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WELL-POSEDNESS OF SOME NON-LINEAR STABLE DRIVEN SDES

NOUFEL FRIKHA, VALENTIN KONAKOV, AND STÉPHANE MENOZZI

ABSTRACT. We prove the well-posedness of some non-linear stochastic differential equations in the sense of McKean-Vlasov driven by non-degenerate symmetric α -stable Lévy processes with values in \mathbb{R}^d under some mild Hölder regularity assumptions on the drift and diffusion coefficients with respect to both space and measure variables. The methodology developed here allows to consider unbounded drift terms even in the so-called super-critical case, i.e. when the stability index $\alpha \in (0, 1)$. New strong well-posedness results are also derived from the previous analysis.

1. INTRODUCTION

In this article, we are interested in some non-linear stochastic differential equations (SDEs for short) with dynamics

$$(1.1) \quad X_t = \xi + \int_0^t b(s, X_s, [X_s])ds + \int_0^t \sigma(s, X_{s-}, [X_s])dZ_s, \quad [\xi] = \mu \in \mathcal{P}(\mathbb{R}^d)$$

driven by a d -dimensional α -stable process Z , $\alpha \in (0, 2)$, with coefficients $b : \mathbb{R}_+ \times \mathbb{R}^d \times \mathcal{P}(\mathbb{R}^d) \rightarrow \mathbb{R}^d$ and $\sigma : \mathbb{R}_+ \times \mathbb{R}^d \times \mathcal{P}(\mathbb{R}^d) \rightarrow \mathbb{R}^d \otimes \mathbb{R}^d$. Here and throughout the article, we will denote by $\mathcal{P}(\mathbb{R}^d)$ the space of probability measures on \mathbb{R}^d and by $[\theta]$ the law of a random variable θ .

This type of dynamics commonly referred to in the literature as McKean-Vlasov, distribution dependent or mean-field SDEs naturally appears as the limit equation of an individual particle evolving within a large system of particles interacting through its empirical measure, as the size of the population grows to infinity, following the well-known propagation of chaos phenomenon. We refer to the original works by Kac [Kac56] in kinetic theory, by McKean [McK66], [McK67] in non-linear parabolic partial differential equations and by Sznitman [Szn91] in the diffusive case $\alpha = 2$.

The well-posedness theory in the weak or strong sense of such dynamics has been an active research topic during the last decades, see e.g. Funaki [Fun84], Oelschlaeger [Oel84], Gärtner [Gär88], [Szn91], Jourdain [Jou97] and more recently, Li and Min [LM16], Chaudru de Raynal [CdR19], Mishura and Veretennikov [MV18], Lacker [Lac18], Chaudru de Raynal and Frikha [CdRF18b], Röckner and Zhang [RZ18] for a short sample in the diffusive setting. Let us also mention the book of Kolokolstov [Kol10] for some related works on kinetic equations. In the current strictly α -stable context $\alpha \in (0, 2)$, we mention the work of Jourdain *et al.* [JMW08] with applications to the fractional porous media equation on the real-line and of Jourdain *et al.* [JMW05b], [JMW05a] for a probabilistic approach to some non-linear equations involving the fractional Laplacian operator. Let us also mention the work of Graham [Gra92] who derives existence and uniqueness results for non-linear diffusions with Lipschitz coefficients and bounded jump rates.

The classical well-posedness theory of such dynamics is now well-understood in the standard Cauchy-Lipschitz framework on the space $\mathbb{R}^d \times \mathcal{P}_p(\mathbb{R}^d)$, $p \geq 1$, $\mathcal{P}_p(\mathbb{R}^d)$ being the space of probability measures with finite p moments equipped with the Wasserstein distance of order p .

Trying to go beyond the previous setting by weakening the regularity assumptions on the coefficients with respect to both space and measure variables remains an interesting and challenging question that has attracted the attention of the research community. In the diffusive setting $\alpha = 2$, let us mention *e.g.* the work [Sch87] by Scheutzow where a counterexample to uniqueness for the SDE (1.1) is exhibited. Namely, when $\sigma \equiv 0$ and $b(t, x, m) = b(m) = \int_{\mathbb{R}} \bar{b}(y)m(dy)$ for some bounded and locally Lipschitz function \bar{b} , the non-linear SDE (1.1) with random initial condition has several solutions. However, when $\sigma \equiv 1$, the noise helps to restore existence and uniqueness. Indeed, the well-posedness of the corresponding martingale problem when $b(t, x, m) = b(x, m) = \int_{\mathbb{R}^d} \bar{b}(x, y)m(dy)$ has been established by Shiga and Tanaka [ST85]. Such a result has been extended by Jourdain [Jou97] where uniqueness is proved under the more general assumption that b is bounded measurable and Lipschitz with respect to the total variation metric on the space of probability measures, the diffusion coefficient being

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Lipschitz in space and independent of the measure argument. Still in the non degenerate diffusive framework, some extensions have been recently obtained in [MV18] who established some weak and strong well-posedness results for a diffusion coefficient depending only on time and space variables, in [Lac18] and also in [RZ18] for dynamics with a singular convolution type interaction kernel in the drift coefficient. We finally mention the recent contribution [CdRF18b] where well-posedness results in the weak and strong sense are established through a fixed point approach on a suitable complete metric space of probability measures under mild Hölder type regularity assumptions on the coefficients with respect to both space and measure arguments. In particular, weak well-posedness is established when the diffusion coefficient is not Lipschitz with respect to the underlying Wasserstein distance. Some new propagation of chaos results for the approximation of the dynamics (1.1) by a system of particles have recently been obtained by the same authors in their companion paper [CdRF18a]. In [JMW08], the authors extend the classical strong well-posedness results of [Szn91] for non-linear SDEs when both initial condition and Lévy driving process are square integrable. When the Lévy process is not square integrable, as it is the case in our current α -stable framework, existence without uniqueness is proved when both coefficients are Lipschitz with respect to some bounded modification of the Wasserstein metric of order 1.

We here revisit the problem of the unique solvability of the SDE (1.1) by tackling the corresponding formulation of the non-linear martingale problem under some rather mild assumptions on the coefficients in both the spatial and measure arguments. Namely, we will consider Hölder continuous in space coefficients with possibly unbounded drift term b and a Lipschitz regularity condition w.r.t. a suitable Hölder type metric (which actually heavily depends on the spatial Hölder exponents of the coefficients) for the measure arguments. We will also assume some *usual* non-degeneracy and boundedness conditions on the Lévy measure ν of the driving symmetric stable process Z and the *diffusion* coefficient σ appearing in (1.1).

In the current setting, we thus aim at finding a unique probability measure \mathbb{P} on $\mathcal{D}(\mathbb{R}^+, \mathbb{R}^d)$ (Skorokhod space of càdlàg functions) such that denoting by $(\mathbb{P}(t))_{t \geq 0}$ its time marginals, $(y_t)_{t \geq 0}$ the associated canonical process and by $(\mathcal{F}_t)_{t \geq 0}$ its filtration, $\mathbb{P}(0) = \mu$ and for any $\varphi \in C_0^{1,2}(\mathbb{R}^+ \times \mathbb{R}^d, \mathbb{R})$ the process

$$(1.2) \quad \varphi(t, y_t) - \varphi(0, y_0) - \int_0^t (\partial_r + \mathcal{A}_r^{\mathbb{P}}) \varphi(r, y_r) dr$$

is an $((\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$ -martingale starting from 0 at time 0. In the above formulation, $\mathcal{A}_r^{\mathbb{P}}$ is the following integro-differential operator

$$(1.3) \quad \mathcal{A}_r^{\mathbb{P}} \varphi(r, x) = \langle b(r, x, \mathbb{P}(r)), D_x \varphi(r, x) \rangle + \int_{\mathbb{R}^d \setminus \{0\}} (\varphi(r, x + \sigma(r, x, \mathbb{P}(r))z) - \varphi(r, x) - \langle \nabla \varphi(r, x), \sigma(r, x, \mathbb{P}(r))z \rangle \mathbf{1}_{|z| \leq 1}) \nu(dz)$$

acting on smooth test functions and in (1.3) ν is a symmetric stable Lévy measure on \mathbb{R}^d , i.e. for $\alpha \in (0, 2)$, writing in polar coordinates $\zeta = r\xi$, $(r, \xi) \in \mathbb{R}_+^* \times \mathbb{S}^{d-1}$:

$$(1.4) \quad \nu(d\zeta) = \frac{dr}{r^{1+\alpha}} \omega(d\xi),$$

where ω is a symmetric measure on the sphere \mathbb{S}^{d-1} satisfying some non-degeneracy assumptions.

Our main idea consists in applying a fixed point argument to a family of *proxy linearized* martingale problems on a suitable complete metric space of probability measures. In the current strictly stable setting $\alpha \in (0, 2)$, for the well-posedness in the weak and strong sense of the corresponding linear martingale problem, *i.e.* when the dependence with respect to the measure argument is frozen, we can mention among others Tanaka *et al.* [TTW74], Bass [Bas03], Bass *et al.* [BBC04], Mikulevicius and Pragarauskas [MP14], Priola [Pri18] or Zhang *et al.* [HWZ19] for the specific (possibly degenerate) stable case. In our previously mentioned Hölder setting, the main idea to complete our fixed point approach, consists in exploiting some parabolic Schauder estimates but under some *rough* regularity conditions on the final condition, whose regularity will be somehow related to the spatial one of the coefficients. Such estimates are obtained through a forward parametrix perturbation argument in the same spirit as those established by Chaudru de Raynal *et al.* [CdRHM18], [CdRMP19] in the Kolmogorov degenerate diffusive setting or in the α -stable supercritical case.

As a by-product of the usual strong uniqueness results on linear stable driven SDEs, the strong well-posedness of the corresponding non-linear SDE is also derived under suitable regularity assumptions on the coefficients and a non-degeneracy assumption on the Lévy measure. Again, such results will be obtained for a bounded non-degenerate σ , with b, σ being Hölder continuous in space and for potentially unbounded in space drift terms b . Moreover, concerning the measure argument, the maps $\mathcal{P}(\mathbb{R}^d) \ni m \mapsto b(t, x, m)$, $\sigma(t, x, m)$ are also assumed to have a bounded linear functional (or flat) derivative denoted by $[\delta b / \delta m](t, x, m)(y)$ and $[\delta \sigma / \delta m](t, x, m)(y)$ (see Section 2 for a precise definition) which is Hölder-continuous w.r.t. both variables x and y .

Compared to the aforementioned references in the strictly α -stable context, our methodology allows to considerably weaken the regularity assumptions on the coefficients b and σ with respect to both space and measure variables. It also covers a large class of interaction in the measure argument with Hölder type regularity. In particular, no Lipschitz regularity with respect to Wasserstein distance is needed which, to the best of our knowledge, appears to be new. Let us indicate as well that, even in the stable Gaussian case, i.e. $\alpha = 2$ for which $Z = B$ (standard Brownian motion), our approach applies and would provide well-posedness for non-linear SDEs with unbounded drift terms and full dependence on the measure in the diffusion coefficient (see Remark 2.4 below).

The article is organized as follows. The basic definitions together with the assumptions, the main results and the strategy of proof are described in Section 2. The sensitivity analysis w.r.t the measure argument of the semigroups generated by the linearized α -stable driven SDE is carried out in Section 3. The proof of some useful technical results are given in Appendix.

2. OVERVIEW: DEFINITIONS, ASSUMPTIONS AND MAIN RESULTS

In order to tackle the non-linear martingale problem as introduced in (1.2), we will proceed through a fixed point procedure on a suitable complete metric space of probability measure valued flows. Let us now describe this functional space.

The metric space of measure valued flows. We first endow $\mathcal{P}(\mathbb{R}^d)$ with the following metric. Fix $\beta \in (0, 1]$ and define for $\mu, \nu \in \mathcal{P}(\mathbb{R}^d)$,

$$(2.1) \quad \mathbf{d}_\beta(\mu, \nu) := \sup_{f: \|f\|_{C^\beta} \leq 1} \int f(z)(\mu - \nu)(dz),$$

where $\|f\|_{C^\beta} := |f|_\infty + \sup_{(x,y) \in (\mathbb{R}^d)^2, x \neq y} \frac{|f(x) - f(y)|}{|x - y|^\beta}$ denotes the usual Hölder norm on $C^\beta(\mathbb{R}^d, \mathbb{R})$, where $|\cdot|$ stands for the usual Euclidean norm of \mathbb{R}^d (see e.g. Krylov [Kry96]). For $\beta = 1$, \mathbf{d}_β is also known as the Fortet-Mourier distance. Let us introduce now the cost function $c_\beta(x, y) = |x - y|^\beta \wedge 1$, $(x, y) \in \mathbb{R}^d$. Denote by $\Pi(\mu, \nu)$ the set of all transport plans from μ to ν w.r.t c_β . Then, for any $\pi \in \Pi(\mu, \nu)$ and any f satisfying $\|f\|_{C^\beta} \leq 1$, it is readily seen

$$\left| \int f(z)(\mu - \nu)(dz) \right| = \left| \int (f(x) - f(y))\pi(dx, dy) \right| \leq 2 \int c_\beta(x, y)\pi(dx, dy)$$

so that optimizing w.r.t π and f

$$\mathbf{d}_\beta(\mu, \nu) \leq 2\widetilde{W}_\beta(\mu, \nu) := 2 \inf_{\pi \in \Pi(\mu, \nu)} \int c_\beta(x, y)\pi(dx, dy).$$

Let us note that $(\mathcal{P}(\mathbb{R}^d), \mathbf{d}_\beta)$ is a complete metric space¹ and that \mathbf{d}_β metrizes the weak convergence of probability measures.

For fixed $0 \leq s < t < \infty$, we then introduce the set of continuous probability measure valued flows $\mathcal{C}([s, t], \mathcal{P}(\mathbb{R}^d))$. We will from now on denote the elements of $\mathcal{C}([s, t], \mathcal{P}(\mathbb{R}^d))$ with a bold capital letter and equip it with the following uniform metric. For all $\mathbf{P}, \mathbf{P}' \in \mathcal{C}([s, t], \mathcal{P}(\mathbb{R}^d))$,

$$(2.2) \quad \mathbf{d}_{\beta, s, t}(\mathbf{P}, \mathbf{P}') := \sup_{r \in [s, t]} \{\mathbf{d}_\beta(\mathbf{P}(r), \mathbf{P}'(r))\}.$$

Since $(\mathcal{P}(\mathbb{R}^d), \mathbf{d}_\beta)$ is a complete metric space, it follows that $\mathcal{C}([s, t], \mathcal{P}(\mathbb{R}^d))$ equipped with $\mathbf{d}_{\beta, s, t}$ is also complete. We refer to Villani [Vil09] for related issues concerning the aforementioned points. There are many topologies with which we can equip the space $\mathcal{C}([s, t], \mathcal{P}(\mathbb{R}^d))$. The above choice of metric strongly derives from the fact that a Hölder-type regularity assumption on both coefficients b and σ appears to be sufficient to establish the well-posedness in the weak sense of standard linear stable driven SDEs (see e.g. [MP14]). Let us also mention that, beyond our Hölder setting, the choice of the metric is very much related to the spatial smoothness of the coefficients. In particular, when little or no regularity is available in space for the coefficients (bounded or $L^q - L^p$, see e.g. [MV18] or [RZ18]) the *natural* metric for the measure appears to be the total variation one.

¹If $(\mu_k)_{k \in \mathbb{N}}$ is a Cauchy sequence in $\mathcal{P}(\mathbb{R}^d)$ equipped with \mathbf{d}_β , then a slight modification of the proof of Lemma 6.12 in [Vil09] allows to conclude that it is tight so that it admits a subsequence (still denoted by $(\mu_k)_{k \in \mathbb{N}}$) which weakly converges to some measure μ . For each k , introduce an optimal transport plan π_k between μ_k and μ (see e.g. Theorem 4.1 in [Vil09] for the existence of such optimal coupling) for the induced related cost function c_β . By Lemma 4.4 in [Vil09], $(\pi_k)_{k \in \mathbb{N}}$ is tight in $\mathcal{P}(\mathbb{R}^d \times \mathbb{R}^d)$. Hence, up to an extraction, we may assume that $\pi_k \rightarrow \pi$ weakly as $k \rightarrow \infty$. Since each π_k is optimal, Theorem 5.19 of [Vil09] guarantees that π is an optimal (trivial) coupling between μ and μ so that $\limsup_{k \rightarrow \infty} \mathbf{d}_\beta(\mu_k, \mu) \leq 2 \limsup_{k \rightarrow \infty} \int_{(\mathbb{R}^d)^2} c_\beta(x, y)\pi_k(dx, dy) = 0$. Since $(\mu_k)_{k \in \mathbb{N}}$ is a Cauchy sequence with a converging subsequence (w.r.t the metric \mathbf{d}_β), it eventually converges to μ .

Now, for a given $\mu \in \mathcal{P}(\mathbb{R}^d)$, we introduce the subset $\mathbf{A}_{s,t,\mu}$ of $\mathcal{C}([s,t], \mathcal{P}(\mathbb{R}^d))$ defined by

$$(2.3) \quad \mathbf{A}_{s,t,\mu} := \{ \mathbf{P} \in \mathcal{C}([s,t], \mathcal{P}(\mathbb{R}^d)) : \mathbf{P}(s) = \mu \}.$$

The subspace $\mathbf{A}_{s,t,\mu}$ equipped with the metric defined in (2.2) can also be viewed as a complete metric space.

In order to rigorously formulate our regularity assumption on the coefficients w.r.t the measure argument, we introduce the notion of flat derivative of a continuous map $U : \mathcal{P}(\mathbb{R}^d) \rightarrow \mathbb{R}$, $\mathcal{P}(\mathbb{R}^d)$ being equipped with the distance \mathbf{d}_β for some $\beta \in (0, 1]$. This notion will play a key role in our analysis. We refer to Chapter 5 of the monograph by Carmona and Delarue [CD18] for a more detailed discussion on the notion of differentiability of functions of probability measures.

Definition 2.1 (Flat derivative of U). The continuous map $U : \mathcal{P}(\mathbb{R}^d) \rightarrow \mathbb{R}$ is said to have a continuous linear functional derivative if there exists a continuous function $\delta U / \delta m : \mathcal{P}(\mathbb{R}^d) \times \mathbb{R}^d \rightarrow \mathbb{R}$ such that $y \mapsto [\delta U / \delta m](m)(y)$ is bounded, uniformly in m for $m \in \mathcal{K}$, \mathcal{K} being any compact subset of $\mathcal{P}(\mathbb{R}^d)$ and such that for any $m, m' \in \mathcal{P}(\mathbb{R}^d)$,

$$\lim_{\varepsilon \downarrow 0} \frac{U((1-\varepsilon)m + \varepsilon m') - U(m)}{\varepsilon} = \int_{\mathbb{R}^d} \frac{\delta U}{\delta m}(m)(y) (m' - m)(dy).$$

The map $y \mapsto [\delta U / \delta m](m)(y)$ being defined up to an additive constant, we will follow the usual normalization convention $\int_{\mathbb{R}^d} [\delta U / \delta m](m)(y) m(dy) = 0$.

Let us illustrate the structure of this notion of differentiation with a couple of commonly encountered examples.

Example 2.2. (1) If the function U is of scalar form, namely, $U(m) = \int_{\mathbb{R}^d} h(y) m(dy)$ for some real-valued bounded and continuous function h defined on \mathbb{R}^d then it is readily seen that

$$\frac{\delta U}{\delta m}(y) = h(y).$$

(2) If the function U is of quadratic type w.r.t the measure, namely, denoting by \star the usual convolution operator, $U(m) = \int_{\mathbb{R}^d} [h \star m](x) m(dx) = \int_{(\mathbb{R}^d)^2} h(x-y) m(dx) m(dy)$ for some bounded and continuous function h defined on \mathbb{R}^d then one has

$$\frac{\delta U}{\delta m}(y) = \int_{\mathbb{R}^d} h(x-y) m(dx) + \int_{\mathbb{R}^d} h(y-x) m(dx).$$

Remark that if U has a linear functional derivative then for any $m, m' \in \mathcal{P}(\mathbb{R}^d)$, it holds

$$(2.4) \quad U(m) - U(m') = \int_0^1 \int_{\mathbb{R}^d} \frac{\delta U}{\delta m}(\lambda m + (1-\lambda)m')(y) (m - m')(dy) d\lambda.$$

The family of inhomogeneous integro-differential operators. In order to solve locally in time the non-linear martingale problem (1.2), we will perform a fixed point argument in the metric space $\mathbf{A}_{s,t,\mu}$. We thoroughly need to consider some inhomogeneous *linearized* integro-differential operators associated with a given measure flow $\mu \in \mathcal{C}([s,t], \mathcal{P}(\mathbb{R}^d))$. Namely, for $(r, x) \in [s,t] \times \mathbb{R}^d$ and all $\varphi \in C_0^{1,2}([s,t] \times \mathbb{R}^d, \mathbb{R})$, we abuse the definition introduced in (1.3)² and write

$$(2.5) \quad \mathcal{A}_r^\mu \varphi(r, x) \equiv \left(\mathcal{A}(r, \cdot, \mu; D_x) \right) \varphi(r, x) = L_r^{\mu, \alpha} \varphi(r, x) + \langle b(r, x, \mu(r)), D_x \varphi(r, x) \rangle,$$

where

$$(2.6) \quad L_r^{\mu, \alpha} \varphi(r, x) = \text{p.v.} \int_{\mathbb{R}^d} (\varphi(x + \sigma(r, x, \mu(r))\zeta) - \varphi(x)) \nu(d\zeta),$$

and ν is the Lévy measure given by (1.4).

We now introduce some non-degeneracy and regularity assumptions on the Lévy measure and coefficients.

²Observe indeed that \mathbb{P} in (1.3) is a probability measure on $\mathcal{D}(\mathbb{R}^+, \mathbb{R})$ whereas $\mu \in \mathcal{C}([s,t], \mathcal{P}(\mathbb{R}^d))$ is a flow of \mathbb{R}^d -valued probability measures.

Non-degeneracy assumptions.

(ND) There exists $\kappa \geq 1$ s.t. for all $z \in \mathbb{R}^d$:

$$(2.7) \quad \kappa^{-1}|z|^\alpha \leq \int_{\mathbb{S}^{d-1}} |\langle z, \xi \rangle|^\alpha \omega(d\xi) \leq \kappa|z|^\alpha.$$

Note that a wide family of spectral measures satisfy condition (2.7), from absolutely continuous ones to (very) singular ones. For instance, taking $\omega = \Lambda_{\mathbb{S}^{d-1}}$ (Lebesgue measure on the sphere) readily yields the previous control. On the other hand, any measure $\omega = \sum_{i=1}^d c_i(\delta_{e_i} + \delta_{-e_i})$ with positive coefficients $(c_i)_{i \in \llbracket 1, d \rrbracket}$, where the $(e_i)_{i \in \llbracket 1, d \rrbracket}$ stand for the canonical vectors, also satisfies (2.7).

