



HAL
open science

FLOCKING: PHASE TRANSITION AND ASYMPTOTIC BEHAVIOUR

Xingyu Li

► **To cite this version:**

Xingyu Li. FLOCKING: PHASE TRANSITION AND ASYMPTOTIC BEHAVIOUR. 2019. hal-02143985v1

HAL Id: hal-02143985

<https://hal.science/hal-02143985v1>

Preprint submitted on 14 Jun 2019 (v1), last revised 10 May 2024 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

FLOCKING: PHASE TRANSITION AND ASYMPTOTIC BEHAVIOUR

XINGYU LI*

Abstract. This paper is devoted to a continuous Cucker-Smale model with noise, which has isotropic and polarized stationary solutions depending on the intensity of the noise. The first result establishes the threshold value of the noise parameter which drives the phase transition. This threshold value is used to classify all stationary solutions and their linear stability properties. Using an entropy, these stability properties are extended to the non-linear regime. The second result is concerned with the asymptotic behaviour of the solutions of the evolution problem. In several cases, we prove that stable solutions attract the other solutions with an optimal exponential rate of convergence determined by the spectral gap of the linearized problem around the stable solutions. The spectral gap has to be computed in a norm adapted to the non-local term.

Keywords. Flocking model; phase transition; symmetry breaking; stability; large time asymptotics; free energy; spectral gap; asymptotic rate of convergence

AMS subject classifications. 35B40; 35P15; 35Q92.

1. Introduction

In many fields such as biology, ecology or economic studies, emerging collective behaviours and self-organization in multiagent interactions have attracted the attention of many researchers. In this paper we consider the Cucker-Smale model in order to describe *flocking*. The original model of [8] describes a population of N birds moving in \mathbb{R}^3 by the equations

$$v_i(t_n + \Delta t) - v_i(t_n) = \frac{\lambda \Delta t}{N} \sum_{j=1}^N a_{ij}(v_j(t_n) - v_i(t_n)), \quad i = 1, 2, \dots, N$$

at discrete times $t_n = n\Delta t$ with $n \in \mathbb{N}$ and $\Delta t > 0$. Here v_i is the velocity of the i th bird, the model is homogeneous in the sense that there is no position variable, and the coefficients a_{ij} model the interaction between pairs of birds as a function of their relative velocities, while λ is an overall coupling parameter. The authors proved that under certain conditions on the parameters, the solution converges to a state in which all birds fly with the same velocity. Another model is the Vicsek model [13] which was derived earlier to study the evolution of a population in which individuals have a given speed but the direction of their velocity evolves according to a diffusion equation with a local alignment term. This model exhibits phase transitions. In [9–12], phase transition has been shown in a continuous version of the model: with high noise, the system is disordered and the average velocity is zero, while for low noise a direction is selected.

Here we consider a model on \mathbb{R}^d , $d \geq 1$ with noise as in [3, 7]. The population is described by a distribution function $f(v, t)$ in which the interaction occurs through a mean-field nonlinearity known as *local velocity consensus* and we also equip the individuals with a so-called *self-propulsion* mechanism which privileges a speed (without a privileged direction) but does not impose a single value to the speed as in the Vicsek model. The

* CEREMADE (CNRS UMR n° 7534), PSL university, Université Paris-Dauphine, Place de Lattre de Tassigny, 75775 Paris 16, France. *E-mail:* li@ceremade.dauphine.fr

distribution function solves

$$\frac{\partial f}{\partial t} = D\Delta f + \nabla \cdot \left((v - \mathbf{u}_f) f + \alpha v (|v|^2 - 1) f \right), \quad f(\cdot, 0) = f_{\text{in}} > 0 \quad (1.1)$$

where $t \geq 0$ denotes the time variable and $v \in \mathbb{R}^d$ is the velocity variable. Here ∇ and Δ are the gradient and the Laplacian with respect to v respectively. The parameter $D > 0$ measures the intensity of the noise, $\alpha > 0$ is the parameter of self-propulsion which tends to force the distribution to be centered on velocities $|v|$ of the order of 1 when α becomes large, and

$$\mathbf{u}_f(t) = \frac{\int_{\mathbb{R}^d} v f(t, v) dv}{\int_{\mathbb{R}^d} f(t, v) dv}$$

is the *mean velocity*. We refer to [2] for more details. Notice that (1.1) is one-homogeneous: from now on, we will assume that the mass satisfies $\int_{\mathbb{R}^d} f(t, v) dv = 1$ for any $t \geq 0$, without loss of generality. In (1.1), the velocity consensus term $v - \mathbf{u}_f$ can be interpreted as a friction force which tends to align v and \mathbf{u}_f . Altogether, individuals are driven to a velocity corresponding to a speed of order 1 and a direction given by \mathbf{u}_f , but this mechanism is balanced by the noise which pushes the system towards an isotropic distribution with zero average velocity. The Vicsek model can be obtained as a limit case in which we let $\alpha \rightarrow +\infty$: see [4]. The competition between the two mechanisms, relaxation towards a non-zero average velocity and noise, is responsible for a phase transition between an ordered state for small values of D , with a distribution function f centered around \mathbf{u} with $\mathbf{u} \neq \mathbf{0}$, and a disordered, *symmetric* state with $\mathbf{u} = \mathbf{0}$. This phase transition can also be interpreted as a *symmetry breaking* mechanism from the *isotropic distribution* to an ordered, asymmetric or *polarized distribution*, with the remarkable feature that nothing but the initial datum determines the direction of \mathbf{u}_f for large values of t and any stationary solution generates a continuum of stationary solutions by rotation. We refer to [12] for more detailed comments and additional references on related models.

So far, a phase transition has been established in [12] when $d=1$ and it has been proved in [1] by A. Barbaro, J. Canizo, J. Carrillo and P. Degond that stationary solutions are isotropic for large values of D while symmetry breaking occurs as $D \rightarrow 0$. The bifurcation diagram showing the phase transition has also been studied numerically in [1] and the phase diagram can be found in [12, Theorem 2.1]. The first purpose of this paper is to classify all stable and unstable stationary solutions and establish a complete description of the phase transition.

THEOREM 1.1. *Let $d \geq 1$ and $\alpha > 0$. There exists a critical intensity of the noise $D_* > 0$ such that*

- (i) *if $D > D_*$ there exists one and only one non-negative stationary distribution which is isotropic and stable,*
- (ii) *if $D < D_*$ there exist one and only one non-negative isotropic stationary distribution which is unstable, and a continuum of stable non-negative non-symmetric stationary distributions, but this non-symmetric stationary solution is unique up to a rotation.*

Under the assumption of mass normalization to 1, it is straightforward to observe that

any stationary solution can be written as

$$f_{\mathbf{u}}(v) = \frac{e^{-\frac{1}{D}(\frac{1}{2}|v-\mathbf{u}|^2 + \frac{\alpha}{4}|v|^4 - \frac{\alpha}{2}|v|^2)}}{\int_{\mathbb{R}^d} e^{-\frac{1}{D}(\frac{1}{2}|v-\mathbf{u}|^2 + \frac{\alpha}{4}|v|^4 - \frac{\alpha}{2}|v|^2)} dv}$$

where $\mathbf{u} = (u_1, \dots, u_d) \in \mathbb{R}^d$ solves $\int_{\mathbb{R}^d} (\mathbf{u} - v) f_{\mathbf{u}}(v) dv = 0$. Up to a rotation, we can assume that $\mathbf{u} = (u, 0, \dots, 0) = u e_1$ and the question of finding stationary solutions to (1.1) is reduced to solve $u \in \mathbb{R}$ such that

$$\mathcal{H}(u) = 0 \tag{1.2}$$

where

$$\mathcal{H}(u) := \int_{\mathbb{R}^d} (v_1 - u) e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)} dv \quad \text{and} \quad \phi_\alpha(v) := \frac{\alpha}{4}|v|^4 + \frac{1-\alpha}{2}|v|^2.$$

Obviously $u = 0$ is always a solution. Moreover, if u is a solution of (1.2), then $-u$ is also a solution. As a consequence, from now on, we always suppose that $u \geq 0$. Theorem 1.1 is proved in Section 2 by analyzing (1.2).

The second purpose of this paper is to study the stability of the stationary states and the rates of convergence of the solutions of the evolution problem. A key tool is the *free energy*

$$\mathcal{F}[f] := D \int_{\mathbb{R}^d} f \log f dv + \int_{\mathbb{R}^d} f \phi_\alpha dv - \frac{1}{2} |\mathbf{u}_f|^2 \tag{1.3}$$

and we shall also consider the *relative entropy* with respect to $f_{\mathbf{u}}$ defined as

$$\mathcal{F}[f] - \mathcal{F}[f_{\mathbf{u}}] = D \int_{\mathbb{R}^d} f \log \left(\frac{f}{f_{\mathbf{u}}} \right) dv - \frac{1}{2} |\mathbf{u}_f - \mathbf{u}|^2$$

where $f_{\mathbf{u}}$ is a stationary solution to be determined. Notice that $f_{\mathbf{u}}$ is a critical point of \mathcal{F} under the mass constraint. Since there is only one stationary solution $f_{\mathbf{u}}$ corresponding to $\mathbf{u} = \mathbf{0}$ if $D > D_*$ and since \mathcal{F} is strictly convex, in that case we know that $f_{\mathbf{0}}$ is the unique minimizer of \mathcal{F} , it is non-linearly stable and in particular we have that $\mathcal{F}[f] - \mathcal{F}[f_{\mathbf{u}}] \geq 0$. See Section 4 for more details.

