



HAL
open science

FROM BOLTZMANN TO INCOMPRESSIBLE NAVIER-STOKES IN SOBOLEV SPACES WITH POLYNOMIAL WEIGHT

Marc Briant, S Merino-Aceituno, C. Mouhot

► **To cite this version:**

Marc Briant, S Merino-Aceituno, C. Mouhot. FROM BOLTZMANN TO INCOMPRESSIBLE NAVIER-STOKES IN SOBOLEV SPACES WITH POLYNOMIAL WEIGHT. *Analysis and Applications*, 2018, 10.1142/S021953051850015X . hal-02307954

HAL Id: hal-02307954

<https://hal.science/hal-02307954v1>

Submitted on 8 Oct 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

FROM BOLTZMANN TO INCOMPRESSIBLE NAVIER-STOKES IN SOBOLEV SPACES WITH POLYNOMIAL WEIGHT

M. BRIANT, S. MERINO-ACEITUNO, C. MOUHOT

ABSTRACT. We study the Boltzmann equation on the d -dimensional torus in a perturbative setting around a global equilibrium under the Navier-Stokes linearisation. We use a recent functional analysis breakthrough to prove that the linear part of the equation generates a C^0 -semigroup with exponential decay in Lebesgue and Sobolev spaces with polynomial weight, independently on the Knudsen number. Finally we show a Cauchy theory and an exponential decay for the perturbed Boltzmann equation, uniformly in the Knudsen number, in Sobolev spaces with polynomial weight. The polynomial weight is almost optimal and furthermore, this result only requires derivatives in the space variable and allows to connect to solutions to the incompressible Navier-Stokes equations in these spaces.

Keywords: Boltzmann equation on the Torus; Incompressible Navier-Stokes hydrodynamical limit; Cauchy theory uniform in Knudsen number, exponential rate of convergence towards global equilibrium.

Acknowledgements: This work was supported by the UK Engineering and Physical Sciences Research Council (EPSRC) grant EP/H023348/1 for the University of Cambridge Centre for Doctoral Training, the Cambridge Centre for Analysis.

CONTENTS

| | |
|---|----|
| 1. Introduction | 1 |
| 2. Main results | 6 |
| 3. The linear part: a C^0 -semigroup in spaces with polynomial weight, proof of Theorem 2.1 | 9 |
| 4. An <i>a priori</i> estimate for the full perturbed equation: proof of Theorem 2.3 | 20 |
| References | 28 |

1. INTRODUCTION

This article deals with the Boltzmann equation in a perturbative setting as the Knudsen number tends to zero. This equation rules the dynamics of rarefied gas particles moving on the flat torus in dimension d , \mathbb{T}^d , when the only interactions taken into account are binary collisions. More precisely, the Boltzmann equation describes the time evolution of the distribution $f = f(t, x, v)$ of particles in position x and velocity v . A formal derivation of the Boltzmann equation from Newton's laws under the rarefied gas assumption can be found in [14], while [15] presents

Lanford's Theorem (see [25] and [16] for detailed proofs) which rigorously proves the derivation in short times.

We denote the Knudsen number by ε and the Boltzmann equation reads

$$\partial_t f + v \cdot \nabla_x f = \frac{1}{\varepsilon} Q(f, f), \text{ on } \mathbb{T}^d \times \mathbb{R}^d,$$

where Q is the Boltzmann collision operator given by

$$Q(f, f) = \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} B(|v - v_*|, \cos \theta) [f' f'_* - f f_*] dv_* d\sigma.$$

The Boltzmann kernel operator B encodes the physics of the collision process and f' , f_* , f'_* and f are the values taken by f at v' , v_* , v'_* and v respectively, where

$$\begin{cases} v' = \frac{v + v_*}{2} + \frac{|v - v_*|}{2} \sigma \\ v'_* = \frac{v + v_*}{2} - \frac{|v - v_*|}{2} \sigma \end{cases}, \text{ and } \cos \theta = \left\langle \frac{v - v_*}{|v - v_*|}, \sigma \right\rangle.$$

The Boltzmann collision operator comes from a symmetric bilinear operator $Q(g, h)$ defined by

$$Q(g, h) = \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} B(|v - v_*|, \cos \theta) [h' g'_* + h'_* g' - h g_* - h_* g] dv_* d\sigma.$$

It is well-known (see [14], [15] or [17] for example) that the global equilibria for the Boltzmann equation are the *Maxwellians*, which are gaussian density functions depending only on the v variable. Without loss of generality we consider only the case of normalized Maxwellians:

$$\mu(v) = \frac{1}{(2\pi)^{\frac{d}{2}}} e^{-\frac{|v|^2}{2}}.$$

In this paper we will assume that the Boltzmann collision kernel is of the following form

$$(1.1) \quad B(|v - v_*|, \cos \theta) = \Phi(|v - v_*|) b(\cos \theta),$$

with Φ and b positive functions. This hypothesis is satisfied for all physical model and is more convenient to work with but do not impede the generality of our results.

We also restrict ourselves to the case of *hard potential* ($\gamma > 0$) or *Maxwellian potential* ($\gamma = 0$), that is to say there is a constant $C_\Phi > 0$ such that

$$(1.2) \quad \Phi(z) = C_\Phi z^\gamma, \quad \gamma \in [0, 1],$$

with a strong form of Grad's *angular cutoff* (see [18]), expressed here by the fact that we assume b to be C^1 with the controls from above

$$(1.3) \quad \forall z \in [-1, 1], \quad b(z), b'(z) \leq C_b.$$

1.1. The problem and its motivations. The Knudsen number is the inverse of the average number of collisions for each particle per unit of time. Therefore, as reviewed in [34], one can expect a convergence, in some sense, from the Boltzmann model towards the acoustics and the fluids dynamics as the Knudsen number tends to 0. However, these different models describe physical phenomena that do not evolve at the same timescale and the right rescaling to approximate the incompressible Navier-Stokes equation (see [4][17][34][31]) is the following equation

$$(1.4) \quad \partial_t f_\varepsilon + \frac{1}{\varepsilon} v \cdot \nabla_x f_\varepsilon = \frac{1}{\varepsilon^2} Q(f_\varepsilon, f_\varepsilon), \text{ on } \mathbb{T}^d \times \mathbb{R}^d,$$

under the linearization $f_\varepsilon(t, x, v) = \mu(v) + \varepsilon h_\varepsilon(t, x, v)$. This leads to the perturbed Boltzmann equation

$$(1.5) \quad \partial_t h_\varepsilon + \frac{1}{\varepsilon} v \cdot \nabla_x h_\varepsilon = \frac{1}{\varepsilon^2} \mathcal{L}(h_\varepsilon) + \frac{1}{\varepsilon} Q(h_\varepsilon, h_\varepsilon),$$

where we defined

$$\mathcal{L}(h) = 2Q(\mu, h).$$

The hydrodynamical limit of the perturbed equation is the system of equations satisfied by the limit, as ε tends to 0, of the hydrodynamical fluctuations that are the following physical observables of h_ε :

$$\begin{aligned} \rho_\varepsilon(t, x) &= \int_{\mathbb{R}^d} h_\varepsilon(t, x, v) dv, \\ u_\varepsilon(t, x) &= \int_{\mathbb{R}^d} v h_\varepsilon(t, x, v) dv, \\ \theta_\varepsilon(t, x) &= \frac{1}{d} \int_{\mathbb{R}^d} (|v|^2 - d) h_\varepsilon(t, x, v) dv. \end{aligned}$$

Note that $(\rho_\varepsilon, u_\varepsilon, \theta_\varepsilon)$ are the linearised fluctuations of the mass, momentum and the thermal energy around the global equilibrium μ .

In our perturbative framework, previous studies [4][5][12] show that the hydrodynamical limits ρ , u and θ are the weak (in the Leray sense [26]) solutions of the linearized incompressible Navier-Stokes equations:

$$(1.6) \quad \begin{aligned} \partial_t u - \nu \Delta u + u \cdot \nabla u + \nabla p &= 0, \\ \nabla \cdot u &= 0, \\ \partial_t \theta - \kappa \Delta \theta + u \cdot \nabla \theta &= 0, \end{aligned}$$

where p is the pressure function and ν and κ are constants determined by L (see [4] or [17] Theorem 5). They also satisfy the Boussineq relation

$$(1.7) \quad \nabla(\rho + \theta) = 0.$$

The aim of the present work is to use a constructive method to obtain existence and exponential decay for solutions to the perturbed Boltzmann equation (1.4), uniformly in the Knudsen number. One will thus be allowed to extract a converging (at least weakly) subsequence of h_ε converging to the incompressible Navier-Stokes equations [5][4][12]. Such uniform results have been obtained on the torus in Sobolev spaces with exponential weight $H_{x,v}^s(\mu^{-1/2})$ in [23][12] and the present work improves

this strong weight to a polynomial weight without the need of derivatives in the velocity variable.

1.2. Existing results. The first part of our work is to prove that the linear part of the Boltzmann equation

$$\mathcal{G}_\varepsilon = \frac{1}{\varepsilon^2} \mathcal{L} - \frac{1}{\varepsilon} v \cdot \nabla_x$$

generates a strongly continuous semigroup with an exponential decay in Lebesgue and Sobolev spaces with polynomial weight, namely $1 + |v|^k$ for some k large enough.

It has been known for long that the linear Boltzmann operator \mathcal{L} is a self-adjoint non positive linear operator in the space $L_v^2(\mu^{-1/2})$. Moreover it has a spectral gap λ_0 . This has been proved in [13][18][19] with non constructive methods for hard potential with cutoff and in [7][8] in the Maxwellian case. These results were made constructive in [2][27] for more general collision operators. One can easily extend this spectral gap to Sobolev spaces $H_v^s(\mu^{-1/2})$ (see for instance [21] Section 4.1).

The next step is to see if the latter properties about \mathcal{L} in the velocity space can be transposed when one adds the skew-symmetric transport operator $-v \cdot \nabla_x$. The first results were obtained in [32] where \mathcal{G}_1 was proven to generate a strong continuous semigroup in $L_v^2 H_x^s(\mu^{-1/2})$ and in $L_v^\infty H_x^s(\mu^{-1/2}(1 + |v|)^k)$, for s and k large enough. Then [29] obtained constructively this result in $H_{x,v}^s(\mu^{-1/2})$ using hypocoercivity properties of the Boltzmann linear operator. Finally, a recent breakthrough proving abstract extension of semigroups [21] showed that \mathcal{G}_1 generates a C^0 -semigroup in all the Sobolev spaces of the form $W_v^{\alpha,q} W_x^{\beta,p}(m)$, for m being an exponential weight (including maxwellian density if $q = p = 2$) or a polynomial weight $(1 + |v|)^k$, as long as $\alpha \leq \beta$ and k is large enough depending on q (with $k > 2$ in the case $q = 1$).

The full Boltzmann equation perturbed around a global equilibrium $\mu(v)$ (1.5) has also been studied in the case $\varepsilon = 1$. The associated Cauchy problem has been worked on over the past fifty years, starting with Grad [20], and it has been studied in different spaces, such as $L_v^2 H_x^s(\mu^{-1/2})$ spaces [32] or $H_{x,v}^s(\mu^{-1/2}(1 + |v|)^k)$ [22][36]. The Cauchy theory was then extended to $H_{x,v}^s(\mu^{-1/2})$ where an exponential trend to equilibrium has also been obtained. This was obtained using hypocoercivity properties of the linear operator [29] or nonlinear estimates on fluid and microscopic parts of the equation [23]. Recently, [21] proved existence and uniqueness for (1.5) in more the general spaces $(W_v^{\alpha,1} \cap W_v^{\alpha,q}) W_x^{\beta,p}(1 + |v|)^k$ for $\alpha \leq \beta$ and β and k large enough with explicit thresholds. This result therefore gets rid of the exponential weight needed in the previous studies.

All the results presented above hold in the case of the torus. We refer the reader interested in the Cauchy problem, both for the torus and the whole space, to the review [33].

For physical purposes, these studies for $\varepsilon = 1$ are relevant since mere rescalings or changes of physical units changes (1.4) to the case where the Knudsen number equals 1. However, if one wants to study the hydrodynamical limits of the Boltzmann equation, one needs to obtain explicit dependencies on the Knudsen number. Using hypocoercivity methods [12] gave a constructive uniform approach on the semigroup generated by \mathcal{G}_ε in $H_{x,v}^s(\mu^{-1/2})$ and its exponential decay. The study of the full perturbed Boltzmann equation (1.5) taking into account the dependencies on the

Knudsen number has been obtained [23][12] in the same spaces $H_{x,v}^s(\mu^{-1/2})$, for s large enough. More precisely, for initial data sufficiently close to μ there exists a unique non-negative solution to (1.4) and it decays exponentially fast towards its equilibrium. The smallness assumption was proven to be independent of the Knudsen number as well as the rate of decay and the methods used in [12] are constructive.

1.3. Our contributions and strategy. The present work brings two major improvements.

In the spirit of [21], we first prove that \mathcal{G}_ε generates a strong continuous semigroup in Sobolev spaces $W_v^{\alpha,1}W_x^{\beta,p}(1+|v|)^k$ for $\alpha \leq \beta$ and β and k large enough with explicit thresholds. It is done by starting from existing results in $H_{x,v}^s(\mu^{-1/2})$ and then decomposing the linear operator \mathcal{G}_ε into a dissipative part and a regularising part that is then treated in more and more regular spaces up to the space where the semigroup properties have been derived in previous articles. We thus improve the existing result [12]. Our main contribution is an adapted version of the abstract extension theorem developed in [21] that takes into account the dependencies on the Knudsen number as well as a careful study of the dissipative and the regularising parts of the operator \mathcal{G}_ε .

The second contribution of this article is the solution to the Cauchy problem with exponential trend to equilibrium, independently on ε , in spaces

$$W_v^{\alpha,1}W_x^{\beta,1}(1+|v|^{2+0}) \quad \text{and} \quad W_v^{\alpha,1}H_x^\beta(1+|v|^{2+0}),$$

for β large enough and all $\alpha \leq \beta$. First, this result makes the recent study [21] uniform in the Knudsen number. Second, it improves the Cauchy theory developed uniformly in ε in [23][12] by dropping the exponential weight and the v -derivatives. Moreover, one can notice that the polynomial weight is almost the optimal one for the Boltzmann equation (conservation of mass and energy).

