

Conditional propagation of chaos for mean field systems of interacting neurons

Xavier Erny, Eva Löcherbach, Dasha Loukianova

▶ To cite this version:

Xavier Erny, Eva Löcherbach, Dasha Loukianova. Conditional propagation of chaos for mean field systems of interacting neurons. Electronic Journal of Probability, 2021, 26. hal-02280882v3

HAL Id: hal-02280882 https://hal.science/hal-02280882v3

Submitted on 2 Mar 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Conditional propagation of chaos for mean field systems of interacting neurons

Xavier Erny*, Eva Löcherbach† and Dasha Loukianova*

Abstract: We study the stochastic system of interacting neurons introduced in De Masi et al. (2015) and in Fournier and Löcherbach (2016) in a diffusive scaling. The system consists of N neurons, each spiking randomly with rate depending on its membrane potential. At its spiking time, the potential of the spiking neuron is reset to 0 and all other neurons receive an additional amount of potential which is a centred random variable of order $1/\sqrt{N}$. In between successive spikes, each neuron's potential follows a deterministic flow. We prove the convergence of the system, as $N \to \infty$, to a limit nonlinear jumping stochastic differential equation driven by Poisson random measure and an additional Brownian motion W which is created by the central limit theorem. This Brownian motion is underlying each particle's motion and induces a common noise factor for all neurons in the limit system. Conditionally on W, the different neurons are independent in the limit system. This is the conditional propagation of chaos property. We prove the well-posedness of the limit equation by adapting the ideas of Graham (1992) to our frame. To prove the convergence in distribution of the finite system to the limit system, we introduce a new martingale problem that is well suited for our framework. The uniqueness of the limit is deduced from the exchangeability of the underlying system.

MSC 2010 subject classifications: 60J75, 60K35, 60G55, 60G09.

Keywords and phrases: Multivariate nonlinear Hawkes processes with variable length memory, Mean field interaction, Piecewise deterministic Markov processes, Interacting particle systems, Propagation of chaos, Exchangeability, Hewitt Savage theorem.

Introduction

This paper is devoted to the study of the Markov process $X_t^N = (X_t^{N,1}, \dots, X_t^{N,N})$ taking values in \mathbb{R}^N and having generator A^N which is defined for any smooth test function $\varphi : \mathbb{R}^N \to \mathbb{R}$ by

$$A^{N}\varphi(x) = -\alpha \sum_{i=1}^{N} \partial_{x^{i}}\varphi(x)x^{i} + \sum_{i=1}^{N} f(x^{i}) \int_{\mathbb{R}} \nu(du) \left(\varphi(x - x^{i}e_{i} + \sum_{j \neq i} \frac{u}{\sqrt{N}}e_{j}) - \varphi(x) \right),$$

where $x=(x^1,\ldots,x^N)$ and where e_j denotes the j-th unit vector in \mathbb{R}^N . In the above formula, $\alpha>0$ is a fixed parameter and ν is a centred probability measure on \mathbb{R} having a second moment. Informally, the process $(X^{N,j})_{1\leq j\leq N}$ solves

$$X_t^{N,i} = X_0^{N,i} - \alpha \int_0^t X_s^{N,i} ds - \int_0^t X_{s-}^{N,i} dZ_s^{N,i} + \frac{1}{\sqrt{N}} \sum_{i \neq i} \int_0^t U^j(s) dZ_s^{N,j}, \tag{1}$$

where $U^j(s)$ are i.i.d. centred random variables distributed according to ν , and where for each $1 \leq j \leq N, Z^{N,j}$ is a simple point process on \mathbb{R}_+ having stochastic intensity $s \mapsto f\left(X_{s-}^{N,j}\right)$.

^{*} Université Paris-Saclay, CNRS, Univ Evry, Laboratoire de Mathématiques et Modélisation d'Evry, 91037, Evry,
France

[†] Statistique, Analyse et Modélisation Multidisciplinaire, Université Paris 1 Panthéon-Sorbonne, EA 4543 et FR FP2M 2036 CNRS

The particle system (1) is a version of the model of interacting neurons considered in De Masi et al. (2015), inspired by Galves and Löcherbach (2013), and then further studied in Fournier and Löcherbach (2016) and Cormier, Tanré and Veltz (2019). The system consists of N interacting neurons. In (1), $Z_t^{N,j}$ represents the number of spikes emitted by the neuron j in the interval [0,t] and $X_t^{N,j}$ the membrane potential of the neuron j at time t. Spiking occurs randomly following a point process of rate f(x) for any neuron of which the membrane potential equals x. Each time a neuron emits a spike, the potentials of all other neurons receive an additional amount of potential. In De Masi et al. (2015), Fournier and Löcherbach (2016) and Cormier, Tanré and Veltz (2019) this amount is of order N^{-1} , leading to classical mean field limits as $N \to \infty$. On the contrary to this, in the present article we study a diffusive scaling where each neuron j receives the amount U/\sqrt{N} at spike times t of neuron t, $t \neq t$, where t = t is a random variable. The variable t is centred modeling the fact that the synaptic weights are balanced. Moreover, right after its spike, the potential of the spiking neuron t is reset to 0, interpreted as resting potential. Finally, in between successive spikes, each neuron has a loss of potential of rate t.

Equations similar to (1) appear also in the frame of multivariate Hawkes processes with mean field interactions. Indeed, if $(Z^{N,i})_{1 \leq i \leq N}$ is a multivariate Hawkes process where the stochastic intensity of each $Z^{N,i}$ is given by $f\left(X^N_{t-}\right)_t$ with

$$X_t^N = e^{-\alpha t} X_0^N + \frac{1}{\sqrt{N}} \sum_{j=1}^N \int_0^t e^{-\alpha(t-s)} U^j(s) dZ_s^{N,j}, \tag{2}$$

then X^N satisfies

$$X_{t}^{N} = X_{0}^{N} - \alpha \int_{0}^{t} X_{s}^{N} ds + \frac{1}{\sqrt{N}} \sum_{i=1}^{N} \int_{0}^{t} U^{j}(s) dZ_{s}^{N,j},$$

which corresponds to equation (1) without the big jumps, i.e. without the reset to 0 after each spike.

The above system of interacting Hawkes processes with intensity given by (2) has been studied in our previous paper Erny, Löcherbach and Loukianova (2019). There we have shown firstly that X^N converges in distribution in $D(\mathbb{R}_+, \mathbb{R})$ to a limit process \bar{X} solving

$$d\bar{X}_t = -\alpha \bar{X}_t dt + \sigma \sqrt{f(\bar{X}_t)} dW_t, \tag{3}$$

and secondly that the sequence of multivariate counting processes $(Z^{N,i})_i$ converges in distribution in $D(\mathbb{R}_+,\mathbb{R})^{\mathbb{N}^*}$ to a limit sequence of counting processes $(\bar{Z}^i)_i$. Every \bar{Z}^i is driven by its own Poisson random measure and has the same intensity $(f(\bar{X}_{t-}))_t$, where \bar{X} is the strong solution of (3) with respect to some Brownian motion W. Consequently, the processes \bar{Z}^i $(i \geq 1)$ are conditionally independent given the Brownian motion W.

In the present paper we add the reset term in (1) that forces the potential $X^{N,i}$ of neuron i to go back to 0 at each jump time of $Z^{N,i}$. This models the well-known biological fact that right after its spike, the membrane potential of the spiking neuron is reset to a resting potential. From a mathematical point of view, this reset to 0 induces a de-synchronization of the processes $X^{N,i}$ ($1 \le i \le N$). In terms of Hawkes processes, it means that in (2), the process X_t^N has been replaced by

$$X_t^{N,i} = \frac{1}{\sqrt{N}} \sum_{i=1}^N \int_{L_t^i}^t e^{-\alpha(t-s)} U^j(s) dZ_s^{N,j} + e^{-\alpha t} X_0^{N,i} \mathbbm{1}_{L_t^i = 0},$$

where $L_t^i = \sup\{0 \leqslant s \leqslant t : \Delta Z_s^{N,i} = 1\}$ is the last spiking time of neuron i before time t, with the convention $\sup \varnothing := 0$. Thus the integral over the past, starting from 0 in (2), is replaced by an integral starting at the last jump time before the present time. Such processes are termed being of variable length memory, in reminiscence of Rissanen (1983). They are the continuous-time analogues of the model considered in Galves and Löcherbach (2013), and we are thus considering multivariate Hawkes processes with mean field interactions and variable length memory. As a consequence, on the contrary to the situation in Erny, Löcherbach and Loukianova (2019), the point processes $Z^{N,i}$ ($1 \leqslant i \leqslant N$) do not share the same stochastic intensity. The reset term in (1) is a jump term that survives in the limit $N \to \infty$.

Before introducing the exact limit equation for the system (1), let us explain informally how the limit particle system associated to $(X^{N,i})_{1 \leq i \leq N}$ should a priori look like. Suppose for the moment that we already know that there exists a process $(\bar{X}^1, \bar{X}^2, \bar{X}^3, \ldots) \in D(\mathbb{R}_+, \mathbb{R})^{\mathbb{N}^*}$ such that for all K > 0, weak convergence $\mathcal{L}(X^{N,1}, \ldots, X^{N,K}) \to \mathcal{L}(\bar{X}^1, \ldots, \bar{X}^K)$ in $D(\mathbb{R}_+, \mathbb{R})^K$, as $N \to \infty$, holds. In equation (1) the only term that depends on N is the martingale term which is approximately given by

$$M_t^N = \frac{1}{\sqrt{N}} \sum_{j=1}^N \int_0^t U^j(s) dZ_s^{N,j}.$$

Then in the infinite neuron model, each process \bar{X}^i should solve the equation (1), where the term M_t^N is replaced by $M_t := \lim_{N \to \infty} M_t^N$. Because of the scaling in $N^{-1/2}$, the limit martingale M_t will be a stochastic integral with respect to some Brownian motion, and its variance the limit of

$$\mathbb{E}\left[(M_t^N)^2\right] = \sigma^2 \int_0^t \mathbb{E}\left[\frac{1}{N} \sum_{j=1}^N f(X_s^{N,j})\right] ds,$$

where σ^2 is the variance of $U^j(s)$. Therefore, the limit martingale (if it exists) must be of the form

$$M_t = \sigma \int_0^t \sqrt{\lim_{N \to \infty} \frac{1}{N} \sum_{j=1}^N f\left(X_s^{N,j}\right)} dW_s = \sigma \int_0^t \sqrt{\lim_{N \to \infty} \mu_s^N(f)} dW_s,$$

where μ_s^N is the empirical measure of the system $(X_s^{N,j})_{1 \le j \le N}$ and W is a one-dimensional standard Brownian motion.

Since the law of the N-particle system $(X^{N,1},\ldots,X^{N,N})$ is symmetric, the law of the limit system $\bar{X}=(\bar{X}^1,\bar{X}^2,\bar{X}^3,\ldots)$ must be exchangeable, that is, for all finite permutations σ , we have that $\mathcal{L}(\bar{X}^{\sigma(1)},\bar{X}^{\sigma(2)},\ldots)=\mathcal{L}(\bar{X})$. In particular, the theorem of Hewitt-Savage, see Hewitt and Savage (1955), implies that the random limit

$$\mu_s := \lim_{N \to \infty} \frac{1}{N} \sum_{i=1}^N \delta_{\bar{X}_s^i} \tag{4}$$

exists. Supposing that μ_s^N converges, it necessarily converges towards μ_s . Therefore, \bar{X} should solve the limit system

$$\bar{X}_t^i = \bar{X}_0^i - \alpha \int_0^t \bar{X}_s^i ds - \int_0^t \bar{X}_{s-}^i d\bar{Z}_s^i + \sigma \int_0^t \sqrt{\mu_s(f)} dW_s, i \in \mathbb{N},$$
 (5)

where each \bar{Z}^i has intensity $t \mapsto f(\bar{X}_{t-}^i)$, and where μ_s is given by (4).

The above arguments are made rigorous in Sections 2.1 and 2.2 below.

Let us briefly discuss the form of the limit equation (5). Analogously to Erny, Löcherbach and Loukianova (2019), the scaling in $N^{-1/2}$ in (1) creates a Brownian motion W in the limit system (5). We will show that the presence of this Brownian motion entails a conditional propagation of chaos, that is the conditional independence of the particles given W. In particular, the limit measure μ_s will be random. This differs from the classical framework, where the scaling is in N^{-1} (see e.g. Delattre, Fournier and Hoffmann (2016), Ditlevsen and Löcherbach (2017) in the framework of Hawkes processes, and De Masi et al. (2015), Fournier and Löcherbach (2016) and Cormier, Tanré and Veltz (2019) in the framework of systems of interacting neurons), leading to a deterministic limit measure μ_s and the true propagation of chaos property implying that the particles of the limit system are independent.

This is not the first time that conditional propagation of chaos is studied in the literature; it has already been considered e.g. in Carmona, Delarue and Lacker (2016), Coghi and Flandoli (2016) and Dermoune (2003). But in these papers the common noise, represented by a common (maybe infinite dimensional) Brownian motion, is already present at the level of the finite particle system, the mean field interactions act on the drift of each particle, and the scaling is the classical one in N^{-1} . On the contrary to this, in our model, this common Brownian motion, leading to conditional propagation of chaos, is only present in the limit, and it is created by the central limit theorem as a consequence of the joint action of the small jumps of the finite size particle system. Moreover, in our model, the interactions survive as a variance term in the limit system due to the diffusive scaling in $N^{-1/2}$.

Now let us discuss the form of μ_s , which is the limit of the empirical measures of the limit system $(\bar{X}_s^i)_{i\geqslant 1}$. The theorem of Hewitt-Savage, Hewitt and Savage (1955), implies that the law of $(\bar{X}_s^i)_{i\geqslant 1}$ is a mixture directed by the law of μ_s . As it has been remarked by Carmona, Delarue and Lacker (2016) and Coghi and Flandoli (2016), this conditioning reflects the dependencies between the particles.

