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# Weighted Spectral Embedding of Graphs

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#### Abstract

We present a novel spectral embedding of graphs that incorporates weights assigned to the nodes, quantifying their relative importance. This spectral embedding is based on the first eigenvectors of some properly normalized version of the Laplacian. We prove that these eigenvectors correspond to the configurations of lowest energy of an equivalent physical system, either mechanical or electrical, in which the weight of each node can be interpreted as its mass or its capacitance, respectively. Experiments on a real dataset illustrate the impact of weighting on the embedding.

#### 1 Introduction

Many types of data can be represented as graphs. Edges may correspond to actual links in the data (e.g., users connected by some social network) or to levels of similarity induced from the data (e.g., users having liked a large common set of movies). The resulting graph is typically *sparse* in the sense that the number of edges is much lower than the total number of node pairs, which makes the data hard to exploit.

A standard approach to the analysis of sparse graphs consists in *embedding* the graph in some vectorial space of low dimension, typically much smaller than the number of nodes [15, 19, 4]. Each node is represented by some vector in the embedding space so that close nodes in the graph (linked either directly or through many short paths in the graph) tend to be represented by close vectors in terms of the Euclidian distance. Standard learning techniques can then be applied to this dense vectorial representation of the graph to recommend new links, rank nodes or find clusters of nodes for instance [7, 3].

The most popular technique for graph embedding is based on the spectral decomposition of the graph Laplacian, each dimension of the embedding space corresponding to an eigenvector of the Laplacian matrix [5, 13, 2, 11, 17, 12]. This spectral embedding can be interpreted in terms of a random walk in the graph [10, 14]. In the full embedding space (including all eigenvectors), the square distance between two vectors is proportional to the mean commute time of the random walk between the two corresponding nodes: close nodes in the graph tend to be close in the embedding space. Viewing this random walk as the path followed by electrons in the corresponding electrical network, with nodes linked by resistors, these square distances can also be interpreted as the effective resistances between pairs of nodes [16].

In this paper, we address the issue of the spectral embedding of graphs including node weights. We shall see that existing spectral embedding techniques implicitly consider either unit weights or so-called internal node weights, depending on the Laplacian used in the spectral decomposition. We here consider the node weights as some additional information representing the relative importance of the nodes, independently of the graph structure. The weight of a node may reflect either its value, its multiplicity if each node represents a category of users or items, or the reliability of the associate data, for instance. Surprisingly, this notion of weight is common for vectoral data (see, e.g., the weighted version of the k-means clustering algorithm [8, 6, 9]) but not for graph data, where weights are typically assigned to edges but not to nodes, apart from those induced from the edges.

Our main contribution is a spectral embedding of the graph, we refer to as the weighted spectral embedding, that incorporates the node weights. It is based on the spectral decomposition of some properly normalized version of the Laplacian. We prove that, when all eigenvectors are used, this embedding is equivalent to the regular spectral embedding shifted so that the origin is the center of mass of the embedding. In practice, only the first eigenvectors are included to get an embedding in low dimension. We show that these eigenvectors can be interpreted as the levels of lowest energy of a physical system, either a mechanical system where nodes are linked by springs (the edges) and have different masses (the node weights), or an electrical network where nodes are linked by resistors (the edges) and connected to the ground by capacitors with different capacitances (the node weights). In particular, the weighted spectral embedding can not be derived from the regular spectral embedding in low dimension. Experiments confirm that these embeddings differ significantly in practice.

The weighted spectral embedding can also be interpreted in terms of a random walk in the graph, where nodes are visited in proportion to their weights. In the full embedding space (including all eigenvectors), the square distances between vectors are proportional to the *mean commute times* of this random walk between the corresponding nodes, as for unit weights. In fact, these mean commute times depend on the weights through their sum only, which explains why the pairwise distances of the embedding are independent of the weights, up to some multiplicative constant. This property is somewhat counter-intuitive as the mean hitting time of one node from another *does* depend on the weights. We shall explain this apparent paradox by some symmetry property of each equivalent physical system.

The rest of the paper is organized as follows. We first introduce the model and the notations. We then present the regular spectral embedding and its interpretation in terms of a random walk in the graph. The weighted version of this random walk is introduced in Section 4. Section 5 presents the weighted spectral embedding and extends the results known for the regular spectral embedding. The analogies with a mechanical system and an electrical network are described in Sections 6 and 7, respectively. Experiments on real data are presented in Section 8. Section 9 concludes the paper.

#### 2 Model

We consider a connected, undirected graph of n nodes, without self-loops. We denote by A its adjacency matrix. In the absence of edge weights, this is a binary, symmetric matrix, with  $A_{ij} = 1$  if and only if there is an edge between nodes i and j. In the presence of edge weights,  $A_{ij}$  is the weight of the edge between nodes i and j, if any, and is equal to 0 otherwise.