In order to solve the non-linear martingale problem, we will also assume that:

(AC) The jump measure writes: $\nu(dy) = f(y)dy$ for a continuously differentiable function $f : \mathbb{R}^d \rightarrow \mathbb{R}$ with bounded and Lipschitz first order derivative. Since ν is a symmetric stable Lévy measure (see (1.4)), we also have that:

$$(2.8) \quad f(y) = \frac{g(\frac{y}{|y|})}{|y|^{d+\alpha}},$$

where g is an even and continuously differentiable spherical function with a bounded and Lipschitz first order derivative.

We also assume that the *diffusion* coefficient σ in (2.6) satisfies the following uniform ellipticity condition.

(UE) There exists $\Lambda \geq 1$ s.t. for all $(t, z, \mu) \in \mathbb{R}_+ \times \mathbb{R}^d \times \mathcal{P}(\mathbb{R}^d)$ and all $\xi \in \mathbb{R}^d$:

$$(2.9) \quad \Lambda^{-1}|\xi|^2 \leq \langle (\sigma^* \sigma)(t, z, \mu) \xi, \xi \rangle \leq \Lambda |\xi|^2.$$

Regularity assumptions on the coefficients. We will also need some regularity assumptions on the coefficients σ and b w.r.t both the spatial and measure arguments.

We assume that the diffusion coefficient σ satisfies the following conditions:

(D_H) For any $(t, \mu) \in \mathbb{R}_+ \times \mathcal{P}(\mathbb{R}^d)$, the mapping $x \mapsto \sigma(t, x, \mu)$ is bounded and 2η -Hölder continuous, for some $\eta \in (0, 1/2]$, uniformly in $(t, \mu) \in \mathbb{R}_+ \times \mathcal{P}(\mathbb{R}^d)$. Namely, there exists $C \geq 1$ s.t., for all $(t, \mu) \in \mathbb{R}_+ \times \mathcal{P}(\mathbb{R}^d)$:

$$(2.10) \quad \|\sigma(t, \cdot, \mu)\|_{\mathcal{C}^{2\eta}} \leq C,$$

where $\mathcal{C}^{2\eta}(\mathbb{R}^d, \mathbb{R}^m) := \{\varphi \in \mathcal{C}(\mathbb{R}^d, \mathbb{R}^m) : \|\varphi\|_\infty + \sup_{x \neq y} \frac{|\varphi(x) - \varphi(y)|}{|x - y|^{2\eta}} < +\infty\}$ with $m \in \{1, d \times d\}$ denotes the usual Hölder space. When there is no ambiguity, as in (2.10), we simply denote $\mathcal{C}^{2\eta} := \mathcal{C}^{2\eta}(\mathbb{R}^d, \mathbb{R}^m)$ to ease the reading.

For any $(t, x) \in \mathbb{R}_+ \times \mathbb{R}^d$, the continuous map $\mathcal{P}(\mathbb{R}^d) \ni m \mapsto \sigma(t, x, m)$ admits a bounded and continuous flat derivative denoted by the map $[\delta\sigma/\delta m](t, x, m)(\cdot)$ such that $(x, y) \mapsto [\delta\sigma/\delta m](t, x, m)(y)$ is 2η -Hölder continuous, uniformly w.r.t the variables t and m .

We suppose that the drift coefficient b fulfills the following conditions:

(B_H) For any $(t, \mu) \in \mathbb{R}_+ \times \mathcal{P}(\mathbb{R}^d)$, the mapping $x \mapsto b(t, x, \mu)$ belongs to the homogeneous Hölder space

$$\mathcal{C}^{2\eta}(\mathbb{R}^d, \mathbb{R}^d) := \{\varphi \in \mathcal{C}(\mathbb{R}^d, \mathbb{R}^d) : \sup_{x \neq y} \frac{|\varphi(x) - \varphi(y)|}{|x - y|^{2\eta}} < +\infty\}, \quad \eta \in (0, \frac{1}{2}]$$

uniformly w.r.t. the variables (t, μ) . In particular, the drift b may be unbounded in space.

For any $(t, x) \in \mathbb{R}_+ \times \mathbb{R}^d$, the continuous map $\mathcal{P}(\mathbb{R}^d) \ni m \mapsto b(t, x, m)$ admits a bounded and continuous flat derivative denoted by the map $[\delta b/\delta m](t, x, m)(\cdot)$ such that $(x, y) \mapsto [\delta b/\delta m](t, x, m)(y)$ is 2η -Hölder continuous, uniformly w.r.t the variables t and m .

We eventually assume that the following condition, linking the spatial Hölder exponent 2η in assumptions **(D_H)** and **(B_H)** and the stability index α is fulfilled:

(L) It holds that:

$$2\eta + \alpha > 1.$$

We importantly point out that the former condition appearing in **(B_H)** is the natural constraint arising in [CdRMP19] to derive Schauder estimates for a general, and potentially singular, non-degenerate spherical measure ω in the sense of (2.7) in the super-critical case.

We will say that assumption **(A_S)** is satisfied if assumptions **(ND)**, **(AC)**, **(UE)**, **(D_H)**, **(B_H)** and **(L)** hold.

Before going further, we first derive an important control that will play a key role in our analysis of the well-posedness of the non-linear martingale problem related to the SDE (1.1). Under the regularity assumptions **(D_H)**

and (\mathbf{B}_H) , denote by U one of the coefficients σ or b and introduce the map $x \mapsto \delta U(t, x, \mu, \nu) := U(t, x, \mu) - U(t, x, \nu)$, for any fixed $\mu, \nu \in \mathcal{P}(\mathbb{R}^d)$. From (2.4) and the very definition (2.1) of $\mathbf{d}_{2\eta}$, one readily gets

$$(2.11) \quad \forall \beta \in [0, 1], \quad |\delta U(t, \cdot, \mu, \nu)|_\infty \leq \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m) \right\|_{C^{2\eta}} \mathbf{d}_{2\eta}(\mu, \nu) \leq \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m) \right\|_{C^{2\eta}} \mathbf{d}_{(1-\beta)2\eta}(\mu, \nu).$$

Similarly,

$$\delta U(t, x, \mu, \nu) - \delta U(t, x', \mu, \nu) = \int_0^1 \int_{\mathbb{R}^d} \left[\frac{\delta}{\delta m} U(t, x, \mu_\lambda) - \frac{\delta}{\delta m} U(t, x', \mu_\lambda) \right] (y) (\mu - \nu)(dy) d\lambda$$

where we introduced the notation $\mu_\lambda := \lambda\mu + (1-\lambda)\nu$ for $\lambda \in [0, 1]$. Using the fact that $(x, y) \mapsto \frac{\delta}{\delta m} U(t, x, m)(y)$ is 2η -Hölder uniformly w.r.t. the variables t and m , for any $\beta \in [0, 1]$, we obtain

$$\begin{aligned} & \left| \left[\frac{\delta}{\delta m} U(t, x, \mu_\lambda) - \frac{\delta}{\delta m} U(t, x', \mu_\lambda) \right] (y) - \left[\frac{\delta}{\delta m} U(t, x, \mu_\lambda) - \frac{\delta}{\delta m} U(t, x', \mu_\lambda) \right] (y') \right| \\ & \leq 2 \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m)(\cdot) \right\|_{C^{2\eta}} |x - x'|^{2\eta} \wedge |y - y'|^{2\eta} \\ & \leq 2 \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m)(\cdot) \right\|_{C^{2\eta}} |x - x'|^{\beta 2\eta} |y - y'|^{(1-\beta)2\eta}. \end{aligned}$$

Now, from the boundedness of $\frac{\delta}{\delta m} U$ and the uniform 2η -Hölder regularity of $x \mapsto \frac{\delta}{\delta m} U(t, x, \mu)$, we similarly get

$$\left| \left[\frac{\delta}{\delta m} U(t, x, \mu_\lambda) - \frac{\delta}{\delta m} U(t, x', \mu_\lambda) \right] (\cdot) \right|_\infty \leq 2 \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m)(\cdot) \right\|_{C^{2\eta}} |x - x'|^{\beta 2\eta}$$

for any $\beta \in [0, 1]$. From the preceding estimates, one thus deduces

$$(2.12) \quad \forall \beta \in [0, 1], \quad \|\delta U(t, \cdot, \mu, \nu)\|_{C^{\beta 2\eta}} \leq 2 \sup_{t, m} \left\| \frac{\delta}{\delta m} U(t, \cdot, m)(\cdot) \right\|_{C^{2\eta}} \mathbf{d}_{(1-\beta)2\eta}(\mu, \nu).$$

Under (\mathbf{A}_S) , for any fixed measure flow \mathbf{Q} in $\mathcal{C}([0, \infty), \mathcal{P}(\mathbb{R}^d))$, with time marginal $\mathbf{Q}(t)$, and any initial conditions $(s, \mu) \in [0, \infty) \times \mathcal{P}(\mathbb{R}^d)$, there exists a unique solution to the linearized martingale problem, that is, the martingale problem associated with the operator $(\mathcal{A}_t^{\mathbf{Q}} \equiv \mathcal{A}(t, \cdot, \mathbf{Q}; D_x))_{t \geq s}$ introduced in (2.5) is well-posed. We can refer to [CdRMP19] which gives the corresponding Schauder estimates which in turn allows to derive the aforementioned well-posedness following the procedure described in e.g. Mikulevicius and Pragarauskas [MP14] or Priola [Pri12].

Moreover, denoting by $\mathbb{P}_{s,x}^{\mathbf{Q}}$ the unique solution to the martingale problem starting from δ_x at time s , we have that the associated canonical process is strong Markov. The corresponding time-inhomogeneous Markov semigroup will be denoted by $(P_{s,t}^{\mathbf{Q}})_{t \geq s}$.

Theorem 2.3 (Well-posedness of the non-linear martingale problem). Under assumption (\mathbf{A}_S) , for any $s \geq 0$ and any initial distribution $\mu \in \mathcal{P}(\mathbb{R}^d)$, the non-linear martingale problem starting from μ at time s is well-posed for any $\eta \in (0, 1/2]$ if $\alpha \geq 1$ and under the condition $\alpha > 2\eta \vee (1 - \eta)$ if $\alpha < 1$.

Observe that, in the super-critical case $\alpha < 1$ the lowest possible stable index α we can attain is strictly greater than $2/3$.

Corollary 2.1 (Strong uniqueness). Under (\mathbf{A}_S) , if $2\eta \geq 1 - \frac{\alpha}{2}$, then equation (1.1) admits a unique strong solution for any $\alpha \geq 1$ and when $\alpha < 1$ under the additional condition $\alpha > 2\eta \vee (1 - \eta)$.

Proof of Theorem 2.3. Let $T > s$ be fixed and such that $T - s$ is *small*. As already mentioned, the central idea is to use the Banach fixed point theorem on the complete metric space $\mathbf{A}_{s,T,\mu}$ equipped with \mathbf{d} .

We consider the map $\mathbf{T} : \mathbf{A}_{s,T,\mu} \rightarrow \mathbf{A}_{s,T,\mu}$ which to the measure flow $\mathbf{Q} \in \mathbf{A}_{s,T,\mu}$ associates the measure flow $\mathbf{T}(\mathbf{Q}) \in \mathbf{A}_{s,T,\mu}$ induced by the probability measure $\mathbb{P}_s^{\mathbf{Q}}$ given by the unique solution to the linear martingale problem associated with the operator $(\mathcal{A}(t, \cdot, \mathbf{Q}; D_x))$ given by (2.5) starting from the initial distribution μ at time s . Namely, one has $(\mathbf{T}(\mathbf{Q}))(t) = [X_t^{\mathbf{Q}}]$, for any $t \in [s, T]$, where $(X_t^{\mathbf{Q}})_{t \in [s, T]}$ is the unique weak solution to the linear SDE with dynamics

$$X_t^{\mathbf{Q}} = \xi + \int_s^t b(r, X_r^{\mathbf{Q}}, \mathbf{Q}(r)) dr + \int_s^t \sigma(r, X_r^{\mathbf{Q}}, \mathbf{Q}(r)) dZ_s, \quad [\xi] = \mu.$$

In particular, one has $\mathbf{T}(\mathbf{Q})(s) = \mu$ and

$$(2.13) \quad \int h(y) \mathbf{T}(\mathbf{Q})(t)(dy) = \int P_{s,t}^{\mathbf{Q}} h(y) \mu(dy),$$

recalling that $(P_{s,t}^{\mathbf{Q}})_{t \geq s}$ denotes the strong Markov semigroup associated with the operator $(\mathcal{A}(t, \cdot, \mathbf{Q}; D_x))_{t \geq s}$.

Let now $\mathbf{P}_i, i = 1, 2$, be two measure flows on $\mathbf{A}_{s,T,\mu}$. Write from (2.13) and (2.2):

$$\begin{aligned} \mathbf{d}_{\eta,s,T}(\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)) &= \sup_{t \in [s,T]} \sup_{h: \|h\|_{\mathcal{C}^\eta} \leq 1} \int h(y) (\mathbf{T}(\mathbf{P}_1)(t) - \mathbf{T}(\mathbf{P}_2)(t)) (dy) \\ (2.14) \quad &= \sup_{t \in [s,T]} \sup_{h: \|h\|_{\mathcal{C}^\eta} \leq 1} \int (P_{s,t}^{\mathbf{P}_1} - P_{s,t}^{\mathbf{P}_2}) h(y) \mu(dy). \end{aligned}$$

Let us now write for notational convenience $P_{s,t}^{\mathbf{P}_i} \equiv P_{s,t}^i, t \in [s, T]$ for $i = 1, 2$. The following Proposition, whose proof is postponed to Section 3 is the key to prove that \mathbf{T} generates a contraction on $\mathbf{A}_{s,T,\mu}$ in small time.

Proposition 2.1 (Sensitivity Analysis of the semi-groups w.r.t. the measure argument). Under (\mathbf{A}_S) and the additional assumption $\alpha \geq 2\eta \vee (1 - \frac{\eta}{2})$ if $\alpha < 1$, for any fixed $T > 0$, there exists a constant $C := C(T, (\mathbf{A}_S), \alpha, \eta) \geq 1$ ($T \mapsto C(T, (\mathbf{A}_S), \alpha, \eta)$ being non-decreasing) and an exponent $\tilde{\zeta} > 0$ s.t. for all $h \in \mathcal{C}^\eta(\mathbb{R}^d, \mathbb{R})$, $\|h\|_{\mathcal{C}^\eta} \leq 1$ and for any $(t, x) \in [s, T] \times \mathbb{R}^d$:

$$(2.15) \quad |(P_{s,t}^1 - P_{s,t}^2)h(x)| \leq C(t-s)^{\tilde{\zeta}} \mathbf{d}_{\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2).$$

From (2.15) and (2.14), recalling as well from (2.2) that $[s, T] \ni t \mapsto \mathbf{d}_{\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2)$ is a non-decreasing function of t , we eventually derive:

$$\mathbf{d}_{\eta,s,T}(\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)) \leq C(T-s)^{\tilde{\zeta}} \mathbf{d}_{\eta,s,T}(\mathbf{P}_1, \mathbf{P}_2).$$

Hence, we see that for T sufficiently close to s , so that $C(T-s)^{\tilde{\zeta}} < 1$, the map \mathbf{T} is a contraction on the complete metric space $\mathbf{A}_{s,T,\mu}$. It thus admits a unique fixed point that we denote by \mathbf{P}^0 . For $n \geq 1$, if \mathbf{P}^{n-1} is constructed, denote by \mathbf{P}^n the unique fixed-point of \mathbf{T} on $\mathbf{A}_{nT, (n+1)T, \mathbf{P}^{n-1}(nT)}$. Now set $\mathbf{P}(t) := \mathbf{P}^n(t)$ if $t \in [nT, (n+1)T]$ and denote by $\mathbb{P}_s^{\mathbf{P}}$ the unique solution to the linear martingale problem starting from the distribution μ at time s in which the measure dependence is given by \mathbf{P} . By the characterization of the martingale problem, it is readily seen that for any integer n and any $t \in [nT, (n+1)T]$, one has $(\mathbb{P}_s^{\mathbf{P}})^{-1} \circ y_t^{\mathbf{P}} = \mathbf{P}^n(t) = \mathbf{P}(t)$. By the uniqueness of the fixed point, it turns out that $\mathbb{P}_s := \mathbb{P}_s^{\mathbf{P}}$ is also the unique solution to the non-linear martingale problem. This completes the proof. \square

At this step, we importantly emphasize that Proposition 2.1 allows to conclude that Theorem 2.3 is valid under the additional stronger condition $\alpha > 2\eta \vee (1 - \frac{\eta}{2})$ which appears to be quite stringent in the supercritical case $\alpha < 1$. Indeed, in this context, the minimal attainable stability index α is strictly greater than $4/5$. In order to establish the well-posedness of the non-linear martingale problem under the sole condition $\alpha > 2\eta \vee (1 - \eta)$, one has to suitably modify the above argument by performing the fixed point argument using the map $\mathbf{T}^2 = \mathbf{T} \circ \mathbf{T}$ still on the complete metric space $\mathbf{A}_{s,T,\mu}$ equipped with the distance $\mathbf{d}_{\beta,s,T}$ for some well-chosen $\beta \in (0, 2\eta)$. For the sake of clarity, the main modifications of the proof of Proposition 2.1 needed to establish that the iterated procedure provides a contraction are presented in Appendix B.

Proof of Corollary 2.1. Theorem 2.3 provides weak uniqueness for the McKean-Vlasov SDE (1.1). The dependence in the law can then classically be viewed as an inhomogeneous time dependence (see e.g. [CdRF18b]). Namely, (1.1) can be rewritten as:

$$X_t = \xi + \int_0^t \tilde{b}(s, X_s) ds + \int_0^t \tilde{\sigma}(s, X_{s-}) dZ_s,$$

where $\tilde{b}(s, X_s) = b(s, X_s, [X_s])$, $\tilde{\sigma}(s, X_s) = \sigma(s, X_s, [X_s])$. To derive strong well-posedness, we can therefore now rely on the classical inhomogeneous linear setting.

If the coefficients are bounded, the result then readily follows from Chen *et al.* [CZZ17] who proved the complete strong well-posedness for stable driven SDEs, including the super-critical case up to $\alpha \in (0, 2)$. It can also be easily checked that the procedure developed in the homogeneous case by [Pri12] for $\alpha \geq 1$ can readily be extended to the current time dependent framework. The methodology therein indeed only depends on *a priori* smoothness controls that have been established under the current assumptions in [CdRMP19]. Hence, the methodology of [Pri12] to derive strong uniqueness still applies and extends to the current framework with unbounded coefficients and some super-critical cases. \square

Remark 2.4 (On possible extensions). Before going to the proofs of our main results, let us discuss some possible extensions of our main existence and uniqueness results.

- Let us first indicate that the global spatial Hölder continuity of the drift assumed in (\mathbf{B}_H) could be weakened to a *local* Hölder continuity condition. Namely, we could only impose that there exists K_0 s.t. uniformly in $(t, \mu) \in \mathbb{R}_+ \times \mathcal{P}(\mathbb{R}^d)$,

$$|b(t, x, \mu) - b(t, y, \mu)| \leq K_0 |x - y|^{2\eta}, |x - y| \leq 1.$$

This is precisely the condition appearing in [CdRMP19]. The drawback of considering the above condition is that it would induce a localization of our arguments and therefore additional technicality. We preferred to avoid it to focus on the specific non-linear aspects. Anyhow, such an assumption would allow to consider drifts which have linear growth in space, like e.g. in the work by Mishura and Veretennikov [MV18].

- In the subcritical case $\alpha \in (1, 2)$, we believe that the approach of [CdRF18b] could be adapted in order to consider for the drift term a Lipschitz continuity condition w.r.t the total variation metric on the space of probability measures uniformly w.r.t. the time and space variables. In this context, one still has to perform the above fixed point procedure on the space $\mathbf{A}_{s,T,\mu}$ now equipped with the total variation distance. This will concern future investigations.
- Let us point out that the condition $\alpha > 2\eta$ appears in Theorem 2.3 only for integrability purposes in order to establish our Schauder estimates in the context of unbounded drift coefficients. In particular, it could also be removed by employing a localization technique³ in the same spirit of the one employed in [CdRMP19] or by assuming that the drift coefficient is globally bounded.
- Even though we only addressed the pure jump case, the methodology developed below (relying mainly on the control of the explosions of the gradient of the solution of a PDE with Hölder terminal condition and coefficients) would also apply in the diffusive stable setting corresponding to $\alpha = 2$. Coupled to the previously mentioned localization procedure, this could allow to consider non-linear drifts, locally Hölder continuous in space and with spatial linear growth as well as a diffusion coefficient also depending on the measure argument.

3. SENSITIVITY OF THE SEMI-GROUPS W.R.T. THE MEASURE ARGUMENT

The main purpose of this section is to prove Proposition 2.1. To derive the key control (2.15) the main point is formally to expand one semi-group around the other and to exploit some Schauder like controls in the same spirit of those established in [CdRMP19] under the current assumption (\mathbf{A}_S) . With the previous notations, defining for any fixed $t \in [0, T]$ with $T \leq 1$ small enough to be specified later on, $(s, x) \in [0, t] \times \mathbb{R}^d$ and $i \in \{1, 2\}$, $u_i(s, x) := P_{s,t}^{\mathbf{P}_i} h(x)$ it is clear that u_i is a mild solution of equation:

$$(3.1) \quad \begin{cases} \left(\partial_s + \mathcal{A}_s^{\mathbf{P}_i} \right) u_i(s, x) = 0, & (s, x) \in [0, t] \times \mathbb{R}^d, \\ u_i(t, x) = h(x), & x \in \mathbb{R}^d. \end{cases}$$

with $\mathcal{A}_s^{\mathbf{P}_i}$ as in (2.5) and $h \in C^\gamma(\mathbb{R}^d, \mathbb{R})$ where $\gamma \in (0, 1)$ corresponds to the index of the underlying distance employed for the fixed point procedure described in the proof of Theorem 2.3. Namely, for the proof of Proposition 2.1, that is, under the assumption $\alpha > 2\eta \vee (1 - \frac{\eta}{2})$, we will take $\gamma = \eta$. This choice notably allows to derive the announced result in the subcritical case $\alpha \geq 1$ and in the super-critical one up to $\alpha > \frac{4}{5}$. In order to handle the supercritical case $\alpha < 1$ under the sole condition $\alpha > 2\eta \vee (1 - \eta)$, we will take $\gamma = 2\eta$ for the first iterate of the map \mathbf{T} and then $\gamma = (1 - \lambda)2\eta$ for a suitable choice of $\lambda \in [0, 1]$ for the second iterate of \mathbf{T} . The main required modification for the latter procedure is precisely described in Appendix B.

