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#### ON SPREADING MODES AND MAGMA SUPPLY AT SLOW AND ULTRASLOW MID-OCEAN RIDGES

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## 10 ABSTRACT

11 The ultraslow eastern Southwest Indian Ridge (SWIR) offers an opportunity to study the effect of 12 magma supply on an ultraslow mid-ocean ridge starting from quasi-melt-free detachment-dominated 13 spreading, and transitioning to volcanic spreading as one nears prominent axial volcanos. Detachments 14 in the quasi-melt-free mode extend along-axis 60 to 95 km and have a lifetime of 0.6 to 1.5 myrs. They 15 cut into their predecessor's footwall with an opposite polarity, causing part of the footwall lithosphere 16 to experience further deformation, hydrothermal alteration, sparse magmatism and possibly thermal 17 rejuvenation, in a hanging wall position. The accretion of the oceanic lithosphere in this context 18 therefore occurs in two distinct stages over the lifetime of two successive detachment faults. We 19 examine the transition from this nearly amagmatic detachment-dominated mode to the more 20 common volcanic mode of spreading, showing that it occurs along-axis over distances  $\leq$  30 km. It 21 involves a significant thinning of the axial lithosphere and a gradual decrease of the amount of tectonic 22 displacement on faults, as the magmatic contribution to the divergence of the two plates increases. 23 We develop a conceptual model of this transition, in which magma plays a double role: it fills the space 24 between the diverging plates, thus reducing the need for displacement along faults, and it modifies 25 the thermal state and the rheology of the plate boundary, affecting its thickness and its tectonic 26 response to plate divergence. Based on a comparison of the ultraslow eastern SWIR, with the faster 27 spreading Mid-Atlantic Ridge, we show that the activation of the volcanic, or of the detachment-28 dominated modes of spreading is connected with the volume of magma supplied per increment of 29 plate separation, over a range of axial lithosphere thickness, and therefore over a range of the M ratio 30 defined by (Buck et al., 2005) as the relative contribution of magma and faults to plate divergence (M 31 is smaller, for a given volume of melt per increment of plate separation, if the plate is thicker). We 32 therefore propose that M does not fully explain the variability in faulting styles observed at slow and 33 ultraslow ridges and propose that rheological changes induced by magma (melt itself is weak, 34 hydrothermally altered gabbro-peridotite mixtures are weak, and melt heat sustains more vigourous 35 hydrothermal circulation) also play a key role, resulting in contrasted potentials for strain localization, 36 footwall flexure on faults and the development of detachment faults.

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38 Keywords: slow and ultraslow mid-ocean ridges, divergent plate boundaries, detachment faults, melt
 39 supply, axial lithosphere, tectonic and magmatic seafloor spreading processes

40

#### 41 INTRODUCTION

42 While all mid-ocean ridges are active volcanic chains, slow (<4 cm/yr) and ultraslow (<2 cm/yr) 43 spreading ridges are also rift zones, where normal faults accommodate part of the plate divergence, 44 interacting with magmatism to shape the newly formed oceanic lithosphere. These ridges are 45 segmented, and significant along-axis variations in melt supply accompany this segmentation, with 46 segment ends receiving less magma than the segments centers (Lin et al., 1990). At segment ends, and 47 in whole segments of melt-poor slow and ultraslow ridge regions, large offset normal faults, also called 48 detachment faults, accommodate a large part of the plate divergence (Cann et al., 1997; Cannat et al., 49 2006; Escartín et al., 2008; Sauter et al., 2013). These detachments bring mantle-derived rocks up 50 through the axial lithosphere, and the resulting seafloor therefore has an ultramafic, non-magmatic 51 component (Cannat, 1993). By contrast, the centers of most slow and ultraslow spreading ridge 52 segments undergo plate divergence through a combination of magma injection and several moderate 53 offset normal faults distributed in the axial domain, and the seafloor exposes volcanic rocks (Smith and 54 Cann, 1999). In the detachment-dominated mode of spreading, specific hydrothermal reactions, 55 associated with the alteration of mantle peridotites, are favored, and feed original microbial 56 communities (Früh-Green et al., 2004). Modes of spreading at slow and ultraslow spreading ridges 57 therefore have a large impact not only on the composition of the oceanic lithosphere, but also on 58 chemical exchanges between the solid Earth and the Ocean, and on the diversity of seafloor biology.

59 Many questions remain about oceanic accretion in the detachment-dominated mode. What controls 60 the geometry and lifetime of detachments, and how do they root into the mantle? How do 61 hydrothermal and magmatic processes operate in this mode and what are their time and space 62 relations to detachment faulting? What is the architecture of the lithosphere that is accreted in the 63 hanging wall plate? Although it is often assumed to be better constrained, there is also a lot we do not 64 yet know about the volcanic spreading mode at slow and ultraslow ridges. While crustal architecture 65 in this mode is commonly proposed to be similar to the layered magmatic crust formed at fast ridges, 66 there is as yet limited direct evidence in support of this hypothesis, and at least one indication that it 67 might not be the case: while volcanic slow spreading ridge regions have a substantial magma supply, 68 they also have a much thicker (up to 7 km-thick at the Mid-Atlantic Ridge; Wolfe et al., 1995) 69 seismogenic lithosphere, compared with the shallow depth to the axial melt lens of faster ridges 70 (Detrick et al., 1987). Finally, while magma supply is identified as a main control on which of the two 71 spreading modes is activated at slow and ultraslow spreading ridges, and although this control has 72 been tested in numerical models (Buck et al., 2005; Olive et al., 2010; Püthe and Gerya, 2014; Tian and 73 Choi, 2017; Tucholke et al., 2008), there is no clear understanding as yet on how it operates in situ: by 74 what melt supply-related mechanism(s) does a slow spreading ridge switch from the volcanic to the 75 detachment-dominated mode of spreading?

76 We address several of these questions based on geological, geophysical and petrological data acquired 77 over the course of several cruises in the eastern region of the Southwest Indian Ridge (SWIR; spreading 78 rate of 14 mm/yr; Cannat et al., 2006), east of the Melville Fracture Zone. This 660 km-long ridge region 79 is a low magma supply end-member of the global mid-ocean ridge system (Cannat et al., 2008). It is 80 characterized by a pronounced focusing of the available melt to axial volcanos, separated by corridors 81 of nearly-amagmatic detachment-dominated spreading where the seafloor exposes almost only 82 ultramafic rocks (Cannat et al., 2006; Sauter et al., 2013). Volcanic areas show numerous and mostly 83 spreading-perpendicular scarps and hummocky volcanic ridges, and correspond to more negative 84 gravity anomalies, consistent with seismic crustal thicknesses up to 8 km (Minshull et al., 2006). Nearly 85 amagmatic spreading corridors have, by contrast, a smooth topography, more positive gravity 86 anomalies (Cannat et al., 2006), and a 2 to 5 km-thick seismic crustal layer that is interpreted as made 87 primarily of serpentinized and fractured mantle-derived rocks (Momoh et al., 2017). These rocks are 88 proposed to have been exhumed by successive axial detachment faults, each cutting into the footwall 89 of the previous fault, with an opposite polarity (Sauter et al., 2013). The eastern SWIR therefore offers 90 opportunities to study the links between magma supply and seafloor spreading modes starting from a 91 quasi melt-free regime, and transitioning to volcanic spreading as one nears the axial volcanos.

