



# Chronic myeloid leukemia model with periodic pulsed treatment

Helal Mohamed, Abdelkader Lakmeche, Fethi Souna

## ► To cite this version:

Helal Mohamed, Abdelkader Lakmeche, Fethi Souna. Chronic myeloid leukemia model with periodic pulsed treatment. 2018. hal-01904650v1

HAL Id: hal-01904650

<https://hal.science/hal-01904650v1>

Preprint submitted on 25 Oct 2018 (v1), last revised 9 May 2019 (v3)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Chronic myeloid leukemia model with periodic pulsed treatment

Helal Mohamed<sup>1</sup>, Abdelkader Lakmeche<sup>2</sup> and Fethi Souna<sup>3</sup>

Biomathematics Laboratory

Univ Sidi-Bel-Abbès

Sidi-Bel-Abbès

Algeria

<sup>1</sup>mhelal\_abbes@yahoo.fr, <sup>2</sup>lakmeche@yahoo.fr and <sup>3</sup>fethiou91@gmail.com

**RÉSUMÉ.** Dans ce travail, nous développons un modèle mathématique de la leucémie myéloïde chronique avec un traitement à effets instantanés. Notre analyse se focalise sur les valeurs du taux de croissance  $\gamma$  pour avoir la stabilité ou l'instabilité de l'équilibre sans maladie. Si le taux de croissance  $\gamma$  des cellules souches leucémiques sensibles est inférieur à un seuil  $\gamma^*$ , nous obtenons la stabilité d'un équilibre sans maladie, ce qui signifie que la maladie sera éradiquée pour toute période de traitement  $\tau_0$ . Sinon, pour  $\gamma$  supérieur à  $\gamma^*$ , la durée du traitement doit être inférieure à une valeur spécifique  $\tau_0^*$ . Dans le cas critique où la période de traitement est égale à  $\tau_0^*$ , nous observons une persistance de la tumeur, ce qui signifie que la maladie est viable.

**ABSTRACT.** In this work we develop a mathematical model of chronic myeloid leukemia including treatment with instantaneous effects. Our analysis focuses on the values of growth rate  $\gamma$  which give either stability or instability of the disease free equilibrium. If the growth rate  $\gamma$  of sensitive leukemic stem cells is less than some threshold  $\gamma^*$ , we obtain the stability of disease free equilibrium which means that the disease is eradicated for any period of treatment  $\tau_0$ . Otherwise, for  $\gamma$  great than  $\gamma^*$ , the period of treatment must be less than some specific value  $\tau_0^*$ . In the critical case when the period of treatment is equal to  $\tau_0^*$ , we observe a persistence of the tumor, which means that the disease is viable.

**MOTS-CLÉS :** Solution périodique, Equation différentielle impulsive, Théorème de point fixe, Bifurcation.

**KEYWORDS :** Positive solution, impulsive differential equation, Fixed point theorem, Bifurcation.

In this paper, we study a mathematical model of chronic myeloid leukemia (CML) under treatment, the model studied here is inspired from [9] and [16].

Several recent works have been developed to study the dynamics of CML under chemotherapy treatment, see ([14], [15], [16] and [17]).

In this work, we consider the following mathematical model which is an extension of the model proposed in [9]. In our model, we assume that normal (resp. leukemic) cells differentiate through two stages of their life cycle, beginning with normal (resp. sensitive leukemic) stem cells which produce normal (resp. leukemic) progenitor cells.

The mathematical form of the system we shall investigate is the following

$$\begin{cases} \dot{x}_0 = (\beta - a_x - \beta_0 x_0 - \lambda(x_1 + y_1))x_0, \\ \dot{x}_1 = a_x x_0 - d_1 x_1, \\ \dot{y}_0 = (\gamma - a_y - \gamma_0 y_0 - \lambda(x_1 + \alpha y_1))y_0, \\ \dot{y}_1 = a_y y_0 - d_2 y_1, \end{cases} \quad [1]$$

with initial conditions

$$x_0(0) \geq 0, x_1(0) \geq 0, y_0(0) \geq 0 \text{ and } y_1(0) \geq 0. \quad [2]$$

Throughout this paper, we assume the following conditions

$$a_y < \gamma, \quad [3]$$

and

$$a_x < \beta. \quad [4]$$

The conditions (3) and (4) are necessary to obtain the biological meanings of the state variables  $x_0$ ,  $x_1$ ,  $y_0$  and  $y_1$ .

In [9], well-posedness of (1), (2) is proved and stability (local and global) of equilibria is investigated. In fact, the disease free equilibrium  $E_f = (x_0^*, \frac{a_x}{d_1}x_0^*, 0, 0)$  is locally asymptotically stable for growth rate of sensitive leukemic stem cells  $\gamma > \gamma^* := a_y + \lambda \frac{a_x}{d_1} x_0^*$  and unstable for  $\gamma < \gamma^*$  where  $x_0^* = \frac{(\beta - a_x)d_1}{\beta_0 d_1 + \lambda a_x}$ .

In this paper, we consider the model above including chemotherapeutic treatment. We obtain the following model

$$\dot{x}_0(t) = F_1(x_0, x_1, y_0, y_1), \quad [5]$$

$$\dot{x}_1(t) = F_2(x_0, x_1, y_0, y_1), \quad [6]$$

$$\dot{y}_0(t) = F_3(x_0, x_1, y_0, y_1), \quad [7]$$

$$\dot{y}_1(t) = F_4(x_0, x_1, y_0, y_1), \quad [8]$$

for  $t > 0$  and  $t \neq t_i$ , where  $t_i$  is the time of the  $i^{th}$  treatment,

$$F_1(x_0, x_1, y_0, y_1) = (\beta - a_x - \beta_0 x_0 - \lambda(x_1 + y_1))x_0, \quad [9]$$

$$F_2(x_0, x_1, y_0, y_1) = a_x x_0 - d_1 x_1, \quad [10]$$

$$F_3(x_0, x_1, y_0, y_1) = (\gamma - a_y - \gamma_0 y_0 - \lambda(x_1 + \alpha y_1))y_0, \text{ and} \quad [11]$$

$$F_4(x_0, x_1, y_0, y_1) = a_y y_0 - d_2 y_1. \quad [12]$$

For  $t = t_i$  we have

$$x_0(t_i^+) = \Theta_1(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)), \quad [9]$$

$$x_1(t_i^+) = \Theta_2(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)), \quad [10]$$

$$y_0(t_i^+) = \Theta_3(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)), \quad [11]$$

$$y_1(t_i^+) = \Theta_4(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)), \quad [12]$$

where  $x_j(t_i^+) = \lim_{\substack{t \rightarrow t_i \\ t > t_i}} x_j(t)$  and  $y_j(t_i^+) = \lim_{\substack{t \rightarrow t_i \\ t > t_i}} y_j(t)$ , ( $j = 0, 1$ ) are the size of  $x_j$  just after the  $i^{th}$  treatment. In our case we have

$$\Theta_1(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)) = x_0(t_i),$$

$$\Theta_2(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)) = x_1(t_i),$$

$$\Theta_3(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)) = T_0 y_0(t_i),$$

$$\Theta_4(x_0(t_i), x_1(t_i), y_0(t_i), y_1(t_i)) = T_1 y_1(t_i).$$

The variables and parameters are

$x_0$  : biomass of the normal stem cells,

$x_1$  : biomass of normal progenitor cells,

$y_0$  : biomass of sensitive leukemic stem cells,

$y_1$  : biomass of sensitive leukemic sensitive progenitor cells,

$\beta_0$  : death rate of the normal stem cells,

$\gamma_0$  : death rate of sensitive leukemic stem cells,

$\beta$  : growth rate of normal stem cells,

$\gamma$  : growth rate of sensitive leukemic stem cells,

$\lambda$  : competitive parameter of the stem and progenitor cells,

$a_x$  : produce rate of the normal stem cells,

$a_y$  : produce rate of the sensitive leukemic stem cells,

$d_1$  : death rate of the normal progenitor cells,

$d_2$  : death rate of the sensitive leukemic progenitor cells,

$\alpha$  : competition parameter ( $0 < \alpha < 1$ ),

$T_0 (< 1)$  : survival fraction of sensitive leukemic stem cells,

$T_1 (< 1)$  : survival fraction of sensitive leukemic progenitor cells, and

$\tau$  : is the time of the first treatment, it's the period between two successive injections, that is  $t_i = i\tau$ ,  $i \in \mathbb{N}$ .

