



HAL
open science

MEAN FIELD LIMIT FOR CUCKER-SMALE MODELS

R. Natalini, Thierry Paul

► **To cite this version:**

R. Natalini, Thierry Paul. MEAN FIELD LIMIT FOR CUCKER-SMALE MODELS. 2019. hal-02008699v1

HAL Id: hal-02008699

<https://hal.science/hal-02008699v1>

Preprint submitted on 5 Feb 2019 (v1), last revised 5 Feb 2021 (v5)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

MEAN FIELD LIMIT FOR CUCKER-SMALE MODELS

R. NATALINI AND THIERRY PAUL

ABSTRACT. In this very short note, we consider the Cucker-Smale dynamical system and we derive rigorously the Vlasov-type equation introduced in [4] in the mean-field limit. The vector field we consider is bounded at infinity in the velocity variables, and Lipschitz continuous in the space variables.

CONTENTS

1. introduction	1
2. The result	2
3. Proof	2
References	5

1. INTRODUCTION

We consider on \mathbf{R}^{2dN} the following Cucker-Smale type vector field [1, 2]

$$(1) \quad \begin{aligned} \dot{x}_i &= v_i \\ \dot{v}_i &= G_i(t, X, V), \quad i = 1, \dots, N \end{aligned}$$

where

$$(2) \quad G_i(t, X, V) = \frac{1}{N} \sum_{j=1}^N \phi(v_i - v_j) \psi(x_i - x_j).$$

Here the function ϕ is defined by

$$\phi(v) = \frac{v}{\sqrt{1 + |v|^2}}$$

and $\psi : \mathbf{R}^d \rightarrow \mathbf{R}$ is supposed to be Lipschitz continuous.

Moreover we used the notation $X = (x_1, \dots, x_N)$, $V = (v_1, \dots, v_N)$.

In fact we are rather interested in the Liouville equation associated to (1) [4], namely

$$(3) \quad \partial_t \rho + v \cdot \nabla_x \rho = \sum_{i=1}^N \nabla_{v_i} (G_i \rho)$$

with $\rho \in \mathcal{P}(\mathbf{R}^{2dN})$.

We want to prove that the marginals of $\rho(t)$ tend, as $N \rightarrow \infty$, to tensorial powers of the solution of a Vlasov type equation.

Such of Vlasov-type equation associated to (3) has been introduced in [4] and reads

$$(4) \quad \partial_t \rho_1(t, x, v) + v \cdot \nabla_x \rho_1 = \nabla_v (G_{\rho_1} \rho_1), \quad \rho_1|_{t=0} = \rho_1^{in} \in L^1(\mathbf{R}^{2d}, dx dv),$$

with

$$(5) \quad G_{\rho_1}(x, v) = \int_{\mathbf{R}^{2d}} \psi(x - y) \phi(v - w) \rho_1(y, w) dy dw.$$

2. THE RESULT

Let us start this section by recalling the definition of the second order Wasserstein distance $\text{dist}_{\text{MK},2}$ (see [5, 6]).

Definition 2.1 (quadratic Wasserstein distance). *The Wasserstein distance of order two between two probability measures μ, ν on \mathbf{R}^m with finite second moments is defined as*

$$\text{dist}_{\text{MK},2}(\mu, \nu)^2 = \inf_{\gamma \in \Gamma(\mu, \nu)} \int_{\mathbf{R}^m \times \mathbf{R}^m} |x - y|^2 \gamma(dx, dy)$$

where $\Gamma(\mu, \nu)$ is the set of probability measures on $\mathbf{R}^m \times \mathbf{R}^m$ whose marginals on the two factors are μ and ν .

We can now state the main result of this little note.

Theorem 2.2. *Let $\rho(t)$ be the solution of (3) with initial condition $\rho^{in} = (\rho_1^{in})^{\otimes N} \in L^1(\mathbf{R}^{2dN})$, $\rho_1^{in} \in L^1(\mathbf{R}^{2d})$, and let $\rho(t)_1 \in L^1(\mathbf{R}^{2d})$ be defined by*

$$(6) \quad \rho(t)_1(x, v) = \int_{\mathbf{R}^{2d(N-1)}} \rho(t)(x, v, x_2, \dots, x_N, v_2, \dots, v_N) dx_2 \dots dx_N dv_2 \dots dv_N.$$

Of course $(\rho^{in})_1 = \rho_1^{in}$.

