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## Annealing-free Si<sub>3</sub>N<sub>4</sub> frequency combs for monolithic integration with Si photonics

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Silicon-nitride-on-insulator (SiNOI) is an attractive platform for optical frequency comb generation in the telecommunication band because of the low two-photon absorption and free carrier induced nonlinear loss when compared with crystalline silicon. However, the high-temperature annealing that has been used so far for demonstrating  $Si_3N_4$ -based frequency combs made the co-integration with silicon-based optoelectronics elusive, thus reducing dramatically its effective complementary metal oxide semiconductor (CMOS) compatibility. We report here on the fabrication and testing of annealing-free SiNOI nonlinear photonic circuits. In particular, we have developed a process to fabricate low-loss, annealing-free, and crack–free  $Si_3N_4$  740-nm-thick films for Kerr-based nonlinear photonics featuring a full process compatibility with front-end silicon photonics. Experimental evidence shows that micro-resonators using such annealing-free silicon nitride films are capable of generating a frequency comb spanning 1300-2100 nm via optical parametrical oscillation based on four-wave mixing. This work constitutes a decisive step toward time-stable power-efficient Kerr-based broadband sources featuring full process compatibility with Si photonic integrated circuits (Si-PICs) on CMOS-lines.

Kerr optical frequency combs constitute a paradigm shift in the development of high-capacity data transmission, integrated spectroscopy, high precision metrology, and frequency synthesis [1]. Since 2010, silicon-nitride-on-insulator (SiNOI) has imposed as an attractive chip-based platform for the generation of wideband frequency combs pumped at telecom wavelengths, because of its relatively high nonlinearity (×10 that of silica and larger than that of highly nonlinear Hydex glass [2]) as well as the absence of two-photon absorption and free carrier generation that plague crystalline silicon. In the meanwhile, silicon photonic integrated circuits (Si-PICs) have demonstrated increasing maturity levels for a wide range of optical functions such as III-V-on-Si integrated lasers [3], high-speed modulators [4], Ge-on-Si photodiodes [5], as well as filters and wavelength (de)multiplexers [6], thus continuously highlighting the potential of silicon optoelectronics integration with cost-effective complementary metal-oxide-semiconductor (CMOS) technology [7, 8]. In this context, the monolithic integration of Kerrbased frequency combs with Si photonics and CMOS micro-nano-electronic nodes holds the promise for on-chip high-capacity transmitters that would benefit from the maturity and low cost of CMOS manufacturing and scalability.

The realization of relatively thick (> 700 nm) stoichiometric Si<sub>3</sub>N<sub>4</sub> films, as required by microring frequency combs, which imply both a tight confinement of light and anomalous group velocity dispersion (GVD), remains challenging. In particular, *all prior works* strictly made use of long high-temperature annealing (~1200 °C for at least 3 h) of the deposited silicon nitride film [9-12]. This extreme annealing step has been accounted for by the need to densify the silicon nitride film through driving out excess hydrogen and break N-H bonds, so as to get closer to a stoichiometric Si<sub>3</sub>N<sub>4</sub> film and reduce the material absorption loss in the C-band. However, this annealing induces thermal stress that eventually leads to cracks during the device processing unless sophisticated pre-patterning strategies are adopted prior to the film deposition [10, 11]. Yet, in the perspective of nonlinear optics-silicon optoelectronic co-integration, these extreme annealing temperatures would severely degrade the frontend-of-the-line (FEOL) silicon optoelectronic circuits underneath. Specifically, doped optical circuits would be unacceptably affected by the undesirable dopant diffusion in junction-based Si modulators and by the hetero-interface degradation of Ge-on-Si photodetectors. As an alternative, optical parametric oscillation has been demonstrated by using the proprietary composition of Hydex waveguides in 2013 [2]. However, the nonlinear parameter  $\gamma$  of Hydex is six times lower than that of silicon nitride, thus increasing the required power for frequency comb generation, hence the risk of waveguide optical damage. Comb generation has also been demonstrated in 2018 using deuterated silicon nitride (SiN:D) deposited at 300 °C [13]. However, although being a CMOS-compatible processing building block, it is worth to point out that deuterated silicon nitride can be exclusively used during the back-end-of-the-line (BEOL) fabrication, thus preventing a straightforward monolithic integration with both silicon photonics and sub-100 nm standard CMOS electronics circuitries during a collective FEOL processing [14]. Furthermore, this material has a relatively strong thermal shift that occurs from the larger residual hydrogen-related absorption coefficient (when compared to purely stoichiometric  $Si_3N_4$  films) making it difficult to stabilize the comb in the soliton state as described in [13]. In addition, for similar resonance quality factors, deuterated silicon nitride microresonators require higher threshold powers for comb generation than stoichiometric  $Si_3N_4$ , which indicates a lower nonlinear index of SiN:D and increases, as for the Hydex, the risk of waveguide optical damage.

