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An exactly solvable continuous-time Derrida–Retaux model

Yueyun Hu* Bastien Mallein† Michel Pain‡

March 27, 2019

Abstract

To study the depinning transition in the limit of strong disorder, Derrida and Retaux [9] introduced a discrete-time max-type recursive model. It is believed that for a large class of recursive models, including Derrida and Retaux’ model, there is a highly non-trivial phase transition. In this article, we present a continuous-time version of Derrida and Retaux model, built on a Yule tree, which yields an exactly solvable model belonging to this universality class. The integrability of this model allows us to study in details the phase transition near criticality and can be used to confirm the infinite order phase transition predicted by physicists. We also study the scaling limit of this model at criticality, which we believe to be universal.

1 Introduction

The problem of the depinning transition has attracted much attention among mathematicians and physicists over the last thirty years, see Giacomin [11, 12] and the references therein. To study the depinning transition in the limit of strong disorder, Derrida and Retaux introduced in [9] a max-type recursive model, that can be defined, up to a simple change of variables, as follows:

$$X_{n+1} \stackrel{(d)}{=} \left(X_n + \widetilde{X}_n - 1 \right)_+, \quad \forall n \geq 0, \quad (1.1)$$

where, for any $x \in \mathbb{R}$, $x_+ := \max(x, 0)$, \widetilde{X}_n denotes an independent copy of X_n , and $\stackrel{(d)}{=}$ stands for the identity in distribution. Note that this model was previously studied by Collet, Eckmann, Glaser and Martin [6] for random variables taking integer values. Through studying the sequence of probability-generating functions, they identified the

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critical manifold of (X_n) , provided that $X_0 \in \mathbb{N}$ a.s. With renormalization group arguments, Derrida and Retaux studied this max-type recursive equation for real-valued random variables.

This recursive model can be viewed as a simplified version of the recursion equation obtained by taking the logarithm of the partition function Z_n of the so-called hierarchical model, which satisfies

$$Z_{n+1} \stackrel{(d)}{=} \frac{Z_n \tilde{Z}_n + b - 1}{b}, \quad \forall n \geq 0, \quad (1.2)$$

where $b > 1$ is a fixed constant and \tilde{Z}_n is an independent copy of Z_n . In other words, Equation (1.1) is a tropicalization of Equation (1.2). The hierarchical model is a pinning model on the diamond lattices, introduced by Derrida, Hakim and Vannimenus [8] and studied in depth by a series of recent works [21, 10, 13, 18, 3].

Despite its simplicity in the definition, the Derrida–Retaux model (which we often abbreviate in the rest of the article as DR model) turns out to exhibit complex behaviours at criticality, which are difficult to prove rigorously, and many fundamental questions remain open. It is believed that for a large class of recursive models, including (1.1) and (1.2), there is an infinite order phase transition from the pinned to the unpinned regime. In this article, we present an exactly solvable version of a continuous-time generalization of the DR model. The integrability of the model allows us to study in details the phase transition of our model near criticality and can be used to clarify some of the conjectures made for the discrete DR models.

Before the introduction of the continuous-time version model, we (re)-present the discrete-time model (1.1) from three different though equivalent viewpoints. In particular, we give a process description of this model, which is best suited for comparison with the continuous-time version model that we are going to introduce.

Notation By $f(t) \sim g(t)$ as $t \rightarrow a$, we mean that $\lim_{t \rightarrow a} f(t)/g(t) = 1$. The constant C is a generic positive constant, which might change from line to line. Moreover, we set $\mathbb{N} := \{0, 1, 2, \dots\}$ and $\mathbb{R}_+ := [0, \infty)$.

1.1 The discrete-time Derrida–Retaux model

At first, one might observe that (1.1) does not describe a stochastic process $(X_n)_{n \geq 0}$, but rather a measure-valued sequence $(\mu_n)_{n \geq 0}$ where μ_n is the law of X_n . This leads us to the following definition.

Definition 1.1 ((Discrete-time) Derrida–Retaux model). A sequence $(\mu_n)_{n \geq 0}$ of probability distributions on \mathbb{R}_+ is called a *DR model* if for all continuous bounded function f , we have

$$\forall n \geq 0, \quad \int_{\mathbb{R}_+} f(x) \mu_{n+1}(dx) = \int_{\mathbb{R}_+} f((x+y-1)_+) \mu_n(dx) \mu_n(dy).$$

The fundamental quantity in a DR model is its associated *free energy*, defined as follows:

$$F_\infty := \lim_{n \rightarrow \infty} 2^{-n} \int_0^\infty x \mu_n(dx). \quad (1.3)$$

It allows the distinction between the pinned and the unpinned regimes in this model. We refer to the model as *unpinned* if $F_\infty = 0$, and as *pinned* if $F_\infty > 0$. Note that as for

$x, y \in \mathbb{R}_+$, $(x + y - 1)_+ \leq x + y$, the sequence $(2^{-n} \int_0^\infty x \mu_n(dx))_{n \geq 0}$ is non-increasing, therefore the limit in (1.3) is always well-defined. One of the main questions for this model is to determine for which starting measure μ_0 the model ends up being pinned or unpinned.

When the support of μ_0 is included in \mathbb{N} , this question was previously answered by Collet et al. [6]. By a careful study of the recursion equation for the generating functions of μ_n , they proved that

$$F_\infty = 0 \iff \int_0^\infty x 2^x \mu_0(dx) \leq \int_0^\infty 2^x \mu_0(dx) < \infty. \quad (1.4)$$

However, this characterization of the unpinned phase, which is valid for integer-valued random variables does not extend to more general initial conditions, even to the half-integers valued random variables. We refer to Chen et al. [5] for further discussions and a list of open questions on this topic.

To discuss the phase transition from the pinned to the unpinned regime, it will be more convenient to specify the mass at 0 of μ_0 . Consider an initial distribution μ_0 of form:

$$\mu_0 = p\delta_0 + (1 - p)\vartheta,$$

where $p \in [0, 1]$ and ϑ denotes a probability measure on $\mathbb{R}_+ \setminus \{0\}$. Denote by $F_\infty(p)$ the associated free energy. Clearly, $p \rightarrow F_\infty(p)$ is non-increasing. Write p_c for the critical parameter distinguishing between the pinned and the unpinned regimes

$$p_c := \sup\{p \in [0, 1] : F_\infty(p) > 0\},$$

with the convention that $\sup \emptyset = 0$.

Assuming that $p_c \in (0, 1)$, are of particular interest to characterize the rate of convergence of $F_\infty(p) \rightarrow 0$ as $p \uparrow p_c$, and to study the critical regime when $p = p_c$. For the rate of convergence, assuming $p_c > 0$ and some integrability conditions on ϑ , Derrida and Retaux [9] conjectured that there exists some constant $C > 0$ such that

$$F_\infty(p) = \exp\left(-\frac{C + o(1)}{\sqrt{p_c - p}}\right), \quad p \uparrow p_c. \quad (1.5)$$

When the support of ϑ is included in $\mathbb{N} \setminus \{0\}$, a weaker form of the above conjecture has been obtained in a recent work by Chen et al. [4].

It is necessary to assume some integrability assumptions on ϑ in Derrida and Retaux' conjecture. Indeed, Collet et al. [6]'s result yields the following characterization:

$$p_c < 1 \iff \int_0^\infty x 2^x \vartheta(dx) < \infty,$$

see Hu and Shi [17]. Furthermore, the latter paper shows some non-universal behaviors of $F_\infty(p)$ as $p \uparrow 1$ for a class of distributions ϑ satisfying $p_c = 1$, in contrast with (1.5).

Now we turn our attention to a second viewpoint of the Derrida and Retaux model. If the support of μ_0 is included in \mathbb{N} , the law μ_n can be constructed as a so-called *parking scheme* on the binary tree, which can be described as follows. For $n \in \mathbb{N}$, we denote by \mathbb{T}_n the binary tree of height n , rooted at \emptyset . To each leaf of the tree at depth n is attached an i.i.d. random variable with law μ_0 , representing the number of cars starting from that position. Each internal node has a single parking spot, initially empty. All cars drive

from the leaves toward the root of the tree, parking as soon as possible. For each $u \in \mathbb{T}_n$, we write $X^n(u)$ the number of cars passing by the parking spot at position u without being allowed to park (because another car is already parked). We call the random map $X^n: \mathbb{T}_n \rightarrow \mathbb{R}_+$ the Derrida–Retaux tree (or DR tree for short). See Figure 1.

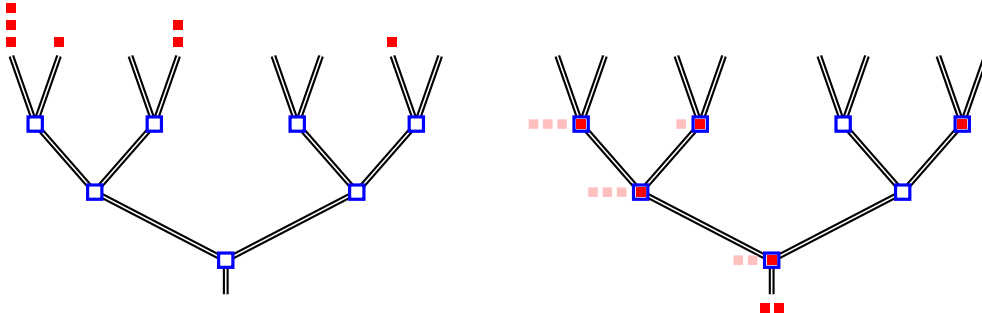


Figure 1: Illustration of the parking scheme on a binary tree. On the left, cars (red squares) are starting from the leaves and each internal node has an empty parking spot. On the right, cars drove down to the root, parking as soon as they could. At each node u , the number $X^n(u)$ of cars that passed by u without being able to park is indicated by pink squares.

Observe that the following equation holds for each internal node $u \in \mathbb{T}_n$:

$$X^n(u) = (X^n(u1) + X^n(u2) - 1)_+, \quad (1.6)$$

where $u1$ and $u2$ are the left and right children of u respectively. As a result, one can observe that for any $0 \leq k \leq n$, the family $X^n(v)$ for all nodes v at the k th generation of \mathbb{T}_n forms a collection of i.i.d. random variables with common law μ_{n-k} . In particular $X^n(\emptyset)$ has law μ_n .

We generalize this definition to any initial law μ_0 on \mathbb{R}_+ . For any $u \in \mathbb{T}_n$, denote by $|u|$ its generation. As such, $\{u : |u| = n\}$ is exactly the set of leaves of \mathbb{T}_n .

Definition 1.2 ((Discrete time) Derrida–Retaux tree). Let μ_0 be a probability distribution on \mathbb{R}_+ . A *DR tree* of height n and initial law μ_0 is a random map $X^n: \mathbb{T}_n \rightarrow \mathbb{R}_+$, with \mathbb{T}_n being the binary tree of height n , which can be constructed as follows:

- The values of the leaves ($X^n(u), |u| = n$) are i.i.d. random variables with law μ_0 .
- Each internal node value is constructed as $X^n(u) = (X^n(u1) + X^n(u2) - 1)_+$, for any $u \in \mathbb{T}_n$ with $|u| < n$.

Equation (1.6) is a max-type recursive equation on trees, that belongs to the class of recursive models analysed in the survey of Aldous and Bandyopadhyay [1]. The parking scheme we describe here has been studied by Goldschmidt and Przykucki [14] on critical Poisson Galton–Watson trees conditioned to be large. Curien and Hénard [7] are also studying parking scheme on uniform trees with n nodes. In these different cases, the key observation is that the models can be analysed through the study of some generating function. However, parking on critical and supercritical Galton–Watson trees appear to belong to different *universality class*, as their behavior, in particular near criticality is very different. Namely, the phase transition of the parking scheme on a critical Galton–Watson tree is of the second order in all known examples ([14, 7]), while it is of infinite order on supercritical Galton–Watson trees (see (1.5)).

The DR tree gives a natural method to generate a random variable X_n with law μ_n . However, this construction does not make easy the consideration of the family $(X_n, n \geq 0)$ as a process. Indeed, making the cars start from generation n and flow back to the root only allows to construct the first n generations of that process. We introduce here a self-contained construction of $(X_n, n \geq 0)$ as a time-inhomogeneous discrete-time Markov chain.

Definition 1.3 (Derrida–Retaux process). Let μ_0 be a probability distribution on \mathbb{R}_+ . A *DR process* with initial law μ_0 is the process constructed from the following recursion equation

$$\forall n \geq 0, \quad \begin{cases} X_{n+1} = (X_n + F_n^{-1}(U_{n+1}) - 1)_+, \\ F_n(x) = \mathbb{P}(X_n \leq x) \text{ for all } x \geq 0. \end{cases} \quad (1.7)$$

where $(U_n, n \geq 1)$ is an i.i.d. family of uniform random variables, and X_0 is an independent random variable with law μ_0 .

Note that the DR process can be seen as a discrete-time version of a solution of a McKean–Vlasov type SDE, see McKean [20], i.e. a Markov process interacting with its distribution. It follows from an easy recursion that for all $n \in \mathbb{N}$, X_n has law μ_n .

The aim of this article is to define continuous-time versions of the DR model, tree and process. In particular, we are looking for a process with a density $r_t(x)$ with respect to the Lebesgue measure which could solve [9, Equation (33)], that we recall here

$$\partial_t r = \partial_x r + r * r - r. \quad (1.8)$$

This differential equation (1.8) plays a key role in the prediction (1.5), and was obtained by Derrida and Retaux as an informal scaling limit of the sequence of measures $(\mu_{[nt]}(ndx))$.

1.2 A continuous-time version of the model, tree and process

By analogy with the DR process, we introduce a continuous-time version as the following Markov process. Starting from an initial distribution X_0 , the process drifts downward at speed 1 until reaching 0, where it stays put. Additionally, at each atom t of a Poisson point process with intensity 1, an independent copy of X_t is added to X_t . More formally, the process is defined in the following way.

Definition 1.4 ((Continuous-time) Derrida–Retaux process). A *DR process* $(X_t)_{t \geq 0}$ is a solution of the McKean–Vlasov type differential equation

$$\begin{cases} X_t = X_0 - \int_0^t \mathbb{1}_{\{X_s > 0\}} ds + \int_0^t \int_{[0,1]} F_s^{-1}(u) N(ds, du) \\ F_t(x) = \mathbb{P}(X_t \leq x), \text{ for all } x \geq 0, t \geq 0, \end{cases} \quad (1.9)$$

where N is a Poisson point process on $\mathbb{R}_+ \times [0, 1]$ with intensity $dsdu$ and F_s^{-1} is the right-continuous inverse of F_s defined by $F_s^{-1}(u) := \inf\{x \geq 0 : F_s(x) > u\}$.

Observe that this continuous-time process has the same features as the discrete-time DR process. It exhibits the similar pattern of losing mass up to hitting zero and being added independent copies of itself. Note that, in order to study a continuous-time DR process Y losing mass at speed a and being added independent copies of x at rate b , one only needs to make the space-time change $Y_t = \frac{a}{b} X_{bt}$.

We prove that for any initial random variable X_0 on \mathbb{R}_+ , there is a unique strong solution of (1.9). Moreover, setting μ_t for the law of X_t , the family $(\mu_t)_{t \geq 0}$ is continuous with respect to the usual topology of the weak convergence. It turns out that the family $(\mu_t)_{t \geq 0}$ will be the solution of a partial differential equation similar to (1.8). In that sense, it will be called a DR model, enticing the following definition.

Definition 1.5. [(Continuous-time) Derrida–Retaux model] A *DR model* is a family $(\mu_t)_{t \geq 0}$ of weakly continuous probability distributions that are weak solution to the partial differential equation

$$\partial_t \mu_t = \partial_x (\mathbb{1}_{\{x > 0\}} \mu_t) + \mu_t * \mu_t - \mu_t, \quad (1.10)$$

where by “weak solution” we mean that for any \mathcal{C}^1 bounded functions f with bounded derivative and any $t > 0$, we have

$$\begin{aligned} & \int_{\mathbb{R}} f(x) \mu_t(dx) - \int_{\mathbb{R}} f(x) \mu_0(dx) \\ &= \int_0^t \left(- \int_{\mathbb{R}} (f'(x) \mathbb{1}_{\{x > 0\}} + f(x)) \mu_s(dx) + \int_{\mathbb{R}^2} f(x+y) \mu_s(dx) \mu_s(dy) \right) ds. \end{aligned} \quad (1.11)$$

Remark 1.6. Note that (1.10) is very similar to Equation (1.8) predicted by physicists. In particular, the behaviors of the PDEs are the same for all $x > 0$. The effect of adding the $\mathbb{1}_{\{x > 0\}}$ term in (1.10) is to specify the evolution of the Dirac mass at 0 separately. Contrarily to what happens in (1.8), this ensures that the total mass of the measures is preserved.

The partial differential equation (1.10) might be related to a Smoluchowski-type coagulation equation [22] with constant kernel. However, we were not able to find references in the PDE literature to partial differential equations exhibiting a behavior similar to (1.10). Nevertheless, it has recently gained focus in the probability literature, for example in the recent article of Lambert and Schertzer [19], where a PDE of the form

$$\partial_t \varrho = \partial_x (\psi \varrho) + a(t) (\varrho * \varrho - \varrho),$$

with ψ the branching mechanism of a continuous-state branching process, was introduced and studied in detail.

The family of law $(\mu_t)_{t \geq 0}$ associated to the DR process can also be constructed on a Yule tree with a method similar to the parking scheme. We recall that a Yule tree is a continuous tree in which each branch has i.i.d. exponential length of parameter 1, and the branching is always binary. For $s \geq 0$, we denote by \mathcal{N}_s the set of individuals alive in the tree at time s .

For some fixed $t > 0$, we define a process $X^t = (X_s^t(u), s \in [0, t], u \in \mathcal{N}_s)$ on a Yule tree truncated at height t by the following *painting scheme*. On each leaf $u \in \mathcal{N}_t$ of the truncated tree, there is a painter with an i.i.d. random amount of paint $X_t^t(u)$, distributed according to the law μ_0 . Then, the procedure is completely deterministic. Each painter climbs down the tree, painting the branches of the tree, using a quantity 1 of paint per unit of branch length. When two painters meet, they put their remaining paint in common. For each $s \in [0, t]$ and $u \in \mathcal{N}_s$, the quantity $X_s^t(u)$ is the amount of paint left when the painters reach level s in the branch u . See Figure 2 for an illustration. This leads to the following definition.

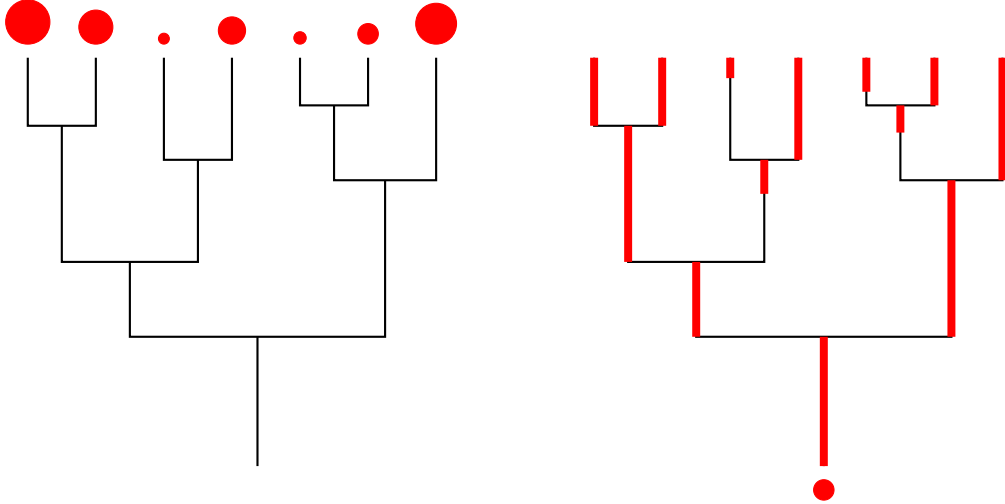


Figure 2: Illustration of the painting scheme on a Yule tree. On the left, painters are starting on the leaves with a random amount of red paint. On the right, painters climbed down to the root, painting the tree as long as they could. The remaining paint at the root is $X_0^t(\emptyset)$.

Definition 1.7 ((Continuous-time) Derrida–Retaux tree). Let μ_0 be a probability distribution on \mathbb{R}_+ . For $t > 0$, a *DR tree* of height t and initial law μ_0 is a random process $X^t = (X_s^t(u), s \in [0, t], u \in \mathcal{N}_s)$, where $X_s^t(v)$ represents the amount of paint remaining at level s on the branch v in the painting scheme starting with $(X_t^t(u), u \in \mathcal{N}_t)$ that are i.i.d. with law μ_0 conditionally on the Yule tree.

Denoting by $\bar{\mu}_t$ the law of $X_0^t(\emptyset)$, we call $(\bar{\mu}_t)_{t \geq 0}$ the *family of laws obtained from the tree-painting scheme* (in particular, $\bar{\mu}_0 = \mu_0$). We observe that by the branching property of the Yule tree, the subtree starting from individual u at time s is a Yule tree of height $t - s$. Therefore for all $s \leq t$, the family $(X_s^t(u), u \in \mathcal{N}_s)$ are i.i.d. random variables with law $\bar{\mu}_{t-s}$. Contrarily to the two previous definitions, the question of existence and uniqueness of $(\bar{\mu}_t)_{t \geq 0}$ is straightforward in Definition 1.7. We prove later that $(\bar{\mu}_t)_{t \geq 0}$ is a DR model, see Theorem 1.8 below.

