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PAINTER: a spatio-spectral image reconstruction algorithm for optical interferometry

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Astronomical optical interferometers sample the Fourier transform of the intensity distribution of a source at the observation wavelength. Because of rapid perturbations caused by atmospheric turbulence, the phases of the complex Fourier samples (visibilities) cannot be directly exploited. Consequently, specific image reconstruction methods have been devised in the last few decades. Modern polychromatic optical interferometric instruments are now paving the way to multiwavelength imaging. This paper is devoted to the derivation of a spatio-spectral (“3D”) image reconstruction algorithm, coined PAINTER (Polychromatic optiAl INTERferometric Reconstruction software). The algorithm relies on an iterative process, which alternates estimation of polychromatic images and of complex visibilities. The complex visibilities are not only estimated from squared moduli and closure phases, but also differential phases, which helps to better constrain the polychromatic reconstruction. Simulations on synthetic data illustrate the efficiency of the algorithm and in particular the relevance of injecting a differential phases model in the reconstruction.

I. INTRODUCTION

Current astronomical Optical Interferometers (OI) do not directly provide images, even if prospective studies exist in this direction [1–3]. Instead, OI data are related to the observed celestial scene at the observation wavelength by a two dimensional Fourier transform (FT) of the corresponding intensity distribution perpendicular to the line of sight. Ideally, the observables measured by an interferometer are the so-called *complex visibilities*, which corresponds to complex samples of the Fourier spectrum. The sampling function is fully defined by the positions of the interfering telescopes and by the observation wavelength. Earth rotation provides additional samples for observations acquired at different epochs, as the telescopes’ positions are modified in time with respect to the line of sight (an effect called super synthesis). However, owing to the small number of telescopes involved (typically three or four), the sampling of the Fourier space that is achieved by OI is always sparse.

In essence, measuring moduli and phases of complex visibilities in OI amounts to measuring contrasts and phases of interference fringes [4]. Atmospheric turbulence randomly shifts these fringes on a timescale of 10 ms. The moduli are thus obtained by estimating fringe contrasts in snapshot mode with short integration times that freeze the atmospheric turbulence. As for the phases, the unknown random phase shifts imply that optical interferometers cannot measure directly the phases (in absence of an artificial or real star, which would provide such a reference [5]). Astronomers thus use a turbulence independent observable related to the phases called *closure*

phase, originally devised for radiointerferometers [6], and from which the phase information can be partially extracted [7–9]. Note that the current situation in optical is however very different from radio, where the numbers of antennas or dipoles is several orders of magnitude greater than in optical (see e.g. [10] for a recent example) and all phases can be estimated.

From an informational viewpoint, the image reconstruction problem posed by OI is highly underdetermined because of the sparse Fourier sampling and of additional missing phase information. This means that formally, infinitely many intensity distributions exist that are consistent with OI data. In this framework, a classical and well-understood strategy for image reconstruction is to adopt an inverse problem approach, where missing information is mitigated, and hopefully compensated for, by *a priori* knowledge [11]. In this case, the image reconstruction algorithm aims at finding an intensity distribution that minimizes a cost function composed of a data fidelity term, which is related to the noise distribution, plus a regularization term and possibly other constraints, which are related to prior knowledge.

Following this path, various algorithms have blossomed in the last twenty years. Most of the proposed algorithms rely on gradient descent methods (*WISARD* [12, 13], *BSMEM* [14], *MiRA* [15], *BBM* [16], *IRBis* [17]). A different approach is used in *MACIM* [18] and in its evolution *SQUEEZE* [19], which rely on Markov Chain Monte Carlo (MCMC) method.

While these algorithms have proved very useful to astronomers, none of them is currently able to tackle polychromatic reconstruction for general sources having variations in their intensity distribution in wavelength. As such, they are monochromatic imagers and are not applicable for multiwavelength image reconstruction.

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Nowadays, modern OI is however polychromatic (see for instance AMBER [20], PIONIER [21], or VEGA [22]) and more powerful polychromatic instruments are in development (like MATISSE [23] and GRAVITY [24]). In such devices, interference fringes are simultaneously recorded in a number of wavelength channels that can reach several hundreds. Indeed, the fundamental justification of polychromatic observations is that the amount and distribution of electromagnetic radiation emitted by astrophysical sources may be very different from one wavelength channel to another, as light-matter interactions are highly variable in wavelength (emission and absorption lines for instance). Multiwavelength OI thus constitutes an extremely rich evolution over monochromatic OI. In order to fully exploit these instruments, the multiwavelength evolution of the mature but monochromatic OI reconstruction algorithms is mandatory.

Building on the monochromatic approaches mentioned above, some first steps have recently been undertaken in the direction of multiwavelength reconstruction. The work [25] implements a semi-parametric algorithm for the image reconstruction of chromatic objects, dedicated to the case of a central object surrounded by an extended structure such as a young star. The approach of [26, 27] uses a sparsity-regularized approach, dedicated to the case where the observed scene is a collection of point-like sources. [28] uses MiRA as a key optimization engine combined to a Self Calibration approach and demonstrates the potential of using the differential phases in an image-reconstruction process.

This paper is devoted to the derivation of a multiwavelength or spatio-spectral images reconstruction (3D) algorithm named PAINTER (for Polychromatic optiCAI INTerferometric Reconstruction software). This approach uses the absolute visibilities and closure phases, which are considered dependent of the wavelength. In addition, we also use the so-called *differential phases*, which are defined as the phases relatively to a reference channel and constitute an additional turbulence independent observable of the phases in multiwavelength observation mode.

The algorithm relies on the alternate estimation of the complex visibilities (from estimated phases and observed noisy moduli) and of the polychromatic intensity distribution (using spatio-spectral regularization and constraints). From a modeling viewpoint, the main originality of the approach is to estimate the unknown phases from both closure phases and differential phases. From an optimization viewpoint, the algorithm is based on ADMM methodology [29]. PAINTER can be seen as an evolution of the *MiRA-3D* algorithm proposed in [26, 27]. An implementation of PAINTER in matlab (octave compatible) with input data in OI-FITS format [30], is available at <https://www-n.oca.eu/aferrari/painter>.

The paper is organized as follows: sections II and III introduce notations and data modeling. We derive here an extended model of phase differences which is specific

to the 3D reconstruction. Section IV tackles the inverse problem approach. We introduce assumptions related to the data fidelity criterion (noise perturbations) as well as prior knowledge in the form of regularizations and constraints. Section V derives the resulting 3D image reconstruction algorithm. Performances of the algorithm are provided and analyzed in section VI.

II. NOTATIONS

$\angle \cdot$	phases of complex numbers (may result in a scalar, a vector or a matrix)
$ \cdot $	moduli of complex numbers (may result in a scalar, a vector or a matrix)
$*$	complex conjugate
$\top$	transpose
$\hat{\cdot}$	Fourier transform
H	complex conjugate transpose
$\mathbf{X}$	matrix
$\mathbf{x}$	vector, with components $\mathbf{x}_n$
$\otimes$	Kronecker product (tensor product)
$\oplus_n^N$	direct sum: $\oplus_n^N \mathbf{X}_n = \text{block diag}(\mathbf{X}_1, \dots, \mathbf{X}_N)$ i.e. $\oplus_{n=1}^2 \mathbf{X}_n = \mathbf{X}_1 \oplus \mathbf{X}_2 = \begin{bmatrix} \mathbf{X}_1 & 0 \\ 0 & \mathbf{X}_2 \end{bmatrix}$
$\odot$	Hadamard product (dot product)
$\mathbf{I}_N$	identity matrix of size $N \times N$
$\mathbf{1}_N$	vector $[1, \dots, 1]^\top$ of length N
$\mathbb{1}_{\mathbb{R}^+}$	indicator function on positive orthant
$\text{tr}(\mathbf{X})$	trace of $\mathbf{X}$
$\text{vec } \mathbf{X}$	matrix vectorization i.e the columns of $\mathbf{X}$ are stacked into one column vector $\mathbf{x}$
$\text{diag}(\mathbf{x})$	diagonal matrix with $\mathbf{x}$ on its diagonal
$\ \mathbf{X}\ _F^2$	squared Frobenius norm
$\ \mathbf{x}\ _2$	ℓ_2 norm
$\mathbf{x}^+$	update of $\mathbf{x}$

III. DATA MODELING

A. Spatio-spectral model

In the absence of atmospheric turbulence the observable measured by an interferometer is the complex visibility [11]. This observable is measured from the fringe pattern obtained by the interference of two beams collected from a pair of telescopes. The spatial position of each such pair defines one of the $N_{\mathbf{b}}$ baselines of the telescope array. Hereafter, the notation for the baseline $\mathbf{b}_{a,b}$ refers to the position vector of a telescope pair (a, b) projected on a plane perpendicular to the line of sight.

In the considered case of polychromatic observations, an astrophysical source is described by an intensity distribution which is a function of wavelength. Because OI instrument always have limited fields of view, we assume that the distribution of interest accounts for the limited spatial response of the interferometer: it is an

apodized version of the intensity distribution around the line of sight. The unknown distribution can be written as $I(\delta, \tau, \lambda) = I_\lambda(\delta, \tau)$, which is a flux density at angular position $\theta = [\delta, \tau]^\top$ of the sky and wavelength λ .

In absence of any perturbation and for purely monochromatic observations, a telescope pair of baseline $\mathbf{b}$ provides a complex visibility defined by $y^\lambda = \hat{I}(\frac{\mathbf{b}}{\lambda})$. Because OI instruments have limited angular and spectral resolutions (respectively set by the maximum distance between two telescopes and by the bandpass of the optical filters), a simple way to represent the unknown spatio-spectral distribution of the sources is to discretize $I(\delta, \tau, \lambda)$ over *voxels*. We consider here for simplicity the same discretization $\Delta\delta = \Delta\tau = \Delta\theta$ in both angular coordinates, resulting in $N_x \times N_x$ parameters per image, and an instrument with N_λ wavelength channels of equal bandwidth $\Delta\lambda$, which is set to the spectral resolution. In this case, all voxels have the same size $\Delta\theta \times \Delta\theta \times \Delta\lambda$. If we further assume a unit transfer function in all channels, one voxel is simply defined as:

$$\mathbf{x}_i^{\lambda_n} = \int_{\delta_l}^{\delta_l + \Delta\theta} \int_{\tau_m}^{\tau_m + \Delta\theta} \int_{\lambda_n}^{\lambda_n + \Delta\lambda} I(\delta, \tau, \lambda) d\theta d\lambda \quad (1)$$

where i refers to a pixel at angular position $[\delta_l, \tau_m]$ and λ_n to the reference wavelength of channel n . The column vector $\mathbf{x}^{\lambda_n}$ collects all voxels and can be organized as the vectorization of a $(N_x \times N_y)$ image of the astrophysical source at wavelength λ_n . In this setting, the goal of the multiwavelength reconstruction algorithm is to estimate the voxels, which represent the unknown parameters of the model.

