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STABILITY OF STANDING WAVES FOR A NONLINEAR KLEIN–GORDON EQUATION WITH DELTA POTENTIALS

ELEK CSOBO, FRANÇOIS GENOUD, MASAHITO OHTA, AND JULIEN ROYER

ABSTRACT. In this paper, we study local well-posedness and orbital stability of standing waves for a singularly perturbed one-dimensional nonlinear Klein–Gordon equation. We first establish local well-posedness of the Cauchy problem by a fixed point argument. Unlike the unperturbed case, a noteworthy difficulty here arises from the possible non-unitarity of the semigroup generating the corresponding linear evolution. We then show that the equation is Hamiltonian and we establish several stability/instability results for its standing waves. Our analysis relies on a detailed study of the spectral properties of the linearization of the equation, and on the well-known ‘slope condition’ for orbital stability.

1. INTRODUCTION

The purpose of this work is to initiate the study of the Cauchy problem for a singularly perturbed one-dimensional nonlinear Klein–Gordon equation, namely

$$\begin{cases} u_{tt} - u_{xx} + m^2u + \gamma\delta u + i\alpha\delta u_t - |u|^{p-1}u = 0, \\ u(t, x) \xrightarrow{|x| \rightarrow \infty} 0, \\ (u(t), \partial_t u(t))|_{t=0} = (u_0, u_1), \end{cases} \quad (1.1)$$

where $u : \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{C}$, $m > 0$, $\alpha, \gamma \in \mathbb{R}$ are parameters and $p > 1$ determines the strength of the nonlinearity. The two coefficients $\delta = \delta(x)$ are singular perturbations both given by a Dirac mass at $x = 0$, often referred to as a ‘delta potential’ in the context of one-dimensional evolution equations. Such space-dependent problems are sometimes termed ‘inhomogeneous’, as they model wave propagation in inhomogeneous media.

The condition that u vanishes at spatial infinity reflects a common physical requirement of having spatially localized waves, sometimes called ‘solitons’. We shall in fact seek solutions of (1.1) in $H^1(\mathbb{R})$.

We will show that the evolution generated by (1.1) admits a peculiar Hamiltonian formalism, with a symplectic structure depending on the coupling constant $\alpha \in \mathbb{R}$. In fact, if $\alpha \in \mathbb{C} \setminus \mathbb{R}$, the Hamiltonian of the system is not constant along the flow, as can be deduced from the symplectic structure in Section 3, or by a direct calculation using a smooth solution. Furthermore, this Hamiltonian system is phase invariant (if u is a solution, so is $e^{i\theta}u$, for any $\theta \in \mathbb{R}$), and thus possesses standing wave solutions, of the form $u(t, x) = e^{i\omega t}\varphi(x)$, with $\omega \in \mathbb{R}$ and $\varphi : \mathbb{R} \rightarrow \mathbb{R}$. The stability of standing waves in Hamiltonian systems with symmetries has attracted a lot of attention since the 1980’s. So far, this issue has been fairly well understood in homogeneous media, while in inhomogeneous media it is still a subject of intense research, both theoretically and experimentally. Inhomogeneous nonlinear dispersive equations appear in various fields of physics such as nonlinear optics, cold quantum gases (e.g. Bose–Einstein condensates), plasma physics, etc. More specifically, our interest in the present problem was initially motivated by [25], where (1.1) arises as an effective model for a superfluid Bose gas.

The nonlinear Klein–Gordon equation in homogeneous media has been extensively studied. A detailed presentation of the local and global well-posedness theory can be found in [7]. Orbital stability of standing wave solutions was first addressed in the classical papers of Shatah [29,30], and Shatah and Strauss [31]. They proved that, in N space dimensions, standing waves of any frequency are orbitally unstable if $p \geq 1 + 4/N$. If $1 < p < 1 + 4/N$, then there exists a critical frequency ω_c such that a standing wave of frequency ω is orbitally stable if $\omega_c < |\omega| < m$ and unstable if $|\omega| < \omega_c$. Strong instability by blow-up in finite time was studied by Liu, Ohta and Todorova [23], and by Ohta and Todorova [26,27]. In [19], Jeanjean and Le Coz introduced a mountain-pass approach to orbital stability for the Klein–Gordon equation, which allowed them to simplify the classical proofs and to obtain new results. In [6], Bellazzini, Ghimenti and Le Coz proved the existence of multi-solitary waves for the nonlinear Klein–Gordon equation.

The effect of a singular potential on the dynamics of the nonlinear Schrödinger equation has recently attracted substantial attention. Well-posedness of the Cauchy problem in the presence of a delta potential was studied in [1],

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while scattering was addressed in [8,14,18]. The stability of standing waves of the nonlinear Schrödinger equation with a delta potential was studied in [11–13,22,24] in various regimes. (Some authors consider δ' interactions as well, see e.g. [2,4].) Stability properties of so-called ‘black solitons’ (standing waves with $|\varphi(x)| \rightarrow 1$ as $|x| \rightarrow \infty$) were also recently addressed in [17].

The present work is a first step in the study of the nonlinear Klein–Gordon equation with delta potentials. Our main goal here is to discuss orbital stability of standing waves of (1.1). Shortly after the seminal works [29–31], a general theory of orbital stability for infinite-dimensional Hamiltonian systems with symmetries was established by Grillakis, Shatah and Strauss [15]. Their approach, based on the so-called ‘energy-momentum’ method of geometric mechanics, was recently revisited by De Bièvre, Rota Nodari and the second author [9], and by Stuart [32]. Under general assumptions on the dynamical system, conditions are given in these papers for orbital stability and instability. Of course, in order to discuss stability of standing waves, an essential preliminary step is to prove that the Hamiltonian system under consideration is locally well-posed. We shall thus start by addressing this issue, which is far from obvious in the context of (1.1).

The singular terms in (1.1) should be interpreted in the sense of distributions. Let us assume that u and u_t are continuous at $x = 0$. δu is then defined by $\langle \delta u, w \rangle = \operatorname{Re} u(0)\bar{w}(0)$, for any function w continuous at $x = 0$. And δu_t is defined similarly. Now, solutions of the equation in (1.1) will be continuous functions satisfying the corresponding unperturbed equation (with $\gamma = \alpha = 0$) pointwise, outside of $x = 0$, together with the jump condition

$$u'(0^+) - u'(0^-) = \gamma u(0) + i\alpha u_t(0). \quad (1.2)$$

Formally, this relation is indeed obtained from the equation with the delta potentials by integrating it over $x \in (-\epsilon, +\epsilon)$ and letting $\epsilon \rightarrow 0^+$. The notion of solution will be made more precise in Section 2, once the appropriate functional setting has been introduced.

Although writing the delta potentials explicitly may be useful for some formal calculations, we now introduce a functional-analytic formulation, based on the jump condition (1.2), which will make our analysis more transparent. It is convenient to reformulate the initial-value problem (1.1) as a first order system for the dependent variables $(u, v) = (u, u_t)$. We will seek solutions to (1.1) with $(u, v) \in \mathcal{H} = H^1(\mathbb{R}) \times L^2(\mathbb{R})$, which we regard as a real Hilbert space, endowed with the inner product

$$\langle (u_1, v_1), (u_2, v_2) \rangle_{\mathcal{H}} = \langle u'_1, u'_2 \rangle_{L^2} + \langle u_1, u_2 \rangle_{L^2} + \langle v_1, v_2 \rangle_{L^2},$$

where the real L^2 inner product is defined as

$$\langle u, v \rangle_{L^2} = \operatorname{Re} \int_{\mathbb{R}} u\bar{v} \, dx.$$

Here and henceforth, $'$ denotes differentiation with respect to $x \in \mathbb{R}$.

We identify $L^2(\mathbb{R}) \times L^2(\mathbb{R})$ with its dual. Then the dual \mathcal{H}^* of \mathcal{H} is $H^{-1}(\mathbb{R}) \times L^2(\mathbb{R})$, and for any $(\varphi, \psi) \in L^2(\mathbb{R}) \times L^2(\mathbb{R}) \subset H^{-1}(\mathbb{R}) \times L^2(\mathbb{R})$, the duality pairing is given by

$$\langle (\varphi, \psi), (u, v) \rangle_{\mathcal{H}^* \times \mathcal{H}} = \langle \varphi, u \rangle_{L^2} + \langle \psi, v \rangle_{L^2}, \quad (u, v) \in \mathcal{H}.$$

We shall merely write $\langle \cdot, \cdot \rangle$ for $\langle \cdot, \cdot \rangle_{\mathcal{H}^* \times \mathcal{H}}$ when no confusion is possible.

The central object in our discussion of the well-posedness of (1.1) in Section 2 is the generator \mathcal{A} of the corresponding linear evolution, defined as

$$\mathcal{A} = \begin{pmatrix} 0 & \operatorname{Id}_{L^2} \\ \partial_x^2 - m^2 & 0 \end{pmatrix}, \quad (1.3)$$

with domain

$$\mathfrak{D} = \{(u, v) \in H^{2,*} \times H^1 : u'(0^+) - u'(0^-) = \gamma u(0) + i\alpha v(0)\} \subset \mathcal{H}, \quad (1.4)$$

where

$$H^{2,*} = H^1(\mathbb{R}) \cap H^2(\mathbb{R} \setminus \{0\}).$$

Note that the effect of the delta potentials is encoded in the domain of the generator.

We will show that the operator \mathcal{A} generates a C^0 -semigroup on \mathcal{H} which, remarkably, may not be a unitary group. In contrast to the classical unperturbed case, it is in general only exponentially bounded. We recall that a skew-adjoint operator on a Hilbert space generates a unitary group (the results which we use in this paper can be found in [7], we refer for instance to [10] for a more detailed presentation of the theory of semigroups). Skew-adjointness of \mathcal{A} depends on the Hilbert structure chosen on \mathcal{H} . For γ non-negative (or negative and close to 0), it will be possible to define on \mathcal{H} an inner product whose corresponding norm is equivalent to the usual one and which makes \mathcal{A} skew-adjoint. However, for other values of γ , we can only consider on \mathcal{H} a Hilbert structure for which \mathcal{A} generates a continuous

(possibly exponentially growing) semigroup. This is in fact enough for our purpose. Notice that the parameter α will not play any role in this discussion.

Using Duhamel's formula and the Banach Fixed Point Theorem, we will then construct, for any initial data in \mathcal{H} , a unique local in time solution. We also prove the blow-up alternative and continuous dependence on the initial data for this solution.

Next, standing waves, which will be our main focus, are solutions of (1.1) of the form

$$u_\omega(t, x) = e^{i\omega t} \varphi_\omega(x),$$

where $\omega \in \mathbb{R}$, and $\varphi_\omega \in H^{2,*}$ is real-valued and satisfies the stationary equation

$$-\varphi'' + (m^2 - \omega^2)\varphi + (\gamma - \alpha\omega)\delta\varphi - |\varphi|^{p-1}\varphi = 0, \quad (1.5)$$

which will be interpreted as

$$-\varphi'' + (m^2 - \omega^2)\varphi - |\varphi|^{p-1}\varphi = 0, \quad \text{a.e. } x \in \mathbb{R}, \quad (1.6)$$

together with the jump condition

$$\varphi'(0^+) - \varphi'(0^-) = (\gamma - \alpha\omega)\varphi(0). \quad (1.7)$$

Non-trivial localized solutions to this problem exist if and only if

$$m^2 - \omega^2 > \frac{(\gamma - \alpha\omega)^2}{4}, \quad (1.8)$$

in which case they are given by the explicit formula (see Proposition 1 and Remark 1 in [22])

$$\varphi_\omega(x) = \left[\frac{(p+1)(m^2 - \omega^2)}{2} \operatorname{sech}^2 \left(\frac{(p-1)\sqrt{m^2 - \omega^2}}{2} |x| + \tanh^{-1} \left(-\frac{\gamma - \alpha\omega}{2\sqrt{m^2 - \omega^2}} \right) \right) \right]^{\frac{1}{p-1}}. \quad (1.9)$$

In particular, there are no standing wave solutions of (1.1) when $m = 0$.

Definition 1.1. For any fixed m , α and γ , we shall say that ω is *admissible* if it satisfies the relation (1.8).

In order to reveal the Hamiltonian structure of the initial-value problem (1.1), we shall follow the notation and terminology of [9]. The Hamiltonian *energy* functional associated with (1.1) is given by

$$E(u, v) = \frac{1}{2} \|u'\|_{L^2}^2 + \frac{m^2}{2} \|u\|_{L^2}^2 + \frac{1}{2} \|v\|_{L^2}^2 + \frac{\gamma}{2} |u(0)|^2 - \frac{1}{p+1} \int_{\mathbb{R}} |u|^{p+1} dx. \quad (1.10)$$

We shall prove in Section 3 that E is a constant of the motion. Another important quantity is conserved along the flow of the solution, namely the *charge*, defined as

$$Q(u, v) = \operatorname{Im} \int_{\mathbb{R}} u\bar{v} dx - \frac{\alpha}{2} |u(0)|^2. \quad (1.11)$$

We will establish in Section 3 that $E, Q \in C^2(\mathcal{H}, \mathbb{R})$. Let us now introduce the *symplector* $\mathcal{J} : \mathcal{H} \rightarrow \mathcal{H}^*$ defined by

$$\mathcal{J}(u, v) = (-i\alpha\delta u - v, u).$$

This notion, somewhat more flexible than that of a symplectic map, is introduced in [9, Sec. 6] to define Hamiltonian systems. It is noteworthy that the coupling constant α appears here in the symplectic structure itself. In this framework, the equation in (1.1) is formulated as the Hamiltonian system

$$\mathcal{J} \frac{d}{dt} U(t) = E'(U(t)), \quad (1.12)$$

where E' denotes the Fréchet derivative of E . A standing wave is now a solution of the form

$$U_\omega(t, x) = e^{i\omega t} \Phi_\omega(x), \quad (1.13)$$

where $\Phi_\omega := (\varphi_\omega, i\omega\varphi_\omega)$ satisfies the stationary equation

$$E'(\Phi_\omega) + \omega Q'(\Phi_\omega) = 0. \quad (1.14)$$

We will study the orbital stability of the standing waves (1.13), for admissible values of $\omega \in \mathbb{R}$, with respect to the symmetry group \mathbb{S}^1 acting on \mathcal{H} through

$$T(\theta)(u, v) = e^{i\theta}(u, v), \quad \theta \in \mathbb{R}. \quad (1.15)$$

This group action leaves (1.12) invariant. The corresponding notion of orbital stability is the following.

Definition 1.2. For a fixed $\omega_0 \in \mathbb{R}$, the standing wave $e^{i\omega_0 t} \Phi_{\omega_0}$ is *orbitally stable* if the following holds: for any $\epsilon > 0$ there is a $\delta > 0$ such that, if $U(t)$ is a solution of (3.3), then we have

$$\|U(0) - \Phi_{\omega_0}\|_{\mathcal{H}} < \delta \implies \inf_{\theta \in \mathbb{R}} \|U(t) - e^{i\theta} \Phi_{\omega_0}\|_{\mathcal{H}} < \epsilon \quad \text{for all } t \in \mathbb{R}. \quad (1.16)$$

Otherwise, Φ_{ω_0} is said to be *orbitally unstable*.