The main issue to obtain uniform results is that the bilinear operator $\varepsilon^{-1}Q$ cannot be treated as a mere perturbation that evolves under the flow of $S_{\mathcal{G}_\varepsilon}$, the semigroup generated by \mathcal{G}_ε , since the latter has an exponential decay of order $O(1)$ that is negligible compared to $O(\varepsilon^{-1})$ as ε tends to zero. We develop an analytic point of view about the extension theorem in [21] and include the bilinear term. More precisely, we decompose the perturbed equation (1.5) into a hierarchy of equations taking place in spaces that have more and more regularity up to $H_{x,v}^s(\mu^{-1/2})$ where estimates had been derived in [12]. At each step we use the dissipative part of the linear operator to control the remainder term $\varepsilon^{-1}Q$ whereas the regularising part is controlled in spaces with higher regularity.

1.4. Organization of the article. Section 2 first introduces the different notations and definitions we will use throughout the paper and then states the precise theorems we prove in this work. Section 2.2 deals with the semigroup generated by the full linear operator $\varepsilon^{-2}\mathcal{L}-\varepsilon^{-1}v\cdot\nabla_x$ whereas Section 2.3 is dedicated to the full Boltzmann equation.

The full linear part of the Boltzmann operator is proven to generate a strongly continuous semigroup in Lebesgue and Sobolev spaces with polynomial weight in Section 3.

We start with Section 3.1, a thorough description of our strategy and a version of the extension theorem of [21] that takes into account the dependencies in ε .

We show in this section that $\varepsilon^{-2}\mathcal{L} - \varepsilon^{-1}v \cdot \nabla_x$ can be decompose into a regularising operator in the velocity variable (Section 3.2) and a dissipative one (Section 3.3).

We then combine the last two properties to gain regularity both in space and velocity (Section 3.4) to finally prove the existence and exponential decay of the associated semigroup (Section 3.5).

The last section, Section 4, proves existence, uniqueness and exponential decay of solutions to the perturbed Boltzmann equation (1.5).

Section 4.1 gives a new point of view on the extension we used to generate the semigroup associated to $\varepsilon^{-2}\mathcal{L} - \varepsilon^{-1}v \cdot \nabla_x$ and how it can be used with the bilinear operator. This strategy is developed through Sections 4.2 and 4.3 and it leads to the proof of the exponential decay towards equilibrium in Section 4.4.

2. MAIN RESULTS

2.1. Notations. We gather here the notations we will use throughout this article.

Function spaces. We first define the following shorthand notation,

$$\langle \cdot \rangle = \sqrt{1 + |\cdot|^2}.$$

The convention we choose is to index the space by the name of the concerned variable so we have, for p in $[1, +\infty]$,

$$L^p_{[0,T]} = L^p([0, T]), \quad L^p_x = L^p(\mathbb{T}^d), \quad L^p_v = L^p(\mathbb{R}^d).$$

Let p and q be in $[1, +\infty)$, α and β in \mathbb{N} and $m : \mathbb{R}^d \rightarrow \mathbb{R}^+$ a strictly positive measurable function. For any multi-indexes $j = (j_1, \dots, j_d)$ and $l = (l_1, \dots, l_d)$ in \mathbb{N}^d we denote the $(j, l)^{th}$ partial derivative by

$$\partial_l^j = \partial_x^l \partial_v^j.$$

We define the space $W_v^{\alpha,q} W_x^{\beta,p}(m)$ by the norm

$$\|f\|_{W_v^{\alpha,q} W_x^{\beta,p}(m)} = \sum_{\substack{|j| \leq \alpha, |l| \leq \beta \\ |l| + |j| \leq \max(\alpha, \beta)}} \|(\partial_l^j f) m\|_{L_v^q L_x^p},$$

where we used the Lebesgue norm

$$\|g\|_{L_v^q L_x^p} = \left[\int_{\mathbb{R}^d} \left(\int_{\mathbb{T}^d} |f(x, v)|^p dx \right)^{q/p} dv \right]^{1/q}.$$

Linear Boltzmann operator. First we use a writing convention. The present work aims at extending results known in a small space E , namely $H_{x,v}^s(\mu^{-1/2})$ with s sufficiently large, into a larger space \mathcal{E} , namely Lebesgue and Sobolev spaces with polynomial weight. We will use curly letters for operators in \mathcal{E} and their non-curly equivalent to denote their restriction to E . For instance, we will denote

$$\mathcal{L}|_E = L.$$

The linear Boltzmann operator L has several properties we will use throughout this paper (see [14][15][35][21] for instance).

L is a closed self-adjoint operator in $L_v^2(\mu^{-1/2})$ with kernel

$$\text{Ker}(L) = \text{Span}\{\phi_0(v), \dots, \phi_{d+1}(v)\} \mu,$$

where $\phi_0(v) = 1$, for $i = 1, \dots, d$ we defined $\phi_i(v) = v_i$ and $\phi_{d+1} = (|v|^2 - d) / \sqrt{2d}$. The family $(\phi_i)_{0 \leq i \leq d+1}$ is an orthonormal basis of $\text{Ker}(L)$ in $L_v^2(\mu^{-1/2})$ and we denote π_L the orthogonal projection onto $\text{Ker}(L)$ in $L_v^2(\mu^{-1/2})$:

$$(2.1) \quad \pi_L(h) = \sum_{i=0}^{d+1} \left(\int_{\mathbb{R}^d} h(u) \phi_i(u) du \right) \phi_i(v) \mu(v),$$

and we define $\pi_L^\perp = \text{Id} - \pi_L$. We will also denote the full linear Boltzmann operator by

$$\mathcal{G}_\varepsilon = \frac{1}{\varepsilon^2} L - \frac{1}{\varepsilon} v \cdot \nabla_x.$$

For s in \mathbb{N} we will use the convention

$$(\mathcal{G}_\varepsilon)|_{H_{x,v}^s(\mu^{-1/2})} = G_\varepsilon.$$

It has been proven ([12] Proposition 3.1) that the kernel of G_ε does not depend on ε and that its generators in $L_{x,v}^2(\mu^{-1/2})$ are the same than the ones of $\text{Ker}(L)$. We therefore have that the orthogonal projection onto $\text{Ker}(G_\varepsilon)$ in $L_{x,v}^2(\mu^{-1/2})$ is given by

$$(2.2) \quad \Pi_G(h) = \Pi_{G_\varepsilon}(h) = \sum_{i=0}^{d+1} \left(\int_{\mathbb{T}^d \times \mathbb{R}^d} h(x, u) \phi_i(u) dx du \right) \phi_i(v) \mu(v),$$

and we define $\Pi_G^\perp = \text{Id} - \Pi_G$.

Note that for a function h in $L_{x,v}^2(\mu^{-1/2})$ we have that

$$\forall (x, v) \in \mathbb{T}^d \times \mathbb{R}^d, \quad \Pi_G(h)(x, v) = \int_{\mathbb{T}^d} \pi_L(h(x_*, \cdot))(v) dx_*.$$

2.2. Results about the full linear part. We first deal with \mathcal{G}_ε , the linear part of the perturbed Boltzmann operator. We prove that it generates a strongly continuous semigroup with an exponential decay in Lebesgue and Sobolev spaces with a weight $\langle v \rangle^k$ as long as k is large enough. The precise statement is the following.

Theorem 2.1. *Let B be a Boltzmann collision kernel satisfying (1.1)-(1.2)-(1.3). There exists $0 < \varepsilon_d \leq 1$ such that for all p, q in $[1, +\infty]$, all α, β in \mathbb{N} with $\alpha \leq \beta$ and all $k > k_q^*$, where*

$$(2.3) \quad k_q^* = \frac{3 + \sqrt{49 - 48/q}}{2} + \gamma \left(1 - \frac{1}{q} \right)$$

with γ defined in (1.2),

- (1) for all $0 < \varepsilon \leq \varepsilon_d$, $\mathcal{G}_\varepsilon = \varepsilon^{-2} \mathcal{L} - \varepsilon^{-1} v \cdot \nabla_x$ generates a C^0 -semigroup, $S_{\mathcal{G}_\varepsilon}(t)$, on $W_v^{\alpha, q} W_x^{\beta, p}(\langle v \rangle^k)$,

(2) for all $\tau > 0$, there exist $C_G(\tau), \lambda_0 > 0$, such that for all $0 < \varepsilon \leq \varepsilon_d$ and for all h_{in} in $W_v^{\alpha,q} W_x^{\beta,p} (\langle v \rangle^k)$, for all $t \geq \tau$

$$\|S_{G_\varepsilon}(t)(h_{in}) - \Pi_G(h_{in})\|_{W_v^{\alpha,q} W_x^{\beta,p} (\langle v \rangle^k)} \leq C_G(\tau) e^{-\lambda_0 t} \|h_{in} - \Pi_G(h_{in})\|_{W_v^{\alpha,q} W_x^{\beta,p} (\langle v \rangle^k)},$$

where Π_G is the spectral projector onto $\text{Ker}(G_\varepsilon)$ which is given, for all ε , by

$$(2.4) \quad \Pi_G(g) = \sum_{i=0}^{d+1} \left(\int_{\mathbb{T}^d \times \mathbb{R}^d} g \phi_i dx dv \right) \phi_i \mu.$$

The constants $\varepsilon_d, C_G(\tau)$ and λ_0 are constructive and only depends on $d, p, q, k, \alpha, \beta$ and the kernel of the Boltzmann operator.

We refer to [24] and [21] Section 2 for definitions and properties of spectral projectors.

Remark 2.2. We can make a couple of remarks about this theorem.

- (1) Of important note is the fact that the exponential decay has a strong dependence on $\tau > 0$. This is inherent to our uniform extension of the enlargement semigroup theory of [21]. However, such a problematic initial layer could not be used to control directly the bilinear operator. That is why, as explained in the introduction, we need a new analytic approach of the enlargement theorem to tackle the non-linearity.
- (2) It has been proven in [12] Section 3, that in $H_{x,v}^1(\mu^{-1/2})$, $\text{Ker}(G_\varepsilon)$ does not depend on ε if ε is positive and we therefore can define $\Pi_G = \Pi_{G_\varepsilon}$. During the proof of Theorem 2.1 we will show that $(\Pi_{G_\varepsilon})|_{H_{x,v}^s(\mu^{-1/2})} = \Pi_{G_\varepsilon}$ and thus Π_G is well-defined and is independent of ε and given by (2.2).
- (3) As noticed in [21], the rate of decay λ_0 can be taken equal to the spectral gap of $\mathcal{L}|_{H_{x,v}^s(\mu^{-1/2})}$ (see [12]), for s as large as wanted, when k is big enough (and we obtained a constructive threshold).
- (4) Finally, we emphasize that in the case $q = 1$, the result holds for all $k > 2$. This is almost the minimal regularity $L_v^2(1 + |v|^2)$ for the Boltzmann equation, that is solutions with bounded mass and energy.

2.3. Existence, uniqueness and trend to equilibrium. A physically relevant requirement for solutions to the Boltzmann equation are that their mass, momentum and energy are preserved with time. This is also an *a priori* property of the Boltzmann equation on the torus (see [35] Chapter 1 Section 2 for instance) which reads

$$\forall t \geq 0, \quad \int_{\mathbb{T}^d \times \mathbb{R}^d} \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} f_\varepsilon(t, x, v) dx dv = \int_{\mathbb{T}^d \times \mathbb{R}^d} \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} f_\varepsilon(0, x, v) dx dv.$$

If one expects trend to the equilibrium $\mu(v)$ for the solutions $f_\varepsilon = \mu + \varepsilon h_\varepsilon$ of the Boltzmann equation (1.4) then it must be that

$$\forall t \geq 0, \quad \int_{\mathbb{T}^d \times \mathbb{R}^d} \begin{pmatrix} 1 \\ v \\ |v|^2 \end{pmatrix} h_\varepsilon(t, x, v) dx dv = 0,$$

that is $\Pi_{\mathcal{G}_\varepsilon}(h_\varepsilon(t, \cdot, \cdot)) = 0$ for all t , which is a property that is indeed preserved along time for solution to the perturbed Boltzmann equation (1.5), see [12] for instance.

We hence state the following theorem answering the Cauchy problem and the exponential convergence towards the equilibrium μ .

Theorem 2.3. *Let B be a Boltzmann collision kernel satisfying (1.1)-(1.2)-(1.3) and let $p = 1$ or $p = 2$. For α, β and $k > 2$ in \mathbb{N} define*

$$\mathcal{E}^p = W_v^{\alpha,1} W_x^{\beta,p} (\langle v \rangle^k).$$

There exists $0 < \varepsilon_d \leq 1$ and β_0 in \mathbb{N} such that for all $\beta \geq \beta_0, \alpha \leq \beta$, for all $k > 2$, for any λ'_0 in $(0, \lambda_0)$ (λ_0 defined in Theorem 2.1) there exist $C_{\alpha,\beta}, \eta_{\alpha,\beta} > 0$ such that for any $0 < \varepsilon \leq \varepsilon_d$, for any distribution $0 \leq f_{in} = \mu + \varepsilon h_{in}$:

If

- (i) h_{in} is in $\text{Ker}(\mathcal{G}_\varepsilon)^\perp$ in \mathcal{E}^p ,
- (ii) $\|h_{in}\|_{\mathcal{E}^p} \leq \eta_{\alpha,\beta}$,

then there exists a unique global solution $f_\varepsilon = f_\varepsilon(t, x, v)$ to (1.4) in \mathcal{E}^p which, moreover, satisfies $f_\varepsilon = \mu + \varepsilon h_\varepsilon \geq 0$ with:

- h_ε belongs to $\text{Ker}(\mathcal{G}_\varepsilon)^\perp$ for all times,
-

$$\|h_\varepsilon\|_{\mathcal{E}^p} \leq C_{\alpha,\beta} \|h_{in}\|_{\mathcal{E}^p} e^{-\lambda'_0 t}.$$

The constants $C_{\alpha,\beta}$ and $\eta_{\alpha,\beta}$ are constructive and depends only on $\alpha, \beta, k, d, \lambda'_0$ and the kernel of the Boltzmann operator.

