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A scanning probe microscopy study of nanostructured TiO₂/poly(3-hexylthiophene) hybrid heterojunctions for photovoltaic applications

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Abstract

The nanoscale morphology of photoactive hybrid heterojunctions plays a key role in the performances of hybrid solar cells. In this work, the heterojunctions consist of a nanocolumnar TiO₂ surface covalently grafted with a monolayer of poly(3-hexylthiophene) (P3HT) functionalized with carboxylic groups (–COOH). Through a joint analysis of the photovoltaic properties at the nanoscale by photoconductive-AFM (PC-AFM) and surface photovoltage imaging, we investigated the physical mechanisms taking place locally during the photovoltaic process and the correlation to the nanoscale morphology. A down-shift of the vacuum level of the TiO₂ surface upon grafting was measured by Kelvin probe force microscopy (KPFM), evidencing the formation of a dipole at the TiO₂/P3HT-COOH interface. Upon in situ illumination, a positive photovoltage was observed as a result of the accumulation of photogenerated holes in the P3HT layer. A positive photocurrent was recorded in PC-AFM measurements, whose spatial mapping was interpreted consistently with the corresponding KPFM analysis, offering a correlated analysis of interest from both a theoretical and material design perspective.

Introduction

Over the past decades, a large range of photovoltaic (PV) technologies have been developed for the production of renewable energy [1]. Inorganic photovoltaic cells are currently the most employed PV devices with a power efficiency ranging from 20 to 40% [2] and a long-term stability up to 20 years [3]. However, a number of drawbacks affect those technologies. Indeed, in addition to high energy consumption for their fabrication, these devices are deposited on rigid substrates and involve relatively heavy and costly materials of possibly low abundance and/or toxicity [4]. New PV technologies, such as organic photovoltaics (OPV) and hybrid solar cells, are now being developed [2] to cope with such issues. In particular, hybrid solar cells can possibly benefit from the low economic and energy costs of production, high absorbance and tailorable absorption spectrum of the organic materials on the one hand, and from the good stability, absorption and electrical properties of the inorganic materials on the other hand.

Hybrid PV devices include various technologies such as perovskite cells, dye sensitized solar cells (DSSC), with power efficiencies up to 13% [5] and hybrid bulk heterojunctions (HBHJ), which combine an organic matrix and inorganic semiconducting nanostructures such as quantum dots. Among the electron acceptor materials commonly used for DSSC and HBHJ, titanium dioxide (TiO₂) is a well-known metal oxide semiconductor [6-8]. Depending on its nanostructure and its crystalline phase, its conductivity varies from 10⁻⁴ Ω⁻¹·cm⁻¹ to 10⁻¹¹ Ω⁻¹·cm⁻¹ [9,10]. TiO₂ is very valuable because it can easily form nanostructures, such as nanoporous layers, nanowires or nanocolumns [5,11,12]. Because of its large band gap (3.2 eV [13]), light absorption is carried out by an organic or inorganic dye. The nanostructuring of the acceptor material is crucial for the cell performance [11], as it allows increasing the specific surface of the layer to enhance the amount of grafted dye, and thereby, the photon absorption yield. Nanostructuring is also likely to improve the conductivity of TiO₂ [14]. Because of the influence of the nanostructuring of TiO₂ on the optoelectronic properties of the device, it is of prime interest to study the photovoltaic properties at the nanoscale. Hybrid heterojunction (HHJ) structures are obtained by impregnation of the porous layer with an absorbing dye or a polymer electron donor. Poly(3-hexylthiophene) (P3HT) is often used, because of its strong absorption, its high hole mobility and its donor-like electronic properties [15]. Upon light absorption by the polymer, excitons are generated and they can be dissociated at the interface with TiO₂, the polymer also acting as the hole-transporting layer.

In this work, we investigated nanostructured TiO₂ layers composed of arrays of nanoscale columns, covalently sensi-

tized with a P3HT-COOH monolayer to form hybrid bulk heterojunctions. The grafting of P3HT on the surface of TiO₂, ensured by the COOH groups, was demonstrated to be beneficial for the photoconversion efficiency of the system [16-18]. The vertically aligned nanostructuring of TiO₂ also makes this system attractive, since it ensures direct percolation paths for the photogenerated electrons from the donor-acceptor interface to the cathode, while providing a simple, controlled and ordered architecture. Furthermore, studies are available in literature regarding the photovoltaic response of TiO₂/P3HT blends [16-23] and can be used as a reference for meaningful interpretations of our measurements, both in terms of photocurrent and photovoltage under illumination. The columnar TiO₂/P3HT-COOH HHJs have been studied by photoconductive-AFM (PC-AFM) and photo-assisted Kelvin probe force microscopy (photo-KPFM) to follow the photovoltaic response, i.e., photocurrent and photovoltage, respectively, at the nanoscale under illumination, in order to understand the local physical processes taking place during the photoconversion of energy, and their correlation with the nanoscale morphology of the active layer. A key aspect of this work consists in the joint analysis of these correlated PC-AFM and KPFM measurements, providing a more fundamental understanding of the photovoltaic mechanisms at stake in the systems. To the best of our knowledge, this joint KPFM/PC-AFM study of such a nanostructured array of TiO₂ columns sensitized with functionalized P3HT-COOH constitutes a novel result of interest from both a theoretical and material design perspectives.

Materials and Methods

The TiO₂ layers were synthesized by magnetron sputtering in grazing mode. A thorough description of the fabrication process can be found in the literature [24], which also identified the optimized fabrication parameters for prospective photovoltaic applications. In compliance with these recommendations, the layers were synthesized without any substrate rotation or bias, while fixing the growth temperature to 450 °C and the tilt angle between the substrate and the cathode axis to 60°. Anatase TiO₂ layers with a 200 nm thick nanocolumnar morphology have been deposited on 85 nm-thick ITO-coated glass substrates (Naranjo B.V., sheet resistance of 15 Ω·sq). The average spacing between the columns is (10 ± 3) nm, with an average width of the columns of (19 ± 4) nm, as determined by SEM measurements [24]. The topography of the deposit is shown in the tapping-mode atomic force microscopy (TM-AFM) image of Figure 1d, where the apex of the columns appears as hemispherical protuberances. Regio-regular P3HT-COOH (5400 g/mol, which corresponds to about 30 monomer units, i.e., a total polymer chain length around 130 Å) was synthesized following a reported procedure [24]. A schematic descrip-

