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# Separable resolution-of-the-identity with all-electron Gaussian bases: Application to cubic-scaling RPA

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## ABSTRACT

We explore a separable resolution-of-the-identity (RI) formalism built on quadratures over limited sets of real-space points designed for all-electron calculations. Our implementation preserves, in particular, the use of common atomic orbitals and their related auxiliary basis sets. The setup of the present density fitting scheme, i.e., the calculation of the system specific quadrature weights, scales cubically with respect to the system size. Extensive accuracy tests are presented for the Fock exchange and MP2 correlation energies. We finally demonstrate random phase approximation (RPA) correlation energy calculations with a scaling that is cubic in terms of operations, quadratic in memory, with a small crossover with respect to our standard RI-RPA implementation.

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## I. INTRODUCTION

The resolution-of-the-identity (RI)<sup>1–7</sup> stands as a central technique in quantum chemistry, relying on the expansion of  $\phi_n\phi_m$  codensities over an auxiliary atomic basis set  $\{\beta\}$  that scales linearly with the number of atoms. Even though the auxiliary basis sets are typically three times larger than the corresponding atomic orbital (AO) basis sets supporting the  $\phi_n$  Hartree-Fock or Kohn-Sham molecular orbitals, this represents a considerable saving both in terms of memory and number of operations to be performed, which comes at the price of a moderate accuracy loss.

While RI was introduced initially to facilitate the calculation of 2-electron 4-center ( $mn|kl$ ) Coulomb integrals, operators that find an exact expression in the product space between the occupied and virtual eigenstates can also be expressed more compactly using auxiliary bases. As such, the independent-electron density-density susceptibility  $\chi_0(\mathbf{r}, \mathbf{r}'; \omega)$  is of central importance to the present study. In particular, the scaling of the related random phase approximation (RPA) approach,<sup>8–10</sup> a popular low-order perturbative approach to correlation energies beyond density functional approximations, can be reduced from  $\mathcal{O}(N^6)$  to  $\mathcal{O}(N^4)$ .<sup>11</sup>

The computational efficiency and accuracy of RI techniques strongly depend on the scheme adopted to build the appropriate coefficients expressing the  $\phi_n\phi_m$  codensities on the auxiliary basis. The original density-fitting (RI-SVS) approach<sup>2,4,5</sup> expresses these coefficients as a direct overlap  $\langle\phi_n\phi_m|\beta\rangle$ , requiring the calculation of the sparse  $\langle\alpha\alpha'|\beta\rangle$  coefficients, with  $\{\alpha\}$  being the AO basis set used to expand the molecular orbitals  $\phi_n$ . The now widely adopted Coulomb-fitting (RI-V) approximation<sup>3,6</sup> requires, on the other hand, the calculation of the denser 3-center Coulomb integrals  $(\alpha\alpha'|\beta)$ , displaying much less sparsity than  $\langle\alpha\alpha'|\beta\rangle$  integrals due to the long-range nature of the Coulomb operator. The Coulomb-fitting formalism is known to be more accurate than the density-fitting scheme for auxiliary basis sets of similar sizes,<sup>12,13</sup> bringing to the standard issue of the trade-off between accuracy and computational/memory costs. The use of short-range or attenuated Coulomb operators,<sup>14–17</sup> or corrective techniques such as “multipole-preserving” constraints to the density-fitting scheme,<sup>4,13</sup> allows one to tune the accuracy-to-cost ratio between these two standard RI approximations.

In conjunction with other powerful techniques, such as the Laplace transform (LT),<sup>18</sup> exploiting the sparsity of the RI-SVS

density-fitting  $\langle \alpha\alpha' | \beta \rangle$  coefficients in the limit of large systems was shown to allow cubic-scaling RPA calculations.<sup>19</sup> As a trade-off between accuracy and efficiency, a Coulomb-attenuated variation of the Coulomb-fitting (RI-V) RPA was recently explored to obtain a low-scaling formulation,<sup>20</sup> exploiting further the decay properties in real-space (RS) of the Laplace-transformed “pseudo” density matrices expressed in the AO basis.<sup>21–24</sup> The efficiency of this latter family of approaches depends on the electronic properties of the system of interest, which are different in nature from the sparsity associated with specific RI formalisms.

In this study, we explore an alternative approach for reducing computational and memory loads by assessing on a large set of molecules the merits of a separable RI formalism relying on a density fitting scheme over compact sets of real-space points  $\{\mathbf{r}_k\}$ . Our approach preserves the use of standard Gaussian atomic orbitals and auxiliary basis sets for all-electron calculations, targeting the accuracy of the Coulomb-fitting (RI-V) formalism. The setup of the fitting procedure scales cubically with system size. The accuracy of our approach is first validated by an extensive benchmark of the Fock exchange and MP2 correlation energies over a large set of molecules. Combined with the Laplace transform technique, and following the so-called space-time approach for calculating the susceptibility operator,<sup>25,26</sup> the calculation of the RPA correlation energy within the present real-space quadrature approach is shown to scale cubically in terms of operations and quadratically in terms of memory, without invoking any system dependent sparsity or weak localization properties.<sup>27</sup> The accuracy of the present real-space RI-RPA formalism is further shown to match that of the standard quartic-scaling Coulomb-fitting RI-RPA calculations for a large set of molecules including the oligoacenes, C<sub>60</sub>, and a larger octapeptide angiotensin II molecule (146 atoms including 71 H atoms) proposed by Eshuis and co-workers in the early days of RI-RPA implementations.<sup>11</sup>

## II. THEORY

In this section, we briefly outline the standard RI-V and RI-SVS approximations, introducing the notations used throughout the paper. We then discuss separable RIs and present our specific implementation preserving the use of standard atomic and auxiliary basis sets. The present approach relies on weighted real-space  $\delta(\mathbf{r} - \mathbf{r}_k)$ -functions to express the density fitting coefficients, relating codensities to auxiliary basis functions through real-space quadratures. The scheme to optimize the distribution of  $\mathbf{r}_k$  and related weights is presented and compared to other real-space quadrature formalisms. We demonstrate, in particular, that the computational cost associated with the setup of the present RI approach scales cubically with the system size. We show then how such a separable RI allows us to obtain cubic-scaling RPA with low crossover with respect to standard quartic RI-RPA formalism when combined with the Laplace transform technique. We conclude this section by presenting the technical details and parameters adopted in this study to perform the calculations illustrating the accuracy and scaling properties of the present approach.

### A. Standard resolution of the identity

The RI approximation<sup>1–7</sup> relies on the expansion of molecular orbital codensities  $\phi\phi'$  over an auxiliary basis set  $\{\beta\}$ , namely,

$$\begin{aligned}\phi(\mathbf{r})\phi'(\mathbf{r}) &\simeq \sum_{\beta} \mathcal{F}_{\beta}(\phi\phi') \beta(\mathbf{r}) \\ &:= \mathcal{F}(\phi\phi'; \mathbf{r}).\end{aligned}\quad (1)$$

The fit  $\mathcal{F}$  is realized through an ensemble of measures  $\{\mathcal{F}_{\beta}\}$ , mapping the  $\phi\phi'$  product-space to the  $\beta$  auxiliary subspace defined so as to scale linearly with the number of atoms. Typical examples of such procedures are the standard RI-V and RI-SVS fitting approaches that use, respectively,

$$\mathcal{F}_{\beta}^V(\phi\phi') = \sum_{\beta'} [V^{-1}]_{\beta\beta'} \langle \beta' | \phi\phi' \rangle, \quad (2)$$

$$\mathcal{F}_{\beta}^{SVS}(\phi\phi') = \sum_{\beta'} [S^{-1}]_{\beta\beta'} \langle \beta' | \phi\phi' \rangle, \quad (3)$$

where  $V$  and  $S$  represent the Coulomb  $\langle \beta | \beta' \rangle$  and overlap  $\langle \beta | \beta' \rangle$  matrices associated with the auxiliary basis set, respectively, and  $[X^{-1}]_{\beta\beta'}$  denotes the  $(\beta, \beta')$  entry of the  $X$  inverse matrix. To explicitly define our  $\langle \cdot | \cdot \rangle$  and  $(\cdot | \cdot)$  notations, we write

$$\begin{aligned}(\beta | \phi\phi') &= \iint d\mathbf{r}d\mathbf{r}' \frac{\beta(\mathbf{r})\phi(\mathbf{r}')\phi'(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}, \\ \langle \beta | \phi\phi' \rangle &= \int d\mathbf{r} \beta(\mathbf{r})\phi(\mathbf{r})\phi'(\mathbf{r}).\end{aligned}$$

