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Homogenization of Bingham Flow in thin porous media

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Abstract

By using dimension reduction and homogenization techniques, we study the steady flow of an incompressible viscoplastic Bingham fluid in a thin porous medium. A main feature of our study is the dependence of the yield stress of the Bingham fluid on the small parameters describing the geometry of the thin porous medium under consideration. Three different problems are obtained in the limit when the small parameter ε tends to zero, following the ratio between the height ε of the porous medium and the relative dimension a_ε of its periodically distributed pores. We conclude with the interpretation of these limit problems, which all preserve the nonlinear character of the flow.

1 Introduction

In this paper we study the asymptotic behavior of the flow of a viscoplastic Bingham fluid in a thin porous medium which contains an array of bodies modeled as vertical cylindrical obstacles (the pores). We refer the reader to the very recent paper [8] and the references therein for the application of our study to problems issued from the real life applications. As a first example one can mention the flow of the volcanic lava through dense forests (see [30]). Another important application is the flow of fresh concrete spreading through networks of steel bars (see [33]).

The model of thin porous medium of thickness much smaller than the distance between the pores was introduced in [34], where a stationary incompressible Navier-Stokes flow was studied. Recently, the model of thin porous medium under consideration in this paper was introduced in [20], where the flow of an incompressible viscous fluid described by the stationary Navier-Stokes equations was studied by the multiscale expansion method, which is a formal but powerful tool to analyse homogenization problems. These results were rigorously proved in [5] using an adaptation introduced in [4] of the periodic unfolding method from [16] and [17]. This adaptation consists of a combination of the unfolding method with a rescaling in the height variable, in order to work with a domain of fixed height, and to use monotonicity arguments to pass to the limit. In [4], in particular, the flow of an incompressible stationary Stokes system with a nonlinear viscosity, being a power law, was studied. For non-stationary incompressible viscous flow in a thin porous medium see [1], where a non-stationary Stokes system is considered, and [2], where a non-stationary non-newtonian Stokes system, where the viscosity obeyed the power law, is studied. For the periodic unfolding method applied to the study of problems stated in other type of thin periodic domains we refer for instance to [23] for crane type structures and to [24], [25] for thin layers with thin beams structures, where elasticity problems are considered. In [32], the homogenization of elasticity problems in thin periodic domains of planar grids type is studied. For problems involving arrays of bodies in high-contrast materials we refer the reader to [6], Chapter 2.

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Keywords: porous medium, thin domain, Bingham fluid

If Π is a three-dimensional domain with smooth boundary $\partial\Pi$ and $f = (f_1, f_2, f_3)$ are external given forces defined on Π , then the velocity $u = (u_1, u_2, u_3)$ of a fluid and its pressure p satisfy the equations of motion

$$-\sum_{j=1}^3 \partial_{x_i}(\sigma(p, u))_{ij} = f_i \text{ in } \Pi, \quad 1 \leq i \leq 3, \quad (1)$$

completed with the fluid's incompressibility condition $\operatorname{div} u = \sum_{i=1}^3 \partial_{x_i} u_i = 0$ in Π , and the no-slip boundary condition $u = 0$ on the boundary $\partial\Pi$. What distinguishes different fluids is the expression of the stress tensor σ . Newtonian fluids are the most encountered ones in real life and as typical examples one can mention the water and the air. For a newtonian fluid, the entries of the stress tensor $\sigma(p, u)$ are given by

$$(\sigma(p, u))_{ij} = -p\delta_{ij} + 2\mu(D(u))_{ij}, \quad 1 \leq i, j \leq 3 \quad (2)$$

where δ_{ij} is the Kronecker symbol, the real positive μ is the viscosity of the fluid and the entries of the strain tensor are $(D(u))_{ij} = (\partial_{x_j} u_i + \partial_{x_i} u_j)/2$. If f belongs to $(L^2(\Pi))^3$ and the space V is defined by $V = \{v \in (H_0^1(\Pi))^3 \mid \operatorname{div} v = 0\}$, then u and p satisfying (1) with (2) are such that (see for instance [22]):

(Stokes) There is a unique $u \in V$ and a unique (up to an additive real constant) $p \in L^2(\Pi)$ such that (if $\langle \cdot, \cdot \rangle$ is the dual pairing between $(H^{-1}(\Pi))^3$ and $(H_0^1(\Pi))^3$)

$$a(u, v) = l(v) - \langle \nabla p, v \rangle, \quad \forall v \in (H_0^1(\Pi))^3, \quad (3)$$

with $a(u, v) = 2\mu \int_{\Pi} D(u) : D(v) dx$ and $l(v) = \int_{\Pi} f \cdot v dx$.

A fluid whose stress is not defined by relation (2) is called a non-newtonian fluid. There are several classes of non-newtonian fluids, as the power law, Carreau, Cross, Bingham fluids. It is on the study of the last type of fluid that we are interested in this paper. We refer to [18] for a review on non-newtonian fluids. For a Bingham fluid, the nonlinear stress tensor is defined by (see [19])

$$(\sigma(p, u))_{ij} = -p\delta_{ij} + 2\mu(D(u))_{ij} + \sqrt{2g} \frac{(D(u))_{ij}}{|D(u)|}, \quad (4)$$

where $|D(u)|^2 = D(u) : D(u) \neq 0$ and the positive number g represents the yield stress of the fluid. If $g = 0$, then (4) becomes (2). Viscoplastic Bingham fluids are quite often encountered in real life. As examples one can mention volcanic lava, fresh concrete, the drilling mud, oils, clays and some paintings. For u_g and p_g satisfying (1) with (4), according to [19], one has the following result:

(Bingham) There is a unique $u_g \in V$ and a (non-unique) $p_g \in L^2(\Pi)/\mathbb{R}$ such that

$$a(u_g, v - u_g) + j(v) - j(u_g) \geq l(v - u_g) - \langle \nabla p_g, v - u_g \rangle, \quad \forall v \in (H_0^1(\Pi))^3. \quad (5)$$

Here $a, l, \langle \cdot, \cdot \rangle$ are as before and

$$j(v) = \sqrt{2g} \int_{\Pi} |D(v)| dx, \quad \forall v \in (H_0^1(\Pi))^3.$$

If the yield stress of the Bingham fluid is of the form $g(\varepsilon)$, with $\varepsilon \in]0, 1[$ and such that $g(\varepsilon)$ tends to zero when ε tends to zero, then, according to [[19], Chapter 6, Théorème 5.1.], the following result holds

When ε tends to zero, one has for the solution u_ε of problem (5) corresponding to $g(\varepsilon)$ the following convergence

$$u_\varepsilon \rightharpoonup u \quad \text{weakly in } V,$$

where u is the solution of problem (3).

This means that, in a fixed domain, the nonlinear character of the Bingham flow is lost in the limit (when the yield stress tends to zero), as it is expected. A natural question that arises is the following: If the yield stress $g(\varepsilon)$

is as before and, moreover, the domain Π itself depends on the small parameter ε , what happens when ε tends to zero? The answer is that, in the limit, the nonlinear character of the flow may be preserved. For instance, if Π_ε is a classical rigid porous medium, it was proven in [29] with the asymptotic expansion method that, in a range of parameters, the nonlinear character of the Bingham flow is preserved in the homogenized problem, which is a nonlinear Darcy equation. The convergence corresponding to the above mentioned result was proven in [10] with the two-scale convergence method and then recovered in [12] with the periodic unfolding method. The case of a doubly periodic rigid porous medium was studied in [11], where a more involved nonlinear Darcy equation is derived. Another class of domains for which the nonlinear character of the flow may be preserved in the limit is those of thin domains. The case of a domain Π_ε which is thin in one direction was addressed in [14] and [15]. We refer to [13] for the asymptotic analysis of a Bingham fluid in a thin T-like shaped domain. In all these cases, a lower-dimensional Bingham-like law was exhibited in the limit. This law was already encountered in engineering (see [31]), but no rigorous mathematical justification was previously known. A first mathematical result combining both periodic and thin domains for the Bingham flow was announced in [3]. For the shallow flow of a viscoplastic fluid we refer the reader to [21], [9], [26], [27] and [28]. For a homogenized non-newtonian viscoelastic model we refer to [7].

The paper is organized as follows. In Section 2. we state the problem: we define in (6) the thin porous medium Ω_ε (see also Figure 1), of height ε and relative dimension a_ε of its periodically distributed pores. In Ω_ε we consider the flow of a viscoplastic Bingham fluid with velocity u_ε and pressure p_ε verifying the nonlinear variational inequality (9). In Section 3. we give some *a priori* estimates for the velocity and for the pressure obtained after the change of variables (10) and verifying (12), and then for the velocity and for the pressure defined in (21). In Section 4. by passing to the limit $\varepsilon \rightarrow 0$, we prove the main convergence results of our paper, stated in Theorems 4.2, 4.4 and 4.6, respectively. Up to our knowledge, problems (36), (57) and (78) are new in the mathematical literature. We conclude in Section 5. with the interpretation of these limit problems, which all three preserve the nonlinear character of the flow; both effects of a nonlinear Darcy equation and a lower dimensional Bingham-like law appear. The paper ends with a list of References.

2 Statement of the problem

A periodic porous medium is defined by a domain ω and an associated microstructure, or periodic cell $Y' = [-1/2, 1/2]^2$, which is made of two complementary parts: the fluid part Y'_f , and the solid part Y'_s ($Y'_f \cup Y'_s = Y'$ and $Y'_f \cap Y'_s = \emptyset$). More precisely, we assume that ω is a smooth, bounded, connected set in \mathbb{R}^2 , and that Y'_s is an open connected subset of Y' with a smooth boundary $\partial Y'_s$, such that $\overline{Y'_s}$ is strictly included in Y' .

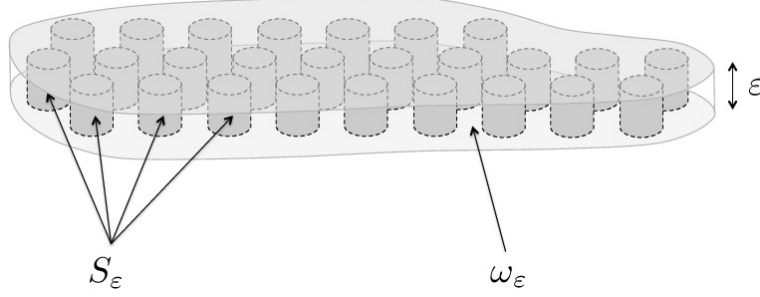
The microscale of a porous medium is a small positive number a_ε . The domain ω is covered by a regular mesh of size a_ε : for $k' \in \mathbb{Z}^2$, each cell $Y'_{k', a_\varepsilon} = a_\varepsilon k' + a_\varepsilon Y'$ is divided in a fluid part $Y'_{f, k', a_\varepsilon}$ and a solid part $Y'_{s, k', a_\varepsilon}$, i.e. is similar to the unit cell Y' rescaled to size a_ε . We define $Y = Y' \times (0, 1) \subset \mathbb{R}^3$, which is divided in a fluid part Y_f and a solid part Y_s , and consequently $Y_{k', a_\varepsilon} = Y'_{k', a_\varepsilon} \times (0, 1) \subset \mathbb{R}^3$, which is also divided in a fluid part Y_{f, k', a_ε} and a solid part Y_{s, k', a_ε} .

