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## Distances and absolute magnitudes from trigonometric parallaxes

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**Abstract.** In astrophysical applications, derived quantities like distances, absolute magnitudes and velocities are used instead of the observed quantities, such as parallaxes and proper motions. As the observed values are affected by random errors and selection effects, the estimates of the astrophysical quantities can be biased if a correct statistical treatment is not used. This paper presents and discusses different approaches to this problem.

We first review the current knowledge of Hipparcos systematic and random errors, in particular small-scale correlations. Then, assuming Gaussian parallax errors and using examples from the recent Hipparcos literature, we show how random errors may be misinterpreted as systematic errors, or transformed into systematic errors.

Finally we summarise how to get unbiased estimates of absolute magnitudes and distances, using either Bayesian or non-parametrical methods. These methods may be applied to get either mean quantities or individual estimates. In particular, we underline the notion of astrometry-based luminosity, which avoids the truncation biases and allows a full use of Hipparcos samples.

### 1. Introduction

Many papers have been devoted along the years to the various biases that can arise in the determination of stellar luminosities from trigonometric parallaxes. The advent of the Hipparcos Catalogue, with its unprecedented accuracy and homogeneous data, could have provided the opportunity to efficiently take these biases into account.

It seems, on the contrary, that, in the majority of recent papers, the sample selections have been mostly based on the parallax relative precision (based on  $\frac{\sigma}{\pi_H}$ , where  $\pi_H$  denotes the Hipparcos parallax and  $\sigma$  its formal precision) while it is well known that sample truncations based on the parallax relative error

lead to biased estimates of quantities derived from the parallax. Furthermore, the various adopted limits on  $\frac{\sigma}{\pi_H}$  are merely a balance between the expected precision on the resulting absolute magnitude and the size of the sample and thus are not based on any statistical criteria. Some illustrative examples are shown in Section 3.3.

The effects of random errors will be thoroughly discussed in the following sections, but it is interesting to summarise here what are the implicit consequences for the resulting sample of a truncation on the “observed” relative error  $\frac{\sigma}{\pi_H}$ :

- the truncation on  $\frac{1}{\pi_H}$  should produce an approximate volume-limited sample, but the error on  $\frac{1}{\pi_H}$  is correlated with the error on the absolute magnitude, implying a bias on this quantity;
- for a given  $\frac{1}{\pi_H}$ , the precision  $\sigma$  is mainly due to photon noise, so that brighter stars will be preferentially selected;
- for a given apparent magnitude,  $\sigma$  also depends on the ecliptic latitude (due to the Hipparcos scanning law) thus adding a spatial selection;
- the values of  $\sigma$  used are not the “real” values of the precision but its estimates. Thus,  $\frac{\sigma}{\pi_H}$  is the combination of two random variables, making the truncation more statistically complex;
- it should also be taken into account that the initial sample (before truncation) represents the content of the Hipparcos Catalogue. Apart from the “Survey”, which is rather well defined, the other selections from which the Catalogue was built are not always clear (there is, e.g., a kinematical bias in favour of nearby stars).

The final effect is that there is in fact *no* knowledge of the representativeness of the sample with respect to the parent population. Any statistics computed from this sample using parallaxes will probably be a biased estimate of the quantity which one would like to obtain for the parent population. Furthermore, if a generic (in the sense of not specifically adapted to the characteristics of the sample) *a posteriori* bias correction is applied, the accuracy of the result is hardly predictable. This is e.g. the case for the Lutz-Kelker (1973) correction, which assumes a uniform stellar density, whereas this assumption may not be realistic. Even when it is adequate, the confidence interval of the correction may be very large (Koen 1992), so that the precision on absolute magnitude will be rather poor.

In general, truncating on parallax relative error in the hope of benefiting from smaller random errors finally gives larger systematic errors. Moreover, the rejection of stars with high relative errors disregards a large amount of data, from which the random errors could have been reduced. Anticipating the conclusions of this paper, it must be noted that no selection on the observed parallaxes should be made. It should also be remembered that the observing list of the brightest stars in the Hipparcos Catalogue (Survey) was prepared on purpose, in order to provide clearly defined samples. When applicable, the selection on

apparent magnitude may then be taken into account in the estimation procedure. The effect of this selection (Malmquist bias) is discussed in Section 3.

## 2. The error law of Hipparcos parallaxes

Since the effect of random parallax errors on derived quantities will be discussed in Section 4, we review in this section the general properties of Hipparcos errors.

### 2.1. Gaussian errors

It has been shown in various papers (e.g. Arenou et al. 1995, 1997) that in general the random errors of the Hipparcos parallaxes may be considered Gaussian. This may be seen for instance when these parallaxes are compared to ground-based values of similar precision, or to distant stars using photometric estimates, from the normalised differences (parallax differences divided by the square root of the quadratic sum of the formal errors). This nice property of the parallax errors may then be fully exploited in parametrical approaches which make use of the conditional law of the observed parallaxes given the true parallax.

The particular case of systematic errors at small angular scale will be discussed in Section 2.3.; for an all-sky sample, one may safely consider that the global systematic error is small ( $< 0.1$  mas), that the formal errors are good estimates of the random error dispersions, and that the random errors are uncorrelated from star to star.

### 2.2. Non-Gaussian errors

Due to their Gaussian behaviour, random errors in the Hipparcos Catalogue are of course expected to produce a number of stars whose astrometric parameters deviate significantly from the  $1\sigma$  error level so that, of course, some hundreds of stars are expected to have an observed parallax which deviates, say, by 3 mas from the true parallax value. This is a logical consequence in a large catalogue like Hipparcos.

In a few cases, however, it may happen that the error on the Hipparcos data is much higher than expected. Although these are probably rare cases, they have been mentioned for the sake of completeness in the Hipparcos documentation and illustrated here.

Apart from the Double and Multiple Star Annex (DMSA, see ESA 1997), most of the Hipparcos Catalogue is constituted by stars assumed to be single. One obvious source of outliers may thus be undetected short period binarity, since in this case the astrometric path of the star will not exactly follow the assumed single star model (5 parameters: position, parallax, linear proper motion).

Two extreme cases of astrometric binaries are discussed below, which may have been biased respectively in parallax and proper motion. It must be stressed that these examples are statistically rare and chosen for the purpose of illustration, and that the duplicity in these few cases had been detected by Hipparcos and flagged in the Catalogue.

The first example concerns HIP 21433, one of the 1561 Hipparcos stochastic solutions (DMSA/X), where an excess scatter of the measurements may be

interpreted as the signature of an unknown orbital motion. Indeed, this star is a spectroscopic binary. The interesting fact is that the period is 330 days, i.e. close to one year, so that there may have been a confusion between the parallactic and orbital motions. Adopting the four known orbital elements  $(P, T, e, \omega_1)$  from Tokovinin et al. (1994), the intermediate astrometric data (ESA, 1997, Vol I, Sect. 2.8) have been re-reduced taking into account the 5 astrometric parameters and the 7 orbital parameters, and the new parallax found is  $30.36 \pm 0.87$  mas, instead of  $34.23 \pm 1.45$  mas as in the published stochastic solution. The parallax from the Hipparcos solution, which did not take into account the binarity, has thus possibly been biased, due to the  $\approx 1$  year orbital motion.