As shown in Ref. 13, both fitting techniques can be combined, preserving the Coulomb-fitting RI-V approach for low angular momentum auxiliary  $\beta$  atomic orbitals. As emphasized here above, the number of  $\langle \beta | \alpha\alpha' \rangle$  overlap matrix elements in the RI-SVS approximation scales linearly with system size, offering a first strategy for reducing computational cost and memory, thanks to sparsity. On the contrary, the number of  $\langle \beta | \alpha\alpha' \rangle$  Coulomb integrals scales quadratically so that sparsity, or sparse tensor algebra, is difficult to exploit within the more accurate Coulomb-fitting (RI-V) approach.

### B. Separable RI

A separable expression for the RI can be obtained through a set of separable measures  $\{f_k\}$  on the codensities, namely,

$$\begin{aligned}\mathcal{F}_{\beta}(\phi\phi') &= \sum_k [M]_{\beta k} \langle f_k | \phi\phi' \rangle \\ &= \sum_k [M]_{\beta k} \langle f_k | \phi \rangle \langle f_k | \phi' \rangle,\end{aligned}\quad (4)$$

where the coefficients of  $M$  have yet to be defined. Trivially separable and generic measures are provided by  $\delta(\mathbf{r} - \mathbf{r}_k)$  distributions centered on a set of  $N_k$  real-space locations  $\{\mathbf{r}_k\}$ . Working with real-space (RS) measures, the  $\mathcal{F}^{RS}$  density fitting procedure then takes the simple form

$$\begin{aligned}\mathcal{F}_{\beta}^{RS}(\phi\phi') &= \sum_k [M]_{\beta k} \langle \delta(\mathbf{r} - \mathbf{r}_k) | \phi\phi' \rangle \\ &= \sum_k [M]_{\beta k} \phi(\mathbf{r}_k) \phi'(\mathbf{r}_k).\end{aligned}\quad (5)$$

The clear advantage of the separability is that the two molecular orbitals  $\phi\phi'$ , originally entangled in, e.g., the  $\mathcal{F}_{\beta}^V(\phi\phi')$  RI-V fitting coefficients through the  $\langle \beta | \phi\phi' \rangle$  Coulomb integrals, are now disentangled. This will prove crucial in the calculation of linear-response

or perturbation theory related quantities where summations over occupied/virtual pairs have to be performed as shown below in the case of the calculation of the RPA correlation energy. We emphasize, however, that, while relying on discrete values of the molecular orbitals in real-space, the present approach remains a RI in the sense that physical continuous quantities such as the codensities, linear-response operators (e.g., susceptibility), etc., are defined everywhere in space in terms of the  $\beta$  auxiliary basis functions.

Amongst existing formalisms adopting real-space quadrature strategies, several studies focused directly on 2-electron Coulomb integrals. The chain-of-sphere (COSX) seminumerical approach to exchange integrals,<sup>28</sup> building on Friesner's pioneering pseudospectral approach,<sup>29</sup> develops only one of the two codensities forming 2-electron Coulomb integrals over a real-space grid. Alternatively, the Least Square Tensor Hypercontraction (LS-THC) formalism<sup>30</sup> fully develops the 2-electron integrals as a quadrature over real-space grid points, with an  $\mathcal{O}(N^4)$  computational complexity associated with the establishment of the quadrature.

On the other hand, the Interpolative Separable Density Fitting (ISDF) from Lu and co-workers provides a  $\mathcal{O}(N^3 \log(N))$  separable fit tensor by using Fourier transform and random projection techniques to select the  $\mathbf{r}_k$  points and define the corresponding auxiliary densities,<sup>31,32</sup> with a recent extension to cubic scaling RPA<sup>33</sup> and a proof-of-concept application to a simple model system. This work adopts as well a separable form for the fit tensor [Eq. (5)] through separable measures along real-space positions, but differs in the way the real-space locations  $\mathbf{r}_k$  and their associated weights  $[M]_{\beta k}$  are constructed, leading to a  $\mathcal{O}(N^3)$  quadrature determination process. We preserve, in particular, the use of the standard auxiliary basis sets optimized by the quantum chemistry community (e.g., cc-pVTZ-RI), targeting the well-documented accuracy of the RI-V associated with such well known basis sets. Furthermore, the flexibility provided by combining self-adjusted real-space grid points with standard auxiliary basis sets allows us to improve computational efficiency while keeping the response operators limited in size for subsequent calculations. The accuracy and efficiency of our approach is further benchmarked on a large number of molecular systems.

### C. Quadrature weight determination

Assuming that the  $\mathbf{r}_k$  locations are known (see below), the determination of the  $[M]_{\beta k}$  coefficients in Eq. (5) is obtained from the following minimization condition:

$$\operatorname{argmin}_M \sum_{\rho\beta} (\mathcal{F}_\beta^{\text{RS}}(\rho) - \mathcal{F}_\beta^{\text{V}}(\rho))^2. \quad (6)$$

Namely, we aim to equate the fitting functions of the present RI-RS approach with those of the Coulomb fitting RI-V scheme. The accuracy of the present RI-RS approach is thus targeting that of the RI-V approach. The set of test codensities  $\{\rho\}$  typically spans the  $\{\alpha\} \otimes \{\alpha\}$  product-space, even though it can be adjusted depending on the problem being addressed. Interestingly, the present fitting scheme allows us to recover at a reduced cost the otherwise  $\mathcal{O}(N^4)$  LS-THC factorization (see the [supplementary material](#)). The solution of Eq. (6) can indeed be achieved with an advantageous  $\mathcal{O}(N^3)$  computational complexity, as demonstrated now.

In order to detail our fitting procedure, we adopt the following matrix notations:  $[D]_{k\rho} = \rho(r_k)$  and  $[F]_{\beta\rho} = \mathcal{F}_\beta^{\text{V}}(\rho)$ . Due to the localization properties of the atomic orbitals, the number of atomic orbital products scales linearly with system size, and thus, the number  $N_\rho$  of test codensities in the test set can be considered  $\propto N_\alpha$ . As a result, the matrices  $D$  ( $N_k \times N_\rho$ ) and  $F$  ( $N_\beta \times N_\rho$ ) as well as the matrix  $M$  ( $N_\beta \times N_k$ ) of Eq. (5) are all  $\mathcal{O}(N^2)$  tensors. Using the Frobenius norm  $\|\cdot\|_F$ , the fit equation (6) can then be formulated as

$$\operatorname{argmin}_M \left\| M \cdot D - F \right\|_F, \quad (7)$$

which leads to the standard least-square estimator

$$M = F \cdot D^\dagger \cdot (D \cdot D^\dagger)^{-1} \quad (8)$$

involving only matrix multiplications and inversions. Computation of  $(D \cdot D^\dagger)^{-1}$  could prove problematic if done explicitly. The term  $D \cdot D^\dagger$  is positive, but has not guaranteed to be definite. On the other hand, due to the large number  $N_\rho$  of test codensities, application of the standard Singular Value Decomposition (SVD) technique to extract the pseudoinverse based estimator leads to rather significant prefactors to the otherwise  $\mathcal{O}(N^3)$  pseudoinverse procedure. We take a side approach by combining simple balancing and Tikhonov  $L_2$  regularization.<sup>34</sup> We first balance the problem by normalizing the rows of  $D$ , writing  $\tilde{D} = d \cdot D$ , where  $d$  is a diagonal matrix and the diagonal terms  $[\tilde{D} \cdot \tilde{D}^\dagger]_{kk} = 1$ . The pseudoinverse is then calculated as