It now follows from the Schauder estimates established in the above reference that, if $h \in C^{\alpha+2\eta}(\mathbb{R}^d, \mathbb{R})$ then $u_i \in L^\infty([0, t], C^{\alpha+2\eta}(\mathbb{R}^d, \mathbb{R}))$. We carefully indicate that the results in [CdRMP19] are stated in the super-critical case $\alpha < 1$ but still hold for $\alpha \in [1, 2)$ (See Remark 6 therein). Also, the estimates of [CdRMP19] are established for a constant *diffusion* coefficient σ . It is anyhow also indicated in Remark 5 of that work that those estimates would extend to dynamics with a variable Hölder in space diffusion coefficient following the usual perturbation approach in that case (see e.g. [MP14] or [ZZ18]). Hence, the estimates are still valid under our current assumption (\mathbf{A}_S) . Furthermore, the techniques we develop below also emphasize how the dependence on the diffusion coefficient must be handled to derive Schauder estimates (see e.g. Lemmas 3.3 and A.3 below which precisely investigate the sensitivity of the non-local term with respect to a diffusion coefficient and the proof of Lemma 3.1).

The point now is that h is here only assumed to be in $C^\gamma(\mathbb{R}^d, \mathbb{R})$ (and not in $C^{\alpha+2\eta}(\mathbb{R}^d, \mathbb{R})$). For the analysis let us now consider a mollification $h^\delta \in C^{\alpha+2\eta}(\mathbb{R}^d, \mathbb{R})$ of $h \in C^\gamma(\mathbb{R}^d, \mathbb{R})$, i.e. $h^\delta = h \star \rho_\delta$ with \star standing for the

³which would actually be the same that the one allowing to consider the local Hölder condition for the drift described in the first point

convolution and for $\rho_\delta(\cdot) = \delta^{-d} \rho(\frac{\cdot}{\delta})$ where $\rho : \mathbb{R}^d \rightarrow \mathbb{R}_+$ is a smooth compactly supported function. In particular, there exists $C \geq 1$ s.t. $\|h^\delta\|_{C^\gamma} \leq C$ and $\|h^\delta - h\|_{C^\gamma} \xrightarrow{\delta \rightarrow 0} 0$. Denoting by $u_i^\delta(s, x) := P_{s,t}^{\mathbf{P}_i} h^\delta(x)$, it is then clear from the uniform Hölder continuity of h that $(u_i^\delta - u_i)(s, x) = P_{s,t}^{\mathbf{P}_i} (h^\delta - h)(x) \xrightarrow{\delta \rightarrow 0} 0$, uniformly in $(s, x) \in [0, t] \times \mathbb{R}^d$. Hence, it suffices to establish (2.15) for the mollified final functions uniformly w.r.t. the smoothing parameter δ .

To this end, observe that $u_i^\delta(s, x)$ satisfies

$$(3.2) \quad \begin{cases} (\partial_s + \mathcal{A}_s^{\mathbf{P}_i}) u_i^\delta(s, x) = 0, & (s, x) \in [0, t] \times \mathbb{R}^d, \\ u_i^\delta(t, x) = h^\delta(x), & x \in \mathbb{R}^d, \end{cases}$$

and $u_i^\delta \in L^\infty([0, t], C^{\alpha+\gamma}(\mathbb{R}^d, \mathbb{R}))$.

To compare both semigroups, we now write the PDE satisfied by $w(s, x) := w^\delta(s, x) = (u_1^\delta - u_2^\delta)(s, x)$. Namely,

$$(3.3) \quad \begin{cases} (\partial_s + \mathcal{A}_s^{\mathbf{P}_1}) w(s, x) = - \left([b(s, x, \mathbf{P}_1) - b(s, x, \mathbf{P}_2)] \cdot D_x u_2^\delta(s, x) \right. \\ \quad \left. + (L_s^{\mathbf{P}_1, \alpha} - L_s^{\mathbf{P}_2, \alpha}) u_2^\delta(s, x) \right) =: -H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(s, x), & (s, x) \in [0, t] \times \mathbb{R}^d, \\ w(t, x) = 0, & x \in \mathbb{R}^d. \end{cases}$$

The important point is now to observe that, proceeding this way we have isolated the terms with the difference of the reference measures in the source (term $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$ in (3.3)). Hence, the above left hand side depends only on one reference measure, here \mathbf{P}_1 , which will serve in order to introduce a proxy associated with the unbounded coefficients of $\mathcal{A}_s^{\mathbf{P}_1}$ similarly to what was previously performed in [CdRMP19] to derive quantitative estimates. Observe that, for the PDE (3.2) with mollified terminal condition, Du_1^δ and $(L_s^{\mathbf{P}_i, \alpha} u_1^\delta)_{i \in \{1, 2\}}$ are indeed well defined pointwise.

The following results will play a central role in our analysis. The first lemma controls the explosions of the gradient and its Hölder modulus of continuity for the solution of (3.2). The second one gives controls of the source $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$ in (3.3) in terms of the distance between the two reference measures, up to multiplicative and integrable time singularities.

Lemma 3.1 (Regularity of the semi-groups for a Hölder source). Assume that $\alpha + \gamma > 1$ and $\gamma \in (0, 2\eta]$. With the notations of (3.2), there exists a constant $C \geq 1$ s.t. for any $\mathbf{P}_i \in \mathcal{C}([0, t], \mathcal{P}(\mathbb{R}^d))$, $i = 1, 2$, for all $(s, x) \in [0, t] \times \mathbb{R}^d$ and all $\delta > 0$:

$$(3.4) \quad |u_i^\delta(s, x)| + (t - s)^{\frac{1}{\alpha} - \frac{\gamma}{\alpha}} |D_x u_i^\delta(s, x)| \leq C.$$

Also, for all $\varepsilon \in (0, 1 - (\frac{1}{\alpha} - \frac{\gamma}{\alpha}))$, there exists C_ε s.t. for all $(s, x, x') \in [0, t] \times (\mathbb{R}^d)^2$:

$$(3.5) \quad (t - s)^{1-\varepsilon} \frac{|D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x')|}{|x - x'|^\vartheta} \leq C_\varepsilon, \quad \vartheta := \vartheta(\alpha, \gamma, \varepsilon) = \alpha \left[1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\gamma}{\alpha} \right) \right] = \alpha(1 - \varepsilon) - 1 + \gamma.$$

Remark 3.1 (About the explosion rates of the gradient). Observe from equation (3.4) that the gradient explodes at an integrable rate in time, precisely from the condition $\alpha + \gamma > 1$ (of course this is meaningful only in the super-critical case $\alpha < 1$, since when $\alpha \geq 1$ it is readily fulfilled by any $\gamma > 0$). Equation (3.5) then specifies how far we can obtain Hölder moduli in space that explodes at an integrable rate. Namely, the gradient can be seen as a function of regularity order $\alpha + \gamma - 1$. We manage to control the corresponding Hölder modulus with an integrable singularity in time up to an additional *small* factor $\varepsilon\alpha$ (exponent ϑ above), where in turn ε quantifies the distance to explosion.

Lemma 3.2 (Controls for the source involving the different measures). Let u_i^δ solve (3.2). Taking then, with the notation introduced in (3.5), $\vartheta > \alpha - 1$, it holds:

$$(3.6) \quad |H_{\mathbf{P}_1, \mathbf{P}_2} u_i^\delta(s, x)| \leq C d_{2\eta}(\mathbf{P}_1(s), \mathbf{P}_2(s)) \left((t - s)^{-1+\varepsilon} \mathbf{1}_{\{\alpha \geq 1\}} + (t - s)^{-\left(\frac{1}{\alpha} - \frac{\gamma}{\alpha}\right)} \right),$$

and for all $(x, x') \in (\mathbb{R}^d)^2$, for any $\lambda \in [0, 1]$

$$(3.7) \quad \begin{aligned} & |H_{\mathbf{P}_1, \mathbf{P}_2} u_i^\delta(s, x) - H_{\mathbf{P}_1, \mathbf{P}_2} u_i^\delta(s, x')| \\ & \leq C d_{(1-\lambda)2\eta}(\mathbf{P}_1(s), \mathbf{P}_2(s)) |x - x'|^{2\lambda\eta} \left((t - s)^{-1+\varepsilon} \mathbf{1}_{\{\alpha \geq 1\}} + (t - s)^{-\left(\frac{1}{\alpha} - \frac{\gamma}{\alpha}\right)} \right) \\ & \quad + C d_{2\eta}(\mathbf{P}_1(s), \mathbf{P}_2(s)) (|x - x'|^\vartheta + |x - x'|^{1+\vartheta-\alpha}) (t - s)^{-1+\varepsilon}. \end{aligned}$$

3.1. Proxy and a priori controls. For given *freezing* parameters $(\tau, \xi) \in [0, T] \times \mathbb{R}^d$ to be specified later on, introduce the flow:

$$(3.8) \quad \theta_{s,\tau}(\xi) = \xi + \int_{\tau}^s b(v, \theta_{v,\tau}(\xi), \mathbf{P}_1(v)) dv, \quad s \geq \tau,$$

$\theta_{s,\tau}(\xi) = \xi$ for $s < \tau$. In the current setting (Hölder drift in space), this object exists from the Peano theorem but is not necessarily unique. Rewrite then (3.3) as follows

$$(3.9) \quad \begin{cases} \left(\partial_v + \tilde{L}_v^{\mathbf{P}_1, \alpha, (\tau, \xi)} \right) w(v, x) + b(v, \theta_{v,\tau}(\xi), \mathbf{P}_1(v)) \cdot D_x w(v, x) \\ = - \left(H_{\mathbf{P}_1, \mathbf{P}_2} u_2^{\delta}(v, x) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, x) \right), \quad (v, x) \in [0, t] \times \mathbb{R}^d, \\ w(t, x) = 0, \quad \text{on } \mathbb{R}^d, \end{cases}$$

where for all $\varphi \in C_0^2(\mathbb{R}^d, \mathbb{R})$

$$(3.10) \quad \tilde{L}_v^{\mathbf{P}_1, \alpha, (\tau, \xi)} \varphi(x) := \text{p.v.} \int_{\mathbb{R}^d} (\varphi(x + \sigma(v, \theta_{v,\tau}(\xi), \mathbf{P}_1(v))\zeta) - \varphi(x)) \nu(d\zeta)$$

and

$$(3.11) \quad R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, x) := (b(v, x, \mathbf{P}_1(v)) - b(v, \theta_{v,\tau}(\xi), \mathbf{P}_1(v))) \cdot D_x w(v, x) + (L_s^{\mathbf{P}_1, \alpha} - \tilde{L}_s^{\mathbf{P}_1, \alpha, (\tau, \xi)}) w(v, x).$$

Precisely, the contribution $R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}$ corresponds, for the measure argument fixed to \mathbf{P}_1 , to the difference of the generator with variable coefficients and the proxy one which is frozen along the deterministic flow (3.8). The following Lemma proved in Appendix A.4 gives quantitative bounds on this remainder term.

Lemma 3.3 (Bounds for the difference of the generators for a fixed measure argument). There exists a constant C s.t. for all $(\tau, \xi), (v, y) \in [0, T] \times \mathbb{R}^d$:

$$(3.12) \quad \begin{aligned} |R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)| &\leq C \left(\|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{C}^{2\eta}} |D_y w(v, y)| + \|\sigma(v, \cdot, \mathbf{P}_1(v))\|_{C^{2\eta}} [(\|w\|_{\infty} + \|D_x w(v, \cdot)\|_{\infty}) \mathbf{1}_{\{\alpha < 1\}} \right. \\ &\quad \left. + \|D_x w(v, \cdot)\|_{C^{\vartheta}} \mathbf{1}_{\{\alpha > 1\}} + \|w(v, \cdot)\|_{C^{1+\vartheta}} \mathbf{1}_{\{\alpha = 1\}} \right] |y - \theta_{v,\tau}(\xi)|^{2\eta}. \end{aligned}$$

Also, under (\mathbf{A}_S) , it is clear that the time-dependent operator $\tilde{L}_v^{\mathbf{P}_1, \alpha, (\tau, \xi)} + b(v, \theta_{v,\tau}(\xi), \mathbf{P}_1(v)) \cdot D_x$ generates a family of transition probability (or two parameter transition semi-group) $(\tilde{P}_{s,v}^{\mathbf{P}_1, \alpha, (\tau, \xi)})_{0 \leq s \leq v \leq t}$. Because of the non-degeneracy conditions of (2.7) and (2.9) in **(ND)** and **(UE)**, the associated heat-kernel exists (see e.g. [Kol00], [Wat07] or Appendix A below) and, for fixed $0 \leq s < v \leq t$, it writes:

$$(3.13) \quad \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) = p_{\Theta_{s,v}, (\tau, \xi), \mathbf{P}_1} \left(y - m_{v,s}^{(\tau, \xi)}(x) \right),$$

where we denoted:

$$(3.14) \quad m_{v,s}^{(\tau, \xi)}(x) := x + \int_s^v b(r, \theta_{r,\tau}(\xi), \mathbf{P}_1(r)) dr,$$

$$(3.15) \quad \Theta_{s,v}, (\tau, \xi), \mathbf{P}_1 := \int_s^v \sigma(r, \theta_{r,\tau}(\xi), \mathbf{P}_1(r)) dZ_r,$$

and where Z is a stable process with Lévy measure ν defined on some probability space $(\Omega, \mathcal{A}, \mathbb{P})$.

Observe that from the definition (3.14) of the *shift* $m_{v,s}^{(\tau, \xi)}(x)$ we have the important property

$$(3.16) \quad m_{v,s}^{(\tau, \xi)}(x)|_{(\tau, \xi) = (s, x)} = \theta_{v,s}(x).$$

Let us first give some bounds on the density function of the stochastic integral (3.15) and its derivatives. They are somehow classic under the absolute continuity condition **(AC)** for the Lévy measure (see e.g. [Kol00]). For the sake of completeness, the proof of the following result is postponed to Appendix A.

Lemma 3.4 (Controls on the proxy density and its derivatives). Assume **(ND)** and **(UE)** hold. Then, there exists a constant C and a probability density $\bar{q} : \mathbb{R}_*^+ \times \mathbb{R}^d \rightarrow \mathbb{R}$ satisfying for all $\gamma < \alpha$ the integrability property

$$(3.17) \quad \int_{\mathbb{R}^d} |z|^\gamma \bar{q}(t-s, z) dz \leq C(t-s)^{\frac{\gamma}{\alpha}},$$

s.t. for all multi-index $\beta, |\beta| \leq 3$ and for any $0 \leq s < t, (\tau, \xi) \in \mathbb{R}_+ \times \mathbb{R}^d, \mu \in \mathcal{C}(\mathbb{R}^+, \mathcal{P}(\mathbb{R}^d)), y \in \mathbb{R}^d$:

$$(3.18) \quad |D_y^\beta p_{\Theta_{s,t}, (\tau, \xi), \mu}(y)| \leq \frac{C}{(t-s)^{\frac{|\beta|}{\alpha}}} \bar{q}(t-s, y).$$

Also, any *diagonal perturbation* does not affect the previous estimate. Namely, for $z \in \mathbb{R}^d$ such that $|z| \leq (t-s)^{\frac{1}{\alpha}}$, it holds:

$$(3.19) \quad |D_y^\beta p_{\Theta_{s,t},(\tau,\xi),\mu}(y+z)| \leq \frac{C}{(t-s)^{\frac{|\beta|}{\alpha}}} \bar{q}(t-s, y).$$

As a very important consequence of this lemma, we also derive from (3.13)-(3.14) that for all $(x, y) \in (\mathbb{R}^d)^2$:

$$(3.20) \quad |D_x^\beta \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y)| := |D_x^\beta p_{\Theta_{s,v},(\tau,\xi),\mathbf{P}_1}(y - m_{v,s}^{(\tau,\xi)}(x))| \leq \frac{C}{(v-s)^{\frac{|\beta|}{\alpha}}} \bar{q}(v-s, y - m_{v,s}^{(\tau,\xi)}(x)).$$

3.2. Proof of Proposition 2.1. Throughout this subsection, we set $\gamma = \eta$, that is, $h \in \mathcal{C}^\eta(\mathbb{R}^d, \mathbb{R})$ with $\|h\|_{\mathcal{C}^\eta} \leq 1$. With the notations and controls of the previous subsection at hand, we can now restart from (3.9) to write for all $(s, x) \in [0, t] \times \mathbb{R}^d$,

$$(3.21) \quad \begin{aligned} w(s, x) &= \int_s^t dv \left(\tilde{p}_{s,v}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \cdot) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, \cdot)) \right)(x) \\ &= \int_s^t dv \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) (H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)) dy. \end{aligned}$$

From (3.21) and the definition of $R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}$ in (3.11), we see that w appears in both sides of the equality. The point is now to perform a circular argument to control w . To this end, we also point out that the a priori controls of Lemma 3.1 will be useful to control $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$ introduced in (3.3). To complete the proof of Proposition 2.1 we aim at showing the following:

$$(3.22) \quad |w(s, x)| \leq C \mathbf{d}_{s,t,\eta}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{\tilde{\zeta}},$$

for some $\tilde{\zeta} > 0$, which precisely gives (2.15) for h replaced by h^δ .

Write first that, from (\mathbf{B}_H) , the definition of $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$ and Lemma 3.1 (see equations (3.5) and (3.6) with $\gamma = \eta$), it holds that there exists a constant C s.t. for all $(v, y) \in [s, t] \times \mathbb{R}^d$:

$$(3.23) \quad |H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y)| \leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \left((t-v)^{-1+\varepsilon} \mathbf{1}_{\{\alpha \geq 1\}} + (t-v)^{-\left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right)} \right),$$

using as well for the above inequality that, from the definitions in (2.1)-(2.2),

$$(3.24) \quad \forall \eta \in (0, \frac{1}{2}], \quad \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) \leq \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)).$$

This bound will be frequently used in the sequel.

On the other hand, from Lemma 3.3 and (3.20):

$$(3.25) \quad \begin{aligned} & \left| \int_s^t dv \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y) dy \right| \\ & \leq C \int_s^t dv \int_{\mathbb{R}^d} dy \left(\|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{\mathcal{C}}^{2\eta}} |D_y w(v, y)| + \|\sigma(v, \cdot, \mathbf{P}_1(v))\|_{\mathcal{C}^{2\eta}} [(\|w\|_\infty + \|D_x w(v, \cdot)\|_\infty) \mathbf{1}_{\{\alpha < 1\}} \right. \\ & \quad \left. + \|D_x w(v, \cdot)\|_{\mathcal{C}^\vartheta} \mathbf{1}_{\{\alpha > 1\}} + \|w(v, \cdot)\|_{\mathcal{C}^{1+\vartheta}} \mathbf{1}_{\{\alpha=1\}}] \right) |y - \theta_{v,\tau}(\xi)|^{2\eta} \bar{q}(v-s, y - m_{v,s}^{(\tau,\xi)}(x)), \end{aligned}$$

denoting here and for the rest of the section with a slight abuse of notation $\|w\|_\infty := \sup_{v \in [s,t]} \|w(v, \cdot)\|_\infty$. Hence, choosing $(\tau, \xi) = (s, x)$ and exploiting (3.16) and Lemma 3.4, we derive:

$$(3.26) \quad \begin{aligned} & \left| \int_s^t dv \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y) dy \right|_{(\tau, \xi) = (s, x)} \\ & \leq C \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left(\|w\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} + \|D_x w(v, \cdot)\|_\infty + \|D_x w(v, \cdot)\|_{\dot{\mathcal{C}}^\vartheta} \mathbf{1}_{\{\alpha \geq 1\}} \right). \end{aligned}$$

From (3.26), (3.23) and (3.21), one derives:

$$(3.27) \quad \begin{aligned} |w(s, x)| &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) [(t-s)^{\varepsilon \wedge (1 - (\frac{1}{\alpha} - \frac{\eta}{\alpha}))} \mathbf{1}_{\{\alpha \geq 1\}} + (t-s)^{(1 - (\frac{1}{\alpha} - \frac{\eta}{\alpha}))}] + (t-s) \|w\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} \right. \\ & \quad \left. + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left(\|D_x w(v, \cdot)\|_\infty + \|D_x w(v, \cdot)\|_{\dot{\mathcal{C}}^\vartheta} \mathbf{1}_{\{\alpha \geq 1\}} \right) \right). \end{aligned}$$

It is now clear that we actually have to estimate the gradient of w , i.e. we need to differentiate (3.21). We will now split the rest of the analysis in function of the index α . Indeed, for $\alpha > 1$, differentiating (3.21) yields an integrable

singularity in time and the previous controls do not change much. The counterpart is that we also need to give a bound on the Hölder modulus of continuity. On the other hand, for $\alpha \leq 1$, the induced time singularity is not integrable and it is crucial to equilibrate it through adequate cancellation techniques.

The case $\alpha \geq 1$. From (3.21), we can differentiate w.r.t x and also use a cancellation technique for the source term $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$. Namely,

$$(3.28) \quad D_x w(s, x) = \int_s^t dv \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) \left([H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))] + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y) \right) dy.$$

Similarly to (3.23), we derive from (\mathbf{B}_H) , equation (3.7) (with $\lambda = 1/2$) of Lemma 3.1 and (3.24) that:

$$(3.29) \quad \begin{aligned} |H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))| &\leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \left[(|y - \theta_{v, \tau}(\xi)|^\eta \wedge 1) \left((t-v)^{-1+\varepsilon} + (t-v)^{-\left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right)} \right) \right. \\ &\quad \left. + (t-v)^{-1+\varepsilon} (|y - \theta_{v, \tau}(\xi)|^\vartheta \wedge 1) + (|y - \theta_{v, \tau}(\xi)|^{1+\vartheta-\alpha} \wedge 1) \right] \\ &\leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \left[(|y - \theta_{v, \tau}(\xi)|^\eta \wedge 1) \left((t-v)^{-1+\varepsilon} + (t-v)^{-\left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right)} \right) \right. \\ &\quad \left. + (t-v)^{-1+\varepsilon} (|y - \theta_{v, \tau}(\xi)|^{1+\vartheta-\alpha} \wedge 1) \right], \end{aligned}$$

recalling that since $\alpha \geq 1$, $1 + \vartheta - \alpha \leq \vartheta$ for the last inequality. From the definition in (3.5) we thus get that, for ε meant to be small,

$$(3.30) \quad 1 + \vartheta - \alpha = 1 - \alpha + \alpha \left[1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\eta}{\alpha} \right) \right] = \eta - \varepsilon \alpha.$$

We thus obtain the following control

$$(3.31) \quad |H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))| \leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) |y - \theta_{v, \tau}(\xi)|^{\eta - \varepsilon \alpha} (t-v)^{-1+\varepsilon}.$$

Hence, from (3.28), (3.31), Lemma 3.3 and (3.18), we get:

$$(3.32) \quad \begin{aligned} &|D_x w(s, x)| \\ &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \int_s^t \frac{dv}{(v-s)^{\frac{1}{\alpha}}} (t-v)^{-1+\varepsilon} \int_{\mathbb{R}^d} dy |y - \theta_{v, \tau}(\xi)|^{\eta - \varepsilon \alpha} \bar{q}(v-s, y - m_{v, s}^{(\tau, \xi)}(x)) \right. \\ &\quad \left. + \int_s^t \frac{dv}{(v-s)^{\frac{1}{\alpha}}} \int_{\mathbb{R}^d} dy \left(\|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{C}^{2\eta}} |D_y w(v, y)| \right. \right. \\ &\quad \left. \left. + \|\sigma(v, \cdot, \mathbf{P}_1(v))\|_{C^{2\eta}} [\|D_x w(v, \cdot)\|_{C^\vartheta} \mathbf{1}_{\alpha > 1} + \|w(v, \cdot)\|_{C^{1+\vartheta}} \mathbf{1}_{\alpha=1}] \right) |y - \theta_{v, \tau}(\xi)|^{2\eta} \bar{q}(v-s, y - m_{v, s}^{(\tau, \xi)}(x)) \right). \end{aligned}$$

Taking $(\tau, \xi) = (s, x)$, it follows again from (3.16) and Lemma 3.4 that:

$$(3.33) \quad \begin{aligned} |D_x w(s, x)| &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{\eta}{\alpha} - \varepsilon} (t-v)^{-1+\varepsilon} \right. \\ &\quad \left. + \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left(\|D_x w(v, \cdot)\|_\infty + \|D_x w(v, \cdot)\|_{\dot{C}^\vartheta} + \|w(v, \cdot)\|_\infty \mathbf{1}_{\alpha=1} \right) \right) \\ &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{-\frac{1}{\alpha} + \frac{\eta}{\alpha}} \right. \\ &\quad \left. + \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left(\|D_x w(v, \cdot)\|_\infty + \|D_x w(v, \cdot)\|_{\dot{C}^\vartheta} + \|w(v, \cdot)\|_\infty \mathbf{1}_{\alpha=1} \right) \right). \end{aligned}$$

From (3.33), we also need to estimate the gradient through a circular argument.

