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ACS APPLIED **ELECTRONIC MATERIALS**

1 Huge Reduction of the Wake-Up Effect in Ferroelectric HZO Thin ² Films

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8 **S** Supporting Information

9 **ABSTRACT:** The wake-up effect is a major issue for ferroelectric HfO_2 -based
10 **memory devices.** Here, two $TiN/HZO/TiN$ structures deposited by magnetron memory devices. Here, two TiN/HZO/TiN structures deposited by magnetron

¹¹ sputtering on silicon are compared. The maximum remanant polarization is

12 higher than 21 μ C/cm² for both samples, but a strong difference is observed in

¹³ the electrical behavior. For the mesa sample, the difference between the

14 maximum and initial remanent polarization is only 3 μ C/cm², whereas it is

15 around 14 μ C/cm² in the non-mesa case. We discuss the root causes of these

¹⁶ behaviors in light of GIXRD results.

Letter

Letter

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17 **KEYWORDS:** HfO₂, ZrO₂, HZO, ferroelectricity, wake-up effect, FRAM

¹⁸ The discovery of ferroelectric Si-doped HfO₂¹ has paved

¹⁹ the way to a wide variety of research interests and

an technical dayslopments. The growing number of papers \blacktriangle the way to a wide variety of research interests and technical developments. The growing number of papers 21 dealing with HfO_2 -based materials² is a relevant indicator of the growing interest for this subject. Undeniably, many 23 applications could benefit from ferroelectric HfO_2 -based materials such as ferroelectric random access memory (FRAM), a ferroelectric field effect transistor (FeFET), a negative capacitance field effect transistor (NCFET), a ferroelectric tunnel junction (FTJ), synaptic devices, pyro- electric energy harvesters, piezoelectric devices, and so on. However, major issues are still hindering the appearance of ferroelectric HfO₂-based materials on the market. These issues are mainly related to the well-known wake-up effect, imprint, fatigue, and endurance. Recent studies on the subject point out crystallographic phase change and oxygen vacancy diffusion during cycling as the root causes of the issues cited above.

³⁵ Of all the numerous dopants that have been tested in 36 ferroelectric HfO_2 -based materials, the most addressed is 37 probably Zr^{4+} , because both $HfO₂$ and $ZrO₂$ are already used 38 in the \dot{CMOS}^3 and DRAM⁴ industries. Moreover, Zr^{4+} allows ³⁹ the lowest annealing temperature of all dopants, making it ⁴⁰ compatible with the silicon semiconductor industry.

 Compared to the hundreds of articles dealing with (Hf, $Zr)O_2$ solid solutions (HZO) deposited by ALD, sputtering 43 has rarely been used.^{5−14} However, the deposition of materials by sputtering shows many benefits compared to ALD:

- (ii) Sputtering can prevent carbon contamination that ⁴⁶ cannot be avoided with ALD precursors.
- (iii) The deposition can be performed at room temperature. ⁴⁸
- (iv) In principle, sputtering can be carried out in one run in ⁴⁹ the same sputtering chamber. $\frac{1}{50}$
- (v) Sputtering is a low-cost reproducible process. 51

Consequently, HZO development could be greatly ⁵² enhanced by using sputtering.

Moreover, most of the articles dealing with HZO sputtering ⁵⁴ use cosputtering.^{7-9,11,13,14} Cosputtering enables more param-ss eters to be tuned, particularly the Zr doping.¹¹ However, for 56 industrial manufacturing, single target sputtering appears to be 57 simpler and less expensive. 58

In this work, two different structures were realized by ⁵⁹ magnetron sputtering with a single target. Their performances ⁶⁰ are among the best in sputtering, but their behaviors are ⁶¹ completely different. Particularly, the mesa structure presents a ⁶² significantly reduced wake-up effect. We discuss the origins of ⁶³ these differences and the consequences for electrical measure- ⁶⁴ ment methods used to characterize ferroelectric $HfO₂$ -based 65 materials. 66

The structure shown in Figure 1a is called NM-sample for 67 f1 "non-mesa" structure and the structure shown in Figure 1b is ⁶⁸ called M-sample for "mesa" structure. 69

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 Si substrates were cleaned with ultrasonication in acetone and ethanol for 1 min. Substrates are Si-doped and (100)- oriented. The native oxide layer of each substrate is removed by a buffer oxide etch (BOE) process. For both processes, a 50 nm-thick bottom electrode (BE) is deposited in the conditions t1 75 detailed in Table 1. This layer blankets the entire Si substrate.