To a distribution function f , we associate the *non-equilibrium Gibbs state*

$$G_f(v) := \frac{e^{-\frac{1}{D}(\frac{1}{2}|v-\mathbf{u}_f|^2 + \frac{\alpha}{4}|v|^4 - \frac{\alpha}{2}|v|^2)}}{\int_{\mathbb{R}^d} e^{-\frac{1}{D}(\frac{1}{2}|v-\mathbf{u}_f|^2 + \frac{\alpha}{4}|v|^4 - \frac{\alpha}{2}|v|^2)} dv}. \tag{1.4}$$

Unless f is a stationary solution of (1.1), let us notice that G_f does not solve (1.1). A crucial observation is that

$$\mathcal{F}[f] = D \int_{\mathbb{R}^d} f \log f dv + \frac{1}{2} \int_{\mathbb{R}^d} |v - \mathbf{u}_f|^2 f dv + \int_{\mathbb{R}^d} \left(\frac{\alpha}{4}|v|^4 + \frac{\alpha}{2}|v|^2 \right) f dv$$

is a Lyapunov function in the sense that

$$\frac{d}{dt} \mathcal{F}[f(t, \cdot)] = -\mathcal{I}[f(t, \cdot)]$$

if f solves (1.1), where $\mathcal{I}[f]$ is the *relative Fisher information* of f defined as

$$\mathcal{I}[f] := \int_{\mathbb{R}^d} \left| D \frac{\nabla f}{f} + \alpha v |v|^2 + (1 - \alpha)v - \mathbf{u}_f \right|^2 f dv = D^2 \int_{\mathbb{R}^d} \left| \nabla \log \left(\frac{f}{G_f} \right) \right|^2 f dv. \quad (1.5)$$

It is indeed clear that $\mathcal{F}[f(t, \cdot)]$ is monotone non-increasing and $\frac{d}{dt} \mathcal{F}[f(t, \cdot)] = 0$ if and only if $f = G_f$ is a stationary solution of (1.1). This is consistent with our first stability result.

PROPOSITION 1.1. *For any $d \geq 1$ and any $\alpha > 0$, f_0 is a linearly stable critical point if and only if $D > D_*$.*

Actually, from the dynamical point of view, we have a better, global result.

THEOREM 1.2. *For any $d \geq 1$ and any $\alpha > 0$, if $D > D_*$, then for any solution f of (1.1) with nonnegative initial datum f_{in} of mass 1 such that $\mathcal{F}[f_{\text{in}}] < \infty$, there are two positive constants C and λ such that, for any time $t > 0$,*

$$0 \leq \mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_0] \leq C e^{-\lambda t}. \quad (1.6)$$

We shall also prove that

$$\int_{\mathbb{R}^d} |f(t, \cdot) - f_0|^2 f_0^{-1} dv \leq C e^{-\lambda t}$$

with same $\lambda > 0$ as in Theorem 1.2, but eventually for a different value of C , and characterize λ as the spectral gap of the linearized evolution operator in an appropriate norm. A characterization of the optimal rate λ is given in Theorem 5.1.

For $D < D_*$, the situation is more subtle. The solution of (1.1) can in principle converge either to the *isotropic stationary solution* f_0 or to a *polarized, non-symmetric stationary solution* $f_{\mathbf{u}}$ with $\mathbf{u} \neq \mathbf{0}$. We will prove that $\mathcal{F}[f] - \mathcal{F}[f_{\mathbf{u}}]$ decays with an exponential rate which is also characterized by a spectral gap in Section 6.

This article is organized as follows. In Section 2, we classify all stationary solutions, prove Theorem 1.1 and deduce that a phase transition occurs at $D = D_*$. Section 3 is devoted to the linearization. The relative entropy and the relative Fisher information provide us with two quadratic forms which are related by the linearized evolution operator. The main result here is to prove a spectral gap property for this operator in the appropriate norm, which is inspired by a similar method used in [5] to study the sub-critical Keller-Segel model: see Proposition 3.1. It is crucial to take into account all terms in the linearization, including the term arising from the non-local mean velocity. The proof of Theorem 1.2 follows using a Grönwall type estimate, in Section 5 (isotropic case). In Section 6, we also give some results in the polarized case.

2. Stationary solutions and phase transition

The aim of this section is to classify all stationary solutions of (1.1) as a first step of the proof of the phase transition result of Theorem 1.1. Our proofs are based on elementary although somewhat painful computations.

2.1. A technical observation

Let us start by the simple observation that

$$-D \frac{\partial}{\partial v_1} \left(e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)} \right) = (v_1 - u + \alpha(|v|^2 - 1)v_1) e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)}$$

can be integrated on \mathbb{R}^d to rewrite \mathcal{H} as

$$\mathcal{H}(u) = \alpha \int_{\mathbb{R}^d} (1 - |v|^2) v_1 e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)} dv$$

and compute

$$\mathcal{H}'(u) = \frac{\alpha}{D} \int_{\mathbb{R}^d} (1 - |v|^2) v_1^2 e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)} dv.$$

We observe that $\mathcal{H}'(0) = \frac{\alpha}{D} |\mathbb{S}^{d-1}| h_d(D)$ where

$$h_d(D) := \int_0^\infty (s^{d+1} - s^{d+3}) e^{-\frac{\varphi_\alpha(s)}{D}} ds \quad \text{and} \quad \phi_\alpha(s) := \frac{\alpha}{4} s^4 + \frac{1-\alpha}{2} s^2.$$

With these notations, we are now in a position to state a key ingredient of the proof.

PROPOSITION 2.1. *For any $d \geq 1$ and any $\alpha > 0$, $h_d(D)$ has a unique positive root D_* . Moreover h_d is positive on $(0, D_*)$ and negative on $(D_*, +\infty)$.*

Proof. Our goal is to prove that $h_d = j_{d+1} - j_{d+3}$ is positive on $(0, D_*)$ and negative on $(D_*, +\infty)$ for some $D_* > 0$, where

$$j_d(D) := \int_0^\infty s^d e^{-\frac{1}{D} \varphi_\alpha(s)} ds. \quad (2.1)$$

Let us start with two useful identities. A completion of the square shows that

$$j_{n+5} - 2j_{n+3} + j_{n+1} = \int_0^\infty s^{n+1} (s^2 - 1)^2 e^{-\frac{\varphi_\alpha}{D}} ds > 0. \quad (2.2)$$

With an integration by parts, we obtain that

$$\alpha j_{n+5} + (1 - \alpha) j_{n+3} = \int_0^\infty s^{n+2} \varphi'_\alpha e^{-\frac{1}{D} \varphi_\alpha} ds = (n+2) D j_{n+1}. \quad (2.3)$$

Next, we split the proof in a series of claims.

- *The function h_d is positive on $(0, 1/(d+2)]$ and negative on $[1/d, +\infty)$.* Let us prove this claim. With $n = d$ and $n = d - 2$, we deduce from (2.2) and (2.3) that

$$h_d > \frac{1 - (d+2)D}{1 + \alpha} j_{d+1} \quad \text{and} \quad h_d < \frac{1 - dD}{1 + \alpha} j_{d-1}.$$

As a consequence, if $h_d(D) = 0$, then $D \in (1/(d+2), 1/d)$.

- *If $\alpha \leq 1$, then $h_d(D) = 0$ has a unique solution.* By a direct computation, we observe that

$$4D^2 h'_d = \alpha h_{d+4} + 2(1 - \alpha) h_{d+2}$$

using (2.3) with $n = d + 2$. If $\alpha \in (0, 1)$, it follows that $h'_d < 0$ on $[1/(d+2), +\infty)$, which proves the claim.

- *If $\alpha > 1$ and $h'_d(D_\circ) = 0$ for some $D_\circ \in (1/(d+2), 1/d)$, then $h_d(D_\circ) > 0$.* Indeed, using

$$4D^2 h'_d = -\alpha j_{d+7} + (3\alpha - 2) j_{d+5} + 2(1 - \alpha) j_{d+3} = 0,$$

combined with (2.3) for $n = d + 2$ and $n = d$, we find that, at $D = D_\circ$,

$$h_d(D_\circ) = \frac{(d+2)D - 1 + \alpha(1 - dD)}{\alpha - 1 + (d+4)D\alpha} j_{d+1}.$$

Collecting our observations concludes the proof. See Fig. A.1 for an illustration. \square

2.2. The one-dimensional case

LEMMA 2.1. *Let us consider a continuous positive function ψ on \mathbb{R}^+ such that the function $s \mapsto \psi(s)e^{s^2}$ is integrable and define*

$$H(u) := \int_0^{+\infty} (1-s^2) \psi(s) \sinh(su) ds \quad \forall u \geq 0.$$

For any $u > 0$, $H''(u) < 0$ if $H(u) \leq 0$. As a consequence, H changes sign at most once on $(0, +\infty)$.

Proof. We first observe that

$$H''(u) - H(u) = \int_0^{+\infty} (1-s^2)(s^2-1) \psi(s) \sinh(su) ds < 0 \quad \forall u > 0. \quad (2.4)$$

Let $u_* > 0$ be such that $H(u_*) = 0$. If $H'(u_*) < 0$, there is a neighborhood of $(u_*)_+$ such that both H and H' are negative. As a consequence, by continuation, $H'(u) < H'(u_*) < 0$ for any $u > u_*$. We also get that $H'(u) < 0$ for any $u > u_*$ if $H'(u_*) = 0$ because we know that $H''(u_*) < 0$. We conclude by observing that $H'(u_*) > 0$ would imply $H'(u) > H'(u_*)$ for any $u \in (0, u_*)$, a contradiction with $H(0) = 0$. \square

PROPOSITION 2.2. *Assume that $d=1$ and $\alpha > 0$. With the notations of Proposition 2.1, Equation (1.2), i.e., $\mathcal{H}(u) = 0$, has as a solution $u = u(D) > 0$ if and only if $D < D_*$ and $\lim_{D \rightarrow (D_*)_-} u(D) = 0$.*

In other words, there exists a solution to (1.2) if and only if $\mathcal{H}'(0) > 0$.

Proof. Since $\mathcal{H}(0) = 0$, for any $D \neq D_*$, $h_d(D)$ and $\mathcal{H}(u)$ have the same sign in a neighborhood of $u = 0_+$. Next we notice that

$$-\frac{1}{\alpha} \mathcal{H}(u) = \int_0^\infty (v^2 - 1) v e^{-\frac{\phi_\alpha(v)}{D}} e^{\frac{uv}{D}} dv - \int_0^\infty (v^2 - 1) v e^{-\frac{\phi_\alpha(v)}{D}} e^{-\frac{uv}{D}} dv.$$

The second term of the right-hand side converges to 0 as $u \rightarrow \infty$ by the dominated convergence theorem. Concerning the first term, let us notice that $|(v^2 - 1)v| e^{-\phi_\alpha(v)/D}$ is bounded on $(0, 3)$, so that

$$\begin{aligned} & \int_0^\infty (v^2 - 1) v e^{-\frac{\phi_\alpha(v)}{D}} e^{\frac{uv}{D}} dv \\ & \geq \int_0^1 (v^2 - 1) v e^{-\frac{\phi_\alpha(v)}{D}} e^{\frac{uv}{D}} dv + \int_2^3 (v^2 - 1) v e^{-\frac{\phi_\alpha(v)}{D}} e^{\frac{uv}{D}} dv \\ & \geq -C_1 e^{u/D} + C_2 e^{2u/D} \rightarrow +\infty \quad \text{as } u \rightarrow +\infty \end{aligned}$$

for some positive constants C_1 and C_2 . This proves that $\lim_{u \rightarrow +\infty} \mathcal{H}(u) = -\infty$ and shows the existence of at least one positive solution of (1.2) if $h_d(D) > 0$.

