A simple turbulent model for compressible flows

Jean-Marc Hérard *

* EDF R&D, MFEE 6 quai Watier, Chatou. France. jean-marc.herard@edf.fr

MODTERCOM - Porquerolles, France. September 08-09, 2014

イロト イポト イヨト イヨト

Outline



- 2 Turbulent compressible model
- Shock and rarefaction waves
- Solution of the 1D Riemann problem
- 5 Summary and perspectives

6 References

イロト イポト イヨト イヨト

Introduction

Turbulent compressible model Shock and rarefaction waves Solution of the 1D Riemann problem Summary and perspectives References

Introduction

This work is a sequel of some previous work on one and two-equation turbulent models for single phase compressible flows (B. Audebert, C. Berthon, F. Coquel, A. Forestier, S.Gavrilyuk, X. Louis, E. Xeuxet, [1,2,3,4,6,8] among others...).

The main objective is to get a meaningful and cheap conservative model for turbulent industrial *unsteady* applications (with little information on boundary and/or initial conditions for the turbulent kinetic energy).

Another goal is to obtain a rough way in order to account for Reynolds stress tensors in two-phase flow two-fluid models.

Introduction

Turbulent compressible model Shock and rarefaction waves Solution of the 1D Riemann problem Summary and perspectives References

Introduction

This work is a sequel of some previous work on one and two-equation turbulent models for single phase compressible flows (B. Audebert, C. Berthon, F. Coquel, A. Forestier, S.Gavrilyuk, X. Louis, E. Xeuxet, [1,2,3,4,6,8] among others...).

The main objective is to get a meaningful and cheap conservative model for turbulent industrial *unsteady* applications (with little information on boundary and/or initial conditions for the turbulent kinetic energy).

Another goal is to obtain a rough way in order to account for Reynolds stress tensors in two-phase flow two-fluid models.

▶ ▲ 同 ▶ ▲ 臣 ▶ ▲ 臣 ▶

Introduction

Turbulent compressible model Shock and rarefaction waves Solution of the 1D Riemann problem Summary and perspectives References

Introduction

This work is a sequel of some previous work on one and two-equation turbulent models for single phase compressible flows (B. Audebert, C. Berthon, F. Coquel, A. Forestier, S.Gavrilyuk, X. Louis, E. Xeuxet, [1,2,3,4,6,8] among others...).

The main objective is to get a meaningful and cheap conservative model for turbulent industrial *unsteady* applications (with little information on boundary and/or initial conditions for the turbulent kinetic energy).

Another goal is to obtain a rough way in order to account for Reynolds stress tensors in two-phase flow two-fluid models.

ト 4 同 ト 4 三 ト 4 三 ト

Governing equations

In the following, $\rho(\mathbf{x}, t)$, $\mathbf{U}(\mathbf{x}, t)$, $P(\mathbf{x}, t)$ and $K(\mathbf{x}, t)$ respectively denote the mean density, the mean velocity, the mean pressure ate the turbulent kinetic energy.

The governing set of equations is:

$$\begin{cases} \partial_t (\rho) + \nabla \cdot (\rho \mathbf{U}) = 0 ;\\ \partial_t (\rho \mathbf{U}) + \nabla \cdot (\rho \mathbf{U} \otimes \mathbf{U} + (P + 2K/3)\mathbf{Id}) = \nabla \cdot (\sigma^{\mathbf{v}}) ;\\ \partial_t (\rho E) + \nabla \cdot (\mathbf{U}(\rho E + P + 2K/3)) = \nabla \cdot (\sigma^{\mathbf{v}} : \mathbf{U}) . \end{cases}$$
(1)

Where we note :

$$\rho E = \rho \ (P, \rho) + \rho \mathbf{U} \cdot \mathbf{U}/2 + K \tag{2}$$

and also:

$$K = K_0 \rho^{5/3} \tag{3}$$

伺き くほき くほき

-

Main properties of this set of PDEs in a 1D framework

We introduce the condensed form of (1) which reads:

$$\partial_t (W) + B(W) \partial_x (W) = 0 \tag{4}$$

and introduce $c(P, \rho)$ such that:

$$c^{2}(\boldsymbol{P},\rho) = \left(\frac{\boldsymbol{P}}{\rho^{2}} - \partial_{\rho}\left(\epsilon(\boldsymbol{P},\rho)\right)\right) / (\partial_{\boldsymbol{P}}\left(\epsilon(\boldsymbol{P},\rho)\right))$$

Property 1 (Hyperbolicity):

• In a 1D framework, system (1) is strictly hyperbolic:

$$\lambda_1 = U$$
 , $\lambda_2 = U - \tilde{c}$, $\lambda_3 = U + \tilde{c}$ (5)

ト (同) (三) (三)

-2

where : $(\tilde{c})^2 = c^2 + 10 K \notin 9$.

