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Adaptive shunted piezoelectric metacomposite: a new integrated technology for vibroacoustic control

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Abstract—In this paper, we present an application of the Floquet-Bloch theorem in the context of electrodynamics for vibroacoustic power flow optimization by mean of distributed and shunted piezoelectric patches. The main purpose of this work is first to propose a dedicated numerical approach able to compute the multi-modal wave dispersions curves into the whole first Brillouin zone for periodically distributed 2D shunted piezomechanical systems. By using two specific indicators evaluating the evanescent part of Bloch's waves and the induced electronic damping, we optimize the piezoelectric shunting electrical impedance for controlling energy diffusion into the proposed semi-active distributed set of cells. Sound radiation efficiency is also analyzed for showing the effects of such smart metamaterial for controlling acoustical noise.

Index Terms—Distributed Active Noise Control, Adaptive Metacomposite, Distributed shunted piezoelectric patches

I. INTRODUCTION

Tailoring the dynamical behavior of wave-guide structures can provide an efficient and physically elegant means to optimize mechanical components with regards to vibration and acoustic criteria, among others. However, achieving this objective may lead to different outcomes depending on the context of the optimization. In the preliminary stages of a product's development, one mainly needs optimization tools capable of rapidly providing global design direction. Such optimization will also depend on the frequency range of interest. One usually discriminates between the low frequency (LF) range and the medium frequency (MF) range, especially if vibration and noise are considered. However, it should be noted that LF optimization of vibration is more common in the literature than MF optimization. For example, piezoelectric materials and other adaptive and smart systems are employed to improve the vibroacoustic quality of structural components, especially in the LF range [1], [2], [3]. Recently, much effort has been spent on developing new multi-functional structures integrating electro-mechanical systems in order to optimize their vibroacoustic behavior over a larger frequency band of interest [4], [5], [6], [7], [8], [9], [10], [11], [12]. However, there is still a lack of studies in the literature for MF optimization of structural vibration. To that end, the focus of this study is to provide a suitable numerical tool for computing wave dispersion in 2D periodic systems incorporating controlling electronics devices. The final aim is to allow their optimization

in order to optimize vibroacoustic diffusion in 2D wave guides and analyze its effect on acoustic radiation.

For optimizing wave dispersion we need to use efficient numerical tools. Two approaches can be distinguished for computing that dispersion: the semi-analytical finite element method (SAFE) and the wave finite element (WFE) method. In the former approach, the displacement field is modeled exactly in the direction of wave propagation by using a harmonic function and approximately in the perpendicular directions by using finite elements (FE). An eigenvalue problem is then formulated by introducing the displacement field into the governing equations. Solving the eigenvalue problem for a given frequency gives the wave numbers of all the propagating modes. The main disadvantage of the SAFE method is that FE used are not standard so they must be specifically defined for each application. Moreover the wave guide needs to be homogeneous. Nevertheless, a large amount of FE has been developed since 1975 to compute dispersion curves of rails [13], stiffened cylinders [14], laminated composite plates [15], [16] and viscoelastic laminated composite plates [17]. To avoid development of specific FE, the WFE method considers the structures as periodic in order to model a single period of the structure with standard FE. By using the periodic structure theory (PST) introduced by Mead [13], an eigenvalue problem can be formulated from the stiffness and mass matrices of the FE model to find wave numbers of all the propagating waves. Contrary to the SAFE method, the displacement field is now approximated in the direction of propagation. The WFE method has been successfully used to deal with wave propagation in two dimensional structures [18], [19]. The approach has also been used for fluid-structure systems using the theory of group representation for computing two dimensional dispersion in undamped vibroacoustic systems [20], [21]. One of the main problem of all these approaches is the difficulty to compute the damped wave numbers in the whole Brillouin domain of multiphysics problem incorporating shunted piezoelectric patches necessary for optimizing vibroacoustic behavior of such periodic smart structures [22], [23].

After recalling the Floquet-Bloch theorems, we introduce a new numerical formulation for computing the multi-modal damped wave numbers dispersion in the whole first Brillouin domain of a periodical smart structure made of periodically

distributed shunted piezoelectric patches. Based on this wave modeling, optimization of the electrical impedance of the shunted circuit is performed in order to decrease the group velocity of flexural waves or to increase damping induced by the electric circuit. The obtained optimal impedances are also tested in controlling the MF response of a semi-distributed infinite system and its effect on sound radiation.

II. PIEZO-ELASTO-DYNAMICAL APPLICATION OF THE FLOQUET-BLOCH THEOREM

In this section the application of the celebrated Floquet-Bloch theorem is presented for piezo-elastodynamic problems. Based on the well known results obtained by Floquet [24] in one-dimensional and later rediscovered by Bloch [25] in multidimensional problems, we propose an original application to bi-dimensional piezo-elastodynamical problem leading to very general numerical implementation for computing waves dispersion for periodically smart distributed mechanical systems incorporating electronic components and damping effects [26].

