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On Exponential Stabilization of Spin- $\frac{1}{2}$ Systems

W. Liang, N. H. Amini, P. Mason

Abstract—In this paper, we study the stabilization problem of quantum spin- $\frac{1}{2}$ systems under continuous-time measurements. In the case without feedback, we show exponential stabilization around the excited and ground state by providing a lower bound of the convergence rate. Based on stochastic Lyapunov techniques, we propose a parametrized feedback controller ensuring exponential convergence toward the target state. Moreover, we provide a lower bound of the convergence rate for this case. Then, we discuss the effect of each parameter appearing in the controller in the convergence rate. Finally, we illustrate the efficiency of such feedback controller through simulations.

I. INTRODUCTION

The evolution of an open quantum system undergoing indirect continuous-time measurements is described by the so-called quantum stochastic master equation, which has been derived by Belavkin in quantum filtering theory [6] (see also [9], [20] for a modern treatment). The solutions of such equation are called quantum trajectories and their properties have been studied in [14], [15].

The deterministic part of quantum stochastic master equation, which corresponds to the average dynamics, is given by the well known Lindblad operator. Its stochastic part represents the back-action effect of continuous-time measurements. Feedback control of open quantum systems has been the subject of many papers [20], [14], [19], [11]. In all of these papers, the control input appears in the Lindblad operator (through the system Hamiltonian). The experiment considered in [?] is a relevant example where a real-time feedback control is applied to stabilize an arbitrary photon number state in a microwave cavity.

The preparation of a pure state is investigated as an essential step towards quantum technologies. According to [1], the stochastic part, unlike the deterministic one, contributes to increase the purity of the quantum state, thus yielding a desirable effect for the preparation of a pure state. On the other hand, if we do not implement a control to the quantum system, the measurement induces a collapse of the quantum state towards either one of the eigenstates of the measurement operator, a phenomenon known as quantum state reduction [2], [20], [14], [18]. Thus, combining the continuous measurement with the feedback control may provide a desirable approach for preparing the target state in practice.

In [20], the authors design for the first time a quantum feedback controller that globally stabilizes a quantum spin- $\frac{1}{2}$

system (which is a special case of quantum angular momentum systems) towards an eigenstate of σ_z in the presence of imperfect measurements. This feedback controller has been designed by looking numerically for an appropriate global Lyapunov function. Then, in [14], by analyzing the stochastic flow and by using stochastic Lyapunov techniques, the authors construct a switching feedback controller which globally stabilizes the N -level quantum angular momentum system, in the presence of imperfect measurements, to the target eigenstate. A continuous version of this feedback controller has been proposed in [19]. The essential ideas in [20], [19] for constructing the continuous feedback controller remain the same: the controllers consist of two parts, the first one making the state converge locally to the target eigenstate, and the second one driving the system away from the antipodal eigenstates. More recently, [11] introduces a non-smooth Lyapunov-like theory to construct a continuous feedback controller that stabilizes globally the generic N -level quantum system at the target eigenstate.

The problem of estimating and optimizing the rate of convergence to the target eigenstate is also relevant in order to make the preparation more robust. For example, in [8], the authors prove the exponential stability of particular target subspaces for general N -level quantum system driven by Wiener processes and Poisson processes with an open-loop control (time-depending Hamiltonian) and in the case of perfect measurements. Also, in the recent paper [?] the authors have proven, by simple Lyapunov arguments, the stochastic exponential stability for a related control system obtained by applying a proportional output feedback.

The contribution of this paper is threefold: Firstly, we prove exponential stabilization around the excited state and ground state for a spin- $\frac{1}{2}$ system without feedback and we provide a lower bound of the convergence rate; Secondly, we construct a parametrized feedback controller which almost surely stabilizes exponentially the spin- $\frac{1}{2}$ system around the excited state (or ground state) providing a lower bound of the convergence rate; Finally, we study the effects of the parameters of the controller on the rate of convergence.

This paper is organized as follows. In Section II, we introduce the stochastic model describing spin- $\frac{1}{2}$ systems containing control inputs in the presence of imperfect measurements and we introduce the notions of stochastic stability needed throughout the paper. In Section III, we analyze the system without feedback and we provide an estimation of the rate of convergence to the excited or ground state in the quantum state reduction phenomenon. In Section IV, we propose a continuous-time feedback controller which almost surely globally exponentially stabilizes the quantum spin- $\frac{1}{2}$ system

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around the target eigenstate. We also analyze the asymptotic behavior of trajectories associated to the quantum dynamics with feedback. We precise that our continuous feedback has a similar form as the ones used in [20], [19]. Simulation results provided in Section V demonstrate the effectiveness of our control design.

II. PRELIMINARIES

A. System description

Consider a filtered probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t), \mathbb{P})$. Let W_t be a standard Wiener process and assume that \mathcal{F}_t is the natural filtration of the process W_t . The stochastic master equation describing the dynamics of a quantum spin- $\frac{1}{2}$ system is given by

$$\begin{aligned} d\rho_t &= \left(-i\frac{\omega_{eg}}{2}[\sigma_z, \rho_t] + \frac{M}{4}(\sigma_z\rho_t\sigma_z - \rho_t) - i\frac{u_t}{2}[\sigma_y, \rho_t] \right) dt \\ &\quad + \frac{\sqrt{\eta M}}{2}[\sigma_z\rho_t + \rho_t\sigma_z - 2\text{Tr}(\sigma_z\rho_t)\rho_t]dW_t \\ &=: F(\rho_t)dt + G(\rho_t)dW_t \end{aligned} \quad (1)$$

where $u_t = u(\rho_t)$ is the feedback controller, ω_{eg} is the difference between the energies of the excited state and the ground state, $\eta \in [0, 1]$ is determined by the efficiency of the photodetectors, and $M > 0$ is the strength of the interaction between the light and the atoms. The matrices σ_x , σ_y , and σ_z correspond to the Pauli matrices. The quantum state is described by the density operator ρ , which belongs to the space \mathcal{S} ,

$$\mathcal{S} := \{\rho \in \mathbb{C}^{2 \times 2} : \rho = \rho^*, \text{Tr}(\rho) = 1, \rho \geq 0\}, \quad (2)$$

where ρ^* denotes Hermitian conjugation. For a 2-level quantum system, ρ can be uniquely characterized by the Bloch sphere coordinates (x, y, z) as

$$\rho = \frac{\mathbb{1} + x\sigma_x + y\sigma_y + z\sigma_z}{2} = \frac{1}{2} \begin{bmatrix} 1+z & x-iy \\ x+iy & 1-z \end{bmatrix}. \quad (3)$$