Table 1. Film Growth Conditions for TiN and HZO

	sputtering	
target-substrate distance	8 cm	
base pressure	$<$ 5 \times 1 ⁻⁷ mbar	
deposited elements	TiN	(Hf, Zr)O ₂
substrate	Si	Si/TiN
target	Ti	ZrO_2/HfO_2 (50/50)
target RF power	300 W	100 W
holder DC bias voltage	60 V	none
gas	$Ar = 50$ sccm	$Ar = 50$ sccm
	$N_2 = 3$ sccm	
working pressure	5×10^{-3} mbar	5×10^{-2} mbar
deposition speed	5.2 nm/min	5 nm/min
rapid thermal annealing (RTA) conditions		
temperature	450 \degree C	
atmosphere	N,	
duration	30 _s	

 Then, an HZO layer covering the whole TiN layer is deposited. The HZO thickness is measured by X-ray reflectivity (XRR). HZO layers are 12.3 nm-thick and 11.0 nm-thick for the M-sample and the NM-sample, respectively. For the NM-sample, UV photolithography is carried out. Then, a 50 nm-thick TiN TE followed by a 100 nm Pt top electrode for mechanical and electrical contact are deposited. A lift-off step is then performed. Finally, the NM-sample is annealed by RTA under the conditions described in Table 1. It means that, for the NM-sample, HZO is only partially covered during the RTA crystallization process.

 For the M-sample, the HZO layer is blanketed with a 50 nm- thick TiN layer. Then, UV photolithography and lift-off are carried out to define the electrodes with the deposition of a 100 nm-thick Pt layer followed by a 50 nm-thick Cr layer. Cr is used as a hard mask for etching. The M-sample is annealed by

RTA in the conditions described in Table 1. Finally, the ⁹² sample is etched by RIE at 300 W (RF) for 20 min with $SF₆$ as 93 a reactive gas. The final structure is presented in Figure 1b. For ⁹⁴ this structure, HZO is fully covered with TiN TE during the ⁹⁵ annealing step. 96

GIXRD is performed on a four-circle Smartlab Rigaku ⁹⁷ diffractometer using a 9 kW copper rotating anode. 98

Electrical characterization is carried out on 20 μ m diameter 99 capacitors using a probe station and a setup composed of an ¹⁰⁰ arbitrary waveform generator (Nippon Factory WF1966), a ¹⁰¹ transimpedance current amplifier (Keithley 428), and a 12 bit ¹⁰² Digital Storage Oscilloscope (Nicolet Integra 40). ¹⁰³

Endurance tests were performed with bipolar square pulses ¹⁰⁴ (commonly called set/reset sequence) until breakdown. The ¹⁰⁵ set/reset sequence is achieved in two different configurations: ¹⁰⁶ 2.5 kHz/3.5 V and 1.0 MHz/3.0 V. The first sequence allows ¹⁰⁷ the stress to be maximized without breakdown. The second ¹⁰⁸ set/reset sequence is used in order to be as close as possible to 109 the microelectronic standard, and it is therefore used as a ¹¹⁰ reference.

P−E curves are obtained using the PUND technique. This 112 consists of applying a negative set pulse followed by two ¹¹³ positive pulses (P and U) and then two negative pulses (N and ¹¹⁴ D). The PUND maximum amplitude voltage equals that of the 115 set/reset sequence. PUND pulses are triangular pulses with a ¹¹⁶ rise/fall time of 100 μ s, corresponding to 2.5 kHz 117 $\left(= \frac{1}{4 \times 100} \mu s^{-1} \right)$.