The formation of many distal continental divergent margins also involves tectonic extension with only incipient magmatism (eg Peron-Pinvidic et al., 2013). Specifically, seismic reflectors imaged in the Ocean Continent Transition (OCT) of the Iberian and Australian margins may also have formed by detachment faults that alternated polarity (Gillard et al., 2015; Reston and McDermott, 2011). Divergence rates and melt supply during the OCT phase, and during the following onset of indisputable seafloor spreading at these margins, are not strongly constrained. Our results and hypothesis on the nearly amagmatic seafloor spreading mode, and on its transition to volcanic seafloor spreading at an

- 99 ultraslow ridge, may nonetheless inform studies of detachment-dominated divergence and mantle
- 100 exhumation at the OCT phase, and of the transition from OCT to true seafloor spreading.
- 101

### 102 **RESULTS**

#### 103 Along-axis extension, and estimated lifetime of successive axial detachments in a nearly amagmatic

## 104 eastern SWIR spreading corridor

105 We focus on the nearly amagmatic corridor between the 64°E and 65.6°E axial volcanos (Figure 1). The 106 smooth seafloor there forms 500 to 1000 m-high broad ridges, parallel to the ridge axis, which is nearly 107 orthogonal to the north-south trending direction of plate divergence to the west of 65°E, and up to 108 25° oblique to the east (Figure 1). Following (Sauter et al., 2013), we interpret these broad ridges 109 (numbered from 1 to 8 going back in time; Figure 1) as formed by successive axial detachments 110 between about 11 myrs-ago (#8) and the present (#1). The along flowline profile in Figure 2 111 summarizes this tectonic interpretation. Axial detachments #6 (which we propose was initiated about 112 5.6 myrs ago) to #1 (presently active) each cut into the footwall of the previous detachment, with an 113 opposite polarity (Sauter et al., 2013). Proposed breakaways B7 and B8 are both located in the 114 northern, African plate (Figures 1 and 2) and the corresponding exhumed fault surfaces locally bear 115 spreading-parallel corrugations (Cannat et al., 2009). These two south-facing detachments face 116 volcanic conjugate seafloor in the Australian plate and therefore formed under more magmatically 117 active axial conditions than detachments #6 to #1 (Cannat et al., 2009).

118 We pick the top of the inward-facing (toward the axis) slope of each ridge as the best approximation 119 (see Appendix Figure A1) for the location of the breakaway of the corresponding detachment fault, 120 and the base of the outward-facing slope of the corresponding broad ridge as the best approximation 121 for the location of the emergence of each fault at the seafloor. The along-axis extension of the eight 122 identified detachments, based on the length of their breakaway (Figure 1 and Table 1), ranges between 123 25 km (B7) and 95 km (B4). While these breakaways are mostly straight, the inferred emergence traces 124 show convex-outward undulations, with a typical wavelength of 15-20 km (Figure 1). These likely 125 reflect along-axis changes in the rheology of the detachment footwall, impacting the fault's emergence 126 angle. Most detachment-controlled ridges extend up to, and in the case of B5 (to the west) and B8 (to 127 the west and east), into, adjacent volcanic areas (Figure 1). Breakaway ridges B4, and to a lesser extend 128 B2, in the Antarctic plate, are long and follow the overall trend of the ridge axis, curving from E-W in 129 the west, to ENE-WSW in the east (Figure 1). Breakaway ridges B5 and B7, in the African plate, are 130 shorter and relayed to the east by broad ridges that trend NE-SW, oblique to the N-S spreading 131 direction but parallel to the local trend of the ridge axis. A shorter ENE-WSW ridge, also interpreted as a detachment fault block, is found in an intermediate position to the east of B3 and B5 (Figure 1). These
 local complexities are interpreted as due to faulting modes specific to oblique nearly amagmatic
 spreading (Sauter et al., 2013) and are beyond the scope of this paper. We thus focus the following
 analysis and discussions on the orthogonal-spreading part of the smooth seafloor corridor, to the west
 of 65°E (Figure 1).

137 The horizontal offset estimated for past detachments is the along-flowline distance between 138 breakaway and emergence (reconstructed at the time of initiation of the next detachment). For 139 detachment #1, which is still active and at an early stage, the proposed estimate takes mass-wasting 140 into account and is explained in Appendix Figure A1. Estimated horizontal offsets for detachments 141 along the 64.6°E flowline (Figure 2) range between 4 km (active detachment #1; Table 1), and ~21 km 142 (detachments #2 and #8). At a spreading rate of 14 mm/yr (Patriat and Segoufin, 1988), these 143 estimated horizontal offsets correspond to faulting durations (Table 1) between 300 kyrs (detachment 144 #1, still active) and 1.5 myrs (detachment #2) for nearly amagmatic detachments 1 to 6, and of about 145 2.8 myrs, for corrugated detachments 7 and 8. For nearly amagmatic, flipping polarity detachments # 146 1 to 6, the horizontal distance between the emergence of a detachment and the location of the next 147 breakaway is 2.8 to 5.6 km (Table 1), indicating that new detachments typically cut into footwall rocks 148 that had been exhumed 200 to 400 kyrs previously.

149

#### 150 Footwall to hanging wall: two stages in the generation of melt-starved oceanic lithosphere

A consequence of having successive detachments of opposite polarity that cut into their predecessor's footwall is that the inward-facing slopes of smooth seafloor breakaway ridges are made of material that was initially accreted and transported off-axis as part of the footwall of a detachment fault (Stage 1), then became part of the hanging wall of the next detachment fault (Stage 2). The oceanic lithosphere forming these inward-facing slopes is therefore predicted to have been generated in two stages and over the lifetime of two successive detachment faults.

157 Stage 2 involves the residence of previously accreted ultramafic seafloor, in the axial valley region as 158 part of the hanging wall of the next axial detachment (Figure 3). Geological observations made in the 159 present-day axial valley near 64.6°E indicate that Stage 2 results in the formation of small offset normal 160 faults and of isolated volcanic patches and ridges, that are distributed in a > 10 km-wide axial valley 161 (Sauter et al., 2013 and Figure 4). Seismic refraction and reflection results obtained in the same area 162 (Momoh et al., 2017) show lower seismic velocities and several sub-horizontal and north-dipping 163 reflectors in the basement of this near-axis domain, consistent with distributed faulting, hydrothermal 164 alteration and magmatic intrusions in the hanging wall of active detachment #1.

165 The outward-facing slopes of smooth seafloor breakaway ridges record only Stage 1: they form in the 166 footwall of a detachment, and are then captured, and moved off-axis into the footwall of the next 167 (antithetic) detachment. Seafloor ages in the outward-facing slopes of detachment ridges therefore 168 probably increase toward the axis (Figure 3). By contrast, most of the seafloor forming the inward-169 facing slopes of these ridges has probably been modified by Stage 2 axial valley tectonics, 170 hydrothermal alteration and sparse volcanism, over the lifetime of the next, antithetic detachment. 171 Seafloor spreading in the near absence of melt is therefore predicted to result in complex seafloor age 172 patterns at scales of a few tens of kilometers or less (Figure 3). At larger scales, seafloor ages do 173 nonetheless increase off-axis: pickings for the two best expressed magnetic anomalies in the area (3A 174 and 5; Figure 1) are accordingly consistent with the ages predicted in Table 1 for breakaways B6 and 175 B8.

