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W Oukil, Ph. Thieullen, A. Kessi. Reduced dimension and Rotation vector formula of ordinary differential equation. 2019. hal-01969255v1

# HAL Id: hal-01969255 https://hal.science/hal-01969255v1

Preprint submitted on 3 Jan 2019 (v1), last revised 3 Apr 2022 (v4)

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# Reduced dimension and Rotation vector formula of ordinary differential equation

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January 3, 2019

#### Abstract

The *leader trajectory* function defined in this article is an approximated solution of a differential equation. It is defined by some independent one-dimensional differential equations. The generalized main result of this article asserts that if the *leader trajectory* exists then it is at finite distance from the solution of the system. However, the differential equation can be controlled by a reduced one-dimensional differential equations. The application of the generalized main result is to control the trajectory of the periodic systems. We prove that for any *periodic* system and any initial condition there exists a *leader* trajectory which is a linear function of the time variable. Otherwise, we find an exact *Rotation vector formula* which is the relation between the rotation vector and the initial condition. In addition, we find easily a necessary and sufficient condition for the existence of a locally constant rotation vector under perturbation of the system, known by the the Arnold Tongue. The mean result answer to many problems in natural sciences.

*Keywords:* Differential equations, bounded solution, periodic system, rotation vector.

### 1 Introduction and main results

One important application of the result is to control the trajectory of the periodic systems as cited in the abstract; The applications of the periodic differential equations are various. The issue is to estimate the value of the *rotation vector*. The list of the related works is very large; for some biological works see [8, 44]. The mean result answer to many problems because we find an explicit relation between the rotation vector and the initial condition of the orbit. We study in this article the following system

$$\dot{x} = f(x), \quad t \in \mathbb{R}, \quad x = (x_i)_{i=1}^n, \quad x(0) = x_0, \quad x_0 \in \mathbb{R}^n.$$
 (1)

where  $f : \mathbb{R}^n \to \mathbb{R}^n$  is a  $C^2$  function. The function  $x(t) = (x_i(t))_{i=1}^n$  is the state of the system. Notice that any non-autonomous system can be written as an autonomous system. We introduce in the present Section some notations, the definition of the leader trajectory and the main results. The Section 2 is the ingredients to study the linearized of system (1) which has the form of the Riccati differential equation. In the rest of Sections we prove the main results

#### **1.1** Periodic differential equations

Many mathematical models are defined by the periodic systems [3, 5, 1]. For example, in biological sciences the components of the rotation vector are called the *frequency* [38, 13, 8]. When the periodic system is stable, in the sens of the *Lyapunov stability*, the rotation vector exists [3, 5, 1, 4].

**Definition 1** (Periodic function). Let be  $q, p \in \mathbb{N}^*$  and let  $g = (g_i)_1^p : \mathbb{R}^q \to \mathbb{R}^p$  be a function. We say that g is *periodic* if it is periodic relative to each variable. Otherwise,

$$g_i(y_1,\ldots,y_j+1,\ldots,y_q) = g_i(y), \ \forall i = \overline{1,q}, \forall j = \overline{1,q}, \ \forall y = (y_i)_{i=1}^q \in \mathbb{R}^q.$$

In this article, we use the usual norm on  $\mathbb{R}^q$ ,  $q \in \mathbb{N}^*$ , defined by  $||z|| = \max_{1 \le i \le q} |z_i|$  for all  $z := (z_i)_{i=1}^q \in \mathbb{R}^q$ . We also denote ||.|| the associated matrix norm.

**Definition 2** (The rotation vector). Let be  $q \in \mathbb{N}^*$  and  $\lambda \in \mathbb{R}^q$ . Let  $g : \mathbb{R} \to \mathbb{R}^q$  be a continuous function. We say that  $\lambda$  is the *rotation vector* of the function g if

$$\sup_{t\in\mathbb{R}}||g(t)-\lambda t||<+\infty.$$

Notice that if the rotation vector exists then it is unique. To formulate the main results let us introduce the following notations; Let  $p, q \in \mathbb{N}^*$ , for every  $q \times p$  matrix  $M = [m_{i,j}]_{i,j}$  we denote

$$\sigma(M) := \sum_{1 \le i \le q} \sum_{1 \le j \le p} m_{i,j}.$$
(2)

Let be  $u \in \mathbb{R}^q$ , we denote Diag(u) the diagonal matrix of diagonal u. Let be  $1 \leq i \leq q$ , define the diagonal matrix  $I_{i,q}$  as

$$I_{i,q} = \text{Diag}(v), \ v = (v_j)_{j=1}^q : \ v_i = -1 \text{ and } v_j = 1, \ \forall 1 \le j \ne i \le q.$$
 (3)

Denote by  $I_q$  the  $q \times q$  identity matrix and put  $I_{0,q} = I_q$ . The particularity of the matrix  $I_{i,q}$  is the relation  $2y_i = \sigma(I_{0,q}y) - \sigma(I_{i,q}y)$  for any vector  $(y_i)_{i=1}^q$ . Let  $g : \mathbb{R}^q \to \mathbb{R}^q$  be a  $C^2$  function we denote dg and  $d^2g$  the first and the second differential of g respectively.

**Main Result.** Let be  $x_0 \in \mathbb{R}^n$  and let x be the solution of the Equation (1). Suppose that f is periodic, then x admits a rotation vector  $\rho := (\rho_j)_{j=1}^n \in \mathbb{R}^n$  which is the unique solution of the following Rotation vector formula

$$\lim_{t \to +\infty} \frac{1}{t} \Big[ \Psi_0[f_{x_0}](\rho, t) - \Psi_i[f_{x_0}](\rho, t) \Big] = 0, \quad \forall i = \overline{1, n}, \tag{4}$$

where the functions  $\Psi_i[f_{x_0}]$  are defined by

$$\Psi_i[f_{x_0}](\rho,t) := \int_0^t A_i(\rho,s) \exp\Big(\int_s^t \sigma\Big(I_{i,q}df(\nu\rho + x_0)I_{i,q}\Big)d\nu\Big)ds,$$
$$A_i(\rho,s) := \sigma\Big(I_{i,q}(\rho - f(s\rho + x_0))\Big).$$

#### 1.2 Generalized Main result

Define the leader trajectory function.

**Definition 3.** [Leader trajectory] Let be  $q \in \mathbb{N}^*$ , let  $g : \mathbb{R}^q \to \mathbb{R}^q$  and  $\mu : \mathbb{R} \to \mathbb{R}^q$  be a  $C^1$  functions. We say that  $\mu$  is a *leader trajectory of the function g* if the following hypotheses are satisfied

• 
$$\sup\left\{\left|\int_{0}^{t}\sigma\left(dg(\mu(\nu))\right)d\nu\right|, \quad t \in \mathbb{R}\right\} < +\infty.$$
  
•  $\forall i = \overline{1,q}, \ \exists \tau_i \in \{-1,1\}: \ \sup_{t \ge 0} \int_{0}^{t}\tau_i\sigma\left(I_{i,q}dg(\mu(\tau_i\nu))I_{i,q}\right)d\nu < +\infty.$ 

•  $\forall i = \overline{1, q}$ :  $\sup_{t \in \mathbb{R}} |\psi_0(\tau_i t) - \psi_i(\tau_i t)| < +\infty.$ 

where the function  $\psi_i : \mathbb{R} \to \mathbb{R}$  satisfies the following one-dimensional differential equations

$$\dot{\psi}_i = \sigma \Big( I_{i,q} \dot{\mu}(t) - I_{i,q} g(\mu(t)) \Big) + \sigma \Big( I_{i,q} dg(\mu(t)) I_{i,q} \Big) \psi_i, \ i = \overline{0, q},$$
$$\psi_i(0) = 0, \quad t \in \mathbb{R}.$$

**Generalized Main Result.** Let be  $x_0 \in \mathbb{R}^n$  and let x be the solution of the Equation (1). Suppose that

$$\|df\|_{\infty} + \|d^2f\|_{\infty} < +\infty.$$

Let  $\mu = (\mu_i)_{i=1}^n : \mathbb{R} \to \mathbb{R}^n$  be a leader trajectory of f such that  $\mu(0) = x_0$ . Then there exists D > 0 such that

$$\|\mu(t) - x(t)\| \le D, \quad \forall t \in \mathbb{R}.$$
(5)