The vector (x, y, z) belongs to the ball

$$B(\mathbb{R}^3) := \{(x, y, z) \in \mathbb{R}^3 : x^2 + y^2 + z^2 \leq 1\}.$$

The stochastic differential equation (1) expressed in the Bloch sphere coordinates takes the following form

$$\begin{aligned} dx_t &= \left(-\omega_{eg}y_t - \frac{Mx_t}{2} + u_t z_t \right) dt - \sqrt{\eta M}x_t z_t dW_t \\ dy_t &= \left(\omega_{eg}x_t - \frac{My_t}{2} \right) dt - \sqrt{\eta M}y_t z_t dW_t \\ dz_t &= -u_t x_t dt + \sqrt{\eta M}(1 - z_t^2) dW_t. \end{aligned} \quad (4)$$

The existence and uniqueness of the solution of (1) have been studied in [15], [14].

B. Stochastic stability

In this subsection, we introduce the notions of stochastic stability needed throughout the paper by adapting classical notions (see e.g. [13], [12]) to our setting. In order to provide them, we first present the definition of Bures distance [7].

Definition 2.1: The Bures distance between two quantum states ρ_a and ρ_b lying in the state space \mathcal{S} is given by

$$d_B(\rho_a, \rho_b) := \sqrt{2 - 2\sqrt{\text{Tr}(\rho_a\rho_b) + 2\sqrt{\det(\rho_a)\det(\rho_b)}}}.$$

Also, the Bures distance between a quantum state ρ_a and a set $E \subseteq \mathcal{S}$ is defined as $d_B(\rho_a, E) = \min_{\rho \in E} d_B(\rho_a, \rho)$. Given $E \subset \mathcal{S}$ and $r > 0$, we define the neighborhood $B_r(E)$ of E as

$$B_r(E) = \{\rho \in \mathcal{S} : d_B(\rho, E) < r\}.$$

Definition 2.2: Let \bar{E} be an invariant set of system (1), then \bar{E} is said to be

1. *locally stable in probability*, if for every $\epsilon \in (0, 1)$ and $r > 0$, there exists a $\delta = \delta(\epsilon, r)$, such that,

$$\mathbb{P}\{\rho_t \in B_r(\bar{E}) \text{ for } t \geq 0\} \geq 1 - \epsilon,$$

whenever $\rho_0 \in B_\delta(\bar{E})$.

2. *exponentially stable in mean*, if for some positive constants α and β ,

$$\mathbb{E}[d_B(\rho_t, \bar{E})] \leq \alpha d_B(\rho_0, \bar{E})e^{-\beta t},$$

whenever $\rho_0 \in \mathcal{S}$.

3. *almost surely exponentially stable*, if

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log d_B(\rho_t, \bar{E}) < 0, \quad a.s.$$

whenever $\rho_0 \in \mathcal{S}$. The left-hand side of the above inequality is called the Lyapunov exponent of the solution and its absolute value describes the rate of convergence.

Note that any equilibrium $\bar{\rho}$ of (1), that is any quantum state satisfying $F(\bar{\rho}) = G(\bar{\rho}) = 0$, corresponds to a special case of invariant set.

Notation 2.3: Suppose that the function $V(\rho, t) : \mathcal{S} \times \mathbb{R}_+ \rightarrow \mathbb{R}_+$ is continuously twice differentiable in ρ and once in t . We denote by \mathcal{L} the infinitesimal generator associated with the equation (1). The latter acts on V in the following way,

$$\begin{aligned} \mathcal{L}V(\rho, t) &:= \frac{\partial V(\rho, t)}{\partial t} + \frac{\partial V(\rho, t)}{\partial \rho} F(\rho) + \frac{1}{2} \frac{\partial^2 V(\rho, t)}{\partial \rho^2} (G(\rho), G(\rho)), \end{aligned}$$

where derivatives must be taken componentwise and the Hessian must be thought as a quadratic form. Then, by Itô's formula, $dV(\rho, t) = \mathcal{L}V(\rho, t)dt + \frac{\partial V(\rho, t)}{\partial \rho} G(\rho)dW_t$.

III. QUANTUM STATE REDUCTION

In this section, we study the dynamics of the quantum spin system (1) with $u_t = 0$. First, we note that in this case the states $\rho_g = \text{diag}(1, 0)$ and $\rho_e = \text{diag}(0, 1)$ are the equilibria of system (1), i.e., $F(\rho_g) = F(\rho_e) = 0$ and $G(\rho_g) = G(\rho_e) = 0$. In the following theorem, we prove that the quantum state reduction of dynamics (1) towards the invariant set $\bar{E} := \{\rho_g, \rho_e\}$ occurs with exponential velocity.

Theorem 3.1 (Quantum state reduction): For system (1), with $u_t = 0$ and $\rho_0 \in \mathcal{S}$, the set \bar{E} is exponentially stable in mean and a.s. with Lyapunov exponent less than $-\frac{\eta M}{2}$. Moreover, the probability of convergence to $\bar{\rho} \in \bar{E}$ is $\text{Tr}(\rho_0 \bar{\rho})$.