A. The Bloch Theorem

The Bloch theorem gives the form of homogeneous states of Schrödinger equation with periodic potential. This theorem can be considered as a multidimensional application of the Floquet theorem [27]. The periodic medium properties satisfy a periodic condition as $M(\mathbf{x} + R \cdot \mathbf{m}) = M(\mathbf{x})$, $\mathbf{m} \in \mathbb{Z}^3$ where $R = [r_1, r_2, r_3] \in \mathbb{R}^{3 \times 3}$ is a matrix grouping the three lattice's basis vectors (in 3D). We can also define the primitive cell as a convex polyhedron of \mathbb{R}^3 called Ω_x . The reciprocal unit cell is denoted by Ω_k limited by the reciprocal lattice vector defined by the three vectors \mathbf{g}_j so that: $\mathbf{r}_i \cdot \mathbf{g}_j = 2\pi \delta_{ij}$ ($\delta_{i,j}$ being the Kronecker index). We note $G = [\mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3]$ the reciprocal lattice matrix in the later. If Ω_x is the irreducible primitive cell, Ω_k corresponds to the first Brillouin zone of the lattice. One can see [28] for details.

The Bloch Theorem stipulates that any functions $\mathbf{u}(\mathbf{x}) \in L^2(\mathbb{R}^3, \mathbb{C}^n)$ can be expressed as

$$\mathbf{u}(\mathbf{x}) = \int_{\Omega_k} e^{i\mathbf{k}\cdot\mathbf{x}} \tilde{\mathbf{u}}(\mathbf{x}, \mathbf{k}) d\mathbf{k} \quad (1)$$

where the Bloch amplitude $\tilde{\mathbf{u}}(\mathbf{x}, \mathbf{k})$ is Ω_x -periodic and has the representations

$$\begin{aligned} \tilde{\mathbf{u}}(\mathbf{x}, \mathbf{k}) &= \sum_{\mathbf{n} \in \mathbb{Z}^3} \hat{\mathbf{u}}(\mathbf{k} + G\mathbf{n}) e^{iG\mathbf{n}\cdot\mathbf{x}}, \\ \mathbf{u}(\mathbf{x}) &= \frac{|\Omega_x|}{(2\pi)^3} \sum_{\mathbf{n} \in \mathbb{Z}^3} \mathbf{u}(\mathbf{x} + R\mathbf{n}) e^{i\mathbf{k}(\mathbf{x} + R\mathbf{n})} \end{aligned} \quad (2)$$

where $\hat{\mathbf{u}}(\mathbf{k})$ stands for the Fourier transform of $\mathbf{u}(\mathbf{x})$. One can also demonstrate that the mean value of the Bloch amplitude is the Fourier amplitude of $\mathbf{u}(\mathbf{x})$ for the corresponding wave vector: $\langle \tilde{\mathbf{u}}(\cdot, \mathbf{k}) \rangle_{\Omega_x} = \hat{\mathbf{u}}(\mathbf{k})$. Using the Bloch theorem to represent the solutions of periodical partial derivative equations implies that all derivatives are shifted by \mathbf{k} in the sense given by the used spatial operator.

Based on that theorem one can define the expansion functions $\mathbf{v}_m(\mathbf{x}, \mathbf{k})$, called the Bloch eigen modes, such that they can be used to represent the Bloch amplitudes of any solution of the corresponding partial derivative equation as

$$\tilde{\mathbf{u}}(\mathbf{x}, \mathbf{k}) = \sum_m \mathbf{u}_m(\mathbf{k}) \mathbf{v}_m(\mathbf{x}, \mathbf{k}) \quad (3)$$

and at the same time diagonalize the partial derivative equations. One notes that the expansion coefficients $\mathbf{u}_m(\mathbf{k})$ depend on the applied disturbance and also on the induced wave vector (see [29] for details).

B. Application to Piezo-Elastodynamic

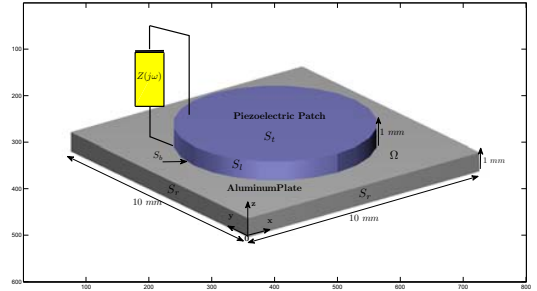


Fig. 1. Generic 3D piezocomposite periodic cell

Let us consider a piezo-elastodynamic problem made of infinite periodic distribution of unitary cell described in figure 1. The harmonic homogeneous dynamical equilibrium of system is driven by the following partial derivative equation:

$$\begin{cases} \rho \ddot{\mathbf{w}}(\mathbf{x}) - \nabla \sigma(\mathbf{x}) = 0 & \forall \mathbf{x} \in \Omega_x \\ -\nabla D(\mathbf{x}) = 0 & \forall \mathbf{x} \in \Omega_x \end{cases} \quad (4)$$

where $\mathbf{w}(\mathbf{x}) \in \mathbb{R}^3(\Omega_x)$ is the displacement vector, σ represents the Cauchy stress tensor, $\epsilon = \nabla_{sym} \mathbf{w} = \frac{1}{2}(\nabla \mathbf{w}^T(\mathbf{x}) + \mathbf{w}(\mathbf{x}) \nabla^T)$ the Green strain tensor, $D(\mathbf{x})$ the electric displacement. The linear constitutive material behavior relationships can be written as

$$\sigma = C_E(\mathbf{x}) \epsilon - e^T(\mathbf{x}) \mathbf{E} \quad (5)$$

$$\mathbf{D} = e(\mathbf{x}) \epsilon + \epsilon_S(\mathbf{x}) \mathbf{E} \quad (6)$$

where $\mathbf{E} = -\nabla V$ the electric field vector (V the voltage), C_E the elasticity tensor at constant electrical field, e^T the piezoelectric coupling tensor and ϵ_S the dielectric permittivity at constant strain. We add to this set of equilibrium equations an output expression

$$q^o = - \int_{S_t} \mathbf{D} \cdot \mathbf{n} dS \quad (7)$$

allowing the introduction of the charge measurement on the piezoelectric's top electrode and hence the dual counterpart of the imposed electrical Dirichlet boundary condition for applying the shunt impedance operator.

