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Direct Observation of Ion Dynamics in Supercapacitor Electrodes using *in situ* Diffusion NMR Spectroscopy

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Ionic transport inside porous carbon electrodes underpins the storage of energy in supercapacitors and the rate at which they can charge and discharge, yet few studies have elucidated the materials properties that influence ion dynamics. Here we use *in situ* pulsed field gradient NMR spectroscopy to measure ionic diffusion in supercapacitors directly. We find that confinement in the nanoporous electrode structures decreases the effective self-diffusion coefficients of ions by over two orders of magnitude compared to neat electrolyte, and in-pore diffusion is modulated by changes in ion populations at the electrode-electrolyte interface during charging. Electrolyte concentration and carbon pore size distributions also affect in-pore diffusion and the movement of ions in and out of the nanopores. In light of our findings we propose that controlling the charging mechanism may allow the tuning of the energy and power performances of supercapacitors for a range of different applications.

As renewable energy and green technologies such as electric vehicles become prevalent, we must develop new ways to store and release energy on a range of timescales. Rechargeable batteries are ideal for timescales of minutes or hours (electric cars, portable electronic devices, grid storage etc.), while supercapacitors are more promising for second or sub-second timescales and are increasingly being used for transport applications where rapid charging and discharging are required. The superior power handling and cycle lifetime of supercapacitors comes at the expense of energy density, with recent materials-driven research aiming to address this issue by fine-tuning the nanoporous structure of the carbon electrodes,^{1,2} and by using ionic liquid electrolytes that are stable at higher voltages.^{3,4} Both approaches have afforded some increases in energy density, though not without sacrificing power density. The delicate balance between energy and power density must be understood if supercapacitors are to be used in a wide range of applications.

Fundamental studies based on spectroscopic,⁵⁻¹⁴ and theoretical,¹⁵⁻¹⁸ methods have recently revealed the complex nature of charging in supercapacitors. Prior to charging, the electrode pores contain a large number of electrolyte ions,^{15,19,20} and as a result charge storage is generally more complex than simple counter-ion adsorption (counter-ions are defined as having charge opposite to the electrode in which they are located).^{5-7,15} A range of different charging mechanisms can operate

depending on the choice of electrode and electrolyte, and the electrode polarisation.^{21,22} For example, charging often proceeds via the exchange of counter-ions and co-ions (co-ions are defined as having charge of the same sign as the electrode in which they are located), while charging by co-ion desorption alone is also a possibility. We introduced the charging mechanism parameter, X , to quantify these different processes, with X taking values of +1, 0 and -1, for the extreme cases of charging by counter-ion adsorption, counter-ion – co-ion exchange, and co-ion desorption, respectively, while intermediate X values indicate contributions from more than one mechanism.²¹

An understanding and control of the charge storage mechanism (X value) may hold the key to optimising the energy and power performance of supercapacitors for different applications. However, a crucial missing part of our understanding is how the electrolyte ions diffuse and migrate in supercapacitor electrodes. Theoretical studies based on molecular dynamics (MD) and mean-field theories have shown that ion-ion interactions, ion-carbon interactions and the electrode pore size all influence the rates of ionic diffusion in supercapacitor electrodes,^{16,18,23,24} though there is not yet a clear consensus on which factors are most important, nor the order of magnitude by which diffusion is influenced by confinement and by charging. Experimental methods to directly probe in-pore motion in an ion-selective and electrode-selective way have until now been lacking.

Here we show how an *in situ* pulsed field gradient (PFG) NMR approach can be used to measure ionic diffusion in the nanopores of supercapacitor electrodes. Confinement results in reductions in self-diffusion coefficients by over two orders of magnitude compared to neat electrolyte solutions, while *in situ* measurements show that changes of in-pore ion populations during charging modulate in-pore ionic diffusion. We also show that the electrolyte concentration and nanopore size have significant effects on in-pore diffusion, and on the exchange of ions between bulk and in-pore sites. Our findings offer detailed insights into diffusion and exchange processes in porous electrodes and bring new opportunities for understanding and controlling the charging dynamics of supercapacitors.