Proof: The proof proceeds in two steps:

Step 1: We take the standard deviation process

$$V(\rho_t) = \sqrt{\text{Tr}(\sigma_z^2 \rho_t) - \text{Tr}^2(\sigma_z \rho_t)} = \sqrt{1 - \text{Tr}^2(\sigma_z \rho_t)}$$

as candidate Lyapunov function. It is twice continuously differentiable with respect to ρ_t , except for the set \bar{E} . By Lemma 1.1 given in the appendix, we can assume that \bar{E} is never attained, which allows us to calculate $\mathcal{L}V(\rho_t) = -\frac{\eta M}{2}V(\rho_t)$. By Itô's formula, for all $\rho_0 \in \mathcal{S}$, $\mathbb{E}[V(\rho_t)] = V(\rho_0) + \int_0^t \mathbb{E}[\mathcal{L}V(\rho_s)]ds$ which implies $\mathbb{E}[V(\rho_t)] = V(\rho_0)e^{-\frac{\eta M}{2}t}$. We have

$$d_B(\rho_t, \bar{E}) = \begin{cases} d_B(\rho_t, \rho_e), & \text{if } \text{Tr}(\sigma_z \rho_t) > 0; \\ d_B(\rho_t, \rho_g), & \text{if } \text{Tr}(\sigma_z \rho_t) \leq 0. \end{cases}$$

Thus, for all $\rho_t \in \mathcal{S}$, $0 \leq d_B(\rho_t, \bar{E}) \leq \sqrt{2 - \sqrt{2}}$. Moreover, $C_1 d_B(\rho_t, \rho_e) \leq V(\rho_t) \leq C_2 d_B(\rho_t, \rho_e)$, if $\text{Tr}(\sigma_z \rho_t) > 0$; $C_1 d_B(\rho_t, \rho_g) \leq V(\rho_t) \leq C_2 d_B(\rho_t, \rho_g)$, if $\text{Tr}(\sigma_z \rho_t) \leq 0$, where $C_1 = \sqrt{1 + \sqrt{2}}/2$ and $C_2 = 2$. Then, we have

$$C_1 d_B(\rho_t, \bar{E}) \leq V(\rho_t) \leq C_2 d_B(\rho_t, \bar{E}). \quad (5)$$

Thus, for all $\rho_0 \in \mathcal{S}$, $\mathbb{E}[d(\rho_t, \bar{E})] \leq C_2/C_1 d(\rho_0, \bar{E})e^{-\frac{\eta M}{2}t}$, which implies that the set \bar{E} is exponentially stable in mean.

Now, we consider the following stochastic process,

$$Q(\rho_t) = e^{\frac{\eta M}{2}t} V(\rho_t) \geq 0,$$

whose infinitesimal generator is given by

$$\mathcal{L}Q(\rho_t) = e^{\frac{\eta M}{2}t} [\eta M/2 V(\rho_t) + \mathcal{L}V(\rho_t)] = 0.$$

Hence, the process $Q(\rho_t)$ is a positive martingale. Due to the Doob's martingale convergence theorem [16], the process $Q(\rho_t)$ converges almost surely to a finite limit A as $t \rightarrow \infty$. Consequently, $\sup_{t \geq 0} Q(\rho_t) = A$ implies $\sup_{t \geq 0} V(\rho_t) = Ae^{-\frac{\eta M}{2}t}$ a.s.. Letting $t \rightarrow \infty$, we obtain: $\limsup_{t \rightarrow \infty} \frac{1}{t} \log V(\rho_t) \leq -\frac{\eta M}{2}$ a.s.. By the inequality (5),

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log d_B(\rho_t, \bar{E}) \leq -\frac{\eta M}{2}, \quad a.s. \quad (6)$$

which means that the set \bar{E} is a.s. exponentially stable. Furthermore, we conclude that almost all paths which never exit the set $B_r(\bar{\rho}) \subset \mathcal{S}$ exponentially converge to $\bar{\rho} \in \bar{E}$.

Step 2: We follow the approach of [4], [2] to calculate the probability of convergence towards $\bar{\rho} \in \bar{E}$. According to the result of Step 1, we denote the terminal state ρ_∞ of the reduction process $\rho_\infty := P_e \rho_e + P_g \rho_g$, where P_e (resp. P_g) is the probability of reducing to ρ_e (resp. ρ_g), as $t \rightarrow \infty$. Since $\mathcal{L}\text{Tr}(\rho_t \bar{\rho}) = 0$, $\text{Tr}(\rho_t \bar{\rho})$ is a positive martingale. By the Doob's martingale convergence theorem, we have

$$\mathbb{E}[\text{Tr}(\rho_\infty \bar{\rho})] = \lim_{t \rightarrow \infty} \mathbb{E}[\text{Tr}(\rho_t \bar{\rho})] = \mathbb{E}[\text{Tr}(\rho_0 \bar{\rho})] = \text{Tr}(\rho_0 \bar{\rho}).$$

Therefore, $P_e = \text{Tr}(\rho_0 \rho_e)$, $P_g = \text{Tr}(\rho_0 \rho_g)$. The proof is complete. \square

IV. EXPONENTIAL STABILIZATION BY CONTINUOUS FEEDBACK

A. Almost sure global exponential stabilization

Lemma 4.1: For system (1) with $\rho_0 \in \mathcal{S}$, assume that under the feedback law u_t the system (1) admits $\bar{\rho} \in \bar{E}$ as

unique equilibrium, then for all $r \in (0, \sqrt{2}]$,

$$\mathbb{P}_{\rho_0} \{\tau_{B_r} < \infty\} = 1,$$

where τ_{B_r} denotes the first entry time of $B_r(\bar{\rho})$.

Proof: In order to prove the lemma, we will make use of the support theorem, in the form considered in [5]. For this purpose, one has to consider the Stratonovich version of the equation (1) and the corresponding controlled deterministic differential equation, which takes the form

$$\dot{\rho} = F(\rho) - \frac{1}{2} \frac{\partial G(\rho)}{\partial \rho} G(\rho) + G(\rho)w(t), \quad (7)$$

where w is the (deterministic) control input. Assume without loss of generality that $\bar{\rho} = \rho_e$, then by the hypothesis of the lemma we have that $F(\rho_g) \neq 0$ and therefore ρ_g is not an equilibrium point of (7). Moreover, since the z component of $G(\rho)$ in the Bloch sphere coordinates is different from zero outside \bar{E} , we deduce that, for any $\rho(0) \in \mathcal{S}$, there exists a control $w(t)$ steering (7) from $\rho(0)$ to $B_r(\bar{\rho})$ in finite time. Thus, the assumptions in [5, Propositions 4.3 and 4.6] are satisfied, in such a way that, for a fixed $T > 0$, there exists $\epsilon_T > 0$ such that, for all $\rho_0 \in \mathcal{S}$, $\mathbb{P}_{\rho_0}(\tau_{B_r} < T) \geq \epsilon_T$.