The equations above are consistent for each kind of material

to the extent that null piezoelectric and permittivity tensors can be used when passive materials are considered. All of these tensors also depend on the spatial location vector \mathbf{x} . The piezo-elastodynamic equilibrium can also be written as:

$$\forall \mathbf{x} \in \Omega_x \quad \rho \omega^2 \mathbf{w}(\mathbf{x}) + \nabla C \nabla_{sym}(\mathbf{w}(\mathbf{x})) + \nabla e^T(\mathbf{x}) \nabla V(\mathbf{x}) = 0 \quad (8)$$

$$- \nabla e(\mathbf{x}) \nabla_{sym}(\mathbf{w}(\mathbf{x})) + \nabla \varepsilon_S(\mathbf{x}) \nabla V(\mathbf{x}) = 0 \quad (9)$$

As the problem is 2D infinitely periodic, only electrostatic boundary conditions have to be considered on each cell:

$$\begin{cases} V = 0 & \forall \mathbf{x} \in S_b \\ V = V^o & \forall \mathbf{x} \in S_t \\ \mathbf{D} \cdot \mathbf{n} = 0 & \forall \mathbf{x} \in S_l \end{cases} \quad (10)$$

where S_b is the grounded bottom electrode of the piezoelectric layer, S_t is the top electrode connected to the external shunt and S_l the lateral electrode less boundary. The top electrode applied feedback voltage V^o depends on the shunt characteristic and on the collected charges q^o (7) and can be expressed in the Fourier space by:

$$V^o(i\omega) = -Z(i\omega)q^o(i\omega) \quad (11)$$

By considering a primitive cell of the periodic problem Ω_x and by using the Bloch theorem, we can compute the associated Bloch eigenmodes (3) and the dispersion functions by searching the eigen solutions of the homogeneous problem (8) and (9) as:

$$\mathbf{u}(\mathbf{x}) = \begin{bmatrix} \mathbf{w}(\mathbf{x}) \\ V(\mathbf{x}) \end{bmatrix} = \mathbf{u}_{n,\mathbf{k}}(\mathbf{x}) e^{i\mathbf{k} \cdot \mathbf{x}} \quad (12)$$

with $\mathbf{u}_{n,\mathbf{k}}(\mathbf{x}) = \begin{bmatrix} \mathbf{w}_{n,\mathbf{k}}(\mathbf{x}) \\ V_{n,\mathbf{k}}(\mathbf{x}) \end{bmatrix}$, Ω_x periodic functions. By introducing expression (12) in the piezo-elastodynamic equations (8), (9), one can demonstrate that $\mathbf{w}_{n,\mathbf{k}}(\mathbf{x})$, $V_{n,\mathbf{k}}(\mathbf{x})$ and $\omega_n(\mathbf{k})$ are solutions of the generalized eigenvalues problem:

$$\begin{aligned} 0 = & \rho \omega_n^2(\mathbf{k}) \mathbf{w}_{n,\mathbf{k}}(\mathbf{x}) + \nabla C \nabla_{sym}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x})) \\ & + ik \{ (C \nabla_{sym}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x}))) \cdot \Phi \\ & + \nabla (C \Xi_{n,\mathbf{k}}(\mathbf{x})) \} - k^2 (C \Xi_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \\ & + \nabla e^T \nabla V_{n,\mathbf{k}}(\mathbf{x}) + ik \{ (\nabla e^T V_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \\ & + (e^T \nabla V_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \} \\ & - k^2 V_{n,\mathbf{k}}(\mathbf{x}) (e^T \Phi) \cdot \Phi \quad \forall \mathbf{x} \in \Omega_x \end{aligned} \quad (13)$$

$$\begin{aligned} 0 = & -\nabla e \nabla_{sym}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x})) - ik \{ \nabla (e \Xi_{n,\mathbf{k}}(\mathbf{x})) \\ & + (e \nabla_{sym}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x}))) \cdot \Phi \} \\ & + k^2 (e \Xi_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \\ & + \nabla \varepsilon_S \nabla V_{n,\mathbf{k}}(\mathbf{x}) + ik \{ (\nabla \varepsilon_S V_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \\ & + (\varepsilon_S \nabla V_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \} \\ & - k^2 (\varepsilon_S \Phi V_{n,\mathbf{k}}(\mathbf{x})) \cdot \Phi \quad \forall \mathbf{x} \in \Omega_x \end{aligned} \quad (14)$$

with the associated boundary conditions:

