Design of finite volume schemes minimizing the grid orientation effect in reservoir simulation

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### The model problem

Dead-oil porous media flow with (a lot of) simplifying assumptions

$$\boldsymbol{u} + \lambda(s)\boldsymbol{\nabla} p = \boldsymbol{0},\tag{1a}$$

$$\operatorname{div} \boldsymbol{u} = q(t)\delta_{\boldsymbol{0}}(\boldsymbol{x}), \tag{1b}$$

$$\phi \,\partial_t s + \operatorname{div}(f(s)\boldsymbol{u}) = q(t)\delta_{\boldsymbol{0}}(\boldsymbol{x}),\tag{1c}$$

where Dirac source terms mimick a water injection at x = 0.
Fractional flux formulation using Bucklet-Leverett's law

$$\begin{split} f(s) &= \frac{\mu_w^{-1} s^2}{\mu_w^{-1} s^2 + \mu_o^{-1} (1-s)^2} = \frac{M s^2}{M s^2 + (1-s)^2},\\ \lambda(s) &= \mu_w^{-1} s^2 + \mu_o^{-1} (1-s)^2 \qquad \text{[total mobility]},\\ M &= \mu_o / \mu_w \qquad \text{[viscosity ratio]}. \end{split}$$

#### The IMPES semi-discretization in time

The time stepping is *implicit* w.r.t. pressure, *explicit* w.r.t. saturation

$$\boldsymbol{u}^{n+1} + \lambda(s^n) \boldsymbol{\nabla} p^{n+1} = \boldsymbol{0}, \tag{2a}$$

$$\operatorname{div} \boldsymbol{u}^{n+1} = q^{n+1} \delta_{\boldsymbol{0}}(\boldsymbol{x}), \qquad (2\mathsf{b})$$

$$\phi \, \frac{s^{n+1} - s^n}{\Delta t} + \operatorname{div}(f(s^n)\boldsymbol{u}^{n+1}) = q^{n+1}\delta_{\boldsymbol{0}}(\boldsymbol{x}). \tag{2c}$$

Decoupling into an elliptic problem in pressure (2a)–(2b)

$$-\operatorname{div}(\lambda(s^n)\boldsymbol{\nabla}p^{n+1}) = q^{n+1}\delta_{\mathbf{0}}(\boldsymbol{x}),$$

followed by a hyperbolic transport in saturation (2c).

#### The 5-point discretization in space

**The domain**  $\Omega$  is divided into rectangles  $\Omega_{ij}$  of *uniform* size



The unknowns at the cell centers (i, j) are  $s_{i,j}^{n+1}$  and  $p_{i,j}^{n+1}$ , while those at the edges  $\sigma = (i + 1/2, j)$  and (i, j + 1/2) are fluxes

$$F_{\sigma} \approx \pm \int_{\sigma} u^{n+1} \cdot n_{\sigma}$$

#### Solving for pressure

Equation  $u^{n+1} + \lambda(s^n) \nabla p^{n+1} = 0$  is discretized as

$$F_{i+1/2,j} = \lambda_{i+1/2,j}^n \frac{\Delta y}{\Delta x} (p_{i,j}^{n+1} - p_{i+1,j}^{n+1}),$$
(3a)

$$F_{i,j+1/2} = \lambda_{i,j+1/2}^{n} \frac{\Delta x}{\Delta y} \left( p_{i,j}^{n+1} - p_{i,j+1}^{n+1} \right), \tag{3b}$$

where  $\lambda_{\sigma}^{n}$  stems from one's favorite averaging procedure. Equation div  $u^{n+1} = q^{n+1}\delta_{0}$  is discretized as

$$F_{i+1/2,j} - F_{i-1/2,j} + F_{i,j+1/2} - F_{i,j-1/2} = q^{n+1} \kappa_{0,0}(i,j).$$
(4)

Plugging (3) into the flux balances (4) results in a linear system.

