



HAL
open science

A Class of Random Walks in Reversible Dynamic Environments: Antisymmetry and Applications to the East Model

Luca Avena, Oriane Blondel, Alessandra Faggionato

► **To cite this version:**

Luca Avena, Oriane Blondel, Alessandra Faggionato. A Class of Random Walks in Reversible Dynamic Environments: Antisymmetry and Applications to the East Model. *Journal of Statistical Physics*, 2016, 165 (1), pp.1-23. 10.1007/s10955-016-1596-7. hal-02072389

HAL Id: hal-02072389

<https://hal.science/hal-02072389v1>

Submitted on 4 Sep 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

A CLASS OF RANDOM WALKS IN REVERSIBLE DYNAMIC ENVIRONMENTS: ANTISYMMETRY AND APPLICATIONS TO THE EAST MODEL

L. AVENA, O. BLONDEL, AND A. FAGGIONATO

ABSTRACT. We introduce via perturbation a class of random walks in reversible dynamic environments having a spectral gap. In this setting one can apply the mathematical results derived in [2]. As first results, we show that the asymptotic velocity is antisymmetric in the perturbative parameter and, for a subclass of random walks, we characterize the velocity and a stationary distribution of the environment seen from the walker as suitable series in the perturbative parameter. We then consider as a special case a random walk on the East model that tends to follow dynamical interfaces between empty and occupied regions. We study the asymptotic velocity and density profile for the environment seen from the walker. In particular, we determine the sign of the velocity when the density of the underlying East process is not $1/2$, and we discuss the appearance of a drift in the balanced setting given by density $1/2$.

1. INTRODUCTION

In [2] we studied continuous-time random walks in dynamic random environments on the d -dimensional integer lattice, $d \geq 1$, in a perturbative regime. More precisely, we considered a stationary Feller Markov process, playing the role of the environment and satisfying the Poincaré inequality¹. In addition, we considered a random walk with transition rates given by functions of the (autonomously) evolving environment. The main assumption required that the random walk is a small perturbation either of an homogeneous or of a “stationary” walk, the latter meaning that the environment viewed from the walker has the same stationary distribution as the environment itself. In this setting, we characterized the ergodic behavior of the environment viewed from the walker, and we derived a law of large numbers (i.e. existence of an asymptotic non random velocity) and an invariance principle (i.e. gaussian fluctuations under diffusive rescaling) for the random walk. One main tool there was the derivation and a careful analysis of a series expansion of Dyson–Phillips type for the semigroup associated with the environment as seen from the walker. We review in Section 2 the main results of [2] that will be used in the rest of the paper (cf. in particular Theorem 2.2 below).

We aim here to illustrate how the results of [2] can give non-trivial information about random walks in dynamic random environments, beyond their diffusive behavior. Environments here will be reversible stochastic particle systems on \mathbb{Z}^d with a positive spectral gap (in particular, the environment at a given time is a configuration in $\{0, 1\}^{\mathbb{Z}^d}$). We first introduce in Section 3 a class of random walks with

¹For a reversible process, the Poincaré inequality is equivalent to the positive spectral gap of the generator.

transition rates satisfying suitable algebraic identities and show a hidden antisymmetry relation in the asymptotic velocity (cf. Theorem 3.1). More precisely, if $v(\varepsilon)$ is the asymptotic velocity at perturbative parameter ε , we have the antisymmetry relation

$$v(\varepsilon) = -v(-\varepsilon).$$

As discussed in Sections 3 and 4, this is not a consequence of obvious symmetries in the system. A special example of random walk in the above mentioned class is given by what we call “ ε -RW”, a one-dimensional random walk with drift 2ε (resp. -2ε) on top of particles (resp. on empty sites). This type of random walk (with different drifts on top of empty/occupied sites) has been recently studied in [3, 4, 5, 6, 19, 20, 21, 22, 27] for different choices of environments. The interest is due to the fact that (in case of opposing drifts) it represents one of the simplest example of a random walk with space-time inhomogeneous random transitions that can give rise to some slow-down or “trapping” effects (cf. [7]) similar to the well known phenomenology in 1-dimensional static random environment [11, 30, 31]. Under our assumptions, the trapping effect does not occur on the diffusive scale, but the ε -RW favors spending time oscillating between a particle and a hole, and so tends to lie at interfaces between occupied/empty regions. Its behavior is therefore connected to space-time correlations in the environment, which can be difficult to grasp. We first derive two main results for generic ε -RW’s: a deeper analysis of the series expansion for its asymptotic velocity $v(\varepsilon)$ (see Proposition 3.2) and for the limiting distribution of the environment viewed from the walker (see Proposition 3.4).

We then study in more detail the ε -RW on the East model. The latter has been introduced in the physics literature as a simplified model for glassy systems [23], and belongs to the class of kinetically constrained model [8]. It has received much attention within the physics and mathematics communities, since it catches some relevant features of glassy dynamics as e.g. aging, dynamical heterogeneity, huge relaxation times (cf. [14, 15, 16, 17, 29] and references therein). Of particular interest to both the physics and mathematics communities is the structure of the space-time correlated “bubbles” of occupied sites (see Figure 2) and tracing the ε -RW on the East model allows to catch some information on these bubbles. We stress that the East model has a positive spectral gap [1, 14] but does not display any uniform mixing property or attractiveness. Therefore we can only use the results of [2] and not for instance those of [4, 5, 28].

For the ε -RW on the East model, we discuss evidence of a negative asymptotic drift in the balanced case of density $1/2$ and give two theoretical results supporting this fact in addition to simulations (cf. Propositions 4.5 and 4.6). It is tempting to interpret the sign of the asymptotic velocity as a signature of the orientation of the East model, but indeed we can show that the velocity remains the same if one replace the East model by the West model, which has the opposite orientation (cf. Corollary 4.4). Finally, in Corollary 4.7 we give a detailed analysis of the density profile of the limiting distribution of the East model viewed from the ε -RW.

Let us notice that the study of the ε -RW on the East model was partially inspired by [24], where the authors consider random walks on the FA1f model (the symmetric version of the East model). An investigation based on the expansion derived in [2] might be performed as well for other types of random walks as e.g. the ones considered in [24]. A further study of random walks in kinetically constrained models is given in [13].

Finally, we mention that the negative drift for the ε -RW on the East model and our qualitative analysis of the density profile of the East model viewed from the ε -RW are supported by numerical simulations performed by Philip Thomann.

2. PRELIMINARIES

2.1. Definitions. We consider a Feller Markov process $(\xi_t)_{t \geq 0}$ on $\Omega := \{0, 1\}^{\mathbb{Z}^d}$ with generator \mathcal{L}_{env} which can be thought of as an interacting particle system, playing the role of dynamic random environment. For $(x, t) \in \mathbb{Z}^d \times [0, \infty)$, if $\xi_t(x) = 1$ we say that there is a particle at time t at position x , else site x is considered empty at time t (equivalently, there is a hole at x at time t). See [26] for a standard reference on this type of Markov processes. We write $\mathbb{E}_\xi^{\text{env}}$ for the expectation w.r.t. the law of the environment started from ξ and τ_x for the translation operator on \mathbb{Z}^d such that $\tau_x \xi(y) = \xi(x + y)$ for $x, y \in \mathbb{Z}^d, \xi \in \Omega$.

Assumption 1. *We assume the following properties for the dynamic environment:*

- (i) *(Reversibility)* $(\xi_t)_{t \geq 0}$ admits a reversible translation invariant probability measure μ on Ω .
- (ii) *(Positive spectral gap)* The generator L_{env} has 0 as simple eigenvalue and the rest of its spectrum is in $[\gamma, +\infty)$ for some $\gamma > 0$.
- (iii) *The Markov semigroup $S_{\text{env}}(t)$, with $S_{\text{env}}(t)f(\xi) := \mathbb{E}_\xi^{\text{env}}[f(\xi_t)]$, commutes with spatial translations, i.e. $S_{\text{env}}(t)(f \circ \tau_x) = (S_{\text{env}}(t)f) \circ \tau_x$ μ -a.s. for any local function f and $x \in \mathbb{Z}^d$.*

We point out that Assumption 1-(ii) is equivalent to the so-called Poincaré inequality: $\gamma \|f\|^2 \leq -\mu(f \mathcal{L}_{\text{env}} f)$ for all $f \in \mathcal{D}(\mathcal{L}_{\text{env}})$ with $\mu(f) = 0$, where $\mathcal{D}(\mathcal{L}_{\text{env}})$ denotes the domain of the operator \mathcal{L}_{env} .

We interpret the process $(\xi_t)_{t \geq 0}$ as a dynamic random environment for a continuous-time random walk $(X_t^{(\varepsilon)})_{t \geq 0}$ on \mathbb{Z}^d which starts at the origin and that we now define. The rate for a jump from x to $x + y$ when the environment is equal to ξ will be denoted by $r_\varepsilon(y, \tau_x \xi)$. Here, ε is a perturbative parameter, whose precise meaning we explicit in Subsection 2.2 below.

Assumption 2. *We assume that for suitable functions $r(y, \eta) \geq 0$, $\hat{r}_\varepsilon(y, \eta)$ with finite support in η and finite range in y^2 , the jump rates admit the decomposition*

$$r_\varepsilon(y, \eta) = r(y, \eta) + \hat{r}_\varepsilon(y, \eta).$$

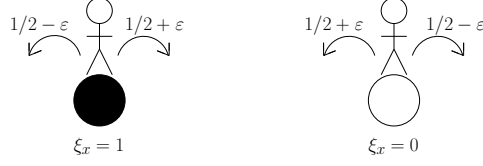
Moreover we assume that

$$r(y, \eta) = r(-y, \tau_y \eta). \tag{1}$$

Note that \hat{r}_ε should be considered as a perturbative contribution to the transition rates r_ε , so that $r(y, \eta)$ can be thought of as the transition rates for an unperturbed random walk. Then, since μ is translation invariant, the last assumption (1) is the detailed balance condition and is equivalent to the reversibility of μ for the environment seen from the unperturbed walker.