Very recently, we reported a method that avoids thermal annealing for growing relatively thick (740 nm) crack-free Si<sub>3</sub>N<sub>4</sub>-based straight nanowaveguides with good linear and nonlinear properties, as measured by self-phase modulation of picosecond pulses propagating along these waveguides [15].

In this paper we report a comb generated from a microresonator made in an *intrinsically* CMOS-compatible annealing-free silicon nitride, following a tailored deposition method and without the typical hour-long 1200 °C post-annealing that is detrimental to the integration with front-end silicon optoelectronic circuits. Our annealing-free and crack-free fabrication process (shown in Fig. 1) provides our devices with the right specification (i.e., microring group velocity dispersion and characteristics) to underpin Kerr frequency combs, thus representing a significant step toward the full compatibility of  $Si_3N_4$ -based Kerr comb sources with the thermal budgets of Si photonics processing. In contrast to all previous approaches, our process does not exceed neither the dopant activation temperature (1030 °C) required for Si modulators [16], nor the H<sub>2</sub> annealing thermal budget used for dislocation control of Ge-on-Si photodiodes (825 °C) [17]. Furthermore, our approach would offer a clear advantage even in the perspective of silicon photonics full integration with sub-100 nm standard micronano-electronics CMOS circuits during collective FEOL processing [18].

In order to control strain and to prevent cracks from appearing, the silicon nitride layer is deposited on a (non-patterned) substrate via low-pressure chemical vapor deposition (LPCVD) in two steps of 370-nm-thick layer each. The deposition is carried out with a tailored ultra-low deposition rate (<2 nm/min) to produce a very high quality film, which is denser optically, and highly nonlinear ( $n_2 = 3.6 \times 10^{-15} \text{ cm}^2 \text{ W}^{-1}$ ) [15].



FIG. 1. Schematics of the annealing-free fabrication process for Si<sub>3</sub>N<sub>4</sub> nonlinear photonics (a)-(f).

Critically, under such a low deposition rate which is 40 % lower than that of standard LPCVD silicon nitride reported in [19] and 30 % lower than the deposition rate mentioned in [20], the thermal activation energy enables silicon and nitrogen to dispose at the silicon nitride film surface via atomic surface migration phenomena, while compelling hydrogen to escape the film. Furthermore, between the two deposition stages, the wafer is rotated by  $45^{\circ}$  in order to distribute the uniaxial strain along the overall film thickness, thus avoiding film cracks upon subsequent subtractive patterning. Each deposition run is carried out at 780 °C with post-deposition cooling to around 630 °C for 20 minutes. Controlled ramp-ups and -downs from/to 780 °C at 10 °C/minute to/from 630 °C are used prior to each deposition which is carried out under a 112 mTorr pressure using NH<sub>3</sub> (200 sccm) and SiH<sub>2</sub>Cl<sub>2</sub> (80 sccm) precursor gases, thus providing a ratio between the two precursor gases (SiH<sub>2</sub>Cl<sub>2</sub>:NH<sub>3</sub>) of 0.4. For comparison, the deposition of silicon-rich silicon nitride requires this ratio to be higher than 2 [21], which typically implies an increase of the SiH<sub>2</sub>Cl<sub>2</sub> gas flow as well as a decrease of the NH<sub>3</sub> gas flow rate. According to [22], a decrease in the NH<sub>3</sub> gas flow inevitably translates into an increase in the deposition rate. Regarding nitride-rich silicon nitride films, the deposition rate is of the order of 17 nm/min [23], which is much higher than our deposition rate too.

The long thermal annealing step adopted in standard LPCVD processes is used to densify the SiN film, through bringing sufficient energy to break up N-H bonds otherwise causing absorption in the material. By reducing the content of residual

hydrogen [11, 19] in the SiN film, the high tensile strain is a clear indication of the film stoichiometry [19, 21]. In order to compare between the stoichiometry of films produced via our annealing-free process and by the standard LPCVD process including a long post-annealing step, we compared their respective tensile strains. The stress of our silicon nitride film is inferred by measuring the wafer bow, before and after removing the silicon nitride from the wafer back side. The material morphological characterization revealed a tensile strain around +1200 MPa, which is 1.5x times higher than the value of 800 MPa reported in [11], employing a long post-annealing step and considered, so far, as the state of the art in stoichiometric silicon nitride processing for nonlinear applications and optical parametric oscillators.