By Dynkin's inequality [10], $\forall \rho_0 \in \mathcal{S}$ and $\forall r \in (0, \sqrt{2}]$,

$$\mathbb{E}_{\rho_0}[\tau_{B_r}] \leq \frac{T}{1 - \sup_{\rho_0 \in \mathcal{S}} \mathbb{P}_{\rho_0}(\tau_{B_r} \geq T)} \leq \frac{T}{\epsilon_T} < \infty.$$

Then by Markov inequality, $\forall \rho_0 \in \mathcal{S}$ and $\forall r \in (0, \sqrt{2}]$, $\mathbb{P}_{\rho_0} \{\tau_{B_r} < \infty\} = 1$. The proof is complete. \square

Theorem 4.2: For system (1) with $\rho_0 \in \mathcal{S}$, denote as $\bar{\rho} \in \bar{E}$ the target eigenstate, and assume that the feedback controller u_t satisfies the condition of Lemma 1.1 given in the appendix and $u_t = 0$ in \bar{E} if and only if $\rho_t = \bar{\rho}$. Suppose that there exists a function $V(\rho)$, which is twice continuously differentiable on the set $B_r(\bar{\rho}) \setminus \bar{\rho}$ with $r \in (0, \sqrt{2}]$, and positive constants C, C_1, C_2 such that, $\forall \rho_t \in B_r(\bar{\rho}) \setminus \bar{\rho}$,

- (i) $C_1 d_B(\rho_t, \bar{\rho}) \leq V(\rho_t) \leq C_2 d_B(\rho_t, \bar{\rho})$,
- (ii) $\mathcal{L}V(\rho_t) \leq -CV(\rho_t)$.

Then $\bar{\rho}$ is a.s. exponentially stable.

Proof: The proof proceeds in 3 steps:

1. $\bar{\rho}$ is locally stable in probability.
2. $\exists T < \infty$ s.t. $\forall t \geq T, \rho_t \in B_r(\bar{\rho})$ a.s..
3. $\bar{\rho}$ is a.s. exponentially stable.

Step 1. Choose $r \in (0, \sqrt{2})$ and let $\epsilon \in (0, 1)$ be arbitrary, by the continuity of $V(\rho_t)$ and the fact that, for all $\rho_t \in \mathcal{S}$, $V(\rho_t) = 0$ if and only if $d_B(\rho_t, \bar{\rho}) = 0$, we can find a $\delta = \delta(\epsilon, r) > 0$ such that,

$$1/\epsilon \sup_{\rho_0 \in B_\delta(\bar{\rho})} V(\rho_0) \leq C_1 r. \quad (8)$$

Let τ be the first exit time of ρ_t from $B_r(\bar{\rho})$. By Itô's formula and the condition (ii),

$$\mathbb{E}[V(\rho_{t \wedge \tau})] \leq V(\rho_0) - C \mathbb{E} \left[\int_0^{t \wedge \tau} V(\rho_s) ds \right] \leq V(\rho_0).$$

For all $t \geq \tau$, $d_B(\rho_{t \wedge \tau}, \bar{\rho}) = d_B(\rho_\tau, \bar{\rho}) = r$. Hence, by the condition (i),

$$\begin{aligned} \mathbb{E}[V(\rho_{t \wedge \tau})] &\geq \mathbb{E}[\mathbb{1}_{\{t \wedge \tau\}} V(\rho_\tau)] \\ &\geq \mathbb{E}[\mathbb{1}_{\{t \wedge \tau\}} C_1 d_B(\rho_\tau, \bar{\rho})] \geq C_1 r \mathbb{P}(\tau \leq t). \end{aligned}$$

Combining with the inequality (8), we have

$$\mathbb{P}(\tau \leq t) \leq \frac{\mathbb{E}[V(\rho_{t \wedge \tau})]}{C_1 r} \leq \frac{V(\rho_0)}{C_1 r} \leq \epsilon.$$

Letting $t \rightarrow \infty$,

$$\mathbb{P}(\tau < \infty) \leq \epsilon \Rightarrow \mathbb{P}\{d_B(\rho_t, \bar{\rho}) < r \text{ for } t \geq 0\} \geq 1 - \epsilon,$$

which concludes the first step.

Step 2. Since $u_t = 0$ in \bar{E} if and only if $\rho_t = \bar{\rho}$, by Lemma 4.1, we obtain, for all $\rho_0 \in \mathcal{S}$, $\mathbb{P}_{\rho_0}\{\tau_{B_\delta} < \infty\} = 1$, where $\tau_{B_\delta} = \inf\{t \geq 0 : \rho_t \in B_\delta(\bar{\rho})\}$ with $\delta \in (0, r)$. It implies that ρ_t enters $B_\delta(\bar{\rho})$ in a finite time almost surely. Then we employ an argument inspired by [14] and suppose $\rho_0 \in \mathcal{S}$. We define two sequences of stopping times $\{\sigma_k\}_{k \geq 0}$ and $\{\tau_k\}_{k \geq 1}$ with $\sigma_0 = 0$ and, for $n \geq 1$, $\tau_n = \inf\{t \geq \sigma_{n-1} : \rho_t \in B_\delta(\bar{\rho})\}$ and $\sigma_n = \inf\{t \geq \tau_n : \rho_t \notin B_r(\bar{\rho})\}$. Then by the strong Markov property and the Bayes' formula, we have

$$\mathbb{P}_{\rho_0}(\tau_n < \infty | \sigma_{n-1} < \infty) = 1, \quad \mathbb{P}_{\rho_0}(\sigma_n < \infty | \tau_n < \infty) \leq \epsilon.$$