We will show that the variables \bar{X}^i are conditionally independent given the Brownian motion W. As a consequence, μ_s will be shown to be the conditional law of the solution given the Brownian motion, that is, P-almost surely,

$$\mu_s(\cdot) = P(\bar{X}_s^i \in \cdot | (W_t)_{0 \le t \le s}) = P(\bar{X}_s^i \in \cdot | W), \tag{6}$$

for any $i \in \mathbb{N}$. Equation (5) together with (6) gives a precise definition of the limit system.

The nonlinear SDE (5) is not clearly well-posed, and our first main result, Theorem 1.2, gives appropriate conditions on the system that guarantee pathwise uniqueness and the existence of a strong solution to (5).

We then prove, in Sections 2.1 and 2.2, our main Theorem 1.7 stating the convergence in distribution of the sequence of empirical measures $\mu^N = N^{-1} \sum_{i=1}^N \delta_{(X_i^{N,i})_{t\geqslant 0}}$, in $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R}))$, to the random limit $\mu = P((\bar{X}_t)_{t\geqslant 0} \in \cdot |W)$. To do so, we first prove that under suitable conditions on the parameters of the system, the sequence μ^N is tight (see Proposition 2.1 below). We then follow a classical road and identify every possible limit as solution of a martingale problem. Since the random limit measure μ will only be the directing measure of the limit system (that is, the conditional law of each coordinate, but not its law), this martingale problem is not a classical one. It is in particular designed to reflect the correlation between the particles and to describe all possible limits of couples of neurons.

Classical representation theorems imply that any coordinate of the limit process must satisfy an equation of the type (5). The fact that our martingale problem describes correlations within couples

of neurons allows to show that each coordinate is driven by its own Poisson random measure and that all coordinates are driven by the same underlying Brownian motion W. But it is not yet clear that μ_s is of the form (6). In other words, it has to be proven that the only common randomness is the one present in the driving Brownian motion W. To prove this last point, we introduce an auxiliary particle system which is a mean field particle version of the limit system, constructed with the same underlying Brownian motion, and we provide an explicit control on the distance (with respect to a particular L^1 -norm) between the two systems.

Let us finally mention that the random limit measure μ satisfies the following nonlinear stochastic PDE in weak form: for any test function $\varphi \in C_b^2(\mathbb{R})$, the set of C^2 -functions on \mathbb{R} such that φ , φ' and φ'' are bounded, for any $t \geq 0$,

$$\int_{\mathbb{R}} \varphi(x) \mu_t(dx) = \int_{\mathbb{R}} \varphi(x) \mu_0(dx) + \int_0^t \left(\int_{\mathbb{R}} \varphi'(x) \mu_s(dx) \right) \sqrt{\mu_s(f)} dW_s$$

$$+ \int_0^t \int_{\mathbb{R}} \left([\varphi(0) - \varphi(x)] f(x) - \alpha \varphi'(x) x + \frac{1}{2} \varphi''(x) \mu_s(f) \right) \mu_s(dx) ds.$$

Organisation of the paper. In Section 1, we state the assumptions and formulate the main results. Section 2 is devoted to the proof of the convergence of $\mu^N := \sum_{j=1}^N \delta_{X^{N,j}}$ (Theorem 1.7). In particular, we introduce our new martingale problem in Section 2.2 and prove the uniqueness of the limit law in Theorem 2.6. Finally, in Appendix, we prove some auxiliary results.

1. Notation, Model and main results

1.1. Notation

We use the following notation throughout the paper.

If E is a metric space, we note:

- $\mathcal{P}(E)$ the space of probability measures on E endowed with the topology of the weak convergence,
- $C_b^n(E)$ the set of the functions g which are n times continuously differentiable such that $g^{(k)}$ is bounded for each $0 \le k \le n$,
- $C_c^n(E)$ the set of functions $g \in C_b^n(E)$ that have a compact support.

In addition, in what follows $D(\mathbb{R}_+, \mathbb{R})$ denotes the space of càdlàg functions from \mathbb{R}_+ to \mathbb{R} , endowed with Skorohod metric, and C and K denote arbitrary positive constants whose values can change from line to line in an equation. We write C_{θ} and K_{θ} if the constants depend on some parameter θ .

1.2. The finite system

We consider, for each $N\geqslant 1$, a family of i.i.d. Poisson measures $(\pi^i(ds,dz,du))_{i=1,\dots,N}$ on $\mathbb{R}_+\times\mathbb$

taking values in \mathbb{R}^N and solving, for i = 1, ..., N, for $t \ge 0$,

$$\begin{cases}
X_t^{N,i} = X_0^{N,i} - \alpha \int_0^t X_s^{N,i} ds - \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} X_{s-}^{N,i} \mathbb{1}_{\{z \le f(X_{s-}^{N,i})\}} \pi^i(ds, dz, du) \\
+ \frac{1}{\sqrt{N}} \sum_{j \ne i} \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} u \mathbb{1}_{\{z \le f(X_{s-}^{N,i})\}} \pi^j(ds, dz, du), \\
X_0^{N,i} \sim \nu_0.
\end{cases} (7)$$

The coefficients of this system are the exponential loss factor $\alpha > 0$, the jump rate function $f : \mathbb{R} \to \mathbb{R}_+$ and the probability measures ν and ν_0 .

In order to guarantee existence and uniqueness of a strong solution of (7), we introduce the following hypothesis.

Assumption 1. The function f is Lipschitz continuous.

In addition, we also need the following condition to obtain a priori bounds on some moments of the process $(X^{N,i})_{1 \le i \le N}$.

Assumption 2. We assume that
$$\int_{\mathbb{R}} x d\nu(x) = 0$$
, $\int_{\mathbb{R}} x^2 d\nu(x) < +\infty$, and $\int_{\mathbb{R}} x^2 d\nu_0(x) < +\infty$.

Under Assumptions 1 and 2, existence and uniqueness of strong solutions of (7) follow from Theorem IV.9.1 of Ikeda and Watanabe (1989), exactly in the same way as in Proposition 6.6 of Erny, Löcherbach and Loukianova (2019).

We now define precisely the limit system and discuss its properties before proving the convergence of the finite to the limit system.

1.3. The limit system

The limit system $(\bar{X}^i)_{i\geq 1}$ is given by

$$\begin{cases}
\bar{X}_{t}^{i} = \bar{X}_{0}^{i} - \alpha \int_{0}^{t} \bar{X}_{s}^{i} ds - \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \bar{X}_{s-}^{i} \mathbb{1}_{\{z \leqslant f(\bar{X}_{s-}^{i})\}} \pi^{i}(ds, dz, du) \\
+ \sigma \int_{0}^{t} \sqrt{\mathbb{E}\left[f\left(\bar{X}_{s}^{i}\right) \middle| \mathcal{W}\right]} dW_{s}, \\
\bar{X}_{0}^{i} \sim \nu_{0}.
\end{cases} \tag{8}$$

In the above equation, $(W_t)_{t\geqslant 0}$ is a standard one-dimensional Brownian motion which is independent of the Poisson random measures, and $W = \sigma\{W_t, t\geqslant 0\}$. Moreover, the initial positions $\bar{X}_0^i, i\geqslant 1$, are i.i.d., independent of W and of the Poisson random measures, distributed according to ν_0 which is the same probability measure as in (7). The common jumps of the particles in the finite system, due to their scaling in $1/\sqrt{N}$ and the fact that they are centred, by the Central Limit Theorem, create this single Brownian motion W_t which is underlying each particle's motion and which induces the common noise factor for all particles in the limit.

The limit equation (8) is not clearly well-posed and requires more conditions on the rate function f. Let us briefly comment on the type of difficulties that one encounters when proving trajectorial uniqueness of (8). Roughly speaking, the jump terms demand to work in an L^1 -framework, while the diffusive terms demand to work in an L^2 -framework. Graham (1992) proposes a unified approach to deal both with jump and with diffusion terms in a non-linear framework, and we shall

rely on his ideas in the sequel. The presence of the random volatility term which involves conditional expectation causes however additional technical difficulties. Finally, another difficulty comes from the fact that the jumps induce non-Lipschitz terms of the form $\bar{X}_s^i f(\bar{X}_s^i)$. For this reason a classical Wasserstein-1—coupling is not appropriate for the jump terms. Therefore we propose a different distance which is inspired by the one already used in Fournier and Löcherbach (2016). To do so, we need to work under the following additional assumption.

Assumption 3. 1. We suppose that inf f > 0.

2. There exists a function $a \in C^2(\mathbb{R}, \mathbb{R}_+)$, strictly increasing and bounded, such that, for a suitable constant C, for all $x, y \in \mathbb{R}$,

$$|a''(x) - a''(y)| + |a'(x) - a'(y)| + |xa'(x) - ya'(y)| + |f(x) - f(y)| \le C|a(x) - a(y)|.$$

Note that Assumption 3 implies Assumption 1 as well as the boundedness of the rate function f.

Proposition 1.1. Suppose that $f(x) = c + d \arctan(x)$, where $c > d \frac{\pi}{2}$, d > 0. Then Assumption 3 holds with a = f.

Proof. We quickly check that $|xa'(x)-ya'(y)| \leq C|a(x)-a(y)|$. We have that $a'(x)=\frac{d}{1+x^2}$, whence $xa'(x)-ya'(y)=d(\frac{x}{1+x^2}-\frac{y}{1+y^2})$. We use that $\left|\frac{d}{dx}\left(\frac{x}{1+x^2}\right)\right|=\left|\frac{1-x^2}{(1+x^2)^2}\right| \leq \frac{1}{1+x^2}$. Suppose w.l.o.g. that $x \leq y$. As a consequence,

$$|xa'(x) - ya'(y)| = d \left| \int_{x}^{y} \frac{1 - t^{2}}{(1 + t^{2})^{2}} dt \right| \le d \int_{x}^{y} \frac{1}{1 + t^{2}} dt = |\arctan(y) - \arctan(x)| = d|a(x) - a(y)|.$$

The other points of Assumption 3 follow immediately.

Under these additional assumptions we obtain the well-posedness of each coordinate of the limit system (8), that is, of the $(\mathcal{F}_t)_t$ adapted process $(\bar{X}_t)_t$ having càdlàg trajectories which is solution of the SDE

$$\begin{cases}
d\bar{X}_{t} = -\alpha \bar{X}_{t}dt - \bar{X}_{t-} \int_{\mathbb{R}_{+} \times \mathbb{R}} \mathbb{1}_{\{z \leq f(\bar{X}_{t-})\}} \pi(dt, dz, du) + \sigma \sqrt{\mu_{t}(f)} dW_{t}, \\
\bar{X}_{0} \sim \nu_{0}, \quad \mu_{t}(f) = \mathbb{E}\left[f\left(\bar{X}_{t}\right) \middle| \mathcal{W}\right] = \mathbb{E}\left[f\left(\bar{X}_{t}\right) \middle| \mathcal{W}_{t}\right].
\end{cases} \tag{9}$$

Here, $\mathcal{F}_t = \sigma\{\pi([0,s] \times A), s \leqslant t, A \in \mathcal{B}(\mathbb{R}_+ \times \mathbb{R})\} \vee \mathcal{W}_t, \mathcal{W}_t = \sigma\{W_s, s \leqslant t\} \text{ and } \mathcal{W} = \sigma\{W_s, s \geqslant 0\}.$

Theorem 1.2. Grant Assumption 3.

- 1. Pathwise uniqueness holds for the nonlinear SDE (9).
- 2. If additionally, $\int_{\mathbb{R}} x^2 d\nu_0(x) < +\infty$, then there exists a unique strong solution $(\bar{X}_t)_{t\geq 0}$ of the nonlinear SDE (9), which is $(\mathcal{F}_t)_t$ adapted with càdlàg trajectories, satisfying for every t>0,

$$\mathbb{E}\left[\sup_{0 \le s \le t} \bar{X}_s^2\right] < +\infty. \tag{10}$$

Remark 1.3. Notice that the stochastic integral $\int_0^t \sqrt{\mu_s(f)} dW_s$ is well-defined since $s \mapsto \sqrt{\mu_s(f)}$ is an $(\mathcal{W}_t)_t$ -progressively measurable process.

In what follows we just give the proof of Item 1. of the above theorem since its arguments are important for the sequel. We postpone the rather classical proof of Item 2. to Appendix.

Proof of Item 1. of Theorem 1.2. Consider two solutions $(\hat{X}_t)_{t\geq 0}$ and $(\check{X}_t)_{t\geq 0}$, $(\mathcal{F}_t)_t$ -adapted, defined on the same probability space and driven by the same Poisson random measure π and the same Brownian motion W, and with $\hat{X}_0 = \check{X}_0$. We consider $Z_t := a(\hat{X}_t) - a(\check{X}_t)$. Denote $\hat{\mu}_s(f) = \mathbb{E}[f(\hat{X}_s)|\mathcal{W}_s]$ and $\check{\mu}_s(f) = \mathbb{E}[f(\check{X}_s)|\mathcal{W}_s]$. Using Ito's formula, we can write

$$\begin{split} Z_t &= -\alpha \int_0^t \left(\hat{X}_s a'(\hat{X}_s) - \check{X}_s a'(\check{X}_s) \right) ds + \frac{1}{2} \int_0^t (a''(\hat{X}_s) \hat{\mu}_s(f) - a''(\check{X}_s) \check{\mu}_s(f)) \sigma^2 ds \\ &+ \int_0^t (a'(\hat{X}_s) \sqrt{\hat{\mu}_s(f)} - a'(\check{X}_s) \sqrt{\check{\mu}_s(f)}) \sigma dW_s \\ &- \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \left[a(\hat{X}_{s-}) - a(\check{X}_{s-}) \right] \mathbb{1}_{\left\{z \leqslant f(\widehat{X}_{s-}) \land f(\widecheck{X}_{s-})\right\}} \pi(ds, dz, du) \\ &+ \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \left[a(0) - a(\hat{X}_{s-}) \right] \mathbb{1}_{\left\{f(\widecheck{X}_{s-}) < z \leqslant f(\widehat{X}_{s-})\right\}} \pi(ds, dz, du) \\ &+ \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \left[a(\widecheck{X}_{s-}) - a(0) \right] \mathbb{1}_{\left\{f(\widehat{X}_{s-}) < z \leqslant f(\widecheck{X}_{s-})\right\}} \pi(ds, dz, du) =: A_t + M_t + \Delta_t, \end{split}$$

where A_t denotes the bounded variation part of the evolution, M_t the martingale part and Δ_t the sum of the three jump terms. Notice that

$$M_t = \int_0^t (a'(\hat{X}_s)\sqrt{\hat{\mu}_s(f)} - a'(\check{X}_s)\sqrt{\check{\mu}_s(f)})\sigma dW_s$$

is a square integrable martingale since f and a' are bounded.