We denote by $\tau(\overline{Y'_{s, k', a_\varepsilon}})$ the set of all translated images of $\overline{Y'_{s, k', a_\varepsilon}}$. The set $\tau(\overline{Y'_{s, k', a_\varepsilon}})$ represents the solids in \mathbb{R}^2 . The fluid part of the bottom $\omega_\varepsilon \subset \mathbb{R}^2$ of the porous medium is defined by $\omega_\varepsilon = \omega \setminus \bigcup_{k' \in \mathcal{K}_\varepsilon} \overline{Y'_{s, k', a_\varepsilon}}$, where $\mathcal{K}_\varepsilon = \{k' \in \mathbb{Z}^2 : Y'_{k', a_\varepsilon} \cap \omega \neq \emptyset\}$. The whole fluid part $\Omega_\varepsilon \subset \mathbb{R}^3$ in the thin porous medium is defined by

$$\Omega_\varepsilon = \{(x_1, x_2, x_3) \in \omega_\varepsilon \times \mathbb{R} : 0 < x_3 < \varepsilon\}. \quad (6)$$

We make the assumption that the solids $\tau(\overline{Y'_{s, k', a_\varepsilon}})$ do not intersect the boundary $\partial\omega$. We define $Y_{s, k', a_\varepsilon}^\varepsilon = Y'_{s, k', a_\varepsilon} \times (0, \varepsilon)$. Denote by S_ε the set of the solids contained in Ω_ε . Then, S_ε is a finite union of solids, i.e. $S_\varepsilon = \bigcup_{k' \in \mathcal{K}_\varepsilon} \overline{Y_{s, k', a_\varepsilon}^\varepsilon}$.

We define $\tilde{\Omega}_\varepsilon = \omega_\varepsilon \times (0, 1)$, $\Omega = \omega \times (0, 1)$, and $Q_\varepsilon = \omega \times (0, \varepsilon)$. We observe that $\tilde{\Omega}_\varepsilon = \Omega \setminus \bigcup_{k' \in \mathcal{K}_\varepsilon} \overline{Y_{s, k', a_\varepsilon}}$, and we define $T_\varepsilon = \bigcup_{k' \in \mathcal{K}_\varepsilon} \overline{Y_{s, k', a_\varepsilon}}$ as the set of the solids contained in $\tilde{\Omega}_\varepsilon$.


 Figure 1: View of the domain Ω_ε

We denote by $:$ the full contraction of two matrices; for $A = (a_{i,j})_{1 \leq i,j \leq 3}$ and $B = (b_{i,j})_{1 \leq i,j \leq 3}$, we have $A : B = \sum_{i,j=1}^3 a_{ij}b_{ij}$.

In order to apply the unfolding method, we will need the following notation. For $k' \in \mathbb{Z}^2$, we define $\kappa : \mathbb{R}^2 \rightarrow \mathbb{Z}^2$ by

$$\kappa(x') = k' \iff x' \in Y'_{k',1}. \quad (7)$$

Remark that κ is well defined up to a set of zero measure in \mathbb{R}^2 (the set $\cup_{k' \in \mathbb{Z}^2} \partial Y'_{k',1}$). Moreover, for every $a_\varepsilon > 0$, we have

$$\kappa\left(\frac{x'}{a_\varepsilon}\right) = k' \iff x' \in Y'_{k',a_\varepsilon}.$$

We denote by C a generic positive constant which can change from line to line.

The points $x \in \mathbb{R}^3$ will be decomposed as $x = (x', x_3)$ with $x' = (x_1, x_2) \in \mathbb{R}^2$, $x_3 \in \mathbb{R}$. We also use the notation x' to denote a generic vector of \mathbb{R}^2 .

In Ω_ε we consider the stationary flow of an incompressible Bingham fluid. As already seen in the Introduction, following Duvaut and Lions [19], the problem is formulated in terms of a variational inequality.

For a vectorial function $v = (v', v_3)$, we define

$$(D(v))_{i,j} = \frac{1}{2} (\partial_{x_j} v_i + \partial_{x_i} v_j), \quad 1 \leq i, j \leq 3, \quad |D(v)|^2 = D(v) : D(v).$$

We introduce the following spaces

$$V(\Omega_\varepsilon) = \{v \in (H_0^1(\Omega_\varepsilon))^3 \mid \operatorname{div} v = 0 \text{ in } \Omega_\varepsilon\}, \quad H(\Omega_\varepsilon) = \{v \in (L^2(\Omega_\varepsilon))^3 \mid \operatorname{div} v = 0 \text{ in } \Omega_\varepsilon, v \cdot n = 0 \text{ on } \partial\Omega_\varepsilon\}.$$

For $u, v \in (H_0^1(\Omega_\varepsilon))^3$, we introduce

$$a(u, v) = 2\mu \int_{\Omega_\varepsilon} D(u) : D(v) dx, \quad j(v) = \sqrt{2}g(\varepsilon) \int_{\Omega_\varepsilon} |D(v)| dx, \quad (u, v)_{\Omega_\varepsilon} = \int_{\Omega_\varepsilon} u \cdot v dx,$$

where the yield stress $g(\varepsilon)$ will be made precise in Section 3.1. Let $f \in (L^2(\Omega))^3$ be given such that $f = (f', 0)$. Let $f_\varepsilon \in (L^2(\Omega_\varepsilon))^3$ be defined by

$$f_\varepsilon(x) = f(x', x_3/\varepsilon), \quad \text{a.e. } x \in \Omega_\varepsilon.$$

The model of the flow is described by the following variational inequality:

Find $u_\varepsilon \in V(\Omega_\varepsilon)$ such that

$$a(u_\varepsilon, v - u_\varepsilon) + j(v) - j(u_\varepsilon) \geq (f_\varepsilon, v - u_\varepsilon)_{\Omega_\varepsilon}, \quad \forall v \in V(\Omega_\varepsilon). \quad (8)$$

From Duvaut and Lions [19], we know that there exists a unique $u_\varepsilon \in V(\Omega_\varepsilon)$ solution of problem (8). Moreover, from Bourgeat and Mikelić [10], we know that if p_ε is the pressure of the fluid in Ω_ε , then problem (8) is equivalent to the following one: Find $u_\varepsilon \in V(\Omega_\varepsilon)$ and $p_\varepsilon \in L_0^2(\Omega_\varepsilon)$ such that

$$a(u_\varepsilon, v - u_\varepsilon) + j(v) - j(u_\varepsilon) \geq (f_\varepsilon, v - u_\varepsilon)_{\Omega_\varepsilon} + (p_\varepsilon, \operatorname{div}(v - u_\varepsilon))_{\Omega_\varepsilon}, \quad \forall v \in (H_0^1(\Omega_\varepsilon))^3. \quad (9)$$

Problem (9) admits a unique solution $u_\varepsilon \in V(\Omega_\varepsilon)$ and a (non) unique solution $p_\varepsilon \in L_0^2(\Omega_\varepsilon)$, where $L_0^2(\Omega_\varepsilon)$ denotes the space of functions belonging to $L^2(\Omega_\varepsilon)$ and of mean value zero.

Our aim is to study the asymptotic behavior of u_ε and p_ε when ε tends to zero. For this purpose, we first use the dilatation of the domain Ω_ε in the variable x_3 , namely

$$y_3 = \frac{x_3}{\varepsilon}, \quad (10)$$

in order to have the functions defined in an open set with fixed height, denoted $\tilde{\Omega}_\varepsilon$.

Namely, we define $\tilde{u}_\varepsilon \in (H_0^1(\tilde{\Omega}_\varepsilon))^3$, $\tilde{p}_\varepsilon \in L_0^2(\tilde{\Omega}_\varepsilon)$ by

$$\tilde{u}_\varepsilon(x', y_3) = u_\varepsilon(x', \varepsilon y_3), \quad \tilde{p}_\varepsilon(x', y_3) = p_\varepsilon(x', \varepsilon y_3) \quad a.e. (x', y_3) \in \tilde{\Omega}_\varepsilon.$$

Let us introduce some notation which will be useful in the following. For a vectorial function $v = (v', v_3)$ and a scalar function w , we will denote $\mathbb{D}_{x'}[v] = \frac{1}{2}(D_{x'}v + D_{x'}^t v)$ and $\partial_{y_3}[v] = \frac{1}{2}(\partial_{y_3}v + \partial_{y_3}^t v)$, where we denote $\partial_{y_3} = (0, 0, \frac{\partial}{\partial y_3})^t$. Moreover, associated to the change of variables (10), we introduce the operators: D_ε , \mathbb{D}_ε , $\operatorname{div}_\varepsilon$ and ∇_ε , defined by

$$(D_\varepsilon v)_{i,j} = \partial_{x_j} v_i \quad \text{for } i = 1, 2, 3, j = 1, 2, \quad (D_\varepsilon v)_{i,3} = \frac{1}{\varepsilon} \partial_{y_3} v_i \quad \text{for } i = 1, 2, 3,$$

$$\mathbb{D}_\varepsilon[v] = \frac{1}{2}(D_\varepsilon v + D_\varepsilon^t v), \quad |\mathbb{D}_\varepsilon[v]|^2 = \mathbb{D}_\varepsilon[v] : \mathbb{D}_\varepsilon[v], \quad \operatorname{div}_\varepsilon v = \operatorname{div}_{x'} v' + \frac{1}{\varepsilon} \partial_{y_3} v_3, \quad \nabla_\varepsilon w = (\nabla_{x'} w, \frac{1}{\varepsilon} \partial_{y_3} w)^t.$$

We introduce the following spaces

$$V(\tilde{\Omega}_\varepsilon) = \{\tilde{v} \in (H_0^1(\tilde{\Omega}_\varepsilon))^3 \mid \operatorname{div}_\varepsilon \tilde{v} = 0 \text{ in } \tilde{\Omega}_\varepsilon\}, \quad H(\tilde{\Omega}_\varepsilon) = \{\tilde{v} \in (L^2(\tilde{\Omega}_\varepsilon))^3 \mid \operatorname{div}_\varepsilon \tilde{v} = 0 \text{ in } \tilde{\Omega}_\varepsilon, \tilde{v} \cdot n = 0 \text{ on } \partial\tilde{\Omega}_\varepsilon\}.$$

For $\tilde{u}, \tilde{v} \in V(\tilde{\Omega}_\varepsilon)$, we introduce

$$a_\varepsilon(\tilde{u}, \tilde{v}) = 2\mu \int_{\tilde{\Omega}_\varepsilon} \mathbb{D}_\varepsilon[\tilde{u}] : \mathbb{D}_\varepsilon[\tilde{v}] dx' dy_3, \quad j_\varepsilon(\tilde{v}) = \sqrt{2}g(\varepsilon) \int_{\tilde{\Omega}_\varepsilon} |\mathbb{D}_\varepsilon[\tilde{v}]| dx' dy_3, \quad (\tilde{u}, \tilde{v})_{\tilde{\Omega}_\varepsilon} = \int_{\tilde{\Omega}_\varepsilon} \tilde{u} \cdot \tilde{v} dx' dy_3.$$

Using the transformation (10), the variational inequality (8) can be rewritten as:

Find $\tilde{u}_\varepsilon \in V(\tilde{\Omega}_\varepsilon)$ such that

$$a_\varepsilon(\tilde{u}_\varepsilon, \tilde{v} - \tilde{u}_\varepsilon) + j_\varepsilon(\tilde{v}) - j_\varepsilon(\tilde{u}_\varepsilon) \geq (f, \tilde{v} - \tilde{u}_\varepsilon)_{\tilde{\Omega}_\varepsilon}, \quad \forall \tilde{v} \in V(\tilde{\Omega}_\varepsilon), \quad (11)$$

and (9) can be rewritten as:

Find $\tilde{u}_\varepsilon \in V(\tilde{\Omega}_\varepsilon)$ and $\tilde{p}_\varepsilon \in L_0^2(\tilde{\Omega}_\varepsilon)$ such that

$$a_\varepsilon(\tilde{u}_\varepsilon, \tilde{v} - \tilde{u}_\varepsilon) + j_\varepsilon(\tilde{v}) - j_\varepsilon(\tilde{u}_\varepsilon) \geq (f, \tilde{v} - \tilde{u}_\varepsilon)_{\tilde{\Omega}_\varepsilon} + (\tilde{p}_\varepsilon, \operatorname{div}_\varepsilon(\tilde{v} - \tilde{u}_\varepsilon))_{\tilde{\Omega}_\varepsilon}, \quad \forall \tilde{v} \in (H_0^1(\tilde{\Omega}_\varepsilon))^3. \quad (12)$$

Our goal now is to describe the asymptotic behavior of this new sequence $(\tilde{u}_\varepsilon, \tilde{p}_\varepsilon)$.

