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Projected gradient descent for non-convex sparse spike estimation

Yann Traonmilin^{1,2}, Jean-François Aujol² and Arthur Leclaire²

Abstract—We propose a new algorithm for sparse spike estimation from Fourier measurements. Based on theoretical results on non-convex optimization techniques for off-the-grid sparse spike estimation, we present a projected gradient descent algorithm coupled with a spectral initialization procedure. Our algorithm permits to estimate the positions of large numbers of Diracs in 2d from random Fourier measurements. We present, along with the algorithm, theoretical qualitative insights explaining the success of our algorithm. This opens a new direction for practical off-the-grid spike estimation with theoretical guarantees in imaging applications.

Index Terms—spike super-resolution, non-convex optimization, projected gradient descent

I. INTRODUCTION

In the space $\mathcal{M} = \mathcal{M}(\mathbb{R}^d)$ (respectively $\mathcal{M} = \mathcal{M}(\mathbb{T}^d)$) of finite signed measures over \mathbb{R}^d (respectively the d -dimensional torus \mathbb{T}^d), we aim at recovering a superposition of impulsive sources $x_0 = \sum_{i=1}^k a_i \delta_{t_i} \in \mathcal{M}$ from the measurements

$$y = Ax_0 + e, \quad (1)$$

where δ_{t_i} is the Dirac measure at position t_i , the operator A is a linear observation operator from \mathcal{M} to \mathbb{C}^m , $y \in \mathbb{C}^m$ are the m noisy measurements and e is a finite energy observation noise. This inverse problem (called spike super-resolution [5], [2], [15], [7], [8]) models many imaging problems found in geophysics, microscopy, astronomy or even (compressive) machine learning [13]. Under a separation assumption on the positions of the Diracs, i.e when x_0 is in a set $\Sigma_{k,\epsilon}$ of sums of k ϵ -separated Diracs with bounded support, it has been shown that x_0 can be estimated by solving a non-convex problem as long as A is an appropriately designed measurement process. This ideal non-convex minimization is:

$$x^* \in \operatorname{argmin}_{x \in \Sigma_{k,\epsilon}} \|Ax - y\|_2^2. \quad (2)$$

Recovery guarantees for this problem are of the form

$$\|x^* - x_0\|_K \leq C\|e\|_2, \quad (3)$$

where $\|\cdot\|_K$ is a kernel norm on \mathcal{M} that measures distances in \mathcal{M} at a given high resolution described by the kernel (in most of the literature K is either a Fejér [5] or Gaussian kernel and $\|\sum_i a_i \delta_{t_i}\|_K^2 = \sum_{i,j} a_i a_j K(t_i - t_j)$ [12], [11]). Mathematically, (3) is guaranteed if the measurement operator A has a restricted isometry property on $\Sigma_{k,\epsilon} - \Sigma_{k,\epsilon}$ (the set of differences of elements of $\Sigma_{k,\epsilon}$) [11]. This property

is typically obtained when the number of measurements is sufficient. For example, recovery guarantees are obtained when $m \geq O(\frac{1}{\epsilon^d})$ for regular low frequency Fourier measurements on the torus [5] and when $m \geq O(k^2 d (\log(k))^2 \log(kd/\epsilon))$ for random Fourier measurements on \mathbb{R}^d [11].

Recent advances in this field proposed a convex relaxation of the problem in the space of measures [5], [7]. While giving theoretical recovery guarantees, these methods are not convex with respect to the parameters due to a polynomial root finding step. Moreover, they rely on a SDP relaxation of a dual formulation, thus squaring the size of the problem (which becomes problematic as the number of parameters increases). Other methods based on greedy heuristics (CL-OMP for compressive k -means [13]) have been proposed. Nevertheless, they still lack theoretical justifications in this context while having a good scaling properties with respect to the number of parameters (amplitudes and positions) even if some first theoretical results are emerging for some particular measurement methods [9].