Moreover, by the construction of the stopping times, we have

$$\mathbb{P}_{\rho_0}(\tau_n < \infty | \sigma_n < \infty) = \mathbb{P}_{\rho_0}(\sigma_{n-1} < \infty | \tau_n < \infty) = 1.$$

Hence by Bayes' formula, we obtain

$$\begin{aligned} & \frac{\mathbb{P}_{\rho_0}(\sigma_n < \infty)}{\mathbb{P}_{\rho_0}(\sigma_{n-1} < \infty)} \\ &= \mathbb{P}_{\rho_0}(\sigma_n < \infty | \tau_n < \infty) \mathbb{P}_{\rho_0}(\tau_n < \infty | \sigma_{n-1} < \infty). \end{aligned}$$

In addition with $\mathbb{P}_{\rho_0}(\sigma_1 < \infty) \leq \epsilon$, we have $\mathbb{P}_{\rho_0}(\sigma_n < \infty) \leq \epsilon^n$ and thus $\sum_{n=1}^{\infty} \mathbb{P}_{\rho_0}(\sigma_n < \infty) \leq \sum_{n=1}^{\infty} \epsilon^n = \frac{\epsilon}{1-\epsilon} < \infty$. By the Borel-Cantelli lemma, we can conclude that, for almost all trajectories, there exists $N < \infty$, such that $\sigma_n = \infty$ for all $n \geq N$. That is, for almost all sample paths, there exists $T < \infty$ such that, for all $t \geq T$, $\rho_t \in B_r(\bar{\rho})$. This completes the proof of Step 2.

Step 3. In this step, we obtain an upper bound of the Lyapunov exponent by employing an argument inspired by [13, Theorem 3.3, p. 121]. For every fixed T consider the event $\Omega_T = \{\rho_t \in B_r(\bar{\rho}) \text{ for all } t \geq T\}$. Conditioning to $\rho_t \in \Omega_T$, from the condition (ii), we get

$$\mathcal{L} \log V(\rho_t) \leq -C - \frac{1}{2} [V_\rho(\rho_t)G(\rho_t)/V(\rho_t)]^2 =: -C - \frac{g^2(t)}{2}$$

for $t \geq T$ which implies

$$\begin{aligned} \log V(\rho_t) &= \log V(\rho_T) - C(t - T) \\ &\quad + \int_T^t g(s) dW_s - \frac{1}{2} \int_T^t g^2(s) ds. \end{aligned}$$

Let $m = 1, 2, 3 \dots$ and take arbitrarily $\epsilon \in (0, 1)$. By the exponential martingale inequality (see e.g. [13]), we have

$$\mathbb{P} \left\{ \sup_{T \leq t \leq T+m} \left[\int_T^t g(s) dW_s - \frac{\epsilon}{2} \int_T^t g^2(s) ds \right] > \frac{2}{\epsilon} \log m \right\} \leq \frac{1}{m^2}$$

Since $\sum_{m=1}^{\infty} \frac{1}{m^2} < \infty$, by Borel-Cantelli lemma, we have that for almost all sample paths, there exists m_0 , such that if $m > m_0$,

$$\sup_{T \leq t \leq T+m} \left[\int_T^t g(s) dW_s - \frac{\epsilon}{2} \int_T^t g^2(s) ds \right] \leq \frac{2}{\epsilon} \log m$$

Thus, for $T \leq t \leq T + m$ and $m > m_0$,

$$\int_T^t g(s) dW_s \leq \frac{2}{\epsilon} \log m + \frac{\epsilon}{2} \int_T^t g^2(s) ds, \quad a.s.$$

If $g^2(t) \geq K$, then we have

$$\log V(\rho_t) \leq \log V(\rho_T) - C(t - T) + \frac{2}{\epsilon} \log m - \frac{1 - \epsilon}{2} K$$

almost surely. It gives

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log V(\rho_t) \leq -C - \frac{1 - \epsilon}{2} K \quad a.s.$$

Letting $\epsilon \rightarrow 0$, we have $\limsup_{t \rightarrow \infty} \frac{1}{t} \log V(\rho_t) \leq -C - \frac{K}{2}$ a.s. In addition, due to (i), we have

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log d_B(\rho_t, \bar{\rho}) \leq -C - \frac{K}{2}, \quad a.s.$$

Since T can be taken arbitrarily large and Step 2 implies that $\lim_{T \rightarrow \infty} \mathbb{P}(\Omega_T) = 1$ we can conclude that $\bar{\rho}$ is almost surely exponentially stable. The proof is complete. \square

B. Feedback controller design

1) *Control design:* For simplicity, we define the following subset of \mathcal{S} ,

$$D_\lambda(\bar{\rho}) := \{\rho \in \mathcal{S} : 0 < \lambda < \text{Tr}(\rho \bar{\rho}) \leq 1\} = B_{r_\lambda}(\bar{\rho})$$

where $r_\lambda = \sqrt{2 - 2\sqrt{\lambda}}$ with $\lambda \in (0, 1)$.

Theorem 4.3: Consider system (1) with $\rho_0 \in \mathcal{S}$, denote as $\bar{\rho} \in \bar{E}$ the target eigenstate and let

$$V(\rho_t) = \sqrt{1 - \text{Tr}(\rho_t \bar{\rho})}, \quad \bar{\rho} \in \bar{E}. \quad (9)$$

Define the feedback control

$$u_t = u_t^{(1)} + u_t^{(2)} = \alpha [V(\rho_t)]^\beta - \gamma \text{Tr}(i[\sigma_y, \rho_t] \bar{\rho}) \quad (10)$$

where $\gamma \geq 0$, $\beta \geq 1$ and $0 < \alpha < \frac{\eta M \lambda^2}{(1-\lambda)^{\frac{\beta-1}{2}}}$ with $\lambda \in (0, 1)$.

Then u_t globally exponentially stabilizes system (1) a.s. to $\bar{\rho}$.