In order to prove the Generalized mains result, we use the following linearized system of (1): Let  $\mu = (\mu_i)_{i=1}^n : \mathbb{R} \to \mathbb{R}^n$  be a  $C^1$  function such that  $\mu(0) = x_0$ . Consider the change of variables

$$\eta(t) = \mu(t) - \Phi^t(x_0), \ t \in \mathbb{R}, \ \eta := (\eta_i)_{i=1}^n,$$

where  $\Phi^t$  is the flow of the system (1). Use Taylor formula to obtain the following linearized system

$$\dot{\eta} = \dot{\mu}(t) - f(\mu(t)) + df(\mu(t))\eta + \eta^T \zeta_{\mu}(t)\eta, \quad \|\eta(0)\| = 0, \quad t \in \mathbb{R},$$
(6)

where  $\zeta_{\mu} : \mathbb{R} \times \mathbb{R}^{n \times n \times n}$  is a continuous function such that  $\|\zeta_{\mu}\|_{\infty} < \|d^2 f\|_{\infty}$ and where  $\eta^T$  is the transpose of the function  $\eta$ . Equation (6) take the form of Riccati differential equation. The aim is to prove that the solution  $\eta : \mathbb{R} \to \mathbb{R}^q$  is uniformly bounded on  $\mathbb{R}$ .

### 2 Riccati differential equation

We denote  $v^T$  the transpose of a vector v and for any  $q \in \mathbb{N}^*$  we denote

$$\mathbb{1}_q := (1, \dots, 1)^T \in \mathbb{R}^q \text{ and } \mathbb{0}_q := (0, \dots, 0)^T \in \mathbb{R}^q.$$

In this Section, we study a generalized Riccati differential equation. Let  $\mathcal{A}: \mathbb{R} \to \mathbb{R}^q$ ,  $\mathcal{B}: \mathbb{R} \to \mathbb{R}^{q \times q}$  and  $\mathcal{H}: \mathbb{R} \to \mathbb{R}^{q \times q \times q}$  be a continuous functions such that

$$\exists \gamma < 1: \max\left\{ \left\| \mathcal{H} \right\|_{\infty}, \left\| \mathcal{B} \right\|_{\infty} \right\} \le \gamma.$$
(7)

4

Consider the following Riccati differential equation

$$\dot{y} = \mathcal{A}(t) + \mathcal{B}(t)y + y^T \mathcal{H}(t)y, \ y(0) = \mathbb{O}_q, \ y = (y_i)_{i=1}^q, \ t \in \mathbb{R}.$$
 (8)

Recall that the notation  $\sigma(.)$  and the diagonal matrix  $I_{i,q}$  are defined by Equations (2) and (3) respectively. Consider the following hypotheses [H1]

$$\sup\left\{\left|\int_0^t \sigma(\mathcal{B}(\nu))d\nu\right|, \quad t\in\mathbb{R}\right\}<+\infty,$$

and for every  $i = \overline{1, q}$  there exists  $\tau_i \in \{-1, 1\}$  such that

$$\sup\left\{\tau_i \int_0^t \sigma(I_{i,q} \mathcal{B}(\tau_i \nu) I_{i,q}) d\nu, \quad t \ge 0\right\} < +\infty,$$

[H2] The following one-dimensional differential equations

$$\dot{\psi}_i = \sigma(I_{i,q}\mathcal{A}(t)) + \sigma(I_{i,q}\mathcal{B}(t)I_{i,q})\psi_i, \quad t \ge 0,$$
(9)

admit a solution  $\psi_i : \mathbb{R}_+ \to \mathbb{R}$  of initial condition  $\psi_i(0) = 0$  such that

$$\sup_{t\in\mathbb{R}}|\psi_0(\tau_i t) - \psi_i(\tau_i t)| < +\infty$$

The following theorem is the main ingredient used to prove the main results.

**Theorem 4.** Let y be the solution of the differential Equation (8) and let be  $\gamma \in (0, 1)$  as defined by Equation (7). Suppose that Hypotheses [H1] and [H2] are satisfied. Then there exists  $\gamma_* > 0$  such that for all  $\gamma \in (0, \gamma_*)$  there exists  $D_{\gamma} > 0$  such that the solution y of Equation (8) satisfies  $\sup_{t \in \mathbb{R}} ||y(t)|| < D_{\gamma}$ .

The strategy it to use an appropriate change of variable in (8) and use the differential equation comparison. In order to simplify, we introduce the following notation. Define the cone  $V^q_+$  of  $\mathbb{R}^q$  by

$$V_{+}^{q} := \{ z = (z_{i})_{i=1}^{q} \in \mathbb{R}^{q} : z_{i} > 0, \quad \forall i = \overline{1, q} \}.$$
(10)

We consider the partial order  $\succeq$  on  $\mathbb{R}^q$  defined by

$$\forall z \in \mathbb{R}^q: \quad z \succ \mathbb{O}_q \iff z \in V^q_+ \quad \text{and} \quad z \succeq \mathbb{O}_q \iff z \in V^q_+.$$

Define the sets  $\mathcal{K}_q^+$  and  $\mathcal{K}_q$  of the  $q \times q$  diagonal matrix by

$$\mathcal{K}_q^+ = \{ I_{i,q}, \ i = \overline{0,q} \} \quad \text{and} \quad \mathcal{K}_q = \{ -I : \ I \in \mathcal{K}_q^+ \} \cup \mathcal{K}_q^+.$$
(11)

We have the following lemma

**Lemma 5.** Let y be the solution of the differential Equation (8). Let be  $\gamma > 0$  given by Equation (7). Let be T > 0 and put U = [-T, T]. For every  $I \in \mathcal{K}_q$  and  $\tau \in \{-1, 1\}$  denote

$$\delta_I(t) := \exp(qt) \Big[ \psi_I(\tau t) - \sigma(Iy(\tau t)) \Big]$$
(12)

where  $\psi_I : \mathbb{R} \to \mathbb{R}$  satisfies

$$\dot{\psi}_I = \sigma(I\mathcal{A}(t)) + \tau\sigma(I\mathcal{B}(t)I)\psi_I, \ t \in \mathbb{R}.$$
(13)

Let be D > 0 and suppose that  $\sup_{t \in U} ||y(t)|| < D/\gamma$ , then

$$\forall t \in \mathbb{R}, \ \forall I \in \mathcal{K}_q: \ \psi_I(\tau t) - \sigma(Iy(\tau t)) = \frac{D^2}{\gamma(1-\gamma)} \implies \dot{\delta}_I(t) > 0.$$

*Proof.* Denote  $\tilde{I}_q := [c_{i,j}]_{1 \le i,j \le q}$  be the  $q \times q$  matrix such that  $c_{i,j} = 1$  for all  $i, j = \overline{1, q}$ . Let be  $I \in \mathcal{K}_q$  and Consider the change of variable

$$z_I(\tau, t) := \exp(t\tilde{I})Iy(\tau t), \quad t \in U.$$
(14)

and put  $\tilde{y}_I(t) := Iy(\tau t)$ . Since the inverse matrix of I is itself the system (8) becomes

$$\dot{\tilde{y}}_I = \tau I \mathcal{A}(\tau t) + \tau I \mathcal{B}(\tau t) I \tilde{y}_I(t) + \tau I y^T(\tau t) \mathcal{H}(\tau t) y(\tau t), \quad \forall t \in U.$$

That follows

$$\dot{z}_{I}(\tau,t) = \exp(t\tilde{I}) \Big[ \tau I \mathcal{A}(\tau t) + [\tilde{I} + \tau I \mathcal{B}(\tau t) I] \tilde{y}_{I} + \tau y^{T}(\tau t) \mathcal{H}(\tau t) y(\tau t) \Big].$$
(15)