The second example concerns HIP 13081, one of the 2622 acceleration solutions (DMSA/G), where the motion has not been linear during the mission, interpreted as a binary of longer period. Taking into account the orbital motion, using data from Tokovinin (1992), a new solution has been computed. The inclination angle  $i$  is close to  $90^\circ$  so that the path on the sky is nearly linear. The proper motion of the barycentre is  $275 \pm 3$  mas/yr instead of the published solution,  $264 \pm 1$  mas/yr. Strictly speaking, this is not a bias, since Hipparcos measured the photocentre of the system, not the barycentre.

In both examples, compared to the “true” value, the published solution is significantly different from what could be expected in the case of Gaussian errors. Although these examples are unfavourable cases, it must be pointed out that they were detected during the Hipparcos data reduction. The same effects may also be present for some other stars where the binarity has not yet been detected, but this implies at the same time that the astrometric perturbation is smaller.

### 2.3. Small-scale systematic errors

The operation mode of the Hipparcos satellite implied that the stars within a given small field were frequently observed together on the same great circles. This introduces correlations between the astrometric parameters of stars within some square degrees but, due to the rather low sky density of Hipparcos, it is only a problem for open star clusters. This effect was studied before the satellite launch by Lindegren (1988) and confirmed using the final results by Lindegren et al. (1997) and Arenou (1997).

A special data reduction process had then to be used for cluster stars. This has been done by van Leeuwen (1997a,b) and Robichon et al. (1997), and for this purpose the angular correlations have been calibrated, as detailed in van Leeuwen & Evans (1998) and Robichon et al. (1999).

Although the correlation effect was known and taken into account, it was possibly not realized that, for a single realisation of a given cluster, this could mean a systematic error for the individual cluster members. It must however be remembered that the Hipparcos data were reduced by two different Consortia, and the systematic error is probably not the same for both, so that the merging of the two sets (Arenou, 1997) probably reduced the effect.

In order to illustrate this correlation, one may take the extreme example of NGC 6231, where all six Hipparcos stars have a negative parallax, whereas the photometric estimate is  $0.71 \pm 0.02$  mas (Dambis, 1998). A straight weighted average of individual parallaxes would give  $-0.71 \pm 0.39$  mas; even taking into

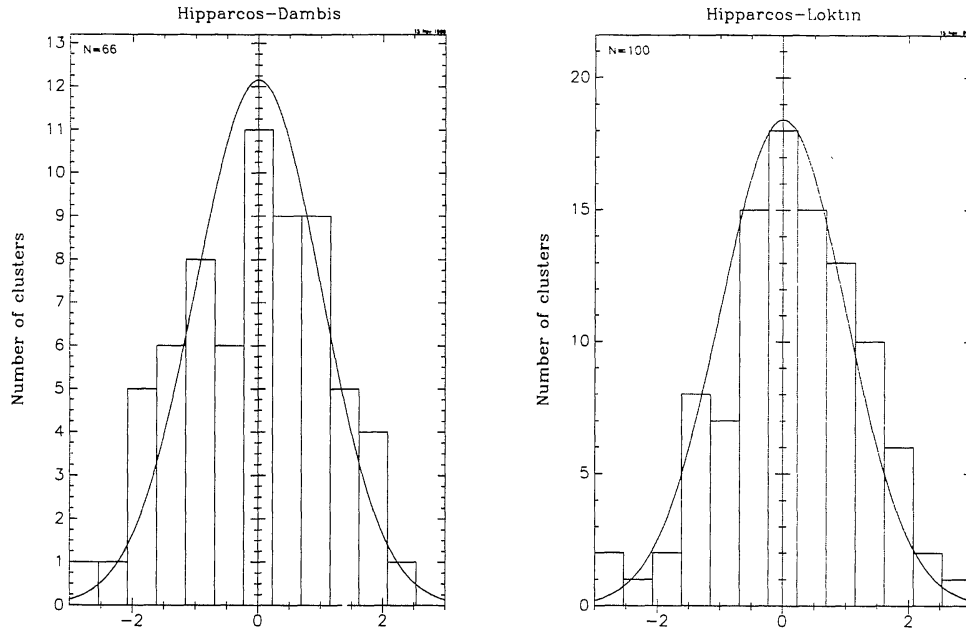


Figure 1. Normalised difference between mean Hipparcos parallaxes of distant clusters, and the parallaxes from Dambis (1998, left) or Loktin & Matkin (1994, right). A Gaussian (0,1) is superimposed.

account the correlations, the mean cluster parallax is  $-0.62 \pm 0.48$  mas, which is still significantly different from the photometric estimate.

Apart from this extreme example, the question is whether the correlation effect has correctly been accounted for in the estimation of the Hipparcos mean parallax of each cluster. Although this is not easy to prove for one given cluster, an indirect statistical evidence may be obtained using a sample of distant clusters. To carry out this test, clusters farther than 300 pc with at least 2 Hipparcos members were used: 66 clusters from Dambis (1998) Catalogue and 102 clusters from Loktin & Matkin (1994). Concerning the latter, a 5% parallax relative error has been adopted. The mean photometric parallax error of these samples is about 0.04 mas, so that the comparison with Hipparcos parallaxes shows mainly the Hipparcos errors.

The normalised differences between the mean Hipparcos parallaxes, taking into account the angular correlations, and the photometric parallaxes are shown Figure 1. It appears that the small-scale systematic errors are 0 on the average, and the unit-weight (dispersion of the normalised differences) about 1.15. If the Loktin & Matkin distance moduli are corrected, taking into account the new Hyades distance modulus (3.33 instead of 3.42), the zero-point, the unit weight (1.17) and the asymmetry are reduced. Since the cluster memberships have not been thoroughly investigated, the 15% underestimation of the formal error on Hipparcos mean parallaxes seems to be an upper limit.