We wish to obtain a control on $|Z_t^*| := \sup_{s \leq t} |Z_s|$. We first take care of the jumps of $|Z_t|$. Notice first that, since f and a are bounded,

$$\Delta(x,y) := (f(x) \land f(y))|a(x) - a(y)| + |f(x) - f(y)| \left| |a(0) - a(y)| + |a(0) - a(x)| \right| \le C|a(x) - a(y)|,$$

implying that

$$\mathbb{E}\sup_{s\leqslant t} |\Delta_s| \leqslant C\mathbb{E}\int_0^t |a(\hat{X}_s) - a(\check{X}_s)| ds \leqslant Ct \, \mathbb{E}|Z_t^*|.$$

Moreover, for a constant C depending on σ^2 , $\|f\|_{\infty}$, $\|a\|_{\infty}$, $\|a'\|_{\infty}$, $\|a''\|_{\infty}$ and α ,

$$\begin{split} \mathbb{E} \sup_{s \leqslant t} |A_s| \leqslant C \int_0^t \mathbb{E} |a'(\hat{X}_s) \hat{X}_s - a'(\check{X}_s) \check{X}_s | ds \\ &+ C \left[\int_0^t |a''(\hat{X}_s) - a''(\check{X}_s)| ds + \int_0^t |\hat{\mu}_s(f) - \check{\mu}_s(f)| ds \right]. \end{split}$$

We know that $|a'(\hat{X}_s)\hat{X}_s - a'(\check{X}_s)\check{X}_s| + |a''(\hat{X}_s) - a''(\check{X}_s)| \leq C|a(\hat{X}_s) - a(\check{X}_s)| = C|Z_s|$. Therefore,

$$\mathbb{E}\sup_{s\leqslant t}|A_s|\leqslant C\mathbb{E}\left[\int_0^t|Z_s|ds+\int_0^t|\widehat{\mu}_s(f)-\widecheck{\mu}_s(f)|ds\right].$$

Moreover,

$$|\hat{\mu}_s(f) - \check{\mu}_s(f)| = \left| \mathbb{E}\left(f(\hat{X}_s) - f(\check{X}_s)|\mathcal{W}\right) \right| \leqslant \mathbb{E}\left(|f(\hat{X}_s) - f(\check{X}_s)||\mathcal{W}\right) \leqslant \mathbb{E}(|Z_s||\mathcal{W}),$$

and thus.

$$\mathbb{E} \int_0^t |\hat{\mu}_s(f) - \widecheck{\mu}_s(f)| ds \leqslant \mathbb{E} \int_0^t |Z_s| ds \leqslant t \mathbb{E} |Z_t^*|.$$

Putting all these upper bounds together we conclude that for a constant C not depending on t,

$$\mathbb{E}\sup_{s\leqslant t}|A_s|\leqslant Ct\mathbb{E}|Z_t^*|.$$

Finally, we treat the martingale part using the Burkholder-Davis-Gundy inequality, and we obtain

$$\mathbb{E}\sup_{s\leqslant t}|M_s|\leqslant C\mathbb{E}\left[\left(\int_0^t (a'(\hat{X}_s)\sqrt{\hat{\mu}_s(f)}-a'(\check{X}_s)\sqrt{\check{\mu}_s(f)})^2ds\right)^{1/2}\right].$$

But

$$(a'(\hat{X}_s)\sqrt{\hat{\mu}_s(f)} - a'(\check{X}_s)\sqrt{\check{\mu}_s(f)})^2 \leqslant C\left[((a'(\hat{X}_s) - a'(\check{X}_s))^2 + (\sqrt{\hat{\mu}_s(f)} - \sqrt{\check{\mu}_s(f)})^2\right]$$

$$\leqslant C|Z_t^*|^2 + C(\sqrt{\hat{\mu}_s(f)} - \sqrt{\check{\mu}_s(f)})^2, \quad (11)$$

where we have used that $|a'(x) - a'(y)| \le C|a(x) - a(y)|$ and that f and a' are bounded. Finally, since $\inf f > 0$,

$$|\sqrt{\widehat{\mu}_s(f)} - \sqrt{\widecheck{\mu}_s(f)}|^2 \leqslant C|\widehat{\mu}_s(f) - \widecheck{\mu}_s(f)|^2 \leqslant C\left(\mathbb{E}(|Z_s^*||\mathcal{W}_s)\right)^2.$$

We use that $|Z_s^*| \leq |Z_t^*|$, implying that $\mathbb{E}(|Z_s^*||\mathcal{W}) \leq \mathbb{E}(|Z_t^*||\mathcal{W})$. Therefore we obtain the upper bound

$$|\sqrt{\hat{\mu}_s(f)} - \sqrt{\check{\mu}_s(f)}|^2 \le C \left(\mathbb{E}(|Z_t^*||\mathcal{W})\right)^2$$

for all $s \leq t$, which implies the control of

$$\mathbb{E}\sup_{s\leqslant t}|M_s|\leqslant C\sqrt{t}\mathbb{E}|Z_t^*|.$$

The above upper bounds imply that, for a constant C not depending on t nor on the initial condition,

$$\mathbb{E}|Z_t^*| \leqslant C(t + \sqrt{t})\mathbb{E}|Z_t^*|,$$

and therefore, for t_1 sufficiently small, $\mathbb{E}|Z_{t_1}^*| = 0$. We can repeat this argument on intervals $[t_1, 2t_1]$, with initial condition \hat{X}_{t_1} , and iterate it up to any finite T because t_1 does only depend on the coefficients of the system but not on the initial condition. This implies the assertion.

Remark 1.4. Theorem 1.2 states the well-posedness of the SDE (9). Under the same hypotheses, with almost the same reasoning, one can prove the well-posedness of the system (8).

In the sequel, we shall also use an important property of the limit system (8), which is the conditional independence of the processes \bar{X}^i $(i \ge 1)$ given the Brownian motion W.

Proposition 1.5. If Assumption 3 holds and $\int_{\mathbb{R}} x^2 d\nu_0(x) < +\infty$, then

- (i) for all $N \in \mathbb{N}^*$ there exists a strong solution $(\bar{X}^i)_{1 \leq i \leq N}$ of (8), and pathwise uniqueness holds.
- (ii) $\bar{X}^1, \dots, \bar{X}^N$ are independent conditionally to W,
- (iii) for all $t \geq 0$, almost surely, the weak limit of $\frac{1}{N} \sum_{i=1}^{N} \delta_{\bar{X}_{[0,t]}^{i}}$ is given by $\lim_{N \to \infty} \frac{1}{N} \sum_{i=1}^{N} \delta_{\bar{X}_{[0,t]}^{i}} = P(\bar{X}_{[0,t]}^{i} \in \cdot | \mathcal{W}_{t}) = P(\bar{X}_{[0,t]}^{i} \in \cdot | \mathcal{W}_{t}).$

Let us finally mention that the random limit measure μ satisfies a nonlinear stochastic PDE in weak form. More precisely,

Corollary 1.6. Grant Assumption 3 and suppose that $\int_{\mathbb{R}} x^2 d\nu_0(x) < +\infty$. Then the measure $\mu = P((\bar{X}_t)_{t \geq 0} \in \cdot | W)$ satisfies the following nonlinear stochastic PDE in weak form: for any $\varphi \in C_b^2(\mathbb{R})$, for any $t \geq 0$,

$$\int_{\mathbb{R}} \varphi(x)\mu_t(dx) = \int_{\mathbb{R}} \varphi(x)\nu_0(dx) + \int_0^t \left(\int_{\mathbb{R}} \varphi'(x)\mu_s(dx) \right) \sqrt{\mu_s(f)}\sigma dW_s$$

$$+ \int_0^t \int_{\mathbb{R}} \left([\varphi(0) - \varphi(x)]f(x) - \alpha\varphi'(x)x + \frac{1}{2}\sigma^2\varphi''(x)\mu_s(f) \right) \mu_s(dx)ds.$$

The proofs of Proposition 1.5 and of Corollary 1.6 are postponed to Appendix.

1.4. Convergence to the limit system

We are now able to state our main result.

Theorem 1.7. Grant Assumptions 1, 2 and 3. Then the empirical measure $\mu^N = \frac{1}{N} \sum_{i=1}^N \delta_{X^{N,i}}$ of the N-particle system $(X^{N,i})_{1 \leq i \leq N}$ converges in distribution in $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R}))$ to $\mu := \mathcal{L}(\bar{X}^1|\mathcal{W})$, where $(\bar{X}^i)_{i \geq 1}$ is solution of (8).

Corollary 1.8. Under the assumptions of Theorem 1.7, $(X^{N,j})_{1 \leq j \leq N}$ converges in distribution to $(\bar{X}^j)_{j \geq 1}$ in $D(\mathbb{R}_+, \mathbb{R})^{\mathbb{N}^*}$.

Proof. Together with the statement of Theorem 1.7, the proof is an immediate consequence of Proposition 7.20 of Aldous (1983).

We will prove Theorem 1.7 in a two step procedure. Firstly we prove the tightness of the sequence of empirical measures, and then in a second step we identify all possible limits as solutions of a martingale problem.

2. Proof of Theorem 1.7

This section is dedicated to prove that the sequence $(\mu^N)_N$ of the empirical measures $\mu^N := \sum_{j=1}^N \delta_{(X_t^j)_{t\geqslant 0}}$ converges in distribution to $\mu := \mathcal{L}(\bar{X}^1|\mathcal{W})$, where $(\bar{X}^j)_{j\geqslant 1}$ is solution of (8).

In a first time, we prove that the sequence $(\mu^N)_N$ is tight on $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R}))$. The main step to prove the convergence of $(\mu^N)_N$ is then to show that each converging subsequence converges to the same limit in distribution. For this purpose, we introduce a new martingale problem, and we show that every possible limit of μ^N is a solution of this martingale problem. Finally, we will show how the uniqueness of the limit law follows from the exchangeability of the system.

2.1. Tightness of $(\mu^N)_N$

Proposition 2.1. Grant Assumptions 1 and 2. For each $N \ge 1$, consider the unique solution $(X_t^N)_{t\geqslant 0}$ to (7) starting from some i.i.d. ν_0 -distributed initial conditions $X_0^{N,i}$.

- (i) The sequence of processes (X_t^{N,1})_{t≥0} is tight in D(ℝ₊, ℝ).
 (ii) The sequence of empirical measures μ^N = N⁻¹ ∑_{i=1}^N δ_{(X_t^{N,i})_{t≥0}} is tight in P(D(ℝ₊, ℝ)).

Proof. First, it is well-known that point (ii) follows from point (i) and the exchangeability of the system, see (Sznitman 1989, Proposition 2.2-(ii)). We thus only prove (i). To show that the family $((X_t^{N,1})_{t\geqslant 0})_{N\geqslant 1}$ is tight in $D(\mathbb{R}_+,\mathbb{R})$, we use the criterion of Aldous, see Theorem 4.5 of Jacod and Shiryaev (2003). It is sufficient to prove that

- (a) for all T > 0, all $\varepsilon > 0$, $\lim_{\delta \downarrow 0} \lim \sup_{N \to \infty} \sup_{(S,S') \in A_{\delta,T}} P(|X_{S'}^{N,1} X_{S}^{N,1}| > \varepsilon) = 0$, where $A_{\delta,T}$ is the set of all pairs of stopping times (S,S') such that $0 \le S \le S' \le S + \delta \le T$ a.s.,
- (b) for all T > 0, $\lim_{K \uparrow \infty} \sup_{N} P(\sup_{t \in [0,T]} |X_t^{N,1}| \ge K) = 0$.