We obtain a special kind of differential equations called impulsive differential equations (see [1]-[6] and [8]).

Our aim is to study (5)-(12), more specifically to study the existence of the healthy steady state and its stability. Moreover, we study the existence of bifurcated periodic positive solutions which represent the persistence of the disease as in [7] and [11]-[13].

## 1. Stability of the disease free equilibrium $E_f$

We can show that  $\zeta(t) := \zeta_0 = (\frac{(\beta - a_x)d_1}{\beta_0 d_1 + \lambda a_x}, \frac{\lambda a_x(\beta - a_x)}{\beta_0 d_1 + \lambda a_x}, 0, 0) = E_f$  is a constant equilibrium of (5)-(12), it is called trivial solution.

To study the stability of  $\zeta$  we use the same approach of fixed point process than in [7] and [11]-[13].

Since solutions of (5)-(8) exist globally in  $\mathbf{R}_+$  and are nonnegative (see [9]) we have

$$X(t) = \Phi(t, X_0), t \geq 0 \quad [13]$$

where  $X(t) = (x_0, x_1, y_0, y_1)(t)$ ,  $X(0) = X_0$  and  $\Phi$  is the flow associated to (5)-(12).

The term  $X(\tau^+)$  denotes the state of the population after the treatment,  $X(\tau^+) = \Theta(X(\tau)) = \Theta(\Phi(\tau, X_0))$ .

To have periodic solution we must have  $X(\tau^+) = X_0$  that is  $X_0 = \Theta(\Phi(\tau, X_0))$ .

Let  $\Psi$  be the operator defined by

$$\Psi(\tau, X_0) = \Theta(\Phi(\tau, X_0)) \quad [14]$$

and denote by  $D_X \Psi$  the derivative of  $\Psi$  with respect to  $X$ . Then  $X = \Phi(., X_0)$  is a  $\tau$ -periodic solution of (5)-(12) if and only if

$$\Psi(\tau, X_0) = X_0, \quad [15]$$

i.e.  $X_0$  is a fixed point of  $\Psi(\tau, .)$ , and it is exponentially stable if and only if the spectral radius  $\rho(D_X \Psi(\tau, .))$  is strictly less than 1 (see [10]).

We have  $D_X \Psi(\tau, X_0) = D_X \Theta(\Phi(\tau, X_0)) \frac{\partial \Phi}{\partial X}(\tau, X_0)$ .

Then, for  $X_0 = \zeta_0$  and  $\tau = \tau_0$  we have

$$\begin{aligned} D_X \Psi(\tau_0, \zeta_0) &= D_X \Theta(\Phi(\tau_0, \zeta_0)) \frac{\partial \Phi}{\partial X}(\tau_0, \zeta_0) \\ &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & T_0 & 0 \\ 0 & 0 & 0 & T_1 \end{pmatrix} \begin{pmatrix} \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_1} & 0 & 0 \\ 0 & 0 & \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_0} & 0 \\ 0 & 0 & 0 & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_0} + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \end{pmatrix} \\ &= \begin{pmatrix} \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_1} & 0 & 0 \\ 0 & 0 & T_0 \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_0} & 0 \\ 0 & 0 & 0 & T_1 \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \end{pmatrix}. \end{aligned}$$

The equilibrium  $\zeta$  is exponentially stable if and only if the spectral radius is less than one.  
We have

$$\det(D_X \Psi(\tau_0, \zeta_0) - \mu I) = \left( T_0 \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_0} - \mu \right) \left( T_1 \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial y_1} - \mu \right) \chi(\mu) \quad [16]$$

where

$$\chi(\mu) = \mu^2 - \left( \frac{\partial \Phi_1}{\partial x_0}(\tau_0, \zeta_0) + \frac{\partial \Phi_2}{\partial x_1}(\tau_0, \zeta_0) \right) \mu + \left( \frac{\partial \Phi_1}{\partial x_0}(\tau_0, \zeta_0) \frac{\partial \Phi_2}{\partial x_1}(\tau_0, \zeta_0) - \frac{\partial \Phi_1}{\partial x_1}(\tau_0, \zeta_0) \frac{\partial \Phi_2}{\partial x_0}(\tau_0, \zeta_0) \right). \quad [17]$$

From (16) and (17), the equilibrium  $\zeta = E_f$  is exponentially stable if and only if  
 $T_0 \left| \frac{\partial \Phi_3}{\partial y_0}(\tau_0, \zeta_0) \right| < 1$ ,  $T_1 \left| \frac{\partial \Phi_4}{\partial y_1}(\tau_0, \zeta_0) \right| < 1$  and  $|\mu_{\pm}| < 1$ . Where  $\mu_{\pm} = \frac{(\frac{\partial \Phi_1}{\partial x_0}(\tau_0, \zeta_0) + \frac{\partial \Phi_2}{\partial x_1}(\tau_0, \zeta_0)) \pm \sqrt{\Delta}}{2}$   
and

$$\Delta = \left( \frac{\partial \Phi_1}{\partial x_0}(\tau_0, \zeta_0) - \frac{\partial \Phi_2}{\partial x_1}(\tau_0, \zeta_0) \right)^2 + 4 \frac{\partial \Phi_1}{\partial x_1}(\tau_0, \zeta_0) \frac{\partial \Phi_2}{\partial x_0}(\tau_0, \zeta_0). \quad [18]$$

From the variational equation  $\frac{d}{dt}(D_X \Phi(t, \zeta_0)) = \frac{\partial F}{\partial X}(\zeta_0) \frac{\partial \Phi}{\partial X}(t, \zeta_0)$ , we have for all  $0 < t \leq \tau_0$

$$\begin{cases} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} &= \frac{d_1+u_2}{u_2-u_1} e^{u_2 t} - \frac{d_1+u_1}{u_2-u_1} e^{u_1 t}, \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} &= \frac{a_x}{u_2-u_1} (e^{u_2 t} - e^{u_1 t}), \\ \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} &= -\frac{d_1+u_2}{u_2-u_1} \frac{d_1+u_1}{a_x} (e^{u_2 t} - e^{u_1 t}), \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} &= -\frac{d_1+u_1}{u_2-u_1} e^{u_2 t} + \frac{d_1+u_2}{u_2-u_1} e^{u_1 t}, \\ \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} &= e^{(\gamma-a_y-\lambda \frac{a_x}{d_1} x_0^*) t}, \text{ and} \\ \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1} &= e^{-d_2 t} \end{cases}$$

where  $u_1 = \frac{-\beta x_0^* - d_1 - \sqrt{\Delta_1}}{2} < 0$ ,  $u_2 = \frac{-\beta x_0^* - d_1 + \sqrt{\Delta_1}}{2} < 0$  and  $\Delta_1 = (\beta_0 x_0^* - d_1)^2 - 4\lambda a_x x_0^* > 0$  for either

$$0 < \beta - a_x < \frac{(\beta_0 d_1 + \lambda a_x)[(\beta_0 d_1 + 2\lambda a_x) - 2\sqrt{\lambda a_x(\beta_0 d_1 + \lambda a_x)}]}{\beta_0^2 d_1}$$

or

$$\beta - a_x > \frac{(\beta_0 d_1 + \lambda a_x)[(\beta_0 d_1 + 2\lambda a_x) + 2\sqrt{\lambda a_x(\beta_0 d_1 + \lambda a_x)}]}{\beta_0^2 d_1}$$

(see Appendix, Subsection 4.1).

