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1 **SCATTERING IN THE WEIGHTED L^2 -SPACE FOR A 2D NONLINEAR**
2 **SCHRÖDINGER EQUATION WITH INHOMOGENEOUS EXPONENTIAL**
3 **NONLINEARITY**

4 ABDELWAHAB BENSOUILAH, VAN DUONG DINH, AND MOHAMED MAJDOUB

ABSTRACT. We investigate the defocusing inhomogeneous nonlinear Schrödinger equation

$$i\partial_t u + \Delta u = |x|^{-b} \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 \right) u, \quad u(0) = u_0, \quad x \in \mathbb{R}^2,$$

with $0 < b < 1$ and $\alpha = 2\pi(2 - b)$. First we show the decay of global solutions by assuming that the initial data u_0 belongs to the weighted space $\Sigma(\mathbb{R}^2) = \{u \in H^1(\mathbb{R}^2) : |x|u \in L^2(\mathbb{R}^2)\}$. Then we combine the local theory with the decay estimate to obtain scattering in Σ when the Hamiltonian is below the value $\frac{2}{(1+b)(2-b)}$.

5 1. INTRODUCTION AND MAIN RESULT

This paper is concerned with the scattering theory for the following initial value problem

$$\begin{cases} i\partial_t u + \Delta u &= |x|^{-b} \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 \right) u, \\ u(0) &= u_0, \end{cases} \quad (1.1)$$

6 where $u = u(t, x)$ is a complex-valued function in space-time $\mathbb{R} \times \mathbb{R}^2$, $0 < b < 1$ and $\alpha =$
7 $2\pi(2 - b)$.

8 The classical nonlinear Schrödinger equation ($b = 0$) with pure power or exponential
9 nonlinearities arises in various physical contexts, as for example the self trapped beams in
10 plasma, the propagation of a laser beam, water waves at the free surface of an ideal fluid and
11 plasma waves (see [21]).

12 From the mathematical point of view, the classical NLS equation, i.e., problem (1.1) with
13 $b = 0$, has attracted considerable attention in the mathematical community and the well-
14 posedness theory as well as the scattering has been extensively studied, see for instance
15 [2, 3, 7, 9, 19, 22]. We refer the reader to [8, 33] and references therein for more properties
16 and information on nonlinear Schrödinger equations.

17 In particular, in [9] a notion of criticality was proposed and the authors established in
18 both subcritical and critical regimes the existence of global solutions in the functional space
19 $C(\mathbb{R}, H^1(\mathbb{R}^2)) \cap L_{loc}^4(\mathbb{R}, W^{1,4}(\mathbb{R}^2))$. Later on in [19], the scattering in the energy space was

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1 obtained in the subcritical case. Note that the critical case was investigated in [3] where the
2 scattering is proved in the radial framework.

3 The situation in the case $b > 0$ is less understood. Recently, in [5] the authors established
4 the global well-posedness in the energy space for $0 < b < 1$. A natural question to ask then
5 is the long time behavior of global solutions, that is the scattering. This means that every
6 global solution of (1.1) approaches solutions to the associated free equation

$$i\partial_t v + \Delta v = 0, \quad (1.2)$$

7 in the energy space H^1 as $t \rightarrow \pm\infty$. The main difficulty is how to obtain the interaction
8 Morawetz inequality? Recall that the interaction Morawetz inequality is nothing but the
9 convolution of the classical one with the mass density. This in particular leads to a priori
10 global bound of the solution in $L_t^4(L_x^8)$ which is the main tool for the scattering in the energy
11 space (see for instance [3, 19, 24]). Note that the interaction Morawetz inequalities were first
12 established for the NLS with power-type nonlinearity, and the proof depends heavily on the
13 form of nonlinearity. Of course the proof can be easily adapted to more general homogeneous
14 nonlinearities. More precisely, for linear combination of powers it suffices that all the powers
15 are quadratic or higher with positive coefficients. The problem with singular weight (or for
16 non-homogeneous nonlinearity) is much more difficult and should be investigated separately.
17 For instance, it was noticed in [11] that the interaction Morawetz inequality for the NLS
18 with singular nonlinearity $N(x, u) = |x|^{-b}|u|^\alpha u$ may not hold due to the lack of momentum
19 conservation law.

20 This is why we restrict ourselves to initial data belonging to the weighted L^2 -space $\Sigma :=$
21 $H^1 \cap L^2(|x|^2 dx)$. Note that the scattering in Σ for the NLS with $N(x, u) = |x|^{-b}|u|^\alpha u$ was
22 considered by the second author in [10].

23 The scattering in the energy space will be investigated in a forthcoming paper, and we
24 believe that some ideas developed in [3] will be helpful.

25 **Remark 1.1.** *We stress that the two-dimensional nonlinear Klein-Gordon equation with pure*
26 *exponential nonlinearity was studied in [16, 18, 17], and a similar trichotomy based on the*
27 *energy was defined. Recently, M. Struwe [30, 31] was able to construct global smooth solution*
28 *for smooth initial data and prove the scattering [29].*

29 Before stating our main result, let us recall that solutions of (1.1) satisfy the conservation
30 of mass and Hamiltonian

$$\mathcal{M}(u(t)) := \|u(t)\|_{L^2}, \quad (1.3)$$

31

$$\mathcal{H}(u(t)) := \int |\nabla u(t, x)|^2 + \frac{1}{\alpha} \int \left(e^{\alpha|u(t, x)|^2} - 1 - \alpha|u(t, x)|^2 - \frac{\alpha^2}{2}|u(t, x)|^4 \right) \frac{dx}{|x|^b}. \quad (1.4)$$

32 Our main result is the following.

1 **Theorem 1.2.** *Let $u_0 \in \Sigma$ be such that $\mathcal{H}(u_0) < \frac{2}{(1+b)(2-b)}$. Then the corresponding global*
 2 *solution u of (1.1) satisfies $u \in L^4(\mathbb{R}, \mathcal{C}^{1/2})$ and there exist $u_0^\pm \in \Sigma$ such that*

$$\lim_{t \rightarrow \pm\infty} \|e^{-it\Delta} u(t) - u_0^\pm\|_\Sigma = 0.$$

3 Let us make some comments. First, we see that $\frac{2}{(1+b)(2-b)} \rightarrow 1$ as $b \rightarrow 0$. Thus our
 4 result extends the one in [19] for initial data in Σ . Second, the condition $\mathcal{H}(u) < \frac{2}{(1+b)(2-b)}$
 5 illustrates the interaction between the wave function u and the potential $|x|^{-b}$. More precisely,
 6 a sufficient condition for scattering is when the energy of the wave is less than a fixed amount
 7 depending on the sole parameter b that characterizes the weight function involved in the
 8 Hamiltonian of (1.1). Finally, a natural question that one could raise is the following: is the
 9 value $\frac{2}{(1+b)(2-b)}$ critical for scattering, in the sense that if the energy of the wave exceeds the
 10 latter quantity, would one get scattering?

11 **Remark 1.3.** *For all $0 < b < 1$, $\frac{8}{9} \leq \frac{2}{(1+b)(2-b)} < 1$.*

The proof of Theorem 1.2 follows a standard strategy for the classical NLS equation. We first derive a decaying property for global solutions by using the pseudo-conformation law. We then show two types of global bounds for the solution u and its weighted variant $(x + 2it\nabla)u$. More precisely, we will show that

$$\|u\|_{S^1(\mathbb{R})} < \infty, \quad \|(x + 2it\nabla)u\|_{S^0(\mathbb{R})} < \infty, \quad (1.5)$$

12 where

$$\|u\|_{S^1(\mathbb{R})} := \|u\|_{L^\infty(\mathbb{R}, H^1)} + \|u\|_{L^4(\mathbb{R}, W^{1,4})}, \quad \|u\|_{S^0(\mathbb{R})} := \|u\|_{L^\infty(\mathbb{R}, L^2)} + \|u\|_{L^4(\mathbb{R}, L^4)}.$$

13 The proof of these global bounds relies on the decaying property, the singular Moser-Trudinger
 14 inequality and the Log estimate. The main difficulty comes from the singular weight $|x|^{-b}$
 15 which does not belong to any Lebesgue space. To overcome this problem, we will take the
 16 advantage of Lorentz spaces. Note that $|x|^{-b} \in L^{\frac{2}{b}, \infty}(\mathbb{R}^2)$, where $L^{p, \infty}$ is the Lorentz space.
 17 Once these global bounds are established, the scattering in weighted L^2 space Σ follows easily.

18 This paper is organized as follows. In Section 2, we recall some useful tools needed in our
 19 problem. The pseudo-conformal law is derived in Section 3. The decaying property of global
 20 solutions in Lebesgue spaces is showed in Section 4. Sections 5 and 6 are devoted to the
 21 proofs of global bounds (1.5). We shall give the proof of our main result in Theorem 1.2 in
 22 Section 7.

2. USEFUL TOOLS

24 In this section, we collect some known and useful tools.

25 **Proposition 2.1** (Moser-Trudinger inequality [1]).

26 *Let $\alpha \in [0, 4\pi)$. A constant c_α exists such that*

$$\|\exp(\alpha|u|^2) - 1\|_{L^1(\mathbb{R}^2)} \leq c_\alpha \|u\|_{L^2(\mathbb{R}^2)}^2, \quad (2.1)$$

1 for all u in $H^1(\mathbb{R}^2)$ such that $\|\nabla u\|_{L^2(\mathbb{R}^2)} \leq 1$. Moreover, if $\alpha \geq 4\pi$, then (2.1) is false.

2 **Remark 2.2.** We point out that $\alpha = 4\pi$ becomes admissible in (2.1) if we require $\|u\|_{H^1(\mathbb{R}^2)} \leq$
 3 1 rather than $\|\nabla u\|_{L^2(\mathbb{R}^2)} \leq 1$. Precisely, we have

$$\sup_{\|u\|_{H^1} \leq 1} \|\exp(4\pi|u|^2) - 1\|_{L^1(\mathbb{R}^2)} < \infty, \quad (2.2)$$

4 and this is false for $\alpha > 4\pi$. See [25] for more details.