Introducing Pulsed Field Gradient NMR Studies of Supercapacitors

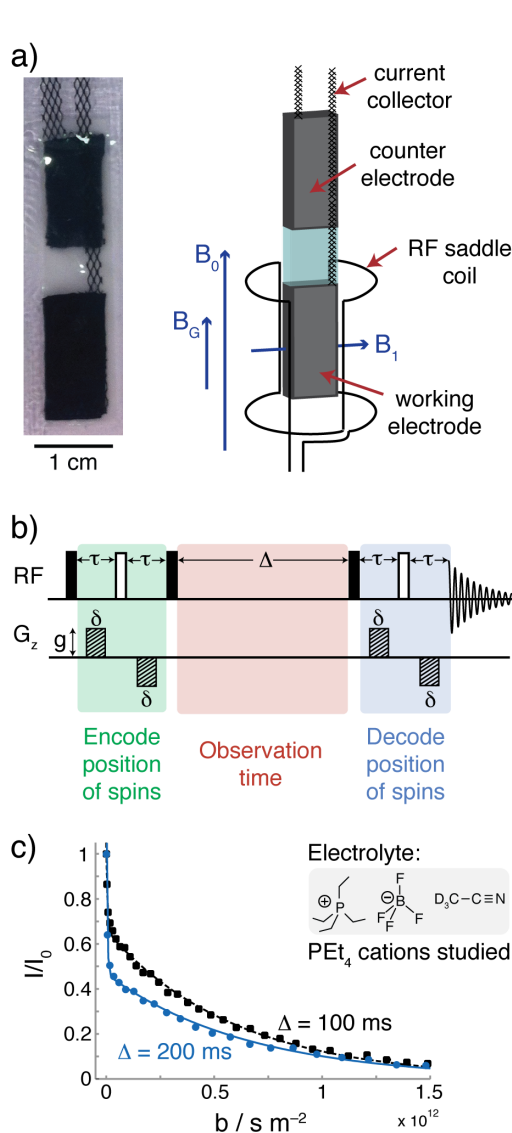


Figure 1: Introducing Pulsed Field Gradient NMR of Supercapacitors. a)

Photograph of pouch cell, and schematic of *in situ* PFG set-up showing the orientation of the magnetic field, B_0 , the magnetic field gradient, B_G , and the radiofrequency (RF) field, B_1 (see Supplementary Discussion of field orientations). **b)** 13-interval sequence used for PFG NMR experiments, RF and magnetic field gradient (G_z) pulses are shown. **c)** Example decays of normalised in-pore cation signal intensity against b for a YP50F supercapacitor (cell voltage = 0 V) with 1.5 M PEt_4BF_4 in deuterated acetonitrile electrolyte (see inset for structural formula). The parameter b is given by: $b = g^2 (\gamma 2\delta)^2 t_{\text{eff}} = g^2 (\gamma 2\delta)^2 (\Delta - 1/6 \delta - 1/2 \tau)$, with g the strength of the magnetic field gradient, γ the gyromagnetic ratio of the nucleus under study, δ the duration of an individual magnetic field gradient pulse, Δ the observation time, τ the separation of the 90 and 180° RF pulses. The time constants of the biexponential fits give self-diffusion coefficients, D .

Two sets of pouch cell supercapacitors²⁵ constructed with electrodes fabricated from two different commercial activated carbons (YP50F and YP80F) with different porosities and 1.5 M tetraethylphosphonium tetrafluoroborate (PEt_4BF_4) electrolyte (in deuterated acetonitrile, D_3CCN) were studied by PFG NMR. In the experiments the cell working-electrode is detected (Figure 1a). Uniquely, NMR allows the separate signals from ions that are inside the carbon nanopores, “in-pore”, and ions that are outside the nanopores in electrolyte reservoirs, “ex-pore”, to be resolved.^{19,26,27} Cations and anions were studied using ^1H and ^{19}F NMR, respectively (see Supplementary Figure 1 for example spectra).

In PFG NMR experiments magnetic field gradient pulses are used to encode and decode the positions of the nuclear spins, with ionic diffusion probed over a given observation time, Δ (Figure 1b). For each diffusion measurement, a series of NMR spectra are acquired with different magnetic field gradient strengths, g , and the normalised signal intensity (I/I_0) is plotted against a parameter, b , that is proportional to g^2 (see Figure 1 caption).^{28,29} Self-diffusion coefficients, D , are then obtained by exponential fits of the form $(I/I_0) = \exp(-D b)$. For in-pore ions biexponential decays were observed (Figure 1c), indicating the presence of two sets of in-pore ions with different D values. The initial steep decay (large D) arises from ions that undergo in-pore \leftrightarrow ex-pore exchange processes during the observation time (see Supplementary Figure 2). 2D exchange NMR measurements confirmed the presence of such exchange (Supplementary Figure 3). The more gradually decaying component in Figure 1c (small D), on the other hand,

arises from ions that remain in the nanopores during the observation time, with a self-diffusion coefficient, $D_{\text{in-pore}}$, and mole fraction, $A_{\text{in-pore}}$. By performing PFG NMR measurements with different Δ times, in-pore diffusion and in-pore \leftrightarrow ex-pore exchange processes can be probed on different timescales (and therefore different lengthscales).

