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**A WAY TO MODEL STOCHASTIC PERTURBATIONS IN
POPULATION DYNAMICS MODELS WITH BOUNDED
REALIZATIONS**

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ABSTRACT. In this paper, we analyze the use of the Ornstein-Uhlenbeck process to model dynamical systems subjected to bounded noisy perturbations. In order to discuss the main characteristics of this new approach we consider some basic models in population dynamics such as the logistic equations and competitive Lotka-Volterra systems. The key is the fact that these perturbations can be ensured to keep inside some interval that can be previously fixed, for instance, by practitioners, even though the resulting model does not generate a random dynamical system. However, one can still analyze the forwards asymptotic behavior of these random differential systems. Moreover, to illustrate the advantages of this type of modeling, we exhibit an example testing the theoretical results with real data, and consequently one can see this method as a realistic one, which can be very useful and helpful for scientists.

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1. INTRODUCTION

In the last years, researchers of many areas in life sciences have been more and more interested in considering non-deterministic parameters in the mathematical models since it allows them to set up models which are much more realistic. However, there are many different ways of introducing random or stochastic disturbances in deterministic models.

First, a decision should be made about the most appropriate kind of stochastic/random perturbation for our model. Amongst the several stochastic processes which can be potential candidates, we need to decide which one fits better the set up of our stochastic/random model. Next, we need to decide about the location in the equations where the disturbances must appear. Later, we should think whether our resulting stochastic/random model is realistic and, of course, we also need to have some tractability in order to carry out a mathematical analysis and effective computations. Due to these facts, it is necessary to find a reasonable balance which make our work both tractable and realistic and, consequently, original and interesting.

The most common stochastic process that is considered when modeling disturbances in real life is the well-known standard Wiener process, see for instance [1, 2] where the authors study random and stochastic modeling for a SIR model, [3, 4, 5] where stochastic prey-predator Lotka-Volterra systems are analyzed or [6, 7, 8, 9] where different ways of modeling stochasticity in the chemostat model are investigated. Nevertheless, this stochastic process has the property of having continuous but not bounded variation paths, which does not suit to the idea of modeling real situations since, in most of cases, the real life is subjected to fluctuations which are known to be bounded.

In some cases analyzed in the previous literature (see [6, 8]), the use of a standard Wiener process to perturb some parameters in deterministic systems can lead to a non-realistic model; for instance, the positiveness of solutions are not necessarily preserved as a consequence of the arbitrary large values and the large fluctuations of this Wiener process. However, in other situations, one can modify the way to perturb the deterministic system, still using standard Wiener processes, and the positiveness of solutions is preserved too (see [7, 9]).

Henceforth, in this paper we consider a noise whose realizations (or sample paths) remains bounded in an interval (previously fixed, for instance, by practitioners), allowing us, in addition, to perform calculations in a simple way. We will be able to obtain realistic mathematical models due to the boundedness of the considered stochastic process, and will be also able to prove the existence of absorbing and attracting sets which will not depend on the realizations of the noise. As a consequence, we will ensure the persistence and coexistence of the species (or population) under some conditions on the parameters in the models.

We remark that every result will be proved forwards in time, unlike the usual pullback convergence, which is characteristic in the theory of random dynamical systems. Although these pullback concepts are very useful for the development of the theory of random dynamical systems, in some cases from applications it might not

provide meaningful information about the forwards behavior of the system. Despite our resulting random differential systems do not generate random dynamical systems, we will be able to investigate the long-time behavior of the random system for every fixed event, which is a relevant and helpful improvement to prove the forwards convergence that we use in our work.

The manuscript is organized as follows. In Section 2 we introduce an Ornstein-Uhlenbeck (O-U) process depending on some parameters whose effects on the dynamics of the stochastic process are detailed explained. Also we recall some essential and useful properties of this process when dealing with the mathematical models. In Sections 3 and 4 we present an example of a logistic differential equation with random disturbances in the environment and the growth rate by means of the O-U process. Therefore, in Section 5 we present an example concerning the parameter estimation in the logistic equation affected by the O-U process by setting up an *observer* which will consist of another differential system providing information about the behavior of the state variables. Then, in Section 6 we analyze a random competitive Lotka-Volterra system where the growth rates of the species are affected by noise, namely by means of the O-U process. In Section 7 we recall the advantages of using the O-U process when modeling reality in the case of the chemostat. Finally, in Section 8 we include some comments and conclusions.

2. THE ORNSTEIN-UHLENBECK PROCESS.

The key in our current work consists of perturbing the deterministic models by means of a suitable O-U process defined as the following random variable

$$(2.1) \quad z_{\beta, \gamma}^*(\theta_t \omega) = -\beta \gamma \int_{-\infty}^0 e^{\beta s} \theta_t \omega(s) ds, \quad t \in \mathbb{R}, \omega \in \Omega, \beta, \gamma > 0,$$

where ω denotes a standard Wiener process in a certain probability space $(\Omega, \mathcal{F}, \mathbb{P})$, β and γ are positive parameters which will be explained in more detail below and θ_t denotes the usual Wiener shift flow given by

$$\theta_t \omega(\cdot) = \omega(\cdot + t) - \omega(t), \quad t \in \mathbb{R}.$$

We note that the O-U process (2.1) can be obtained as the stationary solution of the Langevin equation

$$(2.2) \quad dz + \beta z dt = \gamma d\omega.$$

We would like to highlight that the O-U process is frequently used to transform stochastic models affected by the standard Wiener process into random ones (see [1, 6, 7, 8]), which are much more tractable from the mathematical point of view, but both parameters β and γ are not taken into account or do not play any relevant role. Nevertheless, in the framework that we introduce in this paper we use directly this suitable O-U process depending on the parameters previously mentioned since they will be the key of the advantages provided by this way of modeling, as we will show in the rest of this work.

Due to the importance of those parameters, we introduce them now in a more detailed way and we show the effects that they have in the dynamics of the realizations of the O-U process.

The O-U process given by (2.1) is a stationary mean-reverting Gaussian stochastic process where $\beta > 0$ is a *mean reversion constant* that represents the strength with which the process is attracted by the mean or, in other words, how *strongly* our system reacts under some perturbation, and $\gamma > 0$ is a *volatility constant* which represents the variation or the size of the noise independently of the amount of the noise $\alpha > 0$. In fact, the O-U process can describe the position of some particle by taking into account the friction, which is the main difference with the standard Wiener process and makes our perturbations to be a better approach to the real ones than the ones obtained when using simply the standard Wiener process. In addition, the O-U process could be understood as a generalization of the standard Wiener process as well in the sense that it would correspond to take $\beta = 0$ and $\gamma = 1$ in (2.1). In fact, the O-U also provides a link between the standard Wiener process and no noise at all, as we will see later.

Now, we would like to illustrate the relevant effects caused by both parameters β and γ on the evolution of realization of the O-U process.

Fixed $\beta > 0$. Then, the volatility of the process increases when considering larger values of γ and the evolution of the process is smoother when taking smaller values of γ , which sounds reasonable due to the fact that γ decides the amount of noise introduced to dz , the term which measures the variation of the process. Henceforth, the process will be subject to suffer much more disturbances when taking a larger value of γ . This behavior can be observed in Figure 1, where we simulate two realizations of the O-U process with $\beta = 1$ and we consider $\gamma = 0.1$ (blue) and $\gamma = 0.5$ (orange).

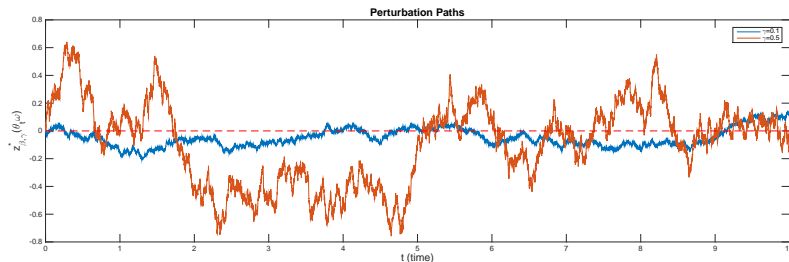


FIGURE 1. Effects of the mean reverting constant on the O-U process

Fixed $\gamma > 0$. In this case the process tends to go further away from the mean value when considering smaller values of β and the attraction of the mean value increases when taking larger values of β . This behavior seems logical since β has a huge influence on the drift of the Langevin equation (2.1). We can observe this behavior in Figure 2, where we simulate two realizations of the O-U process with $\gamma = 0.1$ and we take $\beta = 1$ (blue) and $\beta = 10$ (orange).