Remark 2.4. *Note that the uniform control in \mathcal{E}^p gives a weak convergence of the of the physical observables of $h_\varepsilon(t, x, v)$, that is $\rho_\varepsilon(t, x)$, $u_\varepsilon(t, x)$ and $\theta_\varepsilon(t, x)$. Thanks to the convergence in distribution study made by Bardos, Golse and Levermore [4][3] we know that the weak limit of $(\rho_\varepsilon, u_\varepsilon, \theta_\varepsilon)$ is a solution to (1.6) with (1.7) close to the equilibrium $(1, 0, 1)$.*

3. THE LINEAR PART: A C^0 -SEMGROUP IN SPACES WITH POLYNOMIAL WEIGHT, PROOF OF THEOREM 2.1

In this section we focus on the linear part of the perturbed Boltzmann equation in $W_v^{\alpha,q} W_x^{\beta,p} (\langle v \rangle^k)$. We thus consider the following equation:

$$(3.1) \quad \partial_t h = \mathcal{G}_\varepsilon(h).$$

3.1. Strategy of the proof. If we denote $\mathcal{E} = W_v^{\alpha,q} W_x^{\beta,p} (\langle v \rangle^k)$ and $E = H_{x,v}^s (\mu^{-1/2})$ we have that $E \subset \mathcal{E}$, dense with continuous embedding for s large enough. [12] Theorem 2.1 (with the norm of Theorem 2.4) states that $G_\varepsilon = (\mathcal{G}_\varepsilon)|_E$ generates a strongly continuous semigroup in E with exponential decay. Theorem 2.1 can therefore be understood as the possibility to extend properties of G_ε in a small space E to \mathcal{G}_ε in a larger space \mathcal{E} .

This issue of extending spectral gap properties as well as semigroup properties has been first tackled by Mouhot to obtain constructive rates of convergence to equilibrium for the homogeneous Boltzmann equation [28]. Recently, Gualdani,

Mischler and Mouhot [21] proposed a more abstract approach that allows to deal with the full linear operator. In their work, they proved that if some conditions on \mathcal{G}_ε and G_ε are satisfied then we can pass on some semigroup properties from E to \mathcal{E} . The main argument of the proof of Theorem 2.1 is to show that we can use their result in our setting, independently of ε .

To be more precise, we give below a modified version of their main functional analysis theorem which is combination of Theorem 2.13 and Lemma 2.17 where we added dependencies on ε .

We refer to [21] Section 2 for the definition of hypodissipativity (roughly speaking it is a dissipative property in a different norm on a Banach space) and the definition of the convolution of two semigroups of operators (denoted by the symbol $(*)$). In the sequel we will use $\mathcal{C}(E)$ for the set of closed operators on E and $\mathcal{B}(E)$ for the set of bounded operators on E . For any operator G in $\mathcal{C}(E)$ we denote $R(G)$ its range and $\Sigma(G)$ its spectrum.

Theorem 3.1 (Modified extension theorem from [21]). *Let ε be a parameter such that $0 < \varepsilon \leq 1$.*

Let E, \mathcal{E} be two Banach spaces with $E \subset \mathcal{E}$ dense with continuous embedding, and consider G_ε in $\mathcal{C}(E)$, \mathcal{G}_ε in $\mathcal{C}(\mathcal{E})$ with $(\mathcal{G}_\varepsilon)|_E = G_\varepsilon$ and $a > 0$.

We assume the following

(A1) G_ε generates a semigroup S_{G_ε} on E , $G_\varepsilon + a$ is hypodissipative on $R(\text{Id} - \Pi_{G_\varepsilon, a})$ and

$$\Sigma(G_\varepsilon) \cap \{z \in \mathbb{C}, \text{Re}(z) > -a\} = \{0\} \quad \text{where } 0 \text{ is a discrete eigenvalue.}$$

(A2) *There exists $\mathcal{A}_\varepsilon, \mathcal{B}_\varepsilon$ in $\mathcal{C}(\mathcal{E})$ such that $\mathcal{G}_\varepsilon = \mathcal{A}_\varepsilon + \mathcal{B}_\varepsilon$ (with corresponding restrictions $A_\varepsilon, B_\varepsilon$ on E) and there exist some “intermediate spaces” (not necessarily ordered)*

$$E = \mathcal{E}_J, \mathcal{E}_{J-1}, \dots, \mathcal{E}_2, \mathcal{E}_1 = \mathcal{E}$$

such that, still denoting $\mathcal{B}_\varepsilon := (\mathcal{B}_\varepsilon)|_{\mathcal{E}_j}$ and $\mathcal{A}_\varepsilon := (\mathcal{A}_\varepsilon)|_{\mathcal{E}_j}$

- (i) *For any j in $\{1, \dots, J\}$, $(\mathcal{B}_\varepsilon + a/\varepsilon^2)$ is hypodissipative on \mathcal{E}_j ;*
- (ii) *For any j in $\{1, \dots, J\}$, $\mathcal{A}_\varepsilon \in \mathcal{B}(\mathcal{E}_j)$ with $\|\mathcal{A}_\varepsilon\|_{\mathcal{B}(\mathcal{E}_j)} \leq C_A/\varepsilon^2$;*
- (iii) *there are some constants $l_0, l_1 \in \mathbb{N}^*$, $C \geq 1$, $K \in \mathbb{R}$, $\alpha \in [0, 1)$ such that*

$$\forall t \geq 0, \quad \|T_{l_0}(t)\|_{\mathcal{B}(\mathcal{E}_j, \mathcal{E}_{j+1})} \leq C \frac{e^{Kt/\varepsilon^2}}{\varepsilon^{l_1 t^\alpha}},$$

*for $1 \leq j \leq J-1$, with the notation $T_l := (\mathcal{A}_\varepsilon S_{\mathcal{B}_\varepsilon})^{(*l)}$.*

Then \mathcal{G}_ε is hypodissipative in \mathcal{E} and for all $a' < a$ there exists $n = n(a') \geq 1$ and some positive constants $C_{a'}$ and $C'_{a'}$ such that

$$(3.2) \quad \|T_n(t)\|_{\mathcal{B}(\mathcal{E})} \leq \frac{C_{a'}}{\varepsilon^{nl_1/l_0}} e^{-a't/\varepsilon^2};$$

(3.3)

$$S_{\mathcal{G}_\varepsilon}(t) = S_{G_\varepsilon}(t)\Pi_G + \sum_{l=0}^{n-1} (-1)^l (\text{Id} - \Pi_G) S_{\mathcal{B}_\varepsilon} * T_l(t) + (-1)^n [(\text{Id} - \Pi_G) S_{G_\varepsilon}] * T_n(t);$$

$$(3.4) \quad \left\| S_{\mathcal{G}_\varepsilon}(t) - S_{G_\varepsilon}(t)\Pi_G \right\|_{\mathcal{B}(\mathcal{E})} \leq C'_{a'} \frac{t^n}{\varepsilon^{n(2+l_1/l_0)}} e^{-a't/\varepsilon^2},$$

where $\Pi_{\mathcal{G}}$ has been defined in (2.4).

We will use Theorem 3.1 to directly prove Theorem 2.1. Indeed, [12] Theorem 2.1 states that G_ε generates a strongly continuous semigroup with exponential decay at rate a in $E = H_{x,v}^s(\mu^{-1/2})$, which is the required assumption (A1) (properties about the spectral gap of the spectrum can be found in [2]). Therefore if \mathcal{G}_ε fulfils hypothesis (A2) then it generates a strongly continuous semigroup, with an exponential decay of order λ_0 for all $0 < \lambda_0 < a$. Indeed, take a' in (λ_0, a) in (3.4) and fix $\tau > 0$ and see that

$$(3.5) \quad \forall t \geq \tau, \quad \frac{t^n}{\varepsilon^{n(2+l_1/l_0)}} e^{-a't/\varepsilon^2} \leq \frac{t^n}{\varepsilon^{n(2+l_1/l_0)}} e^{-(a'-\lambda_0)t/\varepsilon^2} e^{-\lambda_0 t/\varepsilon^2} \leq C_{n,l_1/l_0,\tau,a',\lambda_0} e^{-\lambda_0 t},$$

which is the expected result given in Theorem 2.1.

3.2. Decomposition of the operator and assumption (A2)(ii). In this section we find a decomposition $\mathcal{G}_\varepsilon = \mathcal{A}_\varepsilon + \mathcal{B}_\varepsilon$ that will fit the requirements (A1) – (A2) of Theorem 3.1. This decomposition has been found in [21] in the case $\varepsilon = 1$. We will use exactly the same operators but including the dependencies in ε . All the results presented in the rest of this section are true for $\varepsilon = 1$ (see [21] Section 4) so we will try to relate as much as possible our computations with the ones for $\varepsilon = 1$.

For δ in $(0, 1)$, to be chosen later, we consider $\Theta_\delta = \Theta_\delta(v, v_*, \sigma)$ in C^∞ that is bounded by one on the set

$$\{|v| \leq \delta^{-1} \quad \text{and} \quad 2\delta \leq |v - v_*| \leq \delta^{-1} \quad \text{and} \quad |\cos \theta| \leq 1 - 2\delta\}$$

and whose support is included in

$$\{|v| \leq 2\delta^{-1} \quad \text{and} \quad \delta \leq |v - v_*| \leq 2\delta^{-1} \quad \text{and} \quad |\cos \theta| \leq 1 - \delta\}.$$

We define the splitting

$$\mathcal{G}_\varepsilon = \mathcal{A}_\varepsilon^{(\delta)} + \mathcal{B}_\varepsilon^{(\delta)},$$

with

$$\mathcal{A}_\varepsilon^{(\delta)} h(v) = \frac{1}{\varepsilon^2} \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} \Theta_\delta [\mu'_* h' + \mu' h'_* - \mu h_*] b(\cos \theta) |v - v_*|^\gamma d\sigma dv_*$$

and

$$\mathcal{B}_\varepsilon^{(\delta)} h(v) = \mathcal{B}_{2,\varepsilon}^{(\delta)} h(v) - \frac{1}{\varepsilon^2} \nu(v) h(v) - \frac{1}{\varepsilon} v \cdot \nabla_x h(v),$$

where

$$\mathcal{B}_{2,\varepsilon}^{(\delta)} h(v) = \frac{1}{\varepsilon^2} \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} (1 - \Theta_\delta) [\mu'_* h' + \mu' h'_* - \mu h_*] b(\cos \theta) |v - v_*|^\gamma d\sigma dv_*$$

and $\nu(v)$ is the standard collision frequency

$$\nu(v) = \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} b(\cos \theta) |v - v_*|^\gamma \mu_* d\sigma dv_*.$$

Note that there exists $\nu_0, \nu_1 > 0$ such that

$$(3.6) \quad \forall v \in \mathbb{R}^d, \quad \nu_0(1 + |v|^\gamma) \leq \nu(v) \leq \nu_1(1 + |v|^\gamma).$$

We have that

$$\mathcal{A}_\varepsilon^{(\delta)} = \frac{1}{\varepsilon^2} \mathcal{A}_1^{(\delta)} \quad \text{and} \quad \mathcal{B}_{2,\varepsilon}^{(\delta)} = \frac{1}{\varepsilon^2} \mathcal{B}_{2,1}^{(\delta)}.$$

We therefore obtain the following controls on $\mathcal{A}_\varepsilon^{(\delta)}$.

Proposition 3.2. *For all $0 < \varepsilon < \varepsilon_d$, for any q in $[1, +\infty]$ and $\alpha \geq 0$, the operator $\mathcal{A}_\varepsilon^{(\delta)}$ maps L_v^q into $W_v^{\alpha,q}$ with compact support. There exists $C_{\delta,\alpha,q}, R_\delta > 0$ independent of ε such that*

$$\forall h \in L_v^q, \text{supp}(\mathcal{A}_\varepsilon^{(\delta)}h) \subset B(0, R_\delta), \quad \|\mathcal{A}_\varepsilon^{(\delta)}h\|_{W_v^{\alpha,q}} \leq \frac{C_{\delta,\alpha,q}}{\varepsilon^2} \|h\|_{L_v^q}.$$

Moreover, for any p in $[1, +\infty]$ and for all h in $L_v^q L_x^p$,

$$\|\mathcal{A}_\varepsilon^{(\delta)}h\|_{L_v^q L_x^p} \leq \|\mathcal{A}_\varepsilon^{(\delta)}(\|h\|_{L_x^p})\|_{L_v^q}$$

Remark 3.3. *We notice here that this Proposition gives the point (A2)(ii) of Theorem 3.1 if the \mathcal{E}_j are Sobolev spaces.*

Proof of Proposition 3.2. The kernel of the operator $\mathcal{A}_\varepsilon^{(\delta)}$ is of compact support so its Carleman representation (see [13]) gives the existence of $k^{(\delta)}$ in $C_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$ such that

$$(3.7) \quad \mathcal{A}_\varepsilon^{(\delta)}h(v) = \frac{1}{\varepsilon^2} \int_{\mathbb{R}^d} k^{(\delta)}(v, v_*) h(v_*) dv_*,$$

and therefore the control on $\|\mathcal{A}_\varepsilon^{(\delta)}h\|_{W_v^{\alpha,q}}$ is straightforward.

The control of $\|\mathcal{A}_\varepsilon^{(\delta)}h\|_{L_v^q L_x^p}$ comes directly from Minkowski's integral inequality which states

$$\left[\int_{\mathbb{T}^d} \left(\int_{\mathbb{R}^d} k^{(\delta)}(v, v_*) h(x, v_*) dv_* \right)^p dx \right]^{1/p} \leq \int_{\mathbb{R}^d} \left(\int_{\mathbb{T}^d} k^{(\delta)}(v, v_*)^p h(x, v_*)^p dx \right)^{1/p} dv_*.$$

□

3.3. Dissipativity estimates for $\mathcal{B}_\varepsilon^{(\delta)}$, assumption (A2)(i). One can find in [21] proof of Lemma 4.14 case (W2) and (W3) the following estimate on the operator $\mathcal{B}_\varepsilon^{(\delta)}$ in the case $\varepsilon = 1$.