Due to dependence on the environment, such a random walk is not Markovian itself, but the joint process $(\xi_t, X_t^{(\varepsilon)})_{t \geq 0}$ on state space $\Omega \times \mathbb{Z}^d$ is a Markov process

²More precisely, we assume that there exists R such that for $|y| \geq R$, $r(y, \cdot) \equiv 0$ and $\hat{r}_\varepsilon(y, \cdot) \equiv 0$ and for all y , $r(y, \cdot), \hat{r}_\varepsilon(y, \cdot)$ have finite support.

FIGURE 1. A summary of the jump rates for the ε -RW.

with generator

$$\mathcal{L}_{\text{rwe}}^{(\varepsilon)} f(\xi, x) := \mathcal{L}_{\text{env}} f(\xi, x) + \sum_{y \in \mathbb{Z}^d} r_\varepsilon(y, \tau_x \xi) [f(\xi, x + y) - f(\xi, x)], \quad (2)$$

where the operator \mathcal{L}_{env} acts only on the first coordinate of f .

Later, we will consider more closely the following one-dimensional special case, which we call ε -RW.

Definition 2.1. (ε -RW) For $\xi \in \{0, 1\}^{\mathbb{Z}}$ and fixed $\varepsilon \in [-1/2, 1/2]$, set (see Figure 1)

$$r_\varepsilon(y, \xi) := \begin{cases} 1/2 + \varepsilon(2\xi(0) - 1) & \text{if } y = +1, \\ 1/2 - \varepsilon(2\xi(0) - 1) & \text{if } y = -1, \\ 0 & \text{otherwise.} \end{cases} \quad (3)$$

Here, the perturbative role of ε is clear, and in this case, the unperturbed random walk ($\varepsilon = 0$) is the simple symmetric random walk.

2.2. Environment seen by the walker. One of the most common approach to study random motion in random media is to analyze the so called *environment seen by the walker*, that is, the Markov process $(\eta_t^{(\varepsilon)})_{t \geq 0}$ with state space Ω defined by $\eta_t^{(\varepsilon)} := \tau_{X_t^{(\varepsilon)}} \xi_t$, with associated generator given by

$$\mathcal{L}_{\text{ew}}^{(\varepsilon)} f(\eta) := \mathcal{L}_{\text{env}} f(\eta) + \sum_{y \in \mathbb{Z}^d} r_\varepsilon(y, \eta) [f(\tau_y \eta) - f(\eta)], \quad \eta \in \Omega. \quad (4)$$

Note that the jumps of the walker in (2) turn into spatial-shifts for the environment seen by the walker. We write $(S_\varepsilon(t))_{t \geq 0}$ for the semigroup associated with this Markov process $(\eta_t^{(\varepsilon)})_{t \geq 0}$. When $\varepsilon = 0$ we simply write $S(t)$.

In the following theorem we recall some results from [2] that are relevant to our discussion. We set

$$\hat{L}_\varepsilon f(\eta) := \mathcal{L}_{\text{ew}}^{(\varepsilon)} f(\eta) - \mathcal{L}_{\text{ew}}^{(0)} f(\eta) = \sum_{y \in \mathbb{Z}^d} \hat{r}_\varepsilon(y, \eta) [f(\tau_y \eta) - f(\eta)]. \quad (5)$$

Due to our assumptions, \hat{L}_ε has bounded norm $\|\hat{L}_\varepsilon\|$ as operator in $L^2(\mu)$. For example, for the ε -RW the operator \hat{L}_ε is given by

$$\hat{L}_\varepsilon f(\eta) = \varepsilon(2\eta(0) - 1) [f(\tau_1 \eta) - f(\tau_{-1} \eta)]. \quad (6)$$

Note that in this case $\|\hat{L}_\varepsilon\|$ is bounded from above by 2ε .

As the reader will see, our results hold for ε such that $\|\hat{L}_\varepsilon\| < \gamma$ (cf. Assumption 1–(ii)). If, as in the examples discussed below, $\|\hat{L}_\varepsilon\| = C\varepsilon$, this trivially means that we restrict to ε small. Since interesting perturbations are not necessarily explicitly linear in the perturbative parameter, we keep the more general condition $\|\hat{L}_\varepsilon\| < \gamma$.

Theorem 2.2. [2] *Under Assumptions 1 and 2 and further assuming that $\|\hat{L}_\varepsilon\| < \gamma$, the following holds:*

- (i) *The process $(\eta_t^{(\varepsilon)})_{t \geq 0}$ admits a unique probability measure μ_ε which is invariant and absolutely continuous w.r.t. μ . Moreover, μ_ε is time ergodic. The distribution of $\eta_t^{(\varepsilon)}$ converges to μ_ε as $t \rightarrow \infty$ if the distribution of $\eta_0^{(\varepsilon)}$ is absolutely continuous w.r.t. μ .*
- (ii) *If in addition $r(y, \eta) > 0 \Rightarrow r_\varepsilon(y, \eta) > 0$, then μ_ε and μ are mutually absolutely continuous.*
- (iii) *The measure μ_ε admits the following representation: for every $f \in L^2(\mu)$*

$$\mu_\varepsilon(f) = \mu(f) + \sum_{n=0}^{\infty} \int_0^{\infty} \mu \left(\hat{L}_\varepsilon S_\varepsilon^{(n)}(s) f \right) ds, \quad (7)$$

where the integrals and the series are absolutely convergent and the operators $S_\varepsilon^{(n)}(s)$, $n \geq 0$, are defined iteratively as

$$S_\varepsilon^{(0)}(t)f := S(t)f, \quad S_\varepsilon^{(n+1)}(t)f := \int_0^t S(t-s) \hat{L}_\varepsilon S_\varepsilon^{(n)}(s) f ds. \quad (8)$$

Moreover, it holds

$$\left| \int_0^{\infty} \mu \left(\hat{L}_\varepsilon S_\varepsilon^{(n)}(s) f \right) ds \right| \leq (\|\hat{L}_\varepsilon\|/\gamma)^{n+1} \|f - \mu(f)\|. \quad (9)$$

- (iv) *For $\eta \in \Omega$, we introduce the local drift $j^{(\varepsilon)}(\eta) := \sum_y y r_\varepsilon(y, \eta)$ and set $v(\varepsilon) := \mu_\varepsilon(j^{(\varepsilon)})$. Then for μ_ε -a.e. ξ*

$$\frac{1}{t} X_t^{(\varepsilon)} \xrightarrow[t \rightarrow \infty]{} v(\varepsilon) \quad a.s. \quad (10)$$

In particular, $v(\varepsilon)$ can be written as

$$v(\varepsilon) = \mu(j^{(\varepsilon)}) + \sum_{n=0}^{\infty} \int_0^{\infty} \mu(\hat{L}_\varepsilon S_\varepsilon^{(n)}(s) j^{(\varepsilon)}) ds. \quad (11)$$

To obtain the above theorem, not all our assumptions are necessary. We refer the interested reader to Theorem 2 in [2], where a more general statement is given, and to (42) in [2] which allows to get (9). Strictly speaking, in the assumptions of Theorem 2.2 one should include that the environment process has a non pathological generator \mathcal{L}_{env} (see Prop. 3.1 in [2] for a precise statement), anyway this additional technical assumption is satisfied in all standard models. Finally, we point out that a perturbative characterization of the stationary distributions μ_ε is given in [25] in a different form.

3. A CLASS OF RW WITH AN ANTISYMMETRY PROPERTY

3.1. Antisymmetry relation for the velocity. We can now state a first new result given by an antisymmetry relation for the velocity of the random walks in dynamic environment introduced above. To this aim, we introduce some additional assumptions.

Assumption 3. *The following identities are satisfied:*

- (i) $\sum_y y r(y, \cdot) \equiv 0$
- (ii) $\hat{r}_\varepsilon = -\hat{r}_{-\varepsilon}$

(iii) \hat{r}_ε can be factorized into $\hat{r}_\varepsilon(y, \eta) = \alpha(y)\bar{r}_\varepsilon(\eta)$ with α antisymmetric, that is $\alpha(y) = -\alpha(-y)$.

As an example one may think of $\bar{r}_\varepsilon(\eta)$ as ε times the number of particles in a given neighborhood of the origin. In that case, for $\varepsilon > 0$, each particle in the environment that falls in the “vision field” of the walker favors jumps $X_t^{(\varepsilon)} \rightarrow X_t^{(\varepsilon)} + y$ when $\alpha(y) > 0$ and discourages them when $\alpha(y) < 0$ (and vice-versa for $\varepsilon < 0$). Another example (one-dimensional) is given by the ε -RW, which indeed satisfies both Assumptions 2 and 3.

Theorem 3.1. *Assume Assumptions 1, 2, 3 and $\|\hat{L}_\varepsilon\| < \gamma$, then*

$$v(-\varepsilon) = -v(\varepsilon). \quad (12)$$

We prove this theorem in Section 5.1. For the sake of clarity we restrict ourselves there to the case of the ε -RW, but the proof extends easily to the general case.

We point out that $X^{(-\varepsilon)}$ is not a time-reversed version of $X^{(\varepsilon)}$. Indeed the trajectories of these two processes are quite different in general (see Figure 4 for an illustration in the case of the ε -RW on the East model, which will be explained in detail below). In particular, Theorem 3.1 does not follow by taking the time reversion of $X^{(\varepsilon)}$. This is further discussed in Section 4.1.

Let us explain the difficulty behind Theorem 3.1. We first observe that Assumption 3 trivially implies $j^{(\varepsilon)}(\eta) = -j^{(-\varepsilon)}(\eta)$. This identity alone is not enough to prove the antisymmetry relation (12) since, due to Theorem 2.2, $v(\varepsilon) = \mu_\varepsilon(j^{(\varepsilon)})$, while $v(-\varepsilon) = \mu_{-\varepsilon}(j^{(-\varepsilon)})$. One could therefore ask whether the antisymmetry relation could be due to a possible equality of μ_ε and $\mu_{-\varepsilon}$. However, as illustrated for instance in (19) below, the two probability distributions $\mu_\varepsilon, \mu_{-\varepsilon}$ do not coincide in general.