Ion Dynamics in Supercapacitors at 0 V

For the YP50F supercapacitor at 0 V, $D_{\text{in-pore}}$ is reduced by over two orders of magnitude compared to neat electrolyte, for all observation times, Δ (Figure 2a). For example, for a relatively short observation time of $\Delta = 20$ ms, $D_{\text{in-pore}}$ is reduced by a factor of 300 for cations and 180 for anions. These marked reductions could arise from a number of factors: (i) the local reduction in in-pore diffusion due to collisions with the rigid pore walls,^{30,31} (ii) tortuosity arising from the disordered arrangement of pores³²⁻³⁴ in the carbon particles such that ions diffuse in an indirect way, and (iii) the structure and composition of the electrolyte in the pores may differ from neat electrolyte, our previous studies having shown that ions are partially desolvated in the pores.⁶ The difference in diffusion coefficients between anions and cations in neat electrolyte is amplified upon confinement in the pores (Figure 2a), highlighting the important role of ion size, with the larger cations presumably taking more indirect pathways through the pore network (as the smallest pores are inaccessible).³⁰ For a second activated carbon, YP80F, $D_{\text{in-pore}}$ for cations is ~ 4 to 5 times larger compared to YP50F depending on the observation time, showing the crucial influence of carbon structure on in-pore diffusion. This increase in $D_{\text{in-pore}}$ is ascribed to the additional porosity with pore widths between 1 and 3 nm in YP80F (Figure 2b), with these large pores acting as “highways” connecting the smaller pores. Indeed, a previous study showed that the diffusion coefficients of small molecules increased with the pore width of activated carbons.³⁰