Lemma 3.4. *For all p, q in $[1, +\infty]$, for all $k > 2$ and for any δ in $(0, 1)$ and all h in $L_v^q L_x^p(\langle v \rangle^k)$,*

$$\int_{\mathbb{R}^d} \langle v \rangle^{kq} \|h\|_{L_x^p}^{q-p} \left(\int_{\mathbb{T}^d} \text{sgn}(h) |h|^{p-1} \mathcal{B}_1^{(\delta)} h dx \right) dv \leq [\Lambda_{k-\gamma/q',q}(\delta) - 1] \|h\|_{L_v^q L_x^p(\langle v \rangle^{k\nu^{1/q}})}^q,$$

where q' is the conjugate exponent of q and $\Lambda_{k,q}(\delta)$ is a constructive constant such that

$$\lim_{\delta \rightarrow 0} \Lambda_{k,q}(\delta) = \phi_q(k) = \left(\frac{4}{k+2} \right)^{1/q} \left(\frac{4}{k-1} \right)^{1-1/q}.$$

Remark 3.5. As noticed in [21] Remark 4.3, the quantity $\phi_q(k)$ is strictly less than one for k bigger than a constant k_q^{**} . The constant k_q^* we are considering is not optimal and is such that $\phi_q(k - \gamma/q') < 1$, where q' is the conjugate exponent of q . This appearance of $k - \gamma/q'$ is due to a loss of weight of order $\nu^{-1/q'}$ in the estimate of the spectral gap, see proof of Proposition 3.6.

In the case of the Boltzmann operator with hard potential and angular cutoff, point (A2)(i) is fulfilled by $\mathcal{B}_\varepsilon^{(\delta)}$ for δ small enough. This is the purpose of the following lemma. We recall here that $\nu_0 = \inf_{v \in \mathbb{R}^d} (\nu(v)) > 0$ and that we define

$$\|\cdot\|_{W_v^{\alpha,q} W_x^{\beta,p}(\langle v \rangle^k)} = \sum_{\substack{|l|+|j| \leq \max(\alpha,\beta) \\ |j| \leq \alpha, |l| \leq \beta}} \|\partial_l^j \cdot\|_{L_v^q L_x^p(\langle v \rangle^k)}.$$

Proposition 3.6. Consider p, q in $[1, +\infty]$, $k > k_q^*$, defined by (2.3), and α, β in \mathbb{N} such that $\alpha \leq \beta$.

Then there exists $\delta_{k,q}$ in $(0, 1)$ such that for all $0 < \delta \leq \delta_{k,q}$ there exists $\lambda_0 = \lambda_0(k, q, \delta)$ in $(0, \nu_0)$ such that for all $0 < \varepsilon \leq 1$,

- $\lambda_0(k, q, \delta)$ tends to $\lambda_0^*(k, q)$ as δ goes to 0,
- $\lambda_0^*(k, q)$ tends to ν_0 when k goes to $+\infty$,
- $(\mathcal{B}_\varepsilon^{(\delta)} + \lambda_0/\varepsilon^2)$ is hypodissipative in $W_v^{\alpha,q} W_x^{\beta,p}(\langle v \rangle^k)$.

Proof of Proposition 3.6. Let h_0 be in $W_v^{\alpha,q} W_x^{\beta,p}(\langle v \rangle^k)$ and consider h to be a solution to the linear equation

$$(3.8) \quad \partial_t h = \mathcal{B}_\varepsilon^{(\delta)} h = \mathcal{B}_{2,\varepsilon}^{(\delta)} h - \frac{1}{\varepsilon^2} \nu h - \frac{1}{\varepsilon} v \cdot \nabla_x h,$$

with initial value h_0 .

Since the x -derivative commutes with the equation we can consider only the case when $\beta = \alpha$. The proof is split into two parts. First we prove Proposition 3.6 in the case $\alpha = 0$ and then we study the case with v -derivatives.

Step 1: the case $\alpha = 0$. Take p, q in $[1, +\infty)$.

We recall that

$$\|h\|_{L_v^q L_x^p(\langle v \rangle^k)} = \left[\int_{\mathbb{R}^d} (1 + |v|^k)^q \left(\int_{\mathbb{T}^d} |h|^p dx \right)^{q/p} dv \right]^{1/q}.$$

Therefore we can compute

$$(3.9) \quad \begin{aligned} \frac{d}{dt} \|h\|_{L_v^q L_x^p(\langle v \rangle^k)} &= \|h\|_{L_v^q L_x^p(\langle v \rangle^k)}^{1-q} \\ &\times \int_{\mathbb{R}^d} (1 + |v|^k)^q \|h\|_{L_x^p}^{q-p} \left(\int_{\mathbb{T}^d} \operatorname{sgn}(h) |h|^{p-1} \mathcal{B}_\varepsilon^{(\delta)} h dx \right) dv. \end{aligned}$$

Observing that

$$\int_{\mathbb{T}^d} \operatorname{sgn}(h) |h|^{p-1} v \cdot \nabla_x h dx = \frac{1}{p} v \cdot \int_{\mathbb{T}^d} \nabla_x (|h|^p) dx = 0,$$

we deduce

$$\begin{aligned} \frac{d}{dt} \|h\|_{L_v^q L_x^p(\langle v \rangle^k)} &= \|h\|_{L_v^q L_x^p(\langle v \rangle^k)}^{1-q} \\ &\quad \times \frac{1}{\varepsilon^2} \int_{\mathbb{R}^d} \left(1 + |v|^k\right)^q \|h\|_{L_x^p}^{q-p} \left(\int_{\mathbb{T}^d} \operatorname{sgn}(h) |h|^{p-1} \mathcal{B}_1^{(\delta)} h \, dx \right) dv. \end{aligned}$$

We can therefore use Lemma 3.4 which leads to

$$(3.10) \quad \frac{d}{dt} \|h\|_{L_v^q L_x^p(\langle v \rangle^k)} \leq -\frac{1}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})}^q \|h\|_{L_v^q L_x^p(\langle v \rangle^k)}^{1-q},$$

We already noticed that $\Lambda_{k-\gamma/q',q}(\delta)$ is strictly less than 1 for δ smaller than some $\delta_{k,q}$ (see Remark 3.5). Therefore, because $\nu(v) \geq \nu_0$ for all v we have that for all δ smaller than $\delta_{k,q}$ the following holds,

$$\frac{d}{dt} \|h\|_{L_v^q L_x^p(\langle v \rangle^k)} \leq -\frac{\nu_0}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \|h\|_{L_v^q L_x^p(\langle v \rangle^k)},$$

This concludes the proof of Proposition 3.6 for $\alpha = 0$ and $1 \leq p, q < +\infty$. The cases $p = \infty$ and $q = \infty$ are respectively dealt with by taking the limit $p \rightarrow \infty$ and $q \rightarrow \infty$ which is possible since $\delta_{k,q}$ is independent of p and can be chosen to converge to a strictly positive constant when q goes to ∞ , thanks to the definition of $\Lambda_{k,q}(\delta)$.

Step 2: the case with v -derivatives. Take p, q in $[1, +\infty]$ and $\alpha = \beta = 1$.

Since the x -derivative commutes with (3.8) the equation satisfied by h , we have that (3.10) holds for x -derivatives. Notice that $1 - q \leq 0$ gives

$$(3.11) \quad \begin{aligned} \frac{d}{dt} \left(\|h\|_{L_v^q L_x^p(\langle v \rangle^k)} + \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k)} \right) \\ \leq -\frac{\nu_0^{1-1/q}}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \left(\|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} + \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \right). \end{aligned}$$

Applying a v -derivatives to (3.8) yields

$$\begin{aligned} \partial_t \nabla_v h &= \mathcal{B}_\varepsilon^{(\delta)}(\nabla_v h) + (\nabla_v \mathcal{B}_\varepsilon^{(\delta)})(h) \\ &= \mathcal{B}_\varepsilon^{(\delta)}(\nabla_v h) - \frac{1}{\varepsilon} \nabla_x h + \mathcal{R}_\varepsilon^{(\delta)}(h), \end{aligned}$$

where $\mathcal{R}_\varepsilon^{(\delta)}(h) = (\nabla_v \mathcal{B}_{2,\varepsilon}^{(\delta)})(h) - \frac{1}{\varepsilon^2} \nabla_v(\nu)h = \frac{1}{\varepsilon^2} \mathcal{R}_1^{(\delta)}(h)$.

From (3.11), our computations in Step 1 with $\delta \leq \delta_{k,q}$ and the following norm

$$\|h\|_{W_v^{1,q} W_x^{1,p}(\langle v \rangle^k)_\eta} = \|h\|_{L_v^q L_x^p(\langle v \rangle^k)} + \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k)} + \eta \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k)},$$

with $\eta > 0$ to be fixed later, we obtain

$$\begin{aligned}
 & \frac{d}{dt} \|h\|_{W_v^{1,q} W_x^{1,p}(\langle v \rangle^k)_\eta} \\
 & \leq -\frac{\nu_0^{1-1/q}}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \left(\|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} + \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \right) \\
 & \quad - \eta \frac{\nu_0^{1-1/q}}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \\
 & \quad - \frac{\eta}{\varepsilon} \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k)}^{1-q} \int_{\mathbb{R}^d} (\langle v \rangle^k)^q \|\nabla_v h\|_{L_x^p}^{q-p} \left(\int_{\mathbb{T}^d} \text{sgn}(h) |\nabla_v h|^{p-1} \nabla_x h \, dx \right) dv \\
 & \quad + \frac{\eta}{\varepsilon^2} \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k)}^{1-q} \int_{\mathbb{R}^d} (\langle v \rangle^k)^q \|\nabla_v h\|_{L_x^p}^{q-p} \left(\int_{\mathbb{T}^d} \text{sgn}(h) |\nabla_v h|^{p-1} \mathcal{R}_1^{(\delta)}(h) \, dx \right) dv.
 \end{aligned}$$

We take the absolute value and use Hölder inequality twice on the last two terms which makes the terms in $\nabla_v h$ disappear, and this gives

$$\begin{aligned}
 & \frac{d}{dt} \|h\|_{W_v^{1,q} W_x^{1,p}(\langle v \rangle^k)_\eta} \\
 & \leq -\frac{\nu_0^{1-1/q}}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \left(\|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} + \eta \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \right) \\
 & \quad + \frac{1}{\varepsilon^2} \left(\varepsilon \eta \nu_0^{-1/q} - \nu_0^{1-1/q} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \right) \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \\
 & \quad + \frac{\eta}{\varepsilon^2} \left\| \mathcal{R}_1^{(\delta)}(h) \right\|_{L_v^q L_x^p(\langle v \rangle^k)}.
 \end{aligned}$$

One can find in [21] proof of Lemma 4.14 case (W2) and (W3) the following estimate

$$\left\| \mathcal{R}_1^{(\delta)}(h) \right\|_{L_v^q L_x^p(\langle v \rangle^k)} \leq C_\delta \|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})},$$

where $C_\delta > 0$ is a constant only depending on δ .

Because $\varepsilon \leq 1$, this latter estimates yields

$$\begin{aligned}
 (3.12) \quad & \frac{d}{dt} \|h\|_{W_v^{1,q} W_x^{1,p}(\langle v \rangle^k)_\eta} \leq \frac{1}{\varepsilon^2} \left(C_\delta \eta - \nu_0^{1-1/q} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \right) \|h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \\
 & \quad + \frac{1}{\varepsilon^2} \left(\eta \nu_0^{-1/q} - \nu_0^{1-1/q} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \right) \|\nabla_x h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})} \\
 & \quad - \eta \frac{\nu_0^{1-1/q}}{\varepsilon^2} [1 - \Lambda_{k-\gamma/q',q}(\delta)] \|\nabla_v h\|_{L_v^q L_x^p(\langle v \rangle^k \nu^{1/q})},
 \end{aligned}$$

which concludes the proof if we take η small enough in terms of δ , for $\delta \leq \delta_{k,q}$.

The case where $1 < \alpha = \beta$ is dealt with in the same way with the norm

$$\|h\|_{W_v^{\alpha,q} W_x^{\alpha,p}(\langle v \rangle^k)_\eta} = \sum_{0 \leq |j|+|l| \leq \alpha} \eta^{|j|} \|\partial_l^j h\|_{L_v^q L_x^p(\langle v \rangle^k)},$$

with η small enough in terms of δ . □

3.4. Estimates on the iterated convolution product, assumption (A2)(iii). In order to use Theorem 3.1, it remains to show that our equation (3.1) satisfies hypothesis (A2)(iii), that is we need to control the iterated quantities $T_l := \left(\mathcal{A}_\varepsilon^{(\delta)} S_{\mathcal{B}_\varepsilon^{(\delta)}}\right)^{(*l)}$ for some l in \mathbb{N} . The following proposition describes such controls when $p = 1$.

Proposition 3.7. *Consider $k > k_q^*$, defined by (2.3), and s in \mathbb{N} . For any δ in $(0, \delta_{k,q}]$ and any λ'_0 in $(0, \lambda_0)$ ($\delta_{k,q}$ and λ_0 defined in Proposition 3.6), there exists $C_1 = C_1(\lambda'_0, \delta) > 0$ and $R = R(\delta) > 0$ such that for any $t \geq 0$,*

$$\forall n \in \mathbb{N}, \quad \text{supp } T_n(t)h \subset K := B(0, R)$$

and

$$(3.13) \quad \forall s \geq 1, \quad \|T_1(t)h\|_{W_{x,v}^{s+1,1}(K)} \leq C_1 \frac{e^{-\frac{\lambda'_0}{\varepsilon^2}t}}{\varepsilon^2 t} \|h\|_{W_v^{s+1,1} W_x^s(\langle v \rangle^k)},$$

$$(3.14) \quad \forall s \geq 0, \quad \|T_2(t)h\|_{W_{x,v}^{s+1/2,1}(K)} \leq C_1 \frac{e^{-\frac{\lambda'_0}{\varepsilon^2}t}}{\varepsilon^4} \|h\|_{W_{x,v}^{s,1}(\langle v \rangle^k)}.$$

Proof of Proposition 3.7. Most of the proof is an adaptation of [21] proof of Lemma 4.19 to keep track of the dependencies on ε . We will refer to it when we are using some of its computations.