In addition to orbital stability, we will also prove some linear instability results. Writing a solution U of (1.12) in the form $U(t) = e^{i\omega_0 t}(\Phi_{\omega_0} + V(t))$, we have that, at first order, V satisfies the linearized equation

$$\mathcal{J} \frac{d}{dt} V(t) = L''_{\omega_0}(\Phi_{\omega_0}) V(t), \quad (1.17)$$

where L_{ω} is defined in (1.20).

Definition 1.3. The standing wave $e^{i\omega_0 t} \Phi_{\omega_0}$ is *linearly unstable* if 0 is a linearly unstable solution (in the sense of Lyapunov) of (1.17).

In Section 4, we will carry out a stability analysis based on the energy-momentum method developed in [9,15,32]. More precisely, our proofs will make use of the well-known *slope condition* (also known as the ‘Vakhitov–Kolokolov criterion’), which states that the standing wave Φ_{ω_0} is stable/unstable provided

$$\left. \frac{d}{d\omega} \Big|_{\omega=\omega_0} Q(\Phi_{\omega}) > 0 \right/ \left. \frac{d}{d\omega} \Big|_{\omega=\omega_0} Q(\Phi_{\omega}) < 0, \quad (1.18)$$

where the charge of the standing wave (1.13) is explicitly given by

$$Q(\Phi_{\omega}) = Q(\varphi_{\omega}, i\omega\varphi_{\omega}) = -\omega \|\varphi_{\omega}\|_{L^2}^2 - \frac{\alpha}{2} |\varphi_{\omega}(0)|^2. \quad (1.19)$$

The stability/instability of Φ_{ω} in fact relies on a subtle combination of the slope condition (1.18) and suitable spectral properties of the linearization of (1.12) (see e.g. [9, Sec. 10.3] for a detailed discussion in the context of the nonlinear Schrödinger equation). The spectral conditions are conveniently expressed in terms of the Lyapunov functional $L_{\omega} : \mathcal{H} \rightarrow \mathbb{R}$ associated with (1.12), defined by

$$L_{\omega}(u, v) = E(u, v) + \omega Q(u, v). \quad (1.20)$$

Let $\tilde{R} = \text{diag}(R, \text{Id}_{L^2}) : \mathcal{H} \rightarrow \mathcal{H}^*$, where $R = -\partial_x^2 + 1 : H^1(\mathbb{R}) \rightarrow H^{-1}(\mathbb{R})$ is the Riesz isomorphism. It follows from the results of Sections 3 and 4 that, for any $\omega \in \mathbb{R}$, $L_{\omega} \in C^2(\mathcal{H}, \mathbb{R})$, and that $\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega}) : \mathcal{H} \rightarrow \mathcal{H}$ is a bounded selfadjoint operator. Let us denote by $\sigma(\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega})) \subset \mathbb{R}$ its spectrum. The relevant spectral conditions for stability are then formulated as follows.

- (S1) There exists $\lambda_{\omega} \in \mathbb{R}$ such that $\sigma(\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega})) \cap (-\infty, 0) = \{-\lambda_{\omega}^2\}$ and the subspace $\ker(L''_{\omega}(\Phi_{\omega}) + \lambda_{\omega}^2 \tilde{R})$ is one-dimensional.
- (S1') $\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega})$ has two negative eigenvalues (counted with multiplicities): either there exist $\lambda_{\omega}, \mu_{\omega} \in \mathbb{R}$ such that $\sigma(\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega})) \cap (-\infty, 0) = \{-\lambda_{\omega}^2, -\mu_{\omega}^2\}$, $\lambda_{\omega}^2 \neq \mu_{\omega}^2$, and the subspaces $\ker(L''_{\omega}(\Phi_{\omega}) + \lambda_{\omega}^2 \tilde{R})$ and $\ker(L''_{\omega}(\Phi_{\omega}) + \mu_{\omega}^2 \tilde{R})$ are both one-dimensional; or there exists $\lambda_{\omega} \in \mathbb{R}$ such that $\sigma(\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega})) \cap (-\infty, 0) = \{-\lambda_{\omega}^2\}$ and the subspace $\ker(L''_{\omega}(\Phi_{\omega}) + \lambda_{\omega}^2 \tilde{R})$ has dimension 2.
- (S2) $\ker L''_{\omega}(\Phi_{\omega}) = \text{span}\{i\Phi_{\omega}\}$.
- (S3) Apart from the non-positive eigenvalues, $\sigma(\tilde{R}^{-1} L''_{\omega}(\Phi_{\omega}))$ is positive and bounded away from zero.

In the present context, the Cauchy problem (1.1) being locally well posed, the main results of [9,15,32] imply that, if the standing wave Φ_{ω_0} satisfies (S1)–(S3), then it is orbitally stable/unstable provided (1.18) holds. In case (S1) is replaced by (S1'), we will discuss linear instability of the standing waves, by means of results obtained in [16]. We shall therefore carry out a thorough spectral analysis to see when conditions (S1)–(S3) (resp. (S1')–(S3)) are satisfied, depending on the values of the parameters. By discussing the slope condition for some values of the parameters, we will then prove various stability/instability results in \mathcal{H} and in the subspace \mathcal{H}_{rad} of radial functions.

In this analysis, we shall benefit from the explicit dependence of the solution on the parameters, but the calculations required for the slope condition are rather involved. This difficulty is reflected in the intricate form of the results we present in Section 4 and explains why we decided to focus on some regimes and refrained from attempting a comprehensive analysis. Of course, numerics might come in handy to discuss the slope condition outside the scope of our analytical results. We conclude this introduction with the following table, which captures simply what ought to be

checked in order to obtain stability/instability results. The integer n_ω (resp. $n_{\omega,\text{rad}}$) denotes the number of negative eigenvalues (counted with multiplicities) of the operator $\tilde{R}^{-1}L''_\omega(\Phi_\omega)$ in \mathcal{H} (resp. \mathcal{H}_{rad}).¹

	$n_\omega = 1$	$n_\omega = 2$ and $n_{\omega,\text{rad}} = 1$
$\frac{d}{d\omega}Q(\Phi_\omega) > 0$	orbitally stable	linearly unstable
$\frac{d}{d\omega}Q(\Phi_\omega) < 0$	orbitally unstable	orbitally unstable in \mathcal{H}_{rad} , hence in \mathcal{H}

2. LOCAL WELL-POSEDNESS OF THE CAUCHY PROBLEM

In this section we discuss the local well-posedness of the Cauchy problem (1.1). In order to apply the standard theory of operator semigroups, we reformulate (1.1) as a first order system on \mathcal{H} . We consider on \mathcal{H} the operator \mathcal{A} defined by (1.3)–(1.4). Given $f : \mathbb{R} \rightarrow \mathbb{R}$ we set, for $U = (u, v) \in \mathcal{H}$,

$$F(U) = \begin{pmatrix} 0 \\ f(u) \end{pmatrix}.$$

With $f(u) = |u|^{p-1}u$ and $U_0 = (u_0, u_1)$, (1.1) can be rewritten as

$$\begin{cases} U_t(t) - \mathcal{A}U(t) = F(U(t)), \\ U(0) = U_0. \end{cases} \quad (2.1)$$

We will show that \mathcal{A} generates a strongly continuous semigroup on \mathcal{H} , which will allow us to establish the local well-posedness of (1.1).

Definition 2.1. Let $T \in (0, \infty]$.

- A strong solution to (2.1) is a function $U \in C^0([0, T], \mathfrak{D}) \cap C^1([0, T], \mathcal{H})$ such that (2.1) holds on $[0, T]$. We say that u is a strong solution of (1.1) on $[0, T]$ if (u, u_t) is a strong solution of (2.1).
- A weak solution of (2.1) is a function $U \in C^0([0, T], \mathcal{H})$ such that, for all $t \in [0, T]$, there holds

$$U(t) = e^{t\mathcal{A}}U_0 + \int_0^t e^{(t-s)\mathcal{A}}F(U(s)) \, ds. \quad (2.2)$$

We say that u is a weak solution of (1.1) on $[0, T]$ if (u, u_t) is a weak solution of (1.1).

We begin with a lemma which ensures, in particular, that \mathcal{A} is densely defined.

Lemma 2.2. \mathfrak{D} is dense in \mathcal{H} .

Proof. Let $(u, v) \in \mathcal{H}$. We can consider a sequence (v_n) in H^1 such that $v_n(0) = 0$ and $v_n \rightarrow v$ in L^2 . We then choose a sequence (u_n) in H^2 which converges to u in H^1 . For $n \in \mathbb{N}$ and $x \in \mathbb{R}$, we set

$$\zeta_n(x) = 1 + \frac{\gamma|x|}{2}e^{-nx^2}.$$

We have $\zeta'_n(0^+) = -\zeta'_n(0^-) = \frac{\gamma}{2}$, so

$$(u_n\zeta_n)'(0^+) - (u_n\zeta_n)'(0^-) = u_n(0)(\zeta'_n(0^+) - \zeta'_n(0^-)) = \gamma(u_n\zeta_n)(0).$$

This proves that $(u_n\zeta_n, v_n)$ belongs to \mathfrak{D} for all $n \in \mathbb{N}$. Moreover,

$$\|\zeta_n - 1\|_{L^\infty} \xrightarrow{n \rightarrow +\infty} 0 \quad \text{and} \quad \|\zeta'_n\|_{L^\infty} \xrightarrow{n \rightarrow +\infty} 0,$$

so $\|u_n\zeta_n - u_n\|_{H^1} \rightarrow 0$, and hence $(\zeta_n u_n, v_n)$ goes to (u, v) in \mathcal{H} . □

¹It follows from Proposition 4.9 and Remark 4.10 that $n_{\omega,\text{rad}} = 1$ whenever $n_\omega = 2$.

2.1. Linear evolution in the energy space. In this subsection we show that the operator \mathcal{A} generates a strongly continuous group on \mathcal{H} . We know that if \mathcal{A} is skew-adjoint then it generates a one parameter unitary group on \mathcal{H} . Since the notion of skew-adjointness depends on the inner product, we first discuss the choice of a suitable Hilbert structure on \mathcal{H} .

For $\mu \geq 0$ we introduce on \mathcal{H} the quadratic form defined by

$$\|(u, v)\|_{\mathcal{H}, \mu, \gamma}^2 = \|u'\|_{L^2}^2 + \mu^2 \|u\|_{L^2}^2 + \gamma |u(0)|^2 + \|v\|_{L^2}^2. \quad (2.3)$$

We denote by $\langle \cdot, \cdot \rangle_{\mathcal{H}, \mu, \gamma}$ the corresponding bilinear form. With $\mu = m$ we observe that, for $U = (u, v) \in \mathfrak{D}$,

$$\begin{aligned} \langle \mathcal{A}U, U \rangle_{\mathcal{H}, m, \gamma} &= \langle v', u' \rangle_{L^2} + m^2 \langle v, u \rangle_{L^2} + \gamma \operatorname{Re}(v(0)\bar{u}(0)) + \langle u'', v \rangle_{L^2} - m^2 \langle u, v \rangle_{L^2} \\ &= \langle v', u' \rangle_{L^2} + \gamma \operatorname{Re}(v(0)\bar{u}(0)) - \langle u', v' \rangle_{L^2} - \operatorname{Re}(\gamma u(0)\bar{v}(0) + i\alpha |v(0)|^2) \\ &= 0. \end{aligned} \quad (2.4)$$

This makes $\langle \cdot, \cdot \rangle_{\mathcal{H}, m, \gamma}$ a good candidate to be a suitable inner product on \mathcal{H} . However, for negative γ , it may happen that the corresponding quadratic form takes negative values. In this case we have to choose a larger parameter μ .

Lemma 2.3. *Let*

$$\mu_0 = \begin{cases} 0 & \text{if } \gamma \geq 0, \\ \frac{|\gamma|}{2} & \text{if } \gamma < 0. \end{cases}$$

Then for $\mu > \mu_0$ there exists $C_\mu \geq 1$ such that, for all $u \in H^1$, we have

$$C_\mu^{-1} \|u\|_{H^1}^2 \leq \|u'\|_{L^2}^2 + \mu^2 \|u\|_{L^2}^2 + \gamma |u(0)|^2 \leq C_\mu \|u\|_{H^1}^2. \quad (2.5)$$

In particular, the functional $\|\cdot\|_{\mathcal{H}, \mu, \gamma}$ is a norm on \mathcal{H} , equivalent to the usual one.

Proof. For $u \in H^1$ we have

$$|\gamma| |u(0)|^2 = 2|\gamma| \operatorname{Re} \int_{-\infty}^0 u(x)\bar{u}'(x) dx \leq 2|\gamma| \|u\|_{L^2} \|u'\|_{L^2} \leq \frac{\|u'\|_{L^2}^2}{2} + 2\gamma^2 \|u\|_{L^2}^2. \quad (2.6)$$

This gives in particular the second inequality of (2.5). By Theorem I.3.1.4 in [3], we have

$$\|u'\|_{L^2}^2 + \gamma |u(0)|^2 + \mu_0^2 \|u\|_{L^2}^2 \geq 0$$

for all $u \in H^1$. Then for $\epsilon > 0$ we have

$$\begin{aligned} \|u'\|_{L^2}^2 + \mu^2 \|u\|_{L^2}^2 + \gamma |u(0)|^2 &\geq 2\epsilon \|u'\|_{L^2}^2 + 2\epsilon\gamma |u(0)|^2 + (\mu^2 - (1 - 2\epsilon)\mu_0^2) \|u\|_{L^2}^2 \\ &\geq \epsilon \|u'\|_{L^2}^2 + (\mu^2 - (1 - 2\epsilon)\mu_0^2 - 4\epsilon\gamma^2) \|u\|_{L^2}^2. \end{aligned}$$

With $\epsilon > 0$ small enough, this gives the first inequality in (2.5), and the second statement of the proposition follows. \square

We intend to prove the following proposition.

Proposition 2.4. *The operator \mathcal{A} generates a C^0 -semigroup on \mathcal{H} . Moreover, there exist $M \geq 0$ and $\beta \geq 0$ such that, for all $t \in \mathbb{R}$, we have*

$$\|e^{t\mathcal{A}}\|_{\mathcal{L}(\mathcal{H})} \leq M e^{\beta|t|}.$$

For this we need the following lemma.

Lemma 2.5. *Let $\mu \geq 0$ be as in Lemma 2.3 and $\lambda \geq \sqrt{\mu^2 - m^2}$.*

(i) *The bounded operator*

$$-\partial_{xx} + (m^2 + \lambda^2) + (\gamma + i\lambda\alpha)\delta : H^1 \rightarrow H^{-1} \quad (2.7)$$

has a bounded inverse, which we denote by $R(\lambda)$.

