



**HAL**  
open science

# Invariant measure for stochastic Schrödinger equations

Tristan Benoist, Martin Fraas, Yan Pautrat, Clément Pellegrini

► **To cite this version:**

Tristan Benoist, Martin Fraas, Yan Pautrat, Clément Pellegrini. Invariant measure for stochastic Schrödinger equations. *Annales Henri Poincaré*, 2021, 22 (2), pp.347-374. 10.1007/s00023-020-01001-4. hal-02263935

**HAL Id: hal-02263935**

**<https://hal.science/hal-02263935v1>**

Submitted on 23 May 2024

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

## INVARIANT MEASURE FOR STOCHASTIC SCHRÖDINGER EQUATIONS

T. BENOIST, M. FRAAS, Y. PAUTRAT, AND C. PELLEGRINI

ABSTRACT. Quantum trajectories are Markov processes that describe the time-evolution of a quantum system undergoing continuous indirect measurement. Mathematically, they are defined as solutions of the so-called “Stochastic Schrödinger Equations”, which are nonlinear stochastic differential equations driven by Poisson and Wiener processes. This paper is devoted to the study of the invariant measures of quantum trajectories. Particularly, we prove that the invariant measure is unique under an ergodicity condition on the mean time evolution, and a “purification” condition on the generator of the evolution. We further show that quantum trajectories converge in law exponentially fast towards this invariant measure. We illustrate our results with examples where we can derive explicit expressions for the invariant measure.

## CONTENTS

1. Introduction	1
2. Construction of the model	4
3. Invariant measure and exponential convergence in Wasserstein distance	8
4. Set of invariant measures under <b>(Pur)</b>	11
5. <b>(Pur)</b> is not necessary for purification	13
6. Examples	14
References	19

## 1. INTRODUCTION

Under a Markov approximation, the evolution of an open quantum system  $\mathcal{S}$  in interaction with an environment  $\mathcal{E}$  is described by the Gorini–Kossakowski–Sudarshan–Lindblad Master (GKSL) equation [23, 30]. More precisely, assuming that the system is described by the Hilbert space  $\mathbb{C}^k$ , the set of its states is defined as the set  $\mathcal{D}_k$  of density matrices, i.e. positive semidefinite matrices with trace one:

$$\mathcal{D}_k = \{\rho \in M_k(\mathbb{C}) \text{ s.t. } \rho \geq 0, \text{tr } \rho = 1\}.$$

The evolution  $t \in \mathbb{R}_+ \mapsto \bar{\rho}_t \in \mathcal{D}_k$  of states of the system is then determined by the GKSL equation (also called quantum master equation):

$$d\bar{\rho}_t = \mathcal{L}(\bar{\rho}_t) dt, \quad \bar{\rho}_0 \in \mathcal{D}_k, \tag{1.1}$$

where  $\mathcal{L}$  is a bounded linear operator on  $M_k(\mathbb{C})$  of the form

$$\mathcal{L} : \rho \mapsto -i[H, \rho] + \sum_{i \in I} (V_i \rho V_i^* - \frac{1}{2}\{V_i^* V_i, \rho\}), \tag{1.2}$$

with  $I$  a finite set,  $H \in M_k(\mathbb{C})$  self adjoint, and  $V_i \in M_k(\mathbb{C})$  for each  $i \in I$  ( $[\cdot, \cdot]$  and  $\{\cdot, \cdot\}$  are respectively the commutator and anticommutator). Such an  $\mathcal{L}$  is called a Lindblad operator.

Since  $\mathcal{L}$  is linear,  $t \mapsto \bar{\rho}_t$  is given by  $\bar{\rho}_t = e^{t\mathcal{L}}(\bar{\rho}_0)$ . The flow is therefore a semigroup  $(e^{t\mathcal{L}})_t$ , which consists of completely positive, trace-preserving maps (see [38]). In particular,  $\mathcal{L}$  is the generator of a semigroup of contractions, thus  $\text{spec } \mathcal{L} \subset \{\lambda \in \mathbb{C} \text{ s.t. } \text{Re } \lambda \leq 0\}$ . Since  $e^{t\mathcal{L}}$  is trace preserving,  $0 \in \text{spec } \mathcal{L}$ . The following assumption is equivalent to the simplicity of the eigenvalue 0 [38, Proposition 7.6]:

**( $\mathcal{L}$ -erg):** There exists a unique non zero minimal orthogonal projection  $\pi$  such that  $\mathcal{L}(\pi M_k(\mathbb{C})\pi) \subset \pi M_k(\mathbb{C})\pi$ .

Assumption **( $\mathcal{L}$ -erg)** implies directly that there exists a unique  $\rho_{\text{inv}} \in \mathcal{D}_k$  such that  $\mathcal{L}\rho_{\text{inv}} = 0$ . Moreover, one can show that **( $\mathcal{L}$ -erg)** implies the existence of  $\lambda > 0$  such that for any  $\rho \in \mathcal{D}_k$ ,  $e^{t\mathcal{L}}(\rho) = \rho_{\text{inv}} + O(e^{-\lambda t})$  (see [38, Proposition 7.5]).

The above framework generalizes that of continuous-time Markov semigroups on a finite number of sites: density matrices  $\rho$  over  $\mathbb{C}^k$  generalize probability distributions over  $k$  classical states, while Lindbladians  $\mathcal{L}$  generalize generators of Markov jump processes. In Section 6.4, we show how a classical finite state Markov jump process can be encoded in the present formalism.

The family  $(\bar{\rho}_t)_t$  describes the reduced evolution of the system  $\mathcal{S}$  when coupled to an environment  $\mathcal{E}$  in a conservative manner. This evolution can be derived by considering the full Hamiltonian of  $\mathcal{S} + \mathcal{E}$  in relevant limiting regimes, e.g. the *weak coupling* or *fast repeated interactions* regimes, and tracing out the environment degrees of freedom (see [17, 18] and [1] respectively). It can also be described by a *stochastic unravelling*, i.e. a stochastic process  $(\rho_t)_t$  with values in  $\mathcal{D}_k$  such that the expectation  $\bar{\rho}_t$  of  $\rho_t$  satisfies (1.1); this method was developed in [4, 5, 6]. One possible choice of a stochastic unravelling is described by the following stochastic differential equation (SDE), called a *stochastic master equation*:

$$\begin{aligned} d\rho_t = & \mathcal{L}(\rho_{t-}) dt \\ & + \sum_{i \in I_b} \left( L_i \rho_{t-} + \rho_{t-} L_i^* - \text{tr}(\rho_{t-} (L_i + L_i^*)) \rho_{t-} \right) dB_i(t) \\ & + \sum_{j \in I_p} \left( \frac{C_j \rho_{t-} C_j^*}{\text{tr}(C_j \rho_{t-} C_j^*)} - \rho_{t-} \right) \left( dN_j(t) - \text{tr}(C_j \rho_{t-} C_j^*) dt \right), \end{aligned} \quad (1.3)$$

where

- $I = I_b \cup I_p$  is a partition of  $I$  such that  $L_i = V_i$  for  $i \in I_b$  and  $C_j = V_j$  for  $j \in I_p$ ,
- each  $B_i$  is a Brownian motion,
- each  $N_j$  is a Poisson process of intensity  $t \mapsto \int_0^t \text{tr}(C_j \rho_s - C_j^*) ds$ .

*Remark 1.* The processes  $(B_j(t))_t$  and  $(N_j(t) - \int_0^t \text{tr}(C_j \rho_s - C_j^*) ds)_t$  are actually martingales. Then assuming that (1.3) accepts a solution, it is easy to check that for any  $t \geq 0$ , the expectation of  $\rho_t$  is equal to  $\bar{\rho}_t$  whenever  $\rho_0 = \bar{\rho}_0$ .

Proper definitions of these Poisson processes and proofs of existence and uniqueness of the solution to (1.3) can be found in [5, 6, 33, 34, 35]. A solution  $(\rho_t)_t$  of Equation (1.3) is called a *quantum trajectory*.

Equations of the form (1.3) are used to model experiments in quantum optics (photo-detection, heterodyne or homodyne interferometry), particularly for measurement and control (see [15, 24, 37]). They were also introduced as stochastic collapse models (see [19, 22]) and as numerical tools to compute  $\bar{\rho}_t$  (see [16]). Here we are interested in the fact that they model the evolution of the system  $\mathcal{S}$  when continuous measurements are done on the environment  $\mathcal{E}$ . This can be shown starting from quantum stochastic differential equations using quantum filtering [3, 10, 13, 21, 25]. An approach using the notion of a priori and a posteriori states has been also developed using ‘‘classical’’ stochastic calculus (see the reference book by Barchielli and Gregoratti [5], and references therein). Continuous-time limits of discrete-time models can also be considered, see [33, 34, 35].

Equation (1.3) has the property that if  $\rho_0$  is an extreme point of  $\mathcal{D}_k$ , then  $\rho_t$  is almost surely an extreme point of  $\mathcal{D}_k$  for any  $t \in \mathbb{R}_+$ . Since we will extensively use this property, let us make it explicit. The extreme points of  $\mathcal{D}_k$  are the rank-one orthogonal projectors of  $\mathbb{C}^k$ ; for any  $x \in \mathbb{C}^k \setminus \{0\}$ , let  $\hat{x}$  be its equivalence class in  $\text{PC}^k$ , the projective space of  $\mathbb{C}^k$ . For  $\hat{x} \in \text{PC}^k$ , let  $\pi_{\hat{x}}$  be the orthogonal projector onto  $\mathbb{C}x$ . Then  $\hat{x} \in \text{PC}^k \mapsto \pi_{\hat{x}}$  is a bijective map from  $\text{PC}^k$  to the set of extreme points. Assume now that  $\rho_0 = \pi_{\hat{x}_0}$  for some  $\hat{x}_0 \in \text{PC}^k$ . Then it is easy to check that  $\rho_t = \pi_{\hat{x}_t}$  almost surely for any  $t \in \mathbb{R}_+$ , with  $t \mapsto x_t$  the unique solution to the following SDE,

called a *stochastic Schrödinger equation*:

$$\begin{aligned} dx_t &= D(x_{t-})x_{t-} dt \\ &+ \sum_{i \in I_b} (L_i - \frac{1}{2}v_i(t-) \text{Id})x_{t-} dB_i(t) \\ &+ \sum_{j \in I_p} \left( \frac{C_j}{\sqrt{n_j(t-)}} - \text{Id} \right) x_{t-} dN_j(t), \end{aligned} \tag{1.4}$$

for  $x_0 \in \hat{x}_0$  of norm one, where the operator  $D(x_{t-})$  is defined as

$$D(x_{t-}) = -(\mathbf{i}H + \frac{1}{2} \sum_{i \in I_b} L_i^* L_i + \frac{1}{2} \sum_{j \in I_p} C_j^* C_j) + \frac{1}{2} \sum_{i \in I_b} v_i(t-) (L_i - \frac{1}{4} v_i(t-) \text{Id}) + \frac{1}{2} \sum_{j \in I_p} n_j(t-),$$

with

$$v_i(t-) = \langle x_{t-}, (L_i + L_i^*)x_{t-} \rangle, \quad n_j(t-) = \langle x_{t-}, C_j^* C_j x_{t-} \rangle = \|C_j x_{t-}\|^2.$$

The brackets  $\langle \cdot, \cdot \rangle$  denote the scalar product in  $\mathbb{C}^k$ . Without possible confusion, a solution  $(x_t)_t$  will be also called a quantum trajectory. Remark that  $\|x_0\| = 1$  implies  $\|x_t\| = 1$  almost surely for any  $t \in \mathbb{R}_+$ ; remark also that the numerical computation of  $\rho_t$  involves only multiplications of matrices with vectors and not multiplications of matrices (this is the motivation for the use of quantum trajectories as numerical tools mentioned above).

In the physics literature, extreme points of  $\mathcal{D}_k$  are called pure states. In particular, the preceding paragraph shows that the evolution dictated by Eq. (1.3) preserves pure states. It actually has also the property that quantum trajectories (solution of (1.3)) tend to “purify”. This has been formalized by Maassen and Kümmerer in [31] for discrete-time quantum trajectories, and extended to the continuous-time case by Barchielli and Paganoni in [7]. Purification is related to the following assumption (here  $A \propto B$  means there exists  $\lambda \in \mathbb{C}$  such that  $A = \lambda B$  or  $\lambda A = B$ . Particularly we allow for  $\lambda = 0$ ).

**(Pur):** Any non zero orthogonal projector  $\pi$  such that for all  $i \in I_b$ ,  $\pi(L_i + L_i^*)\pi \propto \pi$  and for all  $j \in I_p$ ,  $\pi C_j^* C_j \pi \propto \pi$  has rank one.

As shown in [7], **(Pur)** implies that for any  $\rho_0 \in \mathcal{D}_k$

$$\lim_{t \rightarrow \infty} \inf_{\hat{y} \in \text{PC}^k} \|\rho_t - \pi_{\hat{y}}\| = 0 \quad \text{almost surely.} \tag{1.5}$$

The main goal of this article is to show how the exponential convergence of the solution  $(\bar{\rho}_t)_t$  of Eq. (1.1), induced by **(L-erg)**, translates for its stochastic unravelling  $(\rho_t)_t$  solution of Eq. (1.3). We prove uniqueness of the invariant measure for continuous-time quantum trajectories assuming both **(L-erg)** and **(Pur)**. From (1.5), under these assumptions, the invariant measure will be concentrated on pure states, so we only need to prove uniqueness of the invariant measure for  $(\hat{x}_t)_t$  equivalence class of  $(x_t)_t$  solution of (1.4) (since  $\text{PC}^k$  is compact and the involved process is Feller, the existence of an invariant measure is obvious). The difficulty of this proof lies in the failure of usual techniques like  $\varphi$ -irreducibility. Note that this question has already been partially addressed in the literature: essentially, only diffusive equations have been considered, i.e. equations for which Eq. (1.3) or (1.4) contain no jump term (in our notation,  $I_p = \emptyset$ ). The results of [7] were, to our knowledge, the most advanced ones so far. In that article, algebraic conditions on the vector fields describing the stochastic differential equation are imposed to obtain the uniqueness of the invariant measure. This allows the authors to apply directly standard results from the analysis of stochastic differential equations. Unfortunately their assumptions are hard to check for a given family of matrices  $(L_i)_{i \in I_b}$ .