$$(D \cdot D^\dagger)^{-1} \simeq d \cdot (\tilde{D} \cdot \tilde{D}^\dagger + \epsilon \mathbb{I})^{-1} \cdot d,$$

where the  $L_2$  regularization parameter  $\epsilon$  is adjusted to a small value to maintain definiteness of the problem and ensure numerical stability of the inverse. We identified the value  $\epsilon = 4 \times 10^{-7}$  as a reasonable parameter for double precision arithmetic and kept this value for all the results presented in this work. The resulting final least square estimator is thus

$$M = F \cdot \tilde{D}^\dagger \cdot (\tilde{D} \cdot \tilde{D}^\dagger + \epsilon \mathbb{I})^{-1} \cdot d, \quad (9)$$

which can be computed efficiently through standard numerical inversion techniques. We emphasize that while computing the  $[M]_{\beta k}$  optimal coefficients, there is no need for keeping the 3-center Coulomb integrals or the associated  $\mathcal{F}_\beta^{\text{V}}(\alpha\alpha')$  coefficients: these can be computed once on-the-fly and discarded immediately, avoiding thus any extra memory consumption. In other terms, we never store explicitly the  $F$  and  $D$  matrices but only their  $F \cdot D^\dagger$  and  $D \cdot D^\dagger$  ( $N_k \times N_k$ ) resulting products.

### D. Real-space grids' generation

In the present approach, the optimized  $\{\mathbf{r}_k\}$  sets are generated for isolated atoms, once for every chemical species and their associated atomic basis sets. These atomic grids are then duplicated according to the molecule geometry to form the system-specific quadrature points. With Eq. (9) defining the optimal  $M$  for a given  $\{\mathbf{r}_k\}$  set, locations are adjusted so as to minimize the fit error for single atom test codensities, using the Coulomb metric

$$\operatorname{argmin}_{\{\mathbf{r}_k\}} \sum_{\rho} \left\| \mathcal{F}^{\text{RS}}(\rho) - \mathcal{F}^{\text{V}}(\rho) \right\|_{\text{V}}^2. \quad (10)$$

For the sake of keeping the optimization process relatively simple, we structured the  $\{\mathbf{r}_k\}$  set over four different shells, each one replicated with a different number of radii. The only parameters of the optimization problem are thus the number of radii and their length. The four base shells taken here and denoted  $A_1$ ,  $A_2$ ,  $A_3$ , and  $B_1$  are subsets of the Lebedev quadrature grids<sup>35</sup> (denoted here  $L_i$  for the Lebedev grid of order  $i$ ) in the sense that

$$\begin{aligned}L_3 &= A_1, \\L_5 &= A_1 \cup A_2, \\L_7 &= A_1 \cup A_2 \cup A_3, \\L_{11} &= A_1 \cup A_2 \cup A_3 \cup B_1.\end{aligned}$$

The determination of the number of radii associated with each shell has been done through experimentation until satisfying configurations were obtained. Base shells are provided in Tables S4–S7 of the [supplementary material](#), and the resulting atomic quadrature grids for atoms H, C, N, and O in Tables S8–S11 of the [supplementary material](#). Optimizing freely all  $\mathbf{r}_k$  locations and not only the radii may lead to significant improvement with respect to the grid size/accuracy ratio. Providing such grids falls, however, outside the scope of the present work.

## E. Technical details

Benchmark Hartree-Fock and MP2 calculations were performed on a standard set of 28 medium size organic molecules containing unsaturated aliphatic and aromatic hydrocarbons or heterocycles, aldehydes, ketones, amides, and nucleobases. Such a test set was originally proposed by Thiel and co-workers<sup>36</sup> for reference optical excitations calculations within, e.g., coupled cluster,<sup>36,37</sup> Time Dependent Density Functional Theory (TD-DFT)<sup>38</sup> and more recently Bethe-Salpeter<sup>39–41</sup> formalisms. We adopt the MP2/6-31Gd geometries supplied in Ref. 36.

The assessment of the scaling properties of the present real-space quadrature RPA implementation is further performed on the oligoacene family from benzene to hexacene using the B3LYP cc-pVTZ geometries available in Ref. 42, complemented by the decacene, recently observed,<sup>43</sup> and the (hypothetical) octacene, both relaxed at the B3LYP/6-31Gd level. Finally, we consider the  $C_{60}$  fullerene (B3LYP/6-311Gd geometry provided in the [supplementary material](#)) and the octapeptide angiotensin II molecule originally proposed by Eshuis and co-workers.<sup>11</sup>

All calculations are performed with input molecular orbitals generated at the (spherical) cc-pVTZ<sup>44</sup> Hartree-Fock level using the NWChem package.<sup>45</sup> The corresponding (Cartesian) cc-pVTZ-RI auxiliary basis<sup>46</sup> was adopted in all resolution-of-the-identity (RI) approaches (RI-SVS, RI-V, and real-space quadrature RI-RS). For the sake of comparison, Hartree-Fock exchange and MP2 correlation energies were calculated exactly, namely, without any RI approximation, using the NWChem package as well. All calculations are performed without any frozen-core approximation.

The set  $\{\rho\}$  of test codensities can be adjusted depending on the needs, for example, to match a specific subset of the wave function codensities. In the rest of this work, we adopt the following settings:

$$\{\rho\} = (\{\alpha\} \otimes \{\alpha'\}_{1 \leq 2} \cup \{\beta\}) \quad (11)$$

for both the single atom  $\{\mathbf{r}_k\}$  set problem and the full system optimization of  $M$  coefficients. Limiting the second  $\{\alpha'\}$  atomic orbital

(AO) basis set to  $s$ ,  $p$  and  $d$  orbitals allowed to speed up the computation while no significant change in accuracy was observed. In the minimization process, the weight on the  $s$  and  $p$  AO orbitals have also been stressed with factors 4 and 2, respectively, so as to increase focus on the low-order multipole charge and dipole component of the codensities. Inclusion of the  $\{\beta\}$  auxiliary orbitals within the test set slightly improves regularity of the errors. Such a choice is relevant for the atoms of the present set. The use of products of Kohn-Sham atomic orbitals spanning, e.g., reduced energy range test sets may lead to more accurate fits.

We adopt real-space  $\{\mathbf{r}_k\}$  sets that contain typically 320 points per C, N, and O atom and 180 for hydrogen. This size corresponds to about 3 times the size of the corresponding cc-pVTZ-RI auxiliary basis set. In this study, we do not seek to look for minimal grid sizes, showing here below that excellent accuracy and a small crossover between RI-RS and RI-V can already be obtained with such parameters. Details about the optimized real-space  $\{\mathbf{r}_k\}$  sets, optimized for the cc-pVTZ and cc-pVTZ-RI Gaussian basis sets following Eqs. (6) and (10), are provided in the [supplementary material](#).

