

First-principles calculation of the effects of tetragonal distortions on the Gilbert damping parameter of Co 2 MnSi

Barthélémy Pradines, Rémi Arras, Lionel Calmels

► To cite this version:

Barthélémy Pradines, Rémi Arras, Lionel Calmels. First-principles calculation of the effects of tetragonal distortions on the Gilbert damping parameter of Co 2 MnSi. AIP Advances, 2017, 7 (5), pp.056009. 10.1063/1.4974525 . hal-02084308

HAL Id: hal-02084308 https://hal.science/hal-02084308

Submitted on 18 Jun 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

First-principles calculation of the effects of tetragonal distortions on the Gilbert damping parameter of Co₂MnSi

Cite as: AIP Advances **7**, 056009 (2017); https://doi.org/10.1063/1.4974525 Submitted: 23 September 2016 . Accepted: 28 October 2016 . Published Online: 17 January 2017

B. Pradines, R. Arras, and L. Calmels





ARTICLES YOU MAY BE INTERESTED IN

Effect of Fe substitution on the structural, magnetic and electron-transport properties of half-metallic Co₂TiSi

AIP Advances 7, 055812 (2017); https://doi.org/10.1063/1.4974281

Increased magnetic damping in ultrathin films of Co₂FeAl with perpendicular anisotropy Applied Physics Letters **110**, 252409 (2017); https://doi.org/10.1063/1.4989379

Epitaxial strain controlled magnetocrystalline anisotropy in ultrathin FeRh/MgO bilayers AIP Advances 7, 055914 (2017); https://doi.org/10.1063/1.4974059

AVS Quantum Science

Co-published with AIP Publishing



AIP Advances 7, 056009 (2017); https://doi.org/10.1063/1.4974525 © 2017 Author(s). Coming Soon!



First-principles calculation of the effects of tetragonal distortions on the Gilbert damping parameter of Co₂MnSi

B. Pradines, R. Arras, and L. Calmels

CEMES-CNRS and Université de Toulouse, 29 rue Jeanne Marvig, F-31055 Toulouse, France

(Presented 1 November 2016; received 23 September 2016; accepted 28 October 2016; published online 17 January 2017)

We present an *ab initio* study of the influence of the tetragonal distortion, on the static and dynamic (Gilbert damping parameter) magnetic properties of a Co₂MnSi crystal. This tetragonal distortion can for instance be due to strain, when Co₂MnSi is grown on a substrate with a small lattice mismatch. Using fully relativistic Korringa-Kohn-Rostoker (KKR) calculations, in conjunction with the coherent potential approximation (CPA) to describe atomic disorder and the linear response formalism to compute the Gilbert damping parameter, we show that a tetragonal distortion can substantially change the properties of Co₂MnSi, in a way which depends on the kind of atomic disorder. © 2017 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (http://creativecommons.org/licenses/by/4.0/). [http://dx.doi.org/10.1063/1.4974525]

I. INTRODUCTION

The full-Heusler alloy Co₂MnSi has recently been a subject of intense research activity in the spintronics community. Many theoretical studies have highlighted the half-metallic behavior^{1–3} and the very low Gilbert damping parameter (lower than 1×10^{-3})⁴ of the perfectly ordered phase of this alloy. These properties, with its high measured Curie temperature of 985 K,⁵ make this material a perfect candidate for being used in spintronic applications.

Unfortunately, experimental measurements have shown less impressive magnetic properties than expected from numerical studies: apart from recent measurements reporting a spin-polarization at the Fermi level of 93%⁶ or a damping parameter of 1.5×10^{-4} ,⁷ most of the experimental results only found values of the spin-polarization around 50-60%^{8–11} and α is in general around 4×10^{-3} ,¹² more than 10 times higher than the calculated values.

The presence of crystalline defects in the samples was put forward to explain the qualitative difference between measured and calculated properties. Neutron diffraction experiments on thin films confirmed the presence of atomic disorder with mainly Co/Mn swaps.^{13,14} Theoretical studies have examined the effects of defects on the static magnetic properties.² The torque correlation model (TCM) was used in 2015 to study the impact of atomic disorder on the Gilbert damping parameter for the fully disordered A2 and B2 phases of cubic Co₂MnSi,³ but a disorder resulting from Co/Mn swaps (D0₃-like disorder) has not been considered while it was detected experimentally. Moreover, other defects, like tetragonal deformations, may contribute to the reduction of the spin-polarization at the Fermi level E_F and to the increase of α .