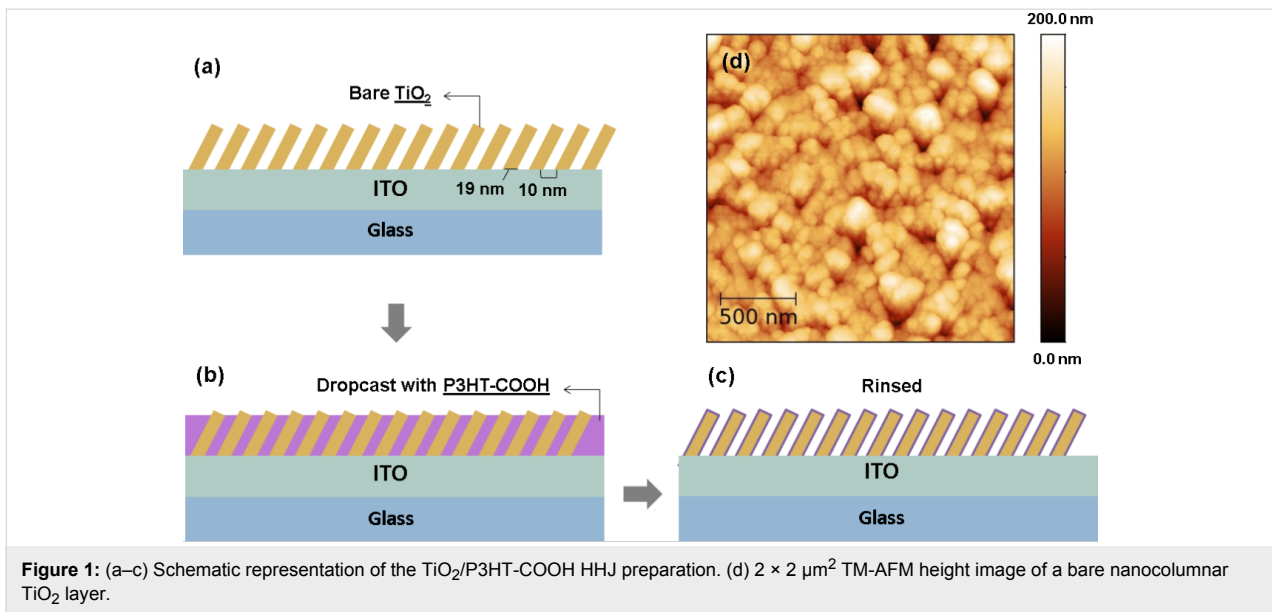


Figure 1: (a–c) Schematic representation of the TiO₂/P3HT-COOH HHJ preparation. (d) 2 × 2 μm² TM-AFM height image of a bare nanocolumnar TiO₂ layer.

tion of the grafting protocol is given in Figure 1a–c. The polymer deposit was obtained by dropcasting a 0.5 mg/mL solution of P3HT-COOH in chlorobenzene on the TiO₂ structure. The covalent grafting of the polymer on the nanoporous TiO₂ surface is ensured by the carboxylic –COOH group. Rinsing with chlorobenzene was then carried out to remove the residual ungrafted polymer chains. The success of the polymer grafting is confirmed by UV–visible optical absorption measurement across a 350–800 nm wavelength range, for which an absorption of light higher by one order of magnitude compared to bare TiO₂ was measured [24]. This indicates a good P3HT impregnation along the columns, the interspacing being sufficient for the polymer infiltration.

The photo-KPFM measurements were carried out in a UHV (<10^{−10} Torr) instrument composed of an Omicron Nanotechnology VT-AFM system with a Nanonis controller. The KPFM electrical excitation used a frequency $\omega_{\text{KPFM}}/2\pi$ of 958 Hz, with a VAC amplitude of 600 mV. The light source for sample irradiation was a green laser diode (wavelength = 500 nm, power density = 1.45 mW/mm²). Photo-assisted KPFM measurements were also performed in ambient conditions, with a Bruker multimode microscope controlled by a Nanoscope III unit coupled to a Nanonis control unit (SPECS Zürich). The KPFM electrical excitation was made at a frequency $\omega_{\text{KPFM}}/2\pi$ of 80 Hz, with a VAC amplitude of 500 mV. The illumination of the sample for photo-KPFM and photovoltage probing was provided by a white light lamp irradiating the sample surface from the top. In both setups, conductive Nanosensors PPP-EFM tips (PtIr-coated Si probes) were used (resonant frequency around 75 kHz). The sample was grounded while the excitation and regulation biases were applied to the tip. The measured

contact potential difference (V_{cpd}) is given by the following expression:

$$|e| V_{\text{cpd}} = \Phi_{\text{tip}} - \Phi_{\text{sample}}, \quad (1)$$

where Φ_{tip} and Φ_{sample} are the workfunction of the tip and the sample, respectively. In this work, no calibration of the tip workfunction was necessary, as only the V_{cpd} variations between the materials constituting the photovoltaic blends and their modifications with incoming light were to be measured. These V_{cpd} variations provide relative but quantitative variations of surface potential at the investigated interfaces.