176 This two-stage evolution has two predictable consequences, in addition to the generation of a complex 177 lithosphere geology. One is the potential for thermal rejuvenation of the deep lithosphere during Stage 178 2, as older portions of the footwall of the first detachment find themselves on-axis again in the hanging 179 wall of the next detachment. The other concerns the potential for nearly amagmatic oceanic 180 lithosphere to bear magnetic anomalies. Ultramafic rocks that may have acquired a remanent 181 magnetization upon serpentinization (eg Oufi et al., 2002) in the footwall of the first detachment, 182 would then be subjected to Stage 2 diking and volcanic eruptions in the hanging wall of the next 183 detachment. For example, this could lead to the overprinting of a serpentinization-related magnetic 184 signature recorded up to 2.9 myrs ago at depth in the footwall of detachment #3, by the magnetic 185 signature recorded in volcanic rocks erupted up to 300 kyrs-ago in the hanging wall of detachment #2 186 (Figure 3). This, in addition to the overall low magnetization of serpentinized peridotites (refs), helps 187 explain why the magnetic anomaly record is poor and ambiguous over the ultramafic seafloor of the 188 study area (Bronner et al., 2014).

189

## 190 Along-axis transition from detachment-dominated to volcanic spreading

Figure 5 shows the transition, in the near-axis region, between nearly amagmatic detachmentdominated spreading (along-flowline profile 1), and more magmatic spreading near the 64°E axial volcano (profiles 3 and 4). Profile 1 has an asymmetrical topography, and the seafloor there exposes almost only serpentinized peridotite (Sauter et al., 2013). Faults in the axial valley floor have vertical offsets <200m; volcanism is sparse, forming isolated ridges and patches, and both rock sampling (Paquet et al., 2016; Sauter et al., 2013) and seismic imaging (Momoh et al., 2017) indicate that magmatic rocks form a volumetrically minor part of the basement. Nearly all the plate divergence is

198 therefore currently accommodated by axial detachment #1. By contrast, profiles 3 and 4, less than 30 199 km to the west, are located fully within volcanic seafloor, with several volcanic ridges, and fault scarps 200 that have a maximum vertical throw of 1000 m (Figure 5c). The gravity anomalies are in the negative 201 range (Figure 5b), consistent with a thicker crust (Cannat et al., 2006), and we see no evidence in these 202 profiles that detachment #1 is continuing at depths beneath rafted blocks of volcanic seafloor (Reston, 203 2018; Reston and Ranero, 2012). The axial topography is symmetrical, and seafloor ages from magnetic 204 anomalies (Cannat et al., 2006) increase monotonically with distance to the axis (Figure 5c). These 205 characteristics are consistent with in situ eruption of the volcanics and are typical of seafloor spreading 206 at moderate melt supply, with plate divergence being partitioned into magma emplacement and fault 207 displacement distributed in the axial domain (eg Buck et al., 2005).

208 Profile 2 is in an intermediate position and cuts the western end of detachment #1 (Figure 5a). The 209 seafloor along this profile shows volcanic morphologies, with spreading perpendicular ridges that are 210 clearly in-situ volcanic constructions (Figure 4 and Sauter et al., 2013), and fault scarps ≤ 200 m-high 211 (Figure 5b); yet the gravity anomalies are in the positive range (Figure 5b). These characteristics are 212 consistent with seafloor spreading at relatively low melt supply, with recent plate divergence being 213 accommodated for a part by detachment #1, and for another part by magma emplacement and smaller 214 offset normal faults distributed in the axial valley. The overall topography of the axial region along 215 profile 2 is distinctly asymmetrical across-axis and similar to that of profile 1. This asymmetric 216 topography is probably mostly due to the broad plate bending effect of the active detachment in the 217 nearly amagmatic eastern domain of profile 1. The axial domain between profiles 2 and 1 is 218 characterized by several volcanic ridges (Figure 4), that formed by local eruptions, on older volcanic 219 seafloor in the west, and directly over ultramafic seafloor to the east (Sauter et al., 2013).

220 Our near-axis observations therefore point to a gradual transition, over an along-axis distance of less 221 than 30 km, from the detachment-dominated, nearly amagmatic mode of spreading, to a melt-222 emplacement dominated, symmetrical, mode of spreading. Such a gradual transition indicates that the 223 amount and the distribution of tectonic displacement (one axial detachment, and/or several 224 distributed smaller offset normal faults) at slow spreading ridges are tuned with the melt supply at 225 scales <10 km, so as to complement the melt emplacement contribution to the divergence of the two 226 plates (Figure 6).

227

### 228 **DISCUSSION**

## 229 Full lithosphere sketches of an ultraslow and magma-poor ridge

230 Figure 6 illustrates the concept of a gradual transition from the magmatically robust volcanic spreading 231 mode, to nearly amagmatic, detachment-dominated spreading, with tectonic displacement on faults 232 being tuned over the width of the axial domain, to complement melt emplacement so that the two 233 plates diverge at the spreading rate that is imposed by far-field plate tectonics forces. This is similar to 234 the interpretation proposed by (MacLeod et al., 2009) for the along-axis transition from detachment-235 dominated spreading to more volcanic ridge domains at 13°N on the Mid-Atlantic Ridge (MAR); and an 236 alternative to the rafting detachment model developed by Reston and Ranero (2012) and by Reston 237 (2018), in which axial detachments are proposed to continue at depth beneath rafted blocks of volcanic 238 seafloor at slow and ultraslow ridges. We discard this alternative interpretation because of clear 239 geological evidence that volcanic rocks in the study area were erupted in situ (Figure 4 and recent 240 submersible observations; Cannat et al., AGU abstract 2017).

241 Figure 6 shows the upper 10 kilometers of the axial domain. Microearthquake depths reported by 242 Schlindwein and Schmid (2016), indicate that the seismogenic lithosphere thickens over an along-axis 243 distance of less than 50 km, from less than 10 km-thick beneath the 65.6°E axial volcano (Figure 1), to 244 about 25 km-thick at the eastern edge of the 64.6°E nearly amagmatic corridor. Assuming that a similar 245 thickening occurs with distance from the 64°E volcano (Figure 1), we infer that the volcanic seafloor at 246 the longitude of profiles 3 and 4, and the nearly amagmatic seafloor at the longitude of profile 1 247 probably have seismogenic lithosphere thicknesses of about 10-15 km, and 20-25 km, respectively. 248 Transitioning from the two spreading modes therefore also involves thinning the axial seismogenic 249 lithosphere by at least 10 km, over an along-axis distance of ~30 km. In the following paragraphs, we 250 discuss full-lithosphere conceptual sketches based on the present-day geology of the eastern SWIR 251 axis, for the nearly amagmatic and for the volcanic spreading modes.