Then, we obtain  $\Delta = (e^{u_2 t} - e^{u_1 t})^2$ ,  $\mu_- = e^{u_1 \tau_0} \in (0, 1)$  and  $\mu_+ = e^{u_2 \tau_0} \in (0, 1)$ .

Let (H1) :  $a_x < \beta < a_x + \frac{(\beta_0 d_1 + \lambda a_x)[(\beta_0 d_1 + 2\lambda a_x) - 2\sqrt{\lambda a_x(\beta_0 d_1 + \lambda a_x)}]}{\beta_0^2 d_1}$  or  
 $\beta > a_x + \frac{(\beta_0 d_1 + \lambda a_x)[(\beta_0 d_1 + 2\lambda a_x) + 2\sqrt{\lambda a_x(\beta_0 d_1 + \lambda a_x)}]}{\beta_0^2 d_1}$ .

We deduce the following results.

### Theorem 1.1

Let (H1) be satisfied.

1) If  $\gamma \leq \gamma^*$ , then the trivial solution  $\zeta$  is exponentially stable for all  $\tau_0 > 0$ .

2) If  $\gamma > \gamma^*$ , then the trivial solution  $\zeta$  is exponentially stable for  $\tau_0 < \tau_0^* := \frac{\ln T_0}{-\gamma + a_y + \lambda \frac{a_x}{d_1} x_0^*}$  and instable for  $\tau > \tau_0^*$ .

REMARK. — I

f(H1) is satisfied and  $\gamma > \gamma^*$  we have  $T_1 \left| \frac{\partial \Phi_4}{\partial y_1}(\tau_0, \zeta_0) \right| < 1$  for  $\tau_0 < \tau_0^*$  and  $T_0 \left| \frac{\partial \Phi_3}{\partial y_0}(\tau_0, \zeta_0) \right| = 1$  for  $\tau_0 = \tau_0^*$ . That is we have a critical case at  $\tau_0 = \tau_0^*$ .

REMARK. — F

rom theorem 1.1, if (H1) is satisfied we show that in case of low growth rate of leukemic sensitive stem cells  $\gamma (\leq \gamma^*)$  we can choose any period  $\tau_0$  of treatment to have eradication of the disease. Otherwise, for high growth rate of leukemic sensitive stem cells  $\gamma (> \gamma^*)$ , the eradication of the disease is acquired only for period  $\tau_0$  less than some threshold  $\tau_0^*$ .

## 2. Bifurcation Analysis of nontrivial periodic solution

In this section, we analyze the bifurcation of nontrivial periodic solutions of system (5) – (12) from  $\zeta$  at  $\tau_0 = \tau_0^*$ . This case is possible if (H1) is satisfied and  $\gamma > \gamma^*$  (see theorem 1.1). The bifurcated solutions means that the disease is installed.

Let  $\bar{\tau}$  and  $\bar{X}$  such that  $\tau = \tau_0^* + \bar{\tau}$  and  $X = \zeta_0 + \bar{X}$ . The equation (15) is equivalent to

$$M(\bar{\tau}, \bar{X}) = 0, \quad [19]$$

where  $M(\bar{\tau}, \bar{X}) = (M_1(\bar{\tau}, \bar{X}), \dots, M_4(\bar{\tau}, \bar{X})) := \zeta_0 + \bar{X} - \Psi(\tau_0 + \bar{\tau}, \zeta_0 + \bar{X})$ . If  $(\bar{\tau}, \bar{X})$  is a zero of  $M$ , then  $(\zeta_0 + \bar{X})$  is a fixed point of  $\Psi(\tau_0^* + \bar{\tau}, .)$ . Let

$$D_X M(\bar{\tau}, \bar{X}) = \begin{pmatrix} a & b & c & d \\ e & f & * & * \\ * & g & * & * \\ * & h & i & * \end{pmatrix}. \quad [20]$$

For  $(\bar{\tau}, \bar{X}) = (0, (0, 0, 0, 0))$ , we have

$$\begin{aligned} D_X M(0, (0, 0, 0, 0)) &= \begin{pmatrix} a_0 & b_0 & c_0 & d_0 \\ e_0 & f_0 & * & * \\ * & g_0 & * & * \\ * & h_0 & i_0 & * \end{pmatrix} \\ &= \begin{pmatrix} 1 - \frac{\partial \Phi_1}{\partial x_0} & -\frac{\partial \Phi_1}{\partial x_1} & -\frac{\partial \Phi_1}{\partial y_0} & -\frac{\partial \Phi_1}{\partial y_1} \\ -\frac{\partial \Phi_2}{\partial x_0} & 1 - \frac{\partial \Phi_2}{\partial x_1} & 0 & 0 \\ 0 & 0 & 1 - T_0 \frac{\partial \Phi_3}{\partial y_0} & 0 \\ 0 & 0 & -T_1 \frac{\partial \Phi_4}{\partial y_0} & 1 - T_1 \frac{\partial \Phi_4}{\partial y_1} \end{pmatrix} (\tau_0^*, \zeta_0). \end{aligned}$$

Then  $a_0 = 1 - \frac{\partial \Phi_1}{\partial x_0}(\tau_0^*, \zeta_0)$ ,  $b_0 = -\frac{\partial \Phi_1}{\partial x_1}(\tau_0^*, \zeta_0)$ ,  $c_0 = -\frac{\partial \Phi_1}{\partial y_0}(\tau_0^*, \zeta_0)$ ,  $d_0 = -\frac{\partial \Phi_1}{\partial y_1}(\tau_0^*, \zeta_0)$ ,  $e_0 = -\frac{\partial \Phi_2}{\partial x_0}(\tau_0^*, \zeta_0)$ ,  $f_0 = 1 - \frac{\partial \Phi_2}{\partial x_1}(\tau_0^*, \zeta_0)$ ,  $g_0 = 1 - T_0 \frac{\partial \Phi_3}{\partial y_0}(\tau_0^*, \zeta_0)$ ,  $h_0 = -T_1 \frac{\partial \Phi_4}{\partial y_0}(\tau_0^*, \zeta_0)$  and  $i_0 = 1 - T_1 \frac{\partial \Phi_4}{\partial y_1}(\tau_0^*, \zeta_0)$ .

We have the critical cases if and only if  $\det D_X M(0, (0, 0, 0, 0)) = (a_0 f_0 - b_0 e_0) g_0 i_0 = 0$ . That is

$$g_0 = 0 \quad [21]$$

since  $i_0 = 1 - T_1 e^{-d_2 \tau_0^*} \in (0, 1)$  and  $a_0 f_0 - b_0 e_0 = (1 - e^{u_2 \tau_0^*})(1 - e^{u_1 \tau_0^*}) \in (0, 1)$ .

We have  $M(0, (0, 0, 0, 0)) = 0$ . Let  $D_X M(0, (0, 0, 0, 0)) = E$ , then  $\dim \ker(E) = \text{co dim } \mathcal{R}(E) = 1$ . Denote by  $P_1$  and  $P_2$  the projectors onto  $\ker(E)$  and  $\mathcal{R}(E)$  respectively, such that  $P_1 + P_2 = Id_{\mathbf{R}^4}$ ,  $P_1 \mathbf{R}^4 = \text{span}\{Y_0\} = \ker(E)$ , with  $Y_0 = (q_1, q_2, 1, q_4)$ ,  $q_1 = \frac{f_0(c_0 i_0 - d_0 h_0)}{i_0(b_0 e_0 - a_0 f_0)}$ ,  $q_2 = -\frac{e_0(c_0 i_0 - d_0 h_0)}{i_0(b_0 e_0 - a_0 f_0)}$ ,  $q_4 = \frac{-h_0}{i_0}$  and

$P_2 \mathbf{R}^4 = \text{span}\{(1, 0, 0, 0), (0, 1, 0, 0), (0, 0, 0, 1)\} = \mathcal{R}(E)$ .