Pinsonneault et al. (1998) suggested that a systematic error affects the mean Hipparcos parallax of the Pleiades, due to the correlations between the right ascension and parallax,  $\rho_{\alpha*\pi}$ . For each star of the distant clusters, the

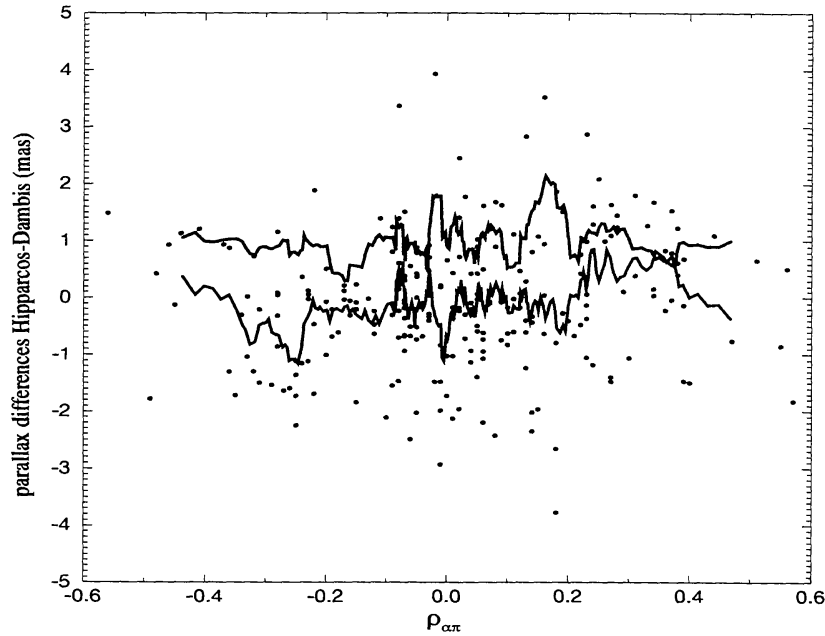


Figure 2. Errors on Hipparcos parallaxes for distant clusters vs the correlation coefficient between right ascension and parallax. The running average (around 0) and standard deviation (around 1) over 10 stars is superimposed.

difference between Hipparcos and Dambis parallaxes is plotted in Figure 2 as a function of  $\rho_{\alpha*\pi}$ . There are significant differences, due to some clusters (NGC 6231 has a  $\rho_{\alpha*\pi} \approx -0.25$ ), but not a linear trend.

From these tests, a 1 mas systematic error for the Pleiades seems unlikely, although one cannot preclude a systematic of a few tenth of mas for some clusters, which would increase the unit-weight as pointed out by Pinsonneault et al. (1998). We may however assume for the following discussion that there is no significant systematic error in the Hipparcos parallaxes, even at small-scale.

### 3. The effect of random parallax errors

The astrometric elements of a star are of not much interest in themselves for an astrophysical purpose. Instead, the quantities of interest are the distance, the absolute magnitude, the radius, the age, or the spatial velocity. Given an observed parallax and proper motion, with its associated errors, unbiased estimates of these quantities are not easy to obtain. For instance, it has been shown by Lutz and Kelker (1973) that a sample selection based on the observed parallax relative error would introduce a bias on the mean absolute magnitude. In fact Lutz and Kelker considered that the bias occurs for all parallaxes, but we will focus on sample selection only. This bias is due to:

- the non-linear relationship between absolute magnitude (or distance, etc.) and parallax,

- the truncation based on the observed parallax, the true parallax distribution not being uniform.

These two points are discussed below, and the influence of parallax errors is shown through the use of examples from the literature. It should be noted that what is true for the absolute magnitude is equally true for the other mentioned quantities. Although obvious, it is worth remarking that  $\frac{\sigma}{\pi_H} \propto \frac{1}{\pi_H}$  so that the “observed” relative error suffers a bias, high dispersion and skewness proportional to those of the “observed” distance. The so-called Lutz-Kelker bias occurs because there is a random error in the observed parallax which is thus present both in the “observed” relative error and in the “observed” absolute magnitude.

### 3.1. Bias from non-linearity

Starting from a symmetric error law for the parallaxes, the error law on derived quantities such as distance or absolute magnitude loses this property. Due to their non-linearity with respect to parallax, a bias is expected, and this is amplified by the fact that the corresponding estimates are not defined when the observed parallax is 0 or negative.

Given the true parallax  $\pi$ , and assuming a Gaussian law for the error on the observed parallax,  $\pi_H \sim \mathcal{N}(\pi, \sigma)$ , the expected bias of the “observed” distance  $r_H = \frac{1}{\pi_H}$  in absence of any truncation is

$$E[r_H|\pi] - \frac{1}{\pi} = \frac{1}{\pi} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \left( \frac{1}{1 + u\frac{\sigma}{\pi}} - 1 \right) e^{-\frac{u^2}{2}} du \quad (1)$$

and the bias for the “observed” absolute magnitude  $M_H = m + 5 \log(\pi_H) + 5 - A$  is

$$E[M_H|\pi] - M = \frac{5}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \log(1 + u\frac{\sigma}{\pi}) e^{-\frac{u^2}{2}} du \quad (2)$$

Apart from the fact that these integrals are not always defined, in both cases a bias will be present when  $\frac{\sigma}{\pi}$  is not negligible. Assuming no truncation on negative or null parallaxes, and for small relative errors, the biases may be approximated by respectively  $B(r_H) \approx \frac{1}{\pi} (\frac{\sigma}{\pi})^2$  and  $B(M_H) \approx -1.09 (\frac{\sigma}{\pi})^2$ , negligible for relative errors smaller than, say, 10%. This bias is due to the asymmetry of the error distribution for  $r_H$  and  $M_H$ , and is what would still be present if an average of these quantities is computed; other statistics (based on the mode or the median) would possibly give a result closer to the true value.

For higher relative errors, the biases and variances are depicted as a function of the true relative error in Brown et al. (1997), Figures 1 and 2 for the distance and the absolute magnitude respectively.

### 3.2. Bias from truncation on observed data

Whereas the bias due to the non-linearity would systematically happen (but with a limited effect for small relative errors), the bias due to the truncation on the “observed” relative error, which is the major effect, could be avoided... if no truncation was done.



An important part of studies in the recent literature based on Hipparcos data have used a truncation procedure, usually based on the relative parallax error and sometimes rejecting only the negative parallaxes. In the hope of selecting only the most precise absolute magnitudes, not only their mean is biased, due to the Lutz-Kelker effect, but moreover the obtained precision on this mean is worse.

A simple —though extreme— simulation may help to understand this fact. We have randomly drawn samples of 1000 stars limited to magnitude 12 of constant absolute magnitude=1<sup>m</sup> with a uniform spatial distribution<sup>1</sup> and a parallax error depending on magnitude. These large samples contain on the average only 10 stars with “observed” relative error better than 30%, and the best weighted mean of the corresponding “observed” absolute magnitudes is  $1.28 \pm 0.28$ , whereas using all stars and the estimate discussed in Section 5.2, the mean absolute magnitude found is  $1.00 \pm 0.08$ . If an unweighted mean had instead been used for the truncated sample, the bias would have reached 0.8 magnitudes.

In this example, the truncation on “observed” relative error gives a 30% systematic error, of the same amount as the mean error, which is itself 3 times larger than what would be obtained without truncation. The truncation thus appears as an efficient way to obtain both biased and unprecise results !

Although the Lutz-Kelker effect is widely known, there seems however to be some confusion about its origin. The confusion was introduced however by Lutz & Kelker themselves who underlined that the systematic “*error is present for ALL stars*”. Had this been true, then every measured quantity whose distribution is not uniform would suffer from a systematic error. This would in particular be the case e.g. for extinction along one line of sight, etc.

