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# Optimal transportation and stationary measures for Randomly Iterated Function Systems 

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In this article, we show how ideas, methods and results from optimal transportation can be used to study various aspects of the stationary measures of Iterated Function Systems equipped with a probability distribution. We recover a classical existence and uniqueness result under a contraction-onaverage assumption, prove moment bounds and generalized moment bounds, consider the convergence of the empirical measure of an associated Markov chain, prove in many cases the Lipschitz continuity of the stationary measure when the system is perturbed, with as a consequence a "linear response formula" at almost every parameter of the perturbation, and prove singularity of the stationary measure in some cases where the classical dimension bound coincides with the dimension of the ambient space.

## 1 Introduction

Let $(X, d)$ be a complete separable metric space (endowed with its Borel $\sigma$-algebra for all measurability purposes) and $\Phi=\left\{\phi_{i}: i \in I\right\}$ be an Iterated Functions System (IFS), i.e. a family of continuous maps $\phi_{i}: X \rightarrow X$ indexed by a set $I$, either countable or endowed with a standard $\sigma$-algebra. The set of probability measures on $X$ is denoted by $\mathcal{P}(X)$.

Hutchinson [Hut81] introduced such IFS to produce fractals: under a contraction hypothesis, there is a unique compact subset $K_{\Phi}$ of $X$ such that

$$
K_{\Phi}=\bigcup_{i \in I} \phi_{i}\left(K_{\Phi}\right) .
$$

[^0]The proof is very simple: one shows that the map $K \mapsto \cup_{i} \phi_{i}(K)$ is a contraction in the Hausdorff metric, and applies the Banach fixed point theorem. Given additionally $\eta \in$ $\mathcal{P}(I)$, one is interested in existence, uniqueness and properties of a measure $\mu \in \mathcal{P}(X)$ such that

$$
\begin{equation*}
\mu=\int\left(\phi_{i}\right)_{*} \mu \mathrm{~d} \eta(i) \tag{1}
\end{equation*}
$$

i.e. $\int f(x) \mathrm{d} \mu(x)=\iint f \circ \phi_{i}(x) \mathrm{d} \mu(x) \mathrm{d} \eta(i)$ for all $f \in C_{b}(X)$, the set of bounded continuous functions $X \rightarrow \mathbb{R}$. Such a measure shall be called a stationary measure for the pair $(\Phi, \eta)$, which we shall call a Randomly Iterated Functions System (RIFS). When the $\phi_{i}$ are contractions, again existence and uniqueness mostly follow from the Banach fixed point theorem; Hutchinson used the now-called Wasserstein distance of exponent 1 (in its dual formulation, restricted to compactly supported measures). Afterward, the contraction hypothesis has been quite relaxed to a form of contraction on average. The now most common technique is to use backward iteration of the system, which to an infinite word $\omega=i_{1} \ldots i_{k} \ldots$ with letters in $I$ and a point $x \in X$ associate the points $x_{k}=\phi_{i_{1}} \circ \phi_{i_{2}} \circ \ldots \phi_{i_{k}}(x)$. One shows (under appropriate hypotheses) that if the $i_{k}$ are drawn randomly independently with law $\eta$, then $x_{k}$ converges almost surely to a random point not depending on $x$, whose law is the stationary measure. As we shall see, using general Wasserstein distances one can use the fixed point theorem approach and get moment estimates at the same time.

More generally, the goal of the present article is to apply tools and ideas from optimal transportation in this context, to show the variety of information they provide with simple (while not always elementary) proofs. General results are stated and proved all along the article, but to give a taste of the main results we start with their applications to specific examples.

### 1.1 Tail estimates for a parabolic-hyperbolic example

For $p, a \in(0,1)$, consider the RIFS given by

$$
\begin{array}{llll}
I=\{0,1\}, & \eta(\{0\})=p & \eta(\{1\})=1-p & \\
& \phi_{0}(x)=a x & \phi_{1}(x)=x+1 \quad \forall x \in \mathbb{R} \tag{2}
\end{array}
$$

When $\phi_{0}$ and $\phi_{1}$ are seen as Möbius transformations (i.e. extended as homography of the real projective line, with $\infty$ as a fixed point) or as hyperbolic isometries (i.e. extended to the Poincaré upper half plane of $\mathbb{C}$ with its hyperbolic metric), $\phi_{0}$ is hyperbolic (one attractive and one repulsive fixed points on the projective line) while $\phi_{1}$ is parabolic (a single fixed point on the projective line, which is repulsive on one side and attractive on the other side). This example is interesting in particular because it is not uniformly contracting, and cannot be made so in any set of coordinates because of the parabolic fixed point at infinity. Let $\mu_{a, p}$ denote the unique stationary measure of this RIFS; it is concentrated on $[0,+\infty)$. Its dimension and absolute continuity have for example been studied in [NSB02]. As an application of the generalized moment estimate obtained from the use of Wasserstein distances, we obtain an exponential tail estimate for $\mu_{a, p}$, sharp up to a polynomial factor.

Corollary A. For all $p, a \in(0,1)$, there exist $c, C>0$ such that for all $t \geq 1$ :

$$
c(1-p)^{t} \leq \mu_{a, p}([t,+\infty)) \leq C t^{\frac{1}{1-a}}(1-p)^{t}
$$

(The main result is the upper bound, the lower bound being trivial.)

### 1.2 Linear response for Bernoulli convolutions

The Bernoulli convolution $\mu_{\lambda}$ (where $\lambda \in(0,1)$ ) is defined as the stationary measure of the following most classical RIFS $\left(\Phi^{\lambda}, \eta\right)$ :

$$
\begin{aligned}
I & =\{0,1\}, & \eta(\{0\}) & =\eta(\{1\})=\frac{1}{2} \\
\phi_{0}^{\lambda}(x) & =\lambda x, & \phi_{1}^{\lambda}(x) & =\lambda x+(1-\lambda)
\end{aligned} \quad \forall x \in \mathbb{R} .
$$

(The precise value $1-\lambda$ of the translation part in $\phi_{1}^{\lambda}$ has no particular relevance -as soon as it is not zero- we chose this value to have a fixed attractor $[0,1]$ but it bears no consequences on the result below.)

We shall prove that the map $\lambda \mapsto \mu_{\lambda}$ is Lipschitz in the Wasserstein distances of all exponents; for exponents larger than 1, thanks to the differentiation theorem of Ambrosio, Gigli and Savaré [AGS08] this implies an almost-everywhere linear response formula ${ }^{1}$ : the map $\lambda \mapsto \mu_{\lambda}$ can be differentiated in some precise sense at almost-all $\lambda$, and while we do not get an explicit expression for the differential we show that it takes a specific form.

Corollary B. The family of Bernoulli convolutions $\left(\mu_{\lambda}\right)_{\lambda \in\left(\frac{1}{2}, 1\right)}$ is differentiable almost everywhere. More specifically, there exist a family $\left(w_{\lambda}\right)_{\lambda \in\left(\frac{1}{2}, 1\right)}$ of measurable functions from $[0,1]$ to $\mathbb{R}$ such that for Lebesgue-almost all $\lambda \in\left(\frac{1}{2}, 1\right)$ and for all $\mathcal{C}^{1}$ compactly supported $f: \mathbb{R} \rightarrow \mathbb{R}$,

$$
\left.\frac{\mathrm{d}}{\mathrm{~d} t} \int f \mathrm{~d} \mu_{t}\right|_{t=\lambda}=\int_{0}^{1} f^{\prime}(x) w_{\lambda}(x) \mathrm{d} x .
$$

Moreover, for Lebesgue-almost all $\lambda>1 / \sqrt{2}$ we have $w_{\lambda} \in L^{q}([0,1])$ for all $q>1$.
In fact, we also have differentiability for almost-all $\lambda \in\left(0, \frac{1}{2}\right)$, but to get the expression above we use that $\mu_{\lambda}$ is absolutely continuous for almost all $\lambda>\frac{1}{2}$, see below.

### 1.3 Singularity of stationary measures via approximate overlaps

The most important question about Bernoulli convolutions is to determine for which $\lambda$ is $\mu_{\lambda}$ absolutely continuous (and also for which $\lambda$ it has Hausdorff dimension 1); we recall a few facts and refer the reader to [PSS00] for more details on these. Important recent works on this topic include [Hoc14, Shm19, Var19b, Var19a].

It is easy to see that for $\lambda<\frac{1}{2}, \mu_{\lambda}$ is supported on a compact Cantor set of dimension $<1$; that $\mu_{\frac{1}{2}}$ is the uniform measure on $[0,1]$, and that the support of $\mu_{\lambda}$ is $[0,1]$ for

[^1]all $\lambda \geq \frac{1}{2}$. Erdős proved [Erd39] that when $\lambda \in\left(\frac{1}{2}, 1\right)$ is the inverse of a Pisot number, then $\mu_{\lambda}$ is singular with respect to Lebesgue; while Solomyak [Sol95] proved that for Lebesgue-almost all $\lambda \in\left(\frac{1}{2}, 1\right), \mu_{\lambda}$ is absolutely continuous. Moreover, for almost all $\lambda \in(1 / \sqrt{2}, 1)$, the density of $\mu_{\lambda}$ is continuous (see (8.1) in [PSS00]).

While Erős used Fourier analysis to prove its singularity result, another common method is to use a dimension upper bound given by the ratio of the entropy of $\eta$ and the opposite to the Lyapunov exponent of $(\Phi, \eta)$. Clearly, if the dimension of a measure on $\mathbb{R}^{d}$ is less than $d$, then the measure cannot be absolutely continuous. Except in the simplest cases, to use the dimension bound effectively one relies on exact overlaps, i.e. equality between two compositions of maps of $\Phi$ with different indices (for example in the above example $\phi_{1}^{\lambda} \phi_{0}^{\lambda} \phi_{0}^{\lambda}=\phi_{0}^{\lambda} \phi_{1}^{\lambda} \phi_{1}^{\lambda}$ when $\lambda=\frac{\sqrt{5}-1}{2}$, the inverse of the golden ratio). We show that one can use approximate overlaps, if not to improve the dimension bound, at least to show in some critical cases that the stationary measure is not absolutely continuous.

Corollary C. Let $\ell \geq 3$ be an integer. There exist a set $T=T(\ell) \subset \mathbb{R}$ with $2^{\aleph_{0}}$ elements such that for all $t_{2} \in T$, for all $\left(t_{3}, \ldots, t_{\ell-1}\right) \in \mathbb{R}^{\ell-3}$, the stationary measure $\mu$ of the RIFS

$$
\begin{array}{r}
I=\{0,1, \ldots, \ell-1\}, \quad \eta(\{i\})=\frac{1}{\ell} \quad \forall i \in I \\
\phi_{0}(x)=\frac{x}{\ell}, \quad \phi_{1}(x)=\frac{x}{\ell}+1, \quad \phi_{i}(x)=\frac{x}{\ell}+t_{i} \quad \forall i \geq 2
\end{array}
$$

is not absolutely continuous with respect to Lebesgue measure.
Note that the dimension bound alluded to above only yields here the useless inequality $\operatorname{dim} \mu \leq 1$ (the entropy is $\log \ell$ and the Lyapunov exponent $\log 1 / \ell$ ): we are in the critical case. What we do is to choose $T$ to ensure many approximate overlaps and use them to show that $\mu$ can be approximated by discrete measures slightly too well for it to be absolutely continuous. The same method can be used in higher dimension as well, see Theorem 5.2.

In the case $\ell=3$, this result does not bring new information since Kenyon proved that for all irrational $t_{2}$, the support of $\mu$ has zero Lebesgue measure [Ken97].

Structure of the article. Section 2 is a very brief introduction to optimal transportation and the Wasserstein distances. In Section 3, we start by using the completeness of Wasserstein spaces to prove existence and uniqueness of stationary measures under a hypothesis of contraction on average; this is a classical result, but we obtain moment estimates for free. Then we deal with generalized moments (Section 3.3, containing the proof of Corollary A), study the convergence toward the stationary measure of the empirical measure of the random walks associated to the RIFS (Section 3.5), and consider stationary measures in the more general setting of skew-products (Section 3.6).

In section 4 we study how much the stationary measure depends on the underlying RIFS, leading to the general linear response formula and the proof of Corollary C.

In Section 5, we explain how the problem of quantization of measures is related to the classical dimension bound, and give a criterion for singularity of a stationary measure leading to Corollary C.

## 2 Notation and definition of Wasserstein distances

We sometimes use $f \lesssim g$ as an alternative to Landau's notation $f=O(g)$, i.e. to express that for some $C$ and all $x, f(x) \leq C g(x)$. Sometimes we use $C$ to denote such constants, and its value can change from line to line.