Control of $T_1(t)h$: The x -derivatives commutes with $T_1(t)$ and therefore it is enough to consider h in $W_v^{s,1} W_x^{1,1}(\langle v \rangle^k)$, with $s \geq 1$, and to control $\|T_1(t)h\|_{W_v^{s+1,1} W_x^{1,1}(K)}$. This yields

$$(3.15) \quad \|T_1(t)h\|_{W_v^{s+1,1} W_x^{1,1}(K)} \leq \|T_1(t)h\|_{W_v^{s+1,1} L_x^1(K)} + \|\nabla_x T_1(t)h\|_{W_v^{s+1,1} L_x^1(K)}.$$

The first term is easily dealt with thanks to the estimate on $\mathcal{A}_\varepsilon^{(\delta)}$, Proposition 3.2, and the dissipativity property of $\mathcal{B}_\varepsilon^{(\delta)}$, Proposition 3.6,

$$(3.16) \quad \|T_1(t)h\|_{W_v^{s+1,1} L_x^1(K)} = \left\| \mathcal{A}_\varepsilon^{(\delta)} S_{\mathcal{B}_\varepsilon^{(\delta)}} h \right\|_{W_v^{s+1,1} L_x^1(K)} \leq \frac{C}{\varepsilon^2} e^{-\frac{\lambda_0}{\varepsilon^2}t} \|h\|_{L_v^1 L_x^1(\langle v \rangle^k)}.$$

For the second term, define $f(t) = S_{\mathcal{B}_\varepsilon^{(\delta)}}(t)h$ and

$$(3.17) \quad D_t = \varepsilon^{-1}t \nabla_x + \nabla_v.$$

By direct computations we have that

$$\varepsilon^{-1}t \nabla_x T_1(t)h = \mathcal{A}_\varepsilon^{(\delta)}(D_t f) - (\nabla_{v_*} \mathcal{A}_\varepsilon^{(\delta)}) f.$$

Above we denote $\nabla_{v_*} \mathcal{A}_\varepsilon^{(\delta)}$ to be the operator $\mathcal{A}_\varepsilon^{(\delta)}$ where the kernel $k(v, v_*)$ is swapped with $-\nabla_{v_*} k(v, v_*)$, see the proof of Proposition 3.2. By Proposition 3.2 we get

$$(3.18) \quad \varepsilon^{-1}t \|\nabla_x T_1(t)h\|_{W_v^{s+1,1} L_x^1(K)} \leq \frac{C}{\varepsilon^2} \left[\|D_t f\|_{L_{x,v}^1(\langle v \rangle^k)} + \|f\|_{L_{x,v}^1(\langle v \rangle^k)} \right].$$

The dissipativity property of $\mathcal{B}_\varepsilon^{(\delta)}$, in particular (3.10) with $q = 1$, yields

$$(3.19) \quad \frac{d}{dt} \|f\|_{L_{x,v}^1(\langle v \rangle^k)} \leq -\frac{\lambda_0}{\varepsilon^2} \|f\|_{L_{x,v}^1(\langle v \rangle^k \nu)}.$$

Direct computations yields

$$\partial_t (D_t f) = \mathcal{B}_\varepsilon^{(\delta)} (D_t f) + \frac{1}{\varepsilon^2} \mathcal{J}^{(\delta)} f,$$

where

$$(3.20) \quad \mathcal{J}^{(\delta)} = \nabla_v (\mathcal{B}^{(\delta)}(\cdot)) - \mathcal{B}^{(\delta)}(\nabla_v(\cdot)) \quad \text{with} \quad \mathcal{B}^{(\delta)} = \mathcal{B}_{2,1}^{(\delta)} - \nu(v)$$

is independent of ε and satisfies (see [21] proof of Lemma 4.19) for all g in $L_v^1(\langle v \rangle^k \nu)$

$$\|\mathcal{J}^{(\delta)} g\|_{L_v^1(\langle v \rangle^k)} \leq C_\delta \|g\|_{L_v^1(\langle v \rangle^k \nu)}.$$

In the same way as proof of Proposition 3.6 we obtain

$$(3.21) \quad \frac{d}{dt} \|D_t f\|_{L_{x,v}^1(\langle v \rangle^k)} \leq -\frac{\lambda_0}{\varepsilon^2} \|D_t f\|_{L_{x,v}^1(\langle v \rangle^k \nu)} + \frac{C_\delta}{\varepsilon^2} \|f\|_{L_v^1(\langle v \rangle^k \nu)}.$$

We then consider λ'_0 in $(0, \lambda_0)$ and define $\eta = (\lambda_0 - \lambda'_0)/C_\delta$. We compute, with (3.19),

$$\frac{d}{dt} \left[e^{\frac{\lambda'_0}{\varepsilon^2} t} \left(\eta \|D_t f\|_{L_{x,v}^1(\langle v \rangle^k)} + \|f\|_{L_{x,v}^1(\langle v \rangle^k)} \right) \right] \leq 0,$$

and thus

$$(3.22) \quad \|D_t f\|_{L_{x,v}^1(\langle v \rangle^k)} + \|f\|_{L_{x,v}^1(\langle v \rangle^k)} \leq \eta^{-1} e^{-\frac{\lambda'_0}{\varepsilon^2} t} \|h\|_{W_v^{1,1} L_x^1(\langle v \rangle^k)}.$$

To conclude we plug (3.22) into (3.18) and we combine it with (3.16) into (3.15). This yields, because $s \geq 1$,

$$\|T_1(t)h\|_{W_v^{s+1,1} W_x^{1,1}(K)} \leq C \frac{e^{-\frac{\lambda'_0}{\varepsilon^2} t}}{\varepsilon^2 t} \|h\|_{W_v^{s,1} L_x^1(\langle v \rangle^k)},$$

which implies the expected result (3.13) because $T_1(t)$ commutes with x -derivatives.

Control of $T_2(t)h$: For $s \geq 0$ we can interpolate (for interpolation theory in Sobolev spaces see [6] Chapters 6) between (3.16) and (3.13) to get

$$\|T_1(t)h\|_{W_{x,v}^{s+1/2,1}(K)} \leq C \frac{e^{-\frac{\lambda'_0}{\varepsilon^2} t}}{\varepsilon^2 \sqrt{t}} \|h\|_{W_v^{s,1} L_x^1(\langle v \rangle^k)}.$$

Then, we firstly use the inequality above and secondly (3.16) to obtain

$$\begin{aligned} \|T_2(t)h\|_{W_{x,v}^{s+1/2,1}(K)} &\leq \int_0^t \|T_1(t-s)T_1(s)h\|_{W_{x,v}^{s+1/2,1}(K)} ds \\ &\leq \frac{C}{\varepsilon^4} e^{-\frac{\lambda'_0}{\varepsilon^2} t} \left(\int_0^t \frac{e^{-\frac{\lambda_0 - \lambda'_0}{\varepsilon^2} s}}{\sqrt{t-s}} ds \right) \|h\|_{W_v^{s,1} L_x^1(\langle v \rangle^k)}, \end{aligned}$$

which is the expected result (3.14).

□

The aim is to link our space $L_v^q L_x^p(\langle v \rangle^k)$ to the space $H_{x,v}^s(\mu^{-1/2})$. We thus state the following control on the iterated convolution in the case $p = 2$.

Proposition 3.8. *Consider $k > k_q^*$, defined by (2.3), and s in \mathbb{N} . For any δ in $(0, \delta_{k,q}]$ there exists $C_2 = C_2(\delta) > 0$ and $R = R(\delta) > 0$ such that for any $t \geq 0$,*

$$\forall n \in \mathbb{N}, \quad \text{supp } T_n(t)h \subset K := B(0, R)$$

and

$$(3.23) \quad \forall s \geq 0, \quad \|T_2(t)h\|_{H_{x,v}^{s+1/2}(K)} \leq \frac{C_T}{\varepsilon^{5/2}} \|h\|_{H_{x,v}^s(\langle v \rangle^k)}.$$

Proof of Proposition 3.8. Consider h in $W_{x,v}^{s,2}(\langle v \rangle^k)$, s in \mathbb{N} .

This Proposition is easier than when $p = 1$ because there exists velocity averaging lemmas in this framework, as discussed in [21] Remark 4.21. The x -derivative commutes with T_1 and therefore we suppose there is no derivative in space.

Define $f(t) = S_{\mathcal{B}_\varepsilon^{(\delta)}}(t)(h)$ so that f is solution to the kinetic equation

$$\partial_t f + \frac{1}{\varepsilon} v \cdot \nabla_x f = s_\varepsilon(t, x, v),$$

with $s_\varepsilon(t, x, v) = -\varepsilon^{-2} \nu f + \varepsilon^{-2} \mathcal{B}_{2,1}^{(\delta)} f$.

Let j be a multi-index such that $|j| \leq s$. We apply ∂_0^j to the latter equation, which gives

$$(3.24) \quad \partial_t (\partial_0^j f) + \frac{1}{\varepsilon} v \cdot \nabla_x (\partial_0^j f) = \partial_0^j s_\varepsilon(t, x, v) + \frac{1}{\varepsilon} \sum_{|i|+|l|=|j|} a_{i,l} \partial_l^i f,$$

where $a_{i,j}$ are non-negative numbers.

A classical averaging lemma (see [9] Lemma 1 and [10] in which we emphasize the dependencies in ε) reads, for (3.24) with $\partial_0^j f(0, x, v) = \partial_0^j h(x, v)$, for all ψ in $\mathcal{D}(\mathbb{R}^d)$

$$(3.25) \quad \left\| \int_{\mathbb{R}^d} \partial_0^j f(t, x, v) \psi(v) dv \right\|_{L_t^2(H_x^{1/2})} \leq \frac{C}{\sqrt{\varepsilon}} \left(\left\| \partial_0^j h(x, v) \right\|_{L_{x,v}^2} + \left\| \partial_0^j s_\varepsilon \right\|_{L_{t,x,v}^2} + \frac{1}{\varepsilon} \left\| \sum_{|i|+|l|=|j|} a_{i,l} \partial_l^i f \right\|_{L_{t,x,v}^2} \right).$$

We use [21], Lemmas 4.4 and 4.7, in order to bound the terms involving $\mathcal{B}_{2,\varepsilon}^{(\delta)} = \varepsilon^{-2} \mathcal{B}_{2,1}^{(\delta)}$ we have that

$$\|s_\varepsilon\|_{H_{x,v}^s(\langle v \rangle^k)} \leq \frac{1}{\varepsilon^2} \|s_1\|_{H_{x,v}^s(\langle v \rangle^k)} \leq \frac{C}{\varepsilon^2} \|f\|_{H_{x,v}^s(\langle v \rangle^k \nu)} \leq \frac{C}{\varepsilon^2} e^{-\frac{\lambda_0}{\varepsilon^2} t} \|h\|_{H_{x,v}^s(\langle v \rangle^k \nu)},$$

where the last inequality comes from the hypodissipativity properties of $S_{\mathcal{B}_\varepsilon}(t)$, see Proposition 3.6.

Using the dissipativity properties of $S_{\mathcal{B}_\varepsilon}(t)$ one more time we deduce that

$$(3.26) \quad \|T_1(t)h\|_{L_t^2(H_{x,v}^{s+1/2}(\langle v \rangle^k))} \leq \frac{C}{\varepsilon^{5/2}} \|h\|_{H_{x,v}^s(\langle v \rangle^{k\nu})}.$$

To conclude we notice that $\int_0^t T_1(t-s)T_1(s) ds$ is a continuous linear operator on the Hilbert space $H_{x,v}^{s+1/2}(K)$ and thus we can see it as an element of $H_{x,v}^{s+1/2}(K)$ by Riesz's representation theorem. Hence, thanks to Cauchy-Schwartz,

$$\begin{aligned} \|T_2(t)h\|_{H_{x,v}^{s+1/2}(K)} &= \left\| \left(\int_0^t T_1(t-s)T_1(s) ds \right) (h) \right\|_{H_{x,v}^{s+1/2}(K)} \\ &\leq \|h\|_{H_{x,v}^s(\langle v \rangle^{k\nu})} \int_0^t \|T_1(t-s)T_1(s)\|_{\mathcal{B}(H_{x,v}^s(\langle v \rangle^{k\nu}), H_{x,v}^{s+1/2}(K))} ds \\ &\leq \|h\|_{H_{x,v}^s(\langle v \rangle^{k\nu})} \left(\int_0^t \|T_1(t-s)\|_{\mathcal{B}(H_{x,v}^s(K), H_{x,v}^{s+1/2}(K))}^2 ds \right)^{1/2} \\ &\quad \times \left(\int_0^t \|T_1(s)\|_{\mathcal{B}(H_{x,v}^s(\langle v \rangle^{k\nu}), H_{x,v}^s(K))}^2 ds \right)^{1/2} \\ &\leq \|h\|_{H_{x,v}^s(\langle v \rangle^{k\nu})} \frac{C}{\varepsilon^{5/2}} \left(\int_0^t \frac{C_A}{\varepsilon^2} e^{-\frac{\lambda'_0}{\varepsilon^2}s} ds \right)^{1/2} \\ &\leq \frac{C}{\varepsilon^{5/2}} \|h\|_{H_{x,v}^s(\langle v \rangle^{k\nu})}, \end{aligned}$$

where we used Proposition 3.2 and the fact that $S_{\mathcal{B}_\varepsilon^{(\delta)}}$ is a contraction semigroup on $H_{x,v}^s$ with spectral gap λ'_0/ε^2 . □

3.5. Proof of Theorem 2.1. As we explained it in Section 3.1, the proof of Theorem 2.1 is direct from the application of Theorem 3.1. This theorem is clearly applicable in our case and we emphasize it through the extreme case of no derivative in space or velocity variables.

Indeed, we consider s in \mathbb{N} to be chosen big enough later. We define $\mathcal{E} = L_v^q L_x^p(\langle v \rangle^k)$ and $E = H_{x,v}^s(\mu^{-1/2})$ and we have $E \subset \mathcal{E}$ for s big enough (dense with continuous embedding). Indeed, in the case $q \geq 2$ and $p \geq 2$, standard Sobolev embeddings (see [11] Section IX.3.) imply $E \subset L_v^q L_x^p(\mu^{-1/2})$. In the case $p < 2$ we have, on the torus, $L_x^2 \subset L_x^p$ and $H_x^s \subset L_x^2$ by the same Sobolev embeddings. Finally, in the case $q < 2$ we have that $L_v^2(\mu^{-1/2}) \subset L_v^q(\langle v \rangle^k)$ (it can be done by a mere Cauchy-Schwarz inequality) and the same Sobolev embeddings give $H_v^s(\mu^{-1/2}) \subset L_v^2(\mu^{-1/2})$.