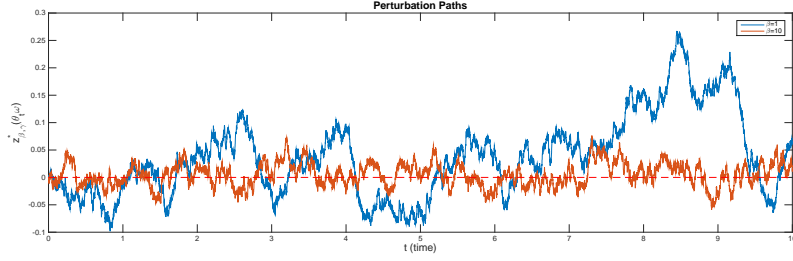


FIGURE 2. Effects of the volatility constant on the O-U process

Once presented the O-U process and the effects that its parameters cause on the behavior of its realization, we state now some essential properties that it satisfies which will be another important key point of the paper.

Proposition 2.1. *There exists a θ_t -invariant set $\tilde{\Omega} \in \mathcal{F}$ of Ω of full \mathbb{P} -measure such that for $\omega \in \tilde{\Omega}$ and $\beta, \gamma > 0$, we have*

(i) *the random variable $|z_{\beta, \gamma}^*(\omega)|$ is tempered.*

(ii) *the mapping*

$$(t, \omega) \rightarrow z_{\beta, \gamma}^*(\theta_t \omega) = -\beta\gamma \int_{-\infty}^0 e^{\beta s} \omega(t+s) ds + \omega(t)$$

is a stationary solution of (2.1) with continuous trajectories;

(iii) *for any $\omega \in \tilde{\Omega}$ one has*

$$\begin{aligned} \lim_{t \rightarrow \pm\infty} \frac{|z_{\beta, \gamma}^*(\theta_t \omega)|}{t} &= 0; \\ \lim_{t \rightarrow \pm\infty} \frac{1}{t} \int_0^t z_{\beta, \gamma}^*(\theta_s \omega) ds &= 0; \\ \lim_{t \rightarrow \pm\infty} \frac{1}{t} \int_0^t |z_{\beta, \gamma}^*(\theta_s \omega)| ds &= \mathbb{E}[z_{\beta, \gamma}^*] < \infty; \end{aligned}$$

(iv) *finally, for any $\omega \in \tilde{\Omega}$,*

$$\lim_{\beta \rightarrow \infty} z_{\beta, \gamma}^*(\theta_t \omega) = 0, \quad \text{for all } t \in \mathbb{R}.$$

The proof of Proposition 2.1 is omitted here. We refer the readers to [10] (Lemma 4.1) for the proof of the last statement and to [11, 12] for the proof of the other points.

To sum up the main idea of this framework, we will have to deal with a random differential system depending on the stationary solution of the Langevin equation (2.2) as follows

$$(2.3) \quad \dot{x} = f(x, z_{\beta, \gamma}^*(\theta_t \omega)),$$

where $z_{\beta, \gamma}^*(\theta_t \omega)$ is the stationary solution of the Langevin equation.

Henceforth, the idea consists of fixing an event $\omega \in \Omega$ and then to solve (2.2) in order to introduce its solution into the differential system (2.4). However, we remark that, for a general choice of the parameter β , it is not possible to ensure that the solution $z_{\beta,\gamma}^*(\theta_t\omega)$ is bounded for every time, which would be the most desirable property since it is well known that real models are expected to be subject to bounded perturbations, in fact, unbounded disturbances has no sense when modeling real systems, as we pointed out previously.

Nevertheless, thanks to the last property (iv) in Proposition 2.1, for every fixed event $\omega \in \Omega$, it is possible to find $\beta > 0$ large enough such that the stationary solution of the Langevin equation (2.2), $z_{\beta,\gamma}^*(\theta_t\omega)$, remain inside any bounded interval previously fixed.

Therefore, for every fixed $\omega \in \Omega$ and for proper $\beta > 0$, we compute $z_{\beta,\gamma}^*(\theta_t\omega)$ such that the realizations of the corresponding perturbed parameter remain inside a strictly positive interval, namely $(b_1, b_2) \subset \mathbb{R}$, which should be previously determined depending on the application or provided by practitioners, where $b_2 > b_1 > 0$. For population dynamics, this property concerning the positiveness of the realizations of the perturbed parameter will be essential in this work, in fact, it will be the main key to be able to guarantee strictly positiveness of all state variables which could mean, for instance, to be able to ensure the persistence of the species involved in the corresponding model.

In most of situations one design a suitable mathematical model with random features for a biological problem. But the main concerns are: is the model realistic? is the model useful? To answer these questions one may try to identify the parameters on sets of real data to test if the choice of stochastic process is able to reproduce satisfactorily real situations.

More precisely, in our case, the parameters β and γ in the O-U process can be estimated from the real data by using a simple mean square regression. As an example of how the process describes the fluctuations in the real models, in Figure 3 we can see time history of values that the input flow of a bioreactor has over a time horizon of one-hundred hours. As it was expected, it seems to be mean-reverting and then it could be described as a realizations of the O-U process. When estimating the parameters by using a mean square regression, we obtain that the mean value is $\mu = 0.2002$, the mean-reverting constant is $\beta = 1.3230$ and the volatility constant is $\gamma = 0.0287$.

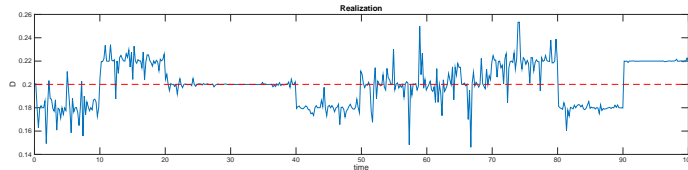


FIGURE 3. Real data: input flow in a bioreactor

Then, in Figure 4 we make some simulations of the O-U process with the previous values of the parameters obtained from Figure 3. As we can observe, every realization remains in the same interval $[0.14, 0.26]$.

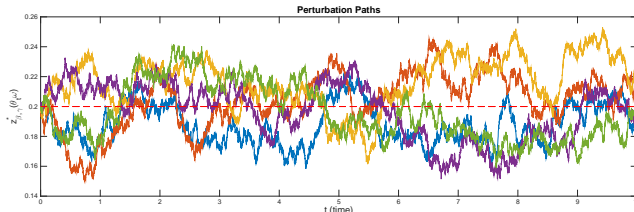


FIGURE 4. Realizations of the O-U process generated with parameters from the real data with β , μ , γ as above

Then, with biological applications in mind, we could consider perturbations of some parameter in a random system of the following kind

$$(2.4) \quad \dot{x} = f(x, z_{\beta, \gamma}^*(\theta_t \omega)).$$

In the real world, one is often interested in analyzing bounded perturbations, thus one can wonder about a suitable way of modeling them. In our method for each ω fixed, we can find values of parameter β_ω such that the random perturbation is bounded in a prefixed interval. This implies that for each realization, we have that our system (2.4) looks like

$$(2.5) \quad \dot{x} = f(x, z_{\beta, \gamma}^*(\theta_t \omega)) = f(x, g(t))$$

where $g(\cdot)$ is a continuous function for each ω , and can be analyzed by the deterministic theory. The main fact is that in some cases the asymptotic behavior of this system may be independent of ω , which then allows us to compare the deterministic and stochastic models and testing how realistic the latter can be.

In other words, the resulting random model is a non-autonomous system perturbed by means of a non-autonomous perturbation which is generated in a random way. As a consequence, in this context, the set of admissible random perturbation coincide with the set of continuous function satisfying properties (i)-(iv) stated in Proposition 2.1.

This kind of non-autonomous perturbations may be obtained by more general stochastic processes, then one could wonder the reason why we are just focusing on the O-U process instead of using another one. The answer is as clear as easy: Indeed, this suitable O-U process given by (2.1) provides us essential ergodic properties which allow us to make calculations when analyzing the resulting random systems, since we will have to deal with integral terms involving such a perturbation.

Notice that real models are subjected to fluctuations which are some what smooth w.r.t. time and it would not make sense to have perturbations which are changing too rapidly from extreme values in short periods of time. This is one of the main reasons to consider the O-U process by taking into account both parameters β and

γ which allow us to have more flexibility to control the noise in order to be able to represent the reality in a better way.

One could also wonder the difference between this way of modeling with the one when considering some random function $a(\theta_t\omega) \in [a_1, a_2]$, which is bounded and continuous respect to the time, as in [1, 13]. On the one hand, in both cases we obtain bounded perturbations, which is the most realistic from the point of view of the applications, as we explained previously. On the other hand, in both cases we are working with continuous functions respect to the time, which is also logical when trying to deal with differential systems as in the current case. Nevertheless, we remark that the random constant $a(\theta_t\omega)$ does not satisfy in general the properties (i)-(iv) stated in Proposition 2.1, as it happens with the O-U process, which are essential when making calculations and allow us to obtain characterizations of absorbing and attracting sets for the solutions of the corresponding random systems as well. Moreover, we would have to assume $a(\theta_t\omega)$ to be continuous and bounded while the O-U process satisfies these properties by definition.