Then  $(I - P_1) \mathbf{R}^4 = \text{span}\{(1, 0, 0, 0), (0, 1, 0, 0), (0, 0, 0, 1)\}$  and

$(I - P_2) \mathbf{R}^4 = \text{span}\{(0, 0, 1, 0)\}$ .

Equation (19) is equivalent to

$$\begin{cases} M_1(\bar{\tau}, \sigma Y_0 + Z) = 0, \\ M_2(\bar{\tau}, \sigma Y_0 + Z) = 0, \\ M_3(\bar{\tau}, \sigma Y_0 + Z) = 0, \\ M_4(\bar{\tau}, \sigma Y_0 + Z) = 0, \end{cases} \quad [22]$$

where  $Z = (z_1, z_2, 0, z_4)$ ,  $(\bar{\tau}, \bar{X}) = (\bar{\tau}, \sigma Y_0 + Z)$  and  $(\sigma, z_1, z_2, z_4) \in \mathbf{R}^4$ .  
From the three equations of (22), we have

$$\det \begin{pmatrix} \frac{\partial M_1(0, (0, 0, 0, 0))}{\partial z_1} & \frac{\partial M_1(0, (0, 0, 0, 0))}{\partial z_2} & \frac{\partial M_1(0, (0, 0, 0, 0))}{\partial z_4} \\ \frac{\partial M_2(0, (0, 0, 0, 0))}{\partial z_1} & \frac{\partial M_2(0, (0, 0, 0, 0))}{\partial z_2} & \frac{\partial M_2(0, (0, 0, 0, 0))}{\partial z_4} \\ \frac{\partial M_4(0, (0, 0, 0, 0))}{\partial z_1} & \frac{\partial M_4(0, (0, 0, 0, 0))}{\partial z_2} & \frac{\partial M_4(0, (0, 0, 0, 0))}{\partial z_4} \end{pmatrix} = \det \begin{pmatrix} a_0 & b_0 & d_0 \\ e_0 & f_0 & 0 \\ 0 & 0 & i_0 \end{pmatrix} = i_0(a_0f_0 - e_0b_0) \neq 0.$$

From the implicit function theorem, there exist a unique continuous function  $Z^*$ , such that

$$Z^*(\bar{\tau}, \sigma) = (z_1^*(\bar{\tau}, \sigma), z_2^*(\bar{\tau}, \sigma), 0, z_4^*(\bar{\tau}, \sigma)), Z^*(0, 0) = (0, 0, 0, 0) \text{ and}$$

$$M_i(\bar{\tau}, (q_1\sigma + z_1^*(\bar{\tau}, \sigma), q_2\sigma + z_2^*(\bar{\tau}, \sigma), \sigma, q_4\sigma + z_4^*(\bar{\tau}, \sigma))) = 0, \quad [23]$$

for  $i = 1, 2, 4$ , with  $\sigma$  and  $\bar{\tau}$  small enough.

We have  $\frac{\partial Z^*}{\partial \bar{\tau}}(0, 0) = 0$  and  $\frac{\partial Z^*}{\partial \sigma}(0, 0) = (-\frac{f_0 c_0}{a_0 f_0 - e_0 b_0}, \frac{c_0 e_0}{a_0 f_0 - e_0 b_0}, 0, 0)$  (see Appendix, subsection 4.3).

Then  $M(\bar{\tau}, \bar{X}) = 0$  if and only if

$$\omega(\bar{\tau}, \sigma) = M_3(\bar{\tau}, (q_1\sigma + z_1^*(\bar{\tau}, \sigma), q_2\sigma + z_2^*(\bar{\tau}, \sigma), \sigma, q_4\sigma + z_4^*(\bar{\tau}, \sigma))) = 0. \quad [24]$$

We find  $\omega(0, 0) = 0$  and  $\frac{\partial \omega(0, 0)}{\partial \bar{\tau}} = \frac{\partial \omega(0, 0)}{\partial \sigma} = 0$  (see Appendix, subsection 4.4).

Let  $\mathcal{A} = \frac{\partial^2 \omega(0, 0)}{\partial \bar{\tau}^2}$ ,  $\mathcal{B} = \frac{\partial^2 \omega(0, 0)}{\partial \bar{\tau} \partial \sigma}$  and  $\mathcal{C} = \frac{\partial^2 \omega(0, 0)}{\partial \sigma^2}$ . It's shown that  $\mathcal{A} = 0$  (see Appendix, subsection 4.5). Hence

$$\omega(\bar{\tau}, \sigma) = \mathcal{B}\bar{\tau}\sigma + \mathcal{C}\frac{\sigma^2}{2} + o(|\sigma|^2 + |\bar{\tau}|^2),$$

where

$$\begin{aligned} \mathcal{B} &= -\frac{\partial \Theta_3}{\partial y_0} \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau} \partial y_0} \\ &= -(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*) < 0 \end{aligned}$$

and

$$\begin{aligned} \mathcal{C} &= -\frac{\partial \Theta_3}{\partial y_0} \left\{ 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial x_0 \partial y_0} \left( q_1 + \frac{\partial z_1^*(0, 0)}{\partial \sigma} \right) + 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial x_1 \partial y_0} \left( q_2 + \frac{\partial z_2^*(0, 0)}{\partial \sigma} \right) \right. \\ &\quad \left. + 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial y_0 \partial y_1} q_4 + \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial y_0^2} \right\} \\ &= 2\lambda \frac{a_x}{u_2 - u_1} \left( \frac{e^{u_2 \tau_0} - 1}{u_2} - \frac{e^{u_1 \tau_0} - 1}{u_1} \right) \frac{f_0(d_0 h_0 - 2c_0 i_0)}{i_0(a_0 f_0 - b_0 e_0)} \\ &\quad + 2\lambda \left( \frac{d_1 + u_2}{u_2 - u_1} \left( \frac{e^{u_2 \tau_0} - 1}{u_2} \right) - \frac{d_1 + u_1}{u_2 - u_1} \left( \frac{e^{u_1 \tau_0} - 1}{u_1} \right) \right) \frac{e_0(d_0 h_0 - 2c_0 i_0)}{i_0(a_0 f_0 - b_0 e_0)} \\ &\quad + 2\lambda \alpha \left( \frac{e^{-d_2 \tau_0} - 1}{d_2} \right) \frac{h_0}{i_0} + 2\gamma_0 \left( \frac{e^{(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*) \tau_0} - 1}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*} \right) \\ &\quad + \lambda \alpha \frac{a_y}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* + d_2} \left( \frac{e^{(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*) \tau_0} - 1}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*} + \frac{e^{-d_2 \tau_0} - 1}{d_2} \right). \end{aligned}$$

For  $\lambda = 0$ , we have

$$\mathcal{C} = 2\gamma_0 \left( \frac{e^{(\gamma-a_y)\tau_0}-1}{\gamma-a_y} \right) > 0.$$

We have the following theorem.

**Theorem 2.1** *Let (H1) be satisfied and  $\gamma > \gamma^*$ . There exist  $\lambda_0 > 0$  such that for  $\lambda \in (0, \lambda_0)$  we have a supercritical bifurcation of nontrivial periodic solutions of (5)-(12) with period  $\tau(\sigma) = \tau_0^* + \bar{\tau}(\sigma)$  starting from  $\zeta_0 + \sigma Y_0 + Z^*(\bar{\tau}(\sigma), \sigma)$  for  $\sigma (> 0)$  small enough where  $\bar{\tau}(\sigma) = -\frac{c}{2B}\sigma + o(\sigma)$ .*

REMARK. — F

rom theorem 2.1, we deduce that for high growth rate of leukemic sensitive stem cells  $\gamma (> \gamma^*)$  and period of treatment dose  $\tau_0 = \tau_0^*$  there is lost of stability of the disease free equilibrium and we note the presence of nontrivial periodic solution which means that the disease is installed for period  $\tau_0(\sigma)$  close to  $\tau_0^*$ .

