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1 2	Are hypervelocity impacts able to produce chondrule-like ejecta?	
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16	Abstract	
17	Chondrules are one of the major components of primitive meteorites. Their sphericity	
18	indicates they formed as molten fragments or droplets but conditions and mechanisms of	
19	chondrule formation remain unknown. A possible scenario is their formation during	
20	hypervelocity impacts and ejections. To challenge this idea, we prepared an experiment that	
21	reproduces analogous of iron metal -rich chondrules by impact between a glassy silicate	
22	projectile and a metallic steel target. The hypervelocity experiment setting allowed an impact	
23	velocity of 5 km/s, and was also designed to collect the ejecta. A scanning electron	
24	microscopy survey shows that silicate ejecta share several similarities with chondrules. They	
25	formed from a population of small melt fragments whose size distribution has the same shape	

as the size distribution of chondrules, with a shift in size: ejecta are about one order of

magnitude smaller than typical chondrules ( $\log(d_{chondrules}/d_{ejecta}) = 1.3^{+0.5}_{-0.7}$ ). We attribute this 27 28 difference in size to the large discrepancy in the size of the impactors (only small 3 mm 29 particle in our experiment versus km-scale planetesimal expected in an impact forming 30 scenario for chondrules). The silicate ejecta formed in the ejecta plume contains numerous 31 small size spherical iron metal beads. Such beads are also observed in numerous chondrules 32 of CO chondrites specifically presented here but also documented in L and LL ordinary 33 chondrites. Size distributions of metal beads in ejecta and chondrules of a carbonaceous 34 chondrite used as reference material (Yamato 81020 CO) display a same shape but with a size 35 shift, quite similar to the one observed between the ejecta droplets and the chondrules: the diameter of metal beads in ejecta is about one order of magnitude smaller than the diameter of 36 the ejecta themselves  $(\log(d_{ejecta}/d_{metal \ beads \ in \ ejecta}) = 1.2^{+0.9}_{-0.8})$ , and the diameter of metal beads 37 38 in chondrules is about one order of magnitude smaller than the diameter of the chondrules themselves  $(\log(d_{chondrules}/d_{metal beads in chondrules}) = 1.4^{+0.6}_{-1.0})$ . We attribute this size differences to 39 40 the blast dynamics: for a same velocity and surface tension, fragments of silicate liquid will be stable when iron liquid fragments of similar size will be separated into smaller droplets. In 41 42 our experiment, the biggest iron metal beads ( $\sim$ 7  $\mu$ m) are within the mean size range of 43 silicate ejecta and can be considered as analogous of the rare large rounded metallic grains 44 (nearly the same size as chondrules) documented in CB chondrites. The textural analogies 45 exposed here provide support for a production of chondrules by impact.

46

#### 47 1. Introduction

Chondrites, the most primitive meteorites, are considered as the main building block of
early planetesimals (e.g. Scott, 2007) and are composed in various proportion – up to 80 vol%
for ordinary chondrites (Jones et al., 2002) – of chondrules. Chondrules are spherules largely
consisting of the silicate minerals olivine and pyroxene formed at high temperatures as

dispersed molten droplets that likely contain much important information on the processes ofthe planetary formation.

Chondrules size distribution is generally quite homogeneous (0.1 to 2 mm in diameter; 54 55 e.g. Teitler et al., 2010). The size distribution of chondrules is typically determined by 56 disaggregation (e.g. Cuzzi et al., 2001), a method that may undercount very small and/or 57 relatively friable chondrule types, and by thin section measurements (e.g. Eisenhour, 1996), a method that requires corrections of the measured size distribution to estimate "true" three-58 59 dimensional grain sizes from random two-dimensional sections. These technics are not accurate in counting very small "microchondrules" (5-40 µm in diameter), found in the matrix 60 or in rare clasts in at least a few chondrites, and "megachondrules" (>1 cm in diameter), 61 discovered as fragments (Nagahara, 1984; Rubin et al., 1982; Ruzicka et al., 1998). 62 Chondrule sizes and masses in various samples have been described as following lognormal 63 64 (King and King, 1979), Rosin or Weibull distributions (Eisenhour, 1996; Hughes, 1978a, 65 1978b) even if Teitler et al. (Teitler et al., 2010) explained that for some chondrule data sets, 66 lognormal or Weibull functions are statistically not good fits of the size distribution. 67 Nevertheless, the size distribution displays a non-zero minimum chondrule size, typically about 10-20% of the mean size (Eisenhour, 1996). Each chemical group of chondrite 68 possesses a distinct size-frequency distribution of chondrules. The mean size is quite uniform 69 70 with ordinary chondrites (L, LL) having the largest, carbonaceous chondrite from Mighei and 71 Ornans groups (CM and CO) having lower average size (King and King, 1978) and 72 carbonaceous chondrite with high metal content (CH) displaying even lower size (e.g. 73 Grossman et al., 1988; Scott, 1988). The process responsible for the formation of chondrules is still a matter of debate and 74 75 it is even not established if a unique or multiple distinct processes are required to form

chondrules. The igneous textures of chondrules indicate that they were produced by melting