Let d = Ae, where e is the n-dimensional vector of ones. The components  $d_1, \ldots, d_n$  of the vector d are equal to the actual node degrees in the absence of edge weights (A is a binary matrix) and to the total weights of incident edges otherwise (A is a non-negative matrix). We refer to  $d_1, \ldots, d_n$  as the *internal* node weights.

Nodes are assigned positive weights  $w_1, \ldots, w_n$  corresponding to their relative importances. These node weights are external parameters, independent of the graph. We denote by w the vector  $(w_1, \ldots, w_n)$ .

# 3 Spectral embedding

We first present the regular spectral embedding, without taking the node weights  $w_1, \ldots, w_n$  into account. Let D = diag(d) be the diagonal matrix of internal node weights. The Laplacian matrix is defined by

$$L = D - A$$
.

This is a symmetric matrix. It is positive semi-definite on observing that:

$$\forall v \in \mathbb{R}^n, \quad v^T L v = \sum_{i < j} A_{ij} (v_j - v_i)^2.$$

The spectral theorem yields

$$L = U\Lambda U^T, \tag{1}$$

where  $\Lambda = \operatorname{diag}(\lambda_1, \dots, \lambda_n)$  is the diagonal matrix of eigenvalues of L, with  $0 = \lambda_1 < \lambda_2 \leq \dots \leq \lambda_n$ , and  $U = (u_1, \dots, u_n)$  is the matrix of corresponding eigenvectors, with  $U^T U = I$  and  $u_1 = e/\sqrt{n}$ .

**Spectral embedding.** Let  $X = \sqrt{\Lambda^+}U^T$ , where  $\Lambda^+ = \text{diag}(0, 1/\lambda_2, \dots, 1/\lambda_n)$  denotes the pseudo-inverse of  $\Lambda$ . The columns  $x_1, \dots, x_n$  of the matrix X define an embedding of the nodes in  $\mathbb{R}^n$ , each dimension corresponding to an eigenvector of the Laplacian. Observe that the first component of each vector  $x_1, \dots, x_n$  is equal to 0, reflecting the fact that the first eigenvector is not informative. Since Xe = 0, the centroid of the n vectors is the origin:

$$\frac{1}{n}\sum_{i=1}^{n}x_{i}=0. (2)$$

The Gram matrix of X is the pseudo-inverse of the Laplacian:

$$X^T X = U \Lambda^+ U^T = L^+$$
.

**Random walk.** Consider a random walk in the graph where the transition rate from node i to node j is  $A_{ij}$ . Specifically, the walker stays at node i an exponential time with parameter  $d_i$ , then moves from node i to node j with probability  $P_{ij} = A_{ij}/d_i$ . This defines a continuous-time Markov chain with generator matrix -L and uniform stationary distribution. The sequence of nodes visited by the random walk forms a discrete-time Markov chain with transition matrix  $P = D^{-1}A$ .

**Hitting times.** Let  $H_{ij}$  be the mean hitting time of node j from node i. Observe that  $H_{ii} = 0$ . The following results, proved in [10, 14], will be extended to weighted spectral embedding in Section 4. We denote by  $e_i$  the n-dimensional unit vector on component i.

**Proposition 1** The matrix H satisfies:

$$LH = ee^T - nI. (3)$$

**Proposition 2** The solution H to equation (3) with zero diagonal entries satisfies:

$$\forall i, j, \quad H_{ij} = n(e_j - e_i)^T L^+ e_j. \tag{4}$$

Since  $L^+ = X^T X$ , we obtain:

$$H_{ij} = n(x_j - x_i)^T x_j.$$

We deduce the mean commute time between nodes i and j:

$$C_{ij} = H_{ij} + H_{ji} = n||x_i - x_j||^2.$$

In view of (2), the mean hitting time of node j from steady state is:

$$h_j = \frac{1}{n} \sum_{i=1}^n H_{ij} = n||x_j||^2.$$

Cosine similarity. Observe that:

$$nx_i^T x_j = h_j - H_{ij} = h_i - H_{ji}.$$

Let  $S_{ij} = \cos(x_i, x_j)$ . This is the cosine-similarity between nodes i and j. We have:

$$S_{ij} = \frac{h_j - H_{ij}}{\sqrt{h_i h_j}} = \frac{h_i - H_{ji}}{\sqrt{h_i h_j}},$$

that is

$$S_{ij} = \frac{h_i + h_j - C_{ij}}{2\sqrt{h_i h_j}}. (5)$$

Thus the cosine-similarity between any vectors  $x_i, x_j$  can be interpreted in terms of the mean commute time between the corresponding nodes i, j relative to their mean hitting times.

