

Seismic reconstruction using FWI with reciprocity misfit functional and dual-sensors data

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AIP 2019, Grenoble, France

July 11th, 2019







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| Intro | Inverse Problem | Reciprocity WI | Experiments | Conclusion |
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- Time-Harmonic Inverse Problem, FWI
 - Dual-sensors data
 - Iterative reconstruction algorithm
- 3 Reconstruction procedure using dual-sensors data

4 Numerical experiments

- Experiments for acoustic media
- Comparison of misfit functions
- Changing the numerical acquisition with $\mathcal{J}_{\mathcal{G}}$
- Extension toward elasticity

5 Conclusion

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Reconstruction of subsurface Earth properties from seismic campaign: collection of **wave** propagation data at the surface.



- Reflection (back-scattered) partial data,
- nonlinear, ill-posed inverse problem.



Reconstruction of subsurface Earth properties from seismic campaign: collection of **wave** propagation data at the surface.





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2 Time-Harmonic Inverse Problem, FWI

- Dual-sensors data
- Iterative reconstruction algorithm



We consider propagation in acoustic media, given by the Euler's equations, heterogeneous medium parameters κ and ρ :

$$\begin{cases} -\mathrm{i}\omega\rho(\boldsymbol{x})\boldsymbol{v}(\boldsymbol{x}) = -\nabla\rho(\boldsymbol{x}), \\ -\mathrm{i}\omega\rho(\boldsymbol{x}) = -\kappa(\boldsymbol{x})\nabla\cdot\boldsymbol{v}(\boldsymbol{x}) + f(\boldsymbol{x}). \end{cases}$$

- *p*: scalar pressure field,
- **v**: vectorial velocity field,
- f: source term,

- κ : bulk modulus,
- $\rho :$ density,
- ω : angular frequency.



We consider propagation in acoustic media, given by the Euler's equations, heterogeneous medium parameters κ and ρ :

$$\begin{cases} -\mathrm{i}\omega\rho(\boldsymbol{x})\boldsymbol{v}(\boldsymbol{x}) = -\nabla p(\boldsymbol{x}), \\ -\mathrm{i}\omega p(\boldsymbol{x}) = -\kappa(\boldsymbol{x})\nabla\cdot\boldsymbol{v}(\boldsymbol{x}) + f(\boldsymbol{x}). \end{cases}$$

p: scalar pressure field, κ : bulk modulus,v: vectorial velocity field, ρ : density,f: source term, ω : angular frequency.

The system reduces to the Helmholtz equation when ρ is constant,

$$(-\omega^2 c(\boldsymbol{x})^{-2} - \Delta) p(\boldsymbol{x}) = 0,$$

with
$$c(\mathbf{x}) = \sqrt{\kappa(\mathbf{x})\rho(\mathbf{x})^{-1}}$$

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The quantitative inverse problem aims the recovery of the physical parameters from surface field measurements.

Dual-sensors record the pressure and vertical velocity:

$$\mathcal{F}(m = (\kappa, \rho)) = \{ p(\mathbf{x}_1), p(\mathbf{x}_2), \dots, p(\mathbf{x}_{n_{rev}}) \}; \\ \{ v_n(\mathbf{x}_1), v_n(\mathbf{x}_2), \dots, v_n(\mathbf{x}_{n_{rev}}) \}.$$



D. Carlson, N. D. Whitmore et al.

Increased resolution of seismic data from a dual-sensor streamer cable – Imaging of primaries and multiples using a dual-sensor towed streamer

SEG, 2007 - 2010

CGG & Lundun Norway (2017 - 2018)

TopSeis acquisition (www.cgg.com/en/What-We-Do/Offshore/Products-and-Solutions/TopSeis)

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FWI provides a **quantitative reconstruction** of the subsurface parameters by solving a minimization problem,

$$\min_{m\in\mathcal{M}} \quad \mathcal{J}(m) = \frac{1}{2} \|\mathcal{F}(m) - d\|^2.$$

d are the observed data,

• $\mathcal{F}(m)$ represents the simulation using an initial model m:



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Misfit functional ${\cal J}$





Numerical methods

- Adjoint-method for the gradient computation, L-BFGS method,
- Hybridizable Discontinuous Galerkin discretization method,
- elasticity, anisotropy, viscosity.