To check (a), consider $(S, S') \in A_{\delta,T}$ and write

$$\begin{split} X_{S'}^{N,1} - X_{S}^{N,1} &= -\int_{S}^{S'} \int_{\mathbb{R}} \int_{0}^{\infty} X_{s-}^{N,1} \mathbbm{1}_{\{z \leqslant f(X_{s-}^{N,1})\}} \pi^{1}(ds,du,dz) - \alpha \int_{S}^{S'} X_{s}^{N,1} ds \\ &+ \frac{1}{\sqrt{N}} \sum_{i=2}^{N} \int_{S}^{S'} \int_{\mathbb{R}} \int_{0}^{\infty} u \mathbbm{1}_{\{z \leqslant f(X_{s-}^{N,j})\}} \pi^{j}(ds,du,dz), \end{split}$$

implying that

$$\begin{split} |X_{S'}^{N,1} - X_{S}^{N,1}| & \leqslant |\int_{S}^{S'} \int_{\mathbb{R}} \int_{0}^{\infty} X_{s-}^{N,1} \mathbbm{1}_{\{z \leqslant f(X_{s-}^{N,1})\}} \pi^{1}(ds,du,dz)| \\ & + \delta \alpha \sup_{0 \leqslant s \leqslant T} \left| X_{s}^{N,1} \right| + |\frac{1}{\sqrt{N}} \sum_{j=2}^{N} \int_{S}^{S'} \int_{\mathbb{R}} \int_{0}^{\infty} u \mathbbm{1}_{\{z \leqslant f(X_{s-}^{N,j})\}} \pi^{j}(ds,du,dz)| \\ & =: |I_{S,S'}| + \delta \alpha \sup_{0 \leqslant s \leqslant T} \left| X_{s}^{N,1} \right| + |J_{S,S'}|. \end{split}$$

We first note that $|I_{S,S'}| > 0$ implies that $\tilde{I}_{S,S'} := \int_S^{S'} \int_{\mathbb{R}} \int_0^\infty \mathbbm{1}_{\{z \leqslant f(X_{s-}^{N,1})\}} \pi^i(ds,du,dz) \geqslant 1$, whence

$$P(|I_{S,S'}| > 0) \leqslant P(\tilde{I}_{S,S'} \geqslant 1) \leqslant \mathbb{E}[\tilde{I}_{S,S'}] \leqslant \mathbb{E}\Big[\int_{S}^{S+\delta} f(X_s^{N,1}) ds\Big] \leqslant ||f||_{\infty} \delta,$$

since f is bounded. We proceed similarly to check that

$$P(|J_{S,S'}| \geqslant \varepsilon) \leqslant \frac{1}{\varepsilon^2} \mathbb{E}[(J_{S,S'})^2] \leqslant \frac{\sigma^2}{N\varepsilon^2} \sum_{j=2}^N \mathbb{E}\Big[\int_S^{S+\delta} f(X_s^{N,j}) ds\Big] \leqslant \frac{\sigma^2}{\varepsilon^2} ||f||_{\infty} \delta.$$

The term $\sup_{0 \le s \le T} |X_s^{N,1}|$ can be handled using Lemma 3.1.(ii).

Finally (b) is a straightforward consequence of Lemma 3.1 (ii) and Markov's inequality.

2.2. Martingale problem

We now introduce a new martingale problem, whose solutions are the limits of any converging subsequence of $\mu^N = \frac{1}{N} \sum_{j=1}^N \delta_{X^{N,j}}$. In this martingale problem, we are interested in couples of trajectories to be able to put hands on the correlations between the particles. In particular, this will allow us to show that, in the limit system (8), the processes \bar{X}^i ($i \ge 1$) share the same Brownian motion, but are driven by Poisson measures π^i ($i \ge 1$) which are independent. The reason why we only need to study the correlation between two particles is the exchangeability of the infinite system.

Let Q be a distribution on $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R}))$. Define a probability measure P on $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R})) \times D(\mathbb{R}_+,\mathbb{R})^2$ by

$$P(A \times B) := \int_{\mathcal{P}(D(\mathbb{R}_+, \mathbb{R}))} \mathbb{1}_A(m) m \otimes m(B) Q(dm). \tag{12}$$

We write any atomic event $\omega \in \Omega := \mathcal{P}(D(\mathbb{R}_+, \mathbb{R})) \times D(\mathbb{R}_+, \mathbb{R})^2$ as $\omega = (\mu, Y)$, with $Y = (Y^1, Y^2)$. Thus, the law of the canonical variable μ is Q, and that of $(Y_t)_{t\geq 0}$ is

$$P_Y = \int_{\mathcal{P}(D(\mathbb{R}_+, \mathbb{R}))} Q(dm)m \otimes m(\cdot).$$

Moreover we have P- almost surely

$$\mu = \mathcal{L}(Y^1|\mu) = \mathcal{L}(Y^2|\mu) \text{ and } \mathcal{L}(Y|\mu) = \mu \otimes \mu.$$

Writing $\mu_t := \int_{D(\mathbb{R}_+,\mathbb{R})} \mu(d\gamma) \delta_{\gamma_t}$ for the projection onto the t-th time coordinate, we introduce the filtration

$$\mathcal{G}_t = \sigma(Y_s, s \leqslant t) \vee \sigma(\mu_s(f), s \leqslant t).$$

Definition 2.2. We say that $Q \in \mathcal{P}(\mathcal{P}(D(\mathbb{R}_+, \mathbb{R})))$ is a solution to the martingale problem (\mathcal{M}) if the following holds.

- (i) $Q-almost\ surely,\ \mu_0=\nu_0.$
- (ii) For all $g \in C_b^2(\mathbb{R}^2)$, $M_t^g := g(Y_t) g(Y_0) \int_0^t Lg(\mu_s, Y_s) ds$ is a $(P, (\mathcal{G}_t)_t)$ -martingale, where

$$Lg(\mu, x) = -\alpha x^{1} \partial_{x^{1}} g(x) - \alpha x^{2} \partial_{x^{2}} g(x) + \frac{\sigma^{2}}{2} \mu(f) \sum_{i,j=1}^{2} \partial_{x^{i}x^{j}}^{2} g(x) + f(x^{1}) (g(0, x^{2}) - g(x)) + f(x^{2}) (g(x^{1}, 0) - g(x)).$$

Remark 2.3. It is not clear if the martingale problem is well-posed, but we are not interested in proving uniqueness for it. However, we will have uniqueness within the class of all possible limits in distribution of μ^N . More precisely, we shall prove that, if μ is a limit in distribution of μ^N such that $\mathcal{L}(\mu)$ is solution to (\mathcal{M}) , then $\mu = \mathcal{L}(\bar{X}|\mathcal{W})$, with \bar{X} the strong solution of (9). Equivalently, defining the problem (\mathcal{M}) for all finite-dimensional distributions, and not only for two coordinates, where Y^i $(i \ge 1)$ are defined as a mixture directed by μ , would lead to uniqueness.

Let $(\bar{X}^i)_{i\geqslant 1}$ be the solution of the limit system (8) and $\mu = \mathcal{L}(\bar{X}^1|\mathcal{W})$. Then we already know that $\mathcal{L}(\mu)$ is a solution of (\mathcal{M}) . Let us now characterise any possible solution of (\mathcal{M}) .

Lemma 2.4. Grant Assumption 3. Let $Q \in \mathcal{P}(\mathcal{P}(D(\mathbb{R}_+,\mathbb{R})))$ be a solution of (\mathcal{M}) . Let (μ,Y) be the canonical variable defined above, and write $Y = (Y^1,Y^2)$. Then there exists a standard $(\mathcal{G}_t)_t$ -Brownian motion W and on an extension $(\tilde{\Omega},(\tilde{\mathcal{G}}_t)_t,\tilde{P})$ of $(\Omega,(\mathcal{G}_t)_t,P)$ there exist $(\tilde{\mathcal{G}}_t)_t$ -Poisson random measures π^1,π^2 on $\mathbb{R}_+ \times \mathbb{R}_+$ having Lebesgue intensity such that W,π^1 and π^2 are independent and

$$dY_t^1 = -\alpha Y_t^1 dt + \sigma \sqrt{\mu_t(f)} dW_t - Y_{t-}^1 \int_{\mathbb{R}_+} \mathbb{1}_{\left\{z \le f(Y_{t-}^1)\right\}} \pi^1(dt, dz),$$

$$dY_t^2 = -\alpha Y_t^2 dt + \sigma \sqrt{\mu_t(f)} dW_t - Y_{t-}^2 \int_{\mathbb{R}_+} \mathbb{1}_{\left\{z \le f(Y_{t-}^2)\right\}} \pi^2(dt, dz).$$

Proof. Item (ii) of of (\mathcal{M}) together with Theorem II.2.42 of Jacod and Shiryaev (2003) imply that Y is a semimartingale with characteristics (B, C, ν) given by

$$\begin{split} B_t^i &= -\alpha \int_0^t Y_s^i ds - \int_0^t Y_s^i f(Y_s^i) ds, \quad 1 \leqslant i \leqslant 2, \\ C_t^{i,j} &= \int_0^t \mu_s(f) ds, \quad 1 \leqslant i, j \leqslant 2, \\ \nu(dt, dy) &= dt (f(Y_{t-}^1) \delta_{(-Y_s^{1-}, 0)}(dy) + f(Y_{t-}^2) \delta_{(0, -Y_s^{2-})}(dy)). \end{split}$$

Then we can use the canonical representation of Y (see Theorem II.2.34 of Jacod and Shiryaev (2003)) with the truncation function h(y) = y for every y: $Y_t - Y_0 - B_t = M_t^c + M_t^d$, where M^c is a continuous local martingale and M^d a purely discontinuous local martingale. By definition of the characteristics, $\langle M^{c,i}, M^{c,j} \rangle_t = C_t^{i,j}$. In particular, $\langle M^{c,i} \rangle_t = \int_0^t \mu_s(f) ds$ (i = 1, 2). Consequently, applying Theorem II.7.1 of Ikeda and Watanabe (1989) to both coordinates, we know that there exist Brownian motions W^1, W^2 such that

$$M_t^{c,i} = \int_0^t \sqrt{\mu_s(f)} dW_s^i, \quad i = 1, 2.$$

We now prove that $W^1=W^2$. Let ρ be the correlation between W^1 and W^2 . Classical computations give $\langle W^1,W^2\rangle_t=\rho t$, implying that $\langle M^{c,1},M^{c,2}\rangle_t=\rho\int_0^t\mu_s(f)ds$. In addition $\langle M^{c,1},M^{c,2}\rangle_t=C_t^{1,2}=\int_0^t\mu_s(f)ds$, and this implies that $\rho=1$ and $W^1=W^2$, since $\int_0^t\mu_s(f)ds>0$ because f is lower-bounded.

We now prove the existence of the independent Poisson measures π^1, π^2 . We know that $M^d = h * (\mu^Y - \nu)$, where $\mu^Y = \sum_s \mathbb{1}_{\{\Delta Y_s \neq 0\}} \delta_{(s,Y_s)}$ is the jump measure of Y and ν is its compensator. We rely on Theorem II.7.4 of Ikeda and Watanabe (1989). Using the notation therein, we introduce $Z = \mathbb{R}_+$, m Lebesgue measure on Z and

$$\theta(t,z) := (-Y_{t-}^1,0) \mathbb{1}_{\left\{z \leqslant f(Y_{t-}^1)\right\}} + (0,-Y_{t-}^2) \mathbb{1}_{\left\{||f||_{\infty} < z \leqslant ||f||_{\infty} + f(Y_{t-}^2)\right\}}.$$

According to Theorem II.7.4 of Ikeda and Watanabe (1989), there exists a Poisson measure π on $\mathbb{R}_+ \times \mathbb{R}_+$ having intensity $dt \cdot dz$ such that, for all $E \in \mathcal{B}(\mathbb{R}^2)$,

$$\mu^{Y}([0,t] \times E) = \int_{0}^{t} \int_{0}^{\infty} \mathbb{1}_{\{\theta(s,z) \in E\}} \pi(ds,dz). \tag{13}$$

Now let us consider two independent Poisson measures $\tilde{\pi}^1, \tilde{\pi}^2$ (independent of everything else) on $[||f||_{\infty}, \infty[$ having Lebesgue intensity. We then define π^1 in the following way: for any $A \in \mathcal{B}(\mathbb{R}_+ \times [0, ||f||_{\infty}]), \ \pi^1(A) = \pi(A), \ \text{and for } A \in \mathcal{B}(\mathbb{R}_+ \times]||f||_{\infty}, \infty[), \ \pi^1(A) = \tilde{\pi}^1(A).$ We define π^2 in a similar way: For $A \in \mathcal{B}(\mathbb{R}_+ \times [0, ||f||_{\infty}]), \ \pi^2(A) = \pi(\{(t, ||f||_{\infty} + z) : (t, z) \in A\}), \ \text{and for } A \in \mathcal{B}(\mathbb{R}_+ \times]||f||_{\infty}, \infty[), \ \pi^2(A) = \tilde{\pi}^2(A).$ By definition of Poisson measures, π^1 and π^2 are independent Poisson measures on \mathbb{R}^2_+ having Lebesgue intensity, and together with (13), we have

$$M_t^{d,i} = -\int_{[0,t]\times\mathbb{R}_+} Y_{s-}^i \mathbb{1}_{\left\{z \leqslant f(Y_{s-}^i)\right\}} \pi^i(ds,dz) + \int_0^t Y_s^i f(Y_s^i) ds, \quad 1 \leqslant i \leqslant 2.$$

Moreover we have the following

Theorem 2.5. Assume that Assumptions 1, 2 and 3 hold. Then the distribution of any limit μ of the sequence $\mu^N := \frac{1}{N} \sum_{j=1}^N \delta_{X^{N,j}}$ is solution of item (ii) of (\mathcal{M}) .