On the torus we have the following embedding: $L_x^p \subset L_x^1$. Thanks to Proposition 3.2 and Proposition 3.6 we obtain (same arguments as (3.16))

$$(3.27) \quad \|T_1(t)h\|_{\mathcal{E}} \leq C \left\| \mathcal{A}_\varepsilon^{(\delta)} S_{\mathcal{B}_\varepsilon^{(\delta)}} h \right\|_{L_v^q L_x^1(K)} \leq \frac{C}{\varepsilon^2} e^{-\frac{\lambda_0}{\varepsilon^2}t} \|h\|_{L_v^1 L_x^1(\langle v \rangle^k)}.$$

We therefore define $\mathcal{E}_2 = L_v^1 L_x^1(\langle v \rangle^k)$.

Then we define by $\mathcal{E}_j = W_{x,v}^{(j-2)/2,1}(\langle v \rangle^k)$ for j from 2 to m with m big enough such that $W_{x,v}^{(m-1)/2,1}(\langle v \rangle^k) \subset L_{x,v}^2(\langle v \rangle^k)$. Then we denote $\mathcal{E}_j = H_{x,v}^{(j-m-1)/2}(\langle v \rangle^k)$ up to $J-1$ where $H_{x,v}^{(J-m-2)/2}(\langle v \rangle^k) \subset E$.

Point (A1) of Theorem 3.1 is satisfied thanks to [2] and [12] Theorem 2.1 (with the norm of Theorem 2.4), point (A2)(i) by Proposition 3.6 and point (A2)(ii) by Proposition 3.2. Finally, point (A2)(iii) is given by (3.27) for \mathcal{E} and \mathcal{E}_1 , then by Proposition 3.7 (3.14) up to \mathcal{E}_m and by Proposition 3.8 from \mathcal{E}_m to \mathcal{E}_J and E .

4. AN *a priori* ESTIMATE FOR THE FULL PERTURBED EQUATION: PROOF OF THEOREM 2.3

In this section we work in $W_v^{\alpha,1}H_x^\beta(\langle v \rangle^k)$ or in $W_v^{\alpha,1}W_x^{\beta,1}(\langle v \rangle^k)$, with $\alpha \leq \beta$ on the full perturbed Boltzmann equation

$$\partial_t h = \mathcal{G}_\varepsilon(h) + \frac{1}{\varepsilon}Q(h, h).$$

4.1. Description of the problem and notations. When $\varepsilon = 1$, the linear part \mathcal{G}_ε has the same order of magnitude than the bilinear term Q in the linearized Boltzmann equation (1.5). In this case, Theorem 2.1 suffices to obtain existence and exponential decay since the contraction property of the semigroup $S_{\mathcal{G}_1}$ controls the bilinear part for small initial data (see [21]).

In the general case, $S_{\mathcal{G}_\varepsilon}$ only generates a semigroup with a spectral gap of order 1, insufficient to control $\varepsilon^{-1}Q$. However, [23][12] show that a careful study of $\varepsilon^{-1}Q$ compared to G_ε yields existence and exponential decay of solutions to (1.5) in $H_{x,v}^s(\mu^{-1/2})$ for s large enough (see Theorem 4.7 for an adapted version of this result). Our strategy is to use the same kind of ideas as when we extended the semigroup properties from $H_{x,v}^\beta(\mu^{-1/2})$ to $W_v^{\alpha,1}H_x^\beta(\langle v \rangle^k)$ and $W_v^{\alpha,1}W_x^{\beta,1}(\langle v \rangle^k)$ but including the bilinear term. Namely, we shall decompose the partial differential equation (1.5) into a system of partial differential equations from $W_v^{\alpha,1}H_x^\beta(\langle v \rangle^k)$ or $W_v^{\alpha,1}W_x^{\beta,1}(\langle v \rangle^k)$ to $H^\beta(\mu^{-1/2})$ and use the perturbative estimates of [12].

As noticed in Remark 2.16 of [21], Theorem 3.1 extending the semigroup generated by \mathcal{G}_ε in $H^s(\mu^{-1/2})$ to $L_v^1L_x^\infty(\langle v \rangle^k)$ can be interpreted as a decomposition of

$$\partial_t f = \mathcal{G}_\varepsilon f,$$

into a system of partial differential equations, involving operators $\mathcal{G}_\varepsilon = \mathcal{A}_\varepsilon + \mathcal{B}_\varepsilon$ (defined in Section 3.2), with $f = f^1 + \dots + f^J$ satisfying

- f^1 is in $L_v^1L_x^\infty(\langle v \rangle^k)$ and $f_{in}^1 = f_{in}$ in $\text{Ker}(\mathcal{G}_\varepsilon)^\perp$,
- for all $2 \leq j \leq J-1$, f^j is in \mathcal{E}_j and $f_{in}^j = 0$,
- f^J is in $H^s(\mu^{-1/2})$, $f_{in}^J = 0$ and in that space we can use the contraction property of $S_{\mathcal{G}_\varepsilon}$.

We will decompose the linearized Boltzmann equation in a similar way than the one explained above. We shall define a sequence of spaces $(\mathcal{E}_j)_{1 \leq j \leq J}$. In each space \mathcal{E}_j , $1 \leq j \leq J-1$, a piece of the bilinear term, of order ε^{-1} , will be added and controlled by the dissipativity property of $\mathcal{B}_\varepsilon^{(\delta)}$, of order ε^{-2} . Contrary to the study in the linear case, the bilinear operator generates terms involving functions in all the spaces \mathcal{E}_j which have to be compared and controlled. This imposes to construct $(\mathcal{E}_j)_{1 \leq j \leq J}$ as a nested sequence.

The difficult part of the linear operator, namely $\mathcal{A}_\varepsilon^{(\delta)}$, enjoys a regularising effect and could therefore be treated in more regular spaces. Of course, our decomposition will be much easier since we solely want to go from an exponential weight into a polynomial weight Sobolev spaces, without losing any derivatives in x or v .

In order to shorten notations we define, for $p = 1, 2$ and k to be defined later,

$$(4.1) \quad \mathcal{E}^p = W_v^{\alpha,1} W_x^{\beta,p} (\langle v \rangle^k) \quad \text{and} \quad E = H_{x,v}^\beta (\mu^{-1/2}).$$

We take h_{in} in \mathcal{E}^p and we decompose the partial differential equation,

$$\partial_t h = \mathcal{G}_\varepsilon(h) + \frac{1}{\varepsilon} Q(h, h) = \mathcal{A}_\varepsilon^{(\delta)}(h) + \mathcal{B}_\varepsilon^{(\delta)}(h) + \frac{1}{\varepsilon} Q(h, h)$$

into an equivalent system of partial differential equations for the following decomposition

$$(4.2) \quad h(t, x, v) = h^0(t, x, v) + h^1(t, x, v),$$

with

(1) In \mathcal{E}^p , $h_{t=0}^0 = h_{in}$ and

$$(4.3) \quad \partial_t h^0 = \mathcal{B}_\varepsilon^{(\delta)}(h^0) + \frac{1}{\varepsilon} Q(h^0, h^0) + \frac{2}{\varepsilon} Q(h^0, h^1),$$

(2) In E , $h_{t=0}^1 = 0$ and

$$(4.4) \quad \partial_t h^1 = \mathcal{G}_\varepsilon(h^1) + \frac{1}{\varepsilon} Q(h^1, h^1) + \mathcal{A}_\varepsilon^{(\delta)}(h^0).$$

The aim of this Section is to establish the following estimate of solutions to the system (4.3) – (4.4).

Theorem 4.1. *Let $p = 1$ or $p = 2$. There exist β_0 in \mathbb{N} and ε_d in $(0, 1]$ depending on d and the kernel of the Boltzmann operator such that:*

For all $\beta \geq \beta_0$, for any δ in $(0, \delta_{k,1}]$ and any λ'_0 in $(0, \lambda_0)$ ($\delta_{k,1}$ and λ_0 defined in Proposition 3.6) there exist $C_\beta, \eta_\beta > 0$ such that for any $0 < \varepsilon \leq \varepsilon_d$ and h_{in} in \mathcal{E}^p , if

- (i) $\|h_{in}\|_{\mathcal{E}^p} \leq \eta_\beta$,
- (ii) (h^0, h^1) is solution to the system (4.3) – (4.4),

then

$$\|h^0 + h^1\|_{\mathcal{E}^p} \leq C_\beta \|h_{in}\|_{\mathcal{E}^p} e^{-\lambda'_0 t}.$$

The constants C_β and η_β are constructive and depends only on $\beta, d, \delta, \lambda'_0$ and the kernel of the Boltzmann operator.

Remark 4.2 (Link with Theorem 2.3). *The existence and uniqueness for the perturbed Boltzmann equation (1.5) in \mathcal{E}^p has been proved for $\varepsilon = 1$, that is equivalent of ε fixed with constant depending on it, in [21] Theorems 5.3 and 5.5 respectively. The constants, as well as the smallness assumption on the initial data, in the theorem above are independent of ε and therefore this a priori result combined with existence and uniqueness developed in [21] and in [12] implies the existence and uniqueness independently of ε which is Theorem 2.3.*

The next subsections deal with the estimates one can get for solutions to the system (4.3) – (4.4). We study each of them independently and the *a priori* exponential decay will be a straightforward application of these results together with a maximum principle argument.

Section 4.2 focuses on the *a priori* study of the equation in \mathcal{E}^p . Section 4.3 deals with (4.4) in E . Finally, Section 4.4 gathers the previous results to prove Theorem 4.1.

4.2. Study of equation (4.3) in \mathcal{E} . In this section we prove the following general proposition about the equation taking place in $\mathcal{E}^p = W_v^{\alpha,1}W_x^{\beta,p}(\langle v \rangle^k)$, for $p = 1$ or $p = 2$. We define the shorthand notation

$$\mathcal{E}_\nu^p = W_v^{\alpha,1}W_x^{\beta,p}(\langle v \rangle^k \nu).$$

Proposition 4.3. *Let $p = 1$ or $p = 2$ and $0 < \varepsilon \leq 1$. Let $k > k_1^* = 2$, $\beta > 2d/p$.*

Let h_{in} be in \mathcal{E}^p and h^1 in \mathcal{E}_ν^p .

For any δ in $(0, \delta_{k,1}]$ and any λ'_0 in $(0, \lambda_0)$ ($\delta_{k,1}$ and λ_0 defined in Proposition 3.6) there exist $C_0, \eta_0 > 0$ such that

if

- (i) $\|h_{in}\|_{\mathcal{E}^p} \leq \eta_0$, $\|h^1\|_{\mathcal{E}_\nu^p} \leq \eta_0$,
- (ii) h^0 satisfies $h^0_{t=0} = h_{in}$ and is solution to

$$\partial_t h^0 = \mathcal{B}_\varepsilon^{(\delta)}(h^0) + \frac{1}{\varepsilon}Q(h^0, h^0) + \frac{2}{\varepsilon}Q(h^0, h^1),$$

then

$$\|h^0\|_{\mathcal{E}^p} \leq C_0 e^{-\frac{\lambda'_0}{\varepsilon^2}t} \|h_{in}\|_{\mathcal{E}^p}.$$

The constant η_0 is constructive and depends only on δ, λ'_0 and the kernel of the Boltzmann operator.

Note that the presence of C_0 is due to the hypodissipativity of the linear operator rather than a direct dissipativity. We need to control the bilinear term Q , which is given by the following lemma.

Lemma 4.4. *For all $p = 1, 2$ and α, β in \mathbb{N} such that $\beta > 2d/p$, there exists $C_{\beta,p} > 0$ such that all f and g*

$$\|Q(f, g)\|_{\mathcal{E}^p} \leq C_{\beta,p} (\|g\|_{\mathcal{E}_\nu^p} \|f\|_{\mathcal{E}^p} + \|g\|_{\mathcal{E}^p} \|f\|_{\mathcal{E}_\nu^p}).$$

This lemma has been proved in Lemma 5.16 in [21], which is adapted from interpolation results in [1] or duality arguments as in [30] Theorem 2.1.

Proof of Proposition 4.3. Consider δ in $(0, \delta_{k,1}]$ and λ'_0 in $(0, \lambda_0)$. Take $p = 1$ or $p = 2$ and $\beta > 2d/p$.

We have that

$$\partial_t h^0 = \mathcal{B}_\varepsilon^{(\delta)}(h^0) + \frac{1}{\varepsilon}Q(h^0, h^0) + \frac{2}{\varepsilon}Q(h^0, h^1).$$

Thanks to the hypodissipativity of property of $\mathcal{B}_\varepsilon^{(\delta)}$, more precisely the proof of Lemma 3.6, we have

$$\begin{aligned} \frac{d}{dt} \|h^0\|_{\mathcal{E}^p} &\leq -\frac{\lambda_0}{\varepsilon^2 \nu_0} \|h^0\|_{\mathcal{E}_v^p} + \frac{1}{\varepsilon} |\langle Q(h^0, h^0) + 2Q(h^0, h^1), h^0 \rangle_{\mathcal{E}^p}| \\ &\leq -\frac{\lambda_0}{\varepsilon^2 \nu_0} \|h^0\|_{\mathcal{E}_v^p} + \frac{1}{\varepsilon} \|Q(h^0, h^0) + 2Q(h^0, h^1)\|_{\mathcal{E}^p}, \end{aligned}$$

where we used the scalar product notation to refer to the product operator appearing in $W_v^{\alpha,1} W_x^{\beta,p}$ when one differentiates $\|h\|_{W_v^{\alpha,1} W_x^{\beta,p}(\langle v \rangle^k)}$ (of the same form as (3.9)).