252

253 Nearly amagmatic, detachment-dominated spreading mode. Detachment #1 (Figure 1) presently 254 accommodates nearly all the plate divergence. It cuts through the thick axial brittle lithosphere and 255 roots into ductile shear zones that are documented in dredged ultramafic samples (Bickert et al., in 256 prep.) at the base of the plate (Figure 7a). This detachment is at an early stage (~300 kyrs of activity 257 for a typical detachment lifetime of 1 to 1.5 myrs; Table 1) and has so far caused an estimated 258 maximum footwall flexure of 15° (Appendix Figure A1). This flexure could at least partially be 259 accommodated by reactivation of former detachment #2, accounting for the change of slope observed 260 across the proposed emergence of detachment #2 (E2; Figures 1 and 7a). Sub-horizontal and dipping 261 seismic reflectors documented in the axial valley basement are interpreted as minor conjugate faults 262 and magmatic intrusions in the hanging wall of detachment #1, some of which fed the sparse basalt 263 patches observed at the seafloor (Momoh et al., 2017). These basalts are compositionally distinct from those erupted in the nearby volcanic seafloor domains and, based on their composition, it is unlikely that they erupted from dikes propagating along-axis from these volcanic domains (Paquet et al., 2016). Seafloor imagery (Sauter et al., 2013 and Figure 4) and submersible observations (Cannat et al., AGU abstract, 2017) indicate that most basalt patches are spatially associated with faults that guide melts in the upper levels of the lithosphere (Figure 7a). Melt infiltrations, and dikes intruded at greater depths and presumably for the most part into the footwall plate (Figure 7a) are also documented in peridotites samples from the smooth seafloor domains (Paquet et al., 2016).

271 Faults and the associated damage also probably channel hydrothermal fluids, so that serpentinization 272 would occur for a good part in and next to the main axial detachment (Andreani et al., 2007; Rouméjon 273 and Cannat, 2014). Sample studies indicate that serpentinization at slow ridges is a multiphase 274 reaction, occurring for the most part in the most kinetically favorable temperature range of 200-350°C, 275 with an initial stage of formation of the serpentine mesh texture at low fluid-rock ratios, followed by 276 several stages of veining and recrystallization of the initial serpentine at higher fluid-rock ratio 277 (Andreani et al., 2007; Früh-Green et al., 2004; Rouméjon et al., 2014). Based on hypocenter depth 278 distribution in melt-poor domains next to the 65.6°E axial volcano, Schlindwein and Schmid (2016) 279 proposed that serpentinization there initiates at depths of ~15 km below seafloor (Figure 7a). 280 Temperatures in the 200-350° C range could, however, also be found at shallower, near seafloor 281 depths in the upflow zone of high temperature, black smoker-type, hydrothermal circulation cells. 282 Such high temperature hydrothermal upflows have not been discovered so far in the area, and given 283 that magmatic rocks are scarce, they should not represent the dominant hydrothermal regime. Yet 284 black smoker-type circulations may develop transiently to cool isolated melt intrusions and trigger 285 episodes of rapid serpentinization in the surrounding ultramafic rocks, both in the footwall and in the 286 hanging wall plates (Figure 7a). Slower, lower temperature serpentinization, as yet not documented 287 by sampling, may also occur in the tectonically damaged footwall and hangingwall domains.

288 At the current early stage of activity of detachment #1, the axial lithosphere, which we infer formed in 289 the footwall of detachment #2 (Figure 3), is 0.6 to 1.5 myrs-old and could thus be thicker than if 290 controlled solely by the active balance between present-day on-axis heat supply and heat loss. Heat in 291 this spreading mode would be supplied primarily by conduction, by tectonic advection of hot mantle, 292 and by serpentinization reactions. Although advected and latent heat from the sparse melt intrusions 293 could support active black smokers, this efficient mode of hydrothermal cooling would be highly 294 episodic and most hydrothermal circulations would occur at lower fluid-rock ratio in tectonically 295 damaged domains. The resulting axial temperature gradient should therefore be nearly conductive 296 (Figure 7c). Over time and as detachment #1 accommodates more displacement, greater footwall 297 flexure will occur, lowering the fault emergence angle and causing enhanced internal deformation and 298 serpentinization. Also, because of the near-zero melt input, the detachment system should migrate in 299 the direction of the hanging wall and away from the axial region of higher heat input and thinner 300 lithosphere, ultimately favoring the initiation of a new detachment (Buck et al., 2005).

301

302 More magmatically robust volcanic spreading mode. The gravity-derived crustal layer on-axis at the 303 longitude of profiles 3 and 4 is 6 to 7 km-thick (Cannat et al., 2006). Plate divergence there thus 304 concerns an axial lithosphere that is thinner than in the nearly amagmatic setting sketched in Figure 305 7a, but still thicker than the crust. Hence, even if magmatic dikes and gabbro bodies do accommodate 306 most plate divergence in the upper lithosphere, plate divergence acting on the deeper axial lithosphere 307 still requires tectonic uplift of mantle material from the asthenosphere, and through-going melts may 308 form mechanically weak intrusions in this axial lithospheric mantle (Cannat, 1993). Melts intruded 309 higher up, into the brittle lithosphere, will also constitute weak zones while they cool, crystallize, and release the heat necessary to fuel vigorous black smoker-type hydrothermal circulation (eg Lowell et 310 311 al., 2013). The axial lithosphere in this volcanic setting (Figure 7b) should therefore not just be thinner 312 than in the nearly amagmatic detachment-dominated setting. It should also have a different type of 313 geotherm (Figure 7c), with an upper brittle part that is more efficiently cooled by hydrothermal 314 convection cells, and a lower part that is kept weak and thin by frequent melt injections. Temperatures 315 in the upper domain are expected to be variable in space and time (Figure 7c), depending on the 316 activation and geometry of black smoker circulations.

317 Whether magmatic intrusions that form in the lithospheric mantle are subsequently uplifted into the 318 crust, or remain at deep levels of the new plate as it is transported off-axis, should depend on the 319 distribution and offset of axial faults. In Figure 7b, we propose that several of the symmetrically-320 distributed axial normal faults observed in profiles 3 and 4 (Figure 5) root into the melt-intruded axial 321 lithospheric mantle, and may thus lift some of its components into the crust. The axial domain of melt 322 emplacement may thus be thinned by faults to produce the off-axis crust. Therefore, although the melt 323 supply is probably such that the crustal layer off-axis is primarily magmatic, we predict that some of 324 these magmatic rocks have initially been intruded in the lithospheric mantle, then tectonically uplifted 325 into the crust by successive moderate offset normal faults, together with small volumes of their host 326 mantle rocks. Exhumed mantle-derived rocks have actually been dredged in the volcanic seafloor 327 domain a few km to the north west of profile 4 (Sauter et al., 2013).

This calls for a discussion of the applicability of the popular M ratio concept to this, and to other thick lithosphere axial settings. The M ratio as defined by Buck et al. (2005) describes the fraction of the total extension of the brittle lithosphere accommodated magmatically; (1-M) is accommodated by 331 faults. If melt emplacement accommodates plate divergence over the full thickness of the brittle plate, 332 as modelled for slow-spreading ridges by Buck et al. (2005) and Tucholke et al. (2008), M is a direct 333 function of the melt flux (m) per increment of plate separation (m can be expressed as the cumulated 334 thickness of melt provided per increment of spreading). On the other hand, the M ratio does not 335 directly reflect m if melts are either also emplaced below the plate on-axis (the common case at fast 336 and many slow ridges), or emplaced only over a portion of a thick axial brittle plate (probably the 337 common case at ultraslow mid-oceanic ridges). Instead, M then also depends on the thickness of the 338 brittle lithosphere on-axis ( $H_B$ ), and on the thickness of the domain of axial melt emplacement ( $H_M$ ). 339 Olive et al. (2010) addressed this for the  $H_M > H_B$  case by defining a second ratio,  $M_D$  for the contribution 340 of melt emplaced in the ductile domain. In the  $H_M < H_B$  case, M can be expressed as m/H<sub>B</sub>. In the 341 volcanic seafloor configuration of Figure 7b, this formula yields a M value of ~0.6, for  $H_B = 12$  km and 342 assuming that the off-axis crustal thickness (~7 km) is a good approximation for m (ie that this off-343 axis crust is mostly free of mantle-derived peridotites, and that there is very little gabbro trapped in 344 the off-axis mantle lithosphere).