The fact that (1.2) has at most one solution on $(0, +\infty)$ follows from Lemma 2.1 applied with $H(u) = \mathcal{H}(Du)$ and $\psi(v) = 2\alpha v e^{-\frac{\phi_\alpha(v)}{D}}$. Finally, as consequence of the regularity of H and of (2.4), the solution $u = u(D)$ of (1.2) is such that $\lim_{D \rightarrow (D_*)_-} u(D) = 0$. \square

2.3. The case of a dimension $d \geq 2$

We extend the result of Proposition 2.2 to higher dimensions.

PROPOSITION 2.3. *Assume that $d \geq 2$ and $\alpha > 0$. With the notations of Proposition 2.1, Equation (1.2), i.e., $\mathcal{H}(u) = 0$, has as a solution $u = u(D) > 0$ if and only if $D < D_*$ and $\lim_{D \rightarrow (D_*)_-} u(D) = 0$.*

Qualitatively, the result is the same as in dimension $d = 1$: there exists a solution to (1.2) if and only if $\mathcal{H}'(0) > 0$. See Fig. A.2.

In radial coordinates such that $s = |v|$ and $v_1 = s \cos \theta$, with $\theta \in [0, \pi]$,

$$\mathcal{H}(u) = \alpha |\mathbb{S}^{d-2}| \int_0^\pi \int_0^{+\infty} (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} \cos \theta (\sin \theta)^{d-2} e^{\frac{us}{D} \cos \theta} ds d\theta$$

written with the convention that $|\mathbb{S}^0| = 2$ can also be rewritten as

$$\mathcal{H}(u) = 2\alpha |\mathbb{S}^{d-2}| \int_0^{\pi/2} \int_0^{+\infty} (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} \cos \theta (\sin \theta)^{d-2} \sinh\left(\frac{us}{D} \cos \theta\right) ds d\theta.$$

Lemma 2.1 does not apply directly. Let us consider

$$\mathfrak{h}(s) := \int_0^{\pi/2} \cos \theta (\sin \theta)^{d-2} \sinh(s \cos \theta) d\theta. \quad (2.5)$$

LEMMA 2.2. *Assume that $d \geq 2$. The function \mathfrak{h} defined by (2.5) is such that $s \mapsto s\mathfrak{h}'(s)/\mathfrak{h}(s)$ is monotone increasing on $(0, +\infty)$.*

Proof. Let s_1 and s_2 be such that $0 < s_1 < s_2$ and consider a series expansion. With

$$P_n := \int_0^\pi (\cos \theta)^{2n} (\sin \theta)^{d-2} d\theta,$$

we know that

$$\begin{aligned} s_2 \mathfrak{h}'(s_2) \mathfrak{h}(s_1) &= \sum_{m=0}^{\infty} \frac{s_2^{2m+1}}{(2m)!} P_{m+1} \sum_{n=0}^{\infty} \frac{s_1^{2n+1}}{(2n+1)!} P_{n+1}, \\ s_1 \mathfrak{h}'(s_1) \mathfrak{h}(s_2) &= \sum_{m=0}^{\infty} \frac{s_1^{2m+1}}{(2m)!} P_{m+1} \sum_{n=0}^{\infty} \frac{s_2^{2n+1}}{(2n+1)!} P_{n+1}. \end{aligned}$$

These series are absolutely converging and we can reindex the difference of the two terms using $i = \min\{m, n\}$ to get

$$\begin{aligned} & s_2 \mathfrak{h}'(s_2) \mathfrak{h}(s_1) - s_1 \mathfrak{h}'(s_1) \mathfrak{h}(s_2) \\ &= \sum_{i=0}^{\infty} \sum_{j=1}^{\infty} \frac{(s_1 s_2)^{2i+1}}{(2i+2j+1)!(2i+1)!} P_{i+1} P_{j+1} \frac{2i+2j+1}{2(i+j+1)} \left(s_2^j - s_1^j\right) > 0. \end{aligned}$$

□

Proof of Proposition 2.3. We prove that $\lim_{u \rightarrow +\infty} \mathcal{H}(u) = -\infty$ as in the case $d = 1$ by considering the domains defined in the coordinates (s, θ) by $0 \leq s \leq 1$ and $\theta \in [0, \pi/2]$ on the one hand, and $2 \leq s \leq 3$ and $0 \leq \theta \leq \theta_*$ for some $\theta_* \in (0, \pi/6)$ on the other hand. The existence of at least one solution $u > 0$ of $\mathcal{H}(u) = 0$ follows from Proposition 2.2 if

$D < D_*$, and if $D > D_*$, we also know that $\mathcal{H}(u) = 0$ has either no positive solution, or at least two.

If there exist u_1 and u_2 such that $\mathcal{H}(u_1) = \mathcal{H}(u_2) = 0$ and $u_1 < u_2$, then

$$\int_0^1 (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) ds = \int_1^\infty (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) ds$$

where $\tilde{u}_1 := u_1/D < u_2/D =: \tilde{u}_2$. We deduce from Lemma 2.2 that the function $s \mapsto k(s) := h(\tilde{u}_2 s)/h(\tilde{u}_1 s)$ is a monotone increasing function on $(0, +\infty)$. Using $\mathcal{H}(u_1) = 0$, we obtain

$$\begin{aligned} \int_0^1 (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_2 s) ds &= \int_0^1 (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) k(s) ds \\ &< \int_0^1 (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) k(1) ds \\ &= \int_1^\infty (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) k(1) ds \\ &< \int_1^\infty (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_1 s) k(s) ds \\ &= \int_1^\infty (1-s^2) s^d e^{-\frac{\varphi_\alpha(s)}{D}} h(\tilde{u}_2 s) ds, \end{aligned}$$

a contradiction with $\mathcal{H}(u_2) = 0$. □

2.4. Classification of the stationary solutions and phase transition

We learn from the expression of \mathcal{I} in (1.5) that any stationary solution of (1.1) is of the form $f_{\mathbf{u}}$ with $\mathbf{u} = u\mathbf{e}_1$ for some u which solves (1.2). Since $\mathcal{H}(0) = 0$, $u = 0$ is always a solution. According to Propositions 2.2 and 2.3, Equation (1.2) has a solution $u = u(D)$ if and only if $D > D_*$ where D_* is obtained as the unique positive root of h_d by Proposition 2.1.

COROLLARY 2.1. *Let $d \geq 1$ and $\alpha > 0$. With the above notations and D_* defined as in Proposition 2.1, we know that*

- (i) *if $D > D_*$ there exists one and only one non-negative stationary distribution $f_{\mathbf{u}}$ given by $\mathbf{u} = \mathbf{0}$, which is isotropic,*
- (ii) *if $D < D_*$ there exists one and only one non-negative isotropic stationary distribution with $\mathbf{u} = \mathbf{0}$, and a continuum of stable non-negative non-symmetric stationary distributions $f_{\mathbf{u}}$ with $\mathbf{u} = u(D)\mathbf{e}$ for any $\mathbf{e} \in \mathbb{S}^{d-1}$, with the convention that $\mathbb{S}^0 = \{-1, 1\}$.*

There are no other stationary solutions. In other words, we have obtained the complete classification of the stationary solutions of (1.1), which shows that there are two phases of stationary solutions: the isotropic ones with $\mathbf{u} = \mathbf{0}$, and the non-isotropic ones with $\mathbf{u} \neq \mathbf{0}$ which are unique up to a rotation and exist only if $D < D_*$. To complete the proof of Theorem 1.1, we have to study the linear stability of these stationary solutions.

2.5. An important estimate

The next result is a technical estimate which is going to play a key role in our analysis.

LEMMA 2.1. *Assume that $d \geq 1$, $\alpha > 0$ and $D > 0$.*

- (i) In the case $\mathbf{u} = \mathbf{0}$, we have that $\int_{\mathbb{R}^d} |v|^2 f_{\mathbf{0}} dv > dD$ if and only if $D < D_*$.
(ii) In the case $D \in (0, D_*)$ and $\mathbf{u} \neq \mathbf{0}$, we have that

$$\int_{\mathbb{R}^d} |(v - \mathbf{u}) \cdot \mathbf{u}|^2 f_{\mathbf{u}} dv < D |\mathbf{u}|^2.$$

- (iii) In the case $d \geq 2$ and $D \in (0, D_*)$ and $\mathbf{u} \neq \mathbf{0}$, we have that

$$\int_{\mathbb{R}^d} |(v - \mathbf{u}) \cdot \mathbf{w}|^2 f_{\mathbf{u}} dv = D |\mathbf{w}|^2 \quad \forall \mathbf{w} \in \mathbb{R}^d \quad \text{such that} \quad \mathbf{u} \cdot \mathbf{w} = 0.$$

Proof. Using Definition (2.1), we observe that $\int_{\mathbb{R}^d} |v|^2 f_{\mathbf{0}} dv - dD$ has the sign of

$$j_{d+1} - dD j_{d-1} = \alpha (j_{d+1} - j_{d+3}) = \alpha h_d(D)$$

by (2.3) with $n = d - 2$. This proves (i) according to Proposition 2.1 and Corollary 2.1.

By integrating $D \mathbf{u} \cdot \nabla((\mathbf{u} \cdot v) f_{\mathbf{u}})$, we obtain that

$$\begin{aligned} 0 &= \int_{\mathbb{R}^d} \left(D |\mathbf{u}|^2 - (\mathbf{u} \cdot v)^2 (\alpha |v|^2 + 1 - \alpha) + u(\mathbf{u} \cdot v) \right) f_{\mathbf{u}} dv \\ &= D |\mathbf{u}|^2 - \int_{\mathbb{R}^d} |(v - \mathbf{u}) \cdot \mathbf{u}|^2 f_{\mathbf{u}} dv + D |\mathbf{u}|^2 \mathcal{H}'(|\mathbf{u}|) \end{aligned}$$

Then (ii) follows from Propositions 2.2 and 2.3 because $\mathcal{H}'(u) < 0$ if $u = u(D) = |\mathbf{u}|$.