We refer to Section 2.1 for a detailed definition of the continuous-time DR tree. However, it is worth noticing that the painting scheme can be obtained straightforwardly as a scaling limit, when $K \rightarrow \infty$, of the parking scheme on a Galton–Watson tree of height Kt , in which particles give birth to two children with probability $1/K$ and to one child with probability $1 - 1/K$, and the starting number of cars in a given leaf has law $\lfloor KX_0 \rfloor$, with X_0 of law μ_0 .

The first main result of the article justifies and unifies the three definitions of the continuous-time model given above.

Theorem 1.8. *Let μ_0 be a probability distribution on \mathbb{R}_+ , and $(\bar{\mu}_t)_{t \geq 0}$ the family of laws obtained from the tree-painting scheme.*

- (i) *There exists a unique strong solution $(X_t)_{t \geq 0}$ to Equation (1.9) with X_0 of law μ_0 . Moreover, for all $t \geq 0$, X_t has law $\bar{\mu}_t$.*
- (ii) *The family $(\bar{\mu}_t)_{t \geq 0}$ is the unique weak solution of the PDE (1.10) with initial condition μ_0 , therefore a DR model with initial law μ_0 .*

Additionally, we obtain preliminary results on the asymptotic behavior of general continuous-time DR models, see Section 6. However, many fundamental questions remain open for this complex dynamics in the general case. In the next subsection, we shall present a simple two-parameter family of measures which turns out to be stable by the dynamics (1.10). Therefore, by restricting ourselves to this family, we are able to introduce an exactly solvable model, which exhibits the researched phase transition at criticality.

1.3 An exactly solvable model

Let $(\mu_t)_{t \geq 0}$ be a continuous-time DR model. We are particularly interested in an exactly solvable case where the initial law μ_0 is a mixture of an exponential distribution and a Dirac mass at 0:

$$\mu_0 = p\delta_0 + (1-p)\lambda e^{-\lambda x} \mathbb{1}_{\{x>0\}} dx, \quad p \in [0, 1], \lambda > 0. \quad (1.12)$$

Denote by $F_\infty(p, \lambda)$ the associated *free energy*:

$$F_\infty(p, \lambda) := \lim_{t \rightarrow \infty} e^{-t} \int_0^\infty x \mu_t(dx). \quad (1.13)$$

To describe the asymptotic behavior of μ_t with respect to the parameters (p, λ) of the initial distribution μ_0 , we introduce three sets:

$$\begin{aligned} \mathcal{P} &:= \{(p, \lambda) \in [0, 1) \times [0, \infty) : \lambda < 1 \text{ or } (\lambda > 1 \text{ and } p < \lambda - \lambda \log \lambda)\}, \\ \mathcal{C} &:= \{(p, \lambda) \in [0, 1) \times [1, \infty) : p = \lambda - \lambda \log \lambda\}, \\ \mathcal{U} &:= \{(p, \lambda) \in [0, 1) \times [1, \infty) : p > \lambda - \lambda \log \lambda\}, \end{aligned}$$

which are represented in Figure 3. The case $p = 1$ is degenerate therefore excluded from the set of parameters.

The following result gives a classification of this DR model in pinned, unpinned and critical behaviors for any starting distribution.

Theorem 1.9. *Let $(\mu_t)_{t \geq 0}$ be the continuous-time DR model with initial distribution μ_0 given by (1.12).*

(i) *If $(p, \lambda) \in \mathcal{P}$, then $F_\infty(p, \lambda) > 0$.*

(ii) *If $(p, \lambda) \in \mathcal{U}$, then there exists some positive constant C such that*

$$\int_0^\infty y \mu_t(dy) \sim C e^{-(x-1)t}, \quad \text{as } t \rightarrow \infty,$$

where x is the unique solution larger than 1 of the equation $Hx - x \log x = 1$ and $H := \frac{p}{\lambda} - \log \lambda$.

(iii) *If $(p, \lambda) \in \mathcal{C}$, then*

$$\int_0^\infty y \mu_t(dy) \sim \frac{2}{t^2}, \quad \text{as } t \rightarrow \infty.$$

Furthermore, conditionally on $(0, \infty)$, μ_t converges weakly to a standard exponential distribution.

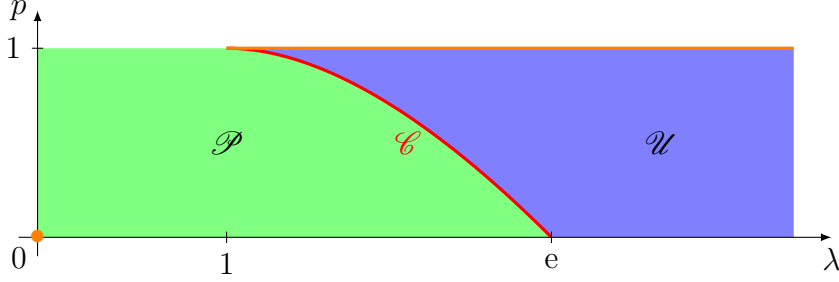


Figure 3: The pinned and unpinned regions of the solvable DR model

For the discrete-time DR model, the counterpart of (iii) was conjectured in Chen et al. [5].

Let us call \mathcal{P} , \mathcal{U} and \mathcal{C} respectively the *supercritical*, *subcritical* and *critical* zones of the dynamics. Note that the critical zone \mathcal{C} is the graph of the function $\lambda \mapsto \lambda - \lambda \log \lambda$ on $(1, \infty)$, which splits the domain $[0, 1) \times \mathbb{R}_+$ into the supercritical and subcritical zones. In the supercritical zone, the dynamics is such that the measure μ_t is attracted by the stable equilibrium $(0, 0)$ (in orange in Figure 3). In the critical zone, μ_t is attracted by the critical point $(1, 1)$ (in red in Figure 3). Finally, in the subcritical zone, μ_t converges toward a point on the orange lane, with $p = 1$ and $\lambda > 1$. In other words, μ_t converges toward δ_0 in probability if $(p, \lambda) \in \mathcal{U} \cup \mathcal{C}$.

It is of interest to give the precise behavior of $p \mapsto F_\infty(p, \lambda)$ as the function hits 0, whenever that point exists. Note that for $\lambda > e$, one has $F_\infty(p, \lambda) = 0$ for all $p \in (0, 1)$, therefore this behavior can be observed only for $\lambda < e$.

Theorem 1.10. *Fix some $\lambda \in (0, e)$ and let $p \in (0, 1)$ vary.*

(i) *If $\lambda \in (1, e)$, setting $p_c = \lambda - \lambda \log \lambda$, there exists $C > 0$ such that*

$$F_\infty(p, \lambda) \sim \mathbb{1}_{\{p < p_c\}} C \exp\left(-\pi\sqrt{2\lambda}(p_c - p)^{-1/2}\right) \quad \text{as } p \uparrow p_c.$$

(ii) *If $\lambda = 1$, we have $F_\infty(1, 1) = 0$ and there exists $C > 0$ such that*

$$F_\infty(p, \lambda) \sim C(1 - p)^{2/3} \exp\left(-\frac{\pi}{\sqrt{2}}(1 - p)^{-1/2}\right) \quad \text{as } p \uparrow 1.$$

(iii) *If $\lambda \in (0, 1)$, we have $F_\infty(1, \lambda) = 0$ and there exists $C > 0$ such that*

$$F_\infty(p, \lambda) \sim C(1 - p)^{1/(1-\lambda)} \quad \text{as } p \uparrow 1.$$

We may observe the difference of a factor 2 in the exponents of the right-hand sides in cases (i) and (ii). Moreover, (i) and (ii) prove the DR conjecture for our model, with enough precision to see the polynomial factor that appears in case (ii). On the other hand, part (iii) gives a continuous version of the polynomial decay obtained in by Hu and Shi [17].

We shall take a particular interest in the DR tree in the critical regime, conditioned on survival. It is believed that the conditioned tree limit is an universal characteristic, and would appear in numerous models related to DR at criticality. For $t > 0$, let X^t denote the continuous-time DR tree with initial distribution μ_0 given by (1.12) (see Definition

1.7). Conditionally on the event $\{X_0^t(\emptyset) > 0\}$, let \mathfrak{X}^t be the restriction of the process X^t to the painted connected component of the root \emptyset . Assuming that the paint is red, the tree \mathfrak{X}^t is called the *red tree*, see Section 4 for a more precise definition and properties.

Theorem 1.11. *Assume $(p, \lambda) \in \mathcal{C}$.*

- (i) *The rescaled process $(\frac{1}{t}\mathfrak{X}_{st}^t)_{0 \leq s < 1}$ converges locally in law as $t \rightarrow \infty$ toward a time-inhomogeneous branching Markov process $(\mathfrak{X}_s)_{0 \leq s < 1}$.*
- (ii) *For each $t \geq 0$, let N_t denote the number of leaves in \mathfrak{X}^t . There exists a positive constant C such that N_t/t^2 converges in law to $C \int_0^1 r^2(s)ds$ as $t \rightarrow \infty$, where $(r(s))_{0 \leq s \leq 1}$ denotes a 4-dimensional Bessel bridge.*

The limiting process \mathfrak{X} is a continuous-time system of particles with masses, that grow continuously and split at random times. As such, it can be called a time-inhomogeneous Markovian growth-fragmentation process. It can be described as satisfying the following properties:

1. It starts at time 0 with a unique particle of mass 0.
2. The mass associated to each particle grows linearly at speed 1.
3. A particle of mass m at time s splits at rate $2m/(1-s)^2$ into two children, the mass m being split uniformly between the two children.
4. Particles behave independently after their splitting time.

The local convergence in the statement (i) can be seen equivalently as the convergence of the particle system on the interval $[0, 1 - \varepsilon]$ for all $\varepsilon > 0$, or the joint convergence in law of the splitting time of the individuals in the n first generations of the process for all $n \geq 1$. Indeed, for fixed $\varepsilon > 0$, the total number of particles in \mathfrak{X}^t at height $(1 - \varepsilon)t$ remains tight as $t \rightarrow \infty$. Conversely, the number of particles at height t explodes as $t \rightarrow \infty$ as shown in statement (ii); see Figure 4. We refer to Theorems 4.5 and 5.1 respectively for a more refined statement of (i) and (ii), respectively.

Remark 1.12. The counterpart of Theorem 1.11(ii) for the discrete DR model is predicted to hold by B. Derrida (personal communication).

1.4 Organization of the paper

The rest of the article is organised as follows. In Section 2, we define properly the three constructions of the continuous version of the DR model. These constructions turn out to be equivalent, as shown by Propositions 2.2 and 2.3. In particular, Theorem 1.8 follows.

In Section 3 we study the exactly solvable model in which the evolution is restricted to the class of mixtures of exponential distributions and Dirac masses at 0. We give at first the phase diagram of the model (Proposition 3.5), then analyze the asymptotic of a system of differential equations (Proposition 3.7), and prove Theorems 1.9 and 1.10 in Sections 3.2 and 3.3 respectively.

Sections 4 and 5 are devoted to the study of the critical regime of our model. Conditioned on survival, the DR tree at any fixed height is a time-inhomogeneous branching Markov process (Proposition 4.1). By letting the height go to infinity, the scaling limit

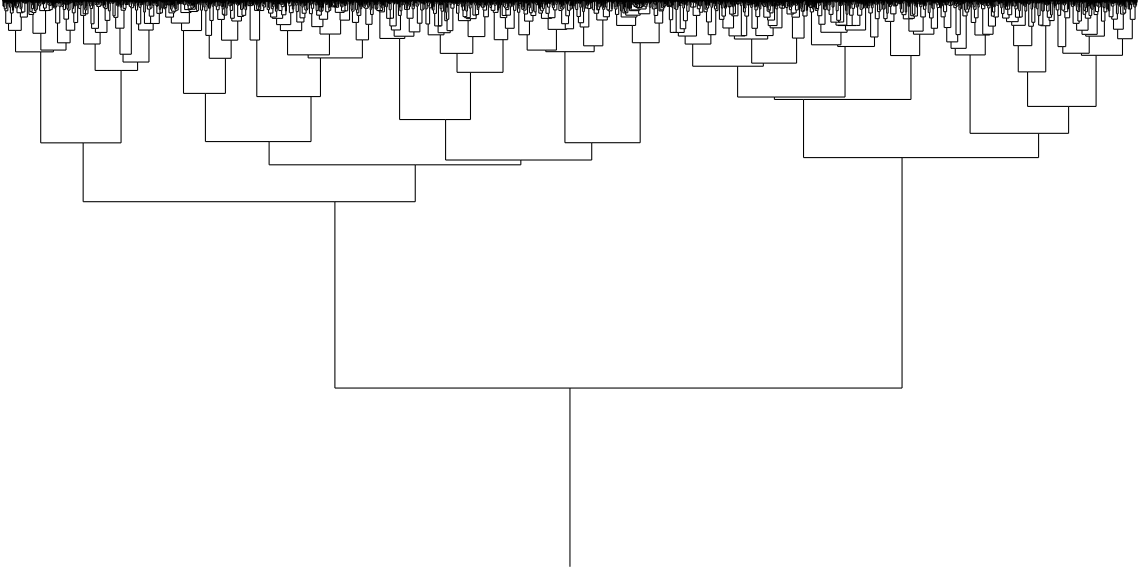


Figure 4: Realization of the red tree \mathfrak{X}^t with $t = 200$ and initial condition $(p, \lambda) = (0, e)$.

of the conditioned tree exists (Theorem 4.5) which gives the part (i) in Theorem 1.11. In Section 5, we are interested in some quantitative characteristics of the conditioned (red) tree, the number N_t and mass M_t of leaves of the red tree of height t . We prove in Lemma 5.3 that the joint Laplace transform of (N_t, M_t) is given by the unique solution of a differential equation whose asymptotics are given in Section 5.3. This allows us to obtain the convergence in law of $(N_t/t^2, M_t/t^2)$ (Theorem 5.1) as $t \rightarrow \infty$. The part (ii) in Theorem 1.11 is a particular case of Theorem 5.1.

In Section 6, we give some results on the continuous-time DR model with a general initial distribution on \mathbb{R}_+ . In particular, we prove that $e^{-t}X_t$ converges in law to an exponential distribution in the pinned case (Proposition 6.3) and give some necessary conditions for the model to being unpinned (Proposition 6.6). We end Section 6 by presenting some open questions.

2 Constructions of the continuous-time Derrida–Re- taux model

In this section, we first give a precise definition of the DR tree and a characterization of the family of laws obtained from the tree-painting scheme. Then, we show that the DR process is well-defined as a solution of the “McKean-Vlasov type” stochastic differential equation (1.9) and that its one-dimensional marginal distributions can be obtained from the tree-painting scheme. Finally, we prove that the DR model is well-defined as the solution of the partial differential equation (1.10): existence follows from Itô formula applied to the DR process and uniqueness from the characterization of the family of laws obtained from the tree-painting scheme. This shows that these constructions are equivalent, as stated in Theorem 1.8.

2.1 Construction from a Yule tree: the Derrida–Retaux tree

We recall that a *Yule tree* is the genealogical tree of a simple population, in which every individual splits after an independent exponential time of parameter one into two children. This tree can be constructed straightforwardly using the standard Ulam–Harris–Neveu notation that we now recall. We denote by $\mathbb{T} = \bigcup_{n \in \mathbb{Z}_+} \{1, 2\}^n$ the space of all finite words on $\{1, 2\}$, where by convention $\{\emptyset\} = \{1, 2\}^0$. The element $u = (u(1), \dots, u(n)) \in \mathbb{T}$ represents the $u(n)$ th child of the $u(n-1)$ th child of the... of the $u(1)$ th child of the root \emptyset . We denote by $|u| = n$ the generation to which u belongs, and for $k \leq n$ by $u_k = (u(1), \dots, u(k))$ its ancestor alive at generation k . We set $u1$ and $u2$ the labels of the two children of u , obtained by adding 1 or 2 to the end of the word u .

The Yule tree is constructed as follows. We denote by $\{e_u, u \in \mathbb{T}\}$ a family of i.i.d. exponential random variables, that represent the time that individual u waits until splitting into its two children $u1$ and $u2$. We then define recursively the birth time b_u of each individual $u \in \mathbb{T}$ by

$$b_u = \sum_{k < |u|} e_{u_k}, \quad \text{with } b_\emptyset = 0.$$

We also set $d_u = b_u + e_u$ the time at which the individual u splits into its children, and $\mathcal{N}_t = \{u \in \mathbb{T} : b_u \leq t < d_u\}$ the set of particles alive at time t .

We now give the detailed definition of the continuous-time DR tree, by analogy with its definition for Galton–Watson trees. Let μ_0 be a probability distribution on \mathbb{R}_+ and $t > 0$. Given the Yule tree, we denote by $(X_t^t(u), u \in \mathcal{N}_t)$ i.i.d. random variables with law μ_0 . For $u \in \mathcal{N}_t$, we define $X_s^t(u) = (X_t^t(u) - (t - s))_+$ for each $s \in [b_u, t]$. Then, we proceed recursively from the leaves to the root by setting, for any $v \in \mathbb{T}$ such that $d_v \leq t$ and any $s \in [b_v, d_v)$,

$$X_s^t(v) = (X_{b_{v1}}^t(v1) + X_{b_{v2}}^t(v2) - (d_v - s))_+, \quad (2.1)$$

noting that $b_{v1} = b_{v2} = d_v$ (see Figure 5). The process $X^t = (X_s^t(u), s \in [0, t], u \in \mathcal{N}_s)$ is a *continuous-time DR tree* of height t and initial law μ_0 , introduced in Definition 1.7. Denoting by μ_t the law of $X_0^t(\emptyset)$, the family $(\mu_t)_{t \geq 0}$ is called *the family of laws obtained from the tree-painting scheme*.

We now give a characterization of family of laws obtained from the tree-painting scheme. For any probability distribution ν on \mathbb{R}_+ and $t > 0$, we introduce the shifted distribution $\tau_t \nu$ satisfying

$$\forall f \in \mathcal{C}_b(\mathbb{R}_+), \quad \int_{\mathbb{R}_+} f(x) (\tau_t \nu)(dx) = f(0) \nu([0, t]) + \int_{\mathbb{R}_+} f(x - t) \nu(dx), \quad (2.2)$$

that represents the law of $(Z - t)_+$ if Z has law ν . We can now describe the law of μ_t by decomposing it with respect to the first branching time in the Yule tree.

Lemma 2.1. *The family $(\mu_t)_{t \geq 0}$ of laws on \mathbb{R}_+ is a family obtained from the tree-painting scheme if and only if*

$$\forall t \geq 0, \quad \mu_t = \int_0^t e^{-s} \tau_s (\mu_{t-s} * \mu_{t-s}) ds + e^{-t} \tau_t (\mu_0). \quad (2.3)$$

Proof. Let \mathbb{T} be a Yule tree, we recall that the first branching time d_\emptyset is distributed according to an exponential random variable with parameter 1. Moreover, the subtrees

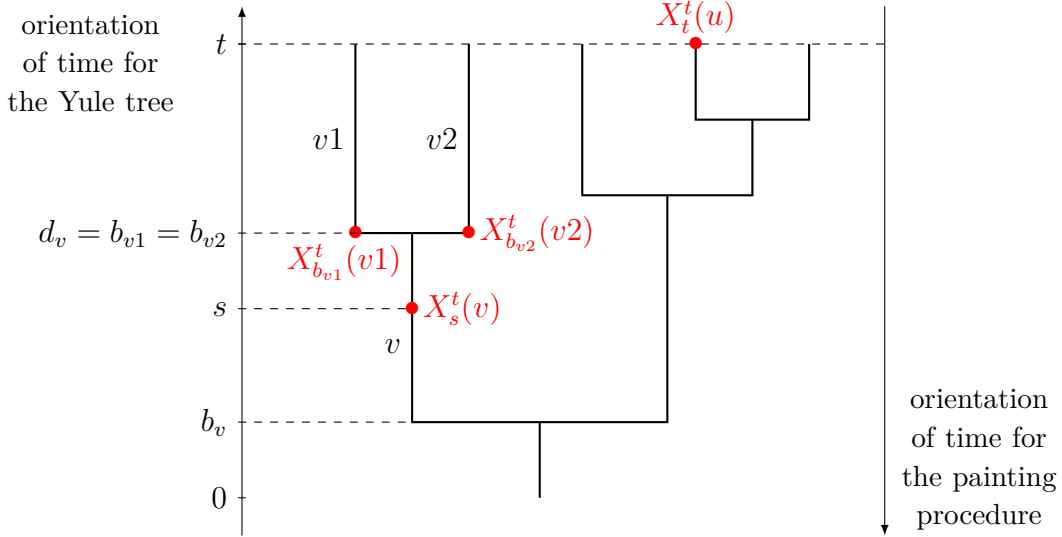


Figure 5: Illustration of the definition of the $X_s^t(u)$'s on the Yule tree.

starting from each child are i.i.d. Yule trees. As a result, conditionally on $\{d_\emptyset = s\}$ with $s \leq t$, $X_s^t(1)$ and $X_s^t(2)$ are two independent random variables with law μ_{t-s} , where as before, 1 and 2 are the two children of the root \emptyset . Note that by construction, for every measurable bounded function f , we have

$$\mathbb{E}\left[f(X_0^t(\emptyset))\right] = \mathbb{E}\left[f((X_{d_\emptyset}^t(1) + X_{d_\emptyset}^t(2) - d_\emptyset)_+) \mathbb{1}_{\{d_\emptyset \leq t\}}\right] + \mathbb{E}\left[f((X_0 - t)_+) \mathbb{1}_{\{d_\emptyset > t\}}\right],$$

where X_0 is a random variable independent of d_\emptyset with law μ_0 . Integrating w.r.t. d_\emptyset , this proves that (μ_t) satisfies (2.3).