Proof: As $\rho_t \geq 0$, we have

$$|\text{Tr}(i[\sigma_y, \rho_t] \bar{\rho})| = |\text{Tr}(\sigma_x \rho_t)| = |x_t| \leq 2V(\rho_t).$$

We easily verify that u_t satisfies the condition of Theorem 4.2. The infinitesimal generator of (9) is given by

$$\mathcal{L}V(\rho_t) = \frac{u_t}{4} \frac{\text{Tr}(i[\sigma_y, \rho_t] \bar{\rho})}{V(\rho_t)} - \frac{\eta M}{2} \text{Tr}^2(\rho_t \bar{\rho}) V(\rho_t) \quad (11)$$

For all $\rho_t \in D_\lambda(\bar{\rho})$,

$$\begin{aligned} \mathcal{L}V(\rho_t) &\leq -1/2[\eta M \lambda^2 - \alpha V^{\beta-1}(\rho_t)] V(\rho_t) \\ &\leq -1/2[\eta M \lambda^2 - \alpha(1-\lambda)^{\frac{\beta-1}{2}}] V(\rho_t) \end{aligned}$$

If $0 < \alpha < \frac{\eta M \lambda^2}{(1-\lambda)^{\frac{\beta-1}{2}}}$, we get

$$\mathcal{L}V(\rho_t) \leq -C_\lambda V(\rho_t) \quad (12)$$

where $C_\lambda = 1/2[\eta M \lambda^2 - \alpha(1-\lambda)^{\frac{\beta-1}{2}}] > 0$. Moreover, we have

$$g^2(t) = [V_\rho(\rho_t)G(\rho_t)/V(\rho_t)]^2 \geq \eta M \lambda^2, \quad \forall \rho_t \in D_\lambda(\bar{\rho})$$

By Theorem 4.2, we can show that u_t exponentially stabilizes the system (1) to $\bar{\rho}$ almost surely. Furthermore, we have

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log d_B(\rho_t, \bar{\rho}) \leq -K_\lambda, \quad a.s.$$

where $K_\lambda = \eta M \lambda^2 - \frac{\alpha}{2}(1-\lambda)^{\frac{\beta-1}{2}}$. The proof is complete. \square

2) *Effect of the parameters of the controller:* We now study the dynamics of the closed-loop system and discuss informally how to choose the coefficients α, β, γ of u_t to stabilize the system more efficiently.

Case 1. Consider the case in which the trajectory ρ_t lies in a small neighborhood of the target eigenstate $\bar{\rho}$ permanently. According to the infinitesimal generator of the Lyapunov function (9) given by (11), we can see that the stochastic part of the stochastic master equation (1) and the deterministic part multiplied by $u_t^{(2)}$ always yield desirable effects on the preparation of the target eigenstate. On the other hand $u_t^{(1)}$ may impact negatively on the convergence to the target state. Choosing a larger β allows us to reduce the negative effect of $u_t^{(1)}$ for any fixed α .

Case 2. Assume that the initial state ρ_0 is contained in a neighborhood of the antipodal eigenstate. It is easy to check that, if such a neighborhood is small enough, then $|\mathcal{L}x_t| \geq \frac{\alpha}{2}$. In particular, choosing a larger α allows one to consider a larger neighborhood such that the inequality holds true. Letting τ be the first exit time from the neighborhood, by Itô's formula, we have $2 \geq \mathbb{E}[|x_\tau - x_0|] \geq \frac{\alpha}{2} \mathbb{E}[\tau]$. By Markov's inequality, we then get $\mathbb{P}[\tau \geq t] \leq \frac{\mathbb{E}[\tau]}{t} \leq \frac{4}{\alpha t}$. Thus taking a larger α permits to exit more quickly from the given neighborhood.

Case 3. Suppose that ρ_0 is far from \bar{E} . In this case, taking V as in (9), the term in $\mathcal{L}V$ corresponding to $u_t^{(2)}$ is always negative outside a neighborhood U of the antipodal eigenstate, and it is proportional to γ . Thus, reasoning as in Case 2 above, restricting to the trajectories that never enter U for $t \geq 0$ and defining τ as the first time when the trajectory enters in a neighborhood U' of $\bar{\rho}$, we get that a larger γ increases the probability of quickly entering U' .

V. NUMERICAL SIMULATIONS

For our simulations, since ρ_t should remain in \mathcal{S} , we consider the following Kraus form [3], [17] associated with (1),

$$\rho_t + d\rho_t = \frac{\mathbb{M}_{dY_t} \rho_t \mathbb{M}_{dY_t}^* + \frac{1-\eta M}{4} \sigma_z \rho_t \sigma_z dt}{\text{Tr} \left(\mathbb{M}_{dY_t} \rho_t \mathbb{M}_{dY_t}^* + \frac{1-\eta M}{4} \sigma_z \rho_t \sigma_z dt \right)},$$

where

$$\mathbb{M}_{dY_t} = \mathbf{1} - \left[\frac{i}{2} (\omega_{eg} \sigma_z + u_t \sigma_y) + \frac{M}{8} \mathbf{1} \right] dt + \frac{\sqrt{\eta M}}{2} \sigma_z dY_t$$

$$dY_t = dW_t + \sqrt{\eta M} \text{Tr}(\sigma_z \rho_t) dt$$

The simulations in the case $u_t = 0$ are shown in Fig. 1. In particular, we observe that the expectation of the Lyapunov function $\mathbb{E}[V(\rho_t)]$ fits precisely the exponential function $V(\rho_0)e^{-\frac{\eta M}{2}t}$, and the expectation of the Bures distance $\mathbb{E}[d_B(\rho_t, \bar{E})]$ is always below the exponential function $\frac{2d_B(\rho_0, \bar{E})}{\sqrt{1+\sqrt{2}/2}} e^{-\frac{\eta M}{2}t}$, which confirms the results of Section III.