Let us now introduce briefly the Wasserstein distances, issued from optimal transportation theory. We only mention here statements that will be used several times or are relevant to several parts of the text. Below, it will happen several times that we use a result of the literature in an crucial way without giving its full statement; we shall only do so when we can give a precise reference, use the result as it is stated without modification, and when restating it would be somewhat redundant with the corollary we get from it. This makes the present article not as self-contained as it could be, but is consistent with the purpose of showing what optimal transportation can bring to the subject and encourage the reader to learn more about it. For details and proofs of the claims made in this section, see for example [Vil09].

Let us fix a reference point $x_{0} \in X$. This choice can be arbitrary and has no conceptual bearing, but can be subject to optimization in some cases. For each $q \in(0,+\infty)$, the $q$-th moment of $\mu \in \mathcal{P}(X)$ is

$$
m_{x_{0}}^{q}(\mu):=\int d\left(x, x_{0}\right)^{q} \mathrm{~d} \mu(x) \in[0,+\infty] .
$$

The set of probability measures $\mu$ of finite $q$-th moment is denoted by $\mathcal{P}_{q}(X)$ and does not depend on $x_{0}$.

Given measures $\mu_{0}, \mu_{1} \in \mathcal{P}(X)$, the set of transport plans or couplings is the set $\Gamma\left(\mu_{0}, \mu_{1}\right)$ of measures $\gamma \in \mathcal{P}(X \times X)$ such that $\gamma(A \times X)=\mu_{0}(A)$ and $\gamma(X \times B)=\mu_{1}(B)$ for all measurable $A, B \subset X$. The number $\gamma(A \times B)$ can be interpreted as the amount of mass moved from $A$ to $B$ under the plan $\gamma$.

One defines the total cost and Wasserstein distance of exponent $q$ between two probability measures by:

$$
\begin{aligned}
\mathrm{C}_{q}\left(\nu_{0}, \nu_{1}\right) & =\inf _{\gamma \in \Gamma\left(\nu_{0}, \nu_{1}\right)} \int d(x, y)^{q} \mathrm{~d} \gamma(x, y) \\
\mathrm{W}_{q}\left(\nu_{0}, \nu_{1}\right) & =\mathrm{C}_{q}\left(\nu_{0}, \nu_{1}\right)^{\min \left(1, \frac{1}{q}\right)}
\end{aligned}
$$

Observe that the Wassertein distance of exponent $q<1$ is actually the Wassertein distance of exponent 1 of $\left(X, d^{q}\right)$. The cost and the Wasserstein distance are finite as soon as $\mu_{0}, \mu_{1} \in \mathcal{P}_{q}(X)$, and $\left(\mathcal{P}_{q}(X), \mathrm{W}_{q}\right)$ is a complete metric space. Convergence in the Wasserstein distance is stronger than weak-* convergence when $X$ is not compact; if $X$ is compact, then $\mathrm{W}_{q}$ metrizes the weak-* topology.

Note that when $\mu_{0}=\delta_{x}$ is a Dirac mass, there is only one possible coupling: $\Gamma\left(\delta_{x}, \mu_{1}\right)=$ $\left\{\delta_{x} \otimes \mu_{1}\right\}$; therefore the $q$-th moment of $\mu$ can be expressed as $m_{x_{0}}^{q}(\mu)=\mathrm{C}_{q}\left(\delta_{x_{0}}, \mu\right)$.

We denote $\operatorname{Lip}(\phi)$ the Lipschitz constant of a map $\phi: X \rightarrow \mathbb{R}$, i.e.

$$
\operatorname{Lip}(\phi)=\inf \{C \geq 0:|\phi(x)-\phi(y)| \leq C d(x, y) \forall x, y \in X\} \in[0,+\infty]
$$

The Kantorovich duality expresses that $\mathrm{W}_{q}$ coincides with the value of a "dual" optimization problem:

$$
\mathrm{W}_{1}\left(\mu_{0}, \mu_{1}\right)=\sup _{\operatorname{Lip}(f) \leq 1}\left|\int f \mathrm{~d} \mu_{0}-\int f \mathrm{~d} \mu_{1}\right| .
$$

Similarly, for all $q<1$ the Wasserstein distance $\mathrm{W}_{q}$ can be expressed as the maximal differences between the integrals of $q$-Hölder functions with Hölder constant 1, since $q$-Hölder functions are precisely the Lipschitz function of the metric $d^{q}$.

## 3 Wasserstein contraction and its consequences

Under a suitable assumptions, one can prove that the dual transfer operator associated to an RIFS is contracting in some Wasserstein distance. The completeness of the Wasserstein spaces thus makes it easy to prove existence and uniqueness of a stationary measure in $\mathcal{P}_{q}(X)$ (some technicalities are needed to prove uniqueness on the whole of $\mathcal{P}(X)$; Huntchinson restricts to compactly supported stationary measures). This is not an original result, but this contraction has several other nice consequences: we have a control on moments of the stationary measure, a good convergence of the empirical measure, and later we shall also obtain a Lipschitz dependency of the stationary measure under perturbation of the RIFS.

### 3.1 Contracting dual operator

We consider the "dual transfer operator" $\mathrm{L}^{*}$ defined on $\mathcal{P}(X)$ by

$$
\mathrm{L}^{*} \mu=\int\left(\phi_{i}\right)_{*} \mu \mathrm{~d} \eta(i),
$$

i.e. $\mathrm{L}^{*} \mu$ is the law of $\mathrm{X}_{n+1}$ if $\left(\mathrm{X}_{n}\right)_{n}$ is a Markov chain jumping from $x$ to $\phi_{i}(x)$ with probability $\mathrm{d} \eta(i)$ and $\mathrm{X}_{n} \sim \mu$. A stationary measure is precisely a $\mu \in \mathcal{P}(X)$ such that $\mathrm{L}^{*} \mu=\mu$.

The "transfer operator" (also known as Markov operator), acting on suitable spaces of functions $X \rightarrow \mathbb{R}$ is defined by

$$
\mathrm{L} f(x)=\int f \circ \phi_{i}(x) \mathrm{d} \eta(x) ;
$$

the transfer operator acts for example on the space of bounded measurable functions. It is a positive operator fixing each constant function, so that when $a \leq f \leq b$ with $a, b \in \mathbb{R}$, then also $a \leq \mathrm{L} f \leq b$. The duality relation between L and $\mathrm{L}^{*}$ is

$$
\int f \mathrm{~d}\left(\mathrm{~L}^{*} \mu\right)=\int \mathrm{L} f \mathrm{~d} \mu
$$

and is a direct consequence of Fubini's theorem.
The dual transfer operator has a natural extension to couplings, which we denote in the same way: given $\gamma \in \Gamma\left(\mu_{0}, \mu_{1}\right)$, we define

$$
\mathrm{L}^{*} \gamma=\int\left(\phi_{i} \times \phi_{i}\right)_{*} \gamma \mathrm{~d} \eta(i) \in \Gamma\left(\mathrm{L}^{*} \mu_{0}, \mathrm{~L}^{*} \mu_{1}\right)
$$

Our first main result is the following, where the control of the moments is our focus. The first hypothesis below asks for contraction on $L^{q}$ average, while the second hypothesis asks the maps not to translate a point too far away.

Theorem 3.1. Let $(\Phi, \eta)$ be a RIFS on a complete metric space $(X, d)$ and fix any $x_{0} \in X$. Assume that for some $q>0, A>0, \lambda \in(0,1)$ we have

$$
\begin{align*}
\int d\left(\phi_{i}(x), \phi_{i}(y)\right)^{q} \mathrm{~d} \eta(i) & \leq \lambda d(x, y)^{q} \quad \forall x, y \in X  \tag{3}\\
\int d\left(x_{0}, \phi_{i}\left(x_{0}\right)\right)^{q} \mathrm{~d} \eta(i) & \leq A \tag{4}
\end{align*}
$$

Then $(\Phi, \eta)$ has a unique stationary measure $\mu \in \mathcal{P}(X)$; moreover $\mu$ has finite $q$-th moment:

$$
m_{x_{0}}^{q}(\mu) \leq \begin{cases}\frac{A}{1-\lambda} & \text { when } q \leq 1 \\ \frac{A}{\left(1-\lambda^{\frac{1}{q}}\right)^{q}} & \text { when } q \geq 1\end{cases}
$$

The proof is split into a few lemmas.
Lemma 3.2. The dual transfer operator preserves $\mathcal{P}_{q}(X)$ and is a contraction of ratio $\bar{\lambda}:=\lambda^{\min \left(1, \frac{1}{q}\right)}$.

Proof. Let $\mu_{0}, \mu_{1} \in \mathcal{P}_{q}(X)$ and choose an optimal coupling $\gamma \in \Gamma\left(\mu_{0}, \mu_{1}\right)$ for $\mathrm{W}_{q}$. Then

$$
\begin{aligned}
\int d(x, y)^{q} \mathrm{~d}\left(\mathrm{~L}^{*} \gamma\right)(x, y) & =\iint d\left(\phi_{i}(x), \phi_{i}(y)\right)^{q} \mathrm{~d} \eta(i) \mathrm{d} \gamma(x, y) \\
& \leq \lambda \int d(x, y)^{q} \mathrm{~d} \gamma(x, y)
\end{aligned}
$$

so that $\mathrm{W}_{q}\left(\mathrm{~L}^{*} \mu_{0}, \mathrm{~L}^{*} \mu_{1}\right) \leq \bar{\lambda} \mathrm{W}_{q}\left(\mu_{0}, \mu_{1}\right)$.
In particular, all elements of $\mathcal{P}_{q}(X)$ are sent at finite $\mathrm{W}_{q}$-distance from $\mathrm{L}^{*} \delta_{x_{0}}$ and we only have left to prove that $\mathrm{L}^{*} \delta_{x_{0}} \in \mathcal{P}_{q}(X)$, which follows from (4):

$$
\mathrm{C}_{q}\left(\delta_{x_{0}}, \mathrm{~L}^{*} \delta_{x_{0}}\right)=\int d(x, y)^{q} \mathrm{~d}\left(\delta_{x_{0}} \otimes \mathrm{~L}^{*} \delta_{x_{0}}\right)(x, y)=\int d\left(x_{0}, \phi_{i}\left(x_{0}\right)\right) \mathrm{d} \eta(i) \leq A
$$

Lemma 3.3. There is a unique stationary measure in $\mathcal{P}_{q}(X)$. Moreover for all $\nu \in$ $\mathcal{P}_{q}(X)$, we have $\mathrm{L}^{* k} \nu \rightarrow \mu$ in the distance $\mathrm{W}_{q}$, exponentially fast.

Proof. Follows from Lemma 3.2 and the Banach fixed point theorem.
Lemma 3.4. For all continuous bounded function $f: X \rightarrow \mathbb{R}$, we have $\mathrm{L}^{k} f(x) \rightarrow \int f \mathrm{~d} \mu$ for all $x \in X$.

Proof. Since $\delta_{x} \in \mathcal{P}_{q}(X)$ we have

$$
\mathrm{L}^{k} f(x)=\int \mathrm{L}^{k} f \mathrm{~d} \delta_{x}=\int f \mathrm{~d}\left(\mathrm{~L}^{* k} \delta_{x}\right) \rightarrow \int f \mathrm{~d} \mu
$$

Lemma 3.5. Every stationary measure has finite $q$-th moment, therefore there is a unique stationary measure in $\mathcal{P}(X)$.