$$\begin{cases} \mathbf{w}_{n,\mathbf{k}}(\mathbf{x} - R \cdot \mathbf{m}) = \mathbf{w}_{n,\mathbf{k}}(\mathbf{x}) & \forall \mathbf{x} \in S_r \quad \mathbf{m} \in \mathbb{Z}^2 \\ V_{n,\mathbf{k}}(\mathbf{x}) = 0 & \forall \mathbf{x} \in S_b \\ V_{n,\mathbf{k}}(\mathbf{x}) V = -Z(i\omega) q_{n,\mathbf{k}}^o & \forall \mathbf{x} \in S_t \\ \mathbf{D} \cdot \mathbf{n} = 0 & \forall \mathbf{x} \in S_l \end{cases}$$

where $\mathbf{k} = k \begin{bmatrix} \cos(\phi) \\ \sin(\phi) \\ 0 \end{bmatrix} = k \Phi$ where ϕ represents the direc-

tion angle into the reciprocal 2D lattice domain and $\Xi_{n,\mathbf{k}}(\mathbf{x}) = \frac{1}{2}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x}) \cdot \Phi^T + \Phi \cdot \mathbf{w}_{n,\mathbf{k}}^T(\mathbf{x}))$ the symmetric dyadic tensor or the dyadic product of the displacement $\mathbf{w}_{n,\mathbf{k}}(\mathbf{x})$ and direction vector Φ . S_r are the interface of the cells continuum, and R the matrix grouping the two lattice's basis vectors (in 2D in the considered problem). In the electrical boundary conditions, $q_{n,\mathbf{k}}^o$ is given by:

$$q_{n,\mathbf{k}}^o = \int_{S_t} [-e(\nabla_{sym}(\mathbf{w}_{n,\mathbf{k}}(\mathbf{x})) + ik \nabla e \Xi_{n,\mathbf{k}}(\mathbf{x})) + \varepsilon_S(\nabla V_{n,\mathbf{k}}(\mathbf{x}) + ik V_{n,\mathbf{k}}(\mathbf{x}) \Phi)] \cdot \mathbf{n} dS \quad (15)$$

where \mathbf{n} is the outpointing unitary normal vector.

The proposed formulation is also based on the computation of the Floquet vectors (13), (14), instead of computing the Floquet propagators commonly used for elastodynamic applications. Our approach allows to obtain the full 2D waves dispersions functions and to clearly introduce damping and electrical impedance into the piezo-elastodynamic operator. The adopted methodology allows the computation of the complete complex map of the dispersion curves including evanescent waves and allowing the introduction of damping and shunt operator if any.

C. Weak formulation and computation of waves dispersion functions in periodical piezo-composite lattice

Let us consider the partial derivative equations (13), (14) on a unit cell Ω . It stands for a generalized eigenvalue problem leading to the computation of the dispersion functions $\omega_n(\mathbf{k})$ and the corresponding Floquet eigenvectors $\mathbf{u}_{n,\mathbf{k}}(\mathbf{x})$. For computing the 2D dispersions curves, we need to introduce a suitable weak formulation.

If $\mathbf{u}_{n,\mathbf{k}}(\mathbf{x})$ is a solution of equations (13), (14), also $\forall \tilde{\mathbf{w}}_{n,\mathbf{k}}(\mathbf{x}) \in \{H_1(\Omega, \mathbb{C}^3) / \tilde{\mathbf{w}}_{n,\mathbf{k}}(\mathbf{x} - R\mathbf{m}) = \tilde{\mathbf{w}}_{n,\mathbf{k}}(\mathbf{x}) \quad \forall \mathbf{x} \in S_r\}$ and $\tilde{V}_{n,\mathbf{k}}(\mathbf{x}) \in \{H_1(\Omega, \mathbb{C}) / \tilde{V}_{n,\mathbf{k}}(\mathbf{x}) = 0 \quad \forall \mathbf{x} \in S_b$ and

$\tilde{V}_{n,k}(\mathbf{x}) = V^o \forall \mathbf{x} \in S_t$ } we have:

$$\begin{aligned}
0 = & \int_{\Omega} \rho \omega_n^2(\mathbf{k}) \tilde{\mathbf{w}}_{n,k}(\mathbf{x}) \mathbf{w}_{n,k}(\mathbf{x}) \\
& - (\tilde{\varepsilon}_{n,k}(\mathbf{x}) - ik \tilde{\Xi}_{n,k}(\mathbf{x})) C(\varepsilon_{n,k}(\mathbf{x}) + ik \Xi_{n,k}(\mathbf{x})) \\
& + (\tilde{\varepsilon}_{n,k}(\mathbf{x}) - ik \tilde{\Xi}_{n,k}(\mathbf{x})) e^T (\nabla V_{n,k}(\mathbf{x}) + ik V_{n,k}(\mathbf{x}) \Phi) \\
& - (\nabla \tilde{V}_{n,k}(\mathbf{x}) - ik \tilde{V}_{n,k}(\mathbf{x}) \Phi) e (\varepsilon_{n,k}(\mathbf{x}) + ik \Xi_{n,k}(\mathbf{x})) \\
& + (\nabla \tilde{V}_{n,k}(\mathbf{x}) - ik \tilde{V}_{n,k}(\mathbf{x}) \Phi) \varepsilon_S (\nabla V_{n,k}(\mathbf{x}) + ik V_{n,k}(\mathbf{x}) \Phi) d\Omega \\
& - \frac{\tilde{V}_{n,k}^t V_{n,k}^t}{Z(i\omega)}
\end{aligned} \tag{16}$$

This weak formulation is simply obtained by integrating equation (13), (14) projected onto any test function $\tilde{\mathbf{u}}_{n,k}(\mathbf{x})$. The boundary integral vanishes as the test functions are chosen so that $\tilde{\mathbf{w}}_{n,k}(\mathbf{x} - R\mathbf{m}) = \tilde{\mathbf{w}}_{n,k}(\mathbf{x})$ on S_r . For a polyhedron cell, each boundary is generally a polyhedral plane sub-domain that can be associated with its parallel opposite one. The symmetry conditions $\tilde{\mathbf{w}}_{n,k}(\mathbf{x} - R\mathbf{m}) = \tilde{\mathbf{w}}_{n,k}(\mathbf{x})$ explicitly link these associated surfaces.

