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Mechanistic Insights into Polyion Complex Associations

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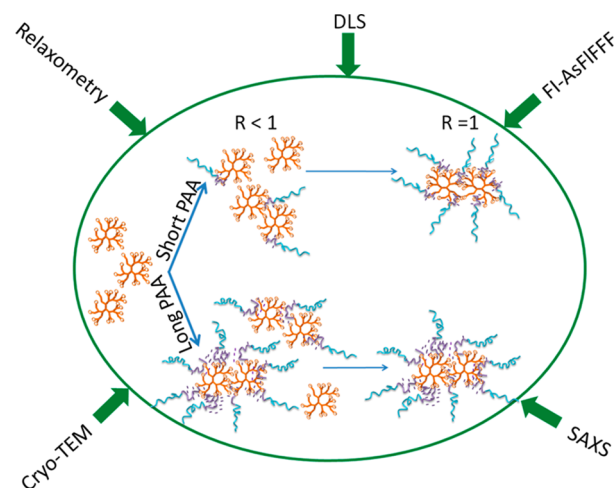
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Supporting Information

ABSTRACT: Polyion complex (PIC) micelles formed from the electrostatic interaction between oppositely charged polymers have been studied for their promising applications in the biomedical field as drug carriers or vectors for gene delivery. In spite of their asset of possible high drug loading, their formation process remains poorly studied. In this work, we investigate the properties of a series of PICs based on poly(ethylene oxide-*b*-acrylic acid) (PEO-PAA)/dendrigrift poly(L-lysine) (DGL3), using PEO-PAA with different compositions and average molecular weights. For each PEO-PAA/DGL3 pair, the complexes were characterized as a function of the ratios between acid and amine moieties combining different techniques: dynamic light scattering (DLS), flow field-flow fractionation (FIFFF), small-angle X-ray scattering (SAXS), and relaxometry. The coupling of batch techniques, i.e., DLS, SAXS, and relaxometry, together with a soft separation technique like FIFFF enabled a finer analysis to elucidate subtle details of the association process and of the polydispersity of the complexes. We show that the formation of PICs is more complex than previously described. In particular, we demonstrate that PICs with stoichiometry 1:1 may form at low ratios provided that the acidic block is long enough to neutralize the cationic dendrigrift with few polymer chains. Moreover, in such conditions, PICs with stoichiometry 1:1 often coexist with free dendritic polymers and other associated complex species.



■ INTRODUCTION

Polyelectrolytes are widely present and essential molecules in living systems. As such, their behavior has been studied for decades.¹ The characterization of the association between different oppositely charged polyelectrolytes has expanded thanks to the extensive work of Kabanov in the early 1990s dealing with so-called interpolyelectrolyte complexes.² Association of a cationic polyelectrolyte to a negatively charged one was shown to lead to soluble complexes only when the ratio between charged groups was different from 1 and the

polyelectrolyte characterized by the higher degree of polymerization (also called the "host") was in excess with respect to the other (also called the "guest").³ At the equimolar stoichiometry, macroscopic precipitation always occurred. For the soluble host-guest IPEC salt addition induced disproportionation; i.e., insoluble IPEC with equimolar stoichiometry precipitated, and

soluble IPEC with a composition different from the parent one, depleted in the short polyelectrolyte, formed. The discovery that such a precipitation could be avoided when using block copolymers instead of homopolymers led Kataoka and co-workers^{4–8} to develop the modern version of such assemblies, namely, polyion complex micelles (PICs) also called block ionomer complexes (BICs), micellar interpolyelectrolyte complexes,⁹ or complex coacervate core micelles.^{10,11} With the development of controlled radical polymerization and polymer engineering, the number of achievable block copolymers has exploded, leading to a burst of new PICs assemblies, thanks also to their potential application in nanomedicine as nanovectors for charged drugs or nanoplexes in DNA in gene therapy.^{12–14}

PIC assemblies are generally formed using at least one double hydrophilic block copolymer with a first neutral block and a second one charged. The other PIC partner consists either in a homopolyelectrolyte or another double hydrophilic block copolymer with a charged block opposite in sign. PICs' assembly is typically monitored with turbidimetry,¹⁵ static (SLS) or dynamic light scattering (DLS),^{5,11,16} cryo-transmission electron microscopy (cryo-TEM),¹¹ and recently two-dimensional proton nuclear magnetic resonance nuclear Overhauser effect spectroscopy.¹⁷

Based on the published studies on the PIC physical–chemical properties, several points of interest for applicative purposes and for the present study can be highlighted. (i) First, the formation and the structural properties of PIC assemblies strongly depend on the balance between the positive and negative charges and the relative length of each segment of the copolymers.⁹ (ii) PICs have been observed to be soluble for any ratio or only for specific ones, depending on the parameters of the copolymers.^{18–21} (iii) A third point is their general instability in the presence of salt due to the screening of the electrostatic interactions that drive PICs' formation.^{20,22} Disassembly is typically observed through static light scattering measurements, where the decrease of the scattered light intensity is associated with a decrease in size of the scattering objects. This drawback was overcome either by cross-linking^{7,16,18,23,24} or through introduction of hydrophobic moieties that provide further favorable contribution to self-assembly.²⁵

Although PICs have been widely studied, some open questions still exist regarding for instance the existence of a critical aggregation concentration.^{5,19,20,23,26,27} The mechanism of association has been rarely addressed, while it is of utmost importance to exert structural control and to fully exploit the applicative potential of PICs. For complexes whose formation is driven by charge compensation, it has been observed that the structure and composition of the aggregates depend on the order of addition, i.e., on the instantaneous composition of the mixture and not on the equilibrium one.^{28–30} For PICs based on poly(ethylene glycol-*b*- α,β -aspartic acid), PEO-P(Asp), and chicken egg white lysozyme, Kataoka suggested they could form through two alternative processes: cooperative (closed association) and noncooperative (open association).^{5,6} In excess of lysozyme, the cooperative process led to PICs with stoichiometry 1:1 of the charged groups even before the macroscopic charge balance was reached; at the molecular level, complete complexation of the charged macromolecular chain to its opposite partner always took place.⁵ In excess of PEO-P(Asp), the noncooperative process led to PICs which decreased in size when the macroscopic stoichiometry 1:1 of

the charged groups was approached; then a partial complexation of the chain with the presence of uncomplexed charged moieties occurred.⁶ Formation through a cooperative process was also claimed by Gohy et al. for poly(2-vinylpyridinium-*b*-ethylene oxide) and sodium poly(4-styrenesulfonate), for all compositions.²² In this example, sodium poly(4-styrenesulfonate) in excess adsorbed on charged balanced micelles.

It is clear that the path followed during formation has a dramatic impact on the final structure of the assemblies. This contribution addresses the above-mentioned structural and mechanistic aspects for a PICs family based on poly(ethylene oxide-*b*-acrylic acid), PEO–PAA, and dendrigraft polylysine (DGL3) through a broad multitechnique approach. Very few examples exist on PICs formed by branched molecules. They mostly dealt with the characterization of the complexes at stoichiometry 1:1,³¹ their resistance to solution conditions.^{32,33} Recently, Nguyen et al.³⁴ described the formation of PICs at pH 7 between insoluble polystyrene-*graft*-poly(2-vinylpyridine), GOPS-*g*-P2VP, and soluble poly(acrylic acid)-*b*-poly(2-hydroxyethyl acrylate), PAA-*b*-PHEA, at different ratios. They formed the assemblies first at pH 4.7 in order to guarantee the protonation of a portion of the amine groups of P2VP ($pK_a = 5$); then the solution pH was adjusted to 7. They observed a better stability of the assemblies for shorter PAA chains in block copolymers with the same length of the neutral block. The increase of the neutral block length was necessary in order to stabilize the assemblies with lower PAA content. Besides, very recently the possible tuning of the self-assemblies properties was nicely demonstrated, i.e., their morphology by varying the number of arms of the star polyelectrolyte³⁵ and their formation kinetics by using thermoresponsive nonionic block.³⁶

Here we selected a dendrigraft polylysine and varied the relative length of the anionic block in a poly(ethylene oxide)-*b*-poly(acrylic acid) (PEO–PAA) set. In a previous work, we characterized PICs formed at equal quantities of acid and amine groups with batch DLS and frit-inlet asymmetric flow field-flow fractionation (FI-AsFIFFF),³⁷ and we demonstrated the importance of a multitechnique approach to gain a clear picture of the structure of the complexes. Batch DLS showed one single broad size distribution, whereas the fractionation provided by FI-AsFIFFF before light scattering analysis revealed the presence of several populations in some of the systems. This analysis highlighted the importance of using FI-AsFIFFF, with its soft injection protocol before separation, in order to properly characterize the self-assemblies. Furthermore, it provided an estimation of the radius of gyration (R_g) and the hydrodynamic radius (R_h), and it opened some questions about the self-assemblies composition and their formation processes. In the present work, in order to address these questions, we focus on the effect of PEO–PAA composition. For each PEO–PAA/DGL3 pair, we varied the ratio between the cationic and anionic groups to follow the formation of the complexes and discern the mechanism of adduct formation. This process was monitored with a variety of techniques, including light scattering, but also flow field flow fractionation, small-angle X-ray scattering, and NMR relaxometry.

■ MATERIALS AND METHODS

Chemicals. Poly(ethylene oxide)-*b*-poly(acrylic acid) (PEO–PAA) was purchased from Polymer Source Inc. (Dorval Montréal, Canada) and was characterized by ¹H NMR and size exclusion chromatography; the obtained number-average molecular weights of the polymers are reported in Table 1. Third-generation dendrigraft polylysine

Table 1. Molar Masses of the Polymers Used As Characterized by ^1H NMR and Size Exclusion Chromatography

name	M_n PEO or dendrigraft DGL3 (g mol^{-1})	DP_n PEO	M_n PAA (g mol^{-1})	DP_n PAA
PEO–PAA 6-3	6000	136	3000	41
PEO–PAA 6-6.5	6000	136	6500	90
PEO–PAA 6-12	6000	136	12000	165
PEO–PAA 5-38	5000	114	38000	524
DGL3 ^a	22200			

^a DP_n or number of NH_2 groups per DGL3 is 123.

trifluoroacetate (DGL3) was kindly provided by Colcom (Montpellier, France). Water was purified through a filter and ion-exchange resin using a Purite device (resistivity $18.2 \text{ M}\Omega \text{ cm}$).

2.2. Formation of PICs. First, a stock solution of DGL3 (5 mg mL^{-1} , pH 7) was diluted with the adequate amount of water and then variable volumes of PEO–PAA (1 mg mL^{-1}) stock solution at pH 7 (adjusted with NaOH) were added. The samples have a final DGL3 concentration of 0.1 wt % and variable PEO–PAA/DGL3 ratios. We define the molar ratio R as the ratio between the concentration of PEO–PAA carboxylic groups and DGL3 amine groups ($[\text{AA}]/[\text{NH}_2]$). In the case of ^1H NMR experiments the polymer stock solutions and NaOH solution were prepared in D_2O .

2.3. Dynamic Light Scattering (DLS). DLS measurements were carried out at 25°C on a Malvern (Orsay, France) Zetasizer NanoZS. Solutions were analyzed in triplicate without being filtered in order to characterize the plain samples. Data were analyzed with a custom-made program named STORMS³⁷ using the general-purpose non-negative least-squares (NNLS) method. This treatment provided residuals lower than 5×10^{-3} . Both intensity- and number-weighted results are presented here in order to provide a more complete overview of the assemblies' populations. The assumptions and hypotheses leading to number-weighted results are further described in the [Supporting Information](#) (SI.1). The typical accuracy for these measurements was 10–20% for systems exhibiting a polydispersity index lower than 0.4. This was determined from repeated experiments associated with different fit treatments. Zeta-potential measurements were carried out using the Smoluchowski model.

2.4. Frit-Inlet Asymmetrical Flow Field-Flow Fractionation (FI-AsFFFF). A FI-AsFFFF channel was linked to an Eclipse 3 system (Wyatt Technology Europe, Dernbach, Germany, [Scheme SI 4.1](#)). The accumulation wall was a 1 kDa cutoff regenerated cellulose ultrafiltration membrane (ConSensuS, Ober-Hilbersheim, Germany). An Agilent 1100 Series isocratic pump (Agilent Technologies,

Waldbronn, Germany) with an in-line vacuum degasser and an Agilent 1100 Autosampler delivered the carrier flow and handled sample injection into the frit inlet AsFFFF channel. A 0.1 mm in-line filter (VVLVP, Millipore, Germany) was installed between the pump and the FFF channel. An aqueous 0.02% sodium azide solution was filtered (vacuum filtration system using 0.1 mm Gelman filters) before use as eluent. The spacer type was 250W or 350W (250 or $350 \mu\text{m}$, respectively). Samples were injected at 0.1 wt % concentration. The elution program used for injection flow V_i of 0.2 mL min^{-1} and for detector flow V_{out} of 1 mL min^{-1} for cross-flow V_x was 0.5 mL min^{-1} for 2 min, 2 mL min^{-1} for 38 min, decreased to 0.5 mL min^{-1} in 10 min followed by 0.5 mL min^{-1} for 40 min. The eluted sample components were detected with a multiangle light scattering (MALS) DAWN Heleos II (Wyatt Technology, Santa Barbara, CA) equipped with QELS (i.e., DLS) at 90° and an OptilabRex Refractometer (Wyatt Technology, Santa Barbara, CA) detectors. The MALS detector was normalized with bovine serum albumin (BSA). Calibration of scattering intensity was performed with HPLC-grade filtered toluene.