77 and solidification in the early solar nebula (e.g. Scott, 2007). There are several hypothesis 78 proposed, including the formation of chondrules by planetesimal collisions (Johnson et al., 79 2018, 2015; Johnson and Melosh, 2014; Urey, 1952), by splashing molten planetesimals 80 where energy for melting is radioactivity (e.g. Asphaug et al., 2011; Sanders and Scott, 2018), by shock waves (Connolly and Love, 1998; Morris and Boley, 2018) or other processes 81 including for instance X-wind model (Shu et al., 1996), lightning (Pilipp et al., 1998), radiant 82 83 heating by magma ocean (Herbst and Greenwood, 2016), or localized magnetized dissipation 84 (Joung et al., 2004).

Krot et al. (2014, 2005) demonstrated that at least some of the latest (~5 Ma post-85 86 Calcium-Aluminium-rich Inclusions (CAI), the oldest substances in the Solar System) ironrich chondrules from Carbonaceous Bencubbin-type (CB) and Carbonaceous "High-metal" 87 chondrites (CH) formed by impact. An impact origin is clearly established, in this particular 88 89 case, and the chondrules and metal grains in the CB chondrites would have formed from a 90 vapor-melt plume produced by a catastrophic disruption between planetary embryos. The 91 remaining question is whether most chondrules are derived from disrupted planetesimals or if 92 CB and CH are exceptions. In other chondrite groups, as in L, LL or CO, chondrules frequently contain iron-nickel metal spherules, distributed in their interiors and/or located at 93 their surface. Where the metal is located at the surface of the chondrules, the shape is not 94 95 spherical (e.g. Uesugi et al., 2008) and the metal forms an asymmetric globule. In most other 96 cases where the metal beads are inside the chondrules, the shape is spherical. Such spherical 97 shape of both chondrules and metal beads is an evidence for a liquid behavior before their 98 solidification (Rambaldi and Wasson, 1981; Skinner and Leenhouts, 1993; Wang et al., 2007) that requires the presence of two distinct and immiscible molten phases. 99

100 Our objective here was to perform a hypervelocity impact experiment to investigate101 the textural properties of ejected material and to compare to the textural properties of iron

metal beads bearing chondrules. Here we aimed at performing impact at a velocity close to 102 103 that expected in the main asteroid belt (Bottke Jr et al., 1994) and a possible order of 104 magnitude of impact velocity in the protoplanetary disk (~ 5.3 km/s). Is it possible to 105 reproduce chondrule-like objects by shooting a silicate projectile into an iron metal target? Is 106 the mechanical process that create a contrasted distribution of metal and silicate liquid 107 droplets in chondrule acting in the hypervelocity impact experiment? This study will help to 108 discuss some physical properties of the spray produced during an impact and the possibility 109 that such chondrules formed from impact during the early Solar System history.

110

#### 111 2. Experimental setup

#### 112 **2.1.** The impact experiment

113 Impact experiment was conducted using a two-stage light-gas gun at Institute of Space and 114 Astronautical Science (ISAS), Japan, where projectile can be accelerated at the impact 115 velocity up to ~7 km/s, which is comparable to typical asteroid-belt collisional velocity 116 (5.3 km/s; Bottke et al., 1994). The projectile was a dunite cylinder, 0.074 g in mass, 3.2 mm 117 in diameter and 3 mm in thickness. The measured impact velocity (4.909 km/s) was close to 118 the one we planned (5 km/s). The target was a steel cylinder used in a previous study (Ganino 119 et al., 2018). The chamber was evacuated to 3 Pa. The trajectory of the projectile and ejecta 120 was captured by a high-speed video camera (Shimazu HPV-X). The camera was operated 121 with frame interval of 2  $\mu$ s with 0.5  $\mu$ s exposure and the spatial resolution is about 0.3 122 mm/pixel. The ejecta plume contains thousands of ejecta that were collected on aluminum 123 witness plates used as "ejecta catchers" (Figure 1). As this unique shoot provides thousands of 124 ejecta to characterize we limited our statistical analysis to one experiment.

125

#### 126 **2.2. Imaging and analyzing major elements using SEM(-EDX)**

The analyses of the ejecta was performed using the Scanning Electron Microscope
Philipps FEI XL30 ESEM LaB6 equipped with a BRUKER Quantax 655 detector, operated at
20 kV and 200 nA beam current at CEMEF-Mines ParisTech.