Reciprocally, let $(\mu_t)_{t \geq 0}$ be a family of probability measures satisfying (2.3). Fix $t > 0$, we shall construct a continuous-time DR tree X^t of height t and initial law μ_0 such that $X_0^t(\emptyset)$ has law μ_t . To this end, we consider on the same probability space three independent random variables T, U, \tilde{U} such that T has the standard exponential distribution, U and \tilde{U} are uniformly distributed on $[0, 1]$. For any $s \geq 0$, let $F_s(x) := \mu_s([0, x])$, $x \geq 0$ and F_s^{-1} be the right-continuous inverse of F_s . Let \emptyset be the root of the DR tree that we are constructing. Define $d_\emptyset := T$ and denote by 1 and 2 the two children of \emptyset . Let

$$X_0^t(\emptyset) := (X_0 - t)_+ \mathbb{1}_{\{T > t\}} + (X_T^t(1) + X_T^t(2) - T)_+ \mathbb{1}_{\{T \leq t\}},$$

where $X_T^t(1) := F_{t-T}^{-1}(U)$, $X_T^t(2) := F_{t-T}^{-1}(\tilde{U})$ and $X_0 = F_0^{-1}(U)$. By (2.3), we observe immediately that $X_0^t(\emptyset)$ has law μ_t .

Now, on the event $\{T \leq t\}$, we consider 1 and 2 as two new roots, and we replace $X_T^t(1)$ and $X_T^t(2)$ by new random variables constructed using new independent exponential and uniform variables on an eventually enlarged probability space, constructed in the same way as $X_0^t(\emptyset)$. The procedure can be iterated until one obtains a Yule tree of height t , with i.i.d. random variables of law μ_0 attached to each leaf of \mathbb{T}_t , in such a way that $X_0^t(\emptyset)$ is the result of the painting scheme applied to that tree, which concludes the proof. \square

2.2 Associated McKean-Vlasov type SDE

In the previous section, we introduced a family $(\mu_t)_{t \geq 0}$ of probability distributions on \mathbb{R}_+ , obtained from the tree-painting scheme on a Yule tree. This procedure allows also to

introduce a Markov process $(X_s)_{s \in [0, t]}$ such that for each $s \in [0, t]$, X_s has law μ_s : for this, consider a Yule tree of height t and let $X_s := X_{t-s}^t(u)$ where u is the leftmost particle alive at time $t - s$. Note that time is running in the opposite direction as for the Yule tree.

This process X can be described as follows: X decreases linearly at speed 1 until reaching 0. At each branching time s of the leftmost particle, X_s is replaced by $X_s + \widetilde{X}_s$, where \widetilde{X}_s is independent of X_s with law μ_s . This description leads to Definition 1.4 which says that a *DR process* is a non-negative Markov process $(X_t)_{t \geq 0}$ solution to the McKean–Vlasov type stochastic differential equation:

$$\begin{cases} X_t = X_0 - \int_0^t \mathbb{1}_{\{X_s > 0\}} ds + \int_0^t \int_0^1 F_s^{-1}(u) N(ds, du) \\ F_t^{-1}(u) = \inf\{y \in [0, \infty) : \mathbb{P}(X_t \leq y) > u\}, \end{cases} \quad (2.4)$$

where N is a Poisson point process on $\mathbb{R}_+ \times [0, 1]$ with intensity $dsdu$. The drift term $-\mathbb{1}_{\{X_s > 0\}} ds$ indicates the amount of paint lost per unit of length, and each atom (s, u) of N corresponds to a time s at which the subtree gains a new branching time, in which case $F_s^{-1}(u)$ is an independent copy of X_s . The following result shows that the one-dimensional marginals of X constitutes a family obtained from the tree-painting scheme.

Proposition 2.2. *Let μ_0 be a probability distribution on \mathbb{R}_+ . There exists a unique strong solution to Equation (2.4) such that X_0 has law μ_0 . Moreover, writing μ_t the law of X_t for each $t \geq 0$, the family $(\mu_t)_{t \geq 0}$ is a family of laws obtained from the tree-painting scheme.*

Proof. We first prove there exists a solution to (2.4). To do so, let $(\mu_t)_{t \geq 0}$ be a family obtained from the tree-painting scheme. For all $t \geq 0$ and $u \in [0, 1]$, we denote by $G_t(u) = \inf\{y \in [0, \infty) : \mu_t([0, y]) > u\}$. We observe easily that there exists a unique strong solution to the SDE

$$X_t = X_0 - \int_0^t \mathbb{1}_{\{X_s > 0\}} ds + \int_0^t \int_0^1 G_s(u) N(ds, du),$$

by decomposing on the atoms of the point measure N . We are going to show that, if X_0 has law μ_0 , then X_t has law μ_t for each $t > 0$. This will prove that X is a strong solution of (2.4), as in that case we will have $G_s = F_s^{-1}$ for all $s \geq 0$.

Let $t > 0$, we denote by $(s_j, u_j)_{1 \leq j \leq n}$ the atoms of N in $[0, t] \times [0, 1]$, ranked in the increasing order of their first coordinate. For each $j \leq n$, we write $\widetilde{X}_j = G_{s_j}(u_j)$. We note that conditionally on $(s_j)_{1 \leq j \leq n}$, (\widetilde{X}_j) are independent random variables, and X_j has law μ_{s_j} . As $(\mu_s)_{s \geq 0}$ is a DR model, one can construct a Yule tree τ_j and a decoration of the leaves with i.i.d. random variables with law μ_0 , such that the amount of paint remaining at the root of τ_j is \widetilde{X}_j .

We then define a tree of height t with a decoration on its leaves as follows: we start with a line of height t , and for each $j \leq n$, we attach the decorated tree τ_j to the line at height $t - s_j$. To the top of the line we attach the random variable X_0 . This procedure constructs a Yule tree of height t , such that each leaf is decorated with an i.i.d. variable with law μ_0 (see Figure 6(b)). Therefore, X_t is obtained as the result of a tree-painting procedure, hence has law μ_t . We conclude that $(X_t, t \geq 0)$ is a strong solution to (2.4).

We now prove strong uniqueness of the solution to (2.4). We first show uniqueness in law, from which we will immediately deduce pathwise uniqueness. Let $(X_t, t \geq 0)$ be a solution to (2.4), for all $t \geq 0$, we denote by ν_t the law of X_t , and we define

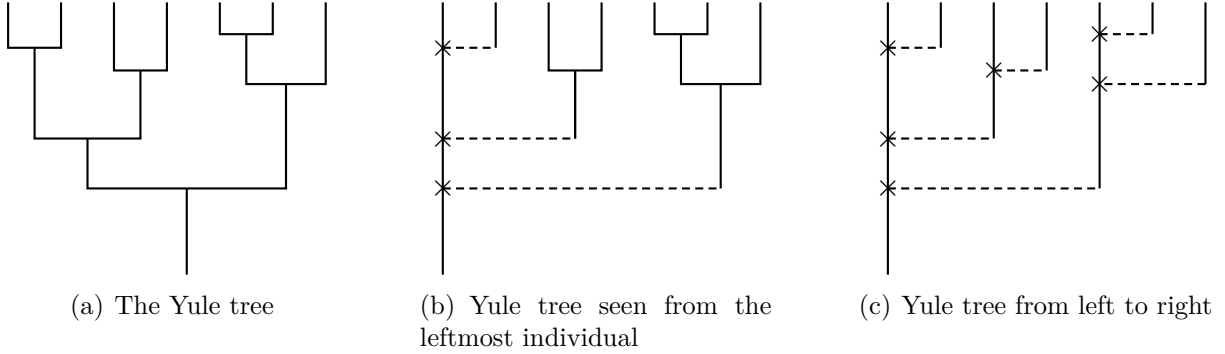


Figure 6: Two alternative constructions of a Yule tree

$H_t(u) = \inf\{y \geq 0 : \nu_t([0, y]) > u\}$. We first remark that there is strong existence and uniqueness of the solution of

$$X_t = X_0 - \int_0^t \mathbb{1}_{\{X_s > 0\}} ds + \int_0^t \int_0^1 H_s(u) N(ds, du), \quad (2.5)$$

and that X is this solution.

Let $t > 0$, our aim is to prove that $\nu_t = \mu_t$. For this, we show that, using the same procedure as above, we can represent X_t as the result of a tree-painting procedure on a Yule tree of height t , with leaves decorated with law μ_0 . Indeed, we denote once again $(s_j, u_j)_{1 \leq j \leq n}$ the atoms of the point process N on $[0, t] \times [0, 1]$ and $\tilde{X}_j := H_{s_j}(u_j)$. Observe that, conditionally on $(s_j)_{1 \leq j \leq n}$, each \tilde{X}_j is a copy of X_{s_j} , independent of $(\tilde{X}_i)_{i \neq j}$ and of X_0 . We construct the tree of height t as follows. We start with a line of height t , and for each $j \leq n$, we add a mark on this line at height $t - s_j$. At each of these marks is added the random variable \tilde{X}_j . As \tilde{X}_j is an independent copy of X_{s_j} , it can be itself constructed as the unique strong solution of Equation (2.5), where the point process N and the initial value X_0 are replaced by i.i.d. copies $N^{(j)}$ and $X_0^{(j)}$.

Hence, we can construct new lines and marks starting from each of the marks in the original line and proceed so on recursively. This constructs a Yule tree, defined in a slightly different way: indeed, in this tree, each line lives until time t , giving independently birth to children at rate 1 (this is the construction pictured in Figure 6(c)). Moreover, on each of the leaves at level t are placed i.i.d. random variables with law μ_0 , corresponding to the starting position for each of the stochastic differential equation. Therefore, X_t is once again constructed as a tree-painting procedure on a Yule tree of height t and we can conclude that $\nu_t = \mu_t$.

In particular, this proves that $H = G$ and, thus, that X is a solution to the SDE

$$X_t = X_0 - \int_0^t \mathbb{1}_{\{X_s > 0\}} ds + \int_0^t \int_0^1 G_s(u) N(ds, du),$$

for which there is pathwise uniqueness of the solution associated to a given X_0 and N . Hence, given X_0 and N , there exists a unique strong solution to (2.4), and its unidimensional marginals form a family of laws obtained from the tree-painting scheme. \square

This result proves there is equivalence between the family of laws defined via the tree-painting scheme, or by the McKean-Vlasov type equation. To complete the proof of Theorem 1.8, it is therefore enough to show that this family of law is the unique solution of the partial differential equation (2.4), which is done in the next section.

2.3 The partial differential equation representation

We observe here that a family of laws obtained from the tree-painting scheme can be rewritten as the weak solution of a partial differential equation, the key being the application of Itô's formula to the solution of the stochastic differential equation (2.4). We then study the existence and uniqueness of the solutions of this partial differential equation using the tree-painting scheme, and in particular Lemma 2.1.

As in Definition 1.5, we say that a family $(\nu_t)_{t \geq 0}$ of finite measures on \mathbb{R} is a weak solution to the partial differential equation

$$\partial_t \nu_t(dx) = \partial_x (\mathbb{1}_{\{x > 0\}} \nu_t(dx)) + (\nu_t * \nu_t)(dx) - \nu_t(dx), \quad (2.6)$$

if $t \mapsto \nu_t$ is measurable and, for all \mathcal{C}^1 bounded functions f with bounded derivative, we have

$$\begin{aligned} & \int_{\mathbb{R}} f(x) \nu_t(dx) - \int_{\mathbb{R}} f(x) \nu_0(dx) \\ &= \int_0^t \left(- \int_{\mathbb{R}} (f'(x) \mathbb{1}_{\{x > 0\}} + f(x)) \nu_s(dx) + \int_{\mathbb{R}^2} f(x+y) \nu_s(dx) \nu_s(dy) \right) ds. \end{aligned} \quad (2.7)$$

In particular, if $(\nu_t, t \geq 0)$ is a family of probability measures, this equation can be rewritten

$$\begin{aligned} & \int_{\mathbb{R}} f(x) \nu_t(dx) - \int_{\mathbb{R}} f(x) \nu_0(dx) \\ &= \int_0^t \left(- \int_{\mathbb{R}} f'(x) \mathbb{1}_{\{x > 0\}} \nu_s(dx) + \int_{\mathbb{R}^2} (f(x+y) - f(x)) \nu_s(dx) \nu_s(dy) \right) ds. \end{aligned} \quad (2.8)$$

The following result proves existence and uniqueness of weak solutions of (2.6) in the case where the initial condition is a probability measure on \mathbb{R}_+ . In other words, there exists a unique (continuous-time) DR model with initial distribution μ_0 , for all probability measure μ_0 on \mathbb{R}_+ .

Proposition 2.3. *Let $(\mu_t)_{t \geq 0}$ be a family of laws obtained from the tree-painting scheme, then $(\mu_t)_{t \geq 0}$ is a weak solution to (2.6).*

Reciprocally, if $(\nu_t)_{t \geq 0}$ is a family of non-negative finite measures on \mathbb{R} which is a weak solution to (2.6) and ν_0 is a probability distribution on \mathbb{R}_+ , then $(\nu_t)_{t \geq 0}$ is a family of laws obtained from the tree-painting scheme.

Proof. From Proposition 2.2, we observe that if $(\mu_t)_{t \geq 0}$ is a family of laws obtained from the tree-painting scheme, we can construct a Markov process $(X_t)_{t \geq 0}$ solution of (2.4) such that the law of X_t is μ_t for all $t \geq 0$. Therefore, for all \mathcal{C}^1 bounded function f , we have

$$\begin{aligned} \mathbb{E}[f(X_t)] &= \mathbb{E} \left[f(X_0) - \int_0^t f'(X_s) \mathbb{1}_{\{X_s > 0\}} ds \right] \\ &\quad + \mathbb{E} \left[\int_0^t \int_0^1 (f(X_{s-} + F_s^{-1}(u)) - f(X_{s-})) N(ds, du) \right], \end{aligned}$$

by the Itô formula. Therefore, using the linearity of the expectation and the Campbell's formula, we obtain

$$\mathbb{E}[f(X_t)] - \mathbb{E}[f(X_0)] = \int_0^t \left(- \mathbb{E}[f'(X_s) \mathbb{1}_{\{X_s > 0\}}] + \mathbb{E}[f(X_s + Y_s) - f(X_s)] \right) ds,$$

with Y_s an independent copy of X_s , proving that $(\mu_t)_{t \geq 0}$ is a weak solution of (2.6).

Reciprocally, let ν_0 be a probability distribution on \mathbb{R} and $(\nu_t)_{t \geq 0}$ be a weak solution of (2.6) (in particular, ν_t is finite). We first observe that for each $t > 0$, the total mass of ν_t remains equal to 1. Indeed, using $f = 1$, we have

$$\nu_t(\mathbb{R}) - \nu_0(\mathbb{R}) = \int_0^t (\nu_s(\mathbb{R})^2 - \nu_s(\mathbb{R})) ds.$$

Hence, the total mass is the solution of the differential equation $y' = -y(1-y)$ with initial condition $y(0) = 1$, which is $y(t) = 1$ for all $t \geq 0$, proving the total mass is preserved in (2.6). We conclude that $(\nu_t, t \geq 0)$ is a family of probability distributions.

We now prove that if ν_0 put no mass on $(-\infty, 0)$ then so does ν_t for all $t > 0$. To do so, we observe that for all \mathcal{C}^1 non-negative bounded function f with support in $(-\infty, 0)$, we have

$$\int f(x) \nu_t(dx) = \int_0^t \int (f(x+y) - f(x)) \nu_s(dx) \nu_s(dy) ds \geq - \int_0^t \int f(x) \nu_s(dx) ds,$$

hence the function $t \mapsto \int f(x) \nu_t(dx)$ remains equal to 0 at all times by Gronwall lemma. We conclude that $\nu_t((-\infty, 0)) = 0$ for all $t > 0$.

Finally, we prove that $(\nu_t)_{t \geq 0}$ satisfies Equation (2.3). To do so, let g be a \mathcal{C}^1 function with compact support in $(0, \infty)$. We define the function

$$G_g(a, b) := \int g((x-a)_+) \nu_b(dx).$$

By dominated convergence and (2.7), we observe that G_g is a \mathcal{C}^1 function, and that

$$\begin{aligned} \partial_a G_g(a, b) &= - \int g'(x-a) \mathbb{1}_{\{x>a\}} \nu_b(dx) \\ \partial_b G_g(a, b) &= - \int (g'(x-a) \mathbb{1}_{\{x>a\}} + g((x-a)_+)) \nu_b(dx) \\ &\quad + \int g((x+y-a)_+) \nu_b(dx) \nu_b(dy). \end{aligned}$$

Combining both equations, it follows that

$$\int g(z) \tau_a(\nu_b * \nu_b)(dz) = \partial_b G_g(a, b) - \partial_a(G_g(a, b)) + \int g((x-a)_+) \nu_b(dx).$$

and, therefore, we get

$$\begin{aligned} G_g(0, t) - e^{-t} G_g(t, 0) &= \int_0^t \partial_s (e^{-(t-s)} G_g(t-s, s)) ds \\ &= \int_0^t \int_0^t e^{-(t-s)} (G_g(t-s, s) - \partial_a G_g(t-s, s) + \partial_b G_g(t-s, s)) ds \\ &= \int_0^t e^{-s} g(z) \tau_s(\nu_{t-s} * \nu_{t-s})(dz) ds, \end{aligned}$$

using the change of variables $s \rightarrow t-s$ in the last line. This proves that $(\nu_t)_{t \geq 0}$ satisfies (2.3), hence by Lemma 2.1, this is a family of laws obtained from the tree-painting scheme, concluding the proof. \square

Remark 2.4. Note that by definition, there exists a unique family $(\mu_t)_{t \geq 0}$ of laws obtained from the tree-painting scheme and associated to a given initial probability measure μ_0 . As a result, Proposition 2.3 shows there exists a unique family of probability measures which are weak solutions of (2.6).

Proof of Theorem 1.8. The proof follows immediately from Propositions 2.2 and 2.3. \square

Thanks to Proposition 2.3, we observe that if we can find a solution to (2.6), then this family of probability measures is a DR model. Despite obtaining few information from (2.6) (for example, we were not able to prove the uniqueness of the solution in the space of distributions, or that a weak solution starting from a smooth initial condition remains smooth at all times), we were able to exhibit a set of measures, which is stable with respect to the continuous-time DR transformation.

We note that strong solutions to (2.6) can be described in the following way. Let p, ϕ be two C^1 functions, the family defined by

$$\mu_t(dx) = p(t)\delta_0(dx) + (1 - p(t))\phi(t, x)dx$$

is a DR model if and only if $p(0) \in [0, 1]$, $\phi(0, x) \geq 0$ and $\int \phi(0, x)dx = 1$, and

$$\begin{cases} p'(t) = (1 - p(t))(\phi(t, 0) - p(t)) \\ \partial_t \phi(t, \cdot) = \partial_x \phi(t, \cdot) + (1 - p(t))\phi(t, \cdot) * \phi(t, \cdot) + (p(t) + \phi(t, 0) - 1)\phi(t, \cdot). \end{cases} \quad (2.9)$$

Finding solutions to this differential equation, or proving existence, uniqueness and regularity for the solutions appears to be an interesting open problem.

3 Exactly solvable Derrida–Retaux model

In this section, we study a continuous-time DR model starting from an initial measure

$$\mu_0(dx) = p(0)\delta_0(dx) + (1 - p(0))\lambda(0)e^{-\lambda(0)x}dx, \quad p(0) \in [0, 1], \lambda(0) > 0, \quad (3.1)$$

which is a mixture of a Dirac mass at 0 and an exponential random variable (on \mathbb{R}_+). It turns out that the evolution of the DR model preserves the class of mixtures of exponential random variables and Dirac masses at 0. In other words, for all $t > 0$, there exists $p(t) \in [0, 1]$ and $\lambda(t) > 0$ such that

$$\mu_t(dx) = p(t)\delta_0(dx) + (1 - p(t))\lambda(t)e^{-\lambda(t)x}dx, \quad (3.2)$$

and the functions (p, λ) solve a system of ordinary differential equations. More precisely, the following result holds.

Lemma 3.1. *The family $(\mu_t)_{t \geq 0}$ is a DR model if and only if (p, λ) is solution of the system of ordinary differential equations*

$$\begin{cases} p' = (1 - p)(\lambda - p) \\ \lambda' = -\lambda(1 - p). \end{cases} \quad (3.3)$$

Proof. By Proposition 2.3, there exists a unique DR model with initial distribution μ_0 given by (3.1). Therefore, letting (p, λ) denote the maximal solution of the system of ordinary differential equations (3.3) with initial condition $(p(0), \lambda(0))$ (unique by Cauchy–Lipschitz theorem), it is enough to prove that this solution is well-defined for all $t \geq 0$ and that $(\mu_t)_{t \geq 0}$ is a DR model.

We first observe that if (p, λ) solve (3.3), then λ is non-increasing, and if $\lambda(0) > 0$ then $\lambda(t) > 0$ as long as the solution is well-defined. Similarly, if $p(0) \leq 1$ then $p(t) \leq 1$, and as $p'(t) \geq -(1 - p(t))p(t)$, $p(t) \geq 0$ for all t such that the solution is well-defined. As a result, we deduce that $(p(t), \lambda(t)) \in [0, 1] \times [0, \lambda_0]$. The solution remaining in a compact set, it cannot explode in finite time, therefore (p, λ) is well-defined for all $t \geq 0$.