Let $y_{a,b}^{\lambda_n}$ be the complex visibility at the spatial frequency $\mathbf{b}/\lambda_n$, and let $\mathbf{y}^{\lambda_n}$ be the column vector collecting the set of complex visibilities corresponding to all available baselines at wavelength λ_n . The complex visibilities can then be related in matrix form to the parameters by the direct model [27, 31]

$$\mathbf{y}^{\lambda_n} = \mathbf{F}^{\lambda_n} \mathbf{x}^{\lambda_n} \quad (2)$$

where $\mathbf{F}^{\lambda_n}$ is obtained from a Non Uniform Discrete Fourier Transform (NuDFT) [32] at the spatial frequencies imposed by the geometry of the telescope array and by the observation wavelength λ_n .

The previous expression describes the complex visibilities by wavelength. A compact notation including all wavelengths and baselines is

$$\begin{aligned} \mathbf{y} &= \mathbf{F} \mathbf{x}, \quad \mathbf{F} = \bigoplus_{n=1}^{N_\lambda} \mathbf{F}^{\lambda_n} \\ \mathbf{x} &= [\mathbf{x}^{\lambda_1 \top}, \dots, \mathbf{x}^{\lambda_{N_\lambda} \top}]^\top \end{aligned} \quad (3)$$

where $\mathbf{F}$ is a block diagonal matrix with each block referring to the NuDFT at a particular wavelength. Vector $\mathbf{y}$ concatenates the complex visibility vectors ($\mathbf{y}^{\lambda_n}$ of Eq. 2) for all wavelengths into a $N_{\mathbf{b}} N_\lambda \times 1$ visibility vector, with

associated moduli γ and phases φ given by

$$\mathbf{y} = [\mathbf{y}^{\lambda_1 \top}, \dots, \mathbf{y}^{\lambda_{N_\lambda} \top}]^\top \quad (4)$$

$$\gamma = [\gamma^{\lambda_1 \top}, \dots, \gamma^{\lambda_{N_\lambda} \top}]^\top = |\mathbf{y}| \quad (5)$$

$$\varphi = [\varphi^{\lambda_1 \top}, \dots, \varphi^{\lambda_{N_\lambda} \top}]^\top = \angle \mathbf{y} \quad (6)$$

In order to analyze the chromatic variation of the visibilities $\mathbf{y}_m^{\lambda_n}$ and of the images $\mathbf{x}^{\lambda_n}$ over the N_λ wavelengths, we also need to introduce the $N_{\mathbf{b}} \times N_\lambda$ matrix $\mathbf{Y}$ and the $N_x^2 \times N_\lambda$ matrix $\mathbf{X}$ defined by

$$\mathbf{Y} = [\mathbf{y}^{\lambda_1}, \dots, \mathbf{y}^{\lambda_{N_\lambda}}] = \text{vec}^{-1} \mathbf{y} \quad (7)$$

$$\mathbf{X} = [\mathbf{x}^{\lambda_1}, \dots, \mathbf{x}^{\lambda_{N_\lambda}}] = \text{vec}^{-1} \mathbf{x}. \quad (8)$$

To clarify the use of a matrix notation note that the n^{th} column of $\mathbf{X}$ corresponds to the vectorization of the image at the wavelength λ_n while the p^{th} line is for the variation of the pixel p along the wavelengths.

B. Model of phase differences

In the presence of atmospheric turbulence, the beams received at each telescope are affected by random and different optical path differences, which corrupt the phases measurements of the complex visibilities. To overcome this difficulty, turbulence independent quantities need to be constructed.

1. Closure phase

The first phase difference information used for image reconstruction in presence of turbulent measurements is the closure phase (operator ψ). It is defined as the phase of the bispectrum [8], i.e., the Fourier transform of the triple correlation. For three baselines $[\mathbf{b}_{a,b}, \mathbf{b}_{b,c}, \mathbf{b}_{a,c}]$ corresponding to a triplet (a, b, c) of telescopes, the ‘‘atmospheric corrupted’’ *instantaneous* visibilities at a given wavelength λ_n can be modeled as

$$y_{a,b}^{\lambda_n} = \gamma_{a,b}^{\lambda_n} \exp(i[\varphi_{a,b}^{\lambda_n} + \eta_a^{\lambda_n} - \eta_b^{\lambda_n}]) \quad (9)$$

$$y_{b,c}^{\lambda_n} = \gamma_{b,c}^{\lambda_n} \exp(i[\varphi_{b,c}^{\lambda_n} + \eta_b^{\lambda_n} - \eta_c^{\lambda_n}]) \quad (10)$$

$$y_{a,c}^{\lambda_n} = \gamma_{a,c}^{\lambda_n} \exp(i[\varphi_{a,c}^{\lambda_n} + \eta_a^{\lambda_n} - \eta_c^{\lambda_n}]) \quad (11)$$

where φ^{λ_n} are the quantities of interest (i.e., the uncorrupted phases) and η^{λ_n} are perturbation terms related to the corresponding telescopes. The closure phase associated with this triplet is defined as

$$\begin{aligned} \psi_{a,b,c}^{\lambda_n} &= \angle y_{a,b}^{\lambda_n} y_{b,c}^{\lambda_n} y_{a,c}^{\lambda_n *} \\ &= \varphi_{a,b}^{\lambda_n} + \varphi_{b,c}^{\lambda_n} - \varphi_{a,c}^{\lambda_n} \end{aligned} \quad (12)$$

$$= \mathbf{h}_{a,b,c}^{\lambda_n} \varphi^{\lambda_n} \quad (13)$$

where φ^{λ_n} is as in Eq. 6 the vector containing all *unperturbed* phases for wavelength λ_n , and $\mathbf{h}_{a,b,c}^{\lambda_n}$ is a sparse row vector with only three non zeros entries which take values $\{1, 1, -1\}$ (as reflected by Eq. 12). Clearly, the phase closure allows to get rid of atmospheric effects for triplets of complex visibilities. If N_t denotes the number of telescopes there are $(N_t - 1)(N_t - 2)/2$ independent closure phases [33].

Let matrix $\mathbf{H}^{\lambda_n}$ concatenate in its rows the independent closure phases of the type (13) that can be obtained for the available triplets of telescopes at wavelength λ_n . For the simplicity of the presentation but without loss of generality, we assume here that all telescope pairs observe the same channels. In this case, the triplets involved in independent closure phases can be taken as the same in each wavelength channel. The global closure phase operator $\mathbf{H}^\psi$ is then simply a block diagonal matrix that replicates $\mathbf{H}^{\lambda_1}$, the closure phase matrix for λ_1

$$\psi = \mathbf{H}^\psi \varphi, \quad \mathbf{H}^\psi = \mathbf{I}_{N_\lambda} \otimes \mathbf{H}^{\lambda_1} \quad (14)$$

where ψ is the vector of all closure phases and φ is the unknown unperturbed phase vector of Eq. 6.

2. Differential Phases

A second phase difference information that is particularly interesting in polychromatic imaging is the differential phase $\Delta\varphi$. For one baseline $\mathbf{b}_{a,b}$, differential phases measure the phase evolution in wavelength with respect to a phase reference channel (see e.g. [34]). They are defined as

$$\begin{aligned} \Delta\varphi_{a,b}^{\lambda_k, \lambda_{\text{ref}}} &= \angle y_{a,b}^{\lambda_k} - \angle y_{a,b}^{\lambda_{\text{ref}}} \\ &= \varphi_{a,b}^{\lambda_k} + \eta_a^{\lambda_k} - \eta_b^{\lambda_k} - \left(\varphi_{a,b}^{\lambda_{\text{ref}}} + \eta_a^{\lambda_{\text{ref}}} - \eta_b^{\lambda_{\text{ref}}} \right) \end{aligned} \quad (15)$$

If the analyzed bandwidth is relatively narrow, which is generally the case, the phase turbulence terms on each telescope $\eta_a^{\lambda_k}$ and $\eta_b^{\lambda_{\text{ref}}}$ are in first approximation independent of the wavelength [35]. Thus the phase difference becomes

$$\Delta\varphi_{a,b}^{\lambda_k, \lambda_{\text{ref}}} = \varphi_{a,b}^{\lambda_k} - \varphi_{a,b}^{\lambda_{\text{ref}}} \quad (16)$$

As closure phases, the differential phases in Eq. 16 are essentially not affected by the atmospheric perturbation.

The reference channel can be chosen as one of the available wavelengths. In this case $N_\lambda - 1$ differential phases are available per baseline. Without loss of generality, we denote below by λ_1 the reference channel. Similarly to Eq. 13 we can write using definition of φ in Eq. 6:

$$\Delta\varphi_{a,b}^{\lambda_k, \lambda_1} = \mathbf{h}_{a,b}^{\lambda_k, \lambda_1} \varphi \quad (17)$$

where $\mathbf{h}_{a,b}^{\lambda_k, \lambda_1}$ is a sparse row vector with only two non zeros entries which take values $\{1, -1\}$. Denoting by

$\Delta\varphi$ the differential phases vector which collects all wavelengths differences for all baselines, and by $\mathbf{H}^{\Delta\varphi}$ the matrix which concatenates all vectors $\mathbf{h}_{a,b}^{\lambda_k, \lambda_1}$ in its rows, we have similarly to Eq. 14

$$\Delta\varphi = \mathbf{H}^{\Delta\varphi} \varphi, \quad \mathbf{H}^{\Delta\varphi} = (-\mathbf{1}_{(N_\lambda - 1)} \otimes \mathbf{I}_{N_b} | \mathbf{I}_{(N_\lambda - 1) \times N_b}) \quad (18)$$

Equations 14 and 18 can be merged in the single measurement equation

$$\xi = \mathbf{H} \varphi \quad (19)$$

where $\mathbf{H} = \begin{bmatrix} \mathbf{H}^\psi \\ \mathbf{H}^{\Delta\varphi} \end{bmatrix}$ models all phase difference information.

C. Practical remarks

In practice, a specific phase reference channel $\varphi_{a,b}^{\lambda_{\text{ref}}}$ is sometimes computed using all the visibilities associated with a baseline. One possibility is to set the phase reference as the angle of the empirical mean over the channels:

$$\varphi_{a,b}^{\lambda_{\text{ref}}} = \angle \frac{1}{N_\lambda} \sum_{n=1}^{N_\lambda} y_{a,b}^{\lambda_n}$$

Another possibility, as for the AMBER instruments [35, 36] is to define a channel dependent reference:

$$\varphi_{a,b}^{\lambda_{\text{ref}}} = \angle \frac{1}{N_\lambda - 1} \sum_{m \neq n}^{N_\lambda} y_{a,b}^{\lambda_m}$$

In these cases, the formalism of Eq. 16 is still valid assuming that all the reference channels for a given baseline are almost equal and will cancel in the difference of differential phases. Equation 16 can then be replaced by:

$$\Delta\varphi_{a,b}^{\lambda_2, \lambda_1} = \Delta\varphi_{a,b}^{\lambda_2, \lambda_{\text{ref}}} - \Delta\varphi_{a,b}^{\lambda_1, \lambda_{\text{ref}}} \quad (20)$$

and the linear Eq. 18 is kept with the same matrix $\mathbf{H}^{\Delta\varphi}$.