3 A Priori Estimates

We start by obtaining some *a priori* estimates for \tilde{u}_ε .

Lemma 3.1. *There exists a constant C independent of ε , such that if $\tilde{u}_\varepsilon \in (H_0^1(\tilde{\Omega}_\varepsilon))^3$ is the solution of problem (11), one has*

i) if $a_\varepsilon \approx \varepsilon$, with $a_\varepsilon/\varepsilon \rightarrow \lambda$, $0 < \lambda < +\infty$, or $a_\varepsilon \ll \varepsilon$, then

$$\|\tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^3} \leq \frac{C}{\mu} a_\varepsilon^2, \quad \|\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \leq \frac{C}{\mu} a_\varepsilon, \quad \|D_\varepsilon \tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \leq \frac{C}{\mu} a_\varepsilon, \quad (13)$$

ii) if $a_\varepsilon \gg \varepsilon$, then

$$\|\tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^3} \leq \frac{C}{\mu} \varepsilon^2, \quad \|\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \leq \frac{C}{\mu} \varepsilon, \quad \|D_\varepsilon \tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \leq \frac{C}{\mu} \varepsilon. \quad (14)$$

Proof. Setting successively $\tilde{v} = 2\tilde{u}_\varepsilon$ and $\tilde{v} = 0$ in (11), we have

$$2\mu \int_{\tilde{\Omega}_\varepsilon} \mathbb{D}_\varepsilon[\tilde{u}_\varepsilon] : \mathbb{D}_\varepsilon[\tilde{u}_\varepsilon] dx' dy_3 + \sqrt{2}g(\varepsilon) \int_{\tilde{\Omega}_\varepsilon} |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]| dx' dy_3 = \int_{\tilde{\Omega}_\varepsilon} f \cdot \tilde{u}_\varepsilon dx' dy_3. \quad (15)$$

Using Cauchy-Schwarz's inequality and the assumption of f , we obtain that

$$\int_{\tilde{\Omega}_\varepsilon} f \cdot \tilde{u}_\varepsilon dx' dy_3 \leq C \|\tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^3},$$

and taking into account that $\int_{\tilde{\Omega}_\varepsilon} |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]| dx' dy_3 \geq 0$, by (15), we have

$$\|\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}}^2 \leq \frac{C}{\mu} \|\tilde{u}_\varepsilon\|_{(L^2(\tilde{\Omega}_\varepsilon))^3}.$$

For the cases $a_\varepsilon \approx \varepsilon$ or $a_\varepsilon \ll \varepsilon$, taking into account Remark 4.3(i) in [4], we obtain the second estimate in (13), and, consequently, from classical Korn's inequality we obtain the last estimate in (13). Now, from the second estimate in (13) and Remark 4.3(i) in [4], we deduce the first estimate in (13). For the case $a_\varepsilon \gg \varepsilon$, proceeding similarly with Remark 4.3(ii) in [4], we obtain the desired result. \square

3.1 The extension of $(\tilde{u}_\varepsilon, \tilde{p}_\varepsilon)$ to the whole domain Ω

We extend the velocity \tilde{u}_ε by zero to the $\Omega \setminus \tilde{\Omega}_\varepsilon$ and denote the extension by the same symbol. Obviously, estimates (13)-(14) remain valid and the extension is divergence free too.

We study in the sequel the following cases for the value of the yield stress $g(\varepsilon)$:

i) if $a_\varepsilon \approx \varepsilon$, with $a_\varepsilon/\varepsilon \rightarrow \lambda$, $0 < \lambda < +\infty$, or $a_\varepsilon \ll \varepsilon$, then $g(\varepsilon) = g a_\varepsilon$,

ii) if $a_\varepsilon \gg \varepsilon$, then $g(\varepsilon) = g \varepsilon$,

where g is a positive number. These choices are the most challenging ones and they answer to the question addressed in the paper, namely they all preserve in the limit the nonlinear character of the flow.

In order to extend the pressure to the whole domain Ω , the mapping R^ε (defined in Lemma 4.5 in [4] as R_2^ε) allows us to extend the pressure p_ε to Q_ε by introducing F_ε in $(H^{-1}(Q_\varepsilon))^3$:

$$\langle F_\varepsilon, w \rangle_{Q_\varepsilon} = \langle \nabla p_\varepsilon, R^\varepsilon w \rangle_{\Omega_\varepsilon}, \quad \text{for any } w \in (H_0^1(Q_\varepsilon))^3. \quad (16)$$

Setting successively $v = u_\varepsilon + R^\varepsilon w$ and $v = u_\varepsilon - R^\varepsilon w$ in (9) we get the inequality

$$|\langle F_\varepsilon, w \rangle_{Q_\varepsilon}| \leq |a(u_\varepsilon, R^\varepsilon w)| + |(f_\varepsilon, R^\varepsilon w)_{\Omega_\varepsilon}| + j(R^\varepsilon w). \quad (17)$$

Moreover, if $\operatorname{div} w = 0$ then $\langle F_\varepsilon, w \rangle_{Q_\varepsilon} = 0$, and the DeRham Theorem gives the existence of P_ε in $L_0^2(Q_\varepsilon)$ with $F_\varepsilon = \nabla P_\varepsilon$.

Using the change of variables (10), we get for any $\tilde{w} \in (H_0^1(\Omega))^3$ where $\tilde{w}(x', y_3) = w(x', \varepsilon y_3)$,

$$\left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega = - \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_\varepsilon \tilde{w} \, dx' \, dy_3 = -\varepsilon^{-1} \int_{Q_\varepsilon} P_\varepsilon \operatorname{div} w \, dx = \varepsilon^{-1} \langle \nabla P_\varepsilon, w \rangle_{Q_\varepsilon}.$$

Then, using the identification (16) of F_ε and the inequality (17),

$$|\left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega| \leq \varepsilon^{-1} (|a(u_\varepsilon, R^\varepsilon w)| + |(f_\varepsilon, R^\varepsilon w)_{\Omega_\varepsilon}| + j(R^\varepsilon w)).$$

and applying the change of variables (10),

$$|\left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega| \leq |a_\varepsilon(\tilde{u}_\varepsilon, \tilde{R}^\varepsilon \tilde{w})| + |(f, \tilde{R}^\varepsilon \tilde{w})_{\tilde{\Omega}_\varepsilon}| + j_\varepsilon(\tilde{R}^\varepsilon \tilde{w}), \quad (18)$$

where $\tilde{R}^\varepsilon \tilde{w} = R^\varepsilon w$ for any $\tilde{w} \in (H_0^1(\Omega))^3$.

Now, we estimate the right-hand side of (18) using the estimates given in Lemma 4.6 in [4].

Lemma 3.2. *There exists a constant C independent of ε , such that the extension $\tilde{P}_\varepsilon \in L_0^2(\Omega)$ of the pressure \tilde{p}_ε satisfies*

$$\left\| \tilde{P}_\varepsilon \right\|_{L_0^2(\Omega)} \leq C. \quad (19)$$

Proof. Let us estimate $\nabla_\varepsilon \tilde{P}_\varepsilon$ in the cases $a_\varepsilon \approx \varepsilon$ or $a_\varepsilon \ll \varepsilon$. We estimate the right-hand side of (18). Using Cauchy-Schwarz's inequality and from the second estimate in (13) we have

$$|a_\varepsilon(\tilde{u}_\varepsilon, \tilde{R}^\varepsilon \tilde{w})| \leq 2\mu \|\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \left\| D_\varepsilon \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} \leq C a_\varepsilon \left\| D_\varepsilon \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}}.$$

Using the assumption made on the function f , we obtain

$$|(f, \tilde{R}^\varepsilon \tilde{w})_{\tilde{\Omega}_\varepsilon}| \leq C \left\| \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^3},$$

and by Cauchy-Schwarz's inequality and taking into account that $|\tilde{\Omega}_\varepsilon| \leq |\Omega|$, we obtain

$$j_\varepsilon(\tilde{R}^\varepsilon \tilde{w}) \leq C a_\varepsilon \left\| D_\varepsilon \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}}.$$

Then, from (18), we deduce

$$\left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| \leq C a_\varepsilon \left\| D_\varepsilon \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} + C \left\| \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^3}.$$

Taking into account the third point in Lemma 4.6 in [4], we have

$$\left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| \leq C a_\varepsilon \left(\frac{1}{a_\varepsilon} \|\tilde{w}\|_{(L^2(\Omega))^3} + \|D_\varepsilon \tilde{w}\|_{(L^2(\Omega))^{3 \times 3}} \right) + C \left(\|\tilde{w}\|_{(L^2(\Omega))^3} + a_\varepsilon \|D_\varepsilon \tilde{w}\|_{(L^2(\Omega))^{3 \times 3}} \right).$$

If $a_\varepsilon \approx \varepsilon$ we take into account that $a_\varepsilon \ll 1$, and if $a_\varepsilon \ll \varepsilon$ we take into account that $a_\varepsilon/\varepsilon \ll 1$ and $a_\varepsilon \ll 1$, and we see that there exists a positive constant C such that

$$\left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| \leq C \|\tilde{w}\|_{(H_0^1(\Omega))^3}, \quad \forall \tilde{w} \in (H_0^1(\Omega))^3,$$

and consequently

$$\left\| \nabla_\varepsilon \tilde{P}_\varepsilon \right\|_{(H^{-1}(\Omega))^3} \leq C.$$

It follows that (see for instance Girault and Raviart [22], Chapter I, Corollary 2.1) there exists a representative of $\tilde{P}_\varepsilon \in L_0^2(\Omega)$ such that

$$\left\| \tilde{P}_\varepsilon \right\|_{L_0^2(\Omega)} \leq C \left\| \nabla \tilde{P}_\varepsilon \right\|_{(H^{-1}(\Omega))^3} \leq C \left\| \nabla_\varepsilon \tilde{P}_\varepsilon \right\|_{(H^{-1}(\Omega))^3} \leq C.$$

Finally, let us estimate $\nabla_\varepsilon \tilde{P}_\varepsilon$ in the case $a_\varepsilon \gg \varepsilon$. Similarly to the previous case, we estimate the right side of (18) by using Cauchy-Schwarz's inequality and from the second estimate in (14), and we have

$$\left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| \leq C\varepsilon \left\| D_\varepsilon \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^{3 \times 3}} + C \left\| \tilde{R}^\varepsilon \tilde{w} \right\|_{(L^2(\tilde{\Omega}_\varepsilon))^3}.$$

Taking into account the proof in Lemma 4.5 in [4], the change of variables (10) and that $a_\varepsilon \gg \varepsilon$, we can deduce

$$\begin{aligned} \left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| &\leq C\varepsilon \left(\frac{1}{\varepsilon} \|\tilde{w}\|_{(L^2(\Omega))^3} + \frac{1}{\varepsilon} \|D_{x'} \tilde{w}\|_{(L^2(\Omega))^{3 \times 2}} + \frac{1}{\varepsilon} \|\partial_{y_3} \tilde{w}\|_{(L^2(\Omega))^3} \right) \\ &+ C \left(\|\tilde{w}\|_{(L^2(\Omega))^3} + a_\varepsilon \|D_{x'} \tilde{w}\|_{(L^2(\Omega))^{3 \times 2}} + \|\partial_{y_3} \tilde{w}\|_{(L^2(\Omega))^3} \right), \end{aligned}$$

and using that $a_\varepsilon \ll 1$, we see that there exists a positive constant C such that

$$\left| \left\langle \nabla_\varepsilon \tilde{P}_\varepsilon, \tilde{w} \right\rangle_\Omega \right| \leq C \|\tilde{w}\|_{(H_0^1(\Omega))^3}, \quad \forall \tilde{w} \in (H_0^1(\Omega))^3,$$

and reasing as the previous case, we have the estimate (19). \square

According to these extensions, problem (12) can be written as:

$$\begin{aligned} &2\mu \int_\Omega \mathbb{D}_\varepsilon[\tilde{u}_\varepsilon] : \mathbb{D}_\varepsilon[\tilde{v} - \tilde{u}_\varepsilon] dx' dy_3 + \sqrt{2}g(\varepsilon) \int_\Omega |\mathbb{D}_\varepsilon[\tilde{v}]| dx' dy_3 - \sqrt{2}g(\varepsilon) \int_\Omega |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]| dx' dy_3 \\ &\geq \int_\Omega f \cdot (\tilde{v} - \tilde{u}_\varepsilon) dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_\varepsilon(\tilde{v} - \tilde{u}_\varepsilon) dx' dy_3, \end{aligned} \quad (20)$$

for every \tilde{v} that is the extension by zero to the whole Ω of a function in $(H_0^1(\tilde{\Omega}_\varepsilon))^3$.