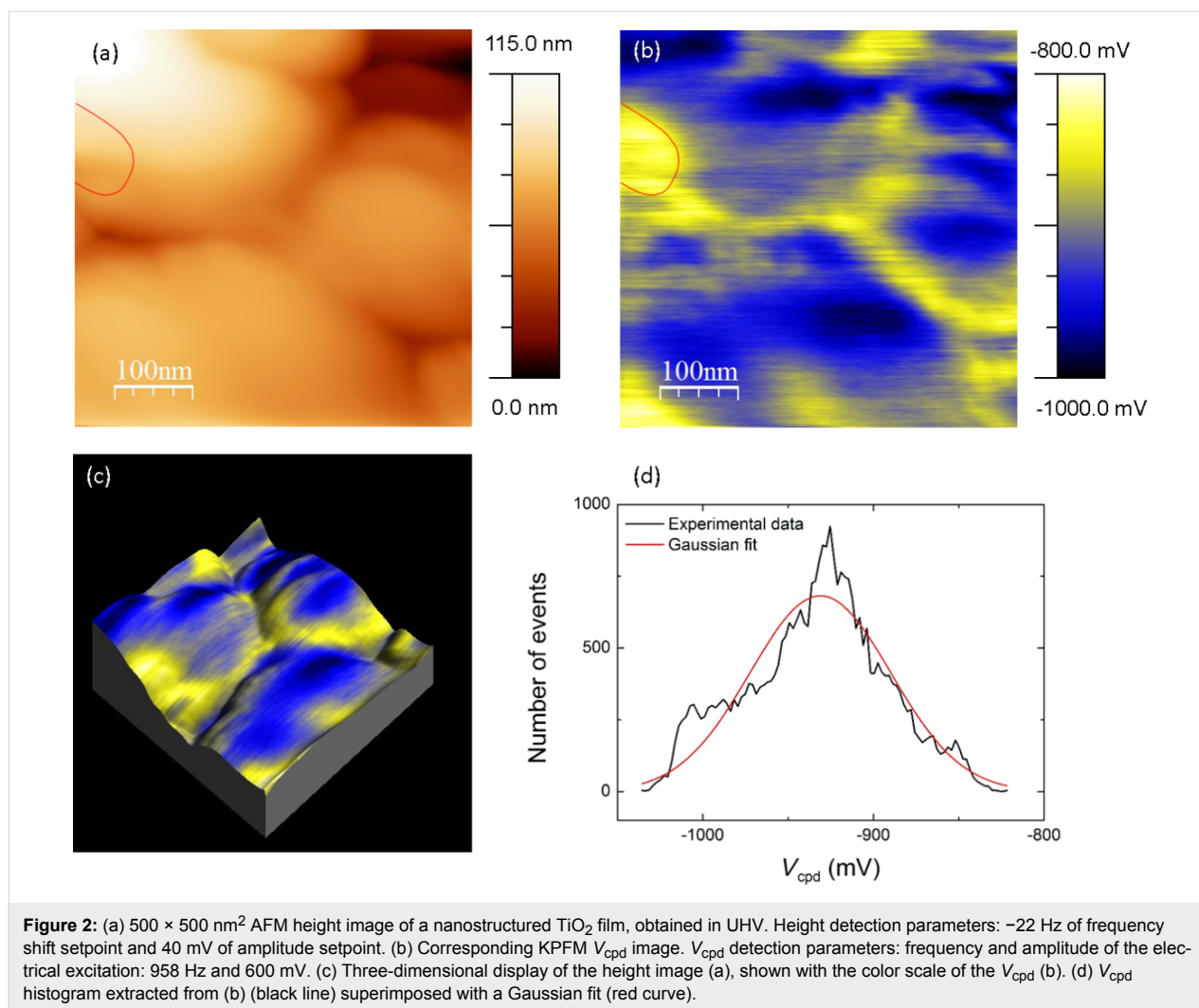
The PC-AFM measurements were carried out in air, using a Bruker Dimension Icon microscope with a Nanoscope V controller. An extended TUNA external module was used for current detection with a detection range within 100 fA to 1 μA. Silicon tips coated with a PtIr conductive alloy (PPP-CONTPt from Nanosensors) were used. The tip and the back-contact were connected while the sample was locally irradiated from the bottom (through the patterned ITO–glass substrates) under AM 1.5 calibrated white light illumination (spot diameter around 200 μm, power density of 100 suns).

Results and Discussion

Photo-KPFM measurements on the TiO₂/P3HT-COOH hybrid heterojunctions

Analysis of the V_{cpd} contrast in the dark

Figure 2a shows a 500 × 500 nm² AFM height image obtained in UHV on a nanocolumnar TiO₂ film deposited over a grounded ITO electrode, where the nanocolumns of TiO₂ are

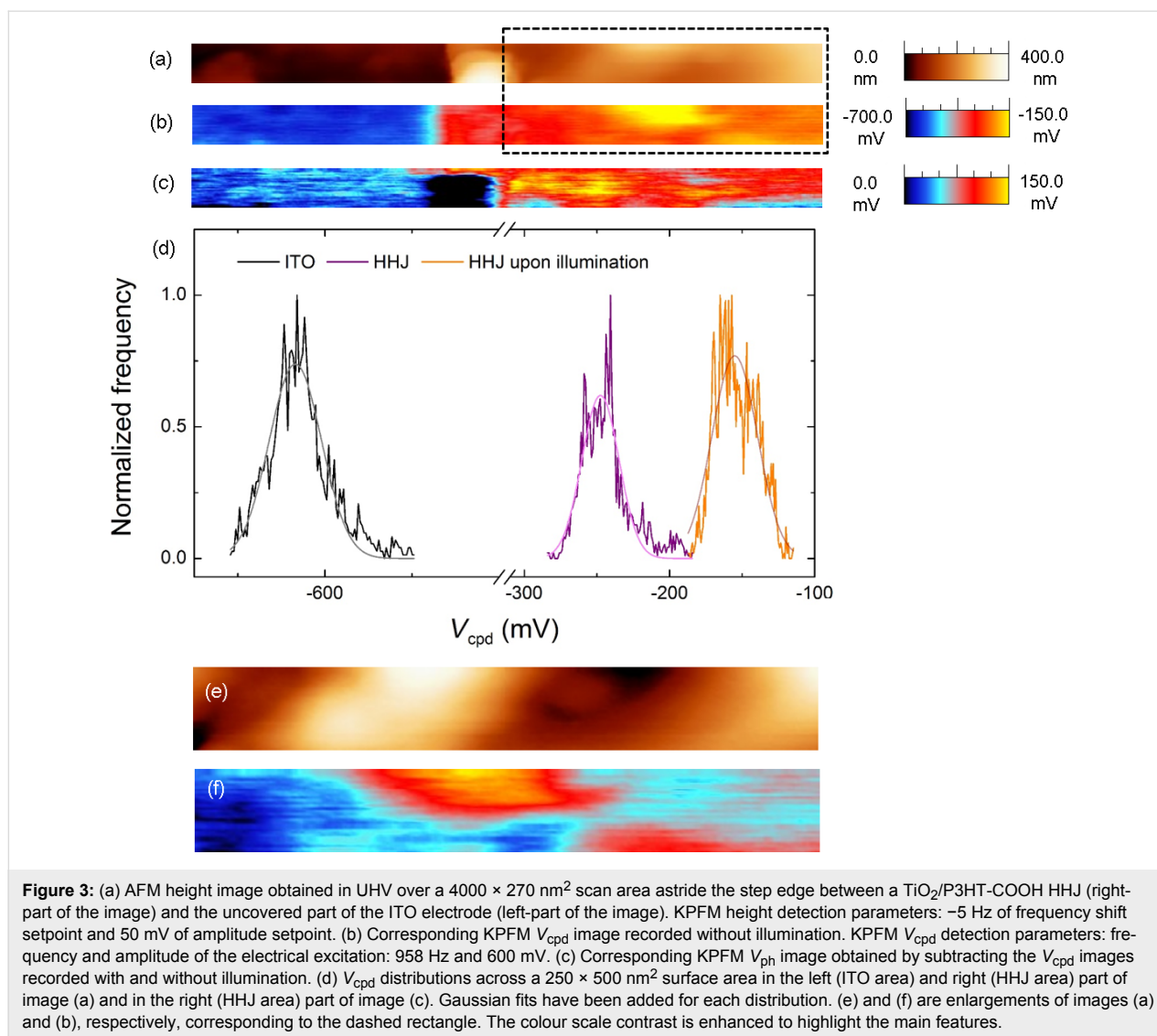


assembled in clumps with a width of several hundred nm. Figure 2b shows the corresponding KPFM V_{cpd} image. Figure 2c presents the three-dimensional display of Figure 2a, where the colour scale refers to the V_{cpd} signal of Figure 2b. The distribution of the V_{cpd} values can be fit with a Gaussian distribution centred at -931 mV with a FWHM of 97 mV (Figure 2d).