345

346 Transition between the detachment-dominated and volcanic modes of spreading. Profile 2 (Figure 5) 347 corresponds to an along-axis transition between the detachment-dominated (to the east) and volcanic 348 (to the west) modes of spreading. Seafloor imagery and sampling (Sauter et al., 2013) suggest that 349 plate divergence at this longitude occurs by a combination of displacement along the westward 350 continuation of detachment #1, along smaller offset normal faults in the axial valley, and by in-situ 351 magmatic injections forming prominent spreading-perpendicular volcanic ridges (Figure 4). As a result, 352 the geophysically-defined crustal layer should be more geologically composite than in the nearly 353 amagmatic smooth seafloor configuration (Figure 7a), with proportions of serpentinized ultramafic 354 rocks and magmatic intrusions that would depend both of recent magmatic input, and of whether the 355 axial configuration in the recent past was more magmatic (in which case the crust offset by detachment 356 #1 would be mostly magmatic), less magmatic (in which case this crust would be mostly ultramafic) or 357 similarly magmatic.

In this intermediate spreading configuration, more melts would transit through and react with the lower lithospheric axial mantle than in the nearly amagmatic configuration. These deep melts would slowly release their intrinsic and latent heat into the surrounding mantle and would therefore keep the lower axial lithosphere thinner and hotter than in Figure 7a. A further increase in the melt supply would eventually allow for axial faults (and for the hydrothermal systems that develop in permeable brittle rocks) to reach down to very near these melt-rich and hot regions, leading to vigorous hydrothermal circulation. This configuration would ultimately resemble the more magmatically robust

365 volcanic mode of Figure 7b. This leads us to conclude that the effect of melts emplaced in thick 366 lithosphere, ultraslow spreading ridge settings goes beyond the mere mechanical accommodation of 367 plate divergence. Melts thin and weaken the lithosphere (thereby modifying the thickness and 368 strength of the rigid domain in which plate separation needs to be accommodated). Melts emplaced 369 in and near the brittle lithosphere also modify the axial thermal regime by fueling vigorous, if transient, 370 hydrothermal circulation, so that the upper and lower lithosphere are maintained at temperatures 371 respectively cooler and hotter than the conductive geotherm. We propose that these combined effects 372 (rheological and thermal) are key to understand how the two highly contrasted ridge settings proposed 373 in Figure 7 transition along-axis over distances <30 km.

374

## 375 Spreading rate, melt supply, and the spreading mode at slow and ultraslow ridges

376 Spreading rate has a large impact on the thermal regime of mid-ocean ridges (eg Chen and Morgan, 377 1990). With a spreading rate about half of that of the MAR, the SWIR is expected to have a thicker axial 378 lithosphere in places where the melt budget per unit of plate separation is the same, and a fortiori also 379 in places where this melt budget is less (the axial lithosphere should be even thicker due to a lesser 380 input of magmatic heat). In thick axial lithosphere settings, lithosphere thickness affects the (M) ratio 381 between the magmatic and tectonic contributions to plate separation and could therefore be expected 382 to be one of the main factors controlling the spreading mode (Buck et al., 2005). We now assess this 383 effect through a comparison of well-studied axial configurations at the MAR and SWIR.

384 The nearly amagmatic flip-flop detachment (after the terminology of Reston and McDermott, 2011) 385 mode documented at the ultraslow SWIR has not so far been identified at slow ridges, an absence that 386 probably results from more extreme degrees of along-axis melt focusing at ultraslow ridges, another 387 consequence of having a very thick axial lithosphere (Cannat et al., 2003; Standish et al., 2008). The 388 other two modes of slow spreading accretion (volcanic and corrugated detachment modes) are 389 common to both slow and ultraslow settings and therefore offer points of comparison. Microseismicity 390 results acquired at 13°N at the Mid-Atlantic Ridge, and in the 49°E region of the SWIR indicate 391 seismogenic lithosphere thicknesses of up to 15 km and 20 km respectively (Parnell-Turner et al., 2017; 392 Yu et al., 2018). Corrugated, dome-shaped, core-complexes that form in the Antarctic plate at 49°39'E 393 on the SWIR (Zhao et al., 2013), therefore result from detachments that cut through some 20 km of 394 brittle lithosphere, compared to up to 15 km for the 13°N MAR detachments (Figure 8). In the two 395 cases, domal exposed detachment surfaces face volcanic seafloor in the other plate (MacLeod et al., 396 2009; Zhao et al., 2013), and seismic crustal thicknesses are similar (with a large 2-5 km range of 397 variation in the two settings). There is therefore no evidence for a significant contrast between the two settings in terms of the volume of melt supplied per increment of plate separation. Along the same line, the volcanic axial configuration in profiles 3 and 4 of Figure 4 resembles that observed at MAR segment centers, with a symmetrical topography, distributed faults and a crustal thickness of 6-7 km, yet it has an axial seismogenic plate thickness that is probably greater (10 to 15 km; Schlindwein and Schmid, 2016) than documented for volcanic domains of the MAR (less than 8 km; Wolfe et al., 1995), and more similar to that of corrugated detachment MAR domains (Figure 8).

404 These comparisons therefore suggest that the activation of the volcanic, or of the corrugated mode of 405 spreading at slow and ultraslow ridges is connected more to the volume of melt per increment of plate 406 separation, than to the thickness of the axial lithosphere. This is not fully consistent with the 407 hypothesis, proposed by (Buck et al., 2005), that spreading modes are determined by the relative 408 contribution of magma and faults to plate divergence, expressed as the M ratio. M is less, for a given 409 magmatic contribution m (the cumulated thickness of melt provided to the ridge per increment of 410 spreading), if the axial brittle plate is thicker. The volcanic and corrugated modes therefore both 411 appear to be activated at lower M at ultraslow than at slow ridges. The explanation we now explore is 412 that the impact of the axial lithosphere on spreading modes at slow and ultraslow ridges lies not just 413 in its actual thickness, but also in its strength, and that this strength is affected by melt supply.