Remark that  $\tilde{I}_q^2 = q\tilde{I}_q$ . We have the relation

$$\exp(t\tilde{I}_q) = I_q + \frac{1}{q}(\exp(qt) - 1)\tilde{I}_q \text{ and } \|\exp(t\tilde{I}_q)\| \le \exp(qt), \ \forall t \ge 0, \quad (16)$$

By hypothesis  $\sup_{t\in U} \|\gamma y(t)\| < D$  and  $\|\mathcal{H}\|_{\infty} \leq \gamma$  then

$$\|\gamma \exp(t\tilde{I})Iy^T(\tau t)\mathcal{H}(\tau t)y(\tau)\| < \exp(qt)D^2, \quad \forall t \in U.$$

From (17), we have

$$\dot{z}_{I}(\tau,t) \prec \gamma \exp(t\tilde{I}) \left[ \tau I \mathcal{A}(\tau t) + [\tilde{I} + \tau I \mathcal{B}(t)I] \tilde{y}_{I} \right] + \frac{1}{\gamma} \exp(qt) D^{2} \mathbb{1}_{q}.$$
 (17)

Suppose that  $\delta_I(t) = D^2/(\gamma(1-\gamma))$  then  $\delta_I(t) > 0$ . Use the fact that for every  $c = (c_i)_{i=1}^q \in \mathbb{R}^q$  and  $v = (v_i)_{i=1}^q \in \mathbb{R}^q$  such that  $\sigma(v) \neq 0$  there exists  $\tilde{c} \in \mathbb{R}$  that satisfies  $|\tilde{c}| \leq ||c||$  and such that

$$\sum_{j=1}^{q} c_j v_i = \tilde{c}\sigma(v).$$

The above vector v play the following role

$$v_i = \frac{1}{q}\psi_I(\tau t) - \tilde{y}_j(t).$$

Use the hypothesis hypothesis  $\|\mathcal{B}(t)\| < \gamma$ , to obtain

$$-\tau I \mathcal{B}(t) I(\frac{1}{q}\psi_I(\tau t) - \tilde{y}_j(t)) \prec \gamma \delta_I(t) \mathbb{1}_q, \quad \forall t \in U,$$
$$-\left[\tilde{I} + \tau I \mathcal{B}(t) I\right] (\frac{1}{q}\psi_I(\tau t) - \tilde{y}_j(t)) \prec -(1-\gamma)\delta_I(t) \mathbb{1}_q, \quad \forall t \in U.$$

Since  $\exp(t\tilde{I})V_q^+ \subset V_q^+$ , then Equation (17) implies that

$$\dot{z}_{I}(\tau,t) \prec \tau \exp(t\tilde{I})I\mathcal{A}(\tau t) + \frac{1}{q}\exp(t\tilde{I})[\tilde{I} + \tau I\mathcal{B}(\tau t)I]\psi_{I}(\tau t) - (1-\gamma)\exp(t\tilde{I})\delta(t)\mathbb{1}_{q} + \frac{1}{\gamma}\exp(qt)D^{2}\mathbb{1}_{q}, \ \forall t \in U.$$

By Equation (16), we have

$$\sigma(I\mathcal{B}(\tau t)I) = \exp(-qt)\sigma(\exp(t\tilde{I})I\mathcal{B}(\tau t)I) \text{ and } \exp(t\tilde{I})\tilde{I} = \exp(qt)\mathbb{1}_q.$$

By addition, we get

$$\frac{d}{dt}[\exp(qt)\psi_I(\tau t) - \sigma(z_I(\tau, t)] < -q(1-\gamma)\exp(qt)\delta(t) + q\frac{1}{\gamma}\exp(qt)D^2,$$

We have proved  $\delta_I(t) = D^2/(\gamma(1-\gamma))$  implies that  $\dot{\delta}_I(t) > 0$ .

Recall that the diagonal matrix  $I_{i,q}$  define the sets  $\mathcal{K}_q^+$  and  $\mathcal{K}_q$  in Equation (11) and recall that  $I_q$  is the  $q \times q$  identity matrix.

**Lemma 6.** Let  $v : \mathbb{R} \to \mathbb{R}^q$  be a function. Let  $\psi : \mathcal{K}_q \times \{-1, 1\} \times \mathbb{R}_+ \to \mathbb{R}$  be a functions such that

• For all 
$$(I, \tau, t) \in \mathcal{K}_q \times \{-1, 1\} \times \mathbb{R}_+$$
  
 $\psi(-I, \tau, t) = -\psi(I, \tau, t).$ 
(18)

• There exists  $\theta, \alpha, T > 0$  such that for all  $I \in \mathcal{K}_q$  there exists  $\tau_I \in \{-1, 1\}$  such that  $\tau_I = \tau_{-I}$  and such that

$$\theta + \psi(I, \tau_I, t) - \sigma(Iv(\tau_I t)) > \alpha > \psi(I, \tau_I, -t) + \sigma(Iv(-\tau_I t)), \ \forall t \in (0, T).$$

• For all  $I_0 \in \{I_q, -I_q\}$  and all  $\tau \in \{-1, 1\}$ 

$$\theta + \psi(I_0, \tau, t) - \sigma(I_0 v(\tau t)) > \alpha > \psi(I_0, \tau, -t) + \sigma(I_0 v(-\tau t)), \ \forall t \in (0, T)$$

Then,

$$\|v(t)\| \le \theta + \alpha + \frac{1}{2}\beta_T, \quad \forall t \in [-T, T],$$

where

$$\beta_T := \sup \left\{ \left| \psi(I_q, \tau_I, t) - \psi(I, \tau_I, t) \right| : \quad t \in [-T, T], \quad I \in \mathcal{K}_q^+ \right\}.$$

*Proof.* Use the fact  $\tau_I = \tau_{-I}$  and use the Equation (18) to obtain

$$\begin{aligned} \theta - \psi(I, \tau_I, t) + \sigma(Iv(\tau_I t)) &> \alpha > -\psi(I, \tau_I, -t) - \sigma(Iv(-\tau_I t)), \\ \theta - \psi(I_0, \tau, t) + \sigma(I_0v(\tau t)) &> \alpha > -\psi(I_0, \tau, -t) - \sigma(I_0v(-\tau t)). \end{aligned}$$

By addition we get:

$$2(\theta - \alpha) + \psi(I_0, \tau_I, t) + \psi(I, \tau_I, t) > [\sigma(I_0v(\tau_I t)) + \sigma(Iv(\tau_I t))], 2\alpha - \psi(I, \tau_I, -t) - \psi(I_0, \tau_I, -t) > \sigma(I_0v(-\tau_I t)) + \sigma(Iv(-\tau_I t)).$$

For  $I_0 = I_q$  and  $I = -I_{i,q}$  we have

$$2(\theta - \alpha) + \psi(I_q, \tau, t) - \psi(I_{i,q}, \tau_I, t) - \alpha > \sigma(I_q v(\tau t)) - \sigma(I_{i,q} v(\tau_I t)),$$
  
$$\alpha - \psi(I_q, \tau, -t) + \psi(I_{i,q}, \tau_I, -t) > \sigma(I_q v(-\tau t)) - \sigma(I_{i,q} v(-\tau_I t)).$$

For  $I_0 = -I_q$  and  $I = I_{i,q}$  we have

$$2(\theta - \alpha) - \psi(I_q, \tau, t) + \psi(I_{i,q}, \tau_I, t) - \alpha > -[\sigma(I_q v(\tau t)) - \sigma(I_{i,q} v(\tau_I t))],\\ \alpha + \psi(I_q, \tau, -t) - \psi(I, \tau_I, -t) > -\sigma(I_q v(-\tau t)) + \sigma(I_{i,q} v(-\tau_I t)).$$

use the fact that for every  $u = (u_i)_{i=1}^q \in \mathbb{R}^q$  we have

$$u_i = \frac{1}{2} [\sigma(I_q u) - \sigma(I_{i,q} u)], \quad \forall i = \overline{1, q}$$

Then for every  $i = \overline{1, q}$  and all  $t \in (0, T)$  we have:

$$\left| v_i(\tau_I t) + \frac{1}{2} \left[ \psi(I_q, \tau_i, t) - \psi(I_{i,q}, \tau_i, t) \right] \right| < \theta + \alpha,$$
  
$$\left| v_i(-\tau_I t) - \frac{1}{2} \left[ \psi(I_q, \tau_i, -t) - \psi(I_{i,q}, \tau_i, -t) \right] \right| < \theta + \alpha.$$

Finally, by definition of the constants  $\beta_T$  in the present Lemma, we get

$$\|v(t)\| \le \theta + \alpha + \frac{1}{2}\beta_T, \quad \forall t \in [-T, T].$$

Proof of Theorem 4.