The Laplace transform (LT) RPA correlation energy calculations are based on time and frequency grids described in the [Appendix](#) together with convergence tests for the benzene correlation energy, other molecules being reported in the [supplementary material](#). The present RI calculations, including the standard RI-V, RI-SVS and the newly developed real-space RI-RS, with and without Laplace transform, are performed with a specific pilot code building on the Coulomb integral libraries implemented in the FIESTA code.<sup>13,39,47</sup>

## III. RESULTS

### A. Assessing the accuracy of the optimized real-space grid: Fock exchange and MP2 correlation energies

As a first accuracy test of the present real-space RI implementation, namely, to assess the quality of the codensity fits, we calculate both the Fock exchange energy,

$$E_{xx} = - \sum_{ij}^{occ} (ij|ij), \quad (12)$$

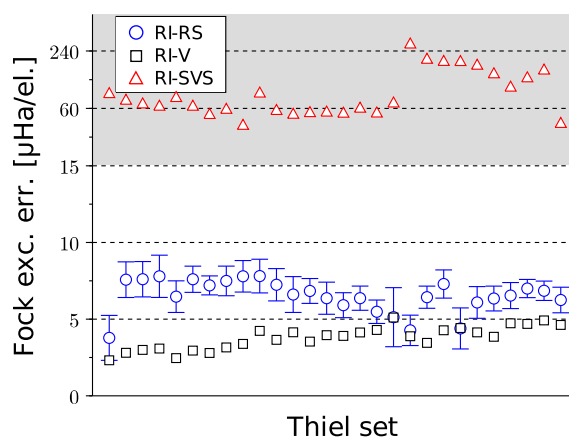
written here for a spin compensated system, and the Møller-Plesset (MP2) correlation energy,

$$E_C^{MP2} = - \sum_{ij}^{occ} \sum_{ab}^{virt} \frac{(ia|jb)[2(ia|jb) - (ib|ja)]}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j}. \quad (13)$$

The Fock exchange and MP2 energies are calculated using the RI expressions of the 2-electron integrals, namely,

$$(ia|jb)^{RI-X} = \sum_{\beta\beta'} \mathcal{F}_{\beta}^X(ia) V_{\beta\beta'} \mathcal{F}_{\beta'}^X(jb), \quad (14)$$

where X = SVS, V, or RS depending on the selected scheme. Since we want to specifically address the accuracy of the fitting technique, we avoid at that stage using extra approximations such as Laplace transform techniques. Also, contrary to the  $\{\beta\}$  Gaussian auxiliary basis, the atomic  $\{\mathbf{r}_k\}$  set is not rotationally invariant. Consequently, all the RI-RS results reported in this section are mean values obtained over 40 random orientations of each molecule, dressed with an



**FIG. 1.** Fock exchange energy error as compared to exact calculations for various RI approximations (RI-RS, RI-V, and RI-SVS). Errors are given in micro-Hartree ( $\mu\text{Ha}$ ) per electron. For RI-SVS data (red triangles), we adopt a log scale (gray shaded area). The blue circles and error bars denote the RI-RS total Fock exchange energy mean error and its associated standard deviation, respectively, computed over 40 random orientations for each molecule of the set (see text). The molecules are ordered from the left to right following the original order provided in Ref. 36 and in Table S1 of the [supplementary material](#).

“error bar” that represents the standard deviation of the corresponding distribution when the molecules are rotated.

The results, namely, the mean errors as compared to exact calculations for the 28 Thiel’s set molecules, are provided in Fig. 1 for the Fock exchange energy calculations. Clearly, the  $\{\mathbf{r}_k\}$  sets adopted in this study provide mean errors that are of the same magnitude as the one from the targeted RI-V approximation, both of which being significantly smaller than those obtained with the RI-SVS scheme. We recall that our real-space quadrature was optimized to reproduce the coefficients  $\mathcal{F}_\beta^V$  of the RI-V formalism [see Eq. (6)] so that it should not be expected that the RI-RS approach yields errors smaller than the RI-V.

The standard deviation per molecule with respect to orientation, as reported in Fig. 1 under the form of an error bar, is found to be well within the corresponding mean error. This ensures an accurate description, regardless of the chosen orientation of the molecule. Along with the per molecule data of Fig. 1 that focus on the sensitivity of our approach with respect to orientation, we also provide the global statistics across the whole Thiel’s molecular set in Table I. More precisely, the statistical analysis is performed over

**TABLE I.** Fock exchange and MP2 correlation error global statistics: global mean error and the corresponding standard deviation  $\sigma$  (both in  $\mu\text{Ha}/\text{electron}$ ) provided by the different RI techniques over the full Thiel’s molecular set. Mean signed and absolute errors are the same in the present case.

	Fock exchange			MP2 correlation	
	RI-SVS	RI-V	RI-RS	RI-V	RI-RS
Err	95.0	3.8	6.5	1.1	1.6
$\sigma$	58.5	0.7	1.1	0.6	0.7

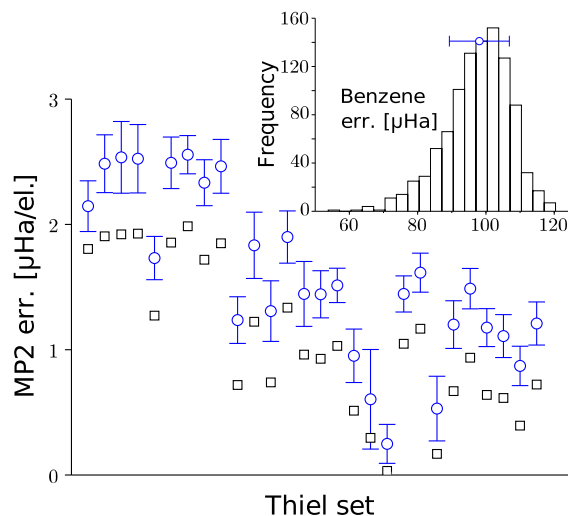
the 40 random orientations for each of the 28 molecules, namely, averaging over a 1000 calculations on medium size molecules representing a large variety of molecular families. In particular, we can see that the global mean errors and standard deviations of the RI-V and RI-RS approaches are comparable.

We now turn to the MP2 correlation energies (Fig. 2), focusing on the comparison between the present RI-RS and the targeted RI-V formalism that we adopt in the following rather than the RI-SVS scheme for the sake of accuracy. Again, we observe that the RI-RS quadrature does not degrade significantly the targeted RI-V MP2 correlation energies, with mean errors remaining lower than a few  $\mu\text{Hartree}$  per electron. Similarly, the corresponding standard deviations remain small. We provide in the inset of Fig. 2 the actual distribution of errors obtained for a 1000 independent random orientations of the benzene molecule, replicating the corresponding standard deviation error bar. We also report the global statistics obtained over the whole molecular set in Table I. Again, we observe comparable performances for RI-V and RI-RS approaches.

We can conclude from the present set of benchmark calculations that the real-space representation (RI-RS), with the adopted  $\{\mathbf{r}_k\}$  distribution size, reproduces accurately the RI-V Coulomb integrals involving codensities (products  $\phi_i\phi_j$  in the Fock exchange expression) and transition densities (products  $\phi_i\phi_a$  in the MP2 energy formula) at the core of all explicitly correlated perturbative techniques.

## B. Laplace transformed RPA

We now turn to the central application of the present study, namely, the calculation of the correlation energy within the random



**FIG. 2.** MP2 correlation energy error, in micro-Hartree ( $\mu\text{Ha}$ ) per electron, as compared to exact calculations for the standard Coulomb-fitting (RI-V) and real-space quadrature (RI-RS) approximations. The blue circles and error bars denote the RI-RS total MP2 energy mean error and its associated standard deviation, respectively, computed over 40 random orientations for each molecule of the set (see text). Inset: details of benzene RI-RS MP2 correlation energy error ( $\mu\text{Ha}$ ) over a 1000 random orientations with the corresponding standard deviation (blue segment). Reference values (no-RI) and errors are provided in Table S1 of the [supplementary material](#).

phase approximation (RPA).<sup>10,48,49</sup> We show, in particular, that the real-space quadrature RI-RS, combined with the standard Laplace transform (LT) approach,<sup>18</sup> allows reducing the scaling with system size down to  $\mathcal{O}(N^3)$ , instead of the  $\mathcal{O}(N^4)$  scaling of standard RI-RPA implementations,<sup>11</sup> without invoking any weak localization or system dependent sparsity arguments.