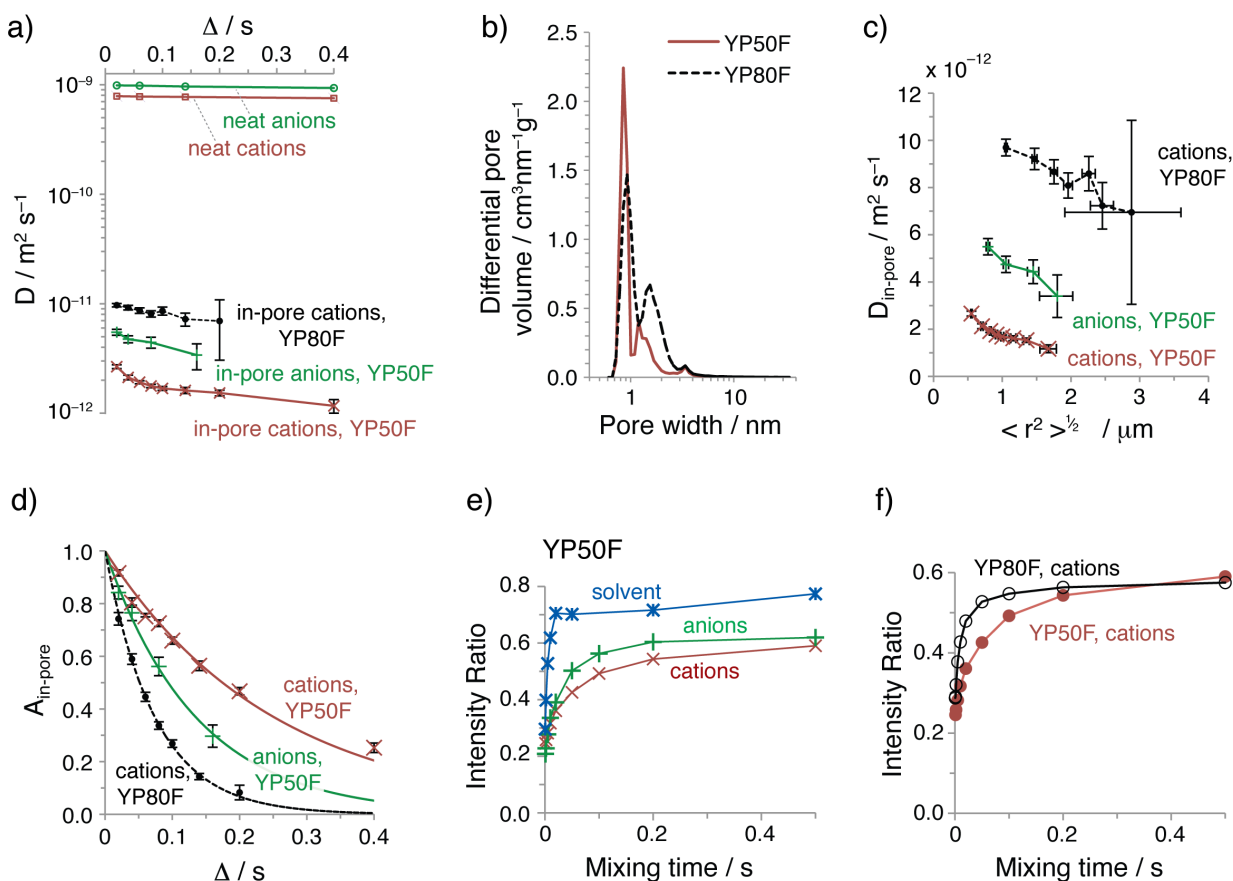


Figure 2: Diffusion measurements on supercapacitors at 0 V. **a)** Diffusion measurements as a function of Δ , for YP50F and YP80F supercapacitors, with 1.5 M PEt_4BF_4 in deuterated acetonitrile electrolyte. Measurements on neat electrolyte are also shown. **b)** Pore size distributions of the two carbons. **c)** $D_{\text{in-pore}}$ as a function of root mean square displacement. **d)** $A_{\text{in-pore}}$ as a function of Δ . Exponential fits allow the mean lifetime of an in-pore ion, $\tau_{\text{in-pore}}$, to be determined. **e)** Intensity ratios from 2D exchange NMR spectroscopy experiments with different mixing times on a YP50F supercapacitor. ^1H , ^{19}F and ^{13}C NMR experiments were used to probe cations, anions and isotopically labelled solvent, respectively. Representative spectra are given in Supplementary Figure 3 and details of intensity ratio calculations are given in Methods. **f)** As in e), but comparing cations in YP50F and YP80F. Error bars in a), c) and d) represent 95% confidence bounds (see Methods). In c) errors in the root mean square displacements arise from uncertainty in $D_{\text{in-pore}}$ values used in their calculation.

Our experiments measure in-pore diffusion on relatively long timescales (20–400 ms), with corresponding lengthscales of several hundreds of nanometres to micrometres (Figure 2c). Root mean square displacements, $\langle r^2 \rangle^{1/2}$, were determined using $D_{\text{in-pore}}$ values measured at different Δ times; $\langle r^2 \rangle^{1/2} = (6D_{\text{in-pore}}t_{\text{eff}})^{1/2}$, with $t_{\text{eff}} = \Delta - 1/6 \delta - 1/2 \tau \approx \Delta$. $D_{\text{in-pore}}$ decreases as larger root mean square displacements, $\langle r^2 \rangle^{1/2}$, are probed (Figure 2c), indicative of restricted diffusion,²⁸ with ions becoming more likely to encounter the smallest pores and bottlenecks for longer observation times (longer $\langle r^2 \rangle^{1/2}$). This shows that the structure of the carbon particles is heterogeneous to ionic motion on lengthscales as large as 2 microns, *i.e.* on the order of the carbon particle sizes (see Supplementary Figure 4). This is consistent with a recent NMR study of solvents in a nanoporous carbon,³⁵ and helps to explain why we observe greater reductions of effective diffusion upon confinement than in MD simulations where much smaller observation times were studied.²³

PFG NMR experiments provide additional dynamic information about the exchange (or swapping) of ions between in- and ex-pore environments (Figure 2d). $A_{\text{in-pore}}$ gives the mole fraction of in-pore ions that have remained in the nanopores during the observation time, Δ . For YP50F, the anions undergo significantly more rapid exchange than the cations, with mean in-pore lifetimes, $\tau_{\text{in-pore}}$,^{29,36} of 135 ± 14 and 252 ± 26 ms, respectively. This suggests that in-pore diffusion limits in-pore \leftrightarrow ex-pore exchange, with the faster anion diffusion resulting in more rapid exchange. Indeed, the larger in-pore cation diffusion coefficient in YP80F electrodes results in even more rapid exchange, with $\tau_{\text{in-pore}} = 74 \pm 3$ ms. A measure of the mean carbon particle radius is obtained by $(6D_{\text{in-pore}}\tau_{\text{in-pore}})^{1/2}$, giving $1.5 \mu\text{m}$ and $1.7 \mu\text{m}$, for YP50F (probed by the cations and anions, respectively) and $1.9 \mu\text{m}$ for YP80F (cations), with the values consistent with typical carbon particle dimensions (see Supplementary Figure 4). The build up of cross peak intensity (plotted as an intensity ratio) in 2D exchange NMR experiments offer a second probe of in-pore \leftrightarrow ex-pore exchange (Figure 2e),^{7,9} and corroborate the $\tau_{\text{in-pore}}$ values obtained from PFG NMR. The steeper build-up of cross peak intensity with mixing time for anions, compared to cations, in YP50F is consistent with their more rapid in-pore \leftrightarrow ex-pore exchange (Figure 2e), while the cations exchange faster in YP80F than in YP50F (Figure 2f). The exchange of the smaller, neutral, acetonitrile solvent molecules (as probed by ^{13}C NMR experiments on isotopically enriched acetonitrile) is more rapid than for the ions, indicating that their in-pore diffusion coefficients are larger (Figure 2e). This is unsurprising since diffusion measurements on neat electrolyte showed that the solvent diffuses 2.0 times faster than the anions, and 2.5 times faster than the cations. MD simulations have also shown that (de)solvation of in-pore ions with acetonitrile takes place as quickly as ~ 1 ps,²³ and a recent study of dimethyl carbonate showed a relatively modest reduction of effective diffusion (by a factor of ~ 0.2) upon confinement in a nanoporous carbon.³⁵ Overall, the slower diffusing ions move through a highly dynamic solvent medium in the nanopores.