We set ρ_e as the target eigenstate. The simulations corresponding to the exponential stabilization with u_t considered as in Section IV-B, are shown in Fig. 2 and 3. In Fig. 2, the system starting at $(-1, 0, 0)$ satisfies the condition of Case 3 of Section IV-B.2. We observe that $V(\rho_t)$ and $d_B(\rho_t, \rho_e)$ decrease quickly, that is ρ_t approaches ρ_e quickly. In Fig. 3, the system starting at $(0, 0, 1)$ satisfies the condition of

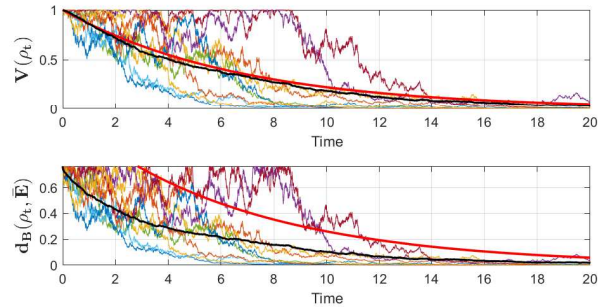


Fig. 1. Quantum state reduction with $u_t = 0$ starting at $(0, 0, 0)$, when $\omega_{eg} = 0, \eta = 0, 3, M = 1$: The black curve represents the mean value of the 10 arbitrary samples, the red curve represents the exponential reference.

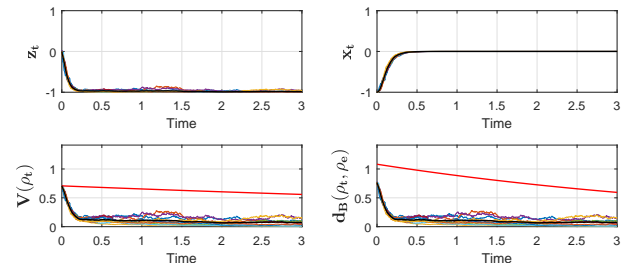


Fig. 2. Exponential stabilization with u_t starting at $(-1, 0, 0)$, when $\omega_{eg} = 0, \eta = 0, 3, M = 1, \alpha = 7.61, \beta = 5, \lambda = 0.9, \gamma = 10$: The black curve represents the mean value of all the samples, the red curve represents the exponential reference.

Case 2 and Case 3 in Section IV-B.2. We observe that $V(\rho_t)$ and $d_B(\rho_t, \rho_e)$ also decrease quickly.

VI. CONCLUSIONS AND FUTURE WORKS

In this paper, we have analyzed the asymptotic behavior of quantum spin- $\frac{1}{2}$ systems both for the cases without and with feedback. For the case without feedback, we have shown that for almost all trajectories ρ_t exponentially converge to the eigenstate ρ_e or ρ_g in mean and almost surely. By combining such asymptotic behavior for the case without feedback and the idea of breaking the symmetry of the state space \mathcal{S} , we designed a parametrized continuous feedback stabilizing exponentially the system toward a predetermined target eigenstate almost surely. Furthermore, we established lower bounds both with and without feedback. Also, we

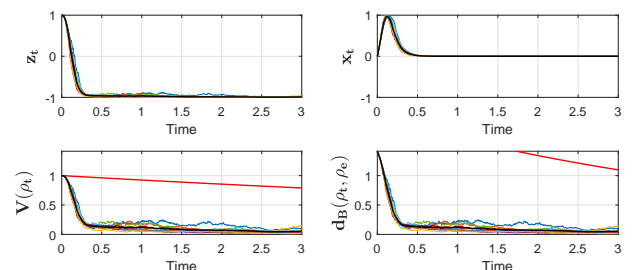


Fig. 3. Exponential stabilization with u_t starting at $(0, 0, 1)$, when $\omega_{eg} = 0, \eta = 0, 3, M = 1, \alpha = 7.61, \beta = 5, \lambda = 0.9, \gamma = 10$: The black curve represents the mean value of all the samples, the red curve represents the exponential reference.

discussed the role that each parameter appearing in the feedback law plays in this convergence speed. Finally, numerical simulations confirm our theoretical results.

Future work will be focused on the extension of the results presented in this paper to the exponential stabilization of an arbitrary eigenstate of N -level angular momentum system. Different directions remain to be explored. For example, optimizing our choice of feedback. Also, the adaptation of our results to the case where there are delays in the feedback loop is included in our future research lines.

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APPENDIX

The following lemma is inspired by analogous results in [12], [13].

Lemma 1.1: Assume that $\rho_0 \neq \bar{\rho}$ with $\bar{\rho} \in \bar{E}$ and u_t is continuous, continuously differentiable in $\mathcal{S} \setminus \bar{\rho}$ and $|u_t| \leq \Gamma\sqrt{1 - \text{Tr}(\rho_t \bar{\rho})}$ for some $\Gamma \in \mathbb{R}_+$. Then,

$$\mathbb{P}\{\rho_t \neq \bar{\rho}, \forall t \geq 0\} = 1.$$

Proof: Due to the fact that $\rho_t \geq 0$, in the Bloch sphere coordinates, if $|z_t| \rightarrow 1$ then $(x_t, y_t) \rightarrow (0, 0)$. Thus, it is sufficient to show the inaccessibility of ρ_e or ρ_g by proving the following two assertions

- 1) For all $z_0 \neq 1$, $|u_t| \leq \Gamma\sqrt{1 - z_t}$ with $\Gamma \in \mathbb{R}_+$,
 $\mathbb{P}\{z_t \neq 1, \forall t \geq 0\} = 1.$
- 2) For all $z_0 \neq -1$, $|u_t| \leq \Gamma\sqrt{1 + z_t}$ with $\Gamma \in \mathbb{R}_+$,
 $\mathbb{P}\{z_t \neq -1, \forall t \geq 0\} = 1.$

We will prove the first assertion. Suppose $\{z_t = 1\}$ is accessible, that is there exists z_0 such that $\mathbb{P}\{\tau < \infty\} > 0$, where $\tau = \inf\{t \geq 0 : z_t = 1\}$. Then, we can find a finite constant T sufficiently large such that $\mathbb{P}(B) > 0$, where $B = \{\tau \leq T\}$. Consider the stochastic process, $Q_t = 1/(1 - z_t)$ for all $z_t \neq 1$ and $t \in [0, T]$. By Itô's formula,

$$\mathcal{L}Q_t = -\frac{u_t x_t}{(1 - z_t)^2} + \frac{\eta M (1 - z_t^2)^2}{(1 - z_t)^2} \leq KQ_t$$

where $K = \sqrt{2}\Gamma + 4\eta M$. The above inequality holds as

$$\left| \sqrt{\eta M}(1 - z_t^2) \right| \leq 2\sqrt{\eta M}(1 - z_t)$$

and, by the assumption of the lemma, $|u_t x_t| \leq \Gamma\sqrt{2}(1 - z_t)$.