Let us now make a short comment before proceeding further. We cannot expect for the gradient of w a better behaviour than the one provided by the *main* term of the perturbative expansion (3.21), i.e. the one associated with $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$. Observe now that we get: $-\frac{1}{\alpha} + \frac{\eta}{\alpha} =: \zeta < 0$. This exponent being negative, we cannot expect to have a pointwise bound for the gradient of w . Having in mind that we want to keep track in the above r.h.s. of the product of the distance $\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)$ and a contribution of type $(t-s)^\zeta$, $\zeta > 0$, which provides a smoothing in time effect, we will therefore investigate the behavior of

$$(3.34) \quad \Phi(v) := (t-v)^\Xi \sup_{x \in \mathbb{R}^d} |Dw(v, x)|, \quad \Xi := \frac{1}{\alpha} - \frac{\eta}{\alpha} + \frac{\varepsilon}{2}, \quad v \in [s, t],$$

To investigate the explosion of the corresponding Hölder modulus, let us introduce as well

$$(3.35) \quad \Psi(v) := (t-v)^{\Xi + \frac{\vartheta}{\alpha}} \|Dw(v, \cdot)\|_{C^\vartheta}, \quad v \in [s, t].$$

Let us note that $\sup_{v \in [s, t]} \Phi(v) < \infty$ and $\sup_{v \in [s, t]} \Psi(v) < \infty$. Indeed, from the definition of $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$ in (3.3), since from [CdRMP19] $u_2^\delta \in L^\infty([0, T], C^{\alpha+2\eta}(\mathbb{R}^d, \mathbb{R}))$ (with corresponding norm which a priori explodes with δ), it holds that, if b is bounded in space, $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta \in L^\infty([0, T], C^{2\eta \wedge (\alpha+2\eta-1)}(\mathbb{R}^d, \mathbb{R}))$. Hence, the previously invoked Schauder estimates still apply for w solving (3.3) (again with corresponding $L^\infty([0, T], C^{\alpha+2\eta \wedge (\alpha+2\eta-1)}(\mathbb{R}^d, \mathbb{R}))$ norm exploding with δ). If now b is not bounded in space, recall that *a priori*, we assumed in (\mathbf{B}_H) that $b(t, \cdot, \mu) \in \dot{C}^{2\eta}(\mathbb{R}^d, \mathbb{R}^d)$, we would then have $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta \in L^\infty([0, T], \dot{C}^{2\eta \wedge (\alpha+2\eta-1)}(\mathbb{R}^d, \mathbb{R}))$. In that case, from the integrability constraint $\alpha > 2\eta$, the arguments of [CdRMP19] could be reproduced to derive the finiteness of $\sup_{v \in [s, t]} \Phi(v), \sup_{v \in [s, t]} \Psi(v)$ which involve the gradient which is in that case bounded whereas the function itself is not (see e.g. Krylov and Priola [KP10] for similar issues in the diffusive setting).

From the previous definitions, a key point, for our circular argument to work is that $\Xi + \frac{\vartheta}{\alpha} < 1$. Indeed, from (3.33), this is necessary in order to make the time singularities generated by Φ and Ψ integrable. We now recall from the definition of ϑ in (3.5) that:

$$(3.36) \quad \Xi + \frac{\vartheta}{\alpha} = \frac{1}{\alpha} - \frac{\eta}{\alpha} + \frac{\varepsilon}{2} + 1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right) = 1 - \frac{\varepsilon}{2} < 1.$$

With these notations at hand, we can from (3.27) get rid of the supremum norm $\|w\|_\infty$ appearing in the case $\alpha = 1$ and upper bound it in terms of the functions Φ, Ψ . Namely, for ε and T small enough (so that $t - s < 1/2$),

$$(3.37) \quad \begin{aligned} \|w\|_\infty &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^\varepsilon + (t-s)\|w\|_\infty \mathbf{1}_{\{\alpha=1\}} \right. \\ &\quad \left. + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left((t-v)^{-\frac{\vartheta}{\alpha}} \Phi(v) + (t-v)^{-(\Xi+\frac{\vartheta}{\alpha})} \Psi(v) \right) \right) \\ &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^\varepsilon + \sup_{v \in [s, t]} \Phi(v) + \sup_{v \in [s, t]} \Psi(v) \right). \end{aligned}$$

up to a modification of the constant C for the last inequality.

We now have to distinguish the *diagonal* and *off-diagonal* regimes w.r.t. the current considered times. Namely, for given $(x, x') \in (\mathbb{R}^d)^2$, if $|x - x'| \geq (t-s)^{\frac{1}{\alpha}}$ we say that the *off-diagonal* regime holds. In this case, we readily get from (3.33) and (3.37):

$$(3.38) \quad \begin{aligned} (t-s)^{\frac{\vartheta}{\alpha}} \frac{|D_x w(s, x) - D_x w(s, x')|}{|x - x'|^\vartheta} &\leq |Dw(s, x) - Dw(s, x')| \leq |Dw(s, x)| + |Dw(s, x')| \\ &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^\zeta \right. \\ &\quad \left. + \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left((t-v)^{-\Xi} \Phi(v) + (t-v)^{-(\Xi+\frac{\vartheta}{\alpha})} \Psi(v) \right) \right), \\ (t-s)^{\Xi+\frac{\vartheta}{\alpha}} \frac{|D_x w(s, x) - D_x w(s, x')|}{|x - x'|^\vartheta} &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{\Xi+\zeta} + (t-s)^{1-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right. \\ &\quad \left. + (t-s)^{1-(\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right). \end{aligned}$$

Let us now turn to the global *diagonal regime*, $|x - x'| \leq (t-s)^{\frac{1}{\alpha}}$. This case is more subtle. Indeed, due to the time integration, i.e. for $v \in [s, t]$, there is a *local off-diagonal regime*, namely for $v \in [s, s + c_0|x - x'|^\alpha]$ for a constant $c_0 \in (0, 1)$ to be specified, and then a *true diagonal regime* w.r.t. the integration variable v which is the one providing the smoothing effects through the heat kernel. The delicate point is that, in order to exploit the off-diagonal bounds, the natural choice of freezing spatial point consists in considering precisely the spatial point in argument of the function. On the other hand, in the *true diagonal regime* the freezing spatial point must be the same for the two quantities to expand. This naturally induces to consider an intermediate quantity for the regime change.

Introduce for a couple $(\tau, \xi') \in \mathbb{R}_+ \times \mathbb{R}^d$ of freezing points, the following *Green kernel* :

$$\forall 0 \leq s < r \leq t, \ x \in \mathbb{R}^d, \ \tilde{G}_{s, r}^{\mathbf{P}_1, \alpha, (\tau, \xi')} f(s, x) := \int_s^r dv \int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, v, x, y) f(v, y).$$

For $x, x' \in \mathbb{R}^d$ s.t. $|x - x'|^\alpha \leq (t - s)$ we can write, similarly to the proof of Proposition 12 in [CdRMP19] (Duhamel formula with change of freezing point), that:

$$\begin{aligned}
& D_x w(s, x') \\
&= (D_x \tilde{G}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(s, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x')} + (D_x \tilde{G}_{\tau_0, t}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(s, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x)} \\
&\quad + (D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x')} - (D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x)} \\
&\quad + \int_s^t dv \int_{\mathbb{R}^d} dy \left(\mathbf{1}_{\{v \leq \tau_0\}} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, v, x', y) R_{\mathbf{P}_1}^{\alpha, (\tau, \xi')}(v, y) \right. \\
&\quad \left. + \mathbf{1}_{\{v > \tau_0\}} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, v, x', y) R_{\mathbf{P}_1}^{\alpha, (\tau, \xi')}(v, y) \right)_{(\tau_0, \tau, \xi', \tilde{\xi}')=(t_0, s, x', x)},
\end{aligned} \tag{3.39}$$

where

$$t_0 = s + c_0 |x - x'|^\alpha. \tag{3.40}$$

The point is now precisely to use the two expansions of the gradient in (3.21) (that we first differentiate w.r.t. x choosing then $(\tau, \xi) = (s, x)$) and (3.39), keeping in mind that the additional term

$$(D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x')} - (D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x)},$$

arising from the change of freezing points precisely needs to be analyzed.

We get:

$$\begin{aligned}
& D_x w(s, x) - D_x w(s, x') \\
&= \left(\int_s^t dv \mathbf{1}_{\{v \leq \tau_0\}} \int_{\mathbb{R}^d} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) (H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)) dy \right. \\
&\quad \left. - \int_s^t dv \mathbf{1}_{\{v \leq \tau_0\}} \int_{\mathbb{R}^d} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, v, x', y) (H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi')}(v, y)) dy \right) \Big|_{(\tau_0, \tau, \xi, \xi')=(t_0, s, x, x')} \\
&\quad - \left((D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x')} - (D_x \tilde{P}_{s, \tau_0}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(w(\tau_0, x'))_{(\tau_0, \tau, \xi')=(t_0, s, x)} \right) \\
&\quad + \int_s^t dv \mathbf{1}_{\{v > \tau_0\}} \int_{\mathbb{R}^d} (D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) - D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x', y)) \\
&\quad \quad \times (H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)) dy \Big|_{(\tau_0, \tau, \xi)=(t_0, s, x)} \\
&=: \Delta w_1(s, x, x') + \Delta w_2(s, x, x') + \Delta w_3(s, x, x').
\end{aligned} \tag{3.41}$$

The term $\Delta w_1(s, x, x')$, corresponding to the local off-diagonal regime within the global diagonal one, can be analyzed as above. Proceeding as in (3.31)-(3.33), and with the bound of (3.37) we get:

$$\begin{aligned}
& |\Delta w_1(s, x, x')| \\
&\leq \left| \int_s^{\tau_0} dv \int_{\mathbb{R}^d} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x, y) ([H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))] + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)) dy \right| \Big|_{(\tau_0, \tau, \xi)=(t_0, s, x)} \\
&\quad + \left| \int_s^{\tau_0} dv \int_{\mathbb{R}^d} D_x \tilde{P}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, v, x', y) ([H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi'))] + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi')}(v, y)) dy \right| \Big|_{(\tau_0, \tau, \xi')=(t_0, s, x')} \\
&\leq C \left(\int_s^{t_0} \frac{dv}{(v - s)^{\frac{1}{\alpha} - (\frac{\eta}{\alpha} - \varepsilon)}} \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) (t - v)^{-1 + \varepsilon} \right. \\
&\quad \left. + \int_s^{t_0} dv (v - s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} [(t - v)^{-\Xi} \Phi(v) + (t - v)^{-(\Xi + \frac{\eta}{\alpha})} \Psi(v)] \frac{|x - x'|^\vartheta}{(v - s)^{\frac{\vartheta}{\alpha}}} \right).
\end{aligned}$$

From (3.40), we then get that for $v \in [s, t_0]$, $(t - v) \geq (1 - c_0)(t - s)$ and therefore, recalling from (3.40) that $|x - x'|^\vartheta = (c_0^{-1}(t_0 - s))^{\frac{\vartheta}{\alpha}}$:

$$\begin{aligned} \frac{|\Delta w_1(s, x, x')|}{|x - x'|^\vartheta} &\leq C \left((t - s)^{-1+\varepsilon} \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t_0 - s)^{1+(\frac{\eta}{\alpha}-\varepsilon)-(\frac{1}{\alpha}+\frac{\vartheta}{\alpha})} \right. \\ &\quad + (t - s)^{-\Xi} \int_s^{t_0} dv (v - s)^{-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} \Phi(v) \frac{1}{(v - s)^{\frac{\vartheta}{\alpha}}} \\ &\quad \left. + (t - s)^{-(\Xi+\frac{\vartheta}{\alpha})} \int_s^{t_0} dv (v - s)^{-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} \Psi(v) \frac{1}{(v - s)^{\frac{\vartheta}{\alpha}}} \right) \\ &\leq C \left((t - s)^{-1+\varepsilon} \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \right. \\ &\quad \left. + (t - s)^{-\Xi+1-\frac{1}{\alpha}+\frac{2\eta}{\alpha}-\frac{\vartheta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t - s)^{-(\Xi+\frac{\vartheta}{\alpha})+1-\frac{1}{\alpha}+\frac{2\eta}{\alpha}-\frac{\vartheta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right), \end{aligned}$$

recalling for the last inequality that:

$$(3.42) \quad 1 + \left(\frac{\eta}{\alpha} - \varepsilon\right) - \left(\frac{1}{\alpha} + \frac{\vartheta}{\alpha}\right) = 1 + \left(\frac{\eta}{\alpha} - \varepsilon\right) - \left(\frac{1}{\alpha} + 1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right)\right) = 0.$$

This eventually gives:

$$\begin{aligned} (t - s)^{\Xi+\frac{\vartheta}{\alpha}} \frac{|\Delta w_1(s, x, x')|}{|x - x'|^\vartheta} &\leq C \left((t - s)^{\Xi+\frac{\vartheta}{\alpha}-1+\varepsilon} \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \right. \\ &\quad \left. + (t - s)^{1-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t - s)^{1-(\frac{1}{\alpha}+\frac{\vartheta}{\alpha})+\frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right). \end{aligned}$$

Recalling from (3.34) that $\Xi := \frac{1}{\alpha} - \frac{\eta}{\alpha} + \frac{\varepsilon}{2}$, from (3.42) we get:

$$(3.43) \quad \Xi + \frac{\vartheta}{\alpha} - 1 + \varepsilon = \left(\frac{1}{\alpha} + \frac{\vartheta}{\alpha}\right) - \frac{\eta}{\alpha} - 1 + \frac{3}{2}\varepsilon = 1 - \varepsilon + \frac{\eta}{\alpha} - \frac{\eta}{\alpha} - 1 + \frac{3}{2}\varepsilon = \frac{\varepsilon}{2} =: \bar{\zeta}.$$

Hence

$$(3.44) \quad \begin{aligned} (t - s)^{\Xi+\frac{\vartheta}{\alpha}} \frac{|\Delta w_1(s, x, x')|}{|x - x'|^\vartheta} &\leq C \left((t - s)^{\bar{\zeta}} \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \right. \\ &\quad \left. + (t - s)^{1-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t - s)^{1-(\frac{1}{\alpha}+\frac{\vartheta}{\alpha})+\frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right). \end{aligned}$$

Turning now to Δw_3 in (3.41), expanding the frozen densities, exploiting as well (3.19) which gives that a diagonal perturbation of the density does not affect the related estimates, we write:

$$\begin{aligned} |\Delta w_3(s, x, x')| &\leq \left| \int_{\tau_0}^t dv \int_0^1 d\lambda \int_{\mathbb{R}^d} D_x^2 \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi)}(s, v, x + \lambda(x' - x), y)(x' - x) \right. \\ &\quad \left. \times ([H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))] + R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)) dy \right|_{(\tau_0, \tau, \xi) = (t_0, s, x)} \\ &\leq C \left(|x - x'| \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \int_{t_0}^t \frac{dv}{(v - s)^{\frac{2}{\alpha} - (\frac{\eta}{\alpha} - \varepsilon)}} (t - v)^{-1+\varepsilon} \right. \\ &\quad \left. + |x - x'|^\vartheta \int_{t_0}^t dv (v - s)^{-\frac{2}{\alpha} + \frac{2\eta}{\alpha}} \left((t - v)^{-\Xi} \Phi(v) + (t - v)^{-(\Xi+\frac{\vartheta}{\alpha})} \Psi(v) + \right) |x - x'|^{1-\vartheta} \right), \end{aligned}$$

where we also used (3.7), for $\gamma = \eta, \lambda = \frac{1}{2}$, and (3.24) for the last inequality. Splitting the first time integral in the right-hand side of the above inequality into the two disjoint intervals $[t_0, \frac{t+s}{2}]$ and $(\frac{t+s}{2}, t]$, noting that in the current diagonal regime, for $c_0 \in (0, \frac{1}{2})$, $t_0 = s + c_0|x - x'|^\alpha \leq (t + s)/2$, we get (recall that we are *far* from the singularity):

$$\int_{t_0}^{\frac{t+s}{2}} dv (v - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon)} (t - v)^{-1+\varepsilon} \leq C(t - s)^{-1+\varepsilon} (t_0 - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon) + 1}$$

and

$$\int_{\frac{t+s}{2}}^t dv (v - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon)} (t - v)^{-1+\varepsilon} \leq C(t - s)^\varepsilon (t - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon)} = C(t - s)^{-\frac{2}{\alpha} + \frac{\eta}{\alpha}}.$$

Recalling that, for $v \in [t_0, t]$, $|x - x'| \leq ((v - s)/c_0)^{\frac{1}{\alpha}}$, we derive:

$$\begin{aligned} |\Delta w_3(s, x, x')| &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) |x - x'| [(t - s)^{-1+\varepsilon} (t_0 - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon) + 1} + (t - s)^{-\frac{2}{\alpha} + \frac{\eta}{\alpha}}] \right. \\ &\quad \left. + |x - x'|^\vartheta c_0^{\frac{\vartheta-1}{\alpha}} \int_{t_0}^t dv (v - s)^{-\frac{1}{\alpha} - \frac{\vartheta}{\alpha} + \frac{2\eta}{\alpha}} [(t - v)^{-\Xi} \Phi(v) + (t - v)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(v)] \right). \end{aligned}$$

Recall, from (3.42), $-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon) + 1 = \frac{\vartheta}{\alpha} - \frac{1}{\alpha}$. This yields $(t_0 - s)^{-\frac{2}{\alpha} + (\frac{\eta}{\alpha} - \varepsilon) + 1} = (c_0^{\frac{1}{\alpha}} |x - x'|)^{-1+\vartheta}$. Also in the current diagonal regime $|x - x'|^{1-\vartheta} (t - s)^{-\frac{2}{\alpha} + \frac{\eta}{\alpha}} \leq (t - s)^{\frac{1-\vartheta}{\alpha} - \frac{2}{\alpha} + \frac{\eta}{\alpha}} = (t - s)^{-\frac{\vartheta}{\alpha} - \frac{1}{\alpha} + \frac{\eta}{\alpha}}$. Eventually, since from (3.42) and (3.43) the other time singularities are integrable, we get:

$$\begin{aligned} (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|\Delta w_3(s, x, x')|}{|x - x'|^\vartheta} &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) [c_0^{\frac{\vartheta-1}{\alpha}} (t - s)^{\Xi + \frac{\vartheta}{\alpha} - 1 + \varepsilon} + (t - s)^{\Xi - \frac{1}{\alpha} + \frac{\eta}{\alpha}}] \right. \\ (3.45) \quad &\quad \left. + c_0^{\frac{\vartheta-1}{\alpha}} \left[(t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t - s)^{1 - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right] \right). \end{aligned}$$

The term $\Delta w_2(s, x, x')$ in (3.41), corresponding to the change of freezing point remains to be analyzed. A key point for this contribution is to note from (3.13)-(3.14) (affine structure in x' of $m_{\tau_0, s}^{(\tau, \xi')}(x')$) that:

$$\begin{aligned} D_x \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y) &= D_x \left(p_{\Theta_{s, \tau_0, (\tau, \xi'), \mathbf{P}_1}}(\tau_0 - s, y - m_{\tau_0, s}^{(\tau, \xi')}(x')) \right) \\ &= D_x \left(p_{\Theta_{s, \tau_0, (\tau, \xi'), \mathbf{P}_1}}(\tau_0 - s, y - x' + \int_s^{\tau_0} b(r, \theta_{r, \tau}(\xi'), \mathbf{P}_1(r)) dr) \right) \\ &= -D_y \left(p_{\Theta_{s, \tau_0, (\tau, \xi'), \mathbf{P}_1}}(\tau_0 - s, y - x' + \int_s^{\tau_0} b(r, \theta_{r, \tau}(\xi'), \mathbf{P}_1(r)) dr) \right) \\ (3.46) \quad &= -D_y \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y). \end{aligned}$$

Hence, with the notation of (3.14), we get:

$$\begin{aligned} &\Delta w_2(s, x, x') \\ &= \left(\int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y) D_y w(\tau_0, y) - \int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y) D_y w(\tau_0, y) \right)_{(\tau, \tau_0, \xi', \tilde{\xi}') = (s, t_0, x', x)} \\ &= \left(\int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y) \left[D_y w(\tau_0, y) - D_y w(\tau_0, \tilde{m}_{s, \tau_0}^{\tau, \xi'}(x')) \right] \right)_{(\tau, \tau_0, \xi') = (s, t_0, x')} \\ &\quad + \left(D_y w(\tau_0, \tilde{m}_{s, \tau_0}^{\tau, \xi'}(x')) - D_y w(\tau_0, \tilde{m}_{s, \tau_0}^{\tau, \tilde{\xi}'}(x')) \right)_{(\tau, \tau_0, \xi', \tilde{\xi}') = (s, t_0, x', x)} \\ &\quad - \left(\int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \tilde{\xi}')}(s, \tau_0, x', y) \left[D_y w(\tau_0, y) - D_y w(\tau_0, \tilde{m}_{s, \tau_0}^{\tau, \tilde{\xi}'}(x')) \right] \right)_{(\tau, \tau_0, \tilde{\xi}') = (s, t_0, x)}, \end{aligned}$$

recalling that $\int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \xi')}(s, \tau_0, x', y) = 1$ and $\int_{\mathbb{R}^d} dy \tilde{p}^{\mathbf{P}_1, \alpha, (\tau, \tilde{\xi}')}(s, \tau_0, x', y) = 1$ for the last equality.