In this paper, we propose a practical method to solve the non-convex minimization problem (2) for a large number of Diracs in imaging problems. Of course, at first sight, it is not possible to solve this problem efficiently. However, we justify qualitatively why our method succeeds. This justification relies on the separation assumption on x_0 and the assumption that there are enough measurements of x_0 . We also give numerical experiments validating the method. One of the main practical advantages of our method is its ability to perform off-the-grid spike estimation from random Fourier measurements with a good scaling with respect to the number of spikes. With this proof of concept, we can estimate many spikes in two dimensions from compressive measurements, yielding potential applications in fields such as astronomy or microscopy where the sum of spikes model is relevant.

Our method, following insights from the literature on non-convex optimization for low-dimensional models [17], [14], [4], [6], relies on two steps:

- Overparametrized spectral initialization: we propose a spectral initialization step for spike estimation that permits a good first estimation of the positions of the Diracs.
- Projected gradient descent algorithm in the parameter space: the idea of projected gradient descent for low-dimensional model recovery has shown its benefits in the finite dimensional case [3], [10]. We adapt this idea to sparse spike recovery. From [16], the global minimizer of (2) can be recovered by unconstrained gradient descent as long as the initialization lies in an explicit basin of attraction of the global minimizer. It was also shown that projecting on the

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separation constraint improves the control on the Hessian of the function we minimize.

Contributions. After recalling that the non-convex sparse spike estimation problem in the space of parameters is a smooth non convex constrained optimization problem, we propose a new projected gradient descent algorithm for sparse spike estimation. The simple gradient descent is already used as a refinement step in greedy algorithms. We show experimentally and justify qualitatively that adding a projection step and using an appropriate initialization leads to a global convergence.

- In Section II, we describe our projected gradient descent algorithm and its implementation details;
- In Section III, we propose a grid based spectral initialization scheme. A qualitative analysis shows that when the number of measurements is large enough, our initialization approximates well the Diracs positions;
- In Section IV, we show the practical benefit of the projection in the descent algorithm and its application to the estimation of large number of Diracs in 2 dimensions.

II. THEORETICAL BACKGROUND AND ALGORITHM DESCRIPTION

A. Measurements and parameter space

The operator A is a linear operator modeling m measurements in \mathbb{C}^m ($\text{Im} A \subset \mathbb{C}^m$) on the space of measures on a domain E (either $E = \mathbb{R}^d$ or $E = \mathbb{T}^d$) defined by:

$$\forall l = 1, \dots, m, \quad (Ax)_l = \int_E \alpha_l(t) dx(t), \quad (4)$$

where $(\alpha_l)_l$ is a collection of (weighted) Fourier measurements: $\alpha_l(t) = c_l e^{-j(\omega_l, t)}$ for some chosen frequencies $\omega_l \in \mathbb{R}^d$ and frequency dependent weights $c_l \in \mathbb{R}$. The model set of ϵ -separated Diracs with $\epsilon > 0$ is:

$$\Sigma_{k, \epsilon} := \left\{ \sum_{r=1}^k a_r \delta_{t_r} : a \in \mathbb{R}^k, t_r \in \mathcal{B}_2(R), \right. \\ \left. \forall r \neq l, \|t_r - t_l\|_2 \geq \epsilon \right\}, \quad (5)$$

where $\mathcal{B}_2(R) = \{t \in \mathbb{R}^d : \|t\|_2 \leq R\}$ is the ℓ^2 ball of radius R centered in 0 in \mathbb{R}^d . We consider the following parametrization of $\Sigma_{k, \epsilon}$: for any $\theta = (a_1, \dots, a_k, t_1, \dots, t_k) \in \mathbb{R}^{k(d+1)}$, we define $\phi(\theta) = \sum_{i=1}^k a_i \delta_{t_i}$, and we set

$$\Theta_{k, \epsilon} := \phi^{-1}(\Sigma_{k, \epsilon}), \quad (6)$$

the reciprocal image of $\Sigma_{k, \epsilon}$ by ϕ . Note that any parametrization of elements of $\Sigma_{k, \epsilon}$ is invariant by permutation of the positions. This is not a problem in practice for the convergence of descent algorithms. We define the parametrized functional

$$g(\theta) := \|A\phi(\theta) - y\|_2^2 \quad (7)$$

and consider the problem

$$\theta^* \in \arg \min_{\theta \in \Theta_{k, \epsilon}} g(\theta). \quad (8)$$