For all  $\alpha \in \mathbb{R}$  and  $i = \overline{0, q}$  let be  $\psi_i(t, \alpha)$  be the solutions of the Equation (9) of initial condition  $\psi_i(0, \alpha) = \alpha$ . Let y be the solution of the differential Equation (8) and put

$$\overline{\delta}_i(t,\alpha) := \psi_i(\tau_i t, \alpha) + \sigma(I_{i,q} y(\tau_i t)) \text{ and } \underline{\delta}_i(\tau_i t, \alpha) := \psi_i(t,\alpha) - \sigma(I_{i,q} y(\tau_i t)).$$

Let be D > 0 and put  $\alpha_* = D^2/(\gamma(1-\gamma))$ . Since  $y(0) = \mathbb{O}_q$  then  $\overline{\delta}_i(0,\alpha_*) = \underline{\delta}_i(0,\alpha_*) = \alpha_*$ . By Lemma 5 we have

$$\overline{\delta}_i(0, \alpha_*) > 0 \text{ and } \underline{\dot{\delta}}_i(0, \alpha_*) > 0 \quad \forall i = \overline{0, q}.$$

Since  $\overline{\delta}_i(0,0) = \underline{\delta}_i(0,0) = 0$ , then there exists  $\epsilon > 0$  such that

$$\begin{split} \overline{\delta}_i(t,\alpha_*) &> \alpha_* > \overline{\delta}_i(-t,0) \quad \forall t \in (0,\epsilon), \quad \forall i = \overline{0,q}, \\ \underline{\delta}_i(t,\alpha_*) &> \alpha_* > \underline{\delta}_i(-t,0) \quad \forall t \in (0,\epsilon), \quad \forall i = \overline{0,q}. \end{split}$$

Define

$$\begin{split} \overline{T} &= \sup\{t \geq 0: \quad \overline{\delta}_i(s, \alpha_*) > \alpha_* > \overline{\delta}_i(-s, 0), \quad \forall s \in (0, t), \quad \forall i = \overline{0, q}\}, \\ \underline{T} &= \sup\{t \geq 0: \quad \underline{\delta}_i(s, \alpha_*) > \alpha_* > \underline{\delta}_i(-s, 0), \quad \forall s \in (0, t), \quad \forall i = \overline{0, q}\} \end{split}$$

Our strategy is to prove that  $\overline{T} = \underline{T} = +\infty$ . By contradiction, suppose that  $\overline{T} = +\infty$  or  $\underline{T} = +\infty$ , then there exists  $i \in \{0, q\}$  such that one of the following cases is true

$$\overline{\delta}_i(\overline{T}, \alpha_*) = \alpha_*, \qquad \overline{\delta}_i(-\overline{T}, 0) = \alpha_*, \\ \underline{\delta}_i(\underline{T}, \alpha_*) = \alpha_*, \qquad \underline{\delta}_i(-\underline{T}, 0) = \alpha_*.$$

Then  $\overline{T}$  and  $\underline{T}$  satisfies

$$\frac{d}{dt} \Big[ \exp(qt)\overline{\delta}_i(\overline{T}, \alpha_*) \Big] \le 0 \quad \text{and} \quad \frac{d}{dt} \Big[ \exp(qt)\overline{\delta}_i(-\overline{T}, 0) \Big] \le 0, \tag{19}$$

$$\frac{d}{dt} \Big[ \exp(qt) \underline{\delta}_i(\underline{T}, \alpha_*) \Big] \le 0 \quad \text{and} \quad \frac{d}{dt} \Big[ \exp(qt) \underline{\delta}_i(-\underline{T}, 0) \Big] \le 0.$$
(20)

Let be  $T_* = \min\{\underline{T}, \overline{T}\}$ , remark that

$$\psi_i(t,\alpha) = \alpha \exp\left(\int_0^t \sigma(I_{i,q}\mathcal{B}(\nu)I_{i,q})d\nu\right) + \psi_i(t,0), \quad \forall t \in \mathbb{R}, \quad \forall I \in \mathcal{K}_q.$$

By hypothesis [H1] there exists  $\omega > 0$  such that for all  $i = \overline{1, q}$  we have

$$\sup\left\{\tau_{i}\int_{0}^{t}\sigma(\mathcal{B}(\tau_{i}\nu))d\nu,\quad\forall t\in\mathbb{R}\right\}<\omega,$$
$$\sup\left\{\tau_{i}\int_{0}^{t}\sigma(I_{i,q}\mathcal{B}(\tau_{i}\nu)I_{i,q})d\nu,\quad\forall t\geq0\right\}<\omega0$$

Then for every  $i = \overline{1, q}$  we obtain the following estimation

$$\begin{split} \psi_i(\tau_i t, \alpha) &\leq \exp(\omega)\alpha + \psi_i(\tau_i t, 0), \quad \forall t \in (0, T_*), \\ \psi_0(\tau_i t, \alpha) &\leq \exp(\omega)\alpha + \psi_0(\tau_i t, 0), \quad \forall t \in (0, T_*), \\ \psi_0(-\tau_i t, \alpha) &\leq \exp(\omega)\alpha + \psi_0(-\tau_i t, 0), \quad \forall t \in (0, T_*). \end{split}$$

Put

$$\psi(I_{i,q}, \tau_i, t) := \psi_i(\tau_i t, 0), \quad \forall i = \overline{1, q}.$$

By definition of the function  $\psi_i(t,0)$  we have  $\psi(-I_{i,q},\tau,t) = -\psi(I_{i,q},\tau,t)$ . By definition of  $T_*$  we have: for all  $I \in \mathcal{K}_q$  there exists  $\tau_I$  such that  $\tau_{I_{i,q}} = \tau_{-I_{i,q}} = \tau_i$  and such that for all  $t \in (0, T_*)$  we have

$$\exp(\omega)\alpha_* + \psi(I,\tau_I,t) + \sigma(Iy(\tau_I t)) > \alpha_* > \psi(I,\tau_I,-t) + \sigma(Iy(\tau_I t))$$

and for all  $\tau \in \{-1, 1\}$  and  $I \in \{I_q, -I_q\}$  we have

$$\exp(\omega)\alpha_* + \psi(I,\tau,t) + \sigma(Iy(\tau t)) > \alpha_* > \psi(I,\tau,-t) + \sigma(Iy(\tau t)).$$

where we recall that  $\alpha_* = D^2/(\gamma(1-\gamma))$ . By Hypothesis [H2] there exists  $\beta > 0$  such that

$$\sup_{t \in \mathbb{R}} |\psi_0(\tau_i t) - \psi_i(\tau_i t)| < \beta, \quad \forall i = \overline{1, q}.$$

By Lemma 6 we have

$$\|\gamma y(t)\| < (\exp(\omega) + 1)\alpha_* + \beta, \quad \forall [-T_*, T_*].$$
 (21)

There exists  $\gamma_* > 0$  such that for all  $\gamma \in (0, \gamma_*)$  there exists  $D_{\gamma} > 0$  such that

$$\frac{\exp(\omega) + 1}{\gamma(1 - \gamma)} D_{\gamma}^2 + \beta < D_{\gamma}, \tag{22}$$

Choose  $\alpha_* = D_{\gamma}^2/(\gamma(1-\gamma))$  and use equation (21) to get

$$\|\gamma y(t)\| < \frac{\exp(\omega) + 1}{\gamma(1 - \gamma)} D_{\gamma}^2 + \beta < D_{\gamma}, \quad \forall [-T_*, T_*].$$