Consider the general case where $N_t > 3$ telescopes configurations are used to observe at $N_\lambda > 1$ wavelengths. For a single exposure the number of spatial visibilities measured and so the number of phases to estimate is $N_\phi = N_b N_\lambda$, where the number of bases $N_b = N_t(N_t - 1)/2$. If the matrix $\mathbf{H}$ involved in the phases to the phases differences transformation is full ranked then there are no phases identifiability ambiguities. In practice, this matrix is the concatenation of two matrices, the first for the phase closures and the second for the differential phases.

A closure is said “closed” if the set of spatial frequencies involved in its construction are the vertex of a polygon. The minimal configuration to measure one closure phase is to use three telescopes, as the addition of several triplets of telescopes can describe all the possible polygon the information is clearly redundant. The

rank of $\mathbf{H}^\psi$, the matrix which connects the phase closures to the polychromatic phases is given by the number of lines which ensure independent phase closures. For a single analyzed wavelength the rank of the matrix $\mathbf{H}^{\lambda_n}$ of section III B 1 is $N_\psi = (N_t - 1)(N_t - 2)/2$ [12]. When polychromatic data are measured the augmented matrix follows the right term of Eq. 14, the phase closures for each wavelength are estimated from the same set of telescopes triplets. Consequently, the rank of $\mathbf{H}^\psi$ is then increased up by a factor N_λ : $\text{rank}(\mathbf{H}^\psi) = N_\psi N_\lambda$. The ratio $\text{rank}(\mathbf{H}^\psi)/N_\phi = 1 - 2/N_t$ depends only on the number of telescopes.

For the differential phases, the construction of non redundant measures is detailed in section III B 2 and is driven by $\mathbf{H}^{\Delta\varphi}$. The rank of this matrix is given by the number of independent differential phases that may result in the transformation from phases and is $\text{rank}(\mathbf{H}^{\Delta\varphi}) = N_b(N_\lambda - 1)$. The ratio $\text{rank}(\mathbf{H}^{\Delta\varphi})/N_\phi = 1 - 1/N_\lambda$ depends only on the number of analyzed wavelengths.

As $\mathbf{H}$ is the concatenation of $\mathbf{H}^\psi$ and $\mathbf{H}^{\Delta\varphi}$ its rank is given by the number of independent relation between these matrices. Consider three telescopes (a, b, c) and two wavelengths (λ_1, λ_k), λ_1 is choised to be the differential phases reference channel, a linear relation which connects the phase closures of Eq. 13 to the differential phases of Eq. 17 is:

$$\psi_{a,b,c}^{\lambda_k} - \psi_{a,b,c}^{\lambda_1} = \Delta\varphi_{a,b}^{\lambda_k,\lambda_1} + \Delta\varphi_{b,c}^{\lambda_k,\lambda_1} - \Delta\varphi_{a,c}^{\lambda_k,\lambda_1} \quad (21)$$

the difference of lines of $\mathbf{H}^\psi$ for a fixed wavelength only defines redundant phase closures. The combination over λ_k and λ_1 of all the triplets of bases, as in Eq. 21, expresses all the differential phases and so the rank of $\mathbf{H}$ is at least the rank of $\mathbf{H}^{\Delta\varphi}$. A more general phase closures differences $\psi_{a,b,c}^{\lambda_k} - \psi_{a,b,c}^{\lambda_m}$ can be expressed in term of (λ_k, λ_1) , leads to redundant information and adds nothing. Consequently, the only remaining independent information is monochromatic and corresponds to all the phase closures related to the reference channel λ_1 so $\text{rank}(\mathbf{H}) = N_b(N_\lambda - 1) + N_\psi$. Finally, the ratio $\text{rank}(\mathbf{H})/N_\phi = 1 - 2/(N_t N_\lambda)$ depends on the number of telescopes and analyzed wavelengths.

The ratio of the number of independent phases difference for the two models (ranks of $\mathbf{H}^\psi$ and $\mathbf{H}^{\Delta\varphi}$) and the augmented model (rank of $\mathbf{H}$) to the number of phases is shown in Fig. 1. The ratio is drawn for $N_t = 3, \dots, 6$ telescopes which corresponds to actual interferometers and up to 100 N_λ . Note as an example that $N_\lambda > 500$ for the high resolution mode of AMBER [34], and $N_\lambda = 3$ for the standard uses of PIONIER [21].

The use of all phases differences information largely increases the rank of the matrix which is always greater than for the phase closures alone.

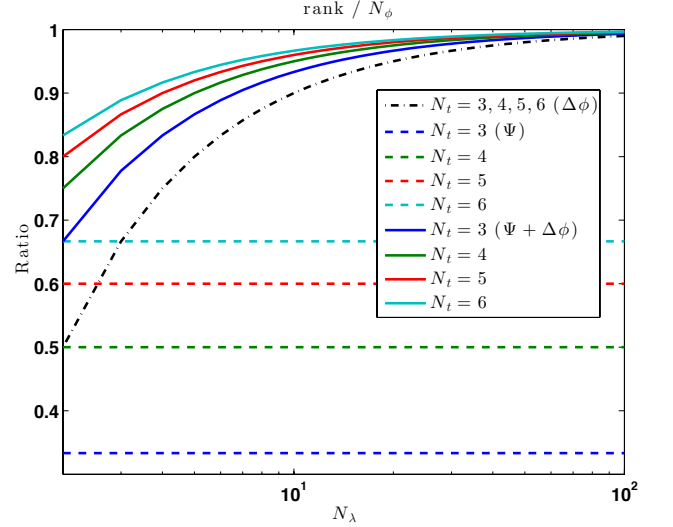


FIG. 1. Ratio of the number of independent phases difference to the total number of phases as a function of the number of telescopes and analyzed wavelengths.

IV. INVERSE PROBLEM APPROACH

A. Image reconstruction problem formulation

According to Eq. 19 and notations defined in Eqs. 5 and 6, a model for measurements of difference phases and squared moduli can be written as

$$\xi = \mathbf{H} \varphi + \eta_\xi \quad (22)$$

$$\zeta = \gamma^2 + \eta_\zeta \quad (23)$$

where η_ξ and η_ζ account for noise and modeling errors. Classical assumptions on their distributions are considered here. The noise η_ζ is assumed to be jointly independent and Gaussian [12, 37], and the noise η_ξ is assumed to be jointly independent and marginally Von Mises distributed [15].

Writing the opposite logarithm of the likelihood related to the squared absolute value and phase differences model, we obtain:

$$g^{\text{data}}(\mathbf{x}) = \alpha g^\zeta(\mathbf{y}^\gamma) + \beta g^\xi(\mathbf{y}^\phi) \quad (24)$$

where, α and β are relative weighting terms and:

$$g^\zeta(\mathbf{y}^\gamma) = \sum_n \frac{1}{\omega_n} (\zeta_n - \gamma_n^2)^2, \quad \gamma = |\mathbf{F}\mathbf{x}| \quad (25)$$

$$g^\xi(\mathbf{y}^\phi) = - \sum_m \kappa_m \cos(\mathbf{h}_m \varphi - \xi_m), \quad \varphi = \angle \mathbf{F}\mathbf{x} \quad (26)$$

where ω_n is the variance of ζ_n . The constants κ_m are related to the variance of ξ_m by

$$\text{var}(\xi_m) = 1 - I_1(\kappa_m)/I_0(\kappa_m)$$

where I_j is the modified Bessel function of order j [38]. In practice κ_n is computed inverting numerically the previous equation. The variance of the closure phase and

differential phases, if not provided by the instrument pipeline, are estimated assuming independence of phase measurements.

Image reconstruction can be seen as an inverse problem [39–41]. The model which connects object to the measured data involves a NuDFT, as described in Eq. 3 and Eqs. 22 and 23. This transformation leads to a poor coverage of the spatial frequency plan and makes the problem ill-conditioned which requires tackling the image reconstruction as a regularized optimization problem [42]. We will adopt here an objective function of the form:

$$\mathbf{x} \leftarrow \underset{\mathbf{x} \in \Pi}{\text{minimize}} (g^{\text{data}}(\mathbf{x}) + f^{\text{reg}}(\mathbf{x})) \quad (27)$$

where the image $\mathbf{x}$ can be constrained to be spatially limited in the support Π , and further constraints such as non negativity can be added in $f^{\text{reg}}(\mathbf{x})$, which contains all the regularization terms. The support constraint is not included in $f^{\text{reg}}(\mathbf{x})$ for technical reasons related to the ADMM methodology described below.

B. Regularizations and constraints

OI images are by nature non negative and sometimes contain sources that are spatially localized. However, specifying the properties of the object parameters $\mathbf{x}$ only in terms of non negativity and spatial support is usually not constraining enough. It follows that the uses of regularization terms to emphasize some inherent a priori knowledge about the image structure is necessary.

Following the matrix notation for the 3D object as defined in Eq. 8, PAINTER in its current form accounts the following classical priors:

- Ridge/Tikhonov regularization, motivated by the poor conditioning of the NuDFT operator.
- Spatial or spectral Total Variation [43].
- Spatial or spectral smoothness [27].

The support constraint is defined by the parameters space Π in Eq. 27 and the non-negativity constraint by the regularization term $\mathbb{1}_{\mathbb{R}^+}(\mathbf{X})$. In this description the regularization function in Eq. 27 writes

$$f^{\text{reg}}(\mathbf{x}) = \mathbb{1}_{\mathbb{R}^+}(\mathbf{X}) + (\mu_\varepsilon/2) \|\mathbf{X}\|_{\text{F}}^2 + \dots \quad (28)$$

$$\mu_s \Omega_{\text{pspat}}(\mathbf{H}^{\text{spat}} \mathbf{X}) + \mu_\lambda \Omega_{\text{pspec}}(\mathbf{X} \mathbf{H}^{\text{spec}})$$

$\mathbf{H}^{\text{spat}}$ and $\mathbf{H}^{\text{spec}}$ are the matrices of finite difference, associated with the spatial and spectral regularizations respectively [27]. $\mathbf{H}^{\text{spat}} \mathbf{X}$ act on the column of $\mathbf{X}$ which are images processed independently while $\mathbf{X} \mathbf{H}^{\text{spec}}$ operates on the line of $\mathbf{X}$ to connects the pixels between wavelengths. $\Omega_{\text{pspat}}(\cdot)$ and $\Omega_{\text{pspec}}(\cdot)$ are matrix regularization terms that can be chosen as:

$$\Omega_1(\mathbf{M}) = \sum_{i,j} |\mathbf{M}_{i,j}| \quad \Omega_2(\mathbf{M}) = \sum_{i,j} \mathbf{M}_{i,j}^2 \quad (29)$$

Finally, μ_ε , μ_s and μ_λ are hyper-parameters which control the weights of the associated regularization terms.

V. 3D RECONSTRUCTION ALGORITHM

Owing to the unavoidable non convexity of the problem as defined by Eq. 27 (see e.g. in [13]), the vast majority of image reconstruction algorithms rely on a descent optimization principle. So does PAINTER by using the flexibility of the Alternate Direction Methods of Multipliers (ADMM) algorithm, which was already used in [26, 27] to reconstruct stellar spectrum of point sources from complex visibilities. Within this framework the reconstruction algorithm will iterate as follow:

1. Update the complex visibilities.
2. Use the estimated complex visibilities and spatio-spectral regularization to reconstruct the polychromatic object (3D-images).
3. Update the Lagrange multipliers.