## 4 Random walk with weights

In this section, we introduce a modified version of the random walk, that takes the node weights into account. Recall that all weights are assumed positive. We denote by  $|w| = \sum_{i=1}^{n} w_i$  the total node weight.

**Random walk.** We modify the random walk as follows: the transition rate from node i to node j is now  $A_{ij}/w_i$ . Thus the walker stays at node i an exponential time with parameter  $d_i/w_i$ , then moves from node i to node j with probability  $P_{ij}$ . Observe that the previously considered random walk corresponds to unit weights. We get a continuous-time Markov chain with generator matrix  $-W^{-1}L$ , with  $W = \operatorname{diag}(w)$ , and stationary distribution  $\pi = w/|w|$ .

**Hitting times.** Let  $H_{ij}$  be the mean hitting time of node j from node i. Observe that  $H_{ii} = 0$ .

**Proposition 3** The matrix H satisfies:

$$LH = we^T - |w|I. (6)$$

*Proof.* We have  $H_{ii} = 0$  while for all  $i \neq j$ ,

$$H_{ij} = \frac{w_i}{d_i} + \sum_{k=1}^n P_{ik} H_{kj}.$$

Thus the matrix  $(I-P)H - D^{-1}we^T$  is diagonal. Equivalently, the matrix  $LH - we^T$  is diagonal. Since  $e^TL = 0$ , this diagonal matrix is -|w|I.

**Proposition 4** The solution H to equation (6) with zero diagonal entries satisfies:

$$\forall i, j, \quad H_{ij} = |w|(e_j - e_i)^T L^+(e_j - \pi).$$

*Proof.* The matrix  $M = we^T - |w|I$  satisfies  $e^T M = 0$  so the seeked solution is of the form:

$$H = L^+(we^T - |w|I) + eh^T,$$

for some n-dimensional vector h. Thus,

$$H_{ij} = |w|e_i^T L^+(\pi - e_j) + h_j.$$

Since  $H_{jj} = 0$ , we get:

$$h_j = -|w|e_j^T L^+(\pi - e_j).$$

Let  $\bar{x}$  be the center of mass of the vectors  $x_1, \ldots, x_n$ :

$$\bar{x} = X\pi = \sum_{i=1}^{n} \pi_i x_i.$$

Since  $L^+ = X^T X$ , we obtain:

$$H_{ij} = |w|(x_j - x_i)^T (x_j - \bar{x}). \tag{7}$$

Thus the mean hitting times are the *same* as without node weights, up to the multiplicative constant |w| and the shift of the origin to the center of mass  $\bar{x}$  of the vectors  $x_1, \ldots, x_n$ .

In particular, the mean commute time between any nodes i and j depends on the weights through their sum only:

$$C_{ij} = H_{ij} + H_{ji} = |w|||x_i - x_j||^2.$$

The mean hitting time of node j from steady state is:

$$h_j = \sum_{i=1}^n \pi_i H_{ij} = |w| ||x_j - \bar{x}||^2.$$

Cosine similarity. Let  $S_{ij} = \cos(x_i - \bar{x}, x_j - \bar{x})$ . This is the cosine-similarity between nodes i and j. We obtain as above:

$$S_{ij} = \frac{h_i + h_j - C_{ij}}{2\sqrt{h_i h_j}}.$$

In particular, the relative mean commute times, as defined by (5), depend on the weights  $w_1, \ldots, w_n$  through the center of mass  $\bar{x}$  of the vectors  $x_1, \ldots, x_n$  only.

## 5 Weighted spectral embedding

We now introduce the weighted spectral embedding. The generator matrix  $-W^{-1}L$  of the weighted random walk is not symmetric in general. Thus we consider the following normalized version of the Laplacian, we refer to as the *weighted* Laplacian:

$$L_W = W^{-\frac{1}{2}} L W^{-\frac{1}{2}}.$$

This matrix is symmetric, positive semi-definite. In the particular case where W = D, this is known as the symmetric normalized Laplacian.

The spectral theorem yields

$$L_W = \hat{U}\hat{\Lambda}\hat{U}^T,\tag{8}$$

where  $\hat{\Lambda} = \operatorname{diag}(\hat{\lambda}_1, \dots, \hat{\lambda}_n)$  is the diagonal matrix of eigenvalues of  $L_W$ , with  $0 = \hat{\lambda}_1 < \hat{\lambda}_2 \leq \dots \leq \hat{\lambda}_n$ , and  $\hat{U} = (\hat{u}_1, \dots, \hat{u}_n)$  is the matrix of corresponding eigenvectors, with  $\hat{U}^T \hat{U} = I$  and  $\hat{u}_1 = \sqrt{\pi}$ .