We would also like to remark that, in the classic random case, the continuous function $a(\theta_t\omega)$ is directly generated thanks to the dynamics of the set of events Ω whereas, in our new framework, every event $\omega \in \Omega$ is fixed and the continuous perturbation is obtained by solving the Langevin equation (2.2). Apart from that, the continuous function $a(\theta_t\omega)$ is an arbitrary continuous function with values in some positive bounded interval $[a_1, a_2]$ whereas, in our new framework, we are simply considering the realizations of the perturbed parameter which are realistic from the point of view of the applications.

Let us underline that the realizations by means of suitable O-U process look quite different from the one generated by a random function $a(\theta_t\omega)$. One can observe time to time realizations of $a(\theta_t\omega)$ which are very unlikely to be observed as generating by the suitable O-U processes. Several kind of such realizations are depicted on Figure 5, for instance, the green one approaches the boundaries of the interval $[a_1, a_2]$ and stay close to it for a while, the orange or violet ones stay close from one of the boundaries and, finally, the brown one switches very rapidly between two values close from the boundaries of the interval.

The realizations generated by the suitable O-U looks more realistic in the sense that its is similar to an agitated particule with a recall force to the mean value.

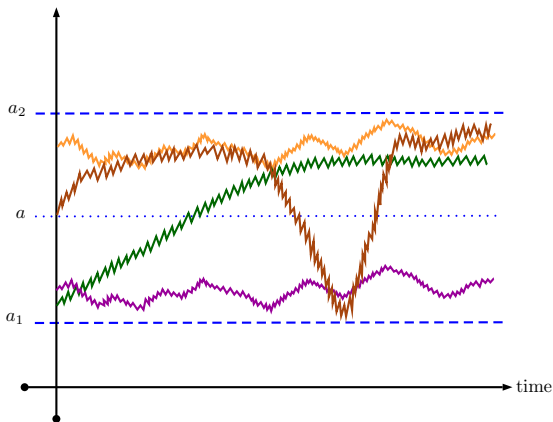


FIGURE 5. Examples of non-realistic realizations of the perturbed parameter

In both cases, when considering a classical random function $a(\theta_t\omega)$ and also when considering a perturbation by means of the O-U process $a + \alpha z_{\beta,\gamma}^*(\theta_t\omega)$, the perturbations are bounded, then one can expect to find bounds for the solutions of the system and, therefore, to be able to provide some conditions under which the persistence of the populations involved in the model can be ensured. Nevertheless, there are important differences between both cases the classic and the new one, for instance, the natural context in the classic random case is to study the pullback convergence whereas, in the new random case involving the new suitable O-U process, the solutions may not generate a random dynamical system, since β in fact depends on ω . However, it does not present any inconvenience since we can analyze the random system for every fixed $\omega \in \Omega$, as we explained before.

In addition, we can prove every mathematical result to hold forwards in time, which is much more realistic than the pullback convergence obtained in the classic case. This improvement concerning the forwards convergence is also very related with the ergodic properties stated Proposition 2.1 (iii) which are proved to hold also forwards in time.

In order to illustrate the modeling approach we propose, we are going to consider in this paper two well known models in the ecology literature: the logistic and the competition Lotka-Volterra model, that we revisit here introducing noise. In addition, we will consider an observer dynamics to show the parameter estimation despite noise in the logistic equation affected by the O-U process, a problem that has not been treated up to now in the literature even though it is object of very interest.

3. ENVIRONMENTAL PERTURBATIONS IN THE LOGISTIC MODEL.

In this section we consider the classical logistic equation given by

$$(3.1) \quad \frac{dx}{dt} = x(a - x),$$

where $x = x(t)$ denotes the number of individual of some population of a certain species and a is the carrying capacity.

As it is very well-known, the carrying capacity of the population whose dynamics are modeled by the logistic equation can be affected by many external factors present on the environment as the climate or the temperature, to name a few. Then it has sense to consider random perturbations on the carrying capacity such that we have the following random equation

$$(3.2) \quad \frac{dx}{dt} = x(a + \alpha z_{\beta, \gamma}^*(\theta_t \omega) - x),$$

where $z_{\beta, \gamma}^*(\theta_t \omega)$ denotes the Ornstein-Uhlenbeck process that we introduced previously and $\alpha > 0$ is the amount of noise.

The solution of the random logistic equation (3.2) exists and its explicit expression is given by

$$(3.3) \quad x(t; 0, \omega, x_0) = \frac{x_0 e^{\int_0^t a + \alpha z_{\beta, \gamma}^*(\theta_s \omega) ds}}{1 + x_0 \int_0^t e^{\int_0^s a + \alpha z_{\beta, \gamma}^*(\theta_\tau \omega) d\tau} ds},$$

for any initial value $x_0 \geq 0$, any $\omega \in \Omega$ and for all $t \geq 0$.

In addition, thanks to a suitable choice of the parameter β in the O-U process presented in the introduction of the paper, we know that $a + \alpha z_{\beta, \gamma}^*(\theta_t \omega) \in [\underline{a}, \bar{a}]$ for every $t \in \mathbb{R}$, where $\bar{a} > \underline{a}$ are positive values. Then from (3.2) we can obtain the following differential inequalities

$$(3.4) \quad x(\underline{a} - x) \leq \frac{dx}{dt} \leq x(\bar{a} - x),$$

whence we can deduce that, as soon as we consider an initial value of the species $x_0 < \underline{a}$, then the dynamics of the population is increasing till it reaches the curve $a + \alpha z_{\beta, \gamma}^*(\theta_t \omega)$ which remains inside the positive interval $[\underline{a}, \bar{a}]$.

Henceforth, from (3.4) it can be deduced that, for every $\varepsilon > 0$, any $\omega \in \Omega$ and any initial value $x_0 < \underline{a}$, there exists some time $T(\varepsilon, \omega) > 0$ such that

$$(3.5) \quad \underline{a} - \varepsilon \leq x(t; 0, \omega, x_0) \leq \bar{a} + \varepsilon,$$

for every $t \geq T(\varepsilon, \omega)$.

From the previous analysis we obtain that, for any $\varepsilon > 0$, $B_\varepsilon = [\underline{a} - \varepsilon, \bar{a} + \varepsilon]$ is a deterministic absorbing set for the solutions of (3.2).

Therefore, $B_0 = [\underline{a}, \bar{a}]$ is a positive attracting set for the solutions of (3.2), i.e.,

$$(3.6) \quad \lim_{t \rightarrow +\infty} \sup_{x_0 \in (0, \underline{a})} \inf_{b_0 \in B_0} |x(t, 0, \omega, x_0) - b_0| = 0.$$

Now, we present some numerical simulations to support the results previously provided and the advantages of using the suitable O-U process presented here when modeling realistic problems. From now on, the blue dashed lines represent the solutions of the deterministic models and the rest are different realizations of the random ones.

In Figure 6 we can see two panels representing several realizations of the solution of the random logistic equation (3.2) for the initial value $x_0 = 2.4$, the nominal carrying capacity is $a = 3$, the amount of noise is $\alpha = 2$ (top) and $\alpha = 2.2$ (bottom), the mean reverting constant is $\beta = 1$ (top) and $\beta = 10$ (bottom) and the volatility constant is $\gamma = 0.1$ (top) and $\gamma = 0.2$ (bottom).

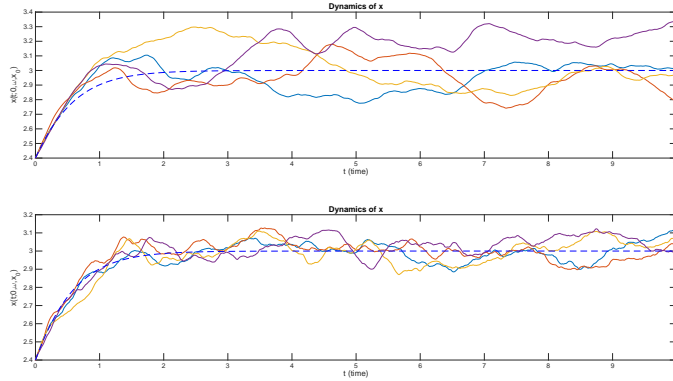


FIGURE 6. Realizations of the solution of the random logistic equation with perturbed carrying capacity for $x_0 = 2.4$. $\alpha = 2$, $\beta = 1$, $\gamma = 0.1$ (top) and $\alpha = 2.2$, $\beta = 10$, $\gamma = 0.2$ (bottom)

In Figure 7 we display two panels representing several realizations of the solution of the random logistic equation (3.2) for the initial value $x_0 = 0.2$, the nominal carrying capacity is $a = 3$, the amount of noise is $\alpha = 2$ (top) and $\alpha = 2.2$ (bottom), the mean reverting constant is $\beta = 1$ (top) and $\beta = 10$ (bottom) and the volatility constant is $\gamma = 0.4$. Compared to Figure 6, now we increase the volatility constant which is significant for small values of the mean reverting constant (as it can be seen in the figure of the top) but the noise can be reduced if we increase the mean reverting constant even though the volatility constant is not decreased.