A direct correlation between the topography and the V_{cpd} signal can be observed, with a higher height corresponding to a more negative V_{cpd} . It is however unlikely that the contrast purely originates from a crosstalk between the topography and the V_{cpd} , as indicated by local mismatching between both contrasts (see red lines in Figure 2a and 2b). Moreover, further measurements (see Figure 3) showed that P3HT grafting barely affects the overall morphology but smooths tremendously the V_{cpd} contrast. Thus, the observed V_{cpd} contrast most probably originates therefore from local variations in the electronic properties of the surface, such as a possibly different free electron

density at the top and at the side of the columns. This explanation is further supported by the PC-AFM measurements presented in the last section. As shown in the Supporting Information File 1 (Figure S1), no ungrounded potential is to be detected at the top of the nanocolumnar TiO₂ film. This can therefore not be the origin of the V_{cpd} contrast observed on the bare TiO₂ columns.

Figure 3a displays a KPFM height image obtained in UHV on a TiO₂ deposit grafted with P3HT-COOH. The left part of the image corresponds to a bare area of the ITO electrode, while the right part shows a TiO₂/P3HT-COOH zone. Figure 3b shows the corresponding V_{cpd} image recorded in the dark. A clear difference between the V_{cpd} intensity over the ITO electrode (-614 ± 18 mV in average) and the TiO₂/P3HT-COOH HHJ (-248 ± 49 mV in average) is observed. This V_{cpd} shift clearly appears in the V_{cpd} distributions of Figure 3d. The more negative V_{cpd} value over the ITO electrode reflects consistently a higher corresponding work function (around 4.7 eV in litera-



ture [6,15]) compared to that of TiO_2 (around 4.3 eV in literature [7]).

The data of Figure 3 were compared with the images obtained on bare nanocolumnar TiO_2 (Figure 2). In both measurements, the ITO electrode was grounded and the same tip was used. The distribution of V_{cpd} values on the $\text{TiO}_2/\text{P3HT-COOH}$ area (right part of Figure 3b) is displayed as the purple curve in Figure 3d; the corresponding Gaussian fit is centred at -248 mV , with a FWHM of 30 mV . As seen in Figure 2d, the V_{cpd} is much more negative on bare nanocolumnar TiO_2 . This indicates that: (i) the P3HT layer induces an up-shift of the V_{cpd} values, and (ii) this up-shift occurs over the entire surface, since no values typical of bare TiO_2 are recorded on the polymer-grafted surface. This indicates that the P3HT covering is complete, with no bare TiO_2 area left. The fact that the V_{cpd} increases upon P3HT grafting indicates that the surface workfunction of $\text{TiO}_2/\text{P3HT-COOH}$ is

lower than that of bare nanocolumnar TiO_2 . This can be understood on the basis of the following discussion, which describes the relative configuration of the electronic levels of the materials within the HHJ.

The covalent bonding between P3HT-COOH and TiO_2 creates a dipole at the interface induced by: (i) the hybridization of the electronic orbitals of the two components, leading to a rearrangement of the charge density at the interface and (ii) the addition of a net dipole intrinsic to the P3HT-COOH molecule itself. The first effect was reported previously [25], evidencing a pinning of the LUMO of P3HT-COOH at the conduction band of the TiO_2 with a net transfer of half an electron per polymer chain from the LUMO of P3HT into the CB of TiO_2 . This results in the formation of a dipole at the P3HT-COOH/ TiO_2 interface, directed away from TiO_2 , where the positive (negative) pole is located in P3HT (TiO_2). Previous KPFM studies

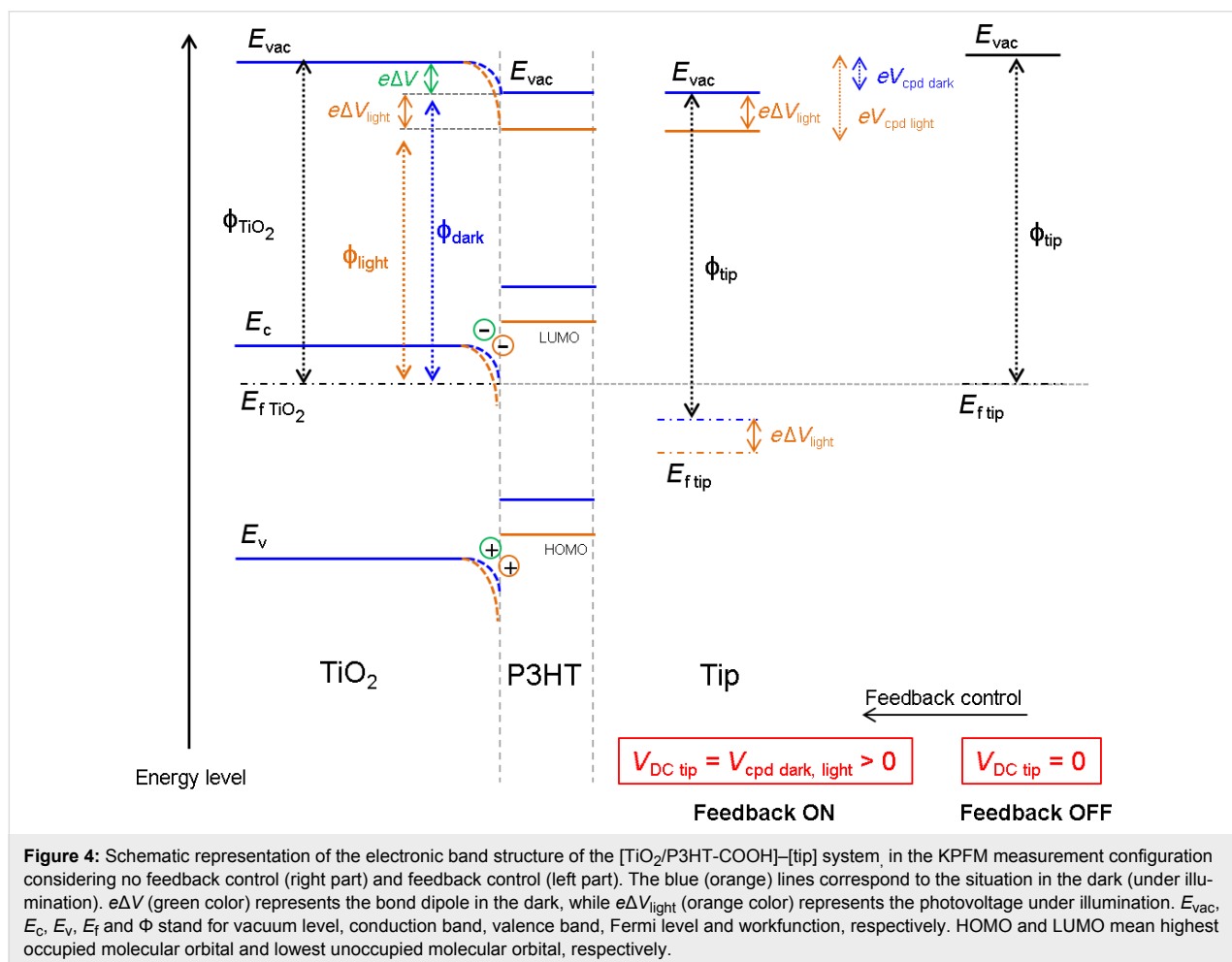
[26–28] confirmed the presence of a dipole directed away from TiO_2 or ITO substrates upon grafting of COOH-containing organic materials. A dipole directed away from the TiO_2 surface (i.e., a negative dipolar moment) means a downshift of the vacuum level upon grafting [29].