However, an observed parallax is individually an unbiased estimate of the true parallax,  $E[\pi_H|\pi] = \pi$ , as indicated Section 2., so the average value for a *random* sample of observed parallaxes will be the same as the average value of the underlying true parallaxes, with no bias

$$\begin{aligned} E[\pi_H] &= \int_{-\infty}^{+\infty} \pi_H f(\pi_H) d\pi_H = \int_{-\infty}^{+\infty} \pi_H \int_0^{+\infty} f(\pi_H|\pi) f(\pi) d\pi d\pi_H \\ &= \int_0^{+\infty} \int_{-\infty}^{+\infty} \pi_H f(\pi_H|\pi) d\pi_H f(\pi) d\pi = \int_0^{+\infty} \pi f(\pi) d\pi = E[\pi] \end{aligned}$$

On the contrary, if a selection is based on the observed parallax distribution (e.g. an integration of  $\pi_H$  from some limit  $\pi_-$  in the previous equations), the mean value will be biased. An example of this problem can be found in Oud-maijer et al. (1998). Although the authors indicate that “*this statistical effect causes measured parallaxes to be too large*”, the bias value will in fact depend on the parallax distribution: for a classical magnitude-limited sample, the measured parallax will be either too large or too small depending on whether the truncation is done on one side or another of a mode of the parallax distribution, as may be deduced from  $E[\pi|\pi_H]$  in Equation 10.

<sup>1</sup>Notice that in this example the samples are not affected by Malmquist bias, even if they are limited in apparent magnitude, because no intrinsic dispersion was introduced on the absolute magnitudes —see Sec. 4.—; thus, any bias will come from non-linearity and parallax truncation.

In fact, for one given star, little can be said when no other information than the observed parallax is available. Let us consider for instance a star with observed parallax of, say, 3 mas, which belongs to two different samples (e.g. with different limiting magnitude), the modes of the distributions of two samples being respectively at e.g. 2 mas and 4 mas: is the observed parallax expected to be too small or too large?

It must also be pointed out that  $E[\pi|\pi_H]$  is *not* the true parallax, but an estimate with also a dispersion. Moreover, this estimate may itself be biased if the *a priori* hypotheses are inappropriate. For example, if an homogeneous spatial distribution ( $f(\pi) \propto \pi^{-4}$ ) is assumed for a large sample of nearby stars, among which stars of a given open cluster are present, then the average cluster parallax computed using the posterior expectation will be biased towards a larger distance.

### 3.3. Examples from the literature

Since the publication of the Hipparcos Catalogue, there have been numerous papers inferring, from samples of Hipparcos stars, the properties of some populations, or comparing the new data with external data. In some cases, the effect of random errors may be misleading, and this is mainly due to the existing correlations between “observed” parallax relative errors, “observed” absolute magnitudes and “observed” distances.

A first example is taken from Tsujimoto et al. (1997), where the absolute magnitudes of RR-Lyræ are calibrated. Although the authors follow a rigorous statistical approach, their Fig. 2 may be misunderstood by the unaware reader. In this Figure, the “observed” absolute magnitude seems to go fainter with increasing (true) distance, the stars with “observed” parallax relative errors greater than 100% being systematically brighter.

What could be interpreted as a systematic error in the parallax is *exactly* what is expected from parallaxes with random errors —and without systematic errors. A simulation of 174 distant stars, assuming a constant absolute magnitude =1, is shown in Fig. 3, excluding obviously those with a negative parallax. The magnitude dispersion increases with distance (due to the increase of true parallax relative errors); the error bars become more and more asymmetrical, shifting some “observed” absolute magnitude towards the brightest end; and a positive random parallax error implies both a fainter “observed” absolute magnitude and a smaller “observed” parallax relative error, producing the correlation between these two data.

A second example is taken from Oudmajer et al. (1998), where the authors discuss the Lutz-Kelker effect and apply the correction to a sample of Cepheids. They first compare the “observed” absolute magnitudes computed using ground-based parallaxes (with a large random error) to those computed with precise Hipparcos parallaxes as a function of the ground-based parallax (their Figure 1, lower panel). The observed effect at small parallaxes is interpreted as a “*completeness effect in the data*”, whereas it is only due to correlations between absolute magnitude errors and parallax errors.

Using a subsample of 26 Cepheids, the difference between “observed” and true absolute magnitudes is plotted in their Figure 4 as a function of “observed” parallax relative error. As expected, the correlation effect is present, and this is

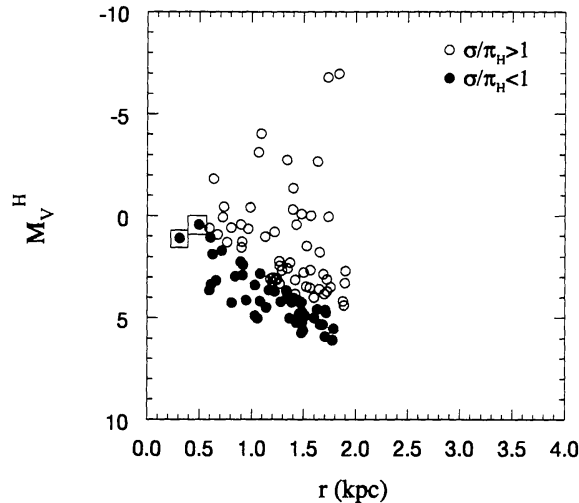


Figure 3. Simulation of distant stars, showing the correlation between “observed” absolute magnitude and “observed” parallax relative error. Only the two points indicated have a *true* relative error smaller than 1. For comparison, see Figure 2 in Tsujimoto et al. (1997)

not due to missing faint stars: a volume-limited simulation will exactly reproduce the correlation effect. Assuming that the initial sample had no selection bias, the rejection of the 3 stars with the higher “observed” relative error creates a bias which is then corrected.

Their result in itself will not be further discussed here. As the authors quote, Koen (1992) showed that for  $\frac{\sigma}{\pi_H} = 0.175$ , the 90% confidence interval of the Lutz-Kelkér correction could span over more than 1.77 magnitudes! Since almost all of the stars used by Oudmaijer et al. have a high parallax relative error one may however wonder how their result may be as precise as 0.02 mag.

The two above examples illustrate the fact that the comparisons should always be done in the plane of the measured quantities (the parallaxes), where the errors may safely be assumed symmetrical, and not in the plane of the derived quantities, where the effect of the random errors is not always clear.

In the following example, however, although the comparison is done in the parallax plane, the effect of asymmetrical errors may be significant, still due to correlations. Figure 2 from Jahreiß et al. (1997) shows the differences between the Hipparcos parallaxes and those deduced from photometric CLLA parallaxes (Carney et al. 1994) versus the CLLA parallaxes. If there is a systematic shift of the photometric absolute magnitudes, then it should be seen as a slope in this graph, and the photometric parallaxes should be corrected by a factor  $(1+\text{slope})$ . This method could also provide an estimate of the Hipparcos zero-point.