Proof. For each $n \in \mathbb{N}$, define a continuous bounded function by

$$
f_{n}(x)=\min \left(d\left(x_{0}, x\right)^{q}, n\right) .
$$

For all stationary measure $\mu^{\prime} \in \mathcal{P}(X)$ and all $n, k \in \mathbb{N}$ we have

$$
\int f_{n} \mathrm{~d} \mu^{\prime}=\int f_{n} \mathrm{~d}\left(\mathrm{~L}^{* k} \mu^{\prime}\right)=\int \mathrm{L}^{k} f_{n} \mathrm{~d} \mu^{\prime}
$$

Since $\mathrm{L}^{k} f_{n}$ is bounded between 0 and $n$ for all $k$, we can apply the dominated convergence theorem as $k \rightarrow \infty$, so that by Lemma 3.4

$$
\int f_{n} \mathrm{~d} \mu^{\prime}=\int \lim _{k \rightarrow \infty} \mathrm{~L}^{k} f_{n} \mathrm{~d} \mu^{\prime}=\int\left(\int f_{n} \mathrm{~d} \mu\right) \mathrm{d} \mu^{\prime}=\int f_{n} \mathrm{~d} \mu \leq \int d\left(x_{0}, x\right)^{q} \mathrm{~d} \mu<\infty
$$

The monotone convergence theorem applied to $f_{n}$ as $n \rightarrow \infty$ then shows that

$$
\int d\left(x_{0}, x\right)^{q} \mathrm{~d} \mu^{\prime} \leq \int d\left(x_{0}, x\right)^{q} \mathrm{~d} \mu<\infty
$$

so that $\mu^{\prime} \in \mathcal{P}_{q}(X)$. Lemma 3.3 shows that $\mu^{\prime}=\mu$.
Lemma 3.6. The unique stationary measure $\mu$ satisfies $m_{x_{0}}^{q}(\mu) \leq A /(1-\bar{\lambda})^{\max (1, q)}$.
Proof. Setting $\bar{A}=A^{\min \left(1, \frac{1}{q}\right)}$ and using $L^{*} \mu=\mu$ we get:

$$
\begin{aligned}
\mathrm{W}_{q}\left(\delta_{x_{0}}, \mu\right) & \leq \mathrm{W}_{q}\left(\delta_{x_{0}}, \mathrm{~L}^{*} \delta_{x_{0}}\right)+\mathrm{W}_{q}\left(\mathrm{~L}^{*} \delta_{x_{0}}, \mathrm{~L}^{*} \mu\right) \\
& \leq \bar{A}+\bar{\lambda} \mathrm{W}_{q}\left(\delta_{x_{0}}, \mu\right) \\
(1-\bar{\lambda}) \mathrm{W}_{q}\left(\delta_{x_{0}}, \mu\right) & \leq \bar{A} .
\end{aligned}
$$

If $q \leq 1, \bar{A}=A$ and $\bar{\lambda}=\lambda$; we get $\int d\left(x_{0}, x\right)^{q} \mathrm{~d} \mu \leq A /(1-\lambda)$. If $q \geq 1 ; \bar{A}=A^{\frac{1}{q}}$ and $\bar{\lambda}=\lambda^{\frac{1}{q}}$ and we get $\left(\int d\left(x_{0}, x\right)^{q} \mathrm{~d} \mu\right)^{\frac{1}{q}} \leq A^{\frac{1}{q}} /\left(1-\lambda^{\frac{1}{q}}\right)$.

Theorem 3.1 follows at once from Lemmas 3.3, 3.5 and 3.6.
As an illustration, let us consider a simple case studied for example in [BMS06]. For each $\omega=\left(i_{1}, \ldots, i_{n}\right) \in\{0,1\}^{n}$ we set $\phi_{\omega}=\phi_{i_{1}} \circ \phi_{i_{2}} \circ \cdots \phi_{i_{n}}$ (here the order of composition, forward or backward, has no particular relevance).
Corollary 3.7. Let $a \in(0,1)$ and $b \in\left(1, \frac{1}{a}\right)$. The RIFS on the line given by

$$
\begin{equation*}
I=\{0,1\}, \quad \eta_{p}(\{0\})=\frac{1}{2} \quad \eta_{p}(\{1\})=\frac{1}{2} \quad l \quad ~\left(\quad \phi_{0}(x)=a x \quad \phi_{1}(x)=b x+1 \quad \forall x \in \mathbb{R}\right. \text { } \tag{5}
\end{equation*}
$$

has a unique stationary measure $\mu$, which has finite moments of all orders $q \in\left(0, q_{0}\right)$ where $q_{0}$ is the unique solution in $(0,+\infty)$ of $a^{q_{0}}+b^{q_{0}}=2$. More precisely

$$
m_{0}^{q}(\mu) \leq \frac{2}{2-\left(a^{q}+b^{q}\right)}
$$

Moreover, for all $q \in\left(0, \min \left(1, q_{0}\right)\right)$, all $q$-Hölder-continuous function $f: \mathbb{R}_{+} \rightarrow \mathbb{R}$ and all $x_{0}$,

$$
\left|\frac{1}{2^{n}} \sum_{\omega \in\{0,1\}^{n}} f\left(\phi_{\omega}\left(x_{0}\right)\right)-\int f \mathrm{~d} \mu\right| \leq C \lambda^{n}
$$

where $\lambda=\frac{1}{2}\left(a^{q}+b^{q}\right)<1$ and $C=\operatorname{Hol}_{q}(f) /(1-\lambda)$.
The notation $\operatorname{Hol}_{q}(f)$ denotes the best Hölder constant of $f$; note that here we ask $f$ to be globally $q$-Hölder, implying it has a growth at infinity of the order of $x^{q}$ at most. Of course, unbalanced versions of this example (i.e. with $\eta(\{0\}) \neq \eta(\{1\}))$ can be studied in the same way.

Proof. First note that the function $q \mapsto a^{q}+b^{q}$ is convex, and the assumptions ensure it is decreasing in some interval $\left(0, q_{1}\right)$ and goes to $+\infty$ when $q \rightarrow+\infty$, so that this function takes the value 2 at exactly two points, 0 and $q_{0}$.

For all $q>0$, we have (4) with $A=1$ and

$$
\int\left|\phi_{i}(x)-\phi_{i}(y)\right|^{q} \mathrm{~d} \eta(i)=\frac{1}{2}\left(a^{q}+b^{q}\right)|x-y|^{q}
$$

so that when $q<q_{0}$, (3) is satisfied with $\lambda=\frac{1}{2}\left(a^{q}+b^{q}\right)$. The claim on existence, uniqueness and moments of $\mu$ thus follows from Theorem 3.1.

The convergence of empirical averages of $f$ toward its integral with respect to $\mu$ follow from Lemma 3.2, observing

$$
\frac{1}{2^{n}} \sum_{\omega \in\{0,1\}^{n}} f\left(\phi_{\omega}\left(x_{0}\right)\right)=\mathrm{L}^{n} f\left(x_{0}\right) .
$$

Indeed as in the proof of Lemma 3.4 we have

$$
\begin{aligned}
\left|\mathrm{L}^{n} f\left(x_{0}\right)-\int f \mathrm{~d} \mu\right| & =\left|\int f \mathrm{~d}\left(\mathrm{~L}^{* n} \delta_{x_{0}}\right)-\int f \mathrm{~d} \mu\right| \\
& \leq \operatorname{Hol}_{q}(f) \mathrm{W}_{q}\left(\mathrm{~L}^{* n} \delta_{x_{0}}, \mu\right) \\
& \leq \operatorname{Hol}_{q}(f) \lambda^{n} \mathrm{~W}_{q}\left(\delta_{x_{0}}, \mu\right) .
\end{aligned}
$$

and $\mathrm{W}_{q}\left(\delta_{x_{0}}, \mu\right)=m_{x 0}^{q}(\mu) \leq \frac{1}{1-\lambda}$.

### 3.2 Two simple tools

To apply Theorem 3.1, it can be convenient to use the following statements.
Proposition 3.8. If $(\Phi, \eta)$ satisfies (3) and (4), then for all $q^{\prime} \in(0, q)$ it also satisfies them with constants $q^{\prime}, \lambda^{\prime}:=\lambda^{q^{\prime} / q}$ and $A^{\prime}:=A^{q^{\prime} / q}$.

Proof. Follows from the Jensen inequality applied to the concave function $r \mapsto r^{q^{\prime} / q}$.
Proposition 3.9. If $(\Phi, \eta)$ satisfies

$$
\int \log \operatorname{Lip}\left(\phi_{i}\right) \mathrm{d} \eta(i)<0 \quad \text { and } \quad \exists p>0, \int \operatorname{Lip}\left(\phi_{i}\right)^{p} \mathrm{~d} \eta(i)<+\infty
$$

then there exists $q>0, \lambda \in(0,1)$ such that (3) holds.
This enables us to recover a more usual 'contracting on average" hypothesis, used for example in [DF99].

Proof. The idea is simply to differentiate $\int \operatorname{Lip}\left(\phi_{i}\right)^{t} \mathrm{~d} \eta(i)$ with respect to $t$ at $t=0$; we shall use truncation to differentiate under the integral sign.

For all $n \in \mathbb{N}$, consider the functions $f_{n}: I \rightarrow \mathbb{R}$ and $F_{n}:(0, p] \rightarrow \mathbb{R}$ defined by

$$
f_{n}(i)=\max \left(\operatorname{Lip}\left(\phi_{i}\right), 1 / n\right) \quad \text { and } \quad F_{n}(t)=\int f_{n}(i)^{t} \mathrm{~d} \eta(i)
$$

Since $-\log n \leq \log f_{n}(i) \leq \frac{1}{p} \operatorname{Lip}\left(\phi_{i}\right)^{p}$ for all $i, n$, the functions $\log f_{n}$ are $\eta$-integrable. The monotone convergence theorem implies that

$$
\int \log f_{n}(i) \mathrm{d} \eta(i) \rightarrow \int \log \operatorname{Lip}\left(\phi_{i}\right) \mathrm{d} \eta(i) \in[-\infty, 0)
$$

so that for some $n \in \mathbb{N}$ we have $\int \log f_{n}(i) \mathrm{d} \eta(i) \in(-\infty, 0)$. Now $F_{n}(0)=1$ and for all $t \in\left[0, \frac{p}{2}\right]:$

$$
-\log (n) \max \left(1, f_{n}(i)^{\frac{p}{2}}\right) \leq \frac{\mathrm{d}}{\mathrm{~d} t} f_{n}(i)^{t}=\log f_{n}(i) \cdot f_{n}(i)^{t} \leq \frac{2}{p} f_{n}(i)^{t+\frac{p}{2}}
$$

is $\eta$-integrable, uniformly so over $t$ ( $n$ being fixed above). The function $F_{n}$ is thus differentiable on $\left[0, \frac{p}{2}\right]$, with $F_{n}^{\prime}(0)=\int \log f_{n}(i) \mathrm{d} \eta_{i}<0$. We conclude that there is some $q \in\left(0, \frac{p}{2}\right)$ such that $F_{n}(q) \in(0,1)$. Now

$$
\int \operatorname{Lip}\left(\phi_{i}\right)^{q} \mathrm{~d} \eta_{i} \leq \int f_{n}(i)^{q} \mathrm{~d} \eta_{i}=F_{n}(q)<1
$$

which readily implies (3).

### 3.3 Generalized moment estimates

In some cases, the stationary measure of an IFS will not be compactly supported, but will have finite moment of all order; it then makes sense to develop tools to estimate exponential, sub-exponential or super-exponential moments. In practice, the following simple result will be quite efficient.

Proposition 3.10. Let $(\Phi, \eta)$ be an RIFS on $X$ and $\varphi, \psi: X \rightarrow[0,+\infty)$ two functions that are bounded on every bounded subset of $X$ (most usually, they will factor through $\left.d\left(x_{0}, \cdot\right)\right)$. Denote by L the transfer operator, i.e. $\mathrm{L} f(x)=\int f\left(\phi_{i}(x)\right) \mathrm{d} \eta(i)$. Assume $\mu$ is a stationary measure for $(\Phi, \eta)$ and $\int \psi \mathrm{d} \mu<+\infty$.

If there exist $\theta \in(0,1)$ and $B \geq 0$ such that for all $x \in X, \mathrm{~L} \varphi(x) \leq \theta \varphi(x)+B \psi$, then $\int \varphi \mathrm{d} \mu<+\infty$ and more precisely

$$
\int \varphi \mathrm{d} \mu \leq \frac{B}{1-\theta} \int \psi \mathrm{d} \mu
$$

Proof. For all $m \in[0,+\infty)$, set $\varphi_{m}(x)=\min (\varphi(x), m)$. We have $L^{*} \mu=\mu$ since $\mu$ is stationary, and by positivity

$$
\mathrm{L} \varphi_{m} \leq \min (\mathrm{L} \varphi, m) \leq \min (\theta \varphi+B \psi, m) \leq \min (\theta \varphi, m)+B \psi=\theta \varphi_{\frac{m}{\theta}}+B \psi .
$$

Applying duality we get:

$$
\int \varphi_{m} \mathrm{~d} \mu=\int \varphi_{m} \mathrm{~d}^{*} \mu=\int \mathrm{L} \varphi_{m} \mathrm{~d} \mu \leq \int\left(\theta \varphi_{\frac{m}{\theta}}+B \psi\right) \mathrm{d} \mu
$$

from which we deduce

$$
\int\left(\varphi_{m}-\theta \varphi_{\frac{m}{\theta}}\right) \mathrm{d} \mu \leq B \int \psi \mathrm{~d} \mu .
$$

For each $x \in X$, the function $m \mapsto \varphi_{m}(x)-\theta \varphi_{\frac{m}{\theta}}(x)$ is non-decreasing and converges to $(1-\theta) \varphi(x)$. The monotone convergence theorem ensures we can pass to the limit in the above inequality, leading precisely to the claimed inequality.

Let us now consider the stationary measures $\mu_{a, p}$ defined in Section 1.1. Theorem 3.1 shows that $\mu_{a, p}$ has finite moment of all orders. Indeed (3) and (4) are satisfied for all $q \geq 1$ with $x_{0}=0, \lambda=1-p+p a^{q}$ (which is less than 1 since $a<1$ ), $A=1-p$. Let us use Proposition 3.10 to show that $\mu_{a, p}$ has some finite exponential moments.