In this paper, we use first-principles methods to calculate, for four crystal phases (nearly- L_1 , B2, D0₃, A2) of Co₂MnSi, the influence on the static magnetic properties and Gilbert damping parameter of a tetragonal distortion, which may be due to the epitaxial growth of this alloy on a substrate with a small lattice mismatch.

II. CALCULATION DETAILS

All the results presented here were obtained using the SPR-KKR package.¹⁵ Based on the KKR multiple scattering formalism, this code enables to calculate the one-electron Green's function of

© Author(s) 2017



defective crystals, describing atomic disorder with the coherent potential approximation (CPA).¹⁶ The main static properties of the system (ground states energy, magnetic moments, density of states) are directly accessible from the calculated Green's function. Moreover, a parameter-free and fully *ab initio* method based on the linear response formalism¹⁷ was implemented to extract the Gilbert damping parameter from the Green's function.

The crystal structure of Co₂MnSi in its most ordered L2₁ phase is described by the space groupe $Fm\overline{3}m$, where the Co, Mn and Si atoms occupy respectively the 8c, 4a and 4b Wyckoff positions. In this paper we have also considered: the B2 phase, where Mn and Si atoms are randomly distributed on the 4a and 4b atomic sites; the D0₃ phase, where Co and Mn atoms are randomly distributed on the 8c and 4a sites; and the A2 phase where the three atomic species are randomly shared between all the atomic sites.

Since the Gilbert damping parameter cannot be calculated at T = 0 K with the linear response formalism for a perfectly ordered phase, the results are not presented for a perfectly ordered L2₁, but for a nearly L2₁ phase containing a 2% amount of B2-like disorder. We have checked that no significant differences between the L2₁ and the nearly L2₁ phases were found for the static magnetic properties.

All the calculations were performed with the local spin-density approximation (LSDA) exchangecorrelation potential of Vosko, Wilk and Nusair¹⁸ with the fully-relativistic mode. We sampled the irreducible part of the first Brillouin zone with 3000 k vectors and set up the selected *l*-cut-off l_{max} of the KKR formalism to 4. A volume optimization for the cubic phase was made and the minimum value of the ground state energy was found for $a_0 = 5.48$ Å. To account for the tetragonal distortions, we constrained the lattice parameters $a_{\parallel} = a_{[100]} = a_{[010]}$ of the tetragonally distorted unit cell of Co₂MnSi to values which can deviate from the calculated lattice constant a_0 of cubic Co₂MnSi: $-5\% \le \varepsilon_{\parallel} = \frac{a_{\parallel}-a_0}{a_0} \le 5\%$, where ε_{\parallel} is the in-plane strain. Then, we calculated the equilibrium lattice parameter $a_{\perp} = a_{[001]}$ by minimizing the total energy. The ratio a_{\perp}/a_{\parallel} decreases almost linearly between 1.12 and 0.91 when ε_{\parallel} varies from -5% to +5%. These values were calculated for the L2₁ phase and were kept for the B2, D0₃ and A2 phases.

III. RESULTS

The magnetization (*M*) of Co₂MnSi is plotted on Fig. 1.a, as a function of ε_{\parallel} and for the different phases of the crystal. Even without compression, chemical disorder causes large variations of the magnetization, especially for disorders involving Co/Mn swaps like in the DO₃ or the A2 phases. It goes from $\approx 120 \ \mu_B$.nm⁻³ for the nearly L2₁ and the B2 phases to respectively $\approx 65 \ \mu_B$.nm⁻³ and $\approx 85 \ \mu_B$.nm⁻³ for the DO₃ and the A2 phases. The effect of the tetragonal deformation generally tends to reduce even more the magnetization. It decreases for example from $\approx 120 \ \mu_B$.nm⁻³ to $\approx 100 \ \mu_B$.nm⁻³ for $\varepsilon_{\parallel} = -5\%$ for the nearly L2₁ phase, mostly because of a decrease of the spin magnetic moment of the Co atoms. For the DO₃ phase, the magnetization goes from $\approx 65 \ \mu_B$.nm⁻³ ($\varepsilon_{\parallel} = -5\%$), this decrease being due to the decrease of the spin magnetic moment of the Co and Mn atoms in 8c sites. For the A2 phase, the decrease of the magnetization resulting from a compression is due to a decrease of the spin magnetic moments of all the chemical species.



FIG. 1. Impact of the tetragonal distortion on the properties of the 4 phases of Co₂MnSi: a) on the magnetization M, b) on the total density of states at the Fermi level Z_F and c) on the Gilbert damping parameter α .