3.2 Adaptation of the Unfolding Method

The change of variable (10) does not provide the information we need about the behavior of \tilde{u}_ε in the microstructure associated to $\tilde{\Omega}_\varepsilon$. To solve this difficulty, we use an adaptation introduced in [4] of the unfolding method from [16] and [17].

Let us recall this adaptation of the unfolding method in which we divide the domain Ω in cubes of lateral length a_ε and vertical length 1. For this purpose, given $(\tilde{u}_\varepsilon, \tilde{P}_\varepsilon) \in (H_0^1(\Omega))^3 \times L_0^2(\Omega)$, we define $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ by

$$\hat{u}_\varepsilon(x', y) = \tilde{u}_\varepsilon \left(a_\varepsilon \kappa \left(\frac{x'}{a_\varepsilon} \right) + a_\varepsilon y', y_3 \right), \quad \hat{P}_\varepsilon(x', y) = \tilde{P}_\varepsilon \left(a_\varepsilon \kappa \left(\frac{x'}{a_\varepsilon} \right) + a_\varepsilon y', y_3 \right), \quad \text{a.e. } (x', y) \in \omega \times Y, \quad (21)$$

where the function κ is defined in (7).

Remark 3.3. For $k' \in \mathcal{K}_\varepsilon$, the restriction of $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ to $Y'_{k', a_\varepsilon} \times Y$ does not depend on x' , whereas as a function of y it is obtained from $(\tilde{u}_\varepsilon, \tilde{P}_\varepsilon)$ by using the change of variables $y' = \frac{x' - a_\varepsilon k'}{a_\varepsilon}$, which transforms Y'_{k', a_ε} into Y .

We are now in position to obtain estimates for the sequences $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$, as in the proof of Lemma 4.9 in [4].

Lemma 3.4. *There exists a constant C independent of ε , such that the couple $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ defined by (21) satisfies*

i) if $a_\varepsilon \approx \varepsilon$, with $a_\varepsilon/\varepsilon \rightarrow \lambda$, $0 < \lambda < +\infty$, or $a_\varepsilon \ll \varepsilon$,

$$\|\hat{u}_\varepsilon\|_{(L^2(\omega \times Y))^3} \leq C a_\varepsilon^2, \quad \|\mathbb{D}_{y'}[\hat{u}_\varepsilon]\|_{(L^2(\omega \times Y))^{3 \times 2}} \leq C a_\varepsilon^2, \quad \|\partial_{y_3}[\hat{u}_\varepsilon]\|_{(L^2(\omega \times Y))^3} \leq C \varepsilon a_\varepsilon,$$

ii) if $a_\varepsilon \gg \varepsilon$,

$$\|\hat{u}_\varepsilon\|_{(L^2(\omega \times Y))^3} \leq C \varepsilon^2, \quad \|\mathbb{D}_{y'}[\hat{u}_\varepsilon]\|_{(L^2(\omega \times Y))^{3 \times 2}} \leq C a_\varepsilon \varepsilon, \quad \|\partial_{y_3}[\hat{u}_\varepsilon]\|_{(L^2(\omega \times Y))^3} \leq C \varepsilon^2,$$

and, moreover, in every cases,

$$\left\| \hat{P}_\varepsilon \right\|_{L_0^2(\omega \times Y)} \leq C.$$

4 Main convergence results

When ε tends to zero, we obtain for problem (20) different behaviors, depending on the magnitude of a_ε with respect to ε . We will analyze them in the next sections.

4.1 Critical case $a_\varepsilon \approx \varepsilon$, with $a_\varepsilon/\varepsilon \rightarrow \lambda$, $0 < \lambda < +\infty$

First, we obtain some compactness results about the behavior of the sequences $(\tilde{u}_\varepsilon, \tilde{P}_\varepsilon)$ and $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ satisfying the *a priori* estimates given in Lemmas 3.1-i) and 3.4-i), respectively.

Lemma 4.1 (Critical case). *For a subsequence of ε still denote by ε , there exist $\tilde{u} \in H^1(0, 1; L^2(\omega)^3)$, where $\tilde{u}_3 = 0$ and $\tilde{u} = 0$ on $y_3 = \{0, 1\}$, $\hat{u} \in L^2(\omega; H_\#^1(Y)^3)$ (“ $\#$ ” denotes Y' -periodicity), with $\hat{u} = 0$ on $\omega \times Y_s$ and $\hat{u} = 0$ on $y_3 = \{0, 1\}$ such that $\int_Y \hat{u}(x', y) dy = \int_0^1 \tilde{u}(x', y_3) dy_3$ with $\int_Y \hat{u}_3 dy = 0$, and $\hat{P} \in L_0^2(\omega \times Y)$, independent of y , such that*

$$\frac{\tilde{u}_\varepsilon}{a_\varepsilon^2} \rightharpoonup (\tilde{u}', 0) \text{ in } H^1(0, 1; L^2(\omega)^3), \quad (22)$$

$$\frac{\hat{u}_\varepsilon}{a_\varepsilon^2} \rightharpoonup \hat{u} \text{ in } L^2(\omega; H^1(Y)^3), \quad \hat{P}_\varepsilon \rightharpoonup \hat{P} \text{ in } L_0^2(\omega \times Y), \quad (23)$$

$$\operatorname{div}_{x'} \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) = 0 \text{ in } \omega, \quad \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) \cdot n = 0 \text{ on } \partial\omega, \quad (24)$$

$$\operatorname{div}_\lambda \hat{u} = 0 \text{ in } \omega \times Y, \quad \operatorname{div}_{x'} \left(\int_Y \hat{u}'(x', y) dy \right) = 0 \text{ in } \omega, \quad \left(\int_Y \hat{u}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega, \quad (25)$$

where $\operatorname{div}_\lambda = \operatorname{div}_{y'} + \lambda \partial_{y_3}$.

Proof. We refer the reader to Lemmas 5.2, 5.3 and 5.4 in [4] for the proof of (22)-(25). Here, we prove that \hat{P} does not depend on the microscopic variable y . To do this, we choose as test function $\tilde{v}(x', y) \in \mathcal{D}(\omega; C_\#^\infty(Y)^3)$ with $\tilde{v}(x', y) = 0 \in \omega \times Y_s$ (thus, $\tilde{v}(x', x'/a_\varepsilon, y_3) \in (H_0^1(\tilde{\Omega}_\varepsilon))^3$). Setting $a_\varepsilon \tilde{v}(x', x'/a_\varepsilon, y_3)$ in (20) (we recall that $g(\varepsilon) = g(a_\varepsilon)$) and using that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$, we have

$$\begin{aligned} & 2\mu a_\varepsilon \int_\Omega \mathbb{D}_\varepsilon[\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'}[\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right) dx' dy_3 - 2\mu \int_\Omega |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]|^2 dx' dy_3 \\ & + \sqrt{2}g a_\varepsilon^2 \int_\Omega \left| \mathbb{D}_{x'}[\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right| dx' dy_3 - \sqrt{2}g a_\varepsilon \int_\Omega |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]| dx' dy_3 \\ & \geq a_\varepsilon \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \int_\Omega f' \cdot \tilde{v}'_\varepsilon dx' dy_3 + a_\varepsilon \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3 + \frac{a_\varepsilon}{\varepsilon} \int_\Omega \tilde{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy_3. \end{aligned} \quad (26)$$

By the change of variables given in Remark 3.3 and by Lemma 3.4, we get for the first term in relation (26)

$$\begin{aligned} & \int_{\Omega} \mathbb{D}_{\varepsilon} [\tilde{u}_{\varepsilon}] : \left(\mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy_3 \\ &= \int_{\omega \times Y} \left(\frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right) : \left(\frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy + O_{\varepsilon}, \end{aligned} \quad (27)$$

and for the second term in relation (26)

$$\int_{\Omega} |\mathbb{D}_{\varepsilon} [\tilde{u}_{\varepsilon}]|^2 dx' dy_3 = \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right|^2 dx' dy = O_{\varepsilon}. \quad (28)$$

Moreover, applying the change of variables given in Remark 3.3 to the fourth term in relation (26), we have

$$\int_{\Omega} |\mathbb{D}_{\varepsilon} [\tilde{u}_{\varepsilon}]| dx' dy_3 = \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right| dx' dy. \quad (29)$$

Therefore, applying the change of variables given in Remark 3.3 to relation (26), we obtain

$$\begin{aligned} & 2\mu a_{\varepsilon} \int_{\omega \times Y} \left(\frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right) : \left(\frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \\ &+ \sqrt{2}g a_{\varepsilon}^2 \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy - \sqrt{2}g a_{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right| dx' dy + O_{\varepsilon} \\ &\geq a_{\varepsilon} \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \int_{\omega \times Y} f' \cdot \hat{u}'_{\varepsilon} dx' dy + a_{\varepsilon} \int_{\omega \times Y} \hat{P}_{\varepsilon} \operatorname{div}_{x'} \tilde{v}' dx' dy + \int_{\omega \times Y} \hat{P}_{\varepsilon} \operatorname{div}_{y'} \tilde{v}' dx' dy \\ &+ \frac{a_{\varepsilon}}{\varepsilon} \int_{\omega \times Y} \hat{P}_{\varepsilon} \partial_{y_3} \tilde{v}_3 dx' dy + O_{\varepsilon}. \end{aligned} \quad (30)$$

According with (23), the first term in relation (30) can be written by the following way

$$2\mu a_{\varepsilon} \int_{\omega \times Y} \left(\frac{1}{a_{\varepsilon}^2} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{a_{\varepsilon}}{\varepsilon} \frac{1}{a_{\varepsilon}^2} \partial_{y_3} [\hat{u}_{\varepsilon}] \right) : \left(\mathbb{D}_{y'} [\tilde{v}] + \frac{a_{\varepsilon}}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (31)$$

In order to pass to the limit in the first nonlinear term, we have

$$\sqrt{2}g a_{\varepsilon} \int_{\omega \times Y} \left| a_{\varepsilon} \mathbb{D}_{x'} [\tilde{v}] + \mathbb{D}_{y'} [\tilde{v}] + \frac{a_{\varepsilon}}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (32)$$