Then  $\|\gamma y(t)\| < D_{\gamma}$  for every  $t \in [-T_*, T^*]$ . In particular  $\|\gamma y(T_*)\| < D_{\gamma}$ and  $\|\gamma y(-T_*)\| < D_{\gamma}$ . By Lemma 5 and Equation (19), we have

$$\frac{d}{dt} \Big[ \exp(qt)\overline{\delta}_i(T_*, \alpha_*) \Big] > 0 \quad \text{and} \quad \frac{d}{dt} \Big[ \exp(qt)\overline{\delta}_i(-T_*, 0) \Big] > 0,$$
$$\frac{d}{dt} \Big[ \exp(qt)\underline{\delta}_i(T_*, \alpha_*) \Big] > 0 \quad \text{and} \quad \frac{d}{dt} \Big[ \exp(qt)\underline{\delta}_i(-T_*, 0) \Big] > 0.$$

We have obtain a contradiction with Equation (19). Then  $T_* = +\infty$  and we have

$$\overline{\delta}_i(t,\alpha_*) > \alpha_* > \overline{\delta}_i(-t,0) \quad \forall t \ge 0, \quad \forall i = \overline{0,q}, \\ \underline{\delta}_i(t,\alpha_*) > \alpha_* > \underline{\delta}_i(-t,0) \quad \forall t \ge 0 \quad \forall i = \overline{0,q}.$$

Finally, use again Equation (22) and Lemma 6 to deduce that

$$\|\gamma y(t)\| < \frac{\exp(\omega) + 1}{\gamma(1 - \gamma)} D_{\gamma}^2 + \beta < D_{\gamma}, \quad \forall t \in \mathbb{R}.$$

## 3 Proof of the Generalized main result

Proof of the Generalized main result.

Consider the Equation (1). By hypothesis of the Generalized main result there exist L such that

$$\max\{\|df\|_{\infty}, \|d^2f\|_{\infty}\} < L$$

Let be  $\epsilon > 0$  and consider the change of variable

$$\eta(t) = \epsilon^2 \Big[ \mu(\epsilon t) - \Phi^{\epsilon t}(x_0) \Big], \ t \in \mathbb{R}, \ \eta := (\eta_i)_{i=1}^n,$$

where  $\Phi^t := (\Phi_i^t)_{i=1}^n$  is the flow of the system (1). Use Taylor formula

$$\dot{\eta} = \epsilon^3 [\dot{\mu}(\epsilon t) - f(\mu(\epsilon t))] + \epsilon^2 df(\mu(\epsilon t))\eta + \eta^T \zeta_{\mu}(t)\eta, \quad \eta(0) = \mathbb{O}_n, \quad t \in \mathbb{R},$$

where he function  $\zeta_{\mu} : \mathbb{R} \times \mathbb{R}^{n \times n \times n}$  is a continuous function and satisfies  $\|\zeta_{\mu}\|_{\infty} \leq \|d^2 f\|_{\infty} < \epsilon L$ . By hypothesis the function  $\mu : \mathbb{R} \to \mathbb{R}^n$  is a leader trajectory of f such that  $\mu(0) = x_0$ . Thanks to Theorem 4 there exists  $\gamma_* > 0$  such that for all  $\epsilon L < \gamma_*$  there exists  $D_{\epsilon,L}$  such that

$$\|\mu(s) - \Phi^s(x_0)\| = \frac{1}{\gamma} \|\eta(t)\| < \frac{1}{\epsilon^2} D_{\epsilon,L}, \quad \forall s \in \mathbb{R}.$$

# 4 Leader trajectory and Rotation vector formula

We show in this Section that for any periodic function g there exists a constant vector  $\rho \in \mathbb{R}^q$  such that the function  $t \mapsto t\mathbb{1}_q$  is a leader trajectory of the function  $t \mapsto I_{\rho}^{-1}g(\rho t)$ . In order to prove the result let us introduce the following notations; Let  $g = (g_i)_{i=1}^q : \mathbb{R}^q \to \mathbb{R}^q$  be a  $C^1$  periodic function. Let  $\Psi_i[g] : \mathbb{R}^q \times \mathbb{R} \to \mathbb{R}$  be the functions defined by

$$\Psi_{i}[g](z,t) := \int_{0}^{t} A_{i}(z,s) \exp\Big(\int_{s}^{t} \sigma\Big(I_{i,q}I_{z}^{-1}dg(\nu z)I_{z}I_{i,q}\Big)d\nu\Big)ds,$$
$$A_{i}(z,s) := \sigma\Big(I_{i,q}(\mathbb{1} - I_{z}^{-1}g(sz))\Big),$$
(23)

where  $I_z$  is the diagonal matrix of diagonal z. For every  $i = \overline{0, q}$  the function  $\Psi_i[g](z, t)$  is the solution of the following one-dimensional differential equation

$$\dot{\psi} := \sigma \Big( I_{i,q} (\mathbb{1}_q - I_z^{-1} g(tz)) \Big) + \sigma \Big( I_{i,q} I_z^{-1} dg(tz) I_z I_{i,q} \Big) \psi$$

$$\psi(0) = 0, \quad t \in \mathbb{R}.$$
(24)

**Proposition 7.** Let  $g = (g_i)_{i=1}^q : \mathbb{R}^q \to \mathbb{R}^q$  be a  $C^1$  periodic function such

$$\exists \beta, \gamma > 0: \inf \{g_i(z), \ z \in \mathbb{R}^q, \ i = \overline{1, q}\} > 1 + \beta \quad and \quad \left\| dg \right\|_{\infty} < \gamma.$$

Then there exits  $\gamma_* > 0$  such that for all  $\gamma \in (0, \gamma_*)$  there exists  $\rho \in \mathbb{R}^q$  such that the function  $t \mapsto t \mathbb{1}_q$  is a leader trajectory of the function  $t \mapsto I_{\rho}^{-1}g(\rho t)$ . Further,  $\rho$  is solution of the following Equation

$$\lim_{t \to +\infty} \frac{1}{t} \Big[ \Psi_0[g](\rho, t) - \Psi_i[g](\rho, t) \Big] = 0, \quad \forall i = \overline{0, q}.$$

For the proof of the following Lemma we refer to [3].

**Lemma 8.** Let be  $q \in \mathbb{N}^*$  and  $g, h : \mathbb{R}^q \to \mathbb{R}$  be a periodic function. Let be  $z \in \mathbb{R}^q$  and suppose that Suppose that

$$\lim_{t \to +\infty} \frac{1}{t} \int_0^t g(zs) ds \le 0,$$

then there exists  $\lambda_z \in \mathbb{R}$  and a bounded function  $p_z : \mathbb{R} \to \mathbb{R}$  such that the following differential equation

$$\dot{\psi} = h(zt) + g(zt)\psi, \quad \psi(0) = 0,$$

admits a solution  $\psi : \mathbb{R} \to \mathbb{R}$  of the form

$$\psi(t) = \lambda_z t + p_z(t), \quad \forall t \in \mathbb{R}.$$

In particular, the Lemma is true when g(z) = 0 for all  $z \in \mathbb{R}^q$ . Further there exists  $\omega > 0$  such that

$$||p_z|| < \omega ||g||_{\infty}, \quad \forall z = (z_i)_{i=1}^q, \quad z_i \ge 1, \quad \forall i = \overline{1, q}.$$

Proof of the Proposition 7. By using Lemma 8, for every  $i = \overline{0, q}$  let  $\tau_i : \mathbb{R}^q \to \{-1, 1\}$  defined by

$$\tau_i(z) = -\operatorname{Sign}(\lambda_{i,z}), \quad \lambda_{i,z} := \lim_{t \to +\infty} \frac{1}{t} \int_0^t \sigma\Big(I_{i,q} I_z^{-1} dg(\nu z) I_z I_{i,q}\Big) d\nu, \quad (25)$$

where the function Sign :  $\mathbb{R} \to \{-1, 1\}$  is defined by Sign(t) = 1 if  $t \ge 0$  and Sign(t) = -1 if t < 0. Remark that for all  $\tau \in \{-1, 1\}$  we have

$$\sigma\left(I_z^{-1}dg(\tau tz)I_z\right) = \tau \frac{d}{dt} \Big[I_z^{-1}g(\tau tz)\Big],$$
  
$$\sigma\left(I_{i,q}I_z^{-1}dg(\tau tz)I_{i,q}I_z\right) = \tau_i(z)\frac{d}{dt} \Big[I_z^{-1}g(\tau tz)\Big] - 2\frac{\partial}{\partial x_i}g(\tau zt).$$