Let $\rho_1(t)$ be the solution of (4) with initial condition ρ_1^{in} . Then, for all $N > 1, t \in \mathbf{R}$,

$$(7) \quad W_2(\rho(t)_1, \rho_1(t)) \leq CN^{-\frac{1}{2}}$$

with

$$(8) \quad C := 4 \|\psi\|_\infty \|\psi\|_\infty \frac{e^{\Lambda t} - 1}{\Lambda} \quad \Lambda := 2(1 + 2(\text{Lip}(\psi) + \text{Lip}(\phi))^2).$$

Remark 2.3. *There exists an equivalent result for higher orders marginals of $\rho(t)$ and other Monge-Kantorovich distances but we prefer in this little note to concentrate on the case of first marginal and quadratic Wasserstein distance. The proof in the more general situation is very close to the one presented here.*

3. PROOF

The proof will follow directly the proof of Theorem 3.1 in [3]. the only difference will be the presence of the external force F and the dependence in velocities of G_i .

Let $\pi_N^{in} \in \Pi((\rho^{in})^{\otimes N}, (\rho^{in})^{\otimes N})$ satisfy

$$(9) \quad T_\sigma \# \pi_N^{in} = \pi_N^{in}, \quad \text{for each } \sigma \in \mathfrak{S}_N,$$

where \mathfrak{S}_N is the group of permutations of N elements and

$$\begin{aligned} & T_\sigma(x_1, v_1, \dots, x_N, v_N, y_1, \xi_1, \dots, y_N, \xi_N) \\ &= (x_{\sigma(1)}, v_{\sigma(1)}, \dots, x_{\sigma(N)}, v_{\sigma(N)}, y_{\sigma(1)}, \xi_{\sigma(1)}, \dots, y_{\sigma(N)}, \xi_{\sigma(N)}). \end{aligned}$$

We will denote $X = (x_1, \dots, x_N), V = (v_1, \dots, v_N), Y = (y_1, \dots, y_N), \Xi = (\xi_1, \dots, \xi_N)$.

The following Lemma will be one of the keys of the proof of our Theorem.

Lemma 3.1. *Let $\pi_N(t)$ be the solution of*

$$(10) \quad \partial_t \pi_N + V \cdot \nabla_X \pi_N + \Xi \cdot \nabla_Y \pi_N = \sum_{i=1}^N \left(\nabla_{\xi_i} \cdot (G_i(Y, \Xi) \pi_N) + \nabla_{v_i} \cdot (G_{\rho_1(t)}(x_i, v_i) \pi_N) \right)$$

with $\pi_N(0) = \pi_N^{\text{in}}$.

Then, for all $t \in \mathbf{R}$, $\pi_N(t)$ is a coupling between $\rho(t)$ and $\rho_1(t)^{\otimes N}$. Moreover $\pi_N(t)$ is invariant by permutations T_σ .

Proof. By taking the two marginals of the two sides of the equality, one get that they satisfy the two Liouville and Vlasov equations. the result is then obtained by uniqueness of the solutions of both equations. \square

Lemma 3.2. *Let*

$$D_N(t) := \int \frac{1}{N} \sum_{j=1}^N (|x_j - y_j|^2 + |\xi_j - \eta_j|^2) ((X - Y)^2 + (V - \Xi)^2) \pi_N(t).$$

Then

$$\frac{dD_N}{dt} \leq \Lambda D_N + \frac{1}{N} \sum_{j=1}^N \int |G_{\rho_1(t)}(x_i, v_i) - G_i(X, V)|^2 d(\rho_1(t))^{\otimes N},$$

with

$$\Lambda := 2(1 + 2 \text{Lip}(\nabla V)^2).$$

Proof. We first notice that

$$\frac{dD_N}{dt} = 2 \int \left((\Xi - V) \cdot (X - Y) + \sum_{i=1}^N (v_i - \xi_i) \cdot (G_i(Y, \Xi) - G_{\rho_1(t)}(x_i, v_i)) \right) d\pi_N$$

Using $2uv \leq u^2 + v^2$ we get

$$\begin{aligned} \frac{dD_N}{dt} &\leq \frac{1}{N} \int \left((X - Y)^2 + 2(V - \Xi)^2 + \frac{1}{N} \sum_{i=1}^N |G_i(Y, \Xi) - G_{\rho_1(t)}(x_i, v_i)|^2 \right) d\pi_N \\ &\leq 2D_N(t) + \frac{1}{N} \int \left(\frac{1}{N} \sum_{i=1}^N |G_i(Y, \Xi) - G_{\rho_1(t)}(x_i, v_i)|^2 \right) d\pi_N. \end{aligned}$$

Let us introduce in the square inside the last integral the nul term

$$\frac{1}{N} \sum_{j=1}^N \frac{v_i - v_j}{\sqrt{1 + |v_i - v_j|^2}} \phi(x_i - x_j) - \frac{1}{N} \sum_{j=1}^N \frac{v_i - v_j}{\sqrt{1 + |v_i - v_j|^2}} \phi(x_i - x_j)$$