3.2. Velocity and density profile for the ε -RW. When applying Theorem 2.2, we get the following more refined results in the case of the ε -RW.

Proposition 3.2. *Under Assumptions 1 and for $2|\varepsilon| < \gamma$, the asymptotic velocity $v(\varepsilon)$ of the ε -RW can be expressed as*

$$v(\varepsilon) = \begin{cases} 2\varepsilon(2\mu(\eta(0)) - 1) + O(\varepsilon^3) & \text{if } \mu(\eta(0)) \neq 1/2 \\ \varepsilon^3\kappa + O(\varepsilon^5) & \text{if } \mu(\eta(0)) = 1/2, \end{cases} \quad (13)$$

with

$$\kappa := -8\mu \left((2\eta(0) - 1) \left\{ \int_0^\infty \mathbb{E}_\eta^{(0)}[\eta_s(1) - \eta_s(-1)] ds \right\}^2 \right), \quad (14)$$

where the expectation $\mathbb{E}_\eta^{(0)}$ refers to the environment viewed from simple random walk, when starting at η . Moreover, the even terms in the expansion (11) equal zero and the antisymmetry relation $v(\varepsilon) = -v(-\varepsilon)$ holds.

When $\mu(\eta(0)) \neq 1/2$, for ε small enough, the sign of the velocity can be read from (13). When $\mu(\eta(0)) = 1/2$, the scenario is more subtle. Since $\nu_{1/2}$ is left invariant by particle-hole exchange and due to the form of transitions of the ε -RW, one may naively guess that the velocity is zero. Despite this guess, the answer seems strongly dependent on the specific dynamics of the underlying environment. In Section 4 we investigate more precisely the case of the East model and we give arguments supporting the negativity of $v(\varepsilon)$ for $\varepsilon > 0$ at density 1/2. Let us conclude this

section by observing that in simple settings it is easy to deduce that $v(\varepsilon) \equiv 0$ at density $1/2$.

Remark 3.3 (Zero velocity for independent spin-flip dynamics). *Suppose that the dynamics of the environment is given by independent spin-flips with generator $\mathcal{L}_{\text{env}}f(\eta) = \gamma \sum_{x \in \mathbb{Z}} [f(\eta^x) - f(\eta)]$, where η^x is the configuration obtained from η by a spin flip at x . Then the product Bernoulli measure with density $1/2$ is ergodic and reversible for this dynamics. Since for this initial distribution the process is invariant by inversion of particles and holes, it is easy to see that $v(\varepsilon) = v(-\varepsilon)$. Consequently, by the antisymmetry relation (12), we conclude that $v(\varepsilon) \equiv 0$.*

Our next result provides some description of the density profile of the environment observed by the ε -RW.

Proposition 3.4. *Under Assumptions 1 and for $2|\varepsilon| < \gamma$, the stationary distribution μ_ε of the environment seen from the ε -RW (i.e. of the process $(\eta_t^{(\varepsilon)})_{t \geq 0}$) admits the following representation for any function $f \in L^2(\mu)$:*

$$\mu_\varepsilon(f) = \mu(f) + 2\varepsilon \int_0^\infty ds \sum_{y \in \mathbb{Z}} p_t(y) \mu(\xi(0) \mathbb{E}_\xi^{\text{env}} [f(\tau_{y+1}\xi_s) - f(\tau_{y-1}\xi_s)]) + O(\varepsilon^2), \quad (15)$$

where $\mathbb{E}_\xi^{\text{env}}$ denotes the expectation w.r.t. the environment with generator \mathcal{L}_{env} starting from $\xi \in \Omega$, and $p_t(y-x)$ is the probability that a simple symmetric random walk jumping at rate 1, started at x , is in position y at time t .

4. THE ε -RW ON THE EAST MODEL

Definition 4.1. (East dynamics) *For $x \in \mathbb{Z}$ and $\xi \in \Omega$, set $c_x^{\text{east}}(\xi) := 1 - \xi(x+1)$. The East model is the Markov process on $\{0, 1\}^{\mathbb{Z}}$ with infinitesimal generator*

$$\mathcal{L}_{\text{east}}f(\xi) := \sum_{x \in \mathbb{Z}} c_x^{\text{east}}(\xi) [\rho(1 - \xi(x)) + (1 - \rho)\xi(x)] [f(\xi^x) - f(\xi)], \quad (16)$$

with $\rho \in (0, 1)$ being a fixed parameter, and ξ^x the configuration obtained by flipping the coordinate of ξ at site x .

The East model can be described as follows: at each site x , after an exponential time of parameter 1 and provided that the kinetic constraint $c_x^{\text{east}} = 1$ is satisfied³, the particle configuration $\xi(x)$ is refreshed and set equal to 1 with probability ρ and equal to 0 with probability $1 - \rho$. It is simple to check that the Bernoulli product measure with density ρ , denoted by ν_ρ , is a reversible probability measure.

Remark 4.2 (West and FA-1f models). *We notice that by definition of c_x^{east} , in order to change the state at site x , the site to its “East”, i.e. at position $x+1$, has to be vacant. This justifies the name of this model. The West model is the process with generator as in (16) when we replace $c_x^{\text{east}}(\xi)$ by $c_x^{\text{west}}(\xi) := 1 - \xi(x-1)$, which means that the constraint has to be satisfied in the other direction. A symmetrized version of the East and West models is the so-called FA1f (Fredrickson-Andersen one spin facilitated) model, that is the process with generator as in (16) with $c_x^{\text{fa}}(\xi) := 1 - \xi(x-1)\xi(x+1)$. In particular, for any $\rho \in (0, 1)$, the Bernoulli product measure ν_ρ is again reversible for the West and the FA1f models. Figure 2 shows space-time realizations of the East and the FA1f particle systems. The West looks like the East reflected w.r.t. the time axis.*

³The East model belongs to the class of kinetically constrained spin models [14].

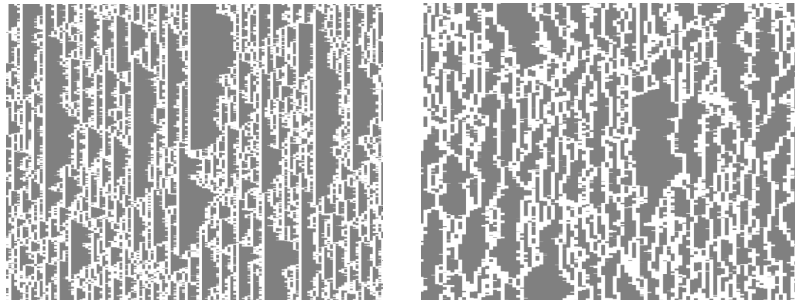


FIGURE 2. Two space-time (horizontal-vertical axis) realizations of the East and the FA1f models, left and right pictures, respectively. Particles (1 's) in gray.

All the models introduced above satisfy Assumption 1 for $\rho \in (0, 1)$ (see [1, 14] for the positivity of the spectral gap), and we can therefore apply the results of [2] to them. However, stronger inequalities of Sobolev type do not hold for the East model [17] and we can hope for no uniform mixing property due to the hardness of the constraint in Definition 4.1. It is also non-attractive, as one can check easily by noticing that more empty sites allow to create more holes, but also to add more particles.

In Figure 2 we present a simulation of the East/FA1f dynamics. One can observe bubbles of occupied sites forming. These are space-time regions with zero activity and a fundamental feature of kinetically constrained dynamics. Rigorous attempts towards an understanding of their structure can be found in [9, 10, 12, 18].

4.1. Asymptotic velocity. Simulations suggest that the ε -RW is ballistic for $\rho = 1/2$, drifting to the left when $\varepsilon > 0$ (see Figure 3). This motivates the following conjecture:

Conjecture 4.3. *When the environment is the East model at density $1/2$, for $\varepsilon > 0$ (resp. $\varepsilon < 0$) we have $v(\varepsilon) < 0$ (resp. $v(\varepsilon) > 0$).*

Below we will give two theoretical arguments supporting the above conjecture, based on Propositions 4.5 and 4.6.

It is tempting to interpret the sign of $v(\varepsilon)$ as a signature of the orientation of the East model and the asymmetry of its dynamics. However, as pointed out in Proposition 3.2, the antisymmetry relation (12) holds for the ε -RW on the East model. Recall the West model mentioned in Remark 4.2, and denote by $P_{\eta,0}^{\text{east}}$ and $P_{\eta,0}^{\text{west}}$, the laws of the ε -RW in the environments East and West, respectively, starting at the origin with environment η . Then, by considering a space reflection at the origin, it is easy to see that $\forall A \subset \mathbb{R}$ and any $\eta \in \Omega$, $P_{\eta,0}^{\text{east}}(X^{(\varepsilon)} \in A) = P_{\eta,0}^{\text{west}}(-X^{(-\varepsilon)} \in A)$. Consequently, at any density, $v_{\text{east}}(\varepsilon) = -v_{\text{west}}(-\varepsilon)$, where $v_{\text{east}}(\varepsilon)$ denotes the asymptotic velocity in (10) in the East environment, and similarly for West. In view of this observation, the following statement is a straightforward consequence of (12) and shows that the orientation of the environment does not determine the sign of the velocity.

