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# A theoretical investigation of Brockett's ensemble optimal control problems 

Jan Bartsch* Alfio Borzì ${ }^{\dagger} \quad$ Francesco Fanelli ${ }^{\ddagger} \quad$ Souvik Roy ${ }^{\S}$

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#### Abstract

This paper is devoted to the analysis of problems of optimal control of ensembles governed by the Liouville (or continuity) equation. The formulation and study of these problems have been put forward in recent years by R.W. Brockett, with the motivation that ensemble control may provide a more general and robust control framework.

Following Brockett's formulation of ensemble control, a Liouville equation with unbounded drift function, and a class of cost functionals that include tracking of ensembles and different control costs is considered. For the theoretical investigation of the resulting optimal control problems, a well-posedness theory in weighted Sobolev spaces is presented for the Liouville and transport equations. Then, a class of non-smooth optimal control problems governed by the Liouville equation is formulated and existence of optimal controls is proved. Furthermore, optimal controls are characterised as solutions to optimality systems; such a characterisation is the key to get (under suitable assumptions) also uniqueness of optimal controls.


Keywords: Liouville and transport equations, well-posedness theory, weighted spaces, optimal control theory, non-smooth optimization, optimality systems.

2010 Mathematics Subject Classification: 49 J 20 (primary); 35L03, 35B65, 49K20, 35Q93 (secondary).

## 1 Introduction

The notion of ensemble control was proposed by R.W. Brockett in [8], and further in [10, 9], while considering the problem of a trade-off between the complexity of implementing a control strategy and the performance of the control system. For the former, Brockett discusses the concept of minimum attention control that results in costs of the control that involve a partial time-derivative of the control function. For the latter, he emphasizes the advantage of considering an ensemble of trajectories, which stem from a distribution of initial conditions, rather than individual trajectories. By these two consideration, Brockett

[^0]concludes that the natural setting for investigating both aspects of the resulting control problem is by means of the Liouville (or continuity) equation that governs the evolution of the ensemble of trajectories.

The Liouville equation is a hyperbolic-type partial differential equation, often used to model the evolution of density functions representing the probability density of multiple trials of a single evolving ordinary differential equation (ODE in brief) system, or the physical (e.g. particle) density of multiple non-interacting systems. In both cases, the function of the dynamics of the ODE model appears as the drift coefficient of the Liouville equation. Therefore the problem of controlling a trajectory of a finite-dimensional dynamical system is lifted to the problem of controlling a continuum of dynamical systems with the same control strategy. Specifically, this setting results in the problem of determining a single closed- or open-loop controller, which applies to a particular system over an infinite number of repeated trials, or to steer a family of finite-dimensional dynamical systems. As discussed by Brockett, this approach represents a new control framework that is able to address a number of issues as uncertainty in initial conditions and the trade-off mentioned above.

The purpose of this paper is to contribute to the development of Brockett's Liouville-based ensemble control with a theoretical investigation of a class of Liouville optimal control problems with unbounded coefficients and cost functionals that are formulated in terms of the density and of different control costs. Those control costs model various requirements on the control functions, including control constraints. To the best of our knowledge, there is no similar investigation available in the literature yet. We will present more details about the formulation of our problem in Subsections 2.1 and 2.2. Thereafter, we give a detailed overview of the main results of the paper in Subsection 2.3.

For the time being, let us point out that our work focuses on two main issues. First of all, we study existence and regularity of solutions to a Liouville initial-value problem with a so-called control-in-thecoefficients, which is related to a differential model with a linear and bilinear control mechanism. In general, this latter model has the following structure: $\dot{x}(t)=\bar{a}+\bar{b} u(t)+\bar{c} v(t) x(t)$, where $\bar{a}, \bar{b}, \bar{c}$ are given functions and $u$ and $v$ denote the linear and bilinear controls, respectively. As we show below, the right-hand side of this equation corresponds to the drift coefficient of the Liouville equation $a(t, x ; u, v)=$ $\bar{a}+\bar{b} u(t)+\bar{c} v(t) x$, which is unbounded for $x \in \mathbb{R}^{d}$. For this reason, we consider the less investigated problem of existence and regularity of solutions to a Liouville equation having a drift which has at most a linear growth in space at infinity: we give a self-contained presentation of the well-posedness theory (due to DiPerna and Lions, see [18]; see also [17]) in Sobolev spaces $H^{m}$, for drifts $a \in L^{1}\left([0, T] ; C^{m+1}\left(\mathbb{R}^{d}\right)\right)$ such that $\nabla a \in L^{1}\left([0, T] ; C_{b}^{m}\left(\mathbb{R}^{d}\right)\right)$. In addition, in order to address continuity and Fréchet differentiability properties of the control-to-state map $G:(u, v) \mapsto \rho$, we extend existence and uniqueness results to the framework of appropriately weighted Sobolev spaces; to the best of our knowledge, such a well-posedness theory, which is natural in our context, seems to be new in the literature.

These results are essential to study ensemble optimal control problems which include Brockett's cost functionals with density of ensemble of trajectories, a $L^{2}$ cost of the control and a $H^{1}$ cost which promotes minimum attention controls. Moreover, we extend this framework including a $L^{1}$ cost of the control, which should promote minimum action control during the time evolution, and include box control-constraints. For the resulting Liouville optimal control problem, we discuss existence and uniqueness of optimal controls and their characterization by optimality systems. Specifically, we prove existence of controls in our general framework with weighted Sobolev spaces; in the case where only $L^{2}$ costs of the control are considered, we prove also their uniqueness. For the characterisation of the optimal solution in our general setting, we use the (sub-)differentiability properties of the cost functional and of the control-to-state map and prove the first-order optimality conditions.

We conclude this brief introduction by remarking that, according to Brockett, in the cost functionals of ensemble controls the density-based tracking and terminal observation terms (i.e. respectively $\theta$ and $\varphi$ in (2.10) below) are formulated by introducing "attracting" potentials, which are quadratic (as discussed in $[10,9])$. Capturing such a framework represents an additional challenge for the theoretical investigation
of these control problems: as shown in the final part of this paper, we are able to address and solve this issue within our weighted Sobolev spaces framework.

The rest of the paper unfolds as follows. In Section 2, we present the problem formulation and an overview of our main results. In Section 3, we state and prove our theoretical results on well-posedness of Liouville and transport equations with unbounded drifts in weighted Sobolev spaces. Section 4 is devoted to the study of the Liouville control-to-state map $G$. Section 5 is devoted to the analysis of the Liouville optimal control problem: we establish existence of optimal controls, their characterisation as solutions of a first-order optimality system, and by use of the previous characterisation, their uniqueness in specific cases. An appendix completes this work, where we postpone the proof of some technical results.

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## Notation

In this section, we present our notation that we use throughout the paper.
Given a domain $\Omega \subset \mathbb{R}^{d}$, the symbol $C_{c}^{\infty}(\Omega)$ denotes the space of infinitely often differentiable functions with compact support in $\Omega$. Given $k \in \mathbb{N}$, we denote by $C^{k}(\Omega)$ the space of all $k$-times continuously differentiable functions defined on $\Omega$, and by $C_{b}^{k}(\Omega)$ the subspace of $C^{k}(\Omega)$ formed by functions which are uniformly bounded together with all their derivatives up to the order $k$. We equip $C_{b}^{k}(\Omega)$ with the $W^{k, \infty}$-norm as follows

$$
\|v\|_{C_{b}^{k}}:=\sum_{|\alpha| \leq k}\left\|D^{\alpha} v\right\|_{L^{\infty}}
$$

For $\alpha \in] 0,1]$, we denote with $C^{0, \alpha}(\Omega)$ the classical Hölder space (Lipschitz space if $\alpha=1$ ), endowed with the norm

$$
\|\Phi\|_{C^{0, \alpha}}:=\sup _{x \in \Omega}|\Phi(x)|+\sup _{\substack{x, y \in \Omega \\ 0<|x-y| \leq 1}} \frac{|\Phi(x)-\Phi(y)|}{|x-y|^{\alpha}}
$$

In particular, $C^{0,1}(\Omega) \equiv W^{1, \infty}(\Omega)$.
For $k \in \mathbb{N}$ and $1 \leq p \leq+\infty$, we denote with $W^{k, p}(\Omega)$ the usual Sobolev space of $L^{p}$ functions with all the derivatives up to the order $k$ in $L^{p}$; we also set $H^{k}(\Omega):=W^{k, 2}(\Omega)$. For $1 \leq p<+\infty$, let $W^{-k, p}(\Omega)$ denote the dual space of $W^{k, p}(\Omega)$. For any $p \in[1,+\infty]$, the space $L_{l o c}^{p}(\Omega)$ is the set formed by all functions which belong to $L^{p}\left(\Omega_{0}\right)$, for any compact subset $\Omega_{0}$ of $\Omega$.

Furthermore, we make use of the so-called Bochner spaces. Given two Banach spaces $X$ and $Y$ and a fixed time $T>0$, we define

$$
X_{T}(Y):=X([0, T] ; Y), \quad \text { with } \quad\|u\|_{X_{T}(Y)}:=\int_{0}^{T}\|u(t)\|_{Y} d t
$$

Given a Banach space $X$ and a sequence $\left(\Phi_{n}\right)_{n}$, we use the notation $\left(\Phi_{n}\right)_{n} \subset X$ meaning that $\Phi_{n} \in X$ for all $n \in \mathbb{N}$ and that this sequence is uniformly bounded in $X$ : there exists some constant $M>0$ such that $\left\|\Phi_{n}\right\|_{X} \leq M \forall n \in \mathbb{N}$.

Given two Banach spaces $X$ and $Y$, the space $X \cap Y$, endowed with the norm $\|\cdot\|_{X \cap Y}:=\|\cdot\|_{X}+\|\cdot\|_{Y}$, is still a Banach space.

For every $p \in[1,+\infty]$, we use the notation $\mathbb{L}_{T}^{p}\left(\mathbb{R}^{d}\right):=L_{T}^{p}\left(\mathbb{R}^{d}\right) \times L_{T}^{p}\left(\mathbb{R}^{d}\right)$. Analogously, $\mathbb{H}_{T}^{1}\left(\mathbb{R}^{d}\right):=$ $H_{T}^{1}\left(\mathbb{R}^{d}\right) \times H_{T}^{1}\left(\mathbb{R}^{d}\right)$. In addition, given two vectors $u$ and $v$ in $\mathbb{R}^{d}$, we write $u \leq v$ if the inequality is satisfied component by component by the two vectors: namely, $u^{i} \leq v^{i}$ for all $1 \leq i \leq d$.

Given two operators $A$ and $B$, we use the standard notation $[A, B]$ for their commutator: $[A, B]:=$ $A B-B A$.

## 2 Problem formulation and overview of the results

In this section, we present the problem formulation and an overview of our results. In Section 2.1, we discuss the Liouville equation and the control mechanism. In Section 2.2, we formulate our Liouville ensemble optimal control problem. Section 2.3 illustrates our main results.

### 2.1 The Liouville equation and a control mechanism

The Liouville model represents the fundamental building block of many important equations in fluid mechanics and related fields, as e.g. the Boltzmann equation, the Fokker-Planck equation, the Vlasov equation. It arises in diverse areas of sciences as biology, finances, mechanics, and physics; see e.g. $[12,13,15,24,19,16,25]$.

Central to our discussion is the relation of the Liouville equation to a given ODE evolution model. To illustrate this fact, consider a smooth vector field $a(t, x)$ over $\mathbb{R}^{d}$, where $(t, x) \in[0, T] \times \mathbb{R}^{d}$, for some time $T>0$. We refer to $a$ as the drift function. It is well-known that, if a scalar function $\rho$, defined on $[0, T] \times \mathbb{R}^{d}$, satisfies the Liouville equation

$$
\begin{equation*}
\partial_{t} \rho(t, x)+\operatorname{div}(a(t, x) \rho(t, x))=0 \tag{2.1}
\end{equation*}
$$

with some (say) smooth initial datum $\rho_{\mid t=0}=\rho_{0}$, then we can represent $\rho$ by the formula

$$
\rho(t, x)=\frac{1}{\operatorname{det} J(t, x)} \rho_{0}\left(\psi_{t}^{-1}(x)\right)
$$

where $\psi_{t}(x)=\psi(t, x)$ denotes the flow map associated to $a, J(t, x)=\nabla_{x} \psi_{t}(x)$ is its Jacobian matrix, and $\psi_{t}^{-1}(x)$ means the inverse with respect to the space variable, at $t$ fixed. By definition of flow map, $\psi$ verifies the following system of ODEs

$$
\begin{equation*}
\partial_{t} \psi(t, x)=a(t, \psi(t, x)), \quad \psi(0, x)=x \tag{2.2}
\end{equation*}
$$

Notice that we can equivalently write this Cauchy problem as $\dot{y}(t)=a(t, y(t)), y(0)=x$, to point out that the (independent) space variable in the Liouville equation corresponds to the (dependent) state variable of the related dynamical system. In the following, we shall use the same symbol for both cases.

Problem (2.1) models the evolution of the ensemble of trajectories of (2.2) for a density distribution of initial conditions given by $\rho_{0}$. Thus, one possible interpretation of (2.1) is that the function $\rho(t, x)$ represents the probability density function of finding the system in $x$ at time $t$, assuming that $\rho_{0}$ prescribes the initial probability density for $\psi_{t}^{-1}(x)$. This interpretation in the space of probability appears in the realm of Kolmogorov-Fokker-Planck and Kramers equations related to stochastic processes, where the Liouville equation corresponds to the first-order differential term of these equations; see e.g. [25]. On the other hand, if $\rho$ represents the material density of non-interacting particles, then (2.1) models the evolution of this density, see e.g. [20]

Notice that the former point of view is the predominant one in Brockett's consideration and we adopt it in the rest of the paper. Therefore, it is natural to assume an initial condition for the Liouville model such that $\rho_{0} \geq 0$, together with the normalization $\int_{\mathbb{R}^{d}} \rho_{0}(x) d x=1$. Consequently, by equation (2.1), it holds for all times $t \geq 0$ that

$$
\rho(t, x) \geq 0 \quad \text { and } \quad \int_{\mathbb{R}^{d}} \rho(t, x) d x=\int_{\mathbb{R}^{d}} \rho_{0}(x) d x=1 .
$$

The first property can be proved by the vanishing viscosity method and the maximum principle, see e.g. [20, 21], or just by solving (2.1) along characteristics; the second property follows from a simple application of the divergence theorem. However, notice that most of our results do not require the latter two assumptions on $\rho_{0}$.

Next, let us discuss our control mechanism. We remark that the focus of ensemble control is the development of a control strategy for the differential model (2.2) augmented with a control mechanism, as follows

$$
\begin{equation*}
\dot{x}=a(t, x ; u), \tag{2.3}
\end{equation*}
$$

where $u$ denotes the control function.
We refer to $[10,9]$ for a discussion on the choice of $u$ as a function of time only, which corresponds to a so-called open-loop control, or as a function of time and of the state variable, which may represent a feedback law. In this paper, while we consider our controlled Liouville model in a general setting that accommodates both choices, we focus our attention on open-loop optimal control problems: this point of view is motivated by the fact that the most used control mechanisms for (2.3) are the linear and bilinear ones, as follows,

$$
\begin{equation*}
a(t, x ; u)=a_{0}(t, x)+u_{1}(t)+x \circ u_{2}(t), \tag{2.4}
\end{equation*}
$$

where $a_{0}$ is a smooth vector field and $u=\left(u_{1}, u_{2}\right)$ is the control, which, for the scope of the present discussion, we assume to be smooth. The control $u_{1}$ represents a linear control mechanism and $u_{2}$ multiplying the state variable $x$ represents the bilinear control term. Both functions $u_{1}$ and $u_{2}$ are defined on the time interval $[0, T]$ with values in $\mathbb{R}^{d}$. Further, with $\circ: \mathbb{R}^{d} \times \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ we denote the Hadamard product of two vectors, i.e. the multiplication component by component.

Now, notice that corresponding to the controlled evolution model (2.3), we have the following controlled Liouville equation

$$
\begin{equation*}
\partial_{t} \rho(t, x)+\operatorname{div}(a(t, x ; u) \rho(t, x))=0 \tag{2.5}
\end{equation*}
$$

In this framework and in the simple case $a_{0}=0$ and $d=1$, we can give a simple interpretation of the role of $u_{1}$ and $u_{2}$ : The control $u_{1}$ represents the driving force of the mean value of the density; $u_{2}$ determines the evolution of the variance of the density. In fact, let $\rho_{0}$ represents a normalized Gaussian centred at $x_{0}$ with variance $v_{0}$ at time $t=0$, and define the following average operator

$$
\mathbb{E}[g](t)=\int_{\mathbb{R}} g(x) \rho(t, x) d x
$$

In particular, we have the mean $m(t)=\mathbb{E}[x](t)$ and the variance $v(t)=\mathbb{E}\left[(x-m(\cdot))^{2}\right](t)$. Then, by taking the average of our controlled ODE model, we obtain the following equations:

$$
\begin{aligned}
\dot{m}(t) & =u_{1}(t)+m(t) u_{2}(t), & m(0) & =x_{0} \\
\dot{v}(t) & =2 v(t) u_{2}(t), & v(0) & =v_{0}
\end{aligned}
$$

However, because of the limiting assumptions, this construction does not provide the degree of generality that the Liouville framework offers, as the latter allows to accommodate any chosen drift with any
control mechanism. The Liouville framework is also very attractive because it allows to consider bimodal distributions of the initial density function.

Finally, we remark that, for the characterization of the solution to our Liouville optimal control problems, we shall deal with (2.5) and with an adjoint Liouville problem, namely a transport problem, given by

$$
\begin{equation*}
\partial_{t} q(t, x)+a(t, x ; u) \cdot \nabla q(t, x)=g(t, x), \quad \text { with } \quad q_{\mid t=0}=q_{0}, \tag{2.6}
\end{equation*}
$$

where $g$ and $q_{0}$ depend on the optimization data.

### 2.2 Formulation of ensemble optimal control problems

In order to discuss the formulation of ensemble control, consider the following ODE optimal control problem

$$
\begin{align*}
& \min j(x, u):=\int_{0}^{T}(\theta(x(t))+\kappa(u(t))) d t+\varphi(x(T))  \tag{2.7}\\
& \text { s.t. } \quad \dot{x}(t)=a(t, x(t) ; u(t)), \quad x(0)=x_{0} \tag{2.8}
\end{align*}
$$

where "s.t." stands for "subject to". Here above, $\theta, \kappa$ and $\varphi$ are usually taken to be continuous convex functions of their arguments; we will better specify their properties later on in the present section.

Further, assume that the optimal control function $u$ is sought in the following set of admissible controls:

$$
\begin{equation*}
U_{a d}:=\left\{u \in \mathbb{L}_{T}^{\infty}\left(\mathbb{R}^{d}\right) \mid \quad u^{a} \leq u(t) \leq u^{b} \quad \text { for a.e. } t \in[0, T]\right\} \tag{2.9}
\end{equation*}
$$

In particular, in the case of (2.4), we have two box constraints $u^{a}=\left(u_{1}^{a}, u_{2}^{a}\right)$ and $u^{b}=\left(u_{1}^{b}, u_{2}^{b}\right)$, where $u_{j}^{a}<u_{j}^{b}, j=1,2$, are given vectors in $\mathbb{R}^{d}$. Clearly, the optimal control function $u$ that solves (2.7)-(2.8) with $u \in U_{a d}$ depends on the initial condition $x_{0}$, which is fixed, and it represents a control strategy that is determined once and for all times for the given $x_{0}$ and the given optimization setting. Therefore any uncertainty on the initial condition is not taken into account in the formulation (2.7)-(2.8) and, hence, the resulting control is not robust. On the other hand, a closed loop control, say, $u=u(t, x)$, would appropriately control the system based on the actual state of the system; however, as pointed out in [9], the cost of implementing such a control mechanism is often prohibitive and may be not justified by real applications.

For this reason, with the purpose to strike a balance between the desired performance of the system and the cost of implementing an effective control, the ensemble control strategy considers a density of initial conditions, and therefore ensemble of trajectories. In this way, it aims at achieving robustness, while choosing control costs which promote controls allowing for easier implementation (see below). Thus, one is led to the formulation of the following ensemble optimal control problem

$$
\begin{align*}
& \min _{u \in U_{a d}} J(\rho, u):=\int_{0}^{T} \int_{\mathbb{R}^{d}} \theta(x) \rho(x, t) d x d t+\int_{\mathbb{R}^{d}} \varphi(x) \rho(x, T) d x+\int_{0}^{T} \kappa(u(t)) d t  \tag{2.10}\\
& \text { s.t. } \quad \partial_{t} \rho+\operatorname{div}(a(t, x ; u) \rho)=0, \quad \rho_{\mid t=0}=\rho_{0} . \tag{2.11}
\end{align*}
$$

This problem is defined on the space-time cylinder $\mathbb{R}^{d} \times[0, T]$, for some $T>0$ fixed. In this formulation, the initial density $\rho_{0}$ represents the probability distribution of the initial condition $x_{0}$ in (2.7)-(2.8), and thus it models the known uncertainty on the initial data.