Proof. Step 1. We first check that for any $t \geq 0$, a.s., $\mu(\{\gamma : \Delta \gamma(t) \neq 0\}) = 0$. We assume by contradiction that there exists t > 0 such that $\mu(\{\gamma : \Delta \gamma(t) \neq 0\}) > 0$ with positive probability. Hence there are a, b > 0 such that the event $E := \{\mu(\{\gamma : |\Delta \gamma(t)| > a\}) > b\}$ has a positive probability. For every $\varepsilon > 0$, we have $E \subset \{\mu(\mathcal{B}_a^{\varepsilon}) > b\}$, where $\mathcal{B}_a^{\varepsilon} := \{\gamma : \sup_{s \in (t-\varepsilon, t+\varepsilon)} |\Delta \gamma(s)| > a\}$, which is an open subset of $D(\mathbb{R}_+, \mathbb{R})$. Thus $\mathcal{P}_{a,b}^{\varepsilon} := \{\mu \in \mathcal{P}(D(\mathbb{R}_+, \mathbb{R})) : \mu(\mathcal{B}_a^{\varepsilon}) > b\}$ is an open subset of $\mathcal{P}(D(\mathbb{R}_+, \mathbb{R}))$. The Portmanteau theorem implies then that for any $\varepsilon > 0$,

$$\liminf_{N \to \infty} P(\mu^N \in \mathcal{P}_{a,b}^{\varepsilon}) \geqslant P(\mu \in \mathcal{P}_{a,b}^{\varepsilon}) \geqslant P(E) > 0.$$
 (14)

Firstly, we can write

$$J^{N,\varepsilon,i} := \sup_{t-\varepsilon < \varepsilon, t+\varepsilon} \left| \Delta X^{N,i}_s \right| = G^{\varepsilon,i}_N \vee S^{\varepsilon}_N,$$

where $G_N^{\varepsilon,i}:=\max_{s\in D_N^{\varepsilon,i}}|X_{s-}^{N,i}|$ is the maximal height of the big jumps of $X^{N,i}$, with $D_N^{\varepsilon,i}:=\{t-\varepsilon\leqslant s\leqslant t+\varepsilon:\pi^i(\{s\}\times[0,f(X_{s-}^{N,i})]\times\mathbb{R}_+)\neq 0\}$. Moreover, $S_N^\varepsilon:=\max\{|U^j(s)|/\sqrt{N}:s\in\bigcup_{1\leqslant j\leqslant N}D_N^{\varepsilon,j}\}$ is the maximal height of the small jumps of $X^{N,i}$, where $U^j(s)$ is defined for $s\in D_N^{\varepsilon,j}$, almost surely, as the only real number that satisfies $\pi^j(\{s\}\times[0,f(X_{s-}^{N,j})]\times\{U^j(s)\})=1$.

We have that

$$\left\{\mu^N(\mathcal{B}_a^\varepsilon)>b\right\}=\left\{\frac{1}{N}\sum_{i=1}^N\mathbbm{1}_{\{J^{N,\varepsilon,j}>a\}}>b\right\}.$$

Consequently, by exchangeability and Markov's inequality,

$$P\left(\mu^{N}(\mathcal{B}_{a}^{\varepsilon}) > b\right) \leqslant \frac{1}{h} \mathbb{E}\left[\mathbb{1}_{\{J^{N,\varepsilon,1} > a\}}\right] = \frac{1}{h} P\left(J^{N,\varepsilon,1} > a\right) \leqslant \frac{1}{h} \left(P\left(G_{N}^{\varepsilon,1} > a\right) + P\left(S_{N}^{\varepsilon} > a\right)\right). \tag{15}$$

The number of big jumps of $X^{N,1}$ in $]t - \varepsilon, t + \varepsilon[$ is smaller than a random variable ξ having Poisson distribution with parameter $2\varepsilon||f||_{\infty}$. Hence

$$P\left(G_{N,1}^{\varepsilon} > a\right) \leqslant P\left(\xi \geqslant 1\right) = 1 - e^{2\varepsilon||f||_{\infty}} \leqslant 2\varepsilon||f||_{\infty}. \tag{16}$$

The small jumps that occur in $]t - \varepsilon, t + \varepsilon[$ are included in $\{U_1/\sqrt{N}, ..., U_K/\sqrt{N}\}$ where K is a \mathbb{N} -valued random variable having Poisson distribution with parameter $2\varepsilon N||f||_{\infty}$, which is independent of the variables U_i ($i \ge 1$) that are i.i.d. with distribution ν . Hence,

$$P\left(S_N^\varepsilon>a\right)\leqslant P\left(\max_{1\leqslant i\leqslant K}\frac{|U_i|}{\sqrt{N}}>a\right)\leqslant \mathbb{E}\left[\left.\mathbb{P}\left(\max_{1\leqslant i\leqslant K}\frac{|U_i|}{\sqrt{N}}>a\right|K\right)\right]=\mathbb{E}\left[\psi(K)\right],$$

where $\psi(k) = P\left(\max_{1 \leq i \leq k} |U_i| > a\sqrt{N}\right) \leq kP\left(|U_1| > a\sqrt{N}\right) \leq ka^{-2}N^{-1}\mathbb{E}\left[U_1^2\right]$. Hence

$$P\left(S_{N}^{\varepsilon} > a\right) \leqslant \frac{\mathbb{E}\left[U_{1}^{2}\right]}{Na^{2}} \mathbb{E}\left[K\right] \leqslant 2||f||_{\infty} \mathbb{E}\left[U_{1}^{2}\right] \frac{1}{a} \varepsilon. \tag{17}$$

Inserting the bounds (16) and (17) in (15), we have

$$P\left(\mu^N(\mathcal{B}_a^{\varepsilon} > b)\right) \leqslant C\varepsilon,$$

where C does not depend on N nor ε . This last inequality is in contradiction with (14) since P(E) does not depend on ε .

Step 2. In the following, we note $\partial^2 \varphi := \sum_{i,j=1}^2 \partial^2_{x^i x^j} \varphi$. For any $0 \leq s_1 < \ldots < s_k < s < t$, any $\varphi_1, \ldots, \varphi_k, \psi_1, \ldots, \psi_k \in C_b(\mathbb{R})$, any $\varphi \in C_b^3(\mathbb{R}^2)$, we introduce

$$F(\mu) := \psi_1(\mu_{s_1}(f)) \dots \psi_k(\mu_{s_k}(f)) \int_{D(\mathbb{R}_+, \mathbb{R})^2} \mu \otimes \mu(d\gamma) \varphi_1(\gamma_{s_1}) \dots \varphi_k(\gamma_{s_k})$$

$$\left[\varphi(\gamma_t) - \varphi(\gamma_s) + \alpha \int_s^t \gamma_r^1 \partial_{x^1} \varphi(\gamma_r) dr + \alpha \int_s^t \gamma_r^2 \partial_{x^2} \varphi(\gamma_r) dr - \frac{\sigma^2}{2} \int_s^t \mu_r(f) \partial^2 \varphi(\gamma_r) dr - \int_s^t f(\gamma_r^1) (\varphi(0, \gamma_r^2) - \varphi(\gamma_r)) dr - \int_s^t f(\gamma_r^2) (\varphi(\gamma_r^1, 0) - \varphi(\gamma_r)) dr \right].$$

To show that $\mathcal{L}(\mu)$ is solution of item (ii) of the martingale problem (\mathcal{M}) , by a classical density argument, it is sufficient to prove that $\mathbb{E}[F(\mu)] = 0$. We have

$$\begin{split} F(\mu^N) &= \psi_1(\mu^N_{s_1}(f)) \dots \psi_k(\mu^N_{s_k}(f)) \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \varphi_1(X^{N,i}_{s_1}, X^{N,j}_{s_1}) \dots \varphi_k(X^{N,i}_{s_k}, X^{N,j}_{s_k}) \cdot \\ & \left[\varphi(X^{N,i}_t, X^{N,j}_t) - \varphi(X^{N,i}_s, X^{N,j}_s) + \alpha \int_s^t X^{N,i}_r \partial_{x^1} \varphi(X^{N,i}_r, X^{N,j}_r) dr \right. \\ & \left. + \alpha \int_s^t X^{N,j}_r \partial_{x^2} \varphi(X^{N,i}_r, X^{N,j}_r) dr - \frac{\sigma^2}{2} \int_s^t \mu^N_r(f) \partial^2 \varphi(X^{N,i}_r, X^{N,j}_r) dr \right. \\ & \left. - \int_s^t f(X^{N,i}_r) (\varphi(0, X^{N,j}_r) - \varphi(X^{N,i}_r, X^{N,j}_r)) dr - \int_s^t f(X^{N,j}_r) (\varphi(X^{N,i}_r, X^{N,j}_r)) - \varphi(X^{N,i}_r, X^{N,j}_r) dr \right]. \end{split}$$

But recalling (7) and using Ito's formula, for any $i \neq j$, we have

$$\begin{split} \varphi(X_t^{N,i},X_t^{N,j}) \\ &= \varphi(X_s^{N,i},X_s^{N,j}) - \alpha \int_s^t X_r^{N,i} \partial_{x^1} \varphi(X_r^{N,i},X_r^{N,j}) dr - \alpha \int_s^t X_r^{N,j} \partial_{x^2} \varphi(X_r^{N,i},X_r^{N,j}) dr \end{split}$$

$$\begin{split} &+ \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,i})\right\}} \left[\varphi\left(0, X_{r-}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \pi^{i}(dr, dz, du) \\ &+ \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,j})\right\}} \left[\varphi\left(X_{r-}^{N,i} + \frac{u}{\sqrt{N}}, 0\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \pi^{j}(dr, dz, du) \\ &+ \sum_{\substack{k=1 \\ k \notin \{i,j\}}}^{N} \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,k})\right\}} \left[\varphi\left(X_{r-}^{N,i} + \frac{u}{\sqrt{N}}, X_{r-}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \pi^{k}(dr, dz, du). \end{split}$$

We use the notation $\tilde{\pi}^j(dr,dz,du)=\pi^j(dr,dz,du)-drdz\nu(du)$ and set

$$\begin{split} M_{s,t}^{N,i,j,1} &:= \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,i})\right\}} \left[\varphi\left(0, X_{r-}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \tilde{\pi}^{i}(dr,dz,du), \\ M_{s,t}^{N,i,j,2} &:= \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,j})\right\}} \left[\varphi\left(X_{r-}^{N,i} + \frac{u}{\sqrt{N}}, 0\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \tilde{\pi}^{j}(dr,dz,du), \\ W_{s,t}^{N,i,j} &:= \sum_{k=1}^{N} \int_{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r-}^{N,k})\right\}} \left[\varphi\left(X_{r-}^{N,i} + \frac{u}{\sqrt{N}}, X_{r-}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \right] \tilde{\pi}^{k}(dr,dz,du), \\ - \varphi(X_{r-}^{N,i}, X_{r-}^{N,j}) \left[\varphi\left(0, X_{r}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(0, X_{r}^{N,j}) \right] d\nu(u) dr, \\ \Delta_{s,t}^{N,i,j,2} &:= \int_{s}^{t} \int_{\mathbb{R}} f(X_{r}^{N,i}) \left[\varphi\left(X_{r}^{N,i} + \frac{u}{\sqrt{N}}, 0\right) - \varphi(X_{r}^{N,i}, 0) \right] d\nu(u) dr, \\ \Gamma_{s,t}^{N,i,j} &:= \sum_{k=1}^{N} \int_{s}^{t} \int_{\mathbb{R}} f(X_{r}^{N,k}) \left[\varphi\left(X_{r}^{N,i} + \frac{u}{\sqrt{N}}, X_{r}^{N,j} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r}^{N,i}, X_{r}^{N,j}) - \frac{u}{\sqrt{N}} \partial_{x^{2}} \varphi(X_{r}^{N,i}, X_{r}^{N,j}) \right] d\nu(u) dr \\ - \int_{s}^{t} \int_{\mathbb{R}} \frac{u^{2}}{2} \partial^{2} \varphi(X_{r}^{N,i}, X_{r}^{N,j}) \frac{1}{N} \sum_{k=1}^{N} f(X_{r}^{N,k}) d\nu(u) dr, \\ R_{s,t}^{N,i,j} &:= \frac{\sigma^{2}}{2} \int_{s}^{t} \delta^{2} \varphi(X_{r}^{N,i}, X_{r}^{N,j}) \left(\frac{1}{N} \sum_{k=1}^{N} f(X_{r}^{N,k}) - \frac{1}{N} \sum_{k=1}^{N} f(X_{r}^{N,k}) \right) dr. \end{split}$$

Finally, for i = j, we have

$$\begin{split} \varphi(X_t^{N,i},X_t^{N,i}) &= \varphi(X_s^{N,i},X_s^{N,i}) \\ &- \alpha \int_s^t X_r^{N,i} \partial_{x^1} \varphi(X_r^{N,i},X_r^{N,i}) dr - \alpha \int_s^t X_r^{N,i} \partial_{x^2} \varphi(X_r^{N,i},X_r^{N,i}) dr \end{split}$$

$$+ \int_{\substack{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}}} \mathbb{1}_{\left\{z \leqslant f(X_{r-}^{N,i})\right\}} \left[\varphi\left(0,0\right) - \varphi(X_{r-}^{N,i},X_{r-}^{N,i}) \right] \pi^{i}(dr,dz,du)$$

$$+ \sum_{\substack{k=1 \\ k \neq i}}^{N} \int_{\substack{]s,t] \times \mathbb{R}_{+} \times \mathbb{R}}} \mathbb{1}_{\left\{z \leqslant f(X_{r-}^{N,k})\right\}} \left[\varphi\left(X_{r-}^{N,i} + \frac{u}{\sqrt{N}},X_{r-}^{N,i} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r-}^{N,i},X_{r-}^{N,i}) \right] \pi^{k}(dr,dz,du).$$

