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► **To cite this version:**

Marie Malbois, Boumediene Nedjar, Stéphane Lavaud, Claude Rospars, Loïc Divet, et al.. On DEF expansion modelling in concrete structures under variable hydric conditions. *Construction and Building Materials*, 2019, 207, pp. 396-402. 10.1016/j.conbuildmat.2019.02.142 . hal-02310874

HAL Id: hal-02310874

<https://hal.science/hal-02310874v1>

Submitted on 25 May 2021

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On DEF expansion modelling in concrete structures under variable hydric conditions

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Abstract

Delayed Ettringite Formation (DEF) in concrete is likely to develop in massive civil engineering structures such as bridges, nuclear plants, and dams with major security issues. In many cases, DEF pathology can lead to swelling and cracking which may significantly impact mass transfer and mechanical properties. It is then of major importance to build predictive tools for engineering conceptions and expertises. In this contribution, the chemical swelling evolution is integrated within the overall constitutive law of concrete that, besides, can experience other phenomena such like damage, plasticity, and long term creep, not all considered here. On another hand, as DEF is activated by environmental humidity above a certain threshold, we introduce the notion of effective time that takes into account the cumulative exposition above this threshold. Hence, a special care is taken with regards to the chem-

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Preprint submitted to Construction and Building Materials

February 15, 2019

ical irreversibility, together with the humidity-drying cycles. Computations are used to calibrate various sets of model parameters with the help of results from the literature, on the one hand, and from an experimental campaign where a calcareous aggregates-based concrete is studied, on the other hand. We show the efficiency of the developed numerical tool through a series of numerical examples.

Keywords: DEF pathology modelling, Variable humidity, Effective time concept

1. Introduction

Delayed ettringite formation (DEF) in concrete can have significant impact on the long term behaviour of massive concrete structures. The swelling induced by the development of this pathology generates cracking and so **can in general affect the mechanical properties, the transfer properties, for instance see [1, 2], and the durability of the concrete structures, for instance [3]**. It is then of major importance to build numerical tools that can simulate this expanding kinematics and its consequences on the structural serviceability. Ettringite formation is delayed under a particular set of conditions, that have been unequally thoroughly studied: cement compositions ([4],[5],[6],[7] among others), aggregate mineralogical nature ([8],[9],[10],[11], among others), curing ([12],[13],[7]) or environmental humidity. Few authors have focused on the study the influence of the ambient humidity of the material. Water has a significant role in the delayed ettringite formation. It has been noticed that the concrete structures with DEF pathology were in an environment with a high ambient relative humidity (submerged parts, inflows of

17 water, or exposition to high humidity). Water seems to be both a reaction
18 agent and environment, as ettringite crystallisation requires 32 molecules of
19 water ([14]), but it also allows the transport of other agents that influence the
20 formation of DEF ([15]). Indeed, it boosts the leaching of alkalines out of the
21 interstitial water in the material, leading to a decrease of pH which favours
22 the formation of ettringite. It is nowadays unanimously accepted that the
23 DEF rate and amplitude is increased by the increase of saturated water, thus
24 environmental humidity, e.g. see for example [16], [15], [17], [18], [19]. For
25 instance in [17] and [19], the authors have even demonstrated experimentally
26 this impact by putting reactive mortar and concrete in different controlled
27 environments from 75%RH to 100%RH. Based on macroscopic observations,
28 they identified a RH threshold of about 92% to initiate DEF.

29 From the mathematical point of view, extensive research has been con-
30 ducted to model expansions caused by DEF, see for example [18, 13, 20, 21,
31 among others] where expansions range around 0.25% to 1.3%, **even, up to**
32 **3% [22]**, and further developments are nowadays still ongoing. The models
33 are mostly based on phenomenological approaches formulated in terms of in-
34 ternal variables. The expansion is in general characterized by two important
35 ingredients:

- 36 (i) An amplitude that depends on the temperature history at early age
37 due to hydration and/or on heating conditions if a curing process is
38 employed such as for pre-casting;
- 39 (ii) A kinetics that depends on the time, more precisely on the cumulated
40 time, of exposure in contact with water.