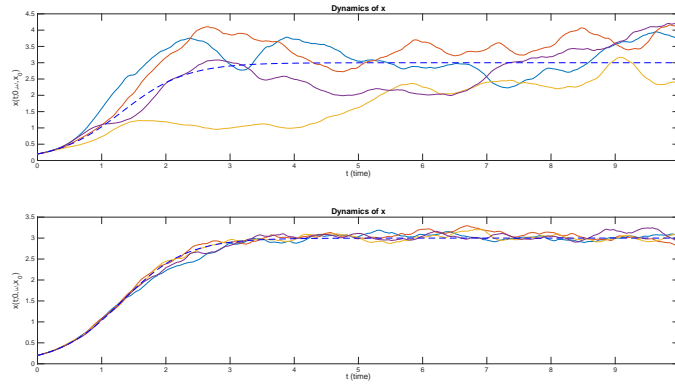


FIGURE 7. Realizations of the solution of the random logistic equation with perturbed carrying capacity for $x_0 = 0.2$. $\alpha = 2$, $\beta = 1$, $\gamma = 0.4$ (top) and $\alpha = 2.2$, $\beta = 10$, $\gamma = 0.4$ (bottom)

We can observe that all the solutions of the random equation (3.2) are fluctuating around the equilibrium of the deterministic case $x = 3$ and these fluctuations remain inside a strictly positive bounded interval which is smaller when taking large values of β and (or) smaller values of γ . Thus, the theoretical results and the advantages of the O-U process are demonstrated on this example.

In Figure 8 we present the behavior of several realizations of the solution of the random logistic equation (3.2) for the initial value $x_0 = 3$, the nominal carrying capacity is $a = 3$, the amount of noise is $\alpha = 2$ (top) and $\alpha = 2.2$ (bottom), the mean reverting constant is $\beta = 1$ (top) and $\beta = 10$ (bottom) and the volatility constant is $\gamma = 0.1$ (top) and $\gamma = 0.2$ (bottom).

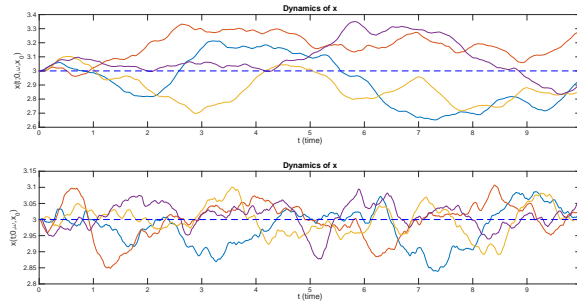


FIGURE 8. Realizations of the solution of the random logistic equation with perturbed carrying capacity for $x_0 = 3$. $\alpha = 2$, $\beta = 1$, $\gamma = 0.1$ (top) and $\alpha = 2.2$, $\beta = 10$, $\gamma = 0.2$ (bottom)

We can observe in this case a similar behavior to the previous one. However there are significant differences when comparing the behavior of the random equation (3.2) and the deterministic one for the initial condition $x_0 = 3$. In the deterministic case,

$x = 3$ is a stable equilibrium and the solution is constant as can be observed in the blue dashed lines. On the random case, one can observe some fluctuations around $x = 3$ which remain inside a strictly positive interval, this later one being deterministic, as proved in the theoretical results. In addition, this interval can be chosen by tuning the parameters of the O-U process as we previously explained.

4. PERTURBATIONS ON THE GROWTH RATE IN THE LOGISTIC EQUATION.

Now, we consider the logistic equation that we rewrite in the following form

$$(4.1) \quad \frac{dx}{dt} = rx \left(1 - \frac{x}{c}\right),$$

where $x = x(t)$ denotes the number of population of some species, r denotes the specific growth rate of the species and c is the carrying capacity of the medium assumed to be constant, both positive.

In this case we are interested in introducing a noise in the reproduction rate by using the O-U process. As a result, we have the following random logistic model

$$(4.2) \quad \frac{dx}{dt} = (r + \alpha z_{\beta, \gamma}^*(\theta_t \omega))x \left(1 - \frac{x}{c}\right),$$

where $z_{\beta, \gamma}^*(\theta_t \omega)$ denotes again the O-U process and $\alpha > 0$ is the amount of noise. We observe that $x = c$ is still an equilibrium for the equation.

As made in the previous case, the solution of equation (4.2) exists and its explicit expression is given by

$$(4.3) \quad x(t; 0, \omega, x_0) = \frac{x_0}{e^{-\int_0^t r + \alpha z_{\beta, \gamma}^*(\theta_s \omega) ds} \left(1 + \frac{x_0}{c}\right) + \frac{x_0}{c}}$$

for every $x_0 \geq 0$, any $\omega \in \Omega$ and $t \geq 0$, whence we observe the property

$$\begin{aligned} \int_0^t r + \alpha z_{\beta, \gamma}^*(\theta_s \omega) ds &= rt + \int_0^t z_{\beta, \gamma}^*(\theta_s \omega) ds \\ &= t \left(r + \frac{1}{t} \int_0^t z_{\beta, \gamma}^*(\theta_s \omega) ds \right) \end{aligned}$$

Thus, thanks to the ergodic properties in Theorem 2.1, we obtain that the dynamics of the population converges to the carrying capacity as in the deterministic case or, in other words, we have that for every $\varepsilon > 0$, any initial value $0 < x_0 < c$ and $\omega \in \Omega$, there exists some time $T(\varepsilon, \omega) > 0$ such that

$$(4.4) \quad c - \varepsilon < x(t; 0, \omega, x_0) < c$$

for all $t \geq T(\varepsilon, \omega)$.

Therefore, we have that $B_\varepsilon = [c - \varepsilon, c]$, for any $\varepsilon > 0$, is a deterministic absorbing set for the solutions of (4.2) whence we have that $B_0 = \{c\}$ defines a positive deterministic attracting set for the solutions of (4.2). As a consequence, every realization of the solution of (4.2) converge to the carrying capacity c as long as the initial value $x_0 > 0$, as in the deterministic case. This is not surprising since, in this second logistic equation, the carrying capacity is still a stable equilibrium even though we are treating a random case.

We would also like to remark that the above calculations are independent of the choice of the parameter β .

Now we perform some numerical simulations to support the results previously stated. In Figure 9 we show the behavior of several realizations of the solution of the random logistic equation (4.2) for the initial values $x_0 = 0.8$ (top), $x_0 = 1.5$ (medium) and $x_0 = 0.2$ (bottom), the growth rate is $r = 2$, the carrying capacity is $c = 1.5$, the amount of noise is $\alpha = 2$, the mean reverting constant is $\beta = 1$ and the volatility constant is $\gamma = 0.4$.

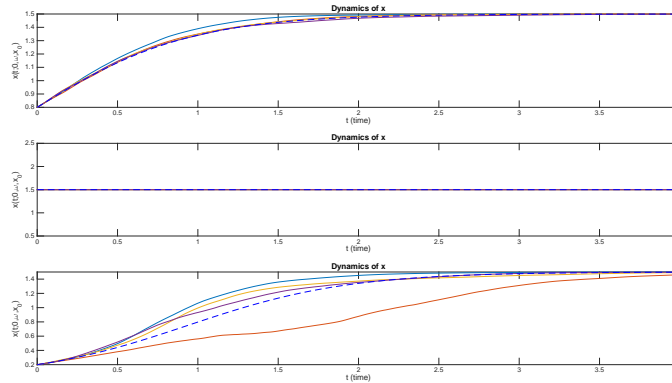


FIGURE 9. Realizations of the solution of the random logistic equation with perturbed growth rate for $x_0 = 0.8$ (top), $x_0 = 1.5$ (medium) and $x_0 = 0.2$ (bottom)

We can observe, differently to the example analyzed in Section 3, that in this case the realizations of the solution of the random equation (4.2) have fluctuations when the population is increasing but these disturbances are not present when the population is close to the carrying capacity, in fact, $x = c$ is an equilibrium of the equation (4.2) as in the deterministic case. In the second plot we can in fact see that the random solution is constant for the initial condition $x_0 = 1.5$.

In Figure 10 we can see two panels representing several realizations of the solution of the random logistic equation (4.2) where the initial values $x_0 = 0.8$ (top), $x_0 = 1.5$ (medium) and $x_0 = 0.2$ (bottom), the growth rate is $r = 2$, the carrying capacity is $c = 1.5$, the amount of noise is $\alpha = 2$, the mean reverting constant is $\beta = 10$ and the volatility constant is $\gamma = 0.4$.