First, the change of the sum of the absolute values of the 118 positive remanent polarization and the negative remanent ¹¹⁹ polarization $|P_{r+}| + |P_{r-}|$ during cycling is studied. Figure 2 120 f2 shows the results for the two different set/reset sequences ¹²¹ described in the experimental section. 122

At 1.0 MHz and 3.0 V (Figure 2a), $|P_{r+}| + |P_{r-}|$ for the NM- 123 sample starts at around 5 μ C/cm². This value increases until 124 10^5 cycles, where it reaches around 29 μ C/cm². Then, it 125 decreases until breakdown. Such an effect is not observable for ¹²⁶ the M-sample. Indeed, $|P_{r+}| + |P_{r-}|$ starts at 25 μ C/cm² (instead 127 of 5 μ C/cm² for the NM-sample) and reaches a maximum of 128 32.8 μ C/cm² at 2 × 10³ cycles. So, the wake-up effect can be 129 considered negligible in these conditions. Subsequently, fatigue ¹³⁰

Figure 2. Change of $|P_{r+}| + |P_{r-}|$ versus the number of cycles with the stress conditions (a) 1.0 MHz and 3.0 V and (b) 2.5 kHz and 3.5 V for the NM-sample (black curve) and M-sample (red curve).

Figure 3. Polarization versus applied electrical field for (a) the NM-sample at 10⁵ cycles and (b) the M-sample at 2×10^3 cycles under the stress conditions: 2.5 kHz and 3.5 V.

131 begins and breakdown occurs after 5×10^5 cycles. In 132 comparison, the NM-sample breaks down after 1×10^7 cycles. 133 More severe stress conditions were also applied (Figure 2b). The same tendencies are observed: here, the wake-up effect is also significantly reduced on the M-sample compared to that on the NM-sample, but it is detrimental for the endurance of the M-sample. In fact, the breakdown occurs at a lower number of cycles for the M-sample compared to for the NM-sample.

139 For both samples, $|P_{r+}| + |P_{r-}|$ values are much higher at 2.5 ¹⁴⁰ kHz/3.5 V and the breakdown occurs earlier. These results are 141 consistent with previous observations.^{15−17}

142 For both stress conditions, the P_r of the NM-sample increases until it exceeds the P_r of the M-sample for a given number of cycles, because fatigue occurs earlier for the Mf3 145 sample. As a consequence, at 1.0 MHz/3.0 V (Figure 3a), between the first cycle and the breakdown, the M-sample has a 147 higher maximum $P_r = 16.4 \mu C/cm^2$ than the NM-sample, whereas at 2.5 kHz/3.5 V (Figure 2b), the M-sample has a 149 higher maximum $P_r = 21.7 \mu C/cm^2$. This is due to the reduced wake-up effect. In fact, with the M-sample, the maximum remanent polarization is always reached very rapidly, whereas with the NM-sample, the higher the stress conditions, the faster you can reach the maximum remanent polarization. Moreover, for the NM-sample, fatigue is always reached at ¹⁵⁴ around 10^6 cycles, which means that under a low-stress 155 sequence, fatigue is reached on the NM-sample before crossing 156 the curve of the M-sample. It follows that the M-sample has a 157 higher remanent polarization under low-stress conditions, ¹⁵⁸ whereas the NM-sample has a higher maximum remanent ¹⁵⁹ polarization under high-stress conditions. Therefore, the ¹⁶⁰ maximum P_r seems to be more dependent on the stress 161 conditions than on the structures themselves, which could have ¹⁶² important consequences on industrial applications. ¹⁶³

The M-sample has a reduced wake-up effect, but it also ¹⁶⁴ shows a sign of fatigue before the NM-sample. It is possible ¹⁶⁵ that the M-sample is already in an awake state, which leads to ¹⁶⁶ the absence of wake-up and an earlier fatigue and breakdown 167 compared to the NM-sample. 168

For the M-sample, the maximum P_r minus the initial P_r is 169 around 3–4 μ C/cm² in both conditions. In comparison, the 170 difference is between 12 and 16 μ C/cm² for the NM-sample. 171 There is 1 order of magnitude between both increases of P_r 172 during the wake-up.

In high-stress conditions, there is no double peak for the M- ¹⁷⁴ sample, and in low-stress conditions, the double peak ¹⁷⁵ disappears at the second measurement cycle. For the NM- ¹⁷⁶

Figure 4. GIXRD measurements comparing the NM-sample (black curve) and the M-sample (red curve) (a) before RTA and (b) after RTA. M stands for monoclinic, O stands for orthorhombic, and T stands for tetragonal phases.

 sample, a double peak is clearly observable during the first measurement cycle, and a small peak (coming from the double peak) is observable until $10³$ cycles for high-stress conditions 180 and until 5×10^4 cycles for low-stress conditions (see the [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acsaelm.9b00367/suppl_file/el9b00367_si_001.pdf) for more details).