With no loss of generality, we can assume that $\mathbf{u} = (u, 0, \dots, 0) \neq \mathbf{0}$. By integrating $\frac{\partial}{\partial v_1} f_{\mathbf{u}}$ on \mathbb{R}^d , we know that $\int_{\mathbb{R}^d} (|v|^2 - 1) v_1 f_{\mathbf{u}} dv = 0$. Let us consider radial coordinates such that $s = |v|$ and $v_1 = s \cos \theta$, with $\theta \in [0, \pi]$. From the integration by parts

$$(d-1)D \int_0^\pi \cos \theta (\sin \theta)^{d-2} e^{\frac{u s}{D} \cos \theta} d\theta = u s \int_0^\pi (\sin \theta)^d e^{\frac{u s}{D} \cos \theta} d\theta,$$

we deduce that $\int_{\mathbb{R}^d} (|v|^2 - 1) (1 - v_1^2) f_{\mathbf{u}} dv = 0$ because $s^2 (\sin \theta)^2 = 1 - v_1^2$ and

$$\int_{\mathbb{R}^d} (|v|^2 - 1) v_i^2 f_{\mathbf{u}} dv = 0 \quad \forall i \geq 2$$

by symmetry among the variables v_2, v_3, \dots, v_d . We conclude by integrating $\frac{\partial}{\partial v_i} f_{\mathbf{u}}$ on \mathbb{R}^d that

$$\int_{\mathbb{R}^d} |v_i|^2 f_{\mathbf{u}} dv = D \quad \forall i \geq 2,$$

which concludes the proof of (iii). \square

COROLLARY 2.2. *Assume that $d \geq 1$, $\alpha > 0$ and $\mathbf{e} \in \mathbb{S}^{d-1}$. There exists a function $D \mapsto \kappa(D)$ on $(0, D_*)$ which is continuous with values in $(0, 1)$ such that, with $\mathbf{u} = u(D)\mathbf{e}$,*

$$\frac{1}{D} \int_{\mathbb{R}^d} |(v - \mathbf{u}) \cdot \mathbf{w}|^2 f_{\mathbf{u}} dv = \kappa(D) (\mathbf{w} \cdot \mathbf{e})^2 + |\mathbf{w}|^2 - (\mathbf{w} \cdot \mathbf{e})^2 \quad \forall \mathbf{w} \in \mathbb{R}^d.$$

With $\kappa(D) := \frac{1}{u(D)^2} \int_{\mathbb{R}^d} |(v - \mathbf{u}) \cdot \mathbf{u}|^2 f_{\mathbf{u}} dv$ and $\mathbf{u} = u(D)\mathbf{e}$ for an arbitrary $\mathbf{e} \in \mathbb{S}^{d-1}$, the proof is a straightforward consequence of Lemma 2.1.

3. The linearized problem: local properties of the stationary solutions

This section is devoted to the quadratic forms associated with the expansion of the free energy \mathcal{F} and the Fisher information \mathcal{I} around the stationary solution $f_{\mathbf{u}}$ studied in Section 2. These quadratic forms are defined for a smooth perturbation g of $f_{\mathbf{u}}$ such that $\int_{\mathbb{R}^d} g f_{\mathbf{u}} dv = 0$ by

$$Q_{1,\mathbf{u}}[g] := \lim_{\varepsilon \rightarrow 0} \frac{2}{\varepsilon^2} \mathcal{F}[f_{\mathbf{u}}(1 + \varepsilon g)] = D \int_{\mathbb{R}^d} g^2 f_{\mathbf{u}} dv - D^2 |\mathbf{v}_g|^2 \quad \text{where } \mathbf{v}_g := \frac{1}{D} \int_{\mathbb{R}^d} v g f_{\mathbf{u}} dv,$$

$$Q_{2,\mathbf{u}}[g] := \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^2} \mathcal{I}[f_{\mathbf{u}}(1 + \varepsilon g)] = D^2 \int_{\mathbb{R}^d} |\nabla g - \mathbf{v}_g|^2 f_{\mathbf{u}} dv.$$

3.1. Stability of the isotropic stationary solution

The first result is concerned with the linear stability of \mathcal{F} around $f_{\mathbf{0}}$.

LEMMA 3.1. *On the space of the functions $g \in L^2(f_{\mathbf{0}} dv)$ such that $\int_{\mathbb{R}^d} g f_{\mathbf{0}} dv = 0$, $Q_{1,\mathbf{0}}$ is a nonnegative (resp. positive) quadratic form if and only if $D \geq D_*$ (resp. $D > D_*$). Moreover, for any $D > D_*$, let $\eta(D) := \alpha \mathcal{C} |h_d(D)$ for some explicit $\mathcal{C} = \mathcal{C}(D) > 0$. Then*

$$Q_{1,\mathbf{0}}[g] \geq \eta(D) \int_{\mathbb{R}^d} g^2 f_{\mathbf{0}} dv \quad \forall g \in L^2(f_{\mathbf{0}} dv) \quad \text{such that} \quad \int_{\mathbb{R}^d} g f_{\mathbf{0}} dv = 0. \quad (3.1)$$

Proof. Let $\mathbf{e} \in \mathbb{S}^{d-1}$. We consider $g(v) = v \cdot \mathbf{e}$ and, using (2.3) with $n = d - 2$, compute

$$Q_{1,\mathbf{0}}[g] = D \int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv - \left(\int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv \right)^2 = \mathcal{C} \int_0^\infty (d D s^{d-1} - s^{d+1}) e^{-\frac{\varphi_\alpha(s)}{D}} ds$$

where the last equality determines the value of \mathcal{C} . If $D < D_*$, this proves that $Q_{1,\mathbf{0}}[g] = -\alpha \mathcal{C} h_d(D) < 0$ and, as a consequence, the linear instability of $f_{\mathbf{0}}$.

On the other hand, let g be a function in $L^2(\mathbb{R}^d, f_{\mathbf{0}} dv)$ such that $\int_{\mathbb{R}^d} g^2 f_{\mathbf{0}} dv = \int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv$. We can indeed normalize g with no loss of generality. With $v_1 = v \cdot \mathbf{e}$, $\mathbf{e} \in \mathbb{S}^{d-1}$ such that $\mathbf{u}_g f_{\mathbf{0}} = u \mathbf{e}$ for some $u \in \mathbb{R}$, we know by the Cauchy-Schwarz inequality that

$$\left(\int_{\mathbb{R}^d} v_1 g f_{\mathbf{0}} dv \right)^2 \leq \int_{\mathbb{R}^d} g^2 f_{\mathbf{0}} dv \int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv = \left(\int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv \right)^2 = \left(\frac{1}{d} \int_{\mathbb{R}^d} |v|^2 f_{\mathbf{0}} dv \right)^2,$$

hence

$$Q_{1,\mathbf{0}}[g] \geq D \int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv - \left(\int_{\mathbb{R}^d} v_1^2 f_{\mathbf{0}} dv \right)^2 = -\alpha \mathcal{C} h_d(D).$$

This proves the linear stability of $f_{\mathbf{0}}$ if $D > D_*$. □

The classification result of Theorem 1.1 is a consequence of Corollary 2.1 and Lemma 3.1.

3.2. A coercivity result

Let us start by recalling the *Poincaré inequality*

$$\int_{\mathbb{R}^d} |\nabla h|^2 f_{\mathbf{u}} dv \geq \Lambda_D \int_{\mathbb{R}^d} |h|^2 f_{\mathbf{u}} dv \quad \forall h \in H^1(\mathbb{R}^d, f_{\mathbf{u}} dv) \quad \text{such that} \quad \int_{\mathbb{R}^d} h f_{\mathbf{u}} dv = 0. \quad (3.2)$$

Here \mathbf{u} is an admissible velocity such that $\mathbf{u} = \mathbf{0}$ if $D \geq D_*$, or $|\mathbf{u}| = u(D)$ if $D < D_*$, and Λ_D denotes the corresponding optimal constant. Since φ_α can be seen as a uniformly strictly convex potential perturbed by a bounded perturbation, it follows from the *carré du champ* method and the Holley-Stroock lemma that Λ_D is a positive constant. Let

$$\begin{aligned} \mathbf{u}[f] &= \mathbf{0} \quad \text{if } D \geq D_* \quad \text{or} \quad \mathbf{u}_f = \mathbf{0} \quad \text{and} \quad D < D_*, \\ \mathbf{u}[f] &= \frac{u(D)}{|\mathbf{u}_f|} \mathbf{u}_f \quad \text{if } D < D_* \quad \text{and} \quad \mathbf{u}_f \neq \mathbf{0}. \end{aligned}$$

Based on (3.2), we have the following coercivity result.

PROPOSITION 3.1. *Let $d \geq 1$, $\alpha > 0$, $D > 0$ and $\mathcal{C}_D = D \Lambda_D$ with Λ_D as in (3.2). Let us consider a nonnegative distribution function $f \in L^1(\mathbb{R}^d)$ with $\int_{\mathbb{R}^d} f dv = 1$, let $\mathbf{u} \in \mathbb{R}^d$ be such that either $\mathbf{u} = \mathbf{0}$ or $|\mathbf{u}| = u(D)$ if $D < D_*$ and consider $g = (f - f_{\mathbf{u}})/f_{\mathbf{u}}$. We assume that $g \in H^1(\mathbb{R}^d, f_{\mathbf{u}} dv)$. If $\mathbf{u} = \mathbf{0}$, then*

$$Q_{2,\mathbf{u}}[g] \geq \mathcal{C}_D Q_{1,\mathbf{u}}[g].$$

Otherwise, if $\mathbf{u} \neq \mathbf{0}$ for some $D \in (0, D_*)$ with D_* as in Corollary 2.1, then we have

$$Q_{2,\mathbf{u}}[g] \geq \mathcal{C}_D (1 - \kappa(D)) \frac{(\mathbf{v}_g \cdot \mathbf{u})^2}{|\mathbf{v}_g|^2 |\mathbf{u}|^2} Q_{1,\mathbf{u}}[g]$$

with $\mathbf{v}_g := \frac{1}{D} \int_{\mathbb{R}^d} (v - \mathbf{u}) g f_{\mathbf{u}} dv$ and $\kappa(D) < 1$ defined as in Corollary 2.2.

As a special case, if $\mathbf{u} = \mathbf{u}[f]$, then $Q_{2,\mathbf{u}}[g] \geq \mathcal{C}_D (1 - \kappa(D)) Q_{1,\mathbf{u}}[g]$.

By construction, \mathbf{v}_g is such that $D \mathbf{v}_g = \int_{\mathbb{R}^d} (v - \mathbf{u}) g f_{\mathbf{u}} dv = \int_{\mathbb{R}^d} v g f_{\mathbf{u}} dv = \mathbf{u}_f - \mathbf{u}$ because $\int_{\mathbb{R}^d} g f_{\mathbf{u}} dv = 0$.