We have the following explicit expression for the pseudo-inverse of the weighted Laplacian:

**Proposition 5** The pseudo-inverse of  $L_W$  is:

$$L_W^+ = W^{\frac{1}{2}} (I - e\pi^T) L^+ (I - \pi e^T) W^{\frac{1}{2}}. \tag{9}$$

*Proof.* Let M be the matrix defined by the right-hand side of (9). Using the fact that Le = 0, we get:

$$L_W M L_W = W^{-\frac{1}{2}} L L^+ L W^{-\frac{1}{2}},$$
  
=  $W^{-\frac{1}{2}} L^+ W^{-\frac{1}{2}} = L_W$ 

which proves that M is the pseudo-inverse of  $L_W$ .

Generalized eigenvalue problem. Let  $V = W^{-\frac{1}{2}}\hat{U}$ . We have:

$$LV = WV\hat{\Lambda},$$

with  $V^TWV = I$ . Thus the columns v of the matrix V are solutions to the generalized eigenvalue problem:

$$Lv = \lambda W v, \tag{10}$$

with corresponding eigenvalues  $\lambda = \hat{\lambda}_1, \dots, \hat{\lambda}_n$ . The first column, associated to the eigenvalue  $\hat{\lambda}_1 = 0$ , satisfies  $v \propto e$ , while the others satisfy  $w^T v = 0$ .

**Spectral embedding.** Let  $Y = \sqrt{\Lambda^+} \hat{U}^T W^{-\frac{1}{2}}$ . The columns  $y_1, \ldots, y_n$  of the matrix Y define an embedding of the nodes, we refer to as the *weighted* spectral embedding. As for the regular spectral embedding, the first component of each vector  $y_1, \ldots, y_n$  is equal to 0. Since Yw = 0, the center of mass of  $y_1, \ldots, y_n$  lies at the origin:

$$\sum_{i=1}^{n} w_i y_i = 0.$$

The Gram matrix of Y is:

$$Y^TY = W^{-\frac{1}{2}} \hat{U} \hat{\Lambda}^+ \hat{U}^T W^{-\frac{1}{2}} = W^{-\frac{1}{2}} L_W^+ W^{-\frac{1}{2}}.$$

In view of Proposition 5,

$$y_i^T y_j = e_i^T W^{-\frac{1}{2}} L_W^+ W^{-\frac{1}{2}} e_j,$$
  
=  $(e_i - \pi)^T L^+ (e_j - \pi)^T,$   
=  $(x_i - \bar{x})^T (x_j - \bar{x}).$ 

Thus the considered weighted spectral embedding is equivalent to the regular spectral embedding, shifted so that the origin is the center of mass  $\bar{x}$ . In particular, the distances between vectors in the embedding space can be interpreted in terms of mean hitting times of the random walk, as shown in Section 4.

# 6 A mechanical system

Consider the mechanical system consisting of n point particles of respective masses  $w_1, \ldots, w_n$  sliding along a bar without friction. Particles i and j are linked by a spring satisfying Hooke's law with stiffness  $A_{ij}$ .

**Eigenmodes.** Assume that the bar has a uniform circular motion with angular velocity  $\omega$  around some fixed axis. We denote by  $v_1, \ldots, v_n$  the locations of the particles along the bar, with the axis taken as the origin. By Newton's second law of motion, the system is in equilibrium if and only if

$$\forall i = 1, \dots, n, \quad \sum_{j=1}^{n} A_{ij}(v_j - v_i) = -w_i v_i \omega^2,$$

that is

$$Lv = \omega^2 W v$$
.

If  $v \neq 0$ , then v is a solution to the generalized eigenvector problem (10) with eigenvalue  $\lambda = \omega^2$ . We refer to these equilibrium states as the *eigenmodes* of the mechanical system.

The first eigenmode  $v \propto e$ , for  $\omega^2 = 0$ , corresponds to the absence of motion. The other eigenmodes give the possible angular velocities of the system, equal to the square roots of the eigenvalues  $\hat{\lambda}_2, \ldots, \hat{\lambda}_n$  of  $L_W$ . Any such eigenmode v satisfies  $w^T v = 0$ , meaning that the center of mass of the system is at the origin.