Specific details are given in the rest of this section.

A. ADMM optimization algorithm

The ADMM framework allows to split complex problem into smaller and easier ones by introducing auxiliary variables. However, in this case additional terms have to be taken into account. The subproblems can be solved independently by means of proximal operators. Standard constraints are taken into account in PAINTER and most proximal operators used in this section are known. The optimization problem of Eq. 27 where $g^{\text{data}}(\mathbf{x})$ is given by Eqs. 24–26 and the regularization term $f^{\text{reg}}(\mathbf{x})$ is given by Eq. 29 is equivalent to:

$$\underset{\mathbf{x} \in \Pi, \mathbf{x}_i, \mathbf{y}, \mathbf{y}^\gamma, \mathbf{y}^\phi}{\text{minimize}} \quad \alpha g^\zeta(\mathbf{y}^\gamma) + \beta g^\xi(\mathbf{y}^\phi) + \frac{\mu_\varepsilon}{2} \|\mathbf{X}\|_{\text{F}}^2 + \mathbb{1}_{\mathbb{R}^+}(\mathbf{X})$$

$$+ \mu_s \Omega_{\text{pspat}}(\mathbf{H}^{\text{spat}} \mathbf{X}) + \mu_\lambda \Omega_{\text{pspec}}(\mathbf{X} \mathbf{H}^{\text{spec}})$$

subject to $\mathbf{y}^\gamma = \mathbf{y}, \mathbf{y}^\phi = \mathbf{y}, \mathbf{X}_i = \mathbf{X}, \mathbf{y} = \mathbf{F} \mathbf{x}$ (30)

In this equation, the vector $\mathbf{y}^\gamma$ denotes the complex visibilities variables whose likelihood is associated with the measurements of visibilities squared absolute values, Eq. 25. The vector $\mathbf{y}^\phi$ denotes the complex visibilities variables whose likelihood is associated with the measurements of the visibilities phases *differences*, Eq. 26.

Using the same approach for the three last regularization terms leads to the introduction of new auxiliary variables $\mathbf{P}$, $\mathbf{T}$ and $\mathbf{V}$ and to replace each regularization term in Eq. 30 by a constrained problem.

- The non-negativity constraint $\mathbb{1}_{\mathbb{R}^+}(\mathbf{X})$ and the support constraint in Eq. 30 are replaced by

$$\underset{\mathbf{P}}{\text{minimize}} \quad \mathbb{1}_{\mathbb{R}^+}(\mathbf{P}) \quad (31)$$

subject to $\mathbf{P} = \mathbf{X}, \mathbf{P} \in \Pi$

- The total variation regularization $\Omega_{\text{Pspat}}(\mathbf{H}^{\text{spat}}\mathbf{X})$ in Eq. 30 is replaced by

$$\begin{aligned} & \underset{\mathbf{T}}{\text{minimize}} \quad \Omega_{\text{Pspat}}(\mathbf{T}) \\ & \text{subject to} \quad \mathbf{T} = \mathbf{H}^{\text{spat}}\mathbf{X} \end{aligned} \quad (32)$$

- The spectral regularization $\Omega_{\text{Pspec}}(\mathbf{X}\mathbf{H}^{\text{spec}})$ in Eq. 30 is replaced by

$$\begin{aligned} & \underset{\mathbf{V}, \mathbf{S}}{\text{minimize}} \quad \Omega_{\text{Pspec}}(\mathbf{V}) \\ & \text{subject to} \quad \mathbf{V} = \mathbf{S}\mathbf{H}^{\text{spec}} \text{ and } \mathbf{S} = \mathbf{X} \end{aligned} \quad (33)$$

the use of two auxiliary variables ensure spectral separability when solving for $\mathbf{X}$.

The final optimization problem obtained by replacing Eqs. 31, 32 and 33 in Eq. 30 is iteratively solved using the ADMM algorithm [29]. Auxiliary variables related to the complex visibility: $\mathbf{y}$, $\mathbf{y}^\gamma$, $\mathbf{y}^\phi$ have a proper Lagrange multiplier: $\mathbf{v}_y$ (or $\mathbf{\Upsilon}_y$, with columns $\mathbf{\Upsilon}_y^{\lambda_n}$, in matrix form), $\mathbf{v}_\gamma$, $\mathbf{v}_\phi$ and share the same augmented Lagrangian parameter ρ_y . The auxiliary variables introduced by the regularization: $\mathbf{P}$, $\mathbf{T}$, $\mathbf{V}$, $\mathbf{S}$ are associated with the Lagrange multipliers $\mathbf{\Upsilon}_P$, $\mathbf{\Upsilon}_T$, $\mathbf{\Upsilon}_V$, $\mathbf{\Upsilon}_S$ and to the augmented Lagrangian parameters ρ_P , ρ_T and ρ_S for $\mathbf{V}$ and $\mathbf{S}$.

Denoting with a “+” superscript updated quantities, alternated minimization of the augmented Lagrangian gives:

$$\mathbf{y}^{\gamma+} = \arg \min_{\mathbf{y}^\gamma} \alpha g^\zeta(\mathbf{y}^\gamma) + \frac{\rho_y}{2} \|\mathbf{y}^\gamma - \tilde{\mathbf{y}}^\gamma\|_2^2 \quad (34)$$

$$\mathbf{y}^{\phi+} = \arg \min_{\mathbf{y}^\phi} \beta g^\xi(\mathbf{y}^\phi) + \frac{\rho_y}{2} \|\mathbf{y}^\phi - \tilde{\mathbf{y}}^\phi\|_2^2 \quad (35)$$

$$\mathbf{y}^+ = \frac{1}{3} (\mathbf{y}^{\gamma+} + \mathbf{y}^{\phi+} + \mathbf{F}\mathbf{x} + \rho_y^{-1}(\mathbf{v}_y - \mathbf{v}_\phi - \mathbf{v}_\gamma)) \quad (36)$$

$$\begin{aligned} \mathbf{X}^{\lambda_n+} &= \mathbf{C}^{\lambda_n-1} \left[\mathbf{F}^{\lambda_n H} \left(\rho_y \mathbf{Y}^{\lambda_n+} - \mathbf{\Upsilon}_y^{\lambda_n} \right) + \dots \right. \\ & \quad \left. \mathbf{H}^{\text{spat} \top} \left(\rho_T \mathbf{T}^{\lambda_n} - \mathbf{\Upsilon}_T^{\lambda_n} \right) + \dots \right. \\ & \quad \left. \left(\rho_P \mathbf{P}^{\lambda_n} - \mathbf{\Upsilon}_P^{\lambda_n} \right) + \left(\rho_S \mathbf{S}^{\lambda_n} - \mathbf{\Upsilon}_S^{\lambda_n} \right) \right] \end{aligned} \quad (37)$$

$$\mathbf{x}^+ = \text{vec } \mathbf{X}^+ \quad (38)$$

$$\mathbf{P}^+ = \arg \min_{\mathbf{P} \in \Pi} \mathbb{1}_{\mathbb{R}^+}(\mathbf{P}) + \frac{\rho_P}{2} \left\| \mathbf{P} - \tilde{\mathbf{P}} \right\|_2^2 \quad (39)$$

$$\mathbf{T}^+ = \arg \min_{\mathbf{T}} \mu_S \Omega_{\text{Pspat}}(\mathbf{T}) + \frac{\rho_T}{2} \left\| \mathbf{T} - \tilde{\mathbf{T}} \right\|_2^2 \quad (40)$$

$$\begin{aligned} \mathbf{S}^+ &= \left[\mathbf{H}^{\text{spec}} \mathbf{H}^{\text{spec} \top} + \mathbf{I}_{N_\lambda} \right]^{-1} \left[(\mathbf{V} - \mathbf{\Upsilon}_V / \rho_S) \mathbf{H}^{\text{spec} \top} + \dots \right. \\ & \quad \left. (\mathbf{X}^+ + \mathbf{\Upsilon}_S / \rho_S) \right] \end{aligned} \quad (41)$$

$$\mathbf{V}^+ = \arg \min_{\mathbf{V}} \mu_\lambda \Omega_{\text{Pspec}}(\mathbf{V}) + \frac{\rho_S}{2} \left\| \mathbf{V} - \tilde{\mathbf{V}} \right\|_2^2 \quad (42)$$

with the definitions:

$$\tilde{\mathbf{y}}^\gamma = \mathbf{y} + \rho_y^{-1} \mathbf{v}_\gamma \quad (43)$$

$$\tilde{\mathbf{y}}^\phi = \mathbf{y} + \rho_y^{-1} \mathbf{v}_\phi \quad (44)$$

$$\begin{aligned} \mathbf{C}^{\lambda_n} &= \rho_y \mathbf{F}^{\lambda_n H} \mathbf{F}^{\lambda_n} + \rho_T \mathbf{H}^{\text{spat} \top} \mathbf{H}^{\text{spat}} + \dots \\ & (\mu_\varepsilon + \rho_P + \rho_S) \mathbf{I}_{N_x^2} \end{aligned} \quad (45)$$

$$\tilde{\mathbf{P}} = \mathbf{X}^+ + \mathbf{\Upsilon}_P / \rho_P \quad (46)$$

$$\tilde{\mathbf{T}} = \mathbf{H}^{\text{spat}} \mathbf{X}^+ + \mathbf{\Upsilon}_T / \rho_T \quad (47)$$

$$\tilde{\mathbf{V}} = \mathbf{S}^+ \mathbf{H}^{\text{spec}} + \mathbf{\Upsilon}_V / \rho_S \quad (48)$$

The update of the Lagrange multipliers are:

$$\mathbf{v}_\gamma^+ = \mathbf{v}_\gamma + \rho_y (\mathbf{y}^+ - \mathbf{y}^{\gamma+}) \quad (49)$$

$$\mathbf{v}_\phi^+ = \mathbf{v}_\phi + \rho_y (\mathbf{y}^+ - \mathbf{y}^{\phi+}) \quad (50)$$

$$\mathbf{v}_y^+ = \mathbf{v}_y + \rho_y (\mathbf{F}\mathbf{x}^+ - \mathbf{y}^+) \quad (51)$$

$$\mathbf{\Upsilon}_P^+ = \mathbf{\Upsilon}_P + \rho_P (\mathbf{X}^+ - \mathbf{P}^+) \quad (52)$$

$$\mathbf{\Upsilon}_T^+ = \mathbf{\Upsilon}_T + \rho_T (\mathbf{H}^{\text{spat}} \mathbf{X}^+ - \mathbf{T}^+) \quad (53)$$

$$\mathbf{\Upsilon}_S^+ = \mathbf{\Upsilon}_S + \rho_S (\mathbf{X}^+ - \mathbf{S}^+) \quad (54)$$

$$\mathbf{\Upsilon}_V^+ = \mathbf{\Upsilon}_V + \rho_S (\mathbf{S}^+ \mathbf{H}^{\text{spec}} - \mathbf{V}^+) \quad (55)$$

The proximal operators [44] in Eqs. 34 and 35 are detailed in subsections VB and VC respectively and the proximal operators for the regularization terms in Eqs. 39, 40 and 42 are detailed in subsections VD–VE.