The main result of the present paper is the following theorem.

**Theorem 1.1.** *Assume that **(Pur)** and **(L-erg)** hold. Then the Markov process  $(\hat{x}_t)_t$  has a unique invariant probability measure  $\mu_{\text{inv}}$ , and there exist  $C > 0$  and  $\lambda > 0$  such that for any initial distribution  $\mu$  of  $\hat{x}_0$  over  $\text{PC}^k$ , for all  $t \geq 0$ , the distribution  $\mu_t$  of  $\hat{x}_t$  satisfies*

$$W_1(\mu_t, \mu_{\text{inv}}) \leq C e^{-\lambda t}$$

where  $W_1$  is the Wasserstein distance of order 1.

This theorem is more general than previous similar results in different ways. First, we consider stochastic Schrödinger equations involving both Poisson and Wiener processes. Second, our assumptions are standard for quantum trajectories and are easy to check for a given family of operators  $(H, (L_i)_{i \in I_b}, (C_j)_{j \in I_p})$ . Last, we prove an exponential convergence towards the invariant measure. As a byproduct, we also provide a simple proof of the purification expressed in Eq. (1.5) (see Proposition 2.5). To complete the picture, assuming only **(Pur)**, we show that **(L-erg)** is necessary. We also provide a complete characterization of the set of invariant measures of  $(\hat{x}_t)$  whenever **(L-erg)** does not hold (see Proposition 4.2). Arguments in Sections 3 and 4 are adaptations of [11], where similar results for discrete-time quantum trajectories are considered.

The paper is structured as follows. In Section 2, we give a precise description of the model of quantum trajectories with a proper definition of the underlying probability space. In particular, we introduce a new martingale which is central to our proofs. In Section 3, we prove Theorem 1.1. In Section 4 we derive the full set of invariant measures assuming only **(Pur)**. In Section 5 we show that **(Pur)** is not necessary even if **(L-erg)** holds. In Section 6, we provide some examples of explicit invariant measures. In Section 6.4 we provide an encoding of any classical finite state Markov jump process into a stochastic master equation.

## 2. CONSTRUCTION OF THE MODEL

**2.1. Construction of quantum trajectories.** In this section we fix the notations and introduce the probability space we use to study  $(\hat{x}_t)_t$ . First, for an element  $x \neq 0$  of  $\mathbb{C}^k$ , and for an operator  $A$  with  $Ax \neq 0$  we denote

$$A \cdot \hat{x} = \widehat{Ax}.$$

We consider the following distance on  $\text{PC}^k$ :

$$d(\hat{x}, \hat{y}) = \sqrt{1 - |\langle x, y \rangle|^2}, \quad (2.1)$$

for all  $\hat{x}, \hat{y} \in \text{PC}^k$ , where  $x$  and  $y$  are norm-one representatives of  $\hat{x}$  and  $\hat{y}$  respectively. We equip  $\text{PC}^k$  with the associated Borel  $\sigma$ -algebra denoted by  $\mathcal{B}$ .

Now we introduce a stochastic process with values in  $M_k(\mathbb{C})$ . Let  $(\Omega, (\mathcal{F}_t)_t, \mathbb{P})$  be a filtered probability space with standard brownian motions  $W_i$  for  $i \in I_b$ , and standard Poisson processes  $N_j$  for  $j \in I_p$ , such that the full family  $(W_i, N_j; i \in I_b, j \in I_p)$  is independent. The filtration  $(\mathcal{F}_t)_t$  is assumed to satisfy the standard conditions, and we denote  $\mathcal{F}_\infty$  by  $\mathcal{F}$  and the processes  $(W_i(t))_t$  and  $(N_j(t) - t)_t$  are  $(\mathcal{F}_t)_t$ -martingales under  $\mathbb{P}$ . We denote by  $\mathbb{E}$  the expectation with respect to  $\mathbb{P}$ .

On  $(\Omega, (\mathcal{F}_t)_t, \mathbb{P})$ , for  $s \in \mathbb{R}_+$ , let  $(S_t^s)_{t \in [s, \infty)}$  be the solution to the following SDE:

$$dS_t^s = \left( K + \frac{\#I_p}{2} \text{Id} \right) S_{t-}^s dt + \sum_{i \in I_b} L_i S_{t-}^s dW_i(t) + \sum_{j \in I_p} (C_j - \text{Id}) S_{t-}^s dN_j(t), \quad S_s^s = \text{Id} \quad (2.2)$$

( $\#I_p$  is the cardinal of  $I_p$ ), where

$$K = -iH - \frac{1}{2} \left( \sum_{i \in I_b} L_i^* L_i + \sum_{j \in I_p} C_j^* C_j \right).$$

Since standard Cauchy–Lipschitz conditions are fulfilled, the SDE defining  $(S_t^s)_t$  has indeed a unique (strong) solution. We denote  $S_t := S_t^0$ . Note that for  $s$  fixed the process  $(S_t^s)_t$  is independent of  $\mathcal{F}_s$ , and we have that for all  $0 \leq r \leq s \leq t$

$$S_t^s S_s^r = S_t^r.$$

In addition, for any  $\rho \in \mathcal{D}_k$ , let  $(Z_t^\rho)_t$  be the positive real-valued process defined by

$$Z_t^\rho = \text{tr}(S_t^* S_t \rho),$$

and let  $(\rho_t)_t$  be the  $\mathcal{D}_k$ -valued process defined by

$$\rho_t = \frac{S_t \rho S_t^*}{\text{tr}(S_t \rho S_t^*)}$$

if  $Z_t^\rho \neq 0$ , taking an arbitrarily fixed value whenever  $Z_t^\rho = 0$  (this value will always appear with probability zero in the sequel).

The following results on the properties of  $(Z_t^\rho)_t$  were proven in [6]. We give short proofs adapted to our restricted setting where the Hilbert space is finite-dimensional, and  $I = I_b \cup I_p$  is a finite set.

**Lemma 2.1.** *For any  $\rho \in \mathcal{D}_k$ , the stochastic process  $(Z_t^\rho)_t$  is the unique solution of the SDE*

$$dZ_t^\rho = Z_{t-}^\rho \left( \sum_{i \in I_b} \text{tr}((L_i + L_i^*)\rho_{t-}) dW_i(t) + \sum_{j \in I_p} (\text{tr}(C_j^* C_j \rho_{t-}) - 1) (dN_j(t) - dt) \right), \quad Z_0^\rho = 1.$$

Moreover,  $(Z_t^\rho)_t$  is a nonnegative martingale under  $\mathbb{P}$ .

*Proof.* The fact that  $(Z_t^\rho)_t$  verifies the given SDE is a direct application of the Itô formula. Since  $(\rho_t)_t$  takes its values in the compact space  $\mathcal{D}_k$ , that SDE verifies standard Cauchy–Lipschitz conditions, ensuring the uniqueness of the solution. Since the processes  $(W_i(t))_t$  and  $(N_j(t) - t)_t$  are  $\mathbb{P}$ -martingales, it follows that  $(Z_t^\rho)_t$  is a  $\mathbb{P}$ -local martingale. Since  $\text{tr}(C_j^* C_j \rho) \geq 0$  for any  $j \in I_p$  and  $\rho \in \mathcal{D}_k$ , and  $(\rho_t)_t$  takes value in the compact space  $\mathcal{D}_k$ , it follows from [27, Theorem 12] that  $(Z_t^\rho)_{t \in [0, T]}$  is a  $\mathbb{P}$ -nonnegative martingale for all  $T$ .  $\square$

For any  $\rho \in \mathcal{D}_k$ , we define a probability  $\mathbb{P}_t^\rho$  on  $(\Omega, \mathcal{F}_t)$ :

$$d\mathbb{P}_t^\rho = Z_t^\rho d\mathbb{P}|_{\mathcal{F}_t}. \quad (2.3)$$

Since  $(Z_t^\rho)_t$  is a  $\mathbb{P}$ -martingale from Lemma 2.1, the family  $(\mathbb{P}_t^\rho)_t$  is consistent, that is  $\mathbb{P}_t^\rho(E) = \mathbb{P}_s^\rho(E)$  for  $t \geq s$  and  $E \in \mathcal{F}_s$ . Kolmogorov's extension theorem defines a unique probability on  $(\Omega, \mathcal{F}_\infty)$ , which we denote by  $\mathbb{P}^\rho$ . We will denote by  $\mathbb{E}^\rho$  the expectation with respect to  $\mathbb{P}^\rho$ .

The following proposition makes explicit the relationship between  $\mathbb{P}$  and  $\mathbb{P}^\rho$ . It follows from a direct application of Girsanov's change of measure Theorem (see [26, Theorems III.3.24 and III.5.19]). For all  $i \in I_b$  and  $t \in \mathbb{R}_+$ , let

$$B_i^\rho(t) = W_i(t) - \int_0^t \text{tr}((L_i + L_i^*)\rho_{s-}) ds.$$

**Proposition 2.2.** *Let  $\rho \in \mathcal{D}_k$ . Then, with respect to  $\mathbb{P}^\rho$ , the processes  $\{B_i^\rho\}_{i \in I_b}$  are independent Wiener processes and the processes  $\{N_j\}_{j \in I_p}$  are point processes of respective stochastic intensity  $\{t \mapsto \text{tr}(C_j^* C_j \rho_{t-})\}_{j \in I_p}$ .*

The process  $(\rho_t)_t$  considered under  $\mathbb{P}^\rho$  models the evolution of a Markov open quantum system subject to indirect measurements. We refer the reader to [5, 14, 15] and references therein for a more detailed discussion of this interpretation.

From Itô calculus,  $(\rho_t)_t$  is solution of the SDE

$$\begin{aligned} d\rho_t &= \mathcal{L}(\rho_{t-}) dt \\ &+ \sum_{i \in I_b} \left( L_i \rho_{t-} + \rho_{t-} L_i^* - \text{tr}(\rho_{t-} (L_i + L_i^*)) \rho_{t-} \right) dB_i^\rho(t) \\ &+ \sum_{j \in I_p} \left( \frac{C_j \rho_{t-} C_j^*}{\text{tr}(C_j \rho_{t-} C_j^*)} - \rho_{t-} \right) (dN_j(t) - \text{tr}(C_j \rho_{t-} C_j^*) dt). \end{aligned} \quad (2.4)$$

Proposition 2.2 then implies that with respect to  $\mathbb{P}^\rho$ , the process  $(\rho_t)_t$  is indeed the unique solution of (1.3) with  $\rho_0 = \rho$ . Similarly, if  $\rho_0 = \pi_{\hat{x}}$  for some  $\hat{x} \in \text{PC}^k$ , then with respect to  $\mathbb{P}^{\pi_{\hat{x}}}$ , the process  $\left( \frac{S_t x}{\|S_t x\|} \right)_t$  is the solution of (1.4) with  $x$  any norm one representative of  $\hat{x}$ .

Remark also that for any  $\rho \in \mathcal{D}_k$ , using (2.3), one has from Remark 1

$$\mathbb{E}(S_t \rho S_t^*) = \mathbb{E}(\rho_t Z_t^\rho) = \mathbb{E}^\rho(\rho_t) = e^{t\mathcal{L}}(\rho). \quad (2.5)$$

Our strategy of proof is based on the study of the joint distribution of  $S_t$  and a random initial state  $\hat{x}$ . To this end, we consider the product space  $\Omega \times \text{PC}^k$  equipped with the filtration  $(\mathcal{F}_t \otimes \mathcal{B})_t$

and the full  $\sigma$ -algebra  $\mathcal{F} \otimes \mathcal{B}$ . For any probability measure  $\mu$  on  $\text{PC}^k$ , and for all  $E \in \mathcal{F}$  and  $A \in \mathcal{B}$ , let

$$\mathbb{Q}_\mu(E \times A) = \int \mathbb{P}^{\pi_{\hat{x}}}(E) \mathbb{1}_{\hat{x} \in A} d\mu(\hat{x}).$$

We will denote by  $\mathbb{E}_\mu$  the expectation with respect to  $\mathbb{Q}_\mu$ . Note that  $d\mathbb{P}_t^{\pi_{\hat{x}}} = \|S_t x\|^2 d\mathbb{P}$  for any  $\hat{x} \in \text{PC}^k$ , so that  $\mathbb{P}^{\pi_{\hat{x}}}(\{S_t x = 0\}) = 0$  for all  $x \in \hat{x}$ . Therefore

$$\mathbb{Q}_\mu(\{S_t x = 0\}) = 0$$

and there exists a process  $(\hat{x}_t)_t$  for which

$$\hat{x}_t = S_t \cdot x$$

holds almost surely. It has the same distribution as the image by the map  $x \mapsto \hat{x}$  of the solution  $(x_t)_t$  to (1.4) with  $x_0 \in \hat{x}$ ,  $\|x_0\| = 1$ .