D. Numerical Computation of the Bloch's waves

The numerical implementation is obtained by using a standard finite elements method to discretize the weak formulation (16). The assembled matrix equation is given by:

$$\begin{aligned}
0 = & (K(Z(i\omega_n(\lambda, \phi)) + \lambda L(\phi, Z(i\omega_n(\lambda, \phi)))) \\
& - \lambda^2 H(\phi, Z(i\omega_n(\lambda, \phi))) - \omega_n^2(\lambda, \phi) M) \mathbf{u}_{n,k}(\phi),
\end{aligned} \tag{17}$$

where $\lambda = ik$, M and $K(Z(i\omega_n(\lambda, \phi)))$ are respectively the standard symmetric semi-definite mass and stiffness matrices (the mass matrix is semi definite because elastostatic equation are condensed into the equation), $L(\phi, Z(i\omega_n(\lambda, \phi)))$ is a skew-symmetric matrix and $H(\phi, Z(i\omega_n(\lambda, \phi)))$ is a symmetric semi-definite positive matrix.

When k and ϕ are fixed and Z does not depend on ω the system (17) is a linear eigen value problem allowing us to compute the dispersion functions $\omega_n^2(\mathbf{k}, \phi)$ and the associated Bloch eigenvector $\mathbf{u}_{n,k}(\phi)$.

This approach has been widely used for developing homogenization techniques and spectral asymptotic analyses like in the work of [30]. It can also be applied for computing wave's dispersion even if Floquet propagators is preferred for 1D or quasi 1D computation, as indicated in [31], [32], [33]. Nevertheless these approaches have been only developed for undamped or lightly damped mechanical systems. In these cases, most of the previously published works present techniques based on the mesh of a real \mathbf{k} -space (i.e \mathbf{k} or λ and ϕ) following the boundary of the first Brillouin zone for obtaining the corresponding frequency dispersion and the associated Floquet vectors. For undamped system, only propagative or evanescent waves exist corresponding to a family of eigen solutions purely real or imaginary. Discrimination between each class of waves is easy. If a highly damped system (K, L, H are complex frequency dependent) and a frequency dependent electrical

shunt impedance are considered, the obtained eigenvalue problem is not quadratic and a complex specific numerical methodology has to be implemented. Furthermore, evanescent part of propagating waves appear as the imaginary part of eigenfrequencies. It then becomes very difficult to distinguish the propagative and evanescent waves as all solution appear complex.

Another much more suitable possibility for computing dispersion in damped systems, dedicated for time/space deconvolution and for computation of diffusion properties as defined by [4], [33], is to consider the following generalized eigen value problem:

$$\begin{aligned}
0 = & (K(Z(\omega) - \omega^2 M + \lambda_n(\omega, \phi) L(\phi, Z(\omega))) \\
& - \lambda_n^2(\omega, \phi) H(\phi, Z(\omega))) \mathbf{u}_n(\omega, \phi).
\end{aligned} \tag{18}$$

In this problem, the pulsation ω is a real parameter corresponding to the harmonic frequency. Wave's numbers and Floquet vectors are then computed. An inverse Fourier transformation in the k -space domain can lead us to evaluate the physical wave's displacements and energy diffusion operator when the periodic distribution is connected to another system as in [4]. Another temporal inverse Fourier transformation can furnish a way to access spatio-temporal response for non-homogeneous initial conditions. As L is skew-symmetric, the obtained eigen values are quadruple $(\lambda, \bar{\lambda}, -\lambda, -\bar{\lambda})$ collapsing into real or imaginary pairs (or a single zero) when all matrices are real (i.e. for an undamped system). In this case a real pair of eigen values correspond to evanescent modes oriented in two opposite directions on the k -space and imaginary values to two traveling waves propagating in opposite direction.

As previously mentioned, the real part of $\mathbf{k} = k\Phi$ vector is restricted to stand inside the first Brillouin zone. In the quadratic eigen value problem (18) nothing restricts computation to only find eigen values satisfying this condition. For direction vector Φ orthogonal to the lattice facelets (i.e. for $\Phi_{p1} = [1, 0]^T$ and $\Phi_{p2} = [0, 1]^T$ in bi-dimensional rectangular cell), the periodical conditions expressed for one dimensional wave guide are still valid: if $\lambda_j(\omega, \Phi_p)$ is an eigen value associated to $\mathbf{w}_j(\omega, \Phi_p)$ then $\forall \mathbf{m} \in \mathbb{Z}^3$, $\lambda + i \Phi_p^T (G \cdot \mathbf{m})$ is also an eigen value associated to $\mathbf{w}_j(\omega, \Phi_p) e^{-i \Phi_p^T (G \cdot \mathbf{m}) x}$. Thus, for undamped systems, all obtained eigenvalues are periodically distributed in the \mathbf{k} -space along its principal directions.