5 **Theorem 2.3.** [26] Let $0 < b < 2$ and $0 < \alpha < 2\pi(2 - b)$. Then, there exists a positive
 6 constant $C = C(b, \alpha)$ such that

$$\int_{\mathbb{R}^2} \frac{e^{\alpha|u(x)|^2} - 1}{|x|^b} dx \leq C \int_{\mathbb{R}^2} \frac{|u(x)|^2}{|x|^b} dx, \quad (2.3)$$

7 for all $u \in H^1(\mathbb{R}^2)$ with $\|\nabla u\|_{L^2(\mathbb{R}^2)} \leq 1$.

8 We point out that $\alpha = 2\pi(2 - b)$ becomes admissible in (2.3) if we require $\|u\|_{H^1(\mathbb{R}^2)} \leq 1$
 9 instead of $\|\nabla u\|_{L^2(\mathbb{R}^2)} \leq 1$. More precisely, we have

10 **Theorem 2.4.** [27] Let $0 < b < 2$. We have

$$\sup_{\|u\|_{H^1(\mathbb{R}^2)} \leq 1} \int_{\mathbb{R}^2} \frac{e^{\alpha|u(x)|^2} - 1}{|x|^b} dx < \infty \quad \text{if and only if} \quad \alpha \leq 2\pi(2 - b). \quad (2.4)$$

11 The following lemma will be very useful.

12 **Lemma 2.5.** Let $0 < b < 2$ and $\gamma \geq 2$. Then, there exists a positive constant $C = C(b, \gamma) > 0$
 13 such that

$$\int_{\mathbb{R}^2} \frac{|u(x)|^\gamma}{|x|^b} dx \leq C \|u\|_{H^1(\mathbb{R}^2)}^\gamma, \quad (2.5)$$

14 for all $u \in H^1(\mathbb{R}^2)$.

15 *Proof.* Note that

$$\| |x|^{-b} \|_{L^r(B)} < \infty \quad \text{if } b < \frac{2}{r}, \quad \| |x|^{-b} \|_{L^r(B^c)} < \infty \quad \text{if } b > \frac{2}{r}, \quad (2.6)$$

16 where $B = B(0, 1)$ is the unit ball in \mathbb{R}^2 and $B^c = \mathbb{R}^2 \setminus B$. Write

$$\int_{\mathbb{R}^2} |x|^{-b} |u(x)|^\gamma dx = \int_B |x|^{-b} |u(x)|^\gamma dx + \int_{B^c} |x|^{-b} |u(x)|^\gamma dx.$$

17 We have from the Sobolev embedding $H^1(\mathbb{R}^2) \subset L^q(\mathbb{R}^2)$ for any $q \in [2, \infty)$ that

$$\int_{B^c} |x|^{-b} |u(x)|^\gamma dx \leq \|u\|_{L^\gamma(\mathbb{R}^2)}^\gamma \lesssim \|u\|_{H^1(\mathbb{R}^2)}^\gamma.$$

18 The first term is estimated as follows. Since $0 < b < 2$, there exists $\varepsilon > 0$ small such that
 19 $b < \frac{2}{1+\varepsilon}$. We apply (2.6) with $r = 1 + \varepsilon$ and get

$$\int_B |x|^{-b} |u(x)|^\gamma dx \leq \| |x|^{-b} \|_{L^{1+\varepsilon}(B)} \| |u|^\gamma \|_{L^{\frac{1+\varepsilon}{\varepsilon}}(\mathbb{R}^2)} \lesssim \|u\|_{L^{\frac{(1+\varepsilon)\gamma}{\varepsilon}}(\mathbb{R}^2)}^\gamma \lesssim \|u\|_{H^1(\mathbb{R}^2)}^\gamma.$$

20 Combining the two terms, we prove the desired estimate. \square

Remark 2.6. *The inequality (2.5) fails for $b \geq 2$. Indeed, let $u \in \mathcal{D}(\mathbb{R}^2)$ (the space of smooth compactly supported functions) be a radial function such that $u(x) \equiv 1$ for $|x| \leq 1$. Then, $u \in H^1(\mathbb{R}^2)$ and*

$$\int_{\mathbb{R}^2} \frac{|u(x)|^\gamma}{|x|^b} dx \geq 2\pi \int_0^1 \frac{r dr}{r^b} = +\infty.$$

1 We also recall the so-called Gagliardo-Nirenberg inequalities and Sobolev embedding.

2 **Proposition 2.7** (Gagliardo-Nirenberg inequalities [12, 23]).

3 *We have*

$$\|u\|_{L^{m+1}} \lesssim \|u\|_{L^{q+1}}^{1-\theta} \|\nabla u\|_{L^p}^\theta, \quad (2.7)$$

where

$$\theta = \frac{pN(m-q)}{(m+1)[N(p-q-1) + p(q+1)]}, \quad 0 \leq q < \sigma - 1, \quad q < m < \sigma,$$

$$\sigma = \begin{cases} \frac{(p-1)N+p}{N-p} & \text{if } p < N \\ \infty & \text{if } p \geq N \end{cases}$$

5 In particular, for $N = 2$, we obtain

$$\|u\|_{L^q} \lesssim \|u\|_{L^2}^{2/q} \|\nabla u\|_{L^2}^{1-2/q}, \quad 2 \leq q < \infty. \quad (2.8)$$

6 **Proposition 2.8** (Sobolev embeddings).

7 *We have*

$$W^{s,p}(\mathbb{R}^N) \hookrightarrow L^q(\mathbb{R}^N), \quad 1 \leq p < \infty, \quad 0 \leq s < \frac{N}{p}, \quad \frac{1}{p} - \frac{s}{N} \leq \frac{1}{q} \leq \frac{1}{p}. \quad (2.9)$$

8

$$W^{1,p}(\mathbb{R}^N) \hookrightarrow \mathcal{C}^{1-\frac{N}{p}}(\mathbb{R}^N), \quad p > N. \quad (2.10)$$

9 The following estimate is an L^∞ logarithmic inequality which enables us to establish
10 the link between $\|e^{4\pi|u|^2} - 1\|_{L_T^1(L^2(\mathbb{R}^2))}$ and dispersion properties of solutions of the linear
11 Schrödinger equation.

12 **Proposition 2.9** (Log estimate [15]).

13 *Let $0 < \beta < 1$. For any $\lambda > \frac{1}{2\pi\beta}$ and any $0 < \mu \leq 1$, a constant $C_\lambda > 0$ exists such that, for
14 any function $u \in H^1(\mathbb{R}^2) \cap \mathcal{C}^\beta(\mathbb{R}^2)$, we have*

$$\|u\|_{L^\infty}^2 \leq \lambda \|u\|_\mu^2 \log \left(C_\lambda + \frac{8^\beta \mu^{-\beta} \|u\|_{\mathcal{C}^\beta}}{\|u\|_\mu} \right), \quad (2.11)$$

15 *where*

$$\|u\|_\mu^2 := \|\nabla u\|_{L^2}^2 + \mu^2 \|u\|_{L^2}^2. \quad (2.12)$$

Recall that $\mathcal{C}^\beta(\mathbb{R}^2)$ denotes the space of β -Hölder continuous functions endowed with the norm

$$\|u\|_{\mathcal{C}^\beta(\mathbb{R}^2)} := \|u\|_{L^\infty(\mathbb{R}^2)} + \sup_{x \neq y} \frac{|u(x) - u(y)|}{|x - y|^\beta}.$$

1 We refer to [15] for the proof of this proposition and more details. We just point out that
 2 the condition $\lambda > \frac{1}{2\pi\beta}$ in (2.11) is optimal.

3 We also recall the so-called Strichartz estimates. We say that (q, r) is an L^2 -admissible
 4 pair if

$$0 \leq \frac{2}{q} = 1 - \frac{2}{r} < 1. \quad (2.13)$$

In particular, note that $(\frac{2}{1-2\sigma}, \frac{1}{\sigma})$ is an admissible pair for any $0 < \sigma < 1/2$ and

$$W^{1, \frac{1}{\sigma}}(\mathbb{R}^2) \hookrightarrow \mathcal{C}^{1-2\sigma}(\mathbb{R}^2).$$

5 **Proposition 2.10** (Strichartz estimates [8]).

6 Let $I \subset \mathbb{R}$ be a time interval and let $t_0 \in I$. Then, for any admissible pairs (q, r) and (\tilde{q}, \tilde{r}) ,
 7 we have

$$\|v\|_{L^q(I, W^{1,r}(\mathbb{R}^2))} \lesssim \|v(t_0)\|_{H^1(\mathbb{R}^2)} + \|i\partial_t v + \Delta v\|_{L^{\tilde{q}}(I, W^{1,\tilde{r}}(\mathbb{R}^2))}. \quad (2.14)$$

8 The following continuity argument (or bootstrap argument) will be useful for our purpose.

Theorem 2.11 (Continuity argument).

Let $X : [0, T] \rightarrow \mathbb{R}$ be a nonnegative continuous function, such that, for every $0 \leq t \leq T$,

$$X(t) \leq a + bX(t)^\theta,$$

where $a, b > 0$ and $\theta > 1$ are constants such that

$$a < \left(1 - \frac{1}{\theta}\right) \frac{1}{(\theta b)^{1/(\theta-1)}} \quad \text{and} \quad X(0) \leq \frac{1}{(\theta b)^{1/(\theta-1)}}.$$

Then, for every $0 \leq t \leq T$, we have

$$X(t) \leq \frac{\theta}{\theta-1} a.$$

9 *Proof.* We sketch the proof for reader's convenience. The function $f : x \mapsto bx^\theta - x + a$ is
 10 decreasing on $[0, (\theta b)^{1/(1-\theta)}]$ and increasing on $[(\theta b)^{1/(1-\theta)}, \infty)$. The assumptions on a and
 11 $X(0)$ imply that $f((\theta b)^{1/(1-\theta)}) < 0$. As $f(X(t)) \geq 0$, $f(0) > 0$ and $X(0) \leq \frac{1}{(\theta b)^{1/(\theta-1)}}$, we
 12 deduce the desired result. \square

13

3. PSEUDO-CONFORMAL LAW

14 In this section, we show a decaying property of global solutions to (1.1). Note that the
 15 conservation laws of mass and Hamiltonian give the boundedness of the L^2 and the H^1 norms
 16 but are insufficient to provide a decay estimate in (more general) Lebesgue spaces. To obtain
 17 such a decay we will take advantage of the pseudo-conformal law.