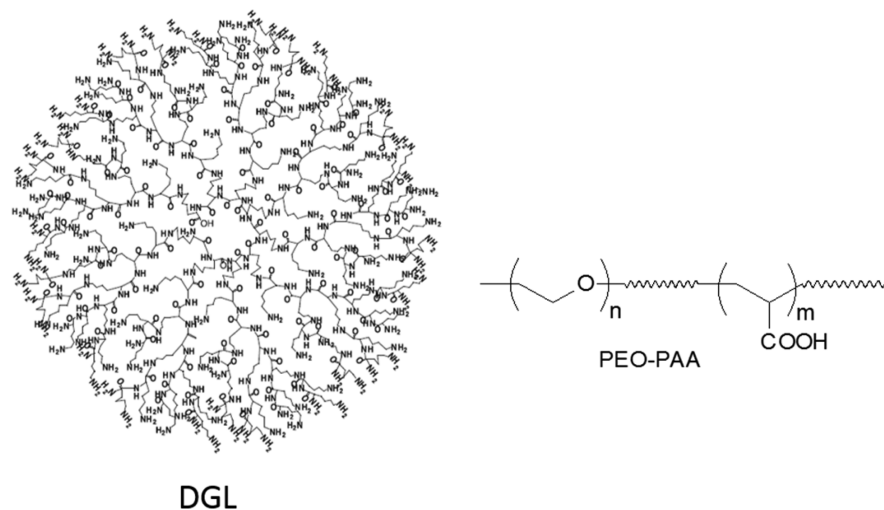
2.5. Small-Angle X-ray Scattering (SAXS). SAXS was performed at the high brilliance beamline ID02 at the European Synchrotron Radiation Facility (ESRF) in Grenoble, France. The 2D SAXS patterns were collected using a Rayonix MX-170HS ccd detector. The measured SAXS profiles were normalized to an absolute scale using the standard procedure reported elsewhere.³⁹ A combination of two sample-to-detector distances (10 and 1 m) was employed, covering a total q -range from 0.01 to 6 nm^{-1} . q is the scattering wave vector defined as $q = (4\pi/\lambda) \sin \theta/2$, λ being the wavelength ($\lambda \sim 1 \text{ \AA}$) and θ the scattering angle. The solutions were loaded in a flow through capillary of 2 mm diameter to ensure an accurate subtraction of the background (water) and to minimize beam damage of the samples.

2.6. ^1H Nuclear Magnetic Relaxation Dispersion (NMRD). NMRD profiles were acquired at 298 K using a high sensitivity fast field cycling relaxometer (Stelar, Mede, Italy). The longitudinal relaxation rates were extracted by single-exponential fitting of magnetization decay or magnetization buildup curves at ^1H Larmor frequencies ranging from 0.01 to 40 MHz. The errors of the collected relaxation rates were below 1%. ^1H NMRD profiles were collected for PEO–PAA 5-38 and PEO–PAA 6-3 based PICs. ^1H NMRD profiles for water solutions of DGL3 0.1 wt %, PEO–PAA 5-38 0.04 wt %, and PEO–PAA 6-3 0.09 wt % were also acquired for reference. These block copolymer concentrations were the maximum values used in the PICs solutions (the same as at $R = 1$).

^1H NMR spectra were recorded in D_2O on a Bruker Avance 500 MHz. The DGL3 concentration was fixed to 0.1 wt %.

2.7. Cryo-Electron Microscopy (Cryo-EM). $3 \mu\text{L}$ of sample was deposited onto glow-discharged lacey carbon grids and placed in the

Scheme 1. Chemical Structure (Idealized for DGL) of the Polymers Used in This Work



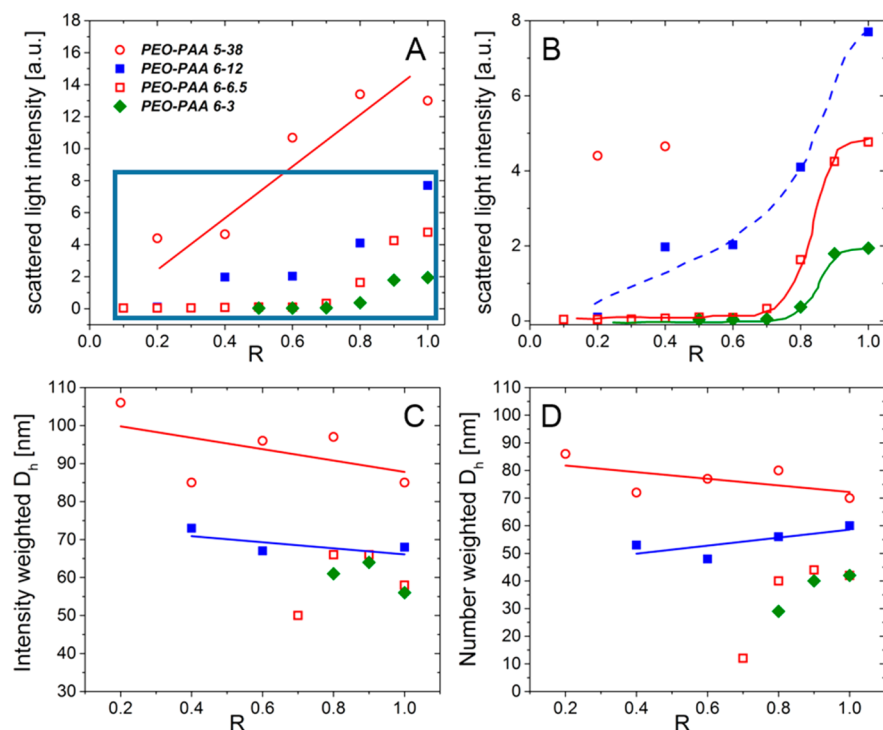


Figure 1. Evolution of the scattering properties as a function of R ($[\text{AA}]/[\text{NH}_2]$) for a solution of DGL3 0.1 wt % titrated with PEO-PAA: (A) scattered light intensity at 173° ; (B) zoom of the low value region of the scattered light intensity highlighted by a rectangle in (A); (C) intensity weighted hydrodynamic diameter; (D) number weighted hydrodynamic diameter. Lines are guides for the eye.

thermostatic chamber of a Leica EM-GP automatic plunge freezer, set at 20°C and 95% humidity. Excess solution was removed by blotting with Whatman no. 1 filter paper for 1 and 2 s, and the grids were immediately flash frozen in liquid ethane at -185°C . The frozen specimens were placed in a Gatan 626 cryo-holder, and cryo-EM was carried out on a Jeol 2100 microscope, equipped with a LaB_6 cathode and operating at 200 kV, under low dose conditions. Images were acquired with SerialEM software, with defocus of $1.5\text{--}3\ \mu\text{m}$, on a Gatan US4000 CCD camera. This device was placed at the end of a GIF Quantum energy filter (Gatan, Inc.), operated in zero-energy-loss mode, with a slit width of 25 eV. Images were recorded at a nominal magnification of 4000 corresponding to calibrated pixel sizes of $1.71\ \text{\AA}$.

Images were analyzed with ImageJ software from which we inferred the average diameter and corresponding standard deviation (confidence intervals were given as average diameter \pm twice standard deviation).

RESULTS

All the PICs presented here are based on poly(ethylene oxide)-*b*-poly(acrylic acid) (PEO-PAA)/dendrigrift polylysine (DGL3) (Scheme 1 and Table 1). Different PEO-PAA copolymers were selected in order to have a variety of molar masses and ratios between the two blocks. On the other hand, only dendrigrift polylysine generation three was used, with a molar mass of $22\,200\ \text{g mol}^{-1}$.

PICs were formed by adding the PEO-PAA stock solution to the diluted DGL3 solution. All PICs were formed at pH 7 with a ratio $R = [\text{AA}]/[\text{NH}_2]$ between 0.1 and 1. A comment is noteworthy regarding the chosen range of R used here. The pK_a of DGL3 α - and ϵ -amine functions can be estimated equal to 7 and 10, respectively,⁴⁰ ensuring that amine groups are mostly positively charged at pH 7. One has, however, to keep in mind that the pK_a value of PAA acid functions is highly dependent on its environment; values between 4.6⁴¹ and 6.5 are reported in the literature. In the case of multilayer assemblies the apparent

pK_a value of PAA diminished to ca. 4 for only two layers and as much as 4 pH units for six layers.⁴² In our case, we can thus reasonably assume that all the PAA units are in the carboxylate form, and therefore they would associate to the available DGL3 ammonium units.

In the following, we present the characterization of mixtures of PEO-PAA and DGL3 varying their ratios. Different complementary batch techniques (DLS, SAXS, NMRD, ^1H NMR) and a soft chromatographic separation method (FI-AsFFFF) are discussed and compared.

Dynamic Light Scattering (DLS). Batch DLS results are presented in Figure 1. We typically titrated DGL3 with PEO-PAA varying R from 0 to 1. For all samples, the DLS analysis was performed only if the scattered light intensity was higher than 1000 counts per second (background level). Indeed, for poorly scattering solutions, the DLS curves were not analyzed.

The evolution of the scattered light intensity (Figure 1A,B) exhibits two distinct types of behavior, either a regular increase over a large ratio range for PEO-PAA 5-38 or a lag phase followed by a sudden increase for all the other systems.

For PEO-PAA 5-38, the size of the objects did not significantly change between ratios 0.2 and 1 (Figure 1C,D), and polydispersity index slightly decreased (Table 3), while the scattered intensity regularly increased with R . This behavior can be associated with an increase of the number density of the scattering objects, the size of which does not evolve. This trend has been already observed in the literature⁵ and is indicative of the presence in solution of complexes with a well-defined stoichiometry irrespectively of the R ratio. In this case, we can assume that objects of a well-defined size and composition are formed even for low ratios: their concentration, but not their nature, depends on R . The fact that the polydispersity index decreases while R increases is probably due to the presence of

free DGL3 molecules in solution whose concentration diminishes.

For PEO–PAA 6-3 and PEO–PAA 6-6.5, the scattered light intensity increased suddenly at ratios close to 0.7–0.8. At this same ratio, the correlation function can be analyzed (see Figure S1.1) to extract hydrodynamic sizes. This behavior could be indicative of complexes formed only close to $R = 1$. Further experimental evidence are required to conclude on the low mixing ratios.

On the other hand, for PEO–PAA 6-12, the scattered light intensity presented two linear regions with a boundary close to 0.7–0.8. The hydrodynamic radius could be extracted starting from $R = 0.4$; the number weighted size moderately increased with increasing R , whereas the intensity weighted size was roughly constant and the polydispersity index slightly decreased. These results suggest that PICs with well-defined stoichiometry formed even at low ratios as in the case of PEO–PAA 5-38.

Zeta Potentials. Zeta potentials (ζ) were also determined by electrophoretic light scattering (ELS) and are shown in Figure 2 as a function of the AA/NH₂ ratio. ζ is defined as the

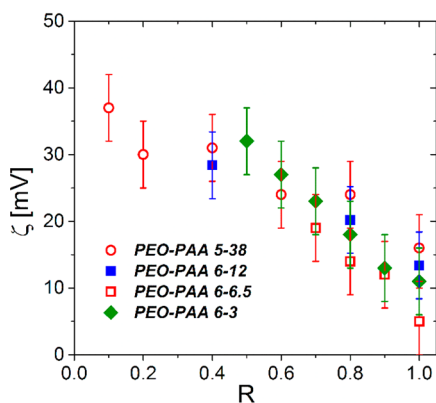


Figure 2. Zeta potentials of PICs self-assemblies. The ratios correspond to that of acrylic acid units over amine units. (○) PEO–PAA 5-38, (■) PEO–PAA 6-12, PEO–PAA 6-6.5, (◆) PEO–PAA 6-3.

potential of the plane where slip between the object and the bulk solution is assumed to occur.⁴⁴ It postulates nonporous objects exhibiting a hard surface and therefore not deformable upon shear. This is naturally doubtful for PICs. Furthermore, even the very existence of a slip plane is questioned for soft objects.⁴⁴

Going beyond this issue, we consider here the qualitative trend of ζ as a function of R . ELS implies the detection of light scattered from the objects. The presence of a scattering intensity roughly 1 order of magnitude higher than that of the solvent is a prerequisite for reliable zeta potentials values. Therefore, we do not report the values for all the low R ratios. For all samples, a regular decrease of ζ is observed (Figure 2), leading to values close to zero when approaching $R = 1$. It means that for all investigated systems at $R = 1$ electrostatic charge is balanced in the nano-objects. For a pure DGL3 solution in the same conditions (0.01 wt % at pH 7) ζ could be measured equal to 42 mV, and this is in agreement with the extrapolation at $R = 0$ of straight lines of Figure 2. On the other hand, the ζ trend cannot alone discriminate PICs formation processes. In the case of PEO–PAA 5-38 and PEO–PAA 6-12, from previously discussed DLS results, we observed that PICs

objects with well-defined stoichiometry formed already below $R = 1$. Indeed, for R below 1, ELS measures an average of free DGL3 molecules and PICs assemblies. For the other systems, the same trend could also originate from objects possessing a net positive charge on their surface, which decreases progressively with increasing the PEO–PAA quantity in an open association mechanism.