41 Of interest in this work is the introduction of a new notion; the concept
42 of *effective* time, that we denote throughout by \tilde{t} , and that is controlled by
43 the amount of water (humidity) during time and, consequently, that influ-
44 ences the kinetics of expansion, i.e. the above ingredient (ii). This extends
45 the applicability of the existing expansion models not only for cases of full
46 saturation, but for varying environmental conditions as well. Among others,
47 the effective time depends on a threshold below which the swelling process
48 stops, and above which swelling strongly depends on humidity as well as on
49 temperature. Within the continuum, the effective time is an internal field
50 variable since each material point \boldsymbol{x} has its own humidity history during the
51 real time t ; $\tilde{t} \equiv \tilde{t}(\boldsymbol{x}, t)$.

52 An outline of the remainder of this paper is as follows: we first recall
53 the basic constitutive equations together with the kinematical assumption
54 we adopt in this work. The notion of effective time is then motivated and
55 detailed with a focus on the most relevant points involved by the present for-
56 mulation. Finally, we present a set of numerical simulations to illustrate the
57 effectiveness of the proposed framework that compare satisfactorily against
58 experimental data from an experimental campaign and from results in the
59 literature.

60 **2. Swelling kinematics**

61 The kinematical choice is as usual based on an additive split of the *total*
62 strain tensor $\boldsymbol{\varepsilon}$ into an elastic part $\boldsymbol{\varepsilon}^e$ and complementary parts, each one
63 corresponding to a phenomenon. To be as clear as possible, let us consider

64 the simplest choice,

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}^e + \boldsymbol{\varepsilon}_{\text{th}} + \boldsymbol{\varepsilon}_{\text{hyd}} + \boldsymbol{\varepsilon}_{\chi} , \quad (1)$$

65 where $\boldsymbol{\varepsilon}_{\text{th}}$ and $\boldsymbol{\varepsilon}_{\text{hyd}}$ are respectively the thermal and hydric dilatations, and
 66 $\boldsymbol{\varepsilon}_{\chi}$ the chemical expansion tensor, in our case due to DEF. We can use the
 67 classical relations for the formers as:

$$\boldsymbol{\varepsilon}_{\text{th}} = \alpha (T - T_0) \mathbf{1} , \quad \boldsymbol{\varepsilon}_{\text{hyd}} = \varpi (S_r - S_{r0}) \mathbf{1} , \quad (2)$$

68 where T is the temperature, S_r is the saturation, α and ϖ are respectively the
 69 thermal and hydric dilatation coefficients, assumed constant for simplicity,
 70 and $\mathbf{1}$ is the second-order identity tensor. Here T_0 and S_{r0} are the initial
 71 temperature and the initial saturation, respectively. Notice that a form based
 72 on the relative humidity can equivalently be used instead of the saturation-
 73 based form (2)₂. **In future contributions, inelastic deformations through a**
 74 **plastic strain $\boldsymbol{\varepsilon}^P$ together with creep through a viscoelastic strain $\boldsymbol{\varepsilon}^V$ will be**
 75 **appended to the decomposition (1) in a classical fashion, i.e. $\boldsymbol{\varepsilon} \equiv \boldsymbol{\varepsilon}^e + \boldsymbol{\varepsilon}_{\text{th}} +$**
 76 **$\boldsymbol{\varepsilon}_{\text{hyd}} + \boldsymbol{\varepsilon}_{\chi} + \boldsymbol{\varepsilon}^P + \boldsymbol{\varepsilon}^V$, see for example [23, 24] for similar couplings.**

77 Now if we consider a purely volumetric chemical expansion, we write:

$$\boldsymbol{\varepsilon}_{\chi} = \varepsilon_{\chi} \mathbf{1} , \quad (3)$$

78 where the scalar functional $\varepsilon_{\chi} \equiv \varepsilon_{\chi}(S_r, T, t \dots)$ is the so-called free chemical
 79 expansion, the expression of which depends on the humidity, the temperature,
 80 and of course on the time as well.