414 Detachment-dominated modes, corrugated or not, develop for substantial lithosphere thicknesses 415 (Figure 8). MAR corrugated detachments are steep at depth (ca 70°; de Martin et al., 2007; Parnell-416 Turner et al., 2017) but, when mature, emerge at very low angles (<15°; Smith et al., 2006), indicating 417 low effective footwall flexural rigidity. Smooth seafloor detachments in nearly amagmatic corridors of 418 the SWIR experience substantial, but probably not as large, footwall flexure (Momoh et al., in prep.). 419 Given that the axial lithosphere is not thin in either case, such large flexures call for mechanisms that 420 weaken the detachment footwall as faulting proceeds, forming weaker domains in an otherwise thick 421 mechanical lithosphere (Buck, 1988; Lavier et al., 1999). Most slow spreading ridge numerical models 422 to date have implemented this effect by making rock strength a function of strain (Lavier et al., 1999). 423 Geological ground truth comes from studies of samples from the exposed footwall of axial 424 detachments. These studies show that strain localization occurs primarily in assemblages of one or 425 several weak hydrous minerals (serpentine, chlorite, amphibole and talc; Boschi et al., 2006; Escartín 426 et al., 2003; Picazo et al., 2012; Schroeder and John, 2004) that form in the hydrothermally altered 427 brittle lithosphere, at greenschists facies temperatures (<500°C). Talc, specifically, is significantly 428 weaker than serpentine, and is not formed in substantial amounts from peridotites without the input 429 of additional silica, via metasomatic hydrothermal fluids that have altered magmatic rocks, or directly 430 from magmatic veins intruded in the peridotites (Boschi et al., 2006; Picazo et al., 2012). Talc is 431 uncommon in samples from the nearly amagmatic smooth seafloor (Rouméjon et al., 2014). It is by 432 contrast common in samples from corrugated seafloor. Large footwall flexure in the corrugated 433 detachment mode is therefore likely facilitated by talc forming from peridotite gabbro mixtures, while 434 nearly amagmatic detachments would operate with more rigid footwalls (Cannat et al., 2009). Higher 435 melt supply, leading to a transition from the corrugated detachment to the volcanic mode, would also 436 likely strengthen the brittle lithosphere (Cannat et al., 2009) because: 1- the proportion of mantle-437 derived peridotites would become too small to form pervasive serpentine and or talc-bearing zones of 438 weakness; and 2- new magmatic intrusions, which initially form weak melt-mush and ductile gabbro, 439 turn, as they cool, into pristine and strong brittle magmatic rocks (Figure 8). More abundant melt, by 440 contrast, weakens the ductile axial lithosphere. The nearly amagmatic detachment mode is therefore 441 expected to be distinct from both the corrugated detachment, and the volcanic mode, in that it 442 operates with a significantly stronger lower axial lithosphere (Figure 8), which would oppose large 443 flexure and may therefore prevent the development of detachments (Lavier, 2002). Field evidence that 444 detachments exist in the SWIR nearly amagmatic spreading corridors indicates that additional strain 445 weakening mechanisms are activated. Microstructures in ductile and semi-brittle shear zones in 446 dredged ultramafic samples show that dynamic grain size reduction in the ductile mantle lithosphere 447 is a plausible candidate (Bickert et al., submitted).

448

#### 449 **CONCLUSIONS**

450 We propose the following conclusions:

451 1- new oceanic lithosphere in the nearly amagmatic detachment-dominated mode is accreted over the 452 lifetime of two successive detachment faults (ie over 1.5 to 3 myrs) and is for the most part made of 453 material that experienced a two-stages tectonic, hydrothermal, and magmatic evolution: first in the 454 footwall of one detachment fault, then in the hanging wall of the next detachment fault.

2- individual active detachments extend over the whole along-axis width of nearly amagmatic seafloor spreading corridors, up to 95 km in the case studied here. They accommodate up to 20 km of plate divergence, and probably cut through a very thick axial lithosphere (≥25 km). Breakaways are straight, cutting into ultramafic seafloor that was exhumed up to 400 kyrs previously in the former detachment fault's footwall. The emergence of mature detachment faults shows 15-20 km-long undulations, which we propose are due to along-axis changes in the rheology of the detachment footwall, impacting the fault's emergence angle.

3- the along-axis transition from nearly amagmatic, detachment-dominated spreading, to more
magmatic spreading next to axial volcanic centers, occurs gradually over distances of < 30 km through</li>
a fine tuning of tectonic displacement on faults to the local rates of melt emplacement, so as to

465 accommodate the plate divergence. We propose that this gradual transition involves two main 466 processes that are both related to the increase of melt supply: 1- more melt reduces the need for 467 displacement on faults as a mean to accommodate plate divergence; 2- more melt heats and weakens 468 the deep axial lithosphere, thus reducing the thickness of the rigid domain over which divergence 469 needs to be accommodated.

470 4- A comparison with the slow spreading MAR further suggests that the activation of the volcanic or 471 detachment-dominated modes of spreading at slow and ultraslow spreading ridges is primarily 472 controlled by the melt budget (the volume of melt supplied per unit of plate separation), over a range 473 of axial lithosphere thickness (the thickness of the plate over which divergence must be 474 accommodated, decreasing with spreading rate at a given melt budget). Spreading modes are 475 therefore not determined solely by the relative contribution of magma and faults to plate divergence 476 (the M ratio of Buck et al., 2005), since this relative contribution also depends on axial lithosphere 477 thickness. We propose that the rheology of the axial lithosphere is a key additional parameter, that 478 also depends on the melt supply. Melt and melt products modify the strength of the brittle lithosphere, 479 resulting in contrasted potential for strain localization and footwall flexure, the strongest configuration 480 being that of the most magmatically active (volcanic) mode, and the weakest and most favorable to 481 long-lasting detachments and large flexure, being that of the moderately magmatic (corrugated) 482 detachment mode. Melt heat can also sustain vigourous, black smoker type hydrothermal circulations 483 that impact the thermal state of the lithosphere and may favor enhanced hydrothermal alteration. 484 Finally, the presence of melt weakens the ductile mantle at the base of the brittle plate, so that the 485 nearly amagmatic detachment mode probably operates on axial domains that are characterized by a 486 strong ductile lithospheric mantle that is lacking in the more magmatically active corrugated 487 detachment and volcanic modes of spreading.

488

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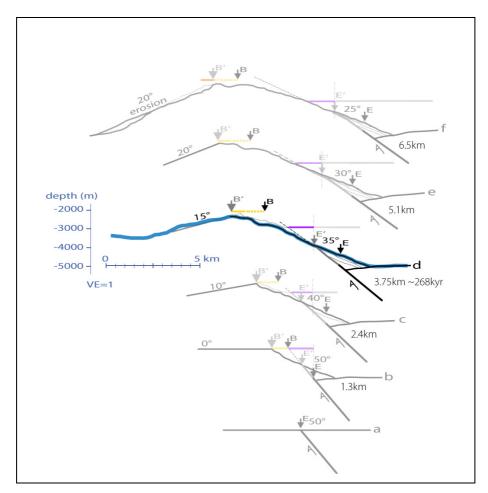
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493

494 APPENDIX

495 Mapping of axial detachment faults and estimating their age and lifetime in melt-poor oceanic496 lithosphere

497 In order to map past detachment faults in the nearly amagmatic spreading corridor of Figure 1, and to 498 estimate their age and lifetime (Table 1), we picked the top of the inward-facing (toward the axis) slope 499 of each detachment-controlled ridge as the best approximation for the breakaway, and the base of 500 their outward-facing slope as the best approximation for the emergence. This approach yields 501 uncertainties, primarily due to the effect of landslides, as illustrated in Figure A1 for presently active 502 detachment #1. For young detachments, the (eroded) breakaways (B in Figure A1) may be a few 503 hundred meters closer to the ridge axis than the top of the inward-facing slope of the detachment 504 ridge, but this discrepancy should diminish and eventually disappear as flexure proceeds so that the 505 outward facing slope of the detachment ridge steepens and also gets subjected to landslides (Figure 506 A1f). Similarly, the emergence of an active detachment may be several hundred meters closer to the 507 breakaway than the base of the detachment-controlled axial relief, which may correspond instead to 508 the front of landslide deposits (Figure A1d). For past detachments, however, this domain has been 509 captured into the footwall of a more recent detachment, probably causing flexure-related deformation 510 and further obscuring the relation between the base of the slope and the past location of the 511 emergence. We estimate that picking the emergence of past detachments at the base of the outward-512 facing slopes is the most straightforward approach, although it may introduce errors of up to 2 km / 513 142 kyrs in the estimated displacement/duration of individual detachments. This is a large error on 514 individual detachments. However, because seafloor spreading in the study area has almost fully been 515 accommodated by detachment faulting from the initiation of detachment #6 (6.1 myrs ago; Table 1), 516 to the present, individual errors made with our method on each detachment must be compensated 517 on the others, within 2 km.