We now prove that $(\mu_t)_{t \geq 0}$ is a DR model. By Proposition 2.3, it is enough to prove that the function is a weak solution to (2.6). Let f be a \mathcal{C}^1 bounded function with a bounded derivative. We define

$$g(t) = \int_{\mathbb{R}_+} f(x) \mu_t(dx) = p(t)f(0) + (1 - p(t))\lambda(t) \int_0^\infty f(x)e^{-\lambda(t)x} dx.$$

We write $E_0(t) := \lambda(t) \int_0^\infty f(x)e^{-\lambda(t)x} dx$ and $E_1(t) := \lambda(t) \int_0^\infty xf(x)e^{-\lambda(t)x} dx$. By dominated convergence, we observe that, on the one hand, the derivative of g is given by

$$\begin{aligned} g' &= p'f(0) + \left(-p' + (1 - p)\frac{\lambda'}{\lambda}\right)E_0 - (1 - p)\lambda'E_1 \\ &= (1 - p)(\lambda - p)f(0) - (1 - p)(1 + \lambda - 2p)E_0 + (1 - p)^2\lambda E_1. \end{aligned}$$

On the other hand, simple computations yield

$$\begin{aligned} \int_{\mathbb{R}_+} (f'(x)\mathbb{1}_{\{x>0\}} + f(x))\mu_t(dx) &= p(t)f(0) + \lambda(t)(1 - p(t)) \int_0^\infty (f'(x) + f(x))e^{-\lambda(t)x} dx \\ &= (p(t) - \lambda(t)(1 - p(t)))f(0) + (1 - p(t))(1 + \lambda(t))E_0(t), \end{aligned}$$

by integration by part, and

$$\begin{aligned} &\int_{(\mathbb{R}_+)^2} f(x + y)\mu_t(dx)\mu_t(dy) \\ &= p(t)^2f(0) + 2(1 - p(t))E_0(t) + (1 - p(t))^2\lambda(t)^2 \int_{(\mathbb{R}_+)^2} f(x + y)e^{-\lambda(t)(x+y)} dx dy \\ &= p(t)^2f(0) + 2p(t)(1 - p(t))E_0(t) + (1 - p(t))^2\lambda(t)E_1(t). \end{aligned}$$

As a result, we obtain that

$$- \int_{\mathbb{R}_+} (f'(x)\mathbb{1}_{\{x>0\}} + f(x))\mu_t(dx) + \int_{(\mathbb{R}_+)^2} f(x + y)\mu_t(dx)\mu_t(dy) = g'(t)$$

showing that $(\mu_t)_{t \geq 0}$ is a DR model, which concludes the proof. \square

As a result of Lemma 3.1, to study the DR model with initial law μ_0 given in (3.1), it is enough to study the system of ordinary differential equations (3.3). This is what is done in the rest of the section, allowing us in particular to prove the DR conjecture in this particular setting.

Remark 3.2. Observe that, similarly, the evolution of a DR model starting by any mixture of Dirac mass at 0 and exponential random variables at time 0 can be represented by a system of differential equations. More precisely, for any $n \in \mathbb{N}$ the measure

$$\nu_t(dx) = p(t)\delta_0 + (1 - p(t)) \sum_{j=1}^n q_j(t)\lambda_j(t)e^{-\lambda_j(t)x} dx$$

is a DR model if and only if the functions p , q_j , λ_j satisfy

$$\begin{cases} p' = (1 - p)(\sum_{k=1}^n q_k \lambda_k - p) \\ q_j' = -q_j \left((1 - p)(1 - q_j + 2 \sum_{k \neq j} \frac{q_k \lambda_k}{\lambda_j - \lambda_k}) + \lambda_j - \sum_{k=1}^n q_k \lambda_k \right) \\ \lambda_j' = -(1 - p)q_j \lambda_j, \end{cases} \quad (3.4)$$

and if $p(0), q_j(0) \in [0, 1]$, $\lambda_j(0) \geq 0$ and $\sum_{j=1}^n q_j(0) = 1$. However, we concentrate on the simple mixture of a Dirac mass at 0 and one exponential random variable in this article, as this simple model exhibits the main features believed to hold for any DR model.

In the next section, we give the phase diagram of the differential system (3.3), with a description of all stable and unstable equilibrium points, as well as their zone of attraction. In Section 3.2, we refine the study by giving asymptotic equivalents for the solutions of (3.3). Finally, we use these results to prove in Section 3.3 the DR conjecture for continuous models with exponential initial condition.

3.1 Phase diagram of the model

We first observe that some quantities are preserved under the dynamics (3.3).

Lemma 3.3. *The function $t \mapsto \frac{p(t)}{\lambda(t)} + \log \lambda(t)$ is a constant.*

Proof. We compute the derivative of the function $H : t \mapsto \frac{p(t)}{\lambda(t)} + \log \lambda(t)$, we have

$$\begin{aligned} H'(t) &= \frac{p'(t)}{\lambda(t)} - \frac{p(t)\lambda'(t)}{\lambda(t)^2} + \frac{\lambda'(t)}{\lambda(t)} \\ &= \frac{(1 - p(t))(\lambda(t) - p(t))}{\lambda(t)} + \frac{p(t)\lambda(t)(1 - p(t))}{\lambda(t)^2} - \frac{\lambda(t)(1 - p(t))}{\lambda(t)} = 0, \end{aligned}$$

which concludes the proof. □

As a straightforward consequence of this remark, we deduce that p can be written as an explicit function of λ , for all initial condition.

Corollary 3.4. *Let $H := \frac{p(0)}{\lambda(0)} + \log \lambda(0)$, for all $t \geq 0$, we have*

$$p(t) = H\lambda(t) - \lambda(t) \log \lambda(t). \quad (3.5)$$

In particular, λ is the solution of the ordinary differential equation

$$\lambda' = -\lambda(1 - H\lambda + \lambda \log \lambda). \quad (3.6)$$

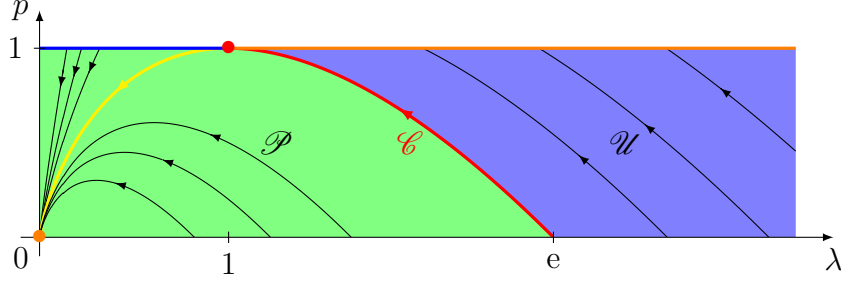


Figure 7: The phase diagram of the differential system (3.3)

Using Lemma 3.3, we can draw the phase diagram of the differential system (3.3), obtain the asymptotic behavior of (p, λ) as $t \rightarrow \infty$ and describe all the equilibrium points of (3.3), with their stability properties.

Recall the three sets \mathcal{P} , \mathcal{C} and \mathcal{U} defined in the introduction, respectively the supercritical, critical and subcritical zones, which we recall on Figure 7. We added on Figure 3, the trajectories of some solutions of the equation (in black). Some interesting additional sets are represented on that figure, such as the supercritical branch (in yellow) of the critical curve

$$\widehat{\mathcal{C}} = \{(p, \lambda) \in [0, 1) \times [0, 1) : p = \lambda - \lambda \log \lambda\},$$

as well as the attractive, critical and repulsive points of the dynamics (in orange, red and dark blue respectively):

$$\mathcal{A}_p = \{(0, 0)\} \cup (\{1\} \times (1, \infty)), \quad \mathcal{C}_p = \{(1, 1)\} \quad \text{and} \quad \mathcal{R}_p = \{1\} \times [0, 1).$$

We now describe the asymptotic behavior of this DR model as $t \rightarrow \infty$.

Proposition 3.5. *Let (p, λ) be a solution of (3.3) on \mathbb{R}_+ . If $(p(0), \lambda(0)) \in [0, 1] \times [0, \infty)$, then the domain of definition of (p, λ) includes $[0, \infty)$, and moreover,*

1. *if $(p(0), \lambda(0)) \in \mathcal{P}$, then $\lim_{t \rightarrow \infty} (p(t), \lambda(t)) = (0, 0)$.*
2. *if $(p(0), \lambda(0)) \in \mathcal{U}$, then $\lim_{t \rightarrow \infty} (p(t), \lambda(t)) = (1, x)$, where x is the unique solution larger than 1 of the equation $Hx - x \log x = 1$.*
3. *if $(p(0), \lambda(0)) \in \mathcal{C}$, then $\lim_{t \rightarrow \infty} (p(t), \lambda(t)) = (1, 1)$.*

Proof. We first remark from (3.3) that the function λ is non-increasing. Moreover, we observe the following behavior for the solutions of (3.3) starting from the boundary of the domain:

- if $p(0) = 1$, then $(p(t), \lambda(t)) = (1, \lambda(0))$ for all $t \in \mathbb{R}$,
- if $\lambda(0) = 0$ and $p(0) \in (0, 1)$, then $(p(t), \lambda(t)) = (\frac{1}{(p(0)^{-1}-1)e^t+1}, 0)$ for all $t \in \mathbb{R}$,
- if $\lambda(0) = p(0) = 0$ then $(p(t), \lambda(t)) = (0, 0)$ for all $t \in \mathbb{R}$,
- if $\lambda(0) > 0$ and $p(0) = 0$, then for all $h > 0$ small enough, we have $(p(h), \lambda(h)) \in (0, 1) \times (0, \infty)$.

The possible values of $\lim_{t \rightarrow \infty} (p(t), \lambda(t))$ then belong to the set of equilibrium points of the dynamics, defined by the system of equation

$$\begin{cases} 0 = (1-p)(\lambda-p) \\ 0 = -\lambda(1-p), \end{cases}$$

which is exactly $\mathcal{A}_p \cup \mathcal{C}_p \cup \mathcal{R}_p$. Moreover, recall that $(p(t), \lambda(t))$ move along the trajectory of $x \mapsto Hx - x \log x$, as $\lambda(t)$ decreases.

Hence, if $(p(0), \lambda(0)) \in \mathcal{P}$, then either $H < 1$, and the only equilibrium point belonging to $x \mapsto Hx - x \log x$ is $(0, 0)$, or $H > 1$ and this is the only equilibrium point left of the starting point. If $(p(0), \lambda(0)) \in \mathcal{U}$, the first equilibrium point met by the trajectory is the one described in the second point. Finally, if $(p(0), \lambda(0)) \in \mathcal{C}$, then the trajectory of (p, λ) would go through the critical point $(1, 1)$ which is impossible, proving the third case. \square

Remark 3.6. Note that $(p(0), \lambda(0)) \in [0, 1] \times [0, 1]$ and $H \geq 1$ is a necessary and sufficient condition for the existence of a meaningful solution of (3.6) defined on \mathbb{R} . From any other starting position, there exists a finite $t \in (-\infty, 0)$ such that $p(t) < 0$, therefore the DR model cannot be extended as a process on \mathbb{R} in general.

3.2 Asymptotic behavior of the solutions

In this section, we study the asymptotic behavior of the solution $(p(t), \lambda(t))$ as $t \rightarrow \infty$, in the three different phases. In the critical case, this corresponds to Conjectures 1.3 and 1.4 of Chen et al. [5] in the discrete setting. Using the properties of the differential system, we were able to study the asymptotic behavior in the critical case beyond the computation of an equivalent.

Proposition 3.7. *Let (p, λ) be a solution of (3.3).*

1. *If $(p(0), \lambda(0)) \in \mathcal{P}$ with $\lambda(0) > 0$, then there exists $K > 0$ such that, as $t \rightarrow \infty$,*

$$\lambda(t) \sim Ke^{-t} \quad \text{and} \quad p(t) \sim Kte^{-t}.$$

2. *If $(p(0), \lambda(0)) \in \mathcal{U}$, then there exists $K > 0$ such that, as $t \rightarrow \infty$,*

$$\lambda(t) - x \sim Ke^{-(x-1)t} \quad \text{and} \quad 1 - p(t) \sim \left(1 - \frac{1}{x}\right)Ke^{-(x-1)t},$$

where x is the unique solution larger than 1 of the equation $Hx - x \log x = 1$.

3. *If $(p(0), \lambda(0)) \in \mathcal{C}$, then we have, as $t \rightarrow \infty$,*

$$\lambda(t) = 1 + \frac{2}{t} - \frac{8 \log t}{3t^2} + o\left(\frac{\log t}{t^2}\right) \quad \text{and} \quad p(t) = 1 - \frac{2}{t^2} + \frac{16 \log t}{3t^3} + o\left(\frac{\log t}{t^3}\right).$$

Proof. Let start with the supercritical case $(p(0), \lambda(0)) \in \mathcal{P}$. Recall from Theorem 1.9 that $\lambda(t) \rightarrow 0$. Using (3.6), we get that $\lambda'(t)/\lambda(t) \rightarrow -1$ and therefore

$$\log \lambda(t) - \log \lambda(0) = \int_0^t \frac{\lambda'(s)}{\lambda(s)} ds \sim -t,$$

hence $\lambda(t) = e^{-t+o(t)}$. Using again (3.6), we get that

$$\frac{\lambda'(t)}{\lambda(t)} + 1 = -\lambda(t) \log \lambda(t) (1 + o(1)) = -e^{-t+o(t)},$$

which is integrable at infinity. Therefore, it follows that

$$\log \lambda(t) - \log \lambda(0) + t = \int_0^t \left(\frac{\lambda'(s)}{\lambda(s)} + 1 \right) ds \xrightarrow{t \rightarrow \infty} \int_0^\infty \left(\frac{\lambda'(s)}{\lambda(s)} + 1 \right) ds,$$

where the last integral is finite. This proves that $\lambda(t) \sim Ke^{-t}$ for some constant $K > 0$. Then, the behavior of $p(t)$ follows directly from (3.5).

We now deal with the subcritical case $(p(0), \lambda(0)) \in \mathcal{U}$. Recall from Theorem 1.9 that $\lambda(t) \rightarrow x$, where $x > 1$ satisfies $Hx - x \log x = 1$. We introduce the auxiliary function $u(t) := \lambda(t) - x$, which is positive for all $t \in \mathbb{R}_+$. Using (3.6), we get

$$\begin{aligned} u' &= -(x+u)(1 - H(x+u) + (x+u) \log(x+u)) \\ &= -(x+u) \left(-Hu + x \log \left(\frac{x+u}{x} \right) + u \log(x+u) \right), \end{aligned}$$

using that $Hx - x \log x = 1$. Since $u(t) \rightarrow 0$, we get the following asymptotic development

$$u'(t) = -x(-H + 1 + \log x)u(t) + O(u(t)^2) = -(x-1)u(t) + O(u(t)^2). \quad (3.7)$$

Proceeding as in the supercritical case, we first deduce from (3.7) that $u(t) = e^{-(x-1)t+o(t)}$ and, then, that $u(t) \sim Ke^{-(x-1)t}$. Plugging this in (3.5), the behavior of $p(t)$ follows.

Finally, we deal with the critical case $(p(0), \lambda(0)) \in \mathcal{C}$. Recall from Theorem 1.9 that $\lambda(t) \rightarrow 1$, so we introduce the auxiliary function $v(t) := \lambda(t) - 1$, which is positive for all $t \in \mathbb{R}_+$. Using (3.6) with here $H = 1$, we get

$$v' = -(1+v)(1 - (1+v) + (1+v) \log(1+v)),$$

and, using that $v(t) \rightarrow 0$, we get

$$v'(t) = -\frac{v(t)^2}{2} - \frac{v(t)^3}{3} + O(v(t)^4). \quad (3.8)$$

First, we deduce from (3.8) that $(1/v)'(t) \rightarrow 1/2$ as $t \rightarrow \infty$ and therefore,

$$\left(\frac{1}{v(t)} - \frac{1}{v(0)} \right) = \int_0^t \left(\frac{1}{v} \right)'(s) ds \sim \frac{t}{2}.$$

It proves that $v(t) \sim 2/t$ as $t \rightarrow \infty$. Then, using again (3.8), we have

$$\left(\frac{1}{v} \right)'(t) - \frac{1}{2} = \frac{v(t)}{3} + o(v(t)) \sim \frac{2}{3t},$$

as $t \rightarrow \infty$, and it follows that

$$\int_1^t \left(\left(\frac{1}{v} \right)'(s) - \frac{1}{2} \right) ds \sim \int_1^t \frac{2}{3s} ds = \frac{2}{3} \log t,$$

and, finally,

$$v(t) = \left(\frac{t}{2} + \frac{2}{3} \log t + o(\log t) \right)^{-1} = \frac{2}{t} - \frac{8 \log t}{3t^2} + o\left(\frac{\log t}{t^2}\right). \quad (3.9)$$

Then, it follows from (3.5) that

$$p(t) = (1 + v(t)) - (1 + v(t)) \log(1 + v(t)) = 1 - \frac{v(t)^2}{2} + O(v(t)^3),$$

and plugging (3.9) in this equation gives the desired result. \square

Proof of Theorem 1.9. Recalling from (3.2) that

$$\mu_t(dx) = p(t)\delta_0(dx) + (1 - p(t))\lambda(t)e^{-\lambda(t)x}dx.$$

Note that $\int_{\mathbb{R}_+} x\mu_t(dx) = \frac{1-p(t)}{\lambda(t)}$. Theorem 1.9 is now an immediate consequence of the asymptotic behavior of $(p(t), \lambda(t))$ stated in Proposition 3.7. \square

3.3 Proof of the Derrida–Retaux conjectures

Using the fact that the DR model we are considering consists in a mixture of exponential random variables and Dirac masses at 0 and is integrable, we obtain a straightforward proof of the DR conjecture for this model, which states that the free energy has an infinite order phase transition between the supercritical and the subcritical phases. We begin by observing that if we write

$$\mu_0(dx) = p(0)\delta_0(dx) + (1 - p(0))\lambda(0)e^{-\lambda(0)x}dx,$$

then the free energy can be rewritten

$$F_\infty(\mu_0) = F_\infty(p(0), \lambda(0)) = \lim_{t \rightarrow \infty} e^{-t} \frac{1 - p(t)}{\lambda(t)}. \quad (3.10)$$

Remark 3.8. The proof of Theorem 1.10 is crucially based on Corollary 3.4, which states that λ satisfies an ordinary differential equation. In the case $\lambda \in [1, e)$, note that as $p(0)$ approaches p_c , the trajectory followed by (p, λ) approaches closer and closer to the critical point $(1, 1)$ (see Figure 7). Moreover, trajectories typically stay a long time around this critical point, as both p' and λ' become small (of order $(1 - p)$). The main asymptotic behavior of the free energy turns out to be driven by the amount of time “lost” by the trajectory (p, λ) in the neighborhood of $(1, 1)$ before eventually converging toward $(0, 0)$.

Proof of Theorem 1.10. We fix some $\lambda \in (0, e)$ and set $p_c := \lambda - \lambda \log \lambda$ if $\lambda > 1$ and, by extension, $p_c := 1$ if $\lambda \leq 1$. We work with p such that $(p, \lambda) \in \mathcal{P}$, which means exactly $p < p_c$. Let (p, λ) be a solution to (3.3) with initial conditions $\lambda(0) = \lambda$ and $p(0) = p$. Recall that setting $H := \frac{p}{\lambda} + \log \lambda$, the function λ is solution to the differential equation

$$\forall t \geq 0, \quad \lambda'(t) = -\lambda(t)(1 - H\lambda(t) + \lambda(t) \log \lambda(t)),$$

which can be rewritten as: for all $t \geq 0$, we have

$$\int_{\lambda(t)}^{\lambda(0)} \frac{dy}{y(1 - Hy + y \log y)} = t.$$

We take interest in the asymptotic behavior, as $x \rightarrow 0$ of the function

$$G_H: x \mapsto \int_x^\lambda \frac{dy}{y(1 - Hy + y \log y)}.$$

Indeed, since $(p, \lambda) \in \mathcal{P}$, we have $p(t) \rightarrow 0$ and $\lambda(t) \rightarrow 0$ by Theorem 1.9 and it follows from (3.10) that

$$F_\infty(p, \lambda) = \lim_{t \rightarrow \infty} \frac{e^{-t}}{\lambda(t)} = \exp\left(-\lim_{x \rightarrow 0} (G_H(x) + \log x)\right), \quad (3.11)$$

by composition of limits. Thus, in order to prove the result, it is enough to compute the asymptotic behavior, as $p \rightarrow p_c$ or equivalently $H \rightarrow H_c := \frac{p_c}{\lambda} + \log \lambda$, of

$$\begin{aligned} \lim_{x \rightarrow 0} (G_H(x) + \log x) &= \log \lambda + \int_0^\lambda \left(\frac{1}{y(1 - Hy + y \log y)} - \frac{1}{y} \right) dy \\ &= \log \lambda + \int_0^\lambda \frac{H - \log y}{1 - Hy + y \log y} dy, \end{aligned} \quad (3.12)$$

in the three different cases $\lambda \in (1, e)$, $\lambda = 1$ and $\lambda \in (0, 1)$.