Corollary 3.11. For all $p, a \in(0,1)$ and all $b<\log \frac{1}{1-p}$ we have

$$
\int e^{b x} \mathrm{~d} \mu_{a, p}(x) \leq\left(\frac{p}{1-(1-p) e^{b}}\right)^{\frac{1}{1-a}}
$$

Proof. Let L* be the dual transfer operator and $\varphi(x)=e^{b x}$. We have

$$
\mathrm{L}^{*} \varphi(x)=p(\varphi(x))^{a}+(1-p) e^{b} \varphi(x) .
$$

Pick any $\theta$ strictly between $(1-p) e^{b}$ and 1 : for some $C>0$ (which could be computed explicitly), $\mathrm{L}^{*} \varphi(x) \leq \theta \varphi(x)$ whenever $x \geq C$. Since $\varphi$ is increasing, we have $\mathrm{L}^{*} \varphi \leq$ $\theta \varphi+e^{b C}$ and we can thus apply Proposition 3.10 with $\psi \equiv 1$, obtaining

$$
\int e^{b x} \mathrm{~d} \mu_{a, p} \leq \frac{e^{b C}}{1-\theta} m_{0}^{1}\left(\mu_{a, p}\right)<\infty
$$

In particular $\int \varphi^{a} \mathrm{~d} \mu_{a, p}<+\infty$ and we can apply Proposition 3.10 again, to the more precise inequality $\mathrm{L}^{*} \varphi \leq(1-p) e^{b} \varphi+p \varphi^{a}$. We get

$$
\begin{aligned}
\int \varphi \mathrm{d} \mu_{a, p} & \leq \frac{p}{1-(1-p) e^{b}} \int \varphi^{a} \mathrm{~d} \mu_{a, p} \\
& \leq \frac{p}{1-(1-p) e^{b}}\left(\int \varphi \mathrm{~d} \mu_{a, p}\right)^{a} \\
\int \varphi \mathrm{~d} \mu_{a, p} & \leq\left(\frac{p}{1-(1-p) e^{b}}\right)^{\frac{1}{1-a}}
\end{aligned}
$$

The bound $\log \frac{1}{1-p}$ is optimal, as shown by Corollary A which we now prove.
Proof of Corollary A. The lower bound follows from the simple observation that for all $n \in \mathbb{N}$ we have $\mu_{a, p}([n+1, n+2]) \geq(1-p) \mu_{a, p}([n, n+1])$, thus $\mu_{a, p}([n, n+1]) \geq$ $(1-p)^{n} \mu_{a, p}([0,1])$. Since $\mu_{a, p}$ is a probability measure, there is some $n \in \mathbb{N}$ such that $\mu_{a, p}\left(\left[0, a^{-n}\right]\right)>0$, and $\mu_{a, p}\left([0,1] \geq p^{n} \mu_{a, p}\left(\left[0, a^{-n}\right]\right)>0\right.$.

The upper bound follows from Corollary 3.11 and Chebyshev's inequality: for all $t>0$ and all $b<\log \frac{1}{1-p}$, we have

$$
\mu_{a, p}([t,+\infty)) \leq e^{-b t}\left(\frac{p}{1-(1-p) e^{b}}\right)^{\frac{1}{1-a}}
$$

Given $t$, we can choose $b$ in order to optimize the above inequality. An elementary computation leads to take

$$
e^{b}=\frac{1}{1-p} \cdot \frac{(1-a) t}{1+(1-a) t},
$$

leading to a bound asymptotically equivalent to

$$
\frac{(p(1-a) t)^{\frac{1}{1-a}}}{(1-p)^{t}}
$$

Problem 3.12. Find an asymptotic for the tail of $\mu_{a, p}$, in the spirit of [Kes73] and [Gol91]. Note these works give (in a slightly different context) a precise asymptotic $\mu([t,+\infty))=f(t)+o(f(t))$ with $f$ a polynomial function; this can thus be written $\mu([t,+\infty))=f(t+o(t))$. While Corollary A already give an estimation of the form $\mu_{a, p}([t,+\infty))=g(t+o(t))$, it might be difficult to get $\mu_{a, p}([t,+\infty))=g(t)+o(g(t))$ since $g$ is exponential and the sensibility on $t$ is thus strong.

### 3.4 A heavy tail of translations

To illustrate the role of assumption (4) in Theorem 3.1 let us consider the following example on $[0,+\infty)$ (simply denoted by $(\Phi, \eta)$ throughout this subsection):

$$
\begin{array}{cc}
I=\mathbb{N}, & \eta(\{n\})=p_{n} \\
\forall n>0, \forall x \in \mathbb{R}: & \phi_{0}(x)=a x \quad \phi_{n}(x)=x+n \tag{6}
\end{array}
$$

where $p_{0}>0$ and, of course, $p_{n} \geq 0$ and $\sum_{n \geq 0} p_{n}=1$.
We have good contraction properties: assumption (3) is satisfied for any $q>0$ with $\lambda=1-\left(1-a^{q}\right) p_{0}$. As in the previous example, $\lambda$ decreases from 1 when $q \rightarrow 0$ to $1-p_{0}$ when $q \rightarrow \infty$, and thus improves with $q$. However the translation part can weight on the moments of the stationary measure. By a direct application of Theorem 3.1, we get:

Proposition 3.13. Let $q>0$; if $\sum n^{q} p_{n}<+\infty$ then ( $\Phi, \eta$ ) has a unique stationary measure $\mu$, and $\mu$ has finite $q$-th moment; if $\sum n^{q} p_{n}=+\infty$, then any stationary measure of $(\Phi, \eta)$ has infinite $q$-th moment.

Giving $\left(p_{n}\right)_{n}$ a heavy tail and taking $a \ll 1$ we get examples with very quick convergence in low-exponent Wasserstein metric but only few finite moments. This begs the question: what happens when $\sum n^{q} p_{n}=\infty$ for all $q>0$, e.g. when $p_{n} \sim 1 / n(\log n)^{2}$ as $n \rightarrow \infty$ ? Does there exist a stationary measure?

### 3.5 Statistical properties

Assume $\left(\mathrm{X}_{k}\right)_{k \in \mathbb{N}}$ is a Markov chain obtained by choosing randomly independently indices $\left(\mathrm{I}_{k}\right)_{k \geq 1}$ with law $\eta$, and setting $\mathrm{X}_{k}=\phi_{\mathrm{I}_{k}}\left(\mathrm{X}_{k-1}\right)$; we shall say that $\left(\mathrm{X}_{k}\right)_{k \in \mathbb{N}}$ is driven by $(\Phi, \eta)$. The fact that the dual transfer operator is a contraction in the Wasserstein distance $\mathrm{W}_{q}$ for some $q \in(0,1]$ is equivalent to $\left(\mathrm{X}_{k}\right)_{k \in \mathbb{N}}$ to have positive Ricci curvature in the sense of Ollivier [Oll09] in the metric space $\left(X, d^{q}\right)$; this for example implies strong concentration properties of the empirical averages

$$
\frac{1}{n} \sum_{k=1}^{n} f\left(\mathrm{x}_{k}\right)=: \hat{\mu}_{n}(f)
$$

whenever $f: X \rightarrow \mathbb{R}$ is a $q$-Hölder function, see [JO10] for effective and completely explicit results (that depends on many specific quantities that may vary between examples).

When it is possible to restrict to a compact domain of Euclidean space, one can also obtain explicit convergence speed for the empirical measure $\hat{\mu}_{n}$ itself toward the stationary measure $\mu$, in terms of $\mathrm{W}_{1}$ or in terms of other distance defined by duality [Klo18a]. Specifically, we shall consider the following metrics between measures:

$$
\left\|\nu_{0}-\nu_{1}\right\|_{\mathcal{C}_{1}^{s}}=\sup _{f \in \mathcal{C}_{1}^{\mathrm{s}}}\left|\int f \mathrm{~d} \nu_{0}-\int f \mathrm{~d} \nu_{1}\right|
$$

where $s$ is any positive integer and $\mathcal{C}_{1}^{\mathrm{s}}$ is the set of $\mathcal{C}^{s-1}$ functions $\mathbb{R}^{n} \rightarrow \mathbb{R}$ with all derivatives not greater than 1 , and with their derivatives of order ( $s-1$ ) 1-Lipschitz. In particular, $\|\cdot\|_{\mathcal{C}_{1}^{1}}$ matches exactly $\mathrm{W}_{1}$.

Theorem 3.14. Let $(\Phi, \eta)$ be an RIFS on $\mathbb{R}^{d}$ such that for some $\lambda \in(0,1)$ every map $\phi_{i}$ of $\Phi$ is $\lambda$-Lipschitz and $\left\{\left\|\phi_{i}(0)\right\|: i \in I\right\}$ is bounded. Let $\left(\mathrm{x}_{k}\right)_{k \in \mathbb{N}}$ be a Markov chain driven by $(\Phi, \eta)$, and consider the empirical measures

$$
\hat{\mu}_{n}:=\frac{1}{n} \sum_{k=1}^{n} \delta_{\mathrm{X}_{k}} .
$$

Then there exists a constant $C>0$ such that for all $n \in \mathbb{N}$,

$$
\mathbb{E}\left[\left\|\hat{\mu}_{n}-\mu\right\|_{C_{1}^{s}}\right] \leq C \begin{cases}\frac{(\log n)^{\frac{d}{2 s+1}}}{\sqrt{n}} & \text { when } 2 s>d  \tag{7}\\ \frac{\log n}{\sqrt{n}} & \text { when } 2 s=d \\ \frac{(\log n)^{d-2 s+\frac{s}{d}}}{n^{\frac{s}{d}}} & \text { when } 2 s<d\end{cases}
$$

These rates cannot be improved, except possibly for the logarithm factors (see [Klo18a] for this and other considerations, including a comparison with concentrations obtained for a fixed observable $f$ ).

Proof. Since Lemma 3.2 with $q=1$ ensures the Markov chain $\left(\mathrm{X}_{k}\right)_{k \in \mathbb{N}}$ is exponentially contracting, this is a direct application of Theorem A in [Klo18a]; the only point to check is that we can restrict to a compact subset of $\mathbb{R}^{d}$. This follows easily from the hypotheses: let $D=\sup \left\{\left\|\phi_{i}(0)\right\|: i \in I\right\}$; then the ball $B$ of center 0 and radius $R=D /(1-\lambda)$ satisfies $\phi_{i}(B) \subset B$ for all $i$. Indeed whenever $\|x\| \leq R$, we have

$$
\left\|\phi_{i}(x)\right\| \leq\left\|\phi_{i}(x)-\phi_{i}(0)\right\|+\left\|\phi_{i}(0)\right\| \leq \lambda\|x\|+D \leq R .
$$

Using relaxed assumptions such as (3) and (4) would be nice, but the rates would necessarily be altered given the stationary measure would not usually have finite moments of all orders (without some moment condition, one cannot expect even an optimal approximation by a discrete measure supported on $n$ point to achieve the rate $1 / n^{\frac{1}{d}}$ in $\mathrm{W}_{1}$ when $d>2$; see [FG15] and [DM19] for rates of convergence of empirical measures under various moment assumptions).

### 3.6 Stationary measures beyond products

While the case of RIFS as defined above, where the randomness is materialized by a sequence of independent random variables of law $\eta \in \mathcal{P}(I)$, is most commonly studied, there has been some interest to generalize this setting. A first generalization is to replace the i.i.d. sequence by a stationary Markov chain; a further generalization is to draw the infinite word $\omega=\omega_{0} \ldots \omega_{k} \ldots$ randomly with law an arbitrary shift-invariant measure $\nu \in \mathcal{P}\left(I^{\mathbb{N}}\right)$ - the case of RIFS corresponding to the independent Bernoulli product $\nu=\eta^{\otimes \mathbb{N}}$; then one can consider a yet further generalization where the shift replaced by an arbitrary measure-preserving dynamical system.

### 3.6.1 Skew-products

We still consider $(X, d)$ a complete metric space, and we additionally fix a standard measured space $(Y, \mathcal{A})$ (i.e. it is isomorphic to $[0,1]$ with its Borel $\sigma$-algebra) equipped with a probability measure $\nu$, and a $\nu$-preserving map $S: Y \rightarrow Y$. A skew-product map over $S$ with fiber $X$ is a map

$$
\begin{aligned}
\Psi: X \times Y & \rightarrow X \times Y \\
(x, y) & \mapsto\left(\psi_{y}(x), S(y)\right)
\end{aligned}
$$

where $(x, y) \mapsto \psi_{y}(x)$ is a measurable map. While, as we have seen above, an RIFS can be studied dynamically by looking at a random orbit $\mathrm{X}_{0}, \mathrm{X}_{n+1}=\phi_{\mathrm{I}_{n}}\left(\mathrm{X}_{n}\right)$ where $\left(\mathrm{I}_{n}\right)_{n \geq 1}$ are i.i.d. random variables of law $\eta$, in the present setting the corresponding random sequence of points is given by $\mathrm{x}_{n+1}=\psi_{S^{n}(\mathrm{Y})}\left(\mathrm{x}_{n}\right)$ where Y is a random element of $Y$ with law $\nu$, taking the place of the whole sequence ( $\mathrm{I}_{1}, \mathrm{I}_{2}, \ldots$ ). In other words, RIFS correspond to the particular case when $Y=I^{\mathbb{N}}, \nu=\eta^{\otimes \mathbb{N}}, S$ is the shift $y=\left(y_{0}, y_{1}, \ldots\right) \mapsto$ $S(y)=\left(y_{1}, y_{2}, \ldots\right)$ and $\psi_{y}(x)=\phi_{y_{0}}(x)$. Note that $\Psi$ carries the information of what are $X, Y$ and $S$; when we refer to this setting, we shall therefore call $(\Psi, \nu)$ a skew product.