 Finally, at 10 cycles, an inflection point can be seen. It is caused by the measurements realized due to the PUND technique, which also wakes the sample. The PUND measurement is negligible after 20 cycles, as compared to the set/reset cycles.

¹⁸⁷ It is to be noted that these results have been reproduced on ¹⁸⁸ many electrodes.

 To assess the quality of our films, Figure 3 shows the 190 polarization versus the electrical field when P_r has reached its maximum value. PUND and PN curves are represented. The 192 value of P_r is higher in PN curves, because PN includes the 193 leakage current. P_r values of the PN measurements have not been discussed in this report, because they can include extrinsic effects. Leakage is very well-compensated, because the PN curves show open cycles, whereas PUND does not. Also, the P−E curves by PUND show very pinched cycles, which means that the samples are very well-saturated.

 Then, grazing incidence X-ray diffraction is carried out to f4 200 enlighten the electrical characterizations. Figure 4 shows GIXRD results before (Figure 4a) and after (Figure 4b) annealing of the NM-sample and the M-sample. ICDD card nos. 00-038-1420 and 00-004-0802 were respectively used for TiN and Pt. For the HZO m-phase, ICDD card no. 04-002- 5428 was used. The peak at around 30.5° for the o/t-phase 206 came from the work of Lee et al., 8 and we knew from the first paper on ferroelectricity in hafnium oxide by Böscke et al.¹ that a peak at around 25° can be mistaken with the peak coming from the m-phase at this angle. In Figure 4a, the M-sample presents a TiN(111) peak, which is not visible in the NM- sample. This is probably caused by the different structures of the TiN TE in each sample. Comparing the NM-sample and the M-sample in Figure 4b, it is clear that the polycrystalline nature of the samples is completely different. The NM-sample shows clear monoclinic phases (m-phases), whereas no monoclinic peak can be seen in the M-sample. On the M-

sample, only the $o(111)$ peak can be seen. Other peaks come 217 from the electrodes. 218

At the microcrystalline scale, the orthorhombic phase (o- ²¹⁹ phase) is recognized as the origin of ferroelectricity. m-Phases ²²⁰ are assumed to be parasitic phases for ferroelectricity. Two ²²¹ hypotheses are generally considered as the causes of the wake- ²²² up effect according to Fengler et al:¹⁸ crystallographic phase 223 change contributions and oxygen vacancy diffusion. Grimley et ²²⁴ al.¹⁹ have observed a phase transition from the m- to the o - 225 phase using scanning transmission electron microscopy ²²⁶ (STEM) and impedance spectroscopy. Therefore, tremendous ²²⁷ reduction in the amount of m-phase in the M-sample could ²²⁸ have led to a drastic decrease in the wake-up effect. Moreover, ²²⁹ the presence of many monoclinic orientations coexisting with ²³⁰ the o-phase on the NM-sample could also play a key role in the ²³¹ oxygen vacancy redistribution, because the number of ²³² interfaces between crystallites has increased on the NM- ²³³ sample. In fact, the crystalline structure influences the ²³⁴ redistribution of charged defects, the mobility of these charges, ²³⁵ charge trapping, and charge screening. For that reason, it can ²³⁶ be concluded that the microcrystalline structure seems to be ²³⁷ the main parameter influencing the two different electrical ²³⁸ behaviors observed in this Letter. 239

One could argue that the different crystalline structures are ²⁴⁰ due to the TiN TE difference between the two samples. In fact, ²⁴¹ in the case of the M-sample, TiN is blanketed over the entire ²⁴² HZO layer, whereas in the case of the NM-sample, it covers ²⁴³ only the areas under the top electrodes. This implies a ²⁴⁴ completely different stress state between the two samples. As a ²⁴⁵ consequence, in the case of the M-sample, it turns the o-phase ²⁴⁶ into the lowest energy phase. It has actually been demonstrated ²⁴⁷ that TiN TE acts as a tensile stressor on the HZO film during ²⁴⁸ the annealing process.^{20−24} The stress inhibits the formation of $_{249}$ the monoclinic phase during HZO crystallization, forming an ²⁵⁰ orthorhombic phase that generates a large ferroelectric ²⁵¹ polarization, even at low process temperatures. ²⁵²