The following proposition shows that the laws of any  $\mathcal{F}$ -measurable random variables are given by a marginal of  $\mathbb{Q}_\mu$ . For a probability measure  $\mu$  on  $\text{PC}^k$ , we define

$$\rho_\mu := \mathbb{E}_\mu(\pi_{\hat{x}}).$$

**Proposition 2.3.** *Let  $\mu$  be a probability measure on  $\text{PC}^k$ , then  $\rho_\mu \in \mathcal{D}_k$  and for any  $E \in \mathcal{F}$ ,*

$$\mathbb{Q}_\mu(E \times \text{PC}^k) = \mathbb{P}^{\rho_\mu}(E).$$

*Proof.* The fact that  $\rho_\mu \in \mathcal{D}_k$  follows from the positivity and linearity of the expectation. Concerning the second part, let  $t \geq 0$  and  $E \in \mathcal{F}_t$ , then

$$\mathbb{Q}_\mu(E \times \text{PC}^k) = \int \mathbb{P}^{\pi_{\hat{x}}}(E) d\mu(x) = \int \int_E \text{tr}(S_t^* S_t \pi_{\hat{x}}) d\mathbb{P} d\mu(x).$$

Fubini's Theorem implies

$$\mathbb{Q}_\mu(E \times \text{PC}^k) = \int_E \text{tr}(S_t^* S_t \rho_\mu) d\mathbb{P} = \int_E Z_t^{\rho_\mu} d\mathbb{P} = \mathbb{P}_t^{\rho_\mu}(E).$$

The uniqueness of the extended measure in Kolmogorov's extension Theorem yields the proposition.  $\square$

*Remark 2.* Any  $\mathcal{F}$ -measurable random variables  $X$  can be extended canonically to a  $\mathcal{F} \otimes \mathcal{B}$ -measurable random variables setting  $X(\omega, \hat{x}) = X(\omega)$ . Proposition 2.3 then implies that the distribution of a  $\mathcal{F}$ -measurable random variable under  $\mathbb{Q}_\mu$  depends on  $\mu$  only through  $\rho_\mu$ . The central idea of our proof is that assumption **(Pur)** will allow us to find a  $\mathcal{F}$ -measurable process approximating  $(\hat{x}_t)_t$ . The  $\mathcal{F}$ -measurability of the process will then imply that it inherits some ergodicity properties from assumption **(L-erg)**.

*Remark 3.* If  $\mu_{\text{inv}}$  is an invariant measure for the Markov chain  $(\hat{x}_t)_t$ , then with the above notation,  $\rho_{\mu_{\text{inv}}}$  is an invariant state for  $(e^{t\mathcal{L}})_t$ . In particular, if **(L-erg)** holds then  $\rho_{\mu_{\text{inv}}} = \rho_{\text{inv}}$ . This follows from the identities

$$e^{t\mathcal{L}}(\rho_{\mu_{\text{inv}}}) = \int e^{t\mathcal{L}}(\pi_{\hat{x}}) d\mu_{\text{inv}} = \int S_t \pi_{\hat{x}} S_t^* d\mathbb{P} d\mu_{\text{inv}}(\hat{x}) = \int \pi_{\hat{x}} d\mu_{\text{inv}}(\hat{x}) = \rho_{\mu_{\text{inv}}}$$

where the second identity uses (2.5).

**2.2. Key martingale.** The following process is the key to construct a  $\mathcal{F}$ -measurable process approximating  $(\hat{x}_t)_t$ . For any  $t \geq 0$ , let

$$M_t = \frac{S_t^* S_t}{\text{tr}(S_t^* S_t)}, \quad (2.6)$$

whenever  $\text{tr}(S_t^* S_t) \neq 0$ , and give  $M_t$  a fixed arbitrary value whenever  $\text{tr}(S_t^* S_t) = 0$ . Since, by definition, for any  $\rho \in \mathcal{D}_k$ ,  $\mathbb{P}^\rho(\{\text{tr} S_t^* S_t = 0\}) = 0$ , the arbitrary definition of  $M_t$  on this set of vanishing probability is irrelevant. It turns out that with respect to  $\mathbb{P}^{\text{Id}/k}$ ,  $(M_t)_t$  is a martingale. For convenience we write  $\mathbb{P}^{\text{ch}} = \mathbb{P}^{\text{Id}/k}$  and similarly for any other  $\rho$ -dependent object, whenever  $\rho = \text{Id}/k$ .

**Theorem 2.4.** *With respect to  $\mathbb{P}^{\text{ch}}$ , the stochastic process  $(M_t)_t$  is a bounded martingale. Therefore, it converges  $\mathbb{P}^{\text{ch}}$ -almost surely and in  $L^1$  to a random variable  $M_\infty$ . Moreover, for any  $\rho \in \mathcal{D}_k$ ,*

$$d\mathbb{P}^\rho = k \operatorname{tr}(\rho M_\infty) d\mathbb{P}^{\text{ch}},$$

and  $(M_t)_t$  converges almost surely and in  $L^1$  to  $M_\infty$  with respect to  $\mathbb{P}^\rho$ .

*Proof.* Expressing  $(S_t)_t$  in terms of  $B_i^{\text{ch}}$  for  $i \in I_b$ , we have that

$$dS_t = \left( K + \frac{\#I_p}{2} \operatorname{Id} + \sum_{i \in I_b} \frac{\operatorname{tr}(S_{t-}^* (L_i + L_i^*) S_{t-})}{\operatorname{tr}(S_{t-}^* S_{t-})} L_i \right) S_{t-} dt + \sum_{i \in I_b} L_i S_{t-} dB_i^{\text{ch}}(t) + \sum_{j \in I_p} (C_j - \operatorname{Id}) S_{t-} dN_j(t).$$

Recall that the distributions of the  $B_i^{\text{ch}}$  and  $N_j$  under  $\mathbb{P}^{\text{ch}}$  are given by Proposition 2.2.

Since  $\operatorname{tr}(S_t^* S_t)$  is  $\mathbb{P}^{\text{ch}}$ -almost surely non zero, we can define  $R_t$  by  $R_t = S_t / \sqrt{\operatorname{tr}(S_t^* S_t)}$  almost surely for  $\mathbb{P}^{\text{ch}}$ , and therefore for  $\mathbb{P}^\rho$  and  $\mathbb{Q}_\mu$ . The Itô formula implies

$$\begin{aligned} dM_t &= \sum_{i \in I_b} \left( R_{t-}^* (L_i + L_i^*) R_{t-} - M_{t-} \operatorname{tr}(R_{t-}^* (L_i + L_i^*) R_{t-}) \right) dB_i^{\text{ch}}(t) \\ &\quad + \sum_{j \in I_p} \left( \frac{R_{t-}^* C_j^* C_j R_{t-}}{\operatorname{tr}(R_{t-}^* C_j^* C_j R_{t-})} - M_{t-} \right) (dN_j(t) - \operatorname{tr}(R_{t-}^* C_j^* C_j R_{t-}) dt). \end{aligned}$$

Hence, with respect to  $\mathbb{P}^{\text{ch}}$ ,  $(M_t)_t$  is a local martingale. By definition, it is positive-semidefinite, and is also bounded since  $\operatorname{tr}(M_t) = 1$  almost surely. Thus  $(M_t)_t$  is a martingale and standard theorems of convergence for martingales imply the convergence almost surely and in  $L^1$ .

By direct computation, we get  $d\mathbb{P}^\rho|_{\mathcal{F}_t} = k \operatorname{tr}(\rho M_t) d\mathbb{P}^{\text{ch}}|_{\mathcal{F}_t}$ . The  $L^1$  convergence of  $(M_t)_t$  with respect to  $\mathbb{P}^{\text{ch}}$  then implies  $d\mathbb{P}^\rho = k \operatorname{tr}(\rho M_\infty) d\mathbb{P}^{\text{ch}}$ . Finally the inequality  $\operatorname{tr}(AB) \leq \|A\| \operatorname{tr}(B)$  for any two positive semidefinite matrices implies  $\mathbb{P}^\rho \leq k \mathbb{P}^{\text{ch}}$ , which yields the  $L^1$  and almost sure convergence with respect to  $\mathbb{P}^\rho$ .  $\square$

Now we are in the position to show that under the assumption **(Pur)** the limit  $M_\infty$  is a rank-one projector. To this end let us introduce the polar decomposition of  $(S_t)_t$ : there exists a process  $(U_t)_t$  with values in the set of  $k \times k$  unitary matrices such that for all  $t \geq 0$

$$S_t = \sqrt{\operatorname{tr}(S_t^* S_t)} U_t M_t^{1/2}.$$

**Proposition 2.5.** *Assume that **(Pur)** holds. Then, for any  $\rho \in \mathcal{D}_k$ ,  $\mathbb{P}^\rho$ -almost surely, the random variable  $M_\infty$  is a rank-one orthogonal projector on  $\mathbb{C}^k$ .*

*Proof.* First, since  $\mathbb{P}^\rho$  is absolutely continuous with respect to  $\mathbb{P}^{\text{ch}}$ , proving the result with  $\rho = \operatorname{Id}/k$  is sufficient. To achieve this, remark that the  $\mathbb{P}^{\text{ch}}$ -almost sure convergence of  $(M_t)_t$  and the  $\mathbb{P}^{\text{ch}}$ -almost sure bound  $\sup_{t \geq 0} \|M_t\| \leq 1$  imply the convergence of  $\mathbb{E}^{\text{ch}}(M_t^2)$ . Now recall that  $R_t = S_t / \sqrt{\operatorname{tr}(S_t^* S_t)}$ . The Itô isometry implies

$$\begin{aligned} \mathbb{E}^{\text{ch}}(M_t^2) &= M_0^2 + \sum_{i \in I_b} \int_0^t \mathbb{E}^{\text{ch}} \left( R_s^* (L_i + L_i^*) R_s - M_s \operatorname{tr}(R_s^* (L_i + L_i^*) R_s) \right)^2 ds \\ &\quad + \sum_{j \in I_p} \int_0^t \mathbb{E}^{\text{ch}} \left( \left( \frac{R_s^* C_j^* C_j R_s}{\operatorname{tr}(R_s^* C_j^* C_j R_s)} - M_s \right)^2 \operatorname{tr}(R_s^* C_j^* C_j R_s) \right) ds. \end{aligned}$$

Therefore, the convergence of  $\mathbb{E}^{\text{ch}}(M_t^2)$  to  $\mathbb{E}^{\text{ch}}(M_\infty^2)$  implies that

$$\int_0^\infty \mathbb{E}^{\text{ch}} \left( R_s^* (L_i + L_i^*) R_s - M_s \operatorname{tr}(R_s^* (L_i + L_i^*) R_s) \right)^2 ds < \infty$$

for all  $i \in I_b$  and

$$\int_0^\infty \mathbb{E}^{\text{ch}} \left( \left( \frac{R_s^* C_j^* C_j R_s}{\operatorname{tr}(R_s^* C_j^* C_j R_s)} - M_s \right)^2 \operatorname{tr}(R_s^* C_j^* C_j R_s) \right) ds < \infty$$



for all  $j \in I_p$ . Since the integrands are nonnegative, their inferior limits at infinity are 0. Hence there exists an unbounded increasing sequence  $(t_n)_n$  such that for any  $i \in I_b$ ,

$$\lim_n \mathbb{E}^{\text{ch}} \left( R_{t_n}^* (L_i + L_i^*) R_{t_n} - M_{t_n} \text{tr} (R_{t_n}^* (L_i + L_i^*) R_{t_n}) \right) = 0$$

and for any  $j \in I_p$ ,

$$\lim_n \mathbb{E}^{\text{ch}} \left( \left( \frac{R_{t_n}^* C_j^* C_j R_{t_n}}{\text{tr}(R_{t_n}^* C_j^* C_j R_{t_n})} - M_{t_n} \right)^2 \text{tr}(R_{t_n}^* C_j^* C_j R_{t_n}) \right) = 0.$$

Since convergence in  $L^1$  implies the almost sure convergence of a subsequence, there exists an unbounded increasing sequence, which we denote also by  $(t_n)_n$ , such that  $\mathbb{P}^{\text{ch}}$ -almost surely,

$$\lim_{n \rightarrow \infty} \left( R_{t_n}^* (L_i + L_i^*) R_{t_n} - M_{t_n} \text{tr}(R_{t_n}^* (L_i + L_i^*) R_{t_n}) \right) = 0$$

and

$$\lim_{n \rightarrow \infty} \left( \left( \frac{R_{t_n}^* C_j^* C_j R_{t_n}}{\text{tr}(R_{t_n}^* C_j^* C_j R_{t_n})} - M_{t_n} \right)^2 \text{tr}(R_{t_n}^* C_j^* C_j R_{t_n}) \right) = 0$$

for all  $i \in I_b$  and  $j \in I_p$ .

Now and for the rest of this paragraph, fix a realization (i.e. an element of  $\Omega$ ) such that  $(M_{t_n})_n$  converges to  $M_\infty$ . The polar decomposition of  $R_t$  is  $R_t = U_t \sqrt{M_t}$ . Since the set of  $k \times k$  unitary matrices is compact, there exists a subsequence  $(s_n)_n$  of  $(t_n)_n$  such that  $(U_{s_n})_n$  converges to  $U_\infty$ . We therefore have

$$\sqrt{M_\infty} U_\infty^* (L_i + L_i^*) U_\infty \sqrt{M_\infty} - M_\infty \text{tr}(M_\infty U_\infty^* (L_i + L_i^*) U_\infty) = 0$$

and

$$\sqrt{M_\infty} U_\infty^* C_j^* C_j U_\infty \sqrt{M_\infty} - M_\infty \text{tr}(M_\infty U_\infty^* C_j^* C_j U_\infty) = 0,$$

for all  $i \in I_b$  and  $j \in I_p$ . Denoting  $P_\infty$  the orthogonal projector onto the range of  $M_\infty$ , it follows that there exist real numbers  $(\alpha_i)_{i \in I_b}$  and  $(\beta_j)_{j \in I_p}$  such that

$$U_\infty P_\infty U_\infty^* (L_i + L_i^*) U_\infty P_\infty U_\infty^* = \alpha_i U_\infty P_\infty U_\infty^*$$

and

$$U_\infty P_\infty U_\infty^* C_j^* C_j U_\infty P_\infty U_\infty^* = \beta_j U_\infty P_\infty U_\infty^*.$$

Assumption **(Pur)** implies that the orthogonal projector  $U_\infty P_\infty U_\infty^*$  has rank one, thus so does  $P_\infty$ . Since  $\text{tr}(M_\infty) = 1$ ,  $M_\infty$  is a rank one orthogonal projector.

Since  $(M_{t_n})_n$  converges  $\mathbb{P}^{\text{ch}}$ -almost surely, the above paragraph and the absolute continuity of  $\mathbb{P}^\rho$  with respect to  $\mathbb{P}^{\text{ch}}$  show that  $M_\infty$  is  $\mathbb{P}^\rho$ -almost surely a rank one orthogonal projector.  $\square$

### 3. INVARIANT MEASURE AND EXPONENTIAL CONVERGENCE IN WASSERSTEIN DISTANCE

This section is devoted to the main result of the paper, which concerns the exponential convergence to the invariant measure for the Markov process  $(\hat{x}_t)_t$ . We first show a convergence result for  $\mathcal{F}$ -measurable random variables. The following theorem is a transcription of [38, Proposition 7.5].