From (3.13), Lemma 3.4 and the definitions in (3.35), we again derive:

$$\begin{aligned} |\Delta w_2(s, x, x')| &\leq C \sum_{z \in \{x, x'\}} \int_{\mathbb{R}^d} dy \bar{q}(t_0 - s, y - \tilde{m}_{s, t_0}^{s, z}(x')) |y - \tilde{m}_{s, t_0}^{s, z}(x')|^\vartheta \Psi(t_0) (t - t_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \\ &\quad + |\tilde{m}_{s, t_0}^{s, x'}(x') - \tilde{m}_{s, t_0}^{s, x}(x')|^\vartheta \Psi(t_0) (t - t_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \\ &\leq \Psi(t_0) (t - t_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \left(C |t_0 - s|^{\frac{\vartheta}{\alpha}} + \left| \int_s^{t_0} (b(v, \theta_{v, s}(x'), \mathbf{P}_1(v)) - b(v, \theta_{v, s}(x), \mathbf{P}_1(v))) dv \right|^\vartheta \right). \end{aligned}$$

Now, from Lemma 11 in [CdRMP19] we have $|\theta_{v, s}(x') - \theta_{v, s}(x)| \leq C(|x - x'| + (v - s)^{\frac{1}{\alpha}})$ and therefore:

$$\begin{aligned} \left| \int_s^{t_0} (b(v, \theta_{v, s}(x'), \mathbf{P}_1(v)) - b(v, \theta_{v, s}(x), \mathbf{P}_1(v))) dv \right| &\leq C \int_s^{t_0} |\theta_{v, s}(x') - \theta_{v, s}(x)|^{2\eta} dv \\ &\leq C(t_0 - s)(|x - x'| + (t_0 - s)^{\frac{1}{\alpha}})^{2\eta} \\ &\leq C c_0 |x - x'|^{\alpha + 2\eta} \leq C c_0 |x - x'|, \end{aligned}$$

recalling that, $\alpha + 2\eta > 1$ and since we are in the diagonal regime and that the time $t - s \leq 1$ we indeed have $|x - x'| \leq 1$. This finally gives:

$$|\Delta w_2(s, x, x')| \leq C\Psi(t_0)(t - t_0)^{-(\Xi + \frac{\vartheta}{\alpha})} c_0^{\frac{\vartheta}{\alpha}} |x - x'|^\vartheta.$$

Recalling that $t - t_0 = t - s - (t_0 - s) \geq t - s - c_0|x - x'|^\alpha \geq (1 - c_0)(t - s)$, this finally yields:

$$(3.47) \quad (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|\Delta w_2(s, x, x')|}{|x - x'|^\vartheta} \leq C c_0^{\frac{\vartheta}{\alpha}} (1 - c_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(t_0).$$

Plugging (3.44), (3.45) and (3.47) into (3.41), we get that in the diagonal case $|x - x'| \leq (t - s)^{\frac{1}{\alpha}}$ and with the definition of $\bar{\zeta} = \frac{\varepsilon}{2}$ in (3.43):

$$(3.48) \quad (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|D_x w(s, x) - D_x w(s, x')|}{|x - x'|^\vartheta} \leq C \left((1 + c_0^{\frac{\vartheta-1}{\alpha}}) \left[\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t - s)^{\bar{\zeta}} + (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right. \right. \\ \left. \left. + (t - s)^{1 + \frac{2\eta}{\alpha} - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha})} \sup_{v \in [s, t]} \Psi(v) + \right] + c_0^{\frac{\vartheta}{\alpha}} (1 - c_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(t_0) \right).$$

Putting together (3.48) and (3.38), corresponding respectively to the control of the normalized Hölder modulus of the gradient in the off-diagonal and diagonal regimes, we eventually derive:

$$(3.49) \quad \Psi(s) \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (1 + c_0^{\frac{\vartheta-1}{\alpha}}) (t - s)^{\bar{\zeta}} + (1 + c_0^{\frac{\vartheta-1}{\alpha}}) (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right. \\ \left. + \sup_{r \in [s, t]} \Psi(r) \left((1 + c_0^{\frac{\vartheta-1}{\alpha}}) (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} + c_0^{\frac{\vartheta}{\alpha}} (1 - c_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \right) \right).$$

Equation (3.49) could be established similarly for s replaced by any $v \in [s, t]$ in the above l.h.s. This in turn yields, taking $T \geq t - s$ and c_0 small enough s.t. $(1 + c_0^{\frac{\vartheta-1}{\alpha}}) T^{\bar{\zeta}}$ is also *small*, up to a modification of C :

$$(3.50) \quad \sup_{v \in [s, t]} \Psi(v) \leq C (1 + c_0^{\frac{\vartheta-1}{\alpha}}) \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{\bar{\zeta}} + (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right).$$

The point is now to plug (3.50) into (3.33) to complete the circular argument. Recalling as well that $1 - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha} = \frac{\eta}{\alpha} - \varepsilon$, this yields:

$$(3.33) \quad |D_x w(s, x)| \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{-\frac{1}{\alpha} + \frac{\eta}{\alpha}} + \sup_{v \in [s, t]} \Phi(v) (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha} - \Xi} + \sup_{v \in [s, t]} \Psi(v) (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha} - (\Xi + \frac{\vartheta}{\alpha})} \right) \\ (3.50) \quad \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \left[(t - s)^{-\frac{1}{\alpha} + \frac{\eta}{\alpha}} + (1 + c_0^{\frac{\vartheta-1}{\alpha}}) (t - s)^{\bar{\zeta} + \frac{\eta}{\alpha} - \varepsilon - \Xi} \right] \right. \\ \left. + \sup_{v \in [s, t]} \Phi(v) \left[(t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha} - \Xi} + (1 + c_0^{\frac{\vartheta-1}{\alpha}}) (t - s)^{(1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}) + \frac{\eta}{\alpha} - \varepsilon - \Xi} \right] \right).$$

Again the previous equation would also hold for any $v \in [s, t]$ instead of s in the previous l.h.s. Thus, from the definition on Ξ in (3.34) and $\bar{\zeta}$ in (3.43),

$$\sup_{v \in [s, t]} \Phi(v) := \sup_{v \in [s, t]} (v - s)^{\Xi} |D_x w(s, x)| \leq C (1 + c_0^{\frac{\vartheta-1}{\alpha}}) \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{\bar{\zeta}} \right. \\ \left. + \sup_{v \in [s, t]} \Phi(v) (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \right).$$

Taking again $T \geq t - s$ and c_0 small enough and s.t. $(1 + c_0^{\frac{\vartheta-1}{\alpha}}) T^{\bar{\zeta}}$ is sufficiently *small*, we derive:

$$(3.51) \quad \sup_{v \in [s, t]} \Phi(v) \leq C (1 + c_0^{\frac{\vartheta-1}{\alpha}}) \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{\bar{\zeta}}.$$

Plugging first (3.51) into (3.50), we get:

$$(3.52) \quad \sup_{v \in [s, t]} \Psi(v) \leq C (1 + c_0^{\frac{\vartheta-1}{\alpha}})^2 \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{\bar{\zeta}}.$$

It now remains to plug (3.51) and (3.52) into (3.27)-(3.37)

$$\begin{aligned}
 |w(s, x)| &\leq \|w\|_\infty \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^\varepsilon \right. \\
 &\quad \left. + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left((t-v)^{-\Xi} \Phi(v) + (t-v)^{-(\Xi+\frac{\eta}{\alpha})} \Psi(v) \right) \right) \\
 (3.53) \quad &\leq C(1 + c_0^{\frac{\vartheta-1}{\alpha}})^2 \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^\varepsilon.
 \end{aligned}$$

Observe indeed carefully that, even though the time contribution in front of the distance in (3.51) and (3.49) appears at a coarser rate, namely $(t-s)^{\tilde{\zeta}} = (t-s)^{\frac{\varepsilon}{2}}$, since we take out the supremum of Φ, Ψ and integrate once again in time in the above first inequality, we eventually get a final contribution in $(t-s)^\varepsilon$ for the distance (from the previous definitions of Ξ, ϑ in (3.34) and (3.5) respectively). This also illustrates the intuitive fact that the remainder terms in the perturbative expansion, those associated with Φ, Ψ , yield somehow negligible contributions.

Equation (3.53) precisely gives the expected bound (3.22) (with $\tilde{\zeta} = \varepsilon$) and completes the proof of Proposition 2.1 in the case $\alpha \geq 1$.

The case $\alpha < 1$. Restarting from (3.27), we write:

$$(3.54) \quad |w(s, x)| \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{\alpha})} + (t-s)\|w\|_\infty + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right)$$

so that

$$(3.55) \quad \|w\|_\infty \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{\alpha})} + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right),$$

assuming that $T \geq t-s$ small enough. Observe first that, for the first term of the above r.h.s., we have assumed that $(\frac{1}{\alpha} - \frac{\eta}{\alpha}) < 1 \iff \alpha + \eta > 1$. This indeed allows to derive the smoothing effect for the contribution $H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta$.

Let us now proceed from the cancellation argument as in (3.28). We get from equations (3.6) and (3.7) with $\lambda = 1/2$ (recalling that $\gamma = \eta$) and (3.24) that in the current case:

$$\begin{aligned}
 |H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))| &\leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \left[|y - \theta_{v, \tau}(\xi)|^\eta (t-v)^{-(\frac{1}{\alpha}-\frac{\eta}{\alpha})} \right. \\
 &\quad \left. + |y - \theta_{v, \tau}(\xi)|^\vartheta (t-v)^{-1+\varepsilon} \right] \\
 (3.56) \quad &\leq C \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) |y - \theta_{v, \tau}(\xi)|^{\alpha-1+\eta-\alpha\varepsilon} (t-v)^{-1+\varepsilon},
 \end{aligned}$$

as soon as $\varepsilon \leq 1 - (\frac{1}{\alpha} - \frac{\eta}{\alpha})$, recalling from (3.5) that $\vartheta = \alpha \left[1 - \varepsilon - (\frac{1}{\alpha} - \frac{\eta}{\alpha}) \right]$ and $\alpha < 1$ for the last inequality. Now, similarly to (3.33), using (3.56) and again Lemma 3.4, we also get:

$$\begin{aligned}
 |D_x w(s, x)| &\leq C \left(\int_s^t dv (v-s)^{-\frac{1}{\alpha}+1-\frac{1}{\alpha}+\frac{\eta}{\alpha}-\varepsilon} \mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) (t-v)^{-1+\varepsilon} \right. \\
 &\quad \left. + \int_s^t dv (v-s)^{-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} (\|w\|_\infty + \|Dw(v, \cdot)\|_\infty) \right) \\
 (3.57) \quad &\stackrel{(3.55)}{\leq} C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{2}{\alpha}-\frac{\eta}{\alpha})} + \int_s^t dv (v-s)^{-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right),
 \end{aligned}$$

provided that

$$(3.58) \quad 2 - \frac{2}{\alpha} + \frac{\eta}{\alpha} > 0 \iff \alpha > 1 - \frac{\eta}{2}$$

in order to obtain integrable time singularities in the first integral (taking as well ε small enough). Recalling that $\eta \in (0, 1/2]$, and that we have also assumed that $\alpha > 2\eta$ for integrability purposes, this means that $\alpha > 2\eta \vee (1 - \frac{\eta}{2})$.

With the notation (3.34) for $(\Phi(v))_{v \in [s, t]}$, taking in our current supercritical case $\Xi = \frac{1}{\alpha} - \frac{\eta}{2\alpha} < 1$ (from (3.58)), we derive from (3.55) and (3.57):

$$\begin{aligned}
 |w(s, x)| &\leq \|w\|_\infty \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{\alpha})} + \sup_{v \in [s, t]} \Phi(v) \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} (t-v)^{-\Xi} \right) \\
 (3.59) \quad &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{\alpha})} + \sup_{v \in [s, t]} \Phi(v) (t-s)^{1+2\frac{\eta}{\alpha}-\Xi} \right),
 \end{aligned}$$

and

$$\begin{aligned} |D_x w(s, x)| &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{2}{\alpha}-\frac{\eta}{\alpha})} + \sup_{v \in [s, t]} \Phi(v) \int_s^t dv (v-s)^{-\frac{1}{\alpha}+\frac{2\eta}{\alpha}} (t-v)^{-\Xi} \right) \\ &\leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{2}{\alpha}-\frac{\eta}{\alpha})} + \sup_{v \in [s, t]} \Phi(v)(t-s)^{1-(\frac{1}{\alpha}-\frac{2\eta}{\alpha})-\Xi} \right), \end{aligned}$$

which in turn gives:

$$(t-s)^{\Xi} |D_x w(s, x)| \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{2}{\alpha}-\frac{\eta}{\alpha})+\Xi} + \sup_{v \in [s, t]} \Phi(v)(t-s)^{1-(\frac{1}{\alpha}-\frac{2\eta}{\alpha})} \right).$$

Observe now that for the previous choice of Ξ and from (3.58), $1-(\frac{2}{\alpha}-\frac{\eta}{\alpha})+\Xi = 1-(\frac{1}{\alpha}-\frac{\eta}{2\alpha}) > 0$, so that we indeed have a regularizing effect. Furthermore, the above equation would still be valid with s replaced by any $\tilde{v} \in [s, t]$ for the l.h.s. This yields that for all $\tilde{v} \in [s, t]$

$$\Phi(\tilde{v}) \leq C \left(\mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{2\alpha})} + \sup_{v \in [s, t]} \Phi(v)(t-s)^{1-(\frac{1}{\alpha}-\frac{2\eta}{\alpha})} \right)$$

which yields

$$(3.60) \quad \sup_{v \in [s, t]} \Phi(v) \leq C \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{2\alpha})}$$

Plugging (3.60) into (3.59) eventually yields:

$$\|w\|_{\infty} \leq C \mathbf{d}_{\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)(t-s)^{1-(\frac{1}{\alpha}-\frac{\eta}{2\alpha})}.$$

This concludes the proof of Proposition 2.1 in the super-critical case under the condition $\alpha > 2\eta \vee (1 - \frac{\eta}{2})$. This naturally gives the constraint on the stability index $\alpha \in (4/5, 1)$. As already mentioned, in Appendix B, we remove this assumption and establish our main result, namely, Theorem 2.3 under the sole condition $\alpha > 2\eta \vee (1 - \eta)$.

APPENDIX A. ADDITIONAL CONTROLS ON THE DENSITY IN THE STABLE CASE

To derive the first main results of Lemma 3.4, the key idea consists in separating the small and the large jumps in the Fourier transform of the proxy $\Theta_{s, v, (\tau, \xi), \mu}$ introduced in (3.15). This *trick*, which is often known as the Itô-Lévy decomposition, has been used in many contexts (see e.g. [Wat07], [Szt10] or [HM16] for density estimates; [HMP19] for L^p bounds related to degenerate stably driven Ornstein-Uhlenbeck processes).

Fix now a final parameter t . Rewrite from (3.15), for any $\zeta \in \mathbb{R}^d$:

$$\begin{aligned} \varphi_{\Theta_{s, t, (\tau, \xi), \mu}}(\zeta) &:= \mathbb{E} \left[\exp(i \langle \zeta, \Theta_{s, t, (\tau, \xi), \mu} \rangle) \right] = \exp \left(\int_s^t dv \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \zeta, \sigma(v, \theta_{v, \tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right) \\ &= \exp \left(\int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \zeta, \sigma(v, \theta_{v, \tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right) \\ &\quad \times \exp \left(\int_s^t dv \int_{(t-s)^{\frac{1}{\alpha}}}^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \zeta, \sigma(v, \theta_{v, \tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right) \\ (A.1) \quad &=: \hat{p}_{M_{s, t, (\tau, \xi), \mu}}(\zeta) \hat{P}_{N_{s, t, (\tau, \xi), \mu}}(\zeta). \end{aligned}$$

It precisely allows to write: $\Theta_{s, t, (\tau, \xi), \mu} := M_{s, t, (\tau, \xi), \mu} + N_{s, t, (\tau, \xi), \mu}$ where $M_{s, t, (\tau, \xi), \mu}$ and $N_{s, t, (\tau, \xi), \mu}$ are independent random variables corresponding respectively to the small and large jumps part of the stochastic integral $\Theta_{s, t, (\tau, \xi), \mu}$. The density of $\Theta_{s, t, (\tau, \xi), \mu}$ writes for all $y \in \mathbb{R}^d$ as

$$(A.2) \quad p_{\Theta_{s, t, (\tau, \xi), \mu}}(y) = \int_{\mathbb{R}^d} p_{M_{s, t, (\tau, \xi), \mu}}(y-z) P_{N_{s, t, (\tau, \xi), \mu}}(dz).$$

It is known, see e.g. Lemma B.1 in [HM16], Lemma A.2 in [HMP19], that the density $p_{M_{s, t, (\tau, \xi), \mu}}$ associated with the small jumps of the stochastic integral is smooth and has a polynomial decay of arbitrary order (i.e. it belongs to the Schwartz class). On the other hand, $P_{N_{s, t, (\tau, \xi), \mu}}$ is a Poisson measure s.t. for all $\beta \in (0, \alpha)$, there exists $C_{\beta, \alpha} \geq 1$ s.t. $\mathbb{E}[|N_{s, t, (\tau, \xi), \mu}|^\beta] = \int_{\mathbb{R}^d} |z|^\beta P_{N_{s, t, (\tau, \xi), \mu}}(dz) \leq C_{\alpha, \beta} (t-s)^{\frac{\beta}{\alpha}}$. Importantly, both components $p_{M_{s, t, (\tau, \xi), \mu}}$ and $P_{N_{s, t, (\tau, \xi), \mu}}$ (when the jump measure ν is absolutely continuous) can be dominated uniformly w.r.t. the freezing parameters $(\tau, \xi), \mu$. For the sake of clarity, and since these steps are also crucial to study further the sensitivities of the quantities $p_{M_{s, t, (\tau, \xi), \mu}}, P_{N_{s, t, (\tau, \xi), \mu}}$ w.r.t. the measure argument, we prove these facts in the next subsection.

A.1. First controls for the two parts of the frozen density. We first give here some useful lemmas concerning the behavior of the laws of the independent random variables $M_{s,t,(\tau,\xi),\mu}$ and $N_{s,t,(\tau,\xi),\mu}$ s.t. $\Theta_{s,t,(\tau,\xi),\mu} = M_{s,t,(\tau,\xi),\mu} + N_{s,t,(\tau,\xi),\mu}$.

Lemma A.1 (Density estimate on the Martingale part and associated derivatives.). For all $m \geq 1$, there exists $C_m \geq 1$ s.t. for all $0 \leq s < t \leq T, y \in \mathbb{R}^d, (\tau, \xi) \in [0, T] \times \mathbb{R}^d$,

$$(A.3) \quad p_{M_{s,t,(\tau,\xi),\mu}}(y) \leq \frac{C_m}{(t-s)^{\frac{d}{\alpha}}} \left(1 + \frac{|y|}{(t-s)^{\frac{1}{\alpha}}}\right)^{-m} =: \tilde{C}_m p_{\bar{M}_{s,t},m}(y),$$

where $\int_{\mathbb{R}^d} p_{\bar{M}_{s,t},m}(y) dy = 1$.

Also, for all $m \geq 1$ and all multi-index β , $|\beta| \leq 3$,

$$|D_y^\beta p_{M_{s,t,(\tau,\xi),\mu}}(y)| \leq \frac{C_m}{(t-s)^{\frac{d+|\beta|}{\alpha}}} \left(1 + \frac{|y|}{(t-s)^{\frac{1}{\alpha}}}\right)^{-m} = \frac{\tilde{C}_m}{(t-s)^{\frac{|\beta|}{\alpha}}} p_{\bar{M}_{s,t},m}(y).$$

Proof. Inverting the Fourier transform in (A.1) write:

$$p_{M_{s,t,(\tau,\xi),\mu}}(y) = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} d\zeta e^{-i\langle \zeta, y \rangle} \exp \left(\int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \zeta, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right).$$

Setting $(t-s)^{1/\alpha} \zeta = \tilde{\zeta}$ yields:

$$(A.4) \quad p_{M_{s,t,(\tau,\xi),\mu}}(y) = \frac{1}{(2\pi)^d} (t-s)^{-\frac{d}{\alpha}} \int_{\mathbb{R}^d} d\tilde{\zeta} e^{-i\langle \tilde{\zeta}, \frac{y}{(t-s)^{\frac{1}{\alpha}}} \rangle} \times \exp \left(\int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}}, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right).$$

Let us now denote

$$\hat{f}_{s,t}(\tilde{\zeta}) := \exp \left(\int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}}, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right).$$

Since the Lévy measure in the above expression has finite support, we get from Theorem 3.7.13 in Jacob [Jac05] that $\hat{f}_{s,t}$ is infinitely differentiable as a function of $\tilde{\zeta}$. Moreover,

$$|D\hat{f}_{s,t}(\tilde{\zeta})| \leq \int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \frac{|\sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda| r}{(t-s)^{\frac{1}{\alpha}}} \left| \sin(\langle \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}}, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) \right| \times \exp \left(\int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}}, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right).$$

Write now:

$$\begin{aligned} & \int_s^t dv \int_0^{(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \frac{|\sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda| r}{(t-s)^{\frac{1}{\alpha}}} \left| \sin(\langle \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}}, \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda r \rangle) \right| \\ & \leq C(t-s) \int_{r \leq (t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \frac{r}{(t-s)^{\frac{1}{\alpha}}} (\mathbf{1}_{\alpha < 1} + \mathbf{1}_{\alpha \geq 1} |\tilde{\zeta}| \frac{r}{(t-s)^{\frac{1}{\alpha}}}) \leq C(1 + |\tilde{\zeta}|). \end{aligned}$$

Thus:

$$\begin{aligned} |D\hat{f}_{s,t}(\tilde{\zeta})| & \leq C(1 + |\tilde{\zeta}|) \exp \left(\int_s^t dv \int_{\mathbb{R}^d} \left\{ \cos(\langle \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) q, \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}} \rangle) - 1 \right\} \nu(dq) \right) \\ & \quad \times \exp(2(t-s) \nu(B(0, (t-s)^{\frac{1}{\alpha}})^c)) \leq C(1 + |\tilde{\zeta}|) \exp(-C^{-1} |\tilde{\zeta}|^\alpha), \quad C \geq 1, \end{aligned}$$

using for the last inequality that from (2.7)

$$\begin{aligned} & \exp \left(\int_s^t dv \int_{\mathbb{R}^d} \left\{ \cos(\langle \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) q, \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}} \rangle) - 1 \right\} \nu(dq) \right) \\ & = \exp \left(-C_{\alpha,d} \int_s^t dv \int_{\mathbb{S}^{d-1}} |\langle \sigma(v, \theta_{v,\tau}(\xi), \mu(v)) \lambda, \frac{\tilde{\zeta}}{(t-s)^{\frac{1}{\alpha}}} \rangle|^\alpha \omega(d\lambda) \right) \\ & \leq C \exp(-C^{-1} |\tilde{\zeta}|^\alpha), \end{aligned}$$

and also that, from the decomposition of ν , $\nu(B(0, (t-s)^{1/\alpha})^c) \leq C/(t-s)$.