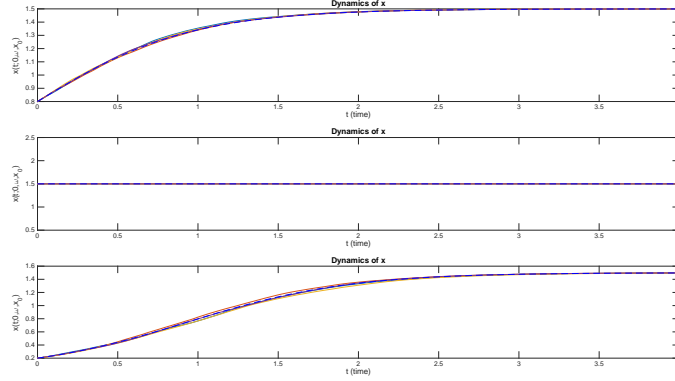


FIGURE 10. Realizations of the solution of the random logistic equation with perturbed growth rate for $x_0 = 0.8$ (top), $x_0 = 1.5$ (medium) and $x_0 = 0.2$ (bottom)

We can observe the same behavior that the one described in the previous simulations. However, the realizations of the solution of the random equation (4.2) are much closer to the deterministic ones since β is larger.

5. PARAMETER ESTIMATION IN THE LOGISTIC MODEL.

We consider again the logistic equation that we write

$$(5.1) \quad \frac{dx}{dt} = ax \left(1 - \frac{x}{K}\right),$$

where we put $K = ac$. In several ecological systems, K represents a *carrying capacity*, which is related to the size of the population when all the sites are colonized. When considering the density of the population or the proportion $p = x/K \in [0, 1]$ of occupied sites, the variable p is solution of the differential equation

$$(5.2) \quad \frac{dp}{dt} = rp(1-p)$$

where the parameter $r = aK$, usually known as the *intrinsic growth rate*, may fluctuate about a nominal value under environmental variations (season, light, temperature...). We consider then random perturbations on r , as in previous section

$$(5.3) \quad \frac{dp}{dt} = (r + \alpha z_{\beta, \gamma}^*(\theta_t \omega))p(1-p),$$

for which there exist positive numbers \underline{r}, \bar{r} with $\underline{r} < \bar{r}$, such that each realization of $r + \alpha z_{\beta, \gamma}^*(\theta_t \omega)$ belongs to the interval (\underline{r}, \bar{r}) . The question under investigation in this section is the estimation of the parameter r , measuring the proportion p over the time, in both deterministic and random frameworks.

From equation (5.2), one obtains an exact expression of r , when the dynamics of p is not at steady state

$$(5.4) \quad r = \frac{1}{t} \int_{p(0)}^{p(t)} \frac{dp}{p(1-p)} = \frac{1}{t} \log \left(\frac{p(t)(1-p(0))}{p(0)(1-p(t))} \right), \quad t > 0,$$

One can then consider

$$(5.5) \quad \hat{r}(t) = \frac{1}{t} \log \left(\frac{p(t)(1-p(0))}{p(0)(1-p(t))} \right), \quad t > 0$$

as an estimator of r in presence of random perturbation. Simulations show that this estimator behaves well as long as $p(t)$ is not too close from its limiting value 1 (see Figure 11, where two panels are presented: the first shows the realizations of the solution of (5.3) and the other overlaps the realization $r + \alpha z_{\beta, \gamma}^*(\theta_i \omega)$ (orange line), its estimator (blue line) and the dashed lines represent the values \underline{r} , r and \bar{r}). Differently to a true observer (see for instance [14] for an introduction to the theory of observers), we have no information to know when to trust this estimator.

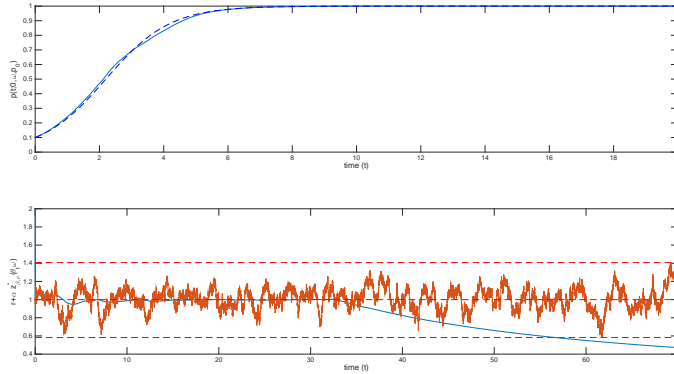


FIGURE 11. Realizations of the solution of (5.3) and estimation of r in presence of random perturbation.

We look instead for observers, i.e. dynamical systems with on-line corrective terms (see Appendix for the main ingredients which are not here to avoid technicalities). After setting up an observer in normal form, it is possible to obtain a practical convergence of the observer, in the sense that for any $\varepsilon > 0$, one can design an observer such that there exists $T > 0$ with $\hat{r}(T) \in [r - \varepsilon, r + \varepsilon]$. One cannot expect a better convergence of this observer as the system is not observable at $p = 1$ and all solutions with non null $p(0)$ converge asymptotically to this singular point. On the simulations depicted on Figure 12, one can see two different panels: the first one shows the solution of (5.3) without the presence of noise and the observer \hat{p} and the second one represents the estimator \hat{r} and r . Moreover, we can observe that the error of the observer has converged much before the system has reached the neighborhood of the steady state $p = 1$, as it is desired for a true observer. Moreover the innovation, that is the difference between the observed variable x and the variable \hat{x} (in blue) of the observer informs on the convergence of the estimator \hat{r} (when \hat{x} stays almost equal to x , we know also that \hat{r} stays close to the unknown value r).

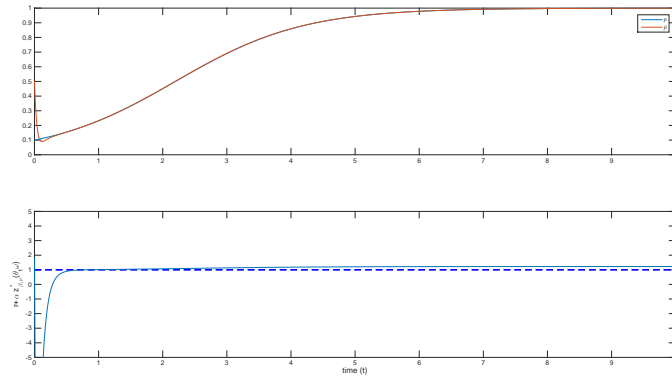


FIGURE 12. Simulation of the observer in normal form for $T = 2$ and $\theta = 15$

Finally, in Figure 13 we present simulations of this observer in presence of random perturbations. The first panel shows the solution of (5.3) in presence of noise by means of the O-U process and the second panel overlaps the realization of $r + \alpha z_{\beta,\gamma}^*(\theta_t\omega)$ (orange line) and the estimator \hat{r} (blue line). The dashed lines represent the values of \bar{r} , r and \underline{r} .

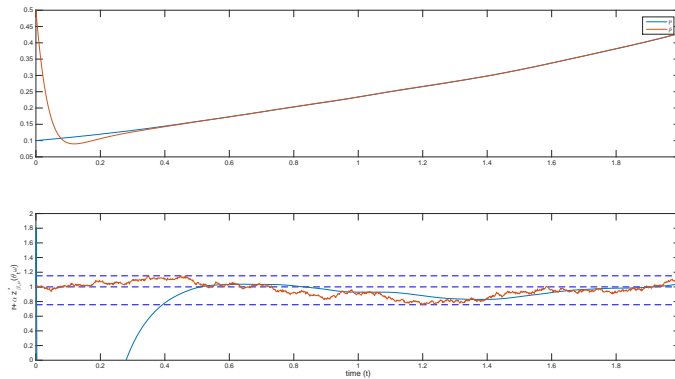


FIGURE 13. Simulation of the observer in normal form for $T = 2$ and $\theta = 15$ with random perturbation on r

6. RANDOM COMPETITIVE LOTKA-VOLTERRA MODELS

In this section we consider a competitive Lotka-Volterra model given by

$$(6.1) \quad \frac{dx}{dt} = x(\lambda - ax - by),$$

$$(6.2) \quad \frac{dy}{dt} = y(\mu - cx - ey)$$

where $x = x(t)$ is the number of population of the *first* species, $y = y(t)$ is the number of population of the *second* species, λ and μ are the specific growth rates

of each species, respectively, a and e are the carrying capacities of each species, respectively, b measures the interaction of the *first* species on the *second* one and c is the interaction that the *second* species has on the *first* ones. We assume that all the parameters in the system are positive.

Concerning the deterministic competitive system (6.1)-(6.2), it is known that coexistence of both populations can be ensured as long as conditions

$$(6.3) \quad \frac{d}{b} > \frac{\mu}{\lambda} > \frac{c}{a} \quad \text{and} \quad ad - bc > 0$$

hold true.

In this case we are interested in studying the previous system where the growth rates are affected by the O-U process. Then, we consider the random competitive model given by

$$(6.4) \quad \frac{dx}{dt} = x(\lambda + \alpha z_{\beta, \gamma}^*(\theta_t \omega) - ax - by),$$

$$(6.5) \quad \frac{dy}{dt} = y(\mu + \alpha z_{\beta, \gamma}^*(\theta_t \omega) - cx - ey)$$

where $z_{\beta, \gamma}^*(\theta_t \omega)$ denotes the O-U process introduced in Section 2 of this work and $\alpha > 0$ represents the intensity of the noise.