B. Proximal operator for squared visibility

The proximal operator of Eq. 34 updates the estimation of the complex visibilities from the measured squared absolute visibilities. Replacing $g^\zeta(\mathbf{y}^\gamma)$ by its expression in Eq. 25, Eq. 34 separates on each component:

$$\mathbf{y}_n^{\gamma+} = \arg \min_{\mathbf{y}_n^\gamma} \alpha \frac{1}{\omega_n} (\zeta_n - \gamma_n^2)^2 + \frac{\rho_y}{2} |\mathbf{y}_n^\gamma - \tilde{\mathbf{y}}_n^\gamma|^2 \quad (56)$$

which separates again on the modulus and phase of $\mathbf{y}^\gamma$ as:

$$\begin{aligned} \gamma_n^+ &= \arg \min_{\gamma_n > 0} \frac{\alpha}{\omega_n} (\zeta_n - \gamma_n^2)^2 + \frac{\rho_y}{2} (\gamma_n - \tilde{\gamma}_n)^2 \\ \varphi_n^+ &= \tilde{\varphi}_n \end{aligned} \quad (57)$$

where $\tilde{\gamma}$ and $\tilde{\varphi}$ denote in this section the moduli and phases of $\tilde{\mathbf{y}}^\gamma$.

The minimization of the previous fourth order polynomial for $\gamma_n > 0$ is obtained by computing the real roots of its derivative:

$$\gamma_n^3 + \left(\rho_y \frac{\omega_n}{4\alpha} - \zeta_n \right) \gamma_n - \rho_y \frac{\omega_n}{4\alpha} \tilde{\gamma}_n = 0 \quad (58)$$

using Cardano’s method [45]. γ_n^+ is the real positive root of Eq. 58 which minimizes the criterium Eq. 56. If Eq. 58 has no positive roots, then the polynomial to minimize is strictly increasing for $\gamma_n > 0$ and $\gamma_n^+ = 0$.

C. Proximal operator for phase differences

The proximal operator of Eq. 35 updates the estimation of the complex visibilities from the measured visibilities phase differences. This problem involves triplets and duets of visibility phases which, contrarily to the moduli, are not separable. Replacing $g^\xi(\mathbf{y}^\phi)$ given in Eq. 26 we obtain:

$$\mathbf{y}^{\phi+} = \arg \min_{\mathbf{y}^\phi} -\beta \sum_m \kappa_m \cos(\mathbf{h}_m \boldsymbol{\varphi} - \boldsymbol{\xi}_m) + \dots \\ \frac{\rho_y}{2} \sum_n \left(\gamma_n^2 + \tilde{\gamma}_n^2 - 2 \gamma_n \tilde{\gamma}_n \cos(\tilde{\varphi}_n - \varphi_n) \right) \quad (59)$$

where $\tilde{\gamma}$ and $\tilde{\varphi}$ denote in this section the modulus and phases of $\tilde{\mathbf{y}}^\phi$. Minimization of Eq. 59 w.r.t. $\gamma_n > 0$ gives:

$$\gamma_n^+(\varphi_n) = \max(0, \tilde{\gamma}_n \cos(\tilde{\varphi}_n - \varphi_n)) \quad (60)$$

Eq. 60 is first replaced in Eq. 59 and the resulting function minimized w.r.t. $\boldsymbol{\varphi}$. This minimization is carried out numerically using a gradient descent algorithm. A compact expression of the gradient of the function to be minimized is:

$$\beta \mathbf{H}^T \text{diag}(\boldsymbol{\kappa})^{-1} \sin(\mathbf{H}\boldsymbol{\varphi} - \boldsymbol{\xi}) - \rho_y \boldsymbol{\Gamma}^+(\boldsymbol{\varphi}) \tilde{\boldsymbol{\Gamma}} \sin(\tilde{\boldsymbol{\varphi}} - \boldsymbol{\varphi}) \quad (61)$$

where $\boldsymbol{\Gamma}^+(\boldsymbol{\varphi}) = \text{diag}(\gamma^+(\boldsymbol{\varphi}))$ and $\tilde{\boldsymbol{\Gamma}} = \text{diag}(\tilde{\gamma})$ are diagonal matrices with elements $\gamma^+(\boldsymbol{\varphi})$ and $\tilde{\gamma}$. Once $\boldsymbol{\varphi}^+$ is obtained, γ^+ is given by $\gamma^+(\boldsymbol{\varphi}^+)$. It is worthy to note that if at an iteration step of the descent algorithm $|\tilde{\varphi}_n - \varphi_n| > \pi/2$, from Eq. 60 the corresponding term in the second part of Eq. 59 reduces to $\tilde{\gamma}_n^2$ and the next estimate of $\boldsymbol{\varphi}$ will not depend on $\tilde{\gamma}_n$ and $\tilde{\varphi}_n$. For this reason the initial estimate of $\boldsymbol{\varphi}$ is set at $\tilde{\boldsymbol{\varphi}}$. The minimization in PAINTER relies on the quasi-Newton algorithm implemented in [46], note that the non convexity of the problem leads to a local minimization to estimate $\boldsymbol{\varphi}$.

D. Positivity and compactness

The solution of the constrained minimization in Eq. 39 [29, 44] is:

$$\mathbf{P}_{k,l}^+ = \begin{cases} \tilde{\mathbf{P}}_{k,l} & \text{if } \tilde{\mathbf{P}}_{k,l} \odot \boldsymbol{\Pi}_{k,l} > 0 \\ 0 & \text{else.} \end{cases} \quad (62)$$

where $\boldsymbol{\Pi}$ is a binary mask matrix (composed of 0 and 1) which represents the support of the object (disk etc.), each column of $\boldsymbol{\Pi}$ is for a wavelength. If the object support is not constrained $\boldsymbol{\Pi}_{k,l} = 1 \forall [k, l]$.

E. Spatial and spectral regularization

The proximal operator of Eqs. 40 and 42 is:

- the *soft thresholding* [29] if $p_{\text{spat}/\text{spec}} = 1$:

$$\mathbf{D}_{k,l}^+ = \begin{cases} \mathbf{U}_{k,l} & \text{if } \mathbf{U}_{k,l} > 0 \\ 0 & \text{else.} \end{cases} \quad (63)$$

where $\mathbf{U}_{k,l} = \left(1 - \frac{\mu_{s/\lambda}}{|\tilde{\mathbf{D}}_{k,l}|} \right) \tilde{\mathbf{D}}_{k,l}$

- if $p_{\text{spat}/\text{spec}} = 2$:

$$\mathbf{D}^+ = \frac{1}{1 + 2 \mu_{s/\lambda}} \tilde{\mathbf{D}} \quad (64)$$

where $\mathbf{D}$ can be $\mathbf{T}$ or $\mathbf{V}$. p_{spat} and p_{spec} can be chosen according to a priori selected on image, if the desired spatial/spectral smoothness regularization must promotes brightness jump then $p_{\text{spat}/\text{spec}} = 1$ otherwise $p_{\text{spat}/\text{spec}} = 2$. As an example, the proximal operator for the ℓ_2 norm spectral smoothness regularization, as expressed in Eq. 42 with $p_{\text{spec}} = 2$, is:

$$\mathbf{V}^+ = \frac{1}{1 + 2 \mu_\lambda} \tilde{\mathbf{V}} \quad (65)$$

Note that according to the structure of matrix $\tilde{\mathbf{T}}$, see Eq. 47, the spatial proximal operator Eq. 63 will apply separately on images at different wavelengths. Similarly, according to the structure of matrix $\tilde{\mathbf{V}}$, the spectral proximal operators operates separately on the voxels. Finally, in order to scale properly the regularization parameters, the μ_s and μ_λ of Eq. 30 are divided by their number of non zero elements in $\mathbf{H}^{\text{spat}}$ and $\mathbf{H}^{\text{spec}}$ respectively. This normalization should makes the parameters independent of the size of the 3D-image.

VI. SIMULATIONS

Computer simulations are presented in order to illustrate the performances of PAINTER and the benefits of a combined use of phase closures and differentials phases. To do so, synthetic data set are used and the algorithm conforms to the complete chain needed in OI image reconstruction i.e going from files to 3D-image.

A. Synthetic data

1. Improvement related to phases differences

The first simulation aims at determining the improvement on the reconstruction due of the combination of phases differences model. The simulated noiseless object consists in two uniform disks with a decreasing radius with wavelengths. The intensity of each disk as a function of the wavelength is constant. The original cube is of size 64×64 pixels $\times 8$ channels. The instrumental

configurations (geometrical position, wavelengths, acquisition time,...) have been generated with the ASPRO2 [47] software which simulates realistic interferometric observation and store the data into OIFITS files [30]. The simulation uses the configuration of the AMBER instrument at the VLTI [34] with three telescopes. Height equi-spaced wavelengths in the range $2.1635 \mu\text{m} - 2.1686 \mu\text{m}$ (high resolution) at five acquisition instants are analyzed. The spatial frequencies coverage, including the earth rotation effect, is shown in Fig. 2 for the spectral channel at $\lambda = 2.16 \mu\text{m}$. This results in the measurement of 129 squared absolute visibilities and 43 phase closures for each wavelength. The differential phases are calculated as in Eq. 20 and corresponds to $128 \times (8 - 1) = 903$ measures. The object is shown in Fig. 3 in a channel per channel view.

2. Resolved stars reconstruction

The second simulation deals with realistic noisy synthetic data with a given signal to noise ratio, $\text{SNR} = 30 \text{ dB}$. More specifically, to generate the data we add zero-mean independent Gaussian random noise to each measurement. For each phase data φ , the standard deviation of the noise is $\sigma_\varphi = 1 \text{ rd}/\text{SNR}$ (in radians). For each amplitude data γ , the standard deviation of the noise is $\sigma_\gamma = \gamma/\text{SNR}$. The simulation consists in two resolved stars of different spectral types: G8V and M3V for the large and small stellar disks respectively. The spectra and their chromatic limb darkening laws are based on Kurucz and van Hamme models [48–50]. A channel of the original 3D-image is shown in Fig. 4. The two disks have a diameters of 22 and 14 milli-arcsecond (mas) and individual brightness distribution is varying among wavelengths. The original cube with a pixel resolution of 0.1 mas is of size $501 \times 501 \text{ pixels} \times 30 \text{ channels}$, the field of view (FOV) is $\approx 50 \text{ mas}$. The instrumental configuration was used in the 2004 International Beauty Contest in Optical Interferometry [51]: Thirty equi-spaced wavelengths in the range $1.45 \mu\text{m} - 1.84 \mu\text{m}$ (high resolution) at 13 acquisition instants are analyzed. The spatial frequencies coverage, including the earth rotation effect, is shown in Fig. 4. This results in the measurement of 195 squared absolute visibilities and 130 phase closures for each wavelength. The differential phases are calculated as in Eq. 20 and corresponds to $130 \times (30 - 1) = 3770$ measures and simulation data are stored into OIFITS files [30].