Since the α_l are smooth, g is a smooth function. Note that performing the minimization (8) allows to recover the minima

of the ideal minimization (2), yielding stable recovery guarantees under a restricted isometry assumption on A which is verified when $m \geq O(k^2 d (\log(k))^2 \log(kd/\epsilon))$ for adequately chosen Gaussian random Fourier measurements (on \mathbb{R}^d) and $m \geq O(\frac{1}{\epsilon^d})$ for regular Fourier measurements on \mathbb{T}^d . In [16], it has been shown that the simple gradient descent converges (without projection) to the global minimum of g as long as the initialization falls into an explicit basin of attraction of this global minimum. It was also shown that the projection on the separation constraint improves the control on the Hessian on g and subsequently the convergence of the descent algorithm.

B. Projected gradient descent in the parameter space

For a user-defined initial number of Diracs k_{in} , we consider the following iterations:

$$\theta_{n+1} = P_{\Theta_{k_{in}, \epsilon}}(\theta_n - \tau_n \nabla g(\theta_n)) \quad (9)$$

where $P_{\Theta_{k_{in}, \epsilon}}$ is a projection on the separation constraint, (notice that there may be several solutions in $\Theta_{k_{in}, \epsilon}$) and τ_n is the step size at iteration n .

The projection $P_{\Theta_{k_{in}, \epsilon}}(\theta)$ could be defined naturally as a solution of the minimization problem $\inf_{\tilde{\theta} \in \Theta_{k_{in}, \epsilon}} \|\phi(\tilde{\theta}) - \phi(\theta)\|_K$. Unfortunately this optimization is not convex. To avoid this, we propose a heuristic (see Algorithm 1) for $P_{\Theta_{k_{in}, \epsilon}}$ that consists in merging Diracs that are not ϵ -separated.

Input: List $\Theta = (a_i, t_i)_i$ of amplitudes and positions ordered by decreasing absolute amplitudes

```

for  $i \geq 1$  do
  for  $j > i$  do
    if  $\|t_i - t_j\| < \epsilon$  then
       $a_i = a_i + a_j$ ;
       $t_i = \frac{|a_i|t_i + |a_j|t_j}{|a_i| + |a_j|}$ 
      Remove  $(a_j, t_j)$  from  $\Theta$ 
    end
  end
end

```

Output: List $\Theta = (a_i, t_i)_i$ of amplitudes and positions of projected spikes

Algorithm 1: Heuristic for the projection $P_{\Theta_{k_{in}, \epsilon}}$

Since we take the barycenter of the positions, if a set of Diracs that are at a distance at most ϵ' of a true position in x_0 is merged, the merged result will be within this distance. After this projection step, we pursue the descent with the remaining number of Diracs. Note that we overparametrize with k_{in} the number of Diracs in the descent to ensure the recovery of all positions in x_0 (see also the next section).

In practice, we implement the projected gradient descent as follows.

- As suggested in [16], to avoid balancing problems between amplitudes and positions, we alternate descent steps between amplitudes and positions.
- To find the step size τ_n , we perform a line search to minimize the value of the function g .

- We start to project after a few iterations of the descent to increase the reduction of dimension of the projection.

From [16], this algorithm will converge as soon as the initialization falls into a basin of attraction of global minimum of g . The basins of attraction get larger as the number of measurements increases (up to a fundamental limit depending on the separation ϵ and the amplitudes in x_0).