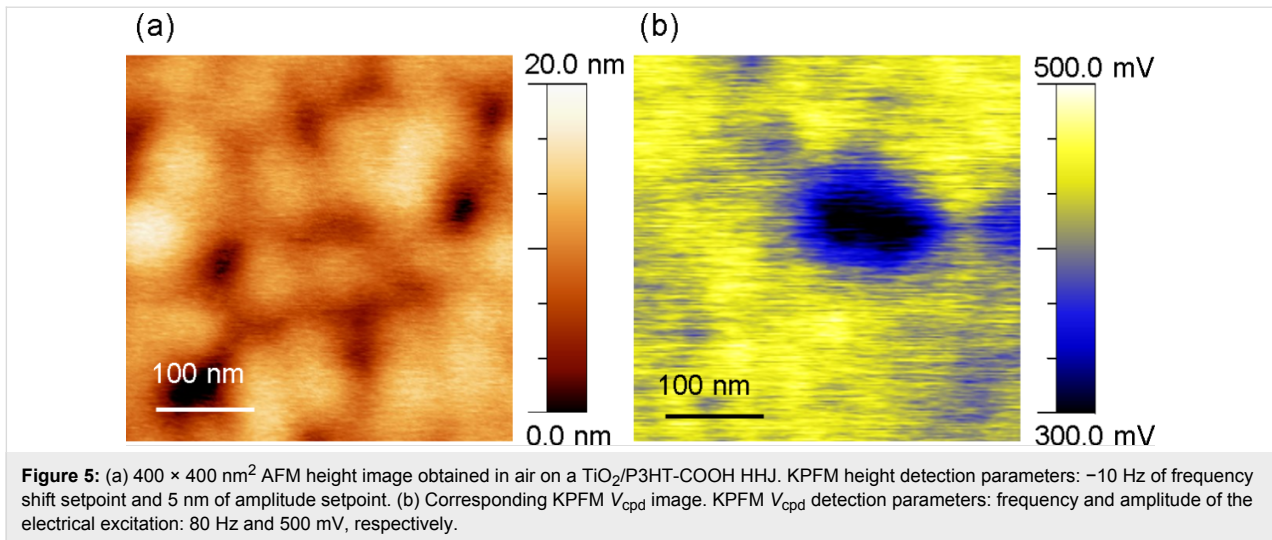
The local variations in the V_{cpd} values (the FWHM of the distribution is about 30 mV) are probably due to slightly different densities of grafted P3HT-COOH chains. Indeed, a homogeneous P3HT covering would induce a homogeneous up-shift of the V_{cpd} across the surface, leading to a variation range of V_{cpd} for the $\text{TiO}_2/\text{P3HT-COOH}$ HHJ having the same origin as that of bare nanocolumnar TiO_2 . Yet, unlike what was observed for bare nanocolumnar TiO_2 , no correlation between the height and V_{cpd} images can be seen between Figure 3e and Figure 3f. The origin of the contrast is therefore not to be linked to the V_{cpd} variations in the TiO_2 surface, but rather to an inhomogeneous contribution of the grafted P3HT-COOH.

Figure 4 shows a schematic representation of the band diagram of the ITO/ TiO_2 /P3HT-COOH/tip electronic system (blue lines)

in a KPFM measurement configuration, i.e., a grounded ITO electrode and the DC and AC bias applied to the tip. Considering no floating potential at the [ITO/ TiO_2 /P3HT-COOH] surface (see Supporting Information File 1, Figure S1), a Fermi level alignment can be assumed across the entire ITO/HHJ structure. The dipole pointing away from the TiO_2 at the $\text{TiO}_2/\text{P3HT-COOH}$ interface, leading to a partial accumulation of e^- (h^+) in the TiO_2 (P3HT), will bend the vacuum level downwards, hence lowering the surface workfunction of the TiO_2 once grafted with P3HT-COOH. The more positive V_{cpd} of $\text{TiO}_2/\text{P3HT-COOH}$ compared to bare TiO_2 confirms this mechanism.

Figure 5a shows a KPFM height image obtained on a $\text{TiO}_2/\text{P3HT-COOH}$ deposit in ambient conditions in the dark. While the top of the TiO_2 columns is visible, topographical features cannot be assigned to the presence of P3HT, probably because the nominal thickness of the P3HT-COOH deposit (13 nm) is similar to the roughness of the columnar assembly. No correlation is observed between the columnar topography and the corresponding surface potential image (Figure 5b),





which shows variations within $[260; 500] \text{ mV}$. By comparison with the data of Figure 2, this confirms that the V_{cpd} contrast is ruled by the presence of P3HT-COOH at the surface of TiO_2 . The V_{cpd} contrast in Figure 5b can be explained on the basis of the bond dipole at the $\text{TiO}_2/\text{P3HT-COOH}$ interface discussed above. V_{cpd} can then be expressed as $V_{\text{cpd}} = V_{\text{cpd TiO}_2} + e\Delta V$, $V_{\text{cpd TiO}_2}$ and $e\Delta V$ being the V_{cpd} of bare TiO_2 and the local bond dipole amplitude, respectively. The lower (higher) V_{cpd} observed in the darker (brighter) zones in Figure 5 (b) corresponds therefore to a lower (higher) $e\Delta V$, which could be related to a lower (higher) P3HT-COOH grafting density.