Following seminal papers,<sup>50–52</sup> we start with the adiabatic-connection fluctuation-dissipation theorem (ACFDT) formula to the RPA correlation energy,

$$E_C^{RPA} = \frac{1}{2\pi} \int_0^\infty d\omega \text{Tr}[\ln(1 - \chi_0(i\omega) \cdot v) + \chi_0(i\omega) \cdot v], \quad (15)$$

where  $v$  is the bare Coulomb operator and  $\chi_0(i\omega)$  is the independent-electron density-density susceptibility at imaginary frequency, that is, for closed-shell systems

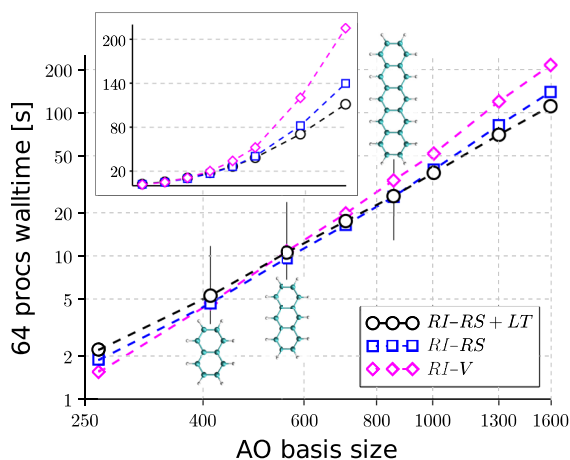
$$\chi_0(\mathbf{r}, \mathbf{r}'; i\omega) = 2 \sum_{ja} \frac{\phi_j^*(\mathbf{r})\phi_a(\mathbf{r})\phi_a^*(\mathbf{r}')\phi_j(\mathbf{r}')}{i\omega - (\epsilon_a - \epsilon_j)} + cc \quad (16)$$

$$\stackrel{RI}{\simeq} \sum_{\beta\beta'} \beta(\mathbf{r})\beta'(\mathbf{r}') \left[ 2 \sum_{ja} \frac{\mathcal{F}_\beta(\phi_j\phi_a)\mathcal{F}_{\beta'}(\phi_a\phi_j)}{i\omega - (\epsilon_a - \epsilon_j)} + cc \right] \quad (17)$$

$$:= \sum_{\beta\beta'} \beta(\mathbf{r})\beta'(\mathbf{r}') [\chi_0^{RI}(i\omega)]_{\beta\beta'}, \quad (18)$$

with  $(j, a)$  indexing (occupied/virtual) molecular eigenstates. The construction of the  $\chi_0^{RI}(i\omega)$  matrix according to Eq. (17) clearly scales as  $\mathcal{O}(N^4)$ .

To discuss such scaling properties, we compare first in Fig. 3 the total computing time for calculating the RI-RPA correlation energy within the standard Coulomb-fitting approach (RI-V) and the novel RI-RS formalism, using the acene family from benzene to



**FIG. 3.** Total walltime for the calculation of the RPA correlation energy over the acene family (benzene to decacene) using several RI schemes. We compare, in particular, the standard Coulomb fitting RI-V approach, the real-space quadrature RI-RS scheme, with and without the Laplace transform (LT) technique. The abscissa provides the size of the (cc-pVTZ) AO basis used to expand the molecular orbitals. Both axes are displayed in log scale. The crossovers between the various RI formalisms are indicated by vertical short segments with the corresponding molecules. Inset: same data points without log scales. Walltimes are given for a run on 64 processors described in Ref. 53.

decacene as a test set. Calculations are performed using the cc-pVTZ AO and cc-pVTZ-RI auxiliary basis sets, together with a 12-point quadrature rule for the imaginary frequency axis integration (see the Appendix). The corresponding correlation energies are provided in the Appendix for benzene and in the supplementary material for other acenes, demonstrating again the accuracy of the real-space approach as compared to the targeted RI-V approach. We further provide for the sake of illustration in the supplementary material the dissociation curve for the benzene dimer (cofacial or sandwich configuration) within the RI-V and RI-RS RPA formalisms, demonstrating that the roughness of the potential energy surface associated with the real-space  $\{\mathbf{r}_k\}$  distribution is marginal even as compared to the small binding energy typical of van der Waals interactions.

We observe that the RI-RS scheme walltime becomes smaller than that of our RI-V implementation for acenes larger than naphthalene. We emphasize, however, that at that stage, both RI-V and RI-RS techniques offer the very same  $\mathcal{O}(N^4)$  scaling, differing only by the expression of the  $\mathcal{F}_\beta(\phi_j\phi_a)$  coefficients. This crossover is related to the effort coming from the  $\mathcal{O}(N^4)$  computation of the full  $\mathcal{F}_\beta(\phi_j\phi_a)$  coefficient set. Without any assumption on the sparsity, the RI-RS scaling relies only on dense algebra techniques and can thus be implemented efficiently.

In order to reduce this scaling, one thus needs to avoid explicit calculation of the full  $\mathcal{F}_\beta(\phi_j\phi_a)$  coefficient set. We achieve this by evaluating the independent-electron susceptibility directly in the real-space representation before transforming it back to the normal auxiliary basis representation,

$$[\chi_0^{RI}(i\omega)]_{\beta\beta'} = \sum_{kk'} M_{\beta k} \chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\omega) M_{\beta' k'}. \quad (19)$$

The second step consists in applying the well known Laplace transform (LT) technique<sup>18</sup> so as to first compute  $\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau)$  in the time domain where its expression is separable<sup>25,26</sup> and transform it back to the frequency domain, using quadrature rules to form  $\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\omega)$  [Eq. (20)]. Such a scheme allows us to work with factorized expression of  $\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau)$  [Eq. (21)],

$$\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\omega) = \sum_{\tau} c_{\tau}(\omega) \chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau), \quad (20)$$

$$\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau) = G^<(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau) G^>(\mathbf{r}_k, \mathbf{r}_{k'}; -i\tau), \quad (21)$$

introducing the propagators of the occupied states and of the unoccupied states, respectively,

$$G^<(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau) = i \sum_j \phi_j(\mathbf{r}_k) \phi_j(\mathbf{r}_{k'}) e^{\epsilon_j \tau}, \quad (22)$$

$$G^>(\mathbf{r}_k, \mathbf{r}_{k'}; -i\tau) = -i \sum_a \phi_a(\mathbf{r}_k) \phi_a(\mathbf{r}_{k'}) e^{-\epsilon_a \tau}, \quad (23)$$

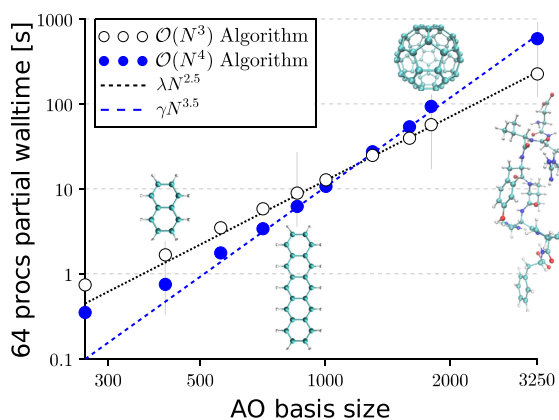
with  $\tau > 0$  and the zero of the occupied/virtual electronic energy levels taken at the Fermi level. As a result of the decoupling of occupied and virtual states, the  $G^<$  and  $G^>$  propagators can be obtained with  $\mathcal{O}(N^3)$  operations. Furthermore, the entire  $[\chi_0^{RI}(i\omega)]_{\beta\beta'}$  matrix stems from a combination of Hadamard product [Eq. (21)] and standard matrix operations [Eqs. (19) and (20)], yielding an overall  $\mathcal{O}(N^3)$  process.

We can now report in Fig. 3 the full calculation walltime associated with the RI-RS+LT approach. Laplace transform quadratures

for each of the 12 imaginary frequencies are performed with a grid of 18 imaginary times, yielding fully converged RPA energies (see the Appendix). With such running parameters, the RI-RS+LT approach becomes more efficient than the standard RI-V formalism for systems larger than anthracene, outperforming the RI-RS approach for molecules larger than pentacene.