III. OVERPARAMETRIZED SPECTRAL INITIALIZATION

The idea of spectral initialization was used for non-convex optimization in the context of phase recovery [17] and blind deconvolution [4]. As we measure the signal x_0 at some frequencies ω_l , a natural way to recover an estimation of the signal is to back-project the irregular spectrum on a grid Γ that samples $\mathcal{B}_2(R)$ at a given precision ϵ_g (to be chosen later). Given a vector of Fourier measurements y at frequencies $(\omega_l)_{l=1,m}$, we calculate $z_\Gamma = B_\Gamma y$ where B_Γ is the linear operator back-projecting the Fourier measurements on a grid in the spatial domain \mathbb{R}^d , it is defined as

$$z_\Gamma = B_\Gamma y := \sum_{s_i \in \Gamma} z_{\Gamma,i} \delta_{s_i} \quad (10)$$

where the $s_i \in \Gamma$ are the grid positions and

$$z_{\Gamma,i} = \sum_l y_l d_l e^{j(\omega_l, s_i)} \quad (11)$$

for some weights d_l to be chosen in the next section. We can show in the noiseless case with Lemmas III.1 and III.2 that when the number of measurements increases and the grid for initialization gets finer, the original positions of Diracs get better approximated. Since the energy of Diracs is well localized by the initialization we then perform overparametrized hard thresholding of the initialization. We propose the initialization θ_{init} defined by

$$\phi(\theta_{init}) := x_{init} = H_{k_{in}}(B_\Gamma y) \quad (12)$$

where for $|z_{\Gamma,j_1}| \geq |z_{\Gamma,j_2}| \geq \dots |z_{\Gamma,j_n}|$, we have $H_{k_{in}}(z_\Gamma) = \sum_{i=1}^{k_{in}} z_{\Gamma,j_i} \delta_{s_{j_i}}$.

A. Ideal spectral initialization and sampling

In the context of Diracs recovery our proposed spectral initialization scheme is a sampling of an ideal back-projection. Let B the operator from \mathbb{C}^m to \mathcal{M} defined for $z = By$ by

$$z(t) = \sum_{l=1}^m d_l y_l e^{j(\omega_l, t)} \quad (13)$$

We call z an ideal spectral initialization because $z_\Gamma = S_\Gamma z$ where S_Γ is the sampling on the grid Γ : for a measure x with a continuous density χ (i.e. $dx(t) = \chi(t) dt$), we define the sampling operation $S_\Gamma(x) = \sum_{t_i \in \Gamma} \chi(t_i) \delta_{t_i}$. Also, the measure $z = By$ has a smooth density as it is a finite sum of complex exponentials.

We first show (all proofs are in the supplementary material) that for the right choice of weights d_l , the energy of $z = BAx$ (where $x = \sum_{i=1}^k a_i \delta_{t_i}$) is localized around the positions t_i in both the regular Fourier sampling on the torus case, and the random Fourier sampling on \mathbb{R}^d case. We show the following

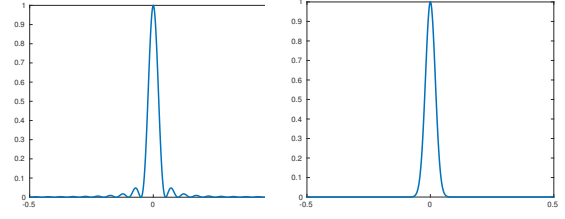


Fig. 1. Kernels used for a separation $\epsilon = 0.1$. Left: On the torus the Fejér kernel of maximum frequency $\frac{2}{\epsilon}$. Right: on \mathbb{R} , the Gaussian kernel of parameter $\sigma = \frac{1}{500\epsilon}$

results for the Fejér and Gaussian kernels as they are typically used in the literature for deterministic [5] and random [11] Fourier sampling.