Variations of V_{cpd} upon illumination

As a preliminary study, KPFM measurements on bare TiO_2 were carried out in the dark and upon illumination (white light). The results are presented in Supporting Information File 1, Figure S2. As expected, no photovoltage is observed, TiO_2 being transparent in the visible spectrum.

Figure 3c shows the KPFM positive photovoltage across the entire $\text{TiO}_2/\text{P3HT-COOH}$ surface (right-part of the image) upon illumination. This up-shift of V_{cpd} upon illumination is better visualized in the corresponding profiles in Figure 3d. This photovoltage confirms locally a complete P3HT covering over the TiO_2 surface. The positive photovoltage means an increase of the V_{cpd} value, i.e., a decrease of the surface workfunction. This effect can be understood on the basis of Figure 4. Upon grafting, it was previously discussed that a dipole is created at the $\text{TiO}_2/\text{P3HT-COOH}$ interface, with positive (negative) charges in the P3HT (TiO_2) layer. This leads to a V_{cpd} value denoted $V_{\text{cpd dark}}$ in Figure 4 and expressed as:

$$eV_{\text{cpd dark}} = \Phi_{\text{tip}} - \Phi_{\text{dark}} = \Phi_{\text{tip}} - \Phi_{\text{TiO}_2} + e\Delta V \quad (2)$$

where Φ_{tip} , Φ_{TiO_2} and Φ_{dark} are the workfunctions of the tip, the TiO_2 layer and the sample surface, respectively. ΔV represents the further voltage compensation needed to cancel the electrostatic forces between the tip and the sample, due to the excess positive charges present in the P3HT layer, i.e., the bond dipole. Upon illumination, it is expected that P3HT-COOH absorbs the incident photons, thus creating excitons. The length of the P3HT-COOH chains being sufficiently small, irrespective of the location where the excitons are generated, they will be able to reach the $\text{TiO}_2/\text{P3HT-COOH}$ interface, and dissociate by transferring an electron from P3HT into the conduction band of TiO_2 . An accumulation of holes in the highest occupied molecular orbital (HOMO) of P3HT and electrons in the conduction band of TiO_2 follows, with the charges remaining close to the interface due to electrostatic attraction. A steady state is then reached between the generation and recombination of charges. The photogeneration of positive charges in the P3HT layer induces an additional V_{DC} that has to be compensated in the KPFM measurement to nullify the tip-sample electrostatic forces. This compensation is denoted ΔV_{light} in Figure 4, and the V_{cpd} value upon illumination, $V_{\text{cpd light}}$, is now expressed as:

$$\begin{aligned} eV_{\text{cpd light}} &= \Phi_{\text{tip}} - \Phi_{\text{light}} \\ &= \Phi_{\text{tip}} - \Phi_{\text{TiO}_2} + e\Delta V + e\Delta V_{\text{light}} \quad (3) \\ &= eV_{\text{cpd dark}} + e\Delta V_{\text{light}} \end{aligned}$$

This provides the following expression for the photovoltage:

$$V_{\text{ph}} = V_{\text{cpd light}} - V_{\text{cpd dark}} = \Delta V_{\text{light}} \quad (4)$$

ΔV_{light} is a positive quantity because the DC bias applied to the tip ($V_{\text{DC tip}}$) (to compensate for positive charges in P3HT) is

necessarily positive. The relation between the surface potential and $V_{\text{DC tip}}$ is given by $V_{\text{cpd}} = V_{\text{DC tip}}$. This leads to a positive value of the photovoltage, as observed experimentally in Figure 3.

Photoconductive-AFM measurements on the $\text{TiO}_2/\text{P3HT-COOH}$ hybrid heterojunctions

A $5 \times 5 \mu\text{m}^2$ height image of a $\text{TiO}_2/\text{P3HT-COOH}$ HHJ is shown in Figure 6a. The corresponding current image in Figure 6b, obtained in short-circuit configuration upon illumination, shows values of photocurrent up to 25 pA. This confirms

light absorption by the P3HT-COOH, followed by the generation of charges at the $\text{TiO}_2/\text{P3HT-COOH}$ interface. The positive sign of the photocurrent means that the charges collected at the tip are holes. The generation and collection of charges upon illumination can be explained on the basis of Figure 6e, which displays the electronic band structure of the ITO/ $\text{TiO}_2/\text{P3HT-COOH}/\text{tip}$ system in short-circuit configuration.

Upon illumination, the photon absorption by P3HT-COOH leads to the creation of excitons in the polymer. The electrons are transferred in the conduction band of TiO_2 at the