Proof. Let us apply (3.2) to $h(v) = g(v) - (v - \mathbf{u}) \cdot \mathbf{v}_g$. Using $\mathbf{v}_g = \frac{1}{D} \int_{\mathbb{R}^d} (v - \mathbf{u}) g f_{\mathbf{u}} dv$ and $\int_{\mathbb{R}^d} g f_{\mathbf{u}} dv = 0$, we obtain

$$\begin{aligned} \frac{1}{D^2} Q_{2,\mathbf{u}}[g] &= \int_{\mathbb{R}^d} |\nabla g - \mathbf{v}_g|^2 f_{\mathbf{u}} dv \\ &\geq \Lambda_D \int_{\mathbb{R}^d} (g^2 + |\mathbf{v}_g \cdot (v - \mathbf{u})|^2 - 2 \mathbf{v}_g \cdot (v - \mathbf{u}) g) f_{\mathbf{u}} dv \\ &= \Lambda_D \left[\int_{\mathbb{R}^d} |g|^2 f_{\mathbf{u}} dv + \int_{\mathbb{R}^d} |\mathbf{v}_g \cdot (v - \mathbf{u})|^2 f_{\mathbf{u}} dv - 2D |\mathbf{v}_g|^2 \right]. \end{aligned}$$

If $\mathbf{u} = \mathbf{0}$, either $\mathbf{v}_g = \mathbf{0}$ and the result is proved, or we know that $\frac{1}{d} \int_{\mathbb{R}^d} |v|^2 f_{\mathbf{0}} dv \geq D$ by Lemma 2.1 because $D \geq D_*$ by assumption. In that case we can estimate the r.h.s. by

$$\int_{\mathbb{R}^d} |g|^2 f_{\mathbf{0}} dv + |\mathbf{v}_g|^2 \left(\frac{1}{d} \int_{\mathbb{R}^d} |v|^2 f_{\mathbf{0}} dv - 2D \right) \geq \int_{\mathbb{R}^d} |g|^2 f_{\mathbf{0}} dv - D |\mathbf{v}_g|^2 = \frac{1}{D} Q_{1,\mathbf{0}}[g],$$

which again proves the result whenever $\mathbf{u} = \mathbf{0}$.

If $\mathbf{u} \neq \mathbf{0}$, let us apply Corollary 2.2 with $\mathbf{w} = \mathbf{v}_g$ and $\kappa = \kappa(D)$:

$$\int_{\mathbb{R}^d} |\mathbf{v}_g \cdot (v - \mathbf{u})|^2 f_{\mathbf{u}} dv = \mathcal{K} D |\mathbf{v}_g|^2 \quad \text{with} \quad \mathcal{K} = 1 - (1 - \kappa) \frac{(\mathbf{v}_g \cdot \mathbf{u})^2}{|\mathbf{v}_g|^2 |\mathbf{u}|^2}.$$

We deduce from the Cauchy-Schwarz inequality

$$D^2 |\mathbf{v}_g|^4 = \left(\int_{\mathbb{R}^d} \mathbf{v}_g \cdot (v - \mathbf{u}) f_{\mathbf{u}} dv \right)^2 \leq \int_{\mathbb{R}^d} |g|^2 f_{\mathbf{0}} dv \int_{\mathbb{R}^d} |\mathbf{v}_g \cdot (v - \mathbf{u})|^2 f_{\mathbf{u}} dv$$

that $D |\mathbf{v}_g|^2 \leq \mathcal{K} \int_{\mathbb{R}^d} |g|^2 f_{\mathbf{0}} dv$. Hence, if $\beta \in (0, 1)$, we obtain

$$\frac{1}{D^2} Q_{2, \mathbf{u}}[g] - \frac{\beta}{D^2} Q_{2, \mathbf{u}}[g] \geq (1 - \beta - (2 - \mathcal{K} - \beta)\mathcal{K}) \int_{\mathbb{R}^d} |g|^2 f_{\mathbf{0}} dv.$$

With $\beta = 1 - \mathcal{K}$, we obtain $1 - \beta - (2 - \mathcal{K} - \beta)\mathcal{K} = 0$, which proves the result. \square

4. Properties of the free energy and consequences

We consider the free energy \mathcal{F} and the Fisher information \mathcal{I} defined respectively by (1.3) and (1.5).

4.1. Basic properties of the free energy

PROPOSITION 4.1. *Assume that f_{in} is a nonnegative function in $L^1(\mathbb{R}^d)$ such that $\mathcal{F}[f_{\text{in}}] < \infty$. Then there exists a solution $f \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^d))$ of (1.1) with initial datum f_{in} such that $\mathcal{F}[f(t, \cdot)]$ is nonincreasing and a.e. differentiable on $[0, \infty)$. Furthermore*

$$\frac{d}{dt} \mathcal{F}[f(t, \cdot)] \leq -\mathcal{I}[f(t, \cdot)], \quad t > 0 \quad \text{a.e.}$$

This result is classical and we shall skip its proof: see for instance [6, Proposition 2.1] for further details. One of the difficulties in the study of \mathcal{F} is that in (1.3), the term $|\mathbf{u}|^2$ has a negative coefficient, so that the functional \mathcal{F} is not convex. A smooth solution realizes the equality, and by approximations, we obtain the result.

PROPOSITION 4.2. *\mathcal{F} is bounded from below on the set*

$$\left\{ f \in L^1_+(\mathbb{R}^d) : \int_{\mathbb{R}^d} f dv = 1 \text{ and } \int_{\mathbb{R}^d} |v|^4 f dv < \infty \right\}$$

and

$$\int_{\mathbb{R}^d} |v|^4 f dv \leq \frac{1}{\alpha^2} \left(D + \alpha + \sqrt{(D + \alpha)^2 + 4\alpha(\mathcal{F}[f] + \frac{d}{2} \log(2\pi)D)} \right)^2.$$

Proof. Let $g = f/\mu$ where $\mu(v) := (2\pi)^{-d/2} e^{-\frac{1}{2}|v|^2}$ and $d\mu = \mu dv$. Since $g \log g \geq g - 1$ and $\int_{\mathbb{R}^d} (g - 1) d\mu = 0$, we have the classical estimate

$$\int_{\mathbb{R}^d} f \log f dv + \frac{1}{2} \int_{\mathbb{R}^d} |v|^2 f dv = \int_{\mathbb{R}^d} g \left(\log g - \frac{d}{2} \log(2\pi) \right) d\mu \geq -\frac{d}{2} \log(2\pi).$$

By the Cauchy-Schwarz inequality, $|\mathbf{u}|^2 \leq \int_{\mathbb{R}^d} |v|^2 f dv$ and $\int_{\mathbb{R}^d} |v|^2 f dv \leq \sqrt{\int_{\mathbb{R}^d} |v|^4 f dv}$, and we deduce that

$$\mathcal{F}[f] \geq -\frac{d}{2} \log(2\pi)D + \frac{\alpha}{4} X^2 - \frac{D + \alpha}{2} X \quad \text{with} \quad X := \sqrt{\int_{\mathbb{R}^d} |v|^4 f dv}.$$

A minimization of the r.h.s. with respect to $X > 0$ shows that $\mathcal{F}[f] \geq -\frac{(D + \alpha)^2}{4\alpha} - \frac{d}{2} \log(2\pi)D$ while the inequality provides the bound on X . \square

4.2. The minimizers of the free energy

COROLLARY 4.1. *Let $d \geq 1$ and $\alpha > 0$. The free energy \mathcal{F} as defined by (1.3) has a unique nonnegative minimizer with unit mass, $f_{\mathbf{0}}$, if $D \geq D_*$. Otherwise, if $D < D_*$, we have*

$$\min \mathcal{F}[f] = \mathcal{F}[f_{\mathbf{u}}] < \mathcal{F}[f_{\mathbf{0}}]$$

for any $u \in \mathbb{R}^d$ such that $|\mathbf{u}| = u(D)$. The above minimum is taken on all nonnegative functions in $L^1(\mathbb{R}^d, (1 + |v|^4) dv)$ such that $\int_{\mathbb{R}^d} f dv = 1$.

Proof. Any minimizing sequence convergence is relatively compact in $L^1(\mathbb{R}^d, dv)$ by the Dunford-Pettis theorem, $f \mapsto \mathbf{u}_f$ is relatively compact and the existence of a minimizer follows by lower semi-continuity. \square

4.3. Proof of Theorem 1.1

By Corollary 4.1, $f_{\mathbf{0}}$ is the unique minimizer if and only if $D \geq D_*$. It is moreover linearly stable by Lemma 3.1. Otherwise $f_{\mathbf{u}}$ with $|\mathbf{u}| = u(D)$ is a minimizer of \mathcal{F} and it is unique up to a rotation. Combined with the results of Corollary 2.1, this completes the proof of Theorem 1.1. \square

4.4. Stability of the polarized stationary solution

Another interesting consequence of Corollary 4.1 is the linear stability of \mathcal{F} around $f_{\mathbf{u}}$ when $D < D_*$.

LEMMA 4.1. *Let $D \in (0, D_*)$ and $\mathbf{u} \in \mathbb{R}^d$ such that $|\mathbf{u}| = u(D)$. On the space of the functions $g \in L^2(f_{\mathbf{u}} dv)$ such that $\int_{\mathbb{R}^d} g f_{\mathbf{u}} dv = 0$, $Q_{1, \mathbf{u}}$ is a nonnegative quadratic form.*

The proof is straightforward as, in the range $D < D_*$, $f_{\mathbf{0}}$ is not a minimizer of \mathcal{F} and the minimum of \mathcal{F} is achieved by any $f_{\mathbf{u}}$ with $|\mathbf{u}| = u(D)$. Details are left to the reader.