By definition of the coefficients  $\tau_i(z)$  in Equation (25) and by Lemma 8, there exists  $\omega > 0$  such that

$$\int_{0}^{t} \sigma \left( I_{z}^{-1} dg(\tau \nu z) I_{z} \right) d\nu < \gamma \omega, \quad \forall z \in \mathbb{R}^{q}, \quad t \in \mathbb{R},$$
$$\int_{0}^{t} \tau_{i}(z) \sigma \left( I_{i,q} I_{z}^{-1} dg(\tau_{i}(z) \nu z) I_{i,q} I_{z} \right) d\nu < \gamma \omega, \quad \forall z \in \mathbb{R}^{q}, \quad \forall t \ge 0.$$
(26)

In order to prove that the function  $t \mapsto t \mathbb{1}_q$  is a leader trajectory of the function  $t \mapsto I_{\rho}^{-1}g(\rho t)$ , we must prove that

$$\sup_{t \in \mathbb{R}} \left| \Psi_0[g](\rho, \tau_i(\rho)t) - \Psi_i[g](\rho, \tau_i(\rho)t) \right| < +\infty, \ \forall i = \overline{0, q}$$

Again, by Lemma 8, it is sufficient to prove that

$$\lim_{t \to +\infty} \frac{1}{t} \Big[ \Psi_0[g](\rho, t) - \Psi_i[g](\rho, t) \Big] = 0, \ \forall i = \overline{0, q}.$$

$$(27)$$

In order to simplify, for every  $i=\overline{0,q},\,z\in\mathbb{R}^q$  and  $t,s\geq 0$  denote

$$\begin{aligned} \theta_i(z,t,s) &:= \exp\left(\gamma \tau_i(z) \int_s^t \sigma \left(I_{i,q} I_z^{-1} dg(\tau_i(z)\nu z) I_z I_{i,q}\right) d\nu\right), \\ \theta_0^i(z,t,s) &:= \exp\left(\gamma \tau_i(z) \int_s^t \sigma \left(I_z^{-1} dg(\tau_i(z)\nu z) I_z\right) d\nu\right), \\ \underline{\theta}_i(z,t,s) &:= \theta_0^i(z,t,s) - \theta_i(z,t,s), \\ \Theta_i(z,t) &:= \int_0^t \theta_0^i(z,t,s) + \frac{1}{2} \underline{\theta}_i(z,t,s) \sigma \left(I_{i,q}(\mathbb{1}_q - I_z^{-1} g(\tau_i(z)sz))\right) ds. \end{aligned}$$

For every fixed  $k \in \mathbb{N}$ , define the function  $\Gamma_k = (\Gamma_{i,k})_{i=1}^q : \mathbb{R}^q \to \mathbb{R}^q$  by

$$\Gamma_{i,k}(z) := \frac{1}{\Theta_i(z,k)} \int_0^k \theta_0^i(z,k,s) g_i(\tau_i(z)sz) ds, \quad z \in \mathbb{R}^q.$$

Equation (27) is satisfied if there exists  $\rho \in \mathbb{R}^q$  and a sequence  $u_k \to +\infty$  such that  $\lim_{k\to+\infty} \Gamma_{u_k}(\rho) = \rho$ . We have

$$\Theta_i(z,k)\Gamma_{i,k}(z) - \int_0^k \theta_0^i(z,k,s)g_i(\tau_i(z)sz)ds = 0.$$

By the Mean value Theorem, for all  $z\in\mathbb{R}^q$  and all  $k\in\mathbb{N}$  there exists  $s^z_{i,k}\geq 0$  such that

$$\theta_0^i(z,k,s_{i,k}^z) + \frac{1}{2}\underline{\theta}_i(z,k,s_{i,k}^z)\sigma\Big(I_{i,q}(\mathbb{1}_q - I_z^{-1}g(s_{i,k}^z\tau_i(z)z))\Big)\Gamma_{i,k}(z)$$

$$= \theta_0^i(z, k, s_{i,k}^z) g_i(s_{i,k}^z \tau_i(z) z).$$
  

$$\Gamma_{i,k}(z) := \frac{g_i(s_{i,k}^z \tau_i(z) z)}{1 + \zeta_i(z, k)},$$
(28)

where

That follows

$$\zeta_{i}(z,k) = \frac{1}{2} \frac{\underline{\theta}_{i}(z,k,s_{i,k}^{z})}{\theta_{0}^{i}(z,k,s_{i,k}^{z})} \sigma \Big( I_{i,q}(\mathbb{1}_{q} - I_{z}^{-1}g(s_{i,k}^{z}\tau_{i}(z)z)) \Big) \\ = \frac{1}{2} \Big[ 1 - \frac{\theta_{i}(z,k,s_{i,k}^{z})}{\theta_{0}^{i}(z,k,s_{i,k}^{z})} \Big] \sigma \Big( I_{i,q}(\mathbb{1}_{q} - I_{z}^{-1}g(s_{i,k}^{z}\tau_{i}(z)z)) \Big).$$
(29)

Let be L > 1 and denote

$$\mathcal{V}_L = \Big\{ z = (z_i)_{i=1}^q \in \mathbb{R}^q : \| z \| \le (L-1) \| g \|_{\infty} \text{ and } z_i > 1, \ \forall i = \overline{1, q} \Big\}.$$

By Equations (26), we have

$$\begin{split} \theta_i(z,t,s) &< \exp(\gamma \omega), \quad \forall i = \overline{1,q}, \quad \forall t \geq s \geq 0, \quad \forall z \in \mathcal{V}_L, \\ \exp(-\gamma \omega) &< \theta_0^i(z,t,s) < \exp(\gamma \omega), \quad \forall i = \overline{1,q}, \quad \forall t \geq s \geq 0, \quad \forall z \in \mathcal{V}_L. \end{split}$$

We have for all  $t \ge 0$  and all  $z \in \mathcal{V}$  the following estimation

$$|\zeta_i(z,t)| \le \frac{1}{2}q[\exp(2\gamma\omega) - 1] \Big[ ||z|| + ||g||_{\infty} \Big] \le \frac{1}{2}qL[\exp(2\omega) - 1] ||g||_{\infty}.$$

Again, by the hypotheses of the present Proposition, there exists  $\beta>0$  such that

$$\inf\{g_i(z), \quad z \in \mathbb{R}^q, \ i = \overline{1,q}\} > 1 + \beta$$

There exists  $\gamma_* > 0$  such that for all  $\gamma \in (0, \gamma_*)$  there exists  $L_{\gamma} > 0$  such that

$$L_{\gamma} > \frac{2}{2 - qL_{\gamma}(\exp(2\gamma\omega) - 1) \|g\|_{\infty}} \text{ and } \frac{2(1+\beta)}{2 + qL_{\gamma}(\exp(2\gamma\omega) - 1) \|g\|_{\infty}} > 1$$

By Equations (28) and (29), for a fixed  $\gamma \in (0, \gamma_*)$  and  $L_{\gamma} > 1$ , the compact and convex set  $\mathcal{V}_{L_{\gamma}}$  satisfies

$$\Gamma_k(\mathcal{V}_{L_{\gamma}}) \subset \mathcal{V}_{L_{\gamma}}, \quad \forall k \in \mathbb{N}.$$

and the function  $\Gamma_k$  is  $C^1$  on  $\mathcal{V}_{L_{\gamma}}$ . By Brouwer fixed-point theorem, for all  $k \in \mathbb{N}$  there exists  $\rho_k \in \mathcal{V}_{L_{\gamma}}$  such that  $\Gamma_k(\rho_k) = \rho_k$ . There exists  $\rho \in \mathcal{V}_{L_{\gamma}}$  and there exists a sequence  $(u_k)_{k \in \mathbb{N}} \subset \mathbb{N}$  such that

$$\lim_{k \to +\infty} \rho_{u_k} = \rho \text{ and } \lim_{k \to +\infty} \frac{1}{u_k} \Gamma_{u_k}(\rho, u_k) = \rho.$$

Otherwise,

$$\lim_{k \to +\infty} \frac{1}{u_k} [\Psi_0[g](\rho, u_k) - \Psi_i[g](\rho, u_k)] = 0.$$

Hence, Equation (27) is satisfied. By Equation (26) of this proof, we deduce that the the function  $t \mapsto t \mathbb{1}_q$  is a leader trajectory of the function  $t \mapsto I_{\rho}^{-1}g(\rho t)$ .