81 In fully saturated conditions ($S_r = 1$ all the time), one has the nowadays
 82 well known expression used for both of DEF and Alkali-Aggregate Reaction

83 (AAR) expansions phenomena, see for example [10, 25, 18, among others] :

$$\varepsilon_\chi = \varepsilon_\infty^0 \frac{1 - e\left(-\frac{t}{\tau_c}\right)}{1 + e\left(-\frac{t-\tau_l}{\tau_c}\right)} \left(1 - \frac{\phi}{t + \delta}\right), \quad (4)$$

84 where,

- 85 • ε_∞^0 is the potential chemical strain that constitutes the amplitude of
- 86 expansion, i.e. ingredient (i) in the Introduction Section,
- 87 • τ_c and τ_l are respectively the characteristic and latent times that char-
- 88 acterize the kinetics of expansion,
- 89 • and ϕ and δ are parameters that control the long term kinetics. They
- 90 are such that $\delta > \phi$. Equation (4) reduces to Larive's law [26] when ϕ
- 91 is set to zero.

92 For the case of a DEF analysis, ε_∞^0 depends on the thermal history at

93 *early-age*. To fix the ideas the following definition can be used for its mod-

94 elling, see for instance [25, 27]:

$$\varepsilon_\infty^0 = \bar{\alpha} \int_0^{t_m} \begin{cases} 0 \, dt, & \text{if } T \leq T_{\text{def}} \\ e\left[-\frac{E_a^{\text{def}}}{R} \frac{1}{T - T_{\text{def}}}\right] dt, & \text{if } T > T_{\text{def}} \end{cases} \quad (5)$$

95 where the constant $\bar{\alpha}$ is a material parameter, T_{def} is the threshold of temper-

96 ature above which DEF can be generated (about 65°C), E_a^{def} is the activation

97 energy relative to DEF expansion, and t_m is the maturation time (few days

98 to few weeks depending on the size of the considered structure).

99 Now to take into account the influences of both of the thermal and hydric

100 conditions, relation (4) must be adapted. To be exhaustive, let us assume

101 the following choices that not all will be considered in this step of the devel-

102 opments due to the lack of experimental results.

103 *2.1. Variable thermal conditions*

104 It has been shown experimentally that the ambient temperature has an
 105 influence on the kinetics of expansion. Among other choices, we can consider
 106 that the two characteristic times be thermo-activated with the forms

$$\tau_l = \bar{\tau}_l e^{\left[\frac{U_l}{R} \left(\frac{1}{T} - \frac{1}{\bar{T}}\right)\right]}, \quad \tau_c = \bar{\tau}_c e^{\left[\frac{U_c}{R} \left(\frac{1}{T} - \frac{1}{\bar{T}}\right)\right]}, \quad (6)$$

107 as suggested in [26, 28] for AAR. Here $\bar{\tau}_l$ and $\bar{\tau}_c$ are reference characteristic
 108 times for a reference temperature \bar{T} , and U_l and U_c are activation energies.
 109 Notice that, even possible, the amplitude ε_∞^0 is not affected by temperature
 110 in our DEF case.

111 *2.2. Variable hydric conditions*

112 Experimental evidences have also shown that the expansion stops below
 113 a certain threshold of humidity, see for example [19]. The relation (4) must
 114 then be adapted to take into account this strong dependency. In terms of
 115 saturation, denoting the threshold below which the expansion reaction stops
 116 by \bar{S}_r , we can formally write:

$$\begin{cases} \dot{\varepsilon}_\chi = 0, & \text{if } S_r \leq \bar{S}_r, \\ \dot{\varepsilon}_\chi > 0, & \text{if } S_r > \bar{S}_r, \\ \text{and } \varepsilon_\chi \text{ is given by (4),} & \text{if } S_r = 1. \end{cases} \quad (7)$$

117 Now for $\bar{S}_r < S_r < 1$, we need to establish a *continuous* link between the
 118 extreme situations (7)₁ and (7)₃. We introduce for this an *effective* time that
 119 we denote by \tilde{t} and such that,

$$\tilde{t} \equiv \tilde{t}(S_r, t, \dots) \in [0, 1]. \quad (8)$$

120 A possible choice would be:

$$\dot{\tilde{t}} = \left(\frac{\langle S_r - \bar{S}_r \rangle_+}{1 - \bar{S}_r} \right)^{\bar{m}} \quad \Longrightarrow \quad \tilde{t} = \int_0^t \left(\frac{\langle S_r - \bar{S}_r \rangle_+}{1 - \bar{S}_r} \right)^{\bar{m}} dt, \quad (9)$$