Then

$$\left| \frac{v_i - v_j}{\sqrt{1 + |v_i - v_j|^2}} \phi(x_i - x_j) - \frac{\xi_i - \xi_j}{\sqrt{1 + |\xi_i - \xi_j|^2}} \phi(y_i - y_j) \right|^2$$

can be approximated by

$$\begin{aligned} &((v_i - v_j - (\xi_i - \xi_j))^2 + (x_i - x_j - (y_i - y_j))^2) \text{Lip}(\phi)^2 \\ &\leq 4 \text{Lip}(\phi)^2 ((x_i - y_i)^2 + (v_i - \xi_i)^2 + (x_j - y_j)^2 + (v_j - \xi_j)^2) \end{aligned}$$

This part gives the $4\text{Lip}(\phi)^2$ part in Λ , and since the remaining part to integrate contains only the (X, V) variable, the integral against π_N can be replaced by the one against $\rho_1(t)^{\otimes N}$. \square

The following result is verbatim Lemma 3.3. in [3] with the special value $p = 2$ and d replaced by $2d$.

Lemma 3.3. *Let F be a bounded vector field on \mathbf{R}^{2d} , and ρ be a probability density on \mathbf{R}^{2d} . For each $j = 1, \dots, N$, one has*

$$\int \left| F \star \rho(x_j, v_j) - \frac{1}{N} \sum_{k=1}^N F(x_j - x_k, v_j - v_k) \right|^2 \prod_{m=1}^N \rho(x_m, v_m) dx_m dv_m \leq \frac{4}{N} (2\|F\|_{L^\infty})^2.$$

Lemma 3.3 with $F(x, v) = \phi(v)\psi(x)$ together with Lemma 3.2 gives immediatly that

$$\frac{dD_N}{dt} \leq \Lambda D_N + 2^4 \frac{2}{N} \|\nabla V\|_{L^\infty}^2$$

and, by Gronwall's inequality,

$$(11) \quad D_N(t) \leq D_N(0)e^{\Lambda t} + 2^5 \|\phi \otimes \psi\|_{L^\infty}^2 \frac{1}{N} \frac{e^{\Lambda t} - 1}{\Lambda}.$$

Lemma 3.4. *$(\pi_N(t))_1$ is a coupling between $\rho_1(t)$ and $(\rho(t))_1$.*

Since $\pi_N(t)$ is symmetric by permutations, one has easily that

$$(12) \quad D_N(t) = \int (|x_1 - y_1|^2 + |v_1 - \xi_1|^2) (\pi_N(t))_1$$

and Lemma 3.4 immediatly implies that

$$(13) \quad D_N(t) \geq W_2(\rho_1(t), (\rho(t))_1)^2,$$

and by Lemma 3.3

$$(14) \quad W_2(\rho_1(t), (\rho(t))_1)^2 \leq D_N(0)e^{\Lambda t} + 2^5 \|\phi\|_{L^\infty}^2 \|\psi\|_{L^\infty}^2 \frac{1}{N} \frac{e^{\Lambda t} - 1}{\Lambda}.$$

Remember that (14) is valid for $D(0)$ as defined in Lemma 3.2 for *any* π_N coupling $(\rho_1^{in})^{\otimes N}$ with itself. Since W_2 is a distance, one has

$$W_2((\rho_1^{in})^{\otimes N}, (\rho_1^{in})^{\otimes N})^2 = \int ((X - Y)^2 + (V - \Xi)^2) \pi_N^{op} = 0$$

for some optimal coupling π_N^{op} .

Choosing now this coupling for the definition of $D_N(0)$ in Lemma 3.2 we get that

$$D_N(0) = \frac{1}{N} W_2((\rho_1^{in})^{\otimes N}, (\rho_1^{in})^{\otimes N})^2 = 0$$

so that (14) gives the result. Theorem 2.2 is proven.

REFERENCES

- [1] F. Cucker and S. Smale, On the mathematics of emergence. *Japan. J. Math.* **2**, 197227, (2007).
- [2] F. Cucker and S. Smale, Emergent behavior in flocks. *IEEE Trans. Automat. Control* **52**, 852-862, 2007.
- [3] F. Golse, C. Mouhot, T. Paul, On the Mean-Field and Classical Limits of Quantum Mechanics, *Communication in Mathematical Physics* **343**, 165-205 (2016).
- [4] S.-Y. Ha, E. Tadmor, From particle to kinetic and hydrodynamic descriptions of flocking, *Kinet. Relat. Models* **1** (2008) 415-435.
- [5] C. Villani: "Topics in Optimal Transportation", Amer. Math. Soc., Providence (RI), 2003.
- [6] C. Villani: "Optimal Transport. Old and New", Springer-Verlag, Berlin, Heidelberg, 2009.

(R.N.) IAC, VIA DEI TAURINI, 19, 00185 ROMA RM ITALY

E-mail address: roberto.natalini@cnr.it

(T.P.) CMLS, ECOLE POLYTECHNIQUE, CNRS, UNIVERSITÉ PARIS-SACLAY, 1128 PALAISEAU CEDEX, FRANCE

E-mail address: thierry.paul@polytechnique.edu