***In situ* Measurements of Ion Dynamics**

In situ pulsed field gradient NMR measurements allow the direct measurement of in-pore diffusion in supercapacitors during charging (Figure 3a, b). Measurements were made under equilibrium conditions at different fixed cell voltages. For negative polarisations, marked decreases of $D_{\text{in-pore}}$ are observed for anions and cations, while for positive polarisations only minor changes are observed, with only slight increases in $D_{\text{in-pore}}$ for anions. For the 1.5 M electrolyte, comparing between 0 and -1.5 V, $D_{\text{in-pore}}$ decreases by a factor of 3.6 for the anions (averaged across observation times, see Supplementary Figures 5 and 6), while the cations show a similar, but slightly smaller, reduction factor of 2.9.

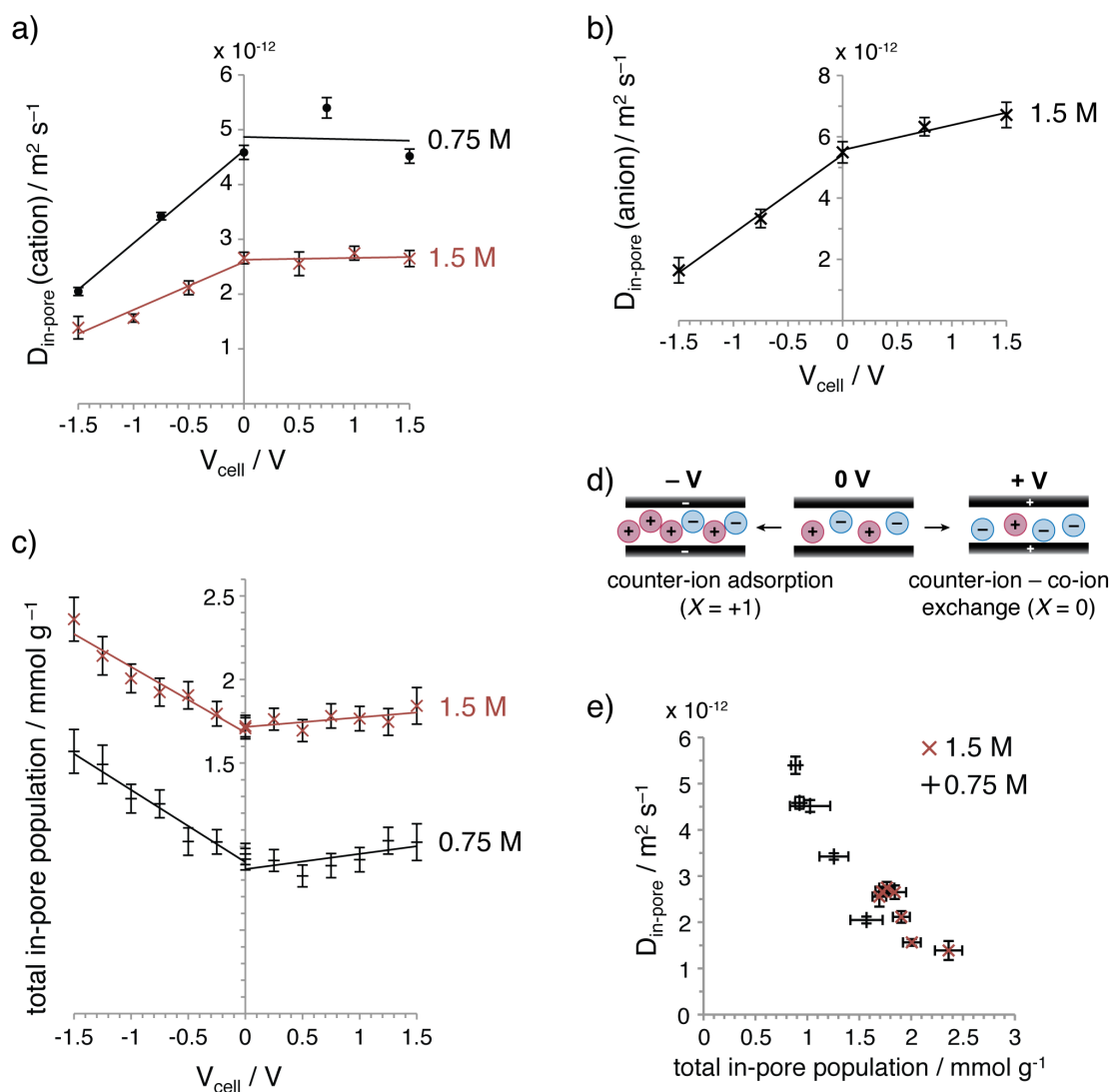


Figure 3: *In situ* measurements of ionic diffusion in supercapacitors. **a)** Cation and **b)** anion diffusion in YP50F – PEt₄BF₄ (1.5 M) in acetonitrile supercapacitors, at different cell voltages. In a) data for a supercapacitor with 0.75 M electrolyte are also shown **c)** Total in-pore ion populations at different potentials, determined from data in Ref. [5]. **d)** Schematic of charge storage mechanism. At 0 V there are in equal number of cations and anions in the pores of the carbon electrodes. For positive potentials charging occurs by the exchange of anions and cations, while for negative potentials charging occurs by cation adsorption. **e)** Correlation between $D_{\text{in-pore}}$ (cations) and total in-pore ion population. All data are for $\Delta = 20$ ms (data recorded with different Δ values show very similar trends, see Supplementary Figures 5 and 6). Error bars in a), b) and d) represent 95% confidence bounds (see Methods). Error bars in c) represent the range of values obtained from four independent fits of each dataset (see Ref. 5).

The variations of $D_{\text{in-pore}}$ with voltage may be understood by examining the changes of in-pore ionic populations at different voltages (Figure 3c). We previously showed that charging in the positive electrode takes place by swapping of anions and cations (counter-ion – co-ion exchange mechanism, $X = 0$), while charging in the negative electrode occurs by counter-ion adsorption ($X = +1$, see Figure 3d).⁵ As a result, for positive polarisations the total in-pore ion population shows only minor changes with voltage, while for negative polarisation a marked increase is observed (Figure 3c). This increase of in-pore ion population correlates with the reduction of $D_{\text{in-pore}}$ for both ions: i.e., it appears that increasing ion-ion interactions lead to a reduction of $D_{\text{in-pore}}$ with voltage for negative polarisations. In principle ion-carbon interactions may also influence $D_{\text{in-pore}}$,²³ though

these appear to be a minor factor, since the anions and cations show such similar variations of $D_{\text{in-pore}}$ with voltage.