(ii) *Let $\varphi \in L^2$ and $\psi \in H^1$. Then $R(\lambda)(\varphi - \delta\psi)$ belongs to $H^{2,*}$. It is the unique solution u in $H^{2,*}$ of the problem*

$$\begin{cases} -u'' + (m^2 + \lambda^2)u = \varphi, \\ u'(0^+) - u'(0^-) = (\gamma + i\lambda\alpha)u(0) + \psi(0). \end{cases} \quad (2.8)$$

Proof. By Lemma 2.3 we have, for all $u \in H^1$,

$$\langle (-\partial_{xx} + (m^2 + \lambda^2) + (\gamma + i\lambda\alpha)\delta)u, u \rangle_{H^{-1}, H^1} \geq \|u'\|_{L^2}^2 + \mu^2 \|u\|_{L^2}^2 + \gamma |u(0)|^2 \gtrsim \|u\|_{H^1}^2.$$

Similarly, for all u and v in H^1 ,

$$\langle (-\partial_{xx} + (m^2 + \lambda^2) + (\gamma + i\lambda\alpha)\delta)u, v \rangle_{H^{-1}, H^1} \lesssim \|u\|_{H^1} \|v\|_{H^1}.$$

Hence, the operator (2.7) has a bounded inverse by the Lax–Milgram Lemma.

Let us assume that (2.8) has a solution $u \in H^{2,*}$. For all $w \in H^1$ we have

$$\langle -u'', w \rangle_{L^2} + (m^2 + \lambda^2) \langle u, w \rangle_{L^2} = \langle \varphi, w \rangle_{L^2}.$$

Integrating by parts and using the jump condition (2.8), we get

$$\langle u', w' \rangle_{L^2} + (m^2 + \lambda^2) \langle u, w \rangle_{L^2} + \operatorname{Re}((\gamma + i\lambda\alpha)u(0)\bar{w}(0)) = \langle \varphi, w \rangle_{L^2} - \operatorname{Re}(\psi(0)\bar{w}(0)). \quad (2.9)$$

This proves that $u = R(\lambda)(\varphi - \delta\psi)$. Conversely, let $u = R(\lambda)(\varphi - \delta\psi) \in H^1$. Then (2.9) holds for all $w \in C_0^\infty(\mathbb{R} \setminus \{0\})$, so u belongs to $H^{2,*}$ and $-u'' + (m^2 + \lambda^2)u = \varphi$. We now write (2.9) with $w \in C_0^\infty(\mathbb{R})$ such that $w(0) = 1$, which yields the jump condition in (2.8). \square

We can now prove Proposition 2.4.

Proof of Proposition 2.4. Consider μ as given by Lemma 2.3. For $U = (u, v) \in \mathfrak{D}$, we have

$$\begin{aligned} \langle \mathcal{A}U, U \rangle_{\mathcal{H}, \mu, \gamma} &= \langle v', u' \rangle_{L^2} + \mu^2 \langle v, u \rangle_{L^2} + \gamma \operatorname{Re}(v(0)\bar{u}(0)) + \langle u'', v \rangle_{L^2} - m^2 \langle u, v \rangle_{L^2} \\ &= \langle v', u' \rangle_{L^2} + (\mu^2 - m^2) \langle v, u \rangle_{L^2} + \gamma \operatorname{Re}(v(0)\bar{u}(0)) - \langle u', v' \rangle_{L^2} - \operatorname{Re}(\gamma u(0)\bar{v}(0) - i\alpha |v(0)|^2) \\ &= (\mu^2 - m^2) \langle v, u \rangle_{L^2}, \end{aligned}$$

and so

$$|\langle \mathcal{A}U, U \rangle_{\mathcal{H}, \mu, \gamma}| = (\mu^2 - m^2) |\langle u, v \rangle_{L^2}| \leq \frac{\mu^2 - m^2}{2} (\|u\|_{L^2}^2 + \|v\|_{L^2}^2).$$

On the other hand, by Lemma 2.3,

$$\|U\|_{\mathcal{H}, \mu, \gamma}^2 \gtrsim \|u\|_{L^2}^2 + \|v\|_{L^2}^2.$$

Hence, fixing $\beta \geq 0$ large enough, we have

$$\langle (\pm\mathcal{A} - \beta)U, U \rangle_{\mathcal{H}, \mu, \gamma} \leq 0. \quad (2.10)$$

Therefore, by [7, Proposition 2.4.2], the operators $\pm\mathcal{A} - \beta$ are dissipative. In particular, for $\lambda > \beta$, we have

$$\|(\pm\mathcal{A} - \lambda)U\|_{\mathcal{H}, \mu, \gamma}^2 \geq \|(\pm\mathcal{A} - \beta)U\|_{\mathcal{H}, \mu, \gamma}^2 + (\lambda - \beta)^2 \|U\|_{\mathcal{H}, \mu, \gamma}^2, \quad (2.11)$$

so that $\pm\mathcal{A} - \lambda$ are injective with closed range. Now, let $F = (f, g) \in \mathcal{H}$. For $U = (u, v) \in \mathfrak{D}$, we have

$$(\mathcal{A} - \lambda)U = F \iff \begin{cases} v = \lambda u + f, \\ u'' - (m^2 + \lambda^2)u = g + \lambda f. \end{cases}$$

By Lemma 2.5, if β is large enough, there exists $U = (u, v) \in \mathfrak{D}$ such that the right-hand side is satisfied. It is given by $u = R(\lambda)(-g - \lambda f - i\alpha\delta f)$ and $v = \lambda u + f$. This proves that $\operatorname{Ran}(\mathcal{A} - \lambda) = \mathcal{H}$. Hence, $(\mathcal{A} - \lambda)$ has a bounded inverse and, by (2.11),

$$\|(\mathcal{A} - \lambda)^{-1}\|_{\mathcal{L}(\mathcal{H})} \leq \frac{1}{\lambda - \beta}.$$

By the Hille–Yosida Theorem, this proves that \mathcal{A} generates a C^0 -semigroup on $(\mathcal{H}, \|\cdot\|_{\mathcal{H}, \mu, \gamma})$. Furthermore, for $t \geq 0$ and $U \in \mathcal{H}$, we have

$$\|e^{t\mathcal{A}}U\|_{\mathcal{H}, \mu, \gamma} \leq e^{\beta t} \|U\|_{\mathcal{H}, \mu, \gamma}.$$

Since the norm $\|\cdot\|_{\mathcal{H}, \mu, \gamma}$ is equivalent to the usual one, there exists $M \geq 1$ such that we also have

$$\|e^{t\mathcal{A}}U\|_{\mathcal{H}} \leq M e^{\beta t} \|U\|_{\mathcal{H}}.$$

Now the same holds true with \mathcal{A} replaced by $-\mathcal{A}$, and the proof is complete. \square

Remark 2.6. If $\gamma > -m$, we can chose $\mu = m$ in Lemma 2.3 and 2.5. As in the proof of Proposition 2.4, we can show that \mathcal{A} is skew-adjoint and, by [7, Theorem 3.2.3], it now generates a one-parameter unitary group:

$$\|e^{t\mathcal{A}}U\|_{\mathcal{H}, m, \gamma} = \|U\|_{\mathcal{H}, m, \gamma}.$$

2.2. Local well-posedness of the nonlinear problem. We are now in a position to prove the local well-posedness of the Cauchy problem (1.1). We suppose that the general nonlinearity $f \in C(\mathbb{C}, \mathbb{C})$ satisfies the following:

$$f(0) = 0, \quad (2.12)$$

$$|f(u) - f(v)| \leq C_f(1 + |u|^{p-1} + |v|^{p-1})|u - v|, \quad (2.13)$$

where $p \geq 1$ and $C_f > 0$.

Lemma 2.7. *For any $R > 0$, there exists $C_R \geq 0$ such that, for $u_1, u_2 \in H^1$ with $\|u_1\|_{H^1} \leq R$ and $\|u_2\|_{H^1} \leq R$, we have*

$$\begin{aligned} \|f(u_1)\|_{L^2} &\leq C_R \|u_1\|_{H^1}, \\ \|f(u_1) - f(u_2)\|_{L^2} &\leq C_R \|u_1 - u_2\|_{H^1}. \end{aligned}$$

Proof. For $j \in \{1, 2\}$ we have $\|u_j\|_{L^\infty} \leq R$, so with C_f we get

$$\begin{aligned} \|f(u_1) - f(u_2)\|_{L^2}^2 &\leq \int_{\mathbb{R}} C_f^2 (1 + |u_1(x)|^{p-1} + |u_2(x)|^{p-1})^2 |u_1(x) - u_2(x)|^2 dx \\ &\leq C_f^2 (1 + 2R^{p-1})^2 \|u_1 - u_2\|_{L^2}^2. \end{aligned}$$

This gives the second inequality. The first one follows by taking $u_2 = 0$. \square

Corollary 2.8. *$F : \mathcal{H} \rightarrow \mathcal{H}$ is Lipschitz continuous on bounded subsets of \mathcal{H} : for any $R > 0$, there is a constant $L(R)$ such that, for $U, V \in \mathcal{H}$ with $\|U\|_{\mathcal{H}} \leq R$ and $\|V\|_{\mathcal{H}} \leq R$, we have*

$$\|F(U) - F(V)\|_{\mathcal{H}} \leq L(R) \|U - V\|_{\mathcal{H}}.$$

Lemma 2.9. *Let $T > 0$, $U_0 \in \mathcal{H}$, and $U, V \in C([0, T], \mathcal{H})$ be two solutions to (2.2). Then $U = V$.*

Proof. Let us set $R = \sup_{t \in [0, T]} \max\{\|U(t)\|_{\mathcal{H}}, \|V(t)\|_{\mathcal{H}}\}$. For $t \in [0, T]$ we have

$$\|U(t) - V(t)\|_{\mathcal{H}} \leq \int_0^t \left\| e^{(t-s)\mathcal{A}} \left(F(U(s)) - F(V(s)) \right) \right\|_{\mathcal{H}} ds \leq M e^{T\beta} L(R) \int_0^t \|U(s) - V(s)\|_{\mathcal{H}} ds.$$

By Gronwall's Lemma we conclude that $U(t) = V(t)$ for all $t \in [0, T]$. \square

In the next proposition we prove the existence of a weak solution to the Cauchy problem.

Proposition 2.10. *Take $R > 0$ and $U_0 \in \mathcal{H}$ such that $\|U_0\|_{\mathcal{H}} \leq R$. Then there exists $T_R > 0$ and a unique solution $U \in C([0, T_R], \mathcal{H})$ of problem (2.2).*

Proof. We only need to prove the existence of the solution, as uniqueness follows from Lemma 2.9. Let M and β be as in Proposition 2.4. Consider $U_0 \in \mathcal{H}$ such that $\|U_0\|_{\mathcal{H}} \leq R$. For $T > 0$ to be determined later, let

$$X := \{U \in C([0, T], \mathcal{H}) : \|U(t)\|_{\mathcal{H}} \leq 3MR \forall t \in [0, T]\}$$

and

$$d(U, V) := \max_{t \in [0, T]} \|U(t) - V(t)\|_{\mathcal{H}}.$$

Then (X, d) is a complete metric space. We now define a map $\Psi : X \rightarrow C([0, T], \mathcal{H})$ by

$$\Psi(U) : t \mapsto e^{t\mathcal{A}} U_0 + \int_0^t e^{(t-s)\mathcal{A}} F(U(s)) ds.$$

Note that, for all $s \in [0, T]$, we have $\|F(U(s))\|_{\mathcal{H}} \leq 3MRL(3MR)$ by Corollary 2.8. Thus,

$$\begin{aligned} \|\Psi(U)(t)\|_{\mathcal{H}} &\leq \|e^{t\mathcal{A}} U_0\|_{\mathcal{H}} + \int_0^t \|e^{(t-s)\mathcal{A}} F(U(s))\|_{\mathcal{H}} ds \\ &\leq M e^{T\beta} R + M e^{T\beta} \int_0^t \|F(U(s))\|_{\mathcal{H}} ds \\ &\leq M e^{T\beta} R + M e^{T\beta} 3MRL(3MR) T. \end{aligned} \quad (2.14)$$

Furthermore, for $U, V \in E$, we have

$$\begin{aligned} \|\Psi(U)(t) - \Psi(V)(t)\|_{\mathcal{H}} &\leq Me^{T\beta} \int_0^t \|F(U(s)) - F(V(s))\|_{\mathcal{H}} \, ds \\ &\leq Me^{T\beta} L(3MR) T d(U, V). \end{aligned}$$

It is now straightforward to check that, if $T = T_R > 0$ is chosen small enough, there holds

$$\begin{aligned} Me^{T\beta} R + Me^{T\beta} 3MR L(3MR) T &\leq 3MR, \\ Me^{T\beta} L(3MR) T &\leq 1/3. \end{aligned}$$

This shows that Ψ maps (X, d) to itself and is a contraction. The result now follows from the Fixed Point Theorem. \square

Theorem 2.11. *There exists a function $T : \mathcal{H} \rightarrow (0, \infty]$ with the following properties. For all $U_0 \in \mathcal{H}$, there exists a function $U \in C([0, T(U_0)), \mathcal{H})$ such that, for all $0 < T < T(U_0)$, U is the unique solution of (2.2) in $C([0, T], \mathcal{H})$. Furthermore, the blow-up alternative holds: if $T(U_0) < \infty$ then $\lim_{t \uparrow T(U_0)} \|U(t)\|_{\mathcal{H}} = \infty$.*

Proof. For all $U_0 \in \mathcal{H}$, we set

$$T(U_0) = \sup\{T > 0 : \exists U \in C([0, T], \mathcal{H}) \text{ solution to (2.2)}\}.$$

From Proposition 2.10, we know that $T(U_0) > T_{\|U_0\|} > 0$ and Lemma 2.9 allows us to extend it to a maximal solution $U \in C([0, T(U_0)], \mathcal{H})$. The blow-up alternative follows from an argument by contradiction. Suppose that $T(U_0) < \infty$ and that there exists a constant C and a sequence t_n in $[0, T(U_0))$ such that $t_n \uparrow T(U_0)$ and $\sup_{n \in \mathbb{N}} \|U(t_n)\|_{\mathcal{H}} \leq C$. Now take a time t_n such that $t_n + T_C > T(U_0)$. Using Lemma 2.9 and Proposition 2.10, we can extend the solution up to $t_n + T_C$ by considering the initial value problem (2.2) with initial value $U(t_n)$. This contradicts the definition of $T(U_0)$ and concludes the proof. \square

Theorem 2.12. *Following the notation of Theorem 2.11, we have the following properties:*

- (i) $T : \mathcal{H} \rightarrow (0, \infty]$ is lower semicontinuous: given the initial conditions $U_0, U_{0,n} \in \mathcal{H}$ such that $U_{0,n}$ converges to U_0 in \mathcal{H} , we have that

$$T(U_0) \leq \liminf_{n \rightarrow \infty} T(U_{0,n}).$$

- (ii) If $U_{0,n} \rightarrow U_0$ and if $T < T(U_0)$, then $U_n \rightarrow U$ in $C([0, T], \mathcal{H})$, where U_n and U are the solutions of (2.2) corresponding to the initial data $U_{0,n}$ and U_0 .

