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Logarithmic Decay of a Wave Equation with Kelvin-Voigt Damping $*$

Luc Robbiano †and Qiong Zhang ‡§

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Abstract

In this paper we analyze the long time behavior of a wave equation with local Kelvin-Voigt Damping. Through introducing proper class symbol and pseudo-differential calculus, we obtain a Carleman estimate, and then establish an estimate on the corresponding resolvent operator. As a result, we show the logarithmic decay rate for energy of the system without any geometric assumption on the subdomain on which the damping is effective.

keywords. Carleman estimate; wave equation; Kelvin-Voigt damping; logarithmic stability

AMS 2010 subject classification: Primary 93B05; Secondary 93B07, 35B37

1 Introduction

In this paper, we consider a wave equation with local Kelvin-Voigt damping and analyze long time behaviour for the -solution of the system. Let $\Omega \subset \mathbb{R}^d$ be a bounded domain with smooth boundary $\Gamma = \partial \Omega$. Denote by ∂_n the unit outward normal vector on boundary Γ. The PDE model is as follows.

$$
\begin{cases}\ny_{tt}(t,x) - \operatorname{div} [\nabla y(t,x) + a(x)\nabla y_t(t,x)] = 0 & \text{in} \quad (0,\infty) \times \Omega, \\
y(t,x) = 0 & \text{on} \quad (0,\infty) \times \Gamma, \\
y(0,x) = y^0, \quad y_t(0,x) = y^1 & \text{in} \quad \Omega,\n\end{cases}
$$
\n(1.1)

where the coefficient function $a(\cdot) \in L^1(\Omega)$ is nonnegative and not identically null.

The natural energy of system (1.1) is

$$
E(t) = \frac{1}{2} \Big[\int_{\Omega} (|\nabla y(t)|^2 + |y_t(t)|^2) dx \Big].
$$
 (1.2)

A direct computation gives that

$$
\frac{d}{dt}E(t) = -\int_{\text{supp }a} a(x) |\nabla y_t(t)|^2 dx.
$$
\n(1.3)

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[†]Laboratoire de Mathématiques Appliquées, UMR 8100 du CNRS, Université Paris–Saclay (site UVSQ), 45 avenue des Etats Unis, 78035 Versailles Cedex, France (luc.robbiano@uvsq.fr)

[‡]School of Mathematics and Statistics, Beijing Key Laboratory on MCAACI, Beijing Institute of Technology, Beijing, 100081, China (zhangqiong@bit.edu.cn).

[§]Corresponding author.

Formula (1.3) shows that the only dissipative mechanism acting on the system is the viscoelastic damping div $[a\nabla y_t]$, which is only effective on supp a.

To rewrite the system as an evolution equation, we set the energy space as

$$
\mathcal{H} = H_0^1(\Omega) \times L^2(\Omega),\tag{1.4}
$$

with norm

$$
||Y||_{\mathcal{H}} = \sqrt{||\nabla y_1||_{L^2(\Omega)}^2 + ||y_2||_{L^2(\Omega)}^2}, \quad \forall Y = (y_1, y_2) \in \mathcal{H}.
$$
 (1.5)

Define an unbounded operator $\mathcal{A}: D(\mathcal{A}) \subset \mathcal{H} \to \mathcal{H}$ by

$$
\mathcal{A}Y = (y_2, \text{div}(\nabla y_1 + a\nabla y_2)), \quad \forall Y = (y_1, y_2) \in D(\mathcal{A}),
$$

and

$$
\mathcal{D}(\mathcal{A}) = \left\{ (y_1, y_2) \in \mathcal{H} : y_2 \in H_0^1(\Omega), \text{div} (\nabla y_1 + a\nabla y_2) \right\} \in L^2(\Omega) \right\}.
$$

Let $Y(t) = (y(t), y_t(t))$. Then system (1.1) can be written as

$$
\frac{d}{dt}Y(t) = \mathcal{A}Y(t), \quad \forall \ t > 0, \quad Y(0) = (y^0, y^1). \tag{1.6}
$$

It is known from [20] that if supp a is non-empty, the operator A generates a contractive C_0 semigroup e^{tA} on H and $i\mathbb{R} \subset \rho(A)$, the resolvent of A. Consequently, the semigroup e^{tA} is strongly stable. Moreover, if the entire medium is of the viscoelastic type (i.e. supp $a = \overline{\Omega}$), the damping for the wave equation not only induces exponential energy decay, but also restricts the spectrum of the associated semigroup generator to a sector in the left half plane, and the associated semigroup is analytic ([13]). When the Kelvin-Voigt damping is localized on a subdomain of Ω , the properties of system is quite complicated. First, it has been proved that properties of regularity and stability of 1-d system (1.1) depend on the continuousness of coefficient function $a(\cdot)$. More precisely, assume that $\Omega = (-1, 1)$ and $a(x)$ behaviours like x^{α} with $\alpha > 0$ in supp $a = [0, 1]$. Then the solution of system (1.1) is eventually differentiable for $\alpha > 1$, exponentially stable for $\alpha \geq 1$, polynomially stable of order $\frac{1}{1-\alpha}$ for $0 < \alpha < 1$, and polynomially stable of optimal decay rate 2 for $\alpha = 0$ (see [19, 22, 25, 27]). For the higher dimensional system, the corresponding semigroup is exponentially stable when $a(\cdot) \in \mathscr{C}^2(\Omega)$ and supp $a \supset \Gamma$ ([21]). However, when the Kelvin-Voigt damping is local and the material coefficient $a(\cdot)$ is a positive constant on supp a, the energy of system (1.1) does not decay exponentially for any geometry of Ω and supp a ([9, 28]). The reason is that the strong damping and non-continuousness of the coefficient function lead to reflection of waves at the interface $\gamma = \partial(\text{supp }a) \setminus \Gamma$, which then fails to be effectively damped because they do not enter the region of damping. It turns out that the Kelvin-Voigt damping does not follow the principle that "geometric optics" condition implies exponential stability, which is true for the wave equation with local viscous damping ([2]).

Recently, [29] proves the polynomial stability of system (1.1) when $a(\cdot) \equiv a_0 > 0$ on supp a and supp α satisfies certain geometry conditions. Then, a natural problem is: how about the decay rate if supp $a \neq \emptyset$ is arbitrary? In [1], a is assumed discontinuous along a $(d-1)$ –manifold, supp a is arbitrary and the rate of the decay of semi-group is estimated by $(\log t)^{-k}$ for a data in $D(A^k)$. In this paper, we analyze the logarithmic decay properties of the solution to (1.1) when a is smooth and supp a is arbitrary. The main result reads as follows.

Theorem 1.1. Suppose that the coefficient function $a(\cdot) \in \mathscr{C}_0^{\infty}(\Omega)$ is nonnegative and supp $a \subset \Omega$ is non-empty. Then the energy of the solution of (1.1) decays at logarithmic speed. More precisely, one has that there exists a positive constant C such that

$$
||e^{t\mathcal{A}}Y_0||_{\mathcal{H}} \le \frac{C}{[\log(t+2)]^{\frac{4k}{5}}} ||Y_0||_{D(\mathcal{A}^k)}, \quad \forall \ t > 0, \ \ Y_0 = (y^0, y^1) \in D(\mathcal{A}^k). \tag{1.7}
$$

Our approach is based on the results duo to [7], which reduced the problem of determining the rate of energy decay to estimating the norm of the resolvent operator along the imaginary axis, see also [10, 18], etc. Our argument divides naturally three steps. First, in Section 2, we show some preliminaries including definitions and classical results about symbol, pseudo-differential calculus, and commutator estimate, etc. Then in Section 3, we prove corresponding Carleman estimates. Finally, in Section 4, we present a resolvent estimate and obtain Theorem 1.1. This theorem is a consequence of a resolvent estimate. The proof is given in Section 4. The method was developed in $(4, 5, 6, 16, 18, 26]$ and the references cited therein).

Throughout this paper, we use $\|\cdot\|_V$ and $(\cdot | \cdot)_V$ to denote the norm and inner product on $L^2(V)$, where $V \subset \mathbb{R}^d$ if there is no further comments. When writing $f \leq g$ (or $f \geq g$), we mean that there exists a positive constant C such that $f \leq Cg$ (or $f \geq Cg$). For $j = 1, 2, \cdots$, define operators $D_j = -i\partial_{x_j}, D = (D_1, \cdots, D_d), D^2 = \sum_{i=1}^d$ $j=1$ D_j^2 and $Da(x)D = \sum_{n=1}^d$ $\sum_{j=1} D_j a(x) D_j.$

2 Preliminaries

We shall prove Theorem 1.1 by Weyl-Hörmander calculus, which was introduced Hörmander ([12, 15]). In this section, some definitions and results on the class of symbol and pseudo-differential calculus are given.

2.1 Symbol and Symbolic calculus

For any $(x,\xi) \in \mathbb{R}^d \times \mathbb{R}^d$, $\lambda \in \mathbb{R}$ and $\tau > 0$, we introduce the metric

$$
g = g_{x,\xi} = \lambda dx^2 + \mu^{-2} d\xi^2, \quad \text{where } \mu^2 = \mu(\tau,\xi)^2 = \tau^2 + |\xi|^2,
$$
 (2.1)

and the weight

$$
\nu = \nu(x, \lambda) = \sqrt{1 + \lambda^2 a(x)^2}.
$$
\n(2.2)

Note that and $g_{x,\xi}(X,\Xi) = \lambda |X|^2 + \mu^{-2}(\tau,\xi) |\Xi|^2$ for all $X,\Xi \in \mathbb{R}^d$. Then we have the following results.