B. Reconstruction parameters

The initial estimate was a Dirac function centered on the image for each wavelength and the algorithm was stopped after 5000 iterations.

The type of regularization was chosen according to the a priori on the object: $p_{\text{spat}} = 1$ and $p_{\text{spec}} = 2$ to enforce

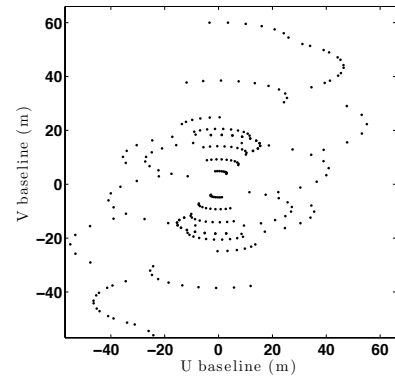


FIG. 2. Spatial frequencies coverage plan. Geometrical configuration the 2004 International Beauty Contest in Optical Interferometry.

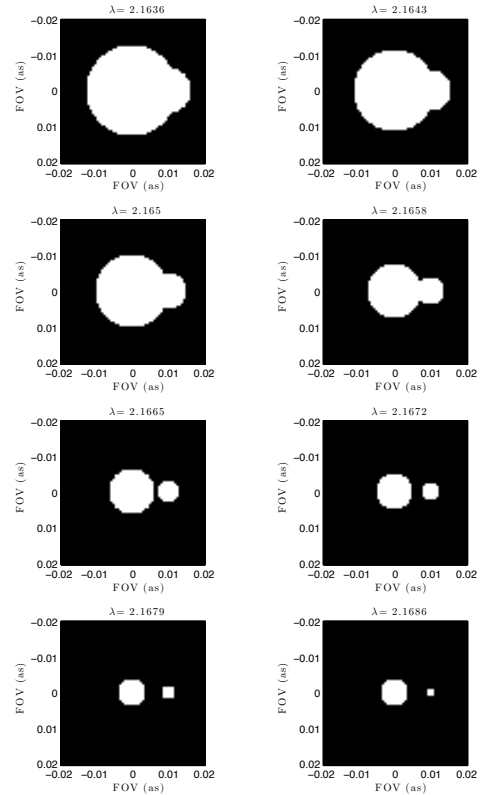


FIG. 3. Synthetic data, two disks with a decreasing radius along wavelength. Per channel view.

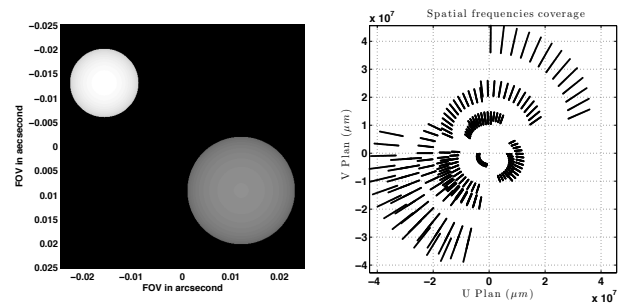


FIG. 4. Left: A channel of the 3D image used in simulation. Right: Spatial frequencies coverage plan. Geometrical configuration the 2004 International Beauty Contest in Optical Interferometry.

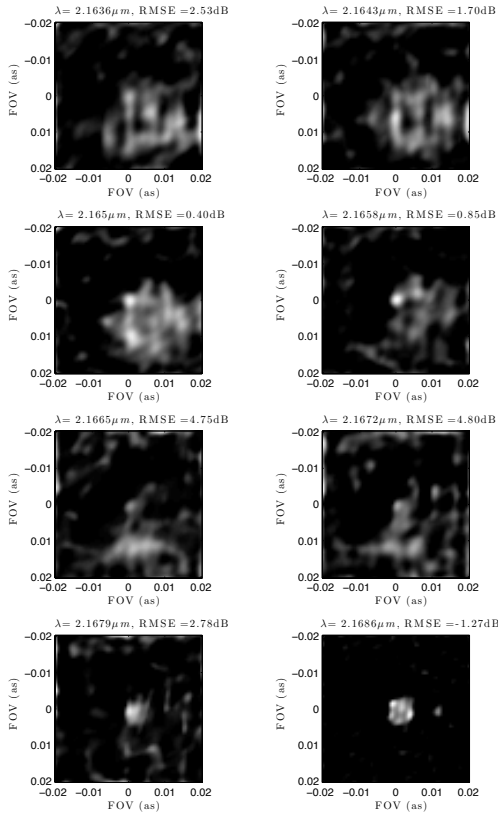


FIG. 5. Estimated 3-D object with closure phases only.

spatial continuity with sharp edges and spectral smoothness. The regularization weights $\mu_s = 2$ and $\mu_\lambda = 2$, were empirically tuned to have a visually acceptable solution (though not the best one). The associated augmented parameters, $\rho_T = 10$ and $\rho_S = 10$, which drive the convergence of the spatial/spectral regularization sub-problems have been tuned to be not too small (to avoid divergence) and not too large (to not slow down the convergence). Without loss of generality the problem can be dimensionless in ρ_y and thus we took $\rho_y = 1$. We tuned the weights of the two terms of Eqs. 34 and 35 so that, at convergences, these terms are of the same order. This yields $\alpha = 10^2$ and $\beta = 10^2$. Finally, we took $\mu_\epsilon = 10^{-6}$. The reconstruct cube is of size 64×64 pixels per channel ($N_\lambda = 8$ and $N_\lambda = 30$ for the first and the second simulation respectively), leading to $4096 \times N_\lambda$ parameters to estimate.

C. Simulation results

1. Improvement related to phases differences

In order to compare the improvement due to the combinations of phase differences we compare the results obtained using the squared visibility with: the closure phases, the differential phases and both phase differences.

The classical observable of phase differences used in

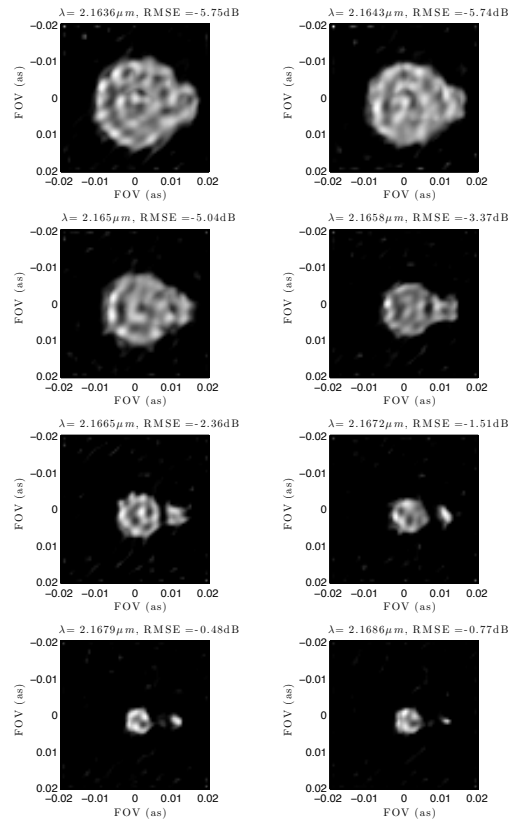


FIG. 6. Estimated 3-D object with differential phases only.

OI images reconstructions is the closure phases. The first estimation uses only this phase information. The result is shown in Fig. 5. Due to the lack of information: 43 phase closures for 129 baselines per wavelength, the reconstruction performances are quite poor. The narrow objects at high wavelength are well reconstructed. The object structure (two disks) is not correctly found at lower wavelengths where artifacts are present.

The second simulation uses the differential phases as the only phase differences measure. It leads to the estimated objects shown in Fig. 6. The object structures are better reconstructed. However some artifacts are present and the object surface are not as smooth as they should be.

Finally, the last simulation uses all the available phase information: phase closures and differential phases. The result in Fig. 7 shows that the object is extremely well reconstructed at all wavelengths. The Relative Mean Squared Error (RMSE) of the reconstructed object ($\hat{x}$) at wavelength λ_n defined as $\|\mathbf{x}^{\lambda_n} - \hat{\mathbf{x}}^{\lambda_n}\|_2^2 / \|\mathbf{x}^{\lambda_n}\|_2^2$ has been computed for each wavelength. The RMSEs in dB are presented in Fig. 5, 6 and 7 to quantify the improvement.

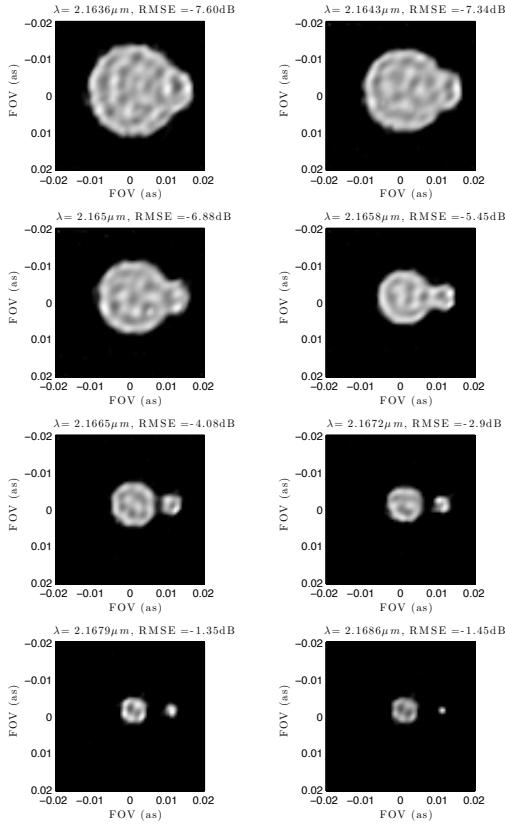


FIG. 7. Synthetic data. Estimated 3-D object with closure phases and differential phases.

2. Resolved stars reconstruction

Images estimated per channel are shown in Fig. 8, values lower than 1% of the maximum of the cube are thresholded. The shape and the size of the two stars are correctly restored. The polychromatic brightness distributions of the two estimated disks are compared to the solution in Fig. 9. Even if the estimated images are not as smooth as the original ones the estimated brightness follows correctly the *true distribution*.

VII. ACKNOWLEDGMENTS

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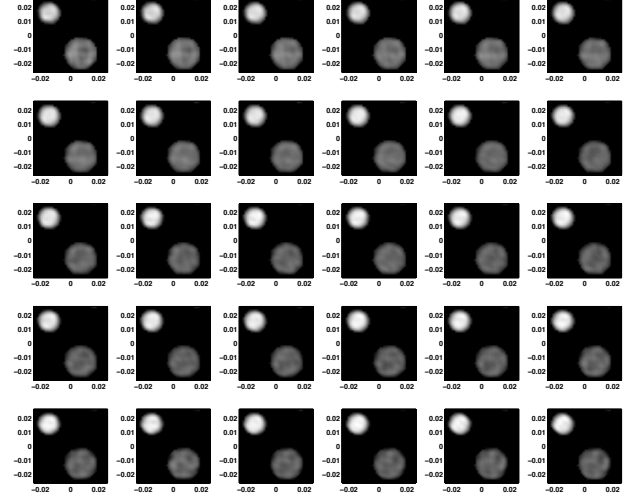


FIG. 8. Estimated Images per channels, the units are in arc second.