Next, we discuss some specific choices of the optimization components in (2.7)-(2.8), and correspondingly in (2.10)-(2.11).

For example, if $x=0$ is a critical point for (2.8), which requires $a(t, 0 ; u)=0$, then the choice $\theta(x)=x^{2}$ appears standard for stabilization purposes. Usually, in this context, the so-called $L^{2}$ cost of the control
is considered, which corresponds to the choice $k(u)=\gamma u^{2}$, where $\gamma>0$ is the weight of the cost of the control. On the other hand, if the purpose of the control in (2.7)-(2.8) is to track a desired and even non-attainable trajectory $x_{d} \in L^{2}\left(0, T ; \mathbb{R}^{d}\right)$, and to come close to a given final configuration $x_{T} \in \mathbb{R}^{d}$ at the final time (possibly with $x_{d}(T) \neq x_{T}$ ), then a natural choice appears to be $\theta(x(t))=\alpha\left(x(t)-x_{d}(t)\right)^{2}$ and $\varphi(x(T))=\beta\left(x(T)-x_{T}\right)^{2}$, with appropriately chosen weights $\alpha, \beta>0$. However, in the context of ensemble control, as in (2.10), the choice of $\theta$ and $\varphi$ as convex functions is problematic because of integrability issues. On the other hand, we remark that the role of this function is to define an attracting potential, that is, to define a well centred at a minimum point such that the minus gradient of the potential is directed towards this minimum. For this purpose, a possible choice is $\theta(x)=1-\exp \left(-x^{2}\right)$, with the minimum at $x=0$. In our analysis, we are able to address both cases in the framework of weighted Sobolev spaces: the case of attracting potentials $\theta$ and $\varphi$ which are both $L^{2}$ integrable, and the case of $\theta$ and $\varphi$ which are quadratic functions. Notice that, in any case, the modelling choice for (2.7)-(2.8) translates without changes to (2.10)-(2.11).

As discussed in $[10,9,8]$, the choice of the cost function $\kappa$ should be such that the effort of implementing the control strategy is as small as possible. In this sense, the cost of implementing a slowly varying control function, and (we add) a control that does not act for all times, should be smaller than that corresponding to a control having large variations. From this perspective, a constant input that controls the system is the cheapest choice, and the next possible choice is a control that slowly changes in time. This requirement leads naturally to a cost of the form

$$
\nu \int_{0}^{T}\left(\frac{d u}{d t}(t)\right)^{2} d t
$$

where $\nu \geq 0$ is a non-negative weight. In fact, as $\nu$ is taken larger, the resulting optimal control will have smaller values of its time derivative, that is, a slowly varying control, which is called "minimum attention control" in [8].

On the other hand, the most common quantification of the cost of a control in the context of (2.7)(2.8) is the $L^{2}$-cost already mentioned above. More recently, there has been a surge of interest in $L^{1}$-costs, originating from signal reconstruction and magnetic resonance imaging [11]. This cost is given by

$$
\delta \int_{0}^{T}|u(t)| d t
$$

where $\delta \geq 0$. The effect of this cost is that it promotes sparsity of the control function, in the sense that, as $\delta>0$ is increased, the $u$ resulting from the minimisation procedure will be zero on open intervals in $] 0, T[$, and these intervals become larger and eventually cover all of $] 0, T[$ as $\delta \rightarrow+\infty$. In the present papaer, we introduce the $L^{1}$-cost in the context of ensemble control and call the resulting sparse control a "minimum action control".

All together, we specify the term $\int_{0}^{T} \kappa(u(t)) d t$ in (2.7) and in (2.10) as follows

$$
\begin{equation*}
\kappa(u(t)):=\frac{\gamma}{2}(u(t))^{2}+\delta|u(t)|+\frac{\nu}{2}\left(\frac{d u}{d t}(t)\right)^{2} \tag{2.12}
\end{equation*}
$$

where $\gamma+\delta+\nu>0$ and the factor $1 / 2$ is chosen for convenience of later calculations.
Notice that different choices of the value of the weights $\gamma, \delta, \nu$ will result in different features of the resulting optimal control function. We investigate the properties of $u$ resulting from these different choices in Section 5.1 and its characterisation in Section 5.2.

### 2.3 Overview of the main results

In this section, we give a short overview of the main results contained in this paper.
The first step in our analysis consists in investigating the well-posedness of the PDE under consideration and its adjoint, namely continuity and transport equations with unbounded drift function, which presents the structure (2.4). We perform such a study in Section 3 below.

We point out that we do not strive for minimal regularity hypotheses on the drift vector field $a$, and frame our work within a setting that can be considered (for bounded coefficients) classical; see e.g. Chapter 3 of [4]. Nonetheless, to the best of our knowledge very few references deal with the case of unbounded drifts, with at most linear growth at infinity, for which the well-posedness theory goes back to the conrnerstone paper [18] (see also [1, 2, 3, 17] and references therein). For this reason, for clarity, we give a self-contained presentation of the well-posedness theory for the Liouville and transport equations with unbounded drifts in classical Sobolev spaces: this part is just a review, and contains no new results.

In addition, for reasons which will become clear later, in Section 3.2 we investigate well-posedness of the Liouville and transport equations in weighted spaces $H_{k}^{m}$, see Definition 3.1 below. Roughly, a tempered distribution $\rho \in H^{m}$ belongs to $H_{k}^{m}$ if it enjoys further integrability at $+\infty$, namely if $\rho$ and all its derivatives up to order $m$ belong to the measurable space $\left(L^{2}\left(\mathbb{R}^{d}\right),(1+|x|)^{k} d x\right)$. On the one hand, existence, uniqueness and regularity properties are derived in this context by standard arguments: the key of the analysis reduces to show suitable a priori estimates on the solutions in weighted norms. On the other hand, the results of this part seem to be completely new in the literature; in addition, this framework reveals to be well-adapted to the investigation of our ensemble optimal control problem, see more details below.

We believe that well-posedness in weighted spaces can be adapted with no special problems to $L^{p}$ based spaces, for any $1 \leq p<+\infty$. Then, ensemble optimal control problems with different integrability conditions can be considered as well. In addition, very likely the $H_{k}^{m}$ theory should generalise also to more general hyperbolic systems which are symmetrizable in the sense of Friedrichs (see e.g. [6, 23]). Consequently, we expect to be able to deal with optimal control problems related to those systems by similar methods as the ones presented in this paper. However, extensions of the present work to both directions go beyond the scope of our paper, and we leave them for further studies.

After establishing well-posedness of the Liouville equation in a suitable framework, we pass to investigating the optimal control problem related to that PDE. For this, we follow a standard scheme.

First of all, in Section 4 we define and study the Liouville control-to-state map $G$, namely the map which associates to any control state $u$ the unique solution $\rho=G(u)$ to the corresponding Liouville equation. A fundamental issue in this part is to prove Fréchet differentiability (in a suitable topology) of $G$. Our method to get that property (see Section 4.2) consists in applying the definition of Fréchet differentibility, and showing the convergence of the limit under consideration in the strong $L_{T}^{\infty}\left(L^{2}\right)$ topology, by performing fine stability estimates on the Liouville equation. Now, dealing with the growth in space of our drift function at $+\infty$ requires the use of weighted norms and weighted spaces in order to carry out those estimates; moreover, due to the hyperbolicity of transport and continuity equations, a loss of regularity occurs, which requires to consider both higher smoothness and higher integrability on the initial data (namely, both $m \geq 2$ and $k \geq 2$ ). Fréchet differentiability reveals to be a fundamental property for the subsequent analysis.

In Section 5, we complete the investigation of the ensemble optimal control problem. First of all, we prove (see Theorem 5.1) existence of optimal controls, in the simpler case of attracting potentials which are moreover $L^{2}$. Then, we characterise these optimal controls as solutions of a first-order optimality system, whose interpretation in terms of the Fréchet differential of the reduced functional $\widehat{J}(u):=J(u, G(u))$ exploits the Fréchet differentiability of $G$, proved above. We also notice that the differentiability properties of $J$ change radically depending on the choice of the optimization weights. For instance, if $\gamma>0$ and
$\delta=\nu=0$, then the optimization space is $L^{2}(0, T)$ and we have Fréchet differentiability of the cost functional. This is the "standard" case. If instead $\delta>0$, then we have a semi-smooth optimal control problem and we have to resort to the use of sub-differentials. Finally, if $\nu>0$, then $H^{1}(0, T)$ is the appropriate control space, and the optimality condition accounts for this fact. If all weights are positive and with control constraints, we have an optimal control problem whose structure (to the best of our knowledge) has never been investigated in PDE optimization. For this general case, we prove (see Theorem 5.2) existence of Lagrange multipliers, and derive the optimality system.

Finally, in Section 5.3 we address the uniqueness of optimal ensemble controls, in the special case $\gamma>0$ and $\delta=\nu=0$. This part of the analysis exploits in a fundamental way the optimality system previously derived, and the characterisation of optimal controls as solutions to it. In Theorem 5.3 we prove uniqueness of the optimal control in the control-unconstrained case; for this, we neeed some further assumptions on the attracting potentials, namely we suppose $\theta$ and $\varphi$ to belong both to $H_{1}^{1}\left(\mathbb{R}^{d}\right)$. Notice that, to the best of our knowledge, our uniqueness result has no counterpart in the literature of PDE control problems. In the case that control constraints are present, in Theorem 5.4 we prove uniqueness of optimal controls, provided a smallness condition is satisfied. Such a condition requires the time $T$ and the size of the data $\rho, g, \theta$ and $\varphi$ to be small enough, or the coefficient $\gamma$ to be sufficiently large.

We recall that all the results of Section 5 that we have summarised above are obtained assuming $\theta, \varphi \in L^{2}\left(\mathbb{R}^{d}\right)$. However, in Section 5.4 we show the necessary modifications to be performed in our arguments in order to address the case $\theta(x)=|x|^{2}$ and $\varphi(x)=|x|^{2}$. In particular, such modifications require to solve the transport equation in weighted Sobolev spaces of negative index, see Lemma 5.1 below.

## 3 Theory of Liouville and transport equations with unbounded drifts

In this section, we present results concerning the well-posedness theory of Liouville and transport equations in the class of Sobolev spaces. In view of formula (2.4), we are especially interested in the case when the drift function $a$ may be unbouded, but has at most a linear growth at infinity.

In Subsection 3.1, we review the well-posedness theory in classical $H^{m}$ spaces, for $m \in \mathbb{N}$ (for simplicity); we will give a self-contained presentation in this framework, and refer to e.g. [18] and [17] for details and more general results. Afterwards in Section 3.2, motivated by the study of our optimal control problem, we extend these results to weighted Sobolev spaces. Notice that, although the statements and their proofs follow the main lines of the classical framework, to the best of our knowledge they are novel.

### 3.1 Classical theory of Liouville and trasport equations

We start our discussion by considering the Liouville equation. Notice that our statements can be repeated in a very similar way (with just slight modifications) also for the adjoint Liouville problem, namely the transport equation: we treat this case in Paragraph 3.1.2, without giving details.

### 3.1.1 Liouville equations in classical Sobolev spaces

Consider the following Liouville initial-value problem

$$
\left\{\begin{array}{lll}
\partial_{t} \rho+\operatorname{div}(a(t, x) \rho)=g(t, x) & \text { in } & {[0, T] \times \mathbb{R}^{d}}  \tag{3.1}\\
\rho_{\mid t=0}=\rho_{0} & \text { on } & \mathbb{R}^{d}
\end{array}\right.
$$

Whenever attempting at solving equation (3.1), we have in mind its weak formulation. Namely, for all $\phi \in C_{c}^{\infty}\left(\mathbb{R}^{d} \times[0, T[)\right.$, we want to verify the following equality

$$
\begin{equation*}
-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho \partial_{t} \phi d x d t-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho a \cdot \nabla \phi d x d t=\int_{0}^{T} \int_{\mathbb{R}^{d}} g \phi d x d t+\int_{\mathbb{R}^{d}} \rho_{0} \phi(0) d x . \tag{3.2}
\end{equation*}
$$

The theory for this equation is classical, at least in the case of a bounded drift function $a$. The following well-posedness result is adapted to our needs from Theorem 3.19 in [4].

Theorem 3.1. Let us fix $T>0$ and $m \in \mathbb{N}$, and let $a \in L^{1}\left([0, T] ; C_{b}^{m+1}\left(\mathbb{R}^{d}\right)\right), \rho_{0} \in H^{m}\left(\mathbb{R}^{d}\right)$ and $g \in L^{1}\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$.

Then there exists a unique weak solution $\rho$ to (3.1), with $\rho \in C\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$. Moreover, there exists a "universal" constant $C>0$, independent of $\rho_{0}, a, g, \rho$ and $T$, such that the following estimate holds true for any $t \in[0, T]$ :

$$
\|\rho(t)\|_{H^{m}} \leq C\left(\left\|\rho_{0}\right\|_{H^{m}}+\int_{0}^{t}\|g(\tau)\|_{H^{m}} d \tau\right) \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{m}} d \tau\right) .
$$

Remark 3.1. In the case $m=0$, one can replace $\|\nabla a\|_{C_{b}^{0}}$ with $\|\operatorname{div} a\|_{L^{\infty}}$ inside the integral in the exponential term.

Motivated by the study of our optimal control problem, see Section 2.1 and especially Definition (2.4), we are rather interested in the case when $a$ may be unbounded, with at most a linear growth at infinity. More precisely, given $m \in \mathbb{N}$, we assume

$$
\left\{\begin{array}{l}
g \in L^{1}\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right) \quad \text { and } \quad \rho_{0} \in H^{m}\left(\mathbb{R}^{d}\right)  \tag{3.3}\\
a \in L^{1}\left([0, T] ; C^{m+1}\left(\mathbb{R}^{d}\right)\right), \quad \text { with } \quad \nabla a \in L^{1}\left([0, T] ; C_{b}^{m}\left(\mathbb{R}^{d}\right)\right)
\end{array}\right.
$$

Remark 3.2. Notice that hypotheses (3.3) imply, in particular, that $a(\cdot, \cdot)$ has at most linear growth in space at infinity: for almost every $(t, x) \in[0, T] \times \mathbb{R}^{d}$, one has

$$
|a(t, x)| \leq C c(t)(1+|x|), \quad \text { for } \quad c=\|\nabla a\|_{L^{\infty}} \in L^{1}([0, T]) .
$$

The condition of at most linear growth at infinity can be proved to be somehow sharp for well-posedness, see e.g. [18], [17] and the references therein.

The main result of this section is the following statement, proved by DiPerna and Lions in [18] (see also [17]). However, we give here a self-contained presentation of its proof.

Theorem 3.2. Let $T>0$ and $m \in \mathbb{N}$ fixed, and let $a, \rho_{0}$ and $g$ satisfy hypotheses (3.3).
Then there exists a unique solution $\rho \in C\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$ to problem (3.1). Moreover, there exists $a$ "universal" constant $C>0$, independent of $\rho_{0}, a, g, \rho$ and $T$, such that the following estimate holds true for any $t \in[0, T]$ :

$$
\begin{equation*}
\|\rho(t)\|_{H^{m}} \leq C\left(\left\|\rho_{0}\right\|_{H^{m}}+\int_{0}^{t}\|g(\tau)\|_{H^{m}} d \tau\right) \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{m}} d \tau\right) . \tag{3.4}
\end{equation*}
$$

We notice that, also in this case, Remark 3.1 applies.
The rest of this paragraph is devoted to the proof of Theorem 3.2. We proceed by truncation, and approximate our problem with a family of Liouville equations with bounded coefficients, to which the classical theory (see Theorem 3.1 above) applies. Then, we pass to the limit in the approximation parameter, proving convergence to a solution of the original equation. We conclude by discussing time regularity and uniqueness issues.

Existence. The first step is to construct a suitable truncation of the drift function. For this purpose, let us introduce a smooth cut-off function $\chi \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ such that $\chi$ is radially decreasing, $\chi(x)=1$ for $|x| \leq 1$ and $\chi(x)=0$ for $|x| \geq 2$. For all real $M>0$, we define

$$
\begin{equation*}
a_{M}(t, x):=\chi\left(\frac{x}{M}\right) a(t, x) . \tag{3.5}
\end{equation*}
$$

Notice that, by assumptions (3.3), we immediately get that $a_{M} \in L_{T}^{1}\left(C_{b}^{m+1}\right)$ for all $M>0$. Moreover, in view of Remark 3.2, an easy computation shows that

$$
\begin{equation*}
\left(\nabla a_{M}\right)_{M} \subset L_{T}^{1}\left(C_{b}^{m}\right), \quad \text { with } \quad\left\|\nabla a_{M}\right\|_{L_{T}^{1}\left(L^{\infty}\right)} \leq C \tag{3.6}
\end{equation*}
$$

for a suitable constant $C>0$ independent of $M$. Indeed, denoting by $\mathbb{1}_{A}$ the characteristic function of a set $A \subset \mathbb{R}^{d}$ and by $B(x, R)$ the ball in $\mathbb{R}^{d}$ of center $x$ and radius $R>0$, we can compute

$$
\begin{aligned}
\left\|\nabla a_{M}\right\|_{L^{\infty}} & =\left\|\frac{1}{M} \nabla \chi\left(\frac{x}{M}\right) a+\chi\left(\frac{x}{M}\right) \nabla a\right\|_{L^{\infty}} \\
& \leq C \frac{1}{M}\left\|a \mathbb{1}_{B(0,2 M)}\right\|_{L^{\infty}}+\|\nabla a\|_{L^{\infty}} \leq C .
\end{aligned}
$$

The bounds for higher order derivatives follow the same lines, after noticing that, at each order of differentiation, we gain a factor $1 / M$ in front of $a$.

At this point, for each fixed $M>0$, we can consider the truncated problem

$$
\left\{\begin{array}{l}
\partial_{t} \rho+\operatorname{div}\left(a_{M} \rho\right)=g  \tag{3.7}\\
\rho_{\mid t=0}=\rho_{0}
\end{array}\right.
$$

which possesses a unique weak solution $\rho_{M} \in C\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$, by virtue of Theorem 3.1. Moreover, each $\rho_{M}$ satisfies the energy estimate (3.4), up to replacing $a$ by $a_{M}$. Thus, we have

$$
\begin{equation*}
\left\|\rho_{M}(t)\right\|_{H^{m}} \leq C\left(\left\|\rho_{0}\right\|_{H^{m}}+\int_{0}^{t}\|g(\tau)\|_{H^{m}} d \tau\right) \exp \left(C \int_{0}^{t}\left\|\nabla a_{M}(\tau)\right\|_{C_{b}^{m}} d \tau\right) \tag{3.8}
\end{equation*}
$$

Thanks to property (3.6), we deduce the uniform bounds

$$
\begin{equation*}
\left(\rho_{M}\right)_{M} \subset L^{\infty}\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right) \tag{3.9}
\end{equation*}
$$

As a consequence of (3.9), we obtain the existence of a $\rho \in L_{T}^{\infty}\left(H^{m}\right)$ such that, up to the extraction of a subsequence, one has

$$
\rho_{M} \stackrel{*}{\rightharpoonup} \rho \quad \text { in } \quad L_{T}^{\infty}\left(H^{m}\right)
$$

Our next goal is to show that $\rho$ actually solves problem (3.1) in the weak form, see equation (3.2). For this purpose, we need to pass to the limit, for $M \rightarrow+\infty$, in the weak formulation of equation (3.7). Of course, it is enough to prove the convergence in the case of minimal regularity, namely for $m=0$. Thus, we restrict to this case in the next argument.

We start by recalling that $\rho_{M}$ is a weak solution to (3.7): given any $\phi \in C_{c}^{\infty}\left(\mathbb{R}^{d} \times[0, T[)\right.$, we have

$$
\begin{equation*}
-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho_{M} \partial_{t} \phi d x d t-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho_{M} a_{M} \cdot \nabla \phi d x d t=\int_{0}^{T} \int_{\mathbb{R}^{d}} g \phi d x d t+\int_{\mathbb{R}^{d}} \rho_{0} \phi(0) d x . \tag{3.10}
\end{equation*}
$$

The only term which presents some difficulties is the term $\rho_{M} a_{M}$, and thus we focus on it. We start by proving the following lemma.

Lemma 3.1. For all compact set $K \subset \mathbb{R}^{d}$, it holds

$$
\left\|a_{M}-a\right\|_{L_{T}^{1}\left(L^{\infty}(K)\right)} \longrightarrow 0 \quad \text { as } \quad M \rightarrow+\infty
$$

Proof. Let $K \subset \mathbb{R}^{d}$ be a compact set, and let $R>0$ such that $K \subset B(0, R)$. The claim of the lemma follows then by noticing that, by definitions, for all $M \geq R+1$ one has $a_{M}(t) \equiv a(t)$ over $K$, for almost every $t \in[0, T]$.