130 2.3. Particle size distribution determination

The contrasted composition of the iron-rich ejecta, the aluminum plates used as ejecta catcher and the  $SiO_2$  –rich ejecta measured using Energy-dispersive X-ray spectroscopy (Ganino et al., 2018) resulted in contrasted gray level on the backscattered electron image (Figure 2), allowing a convenient detection and identification. We used the particle analyser tool of ImageJ software to count and measure the size distribution of both the  $SiO_2$  ejecta and the metal beads contained within the ejecta.

In this article we call melt "fragments" or melt "droplets" the molten material, droplets 137 being the result of break-up of fragments following Johnson and Melosh (Johnson and 138 139 Melosh, 2014), and "beads" the spherical solidified material (e.g., metal spherules in 140 chondrules or in ejecta). As we wanted to estimate the size of the melt fragments and droplets 141 produced during the impact rather than that of the flattened "splashed" ejecta on the 142 aluminum plate, we measured the surface and estimated the volume of observed splashed 143 ejecta and recalculated the diameter of an equivalent sphere. Assumption for this calculation 144 is that most splashed ejecta were 1 µm- thick as inferred from tilted SEM image of the ejecta 145 catcher. 9326 ejecta were measured. Thirteen ejecta were randomly chosen and focused on for 146 determination of the size distribution of the 1625 metal beads they contained. These thirteen 147 ejecta display a large scale range (Figure 2).

#### 148 **2.4. Distribution of metal beads in chondrules**

The presence of metal beads in chondrules is well known, but their size distribution is not
documented and their abundance largely differs from a chondrite family to another (Scott and
Krot, 2003). We used back-scattered electron image of the carbonaceous chondrite

152 Yamato 81020 (CO 3.0) (Figure 3) to describe the size distribution of metal beads. Yamato 153 81020 is a perfect sample for this study because it is a primitive object (carbonaceous 154 chondrite, see Libourel et al, 2017) poorly altered (no sulfidation, etc), and it is a metal rich-155 chondrite with abundant metal beads in the large majority of its chondrules. In the same way 156 to the procedure that we used to determine the size distribution of the ejecta in the 157 experiment, we used the particle analyser tool of ImageJ software to count and measure the 158 2D-size distribution of metal beads contained within the chondrules of Yamato 81020. Here, 159 the data are apparent diameters measured on a two-dimensional surface and include sections 160 of metal beads. When the sample is sectioned, not all particles/grains are sectioned through 161 their maximum diameter, many appear smaller than they actually are. The size distribution is 162 therefore, inherently biased to smaller sizes. The measured size distribution can be corrected with a stereological analysis. We used, a classical Saltykov analysis (Saltykov, 1958) also 163 called "Schwartz-Saltykov", that allows to compute a statistical equivalent 3D size 164 165 distribution from a 2D size distribution considering that the grains are represented by their 166 equivalent spheres.

167

168 **3. Results** 

169 **3.1.** Chondrule size distribution in CO chondrites

170 CO chondrites have different types of chondrules, some of which are non-spherical
171 and present a wide range of shapes (Figure 3). Rubin and Wasson (2005) show that some
172 chondrules in Yamato 81020 are multi-lobate or distended or highly irregular in two
173 dimensions. The circular chondrules are moderately equant with aspect ratios in thin section
174 varying from 1.02 to 1.11.

The size-frequency distribution of 2834 chondrules in CO chondrites was determined
by petrographic analysis of thin sections by Rubin (1989). The data are not corrected and

- apparent diameters were measured on a two-dimensional surface and display a mean of 178  $148^{+132}_{-70} \mu m$ . We used the initial two dimensional chondrule size distribution of Rubin 179 (Rubin, 1989) transformed into a theoretical three dimensional size distribution after a 180 Saltykov analysis (Saltykov, 1958).
- 181 The chondrule theoretical three-dimensional size distribution best fit is a Rosin182 distribution (Hughes, 1978a) expressed as

$$ch\% = 100 * (1 - e^{-bd^n})$$
 (1)

where ch% is the percentage of chondrules, *d* is the diameter in micron and *b* and *n* are constants and are parameters of the distribution (respectively 1/b being the scale parameter and *n* being the shape parameter). For our best fit, b = -2.60 and n = 1.13. The correlation coefficient of the linear regression in the plot log( $-\ln(1-ch\%) = f(\log(d))$ ) is very good (r = -0.96). Following the results from Hughes (Hughes, 1978a) a better correlation (r = 0.99) is obtained using a 3 parameters Weibull law distribution expressed as

$$ch\% = 100 * (1 - e^{-b(d-\gamma)^n})$$
 (2)

189 where  $\gamma$  is the location parameter and with best fit for  $\gamma = 131 \ \mu m$ , b = -1.35 and n = 0.74. 190 The excellent correlation coefficients demonstrate we are in the case where Rosin and three 191 parameters Weibull laws fit the dataset as very often for chondrules (Hughes, 1978a; Teitler 192 et al., 2010).