Small-Angle X-ray Scattering. In order to get further insight into the formation process of the assemblies, we performed a structural characterization study with SAXS. Figure 3 shows the SAXS profiles of DGL3 0.1 wt % solution in the presence of the different PEO–PAA block copolymers at different ratios, and Table 2 reports the fit results.

We first discuss the peculiar case of PEO–PAA 6-3 presented in Figure 3A.

At low PEO–PAA content, the SAXS curves show a pronounced interaction peak at low q due to electrostatic repulsions between the charged molecules of DGL3. Then, the scattering intensity increases with the PEO–PAA content as the interactions between the two polymers diminish the effective DGL3 charge. The evolution of the scattering signal shows a clear transition between $R = 0.5$ and 0.8 when larger structures are present in solution. This is confirmed by the Guinier analysis of the scattering curves, which allows extracting the radius of gyration of the scattering objects for the different ratios (Table 3).

Above $R = 0.8$, the scattering curves could be described using a core–shell sphere form factor (eqs 4, 6, and 7 in SI.2), which includes also a term for the scattering at high q originating from the chain statistics and the interchain interactions of the polymer inside the shell. This intracore contribution is well reproduced by an Ornstein–Zernike type term⁴⁵ with an exponent typical for swollen polymer chains in good solvent. The scattering length densities (SLD) for lysine, PAA, and PEO calculated according to their molecular composition are 0.001 09, 0.001 32, and 0.001 04 nm^{−2}, respectively. Assuming NH₃⁺:AA[−] interactions in the core of the assemblies, we can estimate for the SLD of the core an average starting value of 0.0012 nm^{−2}, which was adjusted in the fitting procedure to 0.001 29 nm^{−2}, in good accordance with theoretical expectations. The resulting value is higher than that of pure lysine molecules, indicating that the latter is not hydrated and mainly formed by DGL3 interacting with AA[−] groups, while the shell is mostly made by the hydrated PEO chains. The average size for the micellar core and shell was found equal to 5.3 and 6.6 nm, respectively (Table 2).

The values of the radius of gyration for $R \leq 0.5$ are comparable with the size of the DGL3 alone. Nevertheless, their scattering curves change with respect to pure DGL3. In order to better understand the origin of the scattering signal, the curves have been modeled with linear combination of the experimental scattering intensity of noninteracting DGL3 and PEO–PAA (Figure SI.2.1 in SI.2 in the case of $R = 0.5$). The curves could be described neither as the sum of noninteracting DGL3 and PEO–PAA nor as pure micelles ($R = 1$) or a mixture of micelles ($R = 1$) and DGL3 alone. This suggests that when PEO–PAA is added in solution, it starts interacting with DGL3. The observed increase of the low- q intensity can be explained with the decrease of the net DGL3 charge (correspondingly less repulsive interactions), the corresponding increase of the SLD of the core with respect to DGL3 alone due to NH₃⁺:AA[−] electrostatic screening, and the increase of the total molar mass.

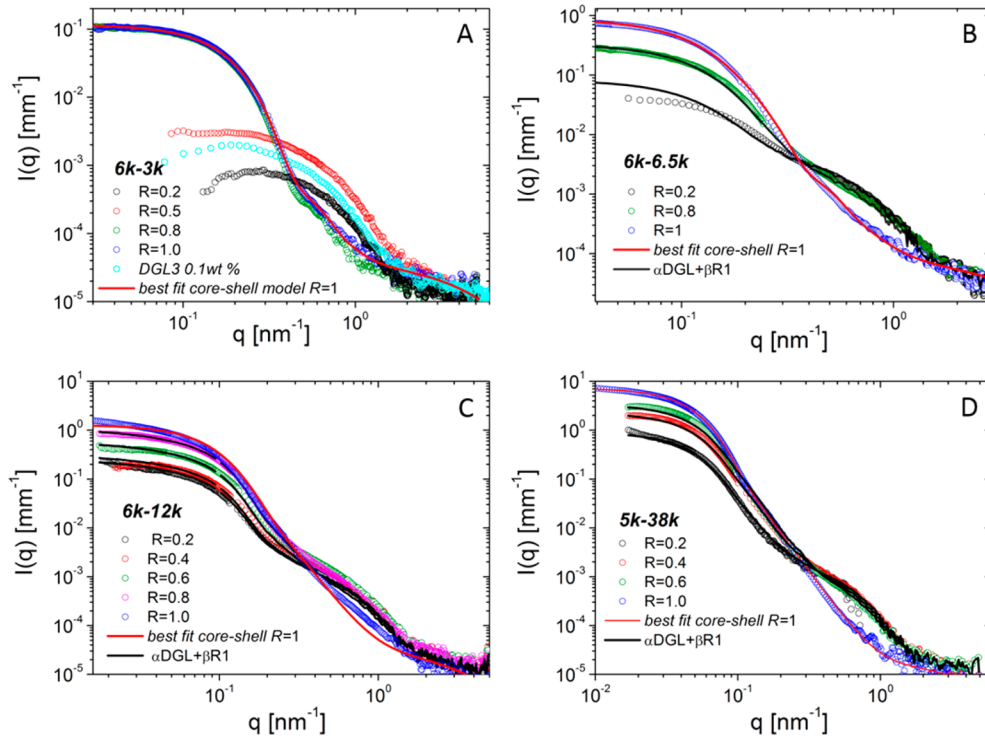


Figure 3. Small-angle X-ray scattering as a function of the scattering vector for PEO–PAA/DGL3 solutions. The DGL3 concentration is constant and equal to 0.1 wt %; PEO–PAA is added at different $R = \text{AA}/\text{NH}_2$ ratios in the case of (A) PEO–PAA 6-3, (B) PEO–PAA 6-6.5, (C) PEO–PAA 6-12, and (D) PEO–PAA 5-38.

Table 2. Fitting Parameters for the Core–Shell Model Described in SI.2 (Eqs 4, 6, and 7) in the Case of $R = 1$ ^a

	R_{core} (nm)	R_{shell} (nm)	R_{sphere}
PEO–PAA 6-3	5.3 ± 1	6.6	
PEO–PAA 6-6.5	7.8 ± 2	9.2	
PEO–PAA 6-12			15 ± 5
PEO–PAA 5-38			30 ± 9

^a R_{core} and R_{shell} are respectively the core radius and the shell radius of the assemblies. ρ_{core} and ρ_{sh} are 0.001 29 nm^{−2} and 0.001 04 nm^{−2} respectively for all assemblies.

We can then conclude that in analogy to “monodendrimer” complexes already observed by Gelissen et al.,⁴⁶ at low R values, DGL3: PAA complexes are present in solution; the PEO block in the shell is highly solvated so only the DGL3 size is clearly resolved by SAXS.

Micelles with a coacervate core and a PEO shell form only approaching $R = 1$. A further experimental proof of the presence of $\text{NH}_3^+ : \text{AA}^-$ interactions at ratios lower than 0.8 was then gained by recording ¹H NMR spectra on solutions containing DGL3 and PEO–PAA 6-3 at different ratios (see subsequent part).

Let us now consider the other block copolymers, PEO–PAA 6-6.5, PEO–PAA 6-12, and PEO–PAA 5-38, whose scattering curves are shown in Figures 3B–D. In the case of $R = 1$ all the objects formed are spherical. The homogeneous sphere model is usually employed in the literature,⁴⁸ and the core contribution is more often the dominant one due to the shell solvation. In our case, the core–shell model works well for the shorter PAA block; for the longer ones the model underestimates both the size and the SLD of the core, toward unphysical values. For this reason, we preferred using a homogeneous sphere model (eqs 5–7 in SI.2) with a value of

the contrast balanced between the one of the $\text{NH}_3^+ : \text{AA}^-$ adducts and PEO. The fit results are reported in Table 2.

For R smaller than 1 (see Figures 3B–D), the scattering curves present two different features in the q range investigated: (a) the scattering signal at low q increases with R , (b) at high q it decreases, and the second shoulder disappears when $R = 1$. The analysis of SAXS data proves that the solutions can be better described as a mixture of micelles with 1:1 stoichiometry and free DGL3 molecules rather than micelles growing when the PEO–PAA content increases. This is well illustrated in the case of PEO–PAA 6-6.5: when using the core–shell model (eq 4 in SI.2) to fit the experimental data at different R values, the scattering signal in the high- q region cannot be only described by the fluctuations of the polymeric chains in the corona and in order to compensate for the high- q scattering signal a too high and unphysical value for the SLD of the core is demanded. On the contrary, the scattering curves are well described by a mixture of charge balanced micelles (stoichiometry 1:1) and free DGL3 molecules. With increasing the amount of PEO–PAA, the micelles number density in solution increases (the scattering intensity increases at low q), and correspondingly the content of free DGL3 decreases (the scattering intensity decreases at high q until disappearing at $R = 1$). We consider that the concentration of the block copolymer free in solution is negligible, and the scattering curves are modeled with a linear combination of the scattering signals from free DGL3 and the core–shell micelles completely formed at $R = 1$ (charge balanced micelles):

$$I(q) = \alpha I(q)_{\text{DGL3}} + \beta I(q)_{R=1}$$

The best values of parameters α and β are reported in Table SI.2.2, and the corresponding curves are showed in Figures 3B–D. In all cases the solution is mainly constituted by

Table 3. Comparison of PICs Size Obtained from Different Techniques

PEO-PAA	R	R_g^b SAXS (nm)	R^c SAXS (nm)	R_g (n) ^c FFF (nm)	R_h (n) ^c FFF (nm)	R_g/R_h FFF (n)	R_h number DLS ^d (nm)	R_h Int DLS ^d (nm)	cumulant PDI	R_g SAXS/ R_h DLS
5-38	0.2	33					43	53	0.16	0.62
	0.4	31.5					36	43	0.14	0.73
	0.6	31.5					39	48	0.10	0.66
	0.8						40	49	0.07	
	1	34	30.6	26/46 ^a	38/53	0.68/0.87	35	46	0.07	0.7
6-12	0.2	17					n.d	n.d	n.d	
	0.4	18.5		31	38	0.55	24	37	0.11	0.5
	0.6	18		16	28	0.57	27	33	0.12	0.55
	0.8	18.7		17/27 ^a	20/41 ^a	0.85/0.66 ^a	27	33	0.09	0.57
	1	19	16.5	25	35	0.71	30	34	0.06	0.56
6-3	0.2	2.4								
	0.4	3.3								
	0.5	3.3								
	0.6	3.3								
	0.8	10	12.6	12	18	0.67	15	30.5	0.3	0.33
	1	9.8	11.9	17	27	0.63	21	28	0.1	0.35

^aBimodal. ^b30% error. ^cThe error for each sample is provided in Table SI.4.1; in most cases, this was below 15%. ^dAs discussed in the Materials and Methods section, the error was estimated in the 10–20% range. ^eThe radius of the object as obtained from the fitting procedure.