**Theorem 3.1.** *Assume that **(L-erg)** holds. Then there exist two constants  $C > 0$  and  $\lambda > 0$  such that for any  $\rho \in \mathcal{D}_k$  and any  $t \geq 0$ ,*

$$\|e^{t\mathcal{L}}(\rho) - \rho_{\text{inv}}\| \leq C e^{-\lambda t}$$

Our next proposition requires the introduction of a shift semigroup. From now on we assume that  $(\Omega, (\mathcal{F}_t)_t, \mathbb{P})$  is a canonical realization of the processes  $W_i$  and  $N_j$ , in particular  $\Omega$  is a subset of  $(\mathbb{R}^{I_b \cup I_p})^{\mathbb{R}_+}$ . We can then define for every  $t \geq 0$  the map  $\theta^t$  on  $\Omega$  by

$$(\theta^t \omega)(s) = \omega(s+t) - \omega(t).$$

From the previous theorem we deduce the following proposition for  $\mathcal{F}$ -measurable random variables.

**Proposition 3.2.** *Assume **( $\mathcal{L}$ -erg)** holds. Then there exist two constants  $C > 0$  and  $\lambda > 0$  such that for any  $\mathcal{F}$ -measurable, essentially bounded function  $f : \Omega \mapsto \mathbb{C}$  with essential bound  $\|f\|_\infty$ , any  $t \geq 0$  and any  $\rho \in \mathcal{D}_k$ ,*

$$|\mathbb{E}^\rho(f \circ \theta^t) - \mathbb{E}^{\rho_{\text{inv}}}(f)| \leq \|f\|_\infty C e^{-\lambda t}. \quad (3.1)$$

*Proof.* Recall that by definition  $\mathbb{P}$  is the law of processes with independent increments. It follows that if  $g$  is  $\mathcal{F}_t$ -measurable and  $h$  is  $\mathcal{F}$ -measurable,  $\mathbb{E}(h \circ \theta^t g) = \mathbb{E}(h \circ \theta^t) \mathbb{E}(g)$ . Then, by definition of  $\mathbb{P}^\rho$ ,

$$\mathbb{E}^\rho(f \circ \theta^t) = \mathbb{E}(f \circ \theta^t Z_{t+s}^\rho).$$

Since  $Z_{t+s}^\rho = \text{tr}(S_{t+s}^{t*} S_{t+s}^t S_t \rho S_t^*)$  where  $S_t^* \rho S_t$  is  $\mathcal{F}_t$ -measurable and  $S_{t+s}^{t*} S_{t+s}^t = S_s^* S_s \circ \theta^t$  by (2.2)

$$\mathbb{E}^\rho(f \circ \theta^t) = \mathbb{E}\left(f \circ \theta^t \text{tr}\left(S_{t+s}^{t*} S_{t+s}^t \mathbb{E}(S_t \rho S_t^*)\right)\right).$$

Then relation (2.5), the  $\theta$ -invariance of  $\mathbb{P}$ , and the definition of the measures  $\mathbb{P}^\rho$  yield

$$\mathbb{E}^\rho(f \circ \theta^t) = \mathbb{E}^{\bar{\rho}_t}(f)$$

with  $\bar{\rho}_t = e^{t\mathcal{L}}(\rho)$ . It follows from Theorem 2.4 that

$$\mathbb{E}^\rho(f \circ \theta^t) - \mathbb{E}^{\rho_{\text{inv}}}(f) = \mathbb{E}^{\text{ch}}\left(f \text{tr}\left(M_\infty(e^{t\mathcal{L}}(\rho) - \rho_{\text{inv}})\right)\right).$$

For any matrix  $A$ , denoting  $\|A\|_1$  its trace norm,  $|\text{tr}(M_\infty A)| \leq \|A\|_1$ . Therefore,

$$|\mathbb{E}^\rho(f \circ \theta^t) - \mathbb{E}^{\rho_{\text{inv}}}(f)| \leq \|f\|_\infty \|e^{t\mathcal{L}}(\rho) - \rho_{\text{inv}}\|_1.$$

Theorem 3.1 then yields the proposition.  $\square$

The main strategy to show Theorem 1.1 is to construct a  $\mathcal{F}$ -measurable process  $(\hat{y}_t)_t$  approximating the process  $(\hat{x}_t)_t$ . Let  $(\hat{z}_t)_t$  be the maximum likelihood process:

$$\hat{z}_t = \underset{\hat{x} \in \text{PC}^k}{\text{argmax}} \|S_t x\| \quad (3.2)$$

where  $x$  is a norm one representative of  $\hat{x}$ . If the largest eigenvalue of  $S_t^* S_t$  is not simple, the choice of  $\hat{z}_t$  may not be unique. However we can always choose an appropriate  $\hat{z}_t$  in an  $(\mathcal{F}_t)_t$ -adapted way. If **(Pur)** holds, Proposition 2.5 ensures that the definition of  $\hat{z}_t$  is almost surely unambiguous for large enough  $t$ : it is the equivalence class of eigenvectors of  $M_t$  corresponding to its largest eigenvalue.

Let now  $(\hat{y}_t)_t$  be the evolution of this maximum likelihood estimate:

$$\hat{y}_t = S_t \cdot \hat{z}_t. \quad (3.3)$$

We shall also use the notation  $\hat{z}_{t+s}^s := \hat{z}_t \circ \theta^s$  and  $\hat{y}_{t+s}^s := \hat{y}_t \circ \theta^s$ , that is, processes defined in the same fashion but substituting  $S_{t+s}^s$  for  $S_t$ . It is worth noticing that these processes are all  $\mathcal{F}$ -measurable.

Our proof that  $(\hat{y}_t)_t$  is an exponentially good approximation of  $(\hat{x}_t)_t$  relies in part on the use of the exterior product of  $\mathbb{C}^k$ . We recall briefly the relevant definitions: for  $x_1, x_2 \in \mathbb{C}^k$  we denote by  $x_1 \wedge x_2$  the alternating bilinear form

$$x_1 \wedge x_2 : (y_1, y_2) \mapsto \det \begin{pmatrix} \langle x_1, y_1 \rangle & \langle x_1, y_2 \rangle \\ \langle x_2, y_1 \rangle & \langle x_2, y_2 \rangle \end{pmatrix}.$$

Then, the set of all  $x_1 \wedge x_2$  is a generating family for the set  $\wedge^2 \mathbb{C}^k$  of alternating bilinear forms on  $\mathbb{C}^k$ . We equip it with a complex inner product by

$$\langle x_1 \wedge x_2, y_1 \wedge y_2 \rangle = \det \begin{pmatrix} \langle x_1, y_1 \rangle & \langle x_1, y_2 \rangle \\ \langle x_2, y_1 \rangle & \langle x_2, y_2 \rangle \end{pmatrix},$$

and denote by  $\|x_1 \wedge x_2\|$  the associated norm (there should be no confusion with the norm on vectors). It is immediate to verify that our metric  $d(\cdot, \cdot)$  on  $\text{PC}^k$  satisfies

$$d(\hat{x}, \hat{y}) = \frac{\|x \wedge y\|}{\|x\| \|y\|}. \quad (3.4)$$

For  $A \in M_k(\mathbb{C})$ , we write  $\wedge^2 A$  for the operator on  $\wedge^2 \mathbb{C}^k$  defined by

$$(\wedge^2 A)(x_1 \wedge x_2) = Ax_1 \wedge Ax_2. \quad (3.5)$$

It follows that  $\wedge^2(AB) = \wedge^2 A \wedge^2 B$ , so that  $\|\wedge^2(AB)\| \leq \|\wedge^2 A\| \|\wedge^2 B\|$ . There exists a useful relationship between the operator norm on  $\wedge^2 M_k(\mathbb{C})$  and singular values of matrices. From e.g. Chapter XVI of [32],

$$\|\wedge^2 A\| = a_1(A) a_2(A), \quad (3.6)$$

where  $a_1(A) \geq a_2(A)$  are the two first singular values of  $A$ , i.e. the square roots of eigenvalues of  $A^*A$ . We recall that the operator norm is defined such that  $\|A\| := a_1(A)$ .

The exponential decrease of  $d(\hat{x}_t, \hat{y}_t)$  is derived from the exponential decay of the following function:

$$f : t \mapsto \mathbb{E}(\|\wedge^2 S_t\|).$$

**Lemma 3.3.** *Assume that **(Pur)** holds. Then there exist two constants  $C > 0$  and  $\lambda > 0$  such that for all  $t \geq 0$*

$$f(t) \leq Ce^{-\lambda t}$$

*Proof.* First, we show that  $f$  converges to zero as  $t$  grows to  $\infty$ . To this end recall that  $R_t = S_t / \sqrt{k Z_t^{\text{ch}}}$ , so that

$$\mathbb{E}(\|\wedge^2 S_t\|) = \mathbb{E}^{\text{ch}}(k \|\wedge^2 R_t\|).$$

Furthermore, since  $R_t^* R_t = M_t$ , we have from Theorem 2.4 and Proposition 2.5 that

$$\lim_{t \rightarrow \infty} \|\wedge^2 R_t\| = \lim_{t \rightarrow \infty} a_1(R_t) a_2(R_t) = 0.$$

Indeed, since  $a_1(R_t)$  and  $a_2(R_t)$  are the largest two eigenvalues of  $\sqrt{M_t}$ , the fact that it converges to a rank-one projector implies that  $a_1(R_t)$  converges to 1 and  $a_2(R_t)$  to zero. The inequality  $\|S_t\|^2 \leq \text{tr}(S_t^* S_t)$  implies  $\|\wedge^2 R_t\| \leq 1$  almost surely. Then Lebesgue's dominated convergence theorem yields  $\lim_{t \rightarrow \infty} f(t) = 0$ .

Second, we show  $f$  is submultiplicative. By the semi-group property,  $S_{t+s} = S_{t+s}^s S_s$  for all  $t, s \geq 0$ . Using that the norm is submultiplicative, for any  $t, s \geq 0$ ,

$$\|\wedge^2 S_{t+s}\| \leq \|\wedge^2 S_{t+s}^s\| \|\wedge^2 S_s\|$$

Since  $\mathbb{P}$  has independent increments,  $\|\wedge^2 S_{t+s}^s\| = \|\wedge^2 S_t\| \circ \theta^s$  is  $\mathcal{F}_s$ -independent and  $\|\wedge^2 S_s\|$  is  $\mathcal{F}_s$ -measurable,

$$\mathbb{E}(\|\wedge^2 S_{t+s}\|) \leq \mathbb{E}(\|\wedge^2 S_{t+s}^s\|) \mathbb{E}(\|\wedge^2 S_s\|)$$

The measure  $\mathbb{P}$  being shift-invariant,

$$\mathbb{E}(\|\wedge^2 S_{t+s}\|) \leq \mathbb{E}(\|\wedge^2 S_t\|) \mathbb{E}(\|\wedge^2 S_s\|)$$

which yields that  $f$  is submultiplicative.

Since  $f$  is measurable, submultiplicative and  $0 \leq f(t) \leq k$  for all  $t$ , Fekete's subadditive lemma ensures that there exists  $\lambda \in (-\infty, \infty]$  such that

$$\lim_{t \rightarrow \infty} \frac{1}{t} \log f(t) = \inf \frac{1}{t} \log f(t) = -\lambda.$$

Since  $f$  converges towards 0, this  $\lambda$  belongs to  $(0, \infty]$ . This yields the lemma.  $\square$

**Proposition 3.4.** *Assume that **(Pur)** holds. Then there exist two constants  $C > 0$  and  $\lambda > 0$  such that for any  $s, t \in \mathbb{R}_+$  and for any probability measure  $\mu$  on  $(\mathbb{P}\mathbb{C}^k, \mathcal{B})$ ,*

$$\mathbb{E}_\mu(d(\hat{x}_{t+s}, \hat{y}_{t+s}^s)) \leq Ce^{-\lambda t}. \quad (3.7)$$

*Proof.* Recall that  $\mathbb{E}_\mu$  is the expectation with respect to  $\mathbb{Q}_\mu$ . Using the Markov property, we have

$$\mathbb{E}_\mu(d(\hat{x}_{t+s}, \hat{y}_{t+s}^s)) = \mathbb{E}_{\mu_s}(d(\hat{x}_t, \hat{y}_t)) \quad (3.8)$$

with  $\mu_s$  the distribution of  $\hat{x}_s$  conditioned on  $\hat{x}_0 \sim \mu$ . Then it is sufficient to prove the proposition for  $s = 0$ . For any  $t \geq 0$ , using the fact that  $\|S_t z_t\| = \|S_t\|$  for  $z_t$  a norm one representative of  $\hat{z}_t$ ,

$$d(\hat{x}_t, \hat{y}_t) = \frac{\|x_t \wedge y_t\|}{\|x_t\| \|y_t\|} = \frac{\|S_t x_0 \wedge S_t z_t\|}{\|S_t x_0\| \|S_t z_t\|} \leq \frac{\|\wedge^2 S_t\|}{\|S_t x_0\| \|S_t\|} \leq \frac{\|\wedge^2 S_t\|}{\|S_t x_0\|^2}$$

Using this inequality and the fact that  $d\mathbb{Q}_\mu|_{\mathcal{F}_t \otimes \mathcal{B}} = \|S_t x_0\|^2 d\mathbb{P} d\mu(\hat{x}_0)$ ,

$$\mathbb{E}_\mu(d(\hat{x}_t, \hat{y}_t)) \leq \int \mathbb{E}(\|\wedge^2 S_t\|) d\mu(\hat{x}_0) \leq f(t).$$

Finally Lemma 3.3 yields the proposition.  $\square$

We turn to the proof of our main theorem, Theorem 1.1. The speed of convergence is expressed in terms of the Wasserstein distance  $W_1$ . Let us recall the definition of this distance for compact metric spaces: for  $X$  a compact metric space equipped with its Borel  $\sigma$ -algebra, the Wasserstein distance of order 1 between two probability measures  $\sigma$  and  $\tau$  on  $X$  can be defined using the Kantorovich–Rubinstein duality Theorem as

$$W_1(\sigma, \tau) = \sup_{f \in \text{Lip}_1(X)} \left| \int_X f d\sigma - \int_X f d\tau \right|,$$

where  $\text{Lip}_1(X) = \{f : X \rightarrow \mathbb{R} \text{ s.t. } |f(x) - f(y)| \leq d(x, y)\}$  is the set of Lipschitz continuous functions with constant one, and  $d$  is the metric on  $X$ . Here we use this for  $X = \mathbb{P}\mathbb{C}^k$  and  $d$  defined in (2.1) (see also (3.4)).