Corollary 4.4. *At any density $\rho \in (0, 1)$*

$$v_{\text{east}}(\varepsilon) = v_{\text{west}}(\varepsilon).$$

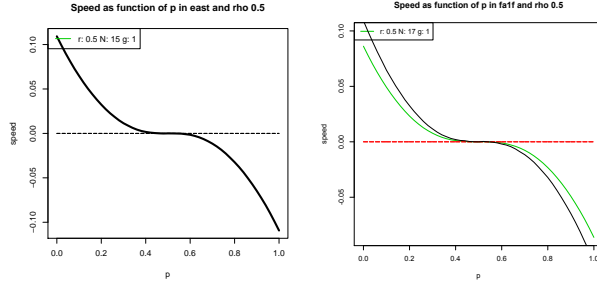


FIGURE 3. Velocity as a function of $p := 1/2 + \varepsilon \in [0, 1]$ of the ε -RW in East (left picture, black curve) and FA1f (right picture, green curve) environments at density $\rho = 1/2$. Note the antisymmetry $v(\varepsilon) = -v(-\varepsilon)$ as in Theorem 3.1 and the non-zero velocity. The curves have been obtained by interpolation over points at distance 0.02 in $[0, 1]$, for each point, the corresponding value of the velocity is given by a sample-mean on ≈ 5000 experiments in which the RW performed 2^{17} jumps.

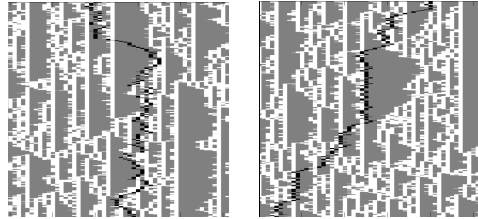


FIGURE 4. Two simulations of ε -RW (in black) on the East model with $\rho = 0.6$. Time goes down, space is horizontal; particles are in gray, holes in white. On the left, $\varepsilon = +0.3$. On the right, $\varepsilon = -0.3$. The two walkers have different-looking trajectories, even up to mirror reflection, due to the fact that one of them sticks to the fluctuating border of the bubbles and the other to the straight border.

Above we use the notation $v_{\text{east}}(\varepsilon)$ and $v_{\text{west}}(\varepsilon)$ to distinguish the velocities in two environments. From now on we consider only the East model and go back to the lighter notation $v(\varepsilon)$.

Besides the numerics in Figure 3, we now show two different results supporting Conjecture 4.3. The following proposition provides a criterion in terms of space-time correlations of the environment implying the negativity of κ defined in Proposition 3.2 (recall that $v(\varepsilon) = \kappa\varepsilon^3 + O(\varepsilon^5)$).

Proposition 4.5. *Set $\rho = 1/2$. If for all $s, t > 0$ and for all $y \geq 1$ it holds*

$$\mathbb{E}_{\nu_{1/2}}^{\text{east}} [\xi_0(0) (2\xi_t(y) - 1) \xi_{t+s}(0)] > 0, \tag{17}$$

then κ in (14) is negative.

Let us explain why we expect (17) to be true. It is clearly equivalent to the following inequality:

$$\mathbb{P}_{\nu_{1/2}}^{\text{east}} (\xi_0(0) = 1, \xi_t(y) = 1, \xi_{t+s}(0) = 1) > \mathbb{P}_{\nu_{1/2}}^{\text{east}} (\xi_0(0) = 1, \xi_t(y) = 0, \xi_{t+s}(0) = 1)$$

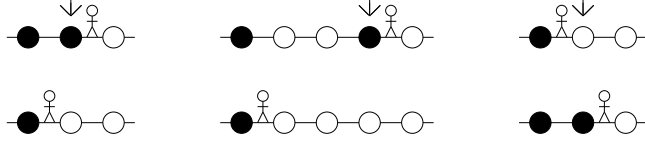


FIGURE 5. Illustration of the different updates in the environment inducing a jump of the degenerate random walker. Particles and holes of the environment are represented by black and white disks respectively. On top, three pre-update situations are depicted, with an arrow pointing the position of next flip occurring in the dynamics of the environment. On the bottom, one can see the result of the spin flip, in particular the induced new position of the walker.

Since we are at density $1/2$ and due to the orientation of the East model (see Appendix A), we have $\mathbb{P}_{\nu_{1/2}}^{\text{east}}(\xi_0(0) = 1, \xi_t(y) = 1) = \mathbb{P}_{\nu_{1/2}}^{\text{east}}(\xi_0(0) = 1, \xi_t(y) = 0) = 1/4$. Therefore the question is whether it is more likely to keep a particle initially present at time zero when $\xi_t(y) = 0$ or $\xi_t(y) = 1$. Intuitively, since zeros can send excitations that allow updates of particles to their left, having a particle at y should work towards conserving a particle on its left, e.g. at the origin, which explains why (17) should hold.

The second argument supporting Conjecture 4.3 is given by Proposition 4.6. More precisely, we introduce below a random walk which is a degenerate version of the ε -RW with $\varepsilon > 0$ on the East model and show that it has negative velocity.

4.1.1. *A degenerate drifting RW.* Let us introduce a degenerate version of the ε -RW. Informally, we introduce a new random walk $(Y_t)_{t \in \mathbb{R}_+}$ living on the edges of \mathbb{Z} with a hole to the right and a particle to the left. Initially, the walker stands on the first edge on the right of the origin satisfying this condition. If the particle to its left flips into a hole, it jumps instantly to the next edge of this type to its left. If instead the hole to its right flips into a particle, the random walk jumps instantly to the next edge of this type to its right. See Figure 5 for an illustration of the different possible jumps for the walker. The latter is a degenerate version of the ε -RW. Indeed it can be thought of as jumping at infinite rate to the left (resp. right) when it is sitting on top of a hole (resp. particle).

Let us give the precise definition of the joint process $(\xi_t, Y_t)_{t \geq 0}$. To this aim, we parametrize the edges of \mathbb{Z} by $1/2 + \mathbb{Z}$ assigning to each $y \in 1/2 + \mathbb{Z}$ the edge $\{y^-, y^+\}$, where $y^+ := y + 1/2$ and $y^- := y - 1/2$. The initial configuration (ξ_0, Y_0) is determined as follows: $\xi_0 \in \Omega$ is sampled from ν_ρ while $Y_0 := k + 1/2$, where k is the lowest non-negative integer such that $\xi_0(k) = 1$, $\xi_0(k+1) = 0$. Then the Markov generator of the joint process (ξ_t, Y_t) is given by

$$L_{deg}f(\xi, y) = \sum_{x \notin \{y^-, y^+\}} c_x^{\text{east}}(\xi) [\rho(1 - \eta(x)) + (1 - \rho)\eta(x)] [f(\eta^x, y) - f(\eta, y)] \\ + (1 - \rho) \left[f(\xi^{y^-}, y - k(\xi, y_-)) - f(\xi, y) \right] \\ + \rho c_{y^+}^{\text{east}}(\xi) \left[f(\xi^{y^+}, y + 1) - f(\xi, y) \right], \quad (18)$$

where $k(\xi, y_-)$ is the first positive integer k such that $\xi(y^- - k) = 1$.

The series in the r.h.s. corresponds to updates of the environment occurring on sites not belonging to the edge where the walker sits. The second line describes what happens when the particle on the left of the walker disappears, and the third when a particle appears on the right of the walker.

Notice that the above joint process is well defined, since at any time there are infinitely many sequences of particle-hole when the initial configuration is sampled from ν_ρ with density $\rho \in (0, 1)$. Moreover, the evolution of ξ_t is the standard East model at equilibrium.

At any density ρ , the random walk $(Y_t)_{t \geq 0}$ has a negative velocity:

Proposition 4.6. *For any $\rho \in (0, 1)$ it holds $\limsup_{t \rightarrow \infty} \frac{Y_t}{t} < 0$ a.s.*

The proof of the above proposition is given in Section 5.4.

4.2. Density profile. By applying Proposition 3.4 to specific environments a more detailed description of the observed density profile can be derived. In the case of the East model, we get the following.

Corollary 4.7. *Suppose that $2|\varepsilon|$ is smaller than the spectral gap of $\mathcal{L}_{\text{east}}$. Then for any $x \in \mathbb{Z}$, it holds*

$$\mu_\varepsilon(\eta(x)) = \rho + 2\varepsilon \int_0^\infty u(s) [p_s(x-1) - p_s(x+1)] ds + O(\varepsilon^2), \quad (19)$$

where

$$u(s) := \rho^2 - \nu_\rho(\xi(0) \mathbb{E}_\xi^{\text{east}}[\xi_s(0)])$$

is a negative increasing function on $[0, +\infty)$, with $u(0) = -\rho(1 - \rho)$ and $|u(s)| \leq \rho(1 - \rho)^{1/2} e^{-\lambda s}$.

In particular, if $x < 0$ (resp. $x > 0$), for $\varepsilon > 0$ small enough,

$$\mu_\varepsilon(\eta(x)) > \rho \quad (\text{resp. } < \rho). \quad (20)$$

Moreover, for any $x, y \in \mathbb{Z}$, we have that

$$\begin{aligned} \mu_\varepsilon(\eta(x)) - \mu_\varepsilon(\eta(x+y)) &= \varepsilon \int_0^\infty u(s) [p_s(x-1) - p_s(x+1) \\ &\quad - p_s(x+y-1) + p_s(x+y+1)] ds + O(\varepsilon^2). \end{aligned} \quad (21)$$

Remark 4.8 (Equilibrium affected around the origin). *From Theorem 3 in [2], we know that, as $|x| \rightarrow \infty$, the density profile of the environment seen by the walker, $(\mu_\varepsilon(\eta(x)))_{x \in \mathbb{Z}}$, approaches the constant profile ρ (corresponding to the equilibrium of the environment process). This reflects the fact that the ε -RW is “sitting at the origin” of the process $(\eta_t^{(\varepsilon)})_{t \geq 0}$. Moreover, recall that for $\varepsilon > 0$ the random walk has a tendency to jump to the right when sitting on top of particles, and vice versa on top of holes. Therefore, heuristically, it should spend more time with a particle to its left and a hole to its right. (20) and (21) confirm this description for ε small enough. Indeed, from (21), it is simple to see that $\mu_\varepsilon(\eta(-1)) > \mu_\varepsilon(\eta(0)) > \mu_\varepsilon(\eta(1))$ for small $\varepsilon > 0$. These observations are summarized in the numerics in Figure 6.*

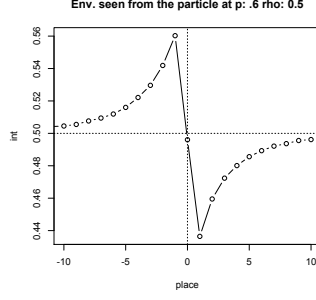


FIGURE 6. Simulation of the density profile $\mu_\varepsilon(\eta(x))$, $x = -10, \dots, 10$ at $\rho = 1/2$ and $\varepsilon = 0.1$.