Another indication of the remaining N-H bonds is provided by the spectral dependence of the intrinsic quality factor ( $Q_i$ ) measured for a Si<sub>3</sub>N<sub>4</sub> micro-ring resonator with 1.4 µm waveguide width (see figure 2). It roughly increases by a factor 2 while moving away from the N-H overtone absorption peak (near 1520 nm), showing the presence of residual N-H bonds in our film. The losses in the ring can be estimated by the expression [19]  $\alpha = \frac{k_0 n_g}{\alpha}$ 

where  $k_0$  is the wavenumber and  $n_g$  is the mode group index of the waveguide with similar width. The difference between the losses at 1550 nm ( $Q_i = 350,000$ ) and the losses at 1520 nm ( $Q_i = 190,000$ ) can thus be estimated to be 0.9 dB/cm. This additional loss due to residual N-H absorption for our annealing-free process is comparable to the value (0.6 dB/cm) inferred for high-temperature annealed Si<sub>3</sub>N<sub>4</sub> waveguides [21], and, as shown below, it does not preclude the oscillation and comb generation in the C-band.



FIG. 2.Intrinsic quality factor as a function of wavelength for an annealing-free silicon nitride microring resonator (waveguide width equal to  $1.4 \mu m$ ). The darker colored line shows the mean value and the brighter shadowed areas illustrates the standard deviation of the measurements.

The Si<sub>3</sub>N<sub>4</sub> circuits were patterned through a combination of deep UV lithography (I-line, 780-nm-thick resist) and fluoride-based (CF<sub>4</sub>-CH<sub>2</sub>F<sub>2</sub>-O<sub>2</sub>) dry etching. They were encapsulated by 3- $\mu$ m-thick SiO<sub>2</sub> cladding layer at 400 °C using high-density plasma-enhanced chemical vapor deposition (HDP-PECVD) to avoid void formation, which is considered to be one of the main causes for device failure. Figure 3(b) shows optical and electron scanning microscope images of the fabricated devices.

The combination between the four-wave mixing processes within the resonator and the optical parametric oscillation results in Kerr frequency comb generation. The optical spectrum out of an annealing-free silicon-nitride-on-insulator microring with a 56-µm radius is shown in Fig 3(a). The cross-section dimensions (1.5-µm-wide × 740-nm-thick) of the ring waveguide ensure that GVD is anomalous at the 1569 nm pump wavelength. The intrinsic quality factor of the microring – which is separated by a 350 nm gap from the bus waveguide - exceeds 580,000. A relatively smooth native line spacing (~406 GHz) frequency comb spanning across about 800 nm between 1300 nm - 2100 nm is obtained for ~1 W TE-polarized continuous-wave power ( $P_{in}$ ) coupled in the bus waveguide. This power value is comparable to the power required (2 W) to fill the entire comb span using high-temperature annealed Si<sub>3</sub>N<sub>4</sub> rings with loaded quality factor of 100,000 (intrinsic quality factor 200,000) [2] and to the pump power (3 W) used to generate an optical frequency comb in similar annealed Si<sub>3</sub>N<sub>4</sub> ring resonators [24]. A spectrum recorded for  $P_{in}$  ~300 mW, i.e., well above the oscillation threshold, is also shown in Fig. 3(c). The comb spacing is wider in this case, as dictated by the spectral gain of degenerate four-wave mixing imposed by the device GVD.



FIG. 3. Comb generation using annealing-free silicon nitride on insulator. (a) A 800-nm-spanning comb generation using a 56-  $\mu$ m-radius Si<sub>3</sub>N<sub>4</sub> microresonator for coupled power  $P_{in} \sim 1$  W. (b) Optical microscope and (insets) scanning electron microscope images of the ring and coupling region. (c) Comb spectrum for  $P_{in} \sim 300$  mW.

In order to draw a comparison between our annealing-free process and the standard LPCVD process using a long post-annealing, some wafers have been annealed at 1050 °C under N<sub>2</sub> atmosphere for 2 hours. This annealing condition was the most extreme we could apply on our CMOS line. Quite surprisingly, we did not observe a clear and significant improvement in the resonance linewidths. One primary reason might be related to our relatively low annealing temperature value as compared with that used by other groups (typically > 1200 °C) which is very close to the silicon wafer fusion limit [12]. A second explanation might originate from the significant value of the mode scattering loss induced by the substantial microring side-wall roughness. The latter has still room for improvement to well below <1 nm RMS via sophisticated pattering strategies which are currently under investigation. Such large scattering loss contribution to the Q-factor of our microring resonances may thus hide the relatively minor gain in *Q* factor brought by the 1050 °C annealing step, which likely falls within our Q factor measurement uncertainty. In any case, our work shows that a sufficient amount of hydrogen can be driven out of the Si<sub>3</sub>N<sub>4</sub> film using our deposition process, so as to get closer to stoichiometric silicon nitride films and to generate a frequency comb.