Finally, Chernikova et al.²⁵ compare the P_r performances of 253 an HZO layer and a La-doped HZO layer called "HZLO". The ²⁵⁴ HZO layer has a very low wake-up effect but also a very low ²⁵⁵ endurance, whereas HZLO has a very high endurance with a ²⁵⁶

 very high wake-up effect. The exact same phenomenon is observed here when measuring the NM-sample and the M- sample. Consequently, the root causes of the huge reduction of the wake-up effect and the decreasing of cycles before breakdown in both articles could be assumed to be the same. However, they observed that "the HZLO film showed a slightly larger grain size compared to the HZO film". Consequently, according the authors of the article, "the bulk free energy difference between the FE o-phase and non-FE m- phase decreased after La doping." Unfortunately, this cannot explain what is happening in our work, because the observation of clear monoclinic peaks in the GIXRD pattern of the NM- sample Figure 4b suggests that the free energy difference decreases. Manufacturing both structures with HZLO instead of HZO may help to shed light on this mystery.

 This Letter reports the fabrication of two samples deposited by magnetron sputtering with excellent performances, quite 274 similar to samples deposited by ALD.²⁵ \tilde{P}_r values are among the highest for samples deposited by sputtering.

 Although the M-sample and the NM-sample show very similar maximum P_r values, both samples demonstrate completely different electrical behavior. During cycling, the 279 increase of the P_r value for the NM-sample is more than an order of magnitude higher than that of the M-sample. It is accompanied by a decrease of the endurance, which is 2 orders of magnitude higher for the NM-sample than for the M-²⁸³ sample.

²⁸⁴ As the electrical behaviors are not the same, for low-stress 285 conditions, the M-sample has a higher maximum P_r value ²⁸⁶ during cycling, whereas for high-stress conditions, the NM-287 sample has a higher P_r value during cycling. As a matter of fact, 288 it has been proven that maximum P_r values are more sensitive ²⁸⁹ to stress conditions than the to structures themselves.

 The origins of the different electrical behaviors come from the microcrystalline structures of the two samples, according to GIXRD results. The crystallization takes place during the annealing step. During annealing, the M-sample is built with a TiN TE fully covering the HZO layer, whereas the TiN only partially covers the HZO layer in the case of the NM-sample. It induces different stress states, which lead to two different microcrystalline patterning. The M-sample shows no mono- clinic peak, whereas the NM-sample shows many monoclinic orientations. This explains the huge reduction of the wake-up ³⁰⁰ effect.

³⁰¹ ■ ASSOCIATED CONTENT

302 Supporting Information

³⁰³ The Supporting Information is available free of charge on the ³⁰⁴ [ACS Publications website](http://pubs.acs.org) at DOI: [10.1021/acsaelm.9b00367.](http://pubs.acs.org/doi/abs/10.1021/acsaelm.9b00367)

 Chronograms for the M-sample and NM-sample describing the presence, the absence, and the mergers of the double peak obtained due to the PUND measurements ([PDF](http://pubs.acs.org/doi/suppl/10.1021/acsaelm.9b00367/suppl_file/el9b00367_si_001.pdf))

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314 Notes

³¹⁵ The authors declare no competing financial interest.

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R REFERENCES 325

(1) Böscke, T. S.; Müller, J.; Bräuhaus, D.; Schröder, U.; Böttger, U. 326 Ferroelectricity in Hafnium Oxide Thin Films. Appl. Phys. Lett. 2011, 327 99, 102903.