### 3. Conclusions

In this work we have analyzed a mathematical model of chronic myeloid leukemia (CML) which is an extension of a model developed in [9] in the case without medical treatment. In our work, we considered the case of a treatment with instantaneous effect described by discrete equations called impulses. We have studied the stability of the healthy equilibrium (trivial solution), it becomes stable if the growth rate of resistant stem cells  $\gamma$  does not exceed a certain threshold  $\gamma^*$ , if it reaches this threshold we obtain a critical case which gives bifurcations what we want say that is the tumor persists and remains viable.

### 4. Appendix

#### 4.1. First derivatives of $\Phi$

For all  $t \in (0, \tau]$ , we have  $\frac{d}{dt} D_X(\Phi(t, \zeta_0)) = \frac{\partial F}{\partial X}(\zeta_0) \frac{\partial \Phi}{\partial X}(t, \zeta_0)$  with the initial condition  $D_X(\Phi(0, \zeta_0)) = I_{\mathbf{R}^4}$ , where

$$\frac{d}{dt} D_X(\Phi(t, \zeta_0)) = \frac{d}{dt} \begin{pmatrix} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_3(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_4(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1} \end{pmatrix},$$

$$\begin{aligned}\frac{\partial F}{\partial X}(\zeta_0) &= \begin{pmatrix} \frac{\partial F_1(\zeta(t))}{\partial x_0} & \frac{\partial F_1(\zeta(t))}{\partial x_1} & \frac{\partial F_1(\zeta(t))}{\partial y_0} & \frac{\partial F_1(\zeta(t))}{\partial y_1} \\ \frac{\partial F_2(\zeta(t))}{\partial x_0} & \frac{\partial F_2(\zeta(t))}{\partial x_1} & \frac{\partial F_2(\zeta(t))}{\partial y_0} & \frac{\partial F_2(\zeta(t))}{\partial y_1} \\ \frac{\partial F_3(\zeta(t))}{\partial x_0} & \frac{\partial F_3(\zeta(t))}{\partial x_1} & \frac{\partial F_3(\zeta(t))}{\partial y_0} & \frac{\partial F_3(\zeta(t))}{\partial y_1} \\ \frac{\partial F_4(\zeta(t))}{\partial x_0} & \frac{\partial F_4(\zeta(t))}{\partial x_1} & \frac{\partial F_4(\zeta(t))}{\partial y_0} & \frac{\partial F_4(\zeta(t))}{\partial y_1} \end{pmatrix} \\ &= \begin{pmatrix} -\beta_0 x_0^* & -\lambda x_0^* & 0 & -\lambda x_0^* \\ a_x & -d_1 & 0 & 0 \\ 0 & 0 & \gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* & 0 \\ 0 & 0 & a_y & -d_2 \end{pmatrix},\end{aligned}$$

and

$$\frac{\partial \Phi}{\partial X}(t, \zeta_0) = \begin{pmatrix} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_2(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_3(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_1} \\ \frac{\partial \Phi_4(t, \zeta_0)}{\partial x_0} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial x_1} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0} & \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1} \end{pmatrix}.$$

From Cauchy Lipschitz theorem (uniqueness of solution) we obtain that  $\frac{\partial \Phi_2(t, \zeta_0)}{\partial y_i} = \frac{\partial \Phi_3(t, \zeta_0)}{\partial x_i} = \frac{\partial \Phi_4(t, \zeta_0)}{\partial x_i} = 0$ ,  $i \in \{0, 1\}$  and  $\frac{\partial \Phi_3(t, \zeta_0)}{\partial y_1} = 0$ . Moreover, we have

$$\frac{d}{dt} \left( \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} \right) = \frac{\partial F_1(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} + \frac{\partial F_1(\zeta(t))}{\partial x_1} \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0}, \quad [25]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} \right) = \frac{\partial F_1(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} + \frac{\partial F_1(\zeta(t))}{\partial x_1} \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1}, \quad [26]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_0} \right) = \frac{\partial F_1(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_0} + \frac{\partial F_1(\zeta(t))}{\partial y_1} \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0}, \quad [27]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_1} \right) = \frac{\partial F_1(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial y_1} + \frac{\partial F_1(\zeta(t))}{\partial y_1} \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1}, \quad [28]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} \right) = \frac{\partial F_2(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} + \frac{\partial F_2(\zeta(t))}{\partial x_1} \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0}, \quad [29]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} \right) = \frac{\partial F_2(\zeta(t))}{\partial x_0} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} + \frac{\partial F_2(\zeta(t))}{\partial x_1} \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1}, \quad [30]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} \right) = \frac{\partial F_3(\zeta(t))}{\partial y_0} \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0}, \quad [31]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0} \right) = \frac{\partial F_4(\zeta(t))}{\partial y_0} \frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} + \frac{\partial F_4(\zeta(t))}{\partial y_1} \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0}, \quad [32]$$

$$\frac{d}{dt} \left( \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1} \right) = \frac{\partial F_4(\zeta(t))}{\partial y_1} \frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1}. \quad [33]$$

From (31) we obtain  $\frac{\partial \Phi_3(t, \zeta_0)}{\partial y_0} = e^{(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*)t}$ .

From (32) we obtain  $\frac{\partial \Phi_4(t, \zeta_0)}{\partial y_0} = \frac{a_y}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* + d_2} (e^{(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^*)t} - e^{-d_2 t})$ .

From (33) we have  $\frac{\partial \Phi_4(t, \zeta_0)}{\partial y_1} = e^{-d_2 t}$ .

From (25) and (29) we have  $\begin{pmatrix} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} \end{pmatrix} = e^{tA} \begin{pmatrix} \frac{\partial \Phi_1(0, \zeta_0)}{\partial x_0} \\ \frac{\partial \Phi_2(0, \zeta_0)}{\partial x_0} \end{pmatrix} = e^{tA} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ , where

$$A = \begin{pmatrix} \frac{\partial F_1(\zeta(t))}{\partial x_0} & \frac{\partial F_1(\zeta(t))}{\partial x_1} \\ \frac{\partial F_2(\zeta(t))}{\partial x_0} & \frac{\partial F_2(\zeta(t))}{\partial x_1} \end{pmatrix} = \begin{pmatrix} -\beta_0 x_0^* & -\lambda x_0^* \\ a_x & -d_1 \end{pmatrix} = P V P^{-1}, \quad V = \begin{pmatrix} u_1 & 0 \\ 0 & u_2 \end{pmatrix},$$

$$P = \begin{pmatrix} \frac{d_1+u_1}{a_x} & \frac{d_1+u_2}{a_x} \\ 1 & 1 \end{pmatrix}, \quad P^{-1} = \begin{pmatrix} -\frac{a_x}{u_2-u_1} & \frac{d_1+u_2}{u_2-u_1} \\ \frac{a_x}{u_2-u_1} & -\frac{d_1+u_1}{u_2-u_1} \end{pmatrix},$$

where  $u_1 = \frac{\beta_0 x_0^* + d_1 - \sqrt{\Delta}}{2}$ ,  $u_2 = \frac{\beta_0 x_0^* + d_1 + \sqrt{\Delta}}{2}$ ,  $\Delta = (\beta_0 x_0^* - d_1)^2 - 4\lambda a_x x_0^*$  and