Lemma 2.1. Assume that there exist positive constants C and λ_0 such that $\lambda \geq \lambda_0$ and $\tau \geq$ $\max\{\mathcal{C}\lambda, 1\}$. It holds

- (i) The metric $g = g_{x,\xi}$ defined by (2.1) is admissible, i.e., it is slowly varying and temperate.
- (ii) The weight $\nu = \nu(x, \lambda)$ defined by (2.2) is admissible, i.e., it is g-continuous and g-temperate.

Proof. (i) From Definition 18.4.1 in [12], the metric $g_{x,\xi}$ defined by (2.1) is slowly varying if there exist $\delta > 0$ and $C > 0$ such that

$$
g_{x,\xi}(y-x,\eta-\xi) \le \delta
$$
 implies $g_{y,\eta}(X,\Xi) \le Cg_{x,\xi}(X,\Xi)$, $\forall x, y, \xi, \eta, X, \Xi \in \mathbb{R}^d$,

where the constants δ and C are independent on the parameters λ and τ .

Suppose $0 < \delta \leq 1/4$ and

$$
g_{x,\xi}(y-x,\eta-\xi) = \lambda |y-x|^2 + (\tau^2 + |\xi|^2)^{-1} |\eta-\xi|^2 \le \delta.
$$

Then, we have

$$
\tau^{2} + |\xi|^{2} \leq \tau^{2} + 2|\xi - \eta|^{2} + 2|\eta|^{2}
$$

$$
\leq \tau^{2} + 2\delta(\tau^{2} + |\xi|^{2}) + 2|\eta|^{2}.
$$

This implies that $\tau^2 + |\xi|^2 \leq 4(\tau^2 + |\eta|^2)$. Consequently,

$$
g_{y,\eta}(X,\Xi) = \lambda |X|^2 + (\tau^2 + |\eta|^2)^{-1} |\Xi|^2 \le \lambda |X|^2 + \frac{1}{4} (\tau^2 + |\xi|^2)^{-1} |\Xi|^2 \le g_{x,\xi}(X,\Xi).
$$

Therefore, g is slowly varying.

For a given metric $g_{x,\xi}$, the associated metric $g_{x,\xi}^{\sigma}$ is defined by $g_{x,\xi}^{\sigma} = (\tau^2 + |\xi|^2)dx^2 + \lambda^{-1}d\xi^2$. The metric $g_{x,\xi}$ is **temperate** if there exist $C > 0$ and $N > 0$, such that

$$
g_{x,\xi}(X,\Xi) \leq C g_{y,\eta}(X,\Xi) \left(1 + g_{x,\xi}^{\sigma}(x-y,\xi-\eta)\right)^N, \ \ \forall \ x,y,\xi,\eta,X,\Xi \in \mathbb{R}^d,
$$
\n
$$
(2.3)
$$

where the constants C and N are independent on the parameters λ and τ (Definition 18.5.1 in [12]).

For the metric $g = g_{x,\xi}$ defined by (2.1), (2.3) is equivalent to

$$
\lambda |X|^2 + (\tau^2 + |\xi|^2)^{-1} |\Xi|^2
$$

\n
$$
\leq C(\lambda |X|^2 + (\tau^2 + |\eta|^2)^{-1} |\Xi|^2) (1 + (\tau^2 + |\xi|^2) |x - y|^2 + \lambda^{-1} |\xi - \eta|^2)^N.
$$
\n(2.4)

First, assume that $\tau^2 + |\eta|^2 \leq 4(\tau^2 + |\xi|^2)$. It follows that

$$
(\tau^2 + |\eta|^2) \le C(\tau^2 + |\xi|^2)(1 + \lambda^{-1}|\xi - \eta|^2)^N, \quad C > 0, N > 0.
$$
 (2.5)

Then it is easy to obtain (2.4) from (2.5) .

Secondly, consider the case $\tau^2 + |\eta|^2 > 4(\tau^2 + |\xi|^2)$. Then

$$
|\eta| > 2 |\xi|, \quad |\eta| > \sqrt{3}\tau,
$$
\n(2.6)

and

$$
|\xi - \eta| > \frac{1}{2} |\eta| > \frac{\sqrt{3}}{2}\tau > \frac{\sqrt{3}}{2}C\lambda.
$$
 (2.7)

It follows from (2.6) and (2.7) that

$$
\lambda^{-1} |\xi - \eta|^2 > \frac{\sqrt{3}}{2} C |\xi - \eta| > \frac{\sqrt{3}}{4} C |\eta|.
$$

Consequently,

$$
(1 + \lambda^{-1} |\xi - \eta|^2)^2 > \frac{3}{16} \mathcal{C}^2 |\eta|^2 > \frac{3}{32} \mathcal{C}^2 (|\eta|^2 + 3\tau^2).
$$

This together with $\tau^2 + |\xi|^2 \ge 1$ yields that there exists a positive constant C such that (2.5) holds with $N = 2$.

(ii) It is known from Definition 18.4.2 in [12] that a weight $\nu(x)$ is g-continuous if there exist $\delta > 0$ and $C > 0$ such that

$$
g_{x,\xi}(y-x, \eta-\xi) \le \delta
$$
 implies $C^{-1}\nu(x) \le \nu(y) \le C\nu(x), \forall x, y, \xi, \eta \in \mathbb{R}^d$.

where the constants δ and C are independent on the parameters λ and τ . Since the weight $\nu(x)$ defined by (2.2) does not depend on ξ , the above condition is reduced to

$$
\lambda |x - y|^2 \le \delta \text{ implies } C^{-1}\nu(x) \le \nu(y) \le C\nu(x), \ \forall x, y \in \mathbb{R}^d.
$$

The weight $\nu(x)$ is g-temperate if there exist $C > 0$ and $N > 0$ such that

$$
\nu(y) \le C\nu(x)\left(1 + g_{y,\eta}^{\sigma}(x - y, \xi - \eta)\right)^N, \ \ \forall \ x, y, \xi, \eta \in \mathbb{R}^d,
$$
\n
$$
(2.8)
$$

where the constants C and N do not depend on the parameters λ and τ (Definition 18.5.1 in [12]). The weight $\nu(x)$ is admissible if it is q-continuous and q-temperate. When a weight is admissible, all the powers of this weight are g -continuous and g -temperate. Therefore, it suffice to prove that $1 + \lambda a(x)$ is admissible.

Let $s \in [0, t]$ and $t \in [0, 1]$. Define $f(s) = \lambda a(x + s(y - x))$ and $F(t) = \sup_{s \in [0, t]} f(s)$ where x, $y \in \Omega$ satisfying $\lambda |x - y|^2 \le \delta$. It is clear that $f'(s) = \lambda a'(x + s(y - x))(y - x)$. Combining this with the following inequality

$$
|a'(x)|^2 \le 2a(x)\|a''\|_{L^{\infty}(\Omega)}, \quad \forall x \in \Omega,
$$
\n(2.9)

we obtain

$$
|f'(s)| \leq \lambda |a'(x+s(y-x))| |y-x| \leq 2\lambda \|a''\|_{\infty}^{\frac{1}{2}} [a(x+s(y-x))]^{\frac{1}{2}} |y-x|.
$$

The proof of (2.9) will be given later. Consequently,

$$
\sup_{s\in[0,t]}|f'(s)| \leq 2\lambda^{\frac{1}{2}} \|a''\|_{L^{\infty}(\Omega)}^{\frac{1}{2}} F(t)^{\frac{1}{2}} \|y-x\|.
$$

Since $f(t) \le f(0) + t \sup_{s \in [0,t]} |f'(s)|$, F is non-decreasing and $\lambda |x - y|^2 \le \delta$, we obtain that for all $t \in [0, 1]$,

$$
f(t) \le f(0) + C\lambda^{\frac{1}{2}} F(t)^{\frac{1}{2}} |y - x| \le f(0) + C\sqrt{\delta} F(t)^{\frac{1}{2}} \le f(0) + C\sqrt{\delta} F(\alpha)^{\frac{1}{2}},
$$

where $C = 2||a''||_{L^{\infty}(\Omega)}^{\frac{1}{2}}$ and $\alpha \in [t, 1]$. Note that $f(0) = F(0)$. It follows that

$$
F(\alpha) = \sup_{t \in [0,\alpha]} f(t) \le F(0) + C\sqrt{\delta}F(\alpha)^{\frac{1}{2}}.
$$

This yields

$$
1 + F(\alpha) \le 1 + F(0) + C\sqrt{\delta} \left(1 + F(\alpha)\right)^{\frac{1}{2}} \le 1 + F(0) + C\sqrt{\delta} \left(1 + F(\alpha)\right). \tag{2.10}
$$

By choosing δ sufficiently small such that $C\sqrt{\delta} \leq 1/2$, one can deduce from (2.10) that

$$
1 + F(\alpha) \le 2(1 + F(0)), \quad \forall \alpha \in [t, 1].
$$

In particular, we have

$$
1 + \lambda a(y) \le 2(1 + \lambda a(x)).
$$

The above inequality remains true if we exchange x and y. Therefore, the weight $1 + \lambda a(x)$ is g-continuous.