The associated martingales and error terms are given by

$$\begin{split} M_{s,t}^{N,i} &:= \int_{]s,t] \times \mathbb{R}_+ \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r^-}^{N,i})\right\}} \left[\varphi\left(0,0\right) - \varphi(X_{r^-}^{N,i},X_{r^-}^{N,i}) \right] \tilde{\pi}^i(dr,dz,du), \\ W_{s,t}^{N,i} &:= \sum_{k=1}^N \int_{]s,t] \times \mathbb{R}_+ \times \mathbb{R}} \mathbbm{1}_{\left\{z \leqslant f(X_{r^-}^{N,k})\right\}} \left[\varphi\left(X_{r^-}^{N,i} + \frac{u}{\sqrt{N}},X_{r^-}^{N,i} + \frac{u}{\sqrt{N}}\right) - \varphi(X_{r^-}^{N,i},X_{r^-}^{N,i}) \right] \tilde{\pi}^k(dr,dz,du), \\ \Delta_{s,t}^{N,i} &:= \int_s^t \int_{\mathbb{R}} f(X_r^{N,i}) \left[\varphi\left(0,0\right) - \varphi(0,X_r^{N,i}) - \varphi(X_r^{N,i},0) + \varphi(X_r^{N,i},X_r^{N,i}) \right] d\nu(u) dr, \\ \Gamma_{s,t}^{N,i} &:= \sum_{k=1}^N \int_s^t \int_{\mathbb{R}} f(X_r^{N,k}) \left[\varphi\left(X_r^{N,i} + \frac{u}{\sqrt{N}},X_r^{N,i} + \frac{u}{\sqrt{N}}\right) - \varphi(X_r^{N,i},X_r^{N,i}) \right] d\nu(u) dr, \\ - \frac{u}{\sqrt{N}} \partial_{x^1} \varphi(X_r^{N,i},X_r^{N,i}) - \frac{u}{\sqrt{N}} \partial_{x^2} \varphi(X_r^{N,i},X_r^{N,i}) \right] d\nu(u) dr \\ - \int_s^t \int_{\mathbb{R}} \frac{u^2}{2} \partial^2 \varphi(X_r^{N,i},X_r^{N,i}) \frac{1}{N} \sum_{k=1 \atop k \neq i}^N f(X_r^{N,k}) d\nu(u) dr, \\ R_{s,t}^{N,i} &:= \frac{\sigma^2}{2} \int_s^t \partial^2 \varphi(X_r^{N,i},X_r^{N,i}) \left(\frac{1}{N} \sum_{k=1 \atop k \neq i}^N f(X_r^{N,k}) - \frac{1}{N} \sum_{k=1}^N f(X_r^{N,k}) \right) dr. \end{split}$$

Then we obtain, since $\int_{\mathbb{R}} u d\nu(u) = 0$, that

$$\begin{split} F(\mu^N) &= \psi_1(\mu^N_{s_1}(f)) \dots \psi_k(\mu^N_{s_k}(f)) \frac{1}{N^2} \sum_{i,j=1,i\neq j}^N \varphi_1(X^{N,i}_{s_1},X^{N,j}_{s_1}) \dots \varphi_k(X^{N,i}_{s_k},X^{N,j}_{s_k}) \\ & \left[M^{N,i,j,1}_{s,t} + M^{N,i,j,2}_{s,t} + W^{N,i,j}_{s,t} + \Delta^{N,i,j,1}_{s,t} + \Delta^{N,i,j,2}_{s,t} + \Gamma^{N,i,j}_{s,t} + R^{N,i,j}_{s,t} \right] \\ & + \psi_1(\mu^N_{s_1}(f)) \dots \psi_k(\mu^N_{s_k}(f)) \frac{1}{N^2} \sum_{i=1}^N \varphi_1(X^{N,i}_{s_1},X^{N,i}_{s_1}) \dots \varphi_k(X^{N,i}_{s_k},X^{N,i}_{s_k}) \\ & \left[M^{N,i}_{s,t} + W^{N,i}_{s,t} + \Delta^{N,i}_{s,t} + \Gamma^{N,i}_{s,t} + R^{N,i}_{s,t} \right]. \end{split}$$

Using exchangeability and the boundedness of the φ_j, ψ_j $(1 \leq j \leq k)$ and the fact that $M^{N,i,j,1}, \ldots, W^{N,i}$ are martingales, this implies

$$|\mathbb{E}\left[F(\mu^{N})\right]| \leqslant C\mathbb{E}\left[|\Delta_{s,t}^{N,i,j,1}| + |\Delta_{s,t}^{N,i,j,2}| + |\Gamma_{s,t}^{N,i,j}| + |R_{s,t}^{N,i,j}| + \frac{|\Delta_{s,t}^{N,i}| + |\Gamma_{s,t}^{N,i}| + |R_{s,t}^{N,i}|}{N}\right].$$

Since f is bounded and $\varphi \in C_h^3(\mathbb{R}^2)$, Taylor-Lagrange's inequality implies then that

$$|\mathbb{E}\left[F(\mu^N)\right]| \leqslant \frac{C}{\sqrt{N}}.$$

Finally, using that F is bounded and almost surely continuous at μ (see Step 1), we have

$$\mathbb{E}\left[F(\mu)\right] = \lim_{N \to \infty} \mathbb{E}\left[F(\mu^N)\right] = 0,$$

concluding our proof.

Now we have all elements to give the proof of the following main result.

Theorem 2.6. Grant Assumptions 1, 2 and 3. Each converging subsequence of $\mu^N := \frac{1}{N} \sum_{j=1}^N \delta_{X^{N,j}}$ converges to the same limit $\mu = \mathcal{L}(\bar{X}|\mathcal{W})$, where \bar{X} is the unique strong solution of (9).

Proof. Let us consider the limit (in distribution) μ of a subsequence of μ^N . By Proposition (7.20) of Aldous (1983), μ is the directing measure of some exchangeable system $(\bar{Y}^i)_{i\geqslant 1}$, and it holds that, for the chosen subsequence, $(X^{N,i})_{1\leqslant i\leqslant N}$ converges in law to $(\bar{Y}^i)_{i\geqslant 1}$. Moreover, we also know that

$$\mu = \mathcal{L}(\bar{Y}^i|\mu)$$
 and $\mu \otimes \mu = \mathcal{L}((\bar{Y}^i, \bar{Y}^j)|\mu)$,

almost surely, for all $i \neq j$. In particular, for all $i \neq j$,

$$\mathcal{L}(\mu, (\bar{Y}^i, \bar{Y}^j)) = P,$$

where P is given by (12), with $Q = \mathcal{L}(\mu)$.

Thanks to Lemma 2.4, together with Theorem 2.5, we know that there exist Brownian motions $W^{(i,j)}$ $(i,j \ge 1)$ and Poisson random measures π^i $(i \ge 1)$ such that for all pairs $(i,j), i \ne j, \pi^i$ is independent of π^j and such that

$$\begin{split} d\bar{Y}_{t}^{i} &= -\alpha \bar{Y}_{t}^{i} dt + \sigma \sqrt{\mu_{t}(f)} dW_{t}^{(i,j)} - \bar{Y}_{t-}^{i} \int_{\mathbb{R}_{+}} \mathbb{1}_{\left\{z \leqslant f(\bar{Y}_{t-}^{i})\right\}} \pi^{i}(dt, dz), \\ d\bar{Y}_{t}^{j} &= -\alpha \bar{Y}_{t}^{j} dt + \sigma \sqrt{\mu_{t}(f)} dW_{t}^{(i,j)} - \bar{Y}_{t-}^{j} \int_{\mathbb{R}_{+}} \mathbb{1}_{\left\{z \leqslant f(\bar{Y}_{t-}^{j})\right\}} \pi^{j}(dt, dz). \end{split}$$

The exchangeability of the system $(\bar{Y}^i)_{i\geqslant 1}$ implies the independence of the $(\pi^i)_{i\geqslant 1}$ and that for all $i,j,k\geqslant 1,\ i\neq j,\ i\neq k,\ W^{(i,j)}=W^{(i,k)}=W.$

The last point to prove is that $\mu_t(f) := \mathbb{E}\left[f(\bar{Y}_t^1)|\mu\right] = \mathbb{E}\left[f(\bar{Y}_t^1)|\mathcal{W}\right]$ a.s.. This would be a consequence of the fact that, conditionally on W, the processes $(\bar{Y}^j)_{j\geqslant 1}$ are i.i.d. (see Lemma (2.12).(a) of Aldous (1983)). But this last assertion is not trivial because we do not know yet that W is the only noise term common to each process \bar{Y}^j $(j\geqslant 1)$. That is why we will introduce an auxiliary particle system which is a mean field version of the limit system and which converges to $(\bar{Y}^j)_{j\geqslant 1}$.

To begin with, Lemma (2.15) of Aldous (1983) implies that $\mu_t(f)$ is the almost sure limit of $N^{-1}\sum_{j=1}^N f(\bar{Y}_t^j)$. Now, let us prove that this sequence converges to $\mathbb{E}\left[f(\bar{Y}_t^1)\big|\mathcal{W}\right]$. For this purpose, we introduce the system $(\tilde{X}^{N,i})_{1\leqslant i\leqslant N}$, driven by the same Brownian motion W and the same Poisson random measures π^i , with $\bar{Y}_0^i = \tilde{X}_0^{N,i}$ $(i \geqslant 1)$, replacing the term $\mu_t(f)$ by the empirical measure:

$$d\widetilde{X}_{t}^{N,i} = -\alpha \widetilde{X}_{t}^{N,i} dt + \sqrt{\frac{1}{N} \sum_{j=1}^{N} f(\widetilde{X}_{t}^{N,j}) dW_{t} - \widetilde{X}_{t-}^{N,i} \int_{\mathbb{R}_{+}} \mathbb{1}_{\left\{z \leq f(\widetilde{X}_{t-}^{N,i})\right\}} \pi^{i}(dt, dz), \ \widetilde{X}_{0}^{N,i} = \bar{Y}_{0}^{i}.$$

Consider finally the system $(\bar{X}^i)_{i\geqslant 1}$ defined in (8), driven by the same Brownian motion W and the same Poisson random measures π^i as $(\bar{Y}^i)_{i\geqslant 1}$, with $\bar{X}^i_0 = \bar{Y}^i_0$ for all $i\geqslant 1$. In this way, $(\bar{X}^i)_{i\geqslant 1}, (\bar{Y}^i)_{i\geqslant 1}$ and $(\tilde{X}^{N,i})_{1\leqslant i\leqslant N}$ are all defined on the same probability space.

It is now sufficient to prove that both for $(\bar{Y}^i)_{i\geq 1}$ and for $(\bar{X}^i)_{i\geq 1}$,

$$\mathbb{E}\left[\left|a(\bar{Y}_t^i) - a(\tilde{X}_t^{N,i})\right|\right] + \mathbb{E}\left[\left|a(\bar{X}_t^i) - a(\tilde{X}_t^{N,i})\right|\right] \leqslant C_t N^{-1/2}.$$
(18)

Indeed, suppose we have already proven the above control (18). Then

$$\mathbb{E}\left[\left|\frac{1}{N}\sum_{j=1}^{N}f(\bar{Y}_{t}^{j}) - \mathbb{E}\left[f(\bar{X}_{t}^{1})|\mathcal{W}\right]\right|\right] \leqslant \frac{1}{N}\sum_{j=1}^{N}\mathbb{E}\left[\left|f(\bar{Y}_{t}^{j}) - f(\tilde{X}_{t}^{N,j})\right|\right] + \frac{1}{N}\sum_{j=1}^{N}\mathbb{E}\left[\left|f(\tilde{X}_{t}^{N,j}) - f(\bar{X}_{t}^{j})\right|\right] + \mathbb{E}\left[\left|\frac{1}{N}\sum_{j=1}^{N}f(\bar{X}_{t}^{j}) - \mathbb{E}\left[f(\bar{X}_{t}^{1})|\mathcal{W}\right]\right|\right].$$

Then, (18) and Assumption 3 imply that the first term and the second one of the sum above are smaller than $C_t N^{-1/2}$ for some $C_t > 0$. In addition, by item (ii) of Proposition 1.5, the variables $(\bar{X}^j)_{1 \leq j \leq N}$ are i.i.d., conditionally on W. Consequently, the third term is also smaller than $C_t N^{-1/2}$. This implies that

$$\mu_t(f) = \mathbb{E}\left[f(\bar{Y}_t^1)|\mu\right] = \mathbb{E}\left[f(\bar{X}_t^1)|\mathcal{W}\right] = \mathbb{E}\left[f(\bar{X}_t^i)|\mathcal{W}\right] \text{ a.s..}$$

As a consequence, $(\bar{Y}^i)_{i \geq 1}$ is solution of the infinite system

$$d\bar{Y}_t^i = -\alpha \bar{Y}_t^i dt + \sigma \sqrt{\mathbb{E}\left[f(\bar{X}_t^i) \middle| \mathcal{W}\right]} dW_t - \bar{Y}_{t-}^i \int_{\mathbb{R}_+} \mathbb{1}_{\left\{z \leqslant f(\bar{Y}_{t-}^i)\right\}} \pi^i (dt, dz),$$

while $(\bar{X}^i)_{i \geq 1}$ in (8) is solution of

$$d\bar{X}_t^i = -\alpha \bar{X}_t^i dt + \sigma \sqrt{\mathbb{E}\left[\left.f(\bar{X}_t^i)\right|\mathcal{W}\right]} dW_t - \bar{X}_{t-}^i \int_{\mathbb{R}_+} \mathbb{1}_{\left\{z \leqslant f(\bar{X}_{t-}^i)\right\}} \pi^i(dt, dz),$$

with $\bar{X}_0^i = \bar{Y}_0^i$, for all $i \ge 1$.