E. Computation of the evanescence and damped power flow criteria

One aim of this paper is to provide a numerical methodology for optimizing the piezoelectric shunt impedance $Z(\omega)$ for controlling energy flow into the periodically distributed piezocomposite structure. For doing this, we need to define suitable criteria.

The first used criterion is based on the computation of the waves group velocities. Indeed, they indicate how energy is transported into the considered system and allow to distinguish the 'propagative' and 'evanescent' waves. If one Bloch eigen

solution (i.e $\mathbf{u}_n(\omega, \phi)$, $k_n(\omega)$) is considered, the associated group velocity vector [34] is given by:

$$\mathbf{C}_{g_n}(\omega, \phi) = \nabla_{\mathbf{k}} \omega = \frac{\langle\langle \mathbf{S} \rangle\rangle}{\langle\langle e_{tot} \rangle\rangle} = \frac{\langle \mathbf{I} \rangle}{\langle E_{tot} \rangle} \quad (19)$$

where $\langle\langle \cdot \rangle\rangle$ is the spatial and time average respectively on one cell and one period, \mathbf{S} is the density of energy flux defined in [34], \mathbf{I} the mean intensity and e_{tot} , E_{tot} the total piezomechanical energy and its time average on a period (see [34] for details). In this problem, we only consider mechanical energy transportation as the electrostatic coupling is decentralized and can be condensed as a mechanical interface as proved in [35] and generally computed in [36]. So we also compute the intensity vector \mathbf{I} by:

$$\langle \mathbf{I}_n \rangle(\omega, \phi) = -\frac{\omega}{2V_{ol}} Re \left(\int_{\Omega_x} C(\varepsilon_n(\mathbf{x}) + ik\Xi_n(\mathbf{x})) \cdot (\mathbf{w}_n^*(\mathbf{x})) d\Omega \right) \quad (20)$$

where $*$ is the complex conjugate and V_{ol} the domain volume. As the spatio-temporal average of the system Lagrangian is null (see [34]), the total energy average is approximated by only computing the kinetic energy average:

$$\langle E_{tot} \rangle(\omega, \phi) = \frac{1}{2V_{ol}} Re \left(\int_{\Omega_x} \rho \omega^2 \mathbf{w}_n(\mathbf{x}) \cdot \mathbf{w}_n^*(\mathbf{x}) d\Omega \right) \quad (21)$$

The group velocity vectors $\mathbf{C}_{g_n}(\omega, \phi)$ is computed for all wave numbers at each frequency. In order to focus our analysis on only flexural modes (S and SH ones) we introduce an indicator allowing to select them by computing the ratio of kinetic energy average on out of plane displacement as:

$$Ind(n, \omega, \phi) = \frac{\frac{1}{2V_{ol}} \left(\int_{\Omega_x} \rho \omega^2 w_{zn}(\mathbf{x}) w_{zn}^*(\mathbf{x}) d\Omega \right)}{\langle E_{tot} \rangle} \quad (22)$$

with $w_{zn}(\mathbf{x}, \omega, \phi)$ being the (Oz) component of vector $\mathbf{w}_n(\mathbf{x}, \omega, \phi)$. Optimization of the shunt impedance $Z(i\omega)$ is based on the minimization of the maximum group velocity collinear to the wave number vector (19) for waves having a ratio of transported flexural kinetic energy (22) greater than 0.8. The used criterion can also be written as:

$$Crit_1(Z(i\omega), \phi) = \max_{n/Ind(n, \omega, \phi) > 0.8} (\mathbf{C}_{g_n}(\omega, \phi) \cdot \Phi) \quad (23)$$

The second used criterion is based on the maximization of the damped electric power directly express as the active electrical power $P_{elec}(n, \omega, \phi) = \frac{1}{2} real(i\omega Z(i\omega) q_{n, \mathbf{k}}^o q_{n, \mathbf{k}}^{o*})$. If one wants to increase damping effect inside the smart metamaterial, this term need to be sufficiently large. In the second case, the used criteria is also

$$Crit_2(Z(i\omega), \phi) = \max_{n/Ind(n, \omega, \phi) > 0.8} \frac{1}{P_{elec}(n, \omega, \phi)} \quad (24)$$

F. Computation of the sound radiation efficiency

The sound radiation efficiency of a plate depends on the coupling of sound waves in the air and flexural waves (vibration) in the plate. Optimum efficiency (maximum energy transfer from vibration to sound or vice versa) is achieved when the plate vibrates such that the wavelength of flexural

Piezoelectric Material		
Symbol	Value	Property
$s_{11}^E = s_{22}^E = s_{33}^E$	$11.6e^{-12} Pa^{-1}$	compliance matrix
$s_{12}^E = s_{13}^E = s_{23}^E$	$-3.33e^{-12} Pa^{-1}$	compliance matrix
$s_{44}^E = s_{55}^E = s_{66}^E$	$45.0e^{-11} Pa^{-1}$	compliance matrix
η	0.1 %	Hysteretic Damping ratio
$d_{31} = d_{32}$	$-6e^{-11} C/N$	piezoelectric matrix
d_{33}	$15.2e^{-11} C/N$	piezoelectric matrix
$d_{24} = d_{15}$	$730e^{-12} C/N$	piezoelectric matrix
ρ	$7600 kg/m^3$	Density
$\varepsilon_{11}^T = \varepsilon_{22}^T$	$504.1 \varepsilon_o C/V/m$	Dielectric Permittivity
ε_{33}^T	$270 \varepsilon_o C/V/m$	Dielectric Permittivity