Now, in order to pass the limit in the second nonlinear term, we are taking into account that

$$\sqrt{2}g a_{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right| dx' dy = \sqrt{2}g a_{\varepsilon}^2 \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}^2} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{a_{\varepsilon}}{\varepsilon} \frac{1}{a_{\varepsilon}^2} \partial_{y_3} [\hat{u}_{\varepsilon}] \right| dx' dy,$$

and using (23) and the fact that the function $E(\varphi) = |\varphi|$ is proper convex continuous, we can deduce that

$$\liminf_{\varepsilon \rightarrow 0} \sqrt{2}g a_{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_{\varepsilon}} \mathbb{D}_{y'} [\hat{u}_{\varepsilon}] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_{\varepsilon}] \right| dx' dy \geq 0. \quad (33)$$

Moreover, using (23) the two first terms in the right hand side of (30) can be written by

$$a_{\varepsilon} \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - a_{\varepsilon}^2 \int_{\omega \times Y} f' \cdot \frac{\hat{u}'_{\varepsilon}}{a_{\varepsilon}^2} dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (34)$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (23), passing to the limit when ε tends to zero, we have

$$\int_{\omega \times Y} \hat{P} \operatorname{div}_{\lambda} \tilde{v} dx' dy. \quad (35)$$

Therefore, taking into account (31)-(35), when we pass to the limit in (30) when ε tends to zero, we have $0 \geq \int_{\omega \times Y} \hat{P} \operatorname{div}_\lambda \tilde{v} dx' dy$. Now, if we choose as test function $-a_\varepsilon \tilde{v}(x', x'/a_\varepsilon, y_3)$ in (20) and we argue similarly, we obtain $\int_{\omega \times Y} \hat{P} \operatorname{div}_\lambda \tilde{v} dx' dy \geq 0$. Thus, we can deduce that $\int_{\omega \times Y} \hat{P} \operatorname{div}_\lambda \tilde{v} dx' dy = 0$, which shows that \hat{P} does not depend on y . \square

Theorem 4.2 (Critical case). *If $a_\varepsilon \approx \varepsilon$, with $a_\varepsilon/\varepsilon \rightarrow \lambda$, $0 < \lambda < +\infty$, then $(\hat{u}_\varepsilon/a_\varepsilon^2, \hat{P}_\varepsilon)$ converges to (\hat{u}, \hat{P}) in $L^2(\omega; H^1(Y)^3) \times L_0^2(\omega \times Y)$, which satisfies the following variational inequality*

$$\begin{aligned} & 2\mu \int_{\omega \times Y} \mathbb{D}_\lambda [\hat{u}] : (\mathbb{D}_\lambda [\tilde{v}] - \mathbb{D}_\lambda [\hat{u}]) dx' dy + \sqrt{2}g \int_{\omega \times Y} |\mathbb{D}_\lambda [\tilde{v}]| dx' dy - \sqrt{2}g \int_{\omega \times Y} |\mathbb{D}_\lambda [\hat{u}]| dx' dy \\ & \geq \int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy - \int_{\omega \times Y} \nabla_{x'} \hat{P} (\tilde{v}' - \hat{u}') dx' dy, \end{aligned} \quad (36)$$

where $\mathbb{D}_\lambda[\cdot] = \mathbb{D}_{y'}[\cdot] + \lambda \partial_{y_3}[\cdot]$ and for every $\tilde{v} \in L^2(\omega; H^1(Y)^3)$ such that

$$\tilde{v}(x', y) = 0 \text{ in } \omega \times Y_s, \quad \operatorname{div}_\lambda \tilde{v} = 0 \text{ in } \omega \times Y, \quad \left(\int_Y \tilde{v}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega.$$

Proof. We choose a test function $\tilde{v}(x', y) \in \mathcal{D}(\omega; C_0^\infty(Y)^3)$ with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ (thus, we have that $\tilde{v}(x', x'/a_\varepsilon, y_3) \in (H_0^1(\tilde{\Omega}_\varepsilon))^3$). We first multiply (20) by a_ε^{-2} and we use that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$. Then, we take as test function $a_\varepsilon^2 \tilde{v}_\varepsilon = a_\varepsilon^2(\tilde{v}'(x', x'/a_\varepsilon, y_3), \lambda \varepsilon/a_\varepsilon v_3(x', x'/a_\varepsilon, y_3))$, with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ and satisfying the incompressibility conditions (25), that is, $\operatorname{div}_\lambda \tilde{v} = 0$ in $\omega \times Y$ and $(\int_Y \tilde{v}'(x', y) dy) \cdot n = 0$ on $\partial\omega$, and we have

$$\begin{aligned} & 2\mu \int_\Omega \mathbb{D}_\varepsilon [\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'} [\tilde{v}_\varepsilon] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right) dx' dy_3 - 2\mu \frac{1}{a_\varepsilon^2} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 \\ & + \sqrt{2}g a_\varepsilon \int_\Omega \left| \mathbb{D}_{x'} [\tilde{v}_\varepsilon] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right| dx' dy_3 - \sqrt{2}g \frac{1}{a_\varepsilon} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]| dx' dy_3 \\ & \geq \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \frac{1}{a_\varepsilon^2} \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \frac{1}{a_\varepsilon} \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3 + \frac{\lambda}{a_\varepsilon} \int_\Omega \tilde{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy_3. \end{aligned} \quad (37)$$

By the change of variables given in Remark 3.3 and by Lemma 3.4, we have (27) for the first term in relation (37), and for the second term in relation (37) we obtain

$$\int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 = \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right|^2 dx' dy. \quad (38)$$

Moreover, applying the change of variables given in Remark 3.3 to the fourth term in relation (37), we have (29). Therefore, applying the change of variables given in Remark 3.3 to relation (37), we obtain

$$\begin{aligned} & 2\mu \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right) dx' dy - 2\mu \frac{1}{a_\varepsilon^2} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right|^2 dx' dy \\ & + \sqrt{2}g a_\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}_\varepsilon] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right| dx' dy - \sqrt{2}g \frac{1}{a_\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\ & \geq \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \frac{1}{a_\varepsilon^2} \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy + \frac{1}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy \\ & + \frac{\lambda}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy + O_\varepsilon. \end{aligned} \quad (39)$$

According with (23), the first term in relation (39) can be written

$$2\mu \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \frac{1}{a_\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right) dx' dy,$$

and, taking into account that $\lambda \varepsilon / a_\varepsilon \rightarrow 1$, this term tends to the following limit

$$2\mu \int_{\omega \times Y} (\mathbb{D}_{y'} [\hat{u}] + \lambda \partial_{y_3} [\hat{u}]) : (\mathbb{D}_{y'} [\tilde{v}] + \lambda \partial_{y_3} [\tilde{v}]) dx' dy. \quad (40)$$

The second term in relation (39) writes

$$2\mu \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \frac{1}{a_\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \frac{1}{a_\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) dx' dy,$$

and, taking into account that the function $B(\varphi) = |\varphi|$ is proper convex continuous and $\lambda \varepsilon / a_\varepsilon \rightarrow 1$, we get that the $\liminf_{\varepsilon \rightarrow 0}$ of this second is greater or equal than

$$2\mu \int_{\omega \times Y} (\mathbb{D}_{y'} [\hat{u}] + \lambda \partial_{y_3} [\hat{u}]) : (\mathbb{D}_{y'} [\hat{u}] + \lambda \partial_{y_3} [\hat{u}]) dx' dy. \quad (41)$$

In order to pass to the limit in the first nonlinear term, we have

$$\begin{aligned} & \left| \sqrt{2g} a_\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}_\varepsilon] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] \right| dx' dy - \sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}] + \lambda \partial_{y_3} [\tilde{v}]| dx' dy \right| \\ & \leq \sqrt{2g} \int_{\omega \times Y} \left| a_\varepsilon \mathbb{D}_{x'} [\tilde{v}_\varepsilon] + \mathbb{D}_{y'} [\tilde{v}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] - \mathbb{D}_{y'} [\tilde{v}] - \lambda \partial_{y_3} [\tilde{v}] \right| dx' dy \\ & \leq \sqrt{2g} \int_{\omega \times Y} |a_\varepsilon \mathbb{D}_{x'} [\tilde{v}_\varepsilon]| dx' dy + \sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}_\varepsilon] - \mathbb{D}_{y'} [\tilde{v}]| dx' dy \\ & + \sqrt{2g} \int_{\omega \times Y} \left| \frac{a_\varepsilon}{\varepsilon} \partial_{y_3} [\tilde{v}_\varepsilon] - \lambda \partial_{y_3} [\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0, \end{aligned}$$

and we can deduce that the first nonlinear term tends to the following limit

$$\sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}] + \lambda \partial_{y_3} [\tilde{v}]| dx' dy. \quad (42)$$

Now, in order to pass the limit in the second nonlinear term, we are taking into account that

$$\sqrt{2g} \frac{1}{a_\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy = \sqrt{2g} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \frac{1}{a_\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy,$$

and using (23) and the fact that the function $E(\varphi) = |\varphi|$ is proper convex continuous, we can deduce that

$$\liminf_{\varepsilon \rightarrow 0} \sqrt{2g} \frac{1}{a_\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy \geq \sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\hat{u}] + \lambda \partial_{y_3} [\hat{u}]| dx' dy. \quad (43)$$

Moreover, using (23) the two first terms in the right hand side of (39) tend to the following limit

$$\int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy. \quad (44)$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (23) the first term of the pressure tends to the following limit $\int_{\omega \times Y} \hat{P} \operatorname{div}_{x'} \tilde{v}' dx' dy$, and using (25) and taking into account that \hat{P} does not depend on y , we have

$$\int_{\omega \times Y} \hat{P} \operatorname{div}_{x'} \tilde{v}' dx' dy = \int_{\omega \times Y} \hat{P} \operatorname{div}_{x'} \tilde{v}' dx' dy - \int_{\omega} \hat{P} \left(\operatorname{div}_{x'} \int_Y \hat{u}' dy \right) dx' = - \int_{\omega \times Y} \nabla_{x'} \hat{P} (\tilde{v}' - \hat{u}') dx' dy \quad (45)$$

Finally, using that $\operatorname{div}_\lambda \tilde{v} = 0$, we have

$$\frac{1}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy + \frac{\lambda}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy = 0. \quad (46)$$

Therefore, taking into account (40)-(46), we have (36). \square

4.2 Subcritical case $a_\varepsilon \ll \varepsilon$ ($\lambda = 0$)

We obtain some compactness results about the behavior of the sequences $(\tilde{u}_\varepsilon, \tilde{P}_\varepsilon)$ and $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ satisfying the *a priori* estimates given in Lemmas 3.1-i) and 3.4-i), respectively.