Thanks to a suitable choice of the parameter β , we know that $\lambda + \alpha z_{\beta, \gamma}^*(\theta_t \omega) \in [\underline{\lambda}, \bar{\lambda}]$ and $\mu + \alpha z_{\beta, \gamma}^*(\theta_t \omega) \in [\underline{\mu}, \bar{\mu}]$ for any $t \in \mathbb{R}$.

Hence, from the random competitive system (6.4)-(6.5), we can obtain the following differential inequalities for the dynamics of the population of both species involved in our model

$$(6.6) \quad \frac{dx}{dt} \leq x(\bar{\lambda} - ax),$$

$$(6.7) \quad \frac{dy}{dt} \leq y(\bar{\mu} - ey).$$

Then, we obtain that the population of the both species are bounded from above

$$(6.8) \quad x(t; 0, \omega, x_0) \leq \frac{\bar{\lambda}}{a} \quad \text{and} \quad y(t; 0, \omega, y_0) \leq \frac{\bar{\mu}}{e}$$

for any initial values $x_0 \geq 0$, $y_0 \geq 0$, any $\omega \in \Omega$ and $t \geq 0$.

In addition, from (6.6) it is possible to obtain the following differential inequality

$$(6.9) \quad \frac{dx}{dt} \geq x \left(\lambda - b \frac{\bar{\mu}}{e} - ax \right)$$

whence we can obtain its explicit solution which is given by

$$(6.10) \quad x(t; 0, \omega, x_0) \geq \frac{x_0}{e^{-(\lambda - b \frac{\bar{\mu}}{e})t} + \frac{x_0 a}{\lambda - b \frac{\bar{\mu}}{e}} \left(1 - e^{-(\lambda - b \frac{\bar{\mu}}{e})t} \right)}$$

for any initial value $x_0 \geq 0$, any $\omega \in \Omega$ and $t \geq 0$, from which, by taking limit when t goes to infinity, we obtain

$$(6.11) \quad \lim_{t \rightarrow +\infty} x(t; 0, \omega, x_0) \geq \frac{\lambda - b\bar{\mu}}{a}.$$

Then, the population of the *first* species persists as long as the condition

$$(6.12) \quad \frac{\bar{\mu}}{\lambda} < \frac{e}{b}$$

is fulfilled.

Concerning the other population, the same argument can be done and we obtain

$$(6.13) \quad \lim_{t \rightarrow +\infty} y(t; 0, \omega, y_0) \geq \frac{\mu - \bar{\lambda}}{e},$$

then the population of the *second* species persists as long as the following condition is fulfilled

$$(6.14) \quad \frac{\mu}{\lambda} > \frac{c}{a}.$$

In conclusion, for any $\varepsilon > 0$, $\omega \in \Omega$ and every initial values $x_0 \geq 0$ and $y_0 \geq 0$, there exists some time $T(\varepsilon, \omega) > 0$ such that the solution of the random system (6.6)-(6.7) can be bounded inside the frame

$$(6.15) \quad B_\varepsilon = \left[\frac{\lambda - b\bar{\mu}}{a} - \varepsilon, \frac{\bar{\lambda}}{a} \right] \times \left[\frac{\mu - \bar{\lambda}}{e} - \varepsilon, \frac{\bar{\mu}}{e} \right]$$

for every $t \geq T(\varepsilon, \omega)$.

Therefore, for any $\varepsilon > 0$, B_ε is a strictly positive deterministic absorbing set for the solutions of the system (6.6)-(6.7), whence we have that

$$(6.16) \quad B_0 = \left[\frac{\lambda - b\bar{\mu}}{a}, \frac{\bar{\lambda}}{a} \right] \times \left[\frac{\mu - \bar{\lambda}}{e}, \frac{\bar{\mu}}{e} \right]$$

is a strictly positive deterministic attracting set for the solutions of the system (6.6)-(6.7).

From the previous analysis, we can observe that, as long as conditions (6.12) and (6.14) are satisfied, we can ensure the coexistence of the population of both species.

Now, we present some numerical simulations to illustrate the results provided in this section. In Figure 14 we present the phase plane with several realizations of the solutions of the random competitive system (6.6)-(6.7) for the initial values $x_0 = 4$ and $y_0 = 3$ and the following values of the rest of the parameters $a = 20$, $b = 2$, $c = 4$, $e = 314$, $\lambda = 5$, $\mu = 7$, the amount of noise is $\alpha = 2$, the mean reverting constant is $\beta = 1$ and the volatility constant is $\gamma = 0.5$. We remark that the right panel shows a zoom of the left one to see the absorbing set of the solutions, which is the box delimited by the dashed lines.

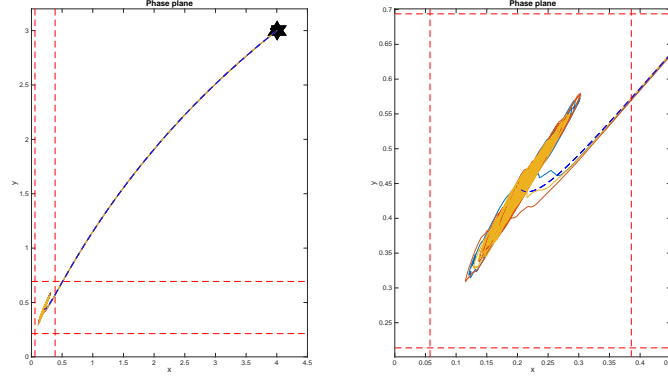


FIGURE 14. Phase plane with realizations of the solutions of the competitive Lotka-Volterra system for $x_0 = 4$ and $y_0 = 3$

We can observe that all the realizations of the solution of the system remain, after some time, inside a rectangle limited by the red dashed lines. This rectangle is the absorbing set B_0 (see (6.16)) obtained in the mathematical results which is deterministic.

In Figure 15 we present the dynamics of both species individually where the red dashed lines represent the bounds guaranteed for the corresponding state variables. We can observe that both species are fluctuating around the deterministic solution inside a strictly positive interval that allows us to guarantee the persistence of both species. In addition, these intervals are deterministic in the sense that they do not depend on the realization on the noise and can be chosen as explained previously.

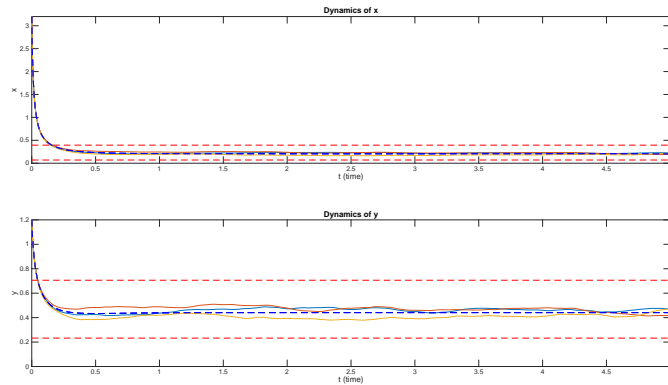


FIGURE 15. Realizations of the solution of the competitive Lotka-Volterra system (both state variables depending on time) for $x_0 = 4$ and $y_0 = 3$

7. RANDOM CHEMOSTAT MODEL

In this section we would like to remark that the O-U process has already also provided a very useful tool when perturbing the input flow in the chemostat model. Let us recall the classical “resource-consumer” model (see for instance [15, 16]):

$$(7.1) \quad \frac{ds}{dt} = D(t)(s_{in} - s) - \frac{1}{Y}\mu(s)x,$$

$$(7.2) \quad \frac{dx}{dt} = \mu(s)x - D(t)x,$$

where $x(t)$ and $s(t)$ denote respectively the consumers and resource densities at time t . The function μ is the specific growth rate of the consumers over the resource (Y is a conversion coefficient). When the system is continuously fed, as in the chemostat device or in ecological situations such as mountain lakes, the resource is brought at a concentration s_{in} and diluted with a dilution rate D . Most often, D is subject to random disturbances but realizations stay bounded.

Every detail about the way of modeling and a complete analysis of the resulting random model can be found in [17, 18] thus we will omit the details in this section. Instead we just give some remarks concerning the work in [17, 18].

As already explained in the introductory section, several drawbacks can be found when perturbing the input flow of the chemostat model by using the standard Wiener process (see [6, 8]). For instance, the input flow could take extremely large values and thus could be negative. Due to this fact, which is unrealistic from the biological point of view since we know that the input flow is fluctuating in a positive bounded interval, we also have that some state variables that describe population size could take negative values which is also unrealistic from the biological point of view. In addition, it is not possible to ensure the persistence of the species as we did and which corresponds to real observations.

However, everything these drawbacks are circumvented when introducing the perturbations on the input flow by means of the O-U process as explained in this paper. The first important improvement is the fact that the perturbed input flow is ensured to be bounded, as in real experiments (see Figure 3 where we presented the real data). In addition, it is possible to prove that there exists absorbing and attracting sets which are deterministic (then they do not depend on the realization of the noise) and moreover, what is essential to prove the persistence of the species, positive. Furthermore, these results are proved in forward sense which suits the point of view of applications..