### Solving for saturation

• Equation 
$$\phi \frac{s^{n+1} - s^n}{\Delta t} + \operatorname{div}(f(s^n)\boldsymbol{u}^{n+1}) = q^{n+1}\delta_0$$
 is discretized as

$$\phi \frac{\Delta x \Delta y}{\Delta t} (s_{i,j}^{n+1} - s_{i,j}^n) + (f(s)F)_{i+1/2,j} - (f(s)F)_{i-1/2,j}$$

$$+ (f(s)F)_{i,j+1/2} - (f(s)F)_{i,j-1/2} = q^{n+1} \kappa_{0,0}(i,j)$$
(5)

The fractional fluxes are upwinded by

$$(f(s)F)_{i+1/2,j} = f(s_{i,j}^n) [F_{i+1/2,j}]^+ + f(s_{i+1,j}^n) [F_{i+1/2,j}]^-$$
$$(f(s)F)_{i,j+1/2} = f(s_{i,j}^n) [F_{i,j+1/2}]^+ + f(s_{i,j+1}^n) [F_{i,j+1/2}]^-$$

where  $[\cdot]^+ = \max(\cdot, 0)$  and  $[\cdot]^- = \min(\cdot, 0)$ .

# The grid orientation effect

# The computed solution propagates faster along the axes of the grid.



M = 200

Water\_saturation 0,00 0,25 0,5 0,75 1,00

For  $M \gg 1$ , mathematical instabilities intrinsic to the model will amplify this numerical error.

### **Objectives**

# The GOE issue is annoying for predicting production at wells.



# Alleviate the GOE. *Make the front round again* in an "cheap" way.

The 9-point scheme: Principle of construction (1)

YANOSIK-MCKRACKEN (1979), COATS (1983), POTEMPA (1985), DING (1990) studied a class of schemes involving 4 diagonal cells



The two diagonal fluxes  $G_{i+1/2,j+1/2}^{\theta}$  and  $H_{i+1/2,j+1/2}^{\theta}$  at vertex (i+1/2, j+1/2) do not correspond to any actual outgoing flux  $\pm \int_{\sigma} \boldsymbol{u} \cdot \boldsymbol{n}$  for any physical "edge"  $\sigma$ .

The 9-point scheme: Principle of construction (2)

To further decrease anisotropy, use two parameters instead of one. This is possible thanks to a "smarter" definition of the diagonal fluxes.



• Every diagonal flux is a linear combination of horizontal and vertical fluxes arising from the 5-point scheme.

### The 9-point discretization in space

### Horizontal and vertical fluxes

$$\widetilde{F}_{i+1/2,j}^{\theta} = (1 - 4\theta_x)F_{i+1/2,j}$$
(6a)
$$\widetilde{F}_{i,j+1/2}^{\theta} = (1 - 4\theta_y)F_{i,j+1/2}$$
(6b)

# Diagonal fluxes

$$\widetilde{G}_{i+1/2,j+1/2}^{\theta} = \theta_y F_{i,j+1/2} + \theta_x F_{i+1/2,j+1} + \theta_x F_{i+1/2,j} + \theta_y F_{i+1,j+1/2}$$
  
$$\widetilde{H}_{i-1/2,j+1/2}^{\theta} = \theta_y F_{i,j+1/2} - \theta_x F_{i-1/2,j+1} - \theta_x F_{i-1/2,j} + \theta_y F_{i-1,j+1/2}$$
(6c)

• The free parameters  $\theta = (\theta_x, \theta_y) \in [0, 1/4]^2$  are to be tuned later.

# Solving for pressure

Equation  $-\operatorname{div}(\lambda(s^n) \nabla p^{n+1}) = q^{n+1} \delta_0$  is discretized as

$$\widetilde{F}_{i+1/2,j}^{\theta} - \widetilde{F}_{i-1/2,j}^{\theta} + \widetilde{F}_{i,j+1/2}^{\theta} - \widetilde{F}_{i,j-1/2}^{\theta} + \widetilde{G}_{i+1/2,j+1/2}^{\theta} - \widetilde{G}_{i-1/2,j-1/2}^{\theta} + \widetilde{H}_{i-1/2,j+1/2}^{\theta} - \widetilde{H}_{i+1/2,j-1/2}^{\theta} = q^{n+1} \kappa_{0,0}(i,j).$$
(7)

Plugging (6) into the flux balances (7) yields a linear system.