5. PROOFS

5.1. Proof of Theorem 3.1. For the sake of clarity, we give the proof for the ε -RW. The proof can be easily generalized to the transition rates satisfying Assumption 3.

We first recall a lemma from [2], Lemma 10.1 therein. The notation is adapted to this specific context.

Lemma 5.1. *Under Assumptions 1, $t \geq 0$, $x \in \mathbb{Z}$, then for all $n \geq 0$, $f : \Omega \rightarrow \mathbb{R}$ local function and $\eta \in \Omega$, it holds*

$$\begin{aligned} \hat{L}_\varepsilon S_\varepsilon^{(n)}(s)f(\eta) &= \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{n-1}} dt_n \sum_{\mathbf{z} \in \{\pm 1\}^{n+1}} \sum_{\delta \in \{0,1\}^{n+1}} (-1)^{|\delta|} \\ &\times \mathbb{E}_\eta^{(0)} \left[\left(\prod_{i=1}^{n+1} \hat{r}_\varepsilon \left(z_i, \tau_{(\delta \cdot \mathbf{z})_{[i-1]}} \eta_{t-t_{i-1}} \right) \right) f \left(\tau_{(\delta \cdot \mathbf{z})_{[n+1]}} \eta_t \right) \right], \end{aligned} \quad (22)$$

where $\hat{r}_\varepsilon(\pm 1, \eta) := \pm \varepsilon(2\eta(0) - 1)$, $|\delta| := \sum_{i=1}^{n+1} (1 - \delta_i)$, $(\delta \cdot \mathbf{z})_{[i]} := \delta_1 z_1 + \dots + \delta_i z_i$, $(\delta \cdot \mathbf{z})_{[0]} := 0$ and $t_0 := t$.

Formula (22) has to be thought with no time integration in the degenerate case $n = 0$. The central ingredient in the proof of Theorem 3.1 is the following lemma.

Lemma 5.2. *Assume Assumptions 1. Then, for all $n \geq 0$, we have*

$$\mu \left(\hat{L}_\varepsilon S_\varepsilon^{(2n)}(t) j^{(\varepsilon)}(\eta) \right) = 0. \quad (23)$$

Proof. For all $n \geq 1$, for all $\eta \in \Omega$, by applying Lemma 5.1 to the function $j^{(\varepsilon)}(\eta) = 2\varepsilon(2\eta(0) - 1) = \pm 2\hat{r}_\varepsilon(\pm 1, \eta)$, we can write

$$\begin{aligned}
\hat{L}_\varepsilon S_\varepsilon^{(n)}(t)j^{(\varepsilon)}(\eta) &= \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{n-1}} dt_n \sum_{\delta \in \{0,1\}^{n+1}} (-1)^{|\delta|} \sum_{\mathbf{z} \in \{\pm 1\}^{n+1}} \\
&\quad \mathbb{E}_\eta^{(0)} \left[\left(\prod_{i=1}^{n+1} \varepsilon z_i f \left(\tau_{(\delta \cdot \mathbf{z})_{[i-1]}} \eta_{t-t_{i-1}} \right) \right) 2\varepsilon f \left(\tau_{(\delta \cdot \mathbf{z})_{[n+1]}} \eta_t \right) \right], \\
&= 2\varepsilon^{n+2} \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{n-1}} dt_n \sum_{\delta \in \{0,1\}^{n+1}} \sum_{\mathbf{z} \in \{\pm 1\}^{n+1}} \\
&\quad (-1)^{|\delta|} \left(\prod_{i=1}^{n+1} z_i \right) \mathbb{E}_\eta^{(0)} \left[\left(\prod_{i=1}^{n+2} f \left(\tau_{(\delta \cdot \mathbf{z})_{[i-1]}} \eta_{t-t_{i-1}} \right) \right) \right]. \quad (24)
\end{aligned}$$

with $f(\eta) := (2\eta(0) - 1)$ and $t_{n+1} := 0, t_0 := t$.

To shorten the notation, write $A_t := \{\mathbf{t} \in \mathbb{R}^n, 0 \leq t_n \leq \dots \leq t_1 \leq t\}$. For $\mathbf{t} \in A_t$, $\delta \in \{0, 1\}^{n+1}$, $\mathbf{z} \in \{\pm 1\}^{n+1}$, let

$$\phi_n(\mathbf{t}, \delta, \mathbf{z}) := \mathbb{E}_\mu^{(0)} \left[\left(\prod_{i=1}^{n+2} f \left(\tau_{(\delta \cdot \mathbf{z})_{[i-1]}} \eta_{t-t_{i-1}} \right) \right) \right]. \quad (25)$$

Then, by (24) and (25), for any $n \geq 1$, we can write that

$$\mu \left(\hat{L}_\varepsilon S_\varepsilon^{(n)}(t)j^{(\varepsilon)}(\eta) \right) = 2\varepsilon^{n+2} \sum_{\delta \in \{0,1\}^{n+1}} (-1)^{|\delta|} \sum_{\mathbf{z} \in \{-1,+1\}^{n+1}} \left(\prod_{i=1}^{n+1} z_i \right) \int_{A_t} d\mathbf{t} \phi_n(\mathbf{t}, \delta, \mathbf{z}), \quad (26)$$

We are now going to show that, for any $n \geq 1$, (26) reduces to

$$\mu \left(\hat{L}_\varepsilon S_\varepsilon^{(n)}(t)j^{(\varepsilon)}(\eta) \right) = 2\varepsilon^{n+2} (-1)^{n+1} \sum_{\mathbf{z} \in \{-1,+1\}^{n+1}} \left(\prod_{i=1}^{n+1} z_i \right) \int_{A_t} d\mathbf{t} \phi_n(\mathbf{t}, \mathbf{1}, \mathbf{z}), \quad (27)$$

Take $\delta \neq \mathbf{1}$ and let $j \in \{1, \dots, n+1\}$ such that $\delta_j = 0$. Then, for every such $\delta \in \{0, 1\}^{n+1}$, we get

$$\begin{aligned}
\sum_{\mathbf{z} \in \{\pm 1\}^{n+1}} \left(\prod_{i=1}^{n+1} z_i \right) \phi_n(\mathbf{t}, \delta, \mathbf{z}) &= \sum_{\substack{\mathbf{z} \in \{\pm 1\}^{n+1} \\ z_j = +1}} \left(\prod_{\substack{i=1, \dots, n+1 \\ i \neq j}} z_i \right) \phi_n(\mathbf{t}, \delta, \mathbf{z}) \\
&\quad - \sum_{\substack{\mathbf{z} \in \{\pm 1\}^{n+1} \\ z_j = -1}} \left(\prod_{\substack{i=1, \dots, n+1 \\ i \neq j}} z_i \right) \phi_n(\mathbf{t}, \delta, \mathbf{z}) = 0, \quad (28)
\end{aligned}$$

where in the last equality we have used that $\phi_n(\mathbf{t}, \delta, \mathbf{z})$ does not depend on z_j since $\delta_j = 0$. Hence, (27) is proven.

In the next steps we will first use reversibility and then that n is even.

On the one hand, by the change of variable,

$$A_t \ni \mathbf{t} = (t_1, t_2, \dots, t_n) \mapsto \mathbf{t}^* := (t - t_n, t - t_{n-1}, \dots, t - t_1) \in A_t,$$

we have

$$\int_{A_t} d\mathbf{t} \phi_n(\mathbf{t}, \mathbf{1}, \mathbf{z}) = \int_{A_t} d\mathbf{t} \phi_n(\mathbf{t}^*, \mathbf{1}, \mathbf{z}). \quad (29)$$

On the other hand, by reversibility, we can show that

$$\phi_n(\mathbf{t}, \mathbf{1}, \mathbf{z}) = \phi_n(\mathbf{t}^*, \mathbf{1}, \bar{\mathbf{z}}), \quad (30)$$

where $\bar{\mathbf{z}} := -(z_{n+1}, \dots, z_1)$.

Let us check (30). Write first

$$\begin{aligned} \phi_n(\mathbf{t}, \mathbf{1}, \mathbf{z}) &= \mathbb{E}_\mu^{(0)} \left[\left(\prod_{i=1}^{n+2} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[i-1]}} \eta_{t-t_{i-1}} \right) \right) \right] = \mathbb{E}_\mu^{(0)} \left[\left(\prod_{i=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[i]}} \eta_{t-t_i} \right) \right) \right] \\ &= \mathbb{E}_\mu^{(0)} \left[\left(\prod_{i=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[i]}} \eta_{t_i} \right) \right) \right] = \mathbb{E}_\mu^{(0)} \left[\left(\prod_{l=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[n+1-l]}} \eta_{t_{n+1-l}} \right) \right) \right] \\ &= \mathbb{E}_\mu^{(0)} \left[\left(\prod_{l=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[n+1-l]}} \eta_{t-t_l^*} \right) \right) \right], \end{aligned} \quad (31)$$

where the third identity follows by reversibility, and the last one by the mapping $\mathbf{t} \mapsto \mathbf{t}^*$ since $t_l^* = t - t_{n+1-l}$.