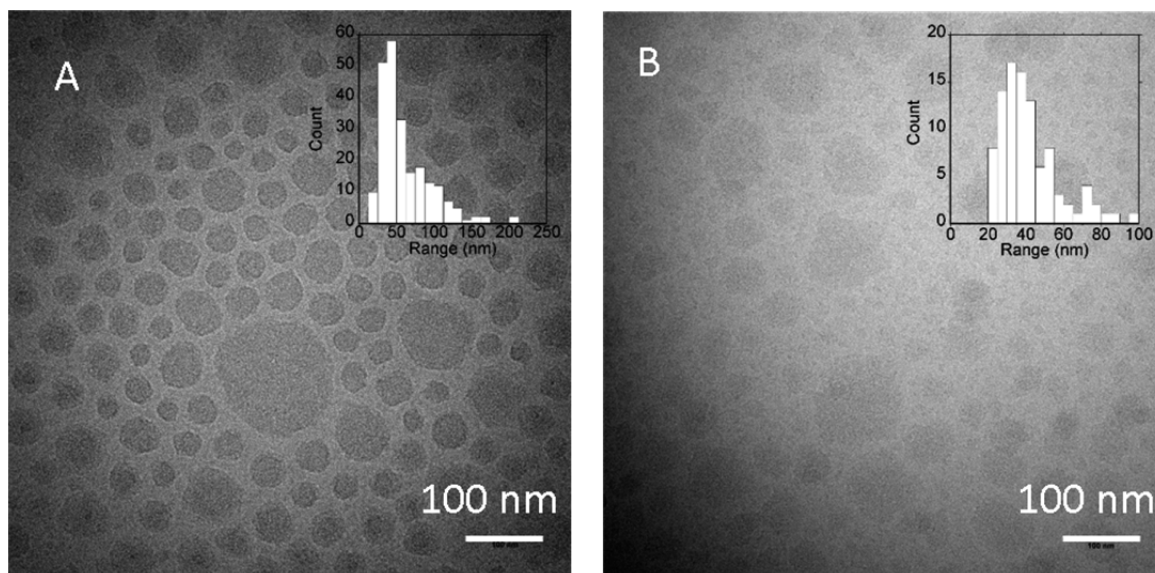


Figure 4. Cryo-EM of PEO-PAA/DGL3 PICs at $R = 1$: (A) PEO-PAA 5-38; (B) PEO-PAA 6-12.

noninteracting DGL3 molecules until $R = 0.8$. This is in agreement with the sudden increase of the light scattering intensity signal observed in Figures 1A,B around this value of R for the three copolymers.

It should be noted that the above-cited model does not perfectly describe the behavior of the scattering curves in the low q range for $R = 0.2$ of PEO-PAA 6-6.5 and the intermediate q region of PEO-PAA 6-12 and PEO-PAA 5-38. Additional structural information might contribute to the overall scattering signal. This will be discussed later in text in comparison with the results from the other experimental techniques. In conclusion, structural investigation with SAXS indicates that for PEO-PAA 5-38, PEO-PAA 6-12 and PEO-PAA 6-6.5 micelles already form at low R values, but all DGL3 molecules are incorporated in micelles only at 1:1 stoichiometric ratio. These results are more in accordance with a close association type mechanism, but we cannot exclude the presence of micelles with stoichiometry different from 1:1.

Cryo-EM. Cryo-EM was also employed to characterize some PICs at $R = 1$ (Figure 4). However, due to poor contrast, only PEO-PAA 5-38 and PEO-PAA 6-12 images could be analyzed. It must be underlined that in cryo-TEM no contrast agent is used, so the contrast of the objects depends on their chemical nature and density. PEO-PAA 6-12 based PICs possessed lower contrast than PEO-PAA 5-38, possibly because they were more hydrated. The statistical analysis provides an average diameter of 60 ± 34 and 42 ± 16 nm for PEO-PAA 5-38 and PEO-PAA 6-12, respectively, consistent with the SAXS (60 ± 19 and 30 ± 10 nm) and DLS (92 ± 18 and 68 ± 15 nm) analyses.

¹H NMR. As discussed in the previous paragraphs, charged balanced micelles are formed at all mixing ratios in the case of all used PEO-PAA with the exception of PEO-PAA 6-3. SAXS experiments suggested that an interaction between the groups of opposite charge already occurred at low mixing ratios. ¹H NMR spectra were then recorded at different mixing ratios.

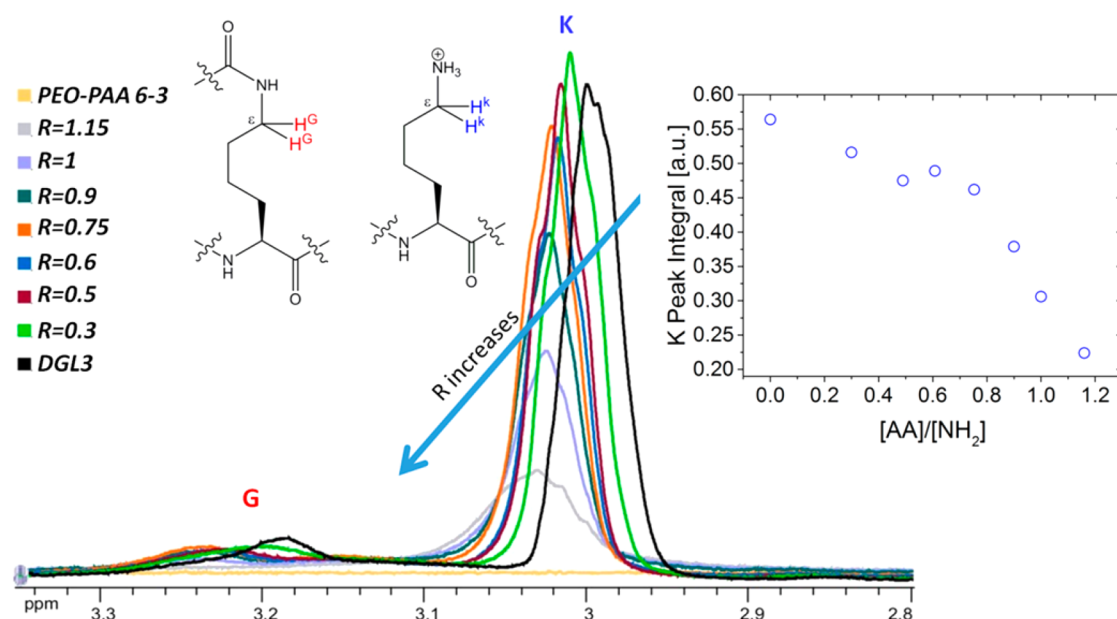


Figure 5. Evolution of ^1H NMR spectra of polylysine protons. Spectra are registered in D_2O with DGL3 0.1 wt % with different PEO–PAA 6-3 amounts of stock solution at pH 7 (adjusted with NaOH). Insets: scheme of lysine residues interested in the NH_3^+ : AA^- interactions; evolution of K polylysine protons integral in ^1H NMR spectra while increasing $R = [\text{AA}]/[\text{NH}_2]$.

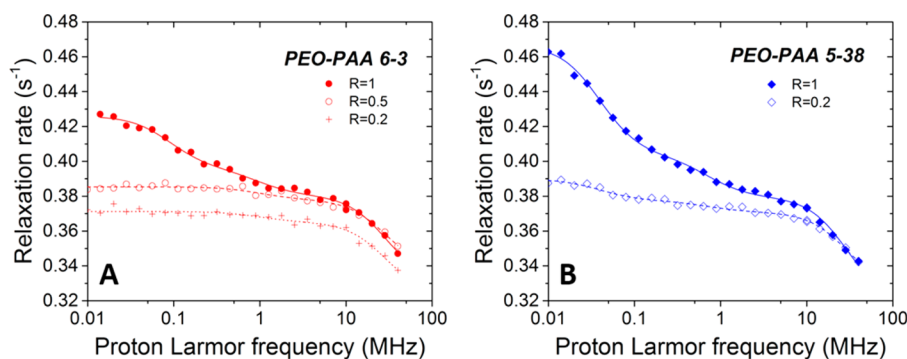


Figure 6. ^1H NMRD profiles for PICs solution at different acid/amine ratios: (A) PEO–PAA 6-3:DGL3 at 0.2:1 (+), 0.5:1 (○), and 1:1 (●) ratios; (B) PEO–PAA 5-38:DGL3 at 0.2:1 (◇) and 1:1 (◆) ratios. In all solutions, the DGL3 concentration was fixed to 0.1 wt %. The lines are the best fit profiles calculated with the parameters reported in Table S1.3.

Figure 5 presents a zoom of protons K and G of the polylysine polymer and shows a downfield shift of the signals by increasing the amount of carboxylic groups. The evolution of the K peak integral is shown in the inset of Figure 5. The signal integral decreases slowly until $R = 0.6$, when a dramatic decrease starts, which can be attributed to the micelle formation.

Therefore, NMR experiments confirm that PEO–PAA 6-3 interacts with DGL3 already at low ratios, even if micelles are not formed yet.

^1H NMRD. Relaxometry profiles were collected at 298 K for water solutions of DGL3, and for PEO–PAA 5-38 and PEO–PAA 6-3 based PICs (Figure 6). The measured ^1H relaxation rates are due to fluctuations in energy of the proton–proton dipole–dipole interactions and depend on the time constants of the motions modulating such interactions. Therefore, they can provide information on the characteristic time constants of such motions.^{49–51} The profiles of DGL3 and the two block copolymers showed a single dispersion at high fields, corresponding to a correlation time of about 3 ns (Figure SI.3.1 in SI.3).

When PEO–PAA 6-3 was added with $R = 0.2$, only a very small difference could be detected with respect to the NMRD profile of DGL3 alone; on the other hand, the relaxation rates increased appreciably at $R = 0.5$, without any dispersion appearing below 1 MHz. Only at $R = 1$, a low field dispersion (at about 0.1 MHz) appeared, indicating the presence of slow reorienting aggregates. The NMRD profiles were fitted using eq 1 described in SI.3.^{52–54} Three correlation times were needed for the fit of the curve obtained at $R = 1$, whereas two correlation times were enough for the fit of the curves at lower ratio (Table SI.3.1 in SI.3). The fastest correlation time was fixed to 3 ns, as obtained from the DGL3 profile, and in agreement with the average value obtained from the best fit τ_3 values obtained from all collected profiles (see Table SI.3.1 in SI.3). The intermediate correlation times ranged from 60 ($R = 0.2$ and $R = 0.5$) to 130 ns ($R = 1$), indicating the presence of motions slower than 3 ns (observed for DGL3 and for the two block copolymers), presumably arising for the presence of $\text{NH}_3^+:\text{AA}^-$ interactions. The slowest correlation time at $R = 1$ was about 1100 ns.

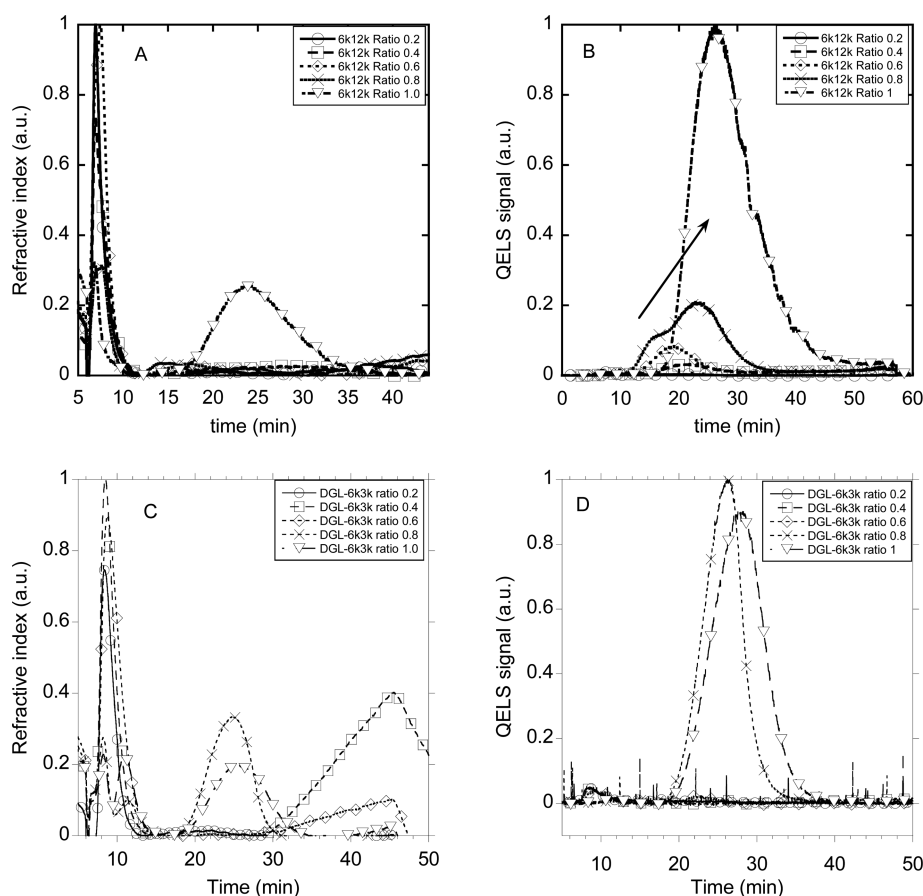


Figure 7. Characterization of PEO-PAA/DGL3 PICs by FI-AsFFFF at different ratios. The ratios correspond to that of acrylic acid units over amine units. (○) ratio 0.2, (□) ratio 0.4, (◇) ratio 0.6, (×) ratio 0.8, (▽) ratio 1.0. (A) PEO-PAA 6-12 RI signal, (B) PEO-PAA 6-12 DLS signal, (C) PEO-PAA 6-3 RI signal, (D) PEO-PAA 6-3 DLS signal.