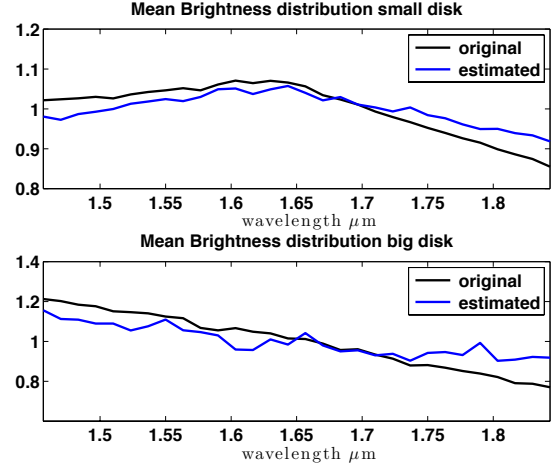


FIG. 9. Original and estimated Brightness distribution overs channels for the two disks.

VIII. CONCLUSIONS

A polychromatic 3-D image reconstruction for optical interferometry namely PAINTER is presented. The main contribution of PAINTER is to uses both phase closure and differential phases information to estimates the original phases of the observed scene. Compared to the use of phases difference independently, the combination leads to use a transformation matrix with a “number of wavelengths/telescopes” dependent rank. The constraint optimization follows the ADMM methodology with spatial and spectral regularization. PAINTER is able to deal with the complete chain of astronomical standard data (OIFITS file to data). Realistic simulation on noisy data are encouraging and shows the potential of PAINTER.

- [1] C. Aime, H. Lantéri, M. Diet, and A. Carlotti. “Strategies for the deconvolution of hypertelescope images,” *A&A*, 543, A42, 1–9, 2012.
- [2] A. Labeyrie. “Resolved imaging of extra-solar planets with future 10-100km optical interferometric arrays,” *A&A Sup. Series*, 118, 517–524, 1996.
- [3] A. Labeyrie, D. Mourard, F. Allouche, R. Chakraborty, J. Dejonghe, A. Surya, Y. Bresson, C. Aime, D. Mary and A. Carlotti. “Concept study of an “extremely Large hyper telescope” (eLhyt) with 1200m sparse aperture for direct imaging at 100 micro-arcsecond resolution,” In *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, July 2012.
- [4] A. Labeyrie. “Interference fringes obtained on VEGA with two optical telescopes,” *ApJ*, 196:L71–L75, March 1975.
- [5] J.D. Monnier. “Optical interferometry in astronomy,” *Reports on Progress in Physics*, 66(789-857), 2003.
- [6] R.C. Jennison. “A phase sensitive interferometer technique for the measurement of the Fourier transforms of spatial brightness distribution of small angular extent,” *MNRAS*, 118(276-284), 1958.
- [7] D. Mourard, I. Tallon-Bosc, F. Rigal, F. Vakili, D. Bonneau, F. Morand and P. Stee. “Estimation of visibility amplitude by optical long-baseline Michelson interferometry with large apertures,” *Astron. Astrophys.*, 288, 675–682, August 1994.
- [8] F. Roddier. “Triple correlation as a phase closure technique,” *Optics Communications*, 60:145–148, November 1986.
- [9] F. Roddier and P. Lena. “Long-baseline Michelson interferometry with large ground-based telescopes operating at optical wavelengths. I - General formalism: Interferometry at visible wavelengths,” *Journal of Optics*, 15:171–182, August 1984.
- [10] M. P. van Haarlem, M. W. Wise, A. W. Gunst, G. Heald, J. P. McKean, J. W. T. Hessels, A. G. de Bruyn, R. Nijboer, J. Swinbank, R. Fallows, M. Brentjens, A. Nelles, R. Beck, H. Falcke, R. Fender, J. Hörandel, L. V. E. Koopmans, G. Mann, G. Miley, H. Röttgering, B. W. Stappers, R. A. M. J. Wijers, S. Zaroubi, M. van den Akker, A. Alexov, J. Anderson, K. Anderson, A. van Ardenne, M. Arts, A. Asgekar, I. M. Avruch, F. Batejat, L. Bähren, M. E. Bell, M. R. Bell, I. van Bemmelen, P. Ben- nema, M. J. Bentum, G. Bernardi, P. Best, L. Birzan, A. Bonafede, A.-J. Boonstra, R. Braun, J. Bregman, F. Breitling, R. H. van de Brink, J. Broderick, P. C. Broekema, W. N. Brouw, M. Brggen, H. R. Butcher, W. van Cappellen, B. Ciardi, T. Coenen, J. Conway, A. Coolen, A. Corstanje, S. Damstra, O. Davies, A. T. Deller, R.-J. Dettmar, G. van Diepen, K. Dijkstra, P. Donker, A. Do- orduin, J. Dromer, M. Drost, A. van Duin, J. Eislöf- fel, J. van Enst, C. Ferrari, W. Frieswijk, H. Gankema, M. A. Garrett, F. de Gasperin, M. Gerbers, E. de Geus, J.-M. Griemeier, T. Grit, P. Gruppen, J. P. Hamaker, T. Hassall, M. Hoeft, H. A. Holties, A. Horneffer, A. van der Horst, A. van Houwelingen, A. Huijgen, M. Iac- obelli, H. Intema, N. Jackson, V. Jelic, A. de Jong, E. Juet- te, D. Kant, A. Karastergiou, A. Koers, H. Kollen, V. I. Kondratiev, E. Kooistra, Y. Koopman, A. Koster, M. Kuniyoshi, M. Kramer, G. Kuper, P. Lambropou- los, C. Law, J. van Leeuwen, J. Lemaitre, M. Loose, P. Maat, G. Macario, S. Markoff, J. Masters, R. A. McFad- den, D. McKay-Bukowski, H. Meijering, H. Meulman, M. Mevius, E. Middelberg, R. Millenaar, J. C. A. Miller- Jones, R. N. Mohan, J. D. Mol, J. Morawietz, R. Mor- ganti, D. D. Mulcahy, E. Mulder, H. Munk, L. Nieuwen- huis, R. van Nieuwpoort, J. E. Noordam, M. Norden, A. Noutsos, A. R. Offringa, H. Olofsson, A. Omar, E. Orr, R. Overeem, H. Paas, M. Pandey-Pommier, V. N. Pandey, R. Pizzo, A. Polatidis, D. Rafferty, S. Rawl- ings, W. Reich, J.-P. de Reijer, J. Reitsma, G. A. Rent- ing, P. Riemers, E. Rol, J. W. Romein, J. Roosjen, M. Ruiter, A. Scaife, K. van der Schaaf, B. Scheers, P. Schel- lart, A. Schoenmakers, G. Schoonderbeek, M. Serylak, A. Shulevski, J. Sluman, O. Smirnov, C. Sobey, H. Spreeuw, M. Steinmetz, C. G. M. Sterks, H.-J. Stiepel, K. Stu- urwold, M. Tagger, Y. Tang, C. Tasse, I. Thomas, S. Thoudam, M. C. Toribio, B. van der Tol, O. Usov, M. van Veelen, A.-J. van der Veen, S. ter Veen, J. P. W. Ver- biest, R. Vermeulen, N. Vermaas, C. Vocks, C. Vogt, M. de Vos, E. van der Wal, R. van Weeren, H. Weggemans, P. Weltevrede, S. White, S. J. Wijnholds, T. Wilhelm- son, O. Wucknitz, S. Yatawatta, P. Zarka, A. Zensus and J. van Zwieten. “LOFAR: The LOw-Frequency ARray”. *Astronomy and Astrophysics*, 556, A2, 1–53, May 2013.
- [11] É. Thiébaud. “Principles of Image Reconstruction in Interferometry,” In D. Mary, C. Theys, and C. Aime, edi- tors, *EAS Publications Series*, volume 59 of *EAS Publi- cations Series*, Published online by Cambridge University Press, pages 157–187, March 2013.
- [12] S. Meimon. “Reconstruction d’images astronomiques en interférométrie optique,” Université Paris Sud, Ph.D 2005.
- [13] S. Meimon, L. M. Mugnier, and G. Le Besnerais. “Recon- struction method for weak-phase optical interferometry,” *Opt. Lett.*, 30(14):1809–1811, Jul 2005.
- [14] D.F. Buscher. “Direct maximum-entropy image recon- struction from the bispectrum,” In J. G. Robertson and W. J. Tango, editors, *Very High Angular Resolu- tion Imaging*, volume 158 of *IAU Symposium*, page 91, 1994.
- [15] É. Thiébaud. “MIRA: an effective imaging algorithm for optical interferometry,” In *Astronomical Telescopes and Instrumentation, Proc. SPIE*, pages 70131I–70131I–12, 2008.
- [16] K.H. Hofmann and G. Weigelt. “Iterative image recon- struction from the bispectrum,” *A&A*, 278:328–339, Oc- tober 1993.
- [17] K.-H. Hofmann, G. Weigelt and D. Schertl. “An image reconstruction method (irbis) for optical/infrared inter- ferometry,” *A&A*, 565, A48; 1–10, 2014.
- [18] M.J. Ireland, J.D. Monnier and N. Thureau. “Monte- Carlo imaging for optical interferometry,” In *Society of Photo-Optical Instrumentation Engineers (SPIE) Con- ference Series*, July 2006.
- [19] F. Baron, B. Kloppenborg, and J. Monnier. “Toward 5d image reconstruction for optical interferometry,” *Proc. SPIE*, 8445:84451D–84451D–9, 2012.
- [20] R. Petrov, F. Malbet, A. Richichi, K.H. Hofmann, D. Mourard, K. Agabi, P. Antonelli, E. Aristidi, C. Baffa, U. Beckmann, P. Béro, Y. Bresson and F.Cassaing. “AM-