We start with the case $\lambda \in [1, e)$ (we will distinguish later between $\lambda = 1$ and $\lambda \neq 1$), in which case $H_c = 1$. We first observe that

$$C_1 := \lim_{H \rightarrow 1} \int_0^\lambda \left(\frac{H - \log y}{1 - Hy + y \log y} - \frac{H + 1 - y}{(1 - H)y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} \right) dy$$

exists and is finite, by dominated convergence¹. Setting $\delta := 1 - H$, it follows that

$$\lim_{x \rightarrow 0} (G_H(x) + \log x) = \log \lambda + C_1 + \int_0^\lambda \frac{2 - \delta - y}{\delta y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} dy + o(1), \quad (3.13)$$

where the $o(1)$ holds as $\delta \rightarrow 0$, or equivalently as $p \rightarrow p_c$. In order to compute the integral, we need to compute the partial fraction decomposition of the rational function in the integral. The polynomial $\delta y + \frac{(y-1)^2}{2} + \frac{(y-1)^3}{6}$ has a real root x_0 and two complex conjugate roots $x_1 + iy_1$ and $x_1 - iy_1$. These roots are explicit and one can check that $x_0 = 4 + O(\delta)$, $x_1 = 1 - O(\delta)$ and $y_1 = \sqrt{2\delta} + O(\delta)$. Then, we have the following decomposition

$$\frac{2 - \delta - y}{\delta y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} = \frac{a}{y - x_0} + \frac{b - ay}{(y - x_1)^2 + y_1^2},$$

with $a = \frac{4}{3} + O(\delta)$ and $b = \frac{10}{3} + O(\delta)$. Therefore, we can finally compute

$$\begin{aligned} &\int_0^\lambda \frac{2 - \delta - y}{\delta y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} dy \\ &= a(\ln(x_0 - \lambda) - \ln(x_0)) + \frac{b - ax_1}{y_1} \left(\arctan\left(\frac{\lambda - x_1}{y_1}\right) - \arctan\left(\frac{-x_1}{y_1}\right) \right) \\ &\quad - \frac{a}{2} \left(\ln((\lambda - x_1)^2 + y_1^2) - \ln(x_1^2 + y_1^2) \right). \end{aligned}$$

¹Note that the main issue to be addressed is the divergence at 1 as $(y-1)^{-2}$ in the limit. For this, we replaced the $\log y$ in the numerator by $y-1$ and the $1+y \log y$ in the denominator by its Taylor expansion $1 + (y-1) + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}$. For the domination, one can use the triangle inequality, in order to make one replacement after the other, and then deal with them separately, using Taylor–Lagrange inequality.

From now on, we distinguish $\lambda > 1$ and $\lambda = 1$. We start with $\lambda > 1$: using the development of the different constants and that $\arctan(t) = \frac{\pi}{2} + O(\frac{1}{t})$ as $t \rightarrow \infty$, we get

$$\int_0^\lambda \frac{2 - \delta - y}{\delta y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} dy = \pi \sqrt{\frac{2}{\delta}} + C_2 + o(1),$$

where C_2 is a constant depending on λ . Recalling that $\delta = 1 - \frac{p}{\lambda} - \log \lambda = \frac{p_c - p}{\lambda}$, coming back to (3.11), (3.12) and (3.13), we get

$$F_\infty(p, \lambda) = \exp\left(-\pi \sqrt{\frac{2\lambda}{p_c - p}} - \log \lambda - C_1 - C_2 + o(1)\right)$$

and it concludes the proof of the first point.

We now consider the case $\lambda = 1$. In that case $\lambda - x_1 = O(\delta)$, so the development of the integral is slightly different:

$$\int_0^\lambda \frac{2 - \delta - y}{\delta y + \frac{(y-1)^2}{2} - \frac{(y-1)^3}{6}} dy = \frac{\pi}{\sqrt{2\delta}} - \frac{2}{3} \ln \delta + C_3 + o(1),$$

and, in the same way,

$$F_\infty(p, \lambda) = \exp\left(-\frac{\pi}{\sqrt{2(p_c - p)}} + \frac{2}{3} \ln(p_c - p) - \log \lambda - C_1 - C_3 + o(1)\right),$$

which proves the second point.

Finally, we deal with $\lambda \in (0, 1)$ and, therefore, $p_c = 1$. Recall from (3.11) and (3.12), that our aim is to compute the asymptotic behavior, as $p \rightarrow 1$, of

$$\int_0^\lambda \frac{H - \log y}{1 - Hy + y \log y} dy = \int_0^\lambda \frac{p - \lambda \log \frac{y}{\lambda}}{\lambda - py + \lambda y \log \frac{y}{\lambda}} dy,$$

where we used that $H = \frac{p}{\lambda} + \log \lambda$. Note that

$$C_4 := \lim_{p \rightarrow 1} \int_0^\lambda \left(\frac{p - \lambda \log \frac{y}{\lambda}}{\lambda - py + \lambda y \log \frac{y}{\lambda}} - \frac{p}{\lambda(1 - \lambda) - y(p - \lambda)} \right) dy$$

exists and is finite, by dominated convergence². Then, we compute

$$\int_0^\lambda \frac{p}{\lambda(1 - \lambda) - y(p - \lambda)} dy = -\frac{p}{p - \lambda} \log\left(\frac{1 - p}{1 - \lambda}\right),$$

and conclude that

$$F_\infty(p, \lambda) = \exp\left(\frac{1}{1 - \lambda} \log\left(\frac{1 - p}{1 - \lambda}\right) - \log \lambda - C_4 + o(1)\right),$$

which proves the third point. □

²Here, the issue is the divergence at λ as $(y - \lambda)^{-1}$ in the limit and, therefore, we replaced the $\log \frac{y}{\lambda}$ in the numerator by 0 and the $y \log \frac{y}{\lambda}$ in the denominator by $\lambda(\frac{y}{\lambda} - 1)$.

4 Critical Derrida–Retaux tree conditioned on survival

As observed in the previous section, a particular trajectory in the DR model exhibits some interesting features. It is the critical line, followed by a DR model with initial law μ_0 given by

$$\mu_0(dx) = p(0)\delta_0 + (1 - p(0))\lambda(0)e^{-\lambda(0)x}dx.$$

with $(p(0), \lambda(0)) \in \mathcal{C}$. From this starting condition, the law μ_t is at all times $t > 0$ given by a mixture of an exponential random variable and a Dirac mass at 0.

In this section, we consider a DR tree $X^t = (X_s^t(u), s \in [0, t], u \in \mathcal{N}_s)$ with initial law μ_0 , as defined in Section 2.1, recalling that \mathcal{N}_s is the set of particles of the Yule tree alive at time s and $X_s^t(u)$ the amount of paint remaining at level s in the branch u . The main objective of the section is to construct, describe and analyze the law of the portion of the Yule tree covered in paint, conditionally on the event that there is “paint” remaining at the origin. The conditioned tree turns out to behave as a time-inhomogeneous growth-fragmentation process, which converges to some limiting process. It is believed that the limiting tree is an universal characteristic, and would appear in numerous models related to DR at criticality. We have some anecdotal evidences of this universality, based on simulations. In Section 5, we take particular interest in the asymptotic behavior of the number and the mass on the leaves of that tree.

On the event $\{X_0^t(\emptyset) > 0\}$, we say that a particle $u \in \mathcal{N}_s$ for $s \leq t$ is *red* if it belongs to the painted connected component of the root: more formally, for any $r \leq s$ and $v \in \mathcal{N}_r$ such that $v \leq u$, $X_r^t(v) > 0$. Let $(\mathfrak{X}_s^t(v), v \in \mathfrak{N}_s^t, s \leq t)$ denote the restriction of the process X^t to the painted connected component of the root. Note that the labeling of particles is modified between the processes X^t and \mathfrak{X}^t . Indeed, if at the first branching time of X^t , all the paint comes from one of the children, the process \mathfrak{X}^t does not “see” this branching event, hence keep calling the line \emptyset . The reindexed portion of the Yule tree, denoted by $(\mathfrak{N}_s^t, s \leq t)$, is called the *red tree* (see Figure 8); note that it depends on t , unlike the Yule tree used to define X^t . By analogy with the definition of the Yule tree in Subsection 2.1, we denote by \mathfrak{e}_u^t , \mathfrak{b}_u^t and \mathfrak{d}_u^t respectively the lifetime, birth time and death time of u in the red tree $(\mathfrak{N}_s^t, s \leq t)$. By convention, all particles die at time t and all their offspring have zero lifetime.

4.1 The red tree as a Markovian branching process

The purpose of this subsection is the description of the law of $(\mathfrak{X}_s^t(v), v \in \mathfrak{N}_s^t, s \leq t)$ conditioned by $X_0^t(\emptyset) = x$ for some $x > 0$. However, the event $\{X_0^t(\emptyset) = x\}$ has zero probability, so the question is not well-posed. But, we give a particular description, as a Markovian branching process, of the law of \mathfrak{X}^t given $X_0^t(\emptyset)$, such that, for this description, it makes sense to consider the process “starting from x ”.

Let $(p(t), \lambda(t), t \geq 0)$ be the solution of (3.3) starting from initial condition $p(0) = p$, $\lambda(0) = \lambda$ for some $(p, \lambda) \in [0, 1) \times (0, \infty)$. Note that in this section, we *do not* restrict ourselves to the critical case. We introduce the auxiliary function

$$\varrho(s) := (1 - p(s))\lambda(s), \quad s \geq 0,$$

which will appear regularly in the sequel.

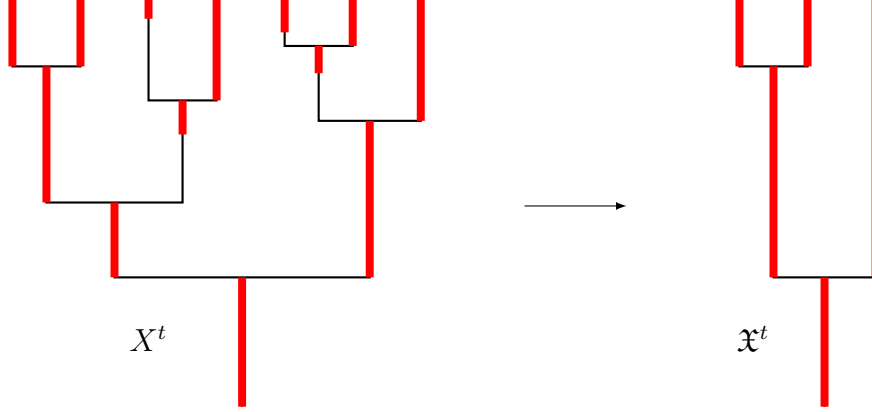


Figure 8: Illustration of the construction of the red tree, by pruning the Yule tree at points where there is no paint.

Proposition 4.1. *Assume that the initial law μ_0 is given by (3.1) with $(p(0), \lambda(0)) \in (0, 1) \times (0, \infty)$. Given $X_0^t(\emptyset)$, on the event $\{X_0^t(\emptyset) > 0\}$, \mathfrak{X}^t has the law of a time-inhomogeneous branching Markov process, which can be described as follows:*

1. *The process starts at time 0 with a unique particle of mass $\mathfrak{X}_0^t(\emptyset) = X_0^t(\emptyset)$.*
2. *The mass associated to each particle grows linearly at speed 1.*
3. *A particle of mass m at time s splits at rate $\rho(t-s)m$ into two children, the mass m being split uniformly between the two children.*
4. *Particles behave independently after their splitting time.*

This proposition allows us to define the distribution of \mathfrak{X}^t “starting from x ”. Indeed, for any $x > 0$, let \mathbb{P}_x denote a probability measure such that, under \mathbb{P}_x , the process \mathfrak{X}^t is the time-inhomogeneous branching Markov process starting with $\mathfrak{X}_0^t(\emptyset) = x$ and satisfying Points 2, 3 and 4 of Proposition 4.1. Then, informally, the law of \mathfrak{X}^t under \mathbb{P} conditioned by $X_0^t(\emptyset) = x$ is the law of \mathfrak{X}^t under \mathbb{P}_x . Before proving Proposition 4.1, the following lemma shows that the whole process X^t satisfies also a branching property.

Lemma 4.2. *Let $s \in (0, t)$. Given the Yule tree and the process X^t until time s , the $(X_{s+r}^t(v), r \in [0, t-s], v \in \mathcal{N}_{s+r}$ s.t. $v \geq u$) for $u \in \mathcal{N}_s$ are independent with respectively the same law as X^{t-s} starting from $X_s^t(u)$.*

Proof. Let $\mathbb{T}_s := (\mathcal{N}_r, r \in [0, s])$ denote the Yule tree until time s . Our aim is to prove that, for any family of nonnegative measurable functions $(F_u)_{u \in \mathbb{T}}$ and any r.v. Υ_s measurable w.r.t. the σ -field generated by \mathbb{T}_s and the process X^t until time s ,

$$\begin{aligned} & \mathbb{E} \left[\Upsilon_s \prod_{u \in \mathcal{N}_s} F_u \left(X_{s+r}^t(v), r \in [0, t-s], v \in \mathcal{N}_{s+r} \text{ s.t. } v \geq u \right) \right] \\ &= \mathbb{E} \left[\Upsilon_s \prod_{u \in \mathcal{N}_s} \tilde{F}_u \left(X_s^t(u) \right) \right]. \end{aligned} \quad (4.1)$$

where, for $u \in \mathbb{T}$, $\tilde{F}_u: \mathbb{R}_+ \rightarrow \mathbb{R}_+$ is the measurable function (unique up to a Lebesgue-negligible set) such that

$$\tilde{G}_u(X_0^{t-s}(\emptyset)) := \mathbb{E}\left[F_u\left(X_r^{t-s}(v), r \in [0, t-s], v \in \mathcal{N}_r\right) \middle| X_0^{t-s}(\emptyset)\right].$$

The process X^t until time s is a deterministic function of \mathbb{T}_s and of the $X_u(s)$ for $u \in \mathcal{N}_s$. Therefore, it suffices to prove (4.1) for Υ_s of the form $\Upsilon_s = g(\mathbb{T}_s) \prod_{u \in \mathcal{N}_s} h_u(X_u(s))$, with g and $(h_u)_{u \in \mathbb{T}}$ measurable functions. But, each factor $h_u(X_u(s))$ can be incorporated in $F_u(\dots)$, with an appropriate change of the function F_u . Hence, it is finally enough to consider $\Upsilon_s = g(\mathbb{T}_s)$.

Given \mathbb{T}_s , the subtrees rooted in the u 's for $u \in \mathcal{N}_s$ are independent Yule trees of height $t-s$ and the labels on their leaves are independent of each other. Therefore, given \mathbb{T}_s , the $(X_{s+r}^t(v), r \in [0, t-s], v \in \mathcal{N}_{s+r} \text{ s.t. } v \geq u)$ for $u \in \mathcal{N}_s$ are independent copies of X^{t-s} and it proves (4.1) in the case $\Upsilon_s = g(\mathbb{T}_s)$. This concludes the proof. \square

Remark 4.3. Note that this proof does not use the fact that we are working with the exponential model. Therefore, for any initial law μ_0 , X^t is a time-inhomogeneous branching Markov process.

Proof of Proposition 4.1. Points 1 and 2 are clearly true. Point 4 follows directly from the branching property for X^t established in Lemma 4.2. We focus now on Point 3. By the branching property, it is sufficient to prove it for the root of the tree: hence, our aim is to prove that, for any $0 \leq s < t$ and any measurable functions $f, g_1, g_2: \mathbb{R} \rightarrow \mathbb{R}_+$,

$$\begin{aligned} & \mathbb{E}\left[f(X_0^t(\emptyset)) \mathbb{1}_{X_0^t(\emptyset) > 0} \mathbb{1}_{\mathbf{e}_{\emptyset}^t \leq s} g_1(\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(1)) g_2(\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(2))\right] \\ &= \int_0^\infty f(x) \varrho(t) e^{-\lambda(t)x} \int_0^s \varrho(t-r)(x+r) \exp\left(-\int_0^r \varrho(t-v)(x+v) dv\right) \\ & \quad \times \int_0^{x+r} g_1(y) g_2(x+r-y) \frac{dy}{x+r} dr dx, \end{aligned} \quad (4.2)$$

recalling that \mathbf{e}_{\emptyset}^t denotes the lifetime of \emptyset in the red tree. First note that $X_0^t(\emptyset) + \mathbf{e}_{\emptyset}^t = \mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(1) + \mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(2)$ and, therefore, the left-hand side of (4.2) is equal to

$$\mathbb{E}\left[f(\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(1) + \mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(2) - \mathbf{e}_{\emptyset}^t) \mathbb{1}_{\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(1) + \mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(2) > \mathbf{e}_{\emptyset}^t} \mathbb{1}_{\mathbf{e}_{\emptyset}^t \leq s} g_1(\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(1)) g_2(\mathfrak{X}_{\mathbf{e}_{\emptyset}^t}^t(2))\right]. \quad (4.3)$$

Then, we distinguish on the first particle u starting from the root, that splits its mass into two *positive* masses between its children. Denoting $\emptyset = u_0 < u_1 < \dots < u_k = u$ the lineage of u and, for $v \in \mathbb{T} \setminus \{\emptyset\}$, \bar{v} the single brother of v , we get that (4.3) is equal to

$$\begin{aligned} & \sum_{k \geq 0} \sum_{|u|=k} \mathbb{E}\left[\mathbb{1}_{\forall i \in [1, k], X_{b_{u_i}}^t(\bar{u}_i) = 0} \mathbb{1}_{X_{d_u}^t(u_1), X_{d_u}^t(u_2) > 0} f(X_{d_u}^t(u_1) + X_{d_u}^t(u_2) - d_u) \right. \\ & \quad \left. \times \mathbb{1}_{X_{d_u}^t(u_1) + X_{d_u}^t(u_2) > d_u} \mathbb{1}_{d_u \leq s} g_1(X_{d_u}^t(u_1)) g_2(X_{d_u}^t(u_2))\right]. \end{aligned} \quad (4.4)$$

Using the branching property, on the event $d_u \leq s$ and given $b_{u_1}, \dots, b_{u_k}, d_u$, the Yule trees rooted in $\bar{u}_1, \dots, \bar{u}_k, u_1, u_2$ and the labels on their leaves are independent and, therefore, $X_{b_{u_1}}^t(\bar{u}_1), \dots, X_{b_{u_k}}^t(\bar{u}_k), X_{d_u}^t(u_1), X_{d_u}^t(u_2)$ are independent with laws $\mu_{b_{u_1}}, \dots, \mu_{b_{u_k}}, \mu_{d_u}, \mu_{d_u}$

respectively. Thus, taking the conditional expectation given $b_{u_1}, \dots, b_{u_k}, d_u$, (4.4) is equal to

$$\begin{aligned} & \sum_{k \geq 0} \sum_{|u|=k} \mathbb{E} \left[\left(\prod_{i=1}^k p(t - b_{u_i}) \right) \mathbb{1}_{d_u \leq s} (1 - p(t - d_u))^2 \lambda(t - d_u)^2 \right. \\ & \quad \left. \times \int_0^\infty \int_0^\infty e^{-\lambda(t-d_u)(y+z)} f(y+z-d_u) \mathbb{1}_{y+z > d_u} g_1(y) g_2(z) dz dy \right]. \end{aligned} \quad (4.5)$$

On the other hand, given d_u , the vector $(b_{u_i})_{1 \leq i \leq k}$ has the same law as the ordered vector $(U_i)_{1 \leq i \leq k}$, where the U_i 's are i.i.d. uniform r.v. on $[0, d_u]$, so, for any measurable positive function h and any $|u| = k$, we have

$$\begin{aligned} \mathbb{E} \left[\left(\prod_{i=1}^k p(t - b_{u_i}) \right) \mathbb{1}_{d_u \leq s} h(d_u) \right] &= \mathbb{E} \left[\left(\int_0^{d_u} p(t - v) \frac{dv}{d_u} \right)^k \mathbb{1}_{d_u \leq s} h(d_u) \right] \\ &= \int_0^s \left(\int_0^r p(t - v) \frac{dv}{r} \right)^k h(r) \frac{r^k e^{-r}}{k!} dr, \end{aligned}$$

using that d_u follows the $\Gamma(k+1, 1)$ distribution. Noting that $\#\{|u| = k\} = 2^k$ and summing over k , we get that (4.5) is equal to

$$\begin{aligned} & \int_{r=0}^s \exp\left(\int_0^r 2p(t-v)dv\right) (1 - p(t-r))^2 \lambda(t-r)^2 \\ & \quad \times \int_{y=0}^\infty \int_{z=0}^\infty e^{-\lambda(t-r)(y+z)} f(y+z-r) \mathbb{1}_{y+z > r} g_1(y) g_2(z) dz dy e^{-r} dr. \end{aligned} \quad (4.6)$$

Changing z into $x = y + z - r$, using Fubini's theorem and recalling that $\varrho(t) = (1 - p(t))\lambda(t)$, (4.6) is equal to

$$\begin{aligned} & \int_{x=0}^\infty f(x) \int_{r=0}^s \exp\left(\int_0^r (2p(t-v) - 1)dv\right) \varrho(t-r)^2 e^{-\lambda(t-r)(x+r)} \\ & \quad \times \int_{y=0}^{x+r} g_1(y) g_2(x+r-y) dy dr dx. \end{aligned}$$

Now, in order to prove (4.2), it is sufficient to prove that

$$\begin{aligned} & \varrho(t-r) e^{-\lambda(t-r)(x+r)} \exp\left(-\int_0^r (1 - 2p(t-v))dv\right) \\ & = \varrho(t) e^{-\lambda(t)x} \exp\left(-\int_0^r \varrho(t-v)(x+v)dv\right). \end{aligned} \quad (4.7)$$

For this, note that $(\log \varrho)' = 2p - 1 - \lambda$ by (3.3) and therefore

$$\frac{\varrho(t-r)}{\varrho(t)} = \exp(\log \varrho(t-r) - \log \varrho(t)) = \exp\left(-\int_0^r (2p(t-v) - 1 - \lambda(t-v))dv\right).$$

On the other hand, using that $\frac{d}{dv}(\lambda(t-v)(x+v)) = \varrho(t-v)(x+v) + \lambda(t-v)$, we get

$$\frac{e^{-\lambda(t-r)(x+r)}}{e^{-\lambda(t)x}} = \exp\left(-\int_0^r (\varrho(t-v)(x+v) + \lambda(t-v))dv\right).$$

Combining the two previous equations, it proves (4.7) and concludes the proof. \square

Remark 4.4. Using [16], one can prove a many-to-one lemma for the red tree. More precisely, the following result holds: For all measurable bounded function f , we have

$$\mathbb{E} \left(\sum_{u \in \mathcal{N}_t} f(Y_s^{t,x}(u), s \leq t) \right) = \mathbb{E} \left(e^{\int_0^t \varrho(t-s) Z_s^{t,x} ds} f(Z_s^{t,x}, s \leq t) \right),$$

where $Z^{t,x}$ is a time-inhomogeneous Markov process with $Z_0^{t,x} = x$, with a positive drift 1, that jumps at rate $2\varrho(t-s)Z_s$ to position $UZ_s^{t,x}$ with U an independent uniform random variable.