Next, note that

$$(\mathbf{1} \cdot \mathbf{z})_{[n+1-l]} = z_1 + \dots + z_{n+1-l} = -(\bar{z}_{n+1} + \bar{z}_n + \dots + \bar{z}_{l+1}) = b + (\mathbf{1} \cdot \bar{\mathbf{z}})_{[l]},$$

where $b := -\sum_{j=1}^{n+1} \bar{z}_j$. Hence, from (31) and translation invariance, we get

$$\begin{aligned} \phi_n(\mathbf{t}, \mathbf{1}, \mathbf{z}) &= \mathbb{E}_\mu^{(0)} \left[\left(\prod_{l=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \mathbf{z})_{[n+1-l]}} \eta_{t-t_l^*} \right) \right) \right] = \mathbb{E}_\mu^{(0)} \left[\left(\prod_{l=0}^{n+1} f \left(\tau_b \tau_{(\mathbf{1} \cdot \bar{\mathbf{z}})_{[l]}} \eta_{t-t_l^*} \right) \right) \right] \\ &= \mathbb{E}_\mu^{(0)} \left[\left(\prod_{l=0}^{n+1} f \left(\tau_{(\mathbf{1} \cdot \bar{\mathbf{z}})_{[l]}} \eta_{t-t_l^*} \right) \right) \right] = \phi_n(\mathbf{t}^*, \mathbf{1}, \bar{\mathbf{z}}), \end{aligned}$$

as claimed in (30).

Finally, for any $n \geq 0$, due to (29) and (27), we can write

$$\begin{aligned} \mu \left(\hat{L}_\varepsilon S_\varepsilon^{(2n)}(t) j^{(\varepsilon)}(\eta) \right) &= 2\varepsilon^{2n+2} (-1)^{2n+1} \sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) \int_{A_t} d\mathbf{t} \phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) \\ &= -\varepsilon^{2n+2} \sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) \int_{A_t} d\mathbf{t} [\phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) + \phi_{2n}(\mathbf{t}^*, \mathbf{1}, \mathbf{z})] \end{aligned}$$

Therefore, to get the claim in (23), it suffices to show that, for any $\mathbf{t} \in A_t$,

$$\sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) [\phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) + \phi_{2n}(\mathbf{t}^*, \mathbf{1}, \mathbf{z})] = 0. \quad (32)$$

In fact, by using that for $2n$ even, $\prod_{i=1}^{2n+1} \bar{z}_i = -\prod_{i=1}^{2n+1} z_i$, we can write

$$\sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) [\phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) + \phi_{2n}(\mathbf{t}^*, \mathbf{1}, \mathbf{z})] = \sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) \phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) \quad (33)$$

$$+ \sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} \bar{z}_i \right) \phi_{2n}(\mathbf{t}^*, \mathbf{1}, \bar{\mathbf{z}}) = \sum_{\mathbf{z} \in \{\pm 1\}^{2n+1}} \left(\prod_{i=1}^{2n+1} z_i \right) [\phi_{2n}(\mathbf{t}, \mathbf{1}, \mathbf{z}) - \phi_{2n}(\mathbf{t}^*, \mathbf{1}, \bar{\mathbf{z}})], \quad (34)$$

and as claimed in (32), the latter equals to zero due to (30). \square

We are now in shape to conclude the proof of Theorem 3.1. Lemma 5.2 implies the cancellation of the even terms in the expansion of the velocity (11), that is,

$$v(\varepsilon) = 2\varepsilon(2\mu(\eta(0)) - 1) + \sum_{n=0}^{\infty} \int_0^{\infty} \mu(\hat{L}_\varepsilon S_\varepsilon^{(2n+1)}(s)j^{(\varepsilon)})ds, \quad (35)$$

and from equation (35) the claim readily follows. Indeed, by using (27) in Lemma 5.2, for any $n \geq 0$, we have that :

$$\begin{aligned} \int_0^{\infty} ds \mu(\hat{L}_\varepsilon S_\varepsilon^{(2n+1)}(s)j^{(\varepsilon)}) &= 2\varepsilon^{2n+1} \sum_{\mathbf{z} \in \{\pm 1\}^{2n}} \left(\prod_{i=1}^{2n} z_i \right) \times \\ &\times \int_0^{\infty} ds \int_{A_s} dt \mathbb{E}_\mu^{(0)} \left[\left(\prod_{i=1}^{2n+1} \left(2\eta_{t-t_{i-1}} \left(\sum_{l=1}^{i-1} z_l \right) - 1 \right) \right) \right] =: \varepsilon^{2n+1} c_{2n}. \end{aligned}$$

By plugging the above expression into (35), we get

$$v(\varepsilon) = 2\varepsilon(2\mu(\eta(0)) - 1) + \varepsilon \sum_{n=1}^{\infty} \varepsilon^{2n} c_{2n},$$

which, as claimed, is an antisymmetric function of ε .

5.2. Proof of Proposition 3.2. By Lemma 5.2 and (35), it suffices to show that

$$\int_0^{\infty} \mu(\hat{L}_\varepsilon S_\varepsilon^{(1)}(s)j^{(\varepsilon)})ds = \varepsilon^3 \kappa, \quad (36)$$

with κ as in (14). To shorten the computations, we abbreviate $f(\eta) := 2\eta(0) - 1$ and $h(\eta) := 2[\eta(1) - \eta(-1)]$. Compute first

$$\begin{aligned} \mu(\hat{L}_\varepsilon S_\varepsilon^{(1)}(s)f) &= \varepsilon \mu(f S_\varepsilon^{(1)}(s)h) = -\varepsilon \mu(h S_\varepsilon^{(1)}(s)f) = -\varepsilon \int_0^s du \mu(h \cdot [S(s-u)\hat{L}_\varepsilon S(u)f]) \\ &= -\varepsilon \int_0^s du \mu([S(s-u)h][\hat{L}_\varepsilon S(u)f]) = -\varepsilon^2 \int_0^s du \mu([S(s-u)h]f[S(u)h]), \end{aligned}$$

where we have used the definition of \hat{L}_ε in (6), translation invariance, the definition of $S_\varepsilon^{(1)}(s)$, reversibility and Assumption 1–(iii).

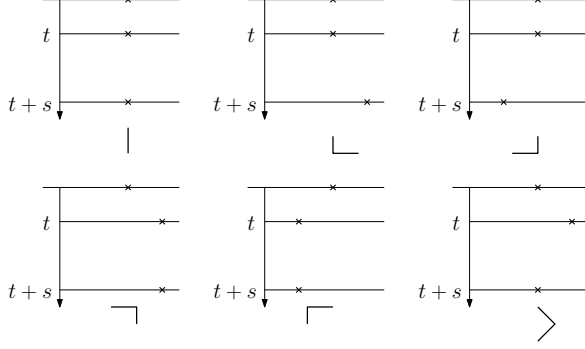


FIGURE 7. The only non-zero terms in A are those where $(1, 0)$, (y, t) and $(y + z + 1, t + s)$ are in one of the six above respective positions (marked by crosses). Similarly for the non-zero terms in B, C, D .

Hence

$$\begin{aligned}
\int_0^\infty \mu(\hat{L}_\varepsilon S_\varepsilon^{(1)}(s)f)ds &= -\varepsilon^2 \int_0^\infty ds \int_0^s du \mu(f[S(s-u)h][S(u)h]) \\
&= -\varepsilon^2 \int_0^\infty du \int_0^\infty dz \mu(fS(z)hS(u)h) = -\varepsilon^2 \mu\left(f\left[\int_0^\infty ds S(s)h\right]^2\right) \\
&= -4\varepsilon^2 \mu\left(f\left[\int_0^\infty ds \mathbb{E}_\eta^{(0)}[\eta_1(s) - \eta_{-1}(s)]\right]^2\right). \tag{37}
\end{aligned}$$

Finally, (36) follows by recalling that $j^{(\varepsilon)}(\eta) = 2\varepsilon f(\eta)$ and (37) above.

5.3. Proof of Proposition 4.5. From (37), we have

$$\kappa = -2 \int_0^\infty dt \int_0^\infty ds \nu_\rho(h \cdot S(t)(fS(s)h)). \tag{38}$$

Recall that $p_t(y)$ is the probability that a continuous time SRW started from 0 is at y at time t . Using that $S(t)g(\eta) = \sum_{y \in \mathbb{Z}} p_t(y) \mathbb{E}_\eta^{\text{east}}[g(\tau_y \eta_t)]$ and the Markov property, we can rewrite the term in the double integral in the above expression as $A - B - C + D$, where

$$A = \sum_{y, z \in \mathbb{Z}} p_t(y) p_s(z) \mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(y) - 1)\eta_{t+s}(z + y + 1)] \tag{39}$$

$$B = \sum_{y, z \in \mathbb{Z}} p_t(y) p_s(z) \mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(-1)(2\eta_t(y) - 1)\eta_{t+s}(z + y + 1)] \tag{40}$$

$$C = \sum_{y, z \in \mathbb{Z}} p_t(y) p_s(z) \mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(y) - 1)\eta_{t+s}(z + y - 1)] \tag{41}$$

$$D = \sum_{y, z \in \mathbb{Z}} p_t(y) p_s(z) \mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(-1)(2\eta_t(y) - 1)\eta_{t+s}(z + y - 1)]. \tag{42}$$