Hybridizable Discontinuous Galerkin (HDG) discretization:

- global matrix with the faces d.o.f. only,
- **local problem** to have the volume solution on DG d.o.f.



- Global matrix needs less memory than FE and DG (order),
- the local problems are small and embarrassingly parallel,
- ▶ 1st order: same accuracy for *p* and *v*,
- topography, sub-surface shapes.

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3 Reconstruction procedure using dual-sensors data



• Compare the pressure and velocity fields separately (L2):

$$\mathcal{J}_{L2} = \sum_{source} \frac{1}{2} \|\mathcal{F}_{p}^{(s)} - d_{p}^{(s)}\|^{2} + \frac{1}{2} \|\mathcal{F}_{v}^{(s)} - d_{v}^{(s)}\|^{2}$$



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Compare the reciprocity-gap:

$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2$$



G. Alessandrini, M.V. de Hoop, F. F., R. Gaburro and E. Sincich Inverse problem for the Helmholtz equation with Cauchy data: reconstruction with conditional well-posedness driven iterative regularization ESAIM: M2AN, 2019.



$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$

Motivated by Green's identity (using variational formulation).



$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$

- Motivated by Green's identity (using variational formulation).
- Reciprocity-gap functional from inverse scattering with Cauchy data.



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Reciprocity Waveform Inversion

| Stability | results | | | Ínva |
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Lipschitz-type stability for the Helmholtz equation with partial data,

 $\|m_1-m_2\| \leq \mathcal{C}\big(\mathcal{J}_{\mathcal{G}}(m_1,m_2)\big)^{1/2},$



Lipschitz-type stability for the Helmholtz equation with partial data, $\|m_1 - m_2\| \leq \mathcal{C} \big(\mathcal{J}_{\mathcal{G}}(m_1, m_2) \big)^{1/2},$

- for piecewise linear parameters.
- Using back-scattered data from one side in a domain with free surface and absorbing conditions,



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G. Alessandrini, M.V. de Hoop, R. Gaburro and E. Sincich

Lipschitz stability for a piecewise linear Schrödinger potential from local Cauchy data arXiv:1702.04222, 2017.

G. Alessandrini, M.V. de Hoop, F. F., R. Gaburro and E. Sincich

Inverse problem for the Helmholtz equation with Cauchy data: reconstruction with conditional well-posedness driven iterative regularization $% \label{eq:construction}$

ESAIM: M2AN, 2019.



It allows the separation of numerical and observational sources:

$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$



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$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$

- ▶ *s*₁ is fixed by the observational setup,
- ► *s*₂ is chosen for the numerical comparisons,
- arbitrary positions of computational source,
- no need for a priori information on the observational source: position and wavelet are not required,
- not possible with the traditional misfit.

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- Experiments for acoustic media
- Comparison of misfit functions
- Changing the numerical acquisition with $\mathcal{J}_{\mathcal{G}}$
- Extension toward elasticity



3D velocity model $2.5 \times 1.5 \times 1.2$ km using dual-sensors data.





We work with time-domain data acquisition.





We work with time-domain data (pressure and velocity).



For the reconstruction, we apply a Fourier transform of the time data.



$$\mathcal{J}_{L2} = \sum_{source} \frac{1}{2} \|\mathcal{F}_{\rho}^{(s)} - d_{\rho}^{(s)}\|^2 + \frac{1}{2} \|\mathcal{F}_{v}^{(s)} - d_{v}^{(s)}\|^2.$$

$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{\text{source source}} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$













But the major advantage of $\mathcal{J}_{\mathcal{G}}$ is the possibility to consider alternative acquisition setup.



Experiment with different obs. and sim. acquisition Unita-

$$\min \mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2$$

Acquisition for the measures s_1

- 160 sources,
- 100 m depth,
- point source,





$$\min \mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| d_v^{(s_1)T} \mathcal{F}_p^{(s_2)} - d_p^{(s_1)T} \mathcal{F}_v^{(s_2)} \|^2.$$

Acquisition for the measures s_1

- 160 sources,
- 100 m depth,
- point source,



Arbitrary numerical acquisition s_2

- ► 5 sources,
- ▶ 80m depth,
- multi-point sources,



No need to known observational source wavelet.