$$\begin{aligned} e^{tA} &= Pe^{tV}P^{-1} \\ &= \begin{pmatrix} \frac{d_1+u_1}{a_x} & \frac{d_1+u_2}{a_x} \\ 1 & 1 \end{pmatrix} \begin{pmatrix} e^{u_1 t} & 0 \\ 0 & e^{u_2 t} \end{pmatrix} \begin{pmatrix} -\frac{a_x}{u_2-u_1} & \frac{d_1+u_2}{u_2-u_1} \\ \frac{a_x}{u_2-u_1} & -\frac{d_1+u_1}{u_2-u_1} \end{pmatrix} \\ &= \begin{pmatrix} \frac{d_1+u_1}{a_x} e^{u_1 t} & \frac{d_1+u_2}{a_x} e^{u_2 t} \\ e^{u_1 t} & e^{u_2 t} \end{pmatrix} \begin{pmatrix} -\frac{a_x}{u_2-u_1} & \frac{d_1+u_2}{u_2-u_1} \\ \frac{a_x}{u_2-u_1} & -\frac{d_1+u_1}{u_2-u_1} \end{pmatrix} \\ &= \begin{pmatrix} \frac{d_1+u_2}{u_2-u_1} e^{u_2 t} - \frac{d_1+u_1}{u_2-u_1} e^{u_1 t} & -\frac{d_1+u_2}{u_2-u_1} \frac{d_1+u_1}{a_x} (e^{u_2 t} - e^{u_1 t}) \\ \frac{d_1+u_1}{u_2-u_1} (e^{u_2 t} - e^{u_1 t}) & -\frac{d_1+u_1}{u_2-u_1} e^{u_2 t} + \frac{d_1+u_2}{u_2-u_1} e^{u_1 t} \end{pmatrix}. \end{aligned}$$

We obtain

$$\begin{cases} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_0} &= \frac{d_1+u_2}{u_2-u_1} e^{u_2 t} - \frac{d_1+u_1}{u_2-u_1} e^{u_1 t}, \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_0} &= \frac{a_x}{u_2-u_1} (e^{u_2 t} - e^{u_1 t}). \end{cases}$$

From (26) and (30) we have  $\begin{pmatrix} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} \end{pmatrix} = e^{tA} \begin{pmatrix} \frac{\partial \Phi_1(0, \zeta_0)}{\partial x_1} \\ \frac{\partial \Phi_2(0, \zeta_0)}{\partial x_1} \end{pmatrix} = e^{tA} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ . We

obtain

$$\begin{cases} \frac{\partial \Phi_1(t, \zeta_0)}{\partial x_1} &= -\frac{d_1+u_2}{u_2-u_1} \frac{d_1+u_1}{a_x} (e^{u_2 t} - e^{u_1 t}), \\ \frac{\partial \Phi_2(t, \zeta_0)}{\partial x_1} &= -\frac{d_1+u_1}{u_2-u_1} e^{u_2 t} + \frac{d_1+u_2}{u_2-u_1} e^{u_1 t}. \end{cases}$$

From (27) we have

$$\frac{\partial \Phi_1(t, \zeta_0)}{\partial y_0} = -\frac{\lambda a_y x_0^* e^{-\beta_0 x_0^* t}}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* + d_2} \left( \frac{e^{(\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* + \beta_0 x_0^*)t} - 1}{\gamma - a_y - \lambda \frac{a_x}{d_1} x_0^* + \beta_0 x_0^*} - \frac{e^{(\beta_0 x_0^* - d_2)t} - 1}{\beta_0 x_0^* - d_2} \right),$$

From (28) we have

$$\frac{\partial \Phi_1(t, \zeta_0)}{\partial y_1} = -\lambda x_0^* e^{-\beta_0 x_0^* t} \left( \frac{e^{(\beta_0 x_0^* - d_2)t} - 1}{\beta_0 x_0^* - d_2} \right).$$









with the initial condition  $\frac{\partial^2 \Phi_3(0, \zeta_0)}{\partial y_1^2} = 0$ , then

$$\frac{d}{dt} \left( \frac{\partial^2 \Phi_3(t, \zeta_0)}{\partial y_1^2} \right) = \frac{\partial F_3(\zeta(t))}{\partial y_0} \frac{\partial^2 \Phi_3(t, \zeta_0)}{\partial y_1^2} \quad [43]$$

with the initial condition  $\frac{\partial^2 \Phi_3(0, \zeta_0)}{\partial y_1^2} = 0$ . From (43) we have

$$\frac{\partial^2 \Phi_3(t, \zeta_0)}{\partial y_1^2} = 0.$$

#### 4.3. First derivatives of $Z^*$

Let  $\eta(\bar{\tau}) = \tau_0 + \bar{\tau}$ ,  $\eta_1(\bar{\tau}, \sigma) = x_0^* + q_1 \sigma + z_1^*(\bar{\tau}, \sigma)$ ,  $\eta_2(\bar{\tau}, \sigma) = \frac{a_x}{d_1} x_0^* + q_2 \sigma + z_2^*(\bar{\tau}, \sigma)$ ,  $\eta_3(\bar{\tau}, \sigma) = \sigma$  and  $\eta_4(\bar{\tau}, \sigma) = q_4 \sigma + z_4^*(\bar{\tau}, \sigma)$ .

From (23) we have

$$\begin{cases} \frac{\partial}{\partial \bar{\tau}} (\eta_1 - \Theta_1 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0, \\ \frac{\partial}{\partial \bar{\tau}} (\eta_2 - \Theta_2 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0, \\ \frac{\partial}{\partial \bar{\tau}} (\eta_4 - \Theta_4 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0. \end{cases}$$

Therefore

$$\begin{cases} \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} - \frac{\partial \Theta_1(\Phi(\tau_0, \zeta_0))}{\partial x_0} \left( \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial \bar{\tau}} + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_0} \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_1} \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} \right. \\ \left. + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} \right) = 0, \\ \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} - \frac{\partial \Theta_2(\Phi(\tau_0, \zeta_0))}{\partial x_1} \left( \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial \bar{\tau}} + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_0} \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_1} \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} \right. \\ \left. + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial y_1} \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} \right) = 0, \\ \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} - \frac{\partial \Theta_4(\Phi(\tau_0, \zeta_0))}{\partial y_1} \left( \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial \bar{\tau}} + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial x_0} \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial x_1} \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} \right. \\ \left. + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial y_1} \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} \right) = 0. \end{cases}$$

Since  $\frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial \bar{\tau}} = \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial \bar{\tau}} = \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial \bar{\tau}} = 0$ , we obtain

$$\begin{cases} a_0 \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} + b_0 \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} + d_0 \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} = 0, \\ e_0 \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} + f_0 \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} = 0, \\ i_0 \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} = 0. \end{cases}$$

That is

$$\begin{cases} \frac{\partial z_1^*(0, 0)}{\partial \bar{\tau}} = 0, \\ \frac{\partial z_2^*(0, 0)}{\partial \bar{\tau}} = 0, \\ \frac{\partial z_4^*(0, 0)}{\partial \bar{\tau}} = 0. \end{cases} \quad [44]$$

In the same way as above, we obtain

$$\begin{cases} \frac{\partial}{\partial \sigma}(\eta_1 - \Theta_1 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0, \\ \frac{\partial}{\partial \sigma}(\eta_2 - \Theta_2 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0, \\ \frac{\partial}{\partial \sigma}(\eta_4 - \Theta_4 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))(0, 0) = 0. \end{cases}$$