8. CONCLUSIONS AND FINAL COMMENTS

In this final section we would like to draw some conclusions and final comments concerning the O-U process introduced in Section 2. We recall that the most important improvement of this way of modeling the noise, compared to other kinds of noise considered previously in the literature is the fact that the O-U process depends of two parameters, the volatility constant γ and the mean reversion constant β which play an important role and allow the noise to have all the expected properties that we have formulated in the introduction.

In order to show the relevance of this new way of modeling we have presented in the previous sections some examples which illustrate the effect of this bounded noise when perturbing some very well-known models such as the logistic growth or the Lotka-Volterra competition. In addition, in [17, 18] the authors consider this noise to model random input flows in the chemostat model where some relevant improvements are also achieved. Finally, this way of modeling is full of advantages from the mathematical analysis point of view but also, which is essential, as a quite realistic modeling from the biological point of view.

In conclusion, we believe that this modeling approach is generic and could be applied in most of the population models, when some aspects or parameters are expected to be subject to randomness with bounded realizations. For instance, it could be very interesting to analyze prey-predator models where, in the deterministic case periodic orbits and limit cycles are present. In this way, it could be possible to define a concept of *random* periodic orbits in the sense that the solutions of the system are fluctuating around the deterministic periodic orbit inside some interval that depends on the parameters of the O-U process, as in the examples of the present paper. Another idea is to analyze the problem of the observer with measurements perturbed by the O-U process. These are some ideas among other ones to carry on the applications of this way of modeling noise.

APPENDIX

Let us give more details about the observer construction used in Section 5. For simplicity, let us first consider by the deterministic framework. For this purpose, we consider the extended dynamics

$$(1) \quad \frac{dp}{dt} = rp(1-p)$$

$$(2) \quad \frac{dr}{dt} = 0$$

with the measured output

$$y(t) = p(t)$$

Notice this system is not observable (see [14] for the definition of observability) at the steady states $p = 0$ or $p = 1$. When the system is not at equilibrium, let us first consider a classical observer of Luenberger form

$$\frac{d\hat{p}}{dt} = \hat{r}y(t)(1-y(t)) + G_1(\hat{p} - y(t))$$

$$\frac{d\hat{r}}{dt} = G_2(\hat{p} - y(t))$$

where the gains parameters G_1, G_2 have to be chosen. The dynamics of the error variables $e_p = \hat{p} - p$, $e_r = \hat{r} - r$ are given by the linear non-autonomous system

$$\frac{de_p}{dt} = G_1e_p + y(t)(1-y(t))e_r$$

$$\frac{de_r}{dt} = G_2e_p$$

Consider then the quadratic function

$$(3) \quad V(e_p, e_r) = \frac{1}{2}(e_p + \gamma e_r)^2 + \frac{1}{2}e_r^2$$

where γ is a parameter. Notice that V is definite positive for any value of γ . One has, along any trajectory:

$$\begin{aligned}\frac{dV}{dt} &= (e_p + \gamma e_r)(G_1 e_p + y(t)(1 - y(t))e_r + \gamma G_2 e_p) + e_r G_2 e_p \\ &= (G_r + \gamma G_2)e_p^2 + \gamma y(t)(1 - y(t))e_r^2 + (\gamma G_1 + \gamma^2 G_2 + G_2)e_p e_r\end{aligned}$$

Take $G_1 < 0$ and $\gamma < 0$ and set

$$G_2 = -\frac{\gamma}{1 + \gamma^2}G_1$$

Notice that for such choice, one has

$$\gamma_p := G_1 + \gamma G_2 = \frac{G_1}{1 + \gamma^2} < 0$$

Equivalently, G_1 and G_2 are defined as

$$\begin{aligned}G_1 &= (1 + \gamma^2)\gamma_p \\ G_2 &= -\gamma\gamma_p\end{aligned}$$

with γ_p and γ negative. For $V > 0$, one has the inequality

$$\frac{dV}{dt} = \gamma_p e_p^2 + \gamma y(t)(1 - y(t))e_r^2 < 0, \quad \forall t > 0.$$

However, we cannot conclude about the convergence of V to 0 because

$$\int_0^{+\infty} y(t)(1 - y(t))dt = \frac{1 - p(0)}{r} < +\infty$$

as this is shown on Figure 16.

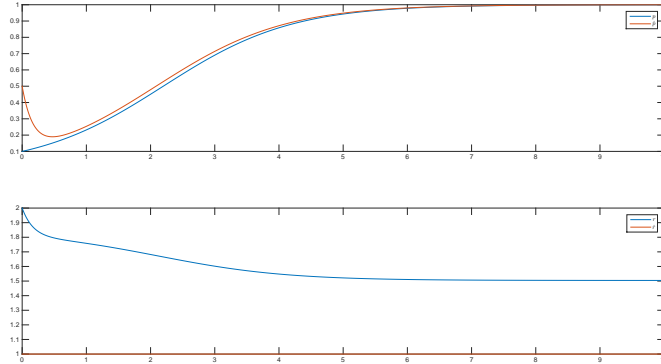


FIGURE 16. Simulation of the Luenberger observer with $\gamma = -5$ and $\gamma_p = -1$

Consider now a second kind of observer, but in *normal form* (see [14, 19]), which consists in applying the change of coordinates $(p, r) \rightarrow (z_1, z_2)$ with

$$\begin{aligned}z_1 &= p \\ z_2 &= rp(1 - p)\end{aligned}$$

When p is different from the steady states $p = 0$ and $p = 1$, parameter r can be reconstructed as

$$r = \varphi(z_1, z_2) := \frac{z_2}{z_1(1 - z_1)}$$

Dynamics (.1)-(2) in these coordinates writes as follows

$$\frac{dz}{dt} = \underbrace{\begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}}_A z + \begin{bmatrix} 0 \\ \psi(z_1, z_2) \end{bmatrix} \quad \text{with} \quad \psi(z_1, z_2) := \varphi(z_1, z_2) \left(1 - \frac{z_1}{2}\right) z_2$$

with the observation

$$y(t) = \underbrace{\begin{bmatrix} 0 & 1 \end{bmatrix}}_C z(t)$$

This leads to consider the following observer

$$\begin{aligned} \frac{d\hat{z}_1}{dt} &= \hat{z}_2 + G_1(\hat{z}_1 - y(t)) \\ \frac{d\hat{z}_2}{dt} &= \psi(y(t), \hat{z}_2) + G_2(\hat{z}_1 - y(t)) \end{aligned}$$

with the estimator

$$\hat{r}(t) = \varphi(y(t), \hat{z}_2(t))$$

Notice that when the system is not at steady state, $\varphi(y(t), \hat{z}_2(t))$ and $\psi(y(t), \hat{z}_2(t))$ are well defined for any $t > 0$. The dynamics of the error $e = \hat{z} - z$ is given by the system

$$\frac{de}{dt} = \begin{bmatrix} G_1 & 1 \\ G_2 & 0 \end{bmatrix} e + \begin{bmatrix} 0 \\ 1 \end{bmatrix} (\psi(y(t) - \hat{z}_2) - \psi(y(t) - z_2))$$

The map $z_2 \mapsto \psi(y(t), z_2)$ is not Lipschitz with respect to z_2 uniformly w.r.t. t . However for any fixed T , it is Lipschitz on any compact set uniformly on $[0, T]$. We can then use the theory of *high-gains observers* [14, 19], which guarantees an exponential decrease of the norm of the error on $[0, T]$, when the gains G_1, G_2 are chosen such that

$$\begin{bmatrix} G_1 \\ G_2 \end{bmatrix} = -S_\theta^{-1} C^\top$$

where S_θ is the symmetric definite positive matrix solution of the Lyapunov equation

$$A^\top S_\theta + S_\theta A - C^\top C + \theta S_\theta = 0$$

and parameter $\theta > 0$ is large enough. One can check that this gives

$$\begin{bmatrix} G_1 \\ G_2 \end{bmatrix} = \begin{bmatrix} -2\theta \\ -\theta^2 \end{bmatrix}$$

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Corrigendum to the paper: A way to model stochastic perturbations in population dynamics models with bounded realizations

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ABSTRACT. In this corrigendum we correct an error in our paper [T. Caraballo, R. Colucci, J. López-de-la-Cruz and A. Rapaport. A way to model stochastic perturbations in population dynamics models with bounded realizations, *Commun Nonlinear Sci Numer Simulat*, **77** (2019) 239–257]. We present the correct way to model real noisy perturbations by considering a slightly different stochastic process based, as in the original paper, on the Ornstein-Uhlenbeck process. Once this modification is made, every result and every example in the initial work remains valid, in fact, the same goal is achieved. The main fact is that we used in our proofs a result in [S. Al-Azzawi, J. Liu and X. Liu, Convergence rate of synchronization of systems with additive noise, *Discrete Contin Dyn Syst Ser B*, **22** (2017) 227–245] which is not true.