Differentiation impossible with least squares types misfit.

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Data from frequency between 3 to 15 Hz, domain size $2.5\times1.5\times1.2$ km, Simulation using 5 sources only.





Frequency from 3 to 15 Hz, $2.5 \times 1.5 \times 1.2$ km, Simulation using 5 sources only. -33% computational time.







Wave propagation in elastic media

$$-\nabla \cdot \underline{\sigma}(\mathbf{x}) - \omega^2 \rho(\mathbf{x}) \mathbf{u}(\mathbf{x}) = \mathbf{g}(\mathbf{x}),$$

 σ is the stress tensor; elastic isotropy, Lamé parameters λ and μ :

 $\underline{\sigma} = \frac{\lambda}{\mathrm{Tr}(\underline{\epsilon})}I_d + 2\underline{\mu}\underline{\epsilon}.$



Wave propagation in elastic media

$$-\nabla \cdot \underline{\sigma}(\mathbf{x}) - \omega^2 \rho(\mathbf{x}) \mathbf{u}(\mathbf{x}) = \mathbf{g}(\mathbf{x}),$$

 σ is the stress tensor; elastic isotropy, Lamé parameters λ and $\mu:$

 $\underline{\sigma} = \frac{\lambda}{\mathrm{Tr}(\underline{\epsilon})}I_d + 2\underline{\mu}\underline{\epsilon}.$

Three (heterogeneous) parameters to characterize the medium:

▶
$$\lambda$$
 and μ , or $c_{\rho} = \sqrt{(\lambda + 2\mu)/\rho}$, $c_{s} = \sqrt{\mu/\rho}$

• Density ρ .

Increased computational requirement and the ill-posedness of the inverse problem.



In elasticity, reciprocity needs measurements of σ and \boldsymbol{u}

$$\mathcal{F}(\boldsymbol{m} := (\lambda, \mu, \rho)) = \{ \boldsymbol{u}(\boldsymbol{x}) \mid_{\Sigma}, (\underline{\sigma}(\boldsymbol{x}) \cdot \boldsymbol{n}) \mid_{\Sigma} \}.$$



In elasticity, reciprocity needs measurements of σ and \boldsymbol{u}

$$\mathcal{F}(\boldsymbol{m} := (\lambda, \mu, \rho)) = \left\{ \boldsymbol{u}(\boldsymbol{x}) \mid_{\Sigma}, (\underline{\sigma}(\boldsymbol{x}) \cdot \boldsymbol{n}) \mid_{\Sigma} \right\}.$$
$$\mathcal{J}_{\mathcal{G}} = \frac{1}{2} \sum_{s_1} \sum_{s_2} \| \boldsymbol{d}_{\boldsymbol{u}}^{(s_1)T} \mathcal{F}_{\sigma \cdot \boldsymbol{n}}^{(s_2)} - \boldsymbol{d}_{\sigma \cdot \boldsymbol{n}}^{(s_1)T} \mathcal{F}_{\boldsymbol{u}}^{(s_2)} \|^2.$$

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2D elastic models of size 31 \times 7km.







2D elastic models of size 31 \times 7km.







- ► The density remains fixed; frequency from 0.50 to 7 Hz,
- Low frequency could be replaced by complex frequency (Laplace domain) or a priori information.



Observational acquisition:

- 150 sources,
- 20 m depth.

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Computational acquisition:

- 30 sources,
- ▶ 10 m depth.

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- ► The density remains fixed; frequency from 0.50 to 7 Hz,
- Low frequency could be replaced by complex frequency (Laplace domain) or a priori information.







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Quantitative time-harmonic inverse wave problem:

- Hybridizable Discontinuous Galerkin discretization, HPC,
- large scale optimization scheme using back-scattered data,
- acoustic, elastic, anisotropy, 2D, 3D, attenuation,
- stability and convergence analysis.

Reciprocity-gap waveform inversion:

- minimal information on the acquisition setup,
- reduced computational cost,
- other applications (elasticity, seismology, helioseismology),
- perspective: design efficient setup; data (rotational seismology).



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THANK YOU

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