## 5 Proof of the Main result

Proof of the Main Result. Let x be the solution of the system 1 and suppose that f is periodic. By periodicity, there exists  $\omega > 0$  such that

$$\max\left\{\left\|f\right\|_{\infty}, \left\|df\right\|_{\infty}, \left\|df\right\|_{\infty}\right\} < \omega.$$
(30)

Let be  $c, \gamma > 0$ , and consider the following change of variable

$$\tilde{x}(t) := ct\mathbb{1}_n + x(\gamma t) - x_0, \quad \forall t \in \mathbb{R}.$$

System 1, becomes

$$\tilde{x} = c + \gamma \tilde{f}_{x_0}(\tilde{x}, t), \quad \tilde{x}(0) = \mathbb{O}_q, \quad t \in \mathbb{R}.$$
  
 $\tilde{f}_{x_0}(z, z_{n+1}) = c + \gamma f(z + x_0 - cz_{n+1}\mathbb{1}_n), \quad \forall z = (z_i)_{i=1}^{n+1}$ 

By Equation (30) we have

$$\begin{aligned} \forall (z,t) \in \mathbb{R}^n \times \mathbb{R} : &\max\left\{ \|\frac{d}{dz} \tilde{f}_{x_0}(z,t)\|, \|\frac{d^2}{dz^2} \tilde{f}_{x_0}(z,t)\| \right\} < \gamma \omega, \\ &\tilde{f}_{x_0}(z,z_{n+1}) > c - \gamma \omega, \quad \forall z = (z_i)_{i=1}^{n+1}. \end{aligned}$$

Le be  $\beta > 0$  fixed then for<br/>a all  $c > 1 + \beta + \gamma \omega$  we have

$$\tilde{f}_{x_0}(z, z_{n+1}) > 1 + \beta, \quad \forall z = (z_i)_{i=1}^{n+1}.$$
 (31)

Thanks to the Proposition 7, there exist  $\gamma > 0$  and  $\rho_* \in \mathbb{R}^n$  such that the function  $\mu(t) := t\mathbb{1}_n$  is a leader trajectory of the function  $I_{\rho_*}^{-1}\tilde{f}_{x_0}$ . Since  $\tilde{x}_0 = \mathbb{O}_n$  and by the Generalized main result we have

$$\sup_{t\in\mathbb{R}}||t\mathbb{1}_n - I_{\rho_*}^{-1}\tilde{x}(t)|| < +\infty,$$

by definition of  $\tilde{x}$  we have

.

$$\sup_{t\in\mathbb{R}}\left\|\frac{1}{\gamma}[(\rho_*-c)t\mathbb{1}_n-x_0+x(t)\|<+\infty,\right.$$

Otherwise, the solution x of the Equation (1) admits a rotation vector  $\rho := (\rho_* - c)/\gamma$ . Again by Proposition 7 and by the uniqueness of the rotation vector the vector we deduce that  $\rho_*$  is the unique solution of the following equation

$$\lim_{t_t \rho \to \infty} \frac{1}{t} \left[ \Psi_0[\tilde{f}_{x_0}](\rho_*, t) - \Psi_i[\tilde{f}_{x_0}](\rho_*, t) \right] = 0, \quad \forall i = \overline{1, n}$$

where the functions  $\Psi_i[\tilde{f}_{x_0}]$  are defined in Equation (23) of the previous Section. Since  $\rho_* = \rho\gamma + c$  and by using the change of variable  $t \to \gamma t$  then it is equivalent

$$\lim_{t_t o + \infty} \frac{1}{t} \Big[ \Psi_0[f_{x_0}](\rho_*, c, t) - \Psi_i[f_{x_0}](\rho_*, c, it) \Big] = 0, \quad \forall i = \overline{1, n},$$

where

$$\Psi_{i}[f_{x_{0}}](\rho, c, t) := \int_{0}^{t} A_{i}(\rho, s) \exp\Big(\int_{s}^{t} \sigma\Big(I_{i,q}I_{\rho_{*}}^{-1}df(\nu\rho + x_{0})I_{\rho_{*}}I_{i,q}\Big)d\nu\Big)ds,$$
$$A_{i}(\rho, c, s) := \sigma\Big(I_{i,q}I_{\rho_{*}}^{-1}(\rho\mathbb{1} - f(s\rho + x_{0}))\Big),$$

By Equation (31) the coefficient c is arbitrary in  $(1 + \beta + \gamma \omega, +\infty)$ , since  $\rho$  is unique, we get

$$\lim_{c \to \infty} \lim_{t \to +\infty} \frac{c}{t} \Big[ \Psi_0[\tilde{f}_{x_0}](\rho_*, c, t) - \Psi_i[\tilde{f}_{x_0}](\rho_*, c, t) \Big] = 0, \quad \forall i = \overline{1, n}.$$

Finally, by Lemma 8 of the previous Section, we deduce that  $\rho$  is the unique solution of the following Rotation vector formula

$$\lim_{t \to +\infty} \frac{1}{t} \Big[ \Psi_0[f_{x_0}](\rho, t) - \Psi_i[f_{x_0}](\rho, t) \Big] = 0, \quad \forall i = \overline{1, n},$$

where

$$\Psi_i[f_{x_0}](\rho,t) := \int_0^t A_i(\rho,s) \exp\Big(\int_s^t \sigma\Big(I_{i,q}df(\nu\rho + x_0)I_{i,q}\Big)d\nu\Big)ds,$$
$$A_i(\rho,s) := \sigma\Big(I_{i,q}(\rho - f(s\rho + x_0))\Big).$$

### 6 Conclusion and open problem

We have proved that for any periodic field f and any initial condition the solution of the system admits a unique rotation vector which is the solution of the Rotation vector formula (4). We conjecture that the solution x of initial condition  $x_0$  and rotation vector  $\rho_x$  is periodic on the torus if and only if there exists T > 0 such that  $T\rho_x \in \mathbb{Q}^n$  and  $\tilde{f}_0(\rho, T) - \tilde{f}_i(\rho, T) = 0$  for every  $i = \overline{1, q}$ . Consider the perturbed system

$$\dot{x} = f(x) + \zeta(t), \quad t \in \mathbb{R}, \quad x(0) = x_0 \in \mathbb{R}^n,$$

where  $\zeta : \mathbb{R} \to \mathbb{R}^n$  is an integrable function. By the definition of the diagonal matrix  $I \in \mathcal{K}_n^+$  we deduce that if f is periodic then the rotation vector is locally constant for a small perturbation if and only if

$$\sup_{t\in\mathbb{R}}\left\|\int_{0}^{t}\exp\left(\int_{s}^{t}\sigma\left(Idf(\rho\nu+x_{0})I\right)d\nu\right)I\zeta(s)ds\right\|<+\infty,\ \forall I\in\mathcal{K}_{n}^{+}.$$

which gives information about the Arnold Tongue. The leader trajectory defined in this article is an approximated solution.