Similarly, for all multi-index β of length $|\beta| = l \in \mathbb{N}$:

$$|D^\beta \hat{f}_{s,t}(\tilde{\zeta})| \leq C_l(1 + |\tilde{\zeta}|^l) \exp(-C^{-1}|\tilde{\zeta}|^\alpha), \quad C_l \geq 1.$$

Thus, $\hat{f}_{s,t}$ belongs the Schwartz space. Denoting by $f_{s,t}$ its inverse Fourier transform, we have:

$$\forall m \geq 0, \forall y \in \mathbb{R}^d, \exists C_m \geq 1 \text{ s.t. : } |f_{s,t}(y)| \leq C_m(1 + |y|)^{-m}.$$

Now, since $p_M(t-s, y) = (t-s)^{-\frac{d}{\alpha}} f_{s,t}(y/(t-s)^{\frac{1}{\alpha}})$, the announced bound follows. The control concerning the derivatives is derived similarly. \square

Lemma A.2 (Controls on the Poisson measure). Let ν be any symmetric stable jump measure satisfying (2.7). Then, for all $\beta \in [0, \alpha)$, there exists $C_{\alpha, \beta} \geq 1$ s.t. for all $0 \leq s < t, z \in \mathbb{R}^d$,

$$(A.5) \quad \int_{\mathbb{R}^d} |y|^\beta P_{N_{s,t}, (\tau, \xi), \mu}(dy) \leq C_{\alpha, \beta} (t-s)^{\frac{\beta}{\alpha}}.$$

Assume now the jump measure ν is absolutely continuous, i.e. Assumption **(AC)** holds. Then, there exists a Poisson measure $P_{\tilde{N}_{\rho_{s,t}}} := \exp(-1) \sum_{n \geq 0} \frac{\rho_{s,t}^{\star n}}{n!}$ where $\rho_{s,t}(dy) = \kappa(t-s) \frac{dr}{r^{1+\alpha}} \mathbf{1}_{|r| \geq c(t-s)^{\frac{1}{\alpha}}} \Lambda_{\mathbb{S}^{d-1}}(d\theta)$, $y = r\theta$, $(r, \theta) \in \mathbb{R}_*^+ \times \mathbb{S}^{d-1}$, with $c > 0$ depending on the non-degeneracy constants in (2.7), (2.9), is a probability measure and $\rho_{s,t}^{\star n}$ denotes its n^{th} fold convolution, for which it holds that:

$$(A.6) \quad P_{N_{s,t}, (\tau, \xi), \mu} \leq C P_{\tilde{N}_{\rho_{s,t}}}.$$

Proof. Introduce $\bar{\nu}_{s,t}(dz) = \nu(dz) \mathbf{1}_{|z| \geq (t-s)^{\frac{1}{\alpha}}}$ which is a finite measure and s.t. $\bar{\nu}_{s,t}(\mathbb{R}^d) \leq C(t-s)^{-1}$. With this notation at hand, write:

$$\begin{aligned} \hat{P}_{N_{s,t}, (\tau, \xi), \mu}(\zeta) &= \exp \left(\int_s^t dv \int_{(t-s)^{\frac{1}{\alpha}}}^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) \left(\cos(\langle \zeta, \sigma(v, \theta_{v, \tau}(\xi), \mu(v)) \lambda r \rangle) - 1 \right) \right) = \\ &= \exp \left(\int_s^t dv \hat{\nu}_{s,t}(\sigma(v, \theta_{v, \tau}(\xi), \mu(v))^* \zeta) - (t-s) \bar{\nu}_{s,t}(\mathbb{R}^d) \right) \\ &= \exp \left(\int_s^t dv \hat{\nu}_{s,t, (v, (\tau, \xi), \mu)}(\zeta) - (t-s) \bar{\nu}_{s,t}(\mathbb{R}^d) \right), \end{aligned}$$

defining $\bar{\nu}_{s,t, (v, (\tau, \xi), \mu)}(A) := \bar{\nu}_{s,t}(\{y \in \mathbb{R}^d : \sigma(v, \theta_{v, \tau}(\xi), \mu(v))y \in A\})$ and where $\hat{\nu}_{s,t}$, $\hat{\nu}_{s,t, (v, (\tau, \xi), \mu)}$ denote the Fourier-Stieltjes transform of the considered measure. Hence, introducing the measure $\zeta_{s,t, (\tau, \xi), \mu} := \int_s^t dv \bar{\nu}_{s,t, (v, (\tau, \xi), \mu)}$ and expanding the previous exponential and by termwise Fourier inversion, we get:

$$(A.7) \quad P_{N_{s,t}, (\tau, \xi), \mu} = \exp(\zeta_{s,t, (\tau, \xi), \mu} - (t-s) \bar{\nu}_{s,t}(\mathbb{R}^d)) = \exp(-(t-s) \bar{\nu}_{s,t}(\mathbb{R}^d)) \sum_{n \geq 0} \frac{(\zeta_{s,t, (\tau, \xi), \mu})^{\star n}}{n!},$$

where in the above equation we again used that, for a finite measure ρ on \mathbb{R}^d , $(\rho)^{\star n} := \rho \star \dots \star \rho$ denotes its n^{th} fold convolution. Observe now from the definition of $\bar{\nu}_{s,t}$ and the non-degeneracy conditions on σ (see eq. (2.9) in **(UE)**) that there exists $C \geq 1$ s.t. for all $0 \leq s < t, z \in \mathbb{R}^d$, $\mu \in \mathcal{C}(\mathbb{R}^+, \mathcal{P}(\mathbb{R}^d))$, $C^{-1} \leq \zeta_{s,t, (\tau, \xi), \mu}(\mathbb{R}^d) = (t-s) \bar{\nu}_{s,t}(\mathbb{R}^d) \leq C$. Introducing the normalized (probability) measure $\tilde{\zeta}_{s,t, (\tau, \xi), \mu} := \frac{\zeta_{s,t, (\tau, \xi), \mu}}{\zeta_{s,t, (\tau, \xi), \mu}(\mathbb{R}^d)} = \frac{\zeta_{s,t, (\tau, \xi), \mu}}{(t-s) \bar{\nu}_{s,t}(\mathbb{R}^d)}$, one then rewrites:

$$\int_{\mathbb{R}^d} |y|^\beta P_{N_{s,t}, (\tau, \xi), \mu}(dy) \leq \exp(-(t-s) \bar{\nu}_{s,t}(\mathbb{R}^d)) \sum_{n=0}^{\infty} \frac{((t-s) \bar{\nu}_{s,t}(\mathbb{R}^d))^n}{n!} \int_{\mathbb{R}^d} |y|^\beta (\tilde{\zeta}_{s,t, (\tau, \xi), \mu})^{\star n}(dy).$$

Introducing now on a some probability space $(\Omega, \mathcal{F}, \mathbb{P})$ a sequence $(X_i)_{i \geq 0}$ of i.i.d. random variables with law $\tilde{\zeta}_{s,t, (\tau, \xi), \mu}$, we can rewrite from the above control:

$$(A.8) \quad \int_{\mathbb{R}^d} |y|^\beta P_{N_{s,t}, (\tau, \xi), \mu}(dy) \leq e^{-C^{-1}} \sum_{n=0}^{\infty} \frac{C^n}{n!} \mathbb{E}[\sum_{i=0}^n |X_i|^\beta] \leq e^{-C^{-1}} \sum_{n=0}^{\infty} \frac{C^n}{n!} (n+1)^{1 \vee \beta} \mathbb{E}[|X_1|^\beta].$$

Now,

$$\begin{aligned}
\mathbb{E}[|X_1|^\beta] &= \int_{\mathbb{R}^d} |y|^\beta \tilde{\zeta}_{s,t,(\tau,\xi),\boldsymbol{\mu}}(dy) \leq C \int_s^t dv \int_{\mathbb{R}^d} |y|^\beta \bar{\nu}_{s,t,(v,(\tau,\xi),\boldsymbol{\mu})}(dy) \\
&= C \int_s^t dv \int_{|y| \geq (t-s)^{\frac{1}{\alpha}}} |\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v)) y|^\beta \nu(dy) \\
&\leq C \int_s^t dv \int_{r \geq (t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \omega(d\lambda) |\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v)) \lambda r|^\beta \leq C \int_s^t dv \int_{r \geq (t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha-\beta}} \\
(A.9) \quad &\leq C_{\alpha,\beta} (t-s)^{\frac{\beta}{\alpha}},
\end{aligned}$$

recalling that $\alpha > \beta$ for the last inequality. Observe as well that the last constant $C_{\alpha,\beta}$ depends on α and β and explodes when $\beta \uparrow \alpha$. Plugging the above bound into (A.8) we derive the first statement (A.5) without any specific assumption on the jump measure.

Assume now that ν is absolutely continuous. Then, still with the previous notations, we can bound the measure $\int_s^t dv \bar{\nu}_{s,t,(v,(\tau,\xi),\boldsymbol{\mu})}$ uniformly in $(v, (\tau, \xi), \boldsymbol{\mu})$. Namely, for all $A \in \mathcal{B}(\mathbb{R}^d)$ and with the notations of assumption **(AC)**:

$$\begin{aligned}
\left(\int_s^t dv \bar{\nu}_{s,t,(v,(\tau,\xi),\boldsymbol{\mu})} \right)(A) &= \int_s^t dv \left(\bar{\nu}_{s,t,(v,(\tau,\xi),\boldsymbol{\mu})}(A) \right) = \int_s^t dv \left(\bar{\nu}_{s,t}(\{y \in \mathbb{R}^d : \sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v)) y \in A\}) \right) \\
&= \int_s^t dv \int_{|y| \geq (t-s)^{\frac{1}{\alpha}}} \mathbf{1}_{\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v)) y \in A} \nu(dy) \\
&= \int_s^t dv \int_{|\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v))^{-1} \tilde{y}| \geq (t-s)^{\frac{1}{\alpha}}} \mathbf{1}_{\tilde{y} \in A} \frac{f(\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v))^{-1} \tilde{y})}{|\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v))^{-1} \tilde{y}|^{d+\alpha}} \frac{d\tilde{y}}{\det(\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v)))} \\
&\leq C \int_s^t dv \int_{|\tilde{y}| \geq c(t-s)^{\frac{1}{\alpha}}} \mathbf{1}_{\tilde{y} \in Ag} \left(\frac{\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v))^{-1} \tilde{y}}{|\sigma(v, \theta_{v,\tau}(\xi), \boldsymbol{\mu}(v))^{-1} \tilde{y}|} \right) \frac{d\tilde{y}}{|\tilde{y}|^{d+\alpha}}.
\end{aligned}$$

Since g is bounded on \mathbb{S}^{d-1} (recall g is Lipschitz continuous on \mathbb{S}^{d-1}), writing $\tilde{y} = r\theta$, $(r, \theta) \in \mathbb{R}_*^+ \times \mathbb{S}^{d-1}$, we derive that there exists a constant C and a probability measure $\rho_{s,t}(d\theta) = \kappa(t-s) \frac{dr}{r^{1+\alpha}} \mathbf{1}_{|r| \geq c(t-s)^{\frac{1}{\alpha}}} \Lambda_{\mathbb{S}^{d-1}}(d\theta)$, s.t. for all $(\tau, \xi, \boldsymbol{\mu}) \in [0, T] \times \mathbb{R}^d \times \mathcal{C}(\mathbb{R}^+, \mathcal{P}(\mathbb{R}^d))$, $A \in \mathcal{B}(\mathbb{R}^d)$:

$$\left(\int_s^t dv \bar{\nu}_{s,t,(v,(\tau,\xi),\boldsymbol{\mu})} \right)(A) \leq C(t-s) \int_{|r| \geq c(t-s)^{\frac{1}{\alpha}}} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) \mathbf{1}_{r\theta \in A} = C\rho_{s,t}(A),$$

up to a modification of C . The above equation and (A.7) eventually give (A.6). \square

Remark A.1 (Integrability of the stochastic integral). As a direct corollary of Lemma A.1 and A.2 we derive that for any non-degenerate jump measure ν satisfying (2.7) and all $\beta \in [0, \alpha)$ there exists $C_\beta \geq 1$ (increasing with β) s.t. for all $0 \leq s < t, z \in \mathbb{R}^d, \boldsymbol{\mu} \in \mathcal{C}(\mathbb{R}^+, \mathcal{P}(\mathbb{R}^d))$:

$$(A.10) \quad \int_{\mathbb{R}^d} |y|^\beta p_{\Theta_{s,t,(\tau,\xi),\boldsymbol{\mu}}}(y) dy \leq C_\beta (t-s)^{\frac{\beta}{\alpha}}.$$

More generally, under **(AC)**, Lemma 3.4 readily follows from Lemmas A.1 and A.2 setting for all $z \in \mathbb{R}^d$, $\bar{q}(t-s, z) := \int_{\mathbb{R}^d} p_{\bar{M}_{s,t,m}}(z - \bar{\zeta}) P_{\bar{N}_{\rho_{s,t}}}(d\bar{\zeta})$, for any arbitrary integer $m \geq d+1$.

A.2. Proof of Lemma 3.1. Quantitative bounds for the solution u_1^δ of the PDE (3.2). To derive the announced estimates, we rewrite the solution $u_i^\delta \in \mathcal{C}^{\alpha+\eta}(\mathbb{R}^d, \mathbb{R})$ of (3.2) as follows:

$$(A.11) \quad \begin{cases} \left(\partial_v + \tilde{L}_v^{\mathbf{P}_i, \alpha, (\tau, \xi)} \right) u_i^\delta(v, x) + b(v, \theta_{v,\tau}(\xi), \mathbf{P}_i(v)) \cdot D_x u_i^\delta(v, x) \\ \quad = -R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, x), \quad (v, x) \in [0, t] \times \mathbb{R}^d, \\ u_i^\delta(t, x) = h^\delta(x), \quad \text{on } \mathbb{R}^d, \end{cases}$$

with the notations of (3.10) and (3.11) (replacing w by u_i^δ in $R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}$), i.e.

$$\tilde{L}_v^{\mathbf{P}_i, \alpha, (\tau, \xi)} \varphi(x) := \text{p.v.} \int_{\mathbb{R}^d} (\varphi(x + \sigma(v, \theta_{v,\tau}(\xi), \mathbf{P}_i(v))) \zeta - \varphi(x)) \nu(d\zeta)$$

and

$$R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, x) := (b(v, x, \mathbf{P}_i(v)) - b(v, \theta_{v,\tau}(\xi), \mathbf{P}_i(v))) \cdot Du_i^\delta(v, x) + \left(L_s^{\mathbf{P}_i, \alpha} - \tilde{L}_s^{\mathbf{P}_i, \alpha, (\tau, \xi)} \right) u_i^\delta(v, x).$$

The PDE (A.11) suggests to rewrite u_i^δ through a Duhamel type expansion (which has been fully justified in Proposition 8 of [CdRMP19]). This yields for all $(s, x) \in [0, t] \times \mathbb{R}^d$,

$$\begin{aligned} u_i^\delta(s, x) &= \tilde{P}_{s,t}^{\mathbf{P}_i, \alpha, (\tau, \xi)} h^\delta(x) + \int_s^t dv \left(\tilde{P}_{s,v}^{\mathbf{P}_i, \alpha, (\tau, \xi)} R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, \cdot) \right)(x) \\ (A.12) \quad &= \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, t, x, y) h^\delta(y) dy + \int_s^t dv \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy. \end{aligned}$$

Now, from (3.11) and (3.20), similarly to (3.25)-(3.26), choosing $(\tau, \xi) = (s, x)$ and exploiting (3.16) and Lemma 3.4, we derive:

$$\begin{aligned} (A.13) \quad & \left| \int_s^t dv \int_{\mathbb{R}^d} \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy \right|_{(\tau, \xi) = (s, x)} \\ & \leq C \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left(\|u_i^\delta\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} + \|Du_i^\delta(v, \cdot)\|_\infty + \|Du_i^\delta(v, \cdot)\|_{\dot{C}^\theta} \mathbf{1}_{\{\alpha \geq 1\}} \right). \end{aligned}$$

Thus, recalling from the definition of $h^\delta = h \star \rho_\delta$ (for a smooth mollifying kernel ρ_δ) that $\|h^\delta\|_\infty \leq \|h\|_\infty$, we get:

$$(A.14) \quad |u_i^\delta(s, x)| \leq C \left(\|h\|_\infty + (t-s) \|u_i^\delta\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left(\|Du_i^\delta(v, \cdot)\|_\infty + \|Du_i^\delta(v, \cdot)\|_{\dot{C}^\theta} \mathbf{1}_{\{\alpha \geq 1\}} \right) \right).$$

From (A.12), we can differentiate w.r.t x and also use a cancellation technique for the final condition. Namely,

$$(A.15) \quad D_x u_i^\delta(s, x) = \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, t, x, y) [h^\delta(y) - h^\delta(\theta_{t,\tau}(\xi))] dy + \int_s^t dv \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy.$$

Hence, from (A.15), recalling that h^δ is γ -Hölder continuous in space uniformly in the regularization parameter δ , with in particular $\|h^\delta\|_{\dot{C}^\gamma} \leq \|h\|_{\dot{C}^\gamma}$, similarly to (A.14), we get:

$$\begin{aligned} (A.16) \quad & |D_x u_i^\delta(s, x)| \\ & \leq C \left(\frac{\|h\|_{\dot{C}^\gamma}}{(t-s)^{\frac{1}{\alpha}}} \int_{\mathbb{R}^d} dy |y - \theta_{v,\tau}(\xi)|^\gamma \bar{q}(v-s, y - m_{v,s}^{(\tau, \xi)}(x)) \right. \\ & + \int_s^t \frac{dv}{(v-s)^{\frac{1}{\alpha}}} \int_{\mathbb{R}^d} dy \left(\|b(v, \cdot, \mathbf{P}_i(v))\|_{C^{2\eta}} |D_y u_i^\delta(v, y)| + \|\sigma(v, \cdot, \mathbf{P}_i(v))\|_{C^{2\eta}} \left(\|u_i^\delta\|_\infty + \|D_x u_i^\delta(v, \cdot)\|_\infty \mathbf{1}_{\{\alpha < 1\}} \right. \right. \\ & \left. \left. + \|D_x u_i^\delta(v, \cdot)\|_{C^\theta} \mathbf{1}_{\{\alpha > 1\}} + \|u_i^\delta(v, \cdot)\|_{C^{1+\theta}} \mathbf{1}_{\{\alpha=1\}} \right) |y - \theta_{v,\tau}(\xi)|^{2\eta} \bar{q}(v-s, y - m_{v,s}^{(\tau, \xi)}(x)) \right). \end{aligned}$$

Taking $(\tau, \xi) = (s, x)$, it follows again from (3.16) and Lemma 3.4 that:

$$\begin{aligned} (A.17) \quad & |D_x u_i^\delta(s, x)| \\ & \leq C \left(\|h\|_{\dot{C}^\gamma} (t-s)^{-\frac{1}{\alpha} + \frac{\gamma}{\alpha}} + \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left(\|u_i^\delta(v, \cdot)\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} + \|D_x u_i^\delta(v, \cdot)\|_\infty + \|D_x u_i^\delta(v, \cdot)\|_{\dot{C}^\theta} \mathbf{1}_{\{\alpha \geq 1\}} \right) \right). \end{aligned}$$

We now use the notations of (3.35) and (3.34) with the exponents modified, namely we consider here

$$(A.18) \quad \Phi(v) := (t-v)^{\Xi} \sup_{x \in \mathbb{R}^d} |D_x u_i^\delta(v, x)|, \quad \Xi := \frac{1}{\alpha} - \frac{\gamma}{\alpha}, \quad v \in [s, t],$$

and

$$(A.19) \quad \Psi(v) := (t-v)^{\Xi + \frac{\theta}{\alpha}} \|D_x u_i^\delta(v, \cdot)\|_{C^\theta}, \quad v \in [s, t].$$

$$(A.20) \quad \Xi + \frac{\theta}{\alpha} = \frac{1}{\alpha} - \frac{\gamma}{\alpha} + 1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\gamma}{\alpha} \right) = 1 - \varepsilon < 1.$$

In particular, with these notations at hand, since $t-s$ is *small*, it follows from (A.14) that

$$(A.21) \quad \|u_i\|_\infty \leq C \left(\|h\|_\infty + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \left((t-v)^{-\Xi} \Phi(v) + (t-v)^{-(\Xi + \frac{\theta}{\alpha})} \Psi(v) \right) \right).$$

As above, we distinguish the *diagonal* and *off-diagonal* regimes w.r.t. the current considered times. In the *off-diagonal* regime holds, we readily get from (A.18):

$$\begin{aligned}
(t-s)^{\frac{\vartheta}{\alpha}} \frac{|D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x')|}{|x - x'|^\vartheta} &\leq |D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x')| \leq |D_x u_i^\delta(s, x)| + |D_x u_i^\delta(s, x')| \\
&\leq C \left(\|h\|_{\dot{C}^\gamma} (t-s)^{-\frac{1}{\alpha} + \frac{\gamma}{\alpha}} \right. \\
&\quad \left. + \int_s^t dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left(\|u_i^\delta\|_\infty \mathbf{1}_{\{\alpha \leq 1\}} + (t-v)^{-\Xi} \Phi(v) + (t-v)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(v) \right) \right), \\
(t-s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x')|}{|x - x'|^\vartheta} &\leq C \left(\|h\|_{\dot{C}^\gamma} + (t-s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right. \\
&\quad \left. + (t-s)^{1 - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right), \tag{A.22}
\end{aligned}$$

using as well (A.21) for the last inequality.