(2) Park, M. H.; Lee, Y. H.; Mikolajick, T.; Schroeder, U.; Hwang, 329 C. S. Review and Perspective on Ferroelectric HfO_2 -based Thin Films 330 for Memory Applications. MRC 2018, 8, 795−808. 331

(3) Kim, K. From the Future Si Technology Perspective: Challenges 332 and Opportunities. 2010 IEEE International Electron Devices Meeting 333 (IEDM) 2010, 1.1.1−1.1.9. 334

(4) James, D. Recent Innovations in DRAM Manufacturing. 2010 335 21st Annual IEEE/SEMI Advanced Semiconductor Manufacturing 336 Conference (ASMC) 2010, 264–269. 337

(5) Kiguchi, T.; Nakamura, S.; Akama, A.; Shiraishi, T.; Konno, T. J. 338 Solid State Epitaxy of $(Hf,Zr)O₂$ Thin Films with Orthorhombic 339 Phase. J. Ceram. Soc. Jpn. 2016, 124, 689-693. 340

(6) Fan, Z.; Xiao, J.; Wang, J.; Zhang, L.; Deng, J.; Liu, Z.; Dong, Z.; 341 Wang, J.; Chen, J. Ferroelectricity and Ferroelectric Resistive 342 Switching in Sputtered $Hf_{0.5}Zr_{0.5}O_2$ Thin Films. Appl. Phys. Lett. 343
2016. 108. 232905. 2016, 108, 232905.

(7) Ambriz-Vargas, F.; Kolhatkar, G.; Broyer, M.; Hadj-Youssef, A.; 345 Nouar, R.; Sarkissian, A.; Thomas, R.; Gomez-Yáñez, C.; Gauthier, M. 346 A.; Ruediger, A. A Complementary Metal Oxide Semiconductor 347 Process-Compatible Ferroelectric Tunnel Junction. ACS Appl. Mater. 348 Interfaces 2017, 9, 13262−13268. 349

(8) Lee, Y. H.; Kim, H. J.; Moon, T.; Kim, K. D.; Hyun, S. D.; Park, 350 H. W.; Lee, Y. B.; Park, M. H.; Hwang, C. S. Preparation and 351 Characterization of Ferroelectric Hf_0 ₅ Zr_0 ₅ O_2 Thin Films Grown by 352 Reactive Sputtering. Nanotechnology 2017, 28, 305703.

(9) Das, K.; Tripathy, N.; Ghosh, S.; Sharma, P.; Singhal, R.; Kar, J. 354 Microstructural, Surface and Interface Properties of Zirconium Doped 355 $HfO₂$ Thin Films Grown by RF Co-sputtering Technique. Vacuum 356 2017, 143, 288−293. ³⁵⁷

(10) Ambriz-Vargas, F.; Kolhatkar, G.; Thomas, R.; Nouar, R.; 358 Sarkissian, A.; Gomez Yáñez, C.; Gauthier, M. A.; Ruediger, A. 359 Tunneling Electroresistance Effect in a $Pt/Hf_{0.5}Zr_{0.5}O_2/Pt$ Structure. 360 Appl. Phys. Lett. 2017, 110, No. 093106. 361

(11) Migita, S.; Ota, H.; Yamada, H.; Shibuya, K.; Sawa, A.; 362 Toriumi, A. Polarization Switching Behavior of Hf Zr O Ferroelectric 363 Ultrathin films Studied Through Coer cive Field Characteristics. Jpn. 364 J. Appl. Phys. 2018, 57, No. 04FB01. 365

(12) Bouaziz, J.; Rojo Romeo, P.; Baboux, N.; Vilquin, B. 366 Characterization of Ferroelectric Hafnium/Zirconium Oxide Solid 367 Solutions Deposited by Reactive Magnetron Sputtering. J. Vac. Sci. 368 Technol., B: Nanotechnol. Microelectron.: Mater., Process., Meas., 369 Phenom. 2019, 37, No. 021203. 370

(13) Migita, S.; Ota, H.; Shibuya, K.; Yamada, H.; Sawa, A.; 371 Matsukawa, T.; Toriumi, A. Phase Transformation Behavior of 372 Ultrathin $Hf_{0.5}Zr_{0.5}O_2$ Films Investigated Through Wide Range 373 Annealing Experiments. Jpn. J. Appl. Phys. 2019, 58, SBBA07. 374 (14) Luo, Q.; Ma, H.; Su, H.; Xue, K.-H.; Cao, R.; Gao, Z.; Yu, J.; 375 Gong, T.; Xu, X.; Yin, J.; Yuan, P.; Tai, L.; Dong, D.; Long, S.; Liu, 376 Q.; Miao, X.-S.; Lv, H.; Liu, M. Composition-Dependent Ferroelectric 377 Properties in Sputtered Hf_xZr_{1−X}O₂ Thin Films. IEEE Electron Device 378 Lett. 2019, 40, 570−573. 379

 (15) Zhou, D.; Xu, J.; Li, Q.; Guan, Y.; Cao, F.; Dong, X.; Müller, J.; Schenk, T.; Schröder, U. Wake-up Effects in Si-doped Hafnium Oxide Ferroelectric Thin Films. Appl. Phys. Lett. 2013, 103, 192904.