For the second inequality we used Hölder inequality between L_x^p and $L_x^{p/(p-1)}$ inside the product operator:

$$\int_{\mathbb{T}^d} \operatorname{sgn}(h^0) |h^0|^{p-1} F(h^0) dx \leq \|h^0\|_{L_x^p}^{p-1} \|F(h^0)\|_{L_x^p}.$$

Then estimating Q using Lemma 4.4 yields

$$(4.5) \quad \frac{d}{dt} \|h^0\|_{\mathcal{E}^p} \leq -\frac{1}{\varepsilon^2} \left[\frac{\lambda_0}{\nu_0} - 2\varepsilon C_{\beta,p} \left(\|h^0\|_{\mathcal{E}^p} + \frac{2}{\nu_0} \|h^1\|_{\mathcal{E}_v^p} \right) \right] \|h^0\|_{\mathcal{E}_v^p},$$

we recall $\nu_0 = \inf_{v \in \mathbb{R}^d} (\nu(v)) > 0$.

Therefore, if

$$\|h^1\|_{\mathcal{E}_v^p} \leq \varepsilon^{-1} \frac{(\lambda_0 - \lambda'_0)}{8C_{\beta,p}} \quad \text{and} \quad \|h_{t=0}\|_{\mathcal{E}^p} \leq \varepsilon^{-1} \frac{(\lambda_0 - \lambda'_0)}{4\nu_0 C_{\beta,p}},$$

then $\|h^0\|_{\mathcal{E}^p}$ is always decreasing in time with

$$\frac{d}{dt} \|h^0\|_{\mathcal{E}^p} \leq -\frac{\lambda'_0}{\varepsilon^2 \nu_0} \|h^0\|_{\mathcal{E}_v^p},$$

which hence yields the expected exponential decay by Grönwall Lemma. \square

4.3. Study of equations (4.4) in E . In the space $E = H_{x,v}^\beta(\mu^{-1/2})$, solutions to the perturbed Boltzmann equation enjoy an exponential decay. More precisely, [12] derived a precise Grönwall that we will now use to obtain estimates on the solution h^1 . We will use the following shorthand notation

$$E_\nu = H_{x,v}^\beta(\mu^{-1/2} \nu^{1/2})$$

In this section we use the previous theorem to obtain exponential decay of h^1 in E . This result is stated in the following proposition, where C_t^0 denotes the space of time-continuous functions.

Proposition 4.5. *Let $p = 1$ or $p = 2$, $0 < \varepsilon \leq \varepsilon^d \leq 1$, $\beta \geq s_0$ and $\alpha \leq \beta$ (ε_d and s_0 being constructive constants that will be defined in Theorem 4.7).*

Let h_{in} be in \mathcal{E}^p and h^0 in $C_t^0 \mathcal{E}^p$.

For any δ in $(0, \delta_{k,1}]$ and any λ'_0 in $(0, \lambda_0)$ ($\delta_{k,1}$ and λ_0 defined in Proposition 3.6) there exist $\eta_1, C_1 > 0$ such that

if

- (i) $\|h_{in}\|_{\mathcal{E}^p} \leq \eta_1$,
- (ii) *there exists $C_0 > 0$ such that $\|h^0\|_{\mathcal{E}^p} \leq C_0 e^{-\frac{\lambda_0 + \lambda'_0}{2\varepsilon^2} t} \|h_{in}\|_{\mathcal{E}^p}$,*
- (iii) h^1 *satisfies $h^1_{t=0} = 0$ and is solution to*

$$\partial_t h^1 = \mathcal{G}_\varepsilon(h^1) + \frac{1}{\varepsilon} Q(h^1, h^1) + \mathcal{A}_\varepsilon^{(\delta)}(h^0)$$

then

$$\|h^1\|_E \leq C_1 e^{-\lambda'_0 t} \|h_{in}\|_{\mathcal{E}^p}$$

The constants C_1 and η_1 are constructive and depends only on δ , λ'_0 and the kernel of the Boltzmann operator.

In order to prove Proposition 4.5 we need a new control on the bilinear term.

For any operator F defined on $E \times E$, we will say that F satisfies the property (H) if the following holds.

Property (H):

- (1) for all g^1, g^2 in E we have $\pi_L(F(g^1, g^2)) = 0$, where π_L is the orthogonal projection on $\text{Ker}(L)$ in $L_v^2(\mu^{-1/2})$ (see (2.1)),
- (2) for all $s' > 0$ there exists $\mathcal{F}_F^{s'} : E \times E \rightarrow \mathbb{R}^+$ such that for all multi-indexes j and l such that $|j| + |l| \leq s'$,

$$\left| \langle \partial_t^j F(g^1, g^2), g^3 \rangle_{L_{x,v}^2(\mu^{-1/2})} \right| \leq \mathcal{F}_F^{s'}(g^1, g^2) \|g^3\|_{L_{x,v}^2(\mu^{-1/2} \nu^{1/2})},$$

with $\mathcal{F}_F^{s'} \leq \mathcal{F}_F^{s'+1}$.

Lemma 4.6. *The Boltzmann linear operator Q satisfies the property (H) with*

$$\forall s > d, \exists C_s > 0, \quad \mathcal{F}_Q^s(f, g) \leq C_s [\|f\|_E \|g\|_{E_\nu} + \|f\|_{E_\nu} \|g\|_E].$$

The latter control on the bilinear part is from [12] Appendix A.2.

Proof of Proposition 4.5. We state below the estimate derived in [12] (note that this is a version of [12] Theorem 2.4 extended by estimates proved in [12] Propositions 2.2 and 7.1).

Theorem 4.7. *There exist $0 < \varepsilon_d \leq 1$ and s_0 in \mathbb{N} such that for any $s \geq s_0$ and any λ'_0 in $(0, \lambda_0)$ there exists $\delta_s, C_s > 0$ such that,*

- *for any h_{in} in $H_{x,v}^s(\mu^{-1/2})$ with*

$$\|h_{in}\|_{H_{x,v}^s(\mu^{-1/2})} \leq \delta_s,$$

- *for any operator F defined on $H_{x,v}^s(\mu^{-1/2}) \times H_{x,v}^s(\mu^{-1/2})$ satisfying the property (H);*

Then for all $0 < \varepsilon \leq \varepsilon_d$ and for all g^1, g^2 in $H_{x,v}^s(\mu^{-1/2})$, if h is a solution to

$$\begin{cases} \partial_t h = G_\varepsilon(h) + \frac{1}{\varepsilon} F(g^1, g^2) \\ h_{t=0} = h_{in}, \end{cases}$$

and h is in $\text{Ker}(G_\varepsilon)$ for all time, then

$$\forall t \in \mathbb{R}^+, \quad \frac{d}{dt} \|h\|_{H_{x,v}^s(\mu^{-1/2})}^2 \leq -\frac{2\lambda_0''}{\nu^2} \|h\|_{H_{x,v}^s(\mu^{-1/2\nu})}^2 + C_s (\mathcal{F}_F^s(g^1, g^2))^2.$$

Remark 4.8. It is of importance to notice that the differential inequality given in Theorem 4.7 is independent of ε . In fact, it holds for a modified $\|\cdot\|_{H_{x,v}^s, \varepsilon}$ that includes ε factors but in a way that $\|\cdot\|_{H_{x,v}^s, \varepsilon} \sim \|\cdot\|_{H_{x,v}^s}$ independently of ε as long as ε is sufficiently small. For clarity purposes, working directly with the $\|\cdot\|_{H_{x,v}^s}$ does not hide any problematic behaviour. If one wants to incorporate the real $\|\cdot\|_{H_{x,v}^s, \varepsilon}$, one just have to do exactly the following computations but with the integrated in time version of Theorem 4.7 where one can use the usual $\|\cdot\|_{H_{x,v}^s}$ thanks to the equivalence of the norms.

Now, let λ'' be in $(0, \lambda_0)$, $s \geq s_0$ and $0 < \varepsilon \leq \varepsilon_d$.

The proof of Proposition 4.5 will be done in two steps. First we study the projection of h^1 onto $\text{Ker}(G_\varepsilon)$ and then its orthogonal part.

Estimate on the projection part. We have that, see the decomposition (4.2), that $h^1 = h - h^0$ with h solution to the perturbed Boltzmann equation and thus satisfying $\Pi_G(h) = 0$. We therefore have that

$$\Pi_G(h^1) = -\Pi_G(h^0).$$

Moreover, Theorem 2.1 tells us that Π_G and Π_G coincide on E and thus

$$\Pi_G(h^1) = -\Pi_G(h^0),$$

and assumption (ii) together with the shape of Π_G (see (2.4)), there exists a constant $C_\Pi > 0$, depending only on the dimension d and s and the constant C_0 , such that

$$(4.6) \quad \|\Pi_G(h^1)\|_{E_\nu} \leq C_\Pi e^{-\frac{\lambda_0 + \lambda_0'}{2\varepsilon^2} t} \|h_{in}\|_{\mathcal{E}^p}.$$

Estimate on the orthogonal part. Applying $\Pi_G^\perp = \text{Id} - \Pi_G$, the orthogonal projection onto $(\text{Ker}(G_\varepsilon))^\perp$ in $L_{x,v}^2(\mu^{-1/2})$, to the differential equation satisfied by h^1 yields

$$\begin{aligned} \partial_t (\Pi_G^\perp(h^1)) &= G_\varepsilon(h^1) + \Pi_G^\perp \left(\frac{1}{\varepsilon} Q(h^1, h^1) + \mathcal{A}_\varepsilon^{(\delta)}(h^0) \right) \\ (4.7) \quad &= G_\varepsilon(\Pi_G^\perp(h^1)) + \Pi_G^\perp \left(\frac{1}{\varepsilon} Q(h^1, h^1) + \mathcal{A}_\varepsilon^{(\delta)}(h^0) \right). \end{aligned}$$

Moreover, we have by definition of Π_G and π_L (see (2.4) and (2.1)) that

$$(\pi_L(h) = 0) \implies (\Pi_G(h) = 0)$$

and therefore

$$\Pi_G^\perp(Q(h^1, h^1)) = Q(h^1, h^1),$$

since Q satisfies property (H).1. by Lemma 4.6. Plugging the latter equality into (4.7) gives

$$\partial_t (\Pi_G^\perp(h^1)) = G_\varepsilon (\Pi_G^\perp(h^1)) + \frac{1}{\varepsilon} Q(h^1, h^1) + \Pi_G^\perp(\mathcal{A}_\varepsilon^{(\delta)}(h^0)).$$

By definition, $\Pi_G^\perp(h^1)$ is in $(\text{Ker}(G_\varepsilon))^\perp$ for all time and thanks to the control on the Boltzmann operator Q in E (Lemma 4.6), we are able to use Theorem 4.7 with $\lambda_0 > \lambda'_0$ to which we have to add the source term $\Pi_G^\perp(\mathcal{A}_\varepsilon^{(\delta)}(h^0))$. This yields the following differential inequality, where we denote by C any positive constant independent of ε ,

$$(4.8) \quad \begin{aligned} \frac{d}{dt} \|\Pi_G^\perp(h^1)\|_E^2 &\leq -\frac{2\lambda_0''}{\nu_0^2} \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 + C (\mathcal{F}_Q^s(h^1, h^1))^2 + |\langle \Pi_G^\perp(\mathcal{A}_\varepsilon^{(\delta)}(h^0)), \Pi_G^\perp(h^1) \rangle_E| \\ &\leq -\frac{2\lambda_0''}{\nu_0^2} \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 + C \|h^1\|_E^2 \|h^1\|_{E_\nu}^2 + \|\Pi_G^\perp(\mathcal{A}_\varepsilon^{(\delta)}(h^0))\|_E \|\Pi_G^\perp(h^1)\|_E, \end{aligned}$$

where we used a Cauchy-Schwarz inequality on the last term on the right-hand side.

Then we can decompose $h^1 = \Pi_G(h^1) + \Pi_G^\perp(h^1)$ to get first

$$\begin{aligned} \|h^1\|_E^2 \|h^1\|_{E_\nu}^2 &\leq 4 \|\Pi_G^\perp(h^1)\|_E^2 \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 + \frac{8}{\nu_0^2} \|\Pi_G(h^1)\|_{E_\nu}^2 \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 \\ &\quad + \frac{4}{\nu_0^2} \|\Pi_G(h^1)\|_{E_\nu}^4, \end{aligned}$$

into which we can plug the control on $\|\Pi_G(h^1)\|_{E_\nu}^2$ we derived in (4.6) to obtain, with $\|h_{in}\| \leq \eta_1$,

$$(4.9) \quad \|h^1\|_E^2 \|h^1\|_{E_\nu}^2 \leq 4 \|\Pi_G^\perp(h^1)\|_E^2 \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 + C\eta_1^2 \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 + C e^{-\frac{2(\lambda_0 + \lambda'_0)}{\varepsilon^2} t} \|h_{in}\|_{\mathcal{E}^p}^4.$$

And finally, this inequality together with assumption (ii) gives the existence of a constant $C_A > 0$ such that

$$(4.10) \quad \|\Pi_G^\perp(\mathcal{A}_\varepsilon^{(\delta)}(h^0))\|_E \|\Pi_G^\perp(h^1)\|_E \leq \frac{C_A}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} e^{-\frac{\lambda_0 + \lambda'_0}{2\varepsilon^2} t} \|\Pi_G^\perp(h^1)\|_E.$$

We plug (4.9) and (4.10) into (4.8) and obtain, with C and C' being positive constants independent of ε ,

$$\begin{aligned} \frac{d}{dt} \|\Pi_G^\perp(h^1)\|_E^2 &\leq -\left[\frac{2\lambda_0''}{\nu_0^2} - \left(4 \|\Pi_G^\perp(h^1)\|_E^2 + C\eta_1^2 \right) \right] \|\Pi_G^\perp(h^1)\|_{E_\nu}^2 \\ &\quad + C' \left(\|h_{in}\|_{\mathcal{E}^p}^4 + \frac{1}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} \|\Pi_G^\perp(h^1)\|_E \right) e^{-\frac{\lambda_0 + \lambda'_0}{2\varepsilon^2} t}. \end{aligned}$$

We now choose η_1 sufficiently small so that

$$C\eta_1^2 \leq \frac{\lambda_0'' - \lambda'_0}{\nu_0^2},$$

which in turns implies

$$(4.11) \quad \begin{aligned} \frac{d}{dt} \|\Pi_G^\perp(h^1)\|_E^2 &\leq - \left[\frac{\lambda_0'' + \lambda_0'}{\nu_0^2} - 4 \|\Pi_G^\perp(h^1)\|_E^2 \right] \|\Pi_G^\perp(h^1)\|_{E\nu}^2 \\ &\quad + C' \left(\|h_{in}\|_{\mathcal{E}^p}^4 + \frac{1}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} \|\Pi_G^\perp(h^1)\|_E \right) e^{-\frac{\lambda_0 + \lambda_0'}{2\varepsilon^2} t}. \end{aligned}$$