We denote by $\pi^{X}$, $\pi^{Y}$ the projection maps from $X \times Y$ to each factor; a measure $\mu \in \mathcal{P}(X)$ is said to be a stationary measure of the skew product $(\Psi, \nu)$ when there exist a measure $\hat{\nu} \in \mathcal{P}(X \times Y)$ such that:

$$
\hat{\nu} \text { is } \Psi \text {-invariant, } \quad \pi_{*}^{Y} \hat{\nu}=\nu, \quad \text { and } \quad \pi_{*}^{X} \hat{\nu}=\mu .
$$

In the case of a RIFS, this coincides with the previous definition of stationary measure. The measure $\hat{\nu}$ as above shall be called a lift of $\nu$. The basic question we want to address under specific assumptions is whether there exist a unique stationary measure; a positive answer will follow from the uniqueness of the lift of $\nu$.

Definition 3.15. We say that $\Psi$ contracts the fibers whenever there exist $\lambda \in(0,1)$ such that for all $y \in Y$, the map $\psi_{y}$ is $\lambda$-Lipschitz.

We say that $\Psi$ has bounded displacement if for some $x_{0} \in X$, there exist an $A>0$ such that the set $d\left(x_{0}, \psi_{y}\left(x_{0}\right)\right) \leq A$ for all $y \in Y$.

Observe that when $\left(\psi_{y}\right)_{y \in Y}$ is an equicontinuous family, e.g. when $\Psi$ contracts the fibers, in the definition of bounded displacement "for some $x_{0}$ " could be equivalently replaced by "for all $x_{0}$ " (up to changing the value of $A$ ).

The main result of this section is the following.
Theorem 3.16. Let $\Psi$ be a skew-product map on $X \times Y$ that contracts the fibers and has bounded displacement. Each $S$-invariant $\nu \in \mathcal{P}(Y)$ has a unique lift, and in particular the skew product $(\Psi, \nu)$ has a unique stationary measure $\mu$, which moreover has bounded support.

Let $\left(\mathrm{X}_{k}\right)_{k \in \mathbb{N}}$ be a stochastic process associated to $(\Psi, \nu)$ as above, with $\mathrm{X}_{0}$ independent from Y and of arbitrary law $\tilde{\mu}_{0} \in \mathcal{P}_{q}(X)$ for some $q>0$, and let $\tilde{\mu}_{k} \in \mathcal{P}(X)$ be the law of $\mathrm{X}_{k}$. Then for all $k \in \mathbb{N}$,

$$
\mathrm{W}_{q}\left(\tilde{\mu}_{k}, \mu\right) \leq D \tilde{\lambda}^{k}
$$

where

$$
\tilde{\lambda}=\lambda^{\min (q, 1)} \in(0,1), \quad D=m_{x_{0}}^{q}\left(\tilde{\mu}_{0}\right)^{\min \left(1, \frac{1}{q}\right)}+\left(\frac{A}{1-\lambda}\right)^{\min (1, q)},
$$

and $A, \lambda$ are the constants in the bounded displacement and fiber contraction hypotheses.

### 3.6.2 Fiber-wise Wasserstein distance

The main tool to prove Theorem 3.16 is a variation of Wasserstein distance that is adapted to a projection map and the inverse images of a given measure on its target space. This notion was at the heart of [Klo18b], from which we adapt the relevant definitions and properties. Theorem A from [Klo18b] is not immediately applicable here since $X$ need not be compact, $\operatorname{diam}\left(\Psi^{n}(X \times\{y\})\right)$ might be infinite for all $n$, and $Y$ is not even a topological space; but the adaptation is relatively straightforward.

Fix any $\nu \in \mathcal{P}(Y)$ and let $\mathcal{P}^{\nu}:=\left(\pi_{*}^{Y}\right)^{-1}(\nu) \subset \mathcal{P}(X \times Y)$ be the fiber of $\nu$, i.e, the set of measures on $X \times Y$ with second marginal equal to $\nu$. Reminding that we fixed a point $x_{0} \in X$, given any $\rho \in \mathcal{P}^{\nu}$ and $q>0$ we define its $q$-th moment by

$$
m_{x_{0}}^{q}(\rho)=\int d\left(x, x_{0}\right)^{q} \mathrm{~d} \rho(x, y)
$$

where the integral is over the whole product $X \times Y$ but distances are recorded only "along the fibers", i.e. over the $X$ factor. We let $\mathcal{P}_{q}^{\nu}$ be the subset of $\mathcal{P}^{\nu}$ consisting of measure of finite $q$-th moment (this set does not depend on $x_{0}$ ).

The product $X^{2} \times Y$ identifies with what was noted $\Delta_{\pi}$ in [Klo18b] (pairs of point in the total space that project to the same point on the base $Y$ ); we consider the maps

$$
\pi_{02}:\left(x, x^{\prime}, y\right) \mapsto(x, y) \quad \pi_{12}:\left(x, x^{\prime}, y\right) \mapsto\left(x^{\prime}, y\right) \quad \pi_{2}:\left(x, x^{\prime}, y\right) \mapsto y
$$

For all $\rho_{0}, \rho_{1} \in \mathcal{P}^{\nu}$ let $\Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right):=\left\{\gamma \in \mathcal{P}\left(X^{2} \times Y\right) \mid\left(\pi_{02 *}\right) \gamma=\rho_{0}\right.$ and $\left.\left(\pi_{12 *}\right) \gamma=\rho_{1}\right\}$ (playing the role of $\Gamma_{\pi}$ in [Klo18b], where we had chosen to emphasize the projection map rather than the image measure) and define

$$
\begin{aligned}
\mathrm{C}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right) & =\inf _{\gamma \in \Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)} \int d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \gamma\left(x, x^{\prime}, y\right) \\
\mathrm{W}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right) & =\left(\mathrm{C}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right)\right)^{\min \left(1, \frac{1}{q}\right)} .
\end{aligned}
$$

The following basic result is proven in the same way as in [Klo18b].
Proposition 3.17. For all $\rho_{0}, \rho_{1} \in \mathcal{P}^{\nu}$, the set $\Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)$ is non-empty. If the moments $m_{x_{0}}^{q}\left(\rho_{i}\right)$ are finite for $i \in\{0,1\}$, then $\mathrm{W}^{\nu}\left(\rho_{0}, \rho_{1}\right)<\infty$. Moreover, if $\left(\xi_{y}\right)_{y \in Y}$ and $\left(\zeta_{y}\right)_{y \in Y}$ are the disintegrations of $\rho_{0}$ and $\rho_{1}$ with respect to $\pi^{Y}$, then

$$
\mathrm{W}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right)= \begin{cases}\left(\int \mathrm{W}_{q}\left(\xi_{y}, \zeta_{y}\right)^{q} \mathrm{~d} \nu(y)\right)^{\frac{1}{q}} & \text { when } q \geq 1  \tag{8}\\ \int \mathrm{~W}_{q}\left(\xi_{y}, \zeta_{y}\right) \mathrm{d} \nu(y) & \text { when } q \leq 1 .\end{cases}
$$

Last, $\mathrm{W}_{q}^{\nu}$ is a complete metric on the set $\mathcal{P}_{q}^{\nu}$.

Proof. Let $\left(\xi_{y}\right)_{y \in Y}$ and $\left(\zeta_{y}\right)_{y \in Y}$ be the disintegrations of $\rho_{0}$ and $\rho_{1}$ with respect to $\pi^{Y}$ (identifying $X$ with the fibers of $\pi^{Y},\left(\xi_{y}\right)_{y \in Y}$ is thus a family of measures on $X$ characterized by $\int f(x, y) \mathrm{d} \xi_{y}(x) \mathrm{d} \nu(y)=\int f(x, y) \mathrm{d} \rho_{0}(x, y)$ for all continuous bounded function $f: X \times Y \rightarrow \mathbb{R}$.)

From any measurable choice of $y \mapsto \gamma_{y} \in \Gamma\left(\xi_{y}, \zeta_{y}\right)$ (e.g. $\left.\gamma_{y}=\xi_{y} \otimes \zeta_{y}\right)$ we can build an element $\gamma$ of $\Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)$ by setting $\int f\left(x, x^{\prime}, y\right) \mathrm{d} \gamma\left(x, x^{\prime}, y\right)=\iint f\left(x, x^{\prime}, y\right) \mathrm{d} \gamma_{y}\left(x, x^{\prime}\right) \mathrm{d} \nu(y)$. In particular $\Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)$ is non-empty.

Conversely, given any $\gamma \in \Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)$ its disintegration with respect to $\pi_{2}$ is a family $\left(\gamma_{y}\right)_{y \in Y}$ of measures on $X \times X$, and by testing $\gamma$ against integrands of the form $f(x) g(y)$ and $f\left(x^{\prime}\right) g(y)$ one sees that $\gamma_{y} \in \Gamma\left(\xi_{y}, \zeta_{y}\right)$ for $\nu$-almost all $y$.

Since $\int d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \gamma\left(x, x^{\prime}, y\right)=\iint d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \gamma_{y}\left(x, x^{\prime}\right) \mathrm{d} \nu(y) \geq \int \mathrm{C}_{q}\left(\xi_{y}, \zeta_{y}\right) \mathrm{d} \check{\mu}(y)$, taking an infimum we get $\mathrm{C}_{q}^{\nu}\left(\mu_{0}, \mu_{1}\right) \geq \int \mathrm{C}_{q}\left(\xi_{y}, \zeta_{y}\right) \mathrm{d} \check{\mu}$.

For each $y$, the set of optimal transport plans from $\xi_{y}$ to $\zeta_{y}$ is compact (see e.g. the proof of Theorem 4.1 in [Vil09]), thus by the measurable selection theorem there is a measurable family $\left(\gamma_{y}\right)_{y \in Y}$ such that for $\nu$-almost all $y \in Y, \int d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \gamma_{y}\left(x, x^{\prime}\right)=$ $\mathrm{C}_{q}\left(\xi_{y}, \zeta_{y}\right)$. It follows $\mathrm{C}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right) \leq \int \mathrm{C}_{q}\left(\xi_{y}, \zeta_{y}\right) \mathrm{d} \nu$ and (8) is proven.

To complete the proof, it remains to be seen that $\mathrm{C}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right)<\infty$ and that $\mathrm{W}_{q}^{\nu}$ is a metric making $\mathcal{P}_{q}^{\nu}$ a complete space. The triangular inequality follows from (8), and then finiteness is obtained by observing

$$
\mathrm{W}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right) \leq \mathrm{W}_{q}^{\nu}\left(\rho_{0}, \delta_{x_{0}} \otimes \nu\right)+\mathrm{W}_{q}^{\nu}\left(\delta_{x_{0}} \otimes \nu, \rho_{1}\right)=m_{x_{0}}^{q}\left(\rho_{0}\right)+m_{x_{0}}^{q}\left(\rho_{1}\right)
$$

Finally, The Riesz-Fischer Theorem for metric-space valued functions ensures that $\mathrm{W}_{q}^{\nu}$ is a complete metric on $\mathcal{P}_{q}^{\nu}$, seen via disintegration as a closed subset of the space of maps $Y \rightarrow \mathcal{P}_{q}(X)$.