519 Figure A1. Conceptual sketches illustrating the evolution of presently active detachment #1. This 520 detachment is at an early stage. Its proposed present-day configuration is shown in A1d, with the 521 measured axial bathymetric profile in blue (no vertical exaggeration). The dip (~35°) and location of 522 the emergence are based on recent submersible observations (Cannat et al., AGU abstract 2017). 523 Panels A1c to A1a are back in time, and panels A1e to A1f into the near future, assuming a progressive 524 footwall flexure, resulting in a decrease of the faults emergence angle from an initial 50° dip (although 525 not shown, we assume that fault dips also increase with depth); and a 20° slope stability angle for 526 serpentinites (Cannat et al., 2013), which is maintained at each step by landslides so that the eroded 527 cross sectionnal area equals that of landslide deposits. B is the location of the actual breakaway, and 528 B' is the location of the top of the ridge-ward slope at each step. The distance between B and B' results 529 from mass-wasting and is expected to decrease with increasing offset and flexure along the 530 detachment. E is the emergence of the fault with respect to the hanging wall plate, and E' is the 531 emergence from the landslide package. Based on this sketch, our best estimate for the horizontal 532 offset/duration of detachment #1 are about 4 km/300 kyrs (Table 1).

- 533
- 534

535

detachment #	facing	along-axis extension at breakaway km	horizontal distance from emergence to breakaway km	horizontal distance from emergence to breakaway of next detachment km	Estimated total plate divergence km	fault	age of breakaway myr	age of seafloor at breakaway myr
1 (active)	south	85	(4)		(4)	(0.3)	(0.3)	(0.6)
2	north	75	21.0	4.2	21	1.5	1.8	2.2
3	south	60	16.2	5.6	16	1.1	2.9	3.1
4	north	95	14.0	2.8	14	1.0	3.9	4.3
5	south	60	15.1	4.7	15	1.1	5.0	5.2
6	north	60	9.3	3.2	9	0.6	5.6	6.1
7 (corrugated)	south	25	19.5	6.3	(39)	(2.8)	8.4	8.4
8 (corrugated)	south	45	20.6		(41)	(2.9)	(11.3)	(11.3)

536 TABLE 1. Characteristics of successive axial detachments as measured and estimated (see Appendix)

along the profile of Figure 2.

539

## 540 FIGURE CAPTIONS

541

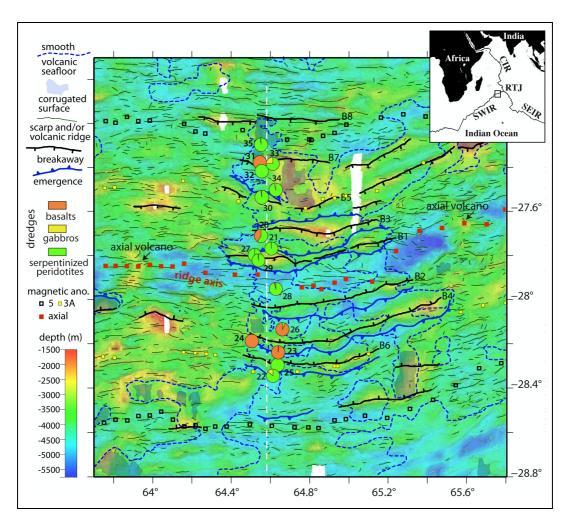


Figure 1. Tectonic sketch map showing the proposed breakaways and emergences of successive detachments in the 64.6°E nearly amagmatic corridor; the volcanic scarps and ridges in the more magmatically active volcanic domains to the east and west; the contours of these volcanic domains, and of several corrugated detachment surfaces and pickings for magnetic anomalies 0, 3A and 5 (Cannat et al., 2006); the recovered lithologies in 16 dredges (Sauter et al., 2013); and the location of the cross-section shown in Figure 2. Detachment faults breakaways (B) and emergences are numbered from 1 (presently active) to 8 (initiated at the time of magnetic anomaly 5; ~11 Myrs ago).

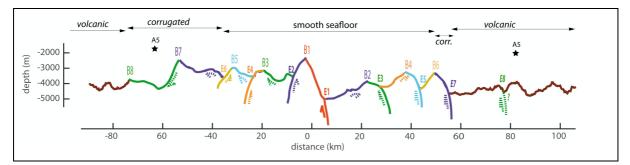
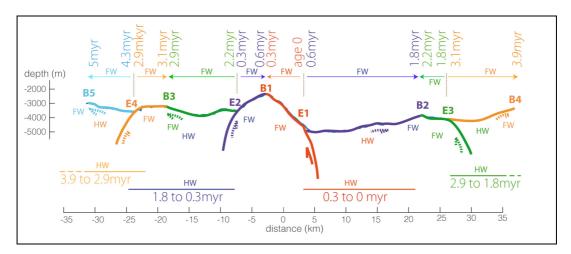


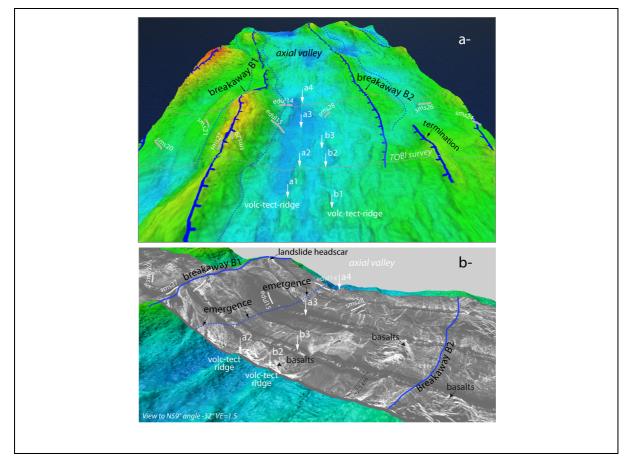
Figure 2. Tectonic interpretation of the along flowline cross-section located in Figure 1, modified after (Sauter et al., 2013), showing the proposed sequence of successive detachments leading to the continuous exposure of ultramafic rocks in the two diverging plates in smooth seafloor domain (detachments #6 to 1), or in the northern, African plate, facing volcanic seafloor in the Antarctic plate (corrugated-volcanic seafloor domains; detachments #8 and 7). Stars: pickings of anomaly A5 in nearest shipboard magnetic profile (Cannat et al., 2006).





557 Figure 3. Tectonic interpretation of the along flowline cross-section located in Figure 1. Zoom to the 558 near-ridge region, showing the ages estimated (text and Table 1) for the initial exposure of ultramafic 559 rocks in the footwall (FW) of each detachment, and for subsequent tectonics, incipient volcanism, and

560 hydrothermal alteration in the hanging wall (HW) of the next detachment.