In PEO-PAA 5-38 based PICs, the low field relaxation rates appreciably increased even for a small ratio ($R = 0.2$), and the presence of a low field dispersion (between 0.01 and 0.1 MHz), corresponding to a slow correlation time, could be clearly appreciated. This indicates the presence of slow reorienting aggregates already at low R values. When R was increased to 1, the increase in the low field relaxation rates became higher as a result of the larger contribution from the slow reorienting aggregates, with an inflection positioned basically at the same field. This indicated that the size of the formed aggregates did not increase with increasing R ; only their concentration did. The NMRD profiles were fitted using eq 1 described in SI.3 and three correlation times, after checking that two correlation times were not sufficient to satisfactorily reproduce the profiles. Since the obtained best fit values of the correlation times τ_i ($i = 1, 2, 3$) were very similar when an unconstrained fit was performed, τ_1 and τ_2 were constrained to have the same values in the fit of the two PEO-PAA 5-38 based PICs profiles. As for PEO-PAA 6-3 based PICs, τ_3 was fixed to 3 ns. The fit was in good agreement with the experimental data and indicated the presence of a fast correlation time of about 2200 ns and an intermediate one of 120 ns for both 0.2 and 1 ratios. The contributions of these correlation times increased with increasing R , as a result of the larger fraction of protons in the aggregates. We can then conclude that the slow reorienting aggregates characterized by the high correlation times are the micelles described in the SAXS section. This indicated the presence of the micelles only at $R = 1$ in the case of PEO-PAA

6-3, even if the sign of interactions are already present for lower R values. Instead, in the case of PEO-PAA 5-38 micelles are present already at $R = 0.2$, and their concentration increases with the copolymer content. It should be stressed that the slowest correlation time obtained for PEO-PAA 5-38 is a factor 2 smaller than the slowest time obtained for PEO-PAA 6-3. This finding is in line with the smallest hydrodynamic radius measured for PEO-PAA 6-3.

Field Flow Fractionation. We performed experiments using field flow fractionation which allows for a separation of nano-objects according to their size. Figure 7 shows the fractograms relative to PEO-PAA 6-3 and PEO-PAA 6-12 which represent well the two different behaviors. The fractograms relative to the other polymers are reported in the Supporting Information (Figure SI.4.2).

Figures 7A,C show the refractive index signal, which is more sensitive to species concentration. For both PEO-PAA 6-3 and PEO-PAA 6-12, a low retention time peak ($t = 6.5$ min), corresponding to DGL3 molecules, was present for all ratios, and it drastically diminished at $R = 0.8$.

For PEO-PAA 6-3 a peak at retention time around 30 min appeared at $R = 0.8$. According to the DLS signal (Figure 7B), this peak belongs to scattering objects with $R_h = 30$ nm. They are the core-shell micelles revealed with SAXS. For PEO-PAA 6-12 a peak at retention time around 25 min appeared at $R = 1$ in the refractive index signal, and it corresponded to nano-objects with a 43 nm hydrodynamic radius; they are the spherical micelles revealed by SAXS (Figure 3). It must be

Table 4. Comparison of the Behavior Observed for All PICs with the Different Methods of Characterization as a Function of R

PEO–PAA	scattered light intensity	size by DLS	FI-AsFIFFF	SAXS	^1H NMRD	^1H NMR
6-3	\uparrow for $R > 0.7$	$R > 0.7$	PICs detected at $R = 0.8$ low retention time peak $R < 1$	\uparrow for $R > 0.7$	evolution of PICs up to $R = 1$	abrupt change $\sim R = 0.7$
6-6.5	\uparrow for $R > 0.7$	\uparrow for $R > 0.7$	increases above $R = 0.8$	PICs detected at $R = 0.2$		
6-12	\uparrow for $R > 0.6$	\sim constant for $R \geq 0.4$	PICs detected at $R = 0.2$ low retention time peak $R < 1$	PICs detected at $R = 0.2$		
5-38	linear	\sim constant		PICs detected at $R = 0.2$	PICs detected at $R = 0.2$	

stressed that PEO–PAA 6-12 SAXS data showed the presence of PICs at 1:1 stoichiometry already at low R values: from 13% to 55% of the scattering signal at $R = 0.2$ and $R = 0.8$, respectively (Table SI.2.2). From FI-AsFIFFF, it looked like the complexes were concentrated enough to be detected by the refractive index signal only at $R = 1$. On the other hand, a very small DLS signal appeared already at $R = 0.2$, and it could be exploited from $R = 0.4$. It indicated objects with size increasing up to ratio 1. In particular, at $R = 0.6$ the DLS signal was a large peak corresponding to a mean $R_h = 28$ nm, and at $R = 0.8$ two populations with R_h equal to 22 and 38 nm appeared, while only the biggest one ($R_h = 43$ nm) was present at $R = 1$. The presence of different populations can account for the difficulties in describing the intermediate- q SAXS spectra as the sum of DGL3 molecules and nano-objects with 1:1 stoichiometry (Figure 3C).

DISCUSSION

PICs Size and Structure: Insights from the Different Experimental Techniques. Table 3 (and Table SI.4.2) compares the size values obtained from the different techniques. As for R_h determination, in most cases, a good agreement between FIFFF and DLS was observed. Regarding R_g values, those obtained from FIFFF are often higher than those obtained from SAXS data. The differences are within the experimental errors and the bigger discrepancies are found in the case of PEO–PAA 6-3, the smallest polymer, whose size is at the limit of light scattering resolution. Nevertheless, the trend of objects' sizes is the same whichever the technique was; i.e., the overall micellar size decreases while the size of the PAA, the core forming block, decreases. Correspondingly, the R_g/R_h ratio diminishes. R_g/R_h ratios for the PICs of this study fall well in the quite wide range, from 0.3 to 1.3, reported in the literature. These values have been attributed to core–shell morphology and the density of the core whose size is generally linked to the R_g .¹⁵ Indeed, inhomogeneity in mass density (a core much denser than the shell, for example) can lead to differences in R_g/R_h values which are difficult to predict. In our case, the more important differences changing the PAA size are found in the R_g values, they diminish more strongly than R_h : the lower R_g/R_h values in the case of PEO–PAA 6-3 based PICs indicate a dense core; instead, for PEO–PAA 5-38 based PICs, the higher R_g/R_h values suggest a more homogeneous structure. Such a trend is indeed in agreement with the angular analysis of the scattering intensity: the nature of the self-assemblies at $R = 1$ depends on the block copolymer composition. SAXS data could be modeled with a core–shell morphology only in the case of the PEO–PAA 6-3 and PEO–PAA 6-6.5, i.e. copolymers with the smaller PAA block. The size of the core extracted for PAA 6-3 and PEO–PAA -6.5 (5.3 and 7.8 nm, respectively) allows reasonable assumption that

more than one DGL3 is present in the core and seems dictated by the number of carboxylic groups on the same polymer chain which can interact with DGL3 molecules. Nevertheless, the absolute number of DGL3 is difficult to extract owing to the expected compaction of the objects core, like in the case of DNA vectors.

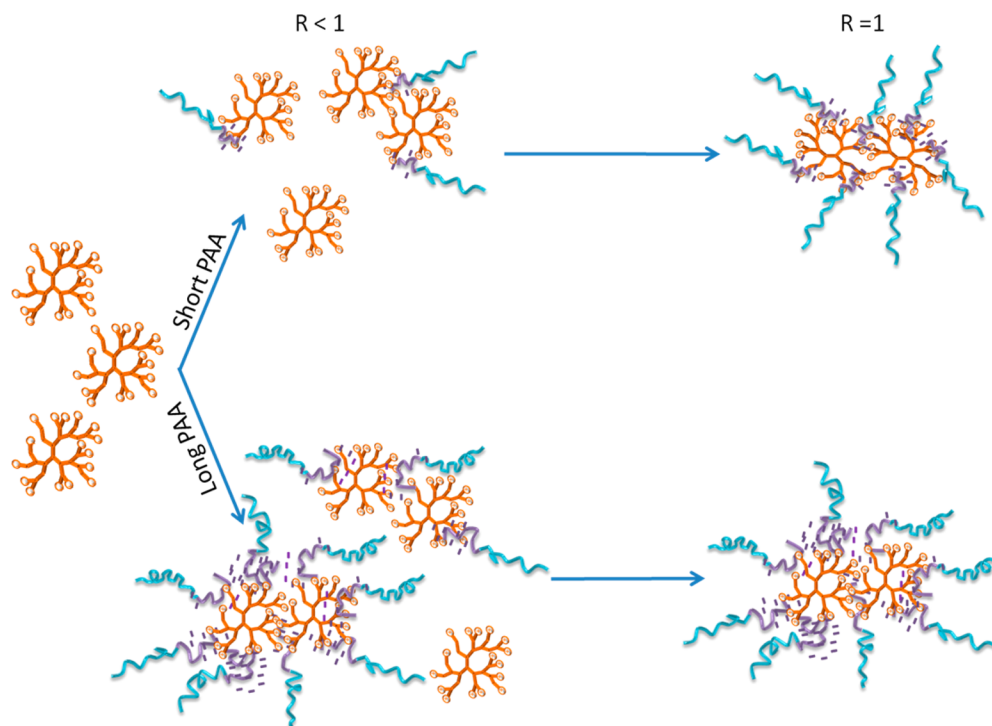
For the two other copolymers, PAA 6-12 and PEO–PAA 5-38, the number of AA groups in a polymer chain is higher than the number of amine groups on a DGL3 molecule; then more than one DGL3 can interact with the same polymer chain. As a consequence, a higher number of block copolymer is necessary in order to stabilize the coacervate complex, which leads to the formation of larger self-assemblies. The complex is probably more hydrated and less compact than in the case of the smaller copolymers; then the self-assemblies are better described as a homogeneous sphere. This is also confirmed by the cryo-EM images where no density gradient could be detected (Figure 4).

Processes of PICs Formation. Table 4 summarizes the observations of the different experimental techniques. In the case of PEO–PAA 6-3 the steep increase of the scattering intensity and the fast decrease of the ^1H NMR peak integrals above $R = 0.7$ indicate the threshold for PIC micelles formation. At this R value, scattering objects with well-defined size are observed with DLS, the SAXS curve can be analyzed with a core–shell sphere model, a large time retention peak is detected with FIFFF, and a third correlation time appears in the ^1H NMRD spectrum: the ensemble of the results indicates the formation of well-defined PICs objects. Below $R = 0.8$, no well-defined R_h value can be determined by DLS, and no retention peak associated with micelle formation could be detected by FIFFF. Nevertheless, SAXS and NMR give clear experimental evidence of the association of carboxylate and ammonium groups even below $R = 0.7$: R_g increases with R (Table 3), the K proton integrals diminish, and a second correlation time appears already at $R \leq 0.5$ (Table SI.3.1 in SI.3). The low retention time observed in the FIFFF fractograms brings the same conclusion: free DGL3 molecules or DGL3 interacting with few, highly hydrated PEO–PAA chains (whose charges are not completely balanced) exist before the appearance of 1:1 stoichiometry complexes at $R = 0.8$.

The observed self-assembling mechanism is then different from the closed association one usually observed for amphiphilic molecules with critical micellar concentration.

Referring to Table 4, for the other block copolymers, SAXS detects the presence of 1:1 stoichiometry PICs already at $R = 0.2$, while the size and concentration of the assemblies are high enough to be detected by DLS at this ratio only in the case of PEO–PAA 5-38. For this latter block copolymer, also ^1H NMRD detects the presence of species with the same correlation times as at $R = 1$. FI-AsFIFFF brings interesting information (Figure SI.4.2): the peak in the fractogram at $R = 1$

Scheme 2. Proposed Mechanism of PIC Formation in the Extreme Cases of Short or Long PAA^a



^aIn both cases at $R = 1$ PIC assemblies characterized by 1:1 stoichiometry are formed. Below $R = 1$ different species are simultaneously present in solution; only in the case of long PAA, PIC with 1:1 stoichiometry are already present in solution.

reveals the presence of two populations. SAXS cannot distinguish between 30% polydisperse spheres with mean R_g of 30 nm and the two populations evidenced by FI-AsFIFFF (R_g 26 and 46 nm). These two populations indicate a mixture of PICs with different $AA^-:NH_3^+$ ratios. The presence of different species may explain the lower agreement of the linear combination approach in the low- q region for $R < 1$ in SAXS spectra (Figure 3D).