Therefore

$$\begin{cases} \frac{\partial z_1^*(0,0)}{\partial \sigma} - \frac{\partial \Theta_1(\Phi(\tau_0, \zeta_0))}{\partial x_0} \left( \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_0} \left( q_1 + \frac{\partial z_1^*(0,0)}{\partial \sigma} \right) + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial x_1} \left( q_2 + \frac{\partial z_2^*(0,0)}{\partial \sigma} \right) \right. \\ \quad \left. + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_0} + \frac{\partial \Phi_1(\tau_0, \zeta_0)}{\partial y_1} \left( q_1 + \frac{\partial z_4^*(0,0)}{\partial \sigma} \right) \right) = 0, \\ \frac{\partial z_2^*(0,0)}{\partial \sigma} - \frac{\partial \Theta_2(\Phi(\tau_0, \zeta_0))}{\partial x_0} \left( \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_0} \left( q_1 + \frac{\partial z_1^*(0,0)}{\partial \sigma} \right) + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial x_1} \left( q_2 + \frac{\partial z_2^*(0,0)}{\partial \sigma} \right) \right. \\ \quad \left. + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial y_0} + \frac{\partial \Phi_2(\tau_0, \zeta_0)}{\partial y_1} \left( q_1 + \frac{\partial z_4^*(0,0)}{\partial \sigma} \right) \right) = 0, \\ \frac{\partial z_4^*(0,0)}{\partial \sigma} - \frac{\partial \Theta_4(\Phi(\tau_0, \zeta_0))}{\partial x_0} \left( \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial x_0} \left( q_1 + \frac{\partial z_1^*(0,0)}{\partial \sigma} \right) + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial x_1} \left( q_2 + \frac{\partial z_2^*(0,0)}{\partial \sigma} \right) \right. \\ \quad \left. + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial y_0} + \frac{\partial \Phi_4(\tau_0, \zeta_0)}{\partial y_1} \left( q_1 + \frac{\partial z_4^*(0,0)}{\partial \sigma} \right) \right) = 0. \end{cases}$$

We obtain

$$\begin{cases} a_0 \frac{\partial z_1^*(0,0)}{\partial \sigma} + b_0 \frac{\partial z_2^*(0,0)}{\partial \sigma} + c_0 + d_0 \frac{\partial z_4^*(0,0)}{\partial \sigma} = 0, \\ e_0 \frac{\partial z_1^*(0,0)}{\partial \sigma} + f_0 \frac{\partial z_2^*(0,0)}{\partial \sigma} = 0, \\ i_0 \frac{\partial z_4^*(0,0)}{\partial \sigma} = 0. \end{cases}$$

That is

$$\begin{cases} \frac{\partial z_1^*(0,0)}{\partial \sigma} = -\frac{f_0 c_0}{a_0 f_0 - b_0 e_0}, \\ \frac{\partial z_2^*(0,0)}{\partial \sigma} = \frac{e_0 c_0}{a_0 f_0 - b_0 e_0}, \\ \frac{\partial z_4^*(0,0)}{\partial \sigma} = 0. \end{cases} \quad [45]$$

#### 4.4. First derivatives of $\omega$

We have

$$\begin{aligned} \frac{\partial \omega}{\partial \bar{\tau}} &= \frac{\partial}{\partial \bar{\tau}}(\eta_3 - \Theta_3 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4)) \\ &= -\frac{\partial \Theta_3}{\partial y_0} \left( \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0} \frac{\partial z_1^*}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1} \frac{\partial z_2^*}{\partial \bar{\tau}} \right. \\ &\quad \left. + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1} \frac{\partial z_4^*}{\partial \bar{\tau}} \right). \end{aligned}$$

At  $(\bar{\tau}, \sigma) = (0, 0)$  we have  $\frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau}} = 0$ , then we obtain

$$\frac{\partial \omega}{\partial \bar{\tau}}(0, 0) = -\frac{\partial \Theta_3}{\partial y_0} \left( \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial x_0} \frac{\partial z_1^*(0,0)}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial x_1} \frac{\partial z_2^*(0,0)}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_1} \frac{\partial z_4^*(0,0)}{\partial \bar{\tau}} \right) = 0.$$

$$\begin{aligned} \frac{\partial \omega}{\partial \sigma} &= \frac{\partial}{\partial \sigma}(\eta_3 - \Theta_3 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4)) \\ &= 1 - \frac{\partial \Theta_3}{\partial y_0} \left( \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \right. \\ &\quad \left. + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0} + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) \right). \end{aligned}$$

At  $(\bar{\tau}, \sigma) = (0, 0)$  we obtain

$$\begin{aligned}\frac{\partial \omega}{\partial \sigma}(0, 0) &= 1 - \frac{\partial \Theta_3}{\partial y_0} \left( \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial 0_0} \left( q_1 + \frac{\partial z_1^*(0, 0)}{\partial \sigma} \right) + \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial x_1} \left( q_2 + \frac{\partial z_2^*(0, 0)}{\partial \sigma} \right) \right. \\ &\quad \left. + \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_0} + \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_1} \left( q_4 + \frac{\partial z_4^*(0, 0)}{\partial \sigma} \right) \right) \\ &= 1 - \frac{\partial \Theta_3}{\partial y_0} \frac{\partial \Phi_3(\tau_0, \zeta_0)}{\partial y_0} \\ &= g_0 = 0\end{aligned}$$

Therefore  $D_{(\bar{\tau}, \sigma)}\omega(0, 0) = (0, 0)$ .

#### 4.5. Second derivatives of $\omega$

Let  $\mathcal{A} = \frac{\partial^2 \omega(0, 0)}{\partial \bar{\tau}^2}$ ,  $\mathcal{B} = \frac{\partial^2 \omega(0, 0)}{\partial \bar{\tau} \partial \sigma}$  and  $\mathcal{C} = \frac{\partial^2 \omega(0, 0)}{\partial \sigma^2}$ .

##### 4.5.1. Calculation of $\mathcal{A}$ .

We have  $\frac{\partial^2 \omega}{\partial \bar{\tau}^2} = \frac{\partial^2}{\partial \bar{\tau}^2} (\eta_3 - \Theta_3 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))$ , then

$$\begin{aligned}\frac{\partial^2 \omega}{\partial \bar{\tau}^2} &= -\frac{\partial \Theta_3}{\partial y_0} \left\{ \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau}^2} + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial x_0} \frac{\partial z_1^*}{\partial \bar{\tau}} + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial x_1} \frac{\partial z_2^*}{\partial \bar{\tau}} \right. \\ &\quad + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial y_1} \frac{\partial z_4^*}{\partial \bar{\tau}} + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0^2} \left( \frac{\partial z_1^*}{\partial \bar{\tau}} \right)^2 + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial x_1} \frac{\partial z_1^*}{\partial \bar{\tau}} \frac{\partial z_2^*}{\partial \bar{\tau}} \\ &\quad + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial y_1} \frac{\partial z_1^*}{\partial \bar{\tau}} \frac{\partial z_4^*}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0} \frac{\partial^2 z_1^*}{\partial \bar{\tau}^2} + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1^2} \left( \frac{\partial z_2^*}{\partial \bar{\tau}} \right)^2 \\ &\quad + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1 \partial y_1} \frac{\partial z_2^*}{\partial \bar{\tau}} \frac{\partial z_4^*}{\partial \bar{\tau}} + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1} \frac{\partial^2 z_2^*}{\partial \bar{\tau}^2} + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1^2} \left( \frac{\partial z_4^*}{\partial \bar{\tau}} \right)^2 \\ &\quad \left. + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1} \frac{\partial^2 z_4^*}{\partial \bar{\tau}^2} \right\}.\end{aligned}$$

At  $(\bar{\tau}, \sigma) = (0, 0)$  we have  $\frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau}^2} = 0$ . Then

$$\mathcal{A} = 0.$$

##### 4.5.2. Calculation of $\mathcal{C}$ .

We have  $\frac{\partial^2 \omega}{\partial \sigma^2} = \frac{\partial^2}{\partial \sigma^2} (\eta_3 - \Theta_3 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4))$ , then

$$\begin{aligned}\frac{\partial^2 \omega}{\partial \sigma^2} &= -\frac{\partial \Theta_3}{\partial y_0} \left( \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0^2} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right)^2 + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial x_1} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \right. \\ &\quad + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial y_1} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial y_0} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0} \left( \frac{\partial^2 z_1^*}{\partial \sigma^2} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1^2} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right)^2 \\ &\quad + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1 \partial y_1} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0 \partial x_1} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1} \left( \frac{\partial^2 z_2^*}{\partial \sigma^2} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1^2} \left( \frac{\partial z_4^*}{\partial \sigma} \right)^2 + 2 \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0 \partial y_1} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) \\ &\quad \left. + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1} \left( \frac{\partial^2 z_4^*}{\partial \sigma^2} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0^2} \right\}.\end{aligned}$$