Due to the Grönwall-Bellman lemma, we construct the following process, $f_t = e^{-Kt}Q_t$ whose infinitesimal generator is given by

$$\mathcal{L}f_t = e^{-Kt}(\mathcal{L}Q_t - KQ_t).$$

Now for any $\epsilon \in (z_0, 1)$, define the stopping time $\tau_\epsilon = \inf\{t \geq 0 : z_t \notin [-1, \epsilon]\}$. By Itô's formula,

$$\mathbb{E}[f_{\tau_\epsilon \wedge t}] = Q_0 + \mathbb{E}\left[\int_0^{\tau_\epsilon \wedge t} \mathcal{L}f_s ds\right] \leq \frac{1}{1 - z_0}.$$

For all events in B , we have $\tau_\epsilon \leq T$ and $z_{\tau_\epsilon} = \epsilon$,

$$\mathbb{E}\left[\frac{e^{-KT}}{1 - \epsilon} \mathbb{1}_B\right] \leq \mathbb{E}[f_{\tau_\epsilon} \mathbb{1}_B] \leq \frac{1}{1 - z_0}$$

which implies $\mathbb{P}(B) \leq (1 - \epsilon)e^{-KT}/(1 - z_0)$. Letting $\epsilon \rightarrow 1$, we get $\mathbb{P}(B) = 0$, which contradicts the definition of B . Hence, the first assertion is proved.

For the second assertion, we have $|u_t x_t| \leq \Gamma\sqrt{2}(1 + z_t)$ and $|\sqrt{\eta M}(1 - z_t^2)| \leq 2\sqrt{\eta M}(1 + z_t)$. We can show similarly to the previous case that $\{z_t = -1\}$ is inaccessible. The proof is then complete. \square

REFERENCES

- [1] T. Abe, T. Sasaki, S. Hara, and K. Tsumura. Analysis on behaviors of controlled quantum systems via quantum entropy. *IFAC Proceedings Volumes*, 41(2):3695–3700, 2008.
- [2] S. L. Adler, D. C. Brody, T. A. Brun, and L. P. Hughston. Martingale models for quantum state reduction. *Journal of Physics A: Mathematical and General*, 34(42):8795, 2001.
- [3] H. Amini, M. Mirrahimi, and P. Rouchon. On stability of continuous-time quantum filters. In *Decision and Control and European Control Conference (CDC-ECC), 2011 50th IEEE Conference on*, pages 6242–6247. IEEE, 2011.
- [4] H. Amini, R. A. Somaraju, I. Dotsenko, C. Sayrin, M. Mirrahimi, and P. Rouchon. Feedback stabilization of discrete-time quantum systems subject to non-demolition measurements with imperfections and delays. *Automatica*, 49(9):2683–2692, 2013.
- [5] P. H. Baxendale. Invariant measures for nonlinear stochastic differential equations. In *Lyapunov Exponents*, pages 123–140. Springer, 1991.
- [6] V. P. Belavkin. Nondemolition measurements, nonlinear filtering and dynamic programming of quantum stochastic processes. In *Modeling and Control of Systems*, pages 245–265. Springer, 1989.
- [7] I. Bengtsson and K. Życzkowski. *Geometry of quantum states: an introduction to quantum entanglement*. Cambridge University Press, 2017.
- [8] T. Benoist, C. Pellegrini, and F. Ticozzi. Exponential stability of subspaces for quantum stochastic master equations. In *Annales Henri Poincaré*, volume 18, pages 2045–2074. Springer, 2017.
- [9] L. Bouten, R. Van Handel, and M. R. James. An introduction to quantum filtering. *SIAM Journal on Control and Optimization*, 46(6):2199–2241, 2007.
- [10] E. B. Dynkin. Markov processes. In *Markov Processes*, pages 77–104. Springer, 1965.
- [11] S. S. Ge, T. L. Vu, and C. C. Hang. Non-smooth lyapunov function-based global stabilization for quantum filters. *Automatica*, 48(6):1031–1044, 2012.
- [12] R. Khasminskii. *Stochastic stability of differential equations*, volume 66. Springer Science & Business Media, 2011.
- [13] X. Mao. *Stochastic differential equations and applications*. Elsevier, 2007.
- [14] M. Mirrahimi and R. Van Handel. Stabilizing feedback controls for quantum systems. *SIAM Journal on Control and Optimization*, 46(2):445–467, 2007.
- [15] C. Pellegrini. Existence, uniqueness and approximation of a stochastic schrödinger equation: the diffusive case. *The Annals of Probability*, pages 2332–2353, 2008.
- [16] D. Revuz and M. Yor. *Continuous martingales and Brownian motion*, volume 293. Springer Science & Business Media, 2013.
- [17] P. Rouchon. Models and feedback stabilization of open quantum systems. *arXiv preprint arXiv:1407.7810*, 2014.
- [18] A. Sarlette and P. Rouchon. Deterministic submanifolds and analytic solution of the quantum stochastic differential master equation describing a monitored qubit. *Journal of Mathematical Physics*, 58(6):062106, 2017.
- [19] K. Tsumura. Global stabilization at arbitrary eigenstates of n -dimensional quantum spin systems via continuous feedback. In *American Control Conference, 2008*, pages 4148–4153. IEEE, 2008.
- [20] R. Van Handel, J. K. Stockton, and H. Mabuchi. Feedback control of quantum state reduction. *IEEE Transactions on Automatic Control*, 50(6):768–780, 2005.