**Potential energy.** The mechanical potential energy of the system in any state v is:

$$\frac{1}{2}v^TLv.$$

If v is an eigenmode of the mechanical system, this is equal to:

$$\frac{\lambda}{2}v^T W v.$$

As  $v^T W v$  is the moment of inertia of the system and  $\lambda = \omega^2$  is the square of the angular velocity, this corresponds to the angular kinetic energy. For a unit moment of inertia  $v^T W v = 1$ , we obtain  $\lambda = v^T L v$ , so that the eigenvalues of the weighted Laplacian  $L_W$  can be viewed as (twice) the levels of energy of the eigenmodes, for unit moments of inertia. The weighted spectral embedding reduced to the k first eigenvectors of  $L_W$  can then be interpreted as that induced by the k eigenmodes of lowest potential energies of the mechanical system, for unit moments of inertia.

**Dirichlet problem.** For any  $i \neq j$ , assume the positions of the point particles i and j are set to 1 and 0, respectively. Let v be the vector of positions at equilibrium, in the absence of rotation. By Newton's first law of motion,

$$Lv = \alpha(e_i - e_j),$$

for some constant  $\alpha$  equal to the force F exerted on both i and j (in opposite directions). This force does not depend on the masses of the point particles.

The solution to this Dirichlet problem is:

$$v = \alpha L^{+}(e_i - e_i) + \beta e,$$

for some constant  $\beta$ . Using the fact that  $v_i = 1$  and  $v_j = 0$ , we get:

$$v_k = \frac{(e_k - e_j)^T L^+(e_i - e_j)}{(e_i - e_j)^T L^+(e_i - e_j)} = \frac{(x_k - x_j)^T (x_i - x_j)}{||x_i - x_j||^2}$$

and

$$\alpha = \frac{1}{||x_i - x_i||^2}.$$

Observe in particular that  $v_k \in [0,1]$  for all nodes k. The mean hitting time of node j from node i by the random walk, given by (7), satisfies:

$$H_{ij} = \frac{|w|}{F}\bar{v},$$

where  $\bar{v}$  is the center of mass of the system:

$$\bar{v} = \sum_{k=1}^{n} \pi_k v_k.$$

By symmetry, we have:

$$H_{ji} = \frac{|w|}{F}(1-\bar{v}).$$

We deduce the mean commute time between nodes i and j:

$$C_{ij} = H_{ij} + H_{ji} = \frac{|w|}{F}.$$

This symmetry in the solutions to the Dirichlet problem explains why, unlike the mean hitting times, the mean commute times depend on the weights through their sum only. The mean commute time between two nodes of the graph is inversely proportional to the force exerted between the two corresponding point particles of the mechanical system, which is independent of the particle masses.

#### 7 An electrical network

Consider the electrical network induced by the graph, with a resistor of conductance  $A_{ij}$  between nodes i and j and a capacitor of capacitance  $w_i$  between node i and the ground. The move of electrons in the network is precisely described by the random walk defined in Section 4.

**Eigenmodes.** Let v(t) be the vector of electric potentials at time t, starting from some arbitrary initial state v. By Ohm's law,

$$\forall i = 1, \dots, n, \quad w_i \frac{dv_i(t)}{dt} = \sum_{i=1}^n A_{ij}(v_j(t) - v_i(t)),$$

that is

$$W\frac{dv(t)}{dt} = -Lv(t).$$

The solution to this differential equation is:

$$v(t) = e^{-W^{-1}Lt}v.$$

If the initial state v is a solution to the generalized eigenvector problem (10), we obtain:

$$v(t) = e^{-\lambda t}v.$$

Since the eigenvectors of  $L_W$  form a basis of  $\mathbb{R}^n$ , any initial state v can be written as a linear combination of solutions to the generalized eigenvector problem, we refer to as the *eigenmodes* of the electrical network. The first eigenmode  $v \propto e$ , for  $\lambda = 0$ , corresponds to a static system, without discharge. The other eigenmodes give the possible discharge rates of the system, equal to the eigenvalues  $\hat{\lambda}_2, \ldots, \hat{\lambda}_n$  of  $L_W$ . Any such eigenmode v satisfies  $w^T v = 0$ , meaning that the total charge accumulated in the capacitors is null.

**Energy dissipation.** The energy dissipated by the resistors for the vector of electric potentials v is:

$$\frac{1}{2}v^TLv.$$

If v be an eigenmode of the electrical network, this is equal to

$$\frac{\lambda}{2}v^T W v.$$

Observing that  $\frac{1}{2}v^TWv$  is the electrical potential energy of the system, we deduce that  $\lambda$  can be viewed as the energy dissipated by the resistors for a unit electric potential energy. In particular, the weighted spectral embedding reduced to the k first eigenvectors of  $L_W$  can be interpreted as that induced by the k eigenmodes of lowest energy dissipation of the electrical network, for unit electric potential energies.