18 More precisely, we define the following quantities

$$\mathbf{V}(t) := \int |x|^2 |u(t, x)|^2 dx, \quad (3.1)$$

$$\mathbf{M}(t) := 2 \int \mathcal{I}(\bar{u}(t, x)x \cdot \nabla u(t, x)) dx, \quad (3.2)$$

$$\mathbf{K}(t) := \|(x + 2it\nabla)u(t)\|_{L^2}^2 + \frac{4t^2}{\alpha} \int \left(e^{\alpha|u(t,x)|^2} - 1 - \alpha|u(t,x)|^2 - \frac{\alpha^2}{2}|u(t,x)|^4 \right) \frac{dx}{|x|^b}, \quad (3.3)$$

$$\begin{aligned} G(t) &:= \frac{4(2-b)}{\alpha} \int \left(e^{\alpha|u(t,x)|^2} - 1 - \alpha|u(t,x)|^2 - \frac{\alpha}{2}|u(t,x)|^4 \right) \frac{dx}{|x|^b} \\ &\quad - \frac{8}{\alpha} \int \left(e^{\alpha|u(t,x)|^2} (\alpha|u(t,x)|^2 - 1) + 1 - \frac{\alpha^2}{2}|u(t,x)|^4 \right) \frac{dx}{|x|^b} \\ &=: \int g(|u(t,x)|^2) \frac{dx}{|x|^b}, \end{aligned} \quad (3.4)$$

where

$$g(\tau) = \frac{4(2-b)}{\alpha} \left(e^{\alpha\tau} - 1 - \alpha\tau - \frac{\alpha^2}{2}\tau^2 \right) - \frac{8}{\alpha} \left(e^{\alpha\tau} (\alpha\tau - 1) + 1 - \frac{\alpha^2}{2}\tau^2 \right). \quad (3.5)$$

1 **Proposition 3.1.** *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Then*

$$\frac{d\mathbf{V}(t)}{dt} = 2\mathbf{M}(t), \quad (3.6)$$

$$\frac{d^2\mathbf{V}(t)}{dt^2} = 8\mathcal{H}(u_0) - G(t), \quad (3.7)$$

$$\frac{d\mathbf{K}(t)}{dt} = tG(t), \quad (3.8)$$

$$G(t) \leq 0, \quad \forall t \in \mathbb{R}. \quad (3.9)$$

Proof. A straightforward computation gives (3.6). Let $N(x, u) := |x|^{-b} \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 \right) u$. Following [32] for instance, we find that

$$\frac{d^2\mathbf{V}(t)}{dt^2} = 8 \int |\nabla u|^2 dx + 4 \int x \cdot \{N(x, u), u\}_p dx,$$

2 where $\{f, g\}_p = \mathcal{R}(f\nabla\bar{g} - g\nabla\bar{f})$ is the momentum bracket.

Now compute the momentum bracket $\{N(x, u), u\}_p$. Expand $N(x, u)$ in a formal series

$$N(x, u) = |x|^{-b} \sum_{k=2}^{\infty} \frac{\alpha^k}{k!} |u|^{2k} u.$$

3 Using the fact

$$\{|x|^{-b}|u|^\beta u, u\}_p = -\frac{\beta}{\beta+2} \nabla(|x|^{-b}|u|^{\beta+2}) - \frac{2}{\beta+2} \nabla(|x|^{-b})|u|^{\beta+2},$$

one gets

$$\begin{aligned} \{N(x, u), u\}_p &= \sum_{k=2}^{\infty} \frac{\alpha^k}{k!} \{|x|^{-b}|u|^{2k} u, u\}_p \\ &= -\sum_{k=2}^{\infty} k \frac{\alpha^k}{(k+1)!} \nabla(|x|^{-b}|u|^{2k+2}) - \sum_{k=2}^{\infty} \frac{\alpha^k}{(k+1)!} \nabla(|x|^{-b})|u|^{2k+2}. \end{aligned}$$

An integration by parts leads

$$\begin{aligned} \int x \cdot \{N(x, u), u\}_p &= 2 \int \left(\sum_{k=2}^{\infty} k \frac{\alpha^k}{(k+1)!} |u|^{2k+2} \right) \frac{dx}{|x|^b} + b \int \left(\sum_{k=2}^{\infty} \frac{\alpha^k}{(k+1)!} |u|^{2k+2} \right) \frac{dx}{|x|^b} \\ &= \frac{2}{\alpha} \int \left(e^{\alpha|u|^2} (\alpha|u|^2 - 1) + 1 - \frac{\alpha^2}{2} |u|^4 \right) \frac{dx}{|x|^b} \\ &\quad + \frac{b}{\alpha} \int \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 - \frac{\alpha^2}{2} |u|^4 \right) \frac{dx}{|x|^b}, \end{aligned}$$

where we have used

$$\begin{aligned} \sum_{k=2}^{\infty} k \frac{\alpha^k}{(k+1)!} |u|^{2k+2} &= \frac{1}{\alpha} \left(e^{\alpha|u|^2} (\alpha|u|^2 - 1) + 1 - \frac{\alpha^2}{2} |u|^4 \right), \\ \sum_{k=2}^{\infty} \frac{\alpha^k}{(k+1)!} |u|^{2k+2} &= \frac{1}{\alpha} \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 - \frac{\alpha^2}{2} |u|^4 \right). \end{aligned}$$

Therefore,

$$\begin{aligned} \frac{d^2 \mathbf{V}(t)}{dt^2} &= 8 \|\nabla u(t)\|_{L^2}^2 + \frac{8}{\alpha} \int \left(e^{\alpha|u|^2} (\alpha|u|^2 - 1) + 1 - \frac{\alpha^2}{2} |u|^4 \right) \frac{dx}{|x|^b} \\ &\quad + \frac{4b}{\alpha} \int \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 - \frac{\alpha^2}{2} |u|^4 \right) \frac{dx}{|x|^b}. \end{aligned}$$

Using the conservation law (1.4), we conclude the proof of (3.7). To prove (3.8), we first remark that

$$\mathbf{K}(t) = \mathbf{V}(t) - t \frac{d\mathbf{V}(t)}{dt} + 4t^2 \mathcal{H}(u_0).$$

Hence

$$\frac{d\mathbf{K}(t)}{dt} = -t \frac{d^2 \mathbf{V}(t)}{dt^2} + 8t \mathcal{H}(u_0),$$

- 1 and the conclusion follows. Finally, for the sign of G , a simple computation shows that (for
2 all $\tau \geq 0$)

$$g'(\tau) = -8(\alpha x e^{\alpha x} - e^{\alpha x} + 1) - 4b(e^{\alpha x} - \alpha x - 1) \leq 0.$$

- 3 Since $g(0) = 0$, we get (3.9). □

- 4 As a consequence of Proposition 3.1, we have

Corollary 3.2. *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Then*

$$\begin{aligned} \|(x + 2it\nabla)u(t)\|_{L^2}^2 + \frac{4t^2}{\alpha} \int \left(e^{\alpha|u(t,x)|^2} - 1 - \alpha|u(t,x)|^2 - \frac{\alpha^2}{2} |u(t,x)|^4 \right) \frac{dx}{|x|^b} \\ = \|xu_0\|_{L^2}^2 + \int_0^t \tau G(\tau) d\tau. \end{aligned}$$

4. DECAY ESTIMATE

Theorem 4.1. *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Then, for all $t \neq 0$ and $2 \leq q < \infty$,*

$$\|u(t)\|_{L^q} \leq C_q \|u_0\|_{\Sigma} |t|^{-(1-\frac{2}{q})},$$

where $C_q > 0$ is a constant depending only on q .

Proof. Set $v(t, x) := e^{-i\frac{|x|^2}{4t}} u(t, x)$. We see that $\|(x + 2it\nabla)u(t)\|_{L^2}^2 = 4t^2 \|\nabla v(t)\|_{L^2}^2$. Hence, by Corollary 3.2,

$$4t^2 \mathcal{H}(v(t)) = \|xu_0\|_{L^2}^2 + \int_0^t \tau G(\tau) d\tau.$$

Using (3.9), we get

$$4t^2 \|\nabla v(t)\|_{L^2}^2 \leq \|xu_0\|_{L^2}^2,$$

or equivalently

$$\|\nabla v(t)\|_{L^2} \lesssim |t|^{-1}.$$

The conservation of mass, the fact that $|u| = |v|$ and the Gagliardo-Nirenberg inequality (2.8), yield, for all $2 \leq q < \infty$,

$$\|u(t)\|_{L^q} = \|v(t)\|_{L^q} \lesssim |t|^{-(1-\frac{2}{q})}.$$

The proof is complete. \square

A natural and useful consequence from the previous theorem is the following bound estimate.

Corollary 4.2. *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Let $1 \leq p < \infty$, $2 \leq q < \infty$ be such that*

$$p \left(1 - \frac{2}{q}\right) > 1. \tag{4.1}$$

Then, for all $T > 0$, we have

$$\|u\|_{L^p([T, \infty); L^q)} \lesssim \frac{T^{\frac{1}{p} + \frac{2}{q} - 1}}{\left(p \left(1 - \frac{2}{q}\right) - 1\right)^{1/p}} < \infty.$$

For bounded time intervals, the local theory allows us to remove the assumption (4.1) to obtain

Corollary 4.3. *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Let $1 \leq p < \infty$, $2 \leq q < \infty$ and $0 < T < S < \infty$. Then*

$$\|u\|_{L^p([T, S]; L^q)} \leq C,$$

where $C > 0$ depends only on p , q , T , S , $\|u_0\|_{\Sigma}$.