Proof. Let $U_0 \in \mathcal{H}$ and $U \in C([0, T(U_0)), \mathcal{H})$ be the solution of (2.2) given by Theorem 2.11. Let $0 < T < T(U_0)$. It suffices to show that, if $U_{0,n} \rightarrow U_0$ then $T(U_{0,n}) > T$ for n large enough, and $U_n \rightarrow U$ in $C([0, T], \mathcal{H})$. We set

$$R = 2 \sup_{t \in [0, T]} \|U(t)\|_{\mathcal{H}},$$

and

$$\tau_n = \sup\{t \in [0, T(U_{0,n})) : \|U_n(s)\|_{\mathcal{H}} \leq 2R \, \forall s \in [0, t]\}.$$

If n is large enough, we have $\|U_{0,n}\|_{\mathcal{H}} \leq R$. Hence, by Proposition 2.10, $0 < T_R < \tau_n$. Now, for all $0 < t \leq \min\{\tau_n, T\}$,

$$\begin{aligned} \|U(t) - U_n(t)\|_{\mathcal{H}} &\leq \|e^{tA}(U_0 - U_{0,n})\|_{\mathcal{H}} + \int_0^t \|e^{(t-s)A}(F(U(s)) - F(U_n(s)))\|_{\mathcal{H}} \, ds \\ &\leq Me^{T\beta} \|U_0 - U_{0,n}\|_{\mathcal{H}} + Me^{T\beta} L(2R) \int_0^T \|U(s) - U_n(s)\|_{\mathcal{H}} \, ds. \end{aligned}$$

Therefore, by Gronwall's lemma,

$$\|U(t) - U_n(t)\|_{\mathcal{H}} \leq Me^{T\beta} \|U_0 - U_{0,n}\|_{\mathcal{H}} e^{Me^{T\beta} L(2R)T} \quad (2.15)$$

for all $t \leq \min\{T, \tau_n\}$. In particular, if n is large enough,

$$\|U_n(t)\|_{\mathcal{H}} \leq R$$

for $t \leq \min\{T, \tau_n\}$. Hence $\tau_n \geq T$, which implies that $T(U_{0,n}) > T$. From (2.15) we also see that $U_n \rightarrow U$ in $C([0, T], \mathcal{H})$, which completes the proof. \square

Theorem 2.13. *Let $U_0 \in \mathfrak{D}$ and $T \in (0, T(U_0))$. Let $U \in C([0, T], \mathcal{H})$ be the corresponding solution of (2.2). Then $U \in C([0, T], \mathfrak{D}) \cap C^1([0, T], \mathcal{H})$.*

Proof. Let $h > 0$ and $t \in [0, T - h)$. By a change of variables it is easy to see that

$$U(t+h) - U(t) = e^{(t+h)\mathcal{A}}U_0 - e^{t\mathcal{A}}U_0 + \int_0^t e^{s\mathcal{A}}(F(U(t+h-s)) - F(U(t-s))) ds + \int_0^h e^{(t+h-s)\mathcal{A}}F(U(s)) ds.$$

Hence,

$$\begin{aligned} \|U(t+h) - U(t)\|_{\mathcal{H}} &\leq \|e^{t\mathcal{A}}(e^{h\mathcal{A}}U_0 - U_0)\|_{\mathcal{H}} + \int_0^t \|e^{s\mathcal{A}}(F(U(t+h-s)) - F(U(t-s)))\|_{\mathcal{H}} ds \\ &\quad + \int_0^h \|e^{(t+h-s)\mathcal{A}}F(U(s))\|_{\mathcal{H}} ds \\ &\leq Me^{T\beta} \|e^{h\mathcal{A}}U_0 - U_0\|_{\mathcal{H}} + Me^{T\beta}L(R) \int_0^t \|U(t+h-s) - U(t-s)\|_{\mathcal{H}} ds \\ &\quad + Me^{T\beta}h \sup_{s \in [0, T]} \|F(U(s))\|_{\mathcal{H}}. \end{aligned}$$

We know that

$$e^{h\mathcal{A}}U_0 - U_0 = \int_0^h e^{s\mathcal{A}}\mathcal{A}U_0 ds,$$

and so $\|e^{h\mathcal{A}}U_0 - U_0\|_{\mathcal{H}} \leq hMe^{\beta T} \|\mathcal{A}U_0\|_{\mathcal{H}}$. Applying Gronwall's Lemma, we get

$$\|U(t+h) - U(t)\|_{\mathcal{H}} \lesssim h$$

for all $0 \leq t < t+h \leq T$. Hence, $U : [0, T] \rightarrow \mathcal{H}$ and $F(U) : [0, T] \rightarrow \mathcal{H}$ are Lipschitz continuous. We conclude by [7, Corollary 1.4.41] and [7, Proposition 4.1.6]. \square

Remark 2.14. It is worth noting that solutions of (1.1) may blow up in finite time. To this aim, let us consider the ordinary differential equation

$$v''(t) + v(t) - v^3(t) = 0. \quad (2.16)$$

For any fixed $T > 0$, this equation has the solution

$$v(t) = \frac{1}{\tanh(T - t/\sqrt{2})}, \quad (2.17)$$

which blows up at time $\sqrt{2}T$. Now, consider (1.1) with $m = 1$, $\gamma = \alpha = 0$ and $p = 3$, and choose the constant initial data $u_0 = 1/\tanh(T)$. By finite speed of propagation (see e.g. [7]), if u_0 is smoothly truncated outside an interval of length $2\sqrt{2}T + 1$, the corresponding solution of (1.1) will blow up like (2.17) at time $\sqrt{2}T$. Again by finite speed of propagation, if the support of the truncated u_0 is chosen far away from $x = 0$, then the solution u will not 'see' the Dirac potentials over the time interval $[0, \sqrt{2}T)$, and will also blow up at time $\sqrt{2}T$, for any values of γ and α .

3. HAMILTONIAN STRUCTURE

In this section we show that (1.1) is a Hamiltonian system, and we establish the relevant conservation laws, namely that the energy and the charge defined in (1.10) and (1.11) are constants of the motion. We shall use the general framework developed in [9] to study orbital stability of standing waves of infinite-dimensional Hamiltonian systems.

We start by showing that, in the terminology of [9, Sec. 6], $(\mathcal{H}, \mathcal{D}, \mathcal{J})$ forms an appropriate symplectic Banach triple for our problem, provided the map $\mathcal{J} : \mathcal{H} \rightarrow \mathcal{H}^*$ defined by

$$\mathcal{J}(u, v) = (-i\alpha\delta u - v, u) \quad (3.1)$$

is a (weak) symplector, in the sense of Definition 6.2 (i) in [9], which we check now.

Lemma 3.1. *The map $\mathcal{J} : \mathcal{H} \rightarrow \mathcal{H}^*$ defined by (3.1) is a symplector, that is:*

- (i) \mathcal{J} is a bounded linear map;
- (ii) \mathcal{J} is one-to-one;
- (iii) \mathcal{J} is anti-symmetric, in the sense that

$$\langle \mathcal{J}(u, v), (w, z) \rangle = -\langle \mathcal{J}(w, z), (u, v) \rangle, \quad (u, v), (w, z) \in \mathcal{H}.$$

Proof. (i) Linearity is obvious and boundedness follows from the Sobolev embedding theorem through the estimate

$$\begin{aligned} |\langle \mathcal{J}(u, v), (w, z) \rangle| &= \left| \operatorname{Im} \alpha u(0) \bar{w}(0) - \operatorname{Re} \int_{\mathbb{R}} v \bar{w} + \operatorname{Re} \int_{\mathbb{R}} u \bar{z} \right| \\ &\leq |\alpha| \|u\|_{H^1} \|w\|_{H^1} + \|v\|_{L^2} \|w\|_{L^2} + \|u\|_{L^2} \|z\|_{L^2} \\ &\lesssim (\|u\|_{H^1} + \|v\|_{L^2}) (\|w\|_{H^1} + \|z\|_{L^2}). \end{aligned}$$

(ii) \mathcal{J} is one-to-one since, clearly, $\mathcal{J}(u, v) = (0, 0)$ if and only if $(u, v) = (0, 0)$.

(iii) The antisymmetry of \mathcal{J} follows by a straightforward calculation, using that $\alpha \in \mathbb{R}$. \square

We now turn our attention to the regularity of the energy and charge functionals respectively introduced in (1.10) and (1.11). In particular, in the terminology of Definition 6.5 in [9], we show that E and Q have \mathcal{J} -compatible derivatives, i.e. that $E'(u, v), Q'(u, v) \in \operatorname{rge} \mathcal{J}$ for all $(u, v) \in \mathfrak{D}$. We write $E, Q \in \operatorname{Dif}(\mathfrak{D}, \mathcal{J})$.

Lemma 3.2. *We have that $E \in C^1(\mathcal{H}, \mathbb{R}) \cap C^2(\mathfrak{D}, \mathbb{R})$ and $Q \in C^2(\mathcal{H}, \mathbb{R})$. For $(\varphi, \psi) \in \mathfrak{D}$ and $(u, v), (w, z) \in \mathcal{H}$, we have*

$$E'(\varphi, \psi) = (-\varphi'' + m^2\varphi - |\varphi|^{p-1}\varphi - i\alpha\delta\psi, \psi),$$

$$\langle E''(\varphi, \psi)(u, v), (w, z) \rangle = \operatorname{Re} \left[\gamma u(0) \bar{w}(0) + \int_{\mathbb{R}} \left(u' \bar{w}' + (m^2 u - |\varphi|^{p-1} u - (p-1)|\varphi|^{p-3} \varphi \operatorname{Re}(\varphi \bar{u})) \bar{w} \right) dx + \int_{\mathbb{R}} v \bar{z} dx \right],$$

and for $(\varphi, \psi), (u, v) \in \mathcal{H}$

$$\begin{aligned} Q'(\varphi, \psi) &= (-\alpha\delta\varphi + i\psi, -i\varphi), \\ Q''(\varphi, \psi)(u, v) &= (-\alpha\delta u + iv, -iu). \end{aligned}$$

Furthermore, $E'(\varphi, \psi) \in \operatorname{rge} \mathcal{J}$ and $Q'(\varphi, \psi) \in \operatorname{rge} \mathcal{J}$ for all $(\varphi, \psi) \in \mathcal{H}$.

Proof. The regularity stated and the expressions obtained for the Fréchet derivatives follow from routine verifications.

Let $(\varphi, \psi) \in \mathfrak{D}$. To see that $E'(\varphi, \psi) \in \operatorname{rge} \mathcal{J}$, one has to find $(w, z) \in \mathcal{H}$ such that

$$-i\alpha\delta w - z = -\varphi'' + m^2\varphi - |\varphi|^{p-1}\varphi - i\alpha\delta\psi \quad \text{and} \quad w = \psi \quad \text{in } \mathcal{H}^*.$$

This yields $(w, z) = (\psi, \varphi'' - m^2\varphi + |\varphi|^{p-1}\varphi)$, which clearly belongs to \mathcal{H} . Similarly, for $(\varphi, \psi) \in \mathcal{H}$,

$$\mathcal{J}(w, z) = Q'(\varphi, \psi) \iff (w, z) = -i(\varphi, \psi). \quad (3.2)$$

This completes the proof. \square

Lemmas 3.1 and 3.2 show that $(\mathcal{H}, \mathfrak{D}, \mathcal{J})$ is a suitable symplectic Banach triple for our problem, with associated Hamiltonian E . For initial conditions $U_0 = (u_0, u_1) \in \mathfrak{D}$, the differential equation in (1.1) can indeed be written as the Hamiltonian system (see Definition 6.6 in [9])

$$\mathcal{J} \frac{d}{dt} U(t) = E'(U(t)). \quad (3.3)$$

Remark 3.3. The well-posedness theory in Section 2 shows that the domain \mathfrak{D} is stable under the flow of (3.3), so that, by Lemma 3.2, $E'(U(t))$ indeed belongs to $\operatorname{rge} \mathcal{J}$ over the lifespan of the solution.

Proposition 3.4. *The energy E and the charge Q are constants of the motion for (3.3), i.e. for any $U_0 = (u_0, u_1) \in \mathcal{H}$, $E(U(t)) = E(U_0)$ and $Q(U(t)) = Q(U_0)$, as long as the solution exists.*

Proof. Following [9, Theorem 5, p. 191], one only needs to check that both E and Q Poisson-commute with E , i.e. that $\{E, E\}(u, v) = \{E, Q\}(u, v) = 0$ for all $(u, v) \in \mathfrak{D}$, where for any $F \in \operatorname{Dif}(\mathfrak{D}, \mathcal{J})$, the Poisson bracket $\{E, F\}$ is defined as

$$\{E, F\}(u, v) = \langle E'(u, v), \mathcal{J}^{-1} F'(u, v) \rangle, \quad (u, v) \in \mathfrak{D}. \quad (3.4)$$

That $\{E, E\}(u, v) = 0$ for all $(u, v) \in \mathfrak{D}$ is a trivial consequence of the anti-symmetry of \mathcal{J} . As for $\{E, Q\}$, using the explicit expression $\mathcal{J}^{-1}(w, z) = (z, -w - i\alpha\delta z)$ (or (3.2)), we have

$$\begin{aligned} \{E, Q\}(u, v) &= \langle (-u'' + m^2 u - |u|^{p-1} u - i\alpha\delta v, v), (-iu, -iv) \rangle \\ &= \operatorname{Re} \left[\int_{\mathbb{R}} (-u'' + m^2 u - |u|^{p-1} u - i\alpha\delta v) \overline{(-iu)} dx + \operatorname{Re} \int_{\mathbb{R}} v \overline{(-iv)} dx \right] \\ &= \operatorname{Re} \left[-i \int_{\mathbb{R}} u'' \bar{u} dx + \alpha v(0) \bar{u}(0) \right] \\ &= \operatorname{Re}[-i\bar{u}(0)(u(0^-) - u(0^+)) + \alpha \bar{u}(0)v(0)] = 0, \end{aligned}$$

which completes the proof. \square

4. STABILITY OF STANDING WAVES

Having established the well-posedness and the Hamiltonian structure of the initial-value problem (1.1), we now investigate the stability of standing waves by applying the energy-momentum method described in the introduction. The criterion for orbital stability of the standing waves (1.13) is the following.