**Dirichlet problem.** For any  $i \neq j$ , assume the electric potentials of nodes i and j are set to 1 and 0, respectively. Let v be the vector of electric potentials at equilibrium. We have:

$$Lv = \alpha(e_i - e_j),$$

for some constant  $\alpha$  equal to the current flowing from i to j. This Dirichlet problem is the same as for the mechanical network, and thus the same results apply. In particular, we get  $v_k \in [0,1]$  for all nodes k, and the current between i to j is:

$$\alpha = \frac{1}{||x_i - x_j||^2}.$$

This is the intensity of the current I generated by a unit potential difference between i and j. Its inverse is known as the *effective resistance* between i and j. The mean hitting time of node j from node i by the random walk satisfies:

$$H_{ij} = \frac{q}{I},$$

where q is the total charge accumulated in the capacitors:

$$q = \sum_{k=1}^{n} w_k v_k.$$

Observe that this expression for the mean hitting time can be seen as a consequence of Little's law, as I is the intensity of positive charges entering the network and q the mean number of positive charges accumulated in the network.

By symmetry, the total charge accumulated in the capacitors when the respective electric potentials of i and j are set to 0 and 1 is:

$$\sum_{k=1}^{n} w_k (1 - v_k) = |w| - q,$$

so that:

$$H_{ji} = \frac{|w| - q}{I}.$$

We deduce the mean commute time between nodes i and j:

$$C_{ij} = H_{ij} + H_{ji} = \frac{|w|}{I}.$$

Again, the symmetry in the solutions to the Dirichlet problem explains why the mean commute times depend on the weights through their sum only. The mean commute time between two nodes of the graph is proportional to the effective resistance between the two corresponding points of the electrical network, which is independent of the capacitors.

# 8 Experiments

We now illustrate the results on a real dataset. The considered graph is that formed by articles of Wikipedia for Schools<sup>1</sup>, a selection of articles of Wikipedia for children [18]. Specifically, we have extracted the largest connected component of this graph, considered as undirected. The resulting graph has 4,589 nodes (the articles) and 106,644 edges (the hyperlinks between these articles). The graph is undirected and unweighted. Both the dataset and the Python code used for the experiments are available online<sup>2</sup>.

**Global clustering.** We first apply k-means clustering<sup>3</sup> to the following embeddings of the graph, each in dimension k = 100:

- the regular spectral embedding, X, restricted to rows  $2, \ldots, k+1$ , say  $\hat{X}$ ;
- the shifted spectral embedding,  $\hat{X} \hat{X}\pi 1^T$ , where  $\pi = w/|w|$  and 1 is the k-dimensional vector of ones;
- the weighted spectral embedding, Y, restricted to rows  $2, \ldots, k+1$ , say  $\hat{Y}$ .

<sup>1</sup>https://schools-wikipedia.org

<sup>&</sup>lt;sup>2</sup>https://github.com/tbonald/spectral\_embedding

<sup>&</sup>lt;sup>3</sup>The algorithm is k-means++ [1] 100 random initializations.

Here the vector of weights w is taken equal to d, the vector of internal node weights. In particular, the weighted spectral embedding is that following from the spectral decomposition of the normalized Laplacian,  $D^{-\frac{1}{2}}LD^{-\frac{1}{2}}$ . Observe that the first row of each embedding X and Y is equal to zero and thus discarded.

Each embedding is normalized so that nodes are represented by k-dimensional unitary vectors. This is equivalent to consider the distance induced by the *cosine similarity* in the original embedding. In particular, the regular spectral embedding and the shifted spectral embedding give different clusterings.

Tables 1 and 2 show the top articles of the clusters found for each embedding, when the number of clusters is set to 20, with the size of each cluster. The selected articles for each cluster correspond to the nodes of highest degrees among the 50% closest nodes from the center of mass of the cluster in the embedding space, with unit weights for the regular spectral embedding and internal node weights d for the shifted spectral embedding and the weighted spectral embedding.

Size	Top articles
1113	Australia, Canada, North America, 20th century
326	UK, England, London, Scotland, Ireland
250	US, New York City, BBC, 21st century, Los Angeles
227	India, Japan, China, United Nations
218	Earth, Sun, Physics, Hydrogen, Moon, Astronomy
200	Mammal, Fish, Horse, Cattle, Extinction
200	France, Italy, Spain, Latin, Netherlands
198	Water, Agriculture, Coal, River, Antarctica
197	Germany, World War II, Russia, World War I
187	Mexico, Brazil, Atlantic Ocean, Argentina
185	Human, Philosophy, Slavery, Religion, Democracy
184	Plant, Rice, Fruit, Sugar, Wine, Maize, Cotton
177	Gold, Iron, Oxygen, Copper, Electron, Color
170	Egypt, Turkey, Israel, Islam, Iran, Middle East
159	English, 19th century, William Shakespeare, Novel
158	Africa, South Africa, Time zone, Portugal
157	Europe, Scientific classification, Animal, Asia
141	Washington, D.C., President of the United States
72	Dinosaur, Fossil, Reptile, Cretaceous, Jurassic
70	Paris, Art, Architecture, Painting, Hist. of painting

Table 1: Global clustering of Wikipedia for Schools by weighted spectral embedding.