Lemma III.1. *On $\mathcal{M}(\mathbb{T}^d)$, we choose A such that $(\omega_l)_{l=1,m}$ is a regular sampling of $[-\omega_{max}, \omega_{max}]^d$ with $\omega_l \in 2\pi \cdot \mathbb{Z}^d$. In (13), take $d_l = \hat{K}_f(\omega_l) / ((2\pi)^d c_l)$ where \hat{K}_f is the Fourier transform of the Fejér kernel K_f on the torus whose Fourier spectrum support is $(\omega_l)_{l=1,m}$, then*

$$z(t) = \sum_{i=1}^k a_i K_f(t - t_i). \quad (14)$$

This immediate lemma just states that on the torus, measuring low frequencies is equivalent to measuring a low-pass filtered signal in the time domain. For example, the low-pass Fejér filter is shown Figure 1. Also, this result holds for any kernel with spectrum supported on the ω_l . Sampling frequencies regularly with maximum frequency $\omega_{max} \geq O(\frac{1}{\epsilon^d})$ guarantees recovery with convex relaxation methods.

For random Fourier sampling, we look at the expected value of z and control its variance with respect to the distribution of the ω_l .

Lemma III.2. *On $\mathcal{M}(\mathbb{R}^d)$, we choose A such that the ω_l are m i.i.d random variables with a Gaussian distribution with density $G(\omega_r) = \frac{\sigma^d}{(\sqrt{2\pi})^d} e^{-\frac{\sigma^2}{2} \|\omega_r\|_2^2}$. Let $K_g(t) = e^{-\frac{\|t\|_2^2}{2\sigma^2}}$. In (13), take $d_l = 1/(m c_l)$ then*

$$E(z(t)) = \sum_{i=1}^k a_i K_g(t - t_i) \quad (15)$$

$$E(|z(t) - E(z(t))|^2) = -\frac{1}{m} |E(z(t))|^2 + \frac{1}{m} \|x_0\|_{K_g}^2 \quad (16)$$

where $\|x_0\|_{K_g}^2$ is the norm associated with the kernel K_g .

Similarly to the regular sampling, the energy of the expected value of z is concentrated around the positions t_i (see Figure 1). In [11] the frequency distribution scales as the inverse of the kernel precision, i.e. the kernel parameter σ of the kernel $h(t) = e^{-\|t\|_2^2/(2\sigma^2)}$ is chosen as $O(1/\epsilon)$.

The control of the variance shows that when the number of measurements increases, the back-projection of these measurements to the space of measures are closer to the ideal initialization which is the expected value of z . In practice we set the number of measurements using a rule $m = \alpha k d$ with a multiplicative parameter α . The quality of the initialization is thus directly linked to α .

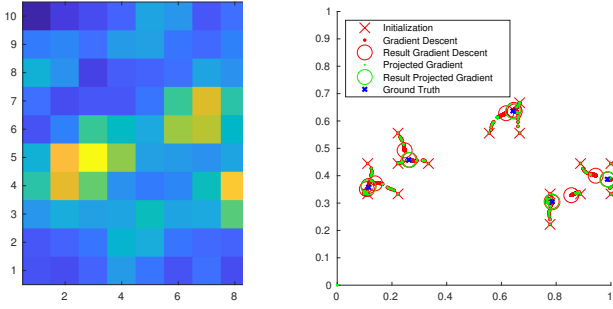


Fig. 2. Result for a few spikes in 2d. Left: back-projection of measurements on a grid. Right: Initialization, gradient descent and projected gradient descent trajectories.

Finally the following lemma makes sure that as the grid gets finer we recover all the energy of the ideal spectral initialization that lies within the domain sampled by the grid.

Lemma III.3. *Let $z_d = (z_{\Gamma,i})_{i \in \Gamma}$. Then $\|\sqrt{\epsilon_g^d} z_d\|_2^2 \rightarrow_{\epsilon_g \rightarrow 0} \|z\|_{L^2(\mathcal{B}_2(R))}^2$.*

We considered the noiseless case. The noisy case just adds a noise term with energy controlled by the noise energy level $\|e\|_2$ because B_Γ is a Fourier back-projection.

IV. NUMERICAL EXPERIMENTS

We first run the algorithm on few Diracs in 2d to illustrate the added benefit of the projection. We then show results with many Diracs in 2d to show the computational feasibility of projected gradient descent for imaging applications. We perform the experiments in the noiseless case and leave the study of the impact of the noise for future work. The Matlab code used to generate these experiments is available at [1].