### References

- [1] A.M. Fink, Almost Periodic Functions. L.N.M., 1974, Springer
- [2] PE. Kloeden., Dynamics of a class of ODEs more general than almost periodic, Nonlinear Analysis: Theory, Methods & Applications, 2011, 74, 7, 2695 - 2719
- [3] D.N. Cheban, Asymptotically Almost Periodic Solutions of Differential Equations, Hindawi Publishing Corporation, New York - Cairo, 2009, 203
- [4] T. Saito, On dynamical systems in n-dimensional torus,

Proceedings of the Symposium on Differential Equations and Dynamical Systems. Lecture Notes in Mathematics, 1971, 206, 18–19

[5] Corduneanu, B., Almost periodic oscillations and waves., New York; London: Springer, ISBN 9780387098180 (hbk.), 2008

- [6] M. Zhang and G. Zong. Almost Periodic Solutions for Stochastic Differential Equations Driven By G-Brownian Motion, Communications in Statistics - Theory and Methods, Taylor & Francis, 2015, 44, 11, 2371-2384
- [7] Y. Yi., On almost automorphic oscillations. Differences and differential equations, Fields Inst. Commun. Amer. Math. Soc., Providence, RI, 2004, 42, 75-99
- [8] Oukil, W. and Kessi, A. and Thieullen, Ph., Synchronization hypothesis in Winfree model, Dynamical Systems, Taylor & Francis, 2016,
- [9] Y. L. Maistrenko and O. V. Popovych and P. A. Tass, Desynchronization and chaos in the Kuramoto model, 2005, Springer Berlin Heidelberg, Berlin, Heidelberg, 285–306, 978-3-540-31520-9,
- [10] G. Giacomin and K. Pakdaman and X. Pellegrin, Global attractor and asymptotic dynamics in the Kuramoto model for coupled noisy phase oscillators, Nonlinearity, 25, 5, 1247, 2012,
- [11] P. Goel and G. B. Ermentrout, Synchrony, stability, and firing patterns in pulse-coupled oscillators,

Physica D, 2002, 163, 191-216

- [12] G. B. Ermentrout, Stable Periodic Solutions to Discrete and Continuum Arrays of Weakly Coupled Nonlinear Oscillators, SIAM Journal on Applied Mathematics, 52, 6, 1665–1687, 1992
- [13] G. B. Ermentrout and M. Pascal and B. S. Gutkin, The Effects of Spike Frequency Adaptation and Negative Feedback on the Synchronization of Neural Oscillators, Neural Computation, 13, 6, 1285–1310, 2001
- [14] R. E. Mirollo and S. H. Strogatz, The spectrum of the locked state for the Kuramoto model of coupled oscillators, Physica D: Nonlinear Phenomena, 205, 1-4, 249 - 266, 2005, 0167-2789,
- [15] R. E. Mirollo and S. H. Strogatz, Synchronization of Pulse-Coupled Biological Oscillators,
   SIAM Journal on Applied Mathematics, 1990, 50, 1645–1662
- [16] J.T. Ariaratnam and S.H. Strogatz, Phase Diagram for the Winfree Model of Coupled Nonlinear Oscillators,

Phys. Rev. Lett, 2001, 86, 4278

- [17] L. Basnarkov and V. Urumov, Critical exponents of the transition from incoherence to partial oscillation death in the Winfree model, Journal of Statistical Mechanics: Theory and Experiment, 2009, 10, P10014, 2009,
- [18] N. Chopra and M. W. Spong, IEEE Transactions on Automatic Control, On Exponential Synchronization of Kuramoto Oscillators, 2009, 54, 2, 353–357
- [19] Cooray, G., The Kuramoto Model, U.U.D.M. project report. Uppsala University, Department of Mathematics., 2008
- [20] F. Giannuzzi and D. Marinazzo and G. Nardulli and M. Pellicoro and S. Stramaglia, Phase diagram of a generalized Winfree model,

Phys. Rev. E, 2007, 75, 051104

[21] A. Jadbabaie and N. Motee and and M. Barahona, On the stability of the Kuramoto model of coupled nonlinear oscillators,

Proceedings of the American Control Conference, 2004, R4296-4301

[22] A. T. Winfree, Biological rhythms and the behavior of populations of coupled oscillators,

J. Theor. Bio, 1967, 16, 15–42

[23] D.D. Quinn and R.H. Rand and S.H. Strogatz, Singular unlocking transition in the Winfree model of coupled oscillators,

Physical Review E, 2007, 75, 036218

[24] G. Nardulli and D. Marinazzo and M. Pellicoro and S. Stramaglia, Phase shifts between synchronized oscillators in the Winfree and Kuramoto models,

InterJournal Complex Systems, 1408, 2004

[25] O. V. Popovych and Y. L. Maistrenko and P. A. Tass, Phase chaos in coupled oscillators,

Phys. Rev. E, 71, 6, 065201, 4, 2005

[26] D. Pazó and E. Montbrió, Low-dimensional Dynamics of Populations of Pulse-Coupled Oscillators,

Phys. Rev. X , 2014, 4, 011009

- [27] S. Louca and F. M. Atay, Spatially structured networks of pulse-coupled phase oscillators on metric spaces, Discrete and Continuous Dynamical Systems, 34, 9, 3703–3745, 2014,
- [28] J. A. Acebrón and L. L. Bonilla and C. J. Pérez Vicente and F. Ritort and R. Spigler, The Kuramoto model: A simple paradigm for synchronization phenomena,

Rev. Mod. Phys., 77, 137-185, 2005

- [29] Y. P. Choil and S.Y Ha and S.B. Yun, Physica D: Nonlinear Phenomena, 1, 32–44, Complete synchronization of Kuramoto oscillators with finite inertia, 240, 2011
- [30] Y. Kuramoto, International Symposium on Mathematical Problems in Theoretical Physics,

Lecture Notes in Physics, Springer, New York, 1975, 39, 420–20

- [31] S.Y. Ha and J. Park and S. W. Ryoo, Emergence of phase-locked states for the Winfree model in a large coupling regime,
   Discrete and Continuous Dynamical Systems, Series A, 35, 2015, 8, 3417– 3436
- [32] S Y. Ha and D. Ko and J. Park and S. W. Ryoo, Emergent dynamics of Winfree oscillators on locally coupled networks,

Journal of Differential Equations, 260, 2016, 5, 4203 – 4236

- [33] Y. Kuramoto and D. Battogtokh, Coexistence of coherence and incoherence in nonlocally coupled phase oscillators, Nonlinear Phenom. Complex Syst, 2002, 5, 380
- [34] O. E. Omel'chenko, Coherence-incoherence patterns in a ring of nonlocally coupled phase oscillators, Nonlinearity, 2013, 26, 2469–2498
- [35] M. J. Panaggio and D. M. Abrams, Chimera states: coexistence of coherence and incoherence in networks of coupled oscillators, Nonlinearity, 2015, 28, R67
- [36] Ermentrout, G. B. and Kopell, N., Multiple pulse interactions and averaging in systems of coupled neural oscillators, Journal of Mathematical Biology, 1991, 29, 3, 195–217, 1432–1416

- [37] G. Wulfram and W. Werner, Spiking Neuron Models: An Introduction, 2002, 0521890799, Cambridge University Press, New York, NY, USA
- [38] Cumin, D. and Unsworth, C.P., Physica D: Nonlinear Phenomena, English, 2, 181–196, Generalising the Kuramoto model for the study of neuronal synchronisation in the brain, 226, 2007
- [39] Kuramoto, Y., Berlin, Heidelberg,
   12977810, 978-3-642-69691-6, 1, Springer Berlin Heidelberg, Chemical Oscillations, Waves, and Turbulence, 19, 1984
- [40] Batista, C. A. S. and Lameu, E. L. and Batista, A. M. and Lopes, S. R. and Pereira, T. and Zamora-López, G. and Kurths, J. and Viana, R. L., Phase synchronization of bursting neurons in clustered small-world networks,

Phys. Rev. E, 86, 1, 016211, 12, 2012

- [41] HODGKIN, A. L. and HUXLEY, A. F., The Journal of physiology, 4, 500–544, The Physiological Society, A quantitative description of membrane current and its application to conduction and excitation in nerve., 117, 1952
- [42] Winfree, A. T., The geometry of biological time, Biomathematics, Springer Verlag, New York, 0-387-09373-7, 1980
- [43] Kolb, B. and Whishaw, I.Q., Cerveau & comportement, 9782744501371, Neurosciences & cognition, 2002, De Boeck
- [44] Françoise, J-P., Oscillations en biologie : analyse qualitative et modèles, Mathématiques et applications, Springer, Berlin, Heidelberg, New York, 3-540-25152-9, 2005