Let now $K$ be the support of $\phi$, where $\phi$ is the test function appearing in (3.10). Thanks to uniform bounds, to the strong convergence of $a_{M}$ to $a$ in $L_{T}^{1}\left(L^{\infty}(K)\right)$ (given by Lemma 3.1) and the weak-* convergence of $\rho_{M}$ to $\rho$ in $L_{T}^{\infty}\left(L^{2}\right)$, we finally deduce that $\left(\rho_{M} a_{M}\right)_{M}$ is uniformly bounded in $L_{T}^{1}\left(L^{2}(K)\right)$, and $\rho_{M} a_{M} \xrightarrow{*} \rho a$ in that space, in the limit when $M \rightarrow+\infty$.

In the end, we have proved that the limit function $\rho$ is a weak solution to (3.1). Observe that, thanks to (3.8), the uniform bounds (3.6) and lower semicontinuity of the norm, we also deduce that $\rho$ verifies the energy estimate (3.4).

Time regularity and uniqueness. It remains to prove uniqueness of solutions and their time regularity. They are both consequences of the next proposition.

Proposition 3.1. Let $T>0$ and take $m \in \mathbb{N}$. Let $\rho \in L_{T}^{\infty}\left(H^{m}\right)$ be a weak solution to equation (3.1) under hypotheses (3.3).

Then $\rho \in C\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$ and it verifies the energy estimate (3.4).
We present the proof of the previous claim just in the minimal regularity case, namely for $m=0$. The general case follows by the same token. To start with, let us state a classical lemma (see e.g. [18] for details), whose proof is recalled in the appendix. For this, we fix a function $s \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$, with $s \equiv 1$ for $|x| \leq 1$ and $s \equiv 0$ for $|x| \geq 2$, $s$ radailly decreasing and such that $\int_{\mathbb{R}^{d}} s=1$. For all $n \in \mathbb{N}$, we then define $s_{n}(x):=n^{d} s(n x)$. We refer to the family $\left(s_{n}\right)_{n}$ as a family of standard mollifiers.
Lemma 3.2. Let $\left(s_{n}\right)_{n}$ be a family of standard mollifiers, as constructed here above. For all $n \in \mathbb{N}$, define the operator $S_{n}$, acting on tempered distributions over $\mathbb{R}_{+} \times \mathbb{R}^{d}$, by the formula

$$
S_{n} \rho:=s_{n} *_{x} \rho
$$

where the symbol $*_{x}$ means that the convolution is taken only with respect to the space variable. For given $\rho \in L_{T}^{\infty}\left(L^{2}\right)$ and $a \in L_{T}^{1}\left(C^{1}\right)$ such that $\nabla a \in L_{T}^{1}\left(C_{b}\right)$, we set, for all $n \in \mathbb{N}$ and $1 \leq j \leq d$,

$$
r_{n}^{j}(\rho):=\partial_{j}\left(\left[a, S_{n}\right] \rho\right)
$$

Then, for all $j$ fixed, we have $\left(r_{n}^{j}\right)_{n} \subset L_{T}^{1}\left(L^{2}\right)$; moreover, for $n \rightarrow+\infty$, we have the strong convergence $r_{n}^{j} \longrightarrow 0$ in $L_{T}^{1}\left(L^{2}\right)$.

Let us also recall the following standard notation. For $X$ a Banach space and $X^{*}$ its predual, we denote by $C_{w}([0, T] ; X)$ the set of measurable functions $f:[0, T] \rightarrow X$ which are continuous with respect to the weak topology. Namely, for any $\phi \in X^{*}$, the function $t \mapsto\langle\phi, f(t)\rangle_{X^{*} \times X}$ is continuous over $[0, T]$.

With this preparation, we are now ready to prove Proposition 3.1.
Proof of Proposition 3.1. With the same notations as in Lemma 3.2, let us define $\rho_{n}:=S_{n} \rho$. Notice that $\left(\rho_{n}\right)_{n} \subset L_{T}^{\infty}\left(L^{2}\right)$. Moreover, $\rho_{n}$ satisfies the equation

$$
\begin{equation*}
\partial_{t} \rho_{n}+\operatorname{div}\left(a \rho_{n}\right)=g_{n}+r_{n}, \quad \text { with } \quad\left(\rho_{n}\right)_{\mid t=0}=S_{n} \rho_{0} \tag{3.11}
\end{equation*}
$$

where we have set $r_{n}:=\operatorname{div}\left(\left[a, S_{n}\right] \rho\right)$. Notice that one has $\left\|S_{n} \rho_{0}\right\|_{L^{2}} \leq C\left\|\rho_{0}\right\|_{L^{2}}$ and $\left\|g_{n}\right\|_{L_{T}^{1}\left(L^{2}\right)} \leq$ $C\|g\|_{L_{T}\left(L^{2}\right)}$. Furthermore, in the limit $n \rightarrow+\infty$, we have the strong convergence properties $g_{n} \longrightarrow g$ in $L_{T}^{1}\left(L^{2}\right)$ and $S_{n} \rho_{0} \longrightarrow \rho_{0}$ in $L^{2}$. In addition, by Lemma 3.2 , we know that $\left\|r_{n}\right\|_{L_{T}^{1}\left(L^{2}\right)} \leq C$ and $r_{n} \longrightarrow 0$ in $L_{T}^{1}\left(L^{2}\right)$.

Now, we remark that an easy inspection of (3.11) implies the property $\left(\partial_{t} \rho_{n}\right)_{n} \subset L_{T}^{1}\left(H_{\text {loc }}^{-1}\right)$, which in turn gives us the uniform embedding $\left(\rho_{n}\right)_{n} \subset C_{T}\left(H_{\text {loc }}^{-1}\right)$. From this latter property, combined with a density argument and the uniform boundedness of $\left(\rho_{n}\right)_{n}$ in $L_{T}^{\infty}\left(L^{2}\right)$, we deduce that $\left(\rho_{n}\right)_{n}$ is uniformly bounded in $C_{w}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)$.

Next, let us take the $L^{2}$ scalar product of equation (3.11) by $\rho_{n}$ : by standard computations we get

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t}\left\|\rho_{n}\right\|_{L^{2}}^{2}+\frac{1}{2} \int \operatorname{div} a\left|\rho_{n}\right|^{2} d x=\int g_{n} \rho_{n} d x \tag{3.12}
\end{equation*}
$$

which implies that, for all $n \in \mathbb{N}$, one has $\left\|\rho_{n}(t)\right\|_{L^{2}} \in C([0, T])$. Thanks to this property, together with the fact that $\rho_{n} \in C_{w}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)$, after writing

$$
\left\|\rho_{n}(t+h)-\rho(t)\right\|_{L^{2}}^{2}=\left\|\rho_{n}(t+h)\right\|_{L^{2}}^{2}-2\left\langle\rho_{n}(t+h), \rho_{n}(t)\right\rangle_{L^{2} \times L^{2}}+\left\|\rho_{n}(t)\right\|_{L^{2}}^{2},
$$

one immediately deduces that, for all $n \in \mathbb{N}$, $\rho_{n}$ belongs to $C_{T}\left(L^{2}\right)$.
On the other hand, by straightforward computations, from relation (3.12) we also infer the following inequality:

$$
\begin{align*}
\left\|\rho_{n}(t)\right\|_{L^{2}} & \leq C \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau)\|_{L^{\infty}} d \tau\right)\left(\left\|S_{n} \rho_{0}\right\|_{L^{2}}+\int_{0}^{t}\left(\left\|g_{n}(\tau)\right\|_{L^{2}}+\left\|r_{n}(\tau)\right\|_{L^{2}}\right) d \tau\right)  \tag{3.13}\\
& \leq C \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau)\|_{L^{\infty}} d \tau\right)\left(\left\|\rho_{0}\right\|_{L^{2}}+\int_{0}^{t}\|g(\tau)\|_{L^{2}} d \tau\right)
\end{align*}
$$

for all $t \in[0, T]$, thanks also to the previous properties on $\left(S_{n} \rho_{0}\right)_{n},\left(g_{n}\right)_{n}$ and $\left(r_{n}\right)_{n}$. In view of this energy estimate, we deduce that $\left(\rho_{n}\right)_{n}$ is uniformly bounded in $C_{T}\left(L^{2}\right)$.

Next, we claim that $\left(\rho_{n}\right)_{n}$ is a Cauchy sequence in $C_{T}\left(L^{2}\right)$. For this, we take $m<n$ and consider the difference $\delta_{m}^{n} \rho:=\rho_{n}-\rho_{m}$. Then, $\delta_{m}^{n} \rho$ fulfils

$$
\partial_{t} \delta_{m}^{n} \rho+\operatorname{div}\left(a \delta_{m}^{n} \rho\right)=\delta_{m}^{n} g+\delta_{m}^{n} r, \quad \text { with } \quad \delta_{m}^{n} \rho_{\mid t=0}=\delta_{m}^{n} \rho_{0}:=\rho_{0}^{n}-\rho_{0}^{m},
$$

where we have defined also $\delta_{m}^{n} g:=g_{n}-g_{m}$ and $\delta_{m}^{n} r:=r_{n}-r_{m}$. To this equation we can also apply the energy estimates, and obtain

$$
\left\|\delta_{m}^{n} \rho\right\|_{L_{T}^{\infty}\left(L^{2}\right)} \leq C \exp \left(C\|\operatorname{div} a\|_{L_{T}^{1}\left(L^{\infty}\right)}\right)\left(\left\|\delta_{m}^{n} \rho_{0}\right\|_{L^{2}}+\left\|\delta_{m}^{n} g\right\|_{L_{T}^{1}\left(L^{2}\right)}+\left\|\delta_{m}^{n} r\right\|_{L_{T}^{1}\left(L^{2}\right)}\right) .
$$

At this point, we can conclude thanks to the fact that $\left(S_{n} \rho_{0}\right)_{n},\left(g_{n}\right)_{n}$ and $\left(r_{n}\right)_{n}$ are strongly convergent in the respective functional spaces, and thus they are, in particular, Cauchy sequences. So, our claim is proved.

Further, we deduce that the limit $\rho$ of the sequence $\left(\rho_{n}\right)_{n}$ belongs to $C_{T}\left(L^{2}\right)$, and the convergence $\rho_{n} \longrightarrow \rho$ is strong in this space. Finally, passing to the limit in the left-hand side of (3.13) we discover that $\rho$ verifies the energy estimate (3.4).

We conclude this part by remarking that stability, and then uniqueness, are easy consequences of Proposition 3.1.

Proposition 3.2. Fix $T>0$ and $m \in \mathbb{N}$, and let a be as in (3.3). For $i=1,2$, take an initial datum $\rho_{0}^{i} \in H^{m}\left(\mathbb{R}^{d}\right)$ and an external force $g^{i} \in L^{1}\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$, and let $\rho^{i} \in L_{T}^{\infty}\left(H^{m}\right)$ be a corresponding solution to (3.1) (whose existence is guaranteed by the previous arguments).

Then, after defining $\delta \rho_{0}:=\rho_{0}^{1}-\rho_{0}^{2}, \delta g:=g^{1}-g^{2}$ and $\delta \rho:=\rho^{1}-\rho^{2}$, the following estimate holds true for all $t \in[0, T]$, for some constant $C$ independent of the data and the respective solutions:

$$
\|\delta \rho(t)\|_{H^{m}} \leq C\left(\left\|\delta \rho_{0}\right\|_{H^{m}}+\int_{0}^{t}\|\delta g(\tau)\|_{H^{m}} d \tau\right) \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{m}} d \tau\right)
$$

Proof. It is enough to remark that, by taking the difference of the equations satisfied by $\rho^{1}$ and $\rho^{2}$, one deduces that $\delta \rho \in L_{T}^{\infty}\left(L^{2}\right)$ is a weak solution to the following equation:

$$
\left\{\begin{array}{l}
\partial_{t} \delta \rho+\operatorname{div}(a \delta \rho)=\delta g \\
\delta \rho_{\mid t=0}=\delta \rho_{0}
\end{array}\right.
$$

Then, Proposition (3.1) applies, and gives us the claimed estimates.

### 3.1.2 The case of the transport equation

The characterization of ensemble controls with the optimality conditions given in Section 5.2, requires the solution of an adjoint Liouville problem, which is given by a linear transport problem. In preparation of that discussion, and to complete the analysis of the present section, we consider the following transport problem

$$
\begin{cases}\partial_{t} q+a \cdot \nabla q+b q=g & \text { in }[0, T] \times \mathbb{R}^{d}  \tag{3.14}\\ q_{\mid t=0}=q_{0} & \text { on } \mathbb{R}^{d} .\end{cases}
$$

We assume that the data $q_{0}, a$ and $g$ verify the assumptions in (3.3), where $\rho_{0}$ is replaced by $q_{0}$. Moreover, we assume that $b$ has the same regularity as div $a$ : that is, $b \in L^{1}\left([0, T] ; C_{b}^{m}\left(\mathbb{R}^{d}\right)\right)$.

We point out that the weak formulation of (3.14) now reads as follows: for all $\phi \in C_{c}^{\infty}\left(\mathbb{R}^{d} \times[0, T[)\right.$, one has the equality

$$
\begin{equation*}
-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho \partial_{t} \phi-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho a \cdot \nabla \phi-\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho \operatorname{div} a \phi+\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho b \phi=\int_{0}^{T} \int_{\mathbb{R}^{d}} g \phi+\int_{\mathbb{R}^{d}} \rho_{0} \phi(0) . \tag{3.15}
\end{equation*}
$$

For (3.14), we have the following well-posedness result, analogous to Theorem 3.2 for the Liouville equation.

Theorem 3.3. Let us fix $T>0$ and $m \in \mathbb{N}$, and let the data $a, b, q_{0}$ and $g$ satisfy the assumptions stated above.

Then there exists a unique solution $q \in C\left([0, T] ; H^{m}\left(\mathbb{R}^{d}\right)\right)$ to equation (3.14). Moreover, there exists a "universal" constant $C>0$, independent of $q_{0}, a, b, g, q$ and $T$, such that the following estimate holds true for any $t \in[0, T]$ :

$$
\begin{equation*}
\|q(t)\|_{H^{m}} \leq C\left(\left\|q_{0}\right\|_{H^{m}}+\int_{0}^{t}\|g(\tau)\|_{H^{m}} d \tau\right) \exp \left(C \int_{0}^{t}\left(\|\nabla a(\tau)\|_{C_{b}^{m}}+\|b(\tau)\|_{C_{b}^{m}}\right) d \tau\right) \tag{3.16}
\end{equation*}
$$

The proof is analogous to the one given for Theorem 3.1, so it is omitted here. In particular, the regularization procedure and the energy estimates work in exactly the same way. The only point which
deserves some attention is passing to the limit in the weak formulation (3.15) at step $n$ of the regularization procedure, especially in the terms

$$
-\int_{0}^{T} \int_{\mathbb{R}^{d}} q^{n} \operatorname{div} a^{n} \phi d x d t+\int_{0}^{T} \int_{\mathbb{R}^{d}} q^{n} b^{n} \phi d x d t
$$

Let us focus on the latter term only; the former can be treated by the same method. Of course, it is enough to treat the case when $m=0$.

Now, we notice that the integral is in fact performed on the compact set $K:=\operatorname{supp} \phi$. Therefore, by Proposition 4.21 and Theorem 4.22 of $[7]$, we deduce that $b^{n} \longrightarrow b$ in $L_{T}^{1}\left(L^{\infty}(K)\right)$ for $n \rightarrow+\infty$. On the other hand, thanks to uniform bounds, $q^{n} \stackrel{*}{\rightharpoonup} q$ in $L_{T}^{\infty}\left(L^{2}\right)$, for some $q$ belonging to that space; so, in particular the weak-* convergence holds true in $L_{T}^{\infty}\left(L^{2}(K)\right)$. Putting these properties together, we deduce that $\left(q^{n} b^{n}\right)_{n}$ is uniformly bounded in $L_{T}^{1}\left(L^{2}(K)\right)$ and it weakly-* converges to $q b$ in that space.

The previous argument shows that we can pass to the limit in the weak formulation of the approximated problems and gather that the limit point $q$ of the sequence $\left(q^{n}\right)_{n}$ solves (3.15).

### 3.2 Well-posedness theory in weighted spaces

In this section, we extend the previous theory to Sobolev spaces with weights. This analysis is especially important for the investigation of the Liouville control-to-state map and of the Liouville ensemble optimal control problem, see the next sections.

Remark 3.3. We limit ourselves to treat the case of the Liouville equation. However, the statements that follow can be proved also for the transport problem, with slight modifications in the proofs.

### 3.2.1 Definition of weighted spaces

For the analysis of the Liouville control-to-state map in Section 4, we need to prove weighted integrability of $\rho$, due to the growth of the drift function. For this purpose, we introduce the following definition.

Definition 3.1. Fix $(m, k) \in \mathbb{N}^{2}$. We define the space $H_{k}^{m}\left(\mathbb{R}^{d}\right)$ in the following way:

$$
H_{k}^{0}\left(\mathbb{R}^{d}\right)=L_{k}^{2}\left(\mathbb{R}^{d}\right):=\left\{\left.f \in L^{2}\left(\mathbb{R}^{d}\right)|\quad| x\right|^{k} f \in L^{2}\left(\mathbb{R}^{d}\right)\right\}
$$

and, for $m \geq 1$, we set

$$
H_{k}^{m}\left(\mathbb{R}^{d}\right):=\left\{\left.f \in H^{m}\left(\mathbb{R}^{d}\right) \cap H_{k}^{m-1}\left(\mathbb{R}^{d}\right)|\quad| x\right|^{k} D^{\alpha} f \in L^{2}\left(\mathbb{R}^{d}\right) \quad \forall|\alpha|=m\right\}
$$

The space $H_{k}^{m}$ is endowed with the following norm:

$$
\|f\|_{H_{k}^{m}}:=\sum_{|\alpha| \leq m}\left\|\left(1+|x|^{k}\right) D^{\alpha} f\right\|_{L^{2}} .
$$

Sometimes, given $m \in \mathbb{N}$, we will use the notation

$$
\left\|\nabla^{m} f\right\|_{L^{2}}=\sum_{|\alpha|=m}\left\|D^{\alpha} f\right\|_{L^{2}}
$$

and analogous writing for weighted norms.

Notice that, for all fixed $m$ and $k$ in $\mathbb{N}$, one has the embedding $H_{k}^{m} \subset H^{m}$. Of course, $H^{m}=H_{0}^{m}$ for all $m \geq 0$. Furthermore, since we want to avoid too singular behaviours close to 0 , we will often focus on the special case (which will be enough for our scopes)

$$
m \leq k
$$

Then, we have a simple characterization of the spaces $H_{k}^{m}$, which will be useful especially in Section 4, when studying the control-to-state map related to our optimal control problem.

Proposition 3.3. (i) Given $k \in \mathbb{N}$, one has $f \in L_{k}^{2}$ if and only if $\left(1+|x|^{k}\right) f \in L^{2}$.
(ii) For $k \in \mathbb{N} \backslash\{0\}$ and $1 \leq m \leq k$, let $f \in H^{m} \cap H_{k}^{m-1}$. Then $f \in H_{k}^{m}$ if and only if $|x|^{k} f \in H^{m}$. In particular, a tempered distribution $f$ belongs to $H_{1}^{1}$ if and only if both $f$ and $|x| f$ belong to $H^{1}$; it belongs to $H_{2}^{2}$ if and only if both $f$ and $|x|^{2} f$ belong to $H^{2}$ and $\nabla f$ belongs to $L_{2}^{2}$.

Before proving this statement, let us state a preliminary result, whose proof is postponed to the Appendix.

Lemma 3.3. Let $(m, k) \in \mathbb{N}^{2}$, with $m \leq k$. If $f \in H_{k}^{m}$, then $\left(1+|x|^{k}\right) f \in H^{m}$.
Thanks to Lemma 3.3, we can prove Proposition 3.3.
Proof of Proposition 3.3. Assertion (i) is elementary. So, let us focus on the proof of (ii).
Suppose that $f \in H^{m} \cap H_{k}^{m-1}$. Then, by Lemma 3.3 above, we have that $|x|^{k} f \in H_{k}^{m-1}$. At this point, for $|\alpha|=m$, we write, using again Leibniz rule,

$$
D^{\alpha}\left(|x|^{k} f\right)=|x|^{k} D^{\alpha} f+\sum_{\beta} D^{\beta}|x|^{k} D^{\alpha-\beta} f
$$

where the sum is performed for all $\beta \leq \alpha$ such that $|\beta| \geq 1$. By the previous arguments, and the fact that $m \leq k$, we have that all the terms in the sum belong to $L^{2}$. Then, the term on the left-hand side belongs to $L^{2}$ if and only if the first term on the right-hand side does.