In the case of PEO–PAA 6-12, the presence of charge-balanced PICs could be detected by DLS only from $R = 0.4$. This is not surprising since from the linear combination of SAXS curves it results that their contribution to scattering signal is still quite low, ca. 16%. As for FI-AsFIFFF, the peak in the fractogram is quite large at $R = 0.6$ and reveals the presence of two populations at $R = 0.8$. In both cases the linear combination of the scattering curve of DGL3 and of micelles at $R = 1$ does not perfectly describe the data in the intermediate q range (Figure 3C). This is indeed related to the presence of more PICs species detected through FI-AsFIFFF analysis of the solutions. Also in this case from SAXS only, it is difficult to distinguish between quite polydisperse micelles and two close populations when even free small species, i.e., DGL3, are present in solution ($R_g = 18.7$ nm with about 30% polydispersity according to SAXS and two populations with most probable values of R_g of 17 and 27 nm according to FI-AsFIFFF, as in the case of $R = 0.8$, Table 3). Therefore, FI-AsFIFFF helps to distinguish among the two circumstances.

In the case of PEO–PAA 6-6.5 at $R = 0.2$ big assemblies clearly exist according to SAXS results, but the linear combination approach does not fully describe the curve, which means that other species different from the charge balanced micelles and DGL3 molecules already form at this ratio. These species exist at very low concentration and their

presence can be detected by DLS only from $R = 0.7$ when their scattering contribution is near 50% (see Table SI.2.1).

Therefore, the complementarity of different techniques indicates a complex scenario where for all ratios smaller than 1, free DGL3 molecules (or DGL3 interacting with few, highly hydrated PEO–PAA chains) coexist with 1:1 stoichiometry micelles and slightly smaller micelles, whose charge is probably not balanced.

The concepts of open and closed association processes are somehow misleading and insufficient to describe such complex scenarios.

CONCLUSION

PICs self-assemblies are complex systems, highly responsive to the external medium conditions like ionic strength and pH, which are inherent to their electrostatic nature. Nevertheless, they are quite interesting and versatile systems easily prepared, directly in water. In this work, we have demonstrated that by combining several experimental approaches, valuable information explaining the formation process and the nature of these objects may be obtained. Self-assemblies with a homogeneous structure were formed when the PAA block size exceeded that of the PEO block, and more than one DGL3 can interact with one block copolymer chain. Instead, core–shell micelles with a dense core were formed when the number of AA units in the block copolymer was lower than or equal to the charged groups in DGL3. The mechanism of formation seems also to depend on the length of the PAA chain in the block copolymer.

Two representative cases of processes for the formation of the PICs (Scheme 2) can be highlighted: For PEO–PAA 6-3, loose interactions between AA^- and NH_3^+ moieties exist even for low R ratios. These interactions are not strong enough to lead to the formation of nano-objects. Well-defined PICs,

however, form as soon as a threshold of 0.7 is passed. In the final PIC, several DGL3 chains are reasonably present in the core as determined by SAXS. The other case is that of the other copolymers for which 1:1 stoichiometry objects are formed even for low *R* ratios and their concentration increases with increasing *R*. These charge-balanced PICs coexist with free DGL3 molecules and/or nano-objects with different stoichiometry (whose charge is not balanced), the size and the concentration of which increase with *R*.

This is reminiscent of “host–guest” soluble IPECs which form when the host chain is larger and in big excess with respect to the guest one (the case of PEO–PAA 6-3 at small *R*).³ On the contrary, insoluble IPEC forms when the concentration of the guest chain (the increase of *R* in the case of PEO–PAA 6-3) or its size (the case of PEO–PAA 5-38 at small *R*) is increased. In our case the presence of PEO chains guarantees colloidal stability so precipitation is avoided in all cases.

The previously described open or closed processes seem then inadequate to fully describe the mechanisms of formation in the case of charge-driven self-assembly, and the processes involved in traditional inter polyelectrolyte complexes have to be considered.

■ ASSOCIATED CONTENT

■ Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: [10.1021/acs.macromol.7b02391](https://doi.org/10.1021/acs.macromol.7b02391).

Characterization of commercial PEO–PAA copolymers, relaxometry results, FI-AsFIFFF (PDF)

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Notes

The authors declare no competing financial interest.

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Supplementary information for Mechanistic insights into Poly Ion Complex association.

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SI.0.1 Block-copolymer characterization

SI.1 DLS characterization

Figure SI.1.1 Typical correlogram functions for PICs at different acid/amine ratios

Figure SI.1.2 Evolution of PEO-PAA 6-3 /DGL 3 size measured by DLS, as a function of acid/amine ratio

Figure SI.1.3 Evolution of PEO-PAA 5-38 /DGL 3 size measured by DLS, as a function of acid/amine ratio

SI.2 SAXS Model. Theoretical Background

Figure SI.2.1. Small Angle X-ray scattering as a function of the scattering vector for DGL3 0.1% wt, PEO-PAA 6-3 0.09% wt (same concentration as in the sample DGL3/*PEO-PAA 6-3* at R=AA/NH₂ =1), DGL3/*PEO-PAA 6-3* solution with R=AA/NH₂ =0.5.

Table SI.2.1. Best values of α and β parameters in order to describe the experimental data with R < 1 as the contribution of micelles (completely formed at R=1) and free DGL3 present in the solutions.

SI.3. ¹H Nuclear Magnetic Relaxation Dispersion

Table SI.3.1 Best fit values obtained from the analysis of the ¹H NMRD profiles reported in Figure 7.

Figure SI.3.1. ¹H NMRD profiles for *PEO-PAA 6-3* 0.09% wt, *PEO-PAA 5-38* 0.04% wt and DGL3 0.1% wt.

SI.4 Field Flow Fractionation

Figure SI.4.1. Experimental set-up for Frit Inlet Asymmetrical FFF

Figure SI.4.2. Characterization of *PEO-PAA 5-38* / DGL3 PIC by FI-AsFIFFF at stoichiometric ratio 1/1

Figure SI.4.3. Characterization of *PEO-PAA 6-6.5*/ DGL3 PICs by FI-AsFIFFF (QELS signal) at different ratios. The ratios correspond to that of acrylic acid units over amine units.

Table SI.4.1. Analysis of FIFFF fractograms

Table SI.4.2 Comparison of PICs size obtained from different techniques in the case of *PEO-PAA 6-6.5* based PICs

SI.0.1 Block-copolymer characterization

Table SI.0.1 Characterization of commercial PEO-PAA copolymers

^a M_n was measured by fixing the PEO block at the value provided by the company, the molecular weight of the PAA block was calculated by comparing the integrals between PEO and PAA blocks.

Announced name from supplier	Announced EO weight fraction	Announced molecular weight	Measured M_w from SEC	Measured M_n from SEC	Measured M_n from ^1H NMR	$f_{\text{PEO exp}}$	Polymer name
PEO-PAA 5-56	0.08	61000	35000	34000	5000-38000	0.12	PEO-PAA 5-38
PEO-PAA 6-3	0.67	9000	8600	7600	6000-2900	0.67	PEO-PAA 6-3
PEO-PAA 6-6.5	0.48	12500	19000	14300	6000-6050	0.5	PEO-PAA 6-6.5
PEO-PAA 6-17.5	0.26	23500	78900	32200	6000-12000	0.33	PEO-PAA 6-12

SI.1 DLS characterization

The Malvern DLS data were analyzed by a custom-made program named STORMS in order to obtain a more precise characterization of the solutions¹. This program has been designed with Matlab and enables the fitting of DLS correlograms using different sets of parameters, corresponding to all hypotheses that have to be made during the treatment. Indeed, going from correlograms to size results implies three levels of hypotheses: the first one consisting in the transformation of autocorrelation data to diffusion coefficient, the second one extracting the size of the scattering object from diffusion coefficient depending on its geometry, and finally using a model enabling the transformation of the intensity-relative population to a number-relative one. The intensity distribution is naturally weighted according to the scattering intensity of each particle fraction or family. As such, the intensity distribution can be somewhat misleading, since a small amount of aggregation or presence of a larger particle species can dominate the distribution. For each step described above, STORMS provides the choice of different parameters.

For the nano-objects presented here, the protocol used a NNLS fitting, assumed a spherical shape for all objects, and the chosen scattering model was that corresponding to a full small sphere (Rayleigh model) for objects with a radius smaller than 100 nm or a sphere (Mie model) for the larger ones. The range of decay rates and the regularization parameter were systematically modified to check the consistency of the results. This treatment provided residuals lower than 5×10^{-3} for all analyses.

From intensity distribution to volume or number size distribution. The intensity distribution can be converted, using Mie theory, to a number distribution describing the relative proportion of multiple components in the sample based on their number rather than based on their scattering intensity.

Given the optical properties of the particle and the scattering angle, Mie theory estimates the scattering intensity $M(x)$ as a function of particle diameter x , dispersant and particle optical properties. The discrete list of Γ_i decay rate associated weighted by α_i could be transformed into a list of radii R_i (assuming spherical particles) through the equation $D_i = \frac{\Gamma_i}{q^2} = k_B T / (6\pi\eta R_i)$ weighted by the coefficient $\alpha_i / M(R_i)$.

Alternatively, conversion can be roughly obtained by assuming that $M(x)$ is proportional to R^6 (in the case of small homogeneous spheres) which is only correct for particle below ca 100

nm of diameter. Note that the Mie theory implies that a particular model has been chosen to describe the particles (homogeneous, spheres, hollow spheres, coated spheres...).

When transforming an intensity distribution to a number distribution, different assumptions are used: *all particles are homogeneous and spherical, the optical properties of the particles are known and intensity distribution is correct*. Moreover DLS technique itself produces distributions with inherent peak broadening, so there will always be some error in the representation of the intensity distribution. As such, number distributions derived from these intensity distributions emphasizes information obtained from a small fraction of the collected data. Therefore they are best used for comparative purposes, or for estimating the relative proportions where there are multiple modes, or peaks, and should never be considered as absolute.

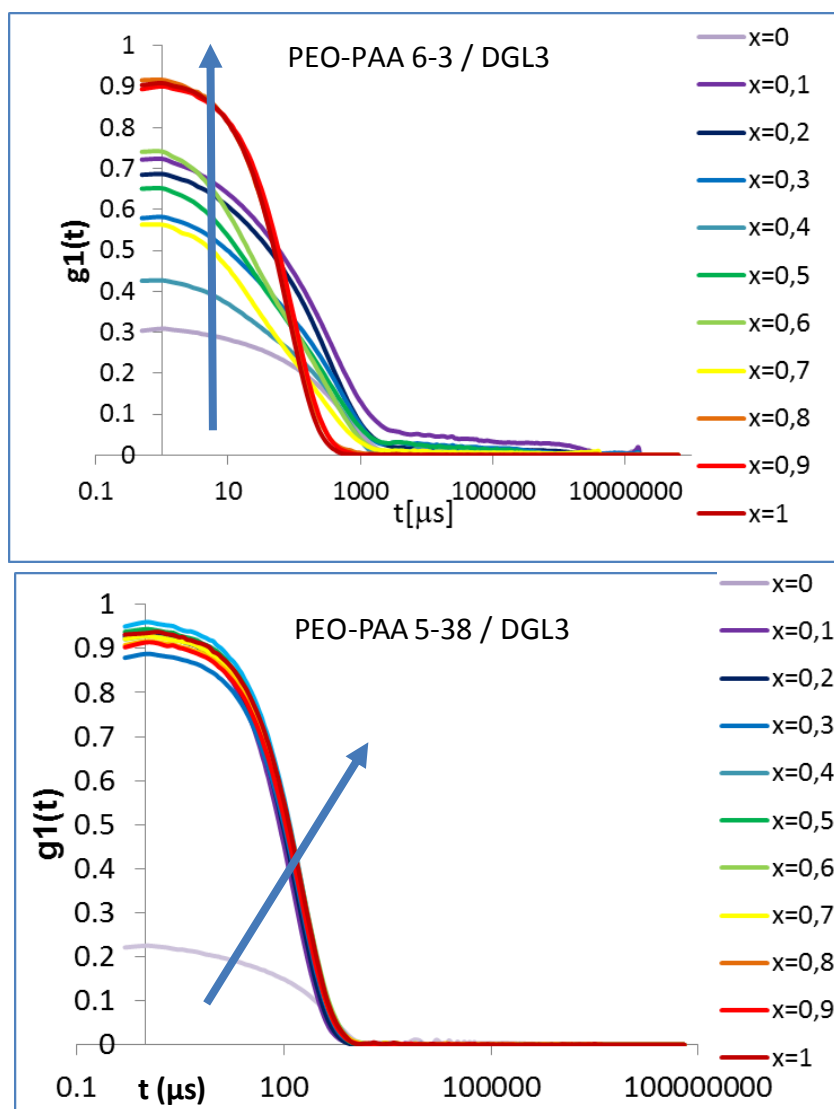


Figure SI.1.1. Typical correlogram functions for PICs at different PEO-PAA carboxylic groups and DGL primary amine groups (AA/NH₂, indicated as x in the graph).