• The 1-field is LD and the 2 – 3- fields are GNL.

Main properties of this set of PDEs in a 1D framework

We introduce the condensed form of (1) which reads:

$$\partial_t (W) + B(W) \partial_x (W) = 0 \tag{4}$$

and introduce $c(P, \rho)$ such that:

$$c^{2}(\boldsymbol{P},\rho) = \left(\frac{\boldsymbol{P}}{\rho^{2}} - \partial_{\rho}\left(\epsilon(\boldsymbol{P},\rho)\right)\right) / (\partial_{\boldsymbol{P}}\left(\epsilon(\boldsymbol{P},\rho)\right))$$

Property 1 (Hyperbolicity):

• In a 1D framework, system (1) is strictly hyperbolic:

$$\lambda_1 = U$$
 , $\lambda_2 = U - \tilde{c}$, $\lambda_3 = U + \tilde{c}$ (5)

ト 4 同 ト 4 三 ト 4 三 ト -

-2

where : $(\tilde{c})^2 = c^2 + 10 K \not= 9$.

• The 1-field is LD and the 2 – 3- fields are GNL.

Entropy inequality

We define $S(P, \rho)$ such that:

$$c^{2}(P,
ho)\partial_{P}\left(\mathcal{S}(P,
ho)
ight)+\partial_{
ho}\left(\mathcal{S}(P,
ho)
ight)=0$$

Property 2 (Entropy inequality):

We note:

$$\eta = \rho Log(S)$$

$$f_{\eta} = \rho Log(S)U$$
(6)

the entropy-entropy flux pair.

Solutions W(x, t) of (1) comply with:

$$\partial_t(\eta) + \nabla_{\cdot}(f_{\eta}) \leq 0$$
 (7)

イロト イポト イヨト イヨト

-20

Jump conditions

Owing to the conservative form, we get jump conditions at once; noting σ the speed of the travelling discontinuity separating two physical states W_l and W_r , we have:

$$-\sigma[W]_{l}^{r} + [F(W)]_{l}^{r} = 0$$
(8)

or:

$$[\rho v]_{l}^{r} = 0 \rho_{l} \rho_{r} ([v]_{l}^{r})^{2} = [P + 2K/3]_{l}^{r} [\rho]_{l}^{r} \rho v \left([\epsilon(P, \rho) + K\tau] l^{r} + (\overline{P}_{l,r} + 2\overline{K}_{l,r}/3)[\tau]_{l}^{r} \right) = 0$$

$$(9)$$

setting :

$$v = U - \sigma$$

and also:

$$[\psi]_{l}^{r} = \psi_{r} - \psi_{l}, \quad (\psi)_{l,r} = (\psi_{r} + \psi_{l})/2.$$

・ロト ・同ト ・ヨト・

We focus on a - γ -perfect gas EOS (with $\gamma > 1$). We recall that in that case, a shock wave separating two states Z_r and Z_l is such that:

$$(\beta)^{-1} \le \frac{\rho *_r}{\rho *_l} \le \beta$$

where: $\beta = \frac{\gamma+1}{\gamma-1}$.

We can derive the exact form of the EOS for the mean variables, that is:

$$\rho (P, \rho) = P/(\gamma - 1).$$
 (10)

What about the shock wave separating mean variables W_I, W_r ?

Actually, the answer is, as expected:

$$(\beta)^{-1} \leq \frac{\rho_r}{\rho_l} \leq \beta$$

We focus on a - γ -perfect gas EOS (with $\gamma > 1$). We recall that in that case, a shock wave separating two states Z_r and Z_l is such that:

$$(\beta)^{-1} \le \frac{\rho *_r}{\rho *_l} \le \beta$$

where: $\beta = \frac{\gamma+1}{\gamma-1}$.

We can derive the exact form of the EOS for the mean variables, that is:

$$\rho (P, \rho) = P/(\gamma - 1).$$
 (10)

What about the shock wave separating mean variables W_I, W_r ?