Let us now turn to the global *diagonal regime*: $|x - x'| \leq (t-s)^{\frac{1}{\alpha}}$. We again need to split as above the analysis in function of the considered running time introducing a change of freezing point. Namely, from Proposition 12 in [CdRMP19], we derive:

$$\begin{aligned}
&D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x') \\
&= \left(\int_s^t dv \mathbf{1}_{v \leq \tau_0} \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy \right. \\
&\quad \left. - \int_s^t dv \mathbf{1}_{v \leq \tau_0} \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi')}(s, v, x', y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi')}(v, y) dy \right) \Big|_{(\tau_0, \tau, \xi, \xi') = (t_0, s, x, x')} \\
&\quad - \left((D_x \tilde{p}_{s, \tau_0}^{\mathbf{P}_i, \alpha, (\tau, \xi')} u_i^\delta(\tau_0, x'))_{(\tau_0, \tau, \xi') = (t_0, s, x')} - (D_x \tilde{p}_{s, \tau_0}^{\mathbf{P}_i, \alpha, (\tau, \xi')} u_i^\delta(\tau_0, x'))_{(\tau_0, \tau, \xi) = (t_0, s, x)} \right) \\
&\quad + \left(\int_{\mathbb{R}^d} dy h^\delta(y) (D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, t, x, y) - D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, t, x', y)) \right) \\
&\quad + \left. \int_s^t dv \mathbf{1}_{v > \tau_0} \int_{\mathbb{R}^d} (D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) - D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x', y)) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy \right) \Big|_{(\tau_0, \tau, \xi) = (t_0, s, x)} \\
&=: \Delta u_{i1}^\delta(s, x, x') + \Delta u_{i2}^\delta(s, x, x') + \Delta u_{i3}^\delta(s, x, x'). \tag{A.23}
\end{aligned}$$

The term $\Delta u_{i1}^\delta(s, x, x')$, corresponding to the local off-diagonal regime within the global diagonal one, can be analyzed as above (see the analysis of the contribution $\Delta w_1(s, x, x')$ introduced in (3.41)). This yields:

$$\begin{aligned}
|\Delta u_{i1}^\delta(s, x, x')| &\leq \left| \int_s^{\tau_0} dv \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x, y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy \right| \Big|_{(\tau_0, \tau, \xi) = (t_0, s, x)} \\
&\quad + \left| \int_s^{\tau_0} dv \int_{\mathbb{R}^d} D_x \tilde{p}^{\mathbf{P}_i, \alpha, (\tau, \xi')}(s, v, x', y) R_{\mathbf{P}_i}^{\alpha, (\tau, \xi')}(v, y) dy \right| \Big|_{(\tau_0, \tau, \xi') = (t_0, s, x')} \\
&\leq C \left(\int_s^{t_0} dv (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} \left[(t-v)^{-\Xi} \Phi(v) + (t-v)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(v) \right] \frac{|x - x'|^\vartheta}{(v-s)^{\frac{\vartheta}{\alpha}}} \right).
\end{aligned}$$

For $v \in [s, t_0]$, $(t-v) \geq (1-c_0)(t-s)$ and therefore:

$$\frac{|\Delta u_{i1}^\delta(s, x, x')|}{|x - x'|^\vartheta} \leq C \left((t-s)^{-\Xi + 1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha} - \frac{\vartheta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t-s)^{-(\Xi + \frac{\vartheta}{\alpha}) + 1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha} - \frac{\vartheta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right),$$

recalling for the last inequality that $-\frac{1}{\alpha} + \frac{2\eta}{\alpha} - \frac{\vartheta}{\alpha} = \frac{2\eta}{\alpha} - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) = \frac{2\eta}{\alpha} - (\frac{1}{\alpha} + 1 - \varepsilon - (\frac{1}{\alpha} - \frac{\gamma}{\alpha})) = -1 + \frac{2\eta}{\alpha} - \frac{\gamma}{\alpha} + \varepsilon > -1$. This eventually gives:

$$\begin{aligned}
(t-s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|\Delta u_{i1}^\delta(s, x, x')|}{|x - x'|^\vartheta} &\leq C \left((t-s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t-s)^{1 - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right). \tag{A.24}
\end{aligned}$$

Turning now to Δu_{13}^δ in (A.23), expanding the frozen densities, exploiting as well (3.19) which gives that a diagonal perturbation of the density does not affect the related estimates, we write:

$$\begin{aligned}
& |\Delta u_{i3}^\delta(s, x, x')| \\
& \leq \left| \int_0^1 d\lambda \int_{\mathbb{R}^d} D_x^2 \bar{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, t, x + \lambda(x' - x), y)(x' - x)[h^\delta(y) - h^\delta(m_{s,t}^{\tau, \xi}(x + \lambda(x' - x)))] dy \right|_{(\tau_0, \tau, \xi) = (t_0, s, x)} \\
& \quad + \left| \int_{\tau_0}^t dv \int_0^1 d\lambda \int_{\mathbb{R}^d} D_x^2 \bar{p}^{\mathbf{P}_i, \alpha, (\tau, \xi)}(s, v, x + \lambda(x' - x), y)(x' - x) \times R_{\mathbf{P}_i}^{\alpha, (\tau, \xi)}(v, y) dy \right|_{(\tau_0, \tau, \xi) = (t_0, s, x)} \\
& \leq C \left(\|h\|_{\dot{C}^\gamma} |x - x'| (t - s)^{-(\frac{2}{\alpha} - \frac{\gamma}{\alpha})} \right. \\
& \quad \left. + |x - x'|^\vartheta \int_{t_0}^t dv (v - s)^{-\frac{2}{\alpha} + \frac{2\eta}{\alpha}} \left((t - v)^{-\Xi} \Phi(v) + (t - v)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(v) \right) |x - x'|^{1-\vartheta} \right).
\end{aligned}$$

Recalling that, for $v \in [t_0, t]$, $|x - x'| \leq ((v - s)/c_0)^{\frac{1}{\alpha}}$, we derive:

$$\begin{aligned}
|\Delta u_{i3}^\delta(s, x, x')| & \leq C \left(\|h\|_{\dot{C}^\gamma} |x - x'|^\vartheta (t - s)^{\frac{\gamma}{\alpha} - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha})} \right. \\
& \quad \left. + |x - x'|^\vartheta \int_{t_0}^t dv (v - s)^{-\frac{1}{\alpha} - \frac{\vartheta}{\alpha} + \frac{2\eta}{\alpha}} \left[(t - v)^{-\Xi} \Phi(v) + (t - v)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(v) \right] \right).
\end{aligned}$$

Therefore, recalling from (A.18)-(A.19) that $\Xi + \frac{\vartheta}{\alpha} = \frac{1-\gamma+\vartheta}{\alpha}$, we get:

$$(A.25) \quad (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|\Delta u_{i3}^\delta(s, x, x')|}{|x - x'|^\vartheta} \leq C \left(\|h\|_{\dot{C}^\gamma} + (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) + (t - s)^{1 - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha}) + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Psi(v) \right).$$

The remaining term Δu_{i2}^δ in (A.23) can be handled exactly as the previous Δw_2 introduced in (3.41), replacing w by u_i^δ . This eventually gives:

$$(A.26) \quad (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|\Delta u_{i2}^\delta(s, x, x')|}{|x - x'|^\vartheta} \leq C c_0^{\frac{\vartheta}{\alpha}} (1 - c_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(t_0).$$

Plugging (A.24), (A.25) and (A.26) into (A.23), we get that in the diagonal case $|x - x'| \leq (t - s)^{\frac{1}{\alpha}}$:

$$\begin{aligned}
(A.27) \quad (t - s)^{\Xi + \frac{\vartheta}{\alpha}} \frac{|D_x u_i^\delta(s, x) - D_x u_i^\delta(s, x')|}{|x - x'|^\vartheta} & \leq C \left(\|h\|_{\dot{C}^\gamma} + (t - s)^{1 - \frac{1}{\alpha} + \frac{2\eta}{\alpha}} \sup_{v \in [s, t]} \Phi(v) \right. \\
& \quad \left. + (t - s)^{1 + \frac{2\eta}{\alpha} - (\frac{1}{\alpha} + \frac{\vartheta}{\alpha})} \sup_{v \in [s, t]} \Psi(v) + c_0^{\frac{\vartheta}{\alpha}} (1 - c_0)^{-(\Xi + \frac{\vartheta}{\alpha})} \Psi(t_0) \right).
\end{aligned}$$

Putting together (A.27) and (A.22), we can conclude the proof of Lemma 3.1 following exactly the procedure for w described in Section 3.1 (using as well (A.14) and (A.18)). This yields the controls in equations (3.4) and (3.5). The proof is now complete. \square

A.3. Proof of Lemma 3.2. We first analyze the sensitivities of the generators w.r.t. the measure arguments or w.r.t to both the spatial points and the measure arguments. The results of Lemma 3.2, i.e. equations (3.6) and (3.7) then readily follow from Lemma A.3 and the a priori bounds of Lemma 3.1.

Lemma A.3 (Sensitivity analysis of the generators w.r.t. the measure and the spatial parameters). Let φ be a function in $\mathcal{C}^{1+\vartheta}(\mathbb{R}^d, \mathbb{R})$ with $\vartheta > \alpha - 1$ if $\alpha \geq 1$. Under (\mathbf{A}_S) it then holds that, there exists a positive constant $C := C((\mathbf{A}_S))$ s.t. for all $v \in [s, t]$:

$$(A.28) \quad |(b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))) \cdot D\varphi(x)| \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \|D\varphi\|_\infty,$$

$$(A.29) \quad \left| (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha}) \varphi(x) \right| \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (\|D\varphi\|_\infty + \|\varphi\|_\infty) \mathbf{1}_{\{\alpha < 1\}} + \|D\varphi\|_{\dot{C}^\vartheta} \mathbf{1}_{\{\alpha > 1\}} + \|\varphi\|_{\mathcal{C}^{1+\vartheta}} \mathbf{1}_{\{\alpha = 1\}}.$$

Furthermore, the following estimates also hold for the differences: there exists a positive constant $C := C((\mathbf{A}_S))$ such that for any $\lambda \in [0, 1]$

$$\begin{aligned}
& \left| (b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))) \cdot D\varphi(x) - (b(v, x', \mathbf{P}_1(v)) - b(v, x', \mathbf{P}_2(v))) \cdot D\varphi(x') \right| \\
(A.30) \quad & \leq C(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^{\lambda 2\eta} \|D\varphi\|_\infty + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^\vartheta \|D\varphi\|_{\dot{C}^\vartheta}), \\
& \left| (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha})\varphi(x) - (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha})\varphi(x') \right| \\
& \leq C\left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^{\lambda 2\eta} (\|D\varphi\|_\infty + \|\varphi\|_\infty) \mathbf{1}_{\{\alpha < 1\}} + \|D\varphi\|_{\dot{C}^\vartheta} \mathbf{1}_{\{\alpha > 1\}} + \|\varphi\|_{C^{1+\vartheta}} \mathbf{1}_{\{\alpha = 1\}}\right) \\
(A.31) \quad & + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^{1+\vartheta-\alpha} \|D\varphi\|_{\dot{C}^\vartheta}.
\end{aligned}$$

Proof. Write first using (\mathbf{B}_H) and (2.11):

$$\left| (b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))) \cdot D\varphi(x) \right| \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \|D\varphi\|_\infty.$$

This gives (A.28).

Write now

$$\begin{aligned}
& (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha})\varphi(x) \\
& = \text{p.v.} \int_{\mathbb{R}^d} [\varphi(x + \sigma(v, x, \mathbf{P}_1(v))z) - \varphi(x)] \nu(dz) - \text{p.v.} \int_{\mathbb{R}^d} [\varphi(x + \sigma(v, x, \mathbf{P}_2(v))z) - \varphi(x)] \nu(dz) \\
& = \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(x + r\theta) - \varphi(x)] \left(\frac{g\left(\frac{(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta}{|(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta|}\right)}{|(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta|^{d+\alpha} \det(\sigma_x^{v, \mathbf{P}_1})} - \frac{g\left(\frac{(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta}{|(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta|}\right)}{|(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta|^{d+\alpha} \det(\sigma_x^{v, \mathbf{P}_2})} \right), \\
(A.32) \quad &
\end{aligned}$$

where we used the absolute continuity condition on the Lévy measure, assumption (\mathbf{AC}) , and the specific structure (2.8) of the density w.r.t. the Lebesgue measure. For simplicity we have also denoted in the above equation for $i \in \{1, 2\}$:

$$(A.33) \quad \sigma_x^{v, \mathbf{P}_i} := \sigma(v, x, \mathbf{P}_i(v)).$$

From now on and for notational convenience, we introduce the following map:

$$(A.34) \quad \mathcal{P}(\mathbb{R}^d) \ni m \mapsto D(v, x, \theta, m) := \frac{g\left(\frac{(\sigma(v, x, m))^{-1}\theta}{|(\sigma(v, x, m))^{-1}\theta|}\right)}{|(\sigma(v, x, m))^{-1}\theta|^{d+\alpha} \det(\sigma(v, x, m))}$$

so that, with our notations

$$\begin{aligned}
& \delta D(v, x, \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) = D(v, x, \theta, \mathbf{P}_1(v)) - D(v, x, \theta, \mathbf{P}_2(v)) \\
(A.35) \quad & = \left(\frac{g\left(\frac{(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta}{|(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta|}\right)}{|(\sigma_x^{v, \mathbf{P}_1})^{-1}\theta|^{d+\alpha} \det(\sigma_x^{v, \mathbf{P}_1})} - \frac{g\left(\frac{(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta}{|(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta|}\right)}{|(\sigma_x^{v, \mathbf{P}_2})^{-1}\theta|^{d+\alpha} \det(\sigma_x^{v, \mathbf{P}_2})} \right).
\end{aligned}$$

Under (\mathbf{D}_H) and (\mathbf{AC}) , using the Lipschitz regularity of g and (2.11) with $\beta = 0$ applied to the diffusion coefficient σ , it is readily seen that:

$$(A.36) \quad |\delta D(v, x, \theta, \mathbf{P}_1(v), \mathbf{P}_2(v))| \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2).$$

Equation (A.29) then readily follows from (A.32)-(A.36) by usual Taylor expansions and the symmetry condition which allows to use cancellation techniques for the small jumps when $\alpha \geq 1$.

Let us now turn to the differences. Write first,

$$\begin{aligned}
& \left| (b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))) \cdot D\varphi(x) - (b(v, x', \mathbf{P}_1(v)) - b(v, x', \mathbf{P}_2(v))) \cdot D\varphi(x') \right| \\
& \leq |b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))| |D\varphi(x) - D\varphi(x')| \\
& \quad + |(b(v, x, \mathbf{P}_1(v)) - b(v, x, \mathbf{P}_2(v))) - (b(v, x', \mathbf{P}_1(v)) - b(v, x', \mathbf{P}_2(v)))| |D\varphi(x')| \\
& \leq C(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^\vartheta \|D\varphi\|_{\dot{C}^\vartheta} + \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2)|x - x'|^{\lambda 2\eta} \|D\varphi\|_\infty),
\end{aligned}$$

using (\mathbf{B}_H) , (2.11) and (2.12) applied to the drift coefficient b , with $\beta = 0$ and $\beta = \lambda$ respectively, for the last inequality. This gives (A.30).

Let us now turn to the difference of the non-local operators. From (A.32) we get with the notation of (A.34):

$$\begin{aligned}
& \left| (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha})\varphi(x) - (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha})\varphi(x') \right| \\
& \leq \left| \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(x + r\theta) - \varphi(x)] \delta D(v, x, \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \right. \\
& \quad \left. - \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(x' + r\theta) - \varphi(x')] \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \right| \\
& \leq \left| \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(x + r\theta) - \varphi(x)] [\delta D(v, x, \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) - \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v))] \right| \\
& \quad + \left| \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(x' + r\theta) - \varphi(x')] - (\varphi(x + r\theta) - \varphi(x)) \right| \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \Big| \\
& =: (\mathcal{D}_1 + \mathcal{D}_2)(v, x, x', \mathbf{P}_1, \mathbf{P}_2).
\end{aligned}
\tag{A.37}$$

Under (\mathbf{AC}) and (\mathbf{D}_H) , using the mean-value theorem and the fact that g has a bounded Lipschitz gradient, we deduce that the continuous map $D(v, x, \theta, \cdot)$ defined by (A.34) admits a bounded and continuous functional derivative such that $(\mathbb{R}^d)^2 \ni (x, y) \mapsto [\delta D(v, x, \theta, m)/\delta m](y)$ is 2η -Hölder uniformly w.r.t. the variables v, θ and m , so that, similarly to (2.12), it is easily seen that there exists C s.t. for all $v \in [s, t], \theta \in \mathbb{S}^{d-1}$ and for any $\lambda \in [0, 1]$

$$|\delta D(v, x, \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) - \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v))| \leq C|x - x'|^{\lambda 2\eta} \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2),$$

which readily gives:

$$\begin{aligned}
& (\text{A.38}) \\
& |\mathcal{D}_1(v, x, x', \mathbf{P}_1, \mathbf{P}_2)| \leq C|x - x'|^{\lambda 2\eta} \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) ((\|D\varphi\|_\infty + \|\varphi\|_\infty) \mathbf{1}_{\{\alpha < 1\}} + \|D\varphi\|_{\dot{C}^\vartheta} \mathbf{1}_{\{\alpha > 1\}} + \|\varphi\|_{C^{1+\vartheta}} \mathbf{1}_{\{\alpha = 1\}}).
\end{aligned}$$

Let us turn now to \mathcal{D}_2 . We will establish:

$$(\text{A.39}) \quad |\mathcal{D}_2(v, x, x', \mathbf{P}_1, \mathbf{P}_2)| \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \|D\varphi\|_{\dot{C}^\vartheta} |x - x'|^{1+\vartheta-\alpha}.$$

Write indeed assuming first that $\alpha < 1$:

$$\begin{aligned}
& |\mathcal{D}_2(v, x, x', \mathbf{P}_1, \mathbf{P}_2)| \\
& \leq \left| \int_0^1 d\lambda \int_{r \leq |x-x'|} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) (D\varphi(x + \lambda r\theta) - D\varphi(x' + \lambda r\theta)) \cdot r\theta \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \right| \\
& \quad + \left| \int_0^1 d\lambda \int_{r \geq |x-x'|} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) \right. \\
& \quad \times (D\varphi(x' + r\theta + \lambda(x - x')) - D\varphi(x' + \lambda(x - x'))) \cdot (x - x') \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \Big| \\
& \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \left(\int_{r \in (0, |x-x'|]} \frac{dr}{r^{1+\alpha}} \|D\varphi\|_{\dot{C}^\vartheta} |x - x'|^\vartheta r + \int_{r \geq |x-x'|} \frac{dr}{r^{1+\alpha}} \|D\varphi\|_{\dot{C}^\vartheta} r^\vartheta |x - x'| \right) \\
& \leq C_{\vartheta, \alpha} \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \|D\varphi\|_{\dot{C}^\vartheta} |x - x'|^{1+\vartheta-\alpha},
\end{aligned}$$

using (A.36) for the last but one inequality and $\alpha > \vartheta$ for the last one (see equation (3.5)).

The only modifications needed for $\alpha \geq 1$ concern the *small jumps*. Indeed, we can introduce the compensator only up to the threshold $|x - x'|$. We are simply led to analyze:

$$\begin{aligned}
& \left| \int_0^1 d\lambda \int_{r \in (0, |x-x'|]} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) ([D\varphi(x + \lambda r\theta) - D\varphi(x)] - [D\varphi(x' + \lambda r\theta) - D\varphi(x')]) \cdot r\theta \right. \\
& \quad \left. \times \delta D(v, x', \theta, \mathbf{P}_1(v), \mathbf{P}_2(v)) \right| \\
& \leq C \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \int_{r \in (0, |x-x'|]} \frac{dr}{r^{1+\alpha}} \|D\varphi\|_{\dot{C}^\vartheta} r^{1+\vartheta} \leq C_{\alpha, \vartheta} \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) \|D\varphi\|_{\dot{C}^\vartheta} |x - x'|^{1+\vartheta-\alpha}.
\end{aligned}$$

This indeed gives (A.39). Plugging (A.38) and (A.39) into (A.37) we derive the statement (A.31). \square

The bound in (3.6) (resp. (3.7)) then readily follows from (A.28), (A.29) (resp. (A.30), (A.31)) and the a priori bounds (3.4), (3.5) which have been established in Lemma 3.1 proven above.

We conclude this section by discussing a bit why in the current McKean-Vlasov setting the absolute continuity condition on the spectral measure (Assumption **(AC)**) seems essential.

Remark A.2 (About the absolute continuity of the spectral measure). The key point is that, when considering the difference of the non-local operators associated with two different measure arguments it seems difficult to make explicitly the distance between the considered measures appear. Let us illustrate this fact considering the elementary following case. Consider $\nu(dz) = \frac{1}{2} \sum_{i=1}^d (\delta_{e_i} + \delta_{-e_i}) \frac{dr}{r^{1+\alpha}}$, where the $(e_i)_{i \in \llbracket 1, d \rrbracket}$ correspond to the canonical vectors, corresponding, up to a suitable normalizing constant to the example of the cylindrical Laplacian. In this case we cannot do the change of variable of (A.32) and directly write for the difference:

$$\begin{aligned} & \left| (L_v^{\mathbf{P}_1, \alpha} - L_v^{\mathbf{P}_2, \alpha}) \varphi(x) \right| \\ & \leq \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right|. \end{aligned} \quad (\text{A.40})$$

As usual, the strategy consists in separating the small and large jumps. Let us first consider the case $\alpha < 1$. In that case we write, for a threshold \mathcal{T} (possibly depending on $(\sigma_x^{v, \mathbf{P}_i})_{i \in \{1,2\}}$):

$$\begin{aligned} & \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_0^{\mathcal{T}} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right| \\ & = \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_0^{\mathcal{T}} \frac{dr}{r^{1+\alpha}} \int_0^1 d\lambda D\varphi(x + (-1)^j ((1-\lambda)\sigma_x^{v, \mathbf{P}_1} + \lambda\sigma_x^{v, \mathbf{P}_2}) e_i r) \cdot (-1)^j (\sigma_x^{v, \mathbf{P}_1} - \sigma_x^{v, \mathbf{P}_2}) r e_i \right| \\ & \leq C_{\alpha, \mathcal{T}} \|D\varphi\|_{\infty} \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)). \end{aligned}$$

Hence, since the required term $\mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v))$ appears this way, it seems natural to consider a threshold \mathcal{T} at a *macro* scale, and in particular independent of $\sigma_x^{v, \mathbf{P}_2} - \sigma_x^{v, \mathbf{P}_1}$. This choice anyhow yields problems for the *large jumps*, i.e. those above the threshold \mathcal{T} . Namely,

$$\begin{aligned} & \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_{\mathcal{T}}^{+\infty} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right| \\ & \leq \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \int_{\mathcal{T}}^{+\infty} \frac{dr}{r^{1+\alpha}} \|\varphi\|_{\dot{C}^{\beta}} |\sigma_x^{v, \mathbf{P}_2} - \sigma_x^{v, \mathbf{P}_1}|^{\beta} r^{\beta}, \end{aligned}$$

for any $\beta \in (0, 1)$ s.t. $\beta < \alpha$ in order to preserve some integrability and where $\|\varphi\|_{\dot{C}^{\beta}}$ again stands for the β -Hölder modulus of φ (homogeneous Hölder norm) which can be estimated from $\|\varphi\|_{\infty}, \|D\varphi\|_{\infty}$ through elementary interpolation. Hence, we cannot recover that way the required control with exactly the distance, we only end up with:

$$\begin{aligned} & \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_{\mathcal{T}}^{+\infty} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right| \\ & \leq C_{\alpha, \beta, \mathcal{T}} \|\varphi\|_{\dot{C}^{\beta}} \mathbf{d}_{\eta}^{\beta}(\mathbf{P}_1(v), \mathbf{P}_2(v)). \end{aligned}$$

One could object that we do not have exploited the *full* regularity of $\varphi \in \mathcal{C}^{1+\vartheta}(\mathbb{R}^d, \mathbb{R})$. But for the small jumps this will not a priori improve the dependence in the distance (see also the small jump part for $\alpha > 1$ below).