Although there is probably such an absolute magnitude zero-point error in that case, it should be pointed out that there are random errors (assumed symmetrical) in the calibrated photometric absolute magnitude, so both the resulting asymmetrical random error of the photometric parallaxes and the correlation between both axes may produce a similar effect (Lindgren 1992).

The way random errors may mimic a systematic error is shown in Fig. 4, where 275 stars have been simulated, assuming a constant density, a linear rela-

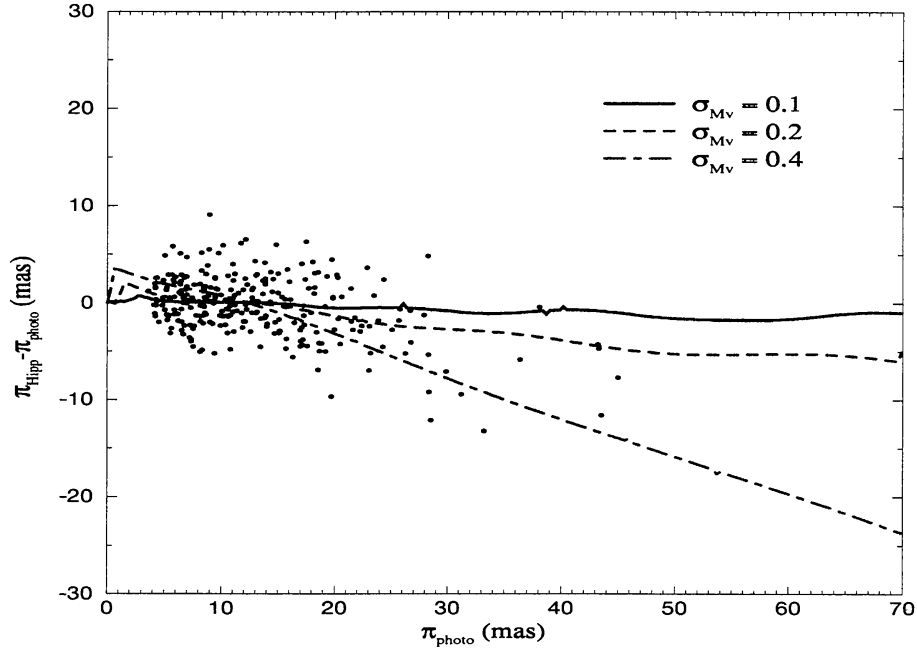


Figure 4. Simulation of differences between trigonometric and photometric parallaxes, with the prediction as a function of the photometric parallax error. For comparison, see Figure 2 in Jahreiß et al. (1997)

tion between colour and absolute magnitude, a 0.4 mag random error on absolute magnitude for the photometric estimate, and an observed parallax computed from the true absolute magnitude.

Denoting  $\pi_P$  the photometric parallax, the theoretical effect may be computed under the assumption of unbiased astrometric and photometric parallaxes:

$$E[\pi_H - \pi_P | \pi_P] = \frac{\int_0^{+\infty} \pi f(\pi_P | \pi) f(\pi) d\pi}{\int_0^{+\infty} f(\pi_P | \pi) f(\pi) d\pi} - \pi_P \quad (3)$$

Assuming a Gaussian law for the distribution of the error on the photometric absolute magnitude, with associated variance  $\sigma_M^2$

$$f(\pi_P | \pi) \propto e^{-\frac{1}{2} \frac{25(\log \pi_P - \log \pi)^2}{\sigma_M^2}}$$

and assuming a magnitude-limited *a priori* distribution for the true parallaxes, the shape of what may be expected from the random errors only is shown in Fig. 4 for different values of the random dispersion of the photometric parallaxes.

In summary, if the random errors on the photometric absolute magnitude are not properly taken into account in the estimation procedure, one could wrongly deduce from such a graph both a systematic error in the absolute magnitude and a systematic error in the trigonometric parallaxes. One can not however infer from this statement that there is no systematic error in the CLLA absolute magnitudes.

#### 4. Malmquist bias

Any finite sample of stars is, by definition, limited in apparent magnitude. In some cases this limit can be ignored if there is another more constraining truncation, like a limit in  $\frac{\sigma}{\pi_H}$ . However, if one wants to avoid as much as possible to introduce any censorship when constructing a sample of stars, one would be left with at least an apparent magnitude limit. It is thus important to understand the effects that such a truncation can have on the estimation of astrophysical quantities.

The simplest case of an apparent magnitude truncation is the case of a sample with a clean apparent magnitude limit. This case was first studied by Malmquist (1936) under some restrictive hypothesis:

- A Gaussian distribution of the individual absolute magnitudes:  $M \sim \mathcal{N}(M_0, \sigma_M)$
- A uniform spatial distribution ( $f(r) \propto r^2$ ).

Under these two hypotheses the joint distribution of absolute magnitudes and distances of the base population has the shape depicted in Fig. 5. However, when the apparent magnitude limit  $m \leq m_{lim}$  is introduced this joint distribution is drastically changed, as depicted in Fig. 6.

While the mean absolute magnitude of the base population is  $M_0$ , the mean absolute magnitude of the truncated sample  $\langle M \rangle$  differs from this value: it is biased. Thus, if one uses such a sample to estimate the absolute magnitude of the base population, even if the use of the trigonometric parallaxes is correct the value obtained will be biased.

This bias in the mean absolute magnitude of a sample due to apparent magnitude truncation is known as the Malmquist bias. Malmquist (1936) calculated its value under the above cited hypothesis:

$$\langle M \rangle \simeq M_0 - 1.38 \sigma_M^2 \quad (4)$$

There is, however, some confusion in the literature when using this correction. As pointed above, the Malmquist correction is valid under the two given hypotheses. If one of the two does not hold, the value of the Malmquist bias may differ from Eq. 4. For instance, in the (rather common) case of an exponential disk spatial distribution the value of the Malmquist bias depends on  $(\sigma_M, m_{lim} - M_0, Z_0)$ , where  $Z_0$  is the scale height of the exponential disk (Luri 1993). An example is given in Fig. 7.

Thus, Malmquist correction should not be blindly applied when an apparent magnitude truncation is present. The correction may vary depending on the absolute magnitude distribution and the spatial distribution of the base population. Furthermore, if the apparent magnitude truncation is not clean-cut the effect will also be different. This is where the Hipparcos Survey may come in handy.

On the other hand, as can be seen in Fig. 5 and 6, the mean distance of the sample is also biased with respect to the base population. This can be important when studying the mean distance of a cluster, for instance.