Lemma 4.3 (Subcritical case). *For a subsequence of ε still denoted by ε , there exist $\tilde{u} \in (L^2(\Omega))^3$, where $\tilde{u}_3 = 0$ and $\tilde{u} = 0$ on $y_3 = \{0, 1\}$, $\hat{u} \in L^2(\Omega; H^1_\#(Y')^3)$ (“ $\#$ ” denotes Y' -periodicity), with $\hat{u} = 0$ in $\omega \times Y_s$ and $\hat{u} = 0$ on $y_3 = \{0, 1\}$ such that $\int_Y \hat{u}(x', y) dy = \int_0^1 \tilde{u}(x', y_3) dy_3$ with $\int_Y \hat{u}_3 dy = 0$ and \hat{u}_3 independent of y_3 , and $\hat{P} \in L^2_0(\omega \times Y)$, independent of y , such that*

$$\frac{\tilde{u}_\varepsilon}{a_\varepsilon^2} \rightharpoonup (\tilde{u}', 0) \text{ in } (L^2(\Omega))^3, \quad (47)$$

$$\frac{\hat{u}_\varepsilon}{a_\varepsilon^2} \rightharpoonup \hat{u} \text{ in } L^2(\Omega; H^1(Y')^3), \quad \hat{P}_\varepsilon \rightharpoonup \hat{P} \text{ in } L^2_0(\omega \times Y), \quad (48)$$

$$\operatorname{div}_{x'} \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) = 0 \text{ in } \omega, \quad \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) \cdot n = 0 \text{ on } \partial\omega, \quad (49)$$

$$\operatorname{div}_{y'} \hat{u}' = 0 \text{ in } \omega \times Y, \quad \operatorname{div}_{x'} \left(\int_Y \hat{u}'(x', y) dy \right) = 0 \text{ in } \omega, \quad \left(\int_Y \hat{u}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega. \quad (50)$$

Proof. See Lemmas 5.2, 5.3 and 5.4 in [4] for the proof of (47)-(50). In order to prove that \hat{P} does not depend on y' we argue as in the proof of Lemma 4.1 using that $a_\varepsilon \ll \varepsilon$, and we obtain $\int_{\omega \times Y} \hat{P} \operatorname{div}_{y'} \tilde{v}' dx' dy = 0$, which shows that \hat{P} does not depend on y' . Now, in order to prove that \hat{P} does not depend on y_3 , setting $\varepsilon \tilde{v} = \varepsilon(0, \tilde{v}_3(x', x'/a_\varepsilon, y_3))$ in (20) (we recall that $g(\varepsilon) = g a_\varepsilon$) and using that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$, we have

$$\begin{aligned} & 2\mu\varepsilon \int_\Omega \mathbb{D}_\varepsilon[\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'}[\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right) dx' dy_3 - 2\mu \int_\Omega |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]|^2 dx' dy_3 \\ & + \sqrt{2} g a_\varepsilon \varepsilon \int_\Omega \left| \mathbb{D}_{x'}[\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right| dx' dy_3 - \sqrt{2} g a_\varepsilon \int_\Omega |\mathbb{D}_\varepsilon[\tilde{u}_\varepsilon]| dx' dy_3 \\ & \geq - \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy_3. \end{aligned} \quad (51)$$

Applying the change of variables given in Remark 3.3 to relation (51) and taking into account (27)-(29), we obtain

$$\begin{aligned} & 2\mu\varepsilon \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3}[\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right) dx' dy \\ & + \sqrt{2} g a_\varepsilon \varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'}[\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3}[\tilde{v}] \right| dx' dy - \sqrt{2} g a_\varepsilon \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'}[\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3}[\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\ & \geq - \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + \int_{\omega \times Y} \hat{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy + O_\varepsilon. \end{aligned} \quad (52)$$

According with (48) and using that $a_\varepsilon \ll \varepsilon$, the first term in relation (52) can be written by the following way

$$2\mu\varepsilon \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon^2} \mathbb{D}_{y'}[\hat{u}_\varepsilon] + \frac{a_\varepsilon}{\varepsilon} \frac{1}{a_\varepsilon^2} \partial_{y_3}[\hat{u}_\varepsilon] \right) : \left(\mathbb{D}_{y'}[\tilde{v}] + \frac{a_\varepsilon}{\varepsilon} \partial_{y_3}[\tilde{v}] \right) dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (53)$$

In order to pass to the limit in the first nonlinear term, we have

$$\sqrt{2} g \varepsilon \int_{\omega \times Y} \left| a_\varepsilon \mathbb{D}_{x'}[\tilde{v}] + \mathbb{D}_{y'}[\tilde{v}] + \frac{a_\varepsilon}{\varepsilon} \partial_{y_3}[\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (54)$$

In order to pass to the limit in the second nonlinear term, we proceed as in Lemma 4.1. Moreover, using (48) the first term in the right hand side of (52) can be written by

$$a_\varepsilon^2 \int_{\omega \times Y} f' \cdot \frac{\hat{u}'_\varepsilon}{a_\varepsilon^2} dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (55)$$

We consider now the term which involves the pressure. Taking into account the convergence of the pressure (48), passing to the limit when ε tends to zero, we have

$$\int_{\omega \times Y} \hat{P} \partial_{y_3} \tilde{v}_3 dx' dy. \quad (56)$$

Therefore, taking into account (33) and (53)-(56), when we pass to the limit in (52) when ε tends to zero, we have $0 \geq \int_{\omega \times Y} \hat{P} \partial_{y_3} \tilde{v}_3 dx' dy$. Now, if we choose as test function $-\varepsilon \tilde{v} = -\varepsilon(0, \tilde{v}_3(x', x'/a_\varepsilon, y_3))$ in (20) and we argue similarly, we can deduce that \hat{P} does not depend on y_3 , so \hat{P} does not depend on y . \square

Theorem 4.4 (Subcritical case). *If $a_\varepsilon \ll \varepsilon$, then $(\hat{u}_\varepsilon/a_\varepsilon^2, \hat{P}_\varepsilon)$ converges to (\hat{u}, \hat{P}) in $L^2(\Omega; H^1(Y')^3) \times L^2_0(\omega \times Y)$, which satisfies the following variational inequality*

$$\begin{aligned} & 2\mu \int_{\omega \times Y} \mathbb{D}_{y'} [\hat{u}'] : (\mathbb{D}_{y'} [\tilde{v}'] - \mathbb{D}_{y'} [\hat{u}']) dx' dy + \sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}']| dx' dy - \sqrt{2g} \int_{\omega \times Y} |\mathbb{D}_{y'} [\hat{u}']| dx' dy \\ & \geq \int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy - \int_{\omega \times Y} \nabla_{x'} \hat{P} (\tilde{v}' - \hat{u}') dx' dy, \end{aligned} \quad (57)$$

for every $\tilde{v} \in L^2(\Omega; H^1(Y')^3)$ such that

$$\tilde{v}(x', y) = 0 \text{ in } \omega \times Y_s, \quad \operatorname{div}_{y'} \tilde{v}' = 0 \text{ in } \omega \times Y, \quad \left(\int_Y \tilde{v}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega.$$

Proof. We choose a test function $\tilde{v}(x', y) \in \mathcal{D}(\omega; C^\infty_\#(Y)^3)$ with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ (thus, we have that $\tilde{v}(x', x'/a_\varepsilon, y_3) \in (H^1_0(\tilde{\Omega}_\varepsilon))^3$). We first multiply (20) by a_ε^{-2} and we use that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$. Then, we take a test function $a_\varepsilon^2 \tilde{v}(x', x'/a_\varepsilon, y_3)$, with \tilde{v}_3 independent of y_3 and with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ and satisfying the incompressibility conditions (50), that is, $\operatorname{div}_{y'} \tilde{v}' = 0$ in $\omega \times Y$ and $(\int_Y \tilde{v}'(x', y) dy) \cdot n = 0$ on $\partial\omega$, and we have

$$\begin{aligned} & 2\mu \int_\Omega \mathbb{D}_\varepsilon [\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy_3 - 2\mu \frac{1}{a_\varepsilon^2} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 \\ & + \sqrt{2g} a_\varepsilon \int_\Omega \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy_3 - \sqrt{2g} \frac{1}{a_\varepsilon} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]| dx' dy_3 \\ & \geq \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \frac{1}{a_\varepsilon^2} \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \frac{1}{a_\varepsilon} \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3. \end{aligned} \quad (58)$$

Applying the change of variables given in Remark 3.3 to relation (58) and taking into account (27), (29) and (38), we obtain

$$\begin{aligned} & 2\mu \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \\ & - 2\mu \frac{1}{a_\varepsilon^2} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right|^2 dx' dy + \sqrt{2g} a_\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy \\ & - \sqrt{2g} \frac{1}{a_\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\ & \geq \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \frac{1}{a_\varepsilon^2} \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy + \frac{1}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy + O_\varepsilon. \end{aligned} \quad (59)$$

In the left-hand side, we only give the details of convergence for the first nonlinear term, the most challenging one.

$$\begin{aligned}
 & \left| \sqrt{2}g a_\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy - \sqrt{2}g \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}]| dx' dy \right| \\
 & \leq \sqrt{2}g \int_{\omega \times Y} \left| a_\varepsilon \mathbb{D}_{x'} [\tilde{v}] + \mathbb{D}_{y'} [\tilde{v}] + \frac{a_\varepsilon}{\varepsilon} \partial_{y_3} [\tilde{v}] - \mathbb{D}_{y'} [\tilde{v}] \right| dx' dy \\
 & \leq \sqrt{2}g \int_{\omega \times Y} |a_\varepsilon \mathbb{D}_{x'} [\tilde{v}]| dx' dy + \sqrt{2}g \int_{\omega \times Y} \left| \frac{a_\varepsilon}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0.
 \end{aligned}$$

Using (48) the two first terms in the right hand side of (59) tend to the following limit

$$\int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy.$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (48) the first term of the pressure tends to the following limit $\int_{\omega \times Y} \hat{P} \operatorname{div}_{x'} \tilde{v}' dx' dy$, and using (50) and taking into account that \hat{P} does not depend on y , we have (45). Finally, using that $\operatorname{div}_{y'} \tilde{v}' = 0$, we have

$$\frac{1}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy = 0. \quad (60)$$

It is straightforward to obtain that $\hat{u}_3 = 0$ and therefore we get (57). \square

4.3 Supercritical case $a_\varepsilon \gg \varepsilon$ ($\lambda = +\infty$)

We obtain some compactness results about the behavior of the sequences $(\tilde{u}_\varepsilon, \tilde{P}_\varepsilon)$ and $(\hat{u}_\varepsilon, \hat{P}_\varepsilon)$ satisfying the *a priori* estimates given in Lemmas 3.1-ii) and 3.4-ii), respectively.

Lemma 4.5 (Supercritical case). *For a subsequence of ε still denote by ε , there exist $\tilde{u} \in H^1(0, 1; L^2(\omega)^3)$, where $\tilde{u}_3 = 0$ and $\tilde{u} = 0$ on $y_3 = \{0, 1\}$, $\hat{u} \in H^1(0, 1; L^2_\#(\omega \times Y')^3)$ (“ $\#$ ” denotes Y' -periodicity), with $\hat{u} = 0$ in $\omega \times Y_s$, $\hat{u} = 0$ on $y_3 = \{0, 1\}$ such that $\int_Y \hat{u}(x', y) dy = \int_0^1 \tilde{u}(x', y_3) dy_3$ with $\int_Y \hat{u}_3 dy = 0$ and \hat{u}_3 independent of y_3 , and $\hat{P} \in L^2_0(\omega \times Y)$, independent of y , such that*

$$\frac{\tilde{u}_\varepsilon}{\varepsilon^2} \rightharpoonup (\tilde{u}', 0) \text{ in } H^1(0, 1; L^2(\omega)^3), \quad (61)$$

$$\frac{\hat{u}_\varepsilon}{\varepsilon^2} \rightharpoonup \hat{u} \text{ in } H^1(0, 1; L^2(\omega \times Y')^3), \quad \hat{P}_\varepsilon \rightharpoonup \hat{P} \text{ in } L^2_0(\omega \times Y), \quad (62)$$

$$\operatorname{div}_{x'} \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) = 0 \text{ in } \omega, \quad \left(\int_0^1 \tilde{u}'(x', y_3) dy_3 \right) \cdot n = 0 \text{ on } \partial\omega, \quad (63)$$

$$\operatorname{div}_{y'} \hat{u}' = 0 \text{ in } \omega \times Y, \quad \operatorname{div}_{x'} \left(\int_Y \hat{u}'(x', y) dy \right) = 0 \text{ in } \omega, \quad \left(\int_Y \hat{u}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega. \quad (64)$$