We recall our main theorem before proving it.

**Theorem 3.5.** *Assume that **(Pur)** and **(L-erg)** hold. Then the Markov process  $(\hat{x}_t)_t$  has a unique invariant probability measure  $\mu_{\text{inv}}$ , and there exist  $C > 0$  and  $\lambda > 0$  such that for any initial distribution  $\mu$  of  $\hat{x}_0$  over  $\mathbb{P}\mathbb{C}^k$ , for all  $t \geq 0$ , the distribution  $\mu_t$  of  $\hat{x}_t$  satisfies*

$$W_1(\mu_t, \mu_{\text{inv}}) \leq C e^{-\lambda t}$$

where  $W_1$  is the Wasserstein distance of order 1.

*Proof.* Let  $f \in \text{Lip}_1(\mathbb{P}\mathbb{C}^k)$ . From the definition of Wasserstein distance, we can restrict ourselves to functions  $f$  that vanish at some point. Remark that since  $\sup_{\hat{x}, \hat{y} \in \mathbb{P}\mathbb{C}^k} d(\hat{x}, \hat{y}) = 1$ , restricting to this set of functions implies  $\|f\|_\infty \leq 1$ . Let  $\mu_{\text{inv}}$  be an invariant probability measure for  $(\hat{x}_t)_t$ . We will prove the exponential convergence of  $(\mu_t)_t$  towards  $\mu_{\text{inv}}$  for any initial  $\mu_0$ , and that will imply that  $(\hat{x}_t)_t$  accepts a unique invariant probability measure. Let  $t \geq 0$ , and recall that  $\hat{y}_t^{t/2} = \hat{y}_{t/2} \circ \theta^{t/2}$ . We have

$$\begin{aligned} \mathbb{E}_\mu(f(\hat{x}_t)) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{x}_t)) &= \mathbb{E}_\mu(f(\hat{x}_t)) - \mathbb{E}_\mu(f(\hat{y}_t^{t/2})) + \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{y}_t^{t/2})) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{x}_t)) \\ &\quad + \mathbb{E}_\mu(f(\hat{y}_t^{t/2})) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{y}_t^{t/2})) \\ &\leq \mathbb{E}_\mu(d(\hat{x}_t, \hat{y}_t^{t/2})) + \mathbb{E}_{\mu_{\text{inv}}}(d(\hat{x}_t, \hat{y}_t^{t/2})) \\ &\quad + \mathbb{E}_\mu(f(\hat{y}_t^{t/2})) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{y}_t^{t/2})). \end{aligned} \tag{3.9}$$

The two terms on the right hand side of line (3.9) are bounded using Proposition 3.4. Using Proposition 2.3, the difference on line (3.10) satisfies

$$\mathbb{E}_\mu(f(\hat{y}_t^{t/2})) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{y}_t^{t/2})) = \mathbb{E}^{\rho\mu}(f(\hat{y}_t^{t/2})) - \mathbb{E}^{\rho\mu_{\text{inv}}}(f(\hat{y}_t^{t/2})).$$

Then bounding the right hand side using Proposition 3.2, it follows there exist  $C > 0$  and  $\lambda > 0$  such that

$$\left| \mathbb{E}_\mu(f(\hat{x}_t)) - \mathbb{E}_{\mu_{\text{inv}}}(f(\hat{x}_t)) \right| \leq 3C e^{-\lambda t/2}.$$

Adapting the two constants yields the theorem.  $\square$

#### 4. SET OF INVARIANT MEASURES UNDER **(Pur)**

The results and proofs of this section are a direct translation of [11, Appendix B]. We reproduce the proofs for the reader's convenience.

Whenever **(L-erg)** does not hold,  $\dim \ker \mathcal{L} > 1$  and the semigroup  $(e^{t\mathcal{L}})_t$  accepts more than one fixed point in  $\mathcal{D}_k$ . The convex set of invariant states can be explicitly classified given the matrices  $(L_i)_{i \in I_b}$  and  $(C_j)_{j \in I_b}$ . Following [9, Theorem 7] (alternatively see Theorem 7.2 and Proposition 7.6 in [38], and [36]), there exists a decomposition

$$\mathbb{C}^k \simeq \mathbb{C}^{n_1} \oplus \dots \oplus \mathbb{C}^{n_d} \oplus \mathbb{C}^D, \quad k = n_1 + \dots + n_d + D$$

with the following properties:

- (1) The range of any invariant states is a subspace of  $V = \mathbb{C}^{n_1} \oplus \dots \oplus \mathbb{C}^{n_d} \oplus \{0\}$ ;
- (2) The restriction of the operators  $L_i$  and  $C_j$  to  $\mathbb{C}^{n_1} \oplus \dots \oplus \mathbb{C}^{n_d}$  are block-diagonal, with

$$\begin{aligned} L_i &= L_{1,i} \oplus \dots \oplus L_{d,i}, \quad i \in I_b, \\ C_j &= C_{1,j} \oplus \dots \oplus C_{d,j}, \quad j \in I_p; \end{aligned} \tag{4.1}$$

- (3) For each  $\ell = 1, \dots, d$  there is a decomposition  $\mathbb{C}^{n_\ell} = \mathbb{C}^{k_\ell} \otimes \mathbb{C}^{m_\ell}$ ,  $n_\ell = k_\ell \times m_\ell$ , a unitary matrix  $U_\ell$  on  $\mathbb{C}^{n_\ell}$  and matrices  $\{\hat{L}_{\ell,i}\}_{i \in I_b}$  and  $\{\hat{C}_{\ell,j}\}_{j \in I_p}$  on  $\mathbb{C}^{k_\ell}$  such that

$$\begin{aligned} L_{\ell,i} &= U_\ell(\hat{L}_{\ell,i} \otimes \text{Id}_{\mathbb{C}^{m_\ell}})U_\ell^*, \quad i \in I_b, \\ C_{\ell,j} &= U_\ell(\hat{C}_{\ell,j} \otimes \text{Id}_{\mathbb{C}^{m_\ell}})U_\ell^*, \quad j \in I_p; \end{aligned} \tag{4.2}$$

- (4) There exists a positive definite matrix  $\rho_\ell$  on  $\mathbb{C}^{k_\ell}$  such that

$$0 \oplus \dots \oplus U_\ell(\rho_\ell \otimes \text{Id}_{\mathbb{C}^{m_\ell}})U_\ell^* \oplus \dots \oplus 0 \tag{4.3}$$

is a fixed point of  $(e^{t\mathcal{L}})_t$ .

Then, the set of fixed points for  $(e^{t\mathcal{L}})$  is

$$U_1(\rho_1 \otimes M_{m_1}(\mathbb{C}))U_1^* \oplus \dots \oplus U_d(\rho_d \otimes M_{m_d}(\mathbb{C}))U_d^* \oplus 0_{M_D(\mathbb{C})}.$$

The decomposition simplifies under the purification assumption.

**Proposition 4.1.** *Assume that **(Pur)** holds. Then there exists a set  $\{\rho_\ell\}_{\ell=1}^d$  of positive definite matrices and an integer  $D$  such that the set of fixed points of  $(e^{t\mathcal{L}})_t$  is*

$$\mathbb{C}\rho_1 \oplus \dots \oplus \mathbb{C}\rho_d \oplus 0_{M_D(\mathbb{C})}.$$

*Proof.* The statement follows from the discussion preceding the proposition if we show that **(Pur)** implies  $m_1 = \dots = m_d = 1$ . Assume that one of the  $m_\ell$ , e.g.  $m_1$ , is greater than 1. Let  $x$  be a norm one vector in  $\mathbb{C}^{k_1}$ . Then  $\pi = U_1(\pi_{\hat{x}} \otimes I_{\mathbb{C}^{m_1}})U_1^* \oplus 0 \oplus \dots \oplus 0$  is an orthogonal projection of rank  $m_1 > 1$ , and

$$\pi(L_i + L_i^*)\pi = \|(L_{1,i} + L_{1,i}^*)x\|^2 \pi \quad \text{and} \quad \pi(C_j^* C_j)\pi = \|C_{1,j}^* C_{1,j}x\|^2 \pi \quad \text{for all } i \in I_b, j \in I_p,$$

and this contradicts **(Pur)**.  $\square$

It is clear from Proposition 4.1 that to each extremal fixed point  $0 \oplus \dots \oplus \rho_\ell \oplus \dots \oplus 0$  corresponds a unique invariant measure  $\mu_\ell$  supported on its range  $\text{ran } \rho_\ell$ . The converse is the subject of the next proposition.

**Proposition 4.2.** *Assume **(Pur)** holds. Then any invariant probability measure of  $(\hat{x}_t)_t$  is a convex combination of the measures  $\mu_\ell$ ,  $\ell = 1, \dots, d$ .*

*Proof.* Let  $\mu$  be an invariant probability measure for  $(\hat{x}_t)_t$  and  $f$  be a continuous function on  $P\mathbb{C}^k$ . Proposition 3.4 implies that

$$\int f \, d\mu = \mathbb{E}_\mu(f(\hat{x}_0)) = \mathbb{E}_\mu(f(\hat{x}_t)) = \lim_{t \rightarrow \infty} \mathbb{E}_\mu(f(\hat{y}_t)).$$

Since  $(\hat{y}_t)_t$  is  $\mathcal{F}$ -measurable, Proposition 2.3 implies  $\int f \, d\mu = \lim_{t \rightarrow \infty} \mathbb{E}^{\rho_\mu}(f(\hat{y}_t))$ , and by Remark 3,  $\rho_\mu \in \mathcal{D}_k$  is a fixed point of  $(e^{t\mathcal{L}})_t$ . Proposition 4.1 ensures that there exist nonnegative numbers  $p_1, \dots, p_d$  summing up to one such that  $\rho_\mu = p_1 \rho_1 \oplus \dots \oplus p_d \rho_d \oplus 0_{M_D(\mathbb{C})}$ . From the definition of  $\mathbb{P}^{\rho_\mu}$ ,

$$\mathbb{P}^{\rho_\mu} = p_1 \mathbb{P}^{\rho_1} + \dots + p_d \mathbb{P}^{\rho_d}$$

with the abuse of notation  $\rho_\ell := 0 \oplus \dots \oplus \rho_\ell \oplus \dots \oplus 0$ , so that

$$\int f \, d\mu = \lim_{t \rightarrow \infty} p_1 \mathbb{E}^{\rho_1}(f(\hat{y}_t)) + \dots + p_d \mathbb{E}^{\rho_d}(f(\hat{y}_t)).$$

The same argument gives  $\int f \, d\mu_\ell = \lim_{t \rightarrow \infty} \mathbb{E}^{\rho_\ell}(f(\hat{y}_t))$ , and we have  $\mu = p_1 \mu_1 + \dots + p_d \mu_d$ .  $\square$

5. **(Pur)** IS NOT NECESSARY FOR PURIFICATION

As shown by the following example, the condition **(Pur)** is sufficient but not necessary for (1.5) to hold.

Let  $k = 3$  and fix an orthonormal basis  $\{e_1, e_2, e_3\}$  of  $\mathbb{C}^3$ . Let  $I_b = \{0, 1\}$ ,  $I_p = \{2\}$ ,  $u = (e_1 + e_2 + e_3)/\sqrt{3}$  and  $v = (e_1 + e_3)/\sqrt{2}$ . Let

$$H = 0, \quad V_0 = L_0 = e_1 u^*, \quad V_1 = L_1 = 2vv^* + e_2 e_2^*, \quad V_2 = C_2 = ue_1^*. \quad (5.1)$$

**Proposition 5.1.** *Let  $\mathcal{L}$  be the Lindblad operator given by (1.2) with  $H, V_1, V_2, V_3$  defined in (5.1). Then **( $\mathcal{L}$ -erg)** holds and the unique invariant state  $\rho_{\text{inv}}$  is positive definite.*

*Proof.* Using [38, Proposition 7.6], it is sufficient to prove that if  $\pi$  is a non null orthogonal projector such that  $(\text{Id} - \pi)L_0\pi = (\text{Id} - \pi)L_1\pi = (\text{Id} - \pi)L_2\pi = 0$ , then  $\pi = \text{Id}$ . Assume  $\text{rank } \pi < 3$ . Since  $\pi \in M_3(\mathbb{C})$ , there exist  $\hat{x} \in \text{PC}^3$  such that either  $\pi = \pi_{\hat{x}}$  or  $\pi = \text{Id} - \pi_{\hat{x}}$ . If the first alternative holds,  $(\text{Id} - \pi)L_0\pi = (\text{Id} - \pi)L_1\pi = (\text{Id} - \pi)L_2\pi = 0$  implies  $\hat{x}$  is the equivalence class of a common eigenvector of  $L_0, L_1$  and  $L_2$ . If the second alternative holds,  $\hat{x}$  is the equivalence class of a common eigenvector of  $L_0^*, L_1^*$  and  $L_2^*$ . The only common eigenvectors of  $L_0$  and  $L_2$  or  $L_0^*$  and  $L_2^*$  are elements of  $\mathbb{C}(e_2 - e_3)$ . Since  $L_1$  is self adjoint, and this eigenspace is not an eigenspace of  $L_1$ , the proposition holds.  $\square$

In the orthonormal basis  $\{e_1, e_2, e_3\}$ ,

$$L_0^* + L_0 = \frac{1}{\sqrt{3}} \begin{pmatrix} 2 & 1 & 1 \\ 1 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \quad L_1^* + L_1 = 2 \begin{pmatrix} 1 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 1 \end{pmatrix} \quad \text{and} \quad C_2^* C_2 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Taking  $\pi$  the orthogonal projector onto the subspace spanned by  $\{e_2, e_3\}$  it follows that **(Pur)** does not hold. Yet we have the following proposition.