At  $(\bar{\tau}, \sigma) = (0, 0)$  we obtain

$$\begin{aligned} \mathcal{C} &= -\frac{\partial \Theta_3}{\partial y_0} \left\{ 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial x_0 \partial y_0} \left( q_1 + \frac{\partial z_1^*(0,0)}{\partial \sigma} \right) + 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial x_1 \partial y_0} \left( q_2 + \frac{\partial z_2^*(0,0)}{\partial \sigma} \right) \right. \\ &\quad \left. + 2 \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial y_0 \partial y_1} q_4 + \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial y_0^2} \right\}. \end{aligned}$$

#### 4.5.3. Calculation of $\mathcal{B}$ .

We have  $\frac{\partial^2 \omega}{\partial \bar{\tau} \partial \sigma} = \frac{\partial}{\partial \bar{\tau}} \left( \frac{\partial}{\partial \sigma} (\eta_3 - \Theta_3 \circ \Phi(\eta, \eta_1, \eta_2, \eta_3, \eta_4)) \right)$ , then

$$\begin{aligned} \frac{\partial^2 \omega}{\partial \bar{\tau} \partial \sigma} &= -\frac{\partial \Theta_3}{\partial y_0} \left\{ \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial x_0} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial x_1} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \right. \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial \bar{\tau} \partial y_0} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0^2} \frac{\partial z_1^*}{\partial \bar{\tau}} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial x_1} \frac{\partial z_1^*}{\partial \bar{\tau}} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0 \partial x_0} \frac{\partial z_1^*}{\partial \bar{\tau}} + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1 \partial x_0} \frac{\partial z_1^*}{\partial \bar{\tau}} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0} \frac{\partial^2 z_1^*}{\partial \bar{\tau} \partial \sigma} + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_0 \partial x_1} \frac{\partial z_2^*}{\partial \bar{\tau}} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1^2} \frac{\partial z_2^*}{\partial \bar{\tau}} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1 \partial y_0} \frac{\partial z_2^*}{\partial \bar{\tau}} \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1 \partial y_1} \frac{\partial z_2^*}{\partial \bar{\tau}} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial x_1} \frac{\partial^2 z_2^*}{\partial \bar{\tau} \partial \sigma} \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1 \partial x_0} \frac{\partial z_4^*}{\partial \bar{\tau}} \left( q_1 + \frac{\partial z_1^*}{\partial \sigma} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1 \partial x_1} \frac{\partial z_4^*}{\partial \bar{\tau}} \left( q_2 + \frac{\partial z_2^*}{\partial \sigma} \right) \\ &\quad + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1^2} \frac{\partial z_1^*}{\partial \bar{\tau}} \left( q_4 + \frac{\partial z_4^*}{\partial \sigma} \right) + \frac{\partial^2 \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_0 \partial y_1} \frac{\partial z_4^*}{\partial \bar{\tau}} \\ &\quad \left. + \frac{\partial \Phi_3(\eta, \eta_1, \eta_2, \eta_3, \eta_4)}{\partial y_1} \frac{\partial^2 z_4^*}{\partial \bar{\tau} \partial \sigma} \right\}. \end{aligned}$$

At  $(\bar{\tau}, \sigma) = (0, 0)$  we have  $\frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau} \partial x_0} = \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau} \partial x_1} = \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau} \partial y_1} = 0$ , then we obtain

$$\mathcal{B} = -\frac{\partial \Theta_3}{\partial y_0} \frac{\partial^2 \Phi_3(\tau_0, \zeta_0)}{\partial \bar{\tau} \partial y_0}.$$

## 5. References

- D. D. BAINOV AND P. S. SIMEONOV, « Oscillation Theory of Impulsive Differential Equations », *International Publications, Orlando, Fla, USA*, 1998.
- D. D. BAINOV AND P. S. SIMEONOV, « Impulsive Differential Equations : Asymptotic Properties of the Solutions », *vol. 24 of Series on Advances in Mathematics for Applied Sciences, World Scientific, Singapore*, 1995.
- D. BAINOV AND V. COVACHEV, « Impulsive Differential Equations with a Small Parameter », *vol. 24 of Series on Advances in Mathematics for Applied Sciences, World Scientific, Singapore*, 1994.

- D. D. BAINOV AND P. S. SIMEONOV, « Impulsive differential equations : periodic solutions and applications », *Longman Scientific and Technical, Essex, England*, 1993.
- D. D. BAINOV, V. LAKSHMIKANTHAM AND P. S. SIMEONOV, « Theory of Impulsive Differential Equations », *vol. 6 of Series in Modern Applied Mathematics, World Scientific, Singapore*, 1989.
- D. D. BAINOV AND P.S. SIMENOV, « Systems with Impulse Effect Stability Theory and Applications », *Ellis Horwood Limited, Chichester*, 1989.
- A. BOUDERME, M. HELAL AND A. LAKMACHE, « Bifurcation of non trivial periodic solutions for pulsed chemotherapy model », *Journal of Mathematical Sciences and Applications, E-Notes*, vol. 2, n° 2, 2014, 22-44.
- A. DISHLIEV AND D. D. BAINOV, « Dependence upon initial conditions and parameters of solutions of impulsive differential equations with variable structure », *International Journal of Theoretical Physics*, vol. 29, 1990, 655-676.
- M. HELAL, M. ADIMY, A LAKMACHE AND L. P. MENJOUET, « Analysis of mathematical model of leukemia », *ITM Web of Conferences*, vol. 4, n° 01005, 2015, 1-14.
- G. IOOSS, « Bifurcation of maps and applications », *Study of mathematics, North Holland*, 1979.
- A. LAKMACHE AND O. ARINO, « Bifurcation of nontrivial periodic solutions of impulsive differential equations arising in therapeutic treatment », *Dynamics Cont. Discr. Impl. Syst.*, vol. 7, 2000, 265-287.
- A. LAKMACHE AND O. ARINO, « Nonlinear mathematical model of pulsed-therapy of heterogeneous tumor », *Nonlinear Anal. Real World Appl.*, vol. 2, 2001, 455-465.
- AH. LAKMACHE, M. HELAL AND A. LAKMACHE, « Pulsed chemotherapy treatment », *Electronic Journal of Mathematical Analysis and Applications*, vol. 2, n° 1, 2014, 127-148.
- P. KIM, P. LEE AND D. LEVY, « Dynamics and potential impact of the immune response to chronic myelogenous leukemia », *PLoS Comput. Biol.*, vol. 4, n° 6, 2008, e1000095.
- N. L. KOMAROVA AND D. WODARZ, « Drug resistance in cancer : Principles of emergence and prevention », *Proc Natl Acad Sci USA*, vol. 102, n° 27, 2005, 9714-9719.
- F. MICHOR, T.P. HUGHES, Y. IWASA, S. BRANFORD, N.P SHAH, C.L. SAWYERS AND M. A. NOWAK, « Dynamic of chronic myeloid leukemia », *Naure* vol. 435, 2005, 1267-1270.
- I. ROEDER, M. HORN, I. GLAUCHE, A. HOCHHAUS, M. C. MUELLER AND M. LOEFFLER, « Dynamic modeling of IMATINIB-treated chronic myeloid leukemia, functional insights and clinical implications », *Nat Med.*, vol. 12, n° 10, 2006, 1181-1184.