On the other hand, note that $1 + \lambda a(x)$ is independent to ξ . Then, to obtain the weight $1 + \lambda a(x)$ is σ -temperate, it is sufficient to prove that

$$
1 + \lambda a(y) \le C(1 + \lambda a(x))(1 + \tau^2 |x - y|^2)^N.
$$
\n(2.11)

In fact, it is clear that $1 + \lambda a(y) \leq 1 + \lambda (a(x) + C |x - y|)$ where $C = ||a'||_{L^{\infty}(\Omega)}$. Therefore, there exists positive constant $C' = C C^{-1}$ such that

$$
1 + \lambda a(y) \le (1 + \lambda a(x))(1 + C' \tau |x - y|) \le (1 + \lambda a(x))(2 + 2(C' \tau |x - y|)^2)^{\frac{1}{2}}.
$$

Thus, we obtain (2.11) with $N = \frac{1}{2}$, $C = 2 \max\{1, C'\}.$

□

Remark 2.1. We claim that (2.9) holds for any compactly supported and nonnegative function $a \in \mathscr{C}^2(\Omega)$. In fact, from the following identity

$$
a(x+h) = a(x) + a'(x)h + \int_0^1 (1-t)a''(x+th)h^2dt, \ \ \forall \ h \in \mathbb{R},
$$

one can get

$$
a(x) + a'(x)h + \frac{1}{2}||a''||_{L^{\infty}(\Omega)}||h||^2 \ge 0.
$$

Let $h = y a'(x)$, where $y \in \mathbb{R}$ and $x \in \Omega$ are arbitrary. It follows from the above inequality that

$$
a(x) + |a'(x)|^2 y + \frac{1}{2} ||a''||_{L^{\infty}(\Omega)} |a'(x)|^2 y^2 \ge 0, \quad \forall x \in \Omega, y \in \mathbb{R}.
$$

Then,

$$
|a'(x)|^4 - 2a(x)||a''||_{L^{\infty}(\Omega)} |a'(x)|^2 \le 0,
$$

and (2.9) is proved.

Definition 2.1. (Section 18.4.2 in [12]) Assume the weight $m(x, \xi)$ is admissible and the metric g is defined by (2.1). Let $q(x,\xi,\lambda,\tau)$ be a \mathscr{C}^{∞} function with respect to (x,ξ) and λ , τ be parameters satisfying conditions in Lemma 2.1. The symbol $q(x,\xi,\lambda,\tau)$ is in **class** $S(m,g)$ if for all $\alpha,\beta \in \mathbb{N}^d$ there exists $C_{\alpha,\beta}$ independent of τ and λ such that

$$
|\partial_x^{\alpha} \partial_{\xi}^{\beta} q(x,\xi,\lambda,\tau)| \leq C_{\alpha,\beta} m(x,\xi) \lambda^{|\alpha|/2} \mu(\tau,\xi)^{-|\beta|}.
$$

Remark 2.2. (i) It is clear that $\mu = \sqrt{\tau^2 + |\xi|^2} \in S(\mu, g)$ since $|\partial_{\xi}^{\beta} \mu(\tau, \xi)| \lesssim \mu^{1-|\beta|}$ for all $\beta \in \mathbb{R}^d$.

(ii) Let v be the weight defined by (2.2). It is easy to get that $\lambda a \in S(\nu, g)$. In fact, if $|\alpha| \geq 2$, it holds that $|\partial_x^{\alpha}(\lambda a(x))| \leq C_{\alpha} \lambda \leq C_{\alpha} \lambda^{|\alpha|/2} \nu(x)$, where $C_{\alpha} > 0$. For the case $|\alpha| = 1$, it follows from (2.9) that

$$
|\partial_x^{\alpha}(\lambda a(x))| \leq \sqrt{2} ||a''||_{L^{\infty}(\Omega)}^{\frac{1}{2}} \lambda^{\frac{1}{2}} (\lambda a(x))^{\frac{1}{2}}
$$

Note that $|\lambda a(x)| < C \nu^2(x)$ for some $C > 0$. This together with the above inequality, we have that

$$
|\partial_x^{\alpha}(\lambda a(x))| < \sqrt{2C} ||a''||_{L^{\infty}(\Omega)}^{\frac{1}{2}} \lambda^{\frac{1}{2}} \nu(x).
$$

(iii) It is known from Lemma 18.4.3 of [12] that if the metric q and weights m_1 , m_2 are admissible, symbols $a \in S(m_1, g)$ and $b \in S(m_2, g)$, then $ab \in S(m_1 m_2, g)$. In particular, $(\lambda a)^j \mu^k \in$ $S(\nu^j \mu^k, g)$ for all j, $k \in \mathbb{N} \cup \{0\}.$

Definition 2.2. Let $b \in S(m, g)$ be a symbol and $u \in \mathscr{S}(\mathbb{R}^d)$, we set

$$
b(x, D, \tau)u(x) = \text{Op}(b)u(x) := (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ix\cdot\xi} b(x, \xi, \tau) \,\widehat{u}(\xi) \,d\xi.
$$

It is known that $Op(b)$: $\mathscr{S}(\mathbb{R}^d) \to \mathscr{S}(\mathbb{R}^d)$ is continuous and $Op(b)$ can be uniquely extended to $\mathscr{S}'(\mathbb{R}^d)$ continuously. The following two lemmas are consequences of Theorem 18.5.4 and 18.5.10 in [12].

Lemma 2.2. Let $b \in S(m, q)$ where m is an admissible weight and g is defined by (2.1). Then there exists $c \in S(m, g)$ such that $Op(b)^* = Op(c)$ and $c(x, \xi) = \overline{b(x, \xi)} + r(x, \xi)$ where the remainder $r \in S(\lambda^{\frac{1}{2}}\mu^{-1}m, g).$

Lemma 2.3. Let $b \in S(m_1, g)$ and $c \in S(m_2, g)$ where m_j are admissible weights for $j = 1, 2$ and g is defined by (2.1). Denote by $[Op(b), Op(c)] = Op(b) \circ Op(c) - Op(c) \circ Op(b)$ and Poisson bracket $\{b, c\}(x, \xi, \tau) = \sum_{1 \leq j \leq d} (\partial_{\xi_j} b \, \partial_{x_j} c - \partial_{x_j} b \, \partial_{\xi_j} c)(x, \xi, \tau).$ Then,

- (i) there exists $d \in S(m_1m_2, g)$ such that $Op(b) Op(c) = Op(d)$ and $d(x, \xi) = b(x, \xi)c(x, \xi) + r(x, \xi)$ where $r \in S(\lambda^{\frac{1}{2}}\mu^{-1}m_1m_2, g)$.
- (ii) for commutator $i[Op(b), Op(c)] = Op(f)$, it holds that $f \in S(\lambda^{\frac{1}{2}}\mu^{-1}m_1m_2, g)$ and $f(x,\xi) =$ ${b, c}(x, \xi) + r(x, \xi)$ where $r \in S(\lambda \mu^{-2}m_1m_2, q)$.

The operators in $S(\nu^j \mu^k, g)$ act on Sobolev spaces adapted to the class of symbol. Let $b \in$ $S(\nu^j \mu^k, g)$, where μ and g are defined by (2.1). Then there exists $C > 0$ such that

 $\|\operatorname{Op}(b)u\|_{\mathbb{R}^d} \leq C\|\nu^j \operatorname{Op}(\mu^k)u\|_{\mathbb{R}^d}, \quad \forall u \in \mathscr{S}(\mathbb{R}^d).$

By symbolic calculus, the above estimate is equivalent to $Op(\mu^{-k}\nu^{-j})Op(b)$ acts on $L^2(\mathbb{R}^d)$ since the operators associated with symbol in $S(1, g)$ act on $L^2(\mathbb{R}^d)$. In particular, if $b \in S(\nu^j \mu, g)$, then for any $\lambda \geq \lambda_0$, $\tau \geq \max\{ \mathcal{C}\lambda, 1 \}$ and $u \in \mathscr{S}(\mathbb{R}^d)$, it holds

$$
\|\operatorname{Op}(b)u\|_{\mathbb{R}^d} \leq C\tau \| \nu^j u\|_{\mathbb{R}^d} + C \| \nu^j Du\|_{\mathbb{R}^d},
$$

where $C > 0$ depends on positive constants λ_0 and \mathcal{C} .

2.2 Commutator estimate

In this subsection, we suppose that $\lambda = 1$ since the symbol does not depend on λ . The metric in (2.1) becomes

$$
\tilde{g} = dx^2 + \mu^{-2} d\xi^2, \quad \text{where } \mu \text{ is defined by (2.1).} \tag{2.12}
$$

To get the commutator estimate, we shall use the following Gårding inequality ([12, Theorem 18.6.7]).

Lemma 2.4. Let $b \in S(\mu^{2k}, \tilde{g})$ be real valued. μ and \tilde{g} are defined by (2.12). We assume there exists $C > 0$ such that $b(x, \xi, \tau) \ge C\mu^{2k}$. Then there exist $\tilde{C} > 0$ and $\tau_0 > 0$ such that

$$
\operatorname{Re}(\operatorname{Op}(b)w \mid w)_{\mathbb{R}^d} \ge \widetilde{C} \|\operatorname{Op}(\mu^k)w\|_{\mathbb{R}^d}^2, \qquad \forall \ w \in \mathscr{S}(\mathbb{R}^d) \text{ and } \tau \ge \tau_0. \tag{2.13}
$$

Definition 2.3. Let V be a bounded open set in \mathbb{R}^d . We say that the weight function $\varphi \in \mathscr{C}^{\infty}(\mathbb{R}^d;\mathbb{R})$ satisfies the **sub-ellipticity** condition in \overline{V} if $|\nabla \varphi| > 0$ in \overline{V} and there exists constant $C > 0$,

$$
\mathbb{p}_{\varphi}(x,\xi,\tau) = 0, \quad \forall (x,\xi) \in \overline{V} \times \mathbb{R}^d, \quad \tau > 0 \quad \Rightarrow \quad \{\mathbb{q}_2, \mathbb{q}_1\}(x,\xi,\tau) \ge C(|\xi|^2 + \tau^2)^{3/2}, \tag{2.14}
$$

where $p_{\varphi}(x,\xi,\tau) = |\xi + i\tau \nabla \varphi(x)|^2 = q_2(x,\xi,\tau) + i q_1(x,\xi,\tau)$ and q_1 , q_2 are real valued.