The first observation is that the three embeddings are very different. The choice of the Laplacian (regular or normalized) matters, because the dimension k is much smaller than the number of nodes n (recall that the shifted spectral embedding and the weighted spectral embedding are equivalent when k = n). The second observation is that the weighted spectral embedding seems to better capture the structure of the graph. Apart from the first cluster, which is significantly larger than the others and thus may contain articles on very different topics, the other clusters look meaningful. The regular spectral embedding puts together articles as different as Meteorology, British English and Number, while the shifted embedding groups together articles about Sanskrit, Light and The Simpsons, that should arguably appear in different clusters.

Size	Top articles
1666	Germany, India, Africa, Spain, Russia, Asia
526	Chordate, Binomial nomenclature, Bird, Mammal
508	Earth, Water, Iron, Sun, Oxygen, Copper, Color
439	Scotland, Ireland, Wales, Manchester, Royal Navy
300	New York City, Los Angeles, California, Jamaica
246	North America, German language, Rome, Kenya
195	English, Japan, Italy, 19th century
131	Language, Mass media, Library, Engineering, DVD
103	Jazz, Piano, Guitar, Music of the United States
91	Microsoft, Linux, Microsoft Windows, Algorithm
67	British monarchy, Bristol, Oxford, Paul of Tarsus
58	Tropical cyclone, Bermuda, Hurricane Andrew
49	Mathematics, Symmetry, Geometry, Algebra, Euclid
48	Fatty acid, List of vegetable oils, Biodiesel
47	Train, Canadian Pacific Railway, Denver, Colorado
38	Eye, Retina, Animation, Glasses, Lego, Retinol
27	Finance, Supply and demand, Stock, Accountancy
19	Meteorology, British English, Number, Vowel
17	Newcastle upon Tyne, Isambard Kingdom Brunel
14	Wikipedia, Jimmy Wales, Wikimedia Foundation
Size	Top articles
452	Chordate, Binomial nomenclature, Bird, Mammal
446	UK, Europe, France, English language, Japan
369	Germany, Spain, Soviet Union, Sweden, Poland
369	Earth, Water, Iron, Sun, Oxygen, Copper, Color
353	India, Africa, Russia, New Zealand, River, Snow
328	Egypt, Greece, Middle Ages, Roman Empire, Nazism
286	United States, New York City, Petroleum, Finland
271	British Empire, 17th century, Winston Churchill
238	Time zone, Turkey, Portugal, Israel, Currency
217	Rice, Fruit, Bacteria, Wine, DNA, Flower, Banana
178	British monarchy, Bristol, Charles II of England
175	Internet, Hebrew language, Language, Mass media
156	Physics, Ancient Egypt, Science, Astronomy, Time
125	Atlantic Ocean, Morocco, Algeria, Barbados
124	Opera, Folk music, Elvis Presley, Bob Dylan
117	Agriculture, Ocean, Geology, Ecology, Pollution
112	Christianity, Switzerland, Judaism, Bible, Deity
108	South America, Pacific Ocean, Tourism, Colombia
94	Film, Sanskrit, Light, The Simpsons, Eye, Shark
71	Mining, Food, Geography, Engineering, Transport

Table 2: Global clustering of Wikipedia for Schools by regular (top) and shifted (bottom) spectral embeddings.

Selective clustering. We are now interested in the clustering of some selection of the articles. We take the 667 articles in the *People* category. A naive approach consists in considering the subgraph induced by these nodes. But this is not satisfactory as the similarity between articles in the People category depends strongly on their connections through articles that not in this category. Indeed, the subgraph of nodes in the People category is not even connected. Our approach consists in assigning a multiplicative factor of 10 to articles in the People category. Specifically, we set  $w_i = 10 \times d_i$  if article i belongs to the People category, and  $w_i = d_i$  otherwise. We compare the three previous embeddings, for this vector of weights.

Tables 3 and 4 show the top articles in the People category for the 20 clusters found for each embedding. The selected articles for each cluster correspond to the nodes of highest degrees in the People category. We also indicate the number of articles in the People category in each cluster. Observe that the last cluster obtained with the regular spectral embedding has no article in the People category. Again, the impact of weighting is significant. The weighted spectral embedding, which is adapted to the People category, seems to better capture the structure of the graph. Except for the first, which is much larger than the others, the clusters tend to group together articles of the People category that are closely related.