TABLE I
Piezoelectric patch characteristics

waves in the plate is equal to the wavelength of acoustic waves in the air (alternatively the flexural wave velocity in the plate is equal to the velocity of acoustic waves in the air). This is more commonly known as the coincidence or critical frequency of radiation. In fact, that phenomenon is based on the computation of the kz_n value of the induced acoustic wave number. If one considers an infinite plate in which an harmonic plane wave is propagating with a wave number kx_n at pulsation ω , as the acoustic pressure radiated from this infinite system is a solution of the Helmholtz equation and is coupled by considering interface continuous normal velocity, one can obtain expression of kz_n by:

$$kz_n = \sqrt{\left(\frac{\omega}{c}\right)^2 - kx_n^2} \quad (25)$$

where c is the sound velocity in the considered semi-infinite acoustic domain.

For analyzing sound radiation effect of our smart metamaterial, acoustic wave numbers kz_n are post-processed from the obtained structural waves given by our previously described computation. As the obtained kx_n are complex and include all modes, we will focus our analysis to flexural modes by using indicator $Ind(n, \omega, \phi) > 0.8$ and consider the complex expression of the out of plane radiation by using equation (25) with a complex kz_n .

III. OPTIMIZATION OF THE FLEXURAL ENERGY FLOW INSIDE THE SHUNTED PERIODIC PIEZO-COMPOSITE

The considered piezo-composite cell is presented in figure 1. The supporting plate material is standard aluminum with 0.1 % of hysteretic damping ratio and the piezoelectric material characteristics are given in table I.

The used numerical optimizations of the criteria are based on a multidimensional unconstrained nonlinear minimization (Nelder-Mead).

A. Optimization of the waves groupes velocities by using $Crit_1$

In a first test, we optimize the waves groups velocities by using $Crit_1$ given in equation (23). We present in figure 2 the obtained real and imaginary parts of the wave number $kx_n(i\omega)$ of plane waves propagating into the smart plate

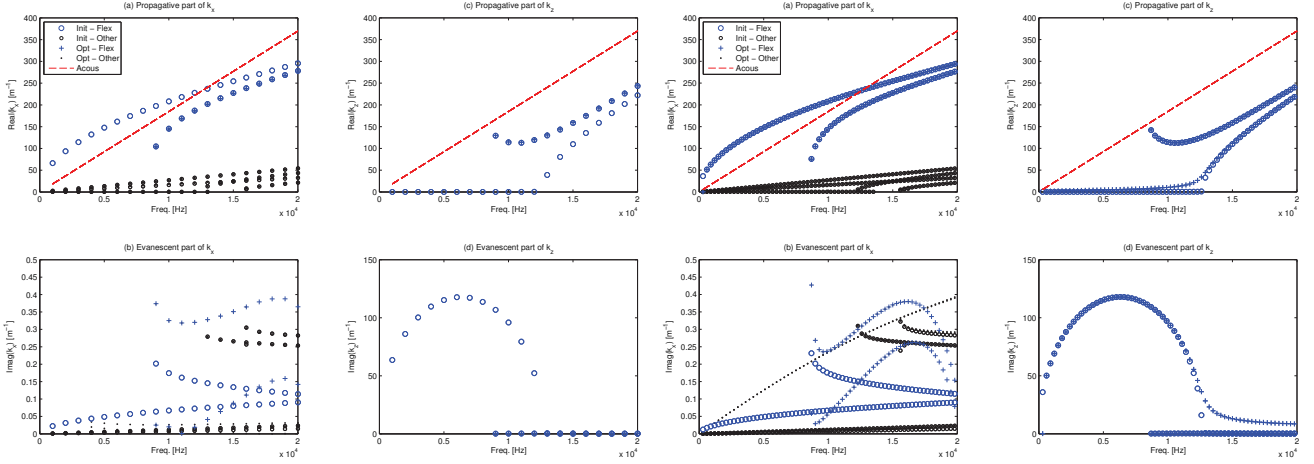


Fig. 2. Real (a) and imaginary (b) parts of the wave number $kx_n(i\omega)$ of plane waves propagating into the smart plate along (Ox) axis and the corresponding real (a) and imaginary (b) parts of the acoustic out plane wave number kz_n .