193

#### 194 **3.2.** Silicate ejecta size distribution in high velocity impact plume

Ejecta we collected appear as stuck patches on the ejecta catcher with a nearly circular
shape, a typical thickness of 1 μm and an apparent diameter that largely varies from 3.7 μm to
61.2 μm. As all ejecta present a circular shape, we infer they were all formed from splashing
melt fragments and droplets.

Despite the relatively low impact energy due to the small size of the projectile, our 199 200 impact experiment provides abundant melted material. In fact, a precise quantification is not 201 straightforward, but at a first glance, most ejecta have rounded shapes and result from 202 solidification of molten material whereas only rare, always metallic, fragments are angular 203 and could be remnants of unmolten pieces of the target. Because of the exposure duration of 204 the high velocity camera (0.5  $\mu$ s), the ejecta particles with a velocity of 5 km/s moved of ~2.5 205 mm during the exposure. In other words, the ejecta particles appear by 2.5 mm longer than 206 they are. Moreover, the spatial resolution of the camera was not sufficient (>0.3 mm/pixel) 207 to determine the size distribution of ejected material during the flight. Therefore, what we can 208 observe on the images (Figure 1) are only slow (< km/s) and relatively large (> mm) 209 fragments. To characterize the ejecta population, we use the ejecta catcher. Our experimental setting was not appropriate to sample the whole plume of ejecta, but we find only rare angular 210 211 fragments, generally of iron-metal (from the target) and relatively large (>100µm) that were 212 ejected during the impact. The remaining ejecta collected on the aluminum plate have a 213 rounded shape compatible with a pronounced melting of silicate (from the projectile) and 214 metal (from the target) material.

The size of the silicate ejecta recovered on the ejecta catcher and that come from the 215 216 major mass of the ejecta plume is related to fragments and droplets with a diameter in the range ~4-17  $\mu$ m (mean = 7.87  $\mu$ m, median = 6.43<sup>+10.31</sup><sub>-2.49</sub>  $\mu$ m; the confidence interval 217 218 calculated to include 90% of the sample data within). The size distribution of the ejecta is 219 very similar in shape with the size distribution of chondrules (Figure 4) and fits a Rosin distribution as in equation (1). In our best fit for the size distribution of silicate ejecta, b = -220 221 0.88 and n = 1.04. The correlation coefficient of the linear regression in the plot log(-ln(1-222 ej% = f(log(d)) where ej% is the percentage of ejecta and d the diameter is very good

223 (r = 0.97). Here again, a better correlation (r = 0.99) is obtained using a three-parameters

224 Weibull law with  $\gamma = 5.3 \ \mu m$ , b = -0.17 and n = 0.62.

225

#### 226 **3.3.** Metal beads size distribution in ejecta and chondrules of Yamato 81020 (CO 3.0)

The silicate ejecta formed in the ejecta plume contains numerous small size (mostly <1  $\mu$ m) spherical iron metal beads (Figure 2). Metal beads were observed and measured in 13 randomly chosen ejecta. All beads are in the range 0.07-7.0  $\mu$ m with a mean size of 0.55  $\mu$ m and a median of  $0.40^{+0.99}_{-0.26}$   $\mu$ m. The size distribution was plotted for the independently analyzed ejecta (thin dotted-lines in Figure 4d) and for the sum of all beads (1625 beads) in that 13 ejecta (thick dotted-line in Figure 4d). The distribution follows a Rosin size distribution (r = 0.98) with b = 0.36 and n = 0.55 or a three-parameter Weibull distribution

234 
$$(r = 0.99)$$
 with  $\gamma = 0.53 \ \mu m, b = 0.50$  and  $n = 0.36$ 

Even if they are generally much larger (up to few tens of micrometers), such spherical 235 236 metal beads are very frequent in chondrules of the carbonaceous chondrite Yamato 81020 237 (CO 3.0). In that meteorite, many chondrules contain metal bead. We observed and measured 238 the size of metal beads in nine randomly chosen chondrules from Yamato 81020. From the 239 two-dimensional size distribution, we calculated the theoretical three-dimensional size 240 distribution after a Saltykov analysis (Saltykov, 1958). Their diameter range is 2.1 – 81.1 µm with a mean of 9.38  $\mu$ m and a median of 5.65<sup>+23.00</sup><sub>-3.21</sub>  $\mu$ m. Here again, the size distribution was 241 242 plotted for the independently analyzed chondrules (thin dotted-lines in Figure 4c) and for the sum of all beads (3517 beads) in that 9 chondrules (thick dotted-line in Figure 4c). The 243 244 distribution follows with an excellent fit (r = 0.99) a Rosin size distribution with b = -0.77 and 245 n = 0.95 or a three-parameter Weibull distribution (r = 0.99) with  $\gamma = 4.71 \,\mu m$ , b = -0.44 and *n* = 0.77. 246