Actually, the answer is, as expected:

We focus on a - γ -perfect gas EOS (with $\gamma > 1$). We recall that in that case, a shock wave separating two states Z_r and Z_l is such that:

$$(\beta)^{-1} \le \frac{\rho *_r}{\rho *_l} \le \beta$$

where: $\beta = \frac{\gamma+1}{\gamma-1}$.

We can derive the exact form of the EOS for the mean variables, that is:

$$\rho (P, \rho) = P/(\gamma - 1). \tag{10}$$

What about the shock wave separating mean variables W_I, W_r ?

Actually, the answer is, as expected:

$$(\beta)^{-1} \leq \frac{\rho_r}{\rho_l} \leq \beta$$

Solution of the 1D Riemann problem

Once more, we focus on a - γ -perfect gas EOS (with $\gamma > 1$).

Property 3 (Solution of the 1D RP):

The Riemann problem associated with initial states $W(x < 0, t = 0) = W_L$ and $W(x > 0, t = 0) = W_R$ admits a unique self-similar solution:

$$W(x,t) = \omega(x/t)$$

with no vacuum occurence, provided that initial left and right states $W_{L,R}$ are such that:

$$U_R - U_L < h(S_L, P_L) + h(S_R, P_R)$$

where:

$$h(S_{L,R}, P_{L,R}) = \int_{0}^{P_{L,R}} \frac{1}{\gamma P} \left(\gamma S_{L,R} (P/S_{L,R})^{\frac{\gamma-1}{\gamma}} + \frac{10K_{0}}{9} (P/S_{L,R})^{\frac{2}{3\gamma}} \right)^{1/2} dP$$

Intermediate states $\rho_{1,2}$ and $P_{1,2}$ are in the physical range.

イロト イポト イヨト イヨト

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

・ロト ・ 同ト ・ ヨト ・ ヨト

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

・ロト ・ 同ト ・ ヨト ・ ヨト

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

・ロト ・ 同ト ・ ヨト ・ ヨト

Summary and perspectives

Summary and perspectives

- A simple turbulent model for single-phase flow applications, even if complex EOS are involved, with an interesting behaviour in shock waves;
- Approximate solutions using Finite Volume methods: exact Godunov scheme or approximate Godunov scheme (VFRoe-ncv using U, P + 2K/3, S variables [5]);
- The solution of the 1DRP enables to derive suitable boundary conditions for computational purposes;
- A comparison with recent models for shallow water flows, proposed by G.L. Richard and S. Gavrilyuk, might be of interest [9,10];
- An extended version of the Baer-Nunziato model (Int. J. Multiphase Flow, 1986) has been recently derived (hyperbolic system, entropy inequality, unique jump conditions in single fields, see [7]).

References

 Contribution à l'analyse des modèles aux tensions de Reynolds pour l'interaction choc turbulence, B. Audebert, *PhD thesis*, Université Paris VI, 2006
 Contribution à l'analyse numérique des équations de Navier-Stokes compressibles à deux entropies spécifiques. Applications à la turbulence compressible, C. Berthon, *PhD thesis*, Université Paris VI, 1999.

3- An approximate Roe-type Riemann solver for a class of realizable second-order closures, G. Brun, J.M.H., D. Jeandel, M. Uhlmann, *223-249*, vol 13(3), 2000 (*IJCFD*).
4- A Godunov type solver to compute turbulent compressible flows, A. Forestier, J.M.H., X. Louis, vol 324, 1997 (*CRAS Paris I*).

5- On the use of symetrizing variables for vacuums, T. Gallouet, J.M. H., N. Seguin *163-194*, vol 40(3), 2003 (*Calcolo*).

6 - Geometric evolution of the Reynolds stress tensor, S.L. Gavriluyk, H. Gouin, *65-73*, vol 59, 2012 (*international Journal of Engineering Science*)

7- A simple turbulent two-phase flow model, J.M.H., H. Lochon, in prepraration.

8- Modélisation numérique de la turbulence compressible, X. Louis, *PhD thesis*, Université Paris VI, 1995.

9 - A new model of roll waves: comparison with Brock's experiments, G.L. Richard, S.L. Gavriluyk, *374-405*, vol 698, 2012 (*Journal of Fluid Mechanics*)

10 - The classical hydraulic jump in a model of shear shallow water flows, G.L. Richard, S.L. Gavriluyk, *492-521*, vol 725, 2013 (*Journal of Fluid Mechanics*)

- A 🗇 🕨 - A 🖻 🕨 - A 🖻 🕨

-