The last sentences follow by straightforward computations. First of all, we have that

$$
\partial_{j}(|x| f)=\partial_{j}|x| f+|x| \partial_{j} f
$$

for all $1 \leq j \leq d$. Furthermore, we also have

$$
\nabla^{2}\left(|x|^{2} f\right) \sim \nabla\left(|x| f+|x|^{2} \nabla f\right) \sim \nabla|x| f+\left(|x|+|x|^{2}\right) \nabla f+|x|^{2} \nabla^{2} f
$$

The equivalence between the two assertions is then apparent. Indeed, arguing as in the beginning of the proof to Lemma 3.3, we gather that, if $f \in H_{2}^{2}$, then $|x|^{j} D^{\alpha} f \in L^{2}$ for all $0 \leq j \leq 2$ and $|\alpha|=0,1$. Hence, all the terms in the right-hand side belong to $L^{2}$, and then so does the one on the left-hand side. On the contrary, if both $f$ and $|x|^{2} f$ belong to $H^{2}$ and $\nabla f$ belongs to $L_{2}^{2}$, then $f \in H^{2} \cap H_{2}^{1}$; finally, by the previous equality, we also discover that $|x|^{2} \nabla^{2} f$ belongs to $L^{2}$, completing the proof of the reverse implication.

The proof of the proposition is hence completed.

### 3.2.2 The Liouville equation in weighted spaces

After the above preliminaries, we are ready to state the main result of this section, which show wellposedness of the Liouville equation in $H_{k}^{m}$ spaces.

Theorem 3.4. Let $T>0$ and $(m, k) \in \mathbb{N}^{2}$ fixed, and let a be a vector field satisfying hypotheses (3.3). Moreover, assume that $\rho_{0} \in H_{k}^{m}\left(\mathbb{R}^{d}\right)$ and $g \in L^{1}\left([0, T] ; H_{k}^{m}\left(\mathbb{R}^{d}\right)\right)$.

Then there exists a unique solution $\rho \in C\left([0, T] ; H_{k}^{m}\left(\mathbb{R}^{d}\right)\right)$ to problem (3.1). Moreover, there exists a "universal" constant $C>0$, independent of $\rho_{0}, a, g, \rho$ and $T$, such that the following estimate holds true for any $t \in[0, T]$ :

$$
\begin{equation*}
\|\rho(t)\|_{H_{k}^{m}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{m}} d \tau\right)\left(\left\|\rho_{0}\right\|_{H_{k}^{m}}+\int_{0}^{t}\|g(\tau)\|_{H_{k}^{m}} d \tau\right) . \tag{3.17}
\end{equation*}
$$

Before proving this statement in its full generality, let us consider its version for simpler cases, which will be needed in the proof of the general case. Moreover, their precise form is important, in view of their application in Section 4.

We start with the case $m=0$.
Lemma 3.4. Assume that the hypotheses of Theorem 3.4 hold true with $m=0$.
Then there exists a unique solution $\rho \in C\left([0, T] ; L_{k}^{2}\left(\mathbb{R}^{d}\right)\right)$ to problem (3.1). Moreover, there exists a "universal" constant $C>0$ such that the following estimate holds true for any $t \in[0, T]$ :

$$
\|\rho(t)\|_{L_{k}^{2}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{L^{\infty}} d \tau\right)\left(\left\|\rho_{0}\right\|_{L_{k}^{2}}+\int_{0}^{t}\|g(\tau)\|_{L_{k}^{2}} d \tau\right)
$$

Proof of Lemma 3.4. Most of the claims of the Lemma follow from Theorem 3.2. We have just to prove propagation of higher integrability (i.e. $k \geq 1$ ) of the initial datum and external force. Omitting a standard regularization procedure for the sake of brevity, we will perform energy estimates directly on equation (3.1).

First of all, for completeness, we consider the case $k=0$. If we take the $L^{2}$ scalar product of equation (3.1) by $\rho$, by standard computations we get

$$
\frac{1}{2} \frac{d}{d t}\|\rho\|_{L^{2}}^{2}+\frac{1}{2} \int \operatorname{div} a|\rho|^{2} d x=\int g \rho d x
$$

From this relation, we easily get

$$
\begin{equation*}
\frac{d}{d t}\|\rho\|_{L^{2}} \leq\|\operatorname{div} a\|_{L^{\infty}}\|\rho\|_{L^{2}}+\|g\|_{L^{2}} . \tag{3.18}
\end{equation*}
$$

Next, let us multiply equation (3.1) by $|x|^{k}$ : we get that $\rho_{k}:=|x|^{k} \rho$ satisfies

$$
\partial_{t} \rho_{k}+\operatorname{div}\left(a \rho_{k}\right)=|x|^{k} g+\rho a \cdot \nabla|x|^{k} .
$$

Taking the $L^{2}$ scalar product of this equation by $\rho_{k}$ and repeating the same computations as above, we find

$$
\begin{equation*}
\frac{d}{d t}\left\|\rho_{k}\right\|_{L^{2}} \leq\|\operatorname{div} a\|_{L^{\infty}}\left\|\rho_{k}\right\|_{L^{2}}+\left\||x|^{k} g\right\|_{L^{2}}+\left\|\rho a \cdot \nabla|x|^{k}\right\|_{L^{2}} \tag{3.19}
\end{equation*}
$$

We need to control the last term on the right-hand side of the previous estimate. For this, we use the fact that $\nabla|x|^{k} \sim|x|^{k-1}$ for all $k \geq 1$, and Remark 3.2, to obtain

$$
\left\|\rho a \cdot \nabla|x|^{k}\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \rho\right\|_{L^{2}} .
$$

Inserting this bound into (3.19) and summing up the resulting expression to (3.18), we have

$$
\begin{equation*}
\frac{d}{d t}\left\|\left(1+|x|^{k}\right) \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) g\right\|_{L^{2}} . \tag{3.20}
\end{equation*}
$$

Hence, an application of Grönwall's lemma gives the desired estimate.
Next, we present results for $m=1$.
Lemma 3.5. Assume that the hypotheses of Theorem 3.4 hold true with $m=1$.
Then there exists a unique solution $\rho \in C\left([0, T] ; H_{k}^{1}\left(\mathbb{R}^{d}\right)\right)$ to problem (3.1). Moreover, there exists a "universal" constant $C>0$ such that the following estimate holds true for any $t \in[0, T]$ :

$$
\|\rho(t)\|_{H_{k}^{1}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{1}} d \tau\right)\left(\left\|\rho_{0}\right\|_{H_{k}^{1}}+\int_{0}^{t}\|g(\tau)\|_{H_{k}^{1}} d \tau\right)
$$

Proof of Lemma 3.5. Once again, it is enough to focus on the proof of the energy estimates. We start by differentiating equation (3.1) with respect to $x^{j}$, for some $1 \leq j \leq d$ : we get

$$
\partial_{t} \partial_{j} \rho+\operatorname{div}\left(a \partial_{j} \rho\right)=\partial_{j} g-\partial_{j} \operatorname{div} a \rho-\partial_{j} a \cdot \nabla \rho
$$

Applying estimate (3.20) to this equation gives

$$
\begin{aligned}
& \frac{d}{d t}\left\|\left(1+|x|^{k}\right) \partial_{j} \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \partial_{j} \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) \partial_{j} g\right\|_{L^{2}}+ \\
&+\left\|\left(1+|x|^{k}\right) \partial_{j} \operatorname{div} a \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) \partial_{j} a \cdot \nabla \rho\right\|_{L^{2}}
\end{aligned}
$$

from which we obtain, for another constant $C>0$, the following bound:

$$
\begin{align*}
\frac{d}{d t}\left\|\left(1+|x|^{k}\right) \nabla \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla \rho\right\|_{L^{2}} & +\left\|\left(1+|x|^{k}\right) \nabla g\right\|_{L^{2}}+  \tag{3.21}\\
& +\left\|\nabla^{2} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \rho\right\|_{L^{2}}
\end{align*}
$$

We can now sum up (3.20) and (3.21) to get

$$
\begin{equation*}
\frac{d}{d t}\|\rho\|_{H_{k}^{1}} \leq C\|\nabla a\|_{C_{b}^{1}}\|\rho\|_{H_{k}^{1}}+\|g\|_{H_{k}^{1}} \tag{3.22}
\end{equation*}
$$

and Grönwall's lemma allows us to get the result.
Now, we can address the proof of the general case, namely of Theorem 3.4.
Proof of Theorem 3.4. We argue by induction on the order of derivatives, i.e. on $m$, the cases $m=0$ and $m=1$ being given by Lemma 3.4 and Lemma 3.5, respectively.

Let $m \geq 2$, and let us assume that, for any $0 \leq \ell \leq m-1$, the following inequality holds true

$$
\begin{align*}
& \frac{d}{d t}\left\|\left(1+|x|^{k}\right) \nabla^{\ell} \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla^{\ell} \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) \nabla^{\ell} g\right\|_{L^{2}}+  \tag{3.23}\\
&+\sum_{0 \leq p \leq \ell-1}\left\|\nabla^{p+1} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla^{p} \rho\right\|_{L^{2}}
\end{align*}
$$

Our goal is to prove an analogous estimate also for $\left\|\left(1+|x|^{k}\right) \nabla^{m} \rho\right\|_{L^{2}}$.

For this purpose, let us take an $\alpha \in \mathbb{N}^{d}$ such that $|\alpha|=m$, and let us apply it to (3.1). Thus, we deduce the following equality

$$
\begin{equation*}
\partial_{t} D^{\alpha} \rho+\operatorname{div}\left(a D^{\alpha} \rho\right)=D^{\alpha} g-\sum_{0<\beta \leq \alpha} D^{\beta} \operatorname{div} a D^{\alpha-\beta} \rho-\sum_{0<\beta \leq \alpha} D^{\beta} a \cdot \nabla D^{\alpha-\beta} \rho, \tag{3.24}
\end{equation*}
$$

where the notation $0<\beta$ means that $\beta \in \mathbb{N}^{d}$ has at least one non-zero component.
Following the computations of Lemma 3.5, we need to estimate the $L_{k}^{2}$ norm of the last two terms in the right-hand side of the previous equation. First of all, we have

$$
\left\|\left(1+|x|^{k}\right) D^{\beta} \operatorname{div} a D^{\alpha-\beta} \rho\right\|_{L^{2}} \leq\left\|\nabla^{|\beta|+1} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) D^{\alpha-\beta} \rho\right\|_{L^{2}} .
$$

Notice that, since $\beta>0$, the terms $D^{\alpha-\beta} \rho$ are lower order. The same can be said of the terms

$$
\left\|\left(1+|x|^{k}\right) D^{\beta} a \cdot \nabla D^{\alpha-\beta} \rho\right\|_{L^{2}} \leq\left\|\nabla^{|\beta|} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla D^{\alpha-\beta} \rho\right\|_{L^{2}},
$$

whenever $|\beta| \geq 2$; on the contrary, when $|\beta|=1$, the terms $\nabla D^{\alpha-\beta} \rho$ contain exactly $m$ derivatives.
Therefore, applying estimate (3.20) to equation (3.24), and using the previous controls, we infer

$$
\begin{align*}
& \frac{d}{d t}\left\|\left(1+|x|^{k}\right) \nabla^{m} \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla^{m} \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) \nabla^{m} g\right\|_{L^{2}}+  \tag{3.25}\\
&+\sum_{0<\beta \leq \alpha}\left\|\nabla^{|\beta|+1} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) D^{\alpha-\beta} \rho\right\|_{L^{2}} \\
& \leq C\|\nabla a\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla^{m} \rho\right\|_{L^{2}}+\left\|\left(1+|x|^{k}\right) \nabla^{m} g\right\|_{L^{2}}+ \\
&+\sum_{0 \leq \ell \leq m-1}\left\|\nabla^{\ell+1} a\right\|_{L^{\infty}}\left\|\left(1+|x|^{k}\right) \nabla^{\ell} \rho\right\|_{L^{2}},
\end{align*}
$$

which proves formula (3.23) at the level $m$. Therefore that formula is true for any $m \in \mathbb{N}$, by induction.
Now, summing up inequality (3.21) for $\ell=0$ to $m$, and by the definition of $H_{k}^{m}$ norms, we get, for some constant also depending on $m$, the following bound:

$$
\frac{d}{d t}\|\rho\|_{H_{k}^{m}} \leq C\|\nabla a\|_{C_{b}^{m}}\|\rho\|_{H_{k}^{m}}+\|g\|_{H_{k}^{m}}
$$

which immediately implies the claimed estimate. Theorem 3.4 is now proved.

## 4 The Liouville control-to-state map

In this section, we define the Liouville control-to-state map and investigate its continuity and differentiability properties. For reasons which will appear clear in the following analysis, we need to resort to weighted spaces $H_{k}^{m}$, as introduced in Section 3.2.

We start by making an important remark.
Remark 4.1. Throughout this section, the data of the Liouville equation has to be thought as fixed. Specifically, for $m \geq 0$ and $k \geq 0$, we take an initial datum $\rho_{0} \in H_{k}^{m}$, a source term $g \in L_{T}^{1}\left(H_{k}^{m}\right)$, and a drift function $a_{0} \in L_{T}^{1}\left(C^{m+1}\right)$, with $\nabla a_{0} \in L_{T}^{1}\left(C_{b}^{m}\right)$.

We are then interested in the dependence of the solution $\rho$ to the Liouville equation (3.1), with drift a given by (2.4), on the control state $u \in U_{a d}$, where $U_{\text {ad }}$ has been defined in (2.9).

### 4.1 Definition and continuity properties

We remark that the statements of Theorems 3.2 and 3.4 cover the case of the Liouville equation with the controlled drift function given by (2.4), where $u \in U_{a d}$. In particular, the next proposition-definition immediately follows.

Proposition 4.1. Fixed data $\rho_{0}, g$ and $a_{0}$ as in Remark 4.1, let us consider drift functions a of the form (2.4), with $u \in U_{\text {ad }}$. Introduce the Liouville control-to-state map $G$, defined by

$$
G: U_{a d} \longrightarrow L^{\infty}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right), \quad u \mapsto \rho:=G(u),
$$

where $\rho$ is the unique solution to the Liouville equation with the given data.
Then $G$ is well-defined.
Let us make an important comment about the previous definition.
Remark 4.2. Notice that the theory developed in Sections 3.1 and 3.2 entails that the solution $\rho$ actually belongs to $C_{T}\left(H_{k}^{m}\right)$. However, for reasons that will appear clear in what follows (namely, a loss of regularity, both in $m$ and $k$, when proving Fréchet differentiability of $G$ ), it is convenient to look at $G$ as a map with values in the space with the weakest topology. Notice that, for any $(m, k) \in \mathbb{N}^{2}$, the space $H_{k}^{m}$ is embedded in $L^{2}$.

Finally, we consider $L^{\infty}$ regularity with respect to time, because it will be convenient also to look at weak continuity properties of $G$, see Proposition 4.2 below.

Next, we study some properties of the map $G$ that are relevant for the analysis of ensemble optimal control problems. We start by establishing that $G$ is weak-weak continuous from $U_{a d}$ into $L_{T}^{\infty}\left(L^{2}\right)$. Notice that we do not need any restriction on $m$ and $k$ (and so, on the initial data) in this case.

Proposition 4.2. Take $m \geq 0$ and $k \geq 0$ and initial data $\rho_{0} \in H_{k}^{m}, g \in L_{T}^{1}\left(H_{k}^{m}\right)$ and $a_{0} \in L_{T}^{1}\left(C^{m+1}\right)$ such that $\nabla a_{0} \in L_{T}^{1}\left(C_{b}^{m}\right)$. Let $u \in U_{\text {ad }}$ and $\left(u^{l}\right)_{l} \subset U_{a d}$ be a sequence of controls, and assume that $u^{l} \xrightarrow{*} u$ in $\mathbb{L}_{T}^{\infty}$.

Then $G\left(u^{l}\right) \stackrel{*}{\rightharpoonup} G(u)$ in the weak-* topology of $L_{T}^{\infty}\left(L^{2}\right)$.
Proof. Of course, it is enough to prove the previous proposition in the case of minimal regularity and integrability, namely for $m=k=0$.

By definition of the set $U_{a d}$, we infer that $\left(u^{l}\right)_{l}$ is uniformly bounded in $\mathbb{L}_{T}^{\infty}$. On the other hand, by hypotheses and Theorem 3.2, for all $l \in \mathbb{N}$ there exists a unique $\rho^{l}:=G\left(u^{l}\right) \in C_{T}\left(L^{2}\right)$ which solves the Liouville equation (3.1). In addition, by inequality (3.4), we deduce that $\left(\rho^{l}\right)_{l}$ is uniformly bounded in $C_{T}\left(L^{2}\right)$; then there exists $\rho \in L_{T}^{\infty}\left(L^{2}\right)$ such that, up to extraction of a subsequence, $\rho^{l} \xrightarrow{*} \rho$ in $L_{T}^{\infty}\left(L^{2}\right)$.

Now, our goal is to prove that $\rho$ is a weak solution to the Liouville equation

$$
\begin{equation*}
\partial_{t} \rho+\operatorname{div}(a(t, x ; u) \rho)=g, \quad \text { with } \quad \rho_{\mid t=0}=\rho_{0} \tag{4.1}
\end{equation*}
$$

Indeed, if this is the case, by uniqueness we get $\rho=G(u)$ and that the whole sequence $\left(\rho^{l}\right)_{l}$ converges, achieving in this way the proof of the proposition.

In order to prove our claim, we need to pass to the limit in the weak formulation of the Liouville equation for $\rho^{l}$, when $l \rightarrow+\infty$. Recalling also our special choice (2.4), it is easy to see that the only term which presents some difficulty is the non-linear term

$$
\begin{equation*}
\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho^{l}\left(u_{1}^{l}(t)+x \circ u_{2}^{l}(t)\right) \cdot \nabla \phi d x d t, \quad \text { for any fixed } \quad \phi \in C_{c}^{\infty}\left(\left[0, T\left[\times \mathbb{R}^{d}\right)\right.\right. \tag{4.2}
\end{equation*}
$$

Therefore, let us focus on the convergence of this integral. First of all, by inspection of the equation $\partial_{t} \rho^{l}=$ $-\operatorname{div}\left(a\left(x, t ; u^{l}\right) \rho^{l}\right)+g$, we discover that $\left(\partial_{t} \rho^{l}\right)_{l} \subset L_{T}^{1}\left(H_{\text {loc }}^{-1}\right)$, which implies that $\left(\rho^{l}\right)_{l} \subset W_{T}^{1,1}\left(H_{\text {loc }}^{-1}\right)$. Then, by the Rellich-Kondrachov theorem and Cantor's diagonal procedure, we discover that, up to an extraction of a subsequence that we do not relabel, $\left(\rho^{l}\right)_{l}$ is compact, and then strongly convergent, in $L_{T}^{1}\left(H_{\mathrm{loc}}^{-2}\right)$. Interpolating this compactness result with the uniform boundedness in $L_{T}^{\infty}\left(L_{\mathrm{loc}}^{2}\right)$, we discover that $\rho^{l} \longrightarrow \rho$ in $L_{T}^{1}\left(H_{\mathrm{loc}}^{-s}\right)$, for any $s>0$.

In view of the uniform boundedness of $\left(u_{1}^{l}\right)_{l}$ and $\left(u_{2}^{l}\right)_{l}$ in $L_{T}^{\infty}$, the previous property is enough to pass to the limit in the integral (4.2), and prove that it converges to

$$
\int_{0}^{T} \int_{\mathbb{R}^{d}} \rho\left(u_{1}(t)+x \circ u_{2}(t)\right) \cdot \nabla \phi d x d t
$$

for all given $\phi \in C_{c}^{\infty}\left(\left[0, T\left[\times \mathbb{R}^{d}\right)\right.\right.$. Thus, we get that (4.1) is satisfied, and then we can conclude the proof as already mentioned above.

For the analysis of our optimal control problem, see Section 5 below, we need stronger regularity properties for $G$. We start by showing Lipschitz continuity, which will be the basis to prove Gâteaux differentiability of $G$, in the next paragraph. The key here is to perform careful estimates in order to identify the right topology: the reason is that, due to hyperbolicity of the Liouville equation, stability estimates involve a loss of regularity.

Lemma 4.1. Let the data $\rho_{0}, g$ and $a_{0}$ be fixed as in Remark 4.1 above, with $m \geq 1$ and $k \geq 1$. Let $u$ and $v$ be in $U_{\text {ad }}$, and denote by $G(u)$ and $G(v)$ the corresponding $C_{T}\left(H_{k}^{m}\right)$ solutions to (3.1), with drift a given by (2.4). Set $\delta G:=G(u)-G(v)$.