4.2 Scaling limit of the red tree

In this subsection, we consider the critical case by taking as initial law μ_0 given by (3.1) with $(p(0), \lambda(0)) \in \mathcal{C}$. In that case, it follows from Proposition 3.7 that $\varrho(s) = \lambda(s)(1-p(s)) \sim 2/s^2$ as $s \rightarrow \infty$. Using this asymptotic result, one can take the limit $t \rightarrow \infty$ in the description of the red tree of Proposition 4.1, as shown in the next result.

Before stating and proving this result, we need some additional formalism. The set of binary continuous trees is $(\mathbb{R}_+)^{\mathbb{T}}$, where a tree is represented by the lengths of its branches, and the set of masses is $\mathcal{C}(\mathbb{R}_+)^{\mathbb{T}}$, where to each particle of the tree is associated the continuous function of its mass during its lifetime. Then, the process $(\mathfrak{X}_s^t(v), v \in \mathfrak{N}_s^t, s \leq t)$ is encoded as the following random variable in $(\mathbb{R}_+)^{\mathbb{T}} \times \mathcal{C}(\mathbb{R}_+)^{\mathbb{T}}$:

$$\left((\mathfrak{e}_u^t)_{u \in \mathbb{T}}, (\mathfrak{X}_{(\mathfrak{b}_u^t + s) \wedge \mathfrak{d}_u^t}(u), s \geq 0)_{u \in \mathbb{T}} \right),$$

where we recall that for any $u \in \mathbb{T}$, \mathfrak{e}_u^t , \mathfrak{b}_u^t and \mathfrak{d}_u^t denotes respectively the lifetime, birth time and death time of u . The set $\mathcal{C}(\mathbb{R}_+)$ is endowed with the topology of uniform convergence on every compact set and $(\mathbb{R}_+)^{\mathbb{T}} \times \mathcal{C}(\mathbb{R}_+)^{\mathbb{T}}$ with the product topology.

Now, for $T \geq 0$, we define the restriction of the tree to the time interval $[0, T]$ by

$$\left(((\mathfrak{d}_u^t \wedge T) - (\mathfrak{b}_u^t \wedge T))_{u \in \mathbb{T}}, (\mathfrak{X}_{(\mathfrak{b}_u^t + s) \wedge \mathfrak{d}_u^t \wedge T}(u) \mathbb{1}_{\mathfrak{b}_u^t \leq T}, s \geq 0)_{u \in \mathbb{T}} \right),$$

which is still an element of $(\mathbb{R}_+)^{\mathbb{T}} \times \mathcal{C}(\mathbb{R}_+)^{\mathbb{T}}$.

Theorem 4.5. *Let $\mathfrak{X} = (\mathfrak{X}_s(u), u \in \mathfrak{N}(s), s \in [0, 1])$ be a time-inhomogeneous branching Markov process such that under \mathbb{P}_x :*

1. *The process starts at time 0 with a unique particle of mass x : $\mathfrak{X}_0(\emptyset) = x$.*
2. *The mass associated to each particle grows linearly at speed 1.*
3. *A particle of mass m at time s splits at rate $2m/(1-s)^2$ into two children, the mass m being split uniformly between the two children.*
4. *Particles behave independently after their splitting time.*

Assume that the initial law μ_0 is given by (3.1) with $(p(0), \lambda(0)) \in \mathcal{C}$. Let $(x_t)_{t \geq 0}$ be a family of positive real numbers such that $x_t/t \rightarrow x \in \mathbb{R}_+$ as $t \rightarrow \infty$. Then, for any $\varepsilon \in (0, 1)$, the restriction of $(\mathfrak{X}_{st}^t(v)/t, v \in \mathfrak{N}_{st}^t, s \in [0, 1])$ to $[0, 1 - \varepsilon]$ under \mathbb{P}_{x_t} (defined below Proposition 4.1) converges in law toward the restriction of \mathfrak{X} to $[0, 1 - \varepsilon]$ under \mathbb{P}_x .

Proof. The first step is to prove that we can deal only with a finite number of particles. Let $m_t := \max_{u \in \mathfrak{N}_{(1-\varepsilon)t}} |u|$ be the maximal generation among the red particles alive at time $(1-\varepsilon)t$, we are going to show that $(M_t)_{t \geq 0}$ is tight in \mathbb{R}_+ . By Proposition 3.7, for any t large enough and $s \in [0, (1-\varepsilon)t]$, we have $\varrho(t-s) \leq 4/(t-s)^2 \leq 4/(\varepsilon t)^2$. On the other hand, note that under \mathbb{P}_{x_t} , the mass of any particle at any time is bounded by $x_t + t$, and therefore by $(x+2)t$ for t large enough. Therefore, by Proposition 4.1, the branching rate of any particle on the time interval $[0, (1-\varepsilon)t]$ can be upper bounded by $4(x+2)/(\varepsilon^2 t)$. We conclude that m_t is stochastically dominated by the maximal generation among the particles alive at time $(1-\varepsilon)$ in a Yule tree with parameter $4(x+2)/\varepsilon^2$ and, thus, $(m_t)_{t \geq 0}$ is tight.

Hence, it is now sufficient to prove, for any $m \in \mathbb{N}$, the following convergence in law

$$\begin{aligned} & \left((\mathbf{e}_u^t/t)_{u \in \mathbb{T}, |u| \leq m}, (\mathfrak{X}_{(b_u^t + st) \wedge \partial_u^t}(u)/t, s \geq 0)_{u \in \mathbb{T}, |u| \leq m} \right) \text{ under } \mathbb{P}_{x_t} \\ & \xrightarrow[t \rightarrow \infty]{(d)} \left((\mathbf{e}_u)_{u \in \mathbb{T}, |u| \leq m}, (\mathfrak{X}_{(b_u + s) \wedge \partial_u}(u), s \geq 0)_{u \in \mathbb{T}, |u| \leq m} \right) \text{ under } \mathbb{P}_x. \end{aligned}$$

By recurrence on m and using the Markov property of \mathfrak{X}^t (Point 4 of Proposition 4.1), it is sufficient to prove that,

$$\left(\mathbf{e}_\emptyset^t/t, (\mathfrak{X}_{(st) \wedge \partial_\emptyset^t}(\emptyset)/t, s \geq 0) \right) \text{ under } \mathbb{P}_{x_t} \xrightarrow[t \rightarrow \infty]{(d)} \left(\mathbf{e}_\emptyset, (\mathfrak{X}_{s \wedge \partial_\emptyset}(\emptyset), s \geq 0) \right) \text{ under } \mathbb{P}_x \quad (4.8)$$

But, $\mathfrak{X}_{(st) \wedge \partial_\emptyset^t}(\emptyset) = x_t + ((st) \wedge \mathbf{e}_\emptyset^t)$ and $\mathfrak{X}_{s \wedge \partial_\emptyset}(\emptyset) = x + (s \wedge \mathbf{e}_\emptyset)$, so (4.8) reduces to the proof of the convergence in distribution of \mathbf{e}_\emptyset^t/t under \mathbb{P}_{x_t} toward \mathbf{e}_\emptyset under \mathbb{P}_x . By Point 3 of Proposition 4.1, for any $r \in (0, 1)$, as $t \rightarrow \infty$,

$$\mathbb{P}_{x_t}(\mathbf{e}_\emptyset^t/t > r) = \exp\left(-\int_0^{rt} \varrho(t-u)(x_t+u)du\right) \xrightarrow[t \rightarrow \infty]{} \exp\left(-\int_0^r \frac{2(x+v)dv}{(1-v)^2}\right), \quad (4.9)$$

using that $\varrho(s) \sim 2/s^2$ as $s \rightarrow \infty$ by Proposition 3.7 and $x_t/t \rightarrow x$ as $t \rightarrow \infty$. Finally, note that, by definition of the limit \mathfrak{X} , the right-hand side of (4.9) is equal to $\mathbb{P}_x(\mathbf{e}_\emptyset > r)$, which completes the proof. \square

Remark 4.6. We believe that the limit \mathfrak{X} should be universal. Defining similarly the red tree for the discrete DR model, we expect the same scaling limit to hold in the critical case. This is supported by computer calculations of the distribution of first branching time in the discrete red tree, which can be compared with the distribution of \mathbf{e}_\emptyset under \mathbb{P}_0 given by

$$\forall r \in [0, 1], \quad \mathbb{P}_0(\mathbf{e}_\emptyset \leq r) = 1 - \exp\left(-\int_0^r \frac{2v dv}{(1-v)^2}\right) = 1 - \frac{1}{(1-r)^2} \exp\left(-\frac{2r}{1-r}\right),$$

see Figure 9. Similarly, computer calculations show that the mass in the red tree of height n is split uniformly among children as $n \rightarrow \infty$.

5 Asymptotic behavior of the number and mass of red leaves

The red tree, studied in the previous section, is a tool for understanding how mass can be brought from the leaves to the root. We focus here on the critical case, with initial law μ_0

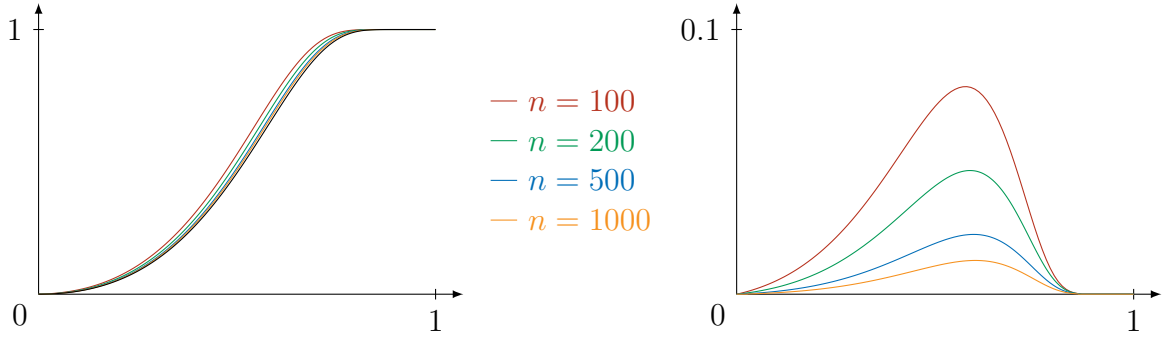


Figure 9: For each $n \in \{100, 200, 500, 1000\}$, the repartition function F_n of the first branching time of the discrete red tree of height n with critical initial law $\frac{4}{5}\delta_0 + \frac{1}{5}\delta_2$ has been computed. On the left, the F_n 's are drawn in color and the repartition function F of ε_\emptyset under \mathbb{P}_0 is drawn in black. On the right, the $F_n - F$'s are drawn.

given by (3.1) with $(p(0), \lambda(0)) \in \mathcal{C}$, and answer the following questions: typically, when mass is brought to the root, how many leaves have contributed to this amount of mass? How much mass do these leaves carry?

In order to answer these questions, let N_t be the number of red leaves and M_t the total mass of the red leaves:

$$N_t := \#\mathfrak{N}_t^t \quad \text{and} \quad M_t := \sum_{u \in \mathfrak{N}_t^t} \mathfrak{X}_t^t(u).$$

The joint asymptotic behavior of N_t and M_t is given by the following result.

Theorem 5.1. *Let $(p(0), \lambda(0)) \in \mathcal{C}$. There exist $\gamma_1, \gamma_2 > 0$ such that, for any family $(x_t)_{t \geq 0}$ of positive real numbers such that $x_t/t \rightarrow x \in \mathbb{R}_+$ as $t \rightarrow \infty$, $(t^{-2}N_t, t^{-2}M_t)$ under \mathbb{P}_{x_t} converges in distribution to $(\gamma_1\eta, \gamma_2\eta)$, where $\eta := \frac{3}{2} \int_0^1 r^2(s)ds$, with $r(\cdot)$ denoting a 4-dimensional Bessel bridge from 0 to $2\sqrt{x}$.*

Constants γ_1 and γ_2 are defined implicitly in Lemma 5.4 and depend on the initial condition $(p(0), \lambda(0))$. This section is devoted to the proof of Theorem 5.1, via the analytic study of the Laplace transform of (N_t, M_t) . In the sequel, $(p(0), \lambda(0)) \in \mathcal{C}$ is fixed and thus also the auxiliary function $\varrho(s) = (1 - p(s))\lambda(s)$.

A further question that we do not answer here is the following: what is the typical mass of a red leaf? Thereupon, we state the following conjecture.

Conjecture 5.2. *For any family $(x_t)_{t \geq 0}$ of positive real numbers such that $x_t/t \rightarrow x \in \mathbb{R}_+$ as $t \rightarrow \infty$, the random measure*

$$\frac{1}{N_t} \sum_{u \in \mathfrak{N}_t^t} \delta_{\mathfrak{X}_t^t(u)} \quad \text{under } \mathbb{P}_{x_t}$$

converges in law for the topology of vague convergence and the limit does not depend on x .

5.1 The Laplace transform of the mass and number of red leaves

For $\varepsilon_1, \varepsilon_2 > 0$, we introduce the following function

$$\phi(t, x) = \phi_{\varepsilon_1, \varepsilon_2}(t, x) := \mathbb{E}_x \left[e^{-\varepsilon_1 N_t - \varepsilon_2 M_t} \right], \quad t, x \geq 0. \quad (5.1)$$

As a function of $(\varepsilon_1, \varepsilon_2)$, $\phi_{\varepsilon_1, \varepsilon_2}(t, x)$ is the Laplace transform of (N_t, M_t) under \mathbb{P}_x . However, in this subsection, we fix $(\varepsilon_1, \varepsilon_2)$ and study analytically $\phi_{\varepsilon_1, \varepsilon_2}$ as a function of (t, x) : therefore, we forget the dependence in $\varepsilon_1, \varepsilon_2$ and write simply ϕ . The aim of this subsection is to prove the following result.

Lemma 5.3. *Let $\varepsilon_1, \varepsilon_2 > 0$. Then, for any $t, x \geq 0$, we have*

$$\phi(t, x) = e^{-(\Theta(t)+x\Theta'(t))}, \quad (5.2)$$

where Θ is the unique solution on \mathbb{R}_+ to the following differential equation

$$\begin{cases} \Theta'' = \varrho(1 - e^{-\Theta}), \\ \Theta(0) = \varepsilon_1, \\ \Theta'(0) = \varepsilon_2. \end{cases} \quad (5.3)$$

Moreover, for any $t \geq 0$, we have $\Theta(t) > 0$ and $\Theta'(t) > 0$.

Proof. Recall that $\mathbf{e}_{\varnothing}^t$ is the first branching time in the red tree. Distinguishing on the cases $\mathbf{e}_{\varnothing}^t = t$ and $\mathbf{e}_{\varnothing}^t < t$ and using Proposition 4.1, we have

$$\phi(t, x) = \mathbb{P}(\mathbf{e}_{\varnothing}^t = t)\phi(0, x+t) + \int_0^t \mathbb{P}(\mathbf{e}_{\varnothing}^t \in ds) \int_0^{x+s} \phi(t-s, y)\phi(t-s, x+s-y) \frac{dy}{x+s}$$

and, moreover, the law of $\mathbf{e}_{\varnothing}^t$ is given by

$$\mathbb{P}(\mathbf{e}_{\varnothing}^t \geq s) = \exp\left(-\int_0^s \varrho(t-r)(x+r)dr\right) = \exp\left(-\int_{t-s}^t \varrho(r)(x+t-r)dr\right),$$

for any $s \in [0, t]$. Therefore, replacing s by $t-s$ in the first integral, we get the following integral equation for ϕ :

$$\begin{aligned} \phi(t, x) &= \exp\left(-\int_0^t \varrho(r)(x+t-r)dr\right)\phi(0, x+t) \\ &\quad + \int_0^t \varrho(s) \exp\left(-\int_s^t \varrho(r)(x+t-r)dr\right) \int_0^{x+t-s} \phi(s, y)\phi(s, x+t-s-y)dyds. \end{aligned} \quad (5.4)$$

Let first check that the function ϕ suggested in the lemma is a solution of Equation (5.4). By Cauchy-Lipschitz theorem, there exists a unique maximal \mathcal{C}^2 solution Θ to Equation (5.3) and (Θ, Θ') has to stay in $(0, \infty) \times (0, \infty)$ at positive times. Therefore, $\varrho(1 - e^{-\Theta})$ remains bounded and it follows that Θ is well-defined on \mathbb{R}_+ . Hence, for $t, x \geq 0$, we can define $\phi(t, x) = e^{-(\Theta(t)+x\Theta'(t))}$ and the second term in the right-hand side of (5.4) is equal to

$$\begin{aligned} &\int_0^t \varrho(s) \exp\left(-\int_s^t \varrho(r)(x+t-r)dr\right)(x+t-s)e^{-(2\Theta(s)+(x+t-s)\Theta'(s))}ds \\ &= \int_0^t \frac{d}{ds} \left[\exp\left(-\int_s^t \varrho(r)(x+t-r)dr\right) e^{-(\Theta(s)+(x+t-s)\Theta'(s))} \right] ds \\ &= \phi(t, x) - \exp\left(-\int_0^t \varrho(r)(x+t-r)dr\right)\phi(0, x+t), \end{aligned}$$

using (5.3) in order to compute the derivative. This proves that $(x, t) \mapsto e^{-(\Theta(t)+x\Theta'(t))}$ is solution of (5.4).

We now prove uniqueness for Equation (5.4) among the set of measurable locally bounded functions from $(\mathbb{R}_+)^2 \rightarrow \mathbb{R}$. Let $\phi_1, \phi_2: (\mathbb{R}_+)^2 \rightarrow \mathbb{R}$ be measurable locally bounded functions satisfying (5.4). We fix some $x_0 \geq 0$ and prove that $\phi_1 = \phi_2$ on the triangle $T := \{(t, x) \in (\mathbb{R}_+)^2 : x + t \leq x_0\}$. It follows from (5.4) that

$$|\phi_1(t, x) - \phi_2(t, x)| \leq C \int_0^t \int_0^{x+t-s} |\phi_1(s, y) - \phi_2(s, y)| dy ds, \quad (5.5)$$

with $C := (\sup_{\mathbb{R}_+} \varrho)(\sup_T |\phi_1| + \sup_T |\phi_2|)$, using that ϱ is bounded and ϕ_1, ϕ_2 are locally bounded. Now, we introduce $T_t := T \cap ([0, t] \times \mathbb{R}_+)$: note that, for any $(x, t) \in T$, the variable (y, s) in the right-hand side of (5.5) remains in T_t . Let $t_0 := 1/(2Cx_0)$, we prove by induction on $k \in \mathbb{N}$ that $\phi_1 - \phi_2 = 0$ on T_{kt_0} . For $k = 0$, it's clearly true. Now, assume that it holds for some $k \in \mathbb{N}$. Then, for any $(t, x) \in T_{(k+1)t_0}$, it follows from (5.5) that

$$|\phi_1(t, x) - \phi_2(t, x)| \leq C(t - kt_0)x_0 \sup_{T_{(k+1)t_0}} |\phi_1 - \phi_2|,$$

and, taking the supremum over $(t, x) \in T_{(k+1)t_0}$, we get

$$\sup_{T_{(k+1)t_0}} |\phi_1 - \phi_2| \leq Ct_0x_0 \sup_{T_{(k+1)t_0}} |\phi_1 - \phi_2| \leq \frac{1}{2} \sup_{T_{(k+1)t_0}} |\phi_1 - \phi_2|,$$

which proves that $\phi_1 - \phi_2 = 0$ on $T_{(k+1)t_0}$. Therefore, we proved that $\phi_1 = \phi_2$ on T and, since it holds for any $x_0 \geq 0$, we get the uniqueness. The function ϕ defined by (5.1) takes value in $[0, 1]$ and the function $(t, x) \mapsto e^{-(\Theta(t)+x\Theta'(t))}$ is continuous on $(\mathbb{R}_+)^2$ and therefore locally bounded. Moreover, they both satisfy Equation (5.4) so they coincide. \square

5.2 An auxiliary differential equation

Recall that our aim is to prove the convergence in law of $(N_t/t^2, M_t/t^2)$ and, for this, the convergence of the Laplace transform $\phi_{\theta_1/t^2, \theta_2/t^2}(t, x_t)$. Therefore, we will need to study the function Θ solution of (5.3) for small initial conditions $\varepsilon_1, \varepsilon_2$ and, when Θ is small, it behaves like the function a , solution to the following linearized equation:

$$a'' = \varrho a. \quad (5.6)$$

In this subsection, we prove a preliminary result concerning the asymptotic behavior of the function a solution of (5.6).

