

Nonlinear Sculpturing of Optical Pulses in Fibre Systems

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Propagation in normally dispersive fiber is governed by the nonlinear Schrödinger equation (NLSE) :

2 \sim 0 $2 \vee \vee$ 2 ∂t^2 17 $i\frac{\partial \mathcal{L}}{\partial \mathcal{L}} - \frac{\mathcal{L}^2}{\partial \mathcal{L}} \frac{\partial \mathcal{L}}{\partial \mathcal{L}} + \gamma \|\psi\|^2 \|\psi = 0$ group ve *z* $2 \theta t^2$ $\theta \psi$ β_2 $\partial^2 \psi$ β_1 β_2 β_3 β_4 β_5 $\gamma \mid \psi \mid^{\mathsf{T}} \psi = 0$ group ve ∂t^2 and ∂t^2 $-\frac{1}{2} - \frac{1}{2} + \gamma |U|^2 |U = 0$ orgun y ∂z 2 ∂t^2

 2 and 2 and 2

 $\partial^2 w$, ∂

group velocity dispersion Kerr nonlinearity

In recent years, there has been a growing interest from the photonics community in the generation of non-conventional optical waveforms because of their applications in all-optical signal processing and microwave signal manipulation. While sinusoidal, Gaussian and hyperbolic secant intensity profiles are now routinely produced by modulators or mode-locked lasers, other signal waveforms such as parabolic, triangular or flat-top pulse shapes remain rather hard to synthesize. The interplay among the effects of dispersion, nonlinearity and gain/loss in optical fibre systems can be efficiently used to shape the pulses and manipulate and control the light dynamics and, hence, lead to different pulse-shaping regimes. In particular, it has been demonstrated that the temporal and spectral properties of a pulse can be tailored by use of such a nonlinear shaping approach to generate ultra-short parabolic or triangular pulses [1, 2].

6 degrees of freedom : (*C*⁰ *, T*⁰ *, P*⁰ *,* b² *, , L*) input pulse

 $\partial \pmb{\mathcal{U}} \qquad \pmb{\beta}_2 \;\; \partial^2 \pmb{\mathcal{U}}$

However, achieving a precise waveform with various prescribed characteristics, such as temporal duration and chirp properties, is a complex issue that requires careful choice of the initial pulse conditions and system parameters. The general problem of optimization towards a target operational regime in a complex multi-parameter space can be intelligently addressed by implementing machine-learning strategies. In this paper, we discuss a novel approach to the characterization and optimization of nonlinear shaping in fibre systems, in which first we perform numerical simulations of the basic propagation model to identify the relevant parameters, and then we use the machine-learning method of neural networks (NNs) to make predictions across a larger range of the data domain.

Due to the longitudinal evolution of g (requiring the solution of steady-state rate equations for an effective two-level system at each longitudinal step), direct simulations are time consuming.

We explore the space of operating states that can be accessed through control of the total pump power $P_T\!=\!P_1\!\!+\!\!P_2$, the fractional pump power $R=P_1/P_T^{}$ and the laser output coupling ratio β .

INTRODUCTION

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[2] S. Boscolo, A. I. Latkin, and S. K. Turitsyn, "Passive nonlinear pulse shaping in normally dispersive fiber systems," IEEE J. Quantum Electron. 44, 1196-1203 (2008).

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[4] S. Kobtsev, A. Ivanenko, A. Kokhanovskiy, and S. Smirnov, "Electronic control of different generation regimes in mode-locked all-fibre F8 laser," Laser Phys. Lett. 15, 045102 (2018).

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NUMERICAL MODEL

NUMERICAL MODEL

A NN regression can reproduce the observed trends.

We numerically characterize the performance of the figure-8 laser demonstrated in [4], based on a dual-pump nonlinear amplifying loop mirror.

Surfaces in the three-dimensional parameter space that enable: (a) the formation of a pulse with a misfit parameter to a parabolic shape $M \leq 0.045$ (b) an output temporal duration of $T_{out} = 10 T_{in}$

IN PASSIVE NORMALLY DISPERSIVE FIBER IN A DUAL-PUMP NONLINEAR AMPLIFYING LOOP MIRROR MODE-LOCKED ALL-FIBRE LASER

- two segments of ytterbium doped fibre are used.
- normally dispersive fibres are included.
- the output coupling ratio can be changed.
- We use a lumped model for the laser, in which propagation in the fibres follows a generalized NLSE:

 $i\psi_z - \frac{1}{2}\beta_2\psi_{tt} + \gamma |\psi|^2 \psi = \frac{i}{2}g(z,\omega)\psi$

The gain *g* includes saturation and spectral response

DIRECT NUMERICAL SIMULATIONS

Using normalized quantities reduces the level of complexity. $i \frac{1}{2} - \frac{1}{2} - \frac{1}{2} + N^2 |u|^2 u$ (*C*, *N*, *ξ*) : normalized input chirp, soliton number, normalized propagation length

2

2

 $- - + N^2 |u|$ $u =$

2 1^{2}

1

 $\rm 0$

 ∂u 1 ∂

OE 20

2

 $\partial \xi = 2 \partial \tau$

u u

Output pulse properties are quantified using the full-width at half maximum temporal and spectral widths, the Strelh ratio, a misfit factor or the kurtosis excess.

SEARCH FOR THE OPTIMUN WORKING PARAMETERS

The three-dimensional space of normalized quantities can be visualized using 3D plots.

The intersections of different surfaces provide the means

The reduced-size data set used to train the NN enables a reduction of the CPU time by a factor of 16 relative to the generation of the full data set.

The predictions from the NN are found to be sufficiently accurate for the targeted application.