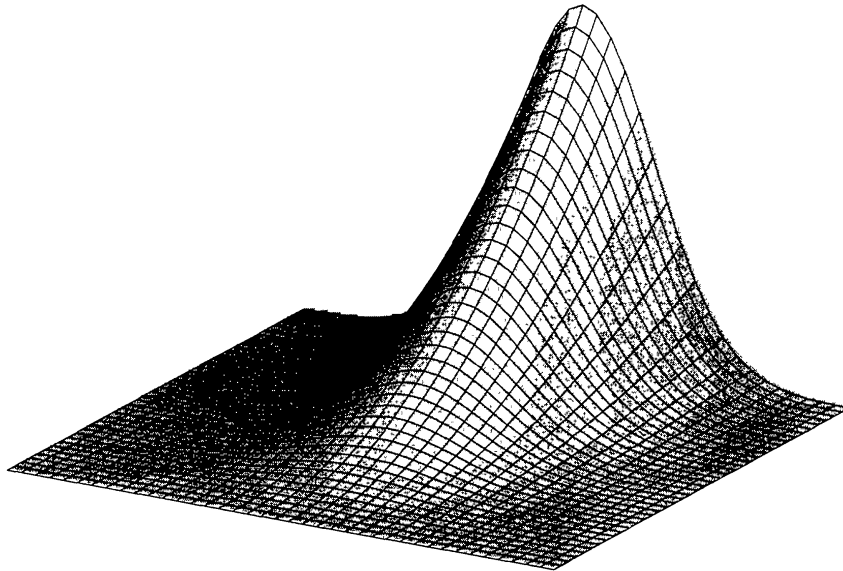


Figure 5. Joint  $(M, r)$  distribution for a base population with a Gaussian distribution in  $M$  and an homogeneous spatial distribution. The figure has been truncated in  $r$  for illustration purposes.

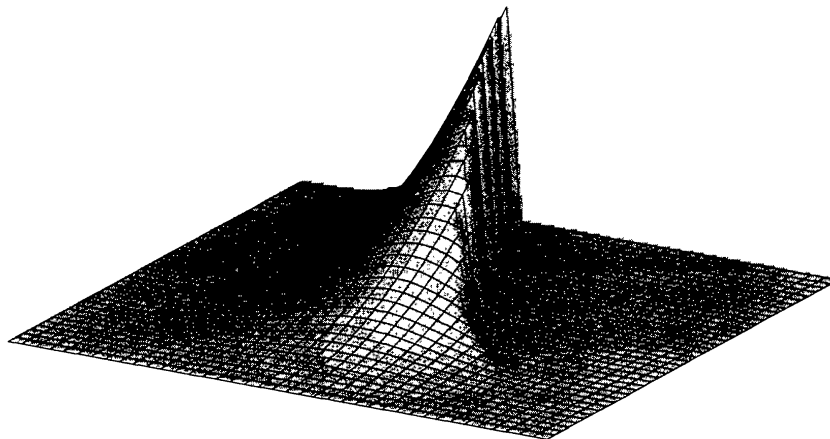


Figure 6. Joint  $(M, r)$  distribution for a sample with a Gaussian distribution in  $M$ , an homogeneous spatial distribution and a truncation in apparent magnitude. The figure has also been truncated in  $r$ .

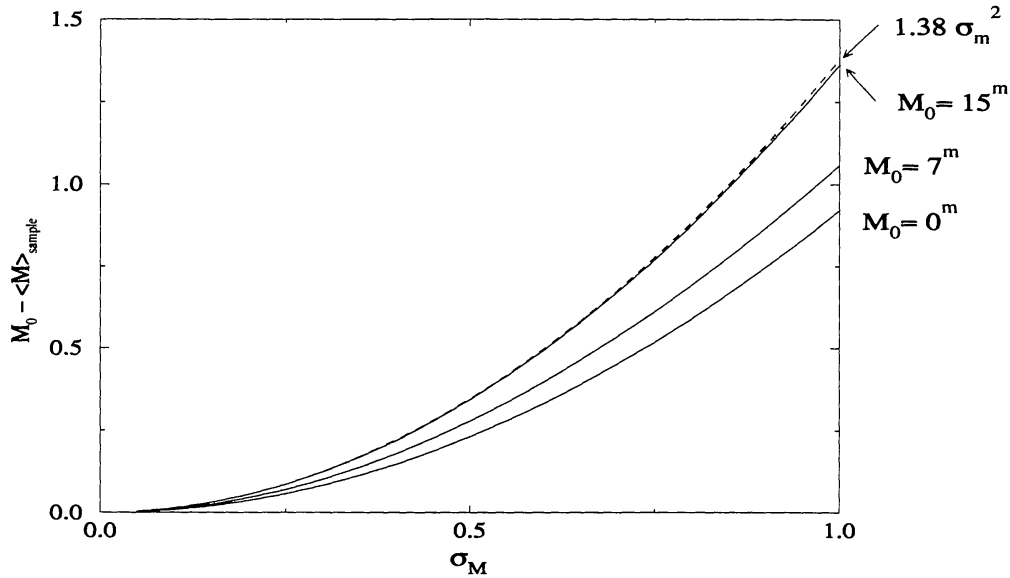


Figure 7. Malmquist bias in the case of a Gaussian distribution of absolute magnitudes, an exponential disk with  $Z_0 = 200$  pc for the spatial distribution and an apparent magnitude limit  $m_{lim} = 15^m$ . The classical Malmquist correction (dashed line) is given for comparison

Finally, a further warning. All the discussion in this section has been centred on the case of a sample truncated *only* in apparent magnitude. In the case of combined truncations the joint effect should be analysed and taken into account. For instance, as pointed out above, a stringent truncation in  $\frac{\sigma}{\pi_H}$  (e.g. 10%) may eliminate the effects of the apparent magnitude truncation, but that may not be the case for a less stringent truncation (e.g. 100%).

## 5. Which distance and absolute magnitude from parallax?

Since Dyson (1926), who corrected the observed parallax distribution in order to get the true absolute magnitude distribution, several methods have been used in order to get unbiased estimates of absolute magnitudes or distances. A first approach uses either a transformation of these quantities, or a correction of the biases. Another approach uses all stars in order to give a smaller bias. Finally a parametrical approach together with supplementary information fits a model to the observed quantities, taking explicitly into account the selection biases. The methods using a galaxy model and simulations (e.g. Bahcall & Soneira 1980, or Robin & Cr ez e 1986) pertain in some sense to this latter approach but will not be discussed here.

### 5.1. Transformation of the distance error law

Recently, Smith & Eichhorn (1996) have tackled the problem of distances derived from trigonometric parallaxes. Assuming Gaussian errors for the parallaxes,

they demonstrate the presence of bias on the “observed” distance and the fact that its variance can be infinite. In this case it would be useless to do a bias correction. Moreover the bias depends on the true parallax relative error, which is unknown. They propose two different methods, using either a transformation based on the observed parallax and its formal error, rendering a positive parallax, or a weighting of these parallaxes, eliminating the zero parallaxes. Each method has advantages and disadvantages depending on whether the bias or the variance of the resulting estimate is considered.

The other problem of the “observed” distance being its asymmetrical error law, Kovalevsky (1998) has proposed a transformation which would give a gaussianized distance error law for small (true) parallax relative errors.