Let us prove that $\bar{X}^i = \bar{Y}^i$ almost surely. For that sake, consider $\tau_M = \inf\{t > 0 : |\bar{X}^i_t| \wedge |\bar{Y}^i_t| > M\}$ for M > 0. We prove that $\mathbb{E}\left[|\bar{X}^i_{t \wedge \tau_M} - \bar{Y}^i_{t \wedge \tau_M}|\right] = 0$ for all M > 0, which implies, by Fatou's lemma, that $\mathbb{E}\left[|\bar{X}^i_t - \bar{Y}^i_t|\right] = 0$, recalling (10), and the fact that we can prove a similar control for \bar{Y}^i . Let $u_M(t) := \mathbb{E}\left[|\bar{X}^i_{t \wedge \tau_M} - \bar{Y}^i_{t \wedge \tau_M}|\right]$. To see that $u_M(t) = 0$, it is sufficient to apply Grönwall's lemma to the following inequality

$$u_{M}(t) \leqslant \alpha \int_{0}^{t} u_{M}(s) ds + \mathbb{E}\left[\int_{[0, t \wedge \tau_{M}] \times \mathbb{R}_{+}} \left| \bar{X}_{s-}^{i} \mathbb{1}_{\left\{z \leqslant f(\bar{X}_{s-}^{i})\right\}} - \bar{Y}_{s-}^{i} \mathbb{1}_{\left\{z \leqslant f(\bar{Y}_{s-}^{i})\right\}} \right| \pi^{i}(ds, dz)\right]$$

implying that

$$u_{M}(t) \leqslant \alpha \int_{0}^{t} u_{M}(s) ds + \mathbb{E}\left[\int_{[0, t \wedge \tau_{M}] \times \mathbb{R}_{+}} \mathbb{1}_{z \in [0, f(\bar{X}_{s-}^{i}) \wedge f(\bar{Y}_{s-}^{i})]} \left| \bar{X}_{s-}^{i} - \bar{Y}_{s-}^{i} \right| \pi^{i}(ds, dz) \right]$$

$$+ \, \mathbb{E}\left[\int_{[0,t \wedge \tau_M] \times \mathbb{R}_+} \mathbb{1}_{z \in]f(\bar{X}^i_{s-}) \wedge f(\bar{Y}^i_{s-}), f(\bar{X}^i_{s-}) \vee f(\bar{Y}^i_{s-})]} |\bar{X}^i_{s-}| \vee |\bar{Y}^i_{s-}| \pi^i(ds,dz) \right],$$

whence

$$u_M(t) \leqslant C(1+M) \int_0^t u_M(s) ds.$$

Hence $(\bar{Y}^i)_{i\geqslant 1}$ is solution of the infinite system (8) and $\mu = \mathcal{L}(\bar{Y}^1|\mathcal{W})$, its directing measure, is uniquely determined. As a consequence, μ^N converges in distribution to $\mathcal{L}(\bar{Y}^1|\mathcal{W})$ in $\mathcal{P}(D(\mathbb{R}_+,\mathbb{R}))$.

Let us now show (18). We only prove it for \bar{Y}^i , the proof for \bar{X}^i is similar. We decompose the evolution of $a(\bar{Y}_t^1)$ in the following way.

$$a(\bar{Y}_{t}^{1}) = a(\bar{Y}_{0}^{1}) - \alpha \int_{0}^{t} a'(\bar{Y}_{s}^{1}) \bar{Y}_{s}^{1} ds + \int_{[0,t] \times \mathbb{R}_{+}} \left(a(0) - a(\bar{Y}_{s-}^{1}) \right) \mathbb{1}_{\{z \leqslant f(\bar{Y}_{s-}^{1})\}} \pi^{1}(ds, dz)$$

$$+ \frac{\sigma^{2}}{2} \int_{0}^{t} a''(\bar{Y}_{s}^{1}) \frac{1}{N} \sum_{j=1}^{N} f(\bar{Y}_{s}^{j}) ds - B_{t}^{N} + \sigma \int_{0}^{t} a'(\bar{Y}_{s}^{1}) \sqrt{\frac{1}{N} \sum_{j=1}^{N} f(\bar{Y}_{s}^{j})} dW_{s} - M_{t}^{N},$$

$$(19)$$

where

$$B^N_t = \frac{\sigma^2}{2} \int_0^t a''(\bar{Y}^1_s) \left(\frac{1}{N} \sum_{j=1}^N f(\bar{Y}^j_s) - \mathbb{E}\left[\left.f(\bar{Y}^1_s)\right|\mu\right]\right) ds$$

and

$$M_t^N = \sigma \int_0^t a'(\bar{Y}_s^1) \left(\sqrt{\frac{1}{N} \sum_{j=1}^N f(\bar{Y}_s^j)} - \sqrt{\mathbb{E}\left[f(\bar{Y}_s^1) | \mu\right]} \right) dW_s.$$

Since

$$< M^N >_t \leqslant \sigma^2 \left(\sup_{x \in \mathbb{R}} |a'(x)|^2 \right) \int_0^t \left(\sqrt{\frac{1}{N} \sum_{j=1}^N f(\bar{Y}_s^j)} - \sqrt{\mathbb{E}\left[f(\bar{Y}_s^1) | \mu\right]} \right)^2 ds,$$

recalling that the variables \bar{Y}_s^j $(1 \leq j \leq N)$ are i.i.d. conditionally to μ , taking conditional expectation $\mathbb{E}(\cdot|\mu)$ implies that

$$\mathbb{E}\left[\langle M^N \rangle_t\right] \leqslant C_t N^{-1} \text{ and } \mathbb{E}\left[B_t^N\right] \leqslant C_t N^{-1}.$$

Then, applying Ito's formula on $\tilde{X}^{N,1}$, we obtain the same equation as (19), but without the terms B_t^N and M_t^N . Introducing

$$u(t) := \sup_{0 \leq s \leq t} \mathbb{E}\left[\left| a(\bar{Y}^1_s) - a(\tilde{X}^{N,1}_s) \right| \right],$$

we can prove with the same reasoning as in the proof of Theorem 1.2 that

$$u(t) \leqslant C(1+t)u(t) + \frac{C_t}{\sqrt{N}},$$

where C and C_t are independent of N. Finally, using the arguments of the proof of Theorem 1.2, this implies (18).

Let us end this section with the

Proof of Theorem 1.7. According to Proposition 2.1, the sequence $(\mu^N)_N$ is tight. Besides, thanks to Theorem 2.6, every converging subsequence of $(\mu^N)_N$ converges to the same limit $\mu = \mathcal{L}(\bar{X}^1|\mathcal{W})$, where $(\bar{X}^j)_{j\geq 1}$ is solution of (8). This implies the result.

3. Appendix

3.1. Well-posedness of the limit equation (9)

Proof of Item 2. of Theorem 1.2. The proof is done using a classical Picard-iteration. For that sake we introduce the sequence of processes $\bar{X}_t^{[0]} \equiv \bar{X}_0$, and

$$\bar{X}_t^{[n+1]} := \bar{X}_0 - \alpha \int_0^t \bar{X}_s^{[n]} ds - \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \bar{X}_{s-}^{[n+1]} \mathbbm{1}_{\{z \leqslant f(\bar{X}_{s-}^{[n]})\}} \pi(ds,dz,du) + \sigma \int_0^t \sqrt{\mu_s^n(f)} dW_s,$$

where

$$\mu_s^n = P(\bar{X}_s^{[n]} \in \cdot | \mathcal{W}).$$

Let us first prove a control on the moments of $\bar{X}^{[n]}$ uniformly in n. Applying Ito's formula we have

$$\begin{split} \mathbb{E}\left[\left(\bar{X}_t^{[n+1]}\right)^2\right] &\leqslant \mathbb{E}\left[\bar{X}_0^2\right] - 2\alpha \int_0^t \mathbb{E}\left[\left(\bar{X}_s^{[n+1]}\right)^2\right] ds + \sigma^2 \int_0^t \mathbb{E}\left[\mu_s^n(f)\right] ds \\ &\leqslant \mathbb{E}\left[\bar{X}_0^2\right] + \sigma^2 \int_0^t \mathbb{E}\left[\mu_s^n(f)\right] ds. \end{split}$$

Using that f is bounded.

$$\mathbb{E}\left[\left(\bar{X}_t^{[n+1]}\right)^2\right]\leqslant \mathbb{E}\left[\bar{X}_0^2\right]+\sigma^2||f||_{\infty}t,$$

implying

$$\sup_{n\in\mathbb{N}} \sup_{0\leqslant s\leqslant t} \mathbb{E}\left[\left(\bar{X}_s^{[n]}\right)^2\right] < +\infty. \tag{20}$$

Now, we prove the convergence of $\bar{X}_t^{[n]}$. The same strategy as the one of the proof of Item 1. of Theorem 1.2 allows to show that

$$\delta^n_t := \mathbb{E} \sup_{s < t} |a(\bar{X}^{[n]}_s) - a(\bar{X}^{[n-1]}_s)| \text{ satisfies } \delta^n_t \leqslant C(t + \sqrt{t}) \delta^{n-1}_t,$$

for all $n \ge 1$, for a constant C only depending on the parameters of the model, but not on n, neither on t. Choose t_1 such that

$$C(t_1 + \sqrt{t_1}) \leqslant \frac{1}{2}.$$

Since $\sup_{s \leq t_1} |a(\bar{X}_s^{[0]})| = a(\bar{X}_0) \leq ||a||_{\infty}$, we deduce from this that

$$\delta_{t_1}^n \leqslant 2^{-n} \|a\|_{\infty}.$$

This implies the almost sure convergence of $a\left(\bar{X}_t^{[n]}\right)_n$ to some random variable Z_t for all $t \in [0, t_1]$.

As a is an increasing function, the almost sure convergence of $\bar{X}_t^{[n]}$ to some (possibly infinite) random variable \bar{X}_t follows from this. The almost sure finiteness of \bar{X}_t is then guaranteed by Fatou's lemma and (20).

Now let us prove that \bar{X} is solution of the limit equation (9) which follows by standard arguments (note that the jump term does not cause troubles because it is of finite activity). The most important point is to notice that

$$\mu_t^n(f) = \mathbb{E}(f(\bar{X}_t^{[n]})|\mathcal{W}) \to \mathbb{E}(f(\bar{X}_t)|\mathcal{W})$$

almost surely, which follows from the almost sure convergence of $f(\bar{X}_t^{[n]}) \to f(\bar{X}_t)$, using dominated convergence.

Once the convergence is proven on the time interval $[0, t_1]$, we can proceed iteratively over successive intervals $[kt_1, (k+1)t_1]$ to conclude that \bar{X} is solution of (9) on \mathbb{R}_+ .

It remains to prove (10). Firstly, by Fatou's lemma and (20), we know that, for all t > 0,

$$\sup_{0 \le s \le t} \mathbb{E}\left[\bar{X}_s^2\right] < \infty. \tag{21}$$

Besides, Ito's formula gives

$$\begin{split} \bar{X}_{t}^{2} &= \bar{X}_{0}^{2} - 2\alpha \int_{0}^{t} \bar{X}_{s}^{2} ds - \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \bar{X}_{s-}^{2} \mathbb{1}_{\{z \leqslant f(\bar{X}_{s-})\}} \pi(ds,dz,du) \\ &+ \sigma^{2} \int_{0}^{t} \mu_{s}(f) ds + 2\sigma \int_{0}^{t} \sqrt{\mu_{s}(f)} \bar{X}_{s} dW_{s} \\ &\leq \bar{X}_{0}^{2} + \sigma^{2} ||f||_{\infty} t + 2\sigma \int_{0}^{t} \sqrt{\mu_{s}(f)} \bar{X}_{s} dW_{s}. \end{split}$$

Inequality (10) is then a straightforward consequence of Burkholder-Davis-Gundy inequality, (21) and the above computation.

We now give the

Proof of Proposition 1.5. (i) Given a Brownian motion W and i.i.d. Poisson measures π^i , the same proof as the one of Theorem 1.2 implies the existence and the uniqueness of the system given in (8) for $1 \le i \le N$.

(ii) The construction of the proof of Item 2. of Theorem 1.2, together with the proof of Theorem 1.1 of Chapter IV.1 and of Theorem 9.1 in Chapter IV.9 of Ikeda and Watanabe (1989), imply the existence of a measurable function Φ that does not depend on $k = 1, \ldots, N$, and that satisfies, for each $1 \leq k \leq N$,

$$\bar{X}^k = \Phi(\bar{X}^k_0, \pi^k, W)$$

and for all $t \ge 0$,

$$\bar{X}_{|[0,t]}^k = \Phi_t(\bar{X}_0^k, \pi_{|[0,t] \times \mathbb{R}_+ \times \mathbb{R}}^k, (W_s)_{s \leqslant t}); \tag{22}$$

in other words, our process is non-anticipative and does only depend on the underlying noise up to time t.

Then we can write, for all continuous bounded functions g, h,

$$\mathbb{E}\left[\left.g(\bar{X}^i)h(\bar{X}^j)\right|\mathcal{W}\right] = \mathbb{E}\left[\left.g(\Phi(\bar{X}^i_0,\pi^i,W))h(\Phi(\bar{X}^j_0,\pi^j,W))\right|\mathcal{W}\right] = \psi(W),$$

where $\psi(w) := \mathbb{E}\left[g(\Phi(\bar{X}_0^i, \pi^i, w))h(\Phi(\bar{X}_0^j, \pi^j, w))\right] = \mathbb{E}\left[g(\Phi(\bar{X}_0^i, \pi^i, w))\right]\mathbb{E}\left[h(\Phi(\bar{X}_0^j, \pi^j, w))\right] =: \psi_i(w)\psi_j(w)$. With the same reasoning, we show that $\mathbb{E}\left[g(\bar{X}^i)\big|\mathcal{W}\right] = \psi_i(W)$ and $\mathbb{E}\left[h(\bar{X}^j)\big|\mathcal{W}\right] = \psi_j(W)$. The same arguments prove the mutual independence of $\bar{X}^1, \dots, \bar{X}^N$ conditionally to W.