Regarding the damage threshold of our devices, throughout the whole measurement time (~6 months), we observed only one single waveguide fusion event, out of a total tested device number of ~40, which occurred under continuouswave coupled power levels above 38 dBm. We think that such fusion event originates from local waveguide imperfections (such a side-walls scattering or other local defects such as encapsulation voids and impurities) as well as waveguide ageing (notably at the upper cladding  $Si_3N_4$ -SiO<sub>2</sub> interface), partly due to the aforementioned residual N-H absorption tail in the annealing-free SiN films. Again though, this low device failure statistics points out that its impact is rather limited.

In order to measure the threshold power for optical parametric oscillation (OPO), the device was pumped with increasing power levels coupled to the bus waveguide. The laser wavelength was scanned across a resonance at 1570 nm. The chip output light was passed through a short pass filter with a cut-off wavelength of 1560 nm that blocks the pump light with an extinction ratio of 60 dB and was then detected by a photodiode. The peak power of the filtered short wavelength OPO signal at each power coupled to the bus waveguide shows that the measured signal exceeds the noise floor (of -50 dBm) above a threshold power of  $\sim 83$  mW.

We measured the dispersion of our devices by scanning an ECDL light while recording the device transmission and two calibration traces from a fiber cavity and a high-finesse free space cavity with FSR of 171MHz and 175MHz,

respectively. The integrated dispersion of the micro-cavity is  $D_{int}(\mu) = \omega_{\mu} - D_{1}\mu = \frac{D_{2}}{2}\mu^{2} + \frac{D_{3}}{6}\mu^{3} + \dots$  where  $\omega_{\mu}$  is

the angular frequency,  $\mu$  is the mode number relative to the pumped mode and  $D_i$  are the dispersion coefficients. Fig. 4(a,b) depict the measured dispersion and the measurement setup. A fit of the measured dispersion gives an anomalous dispersion value of  $D_2/2\pi = 511$ kHz at 1580 nm which is equivalent to a group velocity dispersion of 29 ps/nm/km. This value compares well with our finite element modeling of the dispersion which gave us an anomalous dispersion estimate of 50 ps/nm/km at the same wavelength. Note that the values of simulated and measured dispersion are in agreement with the simulated dispersion of similar annealed Si<sub>3</sub>N<sub>4</sub> waveguides (1.65- $\mu$ m-wide × 730-nm-thick) presented in [2, 19].



FIG. 4. Dispersion measurement of the TE mode of a ring resonator with 200  $\mu$ m radius and 1.5  $\mu$ m width. (a) Plot of the integrated dispersion. (b) Experimental setup used for the dispersion measurement.

In conclusion, these results show that it is possible to generate a wideband comb at telecom wavelengths using annealing-free silicon nitride nonlinear circuits featuring a full FEOL process compatibility with Si photonics. We demonstrated that annealing-free silicon nitride frequency comb microresonators can be made following a tailored deposition method, minimizing the hydrogen content. Our annealing-free and crack-free fabrication process provides our devices with the right specification (microring group velocity dispersion and characteristics) to underpin Kerr frequency combs, thus representing a significant step toward the full compatibility of Si<sub>3</sub>N<sub>4</sub>-based Kerr-comb sources monolithic integration with standard CMOS and Si photonics processing. Through allowing the monolithic integration of broadband comb sources with CMOS-compatible optoelectronics, our work represents a milestone toward the realization of next-generation Petabit/s data transmitters on a chip.

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#### References

<sup>1</sup>P. Marin-Palomo, J. N. Kemal, M. Kaprov, A. Kordts, J. Pfeifle, M. H. Pfeiffer, P. Trocha, S. Wolf, V. Brasch, M. H. Anderson, R. Rosenberger, K. Vijayan, W. Freude, T. J. Kippenberg, and C. Koos, Nature **546**, 247 (2017).

<sup>2</sup>D. J. Moss, R. Morandotti, A. L. Gaeta, and M. Lipson, Nat. Photonics 7, 597 (2013).