It is however important to keep in mind the right use of these corrected distances. For instance, let us assume that we have to compare the distances deduced from Hipparcos parallaxes to the distances deduced from ground-based parallaxes, in order to test if there is a systematic effect in one of the data sets. Whereas the correct comparison would be in the plane of parallaxes, one perverse way to do it would be to compute for the two sets the “observed” distance, then to apply one of the above corrections, and finally to obtain a comparison of distances where biases are unclear and where the high variance may prevent any safe conclusion. . .

## 5.2. Asymptotically unbiased estimates

When one needs to obtain a mean parameter on a sample, such as a mean distance, mean absolute magnitude, etc, all parallaxes may in fact be used, instead of computing biased estimates for each star.

Concerning distance estimation, a simple example is the mean distance of a cluster, neglecting the cluster depth and assuming (which is not the case for Hipparcos) that no correlation exists between individual parallaxes. Two possible estimates,  $\langle \frac{1}{\pi_i} \rangle$  and  $\frac{1}{\langle \pi_i \rangle}$  would look at first sight equivalent. From Equation 1, however, the best of these two distance estimates is obvious: in the first case, the bias will still be present in the average since it occurs for each inverse of parallax, although it would be difficult to know its value since it depends on the true parallax relative error. Whereas, in the second case, the precision of the mean parallax will be  $\frac{\sigma}{\sqrt{n}}$ , so that the bias on the mean will be a factor  $\approx n$  smaller. Asymptotically, the second estimate is thus unbiased and should be preferred over the first, since its variance is also smaller. Although a bias will remain, it is in general very small compared to all the other uncertainties: typically, a cluster at 500 pc with only 9 Hipparcos stars will have a distance bias smaller than 3%, whereas an average of “observed” distances could give a bias larger than 30%.

Concerning the mean absolute magnitude of a star sample, asymptotically unbiased estimates are also used at least since Roman (1952), and detailed in Turon & Cr ez e (1977). This method has recently been used by Feast & Catchpole (1997) or van Leeuwen & Evans (1998) using Hipparcos intermediate astrometric data. The method is summarised at the end of this section.

However, this method concerns the mean absolute magnitude, not individual absolute magnitudes. The question is thus how to handle some individual



stars with poor parallax relative precision. In general, these absolute magnitudes are used in an H-R diagram, e.g. for age determination or luminosity calibrations.

Instead of focusing on the absolute magnitude  $M_V$ , let us consider the quantity

$$a_V = 10^{0.2M_V} = \pi 10^{\frac{m_V+5}{5}} \quad (5)$$

where the apparent magnitude  $m_V$  has been corrected for extinction and the parallax is in arcsec (or  $a_V = \pi 10^{0.2m_V-2}$  with  $\pi$  in mas). Missing a denomination for  $a_V$ , we will refer in what follows to ABL (Astrometry-Based Luminosity). The ABL, equal to the the inverse of the square root of a flux, is much more easy to handle than the absolute magnitude when dealing with stars with a high parallax relative errors or even negative parallaxes (i.e. when the dispersion due to parallax random errors is much larger than the intrinsic dispersion of absolute magnitudes).

In a classical H-R diagram, the absolute magnitude is plotted versus colour; in what we call an “astrometric” H-R diagram, the ABL is plotted versus colour. For illustration purposes, a sample of 1000 stars of age 10 Gy, with  $[\text{Fe}/\text{H}] = -1.4$  and an 0.5 mag dispersion in absolute magnitude, has been simulated. No variations in metallicity or random errors in colours have been added.

The classical H-R diagram for all stars with a 30% truncation on parallax relative error is represented on the left of Fig. 8 (116 stars). The so-called Lutz-Kelker effect appears clearly, showing the trend to get stars below the reference line. Since for each star the parallax relative error is not very large, the error bar asymmetry is not well seen.

Using the ABL, the “astrometric” H-R diagram is represented on the right of Figure 8. For the sake of comparison, the same number of stars has been kept; this has been obtained by using  $\sigma_{a_V} < 3$ . In general, there is however no reason to reject the other stars, as their high number compensates the larger error bars.

Consider for instance a program computing the age and metallicity for a sample of stars through interpolations between isochrones in an H-R diagram: the truncation effect on the parallax relative error may possibly bias the result. On the contrary, we could get unbiased and more precise estimates making use of the “astrometric” H-R diagram. Another application concerns all the luminosity calibrations, the ABL being calibrated as a function of photometric indices.

The use of ABL instead of absolute magnitude has the following advantages:

- the error bars on  $a_V$  due to parallax errors are symmetrical
- there is no Lutz-Kelker bias
- all stars may be used, even those with negative parallaxes
- the higher number of stars allows a gain in precision for mean values

Coming back to the simple case where a mean absolute magnitude has to be computed from a sample of stars, and following Jung (1971) or Turon & Cr ez e

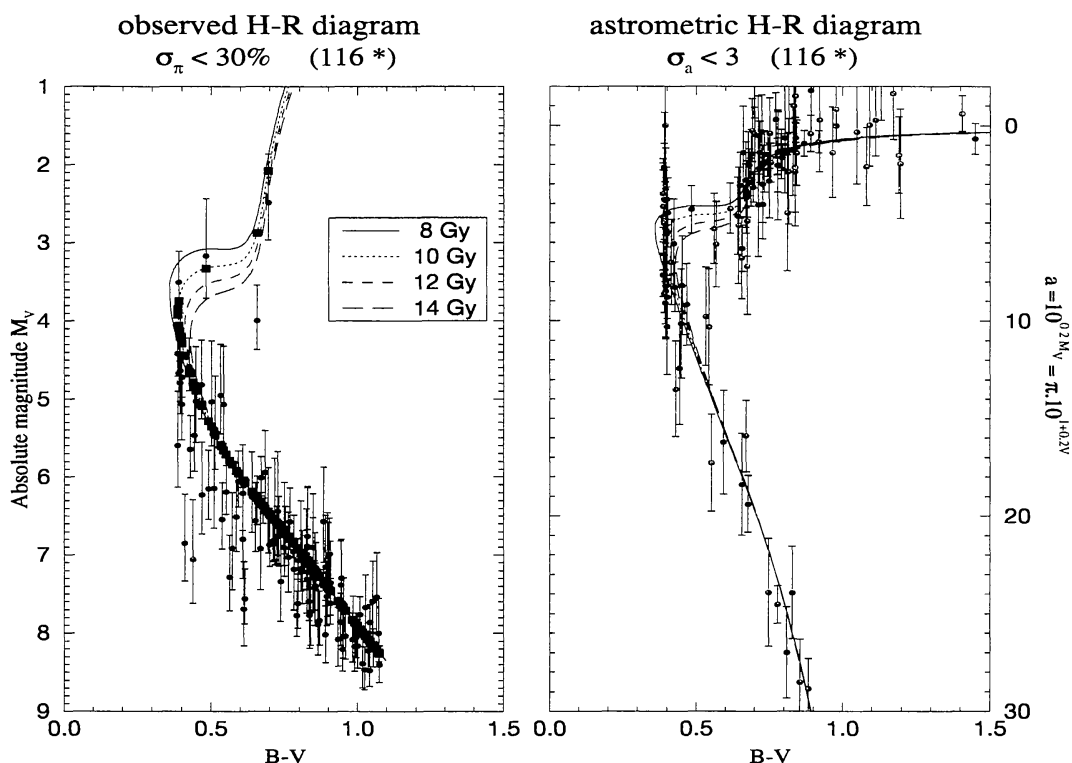


Figure 8. Simulation of a sample of 10 Gy stars with  $[\text{Fe}/\text{H}]=-1.4$  and a 0.5 mag dispersion in absolute magnitude. See text for legend.