121 where the Macauley bracket $\langle \cdot \rangle_+$ denotes the positive part function, and the
 122 exponent parameter most probably depends on the saturation, i.e.

$$\bar{m} \equiv \bar{m}(S_r). \quad (10)$$

123 Hence, by replacing the real time t by the effective time \tilde{t} into the expres-
 124 sion (4), we obtain a free expansion law that covers the requirements (7) for
 125 variable humidity.

126 3. Constitutive equations and mechanical balance

127 The kinematic decomposition, Eq. (1), must now be embedded into a
 128 constitutive relation. At this point, the simplest choice is to consider an
 129 elastic relation for the reversible behaviour:

$$\boldsymbol{\sigma} = \mathbf{C} : \boldsymbol{\varepsilon}^e, \quad (11)$$

130 where $\boldsymbol{\sigma}$ is the stress tensor and the elastic part of the strain tensor $\boldsymbol{\varepsilon}^e$ is
 131 the one that has been used in the decomposition (1). \mathbf{C} is the fourth-order
 132 elasticity tensor that can in turn be affected by chemical damage as

$$\mathbf{C} = (1 - d_\chi) \mathbf{C}_0, \quad (12)$$

133 where \mathbf{C}_0 is the elastic modulus for the undamaged concrete (Hooke's law)
 134 and d_χ is a damage variable in the sense of continuum damage mechanics, see
 135 for example [24, 29]. Intuitively, it can be driven by the chemical expansion

136 itself, then explicitly given as a function of the quantity ε_χ , for instance, the
 137 following form as adopted in [25]:

$$d_\chi = 1 - e^{\left[-\omega \langle \varepsilon_\chi - \varepsilon_{\text{trs}} \rangle_+\right]}, \quad (13)$$

138 where ε_{trs} is the strain-like chemical damage threshold below which no dam-
 139 age occurs, and $\omega \geq 0$ a convenient parameter, i.e. no chemical damage takes
 140 place if we take $\omega = 0$.

141 With (11) we have a minimalist modelling framework with the simplest
 142 possible resolution procedure. The mechanical balance is linear since the
 143 thermal and hydric fields are a priori known at every time step, i.e. assuming
 144 a weak coupling between the thermo-hydric diffusion and the mechanical
 145 equilibrium. *Let us notice however that when inelastic strains (plasticity)*
 146 *will be taken into account together with creep and mechanical damage, the*
 147 *problem clearly becomes nonlinear. Furthermore, the thermal and hydric*
 148 *parts of the modelling are left out of the scope of this paper. Nevertheless,*
 149 *let us stress that for the hydric part at least, care must be taken when solving*
 150 *for the saturation field. Indeed, on the one hand, the hydric diffusion is not*
 151 *constant and, on the other hand, concrete materials are characterized by*
 152 *typical hysteretic sorption-desorption responses that strongly influence the*
 153 *soaking-drying kinetics, to mention a few. Developments on this topic are*
 154 *still in progress.*

155 At the actual time t_{n+1} for instance, the weak form of the mechanical

156 balance at hand is given by:

$$\begin{aligned}
& \int_{\mathcal{B}} \nabla^s \delta \mathbf{u} : \mathbf{C}_{n+1} : \nabla^s \mathbf{u}_{n+1} \, dV = G_{n+1}^{\text{ext}}(\delta \mathbf{u}) \\
& + \int_{\mathcal{B}} \nabla^s \delta \mathbf{u} : \mathbf{C}_{n+1} : (\boldsymbol{\varepsilon}_{\text{th}_{n+1}} + \boldsymbol{\varepsilon}_{\text{hyd}_{n+1}} + \boldsymbol{\varepsilon}_{\chi_{n+1}}) \, dV ,
\end{aligned} \tag{14}$$

157 which must hold for any displacement variation $\delta \mathbf{u}$, and where $\nabla^s(\cdot)$ is the
158 symmetric gradient operator. Here G_{n+1}^{ext} is a short hand notation for the
159 virtual work of the external loads embedding both of the volumetric forces
160 in the body \mathcal{B} and traction forces on part of its boundary $\partial_t \mathcal{B} \subset \partial \mathcal{B}$ applied
161 at time t_{n+1} .