If now $\alpha > 1$ the problems are reversed. Indeed, the large jumps can be readily controlled with the expected bound. Indeed,

$$\begin{aligned} & \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_{\mathcal{T}}^{+\infty} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right| \\ & = \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_{\mathcal{T}}^{+\infty} \frac{dr}{r^{1+\alpha}} \|D\varphi\|_{\infty} |\sigma_x^{v, \mathbf{P}_2} - \sigma_x^{v, \mathbf{P}_1}| r \right| \leq C_{\alpha, \mathcal{T}} \|D\varphi\|_{\infty} \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)). \end{aligned}$$

On the other hand, for the small jumps, we are led to use explicitly the ϑ -Hölder regularity of the gradient $D\varphi$ (recall from (3.5) that $\vartheta = \vartheta(\alpha, \eta, \varepsilon) = \alpha \left[1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\eta}{\alpha}\right)\right]$). Write precisely,

$$\begin{aligned}
& \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} \text{p.v.} \int_0^{\mathcal{T}} \frac{dr}{r^{1+\alpha}} \left([\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - \varphi(x)] - [\varphi(x + (-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) - \varphi(x)] \right) \right| \\
&= \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} (-1)^j \text{p.v.} \int_0^{\mathcal{T}} \frac{dr}{r^{1+\alpha}} \left(\int_0^1 d\lambda [D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) \cdot \sigma_x^{v, \mathbf{P}_1} \pm D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) \cdot \sigma_x^{v, \mathbf{P}_2} \right. \right. \\
&\quad \left. \left. - D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i) \cdot \sigma_x^{v, \mathbf{P}_2}] \right) r e_i \right| \\
&= \left| \sum_{i=1}^d \frac{1}{2} \sum_{j \in \{0,1\}} (-1)^j \text{p.v.} \int_0^{\mathcal{T}} \frac{dr}{r^{1+\alpha}} \left(\int_0^1 d\lambda [(D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - D\varphi(x)) \cdot (\sigma_x^{v, \mathbf{P}_1} - \sigma_x^{v, \mathbf{P}_2}) \right. \right. \\
&\quad \left. \left. + (D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_1} r e_i) - D\varphi(x + \lambda(-1)^j \sigma_x^{v, \mathbf{P}_2} r e_i)) \cdot \sigma_x^{v, \mathbf{P}_2}] \right) r e_i \right| \\
&\leq C \int_0^{\mathcal{T}} \frac{dr}{r^\alpha} \|D\varphi\|_{\dot{C}^\vartheta} r^\vartheta \left(\mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) + \mathbf{d}_\eta^\vartheta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \right) \\
&\leq C_{\alpha, \vartheta} \|D\varphi\|_{\dot{C}^\vartheta} \left(\mathbf{d}_\eta(\mathbf{P}_1(v), \mathbf{P}_2(v)) + \mathbf{d}_\eta^\vartheta(\mathbf{P}_1(v), \mathbf{P}_2(v)) \right),
\end{aligned}$$

where we used a cancellation argument for the first inequality, and exploit as well that $\alpha - \vartheta = 1 - \eta + \varepsilon\alpha < 1$ for ε small enough, for the last one. In any case we see that this is not enough to conclude whereas assumption **(AC)** allows to transfer all the sensitivity analysis to the spectral measure (see again the proof of the previous lemma).

A.4. Proof of Lemma 3.3. The stated controls follow reproducing the *transfer* arguments of the diffusion coefficient to the density of the Lévy measure in equation (A.32) for the integro-differential part. The gradient term is controlled directly. Namely, from (3.11),

$$\begin{aligned}
|R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)| &\leq |(b(v, y, \mathbf{P}_1(v)) - b(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v))) \cdot Dw(v, y)| + |(L_s^{\mathbf{P}_1, \alpha} - \tilde{L}_s^{\mathbf{P}_1, \alpha, (\tau, \xi)})w(v, y)| \\
&\leq \|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{C}^{2\eta}} |y - \theta_{v, \tau}(\xi)|^{2\eta} |Dw(v, y)| + \left| \text{p.v.} \int_{\mathbb{R}^d} [\varphi(y + \sigma(v, y, \mathbf{P}_1(v))z) - \varphi(y)] \nu(dz) \right. \\
&\quad \left. - \text{p.v.} \int_{\mathbb{R}^d} [\varphi(y + \sigma(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v))z) - \varphi(y)] \nu(dz) \right| \\
&\leq \|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{C}^{2\eta}} |y - \theta_{v, \tau}(\xi)|^{2\eta} |Dw(v, y)| + \left| \text{p.v.} \int_0^{+\infty} \frac{dr}{r^{1+\alpha}} \int_{\mathbb{S}^{d-1}} \Lambda_{\mathbb{S}^{d-1}}(d\theta) [\varphi(y + r\theta) - \varphi(y)] \right. \\
&\quad \times \left(\frac{g\left(\frac{(\sigma(v, y, \mathbf{P}_1(v)))^{-1}\theta}{|(\sigma(v, y, \mathbf{P}_1(v)))^{-1}\theta|}\right)}{|(\sigma(v, y, \mathbf{P}_1(v)))^{-1}\theta|^{d+\alpha} \det(\sigma(v, y, \mathbf{P}_1(v)))} \right. \\
&\quad \left. - \frac{g\left(\frac{(\sigma(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v)))^{-1}\theta}{|(\sigma(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v)))^{-1}\theta|}\right)}{|(\sigma(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v)))^{-1}\theta|^{d+\alpha} \det(\sigma(v, \theta_{v, \tau}(\xi), \mathbf{P}_1(v)))} \right) \Big|.
\end{aligned}$$

Again the Lipschitz regularity of g , the non-degeneracy and 2η -Hölder spatial continuity of σ and the symmetry of the Lévy measure yield:

$$\begin{aligned}
|R_{\mathbf{P}_1}^{\alpha, (\tau, \xi)}(v, y)| &\leq C \left(\|b(v, \cdot, \mathbf{P}_1(v))\|_{\dot{C}^{2\eta}} |D_y w(v, y)| + \|\sigma(v, \cdot, \mathbf{P}_1(v))\|_{C^{2\eta}} [(\|w\|_\infty + \|D_x w(v, \cdot)\|_\infty) \mathbf{1}_{\{\alpha < 1\}} \right. \\
&\quad \left. + \|D_x w(v, \cdot)\|_{C^\vartheta} \mathbf{1}_{\{\alpha > 1\}} + \|w(v, \cdot)\|_{C^{1+\vartheta}} \mathbf{1}_{\{\alpha = 1\}}] |y - \theta_{v, \tau}(\xi)|^{2\eta}, \right. \\
\end{aligned} \tag{A.41}$$

which is the required statement.

APPENDIX B. THEOREM 2.3 IN THE SUPERCRITICAL CASE UNDER THE SOLE CONDITION $\alpha > 2\eta \vee (1 - \eta)$

In order to establish Theorem 2.3 under **(A_S)** and the sole condition $\alpha > 2\eta \vee (1 - \eta)$ in the supercritical case $\alpha < 1$, we proceed as follows. Restarting from (3.21) with $h \in C^{\gamma_1}$, with $\alpha + \gamma_1 > 1$ and $\gamma_1 \in (0, 2\eta]$ to be specified later on and employing (3.6) instead of (3.23), similarly to (3.27), we derive:

$$\begin{aligned}
\text{(B.1)} \quad |w(s, x)| &\leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t - s)^{1 - (\frac{1}{\alpha} - \frac{2\eta}{\alpha})} + (t - s) \|w\|_\infty + \int_s^t dv (v - s)^{\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right)
\end{aligned}$$

Observe that, for the first term of the above r.h.s., we have assumed that $(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha}) < 1 \iff \alpha + \gamma_1 > 1$. This seems to be somehow a necessary condition to derive a smoothing effect for the contribution $H_{\mathbf{P}_1, \mathbf{P}_2}^\delta$.

Let us now proceed from the cancellation argument as in (3.28). We get from (3.7) that in the current case:

$$(B.2) \quad |H_{\mathbf{P}_1, \mathbf{P}_2}^\delta u_2^\delta(v, y) - H_{\mathbf{P}_1, \mathbf{P}_2} u_2^\delta(v, \theta_{v, \tau}(\xi))| \leq C \mathbf{d}_{(1-\lambda)2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) |y - \theta_{v, \tau}(\xi)|^{2\lambda\eta} (t-v)^{-(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha})} \\ + C \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) |y - \theta_{v, \tau}(\xi)|^\vartheta (t-v)^{-1+\varepsilon}$$

recalling from (3.5) that $0 < \vartheta := \alpha \left[1 - \varepsilon - \left(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha} \right) \right] < \alpha - 1 + \gamma_1$.

Now, similarly to (3.33), using (B.2) on the time interval $[s, \frac{t+s}{2}]$ and (3.6) on the time interval $[\frac{t+s}{2}, t]$, taking also λ large enough so that $2\lambda\eta + \alpha > 1$, we get:

$$(B.3) \quad |Dw(s, x)| \leq C \left(\int_s^{\frac{t+s}{2}} (v-s)^{-\frac{1}{\alpha} + \frac{2\lambda\eta}{\alpha}} \mathbf{d}_{(1-\lambda)2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) (t-v)^{-\frac{1-\gamma_1}{\alpha}} dv \right. \\ + \int_s^{\frac{t+s}{2}} (v-s)^{-\frac{1}{\alpha} + \frac{\vartheta}{\alpha}} \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) (t-v)^{-1+\varepsilon} dv \\ + \int_{\frac{t+s}{2}}^t (v-s)^{-\frac{1}{\alpha}} \mathbf{d}_{2\eta}(\mathbf{P}_1(v), \mathbf{P}_2(v)) (t-v)^{-\frac{1-\gamma_1}{\alpha}} dv \\ + \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w(v, \cdot)\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \Big) \\ \leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{2\lambda\eta}{\alpha} - \frac{\gamma_1}{\alpha})} + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{\gamma_1}{\alpha})} \right. \\ \left. + \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \right),$$

provided that ϑ is sufficiently large so that

$$(B.4) \quad \vartheta + \alpha > 1.$$

With the notation of (3.34) for $(\Phi(v))_{v \in [s, t]}$, taking in our current supercritical case $\Xi = \frac{1}{\alpha} - \frac{\gamma_1}{\alpha} + \frac{\varepsilon}{2} < 1$, we derive from (B.1) and (B.3):

$$(B.5) \quad |w(s, x)| \leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha})} + (t-s) \|w\|_\infty + \sup_{v \in [s, t]} \Phi(v) \int_s^t (v-s)^{\frac{2\eta}{\alpha}} (t-v)^{-\Xi} dv \right) \\ \leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha})} + (t-s) \|w\|_\infty + \sup_{v \in [s, t]} \Phi(v) (t-s)^{1+\frac{2\eta}{\alpha} - \Xi} \right),$$

and

$$|Dw(s, x)| \leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{2\lambda\eta}{\alpha} - \frac{\gamma_1}{\alpha})} + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{\gamma_1}{\alpha})} \right. \\ + \sup_{v \in [s, t]} \Phi(v) \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (t-v)^{-\Xi} dv \Big) \\ \leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{2\lambda\eta}{\alpha} - \frac{\gamma_1}{\alpha})} + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{\gamma_1}{\alpha})} + \sup_{v \in [s, t]} \Phi(v) (t-s)^{1-(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha}) - \Xi} \right)$$

which in turn gives

$$(t-s)^\Xi |Dw(s, x)| \leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{2\lambda\eta}{\alpha} - \frac{\gamma_1}{\alpha}) + \Xi} + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{\gamma_1}{\alpha}) + \Xi} \right. \\ \left. + \sup_{v \in [s, t]} \Phi(v) (t-s)^{1-(\frac{1}{\alpha} - \frac{\gamma_1}{\alpha})} \right).$$

Taking into account that $0 < t-s \leq T$ is sufficiently small, from the preceding inequality, we obtain

$$\sup_{v \in [s, t]} \Phi(v) \leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{2\lambda\eta}{\alpha} - \frac{\gamma_1}{\alpha}) + \Xi} + \mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{2}{\alpha} - \frac{\gamma_1}{\alpha}) + \Xi} \right)$$

and plugging the previous estimate into (B.5) yields

$$\begin{aligned}
|w(s, x)| &\leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) [(t-s)^{1-(\frac{1}{\alpha}-\frac{\gamma_1}{\alpha})} + (t-s)^{2-(\frac{2}{\alpha}-\frac{2\eta}{\alpha}-\frac{\gamma_1}{\alpha})}] + \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha}-\frac{2\lambda\eta}{\alpha}-\frac{2\eta}{\alpha}-\frac{\gamma_1}{\alpha})} \right. \\
&\quad \left. + (t-s)\|w\|_\infty \right) \\
&\leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{1}{\alpha}-\frac{\gamma_1}{\alpha})} + \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha}-\frac{2\lambda\eta}{\alpha}-\frac{2\eta}{\alpha}-\frac{\gamma_1}{\alpha})} + (t-s)\|w\|_\infty \right)
\end{aligned}$$

where for the last inequality we used the fact that $2\eta + \alpha > 1$.

We now select $\gamma_1 = 2\eta$. We importantly recall the condition on the regularization parameter ϑ : $1 - \alpha < \vartheta < \alpha - 1 + \gamma = \alpha - 1 + 2\eta$ so that a necessary condition is $\alpha + \eta > 1$. Taking into account this condition, we next optimize the previous inequality on w w.r.t $h \in \mathcal{C}^\gamma(\mathbb{R}^d)$. We thus deduce

$$(B.6) \quad \mathbf{d}_{2\eta, s, t}(\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)) \leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{1-(\frac{1}{\alpha}-\frac{2\eta}{\alpha})} + \mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha}-\frac{2\lambda\eta}{\alpha}-\frac{4\eta}{\alpha})} \right)$$

for any $\lambda \in [0, 1]$ s.t. $2\lambda\eta + \alpha > 1$ under the assumption that $\alpha + \eta > 1$.

Our final step consists in iterating the previous analysis. We aim at establishing an estimate similar to (B.6) but for the map \mathbf{T}^2 and a distance $\mathbf{d}_{(1-\lambda)2\eta, s, t}$ for a well-chosen $\lambda \in [0, 1]$. Keeping in mind the inequality (B.6), we proceed in a similar manner but replacing \mathbf{P}_i by $\mathbf{T}(\mathbf{P}_i)$. For $h \in \mathcal{C}^{\gamma_2}$, we consider $u_i^\delta(s, x)$, $i = 1, 2$ satisfying

$$(B.7) \quad \begin{cases} \left(\partial_s + \mathcal{A}_s^{\mathbf{T}(\mathbf{P}_i)} \right) u_i^\delta(s, x) = 0, & (s, x) \in [0, t) \times \mathbb{R}^d, \\ u_i^\delta(t, x) = h^\delta(x), & x \in \mathbb{R}^d, \end{cases}$$

and $u_i^\delta \in L^\infty([0, t], C^{\alpha+\gamma_2}(\mathbb{R}^d, \mathbb{R}))$.

To compare both semigroups, we again write the PDE satisfied by $w(s, x) := w^\delta(s, x) = (u_1^\delta - u_2^\delta)(s, x)$. Namely,

$$(B.8) \quad \begin{cases} \left(\partial_s + \mathcal{A}_s^{\mathbf{T}(\mathbf{P}_1)} \right) w(s, x) = - \left([b(s, x, \mathbf{T}(\mathbf{P}_1)) - b(s, x, \mathbf{T}(\mathbf{P}_2))] \cdot D_x u_2^\delta(s, x) \right. \\ \quad \left. + (L_s^{\mathbf{T}(\mathbf{P}_1), \alpha} - L_s^{\mathbf{T}(\mathbf{P}_2), \alpha}) u_2^\delta(s, x) \right) =: -H_{\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)} u_2^\delta(s, x), & (s, x) \in [0, t) \times \mathbb{R}^d, \\ w(t, x) = 0, & x \in \mathbb{R}^d. \end{cases}$$

The inequality (3.23) now becomes

$$(B.9) \quad |H_{\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)} u_2^\delta(v, y)| \leq C \mathbf{d}_{2\eta}(\mathbf{T}(\mathbf{P}_1)(v), \mathbf{T}(\mathbf{P}_2)(v)) (t-v)^{-\frac{1-\gamma_2}{\alpha}}.$$

This new estimate will be combined with (B.6) in the sequel. Similarly to (B.1), from (B.6), we get

$$\begin{aligned}
|w(s, x)| &\leq C \left(\mathbf{d}_{2\eta, s, t}(\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)) (t-s)^{1-(\frac{1}{\alpha}-\frac{\gamma_2}{\alpha})} + (t-s)\|w\|_\infty + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right) \\
(B.10) \quad &\leq C \left(\mathbf{d}_{(1-\lambda)2\eta, s, t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha}-\frac{2\eta}{\alpha}-\frac{\gamma_2}{\alpha})} + (t-s)\|w\|_\infty + \int_s^t dv (v-s)^{\frac{2\eta}{\alpha}} \|Dw(v, \cdot)\|_\infty \right)
\end{aligned}$$

where we used the direct inequality $\mathbf{d}_{2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2) \leq \mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2)$. Remark also that similarly to (B.3), using directly (B.9), as soon as $\eta + \alpha > 1$ and $\gamma_2 + \alpha > 1$, one has

$$\begin{aligned}
|Dw(s, x)| &\leq C \left(\int_s^t (v-s)^{-\frac{1}{\alpha}} \mathbf{d}_{2\eta}(\mathbf{T}(\mathbf{P}_1)(v), \mathbf{T}(\mathbf{P}_2)(v)) (t-v)^{-\frac{1-\gamma_2}{\alpha}} dv \right. \\
&\quad \left. + \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w(v, \cdot)\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \right) \\
&\leq C \left(\int_s^t (v-s)^{-\frac{1}{\alpha}} \mathbf{d}_{2\eta,s,v}(\mathbf{T}(\mathbf{P}_1), \mathbf{T}(\mathbf{P}_2)) (t-v)^{-\frac{1-\gamma_2}{\alpha}} dv \right. \\
&\quad \left. + \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w(v, \cdot)\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \right) \\
&\leq C \mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2) \int_s^t (v-s)^{1-(\frac{2}{\alpha} - \frac{2\eta}{\alpha})} (t-v)^{-\frac{1-\gamma_2}{\alpha}} dv \\
&\quad + C \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w(v, \cdot)\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \\
(B.11) \quad &\leq C \left(\mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha} - \frac{2\eta}{\alpha} - \frac{\gamma_2}{\alpha})} + \int_s^t (v-s)^{-\frac{1}{\alpha} + \frac{2\eta}{\alpha}} (\|w(v, \cdot)\|_\infty + \|Dw(v, \cdot)\|_\infty) dv \right),
\end{aligned}$$

for any λ s.t. $2\lambda\eta + \alpha > 1$ and where we used (B.6) for the last but one inequality. We now proceed in a completely analogous manner as the first iteration. We thus conclude

$$\|w\|_\infty \leq C \mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha} - \frac{2\eta}{\alpha} - \frac{\gamma_2}{\alpha})}.$$

We now select $\gamma_2 = (1-\lambda)2\eta$ with λ large enough so that $2\lambda\eta + \alpha > 1$ and small enough so that $\alpha + \gamma_2 > 1$. Note that this choice is licit under our current assumption $\alpha + \eta > 1$ and that it also yields $2 - (\frac{2}{\alpha} - \frac{2\eta}{\alpha} - \frac{\gamma_2}{\alpha}) > 0$. We next optimize the previous inequality on w w.r.t $h \in \mathcal{C}^{\gamma_2}(\mathbb{R}^d)$

$$\mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{T}^{(2)}(\mathbf{P}_1), \mathbf{T}^{(2)}(\mathbf{P}_2)) \leq C \mathbf{d}_{(1-\lambda)2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2) (t-s)^{2-(\frac{2}{\alpha} - \frac{2\eta}{\alpha} - \frac{\gamma_2}{\alpha})}.$$

This shows that $\mathbf{T}^{(2)}$ is a contraction on the complete metric space $\mathbf{A}_{s,s+T,\mu}$ w.r.t the distance $\mathbf{d}_{(1-\lambda)2\eta,s,s+T}$ under the condition $\alpha + \eta > 1$ provided T is sufficiently small. According to the Banach fixed point theorem, \mathbf{T} admits a unique fixed point which is the unique solution to the martingale problem on the time interval $[s, s+T]$. We eventually conclude following the same lines of reasonings as those employed at the end of the proof of Theorem 2.3. We omit the remaining technical details.

Remark B.1. We point out that in the previous argument only two iterations of the map \mathbf{T} were needed to derive our result. One may naturally ask if one can do better by performing more iteration of \mathbf{T} . This is not clear since the necessary condition $\alpha + \eta > 1$ (coming from the regularization parameter ϑ) already appears at the first iterate and seems to be the natural one if one wants to be able to compare the distance $\mathbf{d}_{\gamma,s,t}(\mathbf{T}(\mathbf{P}_1, \mathbf{P}_2))$ with the distance $\mathbf{d}_{2\eta,s,t}(\mathbf{P}_1, \mathbf{P}_2)$ in order to prove a contraction property.

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NOUFEL FRIKHA, UNIVERSITÉ DE PARIS, LABORATOIRE DE PROBABILITÉS, STATISTIQUES ET MODÉLISATION, F-75013 PARIS, FRANCE.
E-mail address: `frikha@math.univ-paris-diderot.fr`

VALENTIN KONAKOV, LABORATORY OF STOCHASTIC ANALYSIS, HIGHER SCHOOL OF ECONOMICS, POKROVSKY BOULEVARD 11, MOSCOW, RUSSIAN FEDERATION.
E-mail address: `vkonakov@hse.ru`

STÉPHANE MENOZZI, LABORATOIRE DE MODÉLISATION MATHÉMATIQUE D'EVRY (LAMME), UMR CNRS 8070, UNIVERSITÉ D'EVRY VAL D'ESSONNE, UNIVERSITÉ PARIS-SACLAY, 23 BOULEVARD DE FRANCE 91037 EVRY AND LABORATORY OF STOCHASTIC ANALYSIS, HIGHER SCHOOL OF ECONOMICS, 31 SHABOLOVKA, MOSCOW, RUSSIAN FEDERATION.
E-mail address: `stephane.menozzi@univ-evry.fr`