(1977), the first step is to estimate the best weighted mean ABL for the sample

$$\langle a_V \rangle = \frac{\sum_i \frac{a_i}{\sigma_{a_i}^2}}{\sum_i \frac{1}{\sigma_{a_i}^2}}$$

or possibly a less precise but more robust estimate, then an asymptotically unbiased mean absolute magnitude is obtained with

$$\langle M_V \rangle = 5 \log \langle a_V \rangle$$

In the case where there is an intrinsic dispersion in absolute magnitude (assumed small), it has to be taken into account in the weights of  $\langle a_V \rangle$ . As indicated above, all the stars may (should) be used, although a selection on  $\sigma_{a_V}$  may be applied. However, since this is a selection on luminosity, a Malmquist-type bias should be accounted for. This may be also true for the whole sample. It must be pointed out that a symmetrical error in apparent magnitude will become asymmetrical in  $a_V$ , thus causing a bias. However, given the good photometric precision of Hipparcos, the bias coming from the errors in the apparent magnitudes is negligible and only the errors in the extinction correction may constitute a problem in some cases.

### 5.3. Parametrical approach

The approaches described above make use only of the parallax in order to derive the distance or absolute magnitude. Another approach makes use of all the available information: assuming some parametrical probability density functions (pdf), a maximum likelihood estimation allows one to find the optimal parameters corresponding to the studied sample. An early application of this method may be found in Jung (1971), and in a more modern way by Ratnatunga & Casertano (1991), Arenou et al. (1995) and Luri et al. (1996).

Given the observables  $O = (\pi_H, l, b, m_V, \mu_\alpha, \mu_\delta, V_R)$ , the parameters  $\Theta$  being the coefficients of absolute magnitude as a function of colour, the galactic scale length and scale height, the velocity ellipsoid, etc., are estimated by a maximum likelihood estimation of

$$h(O|\Theta) = \int_0^{+\infty} g(O|\pi, \Theta) p(\pi) d\pi \quad \text{where} \quad (6)$$

$$g(O|\pi, \Theta) = p_1(\pi_H|\pi, \Theta) p_2(m|\pi, \Theta) p_3(\mu_\alpha, \mu_\delta, V_R|\pi, \Theta) p_4(l, b|\pi, \Theta) \quad (7)$$

where each pdf  $p_i$  takes into account a possible censorship, and are assumed to be independent; typically  $p_1$  is chosen Gaussian around the true parallax,  $p_2$  is a Gaussian law for the absolute magnitude around the mean absolute magnitude,  $p_3$  is the velocity ellipsoid, and  $p_4$  is an exponential law in the galactic plane and in  $Z$ . The measurement error on apparent magnitude  $m$  and extinction should be taken into account in  $p_2$ , as for the proper motion  $(\mu_\alpha, \mu_\delta)$  or radial velocity  $V_R$  in  $p_3$ . An application to classical Cepheids is given, in this volume, in the paper by Luri et al. (1999).

As a by-product, the distance and absolute magnitude may be estimated through, e.g., the *a posteriori* expectation:

$$\hat{r} = E[1/\pi|O, \Theta] = \frac{\int_0^{+\infty} \frac{1}{\pi} g(O|\pi, \Theta) p(\pi) d\pi}{h(O|\Theta)} \quad (8)$$

$$\widehat{M} = E[m + 5 + 5 \log \pi|O; \Theta] = m + 5 + 5 \frac{\int_0^{+\infty} \log \pi g(O|\pi, \Theta) p(\pi) d\pi}{h(O|\Theta)} \quad (9)$$

The equations above use all the known information about one given star, and the parameters assumed for it, so that *individual* estimates of distance, absolute magnitude, etc, may be found, even e.g. if the concerned star has a zero or negative observed parallax. As for all Bayesian estimations, the drawback is of course that the *a priori* laws must be adequate, otherwise the final result may be biased.

For completeness, it must be noted that there is one special case where an *a posteriori* estimation may be used without any *a priori* law for the true parallaxes, assuming only a Gaussian error law for the random parallax errors, and making use only of the pdf of the observed parallaxes  $f(\pi_H)$ . This is the expectation of the true parallax given the observed parallax, a result found by Dyson (1926). The precision on the obtained estimate may be also be computed:

$$\hat{\pi} = E[\pi|\pi_H] = \pi_H + \sigma_{\pi_H}^2 \frac{f'(\pi_H)}{f(\pi_H)} \quad (10)$$

$$\sigma_{\hat{\pi}} = \sigma_{\pi_H} \sqrt{1 + \sigma_{\pi_H}^2 \left( \frac{f'(\pi_H)}{f(\pi_H)} \right)^2}$$

which is in general smaller than  $\sigma_{\pi_H}$  for unimodal parallax distributions. A more detailed discussion on the estimation of the true parallax distribution may be found in Lindegren (1995).

## 6. Conclusions

The Hipparcos Catalogue illustrates the various statistical problems one has to face when fundamental parameters have to be deduced from trigonometric parallaxes.

The Hipparcos errors may be considered Gaussian, at least at large scales, with no noticeable bias. At small scales, the correlation effect between measurements must be taken into account. Although the random parallax errors are symmetrical, with zero mean and dispersion as given by the formal error, a few outliers are however expected, e.g. due to binarity, in some rare cases.

The random errors may be misleading if improperly taken into account. In particular the transformation of parallaxes to distance or absolute magnitude should be done with caution. Moreover, truncations based on the observed parallax should be avoided: although corrections to the induced bias exist, they have large confidence intervals.

In order to estimate distances and absolute magnitudes several methods may be used. Either a transformation of the observed parallaxes, the use of asymptotically unbiased estimates, or a Bayesian approach, which takes efficiently into account the selection biases, but which rely on *a priori* laws.

Apart from its numerous astrophysical applications, one of the roles of the Hipparcos Catalogue will be to assess the validity of these *a priori* probability density functions. It will also assess the ground-based trigonometric parallaxes, which will expand our knowledge to fainter stars, until new space astrometry missions such as SIM or GAIA, are launched. In all cases, however, random measurement errors will still have to be taken into account.

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