Then there exists a "universal" constant $C>0$ such that, for all $1 \leq \ell \leq k$, if we set

$$
\begin{equation*}
K_{0}^{(\ell)}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\|u\|_{\mathbb{L}_{T}^{1}}+\|v\|_{\mathbb{L}_{T}^{1}}\right)\right) \times\left(\left\|\rho_{0}\right\|_{H_{\ell}^{1}}+\|g\|_{L_{T}^{1}\left(H_{\ell}^{1}\right)}\right) \tag{4.3}
\end{equation*}
$$

then, for all $t \in[0, T]$, one has

$$
\|\delta G(t)\|_{L_{\ell-1}^{2}} \leq K_{0}^{(\ell)} \int_{0}^{t}|u(\tau)-v(\tau)| d \tau
$$

If moreover $m \geq 2$ and we set

$$
\begin{equation*}
K_{1}^{(\ell)}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)}+\|u\|_{\mathbb{L}_{T}^{1}}+\|v\|_{\mathbb{L}_{T}^{1}}\right)\right) \times\left(\left\|\rho_{0}\right\|_{H_{\ell}^{2}}+\|g\|_{L_{T}^{1}\left(H_{\ell}^{2}\right)}\right) \tag{4.4}
\end{equation*}
$$

we also have

$$
\|\delta G(t)\|_{H_{\ell-1}^{1}} \leq K_{1}^{(\ell)} \int_{0}^{t}|u(\tau)-v(\tau)| d \tau
$$

Proof. By linearity of the Liouville equation, we find that $\delta G$ satisfies

$$
\begin{equation*}
\partial_{t} \delta G+\operatorname{div}(a(t, x ; u) \delta G)=-\operatorname{div}(\bar{a}(t, x ; u-v) G(v)), \quad \delta G_{\mid t=0}=0 \tag{4.5}
\end{equation*}
$$

where we have set

$$
\begin{equation*}
\bar{a}(t, x ; u-v):=a(t, x ; u)-a(t, x ; v)=\left(u_{1}-v_{1}\right)+x \circ\left(u_{2}-v_{2}\right) . \tag{4.6}
\end{equation*}
$$

Applying $L_{\ell-1}^{2}$ estimates of Theorem 3.2 to equation (4.5), we immediately get

$$
\|\delta G(t)\|_{L_{\ell-1}^{2}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau, x ; u)\|_{L^{\infty}} d \tau\right) \int_{0}^{t}\|\operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))\|_{L_{\ell-1}^{2}} d \tau
$$

By explicit computations and using the Leibniz rule, we deduce that

$$
\begin{align*}
& \|\operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))\|_{L_{\ell-1}^{2}} \leq|u(\tau)-v(\tau)|\left(\|G(v)\|_{L_{\ell-1}^{2}}+\|\nabla G(v)\|_{L_{\ell}^{2}}\right)  \tag{4.7}\\
& \quad \leq C|u(\tau)-v(\tau)| \exp \left(C \int_{0}^{\tau}\|\nabla a(s, x ; v)\|_{C_{b}^{1}} d s\right)\left(\left\|\rho_{0}\right\|_{H_{\ell}^{1}}+\int_{0}^{\tau}\|g(s)\|_{H_{\ell}^{1}} d s\right)
\end{align*}
$$

where the second inequality holds true in view of the bound $\|G(v)\|_{L_{\ell-1}^{2}}+\|\nabla G(v)\|_{L_{\ell}^{2}} \leq\|G(v)\|_{H_{\ell}^{1}}$ and Lemma 3.5. This estimate completes the proof of the first inequality, for $L^{2}$-type norms of $\delta G$.

Now, we focus on $H_{\ell-1}^{1}$ bounds for $\delta G$. Thanks to Lemma 3.5, we have

$$
\begin{equation*}
\|\delta G(t)\|_{H_{\ell-1}^{1}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau, x ; u)\|_{C_{b}^{1}} d \tau\right) \int_{0}^{t}\|\operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))\|_{H_{\ell-1}^{1}} d \tau \tag{4.8}
\end{equation*}
$$

By definition, we have that $\|f\|_{H_{\ell-1}^{1}}=\|f\|_{L_{\ell-1}^{2}}+\|\nabla f\|_{L_{\ell-1}^{2}}$. Then, we start with the bound

$$
\begin{align*}
\|\operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))\|_{L_{\ell-1}^{2}} & =\|\operatorname{div} \bar{a} G(v)\|_{L_{\ell-1}^{2}}+\|\bar{a} \cdot \nabla G(v)\|_{L_{\ell-1}^{2}}  \tag{4.9}\\
& \leq C|u(\tau)-v(\tau)|\left(\|G(v)\|_{L_{\ell-1}^{2}}+\|\nabla G(v)\|_{L_{\ell}^{2}}\right) \\
& \leq C|u(\tau)-v(\tau)|\|G(v)\|_{H_{\ell}^{1}} .
\end{align*}
$$

Next, we need to bound in $L_{\ell-1}^{2}$ the quantity $\nabla \operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))$ : we have then to control four terms. First of all, we notice that $\nabla \operatorname{div} \bar{a} \equiv 0$. Moreover, we can write

$$
\begin{equation*}
\|\operatorname{div} \bar{a} \nabla G(v)\|_{L_{\ell-1}^{2}} \leq|u(\tau)-v(\tau)|\|\nabla G(v)\|_{L_{\ell-1}^{2}}, \tag{4.10}
\end{equation*}
$$

and the same estimate holds true also for the term $\nabla \bar{a} \cdot \nabla G(v)$. Finally, we have

$$
\begin{equation*}
\left\|\bar{a} \cdot \nabla^{2} G(v)\right\|_{L_{\ell-1}^{2}} \leq C|u(\tau)-v(\tau)|\left\|\nabla^{2} G(v)\right\|_{L_{\ell}^{2}} \tag{4.11}
\end{equation*}
$$

Putting (4.9), (4.10) and (4.11) together, we infer the bound

$$
\|\operatorname{div}(\bar{a}(\tau, x ; u-v) G(v))\|_{H_{\ell-1}^{1}} \leq C|u(\tau)-v(\tau)|\|G(v)\|_{H_{\ell}^{2}} .
$$

Inserting this last inequality into (4.8) and using the bounds of Theorem 3.4, we finally get the claimed estimate for the $H^{1}$-type norms of $\delta G$.

### 4.2 Differentiability of the control-to-state map

In this section, we investigate differentiability properties of the control-to-state map $G$, defined above.
Now, with Lemma 4.1 at hand, we can prove Gâteaux differentiability of $G$. For any given $u$ in an open set $U_{0} \subset U_{a d}$, let $G(u)$ be the corresponding solution to the Liouville equation, as defined in Proposition 4.1, and let $\delta u=\left(\delta u_{1}, \delta u_{2}\right)$ be an admissible variation of $u$, such that $u+\varepsilon \delta u \in U_{a d}$ for $\varepsilon \in \mathbb{R} \backslash\{0\}$ sufficiently small. Then the Gâteaux derivative of $G$ with respect to the variation $\delta u$ at $u$ is defined as the limit (whenever such a limit exists)

$$
\begin{equation*}
\delta_{\delta u} G(u):=\lim _{\varepsilon \rightarrow 0} \frac{G(u+\varepsilon \delta u)-G(u)}{\varepsilon} . \tag{4.12}
\end{equation*}
$$

The next proposition holds true.

Proposition 4.3. Let $m \geq 2$ and $k \geq 2$. Let the data $\rho_{0}, g$ and $a_{0}$ be fixed as in Remark 4.1 above. Let $u$ belong to int $U_{\text {ad }}$, where int $U_{\text {ad }}$ denotes the interior part of the set $U_{\text {ad }}$.

Then, for any admissible variation $\delta u$ of $u$, the limit (4.12) exists in $L_{T}^{\infty}\left(L^{2}\right)$. In particular, the control-to-state map $G$ is Gâteaux differentiable at $u$. Moreover, $\delta_{\delta u} G$ satisfies the Liouville problem

$$
\begin{equation*}
\partial_{t} \delta_{\delta u} G+\operatorname{div}\left(a(t, x ; u) \delta_{\delta u} G\right)=-\operatorname{div}(\bar{a}(t, x ; \delta u) G(u)), \quad \text { with } \quad \delta_{\delta u} G_{\mid t=0}=0 \tag{4.13}
\end{equation*}
$$

where we have defined $\bar{a}(t, x ; \delta u):=\delta u_{1}+x \circ \delta u_{2}$.
Proof. For any $0<|\varepsilon|<1$ small enough, let us define

$$
\delta G^{\varepsilon}:=\frac{1}{\varepsilon}(G(u+\varepsilon \delta u)-G(u)) .
$$

Following the computations which led to (4.5), we deduce that $\delta G^{\varepsilon}$ solves the equation

$$
\begin{equation*}
\partial_{t} \delta G^{\varepsilon}+\operatorname{div}\left(a(t, x ; u) \delta G^{\varepsilon}\right)=-\operatorname{div}(\bar{a}(t, x ; \delta u) G(u+\varepsilon \delta u)), \tag{4.14}
\end{equation*}
$$

with initial datum $\delta G_{\mid t=0}^{\varepsilon}=0$.
Notice that, by the first part of Lemma 4.1, we can find that, whenever $m \geq 1$ and $k \geq 1$, the sequence $\left(\delta G^{\varepsilon}\right)_{\varepsilon}$ is uniformly bounded in $L_{T}^{\infty}\left(L^{2}\right)$. Then, we can extract a subsequence, which converges weakly-* in this space to some $\rho \in L_{T}^{\infty}\left(L^{2}\right)$. Then, by weak compactness methods (as the ones used in the proof to Proposition 4.2), we can pass to the limit in the weak formulation of the previous equation and gather that $\rho$ solves the problem

$$
\begin{equation*}
\partial_{t} \rho+\operatorname{div}(a(t, x ; u) \rho)=-\operatorname{div}(\bar{a}(t, x ; \delta u) G(u)) \tag{4.15}
\end{equation*}
$$

with initial datum $\rho_{\mid t=0}=0$. Notice that the right-hand side of the previous equation belongs to $L^{2}$ by our assumptions on the initial data. Then, by uniqueness the whole sequence $\left(\delta G^{\varepsilon}\right)_{\varepsilon}$ has to weakly-* converge to $\rho$, and $\rho$ has to coincide with $\delta_{\delta u} G$.

Unfortunately, the previous argument does not prove the Gâteaux differentiability of $G$, because we need that the limit exists in the strong topology, namely in the $L_{T}^{\infty}\left(L^{2}\right)$ norm. In order to get this property, let us write the equation for $\delta G^{\varepsilon}-\rho$ : we find

$$
\partial_{t}\left(\delta G^{\varepsilon}-\rho\right)+\operatorname{div}\left(a(t, x ; u)\left(\delta G^{\varepsilon}-\rho\right)\right)=-\operatorname{div}(\bar{a}(t, x ; \delta u)(G(u+\varepsilon \delta u)-G(u)))
$$

with zero initial datum. For notational simplicity, define $\rho^{\varepsilon}:=\delta G^{\varepsilon}-\rho$; notice also that $G(u+\varepsilon \delta u)-$ $G(u)=\varepsilon \delta G^{\varepsilon}$. Then, an energy estimate for that equation gives us

$$
\begin{aligned}
\left\|\rho^{\varepsilon}(t)\right\|_{L^{2}} & \leq C \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right) \int_{0}^{t}\|\operatorname{div}(\bar{a}(\tau, x ; \delta u)(G(u+\varepsilon \delta u)-G(u)))\|_{L^{2}} d \tau \\
& \leq C \varepsilon \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right) \int_{0}^{t}|\delta u(\tau)|\left(\left\|\delta G^{\varepsilon}\right\|_{L^{2}}+\left\|(1+|x|) \nabla \delta G^{\varepsilon}\right\|_{L^{2}}\right) d \tau \\
& \leq C \varepsilon \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right) \int_{0}^{t}|\delta u(\tau)|\left\|\delta G^{\varepsilon}\right\|_{H_{1}^{1}} d \tau
\end{aligned}
$$

where we have argued as in the first line of (4.7) in order to pass from the first inequality to the second one. At this point, applying the second estimate of Lemma 4.1 to equation (4.14) yields

$$
\left\|\delta G^{\varepsilon}(\tau)\right\|_{H_{1}^{1}} \leq C_{0} \int_{0}^{\tau}|\delta u(s)| d s
$$

for any $\tau \in[0, t], t \leq T$, for a fixed constant $C_{0}$ (depending on $T, u^{a}, u^{b}$, and $\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)},\left\|\rho_{0}\right\|_{H_{2}^{2}}$ and $\left.\|g\|_{L_{T}^{1}\left(H_{2}^{2}\right)}\right)$. Putting this bound in the previous estimate entails

$$
\begin{aligned}
\left\|\rho^{\varepsilon}(t)\right\|_{L^{2}} & \leq C C_{0} \varepsilon \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right)\left(\int_{0}^{t}|\delta u(\tau)| d \tau\right)^{2} \\
& \leq C C_{0} \varepsilon T^{2}\|\delta u\|_{L_{T}^{\infty}}^{2} \exp \left(C \int_{0}^{T}\|\operatorname{div} a(t, x ; u)\|_{L^{\infty}}\right)
\end{aligned}
$$

From this last estimate, we deduce that, in the limit for $\varepsilon \rightarrow 0, \rho^{\varepsilon} \longrightarrow 0$ in $L_{T}^{\infty}\left(L^{2}\right)$. This completes the proof of the proposition.

Next, we tackle the proof of the Fréchet differentiability of $G$. The arguments will follow the proof of Proposition 4.3.
Theorem 4.1. Let $m \geq 2$ and $k \geq 2$. Let the data $\rho_{0}, g$ and $a_{0}$ be fixed as in Remark 4.1 above, and let $u \in \operatorname{int} U_{a d}$. Define $D G(u)[\delta u]$ to be the unique solution to equation (4.13).

Then there exists a constant $C>0$ (depending only on $T, u^{a}, u^{b}$, and $\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)},\left\|\rho_{0}\right\|_{H_{2}^{2}}$ and $\left.\|g\|_{L_{T}^{1}\left(H_{2}^{2}\right)}\right)$ such that

$$
\|G(u+\delta u)-G(u)-D G(u)[\delta u]\|_{L_{T}^{\infty}\left(L^{2}\right)} \leq C\|\delta u\|_{L_{T}^{\infty}}^{2} .
$$

In particular, the map $G$ is Fréchet differentiable from $\operatorname{int} U_{a d}$ into $L_{T}^{\infty}\left(L^{2}\right)$, and its Fréchet differential at any point $u \in \operatorname{int} U_{\text {ad }}$ is given by $D G(u)$.

Proof. In order to prove that $G$ is Fréchet differentiable, with Fréchet differential given by $D G(u)[\delta u]$, we have to show that

$$
\lim _{\|\delta u\|_{L_{T}^{\infty}}^{\infty} \rightarrow 0} \frac{\|G(u+\delta u)-G(u)-D G(u)[\delta u]\|_{L_{T}^{\infty}\left(L^{2}\right)}}{\|\delta u\|_{L_{T}^{\infty}}}=0
$$

We recall also that, if $G$ is Fréchet differentiable at $u$, then it is also Gâteaux differentiable at the same point, and one has $\delta_{\delta u} G=D G(u)[\delta u]$.

For simplicity, let us introduce the notation

$$
\mathcal{G}_{u}(\delta u):=G(u+\delta u)-G(u)-D G(u)[\delta u] .
$$

Remark that, in the proof of Proposition 4.3 above, we have already called $\rho$ the solution to equation (4.13), keep in mind equation (4.15). Therefore, the same computations performed on $\rho^{\alpha}=\delta G^{\alpha}-\rho$ this time lead us to an equation for $\mathcal{G}_{u}(\delta u)$ :

$$
\partial_{t} \mathcal{G}_{u}(\delta u)+\operatorname{div}\left(a(t, x ; u) \mathcal{G}_{u}(\delta u)\right)=-\operatorname{div}(\bar{a}(t, x ; \delta u)(G(u+\delta u)-G(u))),
$$

with initial datum $\mathcal{G}_{u}(\delta u)_{\mid t=0}=0$.
Next, it is just a matter of repeating the estimates performed on $\rho^{\alpha}$ : we easily find, for every $t \in[0, T]$, the inequality

$$
\left\|\mathcal{G}_{u}(\delta u)(t)\right\|_{L^{2}} \leq C \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right) \int_{0}^{t}\|\operatorname{div}(\bar{a}(\tau, x ; \delta u)(G(u+\delta u)-G(u)))\|_{L^{2}} d \tau
$$

$$
\leq C \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right) \int_{0}^{t}|\delta u(\tau)|\|G(u+\delta u)-G(u)\|_{H_{1}^{1}} d \tau
$$

Moreover, owing to the second inequality of Lemma 4.1, we get

$$
\|G(u+\delta u)-G(u)\|_{H_{1}^{1}} \leq C_{0} \int_{0}^{\tau}|\delta u(s)| d s
$$

for a constant $C_{0}$ which depends, as before, only on $T, u^{a}, u^{b},\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)},\left\|\rho_{0}\right\|_{H_{2}^{2}}$ and $\|g\|_{L_{T}^{1}\left(H_{2}^{2}\right)}$. Inserting this relation in the previous estimate, we find

$$
\left\|\mathcal{G}_{u}(\delta u)(t)\right\|_{L^{2}} \leq C C_{0} \exp \left(C \int_{0}^{t}\|\operatorname{div} a(\tau, x ; u)\|_{L^{\infty}}\right)\left(\int_{0}^{t}|\delta u(\tau)| d \tau\right)^{2} \leq K\|\delta u\|_{L_{T}^{\infty}}^{2}
$$

for a new positive constant $K$. From this last inequality, the claims of the theorem follow.

## 5 Analysis of the Liouville optimal control problem

In this section, we investigate our Liouville ensemble optimal control problem. In the first part, after recalling the problem's setting, we prove the existence of optimal controls by means of classical arguments. However, notice that one has to carefully justify that the reduced functional $\widehat{J}$ (see its definition below) is weakly lower semi-continuous. In fact, this property is not obvious, since $\rho=G(u)$ depends non-linearly on $u$. After that, in Section 5.2 we characterise optimal controls as solutions of a related first-order optimality system. In Section 5.3 we discuss uniqueness of optimal controls.

### 5.1 Existence of optimal controls

In this section, we deal with existence of optimal solutions to an ensemble optimal control problem. Our analysis is based on the following assumptions.
(A.1) We fix $(m, k) \in \mathbb{N}^{2}$, and we take an initial datum $\rho_{0} \in H_{k}^{m}\left(\mathbb{R}^{d}\right)$, a force $g \in L^{1}\left([0, T] ; H_{k}^{m}\left(\mathbb{R}^{d}\right)\right)$ and a vector field $a_{0} \in L^{1}\left([0, T] ; C^{m+1}\left(\mathbb{R}^{d}\right)\right)$, with $\nabla a_{0} \in L^{1}\left([0, T] ; C_{b}^{m}\left(\mathbb{R}^{d}\right)\right)$.
(A.2) We fix parameters $(\gamma, \delta, \nu) \in \mathbb{R}^{3}$ such that $\gamma>0, \delta \geq 0$ and $\nu \geq 0$.
(A.3) Chosen $u^{a}=\left(u_{1}^{a}, u_{2}^{a}\right)$ and $u^{b}=\left(u_{1}^{b}, u_{2}^{b}\right)$ in $\mathbb{R}^{2 d}$, with $u^{a} \leq u^{b}$, we define the set of admissible controls to be

$$
\begin{align*}
U_{a d} & :=\left\{u \in \mathbb{L}_{T}^{\infty}\left(\mathbb{R}^{d}\right) \mid \quad u^{a} \leq u(t) \leq u^{b} \quad \text { for a.e. } t \in[0, T]\right\} \quad \text { if } \quad \nu=0  \tag{5.1}\\
U_{a d} & :=\left\{u \in \mathbb{H}_{T}^{1}\left(\mathbb{R}^{d}\right) \mid \quad u^{a} \leq u(t) \leq u^{b} \quad \text { for all } t \in[0, T]\right\} \quad \text { if } \quad \nu>0 \tag{5.2}
\end{align*}
$$

(A.4) Finally, we take two attracting potentials $\theta$ and $\varphi$ in $L^{2}\left(\mathbb{R}^{d}\right)$, in the sense specified in Section 2.2 .

Remark 5.1. We point out that assumption (A.4) (which will be strengthened in Section 5.3 for getting uniqueness, see condition (A.4)* there) is taken for simplicity of presentation, since more general $\theta$ and $\varphi$ can be considered in our framework. For instance, we can allow for $\theta$ to depend on time: $\theta \in L_{T}^{1}\left(L^{2}\right)$, or $\theta \in L_{T}^{1}\left(H_{1}^{1}\right)$ in $(\boldsymbol{A} .4)^{*}$ below. The case $\theta(x)=|x|^{2}$ and $\varphi(x)=|x|^{2}$ is more delicate, and will be matter of further discussions in Section 5.4.