The total density of states (DOS) at the Fermi level Z_F is plotted on Fig. 1.b. As previously explained, the cubic nearly-L2₁ phase being half-metallic, the DOS at E_F is only due to majority spin electrons when $\varepsilon_{\parallel} = 0$, and increases when $\varepsilon_{\parallel} \neq 0$. The half-metallicity of Co₂MnSi is destroyed both by chemical disorder and by tetragonal distortions. Electron states on both sides of the minority spin band gap of the $L2_1$ phase being mainly built from Co d orbitals, disorder-induced modifications of Z_F for cubic crystals ($\varepsilon_{\parallel} = 0$) will be more important for the disordered phases that change the chemical species in the first coordination shell of Co atoms (A2 and DO_3 disorders). This explains why Z_F is higher for disordered cubic crystals (in particular for the A2 and D0₃ phases) than for the L_{2_1} cubic phase. For non-cubic crystals, the tetragonal distortion lifts the degeneracy between the two e_g and the 3 t_{2g} Co atom orbitals. The minority spin d_z^2 orbital of Co atoms, which was unoccupied for the cubic crystal, becomes partly occupied when $\varepsilon_{\parallel} < 0$. Conversely, when $\varepsilon_{\parallel} > 0$, the Co d_z^2 orbital is shifted towards higher energies while the Co $d_{x^2-y^2}$ becomes partly occupied. This explains why Z_F is higher for tetragonally distorted ($\varepsilon_{\parallel} \neq 0$) ordered L2₁ crystals than for the cubic ($\varepsilon_{\parallel} = 0$) ordered L2₁ crystal (we checked this last point by calculating partial DOS-curves for tetragonally distorted crystals in the L2₁ phase, with the code Wien2k¹⁹). The variations of Z_F as a function of ε_{\parallel} are stronger for the nearly-L2₁, the B2 and the D0₃ phases.

A quick analysis of Fig. 1.c shows that any biaxial compression or tension increases the value of α . Except for the A2 phase, this increase is generally higher for $\varepsilon_{\parallel} = 5\%$ than for $\varepsilon_{\parallel} = -5\%$. The effects of the tetragonal distortion are more important for the nearly ordered phase (L2₁+2%B2) for which α is 15 times higher for $\varepsilon_{\parallel} = 5\%$ than for the cubic crystal; they are also rather important for the B2 phase (α is more than 4 times higher for $\varepsilon_{\parallel} = 5\%$), moderate for the A2 phase (α is 2.5 times higher for $\varepsilon_{\parallel} = -5\%$), and negligible for the D0₃ phase.

The comparison of the different panels of Fig. 1 shows that the increase of the Gilbert damping parameter of a strained disordered layer (with respect to the value calculated for the cubic nearly ordered phase), is due both to the tetragonal distortion and to the alloy disorder when Co_2MnSi is in the B2 phase, but mostly to the alloy disorder when it is in the A2 or in the D0₃ phases.

To understand the origin of these variations of α , we used an analytical model proposed by Kamberský, based on the two following contributions:²⁰ $\alpha_{sf} = \frac{\pi\gamma\hbar^2}{\mu_0 M} Z_F (\frac{(g-2)^2}{\tau})^2$ and $\alpha_o = \frac{\pi\gamma}{\mu_0 M} Z_F \lambda_{SO}^2$ $(g-2)^2 \tau$, where α_{sf} accounts for the spin-flip scattering contributions and α_o describes the ordinary scattering. In these equations, γ is the gyromagnetic ratio, λ_{SO} the spin-orbit coupling parameter, g the Landé factor and τ the electron scattering time. M and Z_F have the same meaning than before. The variations of α are roughly the same as those of Z_F for the nearly L2₁ and the B2 phases: Z_F can be considered as a measure of the number of available relaxation channels and if Z_F is high, the magnetization precession will relax faster. For the D0₃ and the A2 cases, we also have to consider the high variations of the magnetization to understand the variations of α ($\alpha \propto \frac{1}{M}$ in the Kamberský model). For the D0₃ phase, Z_F and M both decrease with in-plane compression, balancing each other with nearly constant values of α . For the A2 phase, the DOS at E_F does not change with tetragonal distortions, but the magnetization decreases with compression, resulting in an increase of α .