Lemma A.2 shows that the expectations appearing in A (resp. B , resp. C , resp. D) cancel as soon as $1, y, z + y + 1$ (resp. $-1, y, z + y + 1$, resp. $1, y, z + y - 1$, resp. $-1, y, z + y - 1$) are pairwise distinct. The same holds in the cases where $y < 1$ and $z + y + 1 = 1$ (resp. $y < -1$ and $z + y + 1 = -1$, etc). In fact we only have to deal with terms where the three space-time points involved in the expectation are in one of the six schematic configurations of Figure 7. We gather the terms corresponding

to the different positions in $A_{\downarrow}, A_{\lrcorner}, A_{\ulcorner}, A_{\llcorner}, A_{\lrcorner}, A_{\lrcorner}$ (and similarly for terms in B, C, D) so that

$$A_{\downarrow} = p_t(1)p_s(-1)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(1)-1)\eta_{t+s}(1)], \quad (43)$$

$$A_{\lrcorner} = \sum_{z \geq 0} p_t(1)p_s(z)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(1)-1)\eta_{t+s}(2+z)], \quad (44)$$

$$A_{\ulcorner} = \sum_{z \leq -2} p_t(1)p_s(z)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(1)-1)\eta_{t+s}(2+z)], \quad (45)$$

$$A_{\llcorner} = \sum_{y \geq 2} p_t(y)p_s(-1)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(y)-1)\eta_{t+s}(y)], \quad (46)$$

$$A_{\lrcorner} = \sum_{y \leq 0} p_t(y)p_s(-1)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(y)-1)\eta_{t+s}(y)], \quad (47)$$

$$A_{\lrcorner} = \sum_{y \geq 2} p_t(y)p_s(-y)\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(1)(2\eta_t(y)-1)\eta_{t+s}(1)]. \quad (48)$$

$$(49)$$

Translation invariance and the symmetry of the heat kernel imply that

$$A_{\downarrow} = B_{\downarrow} = C_{\downarrow} = D_{\downarrow}, \quad (50)$$

$$A_{\lrcorner} = B_{\lrcorner}, C_{\lrcorner} = D_{\lrcorner}, \quad (51)$$

$$A_{\ulcorner} = B_{\ulcorner}, C_{\ulcorner} = D_{\ulcorner}, \quad (52)$$

$$A_{\llcorner} + A_{\lrcorner} = B_{\llcorner} + B_{\lrcorner}, C_{\llcorner} + C_{\lrcorner} = D_{\llcorner} + D_{\lrcorner}, \quad (53)$$

so that many terms cancel and the sum reduces to

$$A - B - C + D = A_{\lrcorner} - B_{\lrcorner} - C_{\lrcorner} + D_{\lrcorner}. \quad (54)$$

Rearranging this expression by using translation invariance, we get that $A - B - C + D$ is nothing but

$$\sum_{y \geq 1} (p_t(y+1) - p_t(y-1))(p_s(y+1) - p_s(y-1))\mathbb{E}_{\nu_\rho}^{\text{east}}[\eta(0)(2\eta_t(y)-1)\eta_{t+s}(0)]. \quad (55)$$

The claim follows by noticing that for $s, t > 0, y \geq 1$,

$$(p_t(y+1) - p_t(y-1))(p_s(y+1) - p_s(y-1)) > 0. \quad (56)$$

5.4. Proof of Proposition 4.6. The result is a consequence of Lemma 3.2 in [12]. Let us just recall which process is considered in [12] and show how it can be coupled with our walker. We refer to Appendix A for some standard terminology concerning the East model.

In [12], the author studies the evolution of a single hole, which is called *the front*, through the East dynamics. The “front process” $(F_t)_{t \geq 0}$ is constructed as follows. Start with any configuration $\eta \in \Omega$ with a hole at site k , i.e. $\eta(k) = 0$. Let $F_0 := k$. As long as the Poisson clocks attached to the sites F_t and $F_t - 1$ do not ring, the front process does not jump. If there is a legal ring (see Definition A.1) at F_{t-} at time t and the associated Bernoulli variable is a 1 (i.e. F_{t-} is filled with a particle at time t), set $F_t := F_{t-} + 1$. Note that since the ring is assumed to be legal, the configuration at site F_t is still a hole. If there is a (necessarily legal) ring at $F_{t-} - 1$ at time t and the associated Bernoulli variable is a 0 (i.e. the site $F_{t-} - 1$ has a hole at time t), set $F_t := F_{t-} - 1$. Note that this front process is always on a zero

of the configuration ($\eta_t(F_t) = 0$ for all $t \geq 0$).⁴ Lemma 3.2 in [12], together with Borel-Cantelli lemma, says that the front moves at least with negative linear velocity asymptotically. More precisely, there exists a constant $\underline{v} > 0$ such that

$$\limsup_{t \rightarrow \infty} \frac{F_t}{t} \leq -\underline{v} \quad \text{a.s.} \quad (57)$$

It is not difficult to see, using the same graphical construction of the underlying East dynamics for the degenerate walker and the front process, that if we choose initially

$$F_0 = Y_0 + 1/2,$$

then, for all time $t \geq 0$,

$$F_t \geq Y_t + 1/2.$$

In view of this coupling and (57), the thesis is readily obtained.

5.5. Proof of Proposition 3.4. By Theorem 2.2, for any $f \in L^2(\mu)$, we have that

$$\mu_\varepsilon(f) = \mu(f) + \int_0^\infty \mu \left(\hat{L}_\varepsilon S(s) f \right) ds + O(\varepsilon^2). \quad (58)$$

Now note that by definition $S(t)f(\eta) = \mathbb{E}_\eta^{(0)} [f(\eta_t)] = \sum_{y \in \mathbb{Z}} p_t(y) \mathbb{E}_\eta^{\text{env}} [f(\tau_y \eta_t)]$. By means of this observation, together with the definition of \hat{L}_ε and translation invariance, we can write

$$\begin{aligned} \mu \left(\hat{L}_\varepsilon S(s) f \right) &= \varepsilon \mu \left((2\eta(0) - 1) [S(s)f(\tau_1 \eta) - S(s)f(\tau_{-1} \eta)] \right) \\ &= \varepsilon \mu \left((2\eta(0) - 1) \mathbb{E}_\eta^{(0)} [f(\tau_1 \eta_s) - f(\tau_{-1} \eta_s)] \right) \\ &= 2\varepsilon \sum_{y \in \mathbb{Z}} p_s(y) \mu \left(\eta(0) \mathbb{E}_\eta^{\text{env}} [f(\tau_{y+1} \eta_s) - f(\tau_{y-1} \eta_s)] \right). \end{aligned}$$

By combining the expression above with (58), we get (15).

5.6. Proof of Corollary 4.7. Equation (19) follows by plugging the function $f(\eta) := \eta(x)$ in (15) as we explain. The integrand in (15) then equals

$$\sum_{y \in \mathbb{Z}} p_s(y) \nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(x+y+1) - \xi_s(x+y-1)] \right). \quad (59)$$

To simplify the above expression we first observe that $\nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(y)] \right) = \rho^2$ for any $y \neq 0$ (cf. Lemma A.2 below). So that

$$\begin{aligned} (59) &= \rho^2 \sum_{y \neq -x-1} p_s(y) + p_s(-x-1) \nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(0)] \right) \\ &\quad - \rho^2 \sum_{y \neq -x+1} p_s(y) - p_s(-x+1) \nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(0)] \right) \\ &= [p_s(x-1) - p_s(x+1)] [\rho^2 - \nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(0)] \right)]. \end{aligned}$$

Recalling that $u(s) = \rho^2 - \nu_\rho \left(\xi(0) \mathbb{E}_\xi^{\text{east}} [\xi_s(0)] \right)$ and coming back to (15), one gets (19).

⁴In [12] we actually start with a configuration entirely filled to the left of the initial position of the front. The front is then at any time the left-most zero of the system. Due to the orientation of the East dynamics, however, the above definition gives a process with exactly the same properties.

It remains to analyze the function u . Clearly, $u(0) = -\rho(1-\rho)$. Moreover, setting $g(\eta) := \eta(0) - \rho$, by reversibility we get

$$|u(s)| = |\nu_\rho(g(\xi)\mathbb{E}_\xi^{\text{east}}[g(\xi_s)])| = \nu_\rho([S^{\text{east}}(s/2)g]^2) \leq e^{-s\lambda}\nu_\rho(g^2) = \rho(1-\rho)e^{-s\lambda}$$

where λ and $S^{\text{east}}(\cdot)$ denote respectively the spectral gap and the Markov semigroup of the East process.

Finally, let us focus on the sign and the growth of u . Call T_0 the time of the first *legal ring* at 0, see Definition A.1.

$$\begin{aligned} \nu_\rho(\xi(0)\mathbb{E}_\eta^{\text{east}}[\xi_s(0)]) &= \nu_\rho(\xi(0)\mathbb{E}_\xi^{\text{east}}[\xi_s(0)|T_0 \leq s]\mathbb{P}_\xi^{\text{east}}(T_0 \leq s)) \\ &\quad + \nu_\rho(\xi(0)\mathbb{E}_\xi^{\text{east}}[\xi_s(0)|T_0 > s]\mathbb{P}_\xi^{\text{east}}(T_0 > s)). \end{aligned} \quad (60)$$

On the one hand, if $T_0 > s$, $\xi_s(0) = \xi_0(0)$ (if there has been no legal ring at the origin, by definition the configuration at 0 has not been updated). So that the second term is $\nu_\rho(\xi(0)\mathbb{P}_\xi^{\text{east}}(T_0 > s))$, which in turn is just $\rho\mathbb{P}_{\nu_\rho}^{\text{east}}(T_0 > s)$, since the event $\{T_0 > s\}$ does not depend on $\xi(0)$. On the other hand, the first term can be rewritten as $\nu_\rho(\xi(0)\rho\mathbb{P}_\xi^{\text{east}}(T_0 \leq s))$ by distinguishing the zero that is on site 1 at time T_0 (the fact that T_0 is the time of a legal ring at 0 ensures its existence, cf. [1, 14] for the definition of the distinguished zero and for its properties). Again by orientation of the model, the last mean is equal to $\rho^2\mathbb{P}_{\nu_\rho}^{\text{east}}(T_0 \leq s)$. Due to the previous observations, (60) equals $\rho(1-\rho)\mathbb{P}_\xi^{\text{east}}(T_0 > s) + \rho^2$, hence

$$u(s) = -\rho(1-\rho)\mathbb{P}_{\nu_\rho}^{\text{east}}(T_0 > s) \quad (61)$$

It remains to notice that $\mathbb{P}_{\nu_\rho}^{\text{east}}(T_0 > s)$ is a quantity decreasing with s , so that $u(s)$ is indeed negative and increasing in s .