162 Equation (14) is to be solved for the actual displacement field \mathbf{u}_{n+1} . How-
163 ever, care must be taken in evaluating the chemical strain $\boldsymbol{\varepsilon}_{\chi_{n+1}}$ on the right
164 hand-side of (14). Indeed, this latter is computed with the effective time \tilde{t}_{n+1}
165 that must in turn be updated locally at each time step as:

$$\tilde{t}_{n+1} = \tilde{t}_n + \left(\frac{\langle S_{r_{n+1}} - \bar{S}_r \rangle_+}{1 - \bar{S}_r} \right)^{\bar{m}_{n+1}} \Delta t , \tag{15}$$

166 where $\Delta t = t_{n+1} - t_n$ is the *real* time increment, and the exponent parameter
167 \bar{m}_{n+1} is evaluated with $S_{r_{n+1}}$, i.e. $\bar{m}_{n+1} \equiv \bar{m}(S_{r_{n+1}})$ in the form (10). Conse-
168 quently, the effective time is treated as an internal field variable ; $\tilde{t} \equiv \tilde{t}(\mathbf{x}, t)$.

169 4. Simulations of the influence of the relative humidity on DEF 170 expansions

171 In this section we present simulations with regards to experimental tests
172 given in [19] on the study of the influence of relative humidity on expansion
173 associated with DEF in concrete. A set of 11×22 cm cylindrical samples

174 made with a siliceous-based aggregate concrete were heat-treated at early age
175 so as to potentially trigger DEF expansion, and then separated in different
176 groups, each group submitted to a controlled constant relative humidity RH,
177 see [19] for details:

- 178 - Samples that have been stored at 94%, 96%, 98% and 100% RH, re-
179 spectively denoted by RH94, RH96, RH98 and RH100. ;
- 180 - Samples that have been immersed continuously, denoted by IM ;
- 181 - Samples that have been stored at 91% RH and then immersed in water
182 at the age of 334 days, denoted by RH91-IM.

183 The experimental results are shown in figure 1 where interesting observa-
184 tions are to be pointed out:

- 185 - Expansions under continuous immersion and RH100 (not shown here)
186 are almost identical;
- 187 - With a relative humidity of 91%, the saturation is under the threshold
188 \bar{S}_r ;
- 189 - Above the saturation threshold, the expansion is very sensitive to the
190 ambient humidity, i.e. from 94 to 100% RH.

191 From the point of view of the modelling, the curve obtained at saturation
192 $S_r = 1$ that corresponds to the case of RH100 or IM, is used to identify the
193 parameters of the free-expansion law (4). They are summarized in Table 1
194 where the asymptotic parameters ϕ and δ are used for a better description
195 of the long term expansion kinetics.