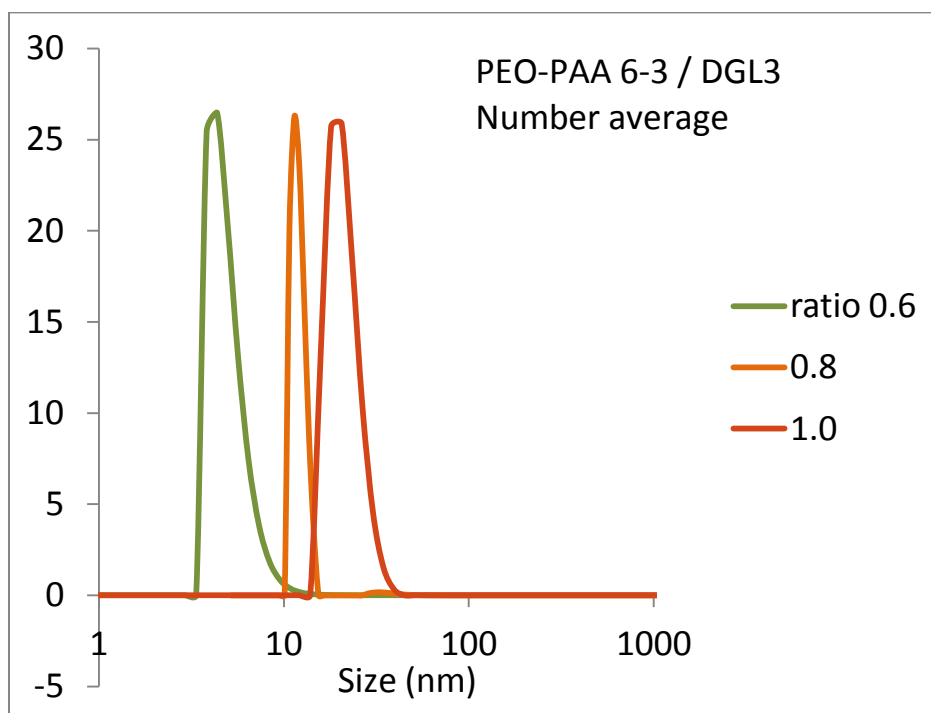
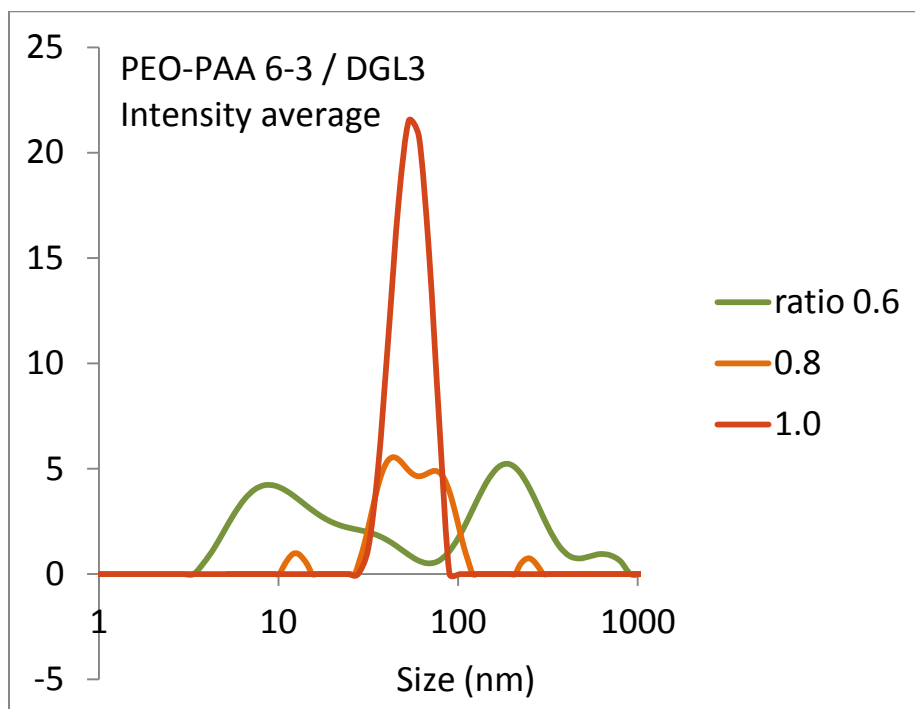


Figure SI.1.2. Evolution of *PEO-PAA 6-3* /DGL3 size measured by DLS, as a function of PEO-PAA carboxylic groups and DGL primary amine groups (AA/NH₂), R.

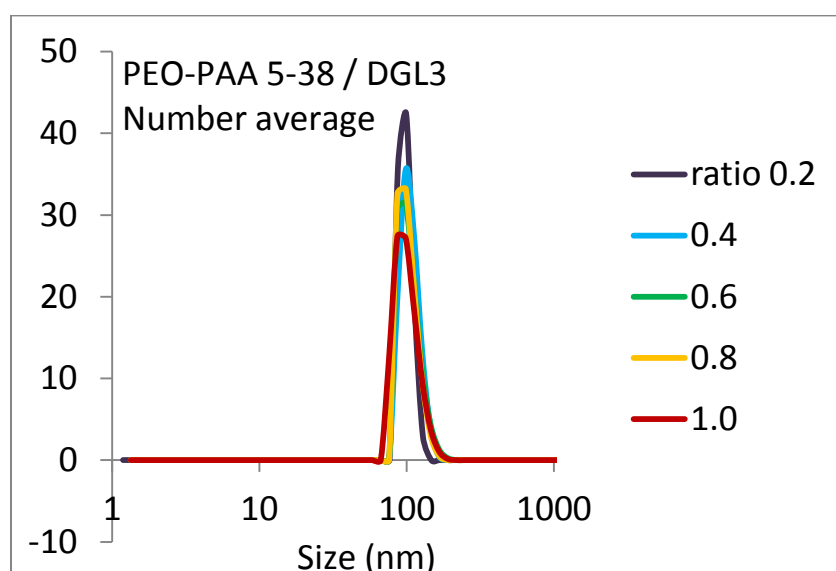
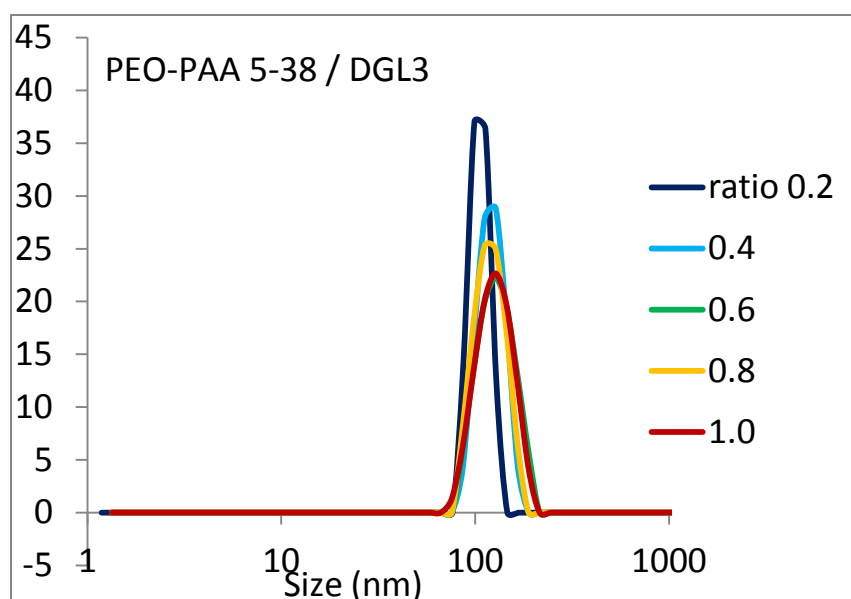


Figure SI.1.3. Evolution of *PEO-PAA 5-38* /DGL3 size measured by DLS, as a function of PEO-PAA carboxylic groups and DGL primary amine groups (AA/NH₂), R.

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SI.2 SAXS Model. Theoretical Background

For a suspension of uniform particles, the normalized scattered intensity is given by

$$I(q) = N P(q) S(q) \quad (1)$$

where N is the number density of scattering objects, $P(q)$ is the single particle scattering function describing the shape of the particles and $S(q)$ is the structure factor describing the inter-particle interactions. For a relatively dilute suspension, $S(q) \approx 1$ and $I(q)$ is governed by the shape of the scattering objects. Polydispersity in real systems require the introduction of a finite size distribution of the scattering objects. The resulting $I(q)$ in the non-interacting case is then given by:

$$I(q) = \int_0^\infty P(q, R) f(R) dR \quad (2)$$

Where the polydispersity can be approximately described by a Schulz size distribution function with corresponding $f(R)$:

$$f(R) = \left(\frac{Z+1}{R_m}\right)^{Z+1} \frac{R^Z}{\Gamma(Z+1)} \exp\left(-\frac{Z+1}{R_m} R\right) \quad (3)$$

with $Z = 1/(1-p^2)$, with p the polydispersity and R_m the mean size.

For spherical core-shell like objects, the single particle scattering function is expressed in terms of the scattering amplitudes of two spheres of radii (R_c) and ($R_c + t_{sh}$):

$$P_0(q, R_c, t_{sh}) = \left(\frac{16}{9} \pi^2\right) \{(\rho_{solv} - \rho_{sh})(R_c + t_{sh})^3 F_0(q(R_c + t_{sh})) - (\rho_{core} - \rho_{sh})R_c^3 F_0(qR_c)\}^2 \quad (4)$$

For spherical homogeneous objects, the single particle scattering function is

$$P_0(q, R) = \left(\frac{16}{9} \pi^2\right) \{(\rho_{core} - \rho_{solv})R^3 F_0(qR)\}^2 \quad (5)$$

$F_0(qR_i)$ in equations (4) and (5) indicates the scattering amplitude of a sphere of radii R_i

$$F_0(qR_i) = \frac{3[\sin(qR_i) - qR_i \cos(qR_i)]}{(qR_i)^3} \quad (6)$$

For block copolymer micelles, it is often necessary to consider an extra contribution at high- q vectors related to fluctuations originating from chain statistics and the inter-chain interactions of PEO chains inside the shell. The intra-corona contribution is well reproduced by the Beaucage function or alternatively by Ornstein-Zernike type term:

$$P(q) = P_0(q) + I_0 / (1 + (q^2 \xi^2)^{d/2}) \quad (7)$$

Where ξ accounts for the length scale of density fluctuations inside the corona (order of nm) and d is related to the fractal dimension of the scattering chains. We fixed d to -1.7 as expected for swollen polymer chains in good solvent. ξ is 0.22 nm in all fits.

Best fits of the form factor $P(q)$ for the measured data were derived using the model of spherical core-shell objects and a Ornstein-Zernike term as implemented in the SAXSutilities analysis package.

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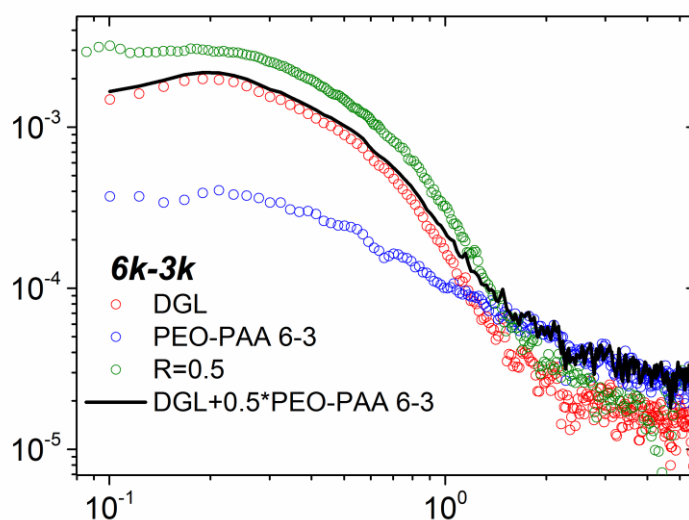


Figure SI.2.1. Small Angle X-ray scattering as a function of the scattering vector for DGL3 0.1%wt, PEO-PAA 6-3 0.09% wt (same concentration as in the sample DGL3/**PEO-PAA 6-3** at $R=AA/NH_2=1$), DGL3/**PEO-PAA 6-3** solution with $R=AA/NH_2=0.5$. The black curve is the linear combination of the experimental data relative to DGL3 plus the scattering contribution of **PEO-PAA 6-3** at the same concentration as in the sample at $R=0.5$.