To better assess the scaling properties associated with the resolution of Eq. (15) within standard and Laplace transformed RI-RPA approaches, we single out the corresponding computation walltimes in Fig. 4. We assume, in particular, that fitted codensities coefficients are already available in the case of standard approaches so that standard RI-V and RI-RS computational loads are equivalent. To probe larger systems, the test set is extended with the  $C_{60}$  fullerene and the original octapeptide angiotensin II molecule proposed by Eshuis and co-workers.<sup>11</sup> For the sake of information, the RI-RS setup and the Laplace transform RI-RS RPA calculation walltimes for the angiotensin II molecule are 418 s and 300 s, respectively.

Without accounting for the fitting of codensity calculation, the Laplace transform RI-RS RPA algorithm supersedes the standard RI-RS (or RI-V) RPA approach for systems larger than hexacene. This delayed crossover (as compared to Fig. 3) can be imputed to both the Laplace transform overheads and the fact that with the  $\{\mathbf{r}_k\}$  real space point set presently used,  $[\chi_0(i\omega)]_{kk'}$  matrices are roughly  $3 \times 3$  times bigger than  $[\chi_0(i\omega)]_{\beta\beta'}$  ones, leading to a  $3^3$  prefactor in the linear algebra operations. This last fact demonstrates the importance of operating on small  $\{\mathbf{r}_k\}$  quadrature sets. The additional cost coming from obtaining the fitting parameters  $\mathcal{F}_\beta(\phi_j\phi_a)$  only adds to the cost of the no-Laplace-transform RI approaches, bringing the overall crossover at the level of a pentacene molecule, as exemplified in Fig. 4. We finally observe  $\mathcal{O}(N^{3.5})$  and  $\mathcal{O}(N^{2.5})$  scaling laws (see dotted lines), indicating that the expected asymptotic  $\mathcal{O}(N^4)$  and  $\mathcal{O}(N^3)$  behaviors are not yet reached for the tested molecule sizes.



**FIG. 4.** Partial walltimes associated with the RI-RS RPA correlation energy calculations with and without Laplace transforms (LTs), removing the extra cost associated with obtaining the  $\mathcal{F}_\beta(\phi_j\phi_a)$  coefficients in the case of the no-LT approach.<sup>53</sup> The test covers the acene family plus the  $C_{60}$  fullerene and octapeptide angiotensin II molecule. The abscissa provides the size of the (cc-pVTZ) AO basis used to expand the molecular orbitals. Both axes are displayed in log scale. Dotted and dashed lines are a schematic guide to the eye for scaling properties. Walltimes are given for a run on 64 processors described in Ref. 53.

We conclude this section by mentioning that the total walltime for the (RI-RS+LT) calculation of the cc-pVTZ RPA correlation energy for  $C_{60}$  takes less than 6500 s on a single processor.

#### IV. DISCUSSION

The present implementation can be compared to the real-space-grid imaginary-time approach introduced originally in the framework of GW calculations by Rojas and co-workers<sup>25</sup> or the recent real-space-grid imaginary-time RPA implementation by Kaltak and co-workers.<sup>26</sup> In such studies, the real-space grid was obtained as the Fourier transform of the planewave basis used to expand the Bloch states in a pseudopotential or Projector Augmented Wave (PAW) framework for periodic systems. Alternatively, the work of Moussa<sup>54</sup> demonstrated as well cubic scaling for a random-phase approximation with the second-order screened exchange formalism, exploiting a real-space grid for both the primary and auxiliary basis sets, combined with nested low-rank approximations to energy denominators. The present RI-RS approach, while targeting all-electron atomic-basis calculations, preserves the use of the standard representation of molecular orbitals, related codensities, and response operators in terms of atomic orbitals and their associated auxiliary bases, adopting further standard Laplace transform techniques.

A central issue in real-space representations, in particular, when performing all-electron calculations, concerns the size of the real-space grid that strongly influences the crossover with standard RI implementations and the memory requirements. This is all the more important in the present study since we aim to calculate and store intermediate nonlocal operators such as the  $\chi^0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau)$  susceptibilities, and not only local functions such as the charge density or the DFT exchange-correlation potential and energy density. Since our real-space  $\{\mathbf{r}_k\}$  distribution must serve in a quadrature reproducing codensities involving molecular orbital products, one may expect that it should be as large as standard grids<sup>55</sup> used to represent the charge density in DFT codes. Taking as an example the Gaussian09 code, the default DFT grid involves about 7000 grid points per atom after pruning. This is consistent with the recommended “Grid3” in the original paper by Trutler and Alrichs<sup>56</sup> yielding 5980 pruned points for elements from Li to Ne and that serves as the default in Turbomole.

Such standard DFT grid sizes are much larger than the number of real-space points used in the present RS approach, roughly 180 for hydrogen and 320 per non-H atom (C, N, and O). The present  $\{\mathbf{r}_k\}$  sets were optimized so that the RI-RS scheme faithfully reproduces the standard RI-V density fitting results. Each set is composed of a number of different shells of high symmetry points, each one associated with a set of different radii. This minimization process results in a nonuniform  $\{\mathbf{r}_k\}$  distribution of real-space points. As emphasized above, we did not seek to explore here in great details the set size to accuracy ratio, for our goal was to demonstrate that one can find an accurate real-space representation yielding a crossover with standard techniques for small system sizes. Smaller optimized sets of real-space points may certainly be explored in the future.

Another advantage is that the RI-RS+LT RPA scheme lies in the related  $\mathcal{O}(N^2)$  memory footprint associated with the underlying matrix algebra. In the case of the  $C_{60}$  molecule, the relevant sizes are 1800 spherical AOs, 6060 Cartesian auxiliary orbitals, and 19 140



real-space quadrature points. On a single processor run, the setup of the RI-RS [Eq. (6)] peaks at about 8 Gb of memory while leaving at exit a memory footprint below 1Gb for the  $M_{\beta k}$  coefficients. Furthermore, each  $\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau)$  requires about 3 Gb. In regards to a parallelization scheme oriented towards central processing unit (CPU) efficiency, one can benefit from storing all the  $(n\tau)$   $\chi_0(\mathbf{r}_k, \mathbf{r}_{k'}; i\tau)$  in memory.

As emphasized here above, the present cubic scaling in terms of floating point operations, and quadratic scaling in terms of memory load, was obtained without invoking weak localization nor system dependent sparsity considerations. In particular, our approach does not require the use of the density-fitting (RI-SVS) approach with its sparse 3-center overlap matrix tensor. However, another class of localization properties, based on the exponential decay in real-space of the one-body Green's function in gaped systems,<sup>57</sup> can be easily combined with the present approach. These localization properties, which strongly depend on the electronic properties of the system of interest, are reminiscent of the low-scaling techniques based on local AO formulations in the treatment of MP2<sup>21,22</sup> or RPA<sup>20,23,24</sup> correlation energies.<sup>58</sup> Such additional considerations, together with the stochastic approach by Neuhauser and co-workers,<sup>59</sup> may be combined and explored in the future to seek further reduction of the memory and computing time.

## V. CONCLUSION

We have introduced a separable RI based on a real-space quadrature of codensities. The efficiency of our approach relies on setting up an optimal and compact distribution of real-space points  $\{\mathbf{r}_k\}$  allowing excellent accuracy, as exemplified in the case of the Fock exchange energy and further the MP2 and RPA correlation energies, taking as a test case a large set of molecular systems. Our approach preserves the use of standard Gaussian atomic orbitals and related auxiliary basis sets for all-electron calculations, the real-space set of points being used as an intermediate representation. We demonstrate that such an approach leads to calculating RPA correlation energies with a cubic scaling in terms of operations, quadratic in memory, without invoking any localization nor sparsity considerations that may be combined in the future. The limited number of needed real-space points allows early crossovers with traditional Coulomb-fitting RI-RPA calculations for systems as small as naphthalene or anthracene (Fig. 3). The application of such a real-space separable RI to other explicitly correlated techniques, such as the GW and Bethe-Salpeter equation (BSE) formalisms for calculating charged and neutral electronic excitations in molecular systems,<sup>60</sup> is currently under exploration. The slow convergence of RPA and GW correlation energy with basis size also invites to explore in a more systematic way sets of real-space points adapted to larger auxiliary basis sets.