247

#### 248 **4. Discussion**

#### 249 4.1. Size of chondrules, metal beads and ejecta

Distinct size-frequency chondrule distributions made some authors propose a sorting mechanism, as for instance aerodynamic drag (Dodd, 1976; Shu et al., 1996). Following the idea of an impact origin for chondrules, their mean size would then be explained by the solidification of ejected melt fragments and droplets produced by impactors <10 km in diameter (Benoit et al., 1999; Melosh and Vickery, 1991).

255 The textural homogeneity observed in our ejecta and chondrules, with silicate and 256 metal droplets having similar pattern of size distribution, suggests that material experienced 257 mechanical sorting. The particles might have been size-sorted under virtually identical 258 conditions and following a similar physical process. The material initially shocked to high 259 pressure could have approached the liquid-vapor phase boundary from the liquid side during 260 decompression: the formation of a condensate layer formed from a gas phase during the 261 experiment was documented by Ganino et al., (Ganino et al., 2018). The decompression 262 caused the break-up into an expanding spray of melt fragments and droplets of the ejecta 263 blast. Following Melosh and Vickery (Melosh and Vickery, 1991), the equilibrium droplet 264 radius (R) depends on the size of the impactor (L) and the impact velocity  $(v_{imp})$  following the 265 equation

266

$$R = 0.11 \frac{\sqrt{L}}{v_{imp}} (3)$$

268

With this model, mean size of chondrules would then be explained by the solidification of
ejected melt fragments produced by impactors <10 km in diameter (Benoit et al., 1999).</li>
When we apply this equation to the ejecta produced during the hypervelocity experiment, we
find that expected radius for a 3 mm impactor shot at 4.909 km/s is about 1.2 µm, equivalent

to a diameter of  $\sim 2.4 \,\mu\text{m}$ . This estimate is too small compared to what we observed. Melosh 273 274 and Vickery (Melosh and Vickery, 1991) modeling corresponds to melt ejection as a half 275 sphere expanding into free space, a geometry that closely represents the geometry of a vapor 276 plume, while the ejecta plume in our experiment presents a significantly different geometry 277 (Figure 1). Johnson and Melosh (Johnson and Melosh, 2014) performed impact simulations using the hydrocode modelling iSALE to constrain the properties of the ejecta curtain. To 278 279 calculate the size (diameter) of ejected melt fragments as a function of ejection velocity, 280 impact velocity, and impactor size they proposed the following equation (Johnson and Melosh, 2014), 281

282 
$$d_{fragments} \sim 0.14 \left(\frac{v_{ej}}{v_{imp}}\right)^{-0.81} \left(\frac{R_{imp}}{v_{imp}}\right)^{\frac{2}{3}}$$
(4)

283

When analyzing the high-speed video images, the leading edge of the ejecta cloud had 284 285 velocity of ~4.6 km/s corresponding to an estimate of the ejection velocity used in equation (4). With this value, the theoretical size of melt fragments calculated from Johnson and 286 Melosh (Johnson and Melosh, 2014) equation would be  $d_{\text{fragments}} = 6.69 \,\mu\text{m}$ . The expected 287 diameter for melt fragment is close to the mean diameter of the ejecta we observed 288 (mean = 7.87  $\mu$ m, median = 6.43<sup>+10.31</sup><sub>-2.49</sub>  $\mu$ m) making us propose that the physical process 289 290 described in this equation could explain the observations from our experiment. In the post-291 impact blast, when fragmentation occurs in the ejecta plume, ejected melt fragments can be 292 accelerated by aerodynamic drag as the surrounding vapor flows past the melt. The balance of 293 the aerodynamic drag force and surface tension determines the size of the melt droplets. In 294 that case, the size of the ejecta we collected and measured on the ejecta-catcher depends on 295 the size of the melt fragments in the impact plume that might be governed by the balance 296 between the surface tension and relative kinetic energy. Johnson and Melosh (Johnson and 297 Melosh, 2014), argued that melt "droplets" also form during the blast, as a result of break-up

of larger melt fragments and proposed an equation to estimate the size (diameter) of meltdroplets:

300 
$$d_{droplets} \sim 0.025 \left(\frac{v_{ej}}{v_{imp}}\right)^{-0.97} R_{imp}^{\frac{1}{2}} v_{imp}^{-1}$$
(5)

For our experiment, the theoretical diameter of melt droplets calculated from Johnson and Melosh (Johnson and Melosh, 2014) equation would be  $d_{droplets} = 0.21 \ \mu\text{m}$ . This size is clearly too low to explain the observed size distribution of silicate ejecta. Nevertheless, it is close to the order of magnitude of the diameter of the metal beads found in the ejecta (mean = 0.55; median =  $0.40^{+0.99}_{-0.26} \ \mu\text{m}$ ) that could be explained, following Johnson and Melosh (Johnson and Melosh, 2014) process, by the break-up of larger melt fragments of metal.