Lemma 2.5. ([12]) Let V be a bounded open set in \mathbb{R}^d and $\psi \in \mathscr{C}^{\infty}(\mathbb{R}^d;\mathbb{R})$ be such that $|\nabla \psi| > 0$ in \overline{V} . Then, for $\gamma > 0$ sufficiently large, $\varphi = e^{\gamma \psi}$ fulfills the sub-ellipticity property in V.

Lemma 2.6. Assume that φ satisfies the sub-ellipticity in Definition 2.3. For all $w \in \mathscr{C}_0^{\infty}(V)$, there exist $C_1, C_2 > 0$ and $\tau_0 > 0$ such that the following inequality holds for all $\tau \geq \tau_0$,

$$
C_1 \tau^3 \|w\|_V^2 + C_1 \tau \|Dw\|_V^2 \leq \text{Re} \left(\text{Op}(\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\}) w \mid w \right)_V
$$

+
$$
C_2 \tau^{-1} \| \text{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) w \|_V^2
$$

+
$$
C_2 \tau^{-1} \| \text{Op}(2\tau \xi \cdot \nabla \varphi(x)) w \|_V^2.
$$
 (2.15)

Proof. First, by homogeneity in (ξ, τ) , compactness arguments and sub-ellipticity condition, we claim that there exist constants $C, \delta > 0$ such that

$$
C[|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2] + \{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\xi \cdot \nabla \varphi(x)\} \ge \delta \mu^2. \tag{2.16}
$$

The proof of (2.16) is classical. In fact, set

$$
\mathcal{K} = \{ (x,\xi,\tau) \in \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R} : x \in \overline{V}, \ |\xi|^2 + \tau^2 = 1, \ \tau \ge 0 \},
$$

and for $(x, \xi, \tau) \in \mathcal{K}, \ \kappa > 0$,

$$
G(x,\xi,\tau,\kappa) = \kappa [|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2] + {|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\xi \cdot \nabla \varphi(x)}.
$$

If $|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2 = 0$ for $(x, \xi, \tau) \in \mathcal{K}$, it is clear that there exists a positive constant δ such that (2.16) holds due to the fact that ϕ is sub-elliptic. When $|2\xi \cdot \nabla \varphi(x)|^2 +$ $\mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2 > 0$, there exists a positive constant $\kappa_{x,\xi,\tau}$ such that $G(x,\xi,\tau,\kappa) > 0$ for every $\kappa \geq \kappa_{x,\xi,\tau}$ since $\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\xi \cdot \nabla \varphi(x)\}\$ is bounded on K. By continuity of $G(x,\xi,\tau,\kappa)$, there exists a neighborhood of (x,ξ,τ) , denoted by $V_{x,\xi,\tau}$, such that $G(x,\xi,\tau,\kappa) > 0$ for all $(x,\xi,\tau) \in V_{x,\xi,\tau}$ and $\kappa \geq \kappa_{x,\xi,\tau}$. Since K is compact, there exist finite sets $V_j = V_{x_j,\xi_j,\tau_j}$ and corresponding constants $\kappa_j = \kappa_{x_j, \xi_j, \tau_j}$ $(j = 1, 2, \cdots, n)$, such that $\mathcal{K} \subset \bigcup_{j=1}^n V_j$ and $G(x, \xi, \tau, \kappa) > 0$ for all $(x,\xi,\tau) \in V_j$ and $\kappa > \kappa_j$. Let $\tilde{\kappa} = \max\{\kappa_j : j = 1, 2, \cdots, n\}$. It follows that $G(x,\xi,\tau,\tilde{\kappa}) > 0$ for all $(x,\xi,\tau) \in \mathcal{K}$ and $\kappa \geq \tilde{\kappa}$. Finally, using the compactness of K again, we conclude that there exists $\delta > 0$ such that $G(x, \xi, \tau, \tilde{\kappa}) \geq \delta$. Thus, (2.16) is reached since g is a homogeneous function of degree 2 with respect to variables (ξ, τ) .

By Gårding inequality (2.13), there exists a constant $C > 0$ such that, for $\tau \geq \tau_0$ with τ_0 sufficiently large,

$$
C \|\operatorname{Op}(\mu)w\|_{V}^{2} \leq \operatorname{Re}\left(\operatorname{Op}\left(|2\xi \cdot \nabla \varphi(x)|^{2} + \mu^{-2} \left(|\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}\right)^{2}\right) + \left\{|\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}, 2\xi \cdot \nabla \varphi(x)\right\} \right) w \, |w|_{V}.
$$
\n
$$
(2.17)
$$

Now we are going to estimate the terms $\tau \text{ Op}(|2\xi \cdot \nabla \varphi(x)|^2)$ and $\tau \text{ Op}(\mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2)$. Firstly, it follows from Lemma 2.3 that

$$
\tau^{-1} \operatorname{Op}(|2\tau\xi \cdot \nabla \varphi(x)|^2) = \tau^{-1} \operatorname{Op}(2\tau\xi \cdot \nabla \varphi(x))^* \operatorname{Op}(2\tau\xi \cdot \nabla \varphi(x)) + \tau \operatorname{Op}(r_1),\tag{2.18}
$$

where $r_1 \in S(\mu, \tilde{g})$ and \tilde{g} is defined by (2.12). Therefore, for any $\varepsilon > 0$, there exists a positive constant C_{ε} such that

$$
\begin{aligned}\n&\left| \left(\tau^{-1} \operatorname{Op}(\Vert 2\tau \xi \cdot \nabla \varphi(x) \Vert^2) w \Vert w \right)_{V} \right| \\
&\leq \tau^{-1} \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x))w\|_{V}^{2} + \tau \|\operatorname{Op}(r_{1})w\|_{V} \|\n\leq \tau^{-1} \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x))w\|_{V}^{2} + \varepsilon \tau \|\operatorname{Op}(\mu)w\|_{V}^{2} + C_{\varepsilon} \tau \|w\|_{V}^{2}.\n\end{aligned} \tag{2.19}
$$

Substituting (2.19) into (2.17) and choosing ε small enough, we have

$$
C\tau \|\operatorname{Op}(\mu)w\|_{V}^{2}
$$
\n
$$
\leq \operatorname{Re}\left(\operatorname{Op}\left(\tau\mu^{-2}\left(|\xi|^{2}-\tau^{2}|\nabla\varphi(x)|^{2}\right)^{2}+\tau\left\{||\xi|^{2}-\tau^{2}|\nabla\varphi(x)|^{2}, 2\xi\cdot\nabla\varphi(x)\right\}\right)w\|w\right)_{V} \qquad (2.20)
$$
\n
$$
+\tau^{-1}\|\operatorname{Op}(2\tau\xi\cdot\nabla\varphi(x))w\|_{V}^{2}+C_{\varepsilon}\tau\|w\|_{V}^{2}.
$$

Secondly, by symbolic calculus, we have that

$$
\tau \operatorname{Op}(\mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2) = \tau \operatorname{Op}(r_0) \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) + \tau \operatorname{Op}(r_2), \tag{2.21}
$$

where $r_0(x,\xi) = \mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \in S(1,\tilde{g})$ and $r_2 \in S(\mu,\tilde{g})$. Therefore, for all $\varepsilon > 0$, there exists $C_{\varepsilon} > 0$ such that

$$
\begin{aligned}\n| \left(\tau \operatorname{Op}(r_0) \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) w \, | \, w \right)_V \\
&\leq C_{\varepsilon} \tau^{-1} \| \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) w \|_V^2 + \varepsilon \tau^3 \| w \|_V^2. \n\end{aligned} \tag{2.22}
$$

We choose ε small enough and combine (2.21)-(2.22) with (2.20) to get

$$
C\tau \|\operatorname{Op}(\mu)w\|_{V}^{2} \leq \operatorname{Re}(\tau \operatorname{Op}\{| \xi |^{2} - \tau^{2} |\nabla \varphi(x)|^{2}, 2\xi \cdot \nabla \varphi(x) \} w \| w)_{V} + C_{\varepsilon} \tau^{-1} \|\operatorname{Op}(| \xi |^{2} - \tau^{2} |\nabla \varphi(x)|^{2}) w \|_{V}^{2} + \tau^{-1} \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x))w\|_{V}^{2} + C_{\varepsilon}(\tau + \varepsilon \tau^{3}) \| w \|_{V}^{2}.
$$
 (2.23)

Finally, it is clear that there exist positive constant C such that

$$
\tau^3 \|w\|_V^2 + \tau \|Dw\|_V^2 \le C\tau \|\operatorname{Op}(\mu)w\|_V^2. \tag{2.24}
$$

Thus, we obtain (2.15) by using (2.23)-(2.24), choosing ε small enough and letting $\tau > \tau_0$ large enough. \Box

3 Carleman Estimate

In this section, we shall prove several Carleman inequalities. Define the operator

$$
P(x, D, \lambda) = D^2 + i\lambda Da(x)D - \lambda^2.
$$

Let the weight function $\varphi \in \mathscr{C}^{\infty}(\mathbb{R}^d;\mathbb{R})$. The associated conjugate operator of $P(x, D, \lambda)$ is $P_{\varphi}(x, D, \lambda) =$ $e^{\tau \varphi} P(x, D, \lambda) e^{-\tau \varphi}$. Then,

$$
P_{\varphi} = (D + i\tau \nabla \varphi(x))^2 + i\lambda (D + i\tau \nabla \varphi(x))a(x)(D + i\tau \nabla \varphi(x)) - \lambda^2.
$$

By setting $Q_2 = \frac{1}{2}(P_{\varphi} + P_{\varphi}^*)$ and $Q_1 = \frac{1}{2i}(P_{\varphi} - P_{\varphi}^*)$, we have $P_{\varphi} = Q_2 + iQ_1$. We denote by $p(x, \xi, \lambda)$, $p_{\varphi}(x, \xi, \lambda)$ the associated symbol of $P(x, D, \lambda)$, $P_{\varphi}(x, D, \lambda)$, respectively.