Table SI.2.1. Best values of α and β parameters in order to describe the experimental data with $R < 1$ as the contribution of micelles (completely formed at $R=1$) and free DGL3 present in the solutions. $I(q)=\alpha*I(q)_{DGL3}+\beta*I(q)_{R=1}$

PEO-PAA	R	α	β
6-6.5	0.2	0.7	0.06
	0.8	0.42	0.58
6-12	0.2	0.87	0.13
	0.4	0.84	0.16
	0.6	0.70	0.30
	0.8	0.45	0.55
5-38	0.2	0.85	0.15
	0.4	0.7	0.3
	0.6	0.55	0.45

SI.3. ^1H Nuclear Magnetic Relaxation Dispersion

The field dependence of ^1H relaxation is described by the expression¹

$$R_1 = a + b \left(\frac{\tau_c}{1 + \omega^2 \tau_c^2} + \frac{4\tau_c}{1 + 4\omega^2 \tau_c^2} \right)$$

where ω is 2π times the proton Larmor frequency, b is a constant related to the dipole-dipole interaction energy, a is a non-dispersive term accounting for the protons diffusing around the complexes, and τ_c , called correlation time, is the characteristic time constant describing the stochastic fluctuations of the ^1H - ^1H dipole-dipole interactions. When internal fast motions are present, they may be unable to average completely the dipole-dipole interaction, so that multiple motional processes, occurring on different time scales, can contribute to R_1 . In the model-free formalism², the resulting relaxation rates are determined by the weighted sum of the contributions corresponding to the different correlation times, according to the following equation³

$$R_1 = a + b \sum_i c_i (J(\omega, \tau_i) + 4J(2\omega, \tau_i)) \quad (1)$$

where c_i represents the molar fraction of protons with correlation time τ_i and $J(\omega, \tau_i)$ is the Lorentzian spectral density function

$$J(\omega, \tau_i) = \frac{\tau_i}{1 + \omega^2 \tau_i^2} \quad (2)$$

Table SI.3.1 Best fit values obtained from the analysis of the ^1H NMRD profiles reported in Figure 6.

	c_1 (%)	τ_1 (ns)	c_2 (%)	τ_2 (ns)	τ_3 (ns)
DGL3	-	-	-	-	3.0 ± 0.4
PEO-PAA 5-38	-	-	-	-	3.0 ± 0.4
PEO-PAA 6-3	-	-	-	-	3.0 ± 0.4
PEO-PAA 5-38:DGL3 R=0.2	0.04 ± 0.01	2250 ± 140	0.3 ± 0.1	120 ± 20	3.0 ± 0.4
PEO-PAA 5-38:DGL3 R=1	0.14 ± 0.01	2250 ± 140	0.9 ± 0.1	120 ± 20	3.0 ± 0.4
PEO-PAA 6-3:DGL3 R=0.2	-	-	0.7 ± 0.1	68 ± 30	3.0 ± 0.4
PEO-PAA 6-3:DGL3 R=0.5	-	-	0.9 ± 0.4	60 ± 20	3.0 ± 0.4
PEO-PAA 6-3:DGL3 R=1	0.14 ± 0.04	1150 ± 210	0.8 ± 0.2	130 ± 30	3.0 ± 0.4

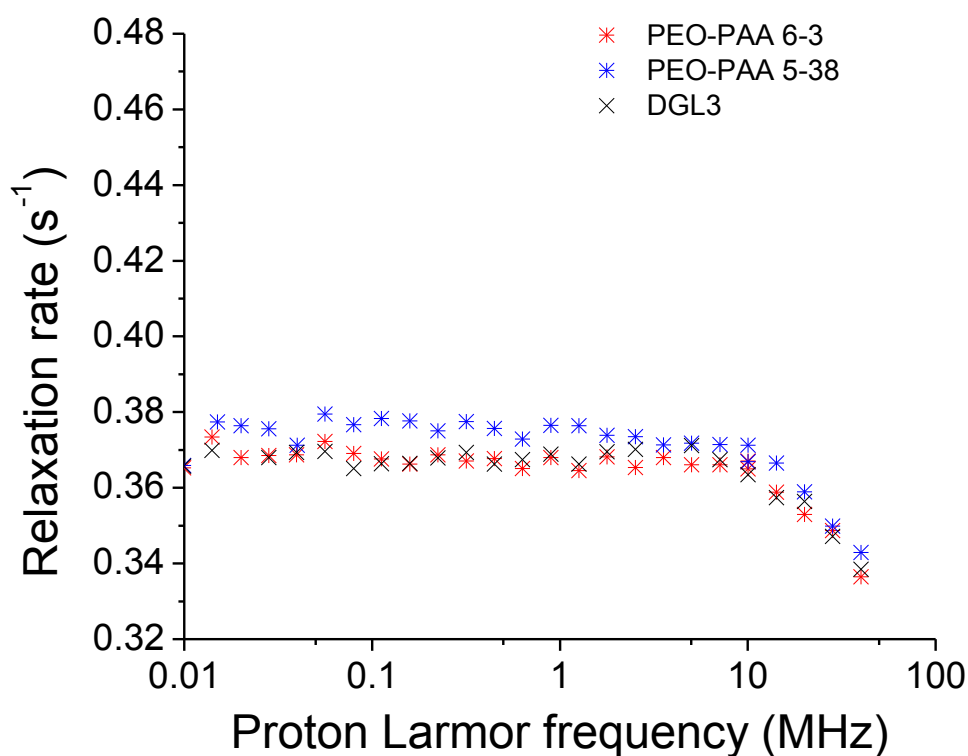


Figure SI.3.1. ^1H NMRD profiles for *PEO-PAA 6-3* 0.09wt%, *PEO-PAA 5-38* 0.04 wt% and DGL3 0.1wt%.

References

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3. Bertini, Fragai, Luchinat, Parigi, *Magn. Reson. Chem.*, **2000**, 38, 543; Ravera, Parigi, Mainz, Religa, Reif, Luchinat, *J. Phys. Chem. B* **2013**, 117, 3548; Ravera, Fragai, Parigi, Luchinat, *ChemPhysChem*, **2015**, 16, 2803.

SI.4 Field Flow Fractionation

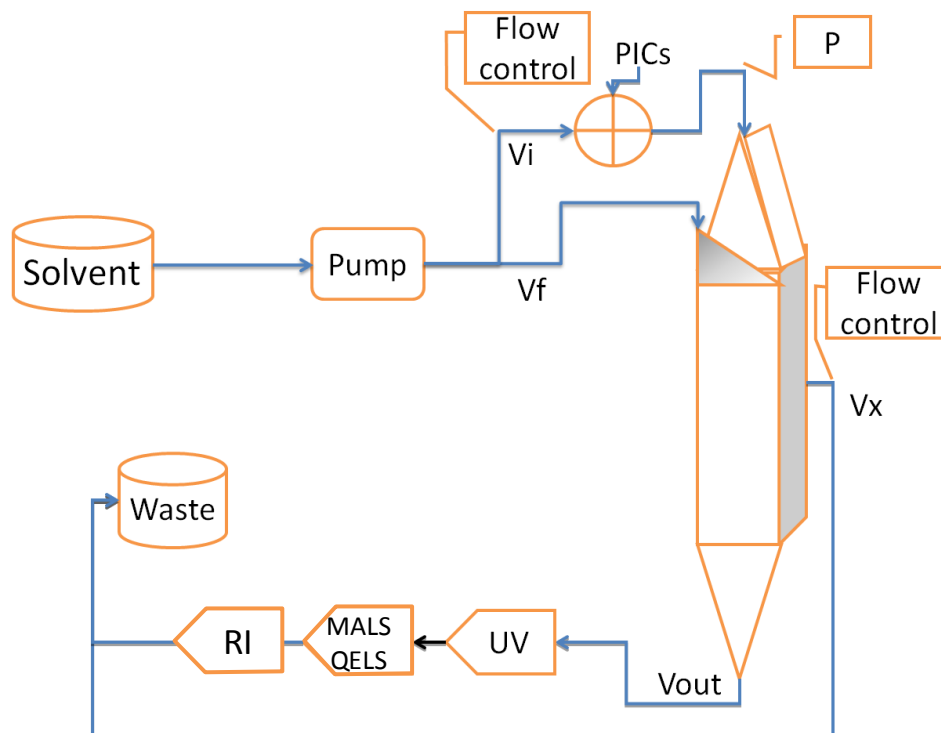


Figure SI.4.1. Experimental set-up for Frit Inlet Asymmetrical FFF

Field Flow Fractionation for *PEO-PAA 5-38* PICs for other ratios than 1 implied specific conditions that showed poor reproducibility. These data could not therefore be used in this study.

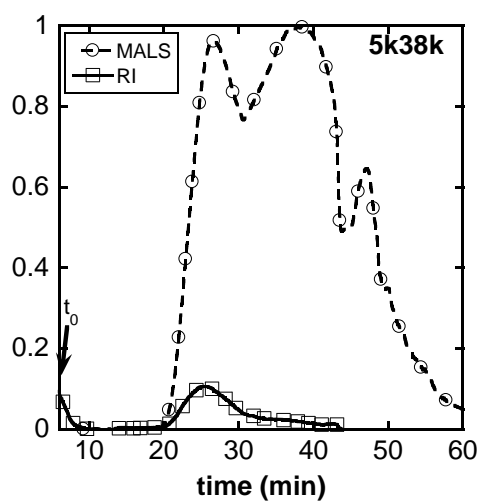


Figure SI.4.2. Characterization of *PEO-PAA 5-38* / DGL3 PIC by FI-AsFIFFF at stoichiometric ratio 1/1

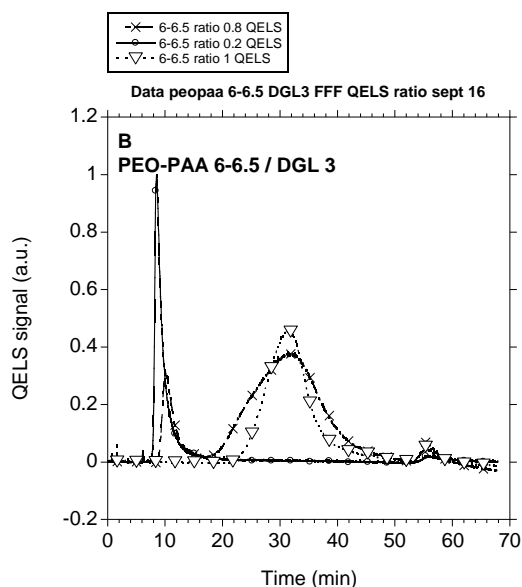


Figure SI.4.3. Characterization of *PEO-PAA 6-6.5* / DGL3 PICs by FI-AsFIFFF (QELS signal) at different ratios. The ratios correspond to that of acrylic acid units over amine units. o ratio 0.2, x ratio 0.8, ∇ ratio 1.0. PEO-PAA 6-6.5

Table SI.4.1. Analysis of FIFFF fractogramms ^a numbers in brackets indicate the accuracy of determination. ^b n number average; w weight average; z stands for z-average

Polymer	Ratio	elution time (min)	R _h (nm) ^a	R _g (nm) ^a	Average type ^b
PEO-PAA 5-38	1.0	25/38	38 (2 %) / 53 (2 %)	26 (8 %) / 46 (3 %)	(n)
			39 (2%) / 55 (2 %)	28 (7%) / 49 (3%)	(w)
			40 (2%) / 58 (2%)	30 (6%) / 53 (4 %)	(z)
PEO-PAA 6-12	0.4	20	38 (20%)	52 (15%)	(n)
			32 (14%)	37 (15/%)	(w)
			26 (10%)	21 (15%)	(z)
	0.6	20	28 (5%)	16 (25%)	(n)
			27 (5%)	15 (19%)	(w)
			28 (4%)	14 (16%)	(z)
	0.8	25	20 (4%) / 41 (4%)	17 (40%)/ 27 (9%)	(n)
			21 (3%) / 39 (3%)	17 (35%) / 26 (9%)	(w)
			22 (2%) / 38 (2%)	18 (30%) / 25 (9%)	(z)
	1.0	25	35 (2%)	25 (10%)	(n)
			39 (2%)	30 (7%)	(w)
			43 (2%)	35 (6%)	(z)
PEO-PAA 6-3	0.8	26	18 (8%)	12 (11%)	(n)
			18 (6%)	12 (36%)	(w)
			19 (6%)	13 (34%)	(z)
	1.0	30	27 (4%)	17 (12%)	(n)
			28 (4%)	17 (11%)	(w)
			30 (4%)	18 (12%)	(z)

Table SI.4.2 Comparison of PICs size obtained from different techniques

PEO-PAA Polymer	Ratio	R _g SAXS (nm)	R _g (z) FFF (nm)	R _h (z) FFF (nm)	R _g /R _h (z)	R _h number DLS (nm)	R _h Int DLS (nm)	Cumulant PDI	R _g SAXS / R _h DLS
6-6.5	0.2	9.2				-	-		
	0.8	12	39/20*	63/30*	0.62/0.66	20	33	0.22	0.36
	1.0	13.8	17	26	0.65	21	29	0.13	0.48