1. INTRODUCTION

In our paper [3] we presented a new way to model real noisy perturbations by making use of the well-known Ornstein-Uhlenbeck (OU) process. In addition, we illustrated this idea with different examples coming from population dynamics and several numerical simulations.

Let us describe firstly the explanations given in [3] in order to let the reader understand the error in [3] and the changes needed to preserve the spirit and the goal of our original work.

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It is well known that most real phenomena are subject to suffer the effects of some kind of noise. Even though there are many works in the literature leading with stochastic models, they mainly focus on making use of the standard Wiener process, in spite of the fact that it is unbounded with probability one. This last property is completely unrealistic from the point of view of applications, since random perturbations in real life are bounded. In order to understand better the drawback caused by the standard Wiener process when modeling noise, we refer the readers to [4–6], where the authors investigate chemostat models by means of the Brownian motion.

Motivated by this fact, we presented in [3] a way to model bounded real perturbations involving the OU process, which has proved to fit in a very loyal way the random disturbances observed in real life.

More precisely, we explained that we could consider perturbations of some parameter in a random system of the following kind

$$(1) \quad \dot{x} = f(x, z_{\beta, \gamma}^*(\theta_t \omega)),$$

where $z_{\beta, \gamma}^*(\theta_t \omega)$ denotes the OU process. This type of systems arises when replacing, for instance, some deterministic parameter, namely a , by $a + \alpha z_{\beta, \gamma}^*(\theta_t \omega)$, where $\alpha > 0$ denotes the intensity of the noise.

In order to deal with system (1), we introduced in [3] the OU process which is defined by the random variable

$$(2) \quad z_{\beta, \gamma}^*(\theta_t \omega) = -\beta\gamma \int_{-\infty}^0 e^{\beta s} \theta_t \omega(s) ds, \quad t \in \mathbb{R}, \omega \in \Omega, \beta, \gamma > 0,$$

and solves the Langevin equation

$$(3) \quad dz + \beta z dt = \gamma d\omega,$$

and then we provided some properties, as the ones in Proposition 2 in [3] that we recall below.

Proposition 1.1. *There exists a θ_t -invariant set $\tilde{\Omega} \in \mathcal{F}$ of Ω of full \mathbb{P} -measure such that for $\omega \in \tilde{\Omega}$ and $\beta, \gamma > 0$, we have*

(i) *the random variable $|z_{\beta, \gamma}^*(\omega)|$ is tempered.*

(ii) *the mapping*

$$(t, \omega) \rightarrow z_{\beta, \gamma}^*(\theta_t \omega) = -\beta\gamma \int_{-\infty}^0 e^{\beta s} \omega(t+s) ds + \omega(t)$$

is a stationary solution of (2) with continuous trajectories;

(iii) *for any $\omega \in \tilde{\Omega}$ one has*

$$\begin{aligned} \lim_{t \rightarrow \pm\infty} \frac{|z_{\beta, \gamma}^*(\theta_t \omega)|}{t} &= 0; \\ \lim_{t \rightarrow \pm\infty} \frac{1}{t} \int_0^t z_{\beta, \gamma}^*(\theta_s \omega) ds &= 0; \\ \lim_{t \rightarrow \pm\infty} \frac{1}{t} \int_0^t |z_{\beta, \gamma}^*(\theta_s \omega)| ds &= \mathbb{E}[z_{\beta, \gamma}^*] < \infty; \end{aligned}$$

(iv) finally, for any $\omega \in \tilde{\Omega}$,

$$\lim_{\beta \rightarrow \infty} z_{\beta, \gamma}^*(\theta_t \omega) = 0, \quad \text{for all } t \in \mathbb{R}.$$

Hence, the idea was the following: (1) we are given an interval, namely $[a_1, a_2]$, (2) we fix any event $\omega \in \tilde{\Omega}$ and (3) for every fixed $\omega \in \tilde{\Omega}$, thanks to Proposition 1.1 (iv), which can be found in [1], Lemma 4.1, we can choose β large enough such that $a + \alpha z_{\beta, \gamma}^*(\theta_t \omega) \in [a_1, a_2]$ for all $t \in \mathbb{R}$.

Then, Proposition 1.1 (iv) was the key of our paper since it allowed us to ensure that the OU process $z_{\beta, \gamma}^*(\theta_t \omega)$ was bounded in some desired interval (typically fixed by practitioners) for every fixed event $\omega \in \tilde{\Omega}$ and β large enough.

As we explained in [3], since β could depend on ω , the resulting random system does not generate a random dynamical system (RDS) and then the theory of RDSs and pullback attractors could not be applied. But this was not a problem at all, since the random system could be analyzed for every fixed event ω . In fact, this allowed us to prove every result of the paper forwards in time, which is a more natural way than the pullback one.

Nevertheless, we noticed that property (iv) in Proposition 1.1 (Proposition 2 (iv) in [3]) is not true. By having a deeper look at the proof of the result in [1] we realized that such a result cannot be proved uniformly in time (as it is asserted in [1]), but for every fixed t .

After this, we found a slightly different way to model bounded noises which allows us to prove every result and present every example in the original paper [3] in a very similar manner. Let us present now the new correct idea.

Assume that we want to consider some random perturbations in a certain parameter a of a deterministic differential system. In addition, as in real life, those disturbances must be bounded (for all the time) in some interval around a , namely $[a - d, a + d]$, where $d > 0$.

Instead of replacing a by $a + \alpha z_{\beta, \gamma}^*(\theta_t \omega)$, as we did in [3], we replace a by $a + \Phi(z^*(\theta_t \omega))$, where $\Phi(z) = \frac{2d}{\pi} \arctan(z)$ and $z^*(\theta_t \omega)$ denotes the OU process (2). In this case, β and γ are fixed through the whole work. Hence, instead of having to deal with system (1), the resulting random system is of the kind

$$(4) \quad \dot{x} = f(x, \Phi(z^*(\theta_t \omega))),$$

and it will be investigated for every fixed event $\omega \in \Omega$, as in the original work [3].

We notice that, in this case, as β and γ are fixed, the resulting random system (4) generates a RDS and then the theory of RDSs and pullback attractors could be applied (which was not the case in the original paper). Nevertheless, instead of using that theory and pullback attraction, we study the random system (4) for every fixed $\omega \in \Omega$ since it allows us to prove every result forwards in time (which is, as explained before, more natural than the pullback one).

It is easy to check that the random perturbations $a + \Phi(z^*(\theta_t \omega))$ are bounded for every $t \in \mathbb{R}$ and any $\omega \in \Omega$, since Φ is bounded. In fact, for every $\omega \in \Omega$, we have

$$a - d \leq a + \Phi(z^*(\theta_t \omega)) \leq a + d, \quad \text{for all } t \in \mathbb{R}.$$

In addition, we can prove the following property.

Proposition 1.2. *Let $\Phi(z) = \frac{2d}{\pi} \arctan(z)$. Then*

$$(5) \quad \lim_{t \rightarrow +\infty} \frac{1}{t} \int_0^t \Phi(z^*(\theta_s \omega)) ds = 0, \quad \text{a.s. in } \Omega.$$

Proof. Since $\Phi(z) = \frac{2d}{\pi} \arctan(z)$, then we have

$$\int_{\Omega} \left| \frac{2d}{\pi} \arctan(z^*(\omega)) \right| d\mathbb{P}(\omega) \leq d|\Omega| = d$$

whence $\Phi(z^*(\cdot)) \in L^1(\Omega, \mathcal{F}, \mathbb{P})$. Then, since \mathbb{P} is invariant by θ (see [2, 7]), from the Birkhoff ergodic theorem we deduce that

$$\lim_{t \rightarrow +\infty} \frac{1}{t} \int_0^t \Phi(z^*(\theta_s \omega)) ds = \mathbb{E}[\Phi(z^*(\omega))], \quad \text{a.s. in } \Omega.$$

Hence, it is enough to show that $\mathbb{E}[\Phi(z^*(\omega))] = 0$. In fact, we have

$$\mathbb{E}[\Phi(z^*(\omega))] = \int_{\mathbb{R}} \Phi(x) f_{\text{OU}}(x) dx = 0,$$

where f_{OU} denotes the density function of the random variable $z^*(\cdot)$, which is an even function since it is Gaussian, and Φ is an odd function. \square

Hence, the stochastic process $\Phi(z^*(\omega))$ satisfies the desired ergodic property (which was false in Proposition 1.1 (iv)) and it is bounded. Then, every result in our original work [3] is true when considering this little modification and every example there can be given to illustrate this new way to model bounded random perturbations.

Concerning the numerical simulations, they can be done again by considering the new function $\Phi(z) = \frac{2d}{\pi} \arctan(z)$ and they are analogous to the ones in the original work [3].

To finish, we would like to remark that we could also use any odd measurable function Φ such that

$$\lim_{z \rightarrow +\infty} \Phi(z) = d < +\infty.$$

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