- BER: the near-infrared focal instrument for the Very Large Telescope Interferometer,” In P. Léna and A. Quirrenbach, editors, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, pages 68–79, July 2000.
- [21] J.B. Le Bouquin, J.P. Berger, B. Lazareff, G. Zins, G. Zins, P. Haguenaue, L. Jocou, P. Kern, R. Millan-Gabet, W. Traub, O. Absil, J.-C. Augereau, M. Benisty, N. Blind, X. Bonfils, P. Bourget, A. Delboulbe, P. Feautrier, M. Germain, P. Gitton, D. Gillier, M. Kiekebusch, J. Kluska, J. Knudstrup, P. Labeye, J.-L. Lizon, J.-L. Monin, Y. Magnard, F. Malbet, D. Maurel, F. Ménard, M. Micallef, L. Michaud, G. Montagnier, S. Morel, T. Moulin, K. Perraut, K. D. Popovic, P. Rabou, S. Rochat, C. Rojas, F. Roussel, A. Roux, E. Stadler, S. Steff, E. Tatulli and N. Ventura. “PIONIER: a 4-telescope visitor instrument at VLTI,” *A & A*, 535:A67, 1–15, November 2011.
- [22] D. Mourard, J. M. Clausse, A. Marcotto, K. Perraut, I. Tallon-Bosc, Ph. Bério, A. Blazit, D. Bonneau, S. Bosio, Y. Bresson, O. Chesneau, O. Delaa, F. Hénault, Y. Hughes, S. Lagarde, G. Merlin, A. Roussel, A. Spang, Ph. Stee, M. Tallon, P. Antonelli, R. Foy, P. Kervella, R. Petrov, E. Thiebaut, F. Vakili, H. McAlister, T. ten Brummelaar, J. Sturmann, L. Sturmann, N. Turner, C. Farrington and P. J. Goldfinger “VEGA: Visible spectrograph and polarimeter for the CHARA array: principle and performance,” *A & A*, 508:1073–1083, December 2009.
- [23] B. Lopez, P. Antonelli, S. Wolf, S. Lagarde, W. Jaffe, R. Navarro, U. Graser, R. Petrov, G. Weigelt, Y. Bresson, K. H. Hofmann, U. Beckman, T. Henning, W. Laun, C. Leinert, S. Kraus, S. Robbe-Dubois, F. and Vakili, A. Richichi, P. Abraham, J. C. Augereau, J. Behrend, P. Bério, N. Berruyer, O. Chesneau, J. M. Clausse, C. Connot, K. Demyk, W. C. Danchi, M. Dugué, G. Finger, S. Flament, A. Glazenborg, H. Hannenburg, M. Heininger, Y. Hugues, J. Hron, S. Jankov, F. Kerschbaum, G. Kroes, H. Linz, J. L. Lizon, P. Mathias, R. Mathar, A. Matter, J. L. Menut, K. Meisenheimer, F. Millour, N. Nardetto, U. Neumann, E. Nussbaum, A. Niedzielski, L. Mosoni, J. Olofsson, Y. Rabbia, T. Ratzka, F. Rigal, A. Roussel, D. Schertl, F. X. Schmider, B. Stecklum, E. Thiebaut, M. Vannier, B. Valat, K. Wagner and L. B. F. M. Waters. “MATISSE, perspective of imaging in the mid-infrared at the VLTI,” In *SPIE*, volume 7013, pages 1–10, 2008.
- [24] F. Eisenhauer, G. Perrin, W. Brandner, C. Straubmeier, A. Böhm, H. Baumeister, F. Cassaing, Y. Clénet, K. Dodds-Eden, A. Eckart, E. Gendron, R. Genzel, S. Gillessen, A. Grter, C. Gueriau, N. Hamaus, X. Haubois, M. Haug, T. Henning, S. Hippler, R. Hofmann, F. Hornuth, K. Houairi, S. Kellner, P. Kervella, R. Klein, J. Kolmeder, W. Laun, P. Léna, R. Lenzen, M. Marteaude, V. Naranjo, U. Neumann, T. Paumard, S. Rabien, J. R. Ramos, J. M. Reess, R.-R. Rohloff, D. Rouan, G. Rousset, B. Ruyet, A. Sevin, M. Thiel, J. Ziegler and D. Ziegler. “Gravity: microarcsecond astrometry and deep interferometric imaging with the VLTI,”. *Proceedings of the International Astronomical Union*, pages 100–101, 9 2007.
- [25] J. Kluska, F. Malbet, J.-P. Berger, F. Baron, B. Lazareff, J.-B. Le Bouquin, J. D. Monnier, F. Soulez and É. Thiébaut. “SPARCO: a semi-parametric approach for image reconstruction of chromatic objects,” *Astron. Astrophys.*, 564, A80–11, 2014.
- [26] F. Soulez and É. Thiébaut. “An image reconstruction framework for polychromatic interferometry,” In *International Colloquium at haute provence observatory, St Michel l’observatoire. France*, sep 2013.
- [27] É. Thiébaut, F. Soulez and L. Denis. “Exploiting spatial sparsity for multiwavelength imaging in optical interferometry,” *Journal of the Optical Society of America A*, 30, 1607170, February 2013.
- [28] F. Millour, A. Meilland, O. Chesneau, P. Stee, S. Kanaan, R. Petrov, D. Mourard and S. Kraus. “Imaging the spinning gas and dust in the disc around the supergiant A[e] star HD 62623,” *Astron. Astrophys.*, vol. 526, no. 14, article A107, 8 pages, 2011.
- [29] S. Boyd, N. Parikh, E. Chu, B. Peleato and J. Eckstein. “Distributed optimization and statistical learning via the alternating direction method of multipliers,” *Found. Trends Mach. Learn.*, 3(1):1–122, January 2011.
- [30] T. A. Pauls, J. S. Young, W. D. Cotton and J. D. Monnier. “A data exchange standard for optical (visible/ir) interferometry,” The Publications of the Astronomical Society of the Pacific, Vol. 117, No. 837, 1255–1262, 2005.
- [31] É. Thiébaut and J.-F. Giovannelli. Image reconstruction in optical interferometry. *IEEE Signal Processing Magazine*, 27:97–109, January 2010.
- [32] J.A. Fessler and B.P. Sutton. “Nonuniform fast fourier transforms using min-max interpolation,” *IEEE Transactions on Signal Processing*, 51(560–574), 2003.
- [33] G. Le Besnerais, S. Lacour, L.M. Mugnier, E. Thiebaut, G. Perrin and S. Meimon. “Advanced imaging methods for long-baseline optical interferometry,” *Selected Topics in Signal Processing, IEEE Journal of*, 2(5):767–780, Oct 2008.
- [34] R. G. Petrov, F. Malbet, G. Weigelt, P. Antonelli, U. Beckmann, Y. Bresson, A. Chelli, M. Dugué, G. Duvert, S. Gennari, L. Glck, P. Kern, S. Lagarde, E. Le Coarer, F. Lisi, F. Millour, K. Perraut, P. Puget, F. Rantakyro, S. Robbe-Dubois, A. Roussel, P. Salinari, E. Tatulli, G. Zins, M. Accardo, B. Acke, K. Agabi, E. Altariba, B. Arezki, E. Aristidi, C. Baffa, J. Behrend, T. Blöcker, S. Bonhomme, S. Busoni, F. Cassaing, J.-M. Clausse, J. Colin, C. Connot, A. Delboulbé, A. Domiciano de Souza, T. Driebe, P. Feautrier, D. Ferruzzi, T. Forveille, E. Foshat, R. Foy, D. Fraix-Burnet, A. Gallardo, E. Giani, C. Gil, A. Glentzlin, M. Heiden, M. Heininger, O. Hernandez Utrera, K.-H. Hofmann, D. Kamm, M. Kiekebusch, S. Kraus, D. Le Contel, J.-M. Le Contel, T. Lesourd, B. Lopez, M. Lopez, Y. Magnard, A. Marconi, G. Mars, G. Martinot-Lagarde, P. Mathias, P. Mge, J.-L. Monin, D. Mouillet, D. Mourard, E. Nussbaum, K. Ohnaka, J. Pacheco, C. Perrier, Y. Rabbia, S. Rebattu, F. Reynaud, A. Richichi, A. Robini, M. Sacchetti, D. Schertl, M. Schöller, W. Solscheid, A. Spang, P. Stee, P. Stefanini, M. Tallon, I. Tallon-Bosc, D. Tasso, L. Testi, F. Vakili, O. von der Lhe, J.-C. Valtier, M. Vannier and N. Ventura “AMBER, the near-infrared spectro-interferometric three-telescope VLTI instrument,” *Astronomy and Astrophysics*, 464:1–12, March 2007.
- [35] R. G. Petrov, F. Millour, S. Lagarde, M. Vannier, S. Rakshit, A. Marconi and G. Weigelt “VLTI/AMBER differential interferometry of the broad-line region of the quasar 3C273,” In *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, volume 8445 of *Society of Photo-Optical Instrumentation Engineers*

- (*SPIE*) *Conference Series*, July 2012.
- [36] F. Millour. “Interférométrie différentielle avec AMBER,”. Université Joseph Fourier (Grenoble), Ph.D 2005.
 - [37] I. Tallon-Bosc, M. Tallon, É. Thiébaud, C. Béchet, G. Mella, S. Lafrasse, O. Chesneau, A. Domiciano de Souza, G. Duvert, D. Mourard, R. Petrov and M. Vannier. “LIT-pro: a model fitting software for optical interferometry,” *Proceedings paper of Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, 7013, July 2008.
 - [38] K. Mardia and P. Jupp. *Directional statistics*. Wiley series in probability and statistics. Wiley, 2000.
 - [39] D. Mary, S. Bourguignon, C. Theys and H. Lanteri. “Interferometric image reconstruction with sparse priors in union of bases,” In *ADA 6 - Sixth Conference on Astronomical Data Analysis*, September 2010.
 - [40] A. Tarantola. *Inverse Problem Theory and Methods for Model Parameter Estimation*. Society for Industrial and Applied Mathematics (SIAM Publications), 2005.
 - [41] E. Thiébaud. “Introduction to Image Reconstruction and Inverse Problems,” In R. Foy and F.C. Foy, editors, *NATO ASIB Proc. 198: Optics in astrophysics*, page 397, January 2005.
 - [42] S. Renard, É. Thiébaud and F. Malbet. “Image reconstruction in optical interferometry: benchmarking the regularization,” *Astron. Astrophys.*, 533:A64, September 2011.
 - [43] S.D. Babacan, R. Molina and A.K. Katsaggelos. “Variational bayesian blind deconvolution using a total variation prior,” *Image Processing, IEEE Transactions on*, 18(1):12–26, Jan 2009.
 - [44] P.L. Combettes and J.C. Pesquet. *Proximal Splitting Methods in Signal Processing*. Fixed-Point Algorithms for Inverse Problems in Science and Engineering, Springer, 185–212, 2011.
 - [45] N. Jacobson. *Basic Algebra I: Second Edition*. Dover Publications, 2009.
 - [46] M. Schmidt. March 2014. <http://www.di.ens.fr/~mschmidt>.
 - [47] G. Duvert, P. Berio and F. Malbet. “Aspro: a software to prepare observations with optical interferometers,” *Proc. SPIE*, 4844:295–299, 2002.
 - [48] A. Belu, É. Thiébaud, M. Ollivier, G. Lagache, F. Selsis and F. Vakili “A scene model of exosolar systems for use in planetary detection and characterisation simulations,” *A&A*, 476(2):1005–1018, 2007.
 - [49] R. L. Kurucz. “Model atmospheres for g, f, a, b, and o stars,” *Astrophysical Journal Supplement Series*, 40:1–340, 1979.
 - [50] W. Van Hamme. “New limb-darkening coefficients for modeling binary star light curves,” *AJ*, 106:2096–2117, November 1993.
 - [51] P.R. Lawson, W.D. Cotton, C.A. Hummel, J.D. Monnier, M. Zhao, J.S. Young, H. Thorsteinsson, S. Meimon, L. Mugnier, G. Le Besnerais, E. Thiébaud and P.G. Tuthill. “An interferometry imaging beauty contest,” *Proc. SPIE*, 5491:886–899, 2004.