**Proposition 5.2.** *Consider the family of processes  $(\rho_t)_t$  defined by (1.3) with  $H, L_0, L_1, C_2$  defined in (5.1). Then for any  $\rho \in \mathcal{D}_k$ ,*

$$\lim_{t \rightarrow \infty} \inf_{\hat{y} \in \text{PC}^k} \|\rho_t - \pi_{\hat{y}}\| = 0 \quad \mathbb{P}^\rho\text{-almost surely.}$$

*Proof.* Proposition 5.1 implies that  $\rho_{\text{inv}}$ , the unique element of  $\mathcal{D}_k$  invariant by  $(e^{t\mathcal{L}})_t$  is positive definite. Then,  $\text{tr}(C_2^* C_2 \rho_{\text{inv}}) > 0$ . The results of [29] thus ensure that for any  $\rho \in \mathcal{D}_k$ ,

$$\lim_{t \rightarrow \infty} N_2(t)/t = \text{tr}(C_2^* C_2 \rho_{\text{inv}}), \quad \mathbb{P}^\rho\text{-almost surely.}$$

Let  $T = \inf\{t \geq 0 : N_2(t) \geq 1\}$ . Then  $\mathbb{P}^\rho(T < \infty) = 1$  and from the definition of  $C_2$ ,

$$\rho_T = \pi_{\hat{u}} \quad \text{and} \quad \rho_t = \pi_{S_T \cdot \hat{u}} \quad \text{for any } t \geq T.$$

Hence  $\inf_{\hat{y} \in \text{PC}^k} \|\rho_t - \pi_{\hat{y}}\| = 0$  for any  $t \geq T$  and  $\mathbb{P}^\rho(T < \infty) = 1$  yield the proposition.  $\square$

**Corollary 5.3.** *Consider the process  $(\hat{x}_t)_t$  defined by (1.4) with  $H, L_0, L_1, C_2$  defined in (5.1). Then  $(\hat{x}_t)_t$  accepts a unique invariant probability measure  $\mu_{\text{inv}}$  and there exist  $C > 0$  and  $\lambda > 0$  such that for any initial distribution  $\mu$  of  $\hat{x}_0$  over  $\text{PC}^3$ , for all  $t \geq 0$ , the distribution  $\mu_t$  of  $\hat{x}_t$  satisfies*

$$W_1(\mu_t, \mu_{\text{inv}}) \leq C e^{-\lambda t}.$$

*Proof.* It is a direct adaptation of our proof of Theorem 1.1. Indeed Theorem 1.1 holds if one substitutes the conclusion of Proposition 2.5 for **(Pur)**. Taking  $\rho_0 = \text{Id}/3$  in the latter proposition yields  $\rho_t = \frac{S_t S_t^*}{\text{tr}(S_t S_t^*)}$  and  $M_t = \frac{S_t^* S_t}{\text{tr}(S_t^* S_t)}$ . Therefore,  $\rho_t$  and  $M_t$  are unitarily equivalent. Following the arguments and notation of the proof of Proposition 5.2 we see that  $\mathbb{P}^{\text{ch}}$ -almost surely,  $M_T$  has rank one, and so does any  $M_t$  for  $t \geq T$ . Hence the conclusion of Proposition 2.5 holds and the corollary is proven.  $\square$

Following the proofs of the discrete-time results of [11], we can prove that the implication in Proposition 2.5 is an equivalence if **(Pur)** is replaced by

**(NSC-Pur):** Any non zero orthogonal projector  $\pi$  that satisfies  $\pi S_t^* S_t \pi \propto \pi$   $\mathbb{P}$ -almost surely for any  $t \geq 0$  has rank one.

Alas, in practice, such a condition is hard to check.

## 6. EXAMPLES

In the following examples  $k = 2$ . We recall the definition of the Pauli matrices:

$$\sigma_x := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y := \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \text{and} \quad \sigma_z := \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

A standard orthonormal basis of  $M_2(\mathbb{C})$  equipped with the Hilbert-Schmidt inner product is

$$\left( \frac{1}{\sqrt{2}} \text{Id}, \frac{1}{\sqrt{2}} \sigma_x, \frac{1}{\sqrt{2}} \sigma_y, \frac{1}{\sqrt{2}} \sigma_z \right).$$

In the basis of Pauli matrices one can write in a unique way any projection  $\pi_{\hat{x}}$  as

$$\pi_{\hat{x}} = \frac{1}{2} (\text{Id} + \mathcal{X} \sigma_x + \mathcal{Y} \sigma_y + \mathcal{Z} \sigma_z)$$

where

$$\mathcal{X} = \text{tr}(\pi_{\hat{x}} \sigma_x), \quad \mathcal{Y} = \text{tr}(\pi_{\hat{x}} \sigma_y), \quad \mathcal{Z} = \text{tr}(\pi_{\hat{x}} \sigma_z).$$

We denote in particular by  $\mathcal{X}_t, \mathcal{Y}_t, \mathcal{Z}_t$  respectively the coordinates associated with  $\pi_{\hat{x}_t}$ .

**6.1. Unitarily perturbed non demolition diffusive measurement.** Our first example consists of a  $\frac{1}{2}$ -spin (or qbit) in a magnetic field oriented along the  $y$ -axis and subject to indirect non demolition measurement along the  $z$ -axis. It is a typical quantum optics experimental situation (see for example [20]). In terms of the parameters defining the related quantum trajectories, we get  $H = \sigma_y$ ,  $I_b = \{0\}$ ,  $I_p = \emptyset$  and  $L_0 = \sqrt{\gamma} \sigma_z$  with  $\gamma > 0$ . Then  $(\pi_{\hat{x}_t})_t$  conditioned on  $\hat{x}_0$  is the solution of

$$d\pi_{\hat{x}_t} = \left( -i[\sigma_y, \pi_{\hat{x}_t}] + \gamma(\sigma_z \pi_{\hat{x}_t} \sigma_z - \pi_{\hat{x}_t}) \right) dt + \sqrt{\gamma} (\sigma_z \pi_{\hat{x}_t} + \pi_{\hat{x}_t} \sigma_z - 2 \text{tr}(\sigma_z \pi_{\hat{x}_t}) \pi_{\hat{x}_t}) dB_t. \quad (6.1)$$

For this quantum trajectory it is immediate to verify **(Pur)**, and solving  $\mathcal{L}(\rho) = 0$  shows that  $\rho_{\text{inv}} = \frac{1}{2} \text{Id}$  is the unique invariant state, so that **(L-erg)** holds. Hence, by Theorem 1.1  $(\hat{x}_t)_t$  has a unique invariant measure. In the following we derive an explicit expression for this invariant measure.

The next Lemma allows us to restrict the state space.

**Lemma 6.1.** *If  $\mu(\mathcal{Y}_0 = 0) = 1$  then  $\mathbb{Q}_\mu(\mathcal{Y}_t = 0) = 1$  for all  $t$  in  $\mathbb{R}$ .*

*Proof.* From equation (6.1),  $(\mathcal{Y}_t)_t$  is the solution of  $d\mathcal{Y}_t = 2\mathcal{Y}_t(\gamma dt - \sqrt{\gamma} \mathcal{Z}_t dB_t)$ . It is therefore a Doléans-Dade exponential:

$$\mathcal{Y}_t = \mathcal{Y}_0 e^{2\gamma t} \exp \left( -2\gamma \int_0^t \mathcal{Z}_s^2 ds - 2\sqrt{\gamma} \int_0^t \mathcal{Z}_s dB_s \right)$$

and the conclusion follows.  $\square$

Now we prove that the invariant measure admits a rotational symmetry.

**Lemma 6.2.** *Assume that the distribution  $\mu$  of  $\hat{x}_0$  is invariant with respect to the mapping  $\hat{x} \mapsto \sigma_y \cdot \hat{x}$ . Then  $\mu_t$  is invariant with respect to the same mapping.*

*Proof.* Since  $\sigma_y$  is unitary and self-adjoint, we have  $\pi_{\sigma_y \cdot \hat{x}} = \sigma_y \pi_{\hat{x}} \sigma_y$ . Since  $\sigma_y \sigma_z = -\sigma_z \sigma_y$ , it follows from (6.1) that

$$\begin{aligned} d(\sigma_y \pi_{\hat{x}_t} \sigma_y) &= \left( -i[\sigma_y, \sigma_y \pi_{\hat{x}_t} \sigma_y] + \gamma(\sigma_z \sigma_y \pi_{\hat{x}_t} \sigma_y \sigma_z - \sigma_y \pi_{\hat{x}_t} \sigma_y) \right) dt \\ &\quad - \sqrt{\gamma} (\sigma_z \sigma_y \pi_{\hat{x}_t} \sigma_y + \sigma_y \pi_{\hat{x}_t} \sigma_y \sigma_z - 2 \text{tr}(\sigma_z \sigma_y \pi_{\hat{x}_t} \sigma_y) \sigma_y \pi_{\hat{x}_t} \sigma_y) dB_t. \end{aligned}$$

Then it follows from  $\sigma_y \cdot \hat{x}_0 \sim \hat{x}_0$  and  $(B_t)_t \sim (-B_t)_t$  that  $(\hat{x}_t)_t$  and  $(\sigma_y \cdot \hat{x}_t)_t$  are both weak solutions to the same SDE with the same initial condition. Since this SDE has a unique solution, they have the same distributions.  $\square$

**Proposition 6.3.** *Let  $(\hat{x}_t)_t$  be the process defined by (6.1). Then its unique invariant measure is the normalized image measure by*

$$\iota : \theta \mapsto \frac{1}{2}(\text{Id} + \sin \theta \sigma_x + \cos \theta \sigma_z) \quad (6.2)$$

of the measure  $\tau(\theta) d\theta$  on  $(-\pi, \pi]$  with

$$\tau(\theta) = \int_{\theta}^{\pi} \exp \frac{\cot x - \cot \theta}{\gamma} \frac{\sin x}{\sin^3 \theta} dx$$

for  $\theta \in [0, \pi]$  and  $\tau(\theta) = \tau(\theta + \pi)$  for  $\theta \in (-\pi, 0]$ .

*Proof.* The convergence results in Theorem 1.1 and Lemma 6.1 imply that the invariant measure  $\mu_{\text{inv}}$  is the image by  $\iota$  of a probability measure  $\tau$  on  $(-\pi, \pi]$ . Let  $(\theta_t)_t$  be the solution of

$$d\theta_t = 2(1 - \gamma \cos \theta_t \sin \theta_t) dt - 2\sqrt{\gamma} \sin \theta_t dB_t \quad (6.3)$$

with initial condition  $\theta_0$ . Remark that  $(\theta_t)_t$  is  $2\pi$ -periodic with respect to its initial condition, namely,  $(\theta_t + 2\pi)_t$  is solution of (6.3) with initial condition  $\theta_0 + 2\pi$ . Now, using the Itô formula,

$$(\cos \theta_t, \sin \theta_t)_t \sim (\text{tr}(\pi_{\hat{x}_t} \sigma_z), \text{tr}(\pi_{\hat{x}_t} \sigma_x))_t$$

for  $(\pi_{\hat{x}_t})_t$  solution of (6.1) with initial condition  $\hat{x}_0 = \frac{1}{2}(\text{Id} + \sin \theta_0 \sigma_x + \cos \theta_0 \sigma_z)$ . Hence  $(\iota(\theta_t))_t$  has the same distribution as  $(\hat{x}_t)_t$ . Therefore  $\tau$  is an invariant measure for the diffusion defined by (6.3); in addition, Theorem 1.1 shows that this invariant measure is unique, and Lemma 6.2 shows that it is  $\pi$ -periodic. Following standard methods (see [28]), one shows that the restriction of  $\tau$  to  $[0, \pi)$  has a density of the form  $\tau(\theta) = C_1 \tau_1(\theta) + C_2 \tau_2(\theta)$  with  $C_1, C_2 \in \mathbb{R}$  and

$$\tau_1(\theta) = \frac{\int_{\theta}^{\pi} \sin x \exp(\frac{1}{\gamma} \cot x) dx}{\sin^3 \theta \exp(\frac{1}{\gamma} \cot \theta)} \quad \tau_2(\theta) = \frac{1}{\sin^3 \theta \exp(\frac{1}{\gamma} \cot \theta)}$$

Now, straightforward analysis shows that  $\int_0^{\pi} \tau_1(\theta) d\theta < \infty$  while  $\int_0^{\pi} \tau_2(\theta) d\theta = \infty$ . Therefore,  $\tau$  is proportional to  $\tau_1$  and the result follows.  $\square$

*Remark 4.* For  $\gamma \rightarrow \infty$ , the invariant measure  $\mu_{\text{inv}}$  in Proposition 6.3 is a Dirac measure at 0 and  $\pi$ . To describe the scaling for  $\gamma$  large we embed  $\tau(\theta)$  into  $L^1(\mathbb{R})$  by defining it to be zero outside the region  $(-\pi, \pi]$ . Then on the positive half line, in the  $L^1$  norm,

$$\lim_{\gamma \rightarrow \infty} \frac{1}{2\gamma^3} \tau\left(\frac{\theta}{\gamma}\right) = \frac{1}{\theta^3} \exp\left(-\frac{1}{\theta}\right).$$

Hence, for large  $\gamma$ , the stationary probability distribution has two peaks of width (of order)  $1/\gamma$  located  $1/\gamma$  radians clockwise from the limit points 0 and  $\pi$ . Furthermore the probability to find the particle around the limit points is exponentially suppressed.

The strong noise limit,  $\gamma \rightarrow \infty$ , was recently studied in various models [2, 8, 12]. This is the first model that allows for an explicit calculation of the shape of the stationary probability measure. The density of the invariant probability distribution is plotted in Figure 1 for three values of  $\gamma$ , and for  $\theta \in [0, \pi]$ .