- (16) Starschich, S.; Menzel, S.; Böttger, U. Pulse Wake-up and 384 Breakdown Investigation of Ferroelectric Yttrium doped HfO₂. J. Appl. Phys. 2017, 121, 154102.
- (17) Starschich, S.; Menzel, S.; Böttger, U. Evidence for Oxygen Vacancies Movement During Wake-up in Ferroelectric Hafnium Oxide. Appl. Phys. Lett. 2016, 108, 032903.
- (18) Fengler, F. P.; Hoffmann, M.; Slesazeck, S.; Mikolajick, T.; Schroeder, U. On The Relationship Between Field Cycling and
- 391 Imprint In Ferroelectric Hf_0 ₅ Zr_0 ₅ O_2 . J. Appl. Phys. 2018, 123, 204101. 392 (19) Grimley, E. D.; Schenk, T.; Sang, X.; Pešić, M.; Schroeder, U.; Mikolajick, T.; LeBeau, J. M. Structural Changes Underlying Field-
- 394 Cycling Phenomena in Ferroelec tric HfO₂ Thin Films. Advanced Electronic Materials 2016, 2, 1600173.
- (20) Park, M. H.; Kim, H. J.; Kim, Y. J.; Jeon, W.; Moon, T.; Hwang, C. S. Ferroelectric Properties and Switching Endurance of 398 $\text{Hf}_{0.5}\text{Zr}_{0.5}\text{O}_2$ Films on TiN Bottom and TiN or RuO₂ Top Electrodes. physica status solidi (RRL). Phys. Status Solidi RRL 2014, 8, 532−535. (21) Hyuk Park, M.; Joon Kim, H.; Jin Kim, Y.; Moon, T.; Seong Hwang, C. The Effects of Crystallographic Orientation and Strain of 402 Thin $\text{Hf}_{0.5}\text{Zr}_{0.5}\text{O}_2$ Film on its Ferroelectricity. Appl. Phys. Lett. 2014,
- 104, No. 072901.
- (22) Hyuk Park, M.; Joon Kim, H.; Jin Kim, Y.; Lee, W.; Kyeom Kim, H.; Seong Hwang, C. Effect of Forming Gas Annealing on the 406 Ferroelectric Properties of $\text{Hf}_{0.5}\text{Zr}_{0.5}\text{O}_2$ Thin Films with and without Pt Electrodes. Appl. Phys. Lett. 2013, 102, 112914.
- (23) Park, M. H.; Kim, H. J.; Kim, Y. J.; Lee, W.; Moon, T.; Kim, K. D.; Hwang, C. S. Study on the Degradation Mechanism of the 410 Ferroelectric Properties of Thin Hf0.5 $Zr_{0.5}O_2$ Films on TiN and Ir Electrodes. Appl. Phys. Lett. 2014, 105, 072902.
- (24) Kim, S. J.; Narayan, D.; Lee, J.-G.; Mohan, J.; Lee, J. S.; Lee, J.; Kim, H. S.; Byun, Y. c.; Lucero, A. T.; Young, C. D.; Summerfelt, S. R.; San, T.; Colombo, L.; Kim, J. Large Ferroelectric Polarization of 415 TiN/Hf_{0.5}Zr0.5O₂/TiN Capacitors Due to Stress-induced Crystal- lization at Low Thermal Budget. Appl. Phys. Lett. 2017, 111, 242901. (25) Chernikova, A. G.; Kozodaev, M. G.; Negrov, D. V.; Korostylev, E. V.; Park, M. H.; Schroeder, U.; Hwang, C. S.; Markeev, A. M. Improved Ferroelectric Switching Endurance of La-420 Doped $\text{Hf}_{0.5}\text{Zr}_{0.5}\text{O}_2$ Thin Films. ACS Appl. Mater. Interfaces 2018, 10, 2701−2708.