Now, consider our cost functional given by

$$
\begin{align*}
J(\rho, u):=\int_{0}^{T} & \int_{\mathbb{R}^{d}} \theta(x) \rho(x, t) d x d t+\int_{\mathbb{R}^{d}} \varphi(x) \rho(x, T) d x  \tag{5.3}\\
& \quad+\frac{\gamma}{2} \int_{0}^{T}|u(t)|^{2} d t+\delta \int_{0}^{T}|u(t)| d t+\frac{\nu}{2} \int_{0}^{T}\left|\frac{d}{d t} u(t)\right|^{2} d t
\end{align*}
$$

Remark that $J$ is well-defined whenever $u \in \mathbb{L}_{T}^{2}$ if $\nu=0$, or $u \in \mathbb{H}_{T}^{1}$ if $\nu>0$, and $\rho \in C\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)$.
Our ensemble optimal control problem requires to find

$$
\begin{equation*}
\min _{u \in U_{a d}} J(\rho, u) \tag{5.4}
\end{equation*}
$$

subject to the differential constraint

$$
\begin{cases}\partial_{t} \rho+\operatorname{div}(a(t, x ; u) \rho)=g & \text { in } \quad[0, T] \times \mathbb{R}^{d}  \tag{5.5}\\ \rho_{\mid t=0}=\rho_{0} & \text { on } \quad \mathbb{R}^{d},\end{cases}
$$

where the drift function $a(t, x ; u)$ is defined as

$$
\begin{equation*}
a(t, x ; u):=a_{0}(t, x)+u_{1}(t)+x \circ u_{2}(t) . \tag{5.6}
\end{equation*}
$$

Under our assumptions, Theorem 3.4 applies. Thus, for every $u \in U_{a d}$, there exists a unique solution solution $\rho \in C\left([0, T] ; H_{k}^{m}\left(\mathbb{R}^{d}\right)\right)$ to the Liouville problem (5.5) corresponding to that $u$. Therefore, resorting to the control-to-state map $G$, as defined in Section 4, we can introduce the so-called reduced cost functional, given by

$$
\begin{equation*}
\widehat{J}(u):=J(G(u), u) . \tag{5.7}
\end{equation*}
$$

Hence, the ensemble optimal control problem (5.4)-(5.5) can be rephrased as follows:

$$
\begin{equation*}
\min _{u \in U_{a d}} \widehat{J}(u) \tag{5.8}
\end{equation*}
$$

Remark 5.2. Recall that we have defined $G$ with values in $L_{T}^{\infty}\left(L^{2}\right)$. However, under our assumptions, we know that the solution to the Liouville equation actually belongs to $C_{T}\left(L^{2}\right)$, so that the $\varphi$-term in (5.3) is well-defined, and thus so is $\widehat{J}$.

In the following, we prove existence of a minimizer to (5.8).
Theorem 5.1. Under assumptions (A.1)-(A.2)-(A.3)-(A.4), the ensemble optimal control problem (5.8) admits at least one solution $u^{*} \in U_{\text {ad }}$. The corresponding state $\rho^{*}:=G\left(u^{*}\right)$ belongs to the space $C\left([0, T] ; H_{k}^{m}\left(\mathbb{R}^{d}\right)\right)$.

Proof. Let us focus on the case $\nu=0$ for simplicity; the case $\nu>0$ follows from the same token.
As already mentioned, the functional $J$ given in (5.3) is well-defined for $(\rho, u) \in C_{T}\left(L^{2}\right) \times \mathbb{L}_{T}^{\infty}$. Now, we remark that $U_{a d}$ is a bounded subset of $\mathbb{L}_{T}^{\infty}$. On the other hand, owing to Theorem 3.4, see especially estimate (3.17), and the embedding $C_{T}\left(H_{K}^{m}\right) \hookrightarrow L_{T}^{\infty}\left(L^{2}\right)$, the map $G$ takes its values in a bounded set of $L_{T}^{\infty}\left(L^{2}\right)$. It follows that $\widehat{J}$ is bounded; in particular, $\widehat{J}$ is a proper map, i.e. $\inf _{U_{a d}} \widehat{J}>-\infty$, and $\widehat{J}$ is not identically equal to $+\infty$.

Next, we claim that $\widehat{J}$ is weakly lower semi-continuous. To prove this fact, it is enough to use the weak-weak continuity of $G$, as stated in Proposition 4.2, and to remark that $J$ is weakly lower semicontinuous. Indeed, the last three terms in (5.3) are norms, so they are weakly lower semi-continuous.

On the other hand, the first two terms are linear in $\rho$, and then they are weakly continuous with respect to the $L_{T}^{\infty}\left(L^{2}\right)$ and $L^{2}$ topologies, respectively. Thus we immediately get that, if $\left(u_{n}\right)_{n} \subset U_{a d}$ is a sequence which converges weakly-* to a $u \in U_{a d}$ in $L_{T}^{\infty}$, we have

$$
\liminf _{n \rightarrow+\infty} \widehat{J}\left(u_{n}\right)=\liminf _{n \rightarrow+\infty} J\left(G\left(u_{n}\right), u_{n}\right) \geq J(G(u), u)=\widehat{J}(u) .
$$

At this point, proving the existence of a minimizer for $\widehat{J}$ is standard. Let us take a minimizing sequence $\left(u_{n}\right)_{n} \subset U_{a d}$. Since $U_{a d}$ is a bounded set in $\mathbb{L}_{T}^{\infty}$, we can extract a weakly-* convergent subsequence, which we do not relabel for simplicity; let us call $u^{*} \in U_{a d}$ its limit-point. Then, by the weak-lower semi-continuity of $\widehat{J}$, we can conclude that $u^{*}$ is a minimizer for $\widehat{J}$.

We discuss uniqueness of the minimizers in Section 5.3 below. For this purpose, we use characterization of minimizers as solutions to a suitable optimality system, which we derive in the next section.

### 5.2 Liouville optimality systems

This section is devoted to the characterization of ensemble optimal controls as solutions of the related first-order optimality system. For this purpose, in addition to hypotheses (A.1)-(A.2)-(A.3)-(A.4) stated above, from now on we take

$$
m \geq 1 \quad \text { and } \quad k \geq 1
$$

In correspondence to (5.3)-(5.4)-(5.5), we consider the Lagrange multipliers framework, see e.g. [22, 27], and introduce the Lagrange functional $\mathcal{L}$ as follows:

$$
\begin{equation*}
\mathcal{L}(\rho, u, q):=J(\rho, u)+\int_{0}^{T} \int_{\mathbb{R}^{d}}\left(\partial_{t} \rho+\operatorname{div}(a(x, t ; u) \rho)-g\right) q d x d t+\int_{\mathbb{R}^{d}}\left(\rho(0, x)-\rho_{0}(x)\right) q_{0}(x) d x \tag{5.9}
\end{equation*}
$$

where, for the sake of generality, we have included a right-hand side $g$. The variable $q$ represents the Lagrange multiplier. Notice that $\mathcal{L}$ is well-defined whenever $u \in \mathbb{L}_{T}^{\infty}$ if $\nu=0, u \in \mathbb{H}_{T}^{1}$ if $\nu>0, q \in L_{T}^{\infty}\left(L^{2}\right)$, $q_{0} \in L^{2}$ and $\rho \in C_{T}\left(L^{2}\right)$ such that both $\partial_{t} \rho$ and $\operatorname{div}(a(x, t ; u) \rho)$ belong to $L_{T}^{1}\left(L^{2}\right)$. In particular, it is enough to have $\rho \in W_{T}^{1,1}\left(L^{2}\right) \cap L_{T}^{\infty}\left(H_{1}^{1}\right)$, recall also Proposition 3.3. Notice that, a posteriori, we will find $q \in C_{T}\left(L^{2}\right)$ and $q_{0}=q(0)$; see the discussion below for details.

In order to derive the optimality system, let us discuss different instances. For clarity, we first discuss the case with $L^{2}$ costs only, then the case with $L^{2}-H^{1}$ costs, and finally the case with $L^{2}-L^{1}-H^{1}$ costs.

The case $\delta=\nu=\mathbf{0}$. If $\delta=0$, then $J$ is Fréchet differentiable over $C_{T}\left(L^{2}\right) \times \operatorname{int} U_{a d}$, since it is linear in $\rho$ and the control costs with $\gamma>0, \nu \geq 0$ are given by differentiable norms. It is then an easy computation to show that $\mathcal{L}$ is Fréchet differentiable over the space

$$
\mathbb{X}_{T}:=\left(W_{T}^{1,1}\left(L^{2}\right) \cap L_{T}^{\infty}\left(H_{1}^{1}\right)\right) \times \mathbb{L}_{T}^{2} \times C_{T}\left(L^{2}\right)
$$

where $\mathbb{L}_{T}^{2}$ has to be replaced by $\mathbb{H}_{T}^{1}$ in the case when $\nu>0$. The Fréchet differential of $\mathcal{L}$ at $(\rho, u, q)$ is given by the linearization of each of its terms at that point.

Now, consider in addition $\nu=0$. The optimality system is obtained by putting to zero the Fréchet derivatives of $\mathcal{L}(\rho, u, q)$ with respect to each of its arguments separately. We obtain

$$
\begin{align*}
& \partial_{t} \rho+\operatorname{div}(a(x, t ; u) \rho)=g, \quad \text { with } \quad \rho_{\mid t=0}=\rho_{0}  \tag{5.10}\\
& -\partial_{t} q-a(x, t ; u) \cdot \nabla q=-\theta,  \tag{5.11}\\
& \text { with } \quad q_{\mid t=T}=-\varphi
\end{align*}
$$

$$
\begin{equation*}
\left(\gamma u_{j}^{r}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u_{j}^{r}} \rho\right) q d x, v_{j}^{r}-u_{j}^{r}\right)_{L^{2}(0, T)} \geq 0 \quad \forall v \in U_{a d}, j=1,2, r=1 \ldots d \tag{5.12}
\end{equation*}
$$

We remark that, denoting by $e^{r}$ the $r$-th unit vector of the canonical basis of $\mathbb{R}^{d}$ and by $x^{r}$ the $r$-th component of the vector $x \in \mathbb{R}^{d}$, by Definition 5.6 we have

$$
\frac{\partial a}{\partial u_{j}^{r}}(t, x ; u)=e^{r} \quad \text { for } j=1, \quad \frac{\partial a}{\partial u_{j}^{r}}(t, x ; u)=x^{r} e^{r} \quad \text { for } j=2
$$

Then, equation (5.12) can be equivalently written in the following form: for any $1 \leq r \leq d$,

$$
\left\{\begin{array}{l}
\left(\gamma u_{1}^{r}+\int_{\mathbb{R}^{d}} \partial_{r} \rho q d x, v_{1}^{r}-u_{1}^{r}\right)_{L^{2}(0, T)} \geq 0 \\
\left(\gamma u_{2}^{r}+\int_{\mathbb{R}^{d}} \partial_{r}\left(x^{r} \rho\right) q d x, v_{2}^{r}-u_{2}^{r}\right)_{L^{2}(0, T)} \geq 0
\end{array}\right.
$$

Further, if we sum up equations (5.12) for all $m$ and all $r$, we can write, in the following compact form

$$
\begin{equation*}
\left(\gamma u+\int_{\mathbb{R}^{d}} \operatorname{div}((e+x) \rho) q d x, v-u\right)_{\mathbb{L}_{T}^{2}} \geq 0 \quad \text { for all } \quad v \in U_{a d} \tag{5.13}
\end{equation*}
$$

where we have defined the vector $e=(1 \ldots 1)$.
Equation (5.10) is our Liouville model (also called the forward equation in this context). The results of Section 3.2 guarantee that, under our assumptions, there exists a unique solution $\rho \in C_{T}\left(H_{1}^{1}\right)$. Moreover, since $u \in U_{a d}$, an inspection of (5.10) reveals that $\partial_{t} \rho \in L_{T}^{1}\left(L^{2}\right)$.

Equation (5.11) is the adjoint Liouville equation; it is obtained by taking the Fréchet derivative of (5.9) with respect to $\rho$. This is a transport equation that evolves backwards in time. By setting $\widetilde{q}(t, x)=q(T-t,-x)$, we obtain a transport problem for $\widetilde{q}$, as in (3.14), with source term $-\theta$ and initial condition $\widetilde{q}_{\mid t=0}=-\varphi$. Thus, the results of Paragraph 3.1.2 guarantee the existence and uniqueness of a Lagrange multiplier $q \in C_{T}\left(L^{2}\right)$, provided that $\theta$ and $\varphi$ are in $L^{2}$.

From the discussion above, we get that any solution to the optimality system (5.10)-(5.11)-(5.12), with $u \in U_{a d}$, belongs indeed to the space $\mathbb{X}_{T}$.

Equation (5.12) represents the optimality condition. To better illustrate this fact, we suppose from now on that

$$
m \geq 2 \quad \text { and } \quad k \geq 2
$$

Then, the reduced cost functional $\widehat{J}$, defined in (5.7), is Fréchet differentiable; in terms of the reduced minimisation problem (5.8), the optimal solution $u^{*}$ in the convex, closed and bounded set $U_{a d}$ is characterized by the optimality condition given by

$$
\left(\nabla_{u} \widehat{J}\left(u^{*}\right), v-u^{*}\right)_{\mathbb{L}_{T}^{2}} \geq 0, \quad \text { for all } v \in U_{a d}
$$

where $\nabla_{u} \widehat{J}$ denotes the $L^{2}$-gradient of $\widehat{J}$ with respect to $u$. In fact, a direct computation of $\nabla_{u} J(G(u), u)$, with the introduction of the auxiliary adjoint variable $q$, gives the optimality system above, and the following relation:

$$
\nabla_{u_{j}^{r}} \widehat{J}(u)=\gamma u_{j}^{r}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u_{j}^{r}} \rho\right) q d x .
$$

The case $\delta=\mathbf{0}, \nu>\mathbf{0}$. Next, assume that $\delta=0$ and $\gamma, \nu>0$. Recall that, in this case, the set $U_{a d}$ is defined by (5.2). Then, the natural Hilbert space where $u^{*}$ is sought is $\widetilde{\mathbb{H}}_{T}^{1}\left(\mathbb{R}^{d}\right):=\widetilde{H}_{T}^{1}\left(\mathbb{R}^{d}\right) \times \widetilde{H}_{T}^{1}\left(\mathbb{R}^{d}\right)$, where $\widetilde{H}_{T}^{1}$ corresponds to the $H_{T}^{1}$ space, endowed with the weighted $H^{1}$-product given by

$$
(u, v)_{\widetilde{H}_{T}^{1}}:=\gamma \int_{0}^{T} u(t) \cdot v(t) d t+\nu \int_{0}^{T} u^{\prime}(t) \cdot v^{\prime}(t) d t
$$

The notation ${ }^{\prime}=d / d t$ stands for the weak time derivative.
Now, let $\mu$ be the $\widetilde{H}^{1}$-Riesz representative of the continuous linear functional

$$
v \mapsto\left(\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u} \rho\right) q d x, v\right)_{\mathbb{L}_{T}^{2}}
$$

Assuming that $u \in U_{a d} \cap H_{0}^{1}\left([0, T] ; \mathbb{R}^{2 d}\right)$, then $\mu$ can be computed by solving the equation

$$
\begin{equation*}
\left(-\nu \frac{d^{2}}{d t^{2}}+\gamma\right) \mu=\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u} \rho\right) q d x, \quad \mu(0)=\mu(T)=0 \tag{5.14}
\end{equation*}
$$

which is understood in a weak sense. Notice that the choice $u \in H_{0}^{1}\left([0, T] ; \mathbb{R}^{2 d}\right)$ corresponds to the modelling requirement that the control is switched on at $t=0$ and switched off at $t=T$. Other initial and final time conditions on $u$ may be required and encoded as boundary conditions in (5.14).

With the setting above, the $\widetilde{H}^{1}$-gradient is given, for $j=1,2$ and all $r=1 \ldots d$, by

$$
\begin{equation*}
\widetilde{\nabla}_{u_{j}^{r}} \widehat{J}(u)=u_{j}^{r}+\mu_{j}^{r} . \tag{5.15}
\end{equation*}
$$

The optimality condition (5.12) then becomes

$$
\begin{equation*}
\left(u_{j}^{r}+\mu_{j}^{r}, v_{j}^{r}-u_{j}^{r}\right)_{\widetilde{H}_{T}^{1}} \geq 0 \tag{5.16}
\end{equation*}
$$

for all $v \in U_{a d}$, where $U_{a d}$ is given in (5.2), $j=1,2$ and $1 \leq r \leq d$.
The case $\delta>\mathbf{0}$. In this case, a $L^{1}$ norm of the control appears in the cost functional. This term is not Gâteaux differentiable and the discussion becomes more involved. By using of the control-to-state map, we start by defining

$$
\begin{aligned}
f(u) & :=\int_{0}^{T} \int_{\mathbb{R}^{d}} \theta(x) G(u)(x, t) d x d t+\int_{\mathbb{R}^{d}} \varphi(x) G(u)(x, T) d x+\frac{\gamma}{2} \int_{0}^{T}|u(t)|^{2} d t+\frac{\nu}{2} \int_{0}^{T}\left|\frac{d}{d t} u(t)\right|^{2} d t \\
g(u) & :=\delta\|u\|_{L_{T}^{1}} .
\end{aligned}
$$

The $L^{1}$-cost, represented by $g$, admits a subdifferential $\partial g(u)=\delta \partial\left(\|u\|_{L^{1}}\right)$, see e.g. Section 2.3 of [5]. If we denote by $\mathbb{L}_{T}^{*}:=\left(\mathbb{L}_{T}^{\infty}\left(\mathbb{R}^{d}\right)\right)^{*}$ and by $\langle\cdot, \cdot\rangle$ the duality product in $\mathbb{L}_{T}^{*} \times \mathbb{L}_{T}^{\infty}$, the following formula holds true:

$$
\begin{align*}
& \partial\left(\|u\|_{L^{1}}\right)=\left\{\phi \in \mathbb{L}_{T}^{*} \mid \quad\|v\|_{L^{1}}-\|u\|_{L^{1}} \geq\langle\phi, v-u\rangle \quad \forall v \in U_{a d}\right\}  \tag{5.17}\\
& =\left\{\begin{array}{lll}
\phi \in \mathbb{L}_{T}^{*} \mid \quad\|\phi\|_{\mathbb{L}_{T}^{*}}=1, \phi(u)=\|u\|_{\mathbb{L}_{T}^{\infty}} & \text { if } & u \not \equiv 0 \\
\text { unit ball in } \mathbb{L}_{T}^{*} & \text { if } & u \equiv 0
\end{array}\right\} .
\end{align*}
$$

Now, the reduced functional can be written as $\widehat{J}(u)=f(u)+g(u)$. In this case, the equations (5.10) and (5.11) in the corresponding optimality system are the same; however, we have a different optimality condition (5.12). In the case $\nu=0$, as in Theorem 2.2 in [14], we have the following result; for its proof, we refer to $[14]$ and $[26]$. Notice that, as for equations (5.10)-(5.11)-(5.12), equation (5.18) below can be written even when $G$, and hence $\widehat{J}$, are not Fréchet differentiable.

Theorem 5.2. Under assumptions (A.1)-(A.2)-(A.3)-(A.4), where we take $m \geq 1$ and $k \geq 1$, suppose moreover that the pair $(\rho, u) \in C_{T}\left(H_{k}^{m}\right) \times U_{a d}$ is a minimizer for (5.8).