### 3.6.3 Proof of Theorem 3.16

Let $\nu$ be any $S$-invariant probability measure on $Y$. We first observe that the fiber contraction and bounded displacement properties ensure that $\Psi_{*}$ preserves $\mathcal{P}_{q}^{\nu}$ for all $q$. These uniform assumptions also ensure that for some bounded set $B \subset X$, the set $B \times Y$ is an absorbing invariant set, i.e. $\Psi(B \times Y) \subset B \times Y$ and for all $(x, y) \in X \times Y$ there is some $k \in \mathbb{N}$ such that $\Psi^{k}(x, y) \in B \times Y$. Let indeed $A>0$ be such that for all $y, d\left(x_{0}, \psi_{y}\left(x_{0}\right)\right) \leq A$, fix any $\varepsilon>0$, set $R=(1+\varepsilon) A /(1-\lambda)$ and let $B=B\left(x_{0}, R\right)$ be the ball of center $x_{0}$ and radius $R$ in $X$; then for all $x, y \in X \times Y$ we have

$$
\begin{aligned}
d\left(x_{0}, \psi_{y}(x)\right) & \leq d\left(x_{0}, \psi_{y}\left(x_{0}\right)\right)+d\left(\psi_{y}\left(x_{0}\right), \psi_{y}(x)\right) \\
& \leq A+\lambda d\left(x_{0}, x\right)
\end{aligned}
$$

When $x \in B$, the right-hand side is at most $A+\lambda R=\frac{1+\varepsilon \lambda}{1-\lambda} A<R$, proving the $B \times Y$ is $\Psi$-invariant. When $x \notin B$, we have $A<\frac{1-\lambda}{1+\varepsilon} d\left(x_{0}, x\right)$ and the right-hand side is at most

$$
\left(\frac{1-\lambda}{1+\varepsilon}+\lambda\right) d\left(x_{0}, x\right)=\frac{1+\varepsilon \lambda}{1+\varepsilon} d\left(x_{0}, x\right)
$$

where $\frac{1+\varepsilon \lambda}{1+\varepsilon}<1$, proving the absorbing property with $k \simeq \log d\left(x_{0}, x\right)$.

Let $\rho_{0}, \rho_{1} \in \mathcal{P}_{q}^{\nu}$. We consider the map $X^{2} \times Y \rightarrow X^{2} \times Y$ defined by

$$
\Psi_{2}\left(x, x^{\prime}, y\right)=\left(\psi_{y}(x), \psi_{y}\left(x^{\prime}\right), S(y)\right)
$$

For $i \in\{0,1\}$ we have $\pi_{i 2} \circ \Psi_{2}=\Psi \circ \pi_{i 2}$; as a consequence, for any $\gamma \in \Gamma^{\nu}\left(\rho_{0}, \rho_{1}\right)$ we have $\Psi_{2 *} \gamma \in \Gamma^{\nu}\left(\Psi_{*} \rho_{0}, \Psi_{*} \rho_{1}\right)$. Observing

$$
\begin{aligned}
\int d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \Psi_{2 *} \gamma\left(x, x^{\prime}, y\right) & =\int d\left(\psi_{y}(x), \psi_{y}\left(x^{\prime}\right)\right)^{q} \mathrm{~d} \gamma\left(x, x^{\prime}, y\right) \\
& \leq \lambda^{q} \int d\left(x, x^{\prime}\right)^{q} \mathrm{~d} \gamma\left(x, x^{\prime}, y\right)
\end{aligned}
$$

and taking an infimum, we see that

$$
\mathrm{W}_{q}^{\nu}\left(\Psi_{*} \rho_{0}, \Psi_{*} \rho_{1}\right) \leq \lambda^{\min (1, q)} \mathrm{W}_{q}^{\nu}\left(\rho_{0}, \rho_{1}\right),
$$

in particular $\Psi_{*}$ induces a contraction on the complete metric space $\left(\mathcal{P}_{q}^{\nu}, \mathrm{W}_{q}^{\nu}\right)$. Therefore, there exists a unique $\Psi$-invariant lift $\hat{\nu}$ of $\nu$ having finite $q$-th moment. By considering different $q$, we already see that the measure $\hat{\nu}$ has finite moments of all orders but, since $B \times Y$ is absorbing, any $\Psi$-invariant measure is concentrated on $B \times Y$. This proves that $\hat{\nu}$ is the unique $\Psi$-invariant lift of $\nu$ on the whole of $\mathcal{P}^{\nu}$, and that its first marginal $\mu$ is supported on a bounded set. Explicitly, by letting $\varepsilon$ above go to 0 , we have that $\mu$ is concentrated on $B\left(x_{0}, A /(1-\lambda)\right)$.

Consider now the stochastic process $\left(\mathrm{x}_{k}\right)_{k \in \mathbb{N}}$. Let $\rho_{0}:=\tilde{\mu}_{0} \otimes \nu$ be the law of ( $\mathrm{x}_{0}, \mathrm{Y}$ ); then the law of $\left(\mathrm{X}_{k}, S^{k}(\mathrm{Y})\right)=\Psi^{k}\left(\mathrm{X}_{0}, \mathrm{Y}\right)$ is $\rho_{k}=\Psi_{*}^{k}\left(\rho_{0}\right)$, by definition has first marginal $\tilde{\mu}_{k}$, and by invariance has second marginal $\nu$. Since $\Psi_{*}$ is a contraction in $\mathcal{P}_{q}^{\nu} \ni \rho_{0}$, we obtain that

$$
\begin{equation*}
\mathrm{W}_{q}^{\nu}\left(\rho_{k}, \hat{\nu}\right) \leq \lambda^{k \min (1, q)} \mathrm{W}_{q}^{\nu}\left(\rho_{0}, \hat{\nu}\right) \tag{9}
\end{equation*}
$$

On the first hand, using the transport plan obtained by projecting an optimal $\gamma \in$ $\Gamma^{\nu}\left(\rho_{k}, \hat{\nu}\right)$ on the first two variables, we get $\mathrm{W}_{q}\left(\tilde{\mu}_{k}, \mu\right) \leq \mathrm{W}_{q}^{\nu}\left(\rho_{k}, \hat{\nu}\right)$. On the other hand, we have

$$
\begin{aligned}
\mathrm{W}_{q}^{\nu}\left(\rho_{0}, \hat{\nu}\right) & \leq \mathrm{W}_{q}^{\nu}\left(\tilde{\mu}_{0} \otimes \nu, \delta_{x_{0}} \otimes \nu\right)+\mathrm{W}_{q}^{\nu}\left(\delta_{x_{0}} \otimes \nu, \hat{\nu}\right) \\
& \leq m_{x_{0}}^{q}\left(\tilde{\mu}_{0}\right)^{\min \left(1, \frac{1}{q}\right)}+(A /(1-\lambda))^{\min (1, q)}
\end{aligned}
$$

since $\hat{\nu}$ is concentrated on $B\left(x_{0}, A /(1-\lambda)\right) \times Y$. Together with (9), this concludes the proof of Theorem 3.16.

## 4 Dependence of the stationary measure on the RIFS and linear response

In this section we seek to quantify how close the stationary measures of two slightly different RIFS must be. To this end, we need to introduce a way to quantify the distance between RIFS; it is both natural and effective to use an adaptation of Wasserstein
distances. There are two points to consider in this adaptation: first, which metric to use for maps; second, how to take into account that we need not consider IFS with the same index set. The second point is easily dealt with, by considering couplings $\gamma \in \mathcal{P}\left(I_{0} \times I_{1}\right)$ of measures $\eta_{0} \in \mathcal{P}\left(I_{0}\right)$ and $\eta_{1} \in \mathcal{P}\left(I_{1}\right)$. There is much flexibility to address the first point; taking the uniform distance $d_{\infty}(\phi, \psi)=\sup _{x \in X} d(\phi(x), \psi(x))$ is ill suited to the non-compact case, as for example the map $x \mapsto a x+b$ acting on $\mathbb{R}$ would not depend continuously on the parameters $a, b$ : changing $a$ the slightest bit would yield a map infinitely far from the original one. We therefore consider a pointed Lipschitz distance, notably suitable for Lipschitz IFS:

$$
\begin{aligned}
d_{x_{0}}(\phi, \psi) & :=\min \left\{\varepsilon \geq 0 \mid \forall x \in X: d(\phi(x), \psi(x)) \leq \varepsilon+\varepsilon d\left(x, x_{0}\right)\right\} \\
& =\sup _{x \in X} \frac{d(\phi(x), \psi(x))}{1+d\left(x, x_{0}\right)}
\end{aligned}
$$

This defines a metric on the space of Lipschitz maps $X \rightarrow X$, and we construct from it the Wasserstein-like distance $\mathrm{W}_{x_{0}, q}$ (possibly taking the value $\infty$ ) between RIFS:

$$
\begin{aligned}
& \mathrm{C}_{x_{0}, q} \\
&\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right):=\inf _{\gamma \in \Gamma\left(\eta_{0}, \eta_{1}\right)} \int d_{x_{0}}\left(\phi_{i}, \psi_{j}\right)^{q} \mathrm{~d} \gamma(i, j) \\
& \mathrm{W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right):=\mathrm{C}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right)^{\min \left(1, \frac{1}{q}\right)}
\end{aligned}
$$

### 4.1 Lipschitz regularity of the stationary measure

Theorem 4.1. Consider two $\operatorname{RIFS}\left(\Phi_{0}, \eta_{0}\right)$ and $\left(\Phi_{1}, \eta_{1}\right)$ such that the first one satisfies (3) and (4) for some $q, \lambda_{0}, A_{0}$ (and thus has a unique stationary measure $\mu_{0}$ ), and such that the second one has at least one stationary measure $\mu_{1}$ with finite $q$-th moment $m_{x_{0}}^{q}\left(\mu_{1}\right)$. Then we have

$$
\begin{aligned}
& \mathrm{W}_{q}\left(\mu_{0}, \mu_{1}\right) \leq C \mathrm{~W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right) \\
& \text { where } C= \begin{cases}\frac{1+m_{x_{0}}^{q}\left(\mu_{1}\right)}{1-\lambda_{0}} & \text { when } q \leq 1 \\
2^{1-\frac{1}{q}} \frac{\left(1+m_{x_{0}}^{q}\left(\mu_{1}\right)\right)^{\frac{1}{q}}}{1-\lambda_{0}^{\frac{1}{q}}} & \text { when } q>1\end{cases}
\end{aligned}
$$

In particular, if we fix $\Phi=\Phi_{0}=\Phi_{1}$ and restrict to measures $\eta$ satisfying (3) and (4) the map $\eta \mapsto \mu$ (which is well-defined by Theorem 3.1) is locally Lipschitz.

Remark 4.2. When the second RIFS satisfy (3) and (4) with constants $q, \lambda_{1}, A_{1}$, we can choose to apply the result after exchanging them to optimize; however using only the moment estimate of Theorem 3.1 this is expected to provide small improvements, since both spectral gaps (i.e. $1-\lambda_{0}$ and $1-\lambda_{1}$ ) are involved in denominators.

Remark 4.3. When the second RISF has several stationary measures with finite $q$-th moment, Theorem 4.1 shows that they all lie within small distance of $\mu_{0}$.

Let us now turn to the proof of Theorem 4.1. We denote by $\mathrm{L}_{k}$ the transfer operator of the $\operatorname{RIFS}\left(\Phi_{k}, \eta_{k}\right)(k \in\{0,1\})$.
Lemma 4.4. For all RIFS $\left(\Phi_{0}, \eta_{0}\right)$, $\left(\Phi_{1}, \eta_{1}\right)$ and all $\nu \in \mathcal{P}_{q}(X)$ we have

$$
\begin{aligned}
& \mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \nu, \mathrm{~L}_{1}^{*} \nu\right) \leq D \mathrm{~W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right) \\
& \text { where } D= \begin{cases}1+m_{x_{0}}^{q}(\nu) & \text { when } q \leq 1 \\
2^{1-\frac{1}{q}}\left(1+m_{x_{0}}^{q}(\nu)\right)^{\frac{1}{q}} & \text { when } q>1\end{cases}
\end{aligned}
$$

Proof. If $\mathrm{C}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right)=\infty$, the statement is emptily true. Assume otherwise, and let $\gamma \in \Gamma\left(\eta_{0}, \eta_{1}\right)$ be an optimal coupling. Let $\bar{\nu}=(\mathrm{Id}, \mathrm{Id})_{*} \nu \in \mathcal{P}(X \times X)$ be the trivial coupling of $\nu$ with itself. Then $\bar{\gamma}:=\int\left(\phi_{i}, \psi_{j}\right)_{*} \bar{\nu} \mathrm{~d} \gamma(i, j)$ is a coupling of $\mathrm{L}_{0}^{*} \nu$ and $\mathrm{L}_{1}^{*} \nu$, so that

$$
\begin{aligned}
\mathrm{C}_{q}\left(\mathrm{~L}_{0}^{*} \nu, \mathrm{~L}_{1}^{*} \nu\right) & \leq \int d(x, y)^{q} \mathrm{~d} \bar{\gamma}(x, y) \\
& =\iint d\left(\phi_{i}(x), \psi_{j}(y)\right)^{q} \mathrm{~d} \bar{\nu}(x, y) \mathrm{d} \gamma(i, j) \\
& =\iint d\left(\phi_{i}(x), \psi_{j}(x)\right)^{q} \mathrm{~d} \nu(x) \mathrm{d} \gamma(i, j) \\
& \leq \iint\left(d_{x_{0}}\left(\phi_{i}, \psi_{j}\right)\left(1+d\left(x, x_{0}\right)\right)\right)^{q} \mathrm{~d} \nu(x) \mathrm{d} \gamma(i, j) \\
& \leq \int d_{x_{0}}\left(\phi_{i}, \psi_{j}\right)^{q} \mathrm{~d} \gamma(i, j) \int\left(1+d\left(x, x_{0}\right)\right)^{q} \mathrm{~d} \nu(x) .
\end{aligned}
$$