1 Another important consequence that will be used to obtain global bounds asserts that one
 2 can decompose any time interval (T, ∞) with $T > 0$ into a finite number of intervals on which
 3 the $L_t^p(L_x^q)$ norm is sufficiently small for every (p, q) satisfying (4.1). More precisely, we have

4 **Corollary 4.4.** *Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). Let $1 \leq p <$
 5 ∞ , $2 \leq q < \infty$, $\varepsilon > 0$ and $T > 0$. Assume that the condition (4.1) is fulfilled. Then there
 6 exists $L \geq 1$ not depending on u and time intervals I_1, I_2, \dots, I_L such that $\bigcup_{\ell=1}^L I_\ell = [T, \infty)$
 7 and*

$$\|u\|_{L^p(I_\ell; L^q)} \leq \varepsilon, \quad \forall \ell = 1, 2, \dots, L. \quad (4.2)$$

Proof. From Corollary 4.2, one can choose $S > T$ sufficiently large (not depending on u) such
 that $\|u\|_{L^p([S, \infty); L^q)} \leq \varepsilon$. Define

$$T_\ell = T + \ell \frac{S - T}{m}, \quad \ell = 0, 1, \dots, m,$$

8 where $m \geq 1$ to be chosen later. Using Hölder's inequality in time, we obtain that

$$\begin{aligned} \|u\|_{L^p(T_\ell, T_{\ell+1}; L^q)} &\leq \left(\frac{S - T}{m} \right)^{\frac{1}{2p}} \|u\|_{L^{2p}([T, S]; L^q)} \\ &\lesssim \left(\frac{S - T}{m} \right)^{\frac{1}{2p}} \\ &\leq \varepsilon, \end{aligned}$$

9 for $m \geq 1$ sufficiently large and for all $\ell = 0, 1, \dots, m - 1$. This finishes the proof of Corollary
 10 4.4. \square

11

5. GLOBAL BOUNDS 1

In this section, we give the proof of the first global bound in (1.5). For a time slab $I \subset \mathbb{R}$,
 we define $S^1(I)$ via

$$\|u\|_{S^1(I)} = \|u\|_{L^\infty(I, H^1)} + \|u\|_{L^4(I, W^{1,4})}.$$

12 By the Strichartz estimates, we have

$$\|u\|_{S^1(I)} \lesssim \|u(T)\|_{H^1} + \|i\partial_t u + \Delta u\|_{L^{\frac{2}{1+2\delta}}(I, W^{1, \frac{1}{1-\delta}})}, \quad (5.1)$$

13 for any $0 < \delta < 1/2$ and $T \in I$. Note that $\left(\frac{2}{1+2\sigma}, \frac{1}{1-\delta} \right)$ is the conjugate pair of the Schrödinger
 14 admissible pair $\left(\frac{2}{1-2\sigma}, \frac{1}{\sigma} \right)$.

15 **Theorem 5.1.** *Let $u_0 \in \Sigma$ be such that $\mathcal{H}(u_0) < \frac{2}{(1+b)(2-b)}$. Then the corresponding global
 16 solution u to (1.1) satisfies $u \in S^1(\mathbb{R})$.*

1 *Proof.* It suffices to estimate the nonlinear term in some dual Strichartz norm as in (5.1). We
2 have

$$|\nabla N(x, u)| \lesssim |x|^{-b} |\nabla u| |u|^2 \left(e^{\alpha|u|^2} - 1 \right) + |x|^{-b-1} |u| \left(e^{\alpha|u|^2} - 1 - \alpha|u|^2 \right) := \mathcal{A} + \mathcal{B}.$$

3 Let $0 < \delta < \frac{1}{2}$ to be chosen adequately, and let I be a time slab. Let us first estimate the
4 norm $\|\mathcal{A}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})}$. By Hölder's inequality,

$$\|\mathcal{A}\|_{L^{\frac{1}{1-\delta}}} \lesssim \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \left\| \frac{e^{\alpha|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1-2\delta}}}.$$

The term $\left\| \frac{e^{\alpha|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1-2\delta}}}$ can be estimated using Lorentz spaces. Indeed, by (A.1), we get

$$\begin{aligned} \left\| \frac{e^{\alpha|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1-2\delta}}} &\lesssim \|e^{\alpha|u|^2} - 1\|_{L^1}^{1-\theta} \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \| |x|^{-b} \|_{L^{\frac{2}{\delta}, \infty}} \\ &\lesssim \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta, \end{aligned}$$

where $\theta := \delta + \frac{1+b}{2}$. Note that we can choose $0 < \delta < \frac{1-b}{2}$ so that $\theta \in (0, 1)$. Here we have used the Moser-Trudinger inequality (2.1) to obtain that $\|e^{\alpha|u|^2} - 1\|_{L^1} \lesssim 1$ since $\|\nabla u\|_{L^2}^2 < \mathcal{H}(u_0) < \frac{2}{(1+b)(2-b)} < 1$. Hence

$$\begin{aligned} \|\mathcal{A}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} &\lesssim \left\| \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \right\|_{L^{\frac{2}{1+2\delta}}(\mathbf{I})} \\ &\quad + \left\| \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \right\|_{L^{\frac{2}{1+2\delta}}(\mathbf{J})}, \end{aligned}$$

where $\mathbf{I} = \{t \in I / \|u(t)\|_{L^\infty} \leq 1\}$ and $\mathbf{J} = \{t \in I / \|u(t)\|_{L^\infty} \geq 1\}$. The first term in the right hand side can be easily estimated as follows

$$\begin{aligned} \left\| \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \right\|_{L^{\frac{2}{1+2\delta}}(\mathbf{I})} &\lesssim \|\nabla u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{2}{1-\delta}})} \|u\|_{L^{\frac{4}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \\ &\lesssim \|u\|_{S^1(I)} \|u\|_{L^{\frac{4}{1+\delta}}(I, L^{\frac{4}{\delta}})}^2, \end{aligned} \quad (5.2)$$

where the following interpolation inequality is used

$$\|\nabla u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{2}{1-\delta}})} \leq \|\nabla u\|_{L^\infty(I, L^2)}^{1-2\delta} \|\nabla u\|_{L^4(I, L^4)}^{2\delta}. \quad (5.3)$$

Let us turn to the second term. For $t \in \mathbf{J}$, we obtain using (2.11) with $\beta = \frac{1}{2} - \frac{\delta}{2}$ that

$$\|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \lesssim \left(1 + \frac{\|u\|_{C^{\frac{1}{2}-\frac{\delta}{2}}}}{\|u\|_\mu} \right)^{\alpha\theta\lambda\|u\|_\mu^2},$$

5 for some $0 < \mu < 1$ and $\lambda > \frac{1}{\pi(1-\delta)}$ to be chosen later. Since $\|u\|_\mu^2 = \|\nabla u\|_{L^2}^2 + \mu^2\|u\|_{L^2}^2 <$
6 $\mathcal{H}(u_0) + \mu^2\mathcal{M}(u_0) =: K^2(\mu)$, we bound

$$\|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \lesssim \left(1 + \frac{\|u\|_{C^{\frac{1}{2}-\frac{\delta}{2}}}}{K(\mu)} \right)^{\alpha\theta\lambda K^2(\mu)}.$$

- 1 Since $K^2(\mu) \rightarrow H(u_0) < \frac{2}{(1+b)(2-b)}$ as $\mu \rightarrow 0$, we can choose $0 < \mu < 1$ sufficiently small so
 2 that $K^2(\mu) < \frac{2}{(1+b)(2-b)}$. Moreover, as $\frac{\theta K^2(\mu)}{1-\delta} \rightarrow \frac{1+b}{2} K^2(\mu) < \frac{1}{2-b}$ as $\delta \rightarrow 0$, we choose $0 <$
 3 $\delta < \frac{1-b}{2}$ sufficiently small such that $\frac{\theta K^2(\mu)}{1-\delta} < \frac{1}{2-b}$. At final, we choose $\frac{1}{\pi(1-\delta)} < \lambda < \frac{2}{\alpha\theta K^2(\mu)}$
 4 so that $\alpha\theta\lambda K^2(\mu) < 2$. It follows that

$$\|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \lesssim (1 + \|u\|_{C^{\frac{1}{2}-\frac{\delta}{2}}})^2 \lesssim \|u\|_{C^{\frac{1}{2}-\frac{\delta}{2}}}^2,$$

where we have used the fact that $\|u(t)\|_{C^{\frac{1}{2}-\frac{\delta}{2}}} \geq \|u(t)\|_{L^\infty} \geq 1$ for all $t \in \mathbf{J}$. Therefore,

$$\begin{aligned} \left\| \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \|e^{\alpha|u|^2} - 1\|_{L^\infty}^\theta \right\|_{L^{\frac{2}{1+2\delta}}(\mathbf{J})} &\lesssim \left\| \|\nabla u\|_{L^{\frac{2}{1-\delta}}} \|u\|_{L^{\frac{4}{\delta}}}^2 \|u\|_{C^{\frac{1}{2}-\frac{\delta}{2}}}^2 \right\|_{L^{\frac{2}{1+2\delta}}(I)} \\ &\lesssim \|\nabla u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{2}{1-\delta}})} \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|u\|_{L^{\frac{4}{1-\delta}}(I, C^{\frac{1}{2}-\frac{\delta}{2}})}^2 \\ &\lesssim \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|u\|_{S^1(I)}^3. \end{aligned} \quad (5.4)$$

The last estimate follows from (5.3) and the fact

$$\begin{aligned} \|u\|_{L^{\frac{4}{1-\delta}}(I, C^{\frac{1}{2}-\frac{\delta}{2}})} &\lesssim \|u\|_{L^{\frac{4}{1-\delta}}(I, W^{1, \frac{4}{1+\delta}})} \\ &\lesssim \|u\|_{L^\infty(I, H^1)}^\delta \|u\|_{L^4(I, W^{1,4})}^{1-\delta} \\ &\lesssim \|u\|_{S^1(I)}. \end{aligned}$$