Then there exists a unique $q \in C_{T}\left(L^{2}\right)$ which solves (5.11), and a $\widehat{\lambda} \in \partial g(u)$ such that the following inequality condition is satisfied:

$$
\begin{equation*}
\left(\gamma u_{j}^{r}+\widehat{\lambda}_{j}^{r}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u_{j}^{r}} \rho\right) q d x, v_{j}^{r}-u_{j}^{r}\right)_{L^{2}(0, T)} \geq 0 \quad \forall v \in U_{a d}, j=1,2, r=1 \ldots d \tag{5.18}
\end{equation*}
$$

Moreover, there exist $\lambda_{+}$and $\lambda_{-}$, belonging to $L_{T}^{\infty}\left(\mathbb{R}^{d}\right)$, such that (5.18) is equivalent to the equations

$$
\left\{\begin{array}{l}
\gamma u_{j}^{r}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u_{j}^{r}} \rho\right) q d x+\left(\lambda_{+}\right)_{j}^{r}-\left(\lambda_{-}\right)_{j}^{r}+\widehat{\lambda}_{j}^{r}=0 \\
\left(\lambda_{+}\right)_{j}^{r} \geq 0, \quad u^{b}-u_{j}^{r} \geq 0, \quad\left(\lambda_{+}\right)_{j}^{r}\left(u^{b}-u_{j}^{r}\right)=0 \\
\left(\lambda_{-}\right)_{j}^{r} \geq 0, \quad u_{j}^{r}-u^{a} \geq 0, \quad\left(\lambda_{-}\right)_{j}^{r}\left(u_{j}^{r}-u^{a}\right)=0 \\
\widehat{\lambda}_{j}^{r}=\delta \quad \text { a.e. in } \quad\left\{t \in[0, T] \mid \quad u_{j}^{r}(t)>0\right\} \\
\left|\widehat{\lambda}_{j}^{r}\right| \leq \delta \quad \text { a.e. in } \quad\left\{t \in[0, T] \mid u_{j}^{r}(t)=0\right\} \\
\widehat{\lambda}_{j}^{r}=\delta \quad \text { a.e.in } \quad\left\{t \in[0, T] \mid \quad u_{j}^{r}(t)<0\right\},
\end{array}\right.
$$

for $j=1,2$ and all $1 \leq r \leq d$.
Remark 5.3. In our case, $\widehat{\lambda}_{j}^{r}$ can be understood to be $\delta \operatorname{sgn}\left(u_{j}^{r}\right)$, where $\operatorname{sgn}(x)$ is the sign function, equal to 1 or -1 depending if $x>0$ or $x<0$ respectively, and equal to 0 if $x=0$.

Furthermore, we notice that the additional Lagrange multipliers $\left(\lambda_{ \pm}\right)_{j}^{r}$ are due to the constraints $u^{a} \leq u(t) \leq u^{b}$ for almost all $t \in[0, T]$.

Finally, the case $\delta>0$ and $\nu>0$ can be treated as done before. After resorting once again to the space $\widetilde{\mathbb{H}}_{T}^{1}$, let $\mu$ be the $\widetilde{H}^{1}$-Riesz representative of the continuous linear functional

$$
v \mapsto\left(\hat{\lambda}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u} \rho\right) q d x, v\right)_{\mathbb{L}_{T}^{2}}
$$

Then, assuming that $u \in U_{a d} \cap H_{0}^{1}\left([0, T] ; \mathbb{R}^{2 d}\right)$, we can compute $\mu$ as above, by solving the equation

$$
\left(-\nu \frac{d^{2}}{d t^{2}}+\gamma\right) \mu=\widehat{\lambda}+\int_{\mathbb{R}^{d}} \operatorname{div}\left(\frac{\partial a}{\partial u} \rho\right) q d x, \quad \mu(0)=\mu(T)=0
$$

which has to be understood again in a weak sense. With this definition, relation (5.15) still holds true, and the optimality condition (5.12) can be expressed once again by equations (5.16).

### 5.3 Uniqueness of optimal controls

In this section, we tackle the problem of uniqueness of optimal controls. Specifically, we are able to prove uniqueness in the situation when $\delta=0$ and $\nu=0$ in (5.3). Indeed, our proof relies on the characterization of optimal controls as solutions to the corresponding optimality system, and especially on the use of the optimality condition (5.12). For this reason, the cases $\delta>0$ or $\nu>0$ read more complicated, recall equation (5.16) and Theorem 5.2 above, and are left aside in our discussion.

We start with discussing uniqueness for the unconstrained-control problem. Then, in Section 5.3.2 we prove uniqueness for the constrained problem, under some additional assumptions.

### 5.3.1 Uniqueness for the unconstrained-control problem

We consider the unconstrained-control optimization problem, with the assumptions (A.1)-(A.2)-(A.4).
The first problem to deal with is then to prove the existence of a global minimizer: it is apparent that the proof of Theorem 5.1 does not work anymore, since it relies on the uniform boundedness of the set $U_{a d}$ in an essential way. For solving this issue, first of all we have to quit the $\mathbb{L}_{T}^{\infty}$ framework, and rather work with controls which are merely in $\mathbb{L}_{T}^{2}$; this is however quite natural, in view of the form of (5.3).

Moreover, we need additional assumptions: we require that both $\rho_{0}$ and $g$ are bounded from below by some constant, that we can suppose, without loss of generality, to be 0 (a case which is physically relevant, recall the discussion in Section 2.1). Moreover, we need also a lower bound on the optimization functions $\theta$ and $\varphi$ : defined $c \in \mathbb{R}$ such a lower bound, up to working with $\theta-c$ and $\varphi-c$ and taking $k$ large enough in (A.1), so that both $\rho_{0}$ and $g$ belong to $L^{1}$ (and so does $\rho$, then), we can assume that also $c=0$.

Our additional hypotheses then read as follows.
(A.5) Suppose that $\rho_{0} \geq 0$ and $g \geq 0$.
(A.6) We assume also that $\theta \geq 0$ and $\varphi \geq 0$.

Under these hypotheses, we are able first of all to prove the existence of a global minimizer.
Proposition 5.1. Under assumptions (A.1)-(A.2)-(A.4)-(A.5)-(A.6) and the conditions $\delta=\nu=0$, the unconstrained ensemble optimal control problem

$$
\min _{u \in \mathbb{L}_{T}^{2}} \widehat{J}(u)
$$

has at least one solution $u^{*} \in \mathbb{L}_{T}^{2}$.
Proof. By assumption (A.5), it is standard to deduce that any $\rho$ solution to (5.5), for any fixed $u$, is positive almost everywhere on $[0, T] \times \mathbb{R}^{d}$. Then, if moreover (A.6) holds, then $\widehat{J}(u) \geq 0$ for any $u \in \mathbb{L}_{T}^{2}$, so $\widehat{J}$ is in particular a proper map.

Next, we claim that $\widehat{J}$ is a coercive map. Indeed, let $\left(u_{n}\right)_{n}$ be a sequence in the space $\mathbb{L}_{T}^{2}$, such that $\left\|u_{n}\right\|_{\mathbb{L}_{T}^{2}} \longrightarrow+\infty$ for $n \rightarrow+\infty$. Then, by (5.3), we have that $\widehat{J}\left(u_{n}\right)$ must explode to $+\infty$.

Finally, proving the weak lower semi-continuity of $\widehat{J}$ requires just small adaptations to the proof of Theorem 5.1. In fact, it is easy to see that Proposition 4.2 holds true even if we replace the weak-* convergence in $\mathbb{L}_{T}^{\infty}$ with the weak convergence in $\mathbb{L}_{T}^{2}$ (this is due to the fact that $W_{T}^{1,1}$ is compactly embedded into $L_{T}^{q}$ for any $\left.1 \leq q<+\infty\right)$.

With these ingredients at hand, and after remarking that the coercivity of $\widehat{J}$ implies that any minimizing sequence has to remain uniformly bounded in $\mathbb{L}_{T}^{2}$, showing the existence of a minimizer also follows the same lines of the proof to Theorem 5.1.

In order to prove uniqueness, we need additional regularity on the cost functions $\theta$ and $\varphi$. We then formulate the following assumption, which strengthen (A.4).
(A.4)* Suppose that both $\theta$ and $\varphi$ belong to $H_{1}^{1}\left(\mathbb{R}^{d}\right)$.

The main result of this section reads as follows.
Theorem 5.3. Under assumptions (A.1)-(A.2)-(A.5)-(A.6)-(A.4)*, suppose also that both $m \geq 2$ and $k \geq 2$. In addition, take $\delta=\nu=0$ in (5.3).

Then, the optimal control $u^{*}$, whose existence is guaranteed by Proposition 5.1, is unique in $\mathbb{L}_{T}^{2}$.

Proof. Let $\left(u_{1}, \rho_{1}, q_{1}\right)$ and $\left(u_{2}, \rho_{2}, q_{2}\right)$ be two optimal controls with corresponding state and adjoint state. Then both triplets have to satisfy the optimality system (5.10)-(5.11), together with the equality

$$
\gamma u+\int_{\mathbb{R}^{d}} \operatorname{div}((e+x) \rho) q d x=0
$$

which replaces the optimality condition (5.12), since now the problem is unconstrained.
We now apply the previous equality to each of the optimal triplets and take the difference: after setting $\delta u=u_{1}-u_{2}$ and analogous notations for $\delta \rho$ and $\delta q$, we find

$$
\gamma \delta u+\int_{\mathbb{R}^{d}} \operatorname{div}((e+x) \delta \rho) q_{1} d x+\int_{\mathbb{R}^{d}} \operatorname{div}\left((e+x) \rho_{2}\right) \delta q d x=0
$$

which immediately implies, for almost every $t \in[0, T]$, the estimate

$$
\begin{equation*}
\gamma|\delta u(t)| \leq \int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho(t)) q_{1}(t)\right|+\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}(t)\right) \delta q(t)\right| . \tag{5.19}
\end{equation*}
$$

Let us now estimate each one of the integral terms. We start with the former term and obtain

$$
\begin{aligned}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho(t)) q_{1}(t)\right| d x & \leq\left\|q_{1}(t)\right\|_{L^{2}}\|\operatorname{div}((e+x) \delta \rho(t))\|_{L^{2}} \\
& \leq C_{1}\left(\|\delta \rho(t)\|_{L^{2}}+\|(1+|x|) \nabla \delta \rho(t)\|_{L^{2}}\right) \leq C_{1}\|\delta \rho(t)\|_{H_{1}^{1}}
\end{aligned}
$$

where, in passing from the first to the second inequality, we have computed explicitly the derivatives in the div term, and we have used Theorem 3.3 applied to the transport equation (5.11) for treating the $q_{1}$ term. Notice that the constant $C_{1}$ can be expressed as

$$
\begin{equation*}
C_{1}:=C \exp \left(C\left(\left\|\operatorname{div} a_{0}\right\|_{L_{T}^{1}\left(L^{\infty}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)\left(\|\varphi\|_{L^{2}}+T\|\theta\|_{L^{2}}\right), \tag{5.20}
\end{equation*}
$$

for a "universal" constant $C>0$ that depends on the space dimension $d$. At this point, we recall that both $\rho_{1}$ and $\rho_{2}$ satisfy equation (5.10), with controls $u_{1}$ and $u_{2}$, respectively. Then, taking their difference and applying Lemma 4.1 finally yields, for a new constant $\widetilde{C}_{1}=C_{1} K_{1}^{(2)}$ just depending on the data of the problem, the following bound:

$$
\begin{equation*}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho(t)) q_{1}(t)\right| d x \leq \widetilde{C}_{1} \int_{0}^{t}|\delta u(\tau)| d \tau \tag{5.21}
\end{equation*}
$$

Next, consider the second integral in (5.19). The computations are similar to the previous ones: first of all, we can estimate

$$
\begin{aligned}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}(t)\right) \delta q(t)\right| d x & \leq\|\delta q(t)\|_{L^{2}}\left\|\operatorname{div}\left((e+x) \rho_{2}(t)\right)\right\|_{L^{2}} \\
& \leq\|\delta q(t)\|_{L^{2}}\left\|\rho_{2}(t)\right\|_{H_{1}^{1}} \leq C_{2}\|\delta q(t)\|_{L^{2}}
\end{aligned}
$$

where this time we have applied Theorem 3.4 to equation (5.10) for $\rho_{2}$ to control its $H_{1}^{1}$ norm. In particular, it follows from that theorem that

$$
\begin{equation*}
C_{2}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)\left(\left\|\rho_{0}\right\|_{H_{1}^{1}}+\|g\|_{L_{T}^{1}\left(H_{1}^{1}\right)}\right) \tag{5.22}
\end{equation*}
$$

for a "universal" constant $C>0$.

As done above, we use now the fact that $q_{1}$ and $q_{2}$ are both solutions of (5.11), related to the controls $u_{1}$ and $u_{2}$ respectively. Hence, taking the difference of those equations and arguing as in the proof of Lemma 4.1 (keep in mind also Remark 3.3), one easily infers the existence of a "universal" constant $C>0$ such that

$$
\begin{aligned}
\|\delta q(t)\|_{L^{2}} & \leq C \exp \left(C\left(\left\|\operatorname{div} a_{0}\right\|_{L_{T}^{1}\left(L^{\infty}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}\right)\right) \int_{0}^{t}|\delta u(\tau)|\left\|(1+|x|) \nabla q_{2}(\tau)\right\|_{L^{2}} d \tau \\
& \leq C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)\left(\|\varphi\|_{H_{1}^{1}}+T\|\theta\|_{H_{1}^{1}}\right) \int_{0}^{t}|\delta u(\tau)| d \tau .
\end{aligned}
$$

After defining the constants

$$
\begin{equation*}
\widetilde{K}_{1}^{(1)}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)\left(\|\varphi\|_{H_{1}^{1}}+T\|\theta\|_{H_{1}^{1}}\right) \tag{5.23}
\end{equation*}
$$

and $\widetilde{C}_{2}:=C_{2} \widetilde{K}_{1}^{(1)}$, we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}(t)\right) \delta q(t)\right| d x \leq \widetilde{C}_{2} \int_{0}^{t}|\delta u(\tau)| d \tau \tag{5.24}
\end{equation*}
$$

At this point, we can insert estimates (5.21) and (5.24) into (5.19), and get, for a new constant $K=\widetilde{C}_{1}+\widetilde{C}_{2}$, the relation

$$
\gamma|\delta u(t)| \leq K \int_{0}^{t}|\delta u(\tau)| d \tau
$$

An application of Grönwall's lemma hence gives that $\delta u \equiv 0$ almost everywhere on $[0, T]$. This concludes the proof of the theorem.

### 5.3.2 The constrained optimization problem

In this section, we consider our optimization problem with constraints on the control $u \in U_{a d}$, where the call $U_{a d}$ is defined by (5.1). For the reasons explained above, we still restrict to the case $\delta=\nu=0$.

In the constrained-control case, the characterization of optimal controls is given by an inequality, see (5.12). This is a very weak information: this is the reason why we are able to prove uniqueness only under a smallness condition, either on the time $T$ or on the size of the data $\rho_{0}, g, \nabla a_{0}, \theta$ and $\varphi$ in their respective functional spaces.

Let us recall that existence of an optimal control has been proved in Theorem 5.1 above. We can now state our uniqueness result.

Theorem 5.4. Under assumptions (A.1)-(A.2)-(A.3)-(A.4)*, suppose that both $m \geq 2$ and $k \geq 2$. Take moreover $\delta=\nu=0$ in (5.3). Finally, define

$$
\widetilde{K}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)}+T \max \left\{\left|u^{a}\right|,\left|u^{b}\right|\right\}\right)\right)\left(\left\|\rho_{0}\right\|_{H_{2}^{2}}+\|g\|_{L_{T}^{1}\left(H_{2}^{2}\right)}\right)\left(\|\varphi\|_{H_{1}^{1}}+T\|\theta\|_{H_{1}^{1}}\right),
$$

where the constant $C>0$ can be taken as the maximum of the constants $C$ appearing in (5.20), (5.22), (5.23) and in the definition (4.4) of $K_{1}^{(2)}$.

If the condition

$$
\frac{\widetilde{K} T}{\gamma}<2
$$

holds true, then there exists at most one optimal control $u^{*}$ in $\operatorname{int} U_{\text {ad }}$.

Proof. As above, let $\left(u, \rho_{1}, q_{1}\right)$ and $\left(v, \rho_{2}, q_{2}\right)$ be two optimal triplets solving the minimization problem (5.8). Then both have to satisfy the optimality system (5.10)-(5.11)-(5.12). In particular, the inequality in (5.12), written in the more explicit form (5.13), gives

$$
\begin{array}{ll}
\left(\gamma u+\int_{\mathbb{R}^{d}} \operatorname{div}\left((e+x) \rho_{1}\right) q_{1}, u-w\right)_{\mathbb{L}_{T}^{2}} \leq 0 & \forall w \in U_{a d} \\
\left(\gamma v+\int_{\mathbb{R}^{d}} \operatorname{div}\left((e+x) \rho_{2}\right) q_{2}, w-v\right)_{\mathbb{L}_{T}^{2}} \geq 0 & \forall w \in U_{a d}
\end{array}
$$

Apply the former inequality with $w=v$ and the latter with $w=u$, and subtract the two obtained relations: if we set $\delta \rho:=\rho_{1}-\rho_{2}$ and $\delta q:=q_{1}-q_{2}$, we get

$$
\left(\gamma(u-v)+\int_{\mathbb{R}^{d}} \operatorname{div}((e+x) \delta \rho) q_{1} d x+\int_{\mathbb{R}^{d}} \operatorname{div}\left((e+x) \rho_{2}\right) \delta q d x, u-v\right)_{\mathbb{L}_{T}^{2}} \leq 0 .
$$

From the previous inequality, straightforward computations allow to deduce that

$$
\begin{align*}
\gamma \int_{0}^{T}|u(t)-v(t)|^{2} d t & \leq-\int_{0}^{T}\left(\int_{\mathbb{R}^{d}} \operatorname{div}((e+x) \delta \rho) q_{1}+\int_{\mathbb{R}^{d}} \operatorname{div}\left((e+x) \rho_{2}\right) \delta q\right)(u(t)-v(t)) d t  \tag{5.25}\\
& \leq \int_{0}^{T}\left(\int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho) q_{1}\right|+\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}\right) \delta q\right|\right)|u(t)-v(t)| d t
\end{align*}
$$

At this point, one may want to apply a Cauchy-Schwarz inequality with respect to time, and a Young inequality to absorb the second term on the left-hand side. This is certainly possible; however, in order to optimize the value of the constants, it is better to directly estimate the integral terms, at any time $t \in[0, T]$. Such a bound can be performed exactly as in the proof of Theorem 5.3 above: we then find again inequalities (5.21) and (5.24). Inserting them into (5.25) yields

$$
\gamma \int_{0}^{T}|u(t)-v(t)|^{2} d t \leq K \int_{0}^{T}|u(t)-v(t)|\left(\int_{0}^{t}|u(s)-v(s)| d s\right) d t,
$$

where $K$ is defined as in the proof of Theorem 5.3 above.
For simplicity of notation, define $\sigma(t):=|u(t)-v(t)|$ for almost every $t \in[0, T]$. The previous estimate becomes then

$$
\gamma \int_{0}^{T}(\sigma(t))^{2} d t \leq K \int_{0}^{T} \sigma(t)\left(\int_{0}^{t} \sigma(s) d s\right) d t=\frac{K}{2}\left(\int_{0}^{T} \sigma(t) d t\right)^{2}
$$

Using a Cauchy-Schwarz inequality for the term on the right-hand side, we finally get

$$
\gamma \int_{0}^{T}(\sigma(t))^{2} d t \leq \frac{K T}{2} \int_{0}^{T}(\sigma(t))^{2} d t
$$

which obviously implies $\sigma \equiv 0$ almost everywhere on $[0, T]$ whenever $K T / \gamma<2$. Then, we conclude the proof remarking that $K \leq \widetilde{K}$.

### 5.4 The case of confining $\theta$ and $\varphi$

As pointed out in Remark 5.1, from the applications viewpoint, it may be desirable to consider the case when both $\theta$ and $\varphi$ are proportional to the function $|x|^{2}$. In this section, we discuss the necessary adaptations to be implemented in our arguments in order to address this case.

Therefore, from now on we assume that

$$
\theta(x)=|x|^{2} \quad \text { and } \quad \varphi(x)=|x|^{2},
$$

although the discussion can be further adapted, in order to treat more general polynomial growths. In order to simplify the presentation, we also assume that $\delta=\nu=0$.

First of all, we notice that, in view of (5.3), for $J$ to be well-defined it is necessary that $|x|^{2} \rho$ belongs to $L^{1}$. Then, we have to assume higher integrability on $\rho$, namely that

$$
\rho \in C\left([0, T] ; L_{k}^{2}\left(\mathbb{R}^{d}\right)\right), \quad \text { for some } \quad k>2+\frac{d}{2}
$$

This of course entails that, in (A.1), one has to take $\rho_{0} \in H_{k}^{m}$ and $g \in L_{T}^{1}\left(H_{k}^{m}\right)$, with the same restriction $k>2+d / 2$. However, the arguments that show existence of an optimal control do not change, so that Theorem 5.1 still holds true.

The main changes pertain Section 5.2, starting from the Definition 5.9 of the functional $\mathcal{L}$. First of all, let us focus on the Lagrangian multiplier $q$. On the one hand, we need it to be in some duality pairing with $\rho$ : then, keeping in mind Definition 3.1, we introduce, for $(m, k) \in \mathbb{N}^{2}$, the spaces

$$
H_{-k}^{m}\left(\mathbb{R}^{d}\right):=\left\{f \in H_{\mathrm{loc}}^{m}\left(\mathbb{R}^{d}\right)\left|\quad(1+|x|)^{-k} D^{\alpha} f \in L^{2}\left(\mathbb{R}^{d}\right) \quad \forall 0 \leq|\alpha| \leq m\right\}\right.
$$

This space is endowed with the natural norm

$$
\|f\|_{H_{-k}^{m}}=\sum_{0 \leq|\alpha| \leq m}\left\|(1+|x|)^{-k} D^{\alpha} f\right\|_{L^{2}}
$$

On the other hand, we still expect $q$ to solve (5.11) to an extent, although the meaning of that equation is now no more clear, owing to the fact that $\theta$ and $\varphi$ do not belong anymore to $L^{2}$. To deal with both issues, we need the following lemma, whose proof can be performed arguing as in the proof of Theorem 3.4 above, using this time the weight $(1+|x|)^{-k}$. We omit to give the details here.