When $q \leq 1$, using $(1+r)^{q} \leq 1+r^{q}$ we obtain

$$
\begin{aligned}
\mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \nu, \mathrm{~L}_{1}^{*} \nu\right) & \leq \mathrm{W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right) \int\left(1+d\left(x, x_{0}\right)^{q}\right) \mathrm{d} \nu(x) \\
& \leq \mathrm{W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right)\left(1+m_{x_{0}}^{q}(\nu)\right)
\end{aligned}
$$

while when $q \geq 1$, using $(1+r)^{q} \leq 2^{q-1}\left(1+r^{q}\right)$ we get

$$
\begin{aligned}
\mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \nu, \mathrm{~L}_{1}^{*} \nu\right) & \leq \mathrm{W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right)\left(2^{q-1} \int\left(1+d\left(x, x_{0}\right)^{q}\right) \mathrm{d} \nu(x)\right)^{\frac{1}{q}} \\
& \leq \mathrm{W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right) \cdot 2^{1-\frac{1}{q}}\left(1+m_{x_{0}}^{q}(\nu)\right)^{\frac{1}{q}}
\end{aligned}
$$

Proof of Theorem 4.1. Apply Lemma 4.4 to $\nu=\mu_{1}$ and use that $\mathrm{L}_{0}^{*}$ is a contraction (Lemma 3.2, recall $\left.\bar{\lambda}_{0}=\lambda_{0}^{\min \left(1, \frac{1}{q}\right)}\right)$ :

$$
\begin{aligned}
\mathrm{W}_{q}\left(\mu_{0}, \mu_{1}\right) & \leq \mathrm{W}_{q}\left(\mu_{0}, \mathrm{~L}_{0}^{*} \mu_{1}\right)+\mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \mu_{1}, \mu_{1}\right) \\
& =\mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \mu_{0}, \mathrm{~L}_{0}^{*} \mu_{1}\right)+\mathrm{W}_{q}\left(\mathrm{~L}_{0}^{*} \mu_{1}, \mathrm{~L}_{1}^{*} \mu_{1}\right) \\
& \leq \bar{\lambda}_{0} \mathrm{~W}_{q}\left(\mu_{0}, \mu_{1}\right)+D \mathrm{~W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right) \\
\mathrm{W}_{q}\left(\mu_{0}, \mu_{1}\right) & \leq \frac{D}{1-\bar{\lambda}_{0}} \mathrm{~W}_{x_{0}, q}\left(\left(\Phi_{0}, \eta_{0}\right),\left(\Phi_{1}, \eta_{1}\right)\right)
\end{aligned}
$$

### 4.2 Linear response

The Rademacher theorem ensures that Lipschitz functions $[a, b] \rightarrow \mathbb{R}$ are differentiable Lebesgue almost-everywhere; a similar result has been proven by Ambrosio, Gigli and Savaré [AGS08] for maps $[a, b] \rightarrow \mathcal{P}_{q}\left(\mathbb{R}^{n}\right)$ for $q>1$, where one has of course to make precise what "differentiable" means. Together with Theorem 4.1, this provides a "linear response formula" in many cases.

Corollary 4.5 (Linear Response). Let $\left(\Phi_{t}, \eta_{t}\right)_{t \in[a, b]}$ be a curve of RIFS on $\mathbb{R}^{n}$ (endowed with the Euclidean metric, the origin $O$ serving as reference point), assume that for some $q>1$,
i. $\left(\Phi_{t}, \eta_{t}\right)_{t}$ is Lipschitz in $\mathrm{W}_{O, q}$, i.e. there exist some $C>0$ such that for all $t, t^{\prime} \in$ $[a, b]$,

$$
\mathrm{W}_{O, q}\left(\left(\Phi_{t}, \eta_{t}\right),\left(\Phi_{t^{\prime}}, \eta_{t^{\prime}}\right)\right) \leq C\left|t-t^{\prime}\right|,
$$

ii. there exist $\lambda_{+} \in(0,1)$ and $A_{+}>0$ such that for all $t \in[a, b]$ the $\operatorname{RIFS}\left(\Phi_{t}, \eta_{t}\right)$ satisfies hypotheses (3) and (4) with parameters $q, \lambda_{t} \leq \lambda_{+}$and $A_{t} \leq A_{+}$
and let $\mu_{t}$ denote the unique stationary measure of $\left(\Phi_{t}, \eta_{t}\right)$. Then there exist a family $\left(v_{t}\right)_{t \in[a, b]}$ of measurable vector fields on $\mathbb{R}^{n}$ such that for Lebesgue almost all $t$ :
i. $\left\|v_{t}\right\| \in L^{q}\left(\mu_{t}\right)$ and $\left|v_{t}\right|^{q-2} v_{t}$ can be approximated by gradients of smooth functions $\mathbb{R}^{n} \rightarrow \mathbb{R}$ in the $L^{q^{\prime}}\left(\mu_{t}\right)$ norm where $q^{\prime}=q /(q-1)$,
ii. $\frac{\mathrm{d}}{\mathrm{d} t} \mu_{t}+\nabla \cdot\left(v_{t} \mu_{t}\right)=0$ weakly, i.e. for almost all $t_{0} \in[a, b]$ and all $C^{1}$ compactly supported functions $f: \mathbb{R}^{n} \rightarrow \mathbb{R}$ :

$$
\left.\frac{\mathrm{d}}{\mathrm{~d} t} \int f \mathrm{~d} \mu_{t}\right|_{t=t_{0}}=\int \nabla f \cdot v_{t_{0}} \mathrm{~d} \mu_{t_{0}} .
$$

Note that here $q>1$ is needed to ensure strict convexity in the optimal transport problem.

Proof. By Theorem 4.1, the family of stationary measures $\left(\mu_{t}\right)_{t \in[a, b]}$ is Lipschitz, in particular absolutely continuous (with bounded derivative). Thus Theorem 8.3.1 of [AGS08] applies, giving precisely the claim (note the formulation (8.1.4) for the interpretation of the continuity equation).

### 4.3 The case of Bernoulli convolutions

Corollary B will follow from Corollary 4.5. We take as reference point $x_{0}=O=0 \in \mathbb{R}$ and recall that the family ( $\Phi^{\lambda}, \eta$ ) of RIFS defining Bernoulli convolutions is given in Section 1.2.

Proof of Corollary B. Fix any $q>1$ and let us check the hypotheses of Corollary 4.5
First, given $\lambda, \lambda^{\prime} \in(0,1)$,

$$
d_{x_{0}}\left(\phi_{0}^{\lambda}, \phi_{0}^{\lambda^{\prime}}\right)=d_{x_{0}}\left(\phi_{1}^{\lambda}, \phi_{1}^{\lambda^{\prime}}\right)=\left|\lambda-\lambda^{\prime}\right|
$$

so that the identity coupling of $\eta$ with itself, defined by $\gamma(\{(0,0)\})=\gamma(\{(1,1)\})=\frac{1}{2}$ and $\gamma(\{(0,1)\})=\gamma(\{(1,0)\})=0$, implies $\mathrm{C}_{x_{0}, q}\left(\left(\Phi^{\lambda}, \eta\right),\left(\Phi^{\lambda^{\prime}}, \eta\right)\right) \leq\left|\lambda-\lambda^{\prime}\right|^{q}$. Second, in restriction to any interval $\left[\frac{1}{2}, 1-\varepsilon\right]$ with $\varepsilon>0$ we have (3) and (4) with uniform bounds $\lambda_{+}=1-\varepsilon$ and $A_{+}=1$.

Corollary 4.5 provides us with a family $\left(v_{\lambda}\right)_{\lambda \in(0,1)}$ of vector fields on $\mathbb{R}$, which can be identified with functions $\mathbb{R} \rightarrow \mathbb{R}$, such that $\left\|v_{\lambda}\right\| \in L^{q}\left(\mu_{\lambda}\right)$ and $\frac{\mathrm{d}}{\mathrm{d} \lambda} \mu_{\lambda}+\left(v_{\lambda} \mu_{\lambda}\right)^{\prime}=0$ weakly for almost all $\lambda$. When $\lambda>\frac{1}{2}$, up to further restricting to a subset of full Lebesgue measure for the parameter $\lambda$, by Solomyak's Theorem [Sol95] we have $\mu_{\lambda}$ absolutely continuous with density denoted by $g_{\lambda}$. For a $\mathcal{C}^{1}$ compactly supported test function $f: \mathbb{R} \rightarrow \mathbb{R}$ we get

$$
\left.\frac{\mathrm{d}}{\mathrm{~d} t} \int f \mathrm{~d} \mu_{t}\right|_{t=\lambda}=\int_{0}^{1} f^{\prime}(x) v_{\lambda}(x) g_{\lambda}(x) \mathrm{d} x
$$

which is the desired formula with $w_{\lambda}=v_{\lambda} g_{\lambda}$. Moreover, for almost-all $\lambda>1 / \sqrt{2}$, as mentioned above $g_{\lambda}$ is continuous, hence bounded. Then $w_{\lambda} \in L^{q}\left(\mu_{\lambda}\right)$ so that $w_{\lambda} \in$ $L^{q}([0,1])$.

Now, $w_{\lambda}$ a priori depends on the choice of $q$. But if $\tilde{w}_{\lambda}$ is another suitable choice for the same $\lambda$ (and possibly different $q$ ), extending both of them by 0 outside $[0,1]$, we have $\int_{\mathbb{R}} f^{\prime} w_{\lambda}=\int_{\mathbb{R}} f^{\prime} \tilde{w}_{\lambda}$ for all test functions $f$. As a consequence the extensions of $w_{\lambda}$ and $\tilde{w}_{\lambda}$ must differ by a constant, and we thus must have $w_{\lambda}=\tilde{w}_{\lambda}$. It follows that we have a single $w_{\lambda}$ belonging to all $L^{q}([0,1])$ simultaneously.

We also got the following Lipschitz estimate along the way.
Proposition 4.6. For all $q \geq 1$ and all $\lambda, \lambda^{\prime} \in(0,1)$ :

$$
\mathrm{W}_{q}\left(\mu_{\lambda}, \mu_{\lambda^{\prime}}\right) \leq \frac{2}{1-\lambda}\left|\lambda-\lambda^{\prime}\right| .
$$

In particular, $\lambda \mapsto \mu_{\lambda}$ is Lipschitz in the metrics $\mathrm{W}_{q}$ on each interval of the form $[0,1-\varepsilon]$ where $\varepsilon>0$.

Proof. On the one hand, we proved above $\mathrm{W}_{0, q}\left(\left(\Phi^{\lambda}, \eta\right),\left(\Phi^{\lambda^{\prime}}, \eta\right)\right) \leq\left|\lambda-\lambda^{\prime}\right|$; on the other hand, the RIFS ( $\Phi^{\lambda}, \eta$ ) obviously satisfies (3) and (4) for any $q>0$ and constants $\lambda^{q}$ and $A=1$; moreover $m_{x_{0}}^{q}\left(\mu_{\lambda^{\prime}}\right) \leq 1$. Theorem 4.1 then ensures that

$$
\mathrm{W}_{q}\left(\mu_{\lambda}, \mu_{\lambda^{\prime}}\right) \leq \frac{2^{1-\frac{1}{q}} 2^{\frac{1}{q}}}{1-\lambda}\left|\lambda-\lambda^{\prime}\right|
$$

## 5 Approximate overlaps and absolute continuity

### 5.1 Dimension, overlaps and approximation

Consider a RIFS $(\Phi, \eta)$ with finite index set $I$; for simplicity, we shall write $\eta_{i}$ for $\eta(\{i\})$. Given a positive integer $n$, for each word $\omega=\left(i_{1}, \ldots, i_{n}\right) \in I^{n}$ one sets $\phi_{\omega}=\phi_{i_{1}} \circ \cdots \circ \phi_{i_{n}}$ and $\eta_{\omega}=\eta_{i_{1}} \cdot \eta_{i_{2}} \cdots \eta_{i_{n}}$; this gives rise to a new IFS $\Phi^{n}:=\left(\phi_{\omega}\right)_{\omega \in I^{n}}$ and a new probability distribution $\eta^{n}=\left(\eta_{\omega}\right)_{\omega \in I^{n}}$. One defines the upper Lyapunov exponent of $(\Phi, \eta)$ as

$$
\chi=\chi(\Phi, \eta):=\lim _{n \rightarrow \infty} \frac{1}{n} \sum_{\omega \in I^{n}} \eta_{\omega} \log \operatorname{Lip}\left(\phi_{\omega}\right)
$$

(the limit exists by sub-additivity) and the entropy of $\eta$ as

$$
h=H(\eta):=\sum_{i \in I} \eta_{i} \log \frac{1}{\eta_{i}} .
$$

Then assuming $h<+\infty$ and $-\infty<\chi<0$, one can prove

$$
\begin{equation*}
\operatorname{dim} \mu \leq \min \left(d, \frac{h}{-\chi}\right) \tag{10}
\end{equation*}
$$

A folklore conjecture (see e.g. [Hoc14]) asserts that when $X=\mathbb{R}$ and the $\phi_{i}$ are similarities (in which case $\chi=\sum \eta_{i} \log r_{i}$, with $r_{i}$ the contraction ratio of $\phi_{i}$ ) then there is equality in (10) unless there are exact overlaps, i.e. there exist $n \in \mathbb{N}$ and words $\omega \neq \omega^{\prime} \in I^{n}$ such that $\phi_{\omega}=\phi_{\omega^{\prime}}$. When exact overlaps occur, $\mu$ can be expressed as the stationary measure for the IFS made of all distinct compositions of $n$ maps among the $\phi_{i}$ with the corresponding probability vector; then the Lyapunov exponent becomes $n \chi$ but the entropy is less than $n h$, and the bound (10) may improve.