Let the metric g and weight ν be defined by (2.1) and (2.2). Due to the results Remark 2.2, we know that $D + i\tau\nabla\varphi(x)$ is an operator with symbol in $S(\mu, g)$ class, λa is in $S(\nu, g)$, and $(1 +$ $i\lambda a(x)$ | $|\xi + i\tau \nabla \varphi(x)|^2$, the principal symbol of P_{φ} belongs to $S(\nu \mu^2, g)$. Furthermore, from Lemma 2.2 and 2.3, one has that

$$
P_{\varphi} = \text{Op}((1 + i\lambda a(x))|\xi + i\tau \nabla \varphi(x)|^2) - \lambda^2 + R_3,
$$

\n
$$
Q_2 = \text{Op}(q_2) - \lambda^2 + R_2,
$$

\n
$$
Q_1 = \text{Op}(q_1) + R_1,
$$
\n(3.1)

where $q_2 = |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2 - 2\lambda \tau a(x) \xi \cdot \nabla \varphi(x), q_1 = 2\tau \xi \cdot \nabla \varphi(x) + \lambda a(x) (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)$ belong to $S(\nu \mu^2, g)$ and the symbols of R_j is in $S(\lambda^{\frac{1}{2}}\nu\mu, g)$ for $j = 1, 2, 3$. It is clear that

$$
||P_{\varphi}v||_{V}^{2} = ||Q_{2}v||_{V}^{2} + ||Q_{1}v||_{V}^{2} + 2\operatorname{Re}(Q_{2}v \,|\, iQ_{1}v)_{V}, \quad \forall \, v \in \mathscr{C}^{\infty}(V). \tag{3.2}
$$

In what follows, several Carleman estimates are introduced. First, we give an estimation on the subdomain which is far away from the boundary Γ.

Theorem 3.1. Suppose φ satisfies sub-ellipticity condition in $V \subset \Omega$. Then, there exist positive constants C, \tilde{K} and λ_0 , such that for every $v \in \mathscr{C}_0^{\infty}(V)$, it holds

$$
\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} + \|Q_{1}v\|_{V}^{2} + \|Q_{2}v\|_{V}^{2} \lesssim \|P_{\varphi}v\|_{V}^{2},\tag{3.3}
$$

where $\lambda \geq \lambda_0$ and $\tau \geq \max{\{\tilde{K}|\lambda|^{\frac{5}{4}}}, 1\}.$

Proof. Let $w = \nu(x)v$ in (2.15). We obtain

$$
C_1 \tau^3 ||\nu(x)v||_V^2 + C_1 \tau ||D(\nu(x)v)||_V^2
$$

\n
$$
\leq \text{Re} \left(\text{Op}(\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\}) \nu(x)v |\nu(x)v \right)_V
$$

\n
$$
+ C_2 \tau^{-1} ||\text{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \nu(x)v||_V^2
$$

\n
$$
+ C_2 \tau^{-1} ||\text{Op}(2\tau \xi \cdot \nabla \varphi(x)) \nu(x)v||_V^2, C_1, C_2 > 0.
$$
\n(3.4)

Since the symbol of $[D, \nu(x)]$ is in $S(\lambda^{\frac{1}{2}}\nu, g)$, we have

$$
\|\nu(x)Dv\|_V^2 \le \|D(\nu(x)v)\|_V^2 + \|[D,\nu(x)]v\|_V^2 \le \|D(\nu(x)v)\|_V^2 + C\lambda \|\nu(x)v\|_V^2.
$$

Consequently, for $\tau \geq C\lambda$ with $C > 0$ sufficiently large, it holds

$$
\tau^3 \|\nu(x)v\|_V^2 + \tau \|\nu(x)Dv\|_V^2 \lesssim \tau^3 \|\nu(x)v\|_V^2 + \tau \|D(\nu(x)v)\|_V^2. \tag{3.5}
$$

It follows from (3.4) and (3.5) that

$$
C'_{1}(\tau^{3}||\nu(x)v||_{V}^{2} + \tau ||\nu(x)Dv||_{V}^{2})
$$

\n
$$
\leq \text{Re}\left(\text{Op}(\{|\xi|^{2} - \tau^{2}\varphi(x)^{2}, 2\tau\xi \cdot \nabla\varphi(x)\})\nu(x)v|\nu(x)v\right)_{V}
$$

\n+*C*₂ $\tau^{-1}||\text{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})\nu(x)v||_{V}^{2}$
\n+*C*₂ $\tau^{-1}||\text{Op}(2\tau\xi \cdot \nabla\varphi(x))\nu(x)v||_{V}^{2}, C'_{1}, C_{2} > 0.$ (3.6)

Now, we estimate the first term on the right side hand of (3.6). Note that

$$
2 \operatorname{Re} (Q_2 v \, | \, iQ_1 v)_V = (Q_2 v \, | \, iQ_1 v)_V + (iQ_1 v \, | \, Q_2 v)_V = (i[Q_2, Q_1] v \, | \, v)_V, \tag{3.7}
$$

where the principal symbol of $i[Q_2, Q_1]$ is $\{q_2, q_1\}$. Due to Lemma 2.3, we obtain that

$$
i[Q_2, Q_1] = \text{Op}(\{q_2, q_1\}) + R_4,\tag{3.8}
$$

where $\{q_2, q_1\} \in S(\lambda^{\frac{1}{2}}\nu^2\mu^3, g)$ and the symbol of R_4 is in $S(\lambda \nu^2\mu^2, g)$. A direct computation gives that

$$
\{q_2, q_1\} = (1 + a^2(x)\lambda^2) \{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\} + \left(\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, \lambda a(x)\} - \lambda a(x) \{2\tau \xi \cdot \nabla \varphi(x), \lambda a(x)\}\right) (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) - 2\tau \xi \cdot \nabla \varphi(x) \left(\lambda a(x) \{\lambda a(x), |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2\} + \{\lambda a(x), 2\tau \xi \cdot \nabla \varphi(x)\}\right).
$$

From the definition of q_1 and q_2 , we have

$$
q_2(x,\xi) + \lambda a(x)q_1(x,\xi) = (1 + \lambda^2 a(x)^2)(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2),
$$

\n
$$
q_1(x,\xi) - \lambda a(x)q_2(x,\xi) = 2\tau\xi \cdot \nabla \varphi(x)(1 + \lambda^2 a^2(x)).
$$
\n(3.9)

Consequently, ϵ

 \sim

$$
\{q_2, q_1\}
$$

= $\nu^2(x)\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, 2\tau\xi \cdot \nabla\varphi(x)\}\$
+ $\nu^{-2}(x)(q_2(x,\xi) + \lambda a(x)q_1(x,\xi))(\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, \lambda a(x)\} - \lambda a(x)\{2\tau\xi \cdot \nabla\varphi(x), \lambda a(x)\})$
- $\nu^{-2}(x)(q_1(x,\xi) - \lambda a(x)q_2(x,\xi))(\lambda a(x)\{\lambda a(x), |\xi|^2 - \tau^2|\nabla\varphi(x)|^2\} + \{\lambda a(x), 2\tau\xi \cdot \nabla\varphi(x)\}).$

Then, it follows from (3.8) and the above equation that

$$
i[Q_2, Q_1] = \nu(x) \operatorname{Op} \left(\{ |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x) \} \right) \nu(x)
$$

$$
+ B_1 \nu^{-1}(x) \operatorname{Op} \left(q_2(x, \xi) + \lambda a(x) q_1(x, \xi) \right)
$$

$$
- B_2 \nu^{-1}(x) \operatorname{Op} \left(q_1(x, \xi) - \lambda a(x) q_2(x, \xi) \right) + R_5,
$$

where

$$
B_1 = \text{Op}\left(\nu^{-1}(x)\left(\left\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, \lambda a(x)\right\} - \lambda a(x)\left\{2\tau\xi \cdot \nabla\varphi(x), \lambda a(x)\right\}\right)\right),
$$

\n
$$
B_2 = \text{Op}\left(\nu^{-1}(x)\left(\lambda a(x)\left\{\lambda a(x), |\xi|^2 - \tau^2|\nabla\varphi(x)|^2\right\} + \left\{\lambda a(x), 2\tau\xi \cdot \nabla\varphi(x)\right\}\right)\right),
$$

symbols of B_1 and B_2 belong to $S(\lambda^{\frac{1}{2}}\nu\mu, g)$ and the symbol of R_5 is in $S(\lambda \nu^2 \mu^2, g)$. Combining this with (3.1) yields

$$
i[Q_2, Q_1] = \nu(x) \operatorname{Op} (\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\}) \nu(x) + B_1 \nu^{-1}(x) [Q_2(x, \xi) + \lambda a(x) Q_1(x, \xi) + \lambda^2] - B_2 \nu^{-1}(x) [Q_1(x, \xi) - \lambda a(x) Q_2(x, \xi) - \lambda^3 a(x)] + R_5,
$$
(3.10)