#### Solving for saturation

• Equation 
$$\phi \frac{s^{n+1} - s^n}{\Delta t} + \operatorname{div}(f(s^n)u^{n+1}) = q^{n+1}\delta_0$$
 is discretized as

$$\phi \frac{\Delta x \Delta y}{\Delta t} (s_{i,j}^{n+1} - s_{i,j}^{n}) + (f(s)\widetilde{F})_{i+1/2,j}^{\theta} - (f(s)\widetilde{F})_{i-1/2,j}^{\theta}$$

$$+ (f(s)\widetilde{F})_{i,j+1/2}^{\theta} - (f(s)\widetilde{F})_{i,j-1/2}^{\theta}$$

$$+ (f(s)\widetilde{G})_{i+1/2,j+1/2}^{\theta} - (f(s)\widetilde{G})_{i-1/2,j-1/2}^{\theta}$$

$$+ (f(s)\widetilde{H})_{i-1/2,j+1/2}^{\theta} - (f(s)\widetilde{H})_{i+1/2,j-1/2}^{\theta} = q^{n+1}\kappa_{0,0}(i,j)$$
(8)

The fractional fluxes are upwinded by

$$\begin{split} (f(s)\widetilde{G})^{\theta}_{i+1/2,j+1/2} &= f(s^n_{i,j}) \, [\widetilde{G}^{\theta}_{i+1/2,j+1/2}]^+ + f(s^n_{i+1,j+1}) \, [\widetilde{G}^{\theta}_{i+1/2,j+1/2}]^- \\ (f(s)\widetilde{H})^{\theta}_{i-1/2,j+1/2} &= f(s^n_{i,j}) \, [\widetilde{H}^{\theta}_{i-1/2,j+1/2}]^+ + f(s^n_{i-1,j+1}) \, [\widetilde{H}^{\theta}_{i-1/2,j+1/2}]^- \end{split}$$

# **Principle of optimization**

- 1. Quantify the anisotropy of the numerical error along each direction, by means of Fourier analysis or equivalent modified equation
- 2. Introduce an ideal behavior of this angular error, declared to be the "least anisotropic" one.
- 3. Minimize the total discrepancy over all direction between the angular error of the scheme and that of the ideal one.
  - Do not seek to achieve higher-order accuracy, but simply to alter the distribution of angular error.

#### Equivalent modified equation

- Neglect the effect of averaging  $\lambda_{\sigma}^n$ , focus on the limit  $\Delta t \to 0$ , carry out Taylor expansions at  $(i\Delta x, j\Delta y)$ , switch to polar coordinates  $(r, \varphi)$ .
- The saturation transport scheme is a second-order approximation of

$$\phi\partial_t s + \operatorname{div}(f(s)\boldsymbol{u}) = \left(r^{-1}\partial_r(r\cdot), r^{-1}\partial_\varphi\right) \left\{ f'(s)\widetilde{\mathbf{D}}^{\boldsymbol{\theta}}_{\Delta x, \Delta y}(\varphi) \begin{pmatrix} \partial_r \\ r^{-1}\partial_\varphi \end{pmatrix} s \right\},\$$

where the diffusion matrix in polar coordinates reads

$$\widetilde{\mathbf{D}}_{\Delta x,\Delta y}^{\boldsymbol{\theta}}(\varphi) = \begin{bmatrix} \widetilde{A}_{\Delta x,\Delta y}^{\boldsymbol{\theta}}(\varphi) & \widetilde{B}_{\Delta x,\Delta y}^{\boldsymbol{\theta}}(\varphi) \\ \widetilde{B}_{\Delta x,\Delta y}^{\boldsymbol{\theta}}(\varphi) & \widetilde{C}_{\Delta x,\Delta y}^{\boldsymbol{\theta}}(\varphi) \end{bmatrix}.$$

The first diagonal entry  $\widetilde{A}^{\theta}_{\Delta x, \Delta y}(\varphi)$  reflects the longitudinal diffusion in the direction of the current location.

#### Properties of the angular error

Let φ<sup>θ</sup><sub>•</sub> = arctan(θ<sub>x</sub>Δy/θ<sub>y</sub>Δx) be the transition angle, where the stencil for saturation changes due to upwinding.
 If 0 ≤ φ ≤ φ<sup>θ</sup><sub>•</sub>,

$$\widetilde{A}^{\boldsymbol{\theta}}_{\Delta x,\Delta y}(\varphi) = \frac{\Delta x}{2} \cos^3 \varphi + \frac{\Delta y}{2} (1 - 4\theta_y) \sin^3 \varphi + 2\Delta y \left( 2\theta_y \frac{\Delta x}{\Delta y} + \theta_x \frac{\Delta y}{\Delta x} \right) \cos \varphi \sin^2 \varphi$$

If  $\varphi_{\bullet}^{\theta} \leq \varphi \leq \pi/2$ ,

$$\widetilde{A}^{\boldsymbol{\theta}}_{\Delta x,\Delta y}(\varphi) = \frac{\Delta y}{2} \sin^3 \varphi + \frac{\Delta x}{2} (1 - 4\theta_x) \cos^3 \varphi + 2\Delta x \left( 2\theta_x \frac{\Delta y}{\Delta x} + \theta_y \frac{\Delta x}{\Delta y} \right) \sin \varphi \cos^2 \varphi$$