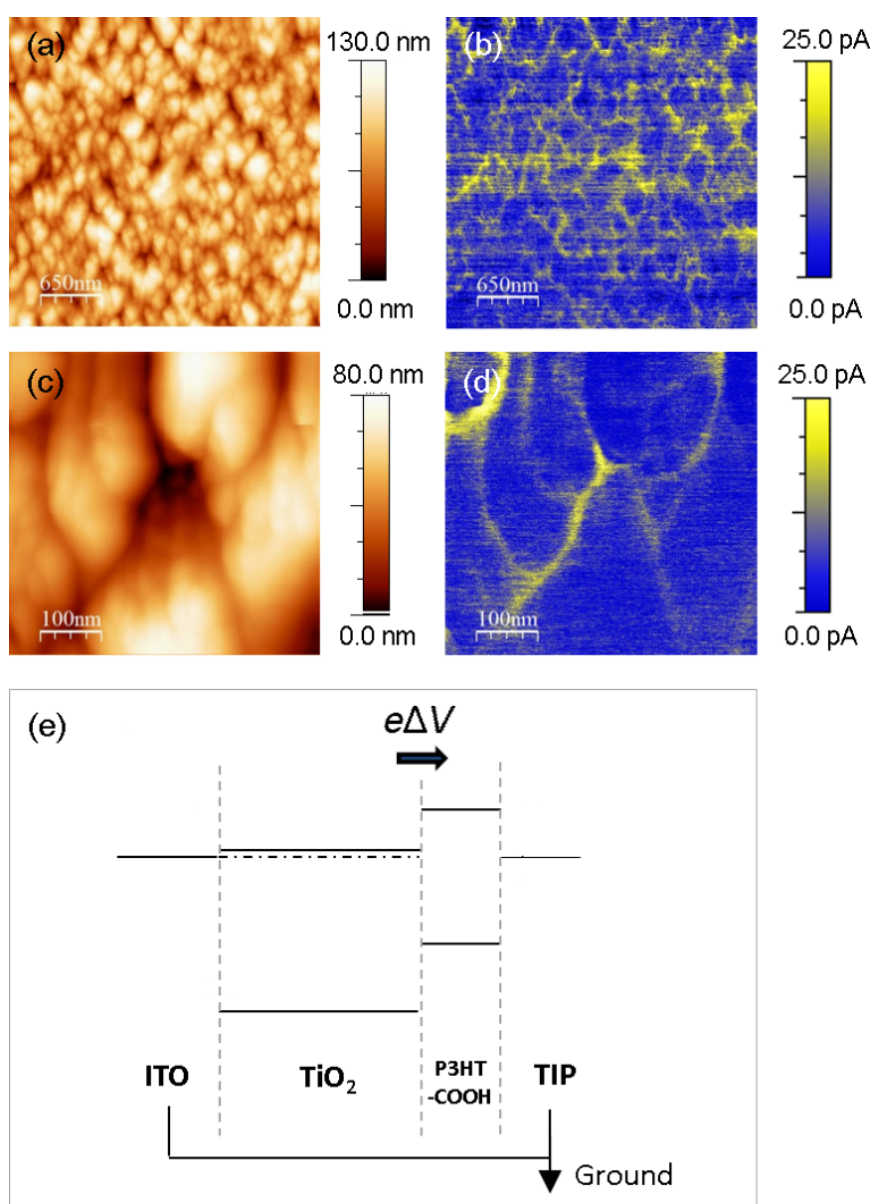


Figure 6: $5 \times 5 \mu\text{m}^2$ (a,b) and $500 \times 500 \text{nm}^2$ (c,d) PC-AFM height and photocurrent images of a $\text{TiO}_2/\text{P3HT-COOH}$ HHJ. The images were recorded upon calibrated illumination (AM 1.5, 100 suns), in short-circuit configuration. (e) Schematic representation of the electronic band structure of the ITO/ $\text{TiO}_2/\text{P3HT-COOH}/\text{tip}$ system in short-circuit configuration. $e\Delta V$ corresponds to the bond dipole.

TiO₂/P3HT-COOH interface. As the COOH group contributes to the LUMO of P3HT-COOH, the transfer of the electron to the conduction band of TiO₂ is favored compared to unsubstituted P3HT [25]. The photogenerated holes (electrons) are collected at the tip (ITO), and a positive photocurrent is measured when probing the P3HT-COOH layer.

However, the photocurrent map of Figure 6b is far from uniform, with a positive photocurrent reaching 25 pA on the regions corresponding to the inter-columnar spaces, while it is 14 pA over the top of the columns. These local variations are highlighted in Figures 6c and 6d.

The origin of those local variations could be due to the difference in tip-sample contact area between the top of the columns and the intercolumnar zones. However, it is observed in Figure 6a and 6b that, while the topographic variations are of similar amplitude across the entire surface, the intensity of photocurrent in the areas between columns varies, and is therefore not impacted solely by the topographic variations.

We note that the I_{ph} contrast is qualitatively similar to that of the V_{cpd} observed in Figure 2, in which the top of the bare TiO₂ nanocolumns displays more negative V_{cpd} values. I_{ph} and V_{cpd} quantify two different physical mechanisms, being the amount of photogenerated charges flowing in the system for the former, and the sample surface workfunction relatively to that of the tip for the latter. However, both quantities are influenced by the electron density in the conduction band of the TiO₂ and the grafting density of P3HT-COOH. These two properties impact the local conductive properties at the tip-sample contact, thus the resulting photocurrent. Φ_{TiO_2} and the P3HT-COOH grafting density are also expected to impact the resulting V_{cpd} since we previously expressed the latter as:

$$V_{cpd} = V_{cpd TiO_2} + e\Delta V,$$

where the first and second terms are directly related to Φ_{TiO_2} and the P3HT-COOH density, respectively.

Due to the small thickness of the P3HT-COOH layer on top of the TiO₂ columns, the photocurrent contrast recorded with the tip in direct contact with the surface is most probably ruled by the TiO₂ electrical properties. This explains why the I_{ph} contrast shows similarities with the V_{cpd} contrast of bare TiO₂, rather than with that of the TiO₂/P3HT-COOH HHJ. In such a configuration, the similarity of contrast between the I_{ph} (Figure 6b,d) and V_{cpd} (Figure 2b,c) images suggests that the lower photocurrent measured on top of the columns might originate from a locally lower initial (i.e., prior to illumination) electron density at the TiO₂ surface. Among various possible factors, this varia-

tion of electron density might be due to the presence of different TiO₂ crystal facets, as the latter are shown to influence the electronic properties of the TiO₂ surface [30,31].

Conclusion

Nanocolumnar TiO₂ layers were sensitized with a layer of P3HT-COOH. KPFM surface potential measurements indicate complete covering of the TiO₂ surface by the polymer. A downshift of the vacuum level of the sample upon grafting, i.e., an increase of the surface potential, was measured, due to the formation of a bond dipole at the TiO₂/P3HT-COOH interface. Upon in situ illumination, a positive photovoltage was observed, which is related to the accumulation of photogenerated holes in the P3HT layer. Along with the surface potential shift, a positive photocurrent was measured by PC-AFM measurements over the TiO₂/P3HT-COOH heterojunction upon illumination, corresponding to a hole collection at the tip. Lower photocurrent values measured on top of the TiO₂ columns can be related to the corresponding more negative V_{cpd} , indicating a locally lower electron density pre-existing the illumination.

Supporting Information

Supporting Information File 1

Supporting Information.

Figure S1 shows a FM-AFM height image obtained in UHV astride the step from a nanostructured TiO₂/P3HT-COOH HHJ to the ITO electrode lying below. The applied DC sample bias was varied during the measurement, without illumination. This result aims at demonstrating the absence of floating potential across the layer composing the sample. Figure S2 shows the superimposition of FM-KPFM height and V_{cpd} profiles over a nanostructured TiO₂ film obtained in UHV and recorded with and without illumination. The result aimed at demonstrating the absence of light-induced artefact during the recording of topography, as well as the negligibility of the photovoltaic effect at the TiO₂/ITO interface. [<https://www.beilstein-journals.org/bjnano/content/supplementary/2190-4286-9-197-S1.pdf>]