An informal way to understand (10) is to first consider the case when $\eta$ is uniform and all $r_{i}$ are equal. A best approximation of a measure of dimension $\alpha$ by a discrete measure supported on $N$ points should be away from the approximated measure by an error of the order of $1 / N^{\frac{1}{\alpha}}$; the discrete measure $\sum_{\omega \in I_{n}} \eta_{\omega} \phi_{\omega}\left(x_{0}\right)$ has at most $|I|^{n}$ points, and is at $\mathrm{W}_{1}$ distance at most $\simeq e^{n \chi}$ from $\mu$, we get $\alpha \leq h /(-\chi)$ in the limit $n \rightarrow \infty$. Then (10) is a natural generalization in the general case. This point of view immediately suggests that approximate overlaps could in principle result in a lower dimension bounds: instead of considering $\sum_{\omega \in I_{n}} \eta_{\omega} \phi_{\omega}\left(x_{0}\right)$ as the approximation measure, one could take advantage of approximate overlaps to merge the points $\phi_{\omega}\left(x_{0}\right)$ that cluster closely enough, reducing the number of support points (and the entropy). It seems difficult to find a RIFS made of similarities where there are enough approximate overlaps to imply a dimension drop (finding one without exact overlap would disprove the folklore conjecture mentioned above). It is the goal of this section to show that this observation can still be used in some particular circumstances to prove that some stationary measures are not absolutely continuous.

### 5.2 Approximate overlaps can prevent absolute continuity

We shall use the following simple definition.
Definition 5.1. The approximate cardinal of the $n$-th iteration of $\Phi$ at $x_{0}$ up to uncertainty $\varepsilon>0$ is defined as the integer

$$
N\left(\Phi, x_{0}, n, \varepsilon\right)=\min \left\{|E|: E \subset \mathbb{R}^{d}, \forall \omega \in I^{n}, \exists y \in E:\left\|\phi_{\omega}\left(x_{0}\right)-y\right\| \leq \varepsilon\right\}
$$

We could instead ask that most (instead of all) points $\phi_{\omega}\left(x_{0}\right)$ can be approximated by some $y \in E$ and it would also be relevant to consider instead the least entropy of an approximation of $\sum_{\omega \in I^{n}} \eta_{\omega} \delta_{\phi_{\omega}\left(x_{0}\right)}$; however the above definition will suffice to treat the case we are interested in.

Theorem 5.2. Let $\ell \geq 3$ be an integer and set $I=\{0,1, \ldots, \ell-1\}$. Let $\Phi=\left(\phi_{i}\right)_{i \in I}$ be an IFS on $\mathbb{R}^{d}$ made of $\ell$ similarities of ratio $\left(\frac{1}{\ell}\right)^{\frac{1}{d}}$, let $\eta$ be the uniform measure on $I$, and let $\mu$ be the stationary measure. The Lyapunov exponent is then $\chi=\frac{1}{d} \log \frac{1}{\ell}$, the entropy is $h=\log \ell$ and $h /(-\chi)=d$.

If there exist an increasing sequence of positive integers $\left(n_{k}\right)_{k \in \mathbb{N}}$ and a sequence of positive reals $\left(\varepsilon_{k}\right)_{k \in \mathbb{N}}$ with $\varepsilon_{k}=O\left(e^{\chi n_{k}}\right)$, such that

$$
N\left(\Phi, x_{0}, n_{k}, \varepsilon_{k}\right)=o\left(e^{-\chi d n_{k}}\right)
$$

then $\mu$ is singular with respect to the Lebesgue measure.
Proof. By assumption, for each $k$ there exists a set $E_{k} \subset \mathbb{R}^{d}$ with $N_{k}=N\left(\Phi, x_{0}, n_{k}, \varepsilon_{k}\right)$ elements and a map $f: I^{n_{k}} \rightarrow E$ such that for each $\omega \in I^{n_{k}},\left\|\phi_{\omega}\left(x_{0}\right)-f(\omega)\right\| \leq \varepsilon_{k} \lesssim e^{\chi n_{k}}$. Let $\mu_{k}=\sum_{\omega \in I^{n_{k}}} \ell^{-n_{k}} \delta_{\phi_{\omega}\left(x_{0}\right)}$ and $\mu_{k}^{\prime}=\sum_{\omega \in I^{n_{k}}} \ell^{-n_{k}} \delta_{f(\omega)}$. On the one hand by using the transport plan naturally defined by $f$, we have $\mathrm{W}_{1}\left(\mu_{k}, \mu_{k}^{\prime}\right) \leq \varepsilon_{k} \lesssim e^{\chi n_{k}}$. On the other hand, since $\mu_{k}=\mathrm{L}^{n_{k}}\left(\delta_{x_{0}}\right)$ where L is the dual transfer operator, $\mathrm{W}_{1}\left(\mu_{k}, \mu\right) \leq e^{\chi n_{k}} m_{x_{0}}^{1}(\mu)$ ( $\mu$ is compactly supported given the assumptions, thus has finite first moment). The triangular inequality then implies $\mathrm{W}_{1}\left(\mu_{k}^{\prime}, \mu\right) \lesssim e^{\chi d n_{k}}$.

We know that $\mu_{k}^{\prime}$ is supported on $N_{k}$ points; by a classical quantization theorem (see [GL00], Theorem 6.2), if $\mu$ where absolutely continuous we would have $\mathrm{W}_{1}\left(\mu, \mu_{k}^{\prime}\right) \gtrsim$ $1 / N_{k}^{\frac{1}{d}}$, i.e. $N_{k} \gtrsim e^{-\chi n_{k}}$. This is precisely contradicted by the assumption.

Remark 5.3. There are several important restrictions in this result: that the $\phi_{i}$ have equal ratio, that $\eta$ is uniform, and that the dimension bound is exactly $d$. They can probably be lifted if we ask $N\left(\Phi, x_{0}, n_{k}, \varepsilon_{k}\right)=O\left(e^{c d n_{k}}\right)$ for some $c<-\chi$, or more precisely an equivalent bound on the minimal discrete entropy of an $\varepsilon_{k}$-approximation of $\mu_{k}$, in which case we expect to get a suitably improved dimension bound for $\mu$. However, we do not know an example where enough approximate overlaps occur. This explains our need for the equality $h /(-\chi)=d$. If the ratios $r_{i}$ of the $\phi_{i}$ where not equal, then we would have $\mathrm{W}_{1}\left(\mu_{k}, \mu\right) \lesssim \bar{r}^{n_{k}}$ with $\bar{r}$ the arithmetic mean of the $r_{i}$, while $e^{\chi}$ is their geometric mean. Therefore the speed of convergence would not be quick enough to conclude. Using $\mathrm{W}_{q}$ with small $q$ would give better results, but would still
be insufficient without a better bound on $N_{k}$. Last, to tackle the case of non-uniform $\eta$ we would have to use "variable rate" quantization, i.e. approximation results where the size of a discrete measure is measured by its entropy and not the cardinal of its support. In this case, bounds of the kind $N_{k}=o\left(e^{-\chi d n_{k}}\right)$ do not prove $\mu$ to be singular, but only that its density (if it exists) is concentrated on small sets (it has infinite entropy relative to Lebesgue).

### 5.3 A case of application

Let us construct an example where Theorem 5.2 applies, proving Corollary C. We set $d=1, \Phi=\left(\phi_{i}\right)_{i \in I}$ with $I=\{0,1, \ldots, \ell-1\}, \phi_{i}(x)=\frac{x}{\ell}+t_{i}, t_{0}=0, t_{1}=1, t_{2}$ to be determined and $t_{i}$ arbitrary for $i \geq 3$, and $\eta(\{i\})=\frac{1}{\ell}$ for all $i$. The Lyapunov exponent is thus $\chi=\log \frac{1}{\ell}$.

Let $\left(d_{k}\right)_{k \geq 1}$ be a sequence of positive integers, assumed to increase very quickly, more precisely $d_{k+1} \geq k\left(d_{k}+1\right) \ell^{d_{k}+1}$ for all $k$ large enough; assume further that for all $k, d_{k+1}$ is a multiple of $d_{k}+1$. Finally, set

$$
t_{2}=\left(\sum_{j=1}^{\infty} \frac{1}{\ell^{d_{j}}}\right)^{-1} \in[\ell-1,+\infty) .
$$

Let $\tilde{\omega}_{k}=100 \cdots 0 \in I^{d_{k}+1}$ and $\omega_{k}=i_{1}^{k} i_{2}^{k} \cdots i_{d_{k}+1}^{k} \in I^{d_{k}+1}$ be defined by $i_{d_{j}+1}^{k}=1$ for all $j \leq k$ and $i_{m}^{k}=0$ for all other indices $m$. Then

$$
\phi_{\omega_{k}}(0)=t_{2} \sum_{j=1}^{k} \frac{1}{\ell^{d_{k}}} \quad \text { and } \quad \phi_{\tilde{\omega}_{k}}(0)=1
$$

so that for all $x$ and for $k$ large enough,

$$
\begin{aligned}
\left|\phi_{\omega_{k}}(x)-\phi_{\tilde{\omega}_{k}}(x)\right| & =\frac{\sum_{j=k+1}^{\infty} \frac{1}{\ell^{d_{j}}}}{\sum_{j=1}^{\infty} \frac{1}{\ell^{d_{j}}}} \\
& =\frac{t_{2}}{\ell^{d_{k+1}}}\left(\sum_{j=k+1}^{\infty} \frac{1}{\ell^{d_{j}-d_{k+1}}}\right) \\
& \lesssim \frac{1}{\ell^{d_{k+1}}}=e^{\chi d_{k+1}} .
\end{aligned}
$$

We choose $n_{k}=d_{k+1}$, and since all $\phi_{i}$ are contracting we deduce that for all pairs of words $\tau_{1}, \tau_{2} \in I^{n_{k}}$ which have the form $\tau_{1}=\alpha \omega_{k} \beta$ and $\tau_{2}=\alpha \tilde{\omega}_{k} \beta$, we also have

$$
\left|\phi_{\tau_{1}}(0)-\phi_{\tau_{2}}(0)\right|=O\left(e^{\chi n_{k}}\right) .
$$

Given $\omega \in I^{n_{k}}$, decompose $\omega$ into $n_{k} / d_{k}$ blocks of length $d_{k}$ and replace each block equal
to $\tilde{\omega}_{k}$ with a copy of $\omega_{k}$; call the result $\omega^{\prime}$ and set $f(\omega)=\phi_{\omega^{\prime}}(0)$. This yields

$$
\begin{aligned}
N\left(\Phi, n_{k}, C e^{\chi n_{k}}\right) & \leq\left(\ell^{d_{k}+1}-1\right)^{\frac{n_{k}}{d_{k}+1}} \\
& \leq \ell^{n_{k}} e^{\frac{n_{k}}{d_{k}+1} \log \left(1-\ell^{-d_{k}-1}\right)} \\
& \leq \ell^{n_{k}} e^{-C \frac{d_{k+1}}{\left(d_{k}+1\right) e^{d_{k}+1}}} \\
& \leq \ell^{n_{k}} e^{-C k}=o\left(\ell^{n_{k}}\right)=o\left(e^{-\chi n_{k}}\right) .
\end{aligned}
$$

Theorem 5.2 applies, and we deduce that the stationary measure of $(\Phi, \eta)$ is not absolutely continuous.

The liberty in the choice of the tail of the sequence, $\left(d_{k}\right)_{k \geq k_{0}}$ enables to produce $2^{\aleph_{0}}$ possible values for $t_{2}$, so that we proved Corollary C. The additional freedom given for the beginning $\left(d_{k}\right)_{1 \leq k \leq k_{0}}$ of the sequence is a coquetry that can be used to cover a wider range of values.

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[^1]:    ${ }^{1}$ This terminology has been coined in dynamical system, see e.g. [Rue98, Rue09, BS12].