562 Figure 4. Perpective views of the near axis region in the 64.6°E nearly amagmatic corridor. a-563 bathymetry (depth scale as in Figure 1) with proposed detachment fault breakaways B1 and B2, dredge 564 tracks (in grey, recovered rocks are > 90% ultramafic; (Sauter et al., 2013)), and outlines of the TOBI 565 (Towed Ocean Bottom Instrument) survey (Sauter et al., 2013). b- TOBI reflectivity image drapped on 566 bathymetry. Hummocky basaltic outcrops rest on less reflective ultramafic seafloor. They tend to form 567 spreading-perpendicular ridges (a1-a2 and b1-b2), that align with small offset faults of the ultramafic 568 basement (a2 to a4 and b2 to b3; these faults are more visible in the bathymetry of panel a). The 569 location of the emergence of detachment #1 is based on recent submersible observations (Cannat et 570 al., AGU abstract 2017).

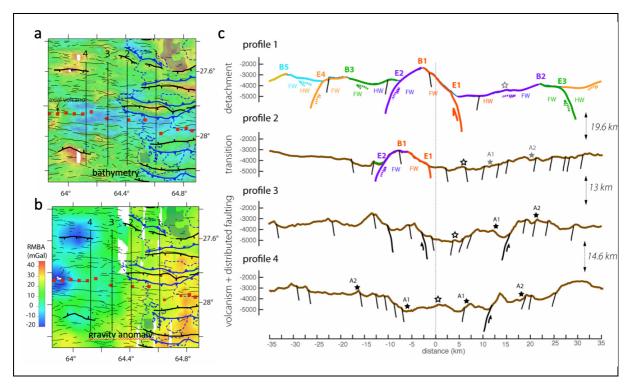
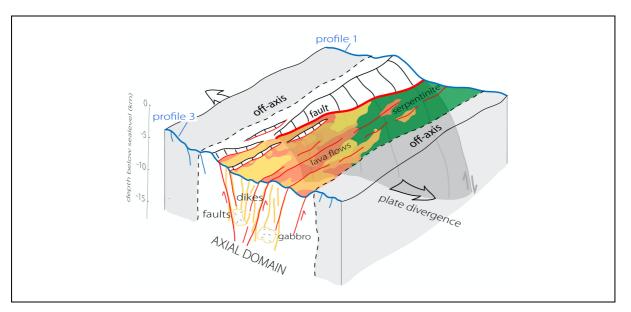
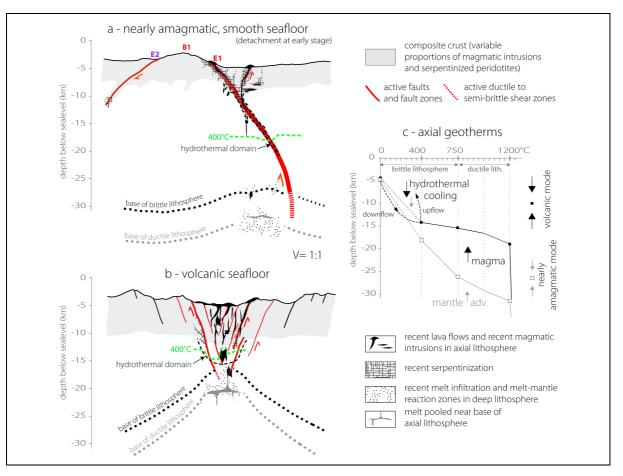


Figure 5. Tectonic interpretation of 4 along flowline cross-sections distributed over the transition from nearly amagmatic seafloor in the east (profile 1), to fully volcanic seafloor in the west (profiles 3 and 4). Panels a and b show the location of these profiles on bathymetry (depth scale, tectonic features and axial magnetic anomaly pickings as in Figure 1) and on a residual mantle Bouguer anomaly map (Cannat et al., 2006). Panel c: seafloor is colored as in Figure 2 for ultramafic seafloor, and in brown for volcanic seafloor; magnetic anomalies 0 (open star), A1 and A2 (closed stars, grey when picking is ambiguous) as in (Cannat et al., 2006).

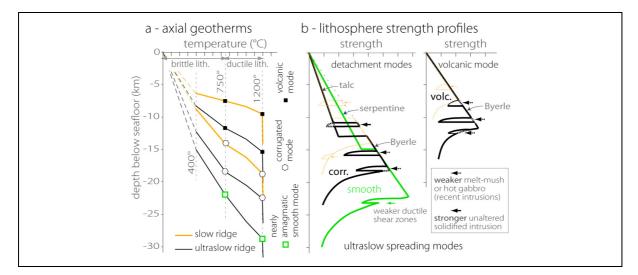


580 Figure 6. 3D perspective sketch of the axial region between across-axis profiles 1 and 3 (in blue; 581 location and tectonic interpretation of these profiles in Figure 5). The concept illustrated here is that

582 tectono-magmatic processes that accommodate plate divergence are tuned so that along-axis 583 variations of melt supply are compensated by greater tectonic displacement on faults. In the nearly 584 amagmatic configuration of profile 1, most plate divergence occurs at a single detachment fault, while 585 in the more magmatically active configuration of profile 3, only 30 km away along-axis, the same total 586 amount of divergence occurs through distributed faulting and magmatic injections. This sketch only 587 goes down ~10 km into basement and therefore does not fully reach to the base of the lithosphere, 588 particularly in the nearly amagmatic configuration, where it probably is >25km-deep (Schlindwein and 589 Schmid, 2016). Full plate 2D sketches for the configurations of profiles 1 and 3 are shown in Figure 7.



591 Figure 7. No vertical exaggeration, full plate 2D conceptual sketches of the ridge axis (a and b) and of 592 axial geotherms (c) for across-axis profiles 1 and 3 (location in Figure 5). The base of the brittle 593 lithosphere in a and b is drawn from constraints on the depth of seismicity (Schlindwein and Schmid, 594 2016). The crust (grey) is a low density, low seismic velocity layer. It is sketched on the basis of gravity 595 and seismic velocity constraints (Cannat et al., 2006; Momoh et al., 2017), and formed by the 596 combination of magmatic flows, intrusions, fault damage, and serpentinization. For clarity, the axial 597 sketches only show the most recent magmatic flows, intrusions, and infiltrations, and the most 598 recently serpentinized domains. See text for explanations.



600 Figure 8. Conceptual sketches of axial geotherms (a) and lithosphere strength profiles (b) for the nearly 601 amagmatic smooth seafloor detachment mode, the weakly magmatic corrugated detachment mode, 602 and the magmatically more robust volcanic spreading mode, at slow and ultraslow ridges. Depth to 603 the base of the brittle lithosphere is based on constraints on the depth of seismicity at the MAR and 604 SWIR (Parnell-Turner et al., 2017; Schlindwein and Schmid, 2016; Wolfe et al., 1995; Yu et al., 2018). 605 Thermal gradients in the lower lithosphere are inferred to be higher at higher melt supply (see Figure 606 7c and text). Lithosphere strength (not scaled) is a function both of lithosphere thickness, and of two 607 types of weakening mechanisms: hydrothermal alteration yielding weaker minerals (simplified here as 608 a choice between serpentine or mechanically weaker talc; see text); and melt injection, which initially 609 weakens the plate as the melt cools, then could make it locally stronger (if cutting though weak, 610 hydrothermally altered exhumed mantle; see text). These weakening mechanisms differ between the 611 3 spreading modes and are proposed to contribute to the transition from one mode to another at 612 varying melt supply (see text).

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