(iii) Using the representation $\bar{X}_{[[0,t]}^k = \Phi_t(\bar{X}_0^k, \pi^k, W)$, we can write for any continuous and bounded function $g: D([0,t],\mathbb{R}) \to \mathbb{R}$,

$$\int_{\mathbb{R}} g d(N^{-1} \sum_{i=1}^{N} \delta_{\bar{X}^{i}_{||0,t|}}) = \frac{1}{N} \sum_{i=1}^{N} g(\bar{X}^{i}_{|[0,t]}) = \frac{1}{N} \sum_{i=1}^{N} g \circ \Phi_{t}(\bar{X}^{i}_{0}, \pi^{i}, W).$$

Using the law of large numbers on the account of the sequence of i.i.d. PRM's and working conditionally on W, we obtain that

$$\lim_{N \to \infty} \int_{\mathbb{R}} g d(N^{-1} \sum_{i=1}^{N} \delta_{\bar{X}^{i}_{|[0,t]}}) = \mathbb{E}\left[g \circ \Phi_{t}(\bar{X}^{1}_{0}, \pi^{1}, W) | \mathcal{W}\right] = \mathbb{E}\left[g(\bar{X}^{1}_{|[0,t]}) | \mathcal{W}\right] = \mathbb{E}\left[g(\bar{X}^{1}_{|[0,t]}) | (W_{s})_{s \leqslant t}\right],$$

where we have used
$$(22)$$
.

3.2. Proof of Corollary 1.6

Applying Ito's formula, we have

$$\varphi(\bar{X}_t) = \varphi(\bar{X}_0) + \int_0^t \left(-\alpha \varphi'(\bar{X}_s) \bar{X}_s + \frac{1}{2} \varphi''(\bar{X}_s) \mu_s(f) \right) ds + \int_0^t \varphi'(\bar{X}_s) \sqrt{\mu_s(f)} dW_s$$

$$+ \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \mathbb{1}_{\{z \le f(\bar{X}_{s-})\}} \left(\varphi(0) - \varphi(\bar{X}_{s-}) \pi(ds, dz, du) \right). \tag{23}$$

Since φ', φ'' and f are bounded, it follows from (10) and Fubini's theorem that

$$\mathbb{E}\left(\int_0^t \left(-\alpha\varphi'(\bar{X}_s)\bar{X}_s + \frac{1}{2}\varphi''(\bar{X}_s)\mu_s(f)\right)ds|W\right) = \int_0^t E\left(-\alpha\varphi'(\bar{X}_s)\bar{X}_s + \frac{1}{2}\varphi''(\bar{X}_s)\mu_s(f)|W\right)ds$$
$$= \int_0^t \int_{\mathbb{R}}\left(-\alpha\varphi'(x)x + \frac{1}{2}\varphi''(x)\mu_s(f)\right)\mu_s(dx)ds.$$

Moreover, by independence of \bar{X}_0 and W, $\mathbb{E}(\varphi(\bar{X}_0)|W) = \int_{\mathbb{R}} \varphi(x)\nu_0(dx)$.

To deal with the martingale part in (23), we use an Euler scheme to approximate the stochastic integral $I_t := \int_0^t \varphi'(\bar{X}_s) \sqrt{\mu_s(f)} dW_s$. For that sake, let $t_k^n := k2^{-n}t, 0 \le k \le 2^n, n \ge 1$, and define

$$I_t^n := \sum_{k=0}^{2^n-1} \varphi'(\bar{X}_{t_k^n}) \Delta_k^n, \ \Delta_k^n = \int_{t_k^n}^{t_{k+1}^n} \sqrt{\mu_s(f)} dW_s,$$

then $\mathbb{E}(|I_t - I_t^n|^2) \to 0$ as $n \to \infty$, and therefore $\mathbb{E}(I_t^n | \mathcal{W}) \to \mathbb{E}(I_t | \mathcal{W})$ in $L^2(P)$, as $n \to \infty$. But

$$\mathbb{E}(I_t^n|\mathcal{W}) = \sum_{k=0}^{2^n - 1} \mathbb{E}(\varphi'(\bar{X}_{t_k^n})|\mathcal{W})\Delta_k^n \to \int_0^t \mathbb{E}(\varphi'(\bar{X}_s)|\mathcal{W})\sqrt{\mu_s(f)}dW_s$$

in $L^2(P)$, since the sequence of processes $Y^n_s := \sum_{k=0}^{2^n-1} \mathbbm{1}_{]t^n_k,t^n_{k+1}]}(s) \mathbb{E}(\varphi'(\bar{X}_{t^n_k})|\mathcal{W}), 0 \leqslant s \leqslant t$, converges in $L^2(\Omega \times [0,t])$ to $\mathbb{E}(\varphi'(\bar{X}_s)|\mathcal{W})$.

We finally deal with the jump part in (23). Since f is bounded, and by independence of W and π , we can rewrite this part in terms of an underlying Poisson process N_t , independent of W and having rate $||f||_{\infty}$, and in terms of i.i.d. variables $(V_n)_{n\geqslant 1}$ uniformly distributed on [0,1], independent of W and of N as follows.

$$\int_{[0,t]\times\mathbb{R}_{+}\times\mathbb{R}} \mathbb{1}_{\{z\leqslant f(\bar{X}_{s-})\}} \left(\varphi(0) - \varphi(\bar{X}_{s-})\pi(ds,dz,du) = \sum_{n=1}^{N_{t}} \mathbb{1}_{\{\|f\|_{\infty}V_{n}\leqslant f(\bar{X}_{T_{n-}})\}} (\varphi(0) - \varphi(\bar{X}_{T_{n-}})).$$

Taking conditional expectation $E(\cdot|\mathcal{W})$, we obtain

$$\mathbb{E}\left(\sum_{n=1}^{N_t} \mathbb{1}_{\{\|f\|_{\infty}V_n \leqslant f(\bar{X}_{T_n-})\}}(\varphi(0) - \varphi(\bar{X}_{T_n-}))|\mathcal{W}\right) = \\ \mathbb{E}\left(\sum_{n=1}^{N_t} \frac{f(\bar{X}_{T_n-})}{\|f\|_{\infty}}(\varphi(0) - \varphi(\bar{X}_{T_n-}))|\mathcal{W}\right) \\ = \int_0^t \mathbb{E}\left(f(\bar{X}_s)(\varphi(0) - \varphi(\bar{X}_s))|\mathcal{W}\right) ds,$$

where we have used the independence properties of $(V_n)_n, N_t$ and W and the fact that conditionally on $\{N_t = n\}$, the jump times (T_1, \ldots, T_n) are distributed as the order statistics of n i.i.d. times which are uniformly distributed on [0, t]. This concludes our proof.

3.3. A priori estimates

In this subsection, we prove useful a priori upper bounds on some moments of the solutions of the SDEs (7).

Lemma 3.1. Assume that 2 holds and that f is bounded:

(i) for all
$$t > 0$$
, $\sup_{N \in \mathbb{N}} \sup_{0 \le s \le t} \mathbb{E}\left[\left(X_s^{N,1}\right)^2\right] < +\infty$,

$$\begin{split} & (i) \ \, \textit{for all} \, \, t > 0, \, \, \sup_{N \in \mathbb{N}^*} \sup_{0 \leqslant s \leqslant t} \mathbb{E} \left[\left(X_s^{N,1} \right)^2 \right] < + \infty, \\ & (ii) \ \, \textit{for all} \, \, t > 0, \, \sup_{N \in \mathbb{N}^*} \mathbb{E} \left[\sup_{0 \leqslant s \leqslant t} \left| X_s^{N,1} \right| \right] < + \infty, \end{split}$$

Proof. Step 1: Let us prove (i).

$$\begin{split} \left(X_{t}^{N,1}\right)^{2} &= \left(X_{0}^{N,1}\right)^{2} - 2\alpha \int_{0}^{t} \left(X_{s}^{N,1}\right)^{2} ds - \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \left(X_{s}^{N,1}\right)^{2} \mathbb{1}_{\left\{z \leqslant f\left(X_{s-}^{N,1}\right)\right\}} d\pi^{j}(s,z,u) \\ &+ \sum_{j=2}^{N} \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \left[\left(X_{s-}^{N,1} + \frac{u}{\sqrt{N}}\right)^{2} - \left(X_{s-}^{N,1}\right)^{2} \right] \mathbb{1}_{\left\{z \leqslant f\left(X_{s-}^{N,j}\right)\right\}} d\pi^{j}(s,z,u) \\ &\leqslant \left(X_{0}^{N,1}\right)^{2} + \sum_{j=2}^{N} \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \left[\left(X_{s-}^{N,1} + \frac{u}{\sqrt{N}}\right)^{2} - \left(X_{s-}^{N,1}\right)^{2} \right] \mathbb{1}_{\left\{z \leqslant f\left(X_{s-}^{N,j}\right)\right\}} d\pi^{j}(s,z,u). \end{split}$$

As f is bounded,

$$\mathbb{E}\left[\left(X_t^{N,1}\right)^2\right] \leqslant \mathbb{E}\left[\left(X_0^{N,1}\right)^2\right] + \frac{\sigma^2}{N} \sum_{j=2}^N \int_0^t \mathbb{E}\left[f\left(X_s^{N,j}\right)\right] ds \leqslant \mathbb{E}\left[\left(X_0^{N,1}\right)^2\right] + \sigma^2 ||f||_{\infty} t.$$

Step 2: Now we prove (ii).

$$\left|X_{t}^{N,1}\right| \leqslant \left|X_{0}^{N,1}\right| + \alpha \int_{0}^{t} \left|X_{s}^{N,1}\right| ds + \int_{[0,t] \times \mathbb{R}_{+} \times \mathbb{R}} \left|X_{s-}^{N,1}\right| \mathbb{1}_{\left\{z \leqslant f(X_{s-}^{N,1})\right\}} d\pi^{1}(s,z,u) + \frac{1}{\sqrt{N}} |M_{t}^{N}|,$$

where M_t^N is the martingale $M_t^N = \sum_{j=2}^N \int_{[0,t]\times\mathbb{R}_+\times\mathbb{R}} u \mathbb{1}_{\{z\leqslant f(X_{s-}^{N,j})\}} d\pi^j(s,z,u)$. Then

$$\begin{split} \sup_{0 \leqslant s \leqslant t} \left| X_s^{N,1} \right| \leqslant \left| X_0^{N,1} \right| + \alpha \int_0^t \left| X_s^{N,1} \right| ds + \int_{[0,t] \times \mathbb{R}_+ \times \mathbb{R}} \left| X_{s-}^{N,1} \right| \mathbbm{1}_{\left\{z \leqslant f(X_{s-}^{N,1})\right\}} d\pi^1(s,z,u) \\ + \frac{1}{\sqrt{N}} \sup_{0 \leqslant s \leqslant t} |M_s^N|. \end{split}$$

To conclude the proof, it is now sufficient to notice that

$$\frac{1}{\sqrt{N}} \mathbb{E} \left[\sup_{0 \leqslant s \leqslant t} |M_s^N| \right] \leqslant \mathbb{E} \left[\frac{1}{N} [M^N]_t \right]^{1/2}$$

is uniformly bounded in N, since f is bounded, and to use the point (i) of the lemma.

Acknowledgements

This research has been conducted as part of FAPESP projectResearch, Innovation and Dissemination Center for Neuromathematics(grant 2013/07699-0).

Bibliography

Aldous, D. (1983). Exchangeability and related topics. In *Ecole d'Eté de Probabilités de Saint-Flour: XIII - 1983*. springer-verlag ed. *Lecture notes in mathematics* 1117.

CARMONA, R., DELARUE, F. and LACKER, D. (2016). Mean field games with common noise. *The Annals of Probability* 44 3740–3803.

Coghi, M. and Flandoli, F. (2016). Propagation of chaos for interacting particles subject to environmental noise. *The Annals of Applied Probability* **26** 1407–1442. MR3513594

CORMIER, Q., TANRÉ, E. and VELTZ, R. (2019). Long time behavior of a mean-field model of interacting neurons. arXiv:1810.08562 [math].

DE MASI, A., GALVES, A., LÖCHERBACH, E. and PRESUTTI, E. (2015). Hydrodynamic Limit for Interacting Neurons. *Journal of Statistical Physics* **158** 866–902.

Delattre, S., Fournier, N. and Hoffmann, M. (2016). Hawkes processes on large networks. The Annals of Applied Probability 26 216–261.

DERMOUNE, A. (2003). Propagation and conditional propagation of chaos for pressureless gas equations. *Probability Theory and Related Fields* **126** 459–476.

- DITLEVSEN, S. and LÖCHERBACH, E. (2017). Multi-class Oscillating Systems of Interacting Neurons. Stochastic Processes and their Applications 127 1840–1869.
- ERNY, X., LÖCHERBACH, E. and LOUKIANOVA, D. (2019). Mean field limits for interacting Hawkes processes in a diffusive regime. arXiv:1904.06985 [math].
- FOURNIER, N. and LÖCHERBACH, E. (2016). On a toy model of interacting neurons. Annales de l'Institut Henri Poincaré Probabilités et Statistiques 52 1844–1876.
- Galves, A. and Löcherbach, E. (2013). Infinite systems of interacting chains with memory of variable length a stochastic model for biological neural nets. *Journal of Statistical Physics* **151** 896–921.
- GRAHAM, C. (1992). McKean-Vlasov Ito-Skorohod equations, and nonlinear diffusions with discrete jump sets. Stochastic Processes and their Applications 40 69–82.
- Hewitt, E. and Savage, L. J. (1955). Symmetric measures on Cartesian products. *Transactions of the American Mathematical Society* **80** 470–501.
- IKEDA, N. and WATANABE, S. (1989). Stochastic Differential Equations and Diffusion Processes, Second ed. North-Holland Publishing Company.
- Jacod, J. and Shiryaev, A. N. (2003). *Limit Theorems for Stochastic Processes*, Second ed. Springer-Verlag BerlinHeidelberg NewYork.
- RISSANEN, J. (1983). A universal data compression system. *IEEE Transactions on Information Theory* **29** 656–664.
- SZNITMAN, A.-S. (1989). Topics in propagation of chaos. In *Ecole d'Eté de Probabilités de Saint-Flour: XIX 1989. Lecture notes in mathematics* **1464** 167–251. Springer, Berlin. OCLC: 23253880.