In our quantitative study, the difference between the size of chondrules in CO, the size 307 308 of ejecta in hypervelocity experiment and the sizes of metal beads in chondrules and in ejecta 309 can be obtained quantifying the horizontal shift between the distribution curves (Figure 5): the diameter of metal beads in ejecta is ~ 1.2 order of magnitude smaller than the diameter of the 310 ejecta themselves  $(\log(d_{ejecta}/d_{metal beads in ejecta}) = 1.2^{+0.9}_{-0.8})$ , and the diameter of metal beads in 311 312 chondrules is ~ 1.4 magnitude smaller than the diameter of the chondrules themselves  $(\log(d_{chondrules}/d_{metal beads in chondrules}) = 1.4^{+0.6}_{-1.0})$ . These results can also be deduced from the 3 313 314 parameters Weibull law distributions we obtained for our data. The shape parameter (n) is very close for all distribution ( $n_{chondrules} = 0.74$ ;  $n_{ejecta} = 0.62$ ;  $n_{beads in chondrules} = 0.77$ ;  $n_{beads in}$ 315  $_{eiecta} = 0.36$ ), and if we focused on the location parameter ( $\gamma$ ), we confirm that there is a one 316 317 order of magnitude shift between the diameter of chondrules and the diameter of metal beads in chondrules  $(\log(\gamma_{chondrules}/\gamma_{metal beads in chondrules}) = 1.4)$  and between the diameter of 318 ejecta and the diameter of metal beads in ejecta  $\left(\log \left(\frac{\gamma_{ejecta}}{\gamma_{metal beads in ejecta}}\right) = 1\right)$ . 319

That discrepancy would suggest that the physical process that created a contrasted distribution 320 321 of metal beads and silicated chondrules in chondrites might govern similar distributions in the 322 hypervelocity impact experiment. The process proposed by Johnson and Melosh (Johnson and Melosh, 2014) could explains the contrasted size between silicate ejecta and metal beads :it 323 324 would be related to the break-up of melt fragments into smaller melt droplets that would 325 occur for metal melt but not for silicate melt. A way to analyze the flow regime in a 326 multiphase fluid flows is to calculate the Weber number (We). This dimensionless parameter 327 is defined as the ratio of the momentum in the vapor layer divided by the surface tension force 328 restraining the liquid following the equation

- 329
- 330

$$We = \frac{\rho v_f^{2l}}{\sigma} (6)$$

331

332 where  $\rho$  is the density of the fluid,  $v_f$  is its velocity, *l* is its characteristic length, for instance 333 the droplet diameter and  $\sigma$  is the surface tension. There is a critical Weber number (We<sub>c</sub>) 334 above which drop break-up occurs. This critical number depends on the fluid properties but is 335 generally close to 12 (e.g. Pilch and Erdman, 1987). In our experiment during the post-impact 336 blast illustrated by Figure 1, the ejecta were propelled as high velocity relatively thick filaments (Figure 1 and Ganino et al., 2018). The intricate non-spherical shapes of the 337 338 filaments show that the Weber number (We) for the fluid flow is high and should drive a first 339 break-up into small spherical fragments by sheet stripping (Chigier and Reitz, 1996). Density 340 being higher in iron liquid than in silicate liquid, for a same velocity and surface tension, 341 fragments of silicate liquid will be stable (We<sub>silicate melt fragments</sub> < We<sub>critical</sub>) when iron liquid 342 fragments of similar size will be separated by a second break-up from "melt fragments" to "melt droplets" (We<sub>iron melt fragments</sub> > We<sub>critical</sub>). If this scenario of droplets formation by sheet 343 344 stripping is true, the size sorting between the silicate droplet and the smaller metal droplet

345 would be acting and efficient in an extremely short scale (the flight length is about10 cm) and 346 timescale (the flight duration is few tens of  $\mu$ s).