## SUPPLEMENTARY MATERIAL

See [supplementary material](#) for (I) MP2 and RPA correlation energies, (II) details of the real-space  $\{\mathbf{r}_k\}$  set of points, (III) geometries for the octacene and decacene (B3LYP 6-31Gd) and the  $C_{60}$  fullerene (B3LYP 6-311Gd), (IV) relation with the LS-THC approach estimator, and (V) RPA potential energy curve for the benzene dimer.

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## APPENDIX: IMAGINARY FREQUENCY AND TIME QUADRATURES

We briefly outline the imaginary frequency and time quadrature strategies adopted for calculating the RPA correlation energy and for setting the Laplace transform. Relying on the pole structure of the susceptibility [Eq. (16)], we seek for a quadrature that reproduces the contribution of all possible poles in the imaginary axis integral, using the exact relation

$$\int_0^\infty d\omega \left[ \frac{1}{E - i\omega} + \frac{1}{E + i\omega} \right] = \pi, \quad (\text{A1})$$

where the pole energy  $E$  can vary from  $E_{min}$ , i.e., the electronic energy gap, to  $E_{max}$  determined by the maximum transition energy between the occupied and virtual energy levels. Even if an optimal solution of this problem could be, in principle, determined through the minimax fitting approach usually applied to Laplace transforms,<sup>26</sup> we prefer here a somewhat numerically simpler least square formulation, namely, we seek for  $n\omega$  frequencies  $z_k$  and weights  $w_k$  corresponding to

$$\underset{w_k, z_k}{\operatorname{argmin}} \left[ \int_{\ln(E_{min})}^{\ln(E_{max})} du \left\| \sum_k w_k \left[ \frac{1}{e^u - iz_k} + \frac{1}{e^u + iz_k} \right] - \pi \right\|^2 \right].$$

Here, the log scale has been preferred so as to obtain uniform oscillations of the error as a function of the integration variable (see Fig. 5), since integrating directly over the energy would favor larger

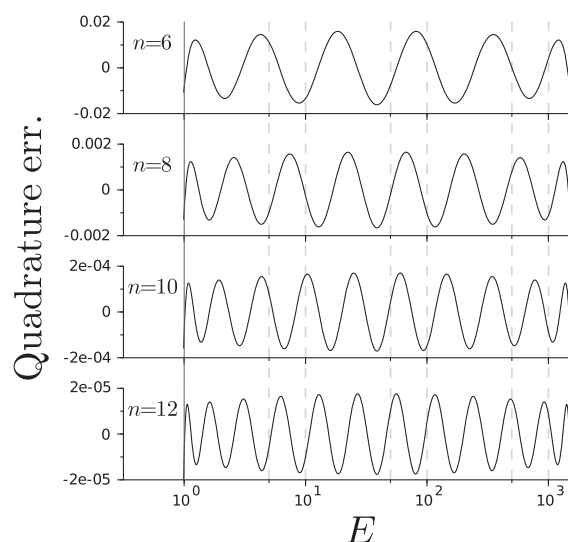


FIG. 5. Imaginary axis quadrature error [see Eq. (A1)] for  $n\omega \in \{6, 8, 10, 12\}$ ,  $E_{min} = 1$ , and  $E_{max} = 1500$ .

**TABLE II.** RPA correlation energy (hartree) for the benzene molecule. The various RI schemes are compared. Calculations are performed at the (spherical) cc-pVTZ level with the corresponding (Cartesian) cc-pVTZ-RI auxiliary basis. The number of imaginary frequencies in the RPA integration is given by  $n\omega$ . The number of times in the Laplace transform (LT) approach is given by  $n\tau$ .

$n\omega$	RI-SVS	RI-V	RI-RS	RI-RS+LT	RI-RS+LT
				( $n\tau = 1.5 \times n\omega$ )	( $n\tau = 2 \times n\omega$ )
6	-1.25169968	-1.25148948	-1.25143651	-1.25148267	-1.25143097
8	-1.25171657	-1.25150621	-1.25145324	-1.25145455	-1.25145327
10	-1.25171791	-1.25150754	-1.25145456	-1.25145456	-1.25145456
12	-1.25171791	-1.25150754	-1.25145456	-1.25145456	-1.25145456

errors in the small energy range.<sup>26</sup> Under such conditions, minimax and least square approaches should provide consistent results. The  $z_k$  frequencies and their corresponding  $w_k$  weights are used for the integration reported in Eq. (15). Convergence tests are provided in Table II for benzene, other acenes being dealt with in the [supplementary material](#). Concerning benzene, our reference RI-V RPA correlation energy falls within less than 8 meV/molecule (0.19 kcal/mol) as compared to that obtained with other codes<sup>61,62</sup> with identified differences in the treatment of preceding Hartree-Fock calculations (RI vs no-RI) and use of the auxiliary basis (Cartesian vs spherical). Together with potential differences in the treatment of the energy integration, these variations result in very small energy differences (see Table S2 of the [supplementary material](#) for details).

For the Laplace transform, the time grid  $\{\tau_p, p \in [1, n\tau]\}$  is set accordingly so as to minimize the Laplace transform errors for the specific quadrature ( $z_k, w_k$ ) points obtained above, namely,

$$\operatorname{argmin}_{w_k^p, \tau_p} \left[ \sum_k \int_{\ln(E_{\min})}^{\ln(E_{\max})} du \left\| \sum_p w_k^p e^{-\tau_p} e^{u} - \left[ \frac{1}{e^u - iz_k} + \frac{1}{e^u + iz_k} \right] \right\|^2 \right],$$

where the weights  $\{w_k^p\}$  depend on the targeted  $z_k$  frequency. Again, the log scale is used so as to allow a regular sampling of the error oscillations. We draw the readers' attention on the fact that the least square approach is clearly at advantage over the minimax fitting approach here since defining an alternant is not possible in the case of simultaneous fits (one fit for each  $z_k$  frequency). The results provided here below for the benzene RPA correlation energy demonstrates already  $<10^{-4}$  Ha convergence for grids containing as few as  $n\omega = 6$  imaginary frequencies and  $n\tau = 9$  times. For  $n\omega = 12$  and  $n\tau = 18$ , the parameters used in this manuscript for actual calculations, the correlation energy is clearly converged.