To explore the role of ion-ion interactions further, a lower concentration electrolyte (0.75 M) was studied (Figure 3a, c). At 0 V, in-pore cation diffusion is faster by a factor of ~ 1.7 for the 0.75 M electrolyte (compared to 1.5 M), and in-pore \leftrightarrow ex-pore exchange is also more rapid, with $\tau_{\text{in-pore}} = 160 \pm 7$ ms, compared to 252 ± 26 ms for 1.5 M (see Supplementary Figure 7). Measurements at different voltages show that in-pore cation diffusion is consistently faster for the 0.75 M electrolyte. This can be rationalised by the smaller total in-pore ion populations for the 0.75 M electrolyte (Figure 3c), and the resulting reduction in ion-ion interactions. The changes of $D_{\text{in-pore}}$ with voltage are very similar for the two concentrations (Figure 3a), which is due to the similar *changes* of in-pore population with voltage (the charging mechanisms are very similar for the two concentrations).⁵ A global plot for the cations in Figure 3e further highlights the inverse correlation between $D_{\text{in-pore}}$ and the total in-pore ion population. It should be noted that the variations of in-pore anion:cation ratios with voltage are different for the two concentrations, so we do not expect the 1.5 and 0.75 M datasets to overlap exactly in Figure 3e. We finally note that the gravimetric electrode capacitances (see Supplementary Table 1) are 107 and 106 F g⁻¹ for the 1.5 and 0.75 M cells, respectively. These are identical within error such that the variation of charge with voltage is very similar for the two cells.

Discussion

Our *in situ* results strongly suggest that changes of in-pore ion populations, and the resulting changes in the number of ion-ion interactions, dominate the variations of in-pore ionic diffusion observed during charging. The ion-ion interactions have at least two contributions: (i) steric effects, with the increased number of ions sterically hindering diffusion, (ii) electrostatic effects, whereby associations between oppositely charged ions can reduce ionic self-diffusion.³⁷⁻³⁹ The latter effect will vary in a complex way as the proportions of anions and cations vary during supercapacitor charging, and may explain the small increases in $D_{\text{in-pore}}$ for the anions, observed for positive polarisations. The apparent dominance of ion-ion interactions on the voltage dependence of $D_{\text{in-pore}}$ is consistent with a recent MD study of an ionic liquid inside slit-pore electrodes,¹⁸ as well as a further MD study on supercapacitors with disordered carbon electrodes.⁴⁰

The dependence of $D_{\text{in-pore}}$ on the in-pore ion population has important consequences for supercapacitor charging rates. Depending on the mechanism that operates (ion exchange, counterion adsorption, co-ion desorption etc.) the charging dynamics will vary. This offers new opportunities to improve charging dynamics in supercapacitors if the charging mechanism (X value)

can be controlled. If $X > 0$, then the in-pore ion population will increase with voltage and $D_{\text{in-pore}}$ is expected to decrease. On the other hand, if $X < 0$ (net ion desorption), $D_{\text{in-pore}}$ would be expected to increase with voltage. For the case of initially filled pores, charging purely by co-ion desorption ($X = -1$) may offer the fastest charging rates at high voltages, though this mechanism has not yet been observed.²¹ Kondrat *et al.* have suggested that ideally the carbon pores would be initially empty of ions (ionophobic pores) to facilitate fast charging,^{16,24,41} though in practice ionophobic pores have not yet been realised experimentally. Ionic liquids with activated carbons represent the opposite extreme as the pores are densely packed with ions in the absence of an applied potential,^{6,15} and our results help to explain their poor power performances.