Lemma 5.1. Let $T>0$ and $(m, k) \in \mathbb{N}^{2}$ fixed, and let a be a vector field satisfying hypotheses (3.3). Moreover, assume that $q_{0} \in H_{-k}^{m}\left(\mathbb{R}^{d}\right)$ and $g \in L^{1}\left([0, T] ; H_{-k}^{m}\left(\mathbb{R}^{d}\right)\right)$.

Then there exists a unique solution $q \in C\left([0, T] ; H_{-k}^{m}\left(\mathbb{R}^{d}\right)\right)$ to the problem

$$
\partial_{t} q+a \cdot \nabla q=g, \quad \text { with } \quad q_{\mid t=0}=q_{0}
$$

Moreover, there exists a "universal" constant $C>0$ such that the following estimate holds true for any $t \in[0, T]$ :

$$
\begin{equation*}
\|q(t)\|_{H_{-k}^{m}} \leq C \exp \left(C \int_{0}^{t}\|\nabla a(\tau)\|_{C_{b}^{m}} d \tau\right)\left(\left\|q_{0}\right\|_{H_{-k}^{m}}+\int_{0}^{t}\|g(\tau)\|_{H_{-k}^{m}} d \tau\right) \tag{5.26}
\end{equation*}
$$

Let us come back to our optimal control problem. In view of Lemma 5.1, we can solve equation (5.11) with $\theta$ and $\varphi$ equal to $|x|^{2}$, getting a unique solution in the space $C_{T}\left(L_{-k}^{2}\right)$ for any $k>2+d / 2$. Let us fix, once for all, the choice ${ }^{1}$

$$
k_{0}=3+\left[\frac{d}{2}\right]
$$

Then, it is easy to see that the functional $\mathcal{L}$ is well-defined on the space

$$
\widetilde{\mathbb{X}}_{T}:=\left(W_{T}^{1,1}\left(L_{k_{0}}^{2}\right) \cap L_{T}^{\infty}\left(H_{k_{0}+1}^{1}\right)\right) \times \mathbb{L}_{T}^{2} \times C_{T}\left(L_{-k_{0}}^{2}\right)
$$

Of course, we also need to take $\rho_{0}$ and $g$ as in assumption (A.1), with $m \geq 1$ and $k \geq k_{0}+1$.
Thereafter, we can write the optimality system (5.10)-(5.11)-(5.12), as done above. In order to characterize equation (5.12) in terms of the gradient of the reduced functional $\widehat{J}$, we need to further assume that $m \geq 2$ and $k \geq k_{0}+2$.

[^1]Finally, notice that also the analysis in Section 5.3 works similarly as above. Of course, we do not need anymore to impose assumption (A.6). On the other hand, assumption (A.4)* is now too strong, and we have to dismiss it.

However, we claim that it is still possible to get results analogous to Theorems 5.3 and 5.4. More precisely, we have the following statement for the unconstrained problem.

Proposition 5.2. Under assumptions (A.1)-(A.2)-(A.5), suppose also that both $m \geq 2$ and $k \geq k_{0}+2$. In addition, take $\delta=\nu=0$ in (5.3), and $\theta(x)=\varphi(x)=|x|^{2}$.

Then, there exists at most one optimal control $u^{*}$ in the class $\mathbb{L}_{T}^{2}$.
Proof. The proof is very similar to the one to Theorem 5.3, therefore we limit ourselves to put in evidence the main changes to be adopted, and to treat the most delicate points of the analysis.

As before, let ( $u_{1}, \rho_{1}, q_{1}$ ) and ( $u_{2}, \rho_{2}, q_{2}$ ) be two optimal controls with corresponding state and adjoint state. Arguing as above, we find that $\delta u=u_{1}-u_{2}$ fulfils estimate (5.19). Let us now focus on the estimate of each integral appearing in that relation.

As for the former integral term, also by use of Lemma 5.1, we can write

$$
\begin{aligned}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho(t)) q_{1}(t)\right| d x & \leq\left\|q_{1}(t)\right\|_{L_{-k_{0}}^{2}}\|\operatorname{div}((e+x) \delta \rho(t))\|_{L_{k_{0}}^{2}} \\
& \leq C_{3}\left(\|\delta \rho(t)\|_{L_{k_{0}}^{2}}+\|\nabla \delta \rho(t)\|_{L_{k_{0}+1}^{2}}\right) \leq C_{3}\|\delta \rho(t)\|_{H_{k_{0}+1}^{1}} .
\end{aligned}
$$

Notice that the constant $C_{3}$ can be expressed as follows

$$
\begin{equation*}
C_{3}:=C(1+T)\left\||x|^{2}(1+|x|)^{-k_{0}}\right\|_{L^{2}} \exp \left(C\left(\left\|\operatorname{div} a_{0}\right\|_{L_{T}^{1}\left(L^{\infty}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}\right)\right), \tag{5.27}
\end{equation*}
$$

for a "universal" constant $C>0$. At this point, the estimate for $\delta \rho$ works as before, finally leading to

$$
\begin{equation*}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}((e+x) \delta \rho(t)) q_{1}(t)\right| d x \leq \widetilde{C}_{3} \int_{0}^{t}|\delta u(\tau)| d \tau \tag{5.28}
\end{equation*}
$$

where we have defined $\widetilde{C}_{3}=C_{3} K_{1}^{\left(k_{0}+2\right)}$, just depending on the data of the problem.
Next, consider the second integral in (5.19). The computations are similar to the previous ones: first of all, we can estimate

$$
\begin{aligned}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}(t)\right) \delta q(t)\right| d x & \leq\|\delta q(t)\|_{L_{-k_{0}}^{2}}\left\|\operatorname{div}\left((e+x) \rho_{2}(t)\right)\right\|_{L_{k_{0}}^{2}} \\
& \leq\|\delta q(t)\|_{L_{-k_{0}}^{2}}\left\|\rho_{2}(t)\right\|_{H_{k_{0}+1}^{1}} \leq C_{4}\|\delta q(t)\|_{L_{-k_{0}}^{2}}
\end{aligned}
$$

where, by Theorem 3.4 applied to equation (5.10) for $\rho_{2}$, we obtain that

$$
\begin{equation*}
C_{4}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)\left(\left\|\rho_{0}\right\|_{H_{k_{0}+1}^{1}}+\|g\|_{L_{T}^{1}\left(H_{k_{0}+1}^{1}\right)}\right), \tag{5.29}
\end{equation*}
$$

for a "universal" constant $C>0$. On the other hand, Lemma 5.1 applied to the equation for $\delta q$ gives, for a constant $C>0$, the estimate

$$
\|\delta q(t)\|_{L_{-k_{0}}^{2}} \leq C \exp \left(C\left(\left\|\operatorname{div} a_{0}\right\|_{L_{T}^{1}\left(L^{\infty}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}\right)\right) \int_{0}^{t}|\delta u(\tau)|\left\|(1+|x|) \nabla q_{2}(\tau)\right\|_{L_{-k_{0}}^{2}} d \tau
$$

Notice that $\left\|(1+|x|) \nabla q_{2}(\tau)\right\|_{L_{-k_{0}}^{2}} \leq\left\|\nabla q_{2}(\tau)\right\|_{L_{-k_{0}+1}^{2}}$. In order to bound this quantity, we can differentiate the equation for $q_{2}$ with respect to $x^{j}$, for $1 \leq j \leq d$, and get (notice that $\partial_{j}|x|^{2}=2 x^{j}$ )

$$
\partial_{t}\left((1+|x|)^{-k_{0}+1} \partial_{j} q_{2}\right)+a\left(t, x ; u_{2}\right) \cdot \nabla\left((1+|x|)^{-k_{0}+1} \partial_{j} q_{2}\right)=
$$

$$
=2 x^{j}(1+|x|)^{-k_{0}+1}-(1+|x|)^{-k_{0}+1} \partial_{j} a\left(t, x ; u_{2}\right) \cdot \nabla q_{2},
$$

with initial datum equal to $2 x^{j}(1+|x|)^{-k_{0}+1}$. Obviously, the latter term in the right-hand side can be absorbed by a Grönwall argument; in addition, an easy computation shows that the former is in $L^{2}$. Therefore, by applying an $L^{2}$ estimate of Theorem 3.3 to the previous equation implies, for a "universal" constant $C>0$, the following bound

$$
\left\|\nabla q_{2}(\tau)\right\|_{L_{-k_{0}+1}^{2}} \leq C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(L^{\infty}\right)}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)(1+T)\left\||x|(1+|x|)^{-k_{0}+1}\right\|_{L^{2}} .
$$

By use of this latter estimate, we finally obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{d}}\left|\operatorname{div}\left((e+x) \rho_{2}(t)\right) \delta q(t)\right| d x \leq \widetilde{C}_{4} \int_{0}^{t}|\delta u(\tau)| d \tau \tag{5.30}
\end{equation*}
$$

where we have defined $\widetilde{C}_{4}:=C_{4} \widetilde{\mathcal{K}}_{1}^{(1)}$ and

$$
\begin{equation*}
\widetilde{\mathcal{K}}_{1}^{(1)}:=C \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{1}\right)}+\left\|u_{1}\right\|_{\mathbb{L}_{T}^{1}}+\left\|u_{2}\right\|_{\mathbb{L}_{T}^{1}}\right)\right)(1+T)\left\||x|(1+|x|)^{-k_{0}+1}\right\|_{L^{2}} \tag{5.31}
\end{equation*}
$$

At this point, we can now insert the estimates (5.28) and (5.30) into (5.19), and get, for a new constant $\mathcal{K}=\widetilde{C}_{3}+\widetilde{C}_{4}$, the relation

$$
\gamma|\delta u(t)| \leq \mathcal{K} \int_{0}^{t}|\delta u(\tau)| d \tau
$$

We then conclude by an application of Grönwall's lemma.
The constrained case can be dealt with similarly, by the use of the new constants $C_{3}, C_{4}, \mathcal{K}_{1}^{(1)}$ and $K_{1}^{\left(k_{0}+2\right)}$. We omit its precise statement and presentation. We just remark that, after defining

$$
\begin{aligned}
\widetilde{\mathcal{K}}:=C(1+ & T)\left\|(1+|x|)^{-k_{0}+2}\right\|_{L^{2}} \times \\
& \times \exp \left(C\left(\left\|\nabla a_{0}\right\|_{L_{T}^{1}\left(C_{b}^{2}\right)}+T \max \left\{\left|u^{a}\right|,\left|u^{b}\right|\right\}\right)\right)\left(\left\|\rho_{0}\right\|_{H_{k_{0}+2}^{2}}+\|g\|_{L_{T}^{1}\left(H_{k_{0}+2}^{2}\right)}\right),
\end{aligned}
$$

the new smallness condition reads $\widetilde{\mathcal{K}} T / \gamma<2$.

## A Appendix - Proof of some technical results

In this appendix, we collect the proof of some technical lemmas which we have used in the course of our investigation. We start with the proof to Lemma 3.2: we limit ourselves to give a sketch of its proof, and refer to [18] for details.

Proof of Lemma 3.2. First of all, let us focus just on the dependence with respect to the space variable, and forget about the time variable in the next computations.

We start by showing that the commutator $r_{n}^{j}:=\partial_{j}\left(\left[a, S_{n}\right] \cdot\right)$ is a bounded operator acting from $L^{2}$ into itself, uniformly in $n$. More precisely, there exists a "universal" constant $C$, dependending on $a$ but independent of $j$ and (more importantly) of $n$, such that

$$
\begin{equation*}
\left\|\partial_{j}\left(\left[a, S_{n}\right] \rho\right)\right\|_{L^{2}} \leq C\|\rho\|_{L^{2}} \quad \forall \rho \in L^{2} \tag{A.1}
\end{equation*}
$$

Indeed, straightforward computations show that

$$
r_{n}^{j}(\rho)=\partial_{j}\left(\left[a, S_{n}\right] \rho\right)=\partial_{j} a S_{n} \rho+\Sigma_{n}^{j}, \quad \text { with } \quad \Sigma_{n}^{j}:=a \partial_{j} S_{n} \rho-\partial_{j} S_{n}(a \rho) .
$$

Of course, $\left\|\partial_{j} a S_{n} \rho\right\|_{L^{2}} \leq C\|\nabla a\|_{L^{\infty}}\|\rho\|_{L^{2}}$, for a constant $C>0$ independent of $n \in \mathbb{N}$. Moreover, we notice that we can write

$$
\Sigma_{n}^{j}=n^{d+1} \int_{\mathbb{R}^{d}} \partial_{j} s(n(x-y)) \rho(y)(a(x)-a(y)) d y
$$

Hence, an application of the mean-value theorem to $a$ implies

$$
\left\|\Sigma_{n}^{j}\right\|_{L^{2}} \leq\|\nabla a\|_{L^{\infty}}\|\rho\|_{L^{2}}\left\|\partial_{j} s(z) \mid z\right\|\left\|_{L^{1}} \leq C\right\| \rho \|_{L^{2}},
$$

as claimed.
On the other hand, whenever $\rho$ is more regular, say $\rho \in H^{1}$, we can write

$$
r_{n}^{j}(\rho)=\partial_{j} a S_{n} \rho-S_{n}\left(\partial_{j} a \rho\right)+a S_{n} \partial_{j} \rho-S_{n}\left(a \partial_{j} \rho\right)=\partial_{j} a S_{n} \rho-S_{n}\left(\partial_{j} a \rho\right)+\left[a, S_{n}\right] \partial_{j} \rho .
$$

Notice now that $\partial_{j} a S_{n} \rho-S_{n}\left(\partial_{j} a \rho\right) \longrightarrow 0$ in $L^{2}$, for $n \rightarrow \infty$ : this holds true, since both terms strongly converge to $\partial_{j} a \rho$ in $L^{2}$. As for the last term on the right-hand side of the previous equality, we can argue similarly as above and write

$$
\left[a, S_{n}\right] \partial_{j} \rho=n^{d} \int_{\mathbb{R}^{d}} s(n(x-y)) \partial_{j} \rho(y)(a(x)-a(y)) d y,
$$

which implies, by use of the mean-value theorem again,

$$
\left\|\left[a, S_{n}\right] \partial_{j} \rho\right\|_{L^{2}} \leq \frac{1}{n}\|\nabla a\|_{L^{\infty}}\|\nabla \rho\|_{L^{2}}\|s(z)|z|\|_{L^{1}}
$$

Therefore, we have proven that, whenever $\rho \in H^{1}$, we have $r_{n}^{j}(\rho) \longrightarrow 0$ strongly in $L^{2}$, whenever $n \rightarrow+\infty$.

This latter property, combined with (A.1) and a standard approximation procedure, implies that, for any $\rho \in L^{2}$, we have $r_{n}^{j}(\rho) \longrightarrow 0$ for $n \rightarrow+\infty$, in the strong topology of $L^{2}$.

Let us now consider the general case when the functions $a$ and $\rho$ also depend on the time variable, verifying the properties assumed in Lemma 3.2. We have then to repeat the previous steps, keeping track of the time regularity.

We just put in evidence that, in the approximation procedure, given $\rho \in L_{T}^{\infty}\left(L^{2}\right)$, we can approximate it with a sequence $\left(\rho_{\delta}\right)_{\delta>0} \subset L_{T}^{\infty}\left(L^{2}\right)$, with $\rho_{\delta} \in L_{T}^{\infty}\left(H^{1}\right)$ for any $\delta>0$ fixed, and such that $\rho_{\delta} \longrightarrow \rho$ strongly in $L_{T}^{p}\left(L^{2}\right)$ for any $1 \leq p<+\infty$.

This entails that, before ahead, we need to regularize $a$ with respect to time, for instance by convolution. In particular, we take a sequence $\left(a_{\eta}\right)_{\eta>0} \subset L_{T}^{1}\left(C^{1}\right)$, with $\left(\nabla a_{\eta}\right)_{\eta>0} \subset L_{T}^{1}\left(L^{\infty}\right)$ such that $\nabla a_{\eta} \longrightarrow \nabla a$ in $L_{T}^{1}\left(L^{\infty}\right)$, and $a_{\eta}$ smooth (say continuous) with respect to time.

Then, after fixing some $1 \leq p<+\infty$ and denoting by $q$ its conjugate exponent (namely, $1 / p+1 / q=$ 1), we can estimate

$$
\begin{aligned}
&\left\|r_{n}^{j}(\rho)\right\|_{L_{T}^{1}\left(L^{2}\right)} \leq\left\|\partial_{j}\left(\left[a-a_{\eta}, S_{n}\right] \rho\right)\right\|_{L_{T}^{1}\left(L^{2}\right)}+\left\|\partial_{j}\left(\left[a_{\eta}, S_{n}\right]\left(\rho-\rho_{\delta}\right)\right)\right\|_{L_{T}^{1}\left(L^{2}\right)}+\left\|\partial_{j}\left(\left[a_{\eta}, S_{n}\right] \rho_{\delta}\right)\right\|_{L_{T}^{1} L^{2}} \\
& \leq C\left(\left\|\nabla\left(a-a_{\eta}\right)\right\|_{L_{T}^{1}\left(L^{\infty}\right)}\|\rho\|_{L_{T}^{\infty}\left(L^{2}\right)}+\right. \\
&\left.\quad+\left\|\nabla a_{\eta}\right\|_{L_{T}^{q}\left(L^{2}\right)}\left\|\rho-\rho_{\delta}\right\|_{L_{T}^{p}\left(L^{2}\right)}+\left\|\partial_{j}\left(\left[a_{\eta}, S_{n}\right]\right) \rho_{\delta}\right\|_{L_{T}^{1}\left(L^{2}\right)}\right) .
\end{aligned}
$$

At this point, for any given $\varepsilon>0$, we can fix first of all $\eta>0$ so that the first term in the right-hand side is smaller than $\varepsilon$; then, in correspondence of that $\eta$, we fix $\delta>0$ so small that also the second term is bounded by $\varepsilon$; finally, in the last term we can pass to the limit with respect to $n$, making it smaller than $\varepsilon$ as well.

Thus the lemma is proved.

Let us now give the details of the proof to Lemma 3.3.
Proof of Lemma 3.3. Of course, by definition of the space $H_{k}^{m}$, we have that $f \in H^{m}$. We only need to prove that $|x|^{k} f$ belongs to $H^{m}$ as well.

First of all, by an easy induction, it follows that $|x|^{k} D^{\alpha} f \in L^{2}$ for all $0 \leq|\alpha| \leq m-1$. In particular $f \in L_{k}^{2}$, i.e. $|x|^{k} f \in L^{2}$. Moreover, since $f \in H^{m}$, we gather also that $|x|^{j} D^{\alpha} f \in L^{2}$ for all $0 \leq j \leq k$ and $0 \leq|\alpha| \leq m-1$.

Now, fixed some $1 \leq \ell \leq m$, let us take a multi-index $\alpha$ such that $|\alpha|=\ell$ : we want to prove that $D^{\alpha}\left(|x|^{k} f\right) \in L^{2}$. For this, we write, by Leibniz rule,

$$
D^{\alpha}\left(|x|^{k} f\right)=|x|^{k} D^{\alpha} f+\sum_{\beta} D^{\beta}|x|^{k} D^{\alpha-\beta} f,
$$

where the sum is performed for all $\beta \leq \alpha$ such that $|\beta| \geq 1$. The first term in the right-hand side belongs to $L^{2}$ by the argument here above. Moreover, one has

$$
\left.\left.\left|D^{\beta}\right| x\right|^{k}|\leq C| x\right|^{k-|\beta|} .
$$

At this point, we observe that $k-|\beta|<k$, since $|\beta| \geq 1$; in addition, $k-|\beta| \geq 0$, owing to the fact that $|\beta| \leq|\alpha|=\ell \leq m \leq k$. Therefore, by the properties previously established, we gather that each term $D^{\beta}|x|^{k} D^{\alpha-\beta} f$ of the sum also belongs to $L^{2}$. In turn, this implies that $D^{\alpha}\left(|x|^{k} f\right) \in L^{2}$, for all $0 \leq|\alpha| \leq m$, that is to say $|x|^{k} f \in H^{m}$.

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[^1]:    ${ }^{1}$ Given $z \in \mathbb{R}$, we denote by $[z]$ its entire part.