4.5. An exponential rate of convergence for radially symmetric solutions

PROPOSITION 4.3. *Let $\alpha > 0$, $D > 0$ and consider a solution $f \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^d))$ of (1.1) with radially symmetric initial datum $f_{\text{in}} \in L^1_+(\mathbb{R}^d)$ such that $\mathcal{F}[f_{\text{in}}] < \infty$. Then (1.6) holds for some $\lambda > 0$.*

Proof. According to Proposition 4.1, we know that

$$\frac{d}{dt} (\mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_{\mathbf{0}}]) \leq -\mathcal{I}[f(t, \cdot)]$$

where \mathcal{I} defined by (1.5) and $\mathbf{u}_f = \mathbf{0}$ because the radial symmetry is preserved by the evolution. We have a logarithmic Sobolev inequality

$$\int_{\mathbb{R}^d} \left| \nabla \log \left(\frac{f}{f_{\mathbf{0}}} \right) \right|^2 f dv \geq \mathcal{K}_{\mathbf{0}} \int_{\mathbb{R}^d} f \log \left(\frac{f}{f_{\mathbf{0}}} \right) dv = \mathcal{F}[f] - \mathcal{F}[f_{\mathbf{0}}] \quad (4.1)$$

for some constant $\mathcal{K}_{\mathbf{0}} > 0$. This inequality holds for the same reason as for the Poincaré inequality (3.2): since φ_{α} can be seen as a uniformly strictly convex potential perturbed by a bounded perturbation, it follows from the carré du champ method and the Holley-Stroock lemma that $\mathcal{K}_{\mathbf{0}}$ is a positive constant. Hence

$$\frac{d}{dt} (\mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_{\mathbf{0}}]) \leq -\frac{\mathcal{K}_{\mathbf{0}}}{D} \int_{\mathbb{R}^d} f \log \left(\frac{f}{f_{\mathbf{0}}} \right) dv = -\frac{\mathcal{K}_{\mathbf{0}}}{D} (\mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_{\mathbf{0}}])$$

and we conclude that

$$\mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_0] \leq (\mathcal{F}[f_{\text{in}}] - \mathcal{F}[f_0]) e^{-\lambda t}$$

with $\lambda = \mathcal{K}_0/D$. The fact that $\mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_0] \geq 0$ is a consequence of Corollary 4.1. \square

4.6. Continuity and convergence of the velocity average

PROPOSITION 4.4. *Let $\alpha > 0$, $D > 0$ and consider a solution $f \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^d))$ of (1.1) with initial datum $f_{\text{in}} \in L^1_+(\mathbb{R}^d)$ such that $\mathcal{F}[f_{\text{in}}] < \infty$. Then $t \mapsto \mathbf{u}_f(t)$ is a Lipschitz continuous function on \mathbb{R}^+ such that $\lim_{t \rightarrow +\infty} \mathbf{u}_f(t) = \mathbf{0}$ if $D \geq D_*$ and $\lim_{t \rightarrow +\infty} \mathbf{u}_f(t) = \mathbf{u}$ with either $\mathbf{u} = \mathbf{0}$ or $|\mathbf{u}| = u(D)$ if $D \in (0, D_*)$.*

Proof. Using (1.1), a straightforward computation shows that

$$\frac{d\mathbf{u}_f}{dt} = -\alpha \int_{\mathbb{R}^d} v (|v|^2 - 1) f dv$$

where the right hand side is bounded by Hölder interpolations using Propositions 4.1 and 4.2. By Proposition 4.2 and Hölder's inequality, we also know that \mathbf{u}_f is bounded.

We have a *logarithmic Sobolev inequality* analogous to (4.1) if we consider the relative entropy with respect to the *non-equilibrium Gibbs state* G_f defined by (1.4) instead of the relative entropy with respect to f_0 : for some constant $\mathcal{K} > 0$,

$$\int_{\mathbb{R}^d} \left| \nabla \log \left(\frac{f}{G_f} \right) \right|^2 f dv \geq \mathcal{K} \int_{\mathbb{R}^d} f \log \left(\frac{f}{G_f} \right) dv = \mathcal{F}[f] - \mathcal{F}[G_f].$$

By the *Csiszár-Kullback inequality*

$$\int_{\mathbb{R}^d} f \log \left(\frac{f}{G_f} \right) dv \geq \frac{1}{4} \|f - G_f\|_{L^1(\mathbb{R}^d)}^2, \quad (4.2)$$

we end up with the fact that $\lim_{t \rightarrow +\infty} \int_t^{+\infty} (\int_{\mathbb{R}^d} |f - G_f| dv)^2 ds = 0$. Using Hölder's inequality

$$\left| \int_{\mathbb{R}^d} v (f - G_f) dv \right| \leq \left(\int_{\mathbb{R}^d} |f - G_f| dv \right)^{3/4} \left(\int_{\mathbb{R}^d} |v|^4 (f + G_f) dv \right)^{1/4}$$

the decay of $\mathcal{F}[f(t, \cdot)]$ and Proposition 4.2, we learn that $\lim_{t \rightarrow +\infty} \int_{\mathbb{R}^d} v (f - G_f) dv = 0$. Let $\mathcal{C}(u) := \int_{\mathbb{R}^d} e^{-\frac{1}{D}(\phi_\alpha(v) - u v_1)} dv$. By definition of \mathcal{H} , we have that

$$\int_{\mathbb{R}^d} v (f - G_f) dv = \mathbf{u}_f - \int_{\mathbb{R}^d} v G_f dv = \int_{\mathbb{R}^d} (\mathbf{u}_f - v) G_f dv = -\frac{\mathcal{H}(u)}{\mathcal{C}(u)} \frac{\mathbf{u}_f}{|\mathbf{u}_f|} \quad \text{with } u = |\mathbf{u}_f|.$$

Since \mathbf{u}_f is bounded, $\mathcal{C}(u)$ is uniformly bounded by some positive constant and we deduce that

$$\lim_{t \rightarrow +\infty} \mathcal{H}(|\mathbf{u}_f|) = 0.$$

\square

5. Large time asymptotic behaviour in the isotropic case

In this section, our main goal is to prove Theorem 1.2. In this section, we shall assume that $D > D_*$.

5.1. A non-local scalar product for the linearized evolution operator

We adapt the strategy of [5] to (1.1). With $\mathbf{v}_g = \frac{1}{D} \int_{\mathbb{R}^d} v g f_0 dv$ as in Section 3,

$$\langle g_1, g_2 \rangle := D \int_{\mathbb{R}^d} g_1 g_2 f_0 dv - D^2 \mathbf{v}_{g_1} \cdot \mathbf{v}_{g_2} \quad (5.1)$$

is a scalar product on the space $\mathcal{X} := \{g \in L^2(f_0 dv) : \int_{\mathbb{R}^d} g f_0 dv = 0\}$ by Lemma 3.1 because $\langle g, g \rangle = Q_{1,0}[g]$. Let us recall that f_0 depends on D and, as a consequence, also $D\mathbf{v}_g$. Equation (1.1) means

$$\frac{\partial f}{\partial t} = \nabla \cdot \left(D \nabla f + (v - \mathbf{u}_f + \nabla \phi_\alpha) f \right)$$

and $D \nabla f_0 = -(v + \nabla \phi_\alpha) f_0$. Hence (1.1) is rewritten in terms of $f = f_0(1 + g)$ as

$$f_0 \frac{\partial g}{\partial t} = D \nabla \cdot \left((\nabla g - \mathbf{v}_g) f_0 - \mathbf{v}_g g f_0 \right)$$

using $\mathbf{u}_f = D\mathbf{v}_g$, that is,

$$\frac{\partial g}{\partial t} = \mathcal{L}g - \mathbf{v}_g \cdot \left(D \nabla g - (v + \nabla \phi_\alpha) g \right) \quad \text{with} \quad \mathcal{L}g = D \Delta g - (v + \nabla \phi_\alpha) \cdot (\nabla g - \mathbf{v}_g) \quad (5.2)$$

and collect some basic properties of \mathcal{X} endowed with the scalar product $\langle \cdot, \cdot \rangle$ and \mathcal{L} considered as an operator on \mathcal{X} .

LEMMA 5.1. *Assume that $D > D_*$ and $\alpha > 0$. Let us consider the scalar product defined by (5.1) on \mathcal{X} . The norm $g \mapsto \sqrt{\langle g, g \rangle}$ is equivalent to the standard norm on $L^2(f_0 dv)$ according to*

$$\eta(D) \int_{\mathbb{R}^d} g^2 f_0 dv \leq \langle g, g \rangle \leq D \int_{\mathbb{R}^d} g^2 f_0 dv \quad \forall g \in \mathcal{X}. \quad (5.3)$$

Here η is as in (3.1). The linearized operator \mathcal{L} is self-adjoint on \mathcal{X} with the scalar product defined by (5.1) in the sense that $\langle g_1, \mathcal{L}g_2 \rangle = \langle \mathcal{L}g_1, g_2 \rangle$ for any $g_1, g_2 \in \mathcal{X}$, and such that

$$-\langle g, \mathcal{L}g \rangle = Q_{2,0}[g]. \quad (5.4)$$

Proof. Inequality (5.3) is a straightforward consequence of Definition (5.1) and (3.1). The self-adjointness of \mathcal{L} is a consequence of elementary computations. By starting with

$$\mathcal{L}g_1 = \left[D \Delta g_1 - (v + \nabla \phi_\alpha) \cdot \nabla g_1 \right] + (v + \nabla \phi_\alpha) \cdot \mathbf{v}_{g_1},$$

we first observe that $\int_{\mathbb{R}^d} [D \Delta g_1 - (v + \nabla \phi_\alpha) \cdot \nabla g_1] g_2 f_0 dv = -D \int_{\mathbb{R}^d} \nabla g_1 \cdot \nabla g_2 f_0 dv$ and, as a consequence (take $g_2 = v_i$ for some $i = 1, 2, \dots, d$), $\mathbf{v}_{\mathcal{L}g_1} = \mathbf{v}_{g_1} - \int_{\mathbb{R}^d} \nabla g_1 f_0 dv$. Hence

$$-\langle \mathcal{L}g_1, g_2 \rangle = D^2 \int_{\mathbb{R}^d} (\nabla g_1 - \mathbf{v}_{g_1}) \cdot (\nabla g_2 - \mathbf{v}_{g_2}) dv,$$

which proves the self-adjointness of \mathcal{L} and Identity (5.4). \square

The scalar product $\langle \cdot, \cdot \rangle$ is well adapted to the linearized evolution operator in the sense that a solution of the *linearized equation*

$$\frac{\partial g}{\partial t} = \mathcal{L}g \quad (5.5)$$

with initial datum $g_0 \in \mathcal{X}$ is such that

$$\frac{1}{2} \frac{d}{dt} Q_{1,0}[g] = \frac{1}{2} \frac{d}{dt} \langle g, g \rangle = \langle g, \mathcal{L}g \rangle = -Q_{2,0}[g]$$

and has exponential decay. According to Proposition 3.1, we know that

$$\langle g(t, \cdot), g(t, \cdot) \rangle = \langle g_0, g_0 \rangle e^{-2\mathcal{C}_D t} \quad \forall t \geq 0.$$