Lemma 5.4. *There exists positive constants $\gamma_1, \gamma_2 > 0$ such that, for any $\theta_1, \theta_2 \in \mathbb{R}$, if a is the unique solution on \mathbb{R}_+ to the differential equation $a'' = \varrho a$ with initial conditions $a(0) = \theta_1$ and $a'(0) = \theta_2$, then we have $a(t) \sim (\gamma_1\theta_1 + \gamma_2\theta_2)t^2$ and $a'(t) \sim 2(\gamma_1\theta_1 + \gamma_2\theta_2)t$ as $t \rightarrow \infty$.*

Proof. First note that, by linearity, it is sufficient to deal with the cases $(\theta_1, \theta_2) = (1, 0)$ and $(\theta_1, \theta_2) = (0, 1)$. From now on, assume that we are in one of this cases. In particular, this implies that $a(t) > 0$ and $a'(t) > 0$ for any $t > 0$. The main idea of the proof is to show that, for large t , we can replace $\varrho(t)$ by $2/t^2$ in the equation: indeed, it follows from Proposition 3.7 that

$$\varrho(t) = \frac{2}{t^2} - \frac{16 \log t}{3t^3} + o\left(\frac{\log t}{t^3}\right),$$

so we can choose some t_0 large enough such that

$$\forall t \geq t_0, \quad 0 < \frac{2}{t^2} - \frac{1}{t^{5/2}} \leq \varrho(t) \leq \frac{2}{t^2}.$$

We first prove an upper bound on $a(t)$. Let a_0 be the solution on $[t_0, \infty)$ of the following ODE:

$$a_0''(t) = \frac{2}{t^2}a_0(t), \quad a_0(t_0) = a(t_0), \quad a_0'(t_0) = a'(t_0).$$

Then, $a_0(t) = c_1 t^2 + c_2 t^{-1}$ for any $t \geq t_0$, for some $c_1, c_2 \in \mathbb{R}$ and, using the initial conditions, we get

$$a_0(t) = \frac{a(t_0) + t_0 a'(t_0)}{3t_0^2} t^2 + \frac{2t_0 a(t_0) - t_0^2 a'(t_0)}{3t}. \quad (5.7)$$

Then, setting

$$C = C(t_0) := \frac{a(t_0) + t_0 a'(t_0)}{t_0^2} > 0,$$

we get that $a_0(t) \leq Ct^2$ for any $t \geq t_0$. On the other hand, for any $t \geq t_0$, we have $a(t) \leq a_0(t)$ and, therefore, we proved that

$$\forall t \geq t_0, \quad a(t) \leq Ct^2, \quad (5.8)$$

which is an upper bound with a rough constant.

Now, we prove that $a(t)$ is of order t^2 and $a'(t)$ of order t as $t \rightarrow \infty$. For this, we consider some $t_1 > t_0$ and let a_1 be the solution on $[t_1, \infty)$ of the following ODE:

$$a_1''(t) = \frac{2}{t^2}a_1(t), \quad a_1(t_1) = a(t_1), \quad a_1'(t_1) = a'(t_1).$$

Similarly to a_0 , a_1 has an exact expression of the form (5.7) with t_0 replaced by t_1 . As before, for any $t \geq t_1$, we have $a(t) \leq a_1(t)$ and

$$\begin{aligned} (a_1 - a)''(t) &= \frac{2}{t^2}a_1(t) - \varrho(t)a(t) = \frac{2}{t^2}(a_1(t) - a(t)) + \left(\frac{2}{t^2} - \varrho(t)\right)a(t) \\ &\leq \frac{2}{t^2}(a_1(t) - a(t)) + \frac{1}{t^{5/2}}Ct^2, \end{aligned} \quad (5.9)$$

using that $a(t) \leq Ct^2$ for any $t \geq t_0$. Then, we introduce the auxiliary function $y(t) := a_1(t) - a(t) + \frac{4C}{5}t^{3/2}$ and it follows from (5.9) that, for any $t \geq t_1$,

$$y''(t) \leq \frac{2}{t^2}y(t).$$

Therefore, we get $y \leq \bar{y}$, where \bar{y} is the solution on $[t_1, \infty)$ of the following ODE:

$$\bar{y}''(t) \leq \frac{2}{t^2}\bar{y}(t), \quad \bar{y}(t_1) = y(t_1) = \frac{4C}{5}t_1^{3/2}, \quad \bar{y}'(t_1) = y'(t_1) = \frac{6C}{5}t_1^{1/2}.$$

One compute as before that, for any $t \geq t_1$,

$$\bar{y}(t) = \frac{2C}{\sqrt{t_1}}t^2 + \frac{2Ct_1^{5/2}}{15t}.$$

Therefore, it follows that

$$a_1(t) - a(t) = y(t) - \frac{4C}{5}t^{3/2} \leq \frac{2C}{\sqrt{t_1}}t^2 + \frac{2Ct_1^{5/2}}{15t}. \quad (5.10)$$

and, using the explicit value of a_1 , we finally proved that, for any $t_1 \geq t_0$, as $t \rightarrow \infty$,

$$\left(\frac{a(t_1) + t_1 a'(t_1)}{3t_1^2} - \frac{2C}{\sqrt{t_1}} \right) t^2 + o(t^2) \leq a(t) \leq \frac{a(t_1) + t_1 a'(t_1)}{3t_1^2} t^2 + o(t^2). \quad (5.11)$$

We apply this result with $t_1 = t_0$ and it gives, as $t \rightarrow \infty$,

$$\left(\frac{C}{3} - \frac{2C}{\sqrt{t_0}} \right) t^2 + o(t^2) \leq a(t) \leq \frac{C}{3} + o(t^2), \quad (5.12)$$

where $\frac{C}{3} - \frac{2C}{\sqrt{t_0}} > 0$ by choosing $t_0 > 36$ at the beginning. Thus, we proved that $a(t)$ is of order t^2 . Combining this with $a'(t) = \int_0^t \varrho(s)a(s)ds$ and $\varrho(t) \sim 2/t^2$, we get

$$2 \left(\frac{C}{3} - \frac{2C}{\sqrt{t_0}} \right) t + o(t) \leq a'(t) \leq \frac{2C}{3} t + o(t), \quad (5.13)$$

so $a'(t)$ is of order t .

Finally, we prove the result. Using (5.12) and (5.13), there exist an increasing sequence $(s_n)_{n \in \mathbb{N}}$ with $s_n \rightarrow \infty$ and two constants $c, c' > 0$ such that

$$\frac{a(s_n)}{s_n^2} \xrightarrow{n \rightarrow \infty} c \quad \text{and} \quad \frac{a'(s_n)}{s_n} \xrightarrow{n \rightarrow \infty} c'. \quad (5.14)$$

Then, for any $n \in \mathbb{N}$, applying (5.11) with $t_1 = s_n$, we get

$$\left(\frac{a(s_n) + s_n a'(s_n)}{3s_n^2} - \frac{2C}{\sqrt{s_n}} \right) t^2 + o(t^2) \leq a(t) \leq \frac{a(s_n) + s_n a'(s_n)}{3s_n^2} t^2 + o(t^2)$$

and, using (5.14), it follows that $a(t)/t^2 \rightarrow (c+c')/3$. In particular, we have $c = (c+c')/3$. Using again that $a'(t) = \int_0^t \varrho(s)a(s)ds$ and $\varrho(t) \sim 2/t^2$, it follows that $a'(t)/t \rightarrow 2c$. Therefore, we proved the result for $(\theta_1, \theta_2) = (1, 0)$ and $(\theta_1, \theta_2) = (0, 1)$ and the result follows. \square

Remark 5.5. Let $\psi(t, x) := \mathbb{E}_x[N_t]$ (or similarly $\mathbb{E}_x[M_t]$), then one can check in the same way as in Lemma 5.3 that $\psi(t, x) = a(t) + a'(t)x$ with a a solution of Equation (5.6). Therefore, Lemma 5.4 implies that $\mathbb{E}_{xt}[N_t] \sim \gamma_1(1 + 2x)t^2$ and $\mathbb{E}_{xt}[M_t] \sim \gamma_2(1 + 2x)t^2$ as $t \rightarrow \infty$.

5.3 Asymptotics of the Laplace transform

In this subsection, we prove Theorem 5.1, using results of the two previous subsections. Recall that $\phi_{\varepsilon_1, \varepsilon_2}(t, x) = \mathbb{E}_x[e^{-\varepsilon_1 N_t - \varepsilon_2 M_t}] = e^{-(\Theta(t) + x\Theta'(t))}$, where Θ is the solution of (5.3). Our aim is to prove the convergence of $\mathbb{E}_{x_t}[e^{-\theta_1 t^{-2} N_t - \theta_2 t^{-2} M_t}]$ as $t \rightarrow \infty$. Hence, we consider Θ_ε solution of

$$\begin{cases} \Theta_\varepsilon'' = \varrho(1 - e^{-\Theta_\varepsilon}), \\ \Theta_\varepsilon(0) = \theta_1 \varepsilon, \\ \Theta_\varepsilon'(0) = \theta_2 \varepsilon, \end{cases} \quad (5.15)$$

for any fixed $\theta_1, \theta_2 > 0$, and study the asymptotics of $\Theta_\varepsilon(1/\sqrt{\varepsilon})$ and $\Theta_\varepsilon'(1/\sqrt{\varepsilon})$ as $\varepsilon \rightarrow 0$. The first step is to control the behavior of Θ_ε and Θ_ε' at large but fixed time, as $\varepsilon \rightarrow 0$. This is done in the following result.

Lemma 5.6. *Let $\theta_1, \theta_2 \in \mathbb{R}$ and $\delta > 0$. For large enough t , there exists $\varepsilon_0 = \varepsilon_0(t, \delta) > 0$ such that for any $0 < \varepsilon < \varepsilon_0$,*

$$\left| \Theta_\varepsilon(t) - \varepsilon c t^2 \right| \leq \delta \varepsilon t^2 \quad \text{and} \quad \left| \Theta_\varepsilon'(t) - 2\varepsilon c t \right| \leq \delta \varepsilon t,$$

where $c = \gamma_1 \theta_1 + \gamma_2 \theta_2$ and γ_1, γ_2 are defined in Lemma 5.4.

Proof. Let a be solution of $a'' = \varrho a$ with initial conditions $a(0) = \theta_1$ and $a'(0) = \theta_2$. The strategy of the proof is to show that $\Theta_\varepsilon(t)$ and $\Theta_\varepsilon'(t)$ are close to $\varepsilon a(t)$ and $\varepsilon a'(t)$ for small ε and then to apply Lemma 5.4. From now on, we fix some $t \geq 0$ and the constants C_i can depend on t , as well as on $\theta_1, \theta_2, \delta$. First note that $\Theta_\varepsilon'' = \varrho(1 - e^{-\Theta_\varepsilon}) \leq \varrho \Theta_\varepsilon$ and, therefore, $\Theta_\varepsilon \leq \varepsilon a$. Then, it follows from Lemma 5.4 that there exists $C > 0$ such that, for any $s \geq 0$,

$$\Theta_\varepsilon(s) \leq \varepsilon a(s) \leq C\varepsilon(s+1)^2. \quad (5.16)$$

In order to prove that $\Theta_\varepsilon(t)$ is close to $\varepsilon a(t)$, we introduce $y := \varepsilon a - \Theta_\varepsilon$, which satisfies $y(0) = 0$, $y'(0) = 0$ and, for any $s \geq 0$,

$$\begin{aligned} y''(s) &= \varrho(s) (\varepsilon a(s) - 1 + e^{-\Theta_\varepsilon(s)}) = \varrho(s) (y(s) + \Theta_\varepsilon(s) - 1 + e^{-\Theta_\varepsilon(s)}) \\ &\leq \varrho(s) \left(y(s) + \frac{\Theta_\varepsilon(s)^2}{2} \right) \leq \varrho(s) \left(y(s) + \frac{C}{2} (s+1)^4 \varepsilon^2 \right), \end{aligned}$$

using (5.16). Since the function ϱ is bounded, we deduce that, for any $s \in [0, t]$, $y''(s) \leq C_1 y(s) + C_2 \varepsilon^2$, where C_2 depends on t . Thus, we get, for any $s \in [0, t]$,

$$y(s) \leq C_2 \varepsilon^2 \left(\cosh\left(\sqrt{C_1} s\right) - 1 \right) \leq C_3 \varepsilon^2.$$

It follows that, for any $s \in [0, t]$,

$$|\Theta_\varepsilon(s) - \varepsilon a(s)| \leq C_3 \varepsilon^2. \quad (5.17)$$

Moreover, since $\Theta_\varepsilon'(t) = \int_0^t \varrho(s) (1 - e^{-\Theta_\varepsilon(s)}) ds$ and $a'(t) = \int_0^t \varrho(s) a(s) ds$, we get

$$\begin{aligned} |\Theta_\varepsilon'(t) - \varepsilon a'(t)| &\leq \int_0^t \varrho(s) |1 - e^{-\Theta_\varepsilon(s)} - \varepsilon a(s)| ds \\ &\leq \int_0^t \varrho(s) \left(\frac{\Theta_\varepsilon(s)^2}{2} + |\Theta_\varepsilon(s) - \varepsilon a(s)| \right) ds \\ &\leq \varepsilon^2 \int_0^t \varrho(s) \left(\frac{C}{2} (s+1)^4 + C_3 \right) ds = C_4 \varepsilon^2, \end{aligned}$$

using (5.16) and (5.17). Therefore, we proved that, for any $\varepsilon < \delta t^2 / (2(C_3 \vee C_4))$, we have

$$|\Theta_\varepsilon(t) - \varepsilon a(t)| \leq \frac{\delta \varepsilon t^2}{2} \quad \text{and} \quad |\Theta'_\varepsilon(t) - \varepsilon a'(t)| \leq \frac{\delta \varepsilon t^2}{2}$$

But, on the other hand, it follows from Lemma 5.4 that there exists $t_0 > 0$ such that, for all $t \geq t_0$,

$$|a(t) - ct^2| \leq \frac{\delta}{2} t^2 \quad \text{and} \quad |a'(t) - 2ct| \leq \frac{\delta}{2} t^2,$$

so the result holds for $t \geq t_0$ and $\varepsilon < \delta t^2 / (2(C_3 \vee C_4))$. \square

Using the previous lemma, we now can control Θ_ε at some large fixed time t . Then, in order to get the asymptotics of $\Theta_\varepsilon(1/\sqrt{\varepsilon})$ as $\varepsilon \rightarrow 0$, we can replace $\varrho(s)$ by $2/s^2$ for $s \in [t, 1/\sqrt{\varepsilon}]$ in Equation (5.15) and this will prove the following result.

Lemma 5.7. *Let $\theta_1, \theta_2 \in \mathbb{R}$. Then, we have*

$$\Theta_\varepsilon(1/\sqrt{\varepsilon}) \xrightarrow{\varepsilon \rightarrow 0} \log \frac{\sinh^2(\sqrt{3c})}{3c} \quad \text{and} \quad \frac{\Theta'_\varepsilon(1/\sqrt{\varepsilon})}{\sqrt{\varepsilon}} \xrightarrow{\varepsilon \rightarrow 0} 2(\sqrt{3c} \coth(\sqrt{3c}) - 1),$$

where $c = \gamma_1 \theta_1 + \gamma_2 \theta_2$ and γ_1, γ_2 are defined in Lemma 5.4.

Proof. By Proposition 3.7, we can choose $t_1 \geq 1$ large enough such that

$$\forall s \geq t_1, \quad 0 < \frac{2}{s^2} - \frac{1}{s^{5/2}} \leq \varrho(s) \leq \frac{2}{s^2}. \quad (5.18)$$

We fix some $t \geq t_1$, and consider z the solution (depending on t and ε) of the following equation:

$$\forall s \geq t, \quad z''(s) = \frac{2}{s^2}(1 - e^{-z(s)}),$$

with initial conditions $z(t) = \Theta_\varepsilon(t)$ and $z'(t) = \Theta'_\varepsilon(t)$. Then, z is explicitly given by

$$\forall s \geq t, \quad z(s) = \log \frac{\sinh^2(\sqrt{\xi}(s-t) + \zeta)}{\xi s^2},$$

where

$$\xi := \frac{1}{t^2} \left(\left(1 + \frac{t}{2} \Theta'_\varepsilon(t) \right)^2 - e^{-\Theta_\varepsilon(t)} \right) \quad \text{and} \quad \zeta := \operatorname{arccosh} \left(e^{\Theta_\varepsilon(t)/2} \left(1 + \frac{t}{2} \Theta'_\varepsilon(t) \right) \right).$$

Using Lemma 5.6, we get that

$$\lim_{t \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \left| \frac{\xi}{\varepsilon} - 3c \right| = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \left| \frac{\zeta}{\sqrt{\varepsilon}} - \sqrt{3ct} \right| = 0,$$

and it follows that

$$\lim_{t \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \left| z(1/\sqrt{\varepsilon}) - \log \frac{\sinh^2(\sqrt{3c})}{3c} \right| = 0, \quad (5.19)$$

$$\lim_{t \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \left| \frac{z'(1/\sqrt{\varepsilon})}{\sqrt{\varepsilon}} - 2(\sqrt{3c} \coth(\sqrt{3c}) - 1) \right| = 0. \quad (5.20)$$

Now, we have to prove that $\Theta_\varepsilon(1/\sqrt{\varepsilon})$ and $\Theta'_\varepsilon(1/\sqrt{\varepsilon})$ are close to $z(1/\sqrt{\varepsilon})$ and $z'(1/\sqrt{\varepsilon})$ when $\varepsilon \rightarrow 0$ and then $t \rightarrow \infty$. Firstly, it follows from (5.18) that $\Theta_\varepsilon(s) \leq z(s)$ for any $s \geq t$. Then, recalling from (5.16) that there exists $C > 0$ such that $\Theta_\varepsilon(s) \leq C\varepsilon s^2$ for any $s \geq 1$, we get that, for any $s \geq t$,

$$\begin{aligned} (z - \Theta_\varepsilon)''(s) &= \frac{2}{s^2} \left(e^{-\Theta_\varepsilon(s)} - e^{-z(s)} \right) + \left(\frac{2}{s^2} - \varrho(s) \right) \left(1 - e^{-\Theta_\varepsilon(s)} \right) \\ &\leq \frac{2}{s^2} (z - \Theta_\varepsilon)(s) + \frac{C\varepsilon}{\sqrt{s}}, \end{aligned} \quad (5.21)$$

using that $\Theta_\varepsilon(s), z(s) \geq 0$ and (5.18). This is the same differential inequality as in (5.9), so we get the same result as in (5.10): for any $s \geq t$,

$$z(s) - \Theta_\varepsilon(s) \leq \frac{2C\varepsilon}{\sqrt{t}} s^2 + \frac{2C\varepsilon t^{5/2}}{15s}. \quad (5.22)$$

Recalling that $z(s) - \Theta_\varepsilon(s) \geq 0$, we get

$$\lim_{t \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} |z(1/\sqrt{\varepsilon}) - \Theta_\varepsilon(1/\sqrt{\varepsilon})| = 0.$$

Coming back to (5.19), the first convergence of the lemma follows. For the second one, note that we have $\Theta'_\varepsilon(1/\varepsilon) \leq z'(1/\varepsilon)$ and, integrating (5.21) and using (5.22),

$$\begin{aligned} (z - \Theta_\varepsilon)'(1/\varepsilon) &\leq \int_t^{1/\sqrt{\varepsilon}} \left(\frac{2}{s^2} \left(\frac{2C\varepsilon}{\sqrt{t}} s^2 + \frac{2C\varepsilon t^{5/2}}{15s} \right) + \frac{C\varepsilon}{\sqrt{s}} \right) ds \\ &\leq \frac{4C\sqrt{\varepsilon}}{\sqrt{t}} + \frac{2C\varepsilon\sqrt{t}}{15} + 2C\varepsilon^{3/4}. \end{aligned}$$

Combining this with (5.20), it proves the second convergence. \square

We can now conclude this subsection by proving Theorem 5.1.

Proof of Theorem 5.1. Let $\theta_1, \theta_2 > 0$. Then, with $\varepsilon = 1/t^2$ and Θ_ε the solution of (5.15), we have by Lemma 5.7

$$\begin{aligned} \mathbb{E}_{x_t} \left[e^{-\theta_1 t^{-2} N_t - \theta_2 t^{-2} M_t} \right] &= \phi_{\theta_1/t^2, \theta_2/t^2}(t, x_t) = \exp \left(-\Theta_\varepsilon(1/\sqrt{\varepsilon}) - \frac{x_t}{t} \frac{\Theta'_\varepsilon(1/\sqrt{\varepsilon})}{\sqrt{\varepsilon}} \right) \\ &\xrightarrow{t \rightarrow \infty} \frac{3c}{\sinh^2(\sqrt{3c})} \exp \left(-2x \left(\sqrt{3c} \coth(\sqrt{3c}) - 1 \right) \right), \end{aligned} \quad (5.23)$$

where $c = \gamma_1 \theta_1 + \gamma_2 \theta_2$. The right-hand side of (5.23), seen as a function of (θ_1, θ_2) , is the Laplace transform of $(\gamma_1 \eta, \gamma_2 \eta)$ by [23, Equation (2.5)], hence it proves the result. \square

Proof of Theorem 1.11. It follows immediately from Theorems 4.5 and 5.1. \square

6 Some properties of Derrida–Retaux models

In this section, we prove some elementary results on general continuous–time DR models, in other words, we do not assume the initial law to be a mixture of exponential laws. We end the paper by some open questions.

In the continuous DR model, the analogous to the free energy in the discrete case is the quantity

$$F_\infty(\mu_0) := \lim_{t \rightarrow \infty} e^{-t} \mathbb{E}(X_t) \in [0, \infty). \quad (6.1)$$

In this section, we prove that the above quantity always exists, and that $e^{-t}X_t$ converges in law, as $t \rightarrow \infty$, toward an exponential random variable with parameter $F_\infty(\mu_0)^{-1}$. We also give some necessary conditions for $F_\infty(\mu_0) = 0$, giving a partial characterization of the supercritical and subcritical laws. The proofs are based on simple coupling arguments, as well as the computation of $\mathbb{E}(f(X_t))$ for various \mathcal{C}^1 function f , using the SDE representation of the process and the Itô formula.