APPENDIX A. USEFUL FACTS ON THE EAST MODEL

Definition A.1. (Graphical representation of the East model) *Starting from a configuration $\eta \in \Omega$, the East dynamics $(\eta_t)_{t \geq 0}$ can be constructed as follows. With every $x \in \mathbb{Z}$ independently we associate a Poisson process with parameter 1 that will be called the (Poisson) clock at x . The process can then be constructed in the following way:*

- *Check the constraint: if the clock at site x rings at time t , look at the constraint at x in η_t , the configuration at time t .*
- *If $c_x^{\text{east}}(\eta_t) = 1$, the constraint is satisfied and the occupation variable at site x is replaced by a Bernoulli variable of parameter ρ independent of all the rest. The ring at time t is said to be a legal ring.*
- *If $c_x^{\text{east}}(\eta_t) = 0$, the constraint is not satisfied and the system is left unchanged.*

The following lemma is a consequence of reversibility and the orientation property of the East model which we use to prove Proposition 4.5. We write $\mathbb{E}_\eta^{\text{east}}$ for the expectation of the East dynamics starting at η , and we define $\mathbb{E}_{\nu_\rho}^{\text{east}}$ similarly.

Lemma A.2. *The following holds:*

- (1) *Let $0 \leq t_1 \leq t_2 \leq \dots \leq t_k$ and let f_1, \dots, f_k be functions on $\{0, 1\}^{\mathbb{Z}}$ such that the convex envelopes of their supports $\text{Conv}(\text{Supp}(f_1)), \dots, \text{Conv}(\text{Supp}(f_k))$*

are disjoint. Then

$$\mathbb{E}_{\nu_\rho}^{\text{east}} [f_1(\eta_{t_1}) \dots f_k(\eta_{t_k})] = \prod_{i=1}^k \nu_\rho(f_i).$$

(2) Let $0 \leq t_1 \leq t_2 \leq \dots \leq t_k$, f_1, \dots, f_k functions on $\{0, 1\}^{\mathbb{Z}}$ and $i_0 \in \{1, \dots, k\}$ such that $\nu_\rho(f_{i_0}) = 0$ and $x < y$ for all $x \in \text{Supp}(f_{i_0})$ and $y \in \cup_{i \neq i_0} \text{Supp}(f_i)$ (i.e. the support of f_{i_0} is to the left of all the other supports). Then

$$\mathbb{E}_{\nu_\rho}^{\text{east}} [f_1(\eta_{t_1}) \dots f_k(\eta_{t_k})] = 0. \quad (62)$$

Proof. The first statement is a consequence of the second one by iteration (let i_0 be the index of the function with left-most support and apply the second statement replacing f_{i_0} by $f_{i_0} - \nu_\rho(f_{i_0})$).

Notice that by reversibility we can construct the process at equilibrium also for negative times by mirroring the graphical construction. The process obtained is invariant by time translation. In particular, we have

$$\begin{aligned} \mathbb{E}_{\nu_\rho}^{\text{east}} [f_1(\eta_{t_1}) \dots f_k(\eta_{t_k})] &= \mathbb{E}_{\nu_\rho}^{\text{east}} \left[f_1(\eta_{t_1-t_{i_0}}) \dots f_k(\eta_{t_k-t_{i_0}}) \right] \\ &= \nu_\rho \left(f_{i_0}(\eta) \mathbb{E}_\eta^{\text{east}} \left[\prod_{i \neq i_0} f_i(\eta_{t_i-t_{i_0}}) \right] \right). \end{aligned}$$

Now notice that $\mathbb{E}_\eta^{\text{east}} \left[\prod_{i \neq i_0} f_i(\eta_{t_i-t_{i_0}}) \right]$ has disjoint support from f_{i_0} thanks to the orientation property of the East model. The two terms in the ν_ρ -mean are therefore decorrelated. Hence the result. \square

Acknowledgements. The authors thank P. Thomann for providing the numerical experiments in Figures 3 and 6. L. Avena has been supported by NWO Gravitation Grant 024.002.003-NETWORKS.

REFERENCES

- [1] D. Aldous and P. Diaconis, *The asymmetric one-dimensional constrained Ising model: rigorous results*. J. Stat. Phys. 107 (2002), 945–975.
- [2] L. Avena, O. Blondel, A. Faggionato, *L^2 -Perturbed Markov processes and applications to random walks in dynamic random environments*. Preprint (2016), arXiv:1602.06322.
- [3] L. Avena, T. Franco, M. Jara, F. Völlering, *Symmetric exclusion as a random environment: hydrodynamic limits*. Ann. Inst. Henri Poincaré 51, (2015) 901–916.
- [4] L. Avena, F. den Hollander, F. Redig, *Large deviation principle for one-dimensional random walk in dynamic random environment: attractive spin-flips and simple symmetric exclusion*. Markov Process. Related Fields 16 (2010), 139–168.
- [5] L. Avena, F. den Hollander, F. Redig, *Law of Large Numbers for a Class of Random Walks in Dynamic Random Environments*. Electron. J. Probab. 16 (2011), 587–617.
- [6] L. Avena, R. dos Santos, F. Völlering, *A transient random walk driven by an exclusion process: regenerations, limit theorems and an Einstein relation*. Latin American Jour. Prob. and Math. Stat. (ALEA) 10 (2), (2013) 693–709.
- [7] L. Avena, P. Thomann, *Continuity and Anomalous Fluctuations in Random Walks in Dynamic Random Environments: Numerics, Phase Diagrams and Conjectures*, J. Stat. Phys. 147 (2012), 1041–1067.
- [8] L. Berthier, G. Biroli, *Theoretical perspective on the glass transition and amorphous materials*, Rev. Mod. Phys. 83 (2011), 587–645.
- [9] T. Bodineau, C. Toninelli, *Activity phase transition for constrained dynamics*. Commun. Math. Phys. 311 (2012), 357–396.
- [10] T. Bodineau, V. Lecomte, C. Toninelli, *Finite Size Scaling of the dynamical free-energy in a kinetically constrained model*. J. Stat. Phys. 147 (2012), 1–17.

- [11] J.-P. Bouchaud, A. Georges, *Anomalous diffusion in disordered media*. Phys. Rep. 195 (1990), 127–293.
- [12] O. Blondel, *Front progression in the East model*. Stochastic Process. Appl. 123 (2013), no. 9, 3430–3465.
- [13] O. Blondel, *Is there a breakdown of the Stokes-Einstein relation in Kinetically Constrained Models at low temperature?* Europhysics Letters 107 (2014) 26005.
- [14] N. Cancrini, F. Martinelli, C. Roberto, C. Toninelli, *Kinetically constrained spin models*. Probab. Theory Related Fields 140 (2008), 459–504 .
- [15] P. Chleboun, A. Faggionato, F. Martinelli, *Time scale separation and dynamic heterogeneity in the low temperature East model*. Comm. Math. Phys. 328 (2014), 955–993.
- [16] A. Faggionato, F. Martinelli, C. Roberto, C. Toninelli; *Aging through hierarchical coalescence in the East model*. Commun. Math. Phys. 309 (2012), 459–495.
- [17] A. Faggionato, F. Martinelli, C. Roberto, C. Toninelli; *The East model: recent results and new progresses*. Markov Processes and Related Fields 19 (2013), 407–452.
- [18] S. Ganguly, E. Lubetzsky, F. Martinelli, *Cutoff for the East process*. Commun. Math. Phys. 335 (2015), 1287–1322.
- [19] M. Hilário, F. den Hollander, V. Sidoravicius, R. S. dos Santos, A. Teixeira, *Random Walk on Random Walks*. Electron. J. Probab. 20:95 (2015) 1-35.
- [20] F. den Hollander, H. Kesten, V. Sidoravicius, *Random walk in a high density dynamic random environment*. Indag. Math. (N.S.) 25 (2014), no. 4, 785–799.
- [21] F. den Hollander, R. dos Santos, *Scaling of a random walk on a supercritical contact process*. Ann. Inst. H. Poincaré Probab. Statist. 50 (2014) 1276–1300.
- [22] F. Huveneers, F. Simenhaus, *Random walk driven by simple exclusion process*. Electron. J. Probab. 20 (2015), no. 105, 42.
- [23] J. Jäckle, S. Eisinger, *A hierarchically constrained Ising model*. Z. Phys. B, 84 (1991), 115–129.
- [24] R. Jack, D. Kelsey, J. P. Garrahan, D. Chandler, *Negative differential mobility of weakly driven particles in models of glass formers*. Phys. Rev. E 78 (1), 011506 (2008).
- [25] T. Komorowski, S. Olla, *On Mobility and Einstein Relation for Tracers in Time-Mixing Random Environments*, J. Stat. Phys. 118 (2005), 407–435.
- [26] T. M. Liggett, *Interacting particle systems*. Grundlehren der Mathematischen Wissenschaften 276, Springer, New York (1985).
- [27] T. Mountford, M. E. Vares, *Random walks generated by equilibrium contact processes*. Electron. J. Probab. 20:3 (2015), 1-17.
- [28] F. Redig, F. Völlering, *Random walks in dynamic random environments: a transference principle*. Ann. Probab. 41 (2013), 3157–3180.
- [29] P. Sollich, M.R. Evans, *Glassy dynamics in the asymmetrically constrained kinetic Ising chain*. Phys. Rev. E 68 (2003), 031504.
- [30] F. Solomon, *Random walks in a random environment*, Ann. Probability 3 (1975), 1–31.
- [31] O. Zeitouni, *Random walks in random environments*, J. Phys. A: Math. Gen. 39, (2006) 433–464.

LUCA AVENA. MATHEMATISCH INSTITUUT UNIVERSITEIT LEIDEN. POSTBUS 9512 2300 RA LEIDEN, THE NETHERLANDS
E-mail address: l.avena@math.leidenuniv.nl

ORIANE BLONDEL. CNRS, UNIV. LYON, CNRS UMR 5208, INSTITUT CAMILLE JORDAN, 43 BLVD. DU 11 NOVEMBRE 1918, F-69622 VILLEURBANNE CEDEX, FRANCE
E-mail address: blondel@math.univ-lyon1.fr

ALESSANDRA FAGGIONATO. DIPARTIMENTO DI MATEMATICA, UNIVERSITÀ DI ROMA LA SAPIENZA. P.LE ALDO MORO 2, 00185 ROMA, ITALY
E-mail address: faggiona@mat.uniroma1.it