347 These results could explain the mean size of metal beads versus silicate ejecta or 348 chondrules. If we now look at the extreme values, i.e. the biggest metal beads produced 349 during the experiment (size =  $7.0 \,\mu$ m), it is of note that they are very close to the mean and median size of silicate ejecta (mean = 7.87  $\mu$ m, median = 6.43<sup>+10.31</sup><sub>-2.49</sub>  $\mu$ m). The largest metal 350 351 beads in our experiment would derived from melt fragments of metal that did not break-up 352 into smaller melt droplets. A possible explanation would be that their We number was too low 353 (We<sub>iron melt fragments</sub> < We<sub>critical</sub>), and it could be the case if these melt fragments were ejected 354 with a velocity below the average, or if their temperature was slightly different, resulting in 355 distinct surface tension, We number, and mechanical behavior during the blast. In a planetary 356 perspective, large metal beads with a diameter close to the typical diameter of chondrules are 357 also documented and are also relatively rare. Large metal beads are described in some 358 chondrites where their size range is similar to the size of chondrules, as for instance the 359 "large" metallic grains found in CB (Krot et al., 2005).

360

361

#### 4.2. Origin of the metal beads in chondrules

362 In natural samples, as in Yamato 81020, not all chondrules display abundant ironmetal beads (see also Libourel and Portail, 2018). An explanation could be that all chondrules 363 364 initially contained iron droplets, but separation operated to various degrees between melted 365 iron globules and melted chondrules (Uesugi et al., 2008). Size sorting of metal and silicate fractions might suggest that sorting of metallic beads is linked to a certain extent to the 366 367 siderophile depletion in chondrites (following Skinner and Leenhouts, 1993). Another explanation would be that the iron-bead content is directly related to the 368

elementary abundance of iron that may strongly differ from one chondrule to another. In the 369

case chondrules were formed by impact, the chemical composition of the planetary bodies 370 involved in the impact are of first importance: impact involving iron-rich differentiated bodies 371 372 would involve larger volume of metal-rich ejecta. With our experimental setting, it is clear 373 that a significant proportion of metal beads may be molten fragments of the steel target 374 (Ganino et al., 2018). The processes of projectile-target interaction and inter-element 375 fractionation during and immediately after the impact process are now well documented (e.g. Ebert et al., 2013; Evans et al., 1994; Hamann et al., 2013; Mittlefehldt et al., 2005, 1992). 376 377 Nevertheless, another process could be involved in the occurrence of iron-metal, as the high 378 temperature reached during the impact may cause the reduction of FeO contained in the 379 silicate projectiles into iron metal following the reaction:

$$\text{FeO}_{\text{(silicate liquid)}} = \text{Fe}_{\text{(metal liquid)}} + \frac{1}{2} O_{2(\text{gas})} (7)$$

Shock induced in-situ reduction of FeO is well known during hypervelocity and
impact experiments (Ganino et al., 2018; Rowan and Ahrens, 1994) and examples are
described in natural craters on Earth where iron metal spheres are also observed disseminated
in silicate impact melt (e.g., Meteor Crater and Wabar craters, respectively, described by
Mittlefehldt et al., 2005 and Hamann et al., 2013).

Hypervelocity impact experiment carried out in this study provides further support to
an impact forming scenario for chondrule formation: similarly, in catastrophic disruption
between large bodies, immiscibility may occur between the silicate and iron liquids – the
latter forming from iron metal initially present in the colliding bodies or forming from FeO
reduction when the temperature rises.

In this article, we underlined the remarkable analogies between chondrules size
distribution in CO carbonaceous chondrites and ejecta formed in the hypervelocity
experiment. We also mentioned the striking analogies in their metal-beads contents, following
the idea that a similar sorting mechanism could act. Important aspects remain to be discussed

395 as the other proxies we can use to challenge such scenario and the possible chronology in the early history of the Solar System. The thermal history of chondrules brings constraints that 396 397 could help to discriminate between chondrule formation models (Jones et al., 2018). 398 Porphyritic textures are produced when a chondrule is heated to a maximum temperature 399 close to its liquidus (typically 1400 to 1700°C), and cooled at initial rates between about 10 400 and 1,000 °C/h. Nonporphyritic chondrules develop when peak temperatures exceed the 401 liquidus and chondrules cool at rates around 500-3,000 °C/h. In our experiment, ejecta 402 remain glassy; such conditions being closed to those of hyper-quench (cooling rate of several  $10^5$ - $10^6$  °C.s<sup>-1</sup>). What would be expected in larger ejecta blast? A modelling of the radiative 403 404 cooling of a ballistically expanding spherical cloud of chondrule droplets ejected from the 405 impact site is proposed by Dullemond et al. (2014). This article confirm its compatibility with 406 the cooling history of chondrules inferred from the texture, and mineral composition (Jones et 407 al., 2018). Radiative cooling of a droplet cloud produces cooling rates that can be related not 408 only to the different types of texture of the chondrules but also to their relative abundance 409 (Delpeyrat et al., 2019). Focusing on jetted material, Johnson et al. (Johnson et al., 2015) 410 demonstrated that melted droplets jetted during large-scale accretionary impacts would also 411 exhibit the observed igneous textures of chondrules. In a chondrule forming process 412 associated to impact involving molten planetesimals (Asphaug et al., 2011) the cooling rate, 413 limited by opacity, is regulated by the expansion timescale and depends on local swarm 414 density and proximity to the boundary. In this case, chondrules can be used as proxy to 415 constrain the swarm dynamics.