- 5 Combining inequalities (5.2) and (5.4), we end up with

$$\|\mathcal{A}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \|u\|_{S^1(I)} \|u\|_{L^{\frac{4}{1+\delta}}(I, L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|u\|_{S^1(I)}^3. \quad (5.5)$$

- 6 Let us now estimate the term $\|\mathcal{B}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})}$. Taking $\frac{1}{1-\delta} < p, q < \infty$ such that $\frac{1}{p} + \frac{1}{q} = 1 - \delta$
 7 and applying Hölder's inequality, we get

$$\|\mathcal{B}\|_{L^{\frac{1}{1-\delta}}} \lesssim \left\| |x|^{-b-1} |u| \left(e^{\alpha|u|^2} - 1 \right) \right\|_{L^{\frac{1}{1-\delta}}} \lesssim \|u\|_{L^q} \left\| \frac{e^{\alpha|u|^2} - 1}{|x|^{b+1}} \right\|_{L^p}.$$

- 8 Clearly,

$$\left\| \frac{e^{\alpha|u|^2} - 1}{|x|^{b+1}} \right\|_{L^p}^p \lesssim e^{\alpha(p-1)\|u\|_{L^\infty}^2} \int \frac{e^{\alpha|u|^2} - 1}{|x|^{p(b+1)}} dx.$$

Since $\frac{1}{1-\delta} < \frac{2}{b+1}$ for $0 < \delta < \frac{1-b}{2}$, we choose $\frac{1}{1-\delta} < p < \frac{2}{b+1}$. Hence we can apply the singular Moser-Trudinger inequality for the term $\int \frac{e^{\alpha|u|^2} - 1}{|x|^{p(b+1)}} dx$ to obtain

$$\|\mathcal{B}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \left\| \|u\|_{L^q} e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \right\|_{L^{\frac{2}{1+2\delta}}(I)} + \left\| \|u\|_{L^q} e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \right\|_{L^{\frac{2}{1+2\delta}}(\mathbf{J})}.$$

- 9 Note that the choice of p leads to $q > \frac{2}{1-b-2\delta}$. Therefore,

$$\|\mathcal{B}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \|u\|_{L^{\frac{2}{1+2\delta}}(I, L^q)} + \|u\|_{L^\gamma(I, L^q)} \left\| e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \right\|_{L^p(\mathbf{J})},$$

- 1 where $\frac{2}{1+2\delta} < \gamma, \rho < \infty$ such that $\frac{1}{\gamma} + \frac{1}{\rho} = \frac{1+2\delta}{2}$. Let $t \in \mathbf{J}$. An application of the Log
 2 estimate (2.11) with $\beta = \frac{1}{2}$ gives

$$e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \lesssim \left(1 + \frac{\|u\|_{C^{\frac{1}{2}}}}{\|u\|_\mu}\right)^{\alpha \frac{p-1}{p} \lambda \|u\|_\mu^2},$$

- 3 for some $0 < \mu < 1$ and $\lambda > \frac{1}{\pi}$ to be chosen shortly. Since $\|u\|_\mu^2 < K^2(\mu)$, it follows that

$$e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \lesssim \left(1 + \frac{\|u\|_{C^{\frac{1}{2}}}}{K(\mu)}\right)^{\alpha \frac{p-1}{p} \lambda K^2(\mu)}.$$

- 4 Choose $0 < \mu < 1$ sufficiently small such that $K^2(\mu) < \frac{2}{(1+b)(2-b)} < 1$. Since $\frac{1}{1-\delta} < \frac{2(2-b)}{2(2-b)-1}$
 5 for $0 < \delta < \frac{1}{2(2-b)}$, we choose $\frac{1}{1-\delta} < p < \frac{2(2-b)}{2(2-b)-1}$. In particular $\frac{p}{2(2-b)(p-1)} > 1 > K^2(\mu)$. At
 6 final, we choose $\frac{1}{\pi} < \lambda < \frac{p}{\alpha(p-1)K^2(\mu)}$ so that $\alpha \frac{p-1}{p} \lambda K^2(\mu) < 1$. We thus get

$$e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2} \lesssim 1 + \|u\|_{C^{\frac{1}{2}}} \lesssim \|u\|_{W^{1,4}},$$

where we have used $\|u(t)\|_{W^{1,4}} \gtrsim \|u(t)\|_{C^{\frac{1}{2}}} \geq \|u(t)\|_{L^\infty} \geq 1$ for all $t \in \mathbf{J}$. Therefore, choosing
 $\frac{2}{1+2\delta} < \rho < 4$, one gets

$$\|e^{\alpha \frac{p-1}{p} \|u\|_{L^\infty}^2}\|_{L^\rho(\mathbf{J})} \lesssim \|u\|_{L^4(I, W^{1,4})}^{\frac{4}{\rho}}.$$

- 7 This finally leads to

$$\|\mathcal{B}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \|u\|_{L^{\frac{2}{1+2\delta}}(I, L^q)} + \|u\|_{L^\gamma(I, L^q)} \|u\|_{S^1(I)}^{\frac{4}{\rho}}.$$

Note also that this choice of p leads to $0 < \delta < \frac{1}{2(2-b)}$ and $q > \frac{2(2-b)}{1-2\delta(2-b)}$. Arguing similarly
 for $\|N(x, u)\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})}$, we conclude that

$$\begin{aligned} \|u\|_{S^1(I)} &\lesssim \|u(T)\|_{H^1} + \|u\|_{S^1(I)} \|u\|_{L^{\frac{4}{1+\delta}}(I, L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|u\|_{S^1(I)}^3 \\ &\quad + \|u\|_{L^{\frac{2}{1+2\delta}}(I, L^q)} + \|u\|_{L^\gamma(I, L^q)} \|u\|_{S^1(I)}^{\frac{4}{\rho}}, \end{aligned} \quad (5.6)$$

- 8 where $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$, $\gamma > \frac{2}{1+2\delta}$ and $q > \max\left\{\frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}\right\}$.

Since $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$ and $\gamma > \frac{2}{1+2\delta}$, it is easy to check that the condition (4.1)
 is satisfied for

$$(p, q) \in \left\{ \left(\frac{4}{1+\delta}, \frac{4}{\delta}\right), \left(\frac{2}{\delta}, \frac{4}{\delta}\right), \left(\frac{2}{1+2\delta}, q\right), (\gamma, q) \right\} \quad (5.7)$$

- 9 provided that q satisfies an additional condition $q > \frac{4}{1-2\delta}$. We thus obtain by Corollary 4.2
 10 that $\|u\|_{L^p((a, \infty), L^q)} < \infty$ for all $a > 0$ and (p, q) in (5.7).

Now let $0 < \delta < \min \left\{ \frac{1-b}{2}, \frac{1}{2(2-b)} \right\}$, $\gamma > \frac{2}{1+2\delta}$ and $q > \max \left\{ \frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}, \frac{4}{1-2\delta} \right\}$. We choose $T > 0$ so that $\|u\|_{L^{\frac{4}{1+\delta}}((T,\infty),L^{\frac{4}{\delta}})}^2 \leq \frac{1}{2}$. Set $A := \|u\|_{L^{\frac{2}{1+2\delta}}((T,\infty),L^q)} < \infty$. We have

$$\begin{aligned} \|u\|_{S^1((T,\infty))} &\lesssim \|u(T)\|_{H^1} + \|u\|_{S^1((T,\infty))} \|u\|_{L^{\frac{4}{1+\delta}}((T,\infty),L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}((T,\infty),L^{\frac{4}{\delta}})}^2 \|u\|_{S^1((T,\infty))}^3 \\ &\quad + \|u\|_{L^{\frac{2}{1+2\delta}}((T,\infty),L^q)} + \|u\|_{L^\gamma((T,\infty),L^q)} \|u\|_{S^1((T,\infty))}^{\frac{4}{\rho}}. \end{aligned}$$

Bounding $\|u(T)\|_{H^1}$ by a constant $C(\mathcal{H}(u_0) + \mathcal{M}(u_0))$ depending only on the mass and energy of the initial data, one infers

$$\begin{aligned} \|u\|_{S^1((T,\infty))} &\lesssim C(\mathcal{H}(u_0) + \mathcal{M}(u_0)) + A + \|u\|_{L^{\frac{2}{\delta}}((T,\infty),L^{\frac{4}{\delta}})}^2 \|u\|_{S^1((T,\infty))}^3 \\ &\quad + \|u\|_{L^\gamma((T,\infty),L^q)} \|u\|_{S^1((T,\infty))}^{\frac{4}{\rho}}. \end{aligned}$$

Using Corollary 4.4, one can pick $\varepsilon > 0$ small (to be determined later) and a finite number of intervals $\{I_\ell\}_{\ell=1,2,\dots,L}$, $I_\ell \subset (T, \infty)$ such that , for all ℓ

$$\|u\|_{L^\gamma(I_\ell, L^q)}, \quad \|u\|_{L^{\frac{2}{\delta}}(I_\ell, L^{\frac{4}{\delta}})} \leq \varepsilon.$$

Thus, by (5.6) and since $\frac{4}{\rho} > 1$, we get

$$\|u\|_{S^1(I_\ell)} \lesssim C(\mathcal{H}(u_0) + \mathcal{M}(u_0)) + A + \varepsilon^2 \|u\|_{S^1(I_\ell)}^3 + \varepsilon \|u\|_{S^1(I_\ell)}^{\frac{4}{\rho}}.$$

- 1 A continuity argument allows us to pick $\varepsilon > 0$ sufficiently small depending only on $C(\mathcal{H}(u_0) +$
- 2 $\mathcal{M}(u_0)) + A$ such that $\|u\|_{S^1(I_\ell)} \leq C(\mathcal{H}(u_0), \mathcal{M}(u_0), A)$. Since the number of intervals is finite
- 3 and the conclusion can be made for all I_ℓ 's, we get $\|u\|_{S^1((T,\infty))} < \infty$. A similar argument
- 4 applies for negative times, and we get $\|u\|_{S^1((-\infty, -S))} < \infty$ for some $S > 0$. We conclude the
- 5 proof by the local theory. \square