5.2. Proof of Theorem 1.2

Let us consider the nonlinear term and prove that a solution g of (5.2) has the same asymptotic decay rate as a solution of the linearized equation (5.5). By rewriting (5.2) as

$$f_0 \frac{\partial g}{\partial t} = D \nabla \cdot \left((\nabla g - \mathbf{v}_g) f_0 \right) - D \mathbf{v}_g \cdot \nabla (g f_0)$$

with $f = f_0(1+g)$ and using $\int_{\mathbb{R}^d} g f_0 dv = 0$, we find that

$$\frac{1}{2} \frac{d}{dt} Q_{1,0}[g] + Q_{2,0}[g] = D^2 \mathbf{v}_g \cdot \int_{\mathbb{R}^d} g (\nabla g - \mathbf{v}_g) f_0 dv.$$

Using $\mathbf{u}_f = D \mathbf{v}_g$, by the Cauchy-Schwarz inequality and (3.1), we obtain

$$\left(\int_{\mathbb{R}^d} g (\nabla g - \mathbf{v}_g) f_0 dv \right)^2 \leq \int_{\mathbb{R}^d} g^2 f_0 dv \int_{\mathbb{R}^d} |\nabla g - \mathbf{v}_g|^2 f_0 dv \leq \frac{Q_{1,0}[g]}{\eta(D)} \frac{Q_{2,0}[g]}{D^2}.$$

After taking into account Proposition 3.1, we have

$$\frac{d}{dt} Q_{1,0}[g] \leq -2 \left(1 - |\mathbf{u}_f(t)| \sqrt{\frac{\mathcal{C}_D}{\eta(D)}} \right) Q_{1,0}[g].$$

By Proposition 4.4, we know that $\lim_{t \rightarrow +\infty} |\mathbf{u}_f(t)| = 0$, which proves that

$$\limsup_{t \rightarrow +\infty} e^{2(\mathcal{C}_D - \varepsilon)t} Q_{1,0}[g(t, \cdot)] < +\infty \quad (5.6)$$

for any $\varepsilon \in (0, \mathcal{C}_D)$. After observing that $f \log(f/f_0) - (f - f_0) \leq \frac{1}{2}(f - f_0)^2/f_0$, this completes the proof of Theorem 1.2. \square

5.3. A sharp rate of convergence

We know from Proposition 3.1 that $Q_{2,0}[g] \geq \mathcal{C}_D Q_{1,0}[g]$ for any $g \in H^1(\mathbb{R}^d, f_0 dv)$ such that $\int_{\mathbb{R}^d} g f_0 dv = 0$. At no cost, we can assume that \mathcal{C}_D is the optimal constant.

THEOREM 5.1. *For any $d \geq 1$ and any $\alpha > 0$, if $D > D_*$, then the result of Theorem 1.2 holds with optimal rate $\lambda = 2\mathcal{C}_D$.*

Proof. We have to prove that 5.6 holds with $\varepsilon = 0$. By definition of \mathbf{u}_f , we have that

$$|\mathbf{u}_f|^2 = \left(\int_{\mathbb{R}^d} v (f - f_0) dv \right)^2 \leq \int_{\mathbb{R}^d} g^2 f_0 dv \int_{\mathbb{R}^d} |v|^2 f_0 dv$$

where $g := (f - f_0)/f_0$. This guarantees that $|\mathbf{u}_f(t)| \leq c\sqrt{\eta(D)}\mathcal{C}_D e^{-\lambda t/2}$. Then the function $y(t) := Q_{1,0}[g(t, \cdot)]$ obeys to the differential inequality

$$y' \leq -2\mathcal{C}_D \left(1 - ce^{-\lambda t/2}\right) y$$

and we deduce as in Section 5.2 that $\limsup_{t \rightarrow +\infty} e^{2\mathcal{C}_D t} y(t)$ is finite by a Grönwall estimate. This rate is optimal as shown by using test functions based on perturbations of f_0 . \square

6. Large time asymptotic behaviour in the polarized case

In this section, we shall assume that $0 < D < D_*$. The situation is more delicate than in the isotropic case $D > D_*$, as several asymptotic behaviours can occur.

6.1. Symmetric and non-symmetric stationary states

By perturbation of f_0 , we know that the set of the functions f such that $\mathcal{F}[f] < \mathcal{F}[f_0]$ is non-empty. Notice that the minimum of \mathcal{F} on radial functions is achieved by f_0 . It follows that any function f such that $\mathcal{F}[f] < \mathcal{F}[f_0]$ is non-radial.

LEMMA 6.1. *For any $d \geq 1$ and any $\alpha > 0$, if $D < D_*$, then for any solution $f \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^d))$ of (1.1) with initial datum $f_{\text{in}} \geq 0$ of mass 1 such that $\mathcal{F}[f_{\text{in}}] < \mathcal{F}[f_0]$. Then $\lim_{t \rightarrow +\infty} |\mathbf{u}_f(t)| = u(D)$ and $\lim_{t \rightarrow +\infty} \mathcal{F}[f(t, \cdot)] = \mathcal{F}[f_{\mathbf{u}}]$ for some $\mathbf{u} \in \mathbb{R}^d$ such that $|\mathbf{u}| = u(D)$ and*

$$f(t+n, \cdot) \longrightarrow f_{\mathbf{u}} \quad \text{in } L^1(\mathbb{R}^+ \times \mathbb{R}^d) \quad \text{as } n \rightarrow +\infty.$$

Proof. We reconsider the proof of Proposition 4.4. Since $\mathbf{u} = \mathbf{0}$ is forbidden by Proposition 4.1 and $t \mapsto \mathbf{u}_f(t)$ is a converging Lipschitz function, there exists a unique limit \mathbf{u} such that $|\mathbf{u}| = u(D)$. The convergence result follows from the logarithmic Sobolev inequality and the Csiszár-Kullback inequality (4.2). \square

6.2. An exponential rate of convergence for partially symmetric solutions

Let us start with a simple case, which is to some extent the analogous of the case of Proposition 4.3 in the polarized case.

PROPOSITION 6.1. *Let $\alpha > 0$, $D > 0$ and consider a solution $f \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^d))$ of (1.1) with initial datum $f_{\text{in}} \in L^1_+(\mathbb{R}^d)$ such that $\mathcal{F}[f_{\text{in}}] < \mathcal{F}[f_0]$ and $\mathbf{u}_{f_{\text{in}}} = (u, 0 \dots 0)$ for some $u \neq 0$. We further assume that $f_{\text{in}}(v_1, v_2, \dots, v_{i-1}, v_i, \dots) = f_{\text{in}}(v_1, v_2, \dots, v_{i-1}, -v_i, \dots)$ for any $i = 2, 3, \dots, d$. Then (1.6) holds with $\lambda = \mathcal{C}_D(1 - \kappa(D)) > 0$, with the notations of Proposition 3.1.*

Here we assume that $f_{\text{in}}(v_1, v_2, \dots, v_{i-1}, v_i, \dots)$ is even with respect to all coordinate of index $i \geq 2$, so that $\mathbf{u}[f] = \mathbf{0}$ or $\mathbf{u}[f] = (\pm u(D), 0 \dots 0)$ at any time $t \geq 0$.

Proof. According to Proposition 4.4, we know that \mathbf{u}_f is continuous. On the other hand, if $\mathbf{u}_f = \mathbf{0}$, then

$$\mathcal{F}[f] - \mathcal{F}[f_0] = \int_{\mathbb{R}^d} f \log \left(\frac{f}{f_0} \right) dv = \int_{\mathbb{R}^d} \frac{f}{f_0} \log \left(\frac{f}{f_0} \right) f_0 dv \geq X \log X \Big|_{X = \int_{\mathbb{R}^d} f dv} = 0$$

by Jensen's inequality, a contradiction with the assumption that $\mathcal{F}[f_{\text{in}}] < \mathcal{F}[f_0]$ and Proposition 4.1. Hence $\mathbf{u} = \mathbf{u}[f]$ is constant and we can reproduce with $Q_{1,\mathbf{u}}[n]$ the proof done for $Q_{1,0}[n]$ in Section 5. \square

6.3. Convergence to a polarized stationary state

To study the rate of convergence towards the stationary solutions $f_{\mathbf{u}}$ with $\mathbf{u} \neq \mathbf{0}$ in the range $D \in (0, D_*)$, we face a severe difficulty if \mathbf{u}_f converges tangentially to the set $u(D)\mathbb{S}^{d-1}$ of admissible velocities for stationary solutions. Otherwise we obtain an exponential rate of convergence as in Theorem 1.2.

PROPOSITION 6.2. *Assume that $d \geq 2$, $\alpha > 0$ and $D \in (0, D_*)$. Let us consider a solution f of (1.1) with nonnegative initial datum f_{in} of mass 1 such that $\mathcal{F}[f_{\text{in}}] < \mathcal{F}[f_0]$ and define $\mathbf{u} = \lim_{t \rightarrow +\infty} \mathbf{u}_f(t)$. If $|(\mathbf{u}_f - \mathbf{u}) \cdot \mathbf{u}| \geq \varepsilon u(D) |\mathbf{u}_f - \mathbf{u}|$ for some $\varepsilon > 0$ and $t > 0$ large enough, then there are two positive constants C , λ and some $\mathbf{u} \in \mathbb{R}^d$ such that*

$$0 \leq \mathcal{F}[f(t, \cdot)] - \mathcal{F}[f_{\mathbf{u}}] \leq C e^{-\lambda t} \quad \forall t \geq 0.$$

Proof. Let us consider $\mathbf{u} = \lim_{t \rightarrow +\infty} \mathbf{u}_f(t)$ as in Proposition 4.4. We adapt the setting of Section 5.2 to $g = (f - f_{\mathbf{u}})/f_{\mathbf{u}}$ and get that

$$\frac{1}{2} \frac{d}{dt} Q_{1,\mathbf{u}}[g] + Q_{2,\mathbf{u}}[g] = D^2 \mathbf{v}_g \cdot \int_{\mathbb{R}^d} g(\nabla g - \mathbf{v}_g) f_{\mathbf{u}} dv.$$

With $Z(t) := \mathcal{C}_D (1 - \kappa(D)) \frac{(\mathbf{v}_g \cdot \mathbf{u})^2}{|\mathbf{v}_g|^2 |\mathbf{u}|^2}$, we can rewrite Proposition 3.1 and the estimate of the nonlinear term as

$$Q_{2,\mathbf{u}}[g] \geq Z(t) Q_{1,\mathbf{u}}[g] \quad \text{and} \quad D^2 \mathbf{v}_g \cdot \int_{\mathbb{R}^d} g(\nabla g - \mathbf{v}_g) f_{\mathbf{u}} dv \leq D |\mathbf{v}_g| \frac{\sqrt{Q_{1,\mathbf{u}}[g] Q_{2,\mathbf{u}}[g]}}{\sqrt{\eta(D)}}$$

By assumption, $Z(t) \geq \mathcal{C}_D (1 - \kappa(D)) \varepsilon^2$. The conclusion follows as in Section 5.2. \square

Appendix A. Some additional properties of D_* .