A. Illustration with few Diracs

As a first proof of concept we run the algorithm with the recovery of 5 Diracs in 2 dimensions from $m = 120$ Gaussian random measurements. The trajectories of 500 iterations of the gradient descent and projected gradient descent are represented in Figure 2. We observe that while the gradient descent with overparametrized initialization might converge with a large number of iterations, the projection step greatly accelerates the convergence.

B. Estimation of 100 Diracs in 2d

We recover 100 Diracs, with a separation 0.01 on the square $[0, 1] \times [0, 1]$ from $m = 2000$ compressive Gaussian measurements (we would need ≈ 10000 regular measurements to obtain a separation 0.01). In practice, the grid Γ must be fine enough to overparametrize the number of Diracs with a good sampling of the ideal spectral initialization. If ϵ_g is too small, the number of initial Diracs needed to sample the energy gets larger, leading to an increased cost in the first iterations of the gradient descent. In this example we use $\epsilon_g = \epsilon$ and use $k_{in} = 4k$. We observe in Figure 3 that with these parameters all the Diracs positions are well estimated after 500 iterations of our algorithm. Similarly to our first example, we observe that spikes that are not separated in the back-projection on the grid are well estimated by our algorithm.

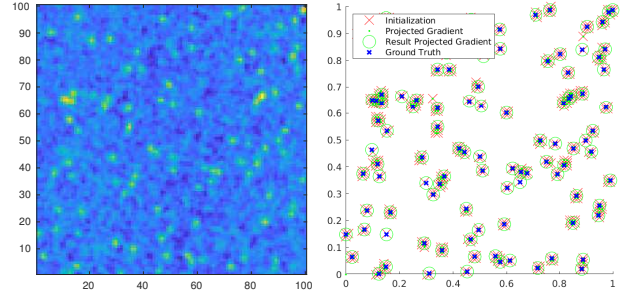


Fig. 3. Result for 100 spikes in 2d. Left: back-projection of measurements on a grid. Right: Initialization, and projected gradient descent trajectories.

C. Complexity

The cost of our algorithm is the sum of the cost of the initialization and the cost of the projected gradient descent algorithm. Of course, the back-projection on the grid scales as $O((1/\epsilon_g)^d)$ (irregular Fourier transform on a grid), but it is done only once and projections using fast transform techniques could be estimated. With our strategy, this cost seems unavoidable as we want to localize Diracs off-the-grid at a precision of the order of ϵ (doing the same on the grid would have this exponential scaling with respect to the dimension and the separation). With this cost the algorithm stays tractable when the dimension d of the domain of the positions of the Diracs is not too large, which is the case in 2d or 3d imaging applications. This cost $O((1/\epsilon_g)^d)$ would also be impossible to avoid in an eventual visualization of the full recovered image over a precise grid on \mathbb{R}^d . Note also that in practice our proposed initialization could be replaced by any satisfying overparametrized initialization, i.e. any state-of-the-art on-the-grid estimation technique could benefit from an added projected gradient descent.

The cost of the projected gradient descent is essentially $O(n_{it} C_\nabla)$ where C_∇ is the cost the calculation of the gradient. This cost is of the order of the calculation of the m Fourier measurements for the current number of Diracs in the descent (which is close to k after a few iterations).

V. CONCLUSION

We gave a practical algorithm to perform off-the-grid sparse spike estimation. This proof-of-concept shows that it is possible to estimate efficiently a large number of Diracs in imaging applications with some strong theoretical insights of success guarantees. Future research directions are:

- Full theoretical convergence proof of the algorithm with sufficient conditions on the number of measurements. One of the main question is to determine if it is possible to have a convergence guarantee without the computational cost $O((1/\epsilon_g)^d)$
- Investigate other methods for reducing the number of parameters after the back-projection on a grid.
- Investigate quasi-Newton schemes to accelerate the descent (some Hessian information is available).
- Study the algorithm stability to noise and modeling error with respect to the number of measurements.

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