Proof. See Lemmas 5.2, 5.3 and 5.4 in [4] for the proof of (61)-(64). Here, we prove that \hat{P} does not depend on the microscopic variable y . To do this, we choose as test function $\tilde{v}(x', y) \in \mathcal{D}(\omega; C^\infty_\#(Y)^3)$ with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ (thus, $\tilde{v}(x', x'/a_\varepsilon, y_3) \in (H^1_0(\tilde{\Omega}_\varepsilon))^3$). In order to prove that \hat{P} does not depend on y_3 , we set $\varepsilon \tilde{v}(x', x'/a_\varepsilon, y_3)$ in (20) (we recall that $g(\varepsilon) = g\varepsilon$) and using that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$, we have

$$\begin{aligned}
 & 2\mu\varepsilon \int_\Omega \mathbb{D}_\varepsilon [\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy_3 - 2\mu \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 \\
 & + \sqrt{2}g\varepsilon^2 \int_\Omega \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy_3 - \sqrt{2}g\varepsilon \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]| dx' dy_3 \\
 & \geq \varepsilon \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + \varepsilon \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \frac{\varepsilon}{a_\varepsilon} \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy_3.
 \end{aligned} \quad (65)$$

Applying the change of variables given in Remark 3.3 to relation (65) and taking into account (27)-(29), we obtain

$$\begin{aligned}
 & 2\mu\varepsilon \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \\
 & + \sqrt{2}g\varepsilon^2 \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy - \sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\
 & \geq \varepsilon \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + \varepsilon \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy + \frac{\varepsilon}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy \\
 & + \int_{\omega \times Y} \hat{P}_\varepsilon \partial_{y_3} \tilde{v}_3 dx' dy + O_\varepsilon.
 \end{aligned} \tag{66}$$

According with (62) and using that $a_\varepsilon \gg \varepsilon$, one has for the first term in relation (66)

$$2\mu\varepsilon \int_{\omega \times Y} \left(\frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] \right) dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{67}$$

We pass to the limit in the first nonlinear term and we have

$$\sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \varepsilon \mathbb{D}_{x'} [\tilde{v}] + \frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{68}$$

In order to pass the limit in the second nonlinear term, we taking into account that

$$\sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy = \sqrt{2}g\varepsilon^2 \int_{\omega \times Y} \left| \frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy,$$

and using (62), with $a_\varepsilon \gg \varepsilon$, and the fact that the function $E(\varphi) = |\varphi|$ is proper convex continuous, we can deduce that

$$\liminf_{\varepsilon \rightarrow 0} \sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy \geq 0. \tag{69}$$

Moreover, using (62) the two first terms in the right hand side of (66) can be written by

$$\varepsilon \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \varepsilon^2 \int_{\omega \times Y} f' \cdot \frac{\hat{u}'_\varepsilon}{\varepsilon^2} dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{70}$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (62) and $a_\varepsilon \gg \varepsilon$, passing to the limit when ε tends to zero, we have

$$\int_{\omega \times Y} \hat{P} \partial_{y_3} \tilde{v}_3 dx' dy. \tag{71}$$

Therefore, taking into account (67)-(71), when we pass to the limit in (66) when ε tends to zero, we have $0 \geq \int_{\omega \times Y} \hat{P} \partial_{y_3} \tilde{v}_3 dx' dy$. Now, if we choose as test function $-\varepsilon \tilde{v}(x', x'/a_\varepsilon, y_3)$ in (20) and we argue similarly, we can deduce that \hat{P} does not depend on y_3 .

Now, in order to prove that \hat{P} does not depend on y' , we set $a_\varepsilon \tilde{v} = a_\varepsilon(\tilde{v}'(x', x'/a_\varepsilon, y_3), 0)$ in (20) and using that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$, we have

$$\begin{aligned}
 & 2\mu a_\varepsilon \int_\Omega \mathbb{D}_\varepsilon [\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy_3 - 2\mu \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 \\
 & + \sqrt{2}g\varepsilon a_\varepsilon \int_\Omega \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy_3 - \sqrt{2}g\varepsilon \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]| dx' dy_3 \\
 & \geq a_\varepsilon \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + a_\varepsilon \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3.
 \end{aligned} \tag{72}$$

Applying the change of variables given in Remark 3.3 to relation (72) and taking into account (27)-(29), we obtain

$$\begin{aligned}
 & 2\mu a_\varepsilon \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \\
 & + \sqrt{2} g \varepsilon a_\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy - \sqrt{2} g \varepsilon \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\
 & \geq a_\varepsilon \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + a_\varepsilon \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy + \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy.
 \end{aligned} \tag{73}$$

According with (62) and using that $a_\varepsilon \gg \varepsilon$, the first term in relation (73) can be written by the following way

$$2\mu a_\varepsilon \int_{\omega \times Y} \left(\frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] \right) dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{74}$$

In order to pass to the limit in the first nonlinear term, we have

$$\sqrt{2} g a_\varepsilon \int_{\omega \times Y} \left| \varepsilon \mathbb{D}_{x'} [\tilde{v}] + \frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] \right| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{75}$$

Moreover, using (62) the two first terms in the right hand side of (73) can be written by

$$a_\varepsilon \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \varepsilon^2 \int_{\omega \times Y} f' \cdot \frac{\hat{u}'_\varepsilon}{\varepsilon^2} dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \tag{76}$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (62), passing to the limit when ε tends to zero, we have

$$\int_{\omega \times Y} \hat{P} \operatorname{div}_{y'} \tilde{v}' dx' dy. \tag{77}$$

Therefore, taking into account (69) and (74)-(77), when we pass to the limit in (73) when ε tends to zero, we have $0 \geq \int_{\omega \times Y} \hat{P} \operatorname{div}_{y'} \tilde{v}' dx' dy$. Now, if we choose as test function $-a_\varepsilon \tilde{v} = -a_\varepsilon (\tilde{v}'(x', x'/a_\varepsilon, y_3), 0)$ in (20) and we argue similarly, we can deduce that \hat{P} does not depend on y' , so \hat{P} does not depend on y . \square

Theorem 4.6 (Supercritical case). *If $a_\varepsilon \gg \varepsilon$, then $(\hat{u}_\varepsilon/\varepsilon^2, \hat{P}_\varepsilon)$ converges to (\hat{u}, \hat{P}) in $H^1(0, 1; L^2(\omega \times Y')^3) \times L^2_0(\omega \times Y)$, which satisfies the following variational equality*

$$\begin{aligned}
 & 2\mu \int_{\omega \times Y} \partial_{y_3} [\hat{u}'] : (\partial_{y_3} [\tilde{v}'] - \partial_{y_3} [\hat{u}']) dx' dy + \sqrt{2} g \int_{\omega \times Y} |\partial_{y_3} [\tilde{v}']| dx' dy - \sqrt{2} g \int_{\omega \times Y} |\partial_{y_3} [\hat{u}']| dx' dy \\
 & \geq \int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy - \int_{\omega \times Y} \nabla_{x'} \hat{P} (\tilde{v}' - \hat{u}') dx' dy,
 \end{aligned} \tag{78}$$

for every $\tilde{v} \in H^1(0, 1; L^2(\omega \times Y')^3)$ such that

$$\tilde{v}(x', y) = 0 \text{ in } \omega \times Y_s, \quad \operatorname{div}_{y'} \tilde{v}' = 0 \text{ in } \omega \times Y, \quad \left(\int_Y \tilde{v}'(x', y) dy \right) \cdot n = 0 \text{ on } \partial\omega.$$

Proof. We choose a test function $\tilde{v}(x', y) \in \mathcal{D}(\omega; C^\infty_\#(Y)^3)$ with $\tilde{v}(x', y) = 0 \in \omega \times Y_s$ (thus, $\tilde{v}(x', x'/a_\varepsilon, y_3) \in (H^1_0(\tilde{\Omega}_\varepsilon))^3$). We first multiply (20) by ε^{-2} and we use that $\operatorname{div}_\varepsilon \tilde{u}_\varepsilon = 0$. Then, we take a test function $\varepsilon^2 \tilde{v}(x', x'/a_\varepsilon, y_3)$, with \tilde{v}_3 independent of y_3 and with $\tilde{v}(x', y) = 0$ in $\omega \times Y_s$ and satisfying the incompressibility conditions (64), that is, $\operatorname{div}_{y'} \tilde{v}' = 0$ in $\omega \times Y$ and $(\int_Y \tilde{v}'(x', y) dy) \cdot n = 0$ on $\partial\omega$, and we have

$$\begin{aligned}
 & 2\mu \int_\Omega \mathbb{D}_\varepsilon [\tilde{u}_\varepsilon] : \left(\mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy_3 - 2\mu \frac{1}{\varepsilon^2} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]|^2 dx' dy_3 \\
 & + \sqrt{2} g \varepsilon \int_\Omega \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy_3 - \sqrt{2} g \frac{1}{\varepsilon} \int_\Omega |\mathbb{D}_\varepsilon [\tilde{u}_\varepsilon]| dx' dy_3 \\
 & \geq \int_\Omega f' \cdot \tilde{v}' dx' dy_3 - \frac{1}{\varepsilon^2} \int_\Omega f' \cdot \tilde{u}'_\varepsilon dx' dy_3 + \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy_3 + \frac{1}{a_\varepsilon} \int_\Omega \tilde{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy_3.
 \end{aligned} \tag{79}$$

Applying the change of variables given in Remark 3.3 to relation (79), arguing as in the critical case, we obtain

$$\begin{aligned}
 & 2\mu \int_{\omega \times Y} \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right) dx' dy \tag{80} \\
 & - 2\mu \frac{1}{\varepsilon^2} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right|^2 dx' dy + \sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy \\
 & - \sqrt{2}g \frac{1}{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy + O_\varepsilon \\
 & \geq \int_{\omega \times Y} f' \cdot \tilde{v}' dx' dy - \frac{1}{\varepsilon^2} \int_{\omega \times Y} f' \cdot \hat{u}'_\varepsilon dx' dy + \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{x'} \tilde{v}' dx' dy + \frac{1}{a_\varepsilon} \int_{\omega \times Y} \hat{P}_\varepsilon \operatorname{div}_{y'} \tilde{v}' dx' dy + O_\varepsilon.
 \end{aligned}$$

According with (62), the first term in relation (80) can be written by the following way

$$2\mu \int_{\omega \times Y} \left(\frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] \right) dx' dy,$$

and, taking into account that $a_\varepsilon \gg \varepsilon$, this term tends to the following limit

$$2\mu \int_{\omega \times Y} \partial_{y_3} [\hat{u}'] : \partial_{y_3} [\tilde{v}'] dx' dy. \tag{81}$$

The second term in relation (80) writes

$$2\mu \int_{\omega \times Y} \left(\frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) : \left(\frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right) dx' dy,$$

and, taking into account that the function $B(\varphi) = |\varphi|$ is proper convex continuous and $a_\varepsilon \gg \varepsilon$, we get that the $\liminf_{\varepsilon \rightarrow 0}$ of this second is greater or equal than

$$2\mu \int_{\omega \times Y} \partial_{y_3} [\hat{u}'] : \partial_{y_3} [\hat{u}'] dx' dy. \tag{82}$$

In order to pass to the limit in the first nonlinear term, using that $a_\varepsilon \gg \varepsilon$, we have

$$\begin{aligned}
 & \left| \sqrt{2}g\varepsilon \int_{\omega \times Y} \left| \mathbb{D}_{x'} [\tilde{v}] + \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \frac{1}{\varepsilon} \partial_{y_3} [\tilde{v}] \right| dx' dy - \sqrt{2}g \int_{\omega \times Y} |\partial_{y_3} [\tilde{v}']| dx' dy \right| \\
 & \leq \sqrt{2}g \int_{\omega \times Y} \left| \varepsilon \mathbb{D}_{x'} [\tilde{v}] + \frac{\varepsilon}{a_\varepsilon} \mathbb{D}_{y'} [\tilde{v}] + \partial_{y_3} [\tilde{v}] - \partial_{y_3} [\tilde{v}] \right| dx' dy \\
 & \leq \sqrt{2}g \int_{\omega \times Y} |\varepsilon \mathbb{D}_{x'} [\tilde{v}]| dx' dy + \sqrt{2}g \frac{\varepsilon}{a_\varepsilon} \int_{\omega \times Y} |\mathbb{D}_{y'} [\tilde{v}]| dx' dy \rightarrow 0, \text{ as } \varepsilon \rightarrow 0.
 \end{aligned}$$