In this appendix, we collect some plots which illustrate Section 2 and state related qualitative properties of D_* .

PROPOSITION A.1. *For any $\alpha > 0$ and $d \geq 1$, the critical value $D_* = D_*(\alpha, d)$ is monotone decreasing as a function of d , such that*

$$\frac{1}{d+2} < D_*(\alpha, d) < \frac{1}{d}$$

with lower and upper bounds achieved respectively as $\alpha \rightarrow 0_+$ and $\alpha \rightarrow +\infty$.

Proof. The monotonicity with respect to d can be read from

$$h_d(D) - h_{d+1}(D) = \int_0^\infty s^{d+1} (s^2 - 1)^2 e^{-\frac{\phi_\alpha}{D}} ds > 0.$$

The lower bound is a consequence of

$$\int_0^\infty (s^{d+1} - s^{d+3}) e^{-\frac{1}{2}(d+2)s^2} ds = 0.$$

As for the upper bound, for any $D > 0$, by considering the derivatives with respect to α of j_{d+1} and j_{d-1} as defined in (2.1), we notice that

$$\frac{j_{d+1}}{j_{d-1}} \sim \frac{2j_{d+3} - j_{d+5}}{2j_{d+1} - j_{d+3}} \sim \frac{\frac{\alpha+1}{\alpha} j_{d+3} - \frac{d+2}{\alpha} D j_{d+1}}{2j_{d+1} - j_{d+3}}$$

by L'Hôpital's rule as $\alpha \rightarrow +\infty$. We recall that $j_{d+1}(D) = j_{d+3}(D)$ at $D = D_*$. By letting $\alpha \rightarrow +\infty$ with $D = D_*$, we conclude that $j_{d+1}/j_{d-1} \rightarrow 1$. On the other hand (2.3) with $n = d - 2$ means that $j_{d+1}(D_*) = dD_*j_{d-1}(D_*)$, from which we conclude that $\lim_{\alpha \rightarrow +\infty} D_*(\alpha, d) = 1/d$. \square

We conclude this appendix by computations of D_* for specific values of the parameters.

- If $d=1$, $\alpha=2$, D_* solves $(1-4D)I_{-1/4}(\frac{1}{16D}) + (1+4D)I_{1/4}(\frac{1}{16D}) + I_{3/4}(\frac{1}{16D}) + I_{5/4}(\frac{1}{16D}) = 0$ where I_γ denotes the modified Bessel function of the first kind. Numerically, we find that $D_* \approx 0.529$ matches [1, Fig. 1, p. 4].
- If $d=2$, $\alpha=2$, we remind that $D_* \approx 0.354$: see Fig. A.2.
- If $d=2$, $\alpha=4$, $D_* \approx 0.398$ solves $(16\Gamma(\frac{3}{2}, \frac{9}{16D}) - 16\sqrt{\pi})D - 8\Gamma(1, \frac{9}{16D})\sqrt{D} + 6\sqrt{\pi} - 3\Gamma(\frac{1}{2}, \frac{9}{16D}) = 0$.

For further numerical examples, we refer the reader to [1, 12].

Acknowledgments This work has been supported by the Project EFI ANR-17-CE40-0030 of the French National Research Agency.

© 2019 by the author. This paper may be reproduced, in its entirety, for non-commercial purposes.

REFERENCES

- [1] BARBARO, A. B. T., CAÑIZO, J. A., CARRILLO, J. A., AND DEGOND, P. Phase transitions in a kinetic flocking model of Cucker-Smale type. *Multiscale Model. Simul.* 14, 3 (2016), 1063–1088. [1, A](#)
- [2] BARBARO, A. B. T., AND DEGOND, P. Phase transition and diffusion among socially interacting self-propelled agents. *Discrete Contin. Dyn. Syst. Ser. B* 19, 5 (2014), 1249–1278. [1](#)
- [3] BOLLEY, F., CAÑIZO, J. A., AND CARRILLO, J. A. Stochastic mean-field limit: non-Lipschitz forces and swarming. *Math. Models Methods Appl. Sci.* 21, 11 (2011), 2179–2210. [1](#)
- [4] BOSTAN, M., AND CARRILLO, J. A. Asymptotic fixed-speed reduced dynamics for kinetic equations in swarming. *Math. Models Methods Appl. Sci.* 23, 13 (2013), 2353–2393. [1](#)
- [5] CAMPOS, J. F., AND DOLBEAULT, J. Asymptotic estimates for the parabolic-elliptic Keller-Segel model in the plane. *Comm. Partial Differential Equations* 39, 5 (2014), 806–841. [1, 5.1](#)
- [6] CARRILLO, J. A., MCCANN, R. J., AND VILLANI, C. Kinetic equilibration rates for granular media and related equations: entropy dissipation and mass transportation estimates. *Rev. Mat. Iberoamericana* 19, 3 (2003), 971–1018. [4.1](#)
- [7] CUCKER, F., AND MORDECKI, E. Flocking in noisy environments. *J. Math. Pures Appl. (9)* 89, 3 (2008), 278–296. [1](#)
- [8] CUCKER, F., AND SMALE, S. Emergent behavior in flocks. *IEEE Trans. Automat. Control* 52, 5 (2007), 852–862. [1](#)
- [9] DEGOND, P., FROUVELLE, A., AND LIU, J.-G. Macroscopic limits and phase transition in a system of self-propelled particles. *J. Nonlinear Sci.* 23, 3 (2013), 427–456. [1](#)
- [10] DEGOND, P., FROUVELLE, A., AND LIU, J.-G. Phase transitions, hysteresis, and hyperbolicity for self-organized alignment dynamics. *Arch. Ration. Mech. Anal.* 216, 1 (2015), 63–115. [1](#)
- [11] FROUVELLE, A., AND LIU, J.-G. Dynamics in a kinetic model of oriented particles with phase transition. *SIAM J. Math. Anal.* 44, 2 (2012), 791–826. [1](#)
- [12] TUGAUT, J. Phase transitions of McKean-Vlasov processes in double-wells landscape. *Stochastics* 86, 2 (2014), 257–284. [1, 1, A](#)
- [13] VICSEK, T., CZIRÓK, A., BEN-JACOB, E., COHEN, I., AND SHOCHET, O. Novel type of phase transition in a system of self-driven particles. *Phys. Rev. Lett.* 75 (Aug 1995), 1226–1229. [1](#)

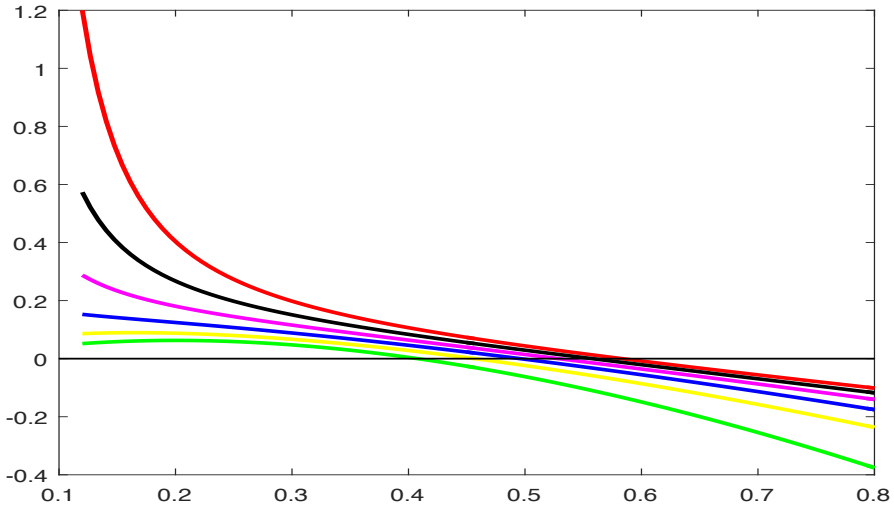


Fig. A.1: Plot of $h_d(D)$ against D when $d=1$ with $\alpha=0.5, 1, \dots, 3.0$.

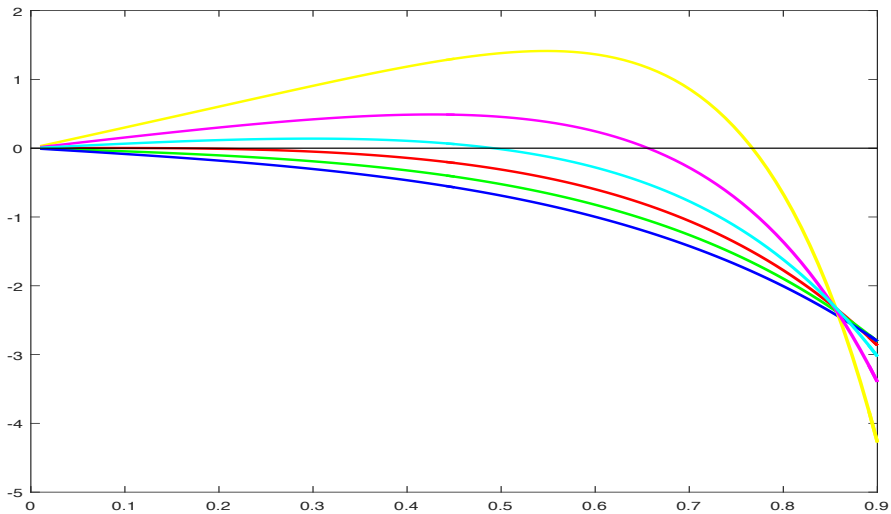


Fig. A.2: Plot of $u \mapsto \mathcal{H}(u)$ when $d=2$, $\alpha=2$, and $D=0.2, 0.25, \dots, 0.45$. In this particular case, $D_* \approx 0.354$ solves $(8\Gamma(\frac{3}{2}, \frac{1}{8D}) - 8\sqrt{\pi})D - \Gamma(\frac{1}{2}, \frac{1}{8D}) + 2\sqrt{\pi} = 0$.