Proposition 4.1. *Suppose the standing wave $e^{i\omega_0 t}\Phi_{\omega_0}(x)$ satisfies the spectral conditions (S1)–(S3). Then it is orbitally stable if*

$$\left. \frac{d}{d\omega} \right|_{\omega=\omega_0} Q(\Phi_\omega) > 0,$$

and orbitally unstable if

$$\left. \frac{d}{d\omega} \right|_{\omega=\omega_0} Q(\Phi_\omega) < 0.$$

Let A be a selfadjoint operator that is bounded below with positive essential spectrum. We shall henceforth denote by $n(A) \in \mathbb{N}$ the number of negative eigenvalues (counted with multiplicities) of A , and we set

$$n_\omega := n(\tilde{R}^{-1}L''_\omega(\Phi_\omega)),$$

for all admissible $\omega \in \mathbb{R}$. In Proposition 4.1, we have $n_{\omega_0} = 1$. For $n_{\omega_0} = 2$, we will exhibit regimes of linear instability using the following criterion, borrowed from [16].

Proposition 4.2. *Let*

$$p(d''(\omega)) = \begin{cases} 1 & \text{if } \left. \frac{d}{d\omega} \right|_{\omega=\omega_0} Q(\Phi_\omega) > 0, \\ 0 & \text{if } \left. \frac{d}{d\omega} \right|_{\omega=\omega_0} Q(\Phi_\omega) < 0. \end{cases}$$

Then the standing wave $e^{i\omega_0 t}\Phi_{\omega_0}$ is linearly unstable if $n_{\omega_0} - p(d''(\omega_0))$ is odd.

Corollary 4.3. *Suppose the standing wave $e^{i\omega_0 t}\Phi_{\omega_0}$ satisfies (S1')–(S3) and*

$$\left. \frac{d}{d\omega} \right|_{\omega=\omega_0} Q(\Phi_\omega) > 0.$$

Then it is linearly unstable.

4.1. Spectral analysis. Our purpose here is to give some spectral properties (in particular the number of negative eigenvalues) of the operator $\tilde{R}^{-1}L''_\omega(\Phi_\omega)$. We will consider α, γ and ω satisfying (1.8). The quantity

$$\beta = \beta(\omega) := \gamma - \alpha\omega,$$

which appears in (1.5), will play an important role below. In view of the admissibility condition (1.8), we shall consider $\beta \in (-\beta_0, \beta_0)$, where $\beta_0 := 2\sqrt{m^2 - \omega^2}$. The main results of this subsection rely on the dependence on β of the key objects entering the spectral analysis. With this in mind (and to avoid a too heavy notation) we shall relabel various quantities by β and temporarily drop the index ω . For instance — with a slight abuse of notation — we will write φ_β instead of φ_ω , that is,

$$\varphi_\beta(x) = \left[\frac{(p+1)(m^2 - \omega^2)}{2} \operatorname{sech}^2 \left(\frac{(p-1)\sqrt{m^2 - \omega^2}}{2} |x| + \tanh^{-1} \left(-\frac{\beta}{2\sqrt{m^2 - \omega^2}} \right) \right) \right]^{\frac{1}{p-1}}. \quad (4.1)$$

One should of course keep in mind the dependence on ω . It will not be relevant for our analysis here, but will come back with full force in the next subsection. For $U = (u, v) \in \mathcal{H}$, we now let

$$L''_\beta U := (-u'' + m^2 u - \varphi_\beta^{p-1} u - (p-1)\varphi_\beta^{p-1} \operatorname{Re}(u) + \beta \delta u + i\omega v, v - i\omega u) \in \mathcal{H}^*. \quad (4.2)$$

In view of Lemma 3.2, this reads $L''_\beta = L''_\omega(\Phi_\omega)$. We shall also use the convenient notation $\tilde{L}''_\beta := \tilde{R}^{-1}L''_\beta$.

Let $\beta \in (-\beta_0, \beta_0)$. We observe that $L''_\beta : \mathcal{H} \rightarrow \mathcal{H}^*$ is a bounded operator and, for $U, W \in \mathcal{H}$, we have

$$\langle \tilde{L}''_\beta U, W \rangle_{\mathcal{H}} = \langle L''_\beta U, W \rangle_{\mathcal{H}^*, \mathcal{H}} = \langle U, L''_\beta W \rangle_{\mathcal{H}, \mathcal{H}^*} = \langle U, \tilde{L}''_\beta W \rangle_{\mathcal{H}}, \quad (4.3)$$

so \tilde{L}''_β is a bounded self-adjoint operator on \mathcal{H} .

Instead of analyzing directly the spectral properties of \tilde{L}''_β , it will be more convenient to work with the operator on $L^2 \times L^2$ associated to the form L''_β . More precisely, we set

$$D_\beta := \{u \in H^{2,*} : u'(0^+) - u'(0^-) = \beta u(0)\} \quad (4.4)$$

and we consider on $L^2 \times L^2$ the operator \mathcal{L}_β defined by $\mathcal{D}(\mathcal{L}_\beta) = D_\beta \times L^2$ and, for $U = (u, v) \in \mathcal{D}(\mathcal{L}_\beta)$,

$$\mathcal{L}_\beta U = (-u'' + m^2 u - \varphi_\beta^{p-1} u - (p-1)\varphi_\beta^{p-1} \operatorname{Re}(u) + i\omega v, v - i\omega u) \in L^2 \times L^2.$$

This defines a (\mathbb{R} -linear) self-adjoint operator which shares the same relevant spectral properties as \tilde{L}''_β :

Proposition 4.4. *The operator \mathcal{L}_β is selfadjoint and bounded from below on $L^2 \times L^2$, and for $U \in D_\beta \times L^2$ we have*

$$\langle \mathcal{L}_\beta U, U \rangle_{L^2 \times L^2} = \langle L''_\beta U, U \rangle_{\mathcal{H}^*, \mathcal{H}} = \langle \tilde{L}''_\beta U, U \rangle_{\mathcal{H}}.$$

We have

$$\ker(\mathcal{L}_\beta) = \ker(\tilde{L}''_\beta). \quad (4.5)$$

Moreover

$$\inf \sigma_{\text{ess}}(\mathcal{L}_\beta) > 0 \iff \inf \sigma_{\text{ess}}(\tilde{L}''_\beta) > 0, \quad (4.6)$$

and, in this case,

$$n(\mathcal{L}_\beta) = n(\tilde{L}''_\beta). \quad (4.7)$$

Proof. • For $U, V \in \mathcal{H}$, we set $Q(U, V) = \langle L''_\beta U, V \rangle_{\mathcal{H}^*, \mathcal{H}}$. This defines on $L^2 \times L^2$ a symmetric bilinear form with domain $\mathcal{D}(Q) = \mathcal{H}$. Using a trace inequality as in (2.6), we can check that Q is bounded from below and closed. We denote by T the corresponding selfadjoint operator given by the Representation Theorem (see for instance [20, VI-Theorem 2.1] for sesquilinear complex forms, the symmetric case being analogous for real bilinear forms). In particular,

$$\mathcal{D}(T) = \{U \in \mathcal{H} : V \mapsto Q(U, V) \text{ is a continuous linear functional on } L^2 \times L^2\},$$

and

$$\langle TU, V \rangle_{L^2 \times L^2} = Q(U, V), \quad \text{for all } U \in \mathcal{D}(T), V \in \mathcal{H}. \quad (4.8)$$

Since $\langle \mathcal{L}_\beta U, V \rangle_{L^2 \times L^2} = Q(U, V)$ for all $U \in \mathcal{D}(T)$ and $V \in \mathcal{H}$, we have $\mathcal{D}(\mathcal{L}_\beta) \subset \mathcal{D}(T)$ and $\mathcal{L}_\beta = T$ on $\mathcal{D}(\mathcal{L}_\beta)$. Now let $U = (u, v) \in \mathcal{D}(T)$. Writing (4.8) with $V = (w, 0)$ for any $w \in C_0^\infty(\mathbb{R}^*)$ proves that $u \in H^{2,*}$. Then, with $w \in C_0^\infty(\mathbb{R})$ such that $w(0) = 1$ or $w(0) = i$, we obtain $u \in D_\beta$, so $U \in \mathcal{D}(\mathcal{L}_\beta)$. This means that $\mathcal{L}_\beta = T$, and the first part of the proposition is proved.

• Let $U = (u, v) \in \ker(\mathcal{L}_\beta)$. In particular we have $U \in \mathcal{H}$ and $\langle \tilde{L}''_\beta U, W \rangle_{\mathcal{H}} = \langle L''_\beta U, W \rangle_{\mathcal{H}^*, \mathcal{H}} = 0$ for all $W \in \mathcal{H}$, so $U \in \ker(\tilde{L}''_\beta)$. Conversely, if $U \in \ker(\tilde{L}''_\beta)$ then $\langle L''_\beta U, W \rangle = \langle 0, W \rangle_{L^2 \times L^2}$ for all $W \in \mathcal{H}$, so $U \in \mathcal{D}(\mathcal{L}_\beta)$ and $\mathcal{L}_\beta U = 0$. This proves (4.5).

• Now suppose that $\inf \sigma_{\text{ess}}(\mathcal{L}_\beta) > 0$. Since \mathcal{L}_β is bounded from below, it has a finite number \tilde{m} of non-positive eigenvalues (counted with multiplicities). We denote by $\tilde{\Theta}$ the subspace of $L^2 \times L^2$ generated by the corresponding eigenvectors, and by $\tilde{\Theta}^\perp$ the orthogonal complement of $\tilde{\Theta}$ in $L^2 \times L^2$. We also set $\tilde{\Theta}_1^\perp = \tilde{\Theta}^\perp \cap \mathcal{H}$. Since $\tilde{\Theta} \subset \mathcal{D}(\mathcal{L}_\beta) \subset \mathcal{H}$, $\tilde{\Theta}$ and $\tilde{\Theta}_1^\perp$ are complementary subspaces of \mathcal{H} .

There exists $\sigma_0 > 0$ such that $\sigma(\mathcal{L}_\beta) \cap (0, \sigma_0) = \emptyset$. Then, for all $U \in \mathcal{D}(\mathcal{L}_\beta) \cap \tilde{\Theta}_1^\perp$, we have

$$\langle \tilde{L}''_\beta U, U \rangle_{\mathcal{H}} = \langle L''_\beta U, U \rangle_{\mathcal{H}^*, \mathcal{H}} = \langle \mathcal{L}_\beta U, U \rangle_{L^2 \times L^2} \geq \sigma_0 \|U\|_{L^2 \times L^2}^2.$$

On the other hand, by the trace inequality there exists $C \geq 0$ such that, for all $U \in \mathcal{H}$,

$$\langle L''_\beta U, U \rangle_{\mathcal{H}^*, \mathcal{H}} \geq \frac{\|U\|_{\mathcal{H}}^2}{2} - C \|U\|_{L^2 \times L^2}.$$

Thus, for $\eta \in (0, 1)$ and $U \in \mathcal{D}(\mathcal{L}_\beta) \cap \tilde{\Theta}_1^\perp$, we have

$$\langle \tilde{L}''_\beta U, U \rangle_{\mathcal{H}} \geq \frac{\eta}{2} \|U\|_{\mathcal{H}}^2 - \eta C \|U\|_{L^2 \times L^2} + (1 - \eta)\sigma_0 \|U\|_{L^2 \times L^2}^2.$$

For $\eta > 0$ small enough, this yields

$$\langle \tilde{L}''_\beta U, U \rangle_{\mathcal{H}} \geq \frac{\eta}{2} \|U\|_{\mathcal{H}}^2, \quad \text{for all } U \in \mathcal{D}(\mathcal{L}_\beta) \cap \tilde{\Theta}_1^\perp.$$

But $\mathcal{D}(\mathcal{L}_\beta) \cap \tilde{\Theta}_1^\perp$ is dense in $\tilde{\Theta}_1^\perp$, so

$$\langle \tilde{L}''_\beta U, U \rangle_{\mathcal{H}} \geq \frac{\eta}{2} \|U\|_{\mathcal{H}}^2, \quad \text{for all } U \in \tilde{\Theta}_1^\perp.$$

Since $\tilde{\Theta}_1^\perp$ is of codimension m in \mathcal{H} , the Min-Max Principle (see, e.g., Theorem XIII.1 in [28]) implies that $\inf \sigma_{\text{ess}}(\tilde{L}_\beta'') > 0$ and that \tilde{L}_β'' has at most \tilde{m} negative eigenvalues (counted with multiplicities).

Conversely, assume that $\inf \sigma_{\text{ess}}(\tilde{L}_\beta'') > 0$. Since \tilde{L}_β'' is bounded, it has a finite number m of non-positive eigenvalues (counted with multiplicities). We denote by Θ the subspace of \mathcal{H} generated by the corresponding eigenvectors, and by Θ^\perp the orthogonal complement of Θ in \mathcal{H} . There exists $\sigma_1 > 0$ such that $\sigma(\tilde{L}_\beta'') \cap (0, \sigma_1) = \emptyset$. Then, for all $U \in \Theta^\perp$, we have

$$\langle L_\beta'' U, U \rangle_{\mathcal{H}^*, \mathcal{H}} = \langle \tilde{L}_\beta'' U, U \rangle_{\mathcal{H}} \geq \sigma_1 \|U\|_{\mathcal{H}}^2 \geq \sigma_1 \|U\|_{L^2 \times L^2}^2.$$

We recall that \mathcal{H} is the form domain of \mathcal{L}_β and that \mathcal{L}_β is associated to the form Q , so by the form version of the Min-Max Principle (see Theorem XIII.2 in [28]), we have $\inf \sigma_{\text{ess}}(\mathcal{L}_\beta) \geq \sigma_1 > 0$ and \mathcal{L}_β has at most m non-positive eigenvalues (counted with multiplicities).

We have thus proved (4.6) and that, in this case, the operators \tilde{L}_β'' and \mathcal{L}_β have the same number of non-positive eigenvalues. With (4.5), this gives (4.7). \square

Since \mathcal{L}_β is not \mathbb{C} -linear, it is usual to split functions into real and imaginary parts. Then, the operator \mathcal{L}_β acting on pairs of complex-valued functions is formally equivalent to the following operator acting on quadruplets of real-valued functions:

$$\begin{pmatrix} L_\beta^+ + \omega^2 & 0 & 0 & -\omega \\ 0 & L_\beta^- + \omega^2 & \omega & 0 \\ 0 & \omega & 1 & 0 \\ -\omega & 0 & 0 & 1 \end{pmatrix},$$

where

$$L_\beta^+ u = -u'' + (m^2 - \omega^2)u - p\varphi_\beta^{p-1}u, \quad (4.9)$$

$$L_\beta^- u = -u'' + (m^2 - \omega^2)u - \varphi_\beta^{p-1}u. \quad (4.10)$$

Here, L_β^+ and L_β^- are \mathbb{R} -linear operators acting on a space of real-valued functions. However, we are going to use some spectral arguments which are more conveniently written with complex operators.