Our measurements on carbons with different pore sizes help to explain previous electrochemical studies of charging dynamics.^{1,42–44} For carbons with smaller average pore sizes and fewer mesopores (< 2 nm), the capacitance decreased more rapidly as the current density was increased. We postulate that this is because slower in-pore transport in the smaller pores results in non-equilibrium charge storage mechanisms when practical current densities are used, whereby the optimum (lowest voltage) configurations of ions cannot be attained. As a result, the voltage shows a larger variation with charge, and the capacitance decreases. Our findings also help to explain the excellent high-current performances of hierarchical meso/micropore structures⁴² and zeolite-templated carbons with pore connectivity in three dimensions.⁴⁵

Conclusion

In conclusion, our *in situ* PFG NMR experiments reveal new insights into ionic diffusion and exchange processes in supercapacitors. Confinement of the ions inside YP50F nanoporous electrodes leads to a large (more than two orders of magnitude) reduction in their effective diffusion coefficients. The important role of ionic size has been demonstrated, with the larger cations showing the largest reductions of diffusion upon confinement. Our measurements also allow the mean lifetime of in-pore ions to be determined ($\tau_{\text{in-pore}}$), and indicate that in-pore diffusion limits the in-pore \leftrightarrow ex-pore exchange process. We have also shown that the structure of the carbon particles is heterogeneous to ionic motion on lengthscales as large as microns.

Going further, the nanoporous structure of the electrode has a significant effect on in-pore ion diffusion and exchange, with the incorporation of a relatively small amount of additional pore volume in the 1 – 3 nm range resulting in a large increase of $D_{\text{in-pore}}$ and decreases of $\tau_{\text{in-pore}}$. Further PFG NMR measurements on a range of different porous carbons will further our understanding and could enable precise control over power and energy performance.

In situ diffusion measurements at different electrolyte concentrations have shown that ion-ion interactions dominate the observed variations of $D_{\text{in-pore}}$ with voltage, with the different charging mechanisms (X values) in the two electrodes resulting in drastically different dynamic behaviour. Charging by counter ion adsorption ($X = +1$) results in decreases of $D_{\text{in-pore}}$ for negative polarisations, while charging by counter-ion – co-ion exchange ($X = 0$) results in much smaller changes of $D_{\text{in-pore}}$ for positive polarisations. Our findings demonstrate the important interplay between charging mechanisms and dynamics, and should offer new opportunities to control the power performance of supercapacitors. One simple way to improve charging rates may be to use lower electrolyte concentrations as this results in faster in-pore diffusion. However, this must be balanced with the availability of electrolyte ions for charging processes, and efforts must also be made to probe any electrolyte concentration gradients⁴⁶ formed during charging. Achieving truly ionophobic pores, or alternatively charging by co-ion desorption, may enhance charging rates. Beyond controlling the charging mechanism, tuning the nature of the inter-ionic interactions⁴⁷ by an informed choice of electrolyte may also offer a route to improved charging rates. We note that our PFG NMR approach can readily be extended to explore the properties of a range of electrolytes including ionic liquids, mixed electrolytes, and those containing multivalent ions.⁴⁸ Computational screening approaches⁴⁹ could also help to identify new carbon-electrolyte combinations for experimental testing.

Our measurements have provided a wealth of information that can bring new opportunities for the design of enhanced supercapacitors. The findings will facilitate modelling of full supercapacitor devices, such that the interplay between structure and ionic transport can be understood across the full range of lengthscales. Finally, future measurements will be made to establish how ionic migration in an electric field differs from equilibrium self-diffusion,⁵⁰⁻⁵² such that the non-equilibrium states reached at the fastest charging/discharging rates can be probed. Only then will we begin to fully unravel the relationships between materials properties and performance in supercapacitors.

Methods

Materials

Freestanding carbon film electrodes (95 wt. % carbon, thickness ~250 μm) were prepared by mixing YP50F (or YP80F) activated carbon (Kuraray Chemical) with polytetrafluoroethylene (Sigma Aldrich, 60 wt. % dispersion in water), as previously described.¹⁹ All supercapacitors used either 1.5 or 0.75 M tetraethylphosphonium tetrafluoroborate (>97% purity, Toyko Chemical

Industries) in deuterated acetonitrile (D_3CCN , 99.8 atom %, Euroisotop) as the electrolyte, with the exception of carbon NMR experiments where $H_3C^{13}CN$ (99 atom %, Isotech) was used.