**6.2. Thermal qubit, diffusive case.** The following second example corresponds to the evolution of a qubit interacting weakly with the electromagnetic field at a fixed temperature. The emission and absorption of photons by the qubit are stimulated by a resonant coherent field (laser). In the limit of a strong stimulating laser, the measurement of emitted photons results in a diffusive signal whose drift depends on the instantaneous average value of the raising and lowering operators of the qubit (see [37, §4.4] for a more detailed physical derivation). We obtain an analytically solvable model if we assume that the unitary rotation of the qubit is compensated for and thus frozen. In terms of the parameters defining the related quantum trajectories, we get  $H = 0$ ,  $I = I_b = \{0, 1\}$ ,  $L_0 = \sqrt{a} \sigma_+$  and  $L_1 = \sqrt{b} \sigma_-$  with  $a, b \in \mathbb{R}_+ \setminus \{0\}$  and  $\sigma_{\pm} = \frac{1}{2}(\sigma_x \pm i\sigma_y)$ , so that  $\sigma_+ = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$  and  $\sigma_- = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$ .



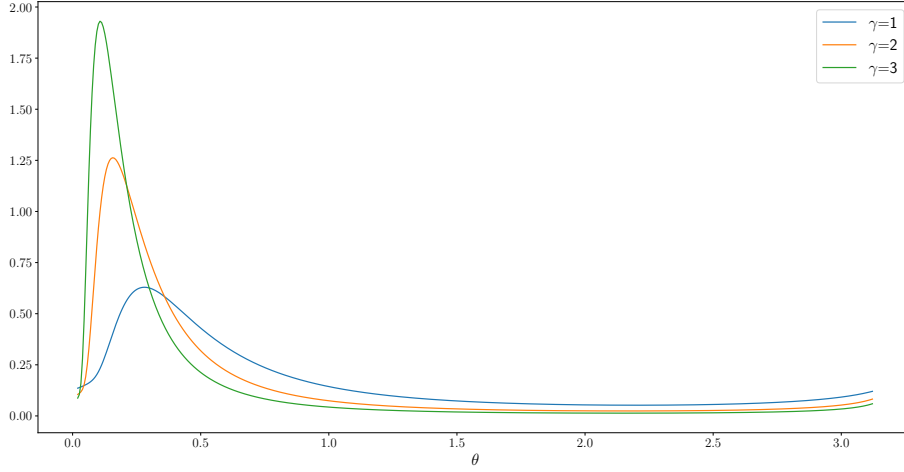


FIGURE 1. The restriction to  $[0, \pi]$  of the density of the invariant probability distribution in Example 6.1.

The stochastic master equation satisfied by  $\pi_{\hat{x}_t}$  is

$$\begin{aligned}
d\pi_{\hat{x}_t} = & a\left(\sigma_+\pi_{\hat{x}_t}\sigma_- - \frac{1}{4}((\text{Id} - \sigma_z)\pi_{\hat{x}_t} + \pi_{\hat{x}_t}(\text{Id} - \sigma_z))\right) dt \\
& + b\left(\sigma_-\pi_{\hat{x}_t}\sigma_+ - \frac{1}{4}((\text{Id} + \sigma_z)\pi_{\hat{x}_t} + \pi_{\hat{x}_t}(\text{Id} + \sigma_z))\right) dt \\
& + \sqrt{a}\left(\sigma_+\pi_{\hat{x}_t} + \pi_{\hat{x}_t}\sigma_- - \text{tr}(\sigma_x\pi_{\hat{x}_t})\pi_{\hat{x}_t}\right) dB_0(t) \\
& + \sqrt{b}\left(\sigma_-\pi_{\hat{x}_t} + \pi_{\hat{x}_t}\sigma_+ - \text{tr}(\sigma_x\pi_{\hat{x}_t})\pi_{\hat{x}_t}\right) dB_1(t)
\end{aligned} \tag{6.4}$$

Again it is immediate to verify **(Pur)**, and solving for  $\mathcal{L}(\rho) = 0$  shows that **(L-erg)** holds.

**Lemma 6.4.** *If  $\mu(\mathcal{Y}_0 = 0) = 1$  then  $\mathbb{Q}_\mu(\mathcal{Y}_t = 0) = 1$  for all  $t$  in  $\mathbb{R}$ .*

*Proof.* From (6.4),  $(\mathcal{X}_t)_t$  and  $(\mathcal{Y}_t)_t$  satisfy

$$d\mathcal{Y}_t = -\mathcal{Y}_t\left(\frac{1}{2}(a+b)dt + \mathcal{X}_t(\sqrt{a}dB_0(t) + \sqrt{b}dB_1(t))\right).$$

Therefore, if one defines

$$M_t = \exp\left(-\frac{1}{2}\int_0^t (a+b)\mathcal{X}_s^2 ds - \int_0^t \mathcal{X}_s(\sqrt{a}dB_0(s) + \sqrt{b}dB_1(s))\right)$$

then one has  $\mathcal{Y}_t = \mathcal{Y}_0 e^{-\frac{1}{2}(a+b)t} M_t$ , and this proves Lemma 6.4.  $\square$

**Proposition 6.5.** *Let  $(\hat{x}_t)_t$  be the process defined by (6.4). Then its unique invariant measure is the normalized image measure by  $\iota$  (defined by (6.2)) of the measure  $\tau(\theta) d\theta$  on  $(-\pi, \pi]$  with*

$$\tau(\theta) = \frac{e^{\varsigma z \arctan(\varsigma(\cos\theta - z))}}{(\cos^2\theta + 1 - 2z\cos\theta)^{3/2}},$$

with  $z = \frac{a-b}{a+b}$  and  $\varsigma = \frac{a+b}{2\sqrt{ab}}$ .

*Proof.* As in the proof of Proposition 6.3, Theorem 1.1 and Lemma 6.4 imply that the invariant measure  $\mu_{\text{inv}}$  is the image by  $\iota$  of a probability measure  $\tau$  on  $(-\pi, \pi]$ . Let  $(\theta_t)_t$  be the solution of

$$d\theta_t = \left((b-a)\sin\theta_t + \frac{1}{2}(a+b)\cos\theta_t\sin\theta_t\right) dt + \sqrt{a}(\cos\theta_t - 1)dB_0(t) + \sqrt{b}(\cos\theta_t + 1)dB_1(t). \tag{6.5}$$

The Itô formula implies once again

$$(\cos \theta_t, \sin \theta_t) \sim (\operatorname{tr}(\pi_{\hat{x}_t} \sigma_z), \operatorname{tr}(\pi_{\hat{x}_t} \sigma_x))_t$$

for  $(\pi_{\hat{x}_t})_t$  solution of (6.4) with initial condition  $\hat{x}_0 = \frac{1}{2}(\operatorname{Id} + \sin \theta_0 \sigma_x + \cos \theta_0 \sigma_z)$ . Hence  $(\iota(\theta_t))$  has the same distribution as  $(\hat{x}_t)$ . As in the proof of Proposition 6.3, standard techniques show that the unique invariant distribution for (6.5) has density proportional to the function  $\tau$  above.  $\square$

The density of the invariant probability distribution for three values of the pair  $(a, b)$  is plotted in Figure 2.

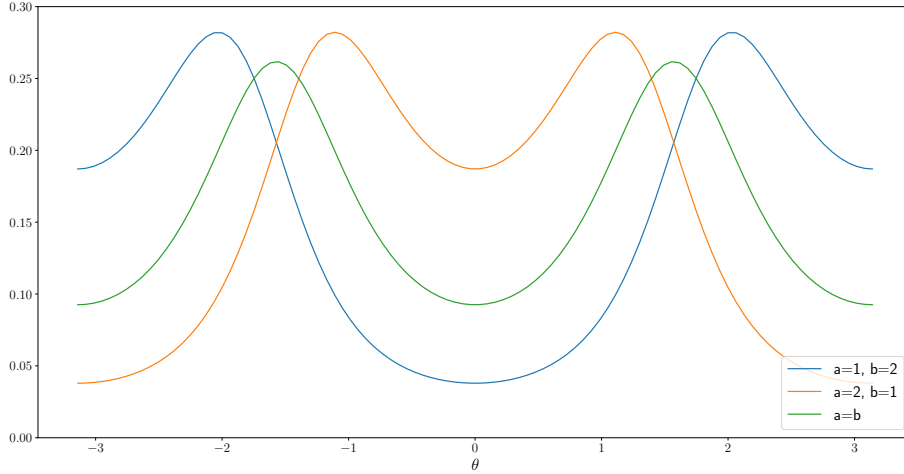


FIGURE 2. The density of the invariant probability distribution for Example 6.2.

**6.3. Thermal qubit, jump case.** Our third example is the second one where the stimulating coherent field has relatively small amplitude and is filtered out. Then, the signal is composed only of the photons absorbed or emitted by the qubit. The resulting trajectory involves only jumps related to these events. The parameters defining the model are then,  $H = 0$ ,  $I_b = \emptyset$  and  $I_p = \{0, 1\}$ ,  $C_0 = \sqrt{a} \sigma_+$  and  $C_1 = \sqrt{b} \sigma_-$  with  $a, b \in \mathbb{R}_+ \setminus \{0\}$ .

The process  $(\pi_{\hat{x}_t})_t$  is solution of

$$\begin{aligned} d\pi_{\hat{x}_t} = & a \left( \sigma_+ \pi_{\hat{x}_{t-}} \sigma_- - \frac{1}{4} ((\operatorname{Id} - \sigma_z) \pi_{\hat{x}_{t-}} + \pi_{\hat{x}_{t-}} (\operatorname{Id} - \sigma_z)) \right) dt \\ & + b \left( \sigma_- \pi_{\hat{x}_{t-}} \sigma_+ - \frac{1}{4} ((\operatorname{Id} + \sigma_z) \pi_{\hat{x}_{t-}} + \pi_{\hat{x}_{t-}} (\operatorname{Id} + \sigma_z)) \right) dt \\ & + \left( \frac{\sigma_+ \pi_{\hat{x}_{t-}} \sigma_-}{\operatorname{tr}(\sigma_- \sigma_+ \pi_{\hat{x}_{t-}})} - \pi_{\hat{x}_{t-}} \right) (dN_0(t) - a \operatorname{tr}(\sigma_- \sigma_+ \pi_{\hat{x}_{t-}}) dt) \\ & + \left( \frac{\sigma_- \pi_{\hat{x}_{t-}} \sigma_+}{\operatorname{tr}(\sigma_+ \sigma_- \pi_{\hat{x}_{t-}})} - \pi_{\hat{x}_{t-}} \right) (dN_1(t) - b \operatorname{tr}(\sigma_+ \sigma_- \pi_{\hat{x}_{t-}}) dt) \end{aligned} \quad (6.6)$$

where  $N_0$  and  $N_1$  are Poisson processes of stochastic intensities

$$t \mapsto \int_0^t a \operatorname{tr}(\sigma_- \sigma_+ \pi_{\hat{x}_{s-}}) ds \quad \text{and} \quad t \mapsto \int_0^t b \operatorname{tr}(\sigma_+ \sigma_- \pi_{\hat{x}_{s-}}) ds.$$

Assumptions **(Pur)** and **(L-erg)** hold as in Example 6.2.

**Proposition 6.6.** *Let  $\{e_1, e_2\}$  denote the canonical basis of  $\mathbb{C}^2$ . The invariant measure for Equation (6.6) is*

$$\mu_{\text{inv}} = \frac{a}{a+b} \delta_{\pi_{e_1}} + \frac{b}{a+b} \delta_{\pi_{e_2}}$$

*Proof.* It is enough to check from (6.6) that, if  $\hat{x}_0$  is either  $\hat{e}_1$  or  $\hat{e}_2$  then  $(\pi_{\hat{x}_t})_t$  is a jump process on  $(\pi_{\hat{e}_1}, \pi_{\hat{e}_2})$  with intensity  $b$  for the jumps from  $\pi_{\hat{e}_1}$  to  $\pi_{\hat{e}_2}$  and intensity  $a$  for the reverse jumps.  $\square$

**6.4. Finite state space Markov process embedding.** In this last example we show how we can recover all the usual continuous-time Markov chains using special quantum trajectories.

Let  $\{e_\ell\}_{\ell=1}^k$  be an orthonormal basis of  $\mathbb{C}^k$ , and  $(X_t)_t$  a  $\{e_1, \dots, e_k\}$ -valued Markov process with generator  $Q$  (we recall that  $Q$  is a  $k \times k$  real matrix such that  $\mathbb{E}(\langle v, X_t \rangle | X_0) = \langle v, e^{tQ} X_0 \rangle$  for any vector  $v \in \mathbb{C}^k$ ). Let  $H$  be diagonal in the basis  $\{e_\ell\}_{\ell=1}^k$ , let  $I_b = \emptyset$  and  $I_p = \{(i, j); i \neq j \text{ in } 1, \dots, k\}$  and for any  $(i, j) \in I_p$  let  $C_{i,j} = \sqrt{Q_{i,j}} e_j e_i^*$ .

**Proposition 6.7.** *Let  $(\hat{x}_t)_t$  be the quantum trajectory defined by Equation (1.4) and the above parameters. Then assumption (Pur) holds. In addition,*

- (i) *Let  $T = \inf\{t \geq 0 : \hat{x}_t \in \{\hat{e}_1, \dots, \hat{e}_k\}\}$ . If for all  $i$  there exists  $j$  with  $Q_{i,j} > 0$  then for any probability measure  $\mu$  over  $\text{PC}^k$ ,  $\mathbb{P}_\mu(T < \infty) = 1$ .*
- (ii) *Conditionally on  $\hat{x}_0 \in \{\hat{e}_\ell\}_{\ell=1}^k$ , the process  $(\hat{x}_t)_t$  has the same distribution as the image by  $x \mapsto \hat{x}$  of  $(X_t)_t$ .*
- (iii) *The assumption (L-erg) holds if and only if  $(X_t)_t$  accepts a unique invariant measure. In that case, the unique invariant measure  $\nu_{\text{inv}}$  for  $(\hat{x}_t)_t$  is the image by  $x \mapsto \hat{x}$  of the unique invariant measure for  $(X_t)_t$ .*

*Proof.* Note first that any  $C_{i,j}^* C_{i,j} = Q_{i,j} e_i e_i^*$ , so that (Pur) holds trivially.