6

6. GLOBAL BOUNDS 2

- 7 In this section, we prove the second global bound in (1.5). For a time slab $I \subset \mathbb{R}$, we define
- 8 $S^0(I)$ by

$$\|u\|_{S^0(I)} = \|u\|_{L^\infty(I, L^2)} + \|u\|_{L^4(I, L^4)}.$$

- 9 **Theorem 6.1.** *Let $u_0 \in \Sigma$ be such that $\mathcal{H}(u_0) < \frac{2}{(1+b)(2-b)}$. Let u the corresponding global*
- 10 *solution to (1.1) and set $w(t) := (x + 2it\nabla)u(t)$. Then it holds that $w \in S^0(\mathbb{R})$.*

- 11 *Proof.* Let $T > 0$ and set $I = (T, \infty)$. Since $x + 2it\nabla$ commutes with $i\partial_t + \Delta$, the Duhamel
- 12 formula implies

$$w(t) = e^{i(t-T)\Delta} w(T) - i \int_T^t e^{i(t-s)\Delta} (x + 2is\nabla) N(x, u) ds.$$

Let $v(t, x) := e^{-i\frac{|x|^2}{4t}} u(t, x)$. It is easy to see that $|v| = |u|$, $|(x+2it\nabla)N(x, u)| = 2|t||\nabla N(x, v)|$ and $2|t||\nabla v| = |w|$. We thus bound

$$|(x+2is\nabla)N(x, u)| \lesssim |x|^{-b}(2|s||\nabla v|)|u|^2 \left(e^{\alpha|u|^2} - 1 \right) + |x|^{-b-1}(2|s||u|^3) \left(e^{\alpha|u|^2} - 1 \right) := \mathcal{C} + \mathcal{D},$$

where we have used the fact that for all $x \geq 0$, $e^x - 1 - x \leq x(e^x - 1)$. It follows from the pseudo-conformal law that $\|w(T)\|_{L^2} \leq \|xu_0\|_{L^2}$. Thus, Strichartz estimate yields

$$\begin{aligned} \|w\|_{S^0(I)} &\lesssim \|xu_0\|_{L^2} + \|(x+2is\nabla)N(x, u)\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \\ &\lesssim \|xu_0\|_{L^2} + \|2|s||\nabla N(x, v)\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})}. \end{aligned}$$

1 As above, one gets

$$\|\mathcal{C}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \|w\|_{S^0(I)} \|u\|_{L^{\frac{4}{1+2\delta}}(I, L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|w\|_{S^0(I)} \|u\|_{S^1(I)}^2$$

2 and

$$\|\mathcal{D}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \| |s||u|^3 \|_{L^{\frac{2}{1+2\delta}}(I, L^q)} + \| |s||u|^3 \|_{L^\gamma(I, L^q)} \|u\|_{S^1(I)}^{\frac{4}{\rho}},$$

3 where $0 < \delta < \min \left\{ \frac{1-b}{2}, \frac{1}{2(2-b)} \right\}$, $\gamma > \frac{2}{1+2\delta}$ and $q > \max \left\{ \frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)} \right\}$. The last
4 inequality can be written as

$$\|\mathcal{D}\|_{L^{\frac{2}{1+2\delta}}(I, L^{\frac{1}{1-\delta}})} \lesssim \| |s|^{\frac{1}{3}} |u|^3 \|_{L^{\frac{6}{1+2\delta}}(I, L^{3q})} + \| |s|^{\frac{1}{3}} |u|^3 \|_{L^{3\gamma}(I, L^{3q})} \|u\|_{S^1(I)}^{\frac{4}{\rho}}.$$

As in Corollary 4.2, we note that for all $a > 0$, the norm $\| |s|^{\frac{1}{3}} u \|_{L^m((a, \infty), L^n)} < \infty$ provided that $1 \leq m < \infty$, $2 \leq n < \infty$ satisfying

$$m \left(\frac{2}{3} - \frac{2}{n} \right) > 1. \quad (6.1)$$

Since $\gamma > \frac{2}{1+2\delta}$, the condition (6.1) is fulfilled for $(m, n) = \left\{ \left(\frac{6}{1+2\delta}, 3q \right), (3\gamma, 3q) \right\}$ provided that $q > \frac{4}{3-2\delta}$. Under the conditions

$$0 < \delta < \min \left\{ \frac{1-b}{2}, \frac{1}{2(2-b)} \right\}, \quad \gamma > \frac{2}{1+2\delta}, \quad q > \max \left\{ \frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}, \frac{4}{3-2\delta} \right\},$$

5 we argue as above to obtain $\|w\|_{S^0((T, \infty))} < \infty$ for some $T > 0$. By the same argument,
6 we prove as well that $\|w\|_{S^0((-\infty, -S))} < \infty$ for some $S > 0$. It remains to show that $w \in$
7 $S^0([-S, T])$. The proof of the latter claim follows the same argument as in [32]. To see this,
8 set $H(t) = x + 2it\nabla$. We are going to prove that $\|Hu\|_{S^0([-S, T])} < \infty$. Divide $[-S, T]$ into a
9 finite number of intervals $J_k = [t_k, t_{k+1}]$ such that $|J_k| \leq \varepsilon$, where $\varepsilon > 0$ is to be chosen later.
10 The Duhamel formula reads

$$H(t)u(t) = e^{i(t-t_k)\Delta} H(t_k)u(t_k) - i \int_{t_k}^t e^{i(t-s)\Delta} H(s)N(x, u) ds.$$

By Strichartz estimates,

$$\begin{aligned} \|Hu\|_{S^0(J_k)} &\lesssim \|H(t_k)u(t_k)\|_{L^2} + \|H(s)N(x, u)\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})} \\ &\lesssim \|H(t_k)u(t_k)\|_{L^2} + \|2|s|\|\nabla N(x, v)\|\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})}. \end{aligned}$$

- 1 Note that in the following all constants involved in \lesssim are independent of k . Using the fact
2 that, for all $x \geq 0$ and all $\eta > 0$, $x(e^x - 1) \leq \frac{e^{(1+\eta)x} - 1}{\eta}$, and that $|v| = |u|$, we bound

$$2|s|\|\nabla N(x, v)\| \lesssim_\eta |x|^{-b}(2|s|\|\nabla v\|) \left(e^{\alpha(1+\eta)|u|^2} - 1 \right) + |x|^{-b-1}(2|s|\|u\|) \left(e^{\alpha(1+\eta)|u|^2} - 1 \right).$$

- 3 The first term in the right hand side is estimated as follows. By Hölder's inequality,

$$\left\| 2|s|\|\nabla v\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{1}{1-\delta}}} \lesssim \|2|s|\|\nabla v\|\|_{L^{\frac{2}{1-\delta}}} \left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1-\delta}}}.$$

- 4 Hence

$$\left\| 2|s|\|\nabla v\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})} \lesssim \|2|s|\|\nabla v\|\|_{L^{\frac{2}{\delta}}(J_k, L^{\frac{2}{1-\delta}})} \left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1+\delta}}(J_k, L^{\frac{2}{1-\delta}})}.$$

- 5 By (5.3), $\|2|s|\|\nabla v\|\|_{L^{\frac{2}{\delta}}(J_k, L^{\frac{2}{1-\delta}})} = \|Hu\|_{L^{\frac{2}{\delta}}(J_k, L^{\frac{2}{1-\delta}})} \lesssim \|Hu\|_{S^0(J_k)}$. Write

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1-\delta}}}^{\frac{2}{1-\delta}} \lesssim \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{1+\delta}{1-\delta}} \int \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{\frac{2b}{1-\delta}}} dx.$$

- 6 Since $\frac{2b}{1-\delta} \rightarrow 2b < 2$ as $\delta \rightarrow 0$ and $\frac{1}{\mathcal{H}(u_0)} > 1$, one can choose $0 < \delta < \frac{1}{2}$ and $\eta > 0$ sufficiently
7 small such that $\frac{2b}{1-\delta} < 2$ and $0 < \eta < \frac{1}{\mathcal{H}(u_0)} - 1$. This guarantees that $\|\nabla(\sqrt{1+\eta}u)\|_{L^2} \leq$
8 $\sqrt{1+\eta}\sqrt{\mathcal{H}(u_0)} < 1$. Hence we can apply the singular Moser-Trudinger inequality for the
9 term $\int \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{\frac{2b}{1-\delta}}} dx$. Thus

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1+\delta}}(J_k, L^{\frac{2}{1-\delta}})} \lesssim \|e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1\|_{L^1(\mathbf{I}_k)}^{\frac{1+\delta}{2}} + \|e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1\|_{L^1(\mathbf{J}_k)}^{\frac{1+\delta}{2}},$$

- 10 where $\mathbf{I}_k := \{t \in J_k / \|u(t)\|_{L^\infty} \leq 1\}$ and $\mathbf{J}_k := \{t \in J_k / \|u(t)\|_{L^\infty} \geq 1\}$. Let $t \in \mathbf{I}_k$. We have

$$e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \lesssim_{\alpha, \eta} 1.$$

- 11 Thus

$$\|e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1\|_{L^1(\mathbf{I}_k)} \lesssim |J_k|.$$

- 12 Let $t \in \mathbf{J}_k$. An application of the Log estimate (2.11) gives

$$\|e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1\|_{L^\infty} \lesssim \left(1 + \frac{\|u\|_{C^{\frac{1}{2}}}}{K(\mu)} \right)^{\alpha(1+\eta)\lambda K^2(\mu)},$$

- 13 where $K^2(\mu) = \mathcal{H}(u_0) + \mu^2 \mathcal{M}(u_0)$, $0 < \mu < 1$ and $\lambda > \frac{1}{\pi}$. Choose $\frac{4}{b+1} < \sigma < 4$. We
14 next choose $0 < \mu < 1$ sufficiently small such that $K^2(\mu) < \frac{\sigma}{2(2-b)}$. This is possible since
15 $K^2(\mu) \rightarrow \mathcal{H}(u_0) < \frac{2}{(1+b)(2-b)} < \frac{\sigma}{2(2-b)}$. Choose $\eta > 0$ sufficiently small such that $1 + \eta <$