<sup>3</sup>B. Ben Bakir, A. Descos, N. Olivier, D. Bordel, P. Grosse, E. Augendre, L. Fulbert, and J.- M. Fedeli Opt. Exp., **19**, 10317, (2011).

<sup>4</sup>L. Liao, D. Samara-Rubio, M. Morse, A. Liu, D. Hodge, D. Rubin, U. Keil, and T. Franck, Opt. Express, 13, 3135, (2005).
<sup>5</sup>L. Vivien, M. Rouvière, J.-M. Fedeli, D. Marris-Morini, J. F. Damlencourt, J. Mangeney, P. Crozat, L. E. Melhaoui, E. Cassan, X. Le Roux, D. Pascal, S. Laval, Opt. Express, 15, 9848, (2007).

<sup>6</sup>W.Bogaerts, S. k. Selvaraja, P. Dumon, J. Brouckaert, K. De Vos, D. Van Thourhout, and R. Baets, IEEE. J. Sel. Topics in quantum Electronics, **16**, 33 (2010).

<sup>7</sup>G. T. Reed, Nature, **427**, 595–596 (2004).

<sup>8</sup>M. Asghari and A. V. Krishnamoorthy, Nat. Photonics, 5, pp. 268-270 (2011).

<sup>9</sup>J. Levy, A. Gondarenko, M. A. Foster, A. C. Turner-Foster, A. L. Gaeta, and M. Lipson, Nat. Photonics 4, 37, (2010).

<sup>10</sup>K. Luke, A. Dutt, C. B. Poitras, and M. Lipson, Opt. Express 21, 22829 (2013).

<sup>11</sup>M. H. Pfeiffer, A. Kordts, V. Brasch, M. Zervas, M. Geiselmann, J. D. Jost, and T. J. Kippenberg, Optica **3**, 20 (2016).

<sup>12</sup>S. Kim, K. Han, C. Wang, J. A. Jaramillo-Villegas, X. Xue, C. Bao, Y. Xuan, D. E. Leaird, A. M. Weiner , and M. Qi, Nat. Communications **8**, 372 (2017).

<sup>13</sup>J. Chiles, N. Nader, D. D. Hickstein, S. P. Yu, T. C. Briles, D. Carlson, H. Jung, J. M. Shainline, S. Diddams, S. B. Papp, S. W. Nam, and R. P. Mirin, Opt. Letters **43**, 1527 (2018).

<sup>14</sup>Y. A. Vlasov, IEEE Communications Magazine, pp. S67-S72, (2012).

<sup>15</sup>H. El Dirani, M. Casale, S. Kerdiles, C. Socquet-Clerc, X. Letartre, C. Monat, and C. Sciancalepore, IEEE Photon. Technol. Letters **30**, 355 (2018).

<sup>16</sup>M. Streshinsky, R. Ding, Y. Liu, A. Novack, Y. Yang, Y. Ma, X. Tu, E. K. S. Chee, A. E. Lim, P. G. Lo, T. Baehr-Jones, and M. Hochberg, Opt. Express **21**, 30350 (2013).

<sup>17</sup>J. Michel, J. Liu, and L. C. Kimerling, Nat. Photonics 4, 527 (2010).

<sup>18</sup>R. C. Johnson, Electronic Engineering Times, published online

https://www.eetimes.com/document.asp?doc id=1326576, (2015).

<sup>19</sup> J. S. Levy, Ph.D. thesis, Cornell University, New York, 2011.

<sup>20</sup> J. Riemensberger, K. Hartinger, T. Herr, V. Brasch, R. Holzwarth, and T.J. Kippenberg, Opt. Express, **20**, 27661 (2012).

<sup>21</sup>C. J. Krückel, A. Fülöp, Z. Ye, P. A. Andrekson, and V. Torres-Company, Opt. Express, **25**, 15370 (2017).

<sup>22</sup>C. Yang and J. Pham, Silicon, 1 (2018). https://doi.org/10.1007/s12633-018-9791-6

<sup>23</sup>L. Wang, H. S. Reehal, F. L. Martimez, E. S. Andres and A. d. Prado, Semiconductor science and technology, **18**, 633 (2003).

<sup>24</sup>T. Herr, K. Hartinger, J. Riemensberger, C. Y. Wang, E. Gavartin, R. Holzwarth, M. L. Gorodetsky and T. J. Kippenberg, Nat. Photonics **6**, 480 (2012).

<sup>25</sup>M. Mitrovic, G. Xiaowei, J. Hua, L. K. Oxenlowe, and L. H. Frandsen, pp. FM1D-6, OSA, (2015).