#### Least anistropic behavior

# For all $\boldsymbol{\theta} \in [0, 1/4]^2$ ,

- $\widetilde{A}^{\boldsymbol{\theta}}_{\overset{\sim}{\sim} x, \Delta y}$  is a continuous, piecewise differentiable function of  $\varphi \in [0, \pi/2]$ .
- $\widetilde{A}^{\theta}_{\Delta x, \Delta y}(\varphi)$  is of first-order in  $\Delta x$  and  $\Delta y$ .
- $\widetilde{A}^{\boldsymbol{\theta}}_{\Delta x,\Delta y}(0) = \Delta x/2$  and  $\widetilde{A}^{\boldsymbol{\theta}}_{\Delta x,\Delta y}(\pi/2) = \Delta y/2$  do *not* depend on  $\boldsymbol{\theta}$ .
- **The ideal behavior is declared to be affine in \sin^2 \varphi, i.e.,**

$$A^*_{\Delta x,\Delta y}(\varphi) = \frac{\Delta x}{2} \mathrm{cos}^2 \varphi + \frac{\Delta y}{2} \mathrm{sin}^2 \varphi$$

#### Minimize the total discrepancy

Exact minimization: find

$$\boldsymbol{\theta}^* = \arg\min_{\boldsymbol{\theta} \in [0, 1/4]^2} \int_0^{\pi/2} |\widetilde{A}^{\boldsymbol{\theta}}_{\Delta x, \Delta y}(\varphi) - A^*_{\Delta x, \Delta y}(\varphi)|^2 \ 2 \sin\varphi \cos\varphi \ \mathrm{d}\varphi.$$

Optimal but slighlty tedious to compute.

Approximate minimization: find  ${m heta}^{\flat}=( heta_x^{\flat}, heta_y^{\flat})\in [0,1/4]^2$  such that

$$\begin{split} \varphi_{\bullet}^{\theta^{\flat}} & \text{is as close to } \pi/4 \text{ as possible}, \\ \widetilde{A}_{\Delta x,\Delta y}^{\theta^{\flat}}(\varphi_{\bullet}^{\theta^{\flat}}) = A_{\Delta x,\Delta y}^{*}(\varphi_{\bullet}^{\theta^{\flat}}). \end{split}$$

The closer  $\varphi_{\bullet}^{\theta^{\flat}}$  is to  $\pi/4$ , the better.

For  $2/7 \le \Delta y / \Delta x \le 7/2$ , it is possible to prescribe  $\varphi_{\bullet}^{\theta^{\flat}} = \pi/4$  without violating the constraint  $\theta^{\flat} \in [0, 1/4]^2$ .

# **Choice of suboptimal parameters**

Let 
$$z = \frac{\Delta y}{\Delta x}$$
. Set  

$$\omega(z) = \begin{cases} \frac{7}{2}z & \text{if } z \leq \frac{2}{7}, \\ 1 & \text{if } \frac{2}{7} \leq z \leq \frac{7}{2}, \\ \frac{2}{7}z & \text{if } z \geq \frac{7}{2} \end{cases}$$



Then, the recommended values for  ${m heta}^{lat}$  are

$$\theta_x^{\flat} = \frac{(1+\omega^2(z)z)\sqrt{1+\omega^2(z)} - (1+\omega^3(z)z)}{8\omega(z)z}$$
$$\theta_y^{\flat} = \frac{(1+\omega^2(z)z)\sqrt{1+\omega^2(z)} - (1+\omega^3(z)z)}{8\omega^2(z)}$$

5-point vs. 9-point (1)

 $\Delta y/\Delta x = 1, M = 200$ 





0,00 0,25 0,5 0,75 1,

5-point vs. 9-point (2)

 $\Delta y/\Delta x = 1/3, M = 200$ 



5-point vs. 9-point (3)

 $\Delta y/\Delta x = 1/5, M = 200$ 



# Conclusion

- Revisited an old scheme with new insights and proposed a quantitative approach by clarifying the notion of least anistropic behavior.
- Incorporated more sophisticated physical effects such as heterogeneous permeabilities, capillary pressure and gravity.
- Extended this optimization strategy to other transport schemes, in particular that of KEILEGAVLEN-KOZDON-MALLISON (2012) where the tuning degree of freedom is a function.