Count	Top articles
228	William Shakespeare, Pope John Paul II
60	Jorge Luis Borges, Rabindranath Tagore
55	Elizabeth II of the UK, Winston Churchill, Tony Blair
49	Charles II of England, Elizabeth I of England
48	Ronald Reagan, Bill Clinton, Franklin Roosevelt
25	Alexander the Great, Genghis Khan, Muhammad
24	Adolf Hitler, Joseph Stalin, Vladimir Lenin
23	Napoleon I, Charlemagne, Louis XIV of France
22	Jesus, Homer, Julius Caesar, Virgil
22	Aristotle, Plato, Charles Darwin, Karl Marx
21	George W. Bush, Condoleezza Rice, Nelson Mandela
18	Elvis Presley, Paul McCartney, Bob Dylan
15	Isaac Newton, Galileo Galilei, Ptolemy
13	Albert Einstein, Gottfried Leibniz, Bertrand Russell
10	Pete Sampras, Boris Becker, Tim Henman
9	William Thomson, 1st Baron Kelvin, Humphry Davy
8	Carolus Linnaeus, James Cook, Gerald Durrell
8	Bill Gates, Richard Stallman, Ralph Nader
5	Floyd Mayweather Jr., Lucy, Lady Duff-Gordon
4	Vasco da Gama, Idit Harel Caperton, Reza Shah

Table 3: Selective clustering of Wikipedia for Schools by weighted spectral clustering.

Count	Top articles
196	Carolus Linnaeus, Adolf Hitler, Jesus, Aristotle
95	Christina Aguilera, Andy Warhol, Auguste Rodin
90	Julius Caesar, Martin Luther King, Jr., Euclid
75	Tony Blair, Victoria of the UK, Charles II of England
45	Ronald Reagan, Franklin D. Roosevelt, Gerald Ford
28	Abraham Lincoln, George Washington, John Adams
25	Igor Stravinsky, Johann Sebastian Bach
24	Albert Einstein, Gottfried Leibniz, Isaac Newton
22	George W. Bush, Napoleon I of France, Plato
10	Pete Sampras, Boris Becker, Tim Henman, Pat Cash
10	Fanny Blankers-Koen, Rosa Parks, Donald Bradman
9	Alexander the Great, Frederick II of Prussia
9	Mahatma Gandhi, Buddha, Muhammad Ali Jinnah
8	Columba, Edwin of Northumbria, Macbeth of Scotland
5	Bill Clinton, Richard Stallman, Linus Torvalds
5	Elizabeth II of the UK, David Beckham, Wayne Rooney
5	Archbishop of Canterbury, Harold Wilson
4	Leonardo da Vinci, Neil Armstrong, Wright brothers
2	George III of the UK, Matthew Brettingham
0	_
Count	Top articles
182	Adolf Hitler, William Shakespeare
113	Elizabeth II of the UK, George W. Bush
69	Tony Blair, Victoria of the UK, Elizabeth I of England
46	Ronald Reagan, Franklin Roosevelt, Jimmy Carter
31	Henry James, Igor Stravinsky, Ezra Pound
26	Paul McCartney, Bob Dylan, Edgar Allan Poe
25	Jesus, Charlemagne, Genghis Khan, Homer, Ptolemy
24	Albert Einstein, Gottfried Leibniz, Isaac Newton
23	Charles Darwin, Galileo Galilei, Nikola Tesla
19	Plato, John Locke, Max Weber, Friedrich Nietzsche
18	Rabindranath Tagore, Mahatma Gandhi, Buddha
17	Margaret Thatcher, David Cameron, George VI
16	Condoleezza Rice, Nelson Mandela, Gerald Ford
12	Dwight D. Eisenhower, Ernest Hemingway
11	Aristotle, Alexander the Great, Fred. II of Prussia
10	Pete Sampras, Boris Becker, Tim Henman
	3.5.1 3.37 53.1 63.1.7.3
8	Muhammad, Norman Borlaug, Osama bin Laden
8 7	Bill Clinton, Bill Gates, Richard Stallman
8	

Table 4: Selective clustering of Wikipedia for Schools by regular (top) and shifted (bottom) spectral clustering.

## 9 Conclusion

We have proposed a novel spectral embedding of graphs that takes node weights into account. We have proved that the dimensions of this embedding correspond to different configurations of an equivalent physical system, either mechanical or electrical, with node weights corresponding to masses or capacitances, respectively.

A practically interesting consequence of our work is in the choice of the Laplacian, when there are no other information on the node weights than the graph itself. Thanks to weighted spectral embedding, we see that the two versions of the Laplacian, regular and normalized, correspond in fact to two relative importance of nodes, given respectively by unit weights and internal node weights.

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