1 $\frac{\sigma}{2(2-b)K^2(\mu)}$. Thus $1 < \frac{\sigma}{2(2-b)(1+\eta)K^2(\mu)}$. One can thus choose $\frac{1}{\pi} < \lambda < \frac{\sigma}{\alpha(1+\eta)K^2(\mu)}$ so that
 2 $\alpha(1+\eta)\lambda K^2(\mu) < \sigma$. As above, one comes to

$$\left\| e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right\|_{L^1(\mathbf{J}_k)} \lesssim |J_k|^{1-\frac{\sigma}{4}} \|u\|_{L^4(J_k, W^{1,4})}^\sigma.$$

3 Therefore,

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1+\delta}}(J_k, L^{\frac{2}{1-\delta}})} \lesssim |J_k|^{\frac{1+\delta}{2}} + |J_k|^{(1-\frac{\sigma}{4})\frac{1+\delta}{2}} \|u\|_{S^1(J_k)}^{\frac{(1+\delta)\sigma}{2}}.$$

4 Conclusion

$$\left\| 2|s|\|\nabla v\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^b} \right\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})} \lesssim \|Hu\|_{S^0(J_k)} \left(|J_k|^{\frac{1+\delta}{2}} + |J_k|^{(1-\frac{\sigma}{4})\frac{1+\delta}{2}} \|u\|_{S^1(J_k)}^{\frac{(1+\delta)\sigma}{2}} \right).$$

5 For the second term, we estimate

$$\left\| |s||u| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})} \lesssim \| |s||u| \|_{L^\gamma(J_k, L^q)} \left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^\rho(J_k, L^p)},$$

6 where $\frac{1}{p} + \frac{1}{q} = 1 - \delta$ and $\frac{1}{\gamma} + \frac{1}{\rho} = \frac{1+2\delta}{2}$. Write

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^p}^p \lesssim \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{p-1} \int \frac{e^{\alpha(1+\delta)|u|^2} - 1}{|x|^{p(b+1)}} dx.$$

7 Since $\frac{1}{1-\delta} \rightarrow 1 < \frac{2}{b+1}$ and $\frac{1}{\mathcal{H}(u_0)} > 1$, one can choose $0 < \delta < \frac{1}{2}$ and $\eta > 0$ sufficiently
 8 small such that $\frac{1}{1-\delta} < \frac{2}{b+1}$ and $0 < \eta < \frac{1}{\mathcal{H}(u_0)} - 1$. This guarantees that $\|\nabla(\sqrt{1+\eta}u)\|_{L^2} \leq$
 9 $\sqrt{1+\eta}\sqrt{\mathcal{H}(u_0)} < 1$. Choose $\frac{1}{1-\delta} < p < \frac{2}{b+1}$. Hence we can apply the singular Moser-
 10 Trudinger inequality for the term $\int \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{p(b+1)}} dx$. Thus

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^\rho(J_k, L^p)} \lesssim \left\| \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{p-1}{p}} \right\|_{L^\rho(\mathbf{I}_k)} + \left\| \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{p-1}{p}} \right\|_{L^\rho(\mathbf{J}_k)}.$$

11 Let $t \in \mathbf{I}_k$. Since

$$\left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{p-1}{p}\rho} \lesssim_{\alpha, \eta, p, \rho} 1,$$

12 we get

$$\left\| \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{p-1}{p}} \right\|_{L^\rho(\mathbf{I}_k)}^\rho \lesssim |J_k|.$$

13 Let $t \in \mathbf{J}_k$. An application of the Log estimate (2.11) with $\beta = \frac{1}{2}$ gives

$$\|e^{\alpha(1+\eta)|u|^2} - 1\|_{L^\infty}^{\frac{p-1}{p}\rho} \lesssim \left(1 + \frac{\|u\|_{\mathcal{C}^{\frac{1}{2}}}}{K(\mu)} \right)^{\alpha(1+\eta)\frac{p-1}{p}\rho\lambda K^2(\mu)},$$

14 for some $0 < \mu < 1$ and $\lambda > \frac{1}{\pi}$. Choose $0 < \mu < 1$ sufficiently small such that $K^2(\mu) < 1$.
 15 Since $\frac{1}{1-\delta} \rightarrow 1 < \frac{2(2-b)}{2(2-b)-1}$ as $\delta \rightarrow 0$, one can choose $0 < \delta < \frac{1}{2}$ sufficiently small such that
 16 $\frac{1}{1-\delta} < \frac{2(2-b)}{2(2-b)-1}$. Choose $\frac{1}{1-\delta} < p < \frac{2(2-b)}{2(2-b)-1}$. In particular $\frac{p}{2(2-b)(p-1)} > 1 > K^2(\mu)$. Choose

- 1 $0 < \eta < 1$ such that $1 + \eta < \frac{p}{2(2-b)(p-1)K^2(\mu)}$. At final, we choose $\frac{1}{\pi} < \lambda < \frac{p}{\alpha(1+\eta)(p-1)K^2(\mu)}$.
 2 Therefore, choosing $\frac{2}{1+2\delta} < \rho < 4$, one gets

$$\|e^{\alpha(1+\eta)|u|^2} - 1\|_{L^\infty}^{\frac{p-1}{p}\rho} \lesssim \left(1 + \|u\|_{C^{\frac{1}{2}}}\right)^4.$$

- 3 Using the fact that $1 \leq \|u(t)\|_{L^\infty} \leq \|u(t)\|_{C^{\frac{1}{2}}} \leq \|u(t)\|_{W^{1,4}}$ for all $t \in \mathbf{J}_k$, one gets

$$\|e^{\alpha(1+\eta)|u|^2} - 1\|_{L^\infty}^{\frac{p-1}{p}\rho} \lesssim \|u\|_{W^{1,4}}^4.$$

- 4 We come to

$$\left\| \left(e^{\alpha(1+\eta)\|u\|_{L^\infty}^2} - 1 \right)^{\frac{p-1}{p}} \right\|_{L^\rho(\mathbf{J}_k)}^\rho \lesssim \|u\|_{L^4(J_k, W^{1,4})}^4.$$

- 5 Thus

$$\left\| \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^\rho(J_k, L^p)} \lesssim \left(|J_k| + \|u\|_{S^1(J_k)}^4 \right)^{\frac{1}{\rho}}.$$

- 6 By the sobolev embedding, one has

$$\| |s|u \|_{L^\gamma(J_k, L^q)} \lesssim \|u\|_{L^\infty(J_k, H^1)} \|s\|_{L^\gamma(J_k)} \lesssim |J_k|^{\frac{\gamma+1}{\gamma}},$$

- 7 where we have used the conservation laws and the fact that $\|s\|_{L^\gamma(J_k)} = \left(\frac{t_k^{\gamma+1} - t_{k-1}^{\gamma+1}}{\gamma+1} \right)^{\frac{1}{\gamma}} \lesssim$

- 8 $(t_{k+1} - t_k)^{\frac{\gamma+1}{\gamma}} = |J_k|^{\frac{\gamma+1}{\gamma}}$. Therefore,

$$\left\| |s|u \frac{e^{\alpha(1+\eta)|u|^2} - 1}{|x|^{b+1}} \right\|_{L^{\frac{2}{1+2\delta}}(J_k, L^{\frac{1}{1-\delta}})} \lesssim |J_k|^{\frac{\gamma+1}{\gamma}} \left(|J_k| + \|u\|_{S^1(J_k)}^4 \right)^{\frac{1}{\rho}}.$$

Collecting the above estimates, we obtain

$$\begin{aligned} \|Hu\|_{S^0(J_k)} &\lesssim \|H(t_k)u(t_k)\|_{L^2} + \|Hu\|_{S^0(J_k)} \left(|J_k|^{\frac{1+\delta}{2}} + |J_k|^{(1-\frac{\sigma}{4})\frac{1+\delta}{2}} \|u\|_{S^1(J_k)}^{\frac{(1+\delta)\sigma}{2}} \right) \\ &\quad + |J_k|^{\frac{\gamma+1}{\gamma}} \left(|J_k| + \|u\|_{S^1(J_k)}^4 \right)^{\frac{1}{\rho}} \\ &\lesssim \|H(t_k)u(t_k)\|_{L^2} + \|Hu\|_{S^0(J_k)} \left(\varepsilon^{\frac{1+\delta}{2}} + \varepsilon^{(1-\frac{\sigma}{4})\frac{1+\delta}{2}} \|u\|_{S^1(J_k)}^{\frac{(1+\delta)\sigma}{2}} \right) \\ &\quad + \varepsilon^{\frac{\gamma+1}{\gamma}} \left(\varepsilon + \|u\|_{S^1(J_k)}^4 \right)^{\frac{1}{\rho}}. \end{aligned}$$

- 9 Since $\|u\|_{S^1(\mathbb{R})} < \infty$, we can choose $\varepsilon > 0$ small enough depending on S, T and $\|u\|_{S^1(\mathbb{R})}$ to
 10 get

$$\|Hu\|_{S^0(J_k)} \leq C \|H(t_k)u(t_k)\|_{L^2} + C,$$

- 11 for some constant $C > 0$ independent of S and T . By induction, we obtain for each k ,

$$\|Hu\|_{S^0(J_k)} \leq C \|H(-S)u(-S)\|_{L^2} + C.$$

- 12 Summing over all subintervals J_k , we prove $\|Hu\|_{S^0([-S, T])} < \infty$. The proof is complete. \square

1

7. SCATTERING IN WEIGHTED L^2 SPACE

2

In this section, we give the proof of our main result in Theorem 1.2.