Now, in order to pass the limit in the second nonlinear term, taking into account that

$$\sqrt{2}g \frac{1}{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy = \sqrt{2}g \int_{\omega \times Y} \left| \frac{\varepsilon}{a_\varepsilon} \frac{1}{\varepsilon^2} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon^2} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy,$$

and using (62) and the fact that the function $E(\varphi) = |\varphi|$ is proper convex continuous and $a_\varepsilon \gg \varepsilon$, we can deduce that

$$\liminf_{\varepsilon \rightarrow 0} \sqrt{2}g \frac{1}{\varepsilon} \int_{\omega \times Y} \left| \frac{1}{a_\varepsilon} \mathbb{D}_{y'} [\hat{u}_\varepsilon] + \frac{1}{\varepsilon} \partial_{y_3} [\hat{u}_\varepsilon] \right| dx' dy \geq \sqrt{2}g \int_{\omega \times Y} |\partial_{y_3} [\hat{u}']| dx' dy. \tag{83}$$

Moreover, using (62) the two first terms in the right hand side of (80) tend to the following limit

$$\int_{\omega \times Y} f' \cdot (\tilde{v}' - \hat{u}') dx' dy. \tag{84}$$

We consider now the terms which involve the pressure. Taking into account the convergence of the pressure (62) the first term of the pressure tends to the following limit $\int_{\omega \times Y} \hat{P} \operatorname{div}_{x'} \tilde{v}' dx' dy$, and using (64) and taking into account that \hat{P} does not depend on y , we have (45). Finally using that $\operatorname{div}_{y'} \tilde{v}' = 0$, we have (60). Therefore, taking into account (45), (60) and (81)-(84), we get (78). \square

5 Conclusions

By using dimension reduction and homogenization techniques, we studied the limiting behavior of the velocity and of the pressure for a nonlinear viscoplastic Bingham flow with small yield stress, in a thin porous medium of small height ε and for which the relative dimension of the pores is a_ε . Three cases are studied following the value of $\lambda = \lim_{\varepsilon \rightarrow 0} a_\varepsilon/\varepsilon$ and, at the limit, they all preserve the nonlinear character of the flow. More precisely, according to [29], each of the limit problems (36), (57) and (78), is written as a nonlinear Darcy equation:

$$\begin{cases} \tilde{U}'(x') = K^\lambda \left(f'(x') - \nabla_{x'} \hat{P}(x') \right) & \text{in } \omega, \\ \operatorname{div}_{x'} \tilde{U}'(x') = 0 & \text{in } \omega, \\ \tilde{U}'(x') \cdot n = 0 & \text{on } \partial\omega. \end{cases} \quad (85)$$

The velocity of filtration $\tilde{U}(x') = \left(\tilde{U}'(x'), \tilde{U}_3(x') \right)$ is defined by

$$\tilde{U}(x') = \int_Y \hat{u}(x', y) dy = \int_0^1 \left(\int_{Y'} \hat{u}(x', y', y_3) dy' \right) dy_3 = \int_0^1 \tilde{u}(x', y_3) dy_3.$$

We remark that in all three cases, the vertical component \tilde{U}_3 of the velocity of filtration equals zero and this result is in accordance with the previous mathematical studies of the flow in this thin porous medium, for newtonian fluids (Stokes and Navier-Stokes equations) and for power law fluids (see [20], [1], [2], [4], [5]). Moreover, despite the fact that the limit pressure is not unique, the velocity of filtration is uniquely determined (see Section 4.3 in [29]). In (85), the function $K^\lambda : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is nonlinear and its expression can not be made explicit for the Bingham flow (see [29]). Nevertheless, in each case, for a given $\xi \in \mathbb{R}^2$, one has $K^\lambda(\xi) = \int_Y \chi_\lambda^\xi(y) dy$, with χ_λ^ξ solution of a local problem stated in the cell Y . If $0 < \lambda < +\infty$, the local problem is a 3-D Bingham problem. If $\lambda = 0$, the local problem is a 2-D Bingham problem (defined for each $y_3 \in]0, 1[$), while if $\lambda = +\infty$ the 1-D local problem (defined for each $y' \in Y'$) corresponds to a lower-dimensional Bingham-like law (see [15]).

We end with the remark that if in the initial problem (9) we take $g = 0$, then the problem under study becomes the Stokes problem. We refer to [4] (case $p = 2$) for the asymptotic analysis of the Stokes problem. If we set $g = 0$ in the limit problems (36), (57) and (78), they become exactly the ones in [4], Theorem 6.1 (case $p = 2$), corresponding to the Stokes case.

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References

- [1] M. Anguiano, Darcy's laws for non-stationary viscous fluid flow in a thin porous medium, *Math. Meth. Appl. Sci.*, 40, No. 8 (2017) 2878-2895.
- [2] M. Anguiano, On the non-stationary non-Newtonian flow through a thin porous medium, *ZAMM-Z. Angew. Math. Mech.* 97, No. 8 (2017) 895-915.
- [3] M. Anguiano, R. Bunoiu, On the Flow of a Viscoplastic Fluid in a Thin Periodic Domain. In: Constanda C., Harris P. (eds) *Integral Methods in Science and Engineering*. Birkhäuser, Cham, 2019.
- [4] M. Anguiano, F.J. Suárez-Grau, Homogenization of an incompressible non-Newtonian flow through a thin porous medium, *Z. Angew. Math. Phys.* (2017) 68:45.

- [5] M. Anguiano, F.J. Suárez-Grau, The transition between the Navier-Stokes equations to the Darcy equation in a thin porous medium, *Mediterr. J. Math.* (2018) 15:45.
- [6] L. Berlyand, A.G.Kolpakov, A. Novikov, Introduction to the Network Approximation Method for Materials Modeling, Cambridge University Press, 2013.
- [7] L. Berlyand, E. Khruslov, Homogenized Non-Newtonian Viscoelastic Rheology of a Suspension of Interacting Particles in a Viscous Newtonian Fluid, *SIAM Journal of Applied Mathematics*, 64(3), (2004) 1002-1034.
- [8] N. Bernabeu, P. Saramito, A. Harris, Laminar shallow viscoplastic fluid flowing through an array of vertical obstacles, *J. Non-Newtonian Fluid Mech.*, 257 (2018) 59-70.
- [9] D. Bresch, E. D. Fernandez-Nieto, I. Ionescu, P. Vigneaux, Augmented Lagrangian Method and Compressible Visco-Plastic Flows : Applications to Shallow Dense Avalanches, *New Directions in Mathematical Fluid Mechanics* (2010) 57-89.
- [10] A. Bourgeat, A. Mikelić, A note on homogenization of Bingham flow through a porous medium, *J. Math. Pures Appl.*, 72 (1993) 405-414.
- [11] R. Bunoiu, G. Cardone, Bingham Flow in Porous Media with Obstacles of Different Size, *Mathematical Methods in the Applied Sciences*, Vol. 40, No. 12, 4514-4528, (2017).
- [12] R. Bunoiu, G. Cardone, C. Perugia, Unfolding Method for the Homogenization of Bingham flow, *Modelling and Simulation in Fluid Dynamics in Porous Media*, Series: Springer Proceedings in Mathematics & Statistics, Vol. 28 (2013) 109-123.
- [13] R. Bunoiu, A. Gaudiello, A. Leopardi, Asymptotic Analysis of a Bingham Fluid in a Thin T-like Shaped Structure, *Journal de Mathématiques Pures et Appliquées*, 123 (2019) 148-166.
- [14] R. Bunoiu, S. Kesavan, Fluide de Bingham dans une couche mince, *Annals of the University of Craiova, Math. Comp. Sci. series 30*, (2003) 1-9.
- [15] R. Bunoiu, S. Kesavan, Asymptotic behaviour of a Bingham fluid in thin layers, *Journal of Mathematical Analysis and Applications*, 293, No. 2 (2004) 405-418.
- [16] D. Cioranescu, A. Damlamian, G. Griso, The periodic Unfolding Method in Homogenization, *SIAM J. Math. Anal.* 40, n. 4, (2008) 1585-1620.
- [17] D. Cioranescu, A. Damlamian, G. Griso, *The Periodic Unfolding Method: Theory and Applications to Partial Differential Problems. Series in Contemporary Mathematics 03*, Springer, 2019.
- [18] D. Cioranescu, V. Girault, K. R. Rajagopal, *Mechanics and mathematics of fluids of the differential type*, vol. 35 of *Advances in Mechanics and Mathematics*, Springer, 2016.
- [19] G. Duvaut, J.L. Lions, *Les inéquations en mécanique et en physique*, Dunod, Paris, 1972.
- [20] J. Fabricius, J.G. I. Hellström, T.S. Lundström, E. Miroshnikova and P. Wall, Darcy's Law for Flow in a Periodic Thin Porous Medium Confined Between Two Parallel Plates, *Transp. Porous Med.*, 115 (2016) 473-493.
- [21] E.D. Fernández-Nieto, P. Noble, J.P. Vila, Shallow water equations for power law and Bingham fluids, *Science China Mathematics* 55 (2) (2012) 277-283.
- [22] V. Girault, P.A. Raviart, *Finite Element Methods for Navier-Stokes Equations, Theory and Algorithms*, Springer Series in Computational Mathematics, 5, Springer-Verlag, 1986.
- [23] G. Griso, Asymptotic behavior of a crane. *C.R.Acad.Sci. Paris, Ser. I*, 338, No. 3, 261-266 (2004).
- [24] G. Griso, L. Merzougui, Junctions between two plates and a family of beams. *Mathematical Methods in the Applied Sciences*, Wiley, (2018), 41 (1).

- [25] G. Griso, A. Migunova, J. Orlik, Asymptotic analysis for domains separated by a thin layer made of periodic vertical beams. *Journal of Elasticity*, 128, 291-331 (2017)
- [26] I. R. Ionescu, Onset and dynamic shallow flow of a viscoplastic fluid on a plane slope, *J. Non-Newtonian Fluid Mech.*, 165 (2010) 1328-1341 .
- [27] I. R. Ionescu, Augmented Lagrangian for shallow viscoplastic flow with topography, *Journal of Computational Physics*, 242 (2013) 544-560.
- [28] I. R. Ionescu, Viscoplastic shallow flow equations with topography, *J. Non-Newtonian Fluid Mech.*, 193 (2013) 116-128.
- [29] J.L. Lions, E. Sánchez-Palencia, Écoulement d'un fluide viscoplastique de Bingham dans un milieu poreux, *J. Math. Pures Appl.*, 60 (1981) 341-360.
- [30] P. Lipman, J. Lockwood, R. Okamura, D. Swanson, K. Yamashita, Ground deformation associated with the 1975 magnitude-7.2 earthquake and resulting changes in activity of kilauea volcano, 1985. Hawaii, Technical report, US Government Printing Office.
- [31] K.F. Liu, C.C. Mei, Approximate equations for the slow spreading of a thin sheet of Bingham plastic fluid, *Phys. Fluids*, A 2, 30-36 (1990).
- [32] S. E. Pastukhova, Asymptotic analysis in elasticity problems on thin periodic structures, *Networks and Heterogeneous Media*, 4 (3), (2009) 577-604.
- [33] K. Vasilic, B. Meng, H.C. Kühne, N. Roussel, Flow of fresh concrete through steel bars: a porous medium analogy, *Cement and Concrete Research* 41 (5) (2011) 496-503.
- [34] Y. Zhengan, Z. Hongxing, Homogenization of a stationary Navier-Stokes flow in porous medium with thin film, *Acta Mathematica Scientia*, 28B(4) (2008) 963-974.

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