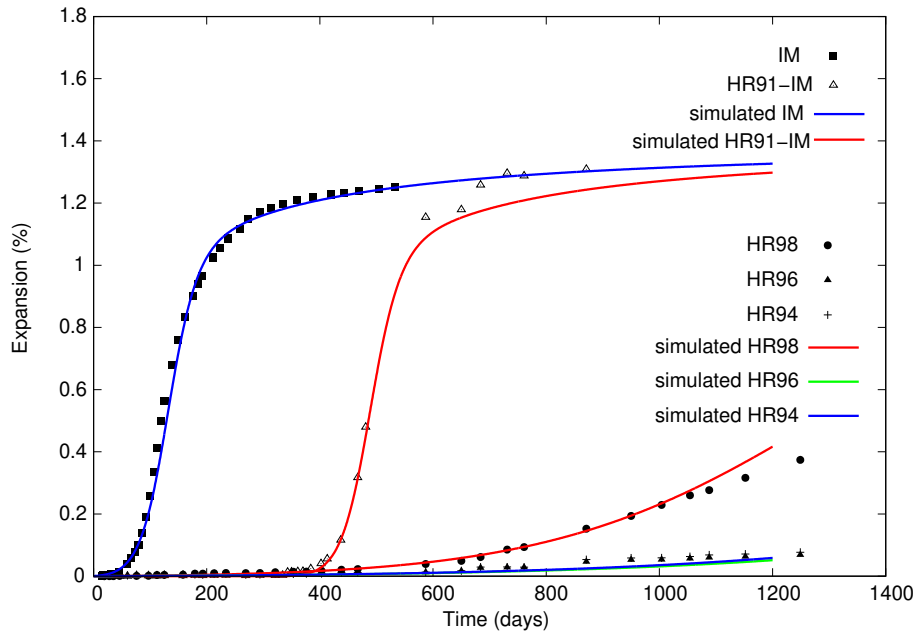


Figure 1: Expansions of all concrete specimens. Experimental results from (Al Shamaa & al., 2015). Superposition with the results of the numerical simulations.

196 Now for the effective time, the property $(7)_2$ is active in this example,
 197 so the expression (9) must be used. We take $\bar{S}_r = 0.93$ for the saturation
 198 threshold, and the exponent parameter \bar{m} has the form (10) chosen here as
 199 a quadratic function given by,

$$\bar{m}(S_r) = 1 + 42.5(S_r - \bar{S}_r) + 1750(S_r - \bar{S}_r)^2. \quad (16)$$

200 With this set of parameters at hand, we show in figure 1 the numerical
 201 results predicted by the model for the different exposures to humidity. One
 202 can notice that the model is able to represent the experimental behaviour.

Table 1: Material parameters with the expansion law (4) for the siliceous concrete of figure 1.

| Definition | identified values |
|--|-------------------|
| expansion amplitude ε_{∞}^0 | 1.4% |
| latent time τ_l | 128 days |
| characteristic time τ_c | 25 days |
| parameter ϕ | 68 days |
| parameter δ | 100 days |

203 5. DEF expansions under cyclic soaking-drying cycles

204 An experimental campaign has been performed in this work as well. Con-
 205 crete samples were designed for their known expansive behaviour when heat-
 206 treated and immersed in water after casting. The concrete formulation used
 207 is detailed in Table 2; A Portland cement CEMII/A L with 6% of calcare-
 208 ous additions (42.5 MPa) is used with a 0.57 water-cement ratio. Calcareous
 209 aggregates were used (Boulonnais aggregates).

210 Now to simulate the heating induced by cement hydration in massive
 211 structures, the samples undergo a thermic treatment as shown in figure 2:
 212 the temperature was maintained at $20^{\circ}C$ for 2h, then increased from $20^{\circ}C$
 213 to $80^{\circ}C$ in 24 h, then maintained at $80^{\circ}C$ for three days, and then decreased
 214 from $80^{\circ}C$ to $20^{\circ}C$ in three days. After this treatment, half of the samples
 215 were immersed in water at $20^{\circ}C$; the other half undergo soaking and drying
 216 cycles. In all cases, 11×22 cm cylindrical samples have been used, and are
 217 monitored following the LCPC method n°66.

Table 2: Concrete composition.

| Materials | Composition |
|-------------------------------|------------------------|
| Cement | 350 kg |
| Water | 210 kg |
| Sand (0-5 mm) | 858 kg |
| Aggregates (5-12.5 mm) | 945 kg |
| Viscosity modifying admixture | 6.6 kg |
| W/C | 0.57 |
| A/S | 1.1 |
| Density | 2354 kg/m ³ |

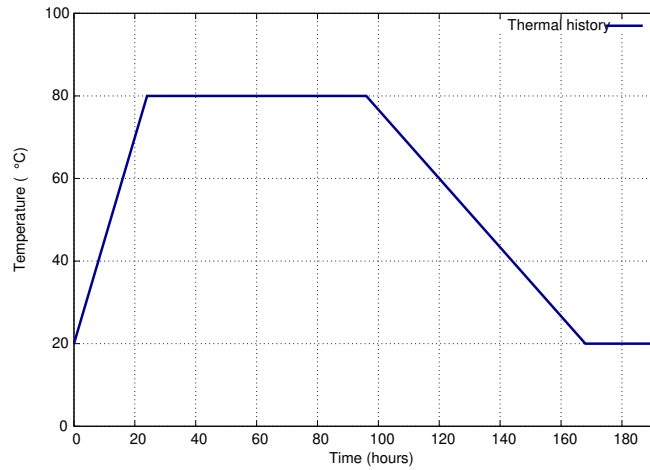


Figure 2: Heating treatment of the 11×22 cm cylindrical samples during first hours after curing.