We refer to Section 2 where the rules on symbolic calculus are given and precise. Therefore, by the continuity of pseudo-differential operator, we have that for $j = 1, 2, k = 0, 1$ and $\ell = 1, 2$,

$$
\begin{split}\n&|\left(B_{j}\nu(x)^{-1}(\lambda a(x))^{k}Q_{\ell}v\,|\,v\right)_{V}| \\
&= \quad |\left(\nu(x)^{-1}(\lambda a(x))^{k}Q_{\ell}v\,|\,B_{j}^{*}v\right)_{V}| \lesssim \|Q_{\ell}v\|_{V}\|B_{j}^{*}v\|_{V} \\
&\leq \quad \frac{1}{10}\|Q_{\ell}v\|_{V}^{2} + C\lambda\tau^{2}\|\nu(x)v\|_{V}^{2} + C\lambda\|\nu(x)Dv\|_{V}^{2}, \quad C > 0, \\
&|\left(B_{j}\nu^{-1}(x)\lambda^{2}(\lambda a(x))^{k}v\,|\,v\right)_{V}| \lesssim \lambda^{\frac{5}{2}}\left(\tau\|\nu(x)v\|_{V} + \|\nu(x)Dv\|_{V}\right)\|\nu(x)v\|_{V} \\
&\leq \lambda^{\frac{5}{2}}\tau\|\nu(x)v\|_{V}^{2} + \lambda^{\frac{5}{2}}\tau^{-1}\|\nu(x)Dv\|_{V}^{2},\n\end{split} \tag{3.12}
$$

and

$$
|(R_5v|v)_V| \lesssim \lambda \tau^2 \|\nu(x)v\|_V^2 + \lambda \|\nu(x)Dv\|_V^2. \tag{3.13}
$$

Due to (3.10) – (3.13) , there exists a positive constant C such that

$$
\begin{aligned} \left(\operatorname{Op}(\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\}) \nu(x) v \, | \, \nu(x) v \right)_V \\ \leq \quad & \left(i[Q_2, Q_1] v \, | \, v \right)_V + \frac{1}{4} \sum_{\ell=1,2} \|Q_\ell v\|_V^2 + C \left(\lambda \tau^2 \| \nu(x) v\|_V^2 \right) \\ + & \lambda \| \nu(x) D v\|_V^2 + \lambda^{\frac{5}{2}} \tau \| \nu(x) v\|_V^2 + \lambda^{\frac{5}{2}} \tau^{-1} \| \nu(x) D v\|_V^2 \right). \end{aligned} \tag{3.14}
$$

Next, we are going to estimate the last two terms on the right side hand of (3.6). It follows from (3.1) and (3.9) that

$$
\|\operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \nu(x) v\|_V^2 + \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x)) \nu(x) v\|_V^2
$$
\n
$$
= \|\operatorname{Op}(\nu(x)^{-2} (q_2 + \lambda a(x) q_1) \nu(x) v\|_V^2 + \|\operatorname{Op}(\nu(x)^{-2} (q_1 - \lambda a(x) q_2) \nu(x) v\|_V^2)
$$
\n
$$
= \|\nu(x)^{-1} [Q_2 + \lambda^2 - R_2 + \lambda a(x) (Q_1 - R_1)] v\|_V^2
$$
\n
$$
+ \|\nu(x)^{-1} [Q_1 - R_1 - \lambda a(x) (Q_2 + \lambda^2 - R_2)] v\|_V^2.
$$
\n(3.15)

Combining this with the fact that the symbols of R_j are in $S(\lambda^{\frac{1}{2}}\nu\mu, g)$ for $j = 1, 2$, we have

$$
\| \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \nu(x) v \|_{V}^2 + \| \operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x)) \nu(x) v \|_{V}^2
$$

\n
$$
\lesssim \sum_{\ell=1,2} (\|Q_{\ell}v\|_{V}^2 + \|R_{\ell}v\|_{V}^2) + \lambda^4 \|v\|_{V}^2
$$

\n
$$
\lesssim \sum_{\ell=1,2} \|Q_{\ell}v\|_{V}^2 + \lambda \tau^2 \|\nu(x)v\|_{V}^2 + \lambda \|\nu(x)Dv\|_{V}^2 + \lambda^4 \|v\|_{V}^2.
$$
\n(3.16)

Consequently, for $\tau \geq C\lambda$ with C large enough, it holds

$$
\tau^{-1} \|\operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \nu(x) v\|_V^2 + \tau^{-1} \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x)) \nu(x) v\|_V^2
$$

$$
\lesssim \tau^{-1} \sum_{\ell=1,2} \|Q_\ell v\|_V^2 + \lambda \tau \|\nu(x)v\|_V^2 + \|\nu(x)Dv\|_V^2 + \lambda^3 \|v\|_V^2.
$$
 (3.17)

Finally, by (3.6), (3.14) and (3.17), one can choose $\tau \ge \max{\{\tilde{K}|\lambda|^{\frac{5}{4}}}, 1\}$ with \tilde{K} sufficiently large such that for some $C > 0$,

$$
C(\tau^3 \|\nu(x)v\|_{V}^2 + \tau \|\nu(x)Dv\|_{V}^2) \leq (i[Q_2, Q_1]v\,|v)_{V} + \frac{1}{2} \sum_{\ell=1,2} \|Q_{\ell}v\|_{V}^2 + \tau^3 \|\nu(x)v\|_{V}^2 + \varepsilon \tau \|\nu(x)Dv\|_{V}^2,
$$
\n(3.18)

 \Box

where $\varepsilon > 0$ is arbitrary. Choosing ε small with respect to C, using (3.2), (3.7) and (3.18), we obtain

$$
C_1(\tau^3 \|\nu(x)v\|_{V}^2 + \tau \|\nu(x)Dv\|_{V}^2) + \frac{1}{2} (||Q_1v\|_{V}^2 + ||Q_2v\|_{V}^2)
$$

\n
$$
\leq (i[Q_2, Q_1]v, v)_V + ||Q_1v\|_{V}^2 + ||Q_2v\|_{V}^2 = ||P_{\varphi}v||_{V}^2,
$$

which implies (3.3).

Remark 3.1. The estimates in (3.12) impose the assumption $\tau \geq \tilde{K} |\lambda|^{\frac{5}{4}}$. The other remainder terms only impose the condition $\tau \geq C\lambda$. This condition is related with the principal normal condition. Indeed for a complex operator, with symbol $p_1 + ip_2$ where p_1, p_2 are both real valued, the Carleman estimate is only true if $\{p_1, p_2\} = 0$ on $p_1 = p_2 = 0$. Here the symbol of operator before conjugation by weight is $|\xi|^2 - \lambda^2 + i\lambda a(x)|\xi|^2$, and the Poisson bracket is $\{|\xi|^2 - \lambda^2, \lambda a(x)|\xi|^2\} = 2\lambda(\xi \cdot \nabla a(x))|\xi|^2$. We can estimate this term, uniformly in a neighborhood of $a(x) = 0$, by $C\lambda a^{\frac{1}{2}}(x) |\xi|^{3}$. This explanation does not justify the power $|\lambda|^{\frac{5}{4}}$ found at the end of computations but shows the difficulties.

Theorem 3.2. Suppose φ satisfies sub-ellipticity condition in $V \subset \Omega$. Then, there exist positive constants \tilde{K} and λ_0 , such that for every $u \in \mathscr{C}_0^{\infty}(V)$, it holds

$$
\tau^3 \|(1+\lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} u\|_V^2 + \tau \|(1+\lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} Du\|_V^2 \lesssim \|e^{\tau \varphi} Pu\|_V^2, \tag{3.19}
$$

where $\lambda \geq \lambda_0$ and $\tau \geq \max{\{\tilde{K}|\lambda|^{\frac{5}{4}}}, 1\}.$

Proof. Set $v = e^{\tau \varphi} u$. From Theorem 3.1, it suffices to prove that (3.19) is equivalent to

$$
\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} \lesssim \|P_{\varphi}v\|_{V}^{2}.
$$
\n(3.20)

First, assume (3.20) holds. Then, $Dv = e^{\tau \varphi} (Du - i\tau \nabla \varphi u)$ and $e^{\tau \varphi} Du = Dv + i\tau \nabla \varphi v$. Then there exist positive constants c_1 , c_2 such that

$$
c_1 \left(\|\nu(x)Dv\|_V + \tau \|\nu(x)v\|_V \right) \leq \|\nu(x)e^{\tau\varphi}Du\|_V + \tau \|\nu(x)e^{\tau\varphi}u\|_V
$$

$$
\leq c_2 \left(\|\nu(x)Dv\|_V + \tau \|\nu(x)v\|_V \right).
$$
 (3.21)

Combining this with (3.20), we conclude that

$$
\tau^{3} \|\nu(x)e^{\tau\varphi}u\|_{V}^{2} + \tau \|\nu(x)e^{\tau\varphi}Du\|_{V}^{2} \lesssim (\tau \|\nu(x)Dv\|_{V}^{2} + \tau^{3} \|\nu(x)v\|_{V}^{2}) \lesssim \|e^{\tau\varphi}Pu\|_{V}^{2}.
$$

On the other hand, (3.19) implies that

$$
\tau^3\|\nu(x)e^{\tau\varphi}u\|_V^2+\tau\|\nu(x)e^{\tau\varphi}Du\|_V^2\lesssim \|P_\varphi v\|_V^2.
$$

Then, we proved (3.20) from the above estimate and (3.21).

Since there is higher order term $div(a(x)\nabla y_t)$ in system (1.1), it is necessary to deal with the term $div(a(x)\nabla f)$ for $f \in H^1(\Omega)$ when proving the resolvent estimate. The following result is analogue to the work by Imanuvilov and Puel ([14]).