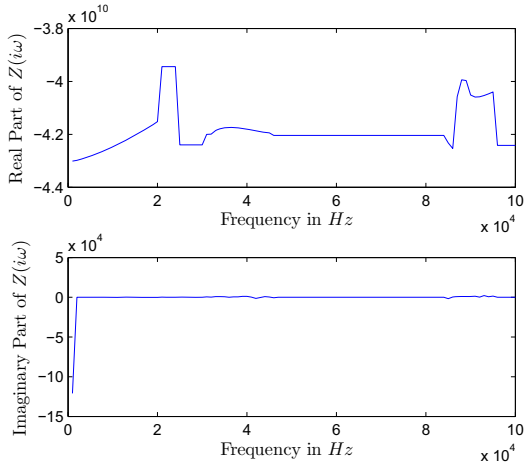


Fig. 3. Real and imaginary parts of the optimal impedance

along (Ox) axis and the corresponding real and imaginary parts of the acoustic out of plane wave number kz_n . The red dashed lines indicate the acoustic wave number $\frac{\omega}{c}$. The circles mark the dispersions curves for $Z = 0$ and the blue crosses the obtained optimal dispersion. The corresponding real and imaginary parts of the optimal impedance are plotted in figure 3. For a sake of clarity, all non propagative waves having a group velocity inferior to 30 m.s^{-1} have been removed from the plotted results.

We immediately observe that the optimization of the shunt impedance leads to greatly modify the group velocity of the A_o mode and to create an evanescent mode. We notice an increase of the imaginary parts of the wave number of A_1 mode indicating an increase of the spatial decay rates even if this mode remains propagative. The acoustic radiation is greatly modified by the cancellation of A_o mode as shown

Fig. 4. Real (a) and imaginary (b) parts of the wave number $kx_n(i\omega)$ of plane waves propagating into the smart plate along (Ox) axis and the corresponding real (a) and imaginary (b) parts of the acoustic out plane wave number kz_n .

in figure 2.c. The imaginary part of kz_n stay unchanged at very low values for radiations induced by A_1 . The optimal impedance values are almost real, and correspond to those obtained if a constant negative capacitance is used. The corresponding average value is $-150.05 \text{ pC.V}^{-1}$. Some imaginary parts of the optimal impedance are negative which indicate that the optimization leads to provide energy to the system for controlling mechanical damping effect introduced with hysteretic damping ratios into the model, and, also, obtain a fully conservative system.

B. Optimization of damped power flow inside the electric shunts by using $Crit_2$

In this second test, we optimize the damped power flow inside the electric shunts by using $Crit_2$ given in equation (24). We present in figure 4 the obtained real and imaginary parts of the wave number $kx_n(i\omega)$ of plane waves propagating into the smart plate along (Ox) axis and the corresponding real and imaginary parts of the acoustic out of plane wave number kz_n . The red dashed lines mark the acoustic wave number $\frac{\omega}{c}$. The circles mark the dispersions curves for $Z = 0$ and the blue crosses the obtained optimal dispersion. The corresponding real and imaginary parts of the optimal impedance are plotted in figure 3. For a sake of clarity, all non propagative waves having a group velocity inferior to 30 m.s^{-1} have been removed from the plotted results.

We immediately observe that the optimization of the shunt impedance leads to greatly modify the imaginary parts of the flexural wave number of A_o and A_1 modes and to create damped propagation while the real parts remain unchanged. The acoustic radiation is greatly modified in improving the spatial decay rates of the out of plane wave number coupled with A_o mode as shown in figure 4.d after 12 kHz . The real parts of kz_n stay unchanged. The optimal impedance values are complex with a positive resistance term that allows energy

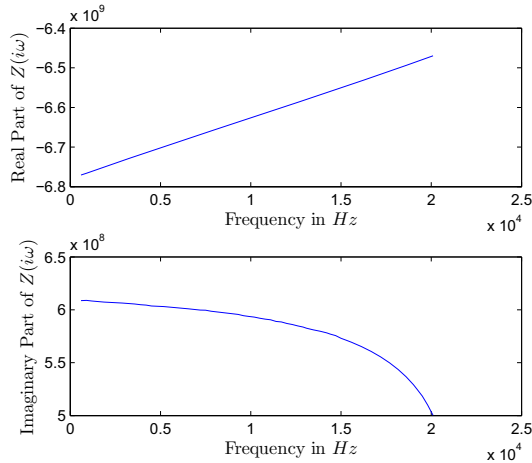


Fig. 5. Real and imaginary parts of the optimal impedance

dissipation inside the electric shunt circuit. The real parts correspond to those obtained if a constant negative capacitance is used. The corresponding average value is also $-151.04 \mu\text{C}\cdot\text{V}^{-1}$ slightly different from the previously considered case.

IV. CONCLUSIONS

This paper presents a numerical procedure able to compute the damped wave's dispersion functions in the whole first Brillouin domain of multi dimensionnal piezo-elastodynamical wave guides. The method was applied for determining the optimal impedance allowing to minimize the group velocities of the flexural waves. Based on this approach, some numerical tests on a finite dimension system incorporating a semi-distributed set of shunted piezo-composite cells have been performed. We underline a strong influence of the designed shunt circuits in the dynamical response of the system and the coupled noise radiation. Even if the link between the obtained wave properties are not clearly established, we also demonstrated that our developed numerical procedures can be used for optimizing the energy diffusion operator of such adaptive mechanical interface. To do so, additional work has to be done for optimizing the complete interface scattering and for controlling the evanescent waves playing an important role in the dynamical response of a finite system incorporating such semi-distributed interface. Another part of this future developments should deal with the complete vibroacoustic optimization incorporating a fully fluid-structure coupling effect. The proposed methodology can also be used for studying particular dissipation phenomenon such as those induced by complex shunted piezoelectric patches as proposed by [6] and [37], or even foams or complex polymers behaviors. The proposed method furnishes an efficient tool for future optimization of distributed smart cells as proposed in the case of 1D wave guide by [4].

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