We define

$$\eta_* = \frac{\lambda_0'' - \lambda_0'}{4\nu_0^2}.$$

We have that $h_{t=0}^1 = 0$ so we can define

$$t_0 = \sup\{t > 0, \quad \|\Pi_G^\perp(h^1)\|_E^2 < \eta_*\}.$$

Suppose that $t_0 < +\infty$, we therefore have for all t in $[0, t_0]$

$$\frac{d}{dt} \|\Pi_G^\perp(h^1)\|_E^2 \leq -\frac{2\lambda_0'}{\nu_0^2} \|\Pi_G^\perp(h^1)\|_{E\nu}^2 + C' \left(\|h_{in}\|_{\mathcal{E}^p}^4 + \frac{\sqrt{\eta_*}}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} \right) e^{-\frac{\lambda_0 + \lambda_0'}{2\varepsilon^2} t},$$

which gives

$$\forall t \in [0, t_0], \quad \frac{d}{dt} \|\Pi_G^\perp(h^1)\|_E^2 \leq -2\lambda_0' \|\Pi_G^\perp(h^1)\|_E^2 + C' \left(\|h_{in}\|_{\mathcal{E}^p}^4 + \frac{\sqrt{\eta_*}}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} \right) e^{-\frac{\lambda_0 + \lambda_0'}{2\varepsilon^2} t},$$

and by Gronwall lemma with $\Pi_G^\perp(h^1)_{(t=0)} = 0$,

$$\begin{aligned} \forall t \in [0, t_0], \quad \|\Pi_G^\perp(h^1)\|_E^2 &\leq C' \left(\|h_{in}\|_{\mathcal{E}^p}^4 + \frac{\sqrt{\eta_*}}{\varepsilon^2} \|h_{in}\|_{\mathcal{E}^p} \right) \left(\int_0^t e^{-\frac{\lambda_0 + \lambda_0'}{2\varepsilon^2} s} e^{2\lambda_0' s} ds \right) e^{-2\lambda_0' t} \\ &\leq C' (\varepsilon^2 \|h_{in}\|_{\mathcal{E}^p}^4 + \sqrt{\eta_*} \|h_{in}\|_{\mathcal{E}^p}) \left(\int_0^{+\infty} e^{-\frac{\lambda_0 - \lambda_0'}{2} u} du \right) e^{-2\lambda_0' t}, \end{aligned}$$

where we used the change of variable $u = \varepsilon^{-2}s$ and we considered $\varepsilon \leq 1/4$ (which only amounts to decreasing ε_d).

Hence, there exists $K > 0$ independent of ε such that

$$\forall t \in [0, t_0], \quad \|\Pi_G^\perp(h^1)\|_E^2 \leq K(\eta_1^4 + \eta_1 \sqrt{\eta_*}).$$

If we thus chose η_1 sufficiently small such that $(\eta_1^4 + \eta_1 \sqrt{\eta_*})K < \eta_*/2$ we reach a contradiction when t goes to t_0 since $\|\Pi_G^\perp(h^1)\|_E^2(t_0) \geq \eta_*$. Therefore, choosing η_1 small enough independently on ε implies first that $t_0 = +\infty$ and second that

$$(4.12) \quad \forall t \in [0, +\infty), \quad \|\Pi_G^\perp(h^1)\|_E^2 \leq C \|h_{in}\|_{\mathcal{E}^p}^2 e^{-2\lambda_0' t}.$$

End of the proof. By just decomposing h^1 into its projection and orthogonal part and using the estimates (4.6) and (4.12) gives the expected exponential decay for h^1 in E . \square

4.4. Proof of Theorem 4.1. Let $p = 1$ or $p = 2$, λ'' be in $(0, \lambda_0)$, $\beta \geq \beta_0 = s_0$ and $0 < \varepsilon \leq \varepsilon_d$. All the constants used in this section are the ones constructed in Proposition 4.3 with $(\lambda_0 + \lambda'_0)/2$ and Proposition 4.5 with λ'_0 .

E is continuously embedded in \mathcal{E}_ν^p because $L_\nu^2(\mu^{-1/2}) \subset L_\nu^2(\langle v \rangle^k)$ (mere Cauchy-Schwarz inequality) and $L_x^2 \subset L_x^1$ because \mathbb{T}^d is bounded. Hence, there exists $C_{E,\varepsilon} > 0$ such that

$$(4.13) \quad \frac{1}{\nu_0} \|\cdot\|_{\mathcal{E}^p} \leq \|\cdot\|_{\mathcal{E}_\nu^p} \leq C_{E,\varepsilon} \|\cdot\|_E.$$

We define

$$\eta = \min\left(\eta_0, \eta_1, \frac{\eta_0}{2C_{E,\varepsilon}C_1}\right),$$

and we assume $\|h_{in}\|_{\mathcal{E}^p} \leq \eta$. Since $h_{t=0}^1 = 0$ we also define

$$t_0 = \sup\{t > 0, \quad \|h^1\|_{\mathcal{E}_\nu^p} < \eta_0\}.$$

Suppose that $t_0 < +\infty$. Then, thanks to Proposition 4.3 we have that

$$\forall t \in [0, t_0], \quad \|h^0\|_{\mathcal{E}^p} \leq \|h_{in}\|_{\mathcal{E}^p} e^{-\frac{\lambda_0 + \lambda'_0}{2\varepsilon^2}t}.$$

We can thus apply Proposition 4.5 and get

$$\forall t \in [0, t_0], \quad \|h^1\|_E \leq C_1 \|h_{in}\|_{\mathcal{E}^p} e^{-\lambda'_0 t} \leq C_1 \eta \leq \frac{\eta_0}{2C_{E,\varepsilon}},$$

which is in contradiction with the definition of t_0 thanks to (4.13). Therefore $t_0 = +\infty$ and we have the expected exponential decay stated in Theorem 4.1 for all time.

REFERENCES

- [1] ARKERYD, L. Stability in L^1 for the spatially homogeneous Boltzmann equation. *Arch. Rational Mech. Anal.* 103, 2 (1988), 151–167.
- [2] BARANGER, C., AND MOUHOT, C. Explicit spectral gap estimates for the linearized Boltzmann and Landau operators with hard potentials. *Rev. Mat. Iberoamericana* 21, 3 (2005), 819–841.
- [3] BARDOS, C., GOLSE, F., AND LEVERMORE, C. D. Fluid dynamic limits of kinetic equations. II. Convergence proofs for the Boltzmann equation. *Comm. Pure Appl. Math.* 46, 5 (1993), 667–753.
- [4] BARDOS, C., GOLSE, F., AND LEVERMORE, D. Fluid dynamic limits of kinetic equations. I. Formal derivations. *J. Statist. Phys.* 63, 1-2 (1991), 323–344.
- [5] BARDOS, C., AND UKAI, S. The classical incompressible Navier-Stokes limit of the Boltzmann equation. *Math. Models Methods Appl. Sci.* 1, 2 (1991), 235–257.
- [6] BERGH, J., AND LÖFSTRÖM, J. *Interpolation spaces. An introduction.* Springer-Verlag, Berlin-New York, 1976. Grundlehren der Mathematischen Wissenschaften, No. 223.
- [7] BOBYLÉV, A. V. The method of the Fourier transform in the theory of the Boltzmann equation for Maxwell molecules. *Dokl. Akad. Nauk SSSR* 225, 6 (1975), 1041–1044.
- [8] BOBYLÉV, A. V. The theory of the nonlinear spatially uniform Boltzmann equation for Maxwell molecules. In *Mathematical physics reviews, Vol. 7*, vol. 7 of *Soviet Sci. Rev. Sect. C Math. Phys. Rev.* Harwood Academic Publ., Chur, 1988, pp. 111–233.
- [9] BOUCHUT, F., AND DESVILLETES, L. Averaging lemmas without time Fourier transform and application to discretized kinetic equations. *Proc. Roy. Soc. Edinburgh Sect. A* 129, 1 (1999), 19–36.
- [10] BOUDIN, L., AND DESVILLETES, L. On the singularities of the global small solutions of the full Boltzmann equation. *Monatsh. Math.* 131, 2 (2000), 91–108.

- [11] BREZIS, H. *Analyse fonctionnelle*. Collection Mathématiques Appliquées pour la Maîtrise. [Collection of Applied Mathematics for the Master's Degree]. Masson, Paris, 1983. Théorie et applications. [Theory and applications].
- [12] BRIANT, M. A constructive method from Boltzmann to incompressible Navier-Stokes on the torus.
- [13] CARLEMAN, T. *Problèmes mathématiques dans la théorie cinétique des gaz*. Publ. Sci. Inst. Mittag-Leffler. 2. Almqvist & Wiksells Boktryckeri Ab, Uppsala, 1957.
- [14] CERCIGNANI, C. *The Boltzmann equation and its applications*, vol. 67 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1988.
- [15] CERCIGNANI, C., ILLNER, R., AND PULVIRENTI, M. *The mathematical theory of dilute gases*, vol. 106 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1994.
- [16] GALLAGHER, I., SAINT-RAYMOND, L., AND TEXIER, B. From newton to boltzmann: the case of short-range potentials.
- [17] GOLSE, F. From kinetic to macroscopic models, 1998.
- [18] GRAD, H. Principles of the kinetic theory of gases. In *Handbuch der Physik (herausgegeben von S. Flügge)*, Bd. 12, *Thermodynamik der Gase*. Springer-Verlag, Berlin, 1958, pp. 205–294.
- [19] GRAD, H. Asymptotic theory of the Boltzmann equation. II. In *Rarefied Gas Dynamics (Proc. 3rd Internat. Sympos., Palais de l'UNESCO, Paris, 1962)*, Vol. I. Academic Press, New York, 1963, pp. 26–59.
- [20] GRAD, H. Asymptotic equivalence of the Navier-Stokes and nonlinear Boltzmann equations. In *Proc. Sympos. Appl. Math., Vol. XVII*. Amer. Math. Soc., Providence, R.I., 1965, pp. 154–183.
- [21] GUALDANI, M. P., MISCHLER, S., AND MOUHOT, C. Factorization for non-symmetric operators and exponential H-theorem.
- [22] GUO, Y. Classical solutions to the Boltzmann equation for molecules with an angular cutoff. *Arch. Ration. Mech. Anal.* 169, 4 (2003), 305–353.
- [23] GUO, Y. Boltzmann diffusive limit beyond the Navier-Stokes approximation. *Comm. Pure Appl. Math.* 59, 5 (2006), 626–687.
- [24] KATO, T. *Perturbation theory for linear operators*. Classics in Mathematics. Springer-Verlag, Berlin, 1995. Reprint of the 1980 edition.
- [25] LANFORD, III, O. E. Time evolution of large classical systems. In *Dynamical systems, theory and applications (Recontres, Battelle Res. Inst., Seattle, Wash., 1974)*. Springer, Berlin, 1975, pp. 1–111. Lecture Notes in Phys., Vol. 38.
- [26] LERAY, J. Sur le mouvement d'un liquide visqueux emplissant l'espace. *Acta Math.* 63, 1 (1934), 193–248.
- [27] MOUHOT, C. Explicit coercivity estimates for the linearized Boltzmann and Landau operators. *Comm. Partial Differential Equations* 31, 7-9 (2006), 1321–1348.
- [28] MOUHOT, C. Rate of convergence to equilibrium for the spatially homogeneous Boltzmann equation with hard potentials. *Communications in Mathematical Physics* 261 (2006), 629–672.
- [29] MOUHOT, C., AND NEUMANN, L. Quantitative perturbative study of convergence to equilibrium for collisional kinetic models in the torus. *Nonlinearity* 19, 4 (2006), 969–998.
- [30] MOUHOT, C., AND VILLANI, C. Regularity theory for the spatially homogeneous Boltzmann equation with cut-off. *Arch. Ration. Mech. Anal.* 173, 2 (2004), 169–212.
- [31] SAINT-RAYMOND, L. *Hydrodynamic limits of the Boltzmann equation*, vol. 1971 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 2009.
- [32] UKAI, S. On the existence of global solutions of mixed problem for non-linear Boltzmann equation. *Proc. Japan Acad.* 50 (1974), 179–184.
- [33] UKAI, S., AND YANG, T. Mathematical theory of the Boltzmann equation. Lecture Notes Series, no. 8, Liu Bie Ju Centre for Mathematical Sciences, City University of Hong Kong, 2006.
- [34] VILLANI, C. Limites hydrodynamiques de l'équation de Boltzmann (d'après C. Bardos, F. Golse, C. D. Levermore, P.-L. Lions, N. Masmoudi, L. Saint-Raymond). *Astérisque*, 282 (2002), Exp. No. 893, ix, 365–405. Séminaire Bourbaki, Vol. 2000/2001.
- [35] VILLANI, C. A review of mathematical topics in collisional kinetic theory. In *Handbook of mathematical fluid dynamics, Vol. I*. North-Holland, Amsterdam, 2002, pp. 71–305.

- [36] YU, H. Global classical solutions of the Boltzmann equation near Maxwellians. *Acta Math. Sci. Ser. B Engl. Ed.* 26, 3 (2006), 491–501.

MARC BRIANT

BROWN UNIVERSITY
DIVISION OF APPLIED MATHEMATICS
182 GEORGE STREET, BOX F PROVIDENCE, RI 02192, USA

E-MAIL: briant.maths@gmail.com

SARA MERINO-ACEITUNO

UNIVERSITY OF CAMBRIDGE
DPMMS, CENTRE FOR MATHEMATICAL SCIENCES
WILBERFORCE ROAD, CAMBRIDGE CB3 0WA, UK

E-MAIL: s.merino-aceituno@maths.cam.ac.uk

CLÉMENT MOUHOT

UNIVERSITY OF CAMBRIDGE
DPMMS, CENTRE FOR MATHEMATICAL SCIENCES
WILBERFORCE ROAD, CAMBRIDGE CB3 0WA, UK

E-MAIL: C.Mouhot@dpms.cam.ac.uk