3

Proof of Theorem 1.2. Let $u_0 \in \Sigma$ and u the corresponding global solution to (1.1). By

4

Duhamel formula, we have

$$e^{-it\Delta}u(t) = u_0 - i \int_0^t e^{-is\Delta}N(x, u)ds.$$

Let $0 < t_1 < t_2 < +\infty$. It follows from Strichartz estimates that

$$\begin{aligned} \|e^{-it_2\Delta}u(t_2) - e^{-it_1\Delta}u(t_1)\|_{H^1} &= \left\| \int_{t_1}^{t_2} e^{-is\Delta}N(x, u)ds \right\|_{H^1} \\ &\lesssim \|N(x, u)\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^{\frac{1}{1-\delta}})} + \|\nabla N(x, u)\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^{\frac{1}{1-\delta}})}. \end{aligned}$$

Arguing as in the proof of (5.6), we obtain

$$\begin{aligned} \|e^{-it_2\Delta}u(t_2) - e^{-it_1\Delta}u(t_1)\|_{H^1} &\lesssim \|u\|_{S^1((t_1, t_2))} \|u\|_{L^{\frac{4}{1+\delta}}((t_1, t_2), L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}((t_1, t_2), L^{\frac{4}{\delta}})}^2 \|u\|_{S^1((t_1, t_2))}^3 \\ &\quad + \|u\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^q)} + \|u\|_{L^\gamma((t_1, t_2), L^q)} \|u\|_{S^1((t_1, t_2))}^{\frac{4}{\delta}}, \end{aligned} \quad (7.1)$$

5

where $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$, $\gamma > \frac{2}{1+2\delta}$ and $q > \max\left\{\frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}\right\}$. By adding

6

an additional condition $q > \frac{4}{1-2\delta}$, we learn from Corollary 4.2 that $\|u\|_{L^p((a, \infty), L^q)} < \infty$

7

for all $a > 0$ and (p, q) in (5.7). In particular, $\|u\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^q)} \rightarrow 0$ as $t_1 \rightarrow +\infty$. Since

8

$\|u\|_{S^1(\mathbb{R})} < \infty$, we infer that the right hand side of (7.1) tends to zero as $t_1, t_2 \rightarrow +\infty$ provided

9

that $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$, $\gamma > \frac{2}{1+2\delta}$ and $q > \max\left\{\frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}, \frac{4}{1-2\delta}\right\}$. This shows

10

that $e^{-it\Delta}u(t)$ is a Cauchy sequence in H^1 as $t \rightarrow +\infty$. There thus exists $u_0^+ \in H^1$ such that

11

$e^{-it\Delta}u(t) \rightarrow u_0^+$ as $t \rightarrow +\infty$. It remains to show that this scattering state u_0^+ belongs to Σ .

12

Since $x + 2it\nabla$ commutes with $i\partial_t + \Delta u$, the Duhamel formula gives

$$(x + 2it\nabla)u(t) = e^{it\Delta}xu_0 - i \int_0^t e^{i(t-s)\Delta}(x + 2is\nabla)N(x, u)ds.$$

13

Using the fact that $x + 2it\nabla = e^{it\Delta}xe^{-it\Delta}$, we write

$$xe^{-it\Delta}u(t) = xu_0 - i \int_0^t e^{-is\Delta}(x + 2is\nabla)N(x, u)ds.$$

By Strichartz estimates, we have

$$\begin{aligned} \|xe^{-it_2\Delta}u(t_2) - xe^{-it_1\Delta}u(t_1)\|_{L^2} &= \left\| \int_{t_1}^{t_2} e^{-is\Delta}(x + 2is\nabla)N(x, u)ds \right\|_{L^2} \\ &\lesssim \|(x + 2is\nabla)N(x, u)\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^{\frac{1}{1-\delta}})} \\ &\lesssim \|2|s|\nabla N(x, u)\|_{L^{\frac{2}{1+2\delta}}((t_1, t_2), L^{\frac{1}{1-\delta}})}, \end{aligned}$$

where $v(t, x) := e^{-i\frac{|x|^2}{4t}} u(t, x)$. Estimating as in the proof of Theorem 6.1, we get

$$\begin{aligned} \|xe^{-it_2\Delta}u(t_2) - xe^{-it_1\Delta}u(t_1)\|_{L^2} &\lesssim \|w\|_{S^0(I)}\|u\|_{L^{\frac{4}{1+\delta}}(I, L^{\frac{4}{\delta}})}^2 + \|u\|_{L^{\frac{2}{\delta}}(I, L^{\frac{4}{\delta}})}^2 \|w\|_{S^0(I)}\|u\|_{S^1(I)}^2 \\ &\quad + \| |s|^{\frac{1}{3}}|u| \|_{L^{\frac{6}{1+2\delta}}(I, L^{3q})}^3 + \| |s|^{\frac{1}{3}}|u| \|_{L^{3\gamma}(I, L^{3q})}^3 \|u\|_{S^1(I)}^{\frac{4}{\rho}}, \end{aligned} \quad (7.2)$$

1 where $w(t) = (x + 2it\nabla)u(t)$, $I = (t_1, t_2)$, $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$, $\gamma > \frac{2}{1+2\delta}$ and
 2 $q > \max\left\{\frac{1}{1-b-2\delta}, \frac{2(2-b)}{1-2\delta(2-b)}\right\}$. Since $0 < \delta < \min\left\{\frac{1-b}{2}, \frac{1}{2(2-b)}\right\}$, Corollary 4.2 implies
 3 that the norms $\|u\|_{L^{\frac{4}{1+\delta}}((a, \infty), L^{\frac{4}{\delta}})}$ and $\|u\|_{L^{\frac{2}{\delta}}((a, \infty), L^{\frac{4}{\delta}})}$ are finite for any $a > 0$. More-
 4 over, adding an additional condition $q > \frac{4}{3-2\delta}$, the condition (6.1) is satisfied for $(m, n) =$
 5 $\left\{\left(\frac{6}{1+2\delta}, 3q\right), (3\gamma, 3q)\right\}$. Thus the norms $\| |s|^{\frac{1}{3}}|u| \|_{L^{\frac{6}{1+2\delta}}((a, \infty), L^{3q})}$ and $\| |s|^{\frac{1}{3}}|u| \|_{L^{3\gamma}((a, \infty), L^{3q})}$
 6 are both finite for any $a > 0$. In particular, $\| |s|^{\frac{1}{3}}|u| \|_{L^{\frac{6}{1+2\delta}}((t_1, t_2), L^{3q})} \rightarrow 0$ as $t_1 \rightarrow +\infty$. Since
 7 $\|u\|_{S^1(\mathbb{R})} < \infty$ and $\|w\|_{S^0(\mathbb{R})} < \infty$, the right hand side of (7.2) tends to zero as $t_1, t_2 \rightarrow +\infty$.
 8 This implies that $xe^{-it\Delta}u(t)$ is a Cauchy sequence in L^2 as $t \rightarrow +\infty$. We thus have $xu_0^+ \in L^2$
 9 and so $u_0^+ \in \Sigma$. Moreover, we have

$$u_0^+(t) = u_0 - i \int_t^{+\infty} e^{-is\Delta} N(x, u) ds.$$

10 Repeating the above argument, we prove that

$$\|e^{-it\Delta}u(t) - u_0^+\|_{\Sigma} \rightarrow 0 \text{ as } t \rightarrow +\infty.$$

11 This completes the proof for positive times, the one for negative times is similar. \square

12

APPENDIX A. LORENTZ SPACES

13 We recall some basic facts about the Lorentz spaces which are relevant to our study. We
 14 refer the reader to [6, 13, 20, 14, 28] and references therein for more properties and information
 15 on Lorentz spaces.

Definition A.1. Let $u : \mathbb{R}^N \rightarrow \mathbb{R}$ be a measurable function. The distribution function of u is given by

$$\mathbf{d}_u(\lambda) := |\{x \in \mathbb{R}^N; |u(x)| > \lambda\}|, \quad \lambda \in (0, \infty).$$

Here, the notation $|E|$ stands for the N -dimensional Lebesgue measure of E . The (unidimensional) decreasing rearrangement of u , denoted by u^* , is defined by

$$u^*(s) = \inf \{ \lambda > 0; \mathbf{d}_u(\lambda) < s \}, \quad s > 0.$$

16 It is clear that \mathbf{d}_u and u^* are non-negative non-increasing functions. The Lorentz spaces
 17 $L^{p,q}(\mathbb{R}^N)$ are defined as follows.

Definition A.2. Let $0 < p < \infty$ and $0 < q \leq \infty$. Then

$$L^{p,q}(\mathbb{R}^N) = \{ u \text{ measurable; } \|u\|_{L^{p,q}} < \infty \},$$

1 where

$$\|u\|_{L^{p,q}} = \begin{cases} \left(\frac{q}{p} \int_0^\infty (s^{1/p} u^*(s))^q \frac{ds}{s} \right)^{1/q} & \text{if } 0 < p, q < \infty \\ \sup_{s>0} (s^{1/p} u^*(s)) & \text{if } 0 < p < \infty, q = \infty \end{cases}$$

We have $L^{p,p} = L^p$ and by convention $L^{\infty,\infty} = L^\infty$. Another way to define the Lorentz space $L^{p,q}$ is via real interpolation theory as follows (see [4])

$$L^{p,q} = \left[L^1, L^\infty \right]_{1-\frac{1}{p}, q}, \quad 1 < p < \infty, \quad 1 \leq q \leq \infty.$$

2 One of the difficulties in our problem is the singular weight $|x|^{-b}$ in the nonlinearity. Since
3 this weight does not belong to any Lebesgue space we have to treat it differently. Fortunately,
4 $|x|^{-b}$ belongs to the Lorentz space $L^{\frac{2}{b},\infty}(\mathbb{R}^2)$ which plays an important role in our proof.

5 The following lemma will be useful.

Lemma A.3. Let $1 < p < \infty$, $1 < p_1 < \infty$ and $1 \leq p_2 \leq \infty$ be such that

$$\frac{1}{p} = \frac{1}{p_1} + \frac{1}{p_2}.$$

6 Then

$$\|fg\|_{L^p} \leq C \|f\|_{L^1}^{1-\theta} \|f\|_{L^\infty}^\theta \|g\|_{L^{p_2,\infty}} \quad (\text{A.1})$$

7 where $\theta = 1 - \frac{1}{p_1}$.

8

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13

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