416 Nevertheless, the formation of chondrules from large impact questions the early
417 evolution of the Solar System: chondrule formation would occur when planetesimals and their
418 collisions were abundant and it is generally believed that planetesimals would only accrete
419 once chondrules formed. A support to this scenario, that seems contradictory at a first glance,

420 comes from isotopic analyses of magmatic iron meteorites (Kruijer et al., 2014). These
421 analyses bring support to a possible synchronicity of planetesimal and chondrules formation,
422 showing that the parent bodies of magmatic iron meteorites greater than 10–100 km in
423 diameter could have accreted ~0.1–0.3 Myr after the formation of CAIs. With a such age,
424 some planetesimals would have formed before most chondrules and chondrites would not be
425 the direct primitive accreting blocks of that planetesimals, but a consequence of their
426 accretion.

427

#### 428 Conclusion

The hypervelocity experiment presented here, involving a silicate projectile on an iron 429 430 metal target, produces ejecta that share several similarities with metal beads rich chondrules. These ejecta appear as largely molten splashed ejecta sampled on a screen, coming from a 431 432 population of small melt fragments and droplets. The size distribution of the ejecta is similar 433 to the size distribution of chondrules in CO chondrites, with a shift in size. The ejecta were 434 about an order of magnitude smaller than CO chondrules, a result that is expected by the fact 435 that the equilibrium melt fragments size depends on is the size of the impactor (small 3 mm particle in our experiment versus km-scale planetesimal in an impact forming scenario for 436 437 chondrules). The silicate ejecta formed in the ejecta plume contains numerous small size 438 spherical iron metal beads as consistent with what is frequently observed in chondrules. The 439 biggest iron metal droplets (~7µm) fall in the size mean size range of silicate ejecta and could 440 be considered as analogous of large (same size as chondrules) rounded metallic grains 441 documented in CB chondrites. The size distributions of metal beads in ejecta and chondrules are similar but with a size shift: iron-metal beads in the ejecta are about an order of magnitude 442 443 smaller than iron-metal beads in the chondrule of Yamato 81020. We attribute the difference 444 in mean size between silicate melt (chondrule or ejecta) and metal melt (beads in chondrules

or beads in ejecta), to the break-up from melt "fragments" to melt "droplets". This process
would be efficient for metal melt but not for silicate melt because of their different density
(higher in iron liquid than in silicate liquid). For a same velocity and surface tension,
fragments of silicate liquid were stable when iron liquid fragments of similar size were
separated. The remarkable textural analogies between chondrules and ejecta exposed here
give support for a production of chondrules during impact, implying that chondrites would not
be the direct primitive accreting blocks of the planets, but a byproduct of their accretion.

453

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611	Figures:
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Figure 1: (a) Aluminum witness plate and target after the impact experiment. (b) High-speed
video image before impact. Projectile is elongated along the trajectory due to exposure
duration (0.5 μs). (c) High-speed video image during the post impact blast (~14 μs after
impact)

Figure 2: BSE-SEM images of the aluminum ejecta catcher showing the 13 ejecta analyzed(iron metal bead in white).

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Figure 3: BSE-SEM images of 9 chondrules of Yamato 81020, including metal beads (inwhite).

623 Figure 4: Cumulative size frequency of (a) chondrules in CO chondrites (data from Rubin,

624 1989 transformed into a theoretical three-dimensional size distribution after a Saltykov

625 analysis; n = 2334); (b) ejecta from hypervelocity experiment (n = 9326), (c) metal

beads in chondrules of Yamato 81020 transformed into a theoretical three-dimensional

627 size distribution after a Saltykov analysis (n = 3517; dotted line = individual analyses in

the nine chondrules shown in Figure 3, bold line = mean), (d) metal beads in ejecta

629 (n = 1625; dotted line = individual analyses in the thirteen ejecta shown in Figure 2, bold
630 line = mean).

Figure 5: Size frequency of chondrules in CO chondrules (data from Rubin, 1989 transformed
into a theoretical three-dimensional size distribution after a Saltykov analysis, n = 2834),
ejecta from hypervelocity experiment (n = 9326), metal beads in chondrules of Yamato
81020 transformed into a theoretical three-dimensional size distribution after a Saltykov

635 analysis (n = 3517) and metal beads in ejecta (n = 1625).