Lemma 6.1. *The limit in (6.1) exists for any initial distribution μ_0 on \mathbb{R}_+ . More precisely, if $\mathbb{E}(X_0) = \infty$, then $\lim_{t \rightarrow \infty} e^{-t}X_t = \infty$ in law and, if $\mathbb{E}(X_0) < \infty$, then*

$$\forall t \geq 0, \quad F_\infty(\mu_0) = e^{-t} \mathbb{E}(X_t) - \int_t^\infty e^{-s} \mathbb{P}(X_s > 0) ds.$$

Proof. We first observe that for all $t \geq 0$, the random variable X_t stochastically dominates

$$S_t = \sum_{j=1}^{\#\mathcal{N}_t} X_0^j - L_t,$$

where $(X_0^j, j \geq 1)$ are i.i.d. copies of X_0 , and L_t is the total length at time t of the Yule tree. By classical results on continuous time branching processes (see Athreya and Ney [2, Section III.7]), there exists a random variable N_∞ such that $e^{-t}\#\mathcal{N}_t \rightarrow N_\infty > 0$ a.s. and it follows easily that

$$e^{-t}L_t = e^{-t} \int_0^t \#\mathcal{N}_s ds \xrightarrow[t \rightarrow \infty]{\text{a.s.}} N_\infty. \quad (6.2)$$

If $\mathbb{E}(X_0) = \infty$, then, by the law of large numbers, we get $\lim_{t \rightarrow \infty} e^{-t}S_t = \infty$ a.s. and therefore $F_\infty(\mu_0) = \infty$.

Reciprocally, we now assume that $\mathbb{E}(X_0) < \infty$, and we prove that $F_\infty(\mu_0)$ exists and is finite. Using the SDE representation of X_t , we have immediately that

$$\mathbb{E}(X_t) = \mathbb{E}(X_0) - \int_0^t \mathbb{P}(X_s > 0) ds + \int_0^t \mathbb{E}(X_s) ds.$$

Solving this differential equation, we obtain that

$$\mathbb{E}(X_t) = \mathbb{E}(X_0)e^t - e^t \int_0^t e^{-s} \mathbb{P}(X_s > 0) ds.$$

As $\mathbb{P}(X_s > 0) \in [0, 1]$, we deduce that $\mathbb{E}(e^{-t}X_t)$ converges as $t \rightarrow \infty$, and that

$$F_\infty(\mu_0) = \mathbb{E}(X_0) - \int_0^\infty e^{-s} \mathbb{P}(X_s > 0) ds.$$

Applying this formula with the starting law X_t , this concludes the proof of the lemma. \square

This lemma proves the existence of the free energy of the continuous-time DR model, and that as soon as $\mathbb{E}(X_t) < \infty$, the free energy can be seen as a preserved quantity of the evolution. In the next lemma, we observe that if X_0 has moments of order n , then the n th moment of X_t converges toward the n th moment of an exponential random variable of parameter $F_\infty(\mu_0)^{-1}$.

Lemma 6.2. *If $\mathbb{E}(X_0^n) < \infty$ for some $n \in \mathbb{N}$, then $\mathbb{E}(X_t^n)e^{-nt}$ converges to $n!F_\infty(\mu_0)^n$ as $t \rightarrow \infty$.*

Proof. We prove this result by recurrence on n . By Lemma 6.1, we know that $\mathbb{E}(X_t)e^{-t} \rightarrow F_\infty(\mu_0)$ as long as $\mathbb{E}(X_0) < \infty$, hence the property is true for $n = 1$.

Let $n \geq 2$ and assume the lemma to be true for initial random variables in L^k for $k < n$. Let X_0 be a random variable in L^n , we note that $X_0 \in L^k$ for all $k < n$, hence by recurrence assumption, we have

$$\lim_{t \rightarrow \infty} \mathbb{E}((e^{-t}X_t)^k) = k!F_\infty(\mu_0)^k. \quad (6.3)$$

Using the SDE representation of X , we now observe that we can write

$$\mathbb{E}(X_t^n) = \mathbb{E}(X_0^n) - n \int_0^t \mathbb{E}(X_s^{n-1})ds + \int_0^t \mathbb{E}((X_s + \widetilde{X}_s)^n - X_s^n)ds,$$

with \widetilde{X}_s an independent copy of X_s . We rewrite this equality as follows:

$$\mathbb{E}(X_t^n) = \mathbb{E}(X_0^n) + \int_0^t \mathbb{E}(X_s^n)ds + \int_0^t \left(-n \mathbb{E}(X_s^{n-1}) + \sum_{k=1}^{n-1} \binom{n}{k} \mathbb{E}(X_s^k) \mathbb{E}(X_s^{n-k}) \right) ds.$$

This first order differential equation can be solved as

$$\mathbb{E}(X_t^n) = \mathbb{E}(X_0^n)e^t + e^t \int_0^t e^{-s} \left(-n \mathbb{E}(X_s^{n-1}) + \sum_{k=1}^{n-1} \binom{n}{k} \mathbb{E}(X_s^k) \mathbb{E}(X_s^{n-k}) \right) ds, \quad (6.4)$$

which can be rewritten

$$\begin{aligned} \mathbb{E}((X_t e^{-t})^n) &= \mathbb{E}(X_0^n) e^{-(n-1)t} - e^{-(n-1)t} \int_0^t e^{(n-2)s} n \mathbb{E}((X_s e^{-s})^{n-1}) ds \\ &\quad + e^{-(n-1)t} \int_0^t e^{(n-1)s} \sum_{k=1}^{n-1} \binom{n}{k} \mathbb{E}((X_s e^{-s})^k) \mathbb{E}((X_s e^{-s})^{n-k}) ds. \end{aligned}$$

Therefore, using (6.3) and the L'Hospital rule, we obtain that $\mathbb{E}((X_t e^{-t})^n)$ converges toward

$$\begin{aligned} \frac{1}{n-1} \sum_{k=1}^{n-1} \binom{n}{k} \lim_{t \rightarrow \infty} \mathbb{E}((X_s e^{-s})^k) \mathbb{E}((X_s e^{-s})^{n-k}) &= \frac{1}{n-1} \sum_{k=1}^{n-1} \binom{n}{k} F_\infty(\mu_0)^n k!(n-k)! \\ &= n!F_\infty(\mu_0)^n \end{aligned}$$

completing the proof. □

A direct consequence of this proof is the convergence, as $t \rightarrow \infty$, of the law of $e^{-t}X_t$, as soon as X_0 has some finite exponential moments. In fact, this convergence holds as soon as $F_\infty(\mu_0)$ is well-defined, which is equivalent to $\mathbb{E}[X_0] < \infty$ by Lemma 6.1.

Proposition 6.3. *Let $(X_t, t \geq 0)$ be a DR model with $\mathbb{E}(X_0) < \infty$. Then $e^{-t}X_t$ converges in law toward an exponential random variable with parameter $F_\infty(\mu_0)^{-1}$.*

Proof. We first observe that by Lemma 6.1, $\sup_{t \geq 0} \mathbb{E}(e^{-t}X_t) < \infty$. Therefore the family $(e^{-t}X_t, t \geq 0)$ is tight. The only thing left to do is thus to identify the limit in distribution.

Remark that for all $t, T \geq 0$, the random variable X_{t+T} can be written as

$$e^{-(t+T)}X_{t+T} = e^{-T} \sum_{u \in \mathcal{N}_T} e^{-t}X_t^{(u)} - e^{-(t+T)}R_T,$$

where, conditionally on \mathcal{N}_T , $(X_t^{(u)}, u \in \mathcal{N}_T)$ are i.i.d. random variables with the same law as X_t , and R_T is a random variable which takes values between 0 and L_T the length of the Yule tree of height T . Recall from (6.2) that $\lim_{T \rightarrow \infty} e^{-T}L_T$ exists a.s., hence if t and T both tend to ∞ , then $e^{-(t+T)}R_T$ converges to 0 in probability.

We now write $(t_n, n \geq 1)$ a sequence along which (X_{t_n}) converges in law toward a random variable Y . Up to extraction, we can assume that $r_n := t_{n+1} - t_n \rightarrow \infty$. Therefore, we have that

$$Y = \lim_{n \rightarrow \infty} e^{-t_{n+1}}X_{t_{n+1}} = \lim_{n \rightarrow \infty} e^{-r_n} \sum_{u \in \mathcal{N}_{r_n}} e^{-t_n}X_{t_n}^{(u)}.$$

Thus, computing the characteristic function of Y for some $\xi \in \mathbb{R}$ and using that $\#\mathcal{N}_t$ is a geometric random variable with parameter e^{-t} , we have

$$\begin{aligned} \mathbb{E}[e^{i\xi Y}] &= \lim_{n \rightarrow \infty} \mathbb{E} \left[\prod_{u \in \mathcal{N}_{r_n}} e^{i\xi e^{-r_n - t_n} X_{t_n}^{(u)}} \right] \\ &= \lim_{n \rightarrow \infty} \frac{e^{-r_n} \mathbb{E}[e^{i\xi e^{-r_n - t_n} X_{t_n}}]}{1 - (1 - e^{-r_n}) \mathbb{E}[e^{i\xi e^{-r_n - t_n} X_{t_n}}]} \\ &= \lim_{n \rightarrow \infty} \frac{\mathbb{E}[e^{i\xi e^{-r_n - t_n} X_{t_n}}]}{\mathbb{E}(e^{i\xi e^{-r_n - t_n} X_{t_n}}) + e^{r_n} \mathbb{E}[1 - e^{i\xi e^{-r_n - t_n} X_{t_n}}]}. \end{aligned}$$

Using the fact that $\mathbb{E}(X_{t_n} e^{-t_n}) \rightarrow F_\infty(\mu_0)$ by Lemma 6.1, we have that

$$\mathbb{E}[e^{i\xi e^{-r_n - t_n} X_{t_n}}] = 1 + i\xi e^{-r_n} F_\infty(\mu_0) + o(e^{-r_n}) \quad \text{as } n \rightarrow \infty.$$

Therefore, this yields

$$\mathbb{E}[e^{i\xi Y}] = \frac{1}{1 - i\xi F_\infty(\mu_0)},$$

which is the characteristic function of the exponential distribution of parameter $F_\infty(\mu_0)^{-1}$, concluding the proof. \square

In the rest of the section, we now focus on the (sub)critical case $F_\infty(\mu_0) = 0$. Previous lemma shows that $\lim_{t \rightarrow \infty} X_t e^{-t} = 0$ in probability. However, an interesting problem is to obtain the precise asymptotic behavior of X_t in this (sub)critical case. It is known that for discrete DR models with integer-valued measures, we have $\lim_{n \rightarrow \infty} X_n = 0$ in probability when the free energy $F_\infty = 0$. Therefore, this yields the following conjecture.

Conjecture 6.4. *If $F_\infty(\mu_0) = 0$, then $X_t \rightarrow 0$ in probability as $t \rightarrow \infty$.*

Assuming the above conjecture, one might be interested in the limiting law of X_t conditioned to be positive. In light of (2.9), if such a limit exists, its density should satisfy $\phi'(x) = -\phi(0)\phi(x)$, $x > 0$. This implies the following conjecture:

Conjecture 6.5. *If $F_\infty(\mu_0) = 0$, then $\mathcal{L}(X_t | X_t > 0)$ converges to an exponential distribution of parameter larger than or equal to 1.*

We were not able to prove these conjectures, and more generally to obtain characterizations, in terms of the law of X_0 , of measures satisfying $F_\infty(\mu_0) = 0$ (as (1.4) in the discrete case with integer-valued measures). However, we obtained the following partial answer, which is notably similar to what is known in the general discrete DR model.

Proposition 6.6. *Let $(X_t, t \geq 0)$ be a DR process, with μ_0 the law of X_0 satisfying $F_\infty(\mu_0) = 0$. The following results hold:*

1. For any $t, x \geq 0$, $\mathbb{P}(X_t \geq x) \leq e^{1-x}$.
2. For any $t \geq 0$, $\mathbb{E}[X_t^2] < \frac{1}{2}$ and, in particular, $\mathbb{E}[X_t] \leq \frac{1}{\sqrt{2}}$.
3. For any $\theta \in (0, 1)$ and $t \geq 0$,

$$\mathbb{E}\left[\left(1 - \theta X_t - 2(1 - \theta)X_t^2\right)e^{\theta X_t}\right] \geq 0.$$

4. We have $\sup_{t \geq 0} \mathbb{E}[X_t e^{X_t}] < \infty$ and, for any $t \geq 0$, $\mathbb{E}[X_t e^{X_t}] \leq \mathbb{E}[e^{X_t}]$.

Note that Point 4 gives a characterization of the subcritical case similar to (1.4) for the discrete model, but here it does not depend only on moments of X_0 . The main reason for this is that the sign of $\mathbb{E}[X_n 2^{X_n}] - \mathbb{E}[2^{X_n}]$ is an invariant of the dynamics for the discrete model with integer valued measures, whereas it is not the case for $\mathbb{E}[X_t e^{X_t}] - \mathbb{E}[e^{X_t}]$ here.

Proof. We first observe that, if $F_\infty(\mu_0) = 0$, then $\mathbb{E}(X_t) \leq 1$ for any $t \geq 0$. Indeed, by Lemma 6.1, we have

$$\mathbb{E}[X_t] = \int_t^\infty e^{-(s-t)} \mathbb{P}(X_s > 0) ds \tag{6.5}$$

and it follows that $\mathbb{E}[X_t] \leq 1$.

Let $s, t \geq 0$, we consider the Yule tree representation of X_{t+s} . One has immediately that

$$X_{t+s} \geq \sum_{u \in \mathcal{N}_s} (X_t^{(u)} - s)_+ \geq \sum_{u \in \mathcal{N}_s} \mathbb{1}_{\{X_t^{(u)} > s+1\}},$$

where, conditionally on \mathcal{N}_s , $(X_t^{(u)}, u \in \mathcal{N}_s)$ are i.i.d. random variables with the same law as X_t . Hence, we have that $\mathbb{E}(X_{t+s}) \geq e^s \mathbb{P}(X_t > s+1)$ and, since $\mathbb{E}(X_{t+s}) \leq 1$, it proves Point 1.

It follows from Point 1. that $\sup_{t \geq 0} \mathbb{E}[X_t^2] < \infty$. Then, using (6.4) with $n = 2$, we have

$$\mathbb{E}[X_t^2] = e^t \mathbb{E}[X_0^2] + e^t \int_0^t 2e^{-s} \mathbb{E}[X_s](\mathbb{E}[X_s] - 1) ds,$$

and, letting $t \rightarrow \infty$,

$$\mathbb{E}[X_0^2] = 2 \int_0^\infty e^{-s} \mathbb{E}[X_s](1 - \mathbb{E}[X_s]) ds \leq \frac{1}{2},$$

using that $\mathbb{E}[X_s] \in [0, 1]$ for any $s \geq 0$. Finally, $\mathbb{E}[X_t^2] \leq \frac{1}{2}$ holds for any $t \geq 0$, because $F_\infty(\mu_t) = e^t F_\infty(\mu_0) = 0$. This proves Point 2.

Let $b \in \mathbb{R}$ be a constant whose value will be fixed later on. Consider the function

$$\phi(t) := \mathbb{E}\left[(1 - \theta X_t - b X_t^2)e^{\theta X_t}\right].$$

By Itô's formula, we have that

$$\begin{aligned} \phi'(t) &= \mathbb{E}\left[\left((\theta^2 + 2b)X_t + \theta b X_t^2\right)e^{\theta X_t}\right] \\ &\quad + \mathbb{E}\left[\left(1 - \theta(X_t + Y_t) - b(X_t + Y_t)^2\right)e^{\theta(X_t + Y_t)}\right] - \mathbb{E}\left[\left(1 - \theta X_t - b X_t^2\right)e^{\theta X_t}\right], \end{aligned}$$

where Y_t is an independent copy of X_t . We introduce the functions

$$g(t) := \mathbb{E}[e^{\theta X_t}] \quad \text{and} \quad h(t) := \mathbb{E}[X_t e^{\theta X_t}],$$

we can rewrite

$$\phi' = (2g - \theta - 1)\phi + 2b(h - h^2) - (g^2 - \theta g).$$

Let $\eta := 2b(h - h^2) - (g^2 - \theta g)$, then, for any $0 \leq r \leq t$, the differential equation can be solved as

$$\phi(t) = e^{\int_r^t (2g(s) - \theta - 1) ds} \left(\phi(r) + \int_r^t \eta(s) e^{\int_r^s (2g(u) - \theta - 1) du} ds \right).$$

On the one hand, we have $g \geq 1$ and therefore $g^2 - \theta g \geq 1 - \theta$ and, on the other hand, $h - h^2 \leq 1/4$. Thus, choosing $b := 2(1 - \theta)$, we have $\eta \leq b/2 - 1 + \theta = 0$. If $\phi(r) < 0$ for some $r \geq 0$, then

$$\phi(t) \leq e^{\int_r^t (2g(s) - \theta - 1) ds} \phi(r) \xrightarrow{t \rightarrow \infty} -\infty,$$

which contradicts the fact that, by Point 1., $\sup_{t \geq 0} \mathbb{E}[e^{\lambda X_t}] < \infty$ for any $\lambda \in (\theta, 1)$. Therefore, $\phi(r) \geq 0$ for any $r \geq 0$ and this proves Point 3.

Now we prove that necessarily $\mathbb{E}[X_0 e^{X_0}] < \infty$. Indeed, if $\mathbb{E}[X_0 e^{X_0}] = \infty$, then $\lim_{n \rightarrow \infty} \mathbb{E}[X_0 e^{X_0} \mathbb{1}_{\{2 \leq X_0 < n\}}] \rightarrow \infty$. Choose and fix an integer n such that

$$\mathbb{E}[X_0 e^{X_0} \mathbb{1}_{\{2 \leq X_0 < n\}}] > 2.$$

Write $Y_0 := X_0 \mathbb{1}_{\{2 \leq X_0 < n\}}$ and $(Y_t)_{t \geq 0}$ the solution of the SDE (2.4) starting from Y_0 . By coupling, a.s. $X_t \geq Y_t$ for all $t \geq 0$. Observe that

$$\mathbb{E}[(1 - Y_0)e^{Y_0}] = \mathbb{P}(Y_0 = 0) + \mathbb{E}[(1 - Y_0)e^{Y_0} \mathbb{1}_{\{Y_0 > 0\}}] \leq 1 - \frac{1}{2} \mathbb{E}[Y_0 e^{Y_0} \mathbb{1}_{\{Y_0 > 0\}}] < 0.$$

By continuity, for $\theta < 1$ sufficiently close to 1, we have $\mathbb{E}[(1 - \theta Y_0)e^{\theta Y_0}] < 0$, which, by Point 3., proves that $\lim_{t \rightarrow \infty} e^{-t} \mathbb{E}[Y_t] > 0$. But we have $F_\infty(\mu_0) \geq \lim_{t \rightarrow \infty} e^{-t} \mathbb{E}[Y_t] > 0$, so it contradicts the assumption $F_\infty(\mu_0) = 0$. This proves that $\mathbb{E}[X_0 e^{X_0}] < \infty$.

The same argument shows that, for any $t > 0$, $\mathbb{E}[X_t e^{X_t}] < \infty$. Therefore, noting that $(1 - \theta)X_t^2 e^{\theta X_t} \leq X_t e^{X_t} \sup_{x \geq 0} x e^{-x}$, we can apply dominated convergence theorem as $\theta \rightarrow 1$ in Point 3. and get $\mathbb{E}[(1 - X_t)e^{X_t}] \geq 0$ for any $t \geq 0$. This proves the second part of Point 4. and can also be re-written as

$$\mathbb{E}[(1 - X_t)e^{X_t} \mathbb{1}_{\{X_t < 2\}}] \geq \mathbb{E}[(X_t - 1)e^{X_t} \mathbb{1}_{\{X_t \geq 2\}}].$$

Note that, on the one hand, $\mathbb{E}[(1 - X_t)e^{X_t} \mathbb{1}_{\{X_t < 2\}}] \leq e^2$ and, on the other hand, $\mathbb{E}[(X_t - 1)e^{X_t} \mathbb{1}_{\{X_t \geq 2\}}] \geq \frac{1}{2} \mathbb{E}[X_t e^{X_t} \mathbb{1}_{\{X_t \geq 2\}}]$. Hence $\mathbb{E}[X_t e^{X_t} \mathbb{1}_{\{X_t \geq 2\}}] \leq 2e^2$, which yields that $\mathbb{E}[X_t e^{X_t}] \leq 4e^2$ and proves the first part of Point 4. \square

In addition to Conjectures 5.2 and 6.4, we list some open questions below. At first, Proposition 6.6 gives some necessary conditions on μ_0 so that the free energy vanishes. Naturally we may ask

Question 6.7. *What are necessary and sufficient condition on μ_0 so that $F_\infty(\mu_0) = 0$?*

It is predicted by physicists that the infinite order phase transition holds for a large classe of recursive model. This motivates our next question:

Question 6.8. *Replace the indicator function $\mathbb{1}_{\{x>0\}}$ in (1.9) by a monotone increasing bounded function f with $f(0) = 0$, such as $\tanh(x)$, $1 - e^{-x}$ or $x \wedge 1$. For regular f , there exists a unique strong solution of the McKean–Vlasov type SDE, see [15]. Are the phase transitions for these models of same nature as for the DR process?*

Concerning our exactly solvable model, a crucial point is the preservation of the mixture of exponential distributions and Dirac masses at 0. This was proved by using the explicit solution of (1.10). However, there is no direct probabilistic proof of this result, yielding the following question.

Question 6.9. *Is there any probabilistic proof that the family of mixtures of exponentials and Dirac masses at 0 is stable under the dynamics of the continuous–time DR model?*

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