IV. CONCLUSION

We have shown that tetragonal distortions and alloy disorder can substantially change the magnetic properties of Co₂MnSi. Tetragonal distortions are responsible for a strong increase of the Gilbert damping parameter for the nearly-L2₁ or the B2 phases, since they have a very low value of α in the cubic structure. It is thus important to avoid any tetragonal distortions to preserve the lowest possible value of α . In comparison, since they already have a high value of α for the cubic cell, the changes induced by the tetragonal deformation are less critical for the D0₃ or A2 phases. The priority must be given to a straight control of atomic disorder during the growth of samples, in order to avoid these last phases. The effects of tetragonal distortion and alloy disorder on the physical properties of Co₂MnSi could be analyzed experimentally by performing photoemission spectroscopy (to measure Z_F), magnetometry (to measure M) and ferromagnetic resonance experiments (to measure α) on thin Co₂MnSi layers strained on substrates with different values of the lattice mismatch, and in which alloy disorder could be changed, for instance by irradiation with light ions.²¹

ACKNOWLEDGMENTS

This work was granted access to the HPC resources of CALMIP supercomputing center under the allocation p1252 (2014-2016).

- ¹S. Fujii, S. Sugimura, S. Ishida, and S. Asano, Journal of Physics: Condensed Matter 2, 8583 (1990).
- ² S. Picozzi, A. Continenza, and A. J. Freeman, Physical Review B 69, 094423 (2004).
- ³ A. Sakuma, Journal of Physics D: Applied Physics 48, 164011 (2015).
- ⁴C. Liu, C. K. A. Mewes, M. Chshiev, T. Mewes, and W. H. Butler, Applied Physics Letters 95, 022509 (2009).
- ⁵ P. J. Webster, Journal of Physics and Chemistry of Solids **32**, 1221 (1971).
- ⁶ M. Jourdan, J. Minr, J. Braun, A. Kronenberg, S. Chadov, B. Balke, A. Gloskovskii, M. Kolbe, H. J. Elmers, G. Schnhense, H. Ebert, C. Felser, and M. Klui, Nature Communications 5, 3974 (2014).
- ⁷ Q. Shi-Zhu, Z. Jie, Q. Yu-Feng, H. Run-Run, Z. Hai, Z. Da-Peng, K. Yun, K. Shi-Shou, Y. Shu-Yun, H. Guang-Bing, Y. Shi-Shen, and M. Liang-Mo, Chinese Physics Letters **32**, 057601 (2015).
- ⁸ Y. Sakuraba, M. Hattori, M. Oogane, Y. Ando, H. Kato, A. Sakuma, T. Miyazaki, and H. Kubota, Applied Physics Letters **88**, 192508 (2006).
- ⁹ M. P. Raphael, B. Ravel, M. A. Willard, S. F. Cheng, B. N. Das, R. M. Stroud, K. M. Bussmann, J. H. Claassen, and V. G. Harris, Applied Physics Letters **79**, 4396 (2001).
- ¹⁰ J. Schmalhorst, A. Thomas, S. Kmmerer, O. Schebaum, D. Ebke, M. D. Sacher, G. Reiss, A. Htten, A. Turchanin, A. Glzhuser, and E. Arenholz, Physical Review B 75, 014403 (2007).
- ¹¹ R. Fetzer, J.-P. Wstenberg, T. Taira, T. Uemura, M. Yamamoto, M. Aeschlimann, and M. Cinchetti, Physical Review B 87, 184418 (2013).
- ¹² M. Oogane, T. Kubota, H. Naganuma, and Y. Ando, Journal of Physics D: Applied Physics 48, 164012 (2015).
- ¹³ M. P. Raphael, B. Ravel, Q. Huang, M. A. Willard, S. F. Cheng, B. N. Das, R. M. Stroud, K. M. Bussmann, J. H. Claassen, and V. G. Harris, *Physical Review B* 66, 104429 (2002).
- ¹⁴ B. Ravel, M. P. Raphael, V. G. Harris, and Q. Huang, Physical Review B 65, 184431 (2002).
- ¹⁵ H. Ebert, "The Munich SPR-KKR package, version 5.4."
- ¹⁶ P. Soven, Physical Review **156**, 809 (1967).
- ¹⁷ H. Ebert, S. Mankovsky, D. Kdderitzsch, and P. J. Kelly, *Physical Review Letters* **107**, 066603 (2011).
- ¹⁸ S. H. Vosko, L. Wilk, and M. Nusair, Canadian Journal of Physics 58, 1200 (1980).
- ¹⁹ P. Blaha, K. Schwarz, P. Sorantin, and S. B. Trickey, Computer Physics Communications 59, 399 (1990).
- ²⁰ V. Kamberský, Canadian Journal of Physics 48, 2906 (1970).
- ²¹ I. Abdallah, N. Ratel-Ramond, C. Magen, B. Pecassou, R. Cours, A. Arnoult, M. Respaud, J. F. Bobo, G. BenAssayag, E. Snoeck, and N. Biziere, Materials Research Express 3, 046101 (2016).