Theorem 3.3. Suppose φ satisfies sub-ellipticity condition on $V \subset \Omega$. Then, there exist C, \tilde{K} , $\lambda_0 >$ 0, such that for all $u \in \mathscr{C}_0^{\infty}(V)$ satisfying

$$
P(x, D, \lambda)u = g_0 + \sum_{j=1}^{d} \partial_{x_j} g_j, \text{ where } g_j \in L^2(V), j = 0, 1, \cdots, d,
$$
 (3.22)

it holds

$$
\tau \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} u \|_{V}^2 + \tau^{-1} \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} Du \|_{V}^2 \le C \sum_{j=0}^d \| e^{\tau \varphi} g_j \|_{V}^2, \tag{3.23}
$$

where $\lambda \geq \lambda_0$ and $\tau \geq \max{\{\tilde{K}|\lambda|^{\frac{5}{4}}}, 1\}.$

Proof. First, from (3.1), we have $D^2 = Q_2 + S_2$ and $\lambda a(x)D^2 = Q_1 + S_1$ where S_1 and S_2 have symbols in $S(\tau \nu \mu, g)$ if $\tau \gtrsim \lambda$. It follows that for any $v \in \mathscr{C}_0^{\infty}(V)$,

$$
||D^{2}v||_{V}^{2} \lesssim ||Q_{2}v||_{V}^{2} + \tau^{2}(\tau^{2}||\nu(x)v||_{V}^{2} + ||\nu(x)Dv||_{V}^{2}),
$$

$$
||\lambda aD^{2}v||_{V}^{2} \lesssim ||Q_{1}v||_{V}^{2} + \tau^{2}(\tau^{2}||\nu(x)v||_{V}^{2} + ||\nu(x)Dv||_{V}^{2}).
$$
\n(3.24)

Using (3.3), (3.24) and the fact that $\|\nu(x)D^2v\|_V^2 \le 2(\|D^2v\|_V^2 + \|\lambda a(x)D^2v\|_V^2)$, we obtain

$$
\tau^3 \|\nu(x)v\|_V^2 + \tau \|\nu(x)Dv\|_V^2 + \tau^{-1} \|\nu(x)D^2v\|_V^2 \lesssim \|P_\varphi v\|_V^2. \tag{3.25}
$$

Let \tilde{u} and χ be in $\mathscr{C}_0^{\infty}(\Omega)$ such that $\chi = 1$ on a neighborhood of supp \tilde{u} . Similarly to (3.21), we obtain

$$
\tau \|\nu(x)\tilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\tilde{u}\|_{V}^{2} \lesssim \tau \|\nu(x)\tilde{u}\|_{V}^{2} + \tau^{-1} \|D(\nu(x)\tilde{u})\|_{V}^{2}.
$$

Then, combining this with Fourier transform and the following inequality

$$
\tau + \frac{|\xi|^2}{\tau} \lesssim \frac{\tau^3}{\tau^2 + |\xi|^2} + \frac{|\xi|^4}{\tau(\tau^2 + |\xi|^2)},
$$

we conclude that

$$
\tau \|\nu(x)\tilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\tilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\operatorname{Op}(\mu^{-1})\nu(x)\chi\tilde{u}\|_{V}^{2} + \tau^{-1} \|D^{2}\operatorname{Op}(\mu^{-1})\nu(x)\chi w\|_{V}^{2}.
$$
 (3.26)

From Lemma 2.3, we have $\text{Op}(\mu^{-1})\nu\chi = \nu\chi \text{Op}(\mu^{-1}) + R_1$, where R_1 has a symbol in $S(\mu^{-2}\nu\lambda^{\frac{1}{2}}, g)$, and $D^2 \text{Op}(\mu^{-1}) \nu \chi = \nu D^2 \chi \text{Op}(\mu^{-1}) + R_2$, where R_2 has a symbol in $S(\nu \lambda^{\frac{1}{2}}, g)$. Then, it follows from (3.26) that

$$
\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\nu(x)\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D^{2}\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2} + \lambda \|\nu(x)\widetilde{u}\|_{V}^{2}.
$$

For $\tau \ge \max\{\mathcal{C}\lambda, 1\}$ with $\mathcal C$ large enough, one has the following result from the above inequality.

$$
\tau \|\nu(x)\tilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\tilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\nu(x)\chi \operatorname{Op}(\mu^{-1})\tilde{u}\|_{V}^{2} + \tau^{-1} \|\nu D^{2}\chi \operatorname{Op}(\mu^{-1})\tilde{u}\|_{V}^{2}.
$$
 (3.27)

Now, we apply (3.25) to $v = \chi \text{Op}(\mu^{-1})\tilde{u}$ to have

$$
\tau^3 \|\nu(x)\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^2 + \tau \|\nu(x)D\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^2 + \tau^{-1} \|\nu(x)D^2\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^2
$$

$$
\lesssim \|P_{\varphi}\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^2.
$$

Thus, combining this with (3.27) yields

$$
\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \|P_{\varphi}\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2}.
$$
\n(3.28)

Finally, note that P_{φ} has a symbol in $S(\nu\mu^2, g)$. Consequently, $P_{\varphi}\chi \operatorname{Op}(\mu^{-1}) = \operatorname{Op}(\mu^{-1})P_{\varphi}\chi + R$, where R has a symbol in $S(\nu \lambda^{\frac{1}{2}}, g)$. Then, we can deduce from (3.28) that

$$
\tau \|\nu(x)\widetilde{u}\|_{V}^2 + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^2 \lesssim \|\operatorname{Op}(\mu^{-1})P_{\varphi}\widetilde{u}\|_{V}^2 + \lambda \|\nu \widetilde{u}\|_{V}^2.
$$

When $\tau \geq C\lambda$ with C large enough, the error term $\lambda ||\nu \tilde{u}||_V^2$ can be absorbed by the left hand side. Consequently,

$$
\tau \|\nu(x)\tilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\tilde{u}\|_{V}^{2} \lesssim \|\operatorname{Op}(\mu^{-1})P_{\varphi}\tilde{u}\|_{V}^{2}.
$$
\n(3.29)

For $\tilde{u} = e^{\tau \varphi} u$, it follows from (3.29) and similar argument as (3.21) that

$$
\tau \|\nu(x)e^{\tau\varphi}u\|_{V}^{2} + \tau^{-1} \|\nu(x)e^{\tau\varphi}Du\|_{V}^{2} \leq \tau \|\nu(x)e^{\tau\varphi}u\|_{V}^{2} + \tau^{-1} \|\nu(x)De^{\tau\varphi}u\|_{V}^{2}
$$
\n
$$
\leq \|\operatorname{Op}(\mu^{-1})P_{\varphi}e^{\tau\varphi}u\|_{V}^{2}.
$$
\n(3.30)

Obviously, one has that

$$
P_{\varphi}\widetilde{u} = e^{\tau \varphi} Pu = e^{\tau \varphi} g_0 + \sum_{j=1}^d e^{\tau \varphi} \partial_{x_j} g_j = e^{\tau \varphi} g_0 + \sum_{j=1}^d \left(\partial_{x_j} (e^{\tau \varphi} g_j) - \tau e^{\tau \varphi} g_j \partial_{x_j} \varphi \right),
$$

which yields

$$
\|\operatorname{Op}(\mu^{-1})P_{\varphi}\tilde{u}\|_{V}^{2} \lesssim \sum_{j=0}^{d} \|e^{\tau\varphi}g_{j}\|_{V}^{2}.
$$
\n(3.31)

Hence, we obtain Theorem 3.3 from (3.30) and (3.31).

 \Box

Remark 3.2. Since $a(\cdot)$ is nonnegative and not identically null, there exists $\delta > 0$ such that $\{x \in$ $\Omega : a(x) > \delta$ $\neq \emptyset$. We introduce several sets as follows.

$$
W_0 = \Omega \setminus \text{supp } a,
$$

\n
$$
W_1 = \overline{\Omega} \setminus \mathcal{O}(\text{supp } a),
$$

\n
$$
W_2 = \Omega \setminus (\{x \in \Omega : a(x) \ge \delta\} \cup \mathcal{O}(\Gamma)),
$$

\n
$$
W_3 = \Omega \setminus \{x \in \Omega : a(x) \ge \frac{\delta}{2}\},
$$

where $\mathcal{O}(\Gamma)$ means the neighborhood of Γ .

It is known that there exists a function $\psi \in \mathscr{C}^{\infty}(\Omega)$ such that ([11])

- 1) $\psi(x) = 0$ for $x \in \partial \Omega$.
- 2) $\partial_n \psi(x) < 0$ for $x \in \partial \Omega$.
- 3) $\nabla \psi(x) \neq 0$ for $x \in \overline{\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}}$.

Let $\varphi = e^{\gamma \psi}$. It follows from Lemma 2.5 that φ satisfies the sub-ellipticity condition on $x \in \Omega \setminus \{x \in \Omega\}$ Ω : $a(x) \geq \delta$ if $\gamma > 0$ is sufficiently large. Then, in Theorem 3.2 and 3.3, one can choose V as $\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}.$

The following result is a classical Carleman estimate and corresponding to the Laplacian with Dirichlet boundary condition ([17], Proposition 2). We shall use it to deal with the terms on Ω \supp a.