We denote by $L_\mathbb{C}^2$ the Lebesgue space $L^2(\mathbb{R}, \mathbb{C})$ endowed with its usual complex structure. Then we define $H_\mathbb{C}^1$ and $H_\mathbb{C}^{2,*}$ accordingly. We also define $D_\mathbb{C}^\beta$ as D_β , with $H^{2,*}$ replaced by $H_\mathbb{C}^{2,*}$. Then we define the operators L_β^+ and L_β^- by $\mathcal{D}(L_\beta^+) = \mathcal{D}(L_\beta^-) = D_\mathbb{C}^\beta \times L_\mathbb{C}^2$ and, for u in $D_\mathbb{C}^\beta \times L_\mathbb{C}^2$, $L_\beta^+ u$ and $L_\beta^- u$ are defined by (4.9) and (4.10). These are in particular \mathbb{C} -linear operators.

For $\lambda \in \mathbb{R} \setminus \{1\}$ we set (see Figure 1)

$$\Lambda(\lambda) := \lambda + \frac{\lambda\omega^2}{1-\lambda}.$$

Proposition 4.5. *The operators L_β^+ and L_β^- are selfadjoint and bounded from below on $L_\mathbb{C}^2$. Moreover, for $\lambda \in \mathbb{R} \setminus \{1\}$,*

(i) $\lambda \in \sigma(\mathcal{L}_\beta)$ if and only if $\Lambda(\lambda) \in \sigma(L_\beta^+) \cup \sigma(L_\beta^-)$,

(ii) we have

$$\dim(\ker(\mathcal{L}_\beta - \lambda)) = \dim(\ker(L_\beta^+ - \Lambda(\lambda))) + \dim(\ker(L_\beta^- - \Lambda(\lambda))), \quad (4.11)$$

and in particular,

$$n(\tilde{L}_\beta'') = n(L_\beta^+) + n(L_\beta^-), \quad (4.12)$$

(iii) $\lambda \in \sigma_{\text{ess}}(\mathcal{L}_\beta)$ if and only if $\Lambda(\lambda) \in \sigma_{\text{ess}}(L_\beta^+) \cup \sigma_{\text{ess}}(L_\beta^-)$.

In (4.11) and (4.12), the left-hand sides are dimensions of real vector spaces while the right-hand sides are dimensions of complex vector spaces.

Proof. • As in the proof of Proposition 4.4, we can check that L_β^+ and L_β^- are the selfadjoint operators corresponding to the sesquilinear forms

$$q_\beta^+ : (u, w) \mapsto \langle u', w' \rangle_{L_\mathbb{C}^2} + (m^2 - \omega^2) \langle u, w \rangle_{L_\mathbb{C}^2} - p \langle \varphi_\beta^{p-1} u, w \rangle_{L_\mathbb{C}^2} + \beta u(0) \bar{w}(0) \quad (4.13)$$

and

$$q_\beta^- : (u, w) \mapsto \langle u', w' \rangle_{L_\mathbb{C}^2} + (m^2 - \omega^2) \langle u, w \rangle_{L_\mathbb{C}^2} - \langle \varphi_\beta^{p-1} u, w \rangle_{L_\mathbb{C}^2} + \beta u(0) \bar{w}(0), \quad (4.14)$$

which are closed, symmetric and bounded from below.

• Let $\lambda \in \mathbb{R}$. Let $U = (u, v)$ and $F = (f, g)$ in $L^2 \times L^2$. We write $u = u_1 + iu_2$ where u_1 and u_2 are real-valued. We use similar notation for v, f and g . Then u belongs to D_β if and only if u_1 and u_2 belong to $D_\mathbb{C}^\beta$ and in this case

$$(\mathcal{L}_\beta - \lambda)U = F \iff \begin{cases} (L_\beta^+ + \omega^2 - \lambda)u_1 - \omega v_2 = f_1, \\ (L_\beta^- + \omega^2 - \lambda)u_2 + \omega v_1 = f_2, \\ (1 - \lambda)v_1 + \omega u_2 = g_1, \\ (1 - \lambda)v_2 - \omega u_1 = g_2. \end{cases}$$

If $\lambda \neq 1$ this gives

$$(\mathcal{L}_\beta - \lambda)U = F \iff \begin{cases} (L_\beta^+ - \Lambda(\lambda))u_1 = f_1 + \frac{\omega g_2}{1 - \lambda}, \\ (L_\beta^- - \Lambda(\lambda))u_2 = f_2 - \frac{\omega g_1}{1 - \lambda}, \\ v_1 = \frac{g_1 - \omega u_2}{1 - \lambda}, \\ v_2 = \frac{g_2 + \omega u_1}{1 - \lambda}. \end{cases} \quad (4.15)$$

• We set $K^+(\lambda) = \ker(L_\beta^+ - \Lambda(\lambda))$ and denote by $K_\mathbb{R}^+(\lambda)$ the \mathbb{R} -linear subspace of real-valued functions in $K^+(\lambda)$. Given $u \in D_\mathbb{C}^\beta$, we have $u \in K^+(\lambda)$ if and only if $\operatorname{Re}(u)$ and $\operatorname{Im}(u)$ are in $K_\mathbb{R}^+(\lambda)$. A family of linearly independent vectors in $K_\mathbb{R}^+(\lambda)$ is also a family of linearly independent vectors in $K^+(\lambda)$, so $\dim_\mathbb{R}(K_\mathbb{R}^+(\lambda)) \leq \dim_\mathbb{C}(K^+(\lambda))$. In particular, if the left-hand side is infinite, then so is the right-hand side. Now assume that $\dim_\mathbb{R}(K_\mathbb{R}^+(\lambda))$ is finite (possibly 0) and consider a basis $e = (e_1, \dots, e_m)$ of $K_\mathbb{R}^+(\mathbb{R})$ (with $m \in \mathbb{N}$). Let $u \in K^+(\lambda)$. Then $\operatorname{Re}(u)$ and $\operatorname{Im}(u)$ belong to $K_\mathbb{R}^+(\lambda)$ and are \mathbb{R} -linear combinations of vectors in e , so u is a \mathbb{C} -linear combination of vectors in e . This proves that

$$\dim_\mathbb{C} K^+(\lambda) = \dim_\mathbb{R} K_\mathbb{R}^+(\lambda).$$

We similarly define $K^-(\lambda)$ and $K_\mathbb{R}^-(\lambda)$ and see that $\dim_\mathbb{C} K^-(\lambda) = \dim_\mathbb{R} K_\mathbb{R}^-(\lambda)$.

If $e_1^+, \dots, e_{m_+}^+$ are linearly independent vectors in $K_\mathbb{R}^+(\lambda)$ and $e_1^-, \dots, e_{m_-}^-$ are linearly independent vectors in $K_\mathbb{R}^-(\lambda)$ (m_\pm may be zero), then

$$\left(e_1^+, \frac{i\omega e_1^+}{1 - \lambda} \right), \dots, \left(e_{m_+}^+, \frac{i\omega e_{m_+}^+}{1 - \lambda} \right), \left(ie_1^-, -\frac{\omega e_1^-}{1 - \lambda} \right), \dots, \left(ie_{m_-}^-, -\frac{\omega e_{m_-}^-}{1 - \lambda} \right) \quad (4.16)$$

is a family of linearly independent vectors in $\ker(\mathcal{L}_\beta - \lambda)$, so

$$\dim(\ker(\mathcal{L}_\beta - \lambda)) \geq \dim(K_\mathbb{R}^+(\lambda)) + \dim(K_\mathbb{R}^-(\lambda)). \quad (4.17)$$

In particular, if the right-hand side is infinite, then so is the left-hand side. Now assume that the right-hand side is finite. If the above families span $K_\mathbb{R}^+(\lambda)$ and $K_\mathbb{R}^-(\lambda)$, then (4.16) span $\ker(\mathcal{L}_\beta - \Lambda(\lambda))$, so the inequality in (4.17) is an equality and (4.11) is proved. Since Λ is a bijection from $(-\infty, 0)$ to itself, (4.12) follows.

• Assume that $\Lambda(\lambda) \in \rho(L_\beta^+) \cap \rho(L_\beta^-)$. Let $F = (f_1 + if_2, g_1 + ig_2) \in L^2 \times L^2$. Let (u_1, u_2, v_1, v_2) be the unique solution of (4.15) and $U = (u_1 + iu_2, v_1 + iv_2)$. Then $U \in \mathcal{D}(\mathcal{L}_\beta)$ and $(\mathcal{L}_\beta - \lambda)U = F$, so $(\mathcal{L}_\beta - \lambda)$ is surjective. We already know that λ is not an eigenvalue of \mathcal{L}_β , so $(\mathcal{L}_\beta - \lambda)$ is bijective. This implies that λ is in the resolvent set of \mathcal{L}_β .

Conversely, assume that $\lambda \in \rho(\mathcal{L}_\beta)$ and let $f = f_1 + if_2 \in L_\mathbb{C}^2$. We denote by u_1 (resp. u_2) the first component of $(\mathcal{L}_\beta - \lambda)^{-1}(f_1, 0, 0, 0)$ (resp. $(\mathcal{L}_\beta - \lambda)^{-1}(f_2, 0, 0, 0)$). Then $u = u_1 + iu_2$ is such that $(L_\beta^+ - \Lambda(\lambda))u = f$, and we deduce that $\Lambda(\lambda) \in \rho(L_\beta^+)$. Similarly, $\Lambda(\lambda) \in \rho(L_\beta^-)$. This proves (i). Then (iii) follows from (i) and (ii). \square

Remark 4.6. If we denote by Λ_- and Λ_+ the restrictions of Λ to $(-\infty, 1)$ and $(1, +\infty)$, then $\Lambda_- : (-\infty, 1) \rightarrow \mathbb{R}$ and $\Lambda_+ : (1, +\infty) \rightarrow \mathbb{R}$ are increasing bijections. Therefore,

$$\sigma_{\text{ess}}(\mathcal{L}_\beta) \setminus \{1\} = \Lambda_-^{-1}(\sigma_{\text{ess}}(L_\beta^+) \cup \sigma_{\text{ess}}(L_\beta^-)) \cup \Lambda_+^{-1}(\sigma_{\text{ess}}(L_\beta^+) \cup \sigma_{\text{ess}}(L_\beta^-)).$$

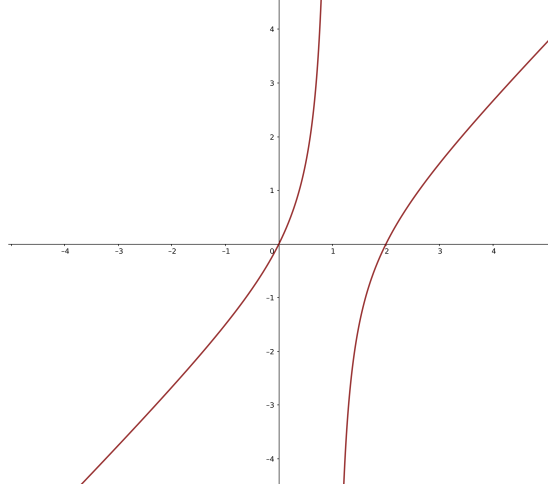
Furthermore, since $\sigma_{\text{ess}}(L_\beta^+) \cup \sigma_{\text{ess}}(L_\beta^-)$ contains a neighborhood of $+\infty$ (see Proposition 4.7 below) and $\sigma_{\text{ess}}(\mathcal{L}_\beta)$ is closed, we have that $1 \in \sigma_{\text{ess}}(\mathcal{L}_\beta)$. More precisely,

$$\sigma_{\text{ess}}(\mathcal{L}_\beta) = [\sigma_1, 1] \cup [\sigma_2, +\infty),$$

with $\sigma_1 = \Lambda_-^{-1}(m^2 - \omega^2) \in (0, 1)$ and $\sigma_2 = \Lambda_+^{-1}(m^2 - \omega^2) > 1$.

With Propositions 4.4 and 4.5, we can deduce the spectral properties of \tilde{L}_β'' from those of L_β^+ and L_β^- , which we now describe.

Proposition 4.7. *Let $\beta \in (-\beta_0, \beta_0)$.*

FIGURE 1. Graph of $\lambda \mapsto \Lambda(\lambda)$

- (i) We have $\sigma_{\text{ess}}(L_\beta^+) = \sigma_{\text{ess}}(L_\beta^-) = [m^2 - \omega^2, +\infty)$.
- (ii) The first eigenvalue of L_β^- is 0, it is simple, and the corresponding eigenspace is spanned by φ_β .
- (iii) 0 is an eigenvalue of L_β^+ if and only if $\beta = 0$. Moreover, $\ker(L_0^+) = \text{span}(\partial_x \varphi_0)$ and the negative spectrum of L_0^+ reduces to a simple eigenvalue.

Proof. It is known from [3, Theorem I-3.1.4] that the essential spectrum of $-\partial_x^2 + (m^2 - \omega^2)$ is $[m^2 - \omega^2, +\infty)$. As L_β^+ and L_β^- are relatively compact perturbations of this operator, (i) follows from the Weyl Theorem (see, e.g., [20, IV-Theorem 5.35]).

Since $\varphi_\beta \in D_\beta$, $L_\beta^- \varphi_\beta = 0$ and $\varphi_\beta > 0$, the first eigenvalue of L_β^- is 0, it is simple and the rest of the spectrum is positive (see, e.g., [5, Chapter 2]). This proves (ii).

As for (iii), we observe that φ_β satisfies

$$-\varphi_\beta'' + (m^2 - \omega^2)\varphi_\beta - \varphi_\beta^p = 0 \quad (4.18)$$

on $(-\infty, 0)$ and on $(0, +\infty)$. When $\beta = 0$, φ_0 is smooth and (4.18) holds on \mathbb{R} . After differentiation, we see that φ_0' belongs to $\ker(L_0^+)$. By Theorem 3.3 in [5, Chapter 2], 0 is a simple eigenvalue of L_0^+ and the corresponding eigenspace is spanned by φ_0' . Moreover, by [21, Lemma 4.16], L_0^+ has one simple negative eigenvalue.