218 *5.1. Simulations for the immersed samples*

219 The experimental evolution of the mean expansions for the samples is
220 presented in figure 3, in square-dots-points. The phase of accelerated degra-
221 dation is finishing and the stabilisation phase is taking place (about 0.23%
222 maximum expansion).

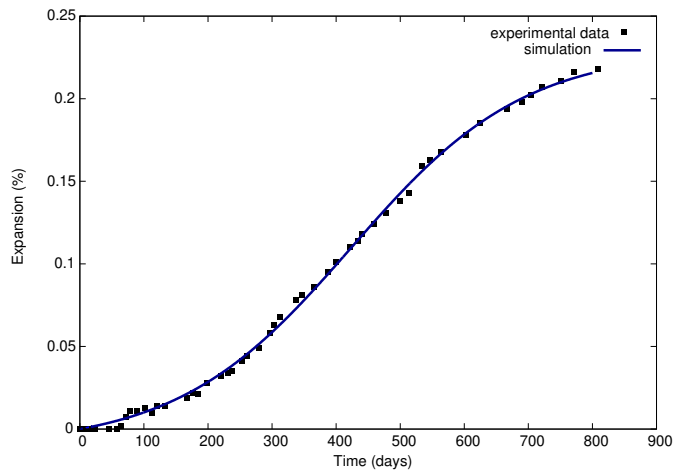


Figure 3: Swelling of concrete samples with calcareous aggregates. Experimental data and numerical simulation.

223 Curve fitting with the free-expansion law (4) gives the parameters sum-
224 marized in Table 3 where Larive’s law is sufficient, i.e. with $\phi = 0$.

225 With these values, the simulated expansion is superposed in figure 3 where
226 a good agreement is to be noticed. As the samples are immersed during
227 the whole time, uniform saturation $S_r = 1$ is taken and, consequently, the
228 effective time is identical to the real time.

Table 3: Material parameters with the expansion law (4) for the calcareous-based concrete.

| Definition | identified values |
|--|-------------------|
| expansion amplitude ε_{∞}^0 | 0.23% |
| latent time τ_l | 423.5 days |
| characteristic time τ_c | 135.48 days |
| parameter ϕ | 0 |
| parameter δ | 0 |

229 *5.2. Simulations of the effect of soaking-drying cycles on DEF*

230 A second set of samples with the above calcareous aggregates was submit-
 231 ted to cycles of drying (at 50 % of relative humidity and 23°C) alternated
 232 with soaking (at 23°C) to represent a more realistic environment of some
 233 structures which are not permanently in contact with water. The cycles are
 234 recorded in Table 4.

235 The swelling monitoring of this new study is plotted in figure 4 (the
 236 dashed line with circle dots), and compared to the previous experimental
 237 curve of figure 3 with the samples kept continuously immersed in water
 238 (square-dots-points curve). One can notice that the drying cycles restrict
 239 the delayed ettringite formation in the material and have an important im-
 240 pact on the expansion kinetics.

241 From the modelling point of view, this falls within the scope of the prop-
 242 erties (7)₁ for drying, and (7)₃ for soaking. For the material parameters of
 243 the expansion law (4), we use the ones of Table 3. And for the effective time
 244 (9), we fix the saturation threshold to $\bar{S}_r = 0.95$. Here the computation is

Table 4: Soaking-drying cycles for the calcareous-based concrete samples.