Fabrication of Supercapacitors

Film electrodes were dried for at least 15 hours at 200 °C *in vacuo*, with supercapacitor devices then prepared in an Ar glovebox (<0.1 ppm H_2O and O_2) in plastic pouch cells, with a design adapted from our previous work.²⁵ These pouch cells were positioned inside either 10 mm or 5 mm NMR tubes for PFG NMR experiments. Electrode masses were equal within 0.2 mg for all cells, and were typically ~20 mg for cells in 10 mm NMR tubes, and ~7 mg for cells in 5 mm tubes.

Pore Size Distributions

Carbon pore size distributions were obtained by analysing N_2 gas sorption isotherms (Supplementary Figure 8) recorded on activated carbon powders at 77 K, using a Quantachrome NOVA 2200e surface and pore volume analyser. Specifically, quenched solid density functional theory (QSDFT) analysis was used,⁵³ assuming slit-shaped carbon pores, and with pore sizes up to 34 nm considered in the analysis.

NMR Measurements

PFG NMR experiments were carried out using a Bruker Avance spectrometer and a Diff-50 probehead at a magnetic field strength of 7.1 T (1H larmor frequency at 300 MHz). Saddle coils (10 mm inner diameter for 1H , 5 mm for ^{19}F and ^{13}C) were used for radiofrequency excitation and detection, with the working electrode positioned in the detection area, with the B_1 field and electrode orientated as indicated in Figure 1a. The 13-interval pulse sequence⁵⁴ (Figure 1b) with bipolar pulsed field gradients was used for diffusion measurements to minimise the effects of any magnetic field gradients inherent to the sample. The maximum g used was $17\ T\cdot m^{-1}$, and the effective gradient pulse duration, δ , (as defined in Figure 1b) was 0.5 ms. Sufficiently long delays of 0.5 ms were inserted between PFG and RF pulses such that eddy currents did not influence diffusion measurements. Experiments on neat electrolytes were performed in 5 mm NMR tubes. The *in situ* PFG NMR methodology is similar to our previous *in situ* NMR studies of supercapacitors, with spectra acquired under equilibrium conditions during the application of constant cell voltages.²⁵ For determination of D and A values, in-pore intensities were first determined by integration in Topspin software and the variation of intensity with the parameter b was analysed using a Matlab script and the curve-fitting tool. Plotted error bars in the diffusion parameters represent upper and lower bounds at 95% confidence limits, as determined using the curve-fitting tool. 2D exchange spectroscopy measurements were performed as in our previous

work,⁷ with a series of measurements made with different mixing times. We consider the “intensity ratio” which is given by the cross peak intensity, normalised by the sum of the cross and diagonal peaks (these peaks are highlighted in the spectra in Supplementary Figure 3). Only these resonances were considered in the analysis as the other resonances are considerably broadened by artefacts arising from truncation of the free induction decay in the indirect dimension. Spectra are referenced as; ¹H - methyl group of ethanol at 1.2 ppm, ¹⁹F – hexafluorobenzene at –164.9 ppm, and ¹³C – ethylene group of ethanol at 58.0 ppm. RF excitation pulses were calibrated on the supercapacitor cells, with powers of between 22-30 kHz used for all nuclei. Recycle delays were set to ~1.2 T_1 to maximise the signal to noise for a given experiment time.

Determination of cell capacitance

The capacitances of *in situ* cells (listed in Supplementary Table 1) were determined by integration of the current, I , vs. time, t , data upon discharge from a cell voltage of 1.5 V to 0 V. Cell capacitances, C_{cell} , were first determined as:

$$C_{cell} = \frac{\int Idt}{\Delta V} \quad [1]$$

with $\Delta V = -1.5$ V. Gravimetric electrode capacitances, $C_{elec.}$, were then calculated as:

$$C_{elec.} = \frac{2C_{cell}}{m_{elec.}} \quad [2]$$

where $m_{elec.}$ is the mass of activated carbon in a single electrode (which is 95% of the total mass of a single electrode). For cyclic voltammograms, see Supplementary Figure 9.

Scanning electronic microscopy

Scanning electron microscopy (SEM) images (Supplementary Figure 4) were recorded on carbon film pieces using a FEI Magellan 400L field emission scanning electron microscope operated at 15 kV.

The data that support the plots within this paper and other findings of this study are available from the corresponding author upon reasonable request.

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Contributions

A.C.F prepared supercapacitor cells, performed PFG NMR experiments and analysed the data. J.M.G. recorded and analysed NMR experiments to determine total in-pore ion populations. N.M.T. performed initial PFG NMR experiments. A.-R.O.R. recorded SEM images and gas sorption isotherms. All authors contributed to design of the research, the discussion of the data and the writing of the paper.

Additional Information

Supplementary information is available in the online version of the paper. Reprints and permissions information is available online at www.nature.com/reprints. Correspondence and requests for materials should be addressed to C.P.G.

Competing financial interests

The authors declare no competing financial interests.