Now assume that $\beta \neq 0$. Let $u \in \ker(L_\beta^+)$. In particular, u satisfies

$$-u'' + (m^2 - \omega^2)u - p\varphi_\beta^{p-1}u = 0 \quad (4.19)$$

on $(-\infty, 0)$ and on $(0, +\infty)$. Since φ_β' is also a solution of (4.19) on $(-\infty, 0)$ and $(0, +\infty)$, there exist $\mu_-, \mu_+ \in \mathbb{R}$ such that $u = \mu_- \varphi_\beta'$ on $(-\infty, 0)$ and $u = \mu_+ \varphi_\beta'$ on $(0, +\infty)$. As u is continuous at 0, we have

$$\mu_- \varphi_\beta'(0^-) = u(0) = \mu_+ \varphi_\beta'(0^+). \quad (4.20)$$

Thus $\mu_- = -\mu_+$ because $\varphi_\beta'(0^+) = -\varphi_\beta'(0^-) \neq 0$. Moreover, φ_β and u both satisfy the jump condition in (4.4), so

$$\beta \varphi_\beta(0) = \varphi_\beta'(0^+) - \varphi_\beta'(0^-) = 2\varphi_\beta'(0^+) \quad (4.21)$$

and

$$\beta u(0) = (u'(0^+) - u'(0^-)) = \mu_+ (\varphi_\beta''(0^+) + \varphi_\beta''(0^-)). \quad (4.22)$$

On the other hand, by (4.18),

$$\varphi_\beta''(0^\pm) = \lim_{x \rightarrow 0^\pm} \varphi_\beta''(x) = (m^2 - \omega^2)\varphi_\beta(0) - \varphi_\beta^p(0). \quad (4.23)$$

Using (4.20)–(4.23) we now have

$$\frac{\mu_+ \beta^2}{2} \varphi_\beta(0) = 2\mu_+ ((m^2 - \omega^2)\varphi_\beta(0) - \varphi_\beta^p(0)).$$

Since $\varphi_\beta(0) \neq 0$, it follows that

$$\mu_+ \varphi_\beta^{p-1}(0) = \mu_+ \left(m^2 - \omega^2 - \frac{\beta^2}{4} \right).$$

But a direct computation using (4.1) gives

$$\varphi_\beta^{p-1}(0) = \frac{p+1}{2} \left(m^2 - \omega^2 - \frac{\beta^2}{4} \right).$$

This proves that $\mu_+ = 0$, whence $u = 0$. Therefore, 0 is not an eigenvalue of L_β^+ when $\beta \neq 0$. The proof is complete. \square

Since $n(L_\beta^-) = 0$, (4.12) now gives $n(\tilde{L}_\beta'') = n(L_\beta^+)$. By Proposition 4.7, L_0^+ has a simple negative eigenvalue, has 0 as a simple eigenvalue, and the rest of its spectrum is positive. We now determine the number of negative eigenvalue of L_β^+ for all $\beta \in (-\beta_0, \beta_0)$ by a perturbation argument.

Proposition 4.8. *Let $\beta \in (\beta_0, \beta_0)$. Then*

$$n(L_\beta^+) = \begin{cases} 1 & \text{if } \beta \leq 0, \\ 2 & \text{if } \beta > 0. \end{cases}$$

Proof. We recall that L_β^+ is associated to the form q_β^+ defined in (4.13). This form is bounded from below and closed with dense domain $H_{\mathbb{C}}^1$. Moreover, for all $u \in H_{\mathbb{C}}^1$, the map $\beta \mapsto q_\beta^+(u, u)$ is analytic (the map $\beta \mapsto \varphi_\beta$ is pointwise analytic and locally bounded in L^∞), so the family of operators L_β^+ is analytic of type (B) in the sense of Kato (see [20, Sec. VII.4.2]). Thus, by analytic perturbation of L_0^+ , there exist $\beta_1 \in (0, \beta_0)$ and an analytic function $\lambda : (-\beta_1, \beta_1) \rightarrow \mathbb{R}$ such that $\lambda(0) = 0$ and, for all $\beta \in (-\beta_1, \beta_1)$, $\lambda(\beta)$ is a simple eigenvalue of L_β^+ , L_β^+ has a simple eigenvalue smaller than $\lambda(\beta)$, and the rest of its spectrum is positive. Furthermore, there exists an analytic function $f : (-\beta_1, \beta_1) \rightarrow L_{\mathbb{C}}^2$ such that, for $\beta \in (-\beta_1, \beta_1)$, $f(\beta)$ belongs to $D_{\mathbb{C}}^\beta$ and is an eigenfunction corresponding to the eigenvalue $\lambda(\beta)$.

In particular, there exist $\lambda_1 \in \mathbb{R}$ and $f_1 \in L_{\mathbb{C}}^2$ such that

$$\lambda(\beta) = \beta \lambda_1 + O(\beta^2), \tag{4.24}$$

$$f(\beta) = \varphi_0' + \beta f_1 + O(\beta^2). \tag{4.25}$$

From (1.9), φ_β is also analytic in $\beta \in (-\beta_0, \beta_0)$ as a function in $H^1(\mathbb{R})$, so there exists $g_1 \in H_{\mathbb{C}}^1$ such that

$$\varphi_\beta = \varphi_0 + \beta g_1 + O(\beta^2). \tag{4.26}$$

For β small, the sign of $\lambda(\beta)$ is given by the sign of λ_1 , which we now compute. We have

$$\langle L_\beta^+ f(\beta), \varphi_0' \rangle_{L_{\mathbb{C}}^2} = \langle \lambda(\beta) f(\beta), \varphi_0' \rangle_{L_{\mathbb{C}}^2} = \lambda_1 \beta \|\varphi_0'\|_{L_{\mathbb{C}}^2}^2 + O(\beta^2). \tag{4.27}$$

On the other hand, since L_β^+ is selfadjoint and $f(\beta), \varphi_0' \in D_\beta$,

$$\langle L_\beta^+ f(\beta), \varphi_0' \rangle_{L_{\mathbb{C}}^2} = \langle f(\beta), L_\beta^+ \varphi_0' \rangle_{L_{\mathbb{C}}^2}.$$

Then, by (4.26),

$$L_\beta^+ \varphi_0' = (L_\beta^+ - L_0^+) \varphi_0' = -p(\varphi_\beta^{p-1} - \varphi_0^{p-1}) \varphi_0' = -\beta p(p-1) \varphi_0^{p-2} \varphi_0' g_1 + O(\beta^2).$$

With (4.25), this yields

$$\langle L_\beta^+ f(\beta), \varphi_0' \rangle_{L_{\mathbb{C}}^2} = -\beta \langle \varphi_0', p(p-1) \varphi_0^{p-2} \varphi_0' g_1 \rangle_{L_{\mathbb{C}}^2} + O(\beta^2). \tag{4.28}$$

A straightforward calculation using that $L_\beta^- \varphi_0 = 0$ gives

$$p(p-1) \varphi_0^{p-2} (\varphi_0')^2 = L_0^+ ((m^2 - \omega^2) \varphi_0 - \varphi_0^p). \tag{4.29}$$

Now consider an arbitrary $\psi \in H^1(\mathbb{R}, \mathbb{R})$. Differentiating the identity $q_\beta^-(\psi, \varphi_\beta) = 0$ with respect to β at $\beta = 0$ yields

$$q_0^+(\psi, g_1) = -\psi(0) \varphi_0(0). \tag{4.30}$$

In view of (4.29) and (4.30), (4.28) then becomes

$$\langle L_\beta^+ f(\beta), \varphi_0' \rangle_{L_{\mathbb{C}}^2} = -\beta q_0^+ ((m^2 - \omega^2) \varphi_0 - \varphi_0^p, g_1) + O(\beta^2) = \beta [(m^2 - \omega^2) \varphi_0(0)^2 - \varphi_0(0)^{p+1}] + O(\beta^2).$$

Combining this with (4.27), we obtain

$$\lambda_1 = \frac{(m^2 - \omega^2)\varphi_0(0)^2 - \varphi_0(0)^{p+1}}{\|\varphi_0'\|_{L^2_c}^2}.$$

But from (1.9) we have

$$\varphi_0(0)^{p-1} = \frac{p+1}{2}(m^2 - \omega^2) > (m^2 - \omega^2),$$

hence $\lambda_1 < 0$. It follows that there exists $\beta_2 \in (0, \beta_1)$ such that L_β^+ has exactly one negative eigenvalue for all $\beta \in [-\beta_2, 0)$ and exactly two negative eigenvalues for all $\beta \in (0, \beta_2]$.

Finally, there exists $\kappa \in \mathbb{R}$ such that $L_\beta^+ \geq \kappa$ for all $\beta \in (-\beta_0, \beta_0)$, so the negative eigenvalues of L_β^+ are in $[\kappa, 0)$. Moreover, we know from Proposition 4.7 that 0 is not an eigenvalue of L_β^+ if $\beta \neq 0$. We define the contour Γ as the boundary of the set $\Omega = (\kappa - 1, 0) + i(-1, 1)$. Since Γ is in the resolvent set of L_β^+ for $\beta \neq 0$, we know from the analytic perturbation theory (see [20, VII-Section 1.3]) that the number of eigenvalues of L_β^+ in Ω does not depend on $\beta \in (-\beta_0, 0)$ or $\beta \in (0, \beta_0)$. Since L_{ω, β_2}^+ has two negative eigenvalues and $L_{\omega, -\beta_2}^+$ has exactly one, we have $n(L_\beta^+) = 1$ for all $\beta \in (-\beta_0, 0)$ and $n(L_\beta^+) = 2$ for all $\beta \in (0, \beta_0)$. \square

Combining Propositions 4.4, 4.5, 4.7 and 4.8 with $\beta = \gamma - \alpha\omega$, we finally obtain the following result.

Proposition 4.9. *Suppose $m^2 - \omega^2 > (\gamma - \alpha\omega)^2/4$. Then*

$$n_\omega = \begin{cases} 1 & \text{if } \gamma - \alpha\omega \leq 0, \\ 2 & \text{if } \gamma - \alpha\omega > 0. \end{cases}$$

Remark 4.10. For $\gamma - \alpha\omega > 0$, the operator \tilde{L}_β'' restricted to \mathcal{H}_{rad} has only one negative eigenvalue (see [22, Lemma 21]). Hence, $\frac{d}{d\omega}Q(\Phi_\omega) < 0$ implies orbital instability in \mathcal{H}_{rad} , and so orbital instability in \mathcal{H} .

4.2. Slope condition. We shall now turn our attention to the slope condition in order to classify various stability/instability regimes. We still consider α, γ and ω satisfying (1.8), and we now restore the dependence on ω in the notation — which was dropped in the previous subsection, where the parameter $\beta = \gamma - \alpha\omega$ played the key role.

From (1.9) and (1.11), we get

$$\begin{aligned} Q(\Phi_\omega) &= -\omega\|\varphi_\omega\|_{L^2}^2 - \frac{\alpha}{2}|\varphi_\omega(0)|^2 \\ &= -C(\omega)\frac{4\omega}{(p-1)\sqrt{m^2 - \omega^2}}\int_{\tau(\omega)}^\infty \operatorname{sech}^{\frac{4}{p-1}}(y) dy - \frac{\alpha}{2}C(\omega)\left(1 - \frac{(\gamma - \alpha\omega)^2}{4(m^2 - \omega^2)}\right)^{\frac{2}{p-1}}, \end{aligned} \quad (4.31)$$

where

$$C(\omega) = \left(\frac{(p+1)(m^2 - \omega^2)}{2}\right)^{\frac{2}{p-1}} \quad \text{and} \quad \tau(\omega) = \tanh^{-1}\left(\frac{-(\gamma - \alpha\omega)}{2\sqrt{m^2 - \omega^2}}\right).$$

We first investigate the stability of standing waves when $p = 3$, in which case (4.31) reduces to

$$\begin{aligned} Q(\Phi_\omega) &= 2(m^2 - \omega^2)\left[\frac{-2\omega}{\sqrt{m^2 - \omega^2}}\left(1 + \frac{\gamma - \alpha\omega}{2\sqrt{m^2 - \omega^2}}\right) - \frac{\alpha}{2}\left(1 - \frac{(\gamma - \alpha\omega)^2}{4(m^2 - \omega^2)}\right)\right] \\ &= -4\omega\sqrt{m^2 - \omega^2} - 2\omega(\gamma - \alpha\omega) - \alpha(m^2 - \omega^2) + \frac{\alpha}{4}(\gamma - \alpha\omega)^2. \end{aligned}$$

We shall inspect the derivative of $Q(\Phi_\omega)$ with respect to ω , which is given by

$$\frac{d}{d\omega}Q(\Phi_\omega) = \frac{4\omega^2}{\sqrt{m^2 - \omega^2}} - 4\sqrt{m^2 - \omega^2} + \left(\frac{\alpha^3}{2} + 6\alpha\right)\omega - \gamma\left(2 + \frac{\alpha^2}{2}\right). \quad (4.32)$$

In the following theorem we address the case when either $\alpha = 0$ or $\gamma = 0$. Let us first remark that in these cases there exists an H^1 solution of (1.5). Indeed, if

$$\alpha = 0, \quad |\gamma| < 2m \quad \text{and} \quad \omega \in (-\omega_\gamma, \omega_\gamma), \quad \text{with} \quad \omega_\gamma = \sqrt{m^2 - \frac{\gamma^2}{4}},$$

or if

$$\gamma = 0 \quad \text{and} \quad \omega \in (-\omega_\alpha, \omega_\alpha), \quad \text{with} \quad \omega_\alpha = \frac{2m}{\sqrt{4 + \alpha^2}},$$

then the admissibility relation (1.8) is satisfied.

Theorem 4.11. *Let $p = 3$ and $m > 0$.*

(i) *Suppose that $\alpha = 0$, $|\gamma| < 2m$ and $|\omega| \leq \omega_\gamma$. Let*

$$\tilde{\omega}_\gamma = \sqrt{\frac{16m^2 - \gamma^2 + \gamma\sqrt{\gamma^2 + 32m^2}}{32}} = m\sqrt{\frac{1}{2} + \frac{\gamma}{\sqrt{\gamma^2 + 32m^2} + \gamma}}.$$

- *For $\gamma < 0$, $e^{i\omega t}\Phi_\omega$ is orbitally stable if $|\omega| > \tilde{\omega}_\gamma$ and orbitally unstable if $|\omega| < \tilde{\omega}_\gamma$.*
- *For $\gamma > 0$, $e^{i\omega t}\Phi_\omega$ is linearly unstable if $|\omega| > \tilde{\omega}_\gamma$ and orbitally unstable on \mathcal{H}_{rad} if $|\omega| < \tilde{\omega}_\gamma$.*

(ii) *Suppose that $\gamma = 0$ and $|\alpha| < 2\sqrt{\sqrt{5} - 2}$. We set*

$$\omega_\alpha^\pm = \frac{m}{\sqrt{2}}\sqrt{1 \mp \frac{|\kappa|}{\sqrt{4 + \kappa^2}}}, \quad \text{where } \kappa = \frac{1}{4}\left(\frac{\alpha^3}{2} + 6\alpha\right).$$

Suppose $\alpha < 0$.

- *If $\omega \in (-\omega_\alpha, -\omega_\alpha^-)$ then $e^{i\omega t}\Phi_\omega$ is orbitally stable.*
- *If $\omega \in (-\omega_\alpha^-, 0)$ then $e^{i\omega t}\Phi_\omega$ is orbitally unstable.*
- *If $\omega \in (0, \omega_\alpha^+)$ then $e^{i\omega t}\Phi_\omega$ is orbitally unstable on \mathcal{H}_{rad} .*
- *If $\omega \in (\omega_\alpha^+, \omega_\alpha)$ then $e^{i\omega t}\Phi_\omega$ is linearly unstable.*

Suppose $\alpha > 0$.