| Cycle | Soaking time | Drying time |
|-----------|--------------|-------------|
| n° | (days) | (days) |
| 1 | 9 | 14 |
| 2 | 22 | 20 |
| 3 | 7 | 7 |
| 4 | 10 | 12 |
| 5 | 20 | 12 |
| 6 | 35 | 9 |
| 7 | 22 | 22 |
| 8 | 10 | 7 |
| 9 | 23 | 19 |
| 10 | 16 | 7 |
| 11 | 34 | 10 |
| 12 | 19 | 21 |
| 13 | 13 | 22 |
| 14 | 12 | 7 |
| 15 | 37 | 22 |
| 16 | 13 | 21 |
| 17 | 13 | 17 |
| 18 | 38 | 23 |
| 19 | 41 | 24 |
| 20 | 14 | 17 |
| 21 | 30 | 20 |
| 22 | 38 | 38 |
| 23 | 42 | 24 |

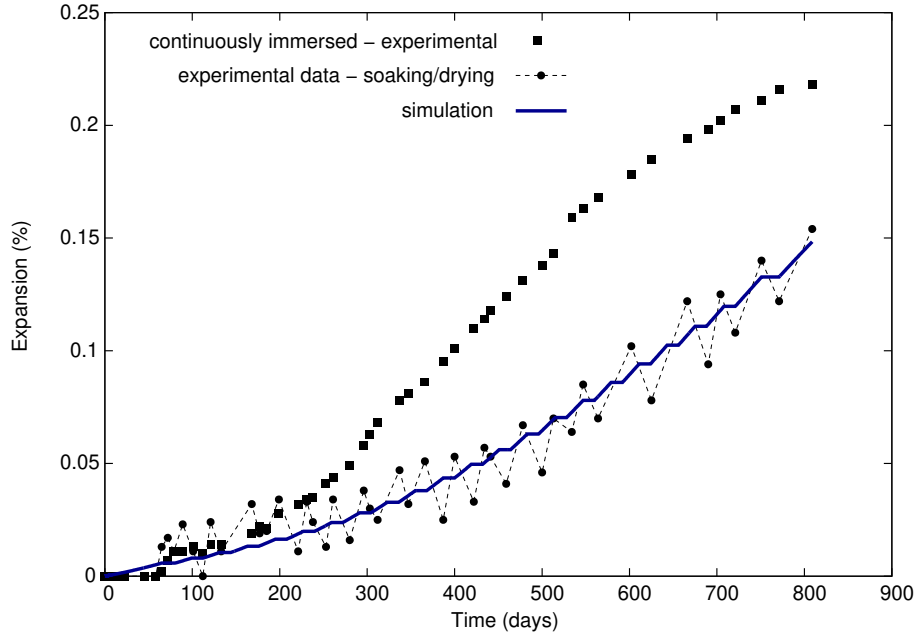


Figure 4: Impact of drying cycles on the swelling of concrete samples with calcareous aggregates. Experimental data and numerical simulation.

245 independent of the exponent parameter \bar{m} since:

$$\tilde{t} = \begin{cases} 1 & \text{for soaking, with } S_r = 1, \\ 0 & \text{for drying, since } S_r \ll \bar{S}_r. \end{cases} \quad (17)$$

246 In the numerical computation, uniform saturation $S_r = 1$ is taken during
 247 the first 72 days, then cycles of drying for 13 days at $S_r = 0.48$ followed by
 248 soaking for 19 days at $S_r = 1$ are prescribed. In total, the sample is immersed
 249 514 days and dried 295 days.

250 The simulated expansion is superposed in figure 4 where good agreement
 251 is again observed. Notice further that the expansion after 809 days of alter-
 252 nate soaking/drying cycles is almost the same as the 514 days' expansion in
 253 continuous immersion conditions.

254 **6. Conclusions**

255 In this study, we have presented and numerically tested an efficient tool
256 for the modelling of the influence of relative humidity on the development
257 of DEF expansions. The simple notion of effective time implicitly takes into
258 account the amount of water during time in variable environmental conditions
259 that strongly influence the kinetics of DEF expansions.

260 From the numerical point of view, we have briefly presented a numerical
261 design for the integration of the chemical strain and the update of the effective
262 time. In particular, this latter must be performed at the integration points
263 level when use is made of the finite element method.

264 Numerical simulations have shown good agreements with the experimen-
265 tal results issued from an experimental compaign, and others from the lit-
266 erature. However, further studies are ongoing **to include this notion within**
267 **more detailed and complex formulations, i.e. including creep, plasticity and**
268 **mechanical damage among others.**

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