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References

- Parida, B.; Iniyani, S.; Goic, R. *Renewable Sustainable Energy Rev.* **2011**, *15*, 1625–1636. doi:10.1016/j.rser.2010.11.032
- Polman, A.; Knight, M.; Garnett, E. C.; Ehrler, B.; Sinke, W. C. *Science* **2016**, *352*, aad4424. doi:10.1126/science.aad4424
- Branker, K.; Pathak, M.; Pearce, J. *Renewable Sustainable Energy Rev.* **2011**, *15*, 4470–4482. doi:10.1016/j.rser.2011.07.104
- Wright, M.; Uddin, A. *Sol. Energy Mater. Sol. Cells* **2012**, *107*, 87–111. doi:10.1016/j.solmat.2012.07.006
- Mathew, S.; Yella, A.; Gao, P.; Humphry-Baker, R.; Curchod, B. F. E.; Ashari-Astani, N.; Tavernelli, I.; Rothlisberger, U.; Nazeeruddin, M. K.; Grätzel, M. *Nat. Chem.* **2014**, *6*, 242–247. doi:10.1038/nchem.1861
- Tengstedt, C.; Osikowicz, W.; Salaneck, W. R.; Parker, I. D.; Hsu, C.-H.; Fahlman, M. *Appl. Phys. Lett.* **2006**, *88*, 053502. doi:10.1063/1.2168515
- Chung, I.; Lee, B.; He, J.; Chang, R. P. H.; Kanatzidis, M. G. *Nature* **2012**, *485*, 486–489. doi:10.1038/nature11067
- Hagfeldt, A.; Boschloo, G.; Sun, L.; Kloo, L.; Pettersson, H. *Chem. Rev.* **2010**, *110*, 6595–6663. doi:10.1021/cr900356p
- Othman, M. A.; Amat, N. F.; Ahmad, B. H.; Rajan, J. *J. Phys.: Conf. Ser.* **2014**, *495*, 012027. doi:10.1088/1742-6596/495/1/012027
- Pomoni, K.; Sofianou, M.; Georgakopoulos, T.; Boukos, N.; Trapalis, C. *J. Alloys Compd.* **2013**, *548*, 194–200. doi:10.1016/j.jallcom.2012.08.136
- Tian, J.; Zhao, Z.; Kumar, A.; Boughton, R. I.; Liu, H. *Chem. Soc. Rev.* **2014**, *43*, 6920–6937. doi:10.1039/C4CS00180J
- Weng, Z.; Guo, H.; Liu, X.; Wu, S.; Yeung, K. W. K.; Chu, P. K. *RSC Adv.* **2013**, *3*, 24758–24775. doi:10.1039/c3ra44031a
- Detle, C.; Pérez-Osorio, M. A.; Kley, C. S.; Punke, P.; Patrick, C. E.; Jacobson, P.; Giustino, F.; Jung, S. J.; Kern, K. *Nano Lett.* **2014**, *14*, 6533–6538. doi:10.1021/nl503131s
- Reckers, P.; Dimamay, M.; Klett, J.; Trost, S.; Zilberberg, K.; Riedl, T.; Parkinson, B. A.; Brötz, J.; Jaegermann, W.; Mayer, T. *J. Phys. Chem. C* **2015**, *119*, 9890–9898. doi:10.1021/acs.jpcc.5b01264
- Moerman, D.; Lazzaroni, R.; Douhéret, O. *Appl. Phys. Lett.* **2011**, *99*, 093303. doi:10.1063/1.3628658
- Weickert, J.; Auras, F.; Bein, T.; Schmidt-Mende, L. *J. Phys. Chem. C* **2011**, *115*, 15081–15088. doi:10.1021/jp203600z
- Huang, Y.-C.; Yen, W.-C.; Liao, Y.-C.; Yu, Y.-C.; Hsu, C.-C.; Ho, M.-L.; Chou, P.-T.; Su, W.-F. *Appl. Phys. Lett.* **2010**, *96*, 123501. doi:10.1063/1.3357425
- Wu, M.-C.; Wu, Y.-J.; Yen, W.-C.; Lo, H.-H.; Lin, C.-F.; Su, W.-F. *Nanoscale* **2010**, *2*, 1448–1454. doi:10.1039/b9nr00385a
- Zeng, T.-W.; Ho, C.-C.; Tu, Y.-C.; Tu, G.-Y.; Wang, L.-Y.; Su, W.-F. *Langmuir* **2011**, *27*, 15255–15260. doi:10.1021/la203533u
- Phan, H.; Jahnke, J. P.; Chmelka, B. F.; Nguyen, T.-Q. *Appl. Phys. Lett.* **2014**, *104*, 233305. doi:10.1063/1.4883001
- Batra, Y.; Rai, D.; Mehta, B. R. *Appl. Phys. Express* **2013**, *6*, 041602. doi:10.7567/APEX.6.041602
- Zeng, T.-W.; Hsu, F.-C.; Tu, Y.-C.; Lin, T.-H.; Su, W.-F. *Chem. Phys. Lett.* **2009**, *479*, 105–108. doi:10.1016/j.cplett.2009.07.104
- Lin, Y.-Y.; Chu, T.-H.; Li, S.-S.; Chuang, C.-H.; Chang, C.-H.; Su, W.-F.; Chang, C.-P.; Chu, M.-W.; Chen, C.-W. *J. Am. Chem. Soc.* **2009**, *131*, 3644–3649. doi:10.1021/ja8079143
- Dervaux, J.; Cormier, P.-A.; Konstantinidis, S.; Di Ciuccio, R.; Coulembier, O.; Dubois, P.; Snyders, R. *Vacuum* **2015**, *114*, 213–220. doi:10.1016/j.vacuum.2014.10.016
- Lasser, L.; Ronca, E.; Pastore, M.; De Angelis, F.; Cornil, J.; Lazzaroni, R.; Beljonne, D. *J. Phys. Chem. C* **2015**, *119*, 9899–9909. doi:10.1021/acs.jpcc.5b01267
- Nüesch, F.; Carrara, M.; Zuppiroli, L. *Langmuir* **2003**, *19*, 4871–4875. doi:10.1021/la026962w
- Krüger, J.; Bach, U.; Grätzel, M. *Adv. Mater.* **2000**, *12*, 447–451. doi:10.1002/(SICI)1521-4095(200003)12:6<447::AID-ADMA447>3.0.CO;2-8
- Henning, A.; Günzburger, G.; Jöhr, R.; Rosenwaks, Y.; Bozic-Weber, B.; Housecroft, C. E.; Constable, E. C.; Meyer, E.; Glatzel, T. *Beilstein J. Nanotechnol.* **2013**, *4*, 418–428. doi:10.3762/bjnano.4.49
- Goh, C.; Scully, S. R.; McGehee, M. D. *J. Appl. Phys.* **2007**, *101*, 114503. doi:10.1063/1.2737977
- Wang, X.; Li, T.; Yu, R.; Yu, H.; Yu, J. *J. Mater. Chem. A* **2016**, *4*, 8682–8689. doi:10.1039/C6TA02039A
- Selçuk, S.; Selloni, A. *J. Phys. D: Appl. Phys.* **2017**, *50*, 273002. doi:10.1088/1361-6463/aa7540

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