- *If $\omega \in (-\omega_\alpha, -\omega_\alpha^+)$ then $e^{i\omega t}\Phi_\omega$ is linearly unstable.*
- *If $\omega \in (-\omega_\alpha^+, 0)$ then $e^{i\omega t}\Phi_\omega$ is orbitally unstable on \mathcal{H}_{rad} .*
- *If $\omega \in (0, \omega_\alpha^-)$ then $e^{i\omega t}\Phi_\omega$ is orbitally unstable.*
- *If $\omega \in (\omega_\alpha^-, \omega_\alpha)$ then $e^{i\omega t}\Phi_\omega$ is orbitally stable.*

Let us remark that orbitally instability on \mathcal{H}_{rad} implies orbitally instability on \mathcal{H} .

Proof. (i) Since $\alpha = 0$, we have $\frac{d}{d\omega}Q(\Phi_\omega) = 0$ if and only if

$$4\omega^2 - 2m^2 = \gamma\sqrt{m^2 - \omega^2},$$

that is,

$$16\omega^4 + (\gamma - 16m^2)\omega^2 + (4m^4 - \gamma^2m^2) = 0 \quad \text{and} \quad \text{sgn}(2\omega^2 - m^2) = \text{sgn}(\gamma).$$

The only possibility is $\omega^2 = \tilde{\omega}_\gamma^2$. Since $\frac{d}{d\omega}Q(\Phi_\omega)$ is negative when $\omega = 0$ and goes to $+\infty$ when ω goes to $\pm\mu$, we deduce that $\frac{d}{d\omega}Q(\Phi_\omega) < 0$ if and only if $|\omega| \leq \tilde{\omega}_\gamma$. Furthermore, Proposition 4.9 with $\alpha = 0$ yields $n_\omega = 1$ if $\gamma < 0$ and $n_\omega = 2$ if $\gamma > 0$. Hence, the conclusions in (i) follow from Proposition 4.1, Proposition 4.2 and Remark 4.10.

(ii) We now consider the case $\gamma = 0$. We have that $\frac{d}{d\omega}Q(\Phi_\omega) = 0$ if and only if

$$\omega^4 - m^2\omega^2 + \frac{m^4}{4 + \kappa^2} = 0 \quad \text{and} \quad \text{sgn}(2\omega^2 - m^2) = \text{sgn}(-\kappa\omega).$$

The solutions are $-\omega_\alpha^-$ and ω_α^+ . Then, for $|\omega| < \omega_\alpha$, we have $\frac{d}{d\omega}Q(\Phi_\omega) < 0$ if and only if $\omega \in (-\omega_\alpha^-, \omega_\alpha^+)$. Furthermore, $n_\omega = 1$ if $\alpha\omega \geq 0$ and $n_\omega = 2$ if $\alpha\omega < 0$. Hence, the conclusions again follow from Proposition 4.1, Proposition 4.2 and Remark 4.10. \square

Remark 4.12. Notice that $\tilde{\omega}_\gamma > \omega_\gamma$ when $\gamma > 2m/\sqrt{3}$. In this case, $e^{i\omega t}\Phi_\omega$ is orbitally unstable for all $\omega \in (-\omega_\gamma, \omega_\gamma)$.

Similarly, if $|\alpha| \geq 2\sqrt{\sqrt{5} - 2}$ then $\omega_\alpha^\pm \geq \omega_\alpha$, so the set of ω for which we have linear instability is empty.

We next give some results with non-zero coupling constants, $\gamma \neq 0$ and $\alpha \neq 0$. We first observe that the right-hand side of (4.32) vanishes for

$$\gamma = \tilde{\gamma}(\alpha, \omega) := \frac{2}{4 + \alpha^2} \left[\frac{4\omega^2}{\sqrt{m^2 - \omega^2}} - 4\sqrt{m^2 - \omega^2} + \left(\frac{\alpha^3}{2} + 6\alpha\right)\omega \right].$$

It follows that

$$\text{sgn} \frac{d}{d\omega}Q(\Phi_\omega) = -\text{sgn}(\gamma - \tilde{\gamma}).$$

The following theorem is then proved using Proposition 4.1, Proposition 4.2 and Remark 4.10, similarly to the proof of Theorem 4.11.

Theorem 4.13. *Let $p = 3$ and consider $\omega \in (-m, m)$, $\alpha \in \mathbb{R}$ and $\gamma \in \mathbb{R}$ satisfying (1.8).*

- (i) *Suppose $\gamma - \alpha\omega < 0$. Then $e^{i\omega t}\Phi_\omega$ is orbitally stable if $\gamma < \tilde{\gamma}$ and orbitally unstable if $\gamma > \tilde{\gamma}$.*
- (ii) *Suppose $\gamma - \alpha\omega > 0$. Then $e^{i\omega t}\Phi_\omega$ is linearly unstable if $\gamma < \tilde{\gamma}$ and orbitally unstable in \mathcal{H}_{rad} if $\gamma > \tilde{\gamma}$.*

Remark 4.14. For any fixed $\alpha \in \mathbb{R}$, there always exist values of the parameters ω, γ satisfying the above conditions for stability/instability. For instance, conditions (1.8), $\gamma < \alpha\omega$ and $\gamma < \tilde{\gamma}$ are all satisfied provided

$$0 < \tilde{\gamma} - \alpha\omega + 2\sqrt{m^2 - \omega^2} = \frac{8}{4 + \alpha^2} \frac{2\omega^2 - m^2}{\sqrt{m^2 - \omega^2}} + 2\sqrt{m^2 - \omega^2} + \frac{8\alpha\omega}{4 + \alpha^2}.$$

Clearly, this condition is satisfied for $|\omega| \approx m$. The other cases follow by similar arguments.

We now consider more general values of the power $1 < p < 5$. To keep the exposition simple enough, we focus on the cases where the coupling constants α, γ have the same sign. Of course, mixed cases could also be considered.

Lemma 4.15. *Let $1 < p < 5$.*

- (i) *Suppose $\alpha, \gamma > 0$. We have*

$$\frac{d}{d\omega}Q(\Phi_\omega) < 0 \quad \text{for } \omega \in \left(-\frac{m}{2}\sqrt{p-1}, 0\right)$$

and

$$\frac{d}{d\omega}Q(\Phi_\omega) > 0 \quad \text{for } \omega \in \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha m^2}{\gamma}\right) \cap \left(\frac{m}{2}\sqrt{p-1}, m\right),$$

whenever these intervals are not empty.

- (ii) *Suppose $\alpha, \gamma < 0$. We have*

$$\frac{d}{d\omega}Q(\Phi_\omega) < 0 \quad \text{for } \omega \in \left(0, \frac{m}{2}\sqrt{p-1}\right) \cap \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha m^2}{\gamma}\right)$$

and

$$\frac{d}{d\omega}Q(\Phi_\omega) > 0 \quad \text{for } \omega \in \left(\frac{\alpha m^2}{\gamma}, \frac{\alpha\gamma}{4 + \alpha^2}\right) \cap \left(\frac{m}{2}\sqrt{p-1}, m\right),$$

whenever these intervals are not empty.

Proof. We only prove (i), as (ii) is proved by similar calculations. We rewrite (4.31) as

$$Q(\Phi_\omega) = C_1(\omega)I(\omega) + C_2(\omega),$$

where

$$C_1(\omega) = -\frac{4}{p-1} \left(\frac{p+1}{2}\right)^{\frac{2}{p-1}} \omega(m^2 - \omega^2)^{\frac{2}{p-1} - \frac{1}{2}},$$

$$I(\omega) = \int_{\tau(\omega)}^{\infty} \operatorname{sech}^{\frac{4}{p-1}}(y) dy,$$

$$C_2(\omega) = -\frac{\alpha}{2} \left(\frac{p+1}{8}\right)^{\frac{2}{p-1}} (4(m^2 - \omega^2) - (\gamma - \alpha\omega)^2)^{\frac{2}{p-1}}.$$

We first find

$$\frac{\partial C_2}{\partial \omega} = -\frac{\alpha}{p-1} \left(\frac{p+1}{8}\right)^{\frac{2}{p-1}} (4(m^2 - \omega^2) - (\gamma - \alpha\omega)^2)^{\frac{2}{p-1} - 1} (-8\omega + 2\alpha(\gamma - \alpha\omega)).$$

It follows that $\frac{\partial C_2}{\partial \omega} > 0$ if

$$\omega > \frac{\alpha\gamma}{4 + \alpha^2}. \tag{4.33}$$

We next determine the sign of $\frac{\partial}{\partial \omega}(C_1 I)$ assuming $\omega > 0$. Since I is positive and C_1 is negative for $\omega > 0$, we will have

$$\frac{\partial}{\partial \omega}(C_1 I) = \frac{\partial C_1}{\partial \omega} I + C_1 \frac{\partial I}{\partial \omega} > 0,$$

provided $\frac{\partial C_1}{\partial \omega} > 0$ and $\frac{\partial I}{\partial \omega} < 0$. On the one hand, we have

$$\frac{\partial C_1}{\partial \omega} = -\frac{4}{p-1} \left(\frac{p+1}{2}\right)^{\frac{2}{p-1}} (m^2 - \omega^2)^{\frac{2}{p-1} - \frac{1}{2}} + \frac{8}{p-1} \left(\frac{p+1}{2}\right)^{\frac{2}{p-1}} \left(\frac{2}{p-1} - \frac{1}{2}\right) \omega^2 (m^2 - \omega^2)^{\frac{2}{p-1} - \frac{3}{2}},$$

which is positive if

$$|\omega| > \frac{m}{2} \sqrt{p-1}. \quad (4.34)$$

On the other,

$$\frac{\partial I}{\partial \omega} = -\operatorname{sech}^{\frac{4}{p-1}}(\tau) \frac{\partial}{\partial \omega} \tau = -\frac{1}{\sqrt{m^2 - \omega^2}} \left(1 - \frac{(\gamma - \alpha\omega)^2}{4(m^2 - \omega^2)} \right)^{\frac{2}{p-1}} \frac{2\alpha m^2 - 2\gamma\omega}{4(m^2 - \omega^2) - (\gamma - \alpha\omega)^2},$$

which is negative if

$$\omega < \frac{\alpha m^2}{\gamma}. \quad (4.35)$$

Hence, it follows from (4.33)–(4.35) that $\frac{d}{d\omega}Q(\Phi_\omega) > 0$ if

$$\omega \in \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha m^2}{\gamma} \right) \cap \left(\frac{m}{2} \sqrt{p-1}, m \right).$$

We now show that $\frac{d}{d\omega}Q(\Phi_\omega) < 0$ for $\omega \in (-\frac{m}{2}\sqrt{p-1}, 0)$. As C_1 and I are both positive for $\omega \in (-m, 0)$, we will have

$$\frac{\partial}{\partial \omega}(C_1 I) = \frac{\partial C_1}{\partial \omega} I + C_1 \frac{\partial I}{\partial \omega} < 0,$$

provided $\frac{\partial C_1}{\partial \omega} < 0$ and $\frac{\partial I}{\partial \omega} < 0$. From the previous calculations, we know that $\frac{\partial C_1(\omega)}{\partial \omega} < 0$ if

$$|\omega| < \frac{m}{2} \sqrt{p-1} \quad (4.36)$$

and $\frac{\partial I}{\partial \omega} < 0$ if

$$\omega < \frac{\alpha m^2}{\gamma}. \quad (4.37)$$

Finally, $\frac{\partial C_2}{\partial \omega} < 0$ if

$$\omega < \frac{\alpha\gamma}{4 + \alpha^2}. \quad (4.38)$$

Since $\alpha > 0$ and $\gamma > 0$, we conclude from (4.36)–(4.38) that $\frac{d}{d\omega}Q(\Phi_\omega) < 0$ for all $\omega \in (-\frac{m}{2}\sqrt{p-1}, 0)$. \square

We finally combine Lemma 4.15 with the spectral conditions in Proposition 4.9 to get the following result.

Theorem 4.16. *Let $1 < p < 5$ and $m = 1$.*

(i) *Let $\alpha, \gamma > 0$. If $\frac{\alpha\gamma}{4 + \alpha^2} < 1$ and $1 < \frac{\gamma^2}{4 + \alpha^2} + \frac{2\gamma}{\alpha(4 + \alpha^2)} \sqrt{4 + \alpha^2 - \gamma^2}$, then there exists ω satisfying (1.8) and*

$$\omega \in \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha}{\gamma} \right) \cap \left(\frac{1}{2} \sqrt{p-1}, 1 \right).$$

For such ω , the standing wave $e^{i\omega t} \Phi_\omega$ is orbitally stable.

If $\alpha\gamma < 2\sqrt{4 + \alpha^2 - \gamma^2}$, then there exists ω satisfying (1.8) and

$$\omega \in \left(-\frac{1}{2} \sqrt{p-1}, 0 \right).$$

For such ω , $e^{i\omega t} \Phi_\omega$ is orbitally unstable.

(ii) *Let $\alpha, \gamma < 0$. If $1 < \frac{\gamma}{\alpha}$, then there exists $\omega \in \mathbb{R}$ satisfying (1.8) and*

$$\omega \in \left(\frac{\alpha}{\gamma}, \frac{\alpha\gamma}{4 + \alpha^2} \right) \cap \left(\frac{1}{2} \sqrt{p-1}, 1 \right).$$

For such ω , $e^{i\omega t} \Phi_\omega$ is orbitally stable.

If $\frac{\gamma}{\alpha} < \frac{1}{2} \sqrt{p-1}$, then there exists $\omega \in \mathbb{R}$ satisfying (1.8) and

$$\omega \in \left(0, \frac{1}{2} \sqrt{p-1} \right) \cap \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha}{\gamma} \right).$$

For such ω , $e^{i\omega t} \Phi_\omega$ is orbitally unstable.

Proof. We only prove (i), the proof of (ii) being similar. The hypotheses $m = 1$ and $\gamma < \alpha$ imply that there exists $\omega \in \mathbb{R}$ satisfying

$$\omega \in \left(\frac{\alpha\gamma}{4 + \alpha^2}, \frac{\alpha}{\gamma} \right) \cap \left(\frac{1}{2}\sqrt{p-1}, 1 \right).$$

In particular, since $\frac{\alpha\gamma}{4 + \alpha^2} < \frac{\gamma}{\alpha}$, we also have $\omega < \frac{\gamma}{\alpha}$. Furthermore, if $1 < \frac{\gamma^2}{4 + \alpha^2} + \frac{2\gamma}{\alpha(4 + \alpha^2)}\sqrt{4 + \alpha^2 - \gamma^2}$, then ω satisfies the admissibility condition (1.8). Orbital stability then follows from Proposition 4.9 and Lemma 4.15.

The condition $\alpha\gamma < 2\sqrt{4 + \alpha^2 - \gamma^2}$ implies that there exists a $\omega \in (-\frac{1}{2}\sqrt{p-1}, 0)$ satisfying (1.8). Orbital instability follows from Proposition 4.9, Remark 4.10 and Lemma 4.15. \square

Remark 4.17. The condition $1 < \frac{\gamma^2}{4 + \alpha^2} + \frac{2\gamma}{\alpha(4 + \alpha^2)}\sqrt{4 + \alpha^2 - \gamma^2}$ is satisfied if $0 < \alpha < \gamma < \sqrt{4 + \alpha^2}$.

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