## REFERENCES

- J. L. Whitten, *J. Chem. Phys.* **58**, 4496 (1973).
- B. I. Dunlap, J. W. D. Connolly, and J. R. Sabin, *J. Chem. Phys.* **71**, 3396 (1979).
- J. W. Mintmire and B. I. Dunlap, *Phys. Rev. A* **25**, 88 (1982).
- C. Van Alsenoy, *J. Comput. Chem.* **9**, 620 (1988).
- M. Feyereisen, G. Fitzgerald, and A. Komornicki, *Chem. Phys. Lett.* **208**, 359 (1993).
- O. Vahtras, J. Almlöf, and M. Feyereisen, *Chem. Phys. Lett.* **213**, 514 (1993).
- W. Klopper and C. C. M. Samson, *J. Chem. Phys.* **116**, 6397 (2002).
- D. Bohm and D. Pines, *Phys. Rev.* **92**, 609 (1953).
- O. Gunnarsson and B. I. Lundqvist, *Phys. Rev. B* **13**, 4274 (1976).
- D. C. Langreth and J. P. Perdew, *Phys. Rev. B* **15**, 2884 (1977).
- H. Eshuis, J. Yarkony, and F. Furche, *J. Chem. Phys.* **132**, 234114 (2010).
- X. Ren, P. Rinke, V. Blum, J. Wierfink, A. Tkatchenko, A. Sanfilippo, K. Reuter, and M. Scheffler, *New J. Phys.* **14**, 053020 (2012).
- I. Duchemin, J. Li, and X. Blase, *J. Chem. Theory Comput.* **13**, 1199 (2017).
- Y. Jung, A. Sodt, P. M. W. Gill, and M. Head-Gordon, *Proc. Natl. Acad. Sci. U. S. A.* **102**, 6692 (2005).
- S. Reine, E. Tellgren, A. Krapp, T. Kjrgaard, T. Helgaker, B. Jansik, S. Hst, and P. Salek, *J. Chem. Phys.* **129**, 104101 (2008).
- A. Sodt and M. Head-Gordon, *J. Chem. Phys.* **128**, 104106 (2008).
- A. C. Ihrig, J. Wierfink, I. Y. Zhang, M. Ropo, X. Ren, P. Rinke, M. Scheffler, and V. Blum, *New J. Phys.* **17**, 093020 (2015).
- J. Almlöf, *Chem. Phys. Lett.* **181**, 319 (1991).
- J. Wilhelm, P. Seewald, M. Del Ben, and J. Hutter, *J. Chem. Theory Comput.* **12**, 5851 (2016).
- A. Luenser, H. F. Schurkus, and C. Ochsenfeld, *J. Chem. Theory Comput.* **13**, 1647 (2017).
- M. Häser, *Theor. Chim. Acta* **87**, 147 (1993).
- P. Y. Ayala and G. E. Scuseria, *J. Chem. Phys.* **110**, 3660 (1999).
- M. Kállay, *J. Chem. Phys.* **142**, 204105 (2015).
- H. F. Schurkus and C. Ochsenfeld, *J. Chem. Phys.* **144**, 031101 (2016).
- H. N. Rojas, R. W. Godby, and R. J. Needs, *Phys. Rev. Lett.* **74**, 1827 (1995).
- M. Kaltak, J. Klimeš, and G. Kresse, *J. Chem. Theory Comput.* **10**, 2498 (2014).
- We will distinguish strong localization properties that are independent of the system and weak localization that depends on the system characteristics (e.g., its dimensionality, size, electronic HOMO-LUMO gap, etc.).
- F. Neese, F. Wennmohs, A. Hansen, and U. Becker, *Chem. Phys.* **356**, 98 (2009), moving Frontiers in Quantum Chemistry.
- R. A. Friesner, *Annu. Rev. Phys. Chem.* **42**, 341 (1991).
- R. M. Parrish, E. G. Hohenstein, T. J. Martinez, and C. D. Sherrill, *J. Chem. Phys.* **137**, 224106 (2012).
- J. Lu and L. Ying, *J. Comput. Phys.* **302**, 329 (2015).
- J. Lu and L. Ying, *Ann. Math. Sci. Appl.* **1**, 321 (2016).
- J. Lu and K. Thicke, *J. Comput. Phys.* **351**, 187 (2017).
- A. Tikhonov, V. Goncharov, A. Stepanov, and A. Yagola, *Numerical Methods for the Solution of Ill-Posed Problems* (Springer Netherlands, 1995).
- V. Lebedev, *USSR Comput. Math. Math. Phys.* **15**, 44 (1975).
- M. Schreiber, M. R. Silva-Junior, S. P. A. Sauer, and W. Thiel, *J. Chem. Phys.* **128**, 134110 (2008).
- M. R. Silva-Junior, S. P. A. Sauer, M. Schreiber, and W. Thiel, *Mol. Phys.* **108**, 453 (2010).
- D. Jacquemin, V. Wathelet, E. A. Perpète, and C. Adamo, *J. Chem. Theory Comput.* **5**, 2420 (2009).

- <sup>39</sup>D. Jacquemin, I. Duchemin, and X. Blase, *J. Chem. Theory Comput.* **11**, 3290 (2015).
- <sup>40</sup>F. Bruneval, S. M. Hamed, and J. B. Neaton, *J. Chem. Phys.* **142**, 244101 (2015).
- <sup>41</sup>K. Krause and W. Klopper, *J. Comput. Chem.* **38**, 383 (2017).
- <sup>42</sup>T. Rangel, S. M. Hamed, F. Bruneval, and J. B. Neaton, *J. Chem. Theory Comput.* **12**, 2834 (2016).
- <sup>43</sup>J. Krüger, F. García, F. Eisenhut, D. Skidin, J. Alonso, E. Guitián, D. Pérez, G. Cuniberti, F. Moresco, and D. Peña, *Angew. Chem., Int. Ed.* **56**, 11945 (2017).
- <sup>44</sup>T. H. Dunning, *J. Chem. Phys.* **90**, 1007 (1989).
- <sup>45</sup>M. Valiev, E. Bylaska, N. Govind, K. Kowalski, T. Straatsma, H. V. Dam, D. Wang, J. Nieplocha, E. Apra, T. Windus, and W. de Jong, *Comput. Phys. Commun.* **181**, 1477 (2010).
- <sup>46</sup>F. Weigend, A. Köhn, and C. Hättig, *J. Chem. Phys.* **116**, 3175 (2002).
- <sup>47</sup>J. Li, G. Davino, I. Duchemin, D. Beljonne, and X. Blase, *J. Phys. Chem. Lett.* **7**, 2814 (2016).
- <sup>48</sup>D. Pines and D. Bohm, *Phys. Rev.* **85**, 338 (1952).
- <sup>49</sup>P. Nozières and D. Pines, *Phys. Rev.* **111**, 442 (1958).
- <sup>50</sup>F. Furche and T. V. Voorhis, *J. Chem. Phys.* **122**, 164106 (2005).
- <sup>51</sup>M. Fuchs, Y.-M. Niquet, X. Gonze, and K. Burke, *J. Chem. Phys.* **122**, 094116 (2005).
- <sup>52</sup>F. Furche, *J. Chem. Phys.* **129**, 114105 (2008).
- <sup>53</sup>All calculations were performed on 2 nodes of  $2 \times 16$  cores of Intel Xeon Gold 6130 CPU @ 2.10 GHz, with Intel Omni-Path Interconnect, except for the RI-V RPA correlation energy of the octapeptide angiotensin II molecule where 4 half filled nodes were used in order to circumvent memory issues.
- <sup>54</sup>J. E. Moussa, *J. Chem. Phys.* **140**, 014107 (2014).
- <sup>55</sup>P. M. Gill, B. G. Johnson, and J. A. Pople, *Chem. Phys. Lett.* **209**, 506 (1993).
- <sup>56</sup>O. Treutler and R. Ahlrichs, *J. Chem. Phys.* **102**, 346 (1995).
- <sup>57</sup>A. Schindlmayr, *Phys. Rev. B* **62**, 12573 (2000).
- <sup>58</sup>Green's functions in real-space and imaginary-time are reminiscent of the so-called auxiliary (pseudo)density matrices ( $\bar{P}, \bar{D}$ ) of Eqs. (15) and (16) in Ref. 24, labeled as well ( $\mathbf{X}, \mathbf{Y}$ ) in Eq. (3.1) of Ref. 22, expressed in the AO basis set.
- <sup>59</sup>D. Neuhauser, E. Rabani, and R. Baer, *J. Phys. Chem. Lett.* **4**, 1172 (2013).
- <sup>60</sup>X. Blase, I. Duchemin, and D. Jacquemin, *Chem. Soc. Rev.* **47**, 1022 (2018).
- <sup>61</sup>F. Bruneval, T. Rangel, S. M. Hamed, M. Shao, C. Yang, and J. B. Neaton, *Comput. Phys. Commun.* **208**, 149 (2016).
- <sup>62</sup>TURBOMOLE V6.2 2010, a development of University of Karlsruhe and Forschungszentrum Karlsruhe GmbH, 1989-2007, TURBOMOLE GmbH, since 2007; available from <http://www.turbomole.com>.