To prove (i), let  $T_1 = \inf\{t > 0; \exists (i, j) \in I_p \text{ such that } N_{i,j}(t) > 0\}$ . Remark that because  $\text{tr}(C_{i,j} \pi_{\hat{x}_{s-}} C_{i,j}^*) = Q_{i,j} |\langle e_i, x_{s-} \rangle|^2$ , the sum  $\sum_{i,j} N_{i,j}$  of independent Poisson processes has intensity

$$\sum_{i,j} \int_0^t Q_{i,j} |\langle e_i, x_{s-} \rangle|^2 ds \geq t \min_i Q_i$$

where  $Q_i = \sum_j Q_{i,j}$  is positive by assumption, so that  $T_1$  is almost surely finite. Now consider the almost surely unique  $(i, j)$  in  $I_p$  such that  $N_{i,j}(T_1) > 0$ ; necessarily  $\text{tr}(C_{i,j} \pi_{\hat{x}_{T_1-}} C_{i,j}^*) > 0$ , and then  $\frac{C_{i,j} \pi_{\hat{x}_{T_1-}} C_{i,j}^*}{\text{tr}(C_{i,j} \pi_{\hat{x}_{T_1-}} C_{i,j}^*)} = \pi_{\hat{e}_j}$ , so that  $T \leq T_1$ . This proves (i).

Now, to prove (ii), remark that equation (1.3) can be rewritten in the form

$$\begin{aligned} d\pi_{\hat{x}_t} &= \sum_{(i,j) \in I_p} (\text{tr}(C_{i,j} \pi_{\hat{x}_{t-}} C_{i,j}^*) \pi_{\hat{x}_{t-}} - \frac{1}{2} \{C_{i,j}^* C_{i,j}, \pi_{\hat{x}_{t-}}\}) dt \\ &\quad + \sum_{(i,j) \in I_p} \left( \frac{C_{i,j} \pi_{\hat{x}_{t-}} C_{i,j}^*}{\text{tr}(C_{i,j} \pi_{\hat{x}_{t-}} C_{i,j}^*)} - \pi_{\hat{x}_{t-}} \right) dN_{i,j}(t). \end{aligned}$$

Let  $T_1$  be defined as above; then for  $t < T_1$  the process  $(\pi_{\hat{x}_t})_t$  satisfies

$$\pi_{\hat{x}_t} = \pi_{\hat{x}_0} + \sum_{(i,j) \in I_p} \int_0^t (\text{tr}(C_{i,j} \pi_{\hat{x}_{s-}} C_{i,j}^*) \pi_{\hat{x}_{s-}} - \frac{1}{2} \{C_{i,j}^* C_{i,j}, \pi_{\hat{x}_{s-}}\}) ds. \quad (6.7)$$

Starting with an initial condition  $\hat{x}_0 \in \{\hat{e}_\ell\}_{\ell=1}^k$ , one proves easily that the integrand is zero, which means that  $\pi_{\hat{x}_t} = \pi_{\hat{x}_0}$  for  $t < T_1$ . This shows in addition that for  $t < T_1$ , the intensity of  $N_{i,j}$  is

$$\int_0^t \text{tr}(C_{i,j} \pi_{\hat{x}_{s-}} C_{i,j}^*) ds = \begin{cases} Q_{i,j} t & \text{if } x_0 = e_i \\ 0 & \text{otherwise.} \end{cases}$$

Therefore, conditionally on  $x_0 = e_i$ ,  $T_1 = \inf\{t > 0; \exists j \neq i \text{ such that } N_{i,j}(t) > 0\}$  and there exists an almost surely unique  $j$  such that  $N_{i,j}(T_1) > 0$ . One then has

$$\pi_{\hat{x}_{T_1}} = \frac{C_{i,j} \pi_{\hat{x}_{T_1-}} C_{i,j}^*}{\text{tr}(C_{i,j} \pi_{\hat{x}_{T_1-}} C_{i,j}^*)} = \pi_{\hat{e}_j}.$$

This shows that for  $t \in [0, T_1]$  the process  $(\hat{x}_t)_t$  has the same distribution as the process of equivalence classes of  $X_t$ . This extends to all  $t$  by the Markov property of the Poisson processes. This proves (ii).

Points (i) and (ii) show that for  $t > T_1$ , the process  $(\hat{x}_t)_t$  has the same distribution as  $(X_t)_t$  with initial condition  $X_{T_1}$  satisfying  $\hat{X}_{T_1} = \hat{x}_{T_1}$ . Therefore any invariant measure for  $(\hat{x}_t)_t$  is the image by  $x \mapsto \hat{x}$  of an invariant measure for  $(X_t)_t$ . Theorem 1.1 and Section 4 show that  $(\hat{x}_t)_t$  admits at least one invariant measure, and that the invariant measure is unique if and only if **( $\mathcal{L}$ -erg)** holds. This implies that  $(X_t)_t$  has a unique invariant measure if and only if **( $\mathcal{L}$ -erg)** holds.  $\square$

**Acknowledgments.** The research of T.B., Y.P. and C.P. has been supported by the ANR project StoQ ANR-14-CE25-0003-01. The research of T.B. has been supported by ANR-11-LABX-0040-CIMI within the program ANR-11-IDEX-0002-02. The research of M.F. was supported in part by funding from the Simons Foundation and the Centre de Recherches Mathématiques, through the Simons-CRM scholar-in-residence program. Y.P. acknowledges the support of ANR project Non-Stops ANR-17-CE40-0006, and of the Cantab Capital Institute for the Mathematics of Information at the University of Cambridge.

## REFERENCES

- [1] S. Attal and Y. Pautrat. From repeated to continuous quantum interactions. *Annales Henri Poincaré*, 7(1):59–104, 2006. [2](#)
- [2] M. Ballesteros, N. Crawford, M. Fraas, J. Fröhlich, and B. Schubnel. Perturbation theory for weak measurements in quantum mechanics, systems with finite-dimensional state space. *Annales Henri Poincaré*, 20(1):299–335, 2019. [15](#)
- [3] A. Barchielli. Continual measurements in quantum mechanics and quantum stochastic calculus. In *Open quantum systems. III*, volume 1882 of *Lecture Notes in Math.*, pages 207–292. Springer, Berlin, 2006. [2](#)
- [4] A. Barchielli and V. P. Belavkin. Measurements continuous in time and a posteriori states in quantum mechanics. *J. Phys. A*, 24(7):1495–1514, 1991. [2](#)
- [5] A. Barchielli and M. Gregoratti. *Quantum trajectories and measurements in continuous time*, volume 782 of *Lecture Notes in Physics*. Springer, Heidelberg, 2009. The diffusive case. [2](#), [5](#)
- [6] A. Barchielli and A. S. Holevo. Constructing quantum measurement processes via classical stochastic calculus. *Stochastic Process. Appl.*, 58(2):293–317, 1995. [2](#), [5](#)
- [7] A. Barchielli and A. M. Paganoni. On the asymptotic behaviour of some stochastic differential equations for quantum states. *Infin. Dimens. Anal. Quantum Probab. Relat. Top.*, 6(2):223–243, 2003. [3](#)
- [8] M. Bauer, D. Bernard, and A. Tilloy. Computing the rates of measurement-induced quantum jumps. *Journal of Physics A: Mathematical and Theoretical*, 48(25):25FT02, 2015. [15](#)
- [9] B. Baumgartner and H. Narnhofer. The structures of state space concerning quantum dynamical semigroups. *Rev. Math. Phys.*, 24(02):1250001, 2012. [11](#)
- [10] V. P. Belavkin. Quantum stochastic calculus and quantum nonlinear filtering. *J. Multivariate Anal.*, 42(2):171–201, 1992. [2](#)
- [11] T. Benoist, M. Fraas, Y. Pautrat, and C. Pellegrini. Invariant measure for quantum trajectories. *Probability Theory and Related Fields*, Jul 2018. [4](#), [11](#), [13](#)
- [12] C. Bernardin, R. Chetrite, R. Chhaibi, J. Najnudel, and C. Pellegrini. Spiking and collapsing in large noise limits of sde’s. *arXiv preprint arXiv:1810.05629*, 2018. [15](#)
- [13] L. Bouten and R. van Handel. Quantum filtering: a reference probability approach. arXiv:math-ph/0508006, 2006. [2](#)
- [14] H.-P. Breuer and F. Petruccione. *The theory of open quantum systems*. Oxford University Press, New York, 2002. [5](#)
- [15] H. Carmichael. *An Open Systems Approach to Quantum Optics: Lectures Presented at the Université Libre de Bruxelles, October 28 to November 4, 1991*. Springer Science, Jan. 1993. [2](#), [5](#)
- [16] J. Dalibard, Y. Castin, and K. Mølmer. Wave-function approach to dissipative processes in quantum optics. *Physical Review Letters*, 68(5):580–583, Feb. 1992. [2](#)
- [17] E. B. Davies. Markovian master equations. *Communications in Mathematical Physics*, 39(2):91–110, June 1974. [2](#)
- [18] E. B. Davies. Markovian master equations. II. *Mathematische Annalen*, 219(2):147–158, June 1976. [2](#)
- [19] L. Diosi. Quantum stochastic processes as models for state vector reduction. *Journal of Physics A: Mathematical and General*, 21(13):2885–2898, July 1988. [2](#)
- [20] Q. Ficheux, S. Jezouin, Z. Leghtas, and B. Huard. Dynamics of a qubit while simultaneously monitoring its relaxation and dephasing. *Nature communications*, 9(1):1926, 2018. [14](#)
- [21] C. W. Gardiner and P. Zoller. *Quantum noise*. Springer Series in Synergetics. Springer-Verlag, Berlin, third edition, 2004. A handbook of Markovian and non-Markovian quantum stochastic methods with applications to quantum optics. [2](#)
- [22] N. Gisin. Quantum Measurements and Stochastic Processes. *Physical Review Letters*, 52(19):1657–1660, May 1984. [2](#)
- [23] V. Gorini, A. Kossakowski, and E. C. G. Sudarshan. Completely positive dynamical semigroups of  $N$ -level systems. *J. Mathematical Phys.*, 17(5):821–825, 1976. [1](#)

- [24] S. Haroche and J.-M. Raimond. *Exploring the quantum*. Oxford Graduate Texts. Oxford University Press, Oxford, 2006. Atoms, cavities and photons. 2
- [25] R. L. Hudson and K. R. Parthasarathy. Quantum Ito's formula and stochastic evolutions. *Comm. Math. Phys.*, 93(3):301–323, 1984. 2
- [26] J. Jacod and A. N. Shiryaev. *Limit theorems for stochastic processes*, volume 288 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer-Verlag, Berlin, second edition, 2003. 5
- [27] J. M. Kabanov, R. Š. Lipcer, and A. Širjaev. Absolute continuity and singularity of locally absolutely continuous probability distributions. i. *Mathematics of the USSR-Sbornik*, 35(5):631, 1979. 5
- [28] S. Karlin and H. M. Taylor. *A second course in stochastic processes*. Academic Press, Inc. [Harcourt Brace Jovanovich, Publishers], New York-London, 1981. 15
- [29] B. Kümmerner and H. Maassen. A pathwise ergodic theorem for quantum trajectories. *J. Phys. A*, 37(49):11889–11896, 2004. 13
- [30] G. Lindblad. On the generators of quantum dynamical semigroups. *Comm. Math. Phys.*, 48(2):119–130, 1976. 1
- [31] H. Maassen and B. Kümmerner. Purification of quantum trajectories. *Lecture Notes-Monograph Series*, 48:252–261, 2006. 3
- [32] S. Mac Lane and G. Birkhoff. *Algebra*. Chelsea Publishing Co., New York, third edition, 1988. 10
- [33] C. Pellegrini. Existence, uniqueness and approximation of a stochastic Schrödinger equation: the diffusive case. *Ann. Probab.*, 36(6):2332–2353, 2008. 2
- [34] C. Pellegrini. Poisson and diffusion approximation of stochastic master equations with control. *Ann. Henri Poincaré*, 10(5):995–1025, 2009. 2
- [35] C. Pellegrini. Markov chains approximation of jump-diffusion stochastic master equations. *Ann. Inst. Henri Poincaré Probab. Stat.*, 46(4):924–948, 2010. 2
- [36] F. Ticozzi and L. Viola. Quantum markovian subsystems: invariance, attractivity, and control. *IEEE Transactions on Automatic Control*, 53(9):2048–2063, 2008. 11
- [37] H. M. Wiseman and G. J. Milburn. *Quantum Measurement and Control*. Cambridge University Press, Jan. 2010. 2, 15
- [38] M. M. Wolf. Quantum channels & operations: Guided tour. <http://www-m5.ma.tum.de/foswiki/pub/M5/Allgemeines/MichaelWolf/QChannelLecture.pdf>, 2012. Lecture notes based on a course given at the Niels-Bohr Institute. 1, 2, 8, 11, 13

INSTITUT DE MATHÉMATIQUES DE TOULOUSE, UMR5219, UNIVERSITÉ DE TOULOUSE, CNRS, UPS IMT, F-31062 TOULOUSE CEDEX 9, FRANCE  
*E-mail address:* `tristan.benoist@math.univ-toulouse.fr`

DEPARTMENT OF MATHEMATICS, VIRGINIA TECH, BLACKSBURG, VA 24061, UNITED STATES OF AMERICA  
*E-mail address:* `fraas@vt.edu`

LABORATOIRE DE MATHÉMATIQUES D'ORSAY, UNIV. PARIS-SUD, CNRS, UNIVERSITÉ PARIS-SACLAY, 91405 ORSAY, FRANCE  
*E-mail address:* `yan.pautrat@math.u-psud.fr`

INSTITUT DE MATHÉMATIQUES DE TOULOUSE, UMR5219, UNIVERSITÉ DE TOULOUSE, CNRS, UPS IMT, F-31062 TOULOUSE CEDEX 9, FRANCE  
*E-mail address:* `clement.pellegrini@math.univ-toulouse.fr`