Lemma 3.1. Suppose φ is chosen as in Remark 3.2. Then, there exist $C > 0$, $\lambda_0 > 0$, such that for all $u \in \mathscr{C}^{\infty}(\overline{\Omega})$ satisfying supp $u \subset W_1$ and $u = 0$ on Γ , it holds

$$
\tau^3 \|e^{\tau \varphi} u\|_{\Omega}^2 + \tau \|e^{\tau \varphi} Du\|_{\Omega}^2 \lesssim \|e^{\tau \varphi} (D^2 - \lambda^2) u\|_{\Omega}^2,\tag{3.32}
$$

where $\lambda \geq \lambda_0$ and $\tau \geq \max\{\mathcal{C}\lambda, 1\}.$

Theorem 3.4. Suppose φ is chosen as in Remark 3.2. Let $u \in \mathscr{C}^{\infty}(\overline{\Omega})$ and satisfy

$$
P(x, D, \lambda)u = f_0 + \sum_{j=1}^{d} \partial_{x_j} f_j \quad in \quad \Omega,
$$

$$
u = 0 \qquad on \quad \Gamma,
$$

where $f_j \in L^2(\Omega)$, supp $f_0 \subset \Omega$ and supp $f_j \subset \mathcal{O}(\text{supp } a)$ for $j = 1, \dots, d$. Then, there exist $\tilde{K} > 0$, $\lambda_0 > 0$, such that for all $\lambda \geq \lambda_0$ and $\tau \geq \max{\{\tilde{K}|\lambda|^{\frac{5}{4}}}, 1\}$, it holds

$$
\tau \|e^{\tau \varphi} u\|_{W_3}^2 + \tau^{-1} \|e^{\tau \varphi} Du\|_{W_3}^2 \lesssim \sum_{j=0}^d \|e^{\tau \varphi} f_j\|_{\Omega}^2 + \lambda \|e^{\tau \varphi} u\|_{\{x \in \Omega : \, a(x) \ge \delta/2\}}^2,
$$

where the positive constant δ is defined as in Remark 3.2.

Proof. Let $\chi_1, \chi_2 \in \mathscr{C}^\infty(\Omega)$ be non-negative and satisfy the following assumption

- (i) $0 \leq \chi_1, \chi_2 \leq 1$; χ_1 and χ_2 are supported on W_1 and W_2 , respectively.
- (ii) $\chi_1 + \chi_2 \geq 1$ in W_3 . In particular, $\chi_1 \equiv 1$ on $[\mathcal{O}(\partial \Omega) \cap \Omega] \setminus \mathcal{O}(\text{supp }a)$, and $\chi_2 \equiv 1$ on $\mathcal{O}(\text{supp }a) \setminus \{x \in \Omega : a(x) \geq \frac{\delta}{2}\}.$

First, it is clear that

$$
P\chi_2 u = \chi_2 f_0 + \sum_{j=1}^d \left(\partial_{x_j} (\chi_2 f_j) - f_j \partial_{x_j} \chi_2 \right) + [P, \chi_2] u.
$$

Since $[P, \chi_2]$ is a first order operator, we have there exist a_0, a_1, \dots, a_d and b_0, b_1, \dots, b_d such that

$$
[P, \chi_2]u = \sum_{j=0}^d a_j \partial_{x_j} u + \lambda \sum_{j=0}^d b_j \partial_{x_j} u,
$$

where supp $a_j \subset \{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\} \cup [\Omega \setminus (\mathcal{O}(\text{supp } a) \cup \mathcal{O}(\text{supp } \partial \Omega))]$, supp $b_j \subset \{x \in \Omega \}$ Ω : $\frac{\delta}{2} < a(x) < \delta$ for $j = 0, 1, \dots, d$. Then, applying Theorem 3.3 with $\chi_2 u$ instead of u, $\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}$ instead of V, we obtain

$$
\tau \|(1+\lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} \chi_2 u\|_{\Omega}^2 + \tau^{-1} \|(1+\lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} D \chi_2 u\|_{\Omega}^2
$$

$$
\lesssim \sum_{j=0}^d \|e^{\tau \varphi} f_j\|_{\Omega}^2 + \|e^{\tau \varphi} u\|_{\{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\} \cup (\Omega \setminus \text{supp } a)}^2 + \lambda \|e^{\tau \varphi} u\|_{\{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\}}^2.
$$

Consequently, due to $(1 + \lambda^2 a(x)^2) \ge 1$, $\tau \ge {\mathcal C \lambda}$, 1} and $\lambda \ge \lambda_0$, we obtain

$$
\tau \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} \chi_2 u \|_{\Omega}^2 + \tau^{-1} \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} D \chi_2 u \|_{\Omega}^2
$$
\n
$$
\lesssim \sum_{j=0}^d \| e^{\tau \varphi} f_j \|_{\Omega}^2 + \| e^{\tau \varphi} u \|_{\Omega \backslash \text{supp } a}^2 + \lambda \| e^{\tau \varphi} u \|_{\{x \in \Omega : a(x) \ge \frac{\delta}{2}\}}^2.
$$
\n(3.33)

On the other hand, by using $\chi_1 u$ instead of u in Lemma 3.1, we obtain

$$
\tau^3 \|e^{\tau \varphi} \chi_1 u\|_{\Omega}^2 + \tau \|e^{\tau \varphi} D\chi_1 u\|_{\Omega}^2 \lesssim \|e^{\tau \varphi} (P - i\lambda D a(x) D)\chi_1 u\|_{\Omega}^2. \tag{3.34}
$$

Since $\chi_1 \partial_{x_j} f_j = 0$ for $j = 1, \dots, d$, we have

$$
P\chi_1 u = \chi_1 f_0 + [P, \chi_1]u.
$$

Therefore, combining these with (3.34) and $\mathcal{O}(\text{supp } a) \cap \text{supp } \chi_1 = \emptyset$ yields

$$
\tau \|e^{\tau \varphi} \chi_1 u\|_{\Omega}^2 + \tau^{-1} \|e^{\tau \varphi} D\chi_1 u\|_{\Omega}^2 \lesssim \tau^{-2} (\|e^{\tau \varphi} f_0\|_{\Omega}^2 + \|e^{\tau \varphi} [P, \chi_1] u\|_{\Omega}^2).
$$
 (3.35)

Adding up (3.35) and (3.33), using $(1 + \lambda^2 a(x)^2) \ge 1$ and $\tau \ge 1$, we obtain

$$
\tau \|e^{\tau\varphi}(\chi_1 + \chi_2)u\|_{\Omega}^2 + \tau^{-1} \|e^{\tau\varphi}D(\chi_1 + \chi_2)u\|_{\Omega}^2
$$

$$
\lesssim \sum_{j=0}^d \|e^{\tau\varphi}f_j\|_{\Omega}^2 + \|e^{\tau\varphi}u\|_{\Omega\backslash \mathrm{supp}\ a}^2 + \lambda \|e^{\tau\varphi}u\|_{\{x \in \Omega : a(x) \ge \frac{\delta}{2}\}}^2 + \tau^{-2} \|e^{\tau\varphi}[P, \chi_1]u\|_{\Omega}^2.
$$

Note that $[P, \chi_1]$ is a first order operator and supported on supp $\chi_1 \setminus \{x \in \Omega : \chi_1 \equiv 1\}$, which is an subset of $\{x \in \Omega : \chi_1 + \chi_2 = 1\}$. Then, for τ sufficiently large, we have

$$
\tau \|e^{\tau\varphi}(\chi_1+\chi_2)u\|_{\Omega}^2 + \tau^{-1} \|e^{\tau\varphi}(\chi_1+\chi_2)Du\|_{\Omega}^2 \lesssim \sum_{j=0}^d \|e^{\tau\varphi}f_j\|_{\Omega}^2 + \lambda \|e^{\tau\varphi}u\|_{\{x\in\Omega\,:\,a(x)\ge\frac{\delta}{2}\}}^2.
$$

Then as $1 \leq \chi_1 + \chi_2 \leq 2$ on W_3 , we obtain the statement of Theorem 3.4.

 \Box

4 Resolvent estimate

In this section, we shall prove the main result. From the results in [3, 7, 10], the logarithmic decay of the energy in Theorem 1.1 can be obtained through the following theorem.

Theorem 4.1. Suppose conditions in Theorem 1.1 hold. Then, for every $\lambda \in \mathbb{R}$ with $|\lambda|$ large enough, there exists $C > 0$ such that

$$
\left\| (\mathcal{A} - i\lambda)^{-1} \right\|_{\mathcal{L}(\mathcal{H})} \lesssim e^{C \|\lambda\|^{\frac{5}{4}}}.
$$
\n(4.1)

Let λ be a real number such that $|\lambda|$ is large enough. Consider the resolvent equation:

$$
F = (\mathcal{A} - i\lambda)Y, \quad \text{where } Y = (y_1, y_2) \in D(\mathcal{A}), \quad F = (f_1, f_2) \in \mathcal{H}, \tag{4.2}
$$

which yields

$$
\begin{cases}\n\operatorname{div} (\nabla y_1 + a \nabla y_2) + \lambda^2 y_1 = i \lambda f_1 + f_2, & \text{in } \Omega, \\
y_2 = i \lambda y_1 + f_1 & \text{in } \Omega, \\
y_1 |_{\Gamma} = 0\n\end{cases}
$$
\n(4.3)

In what follows, we shall prove that there exists a positive constant C such that

$$
||(y_1, y_2)||_{\mathcal{H}} \lesssim e^{C \, |\,\lambda|^{\frac{5}{4}}} \|(f_1, f_2)||_{\mathcal{H}}.
$$

First, let $\eta > 0$ and $\chi \in \mathscr{C}_0^{\infty}(\Omega)$ be real valued function such that supp $\chi = \{x \in \Omega : a(x) > \eta\}$ and $\chi \equiv 1$ in $\{x \in \Omega : a(x) > 2\eta\}$. The following lemma is helpful.

Lemma 4.1. For y_1 , f_1 , f_2 satisfying (4.3) and $|\lambda|$ large enough, it holds

- (i) $\|\chi y_1\|_{\Omega} \lesssim \|\nabla y_1\|_{\{x \in \Omega : a(x) > \eta\}} + \|f_1\|_{H^1(\Omega)} + \|f_2\|_{\Omega}.$
- (ii) $\int_{\Omega} a(x) |\nabla y_1|^2 dx \lesssim ||f_1||_{H^1(\Omega)}^2 + (||f_1||_{\Omega} + ||f_2||_{\Omega}) ||y_1||_{\Omega}.$

Proof. (i) Multiplying the first equation in (4.3) by $\chi^2 \overline{y}_1$ and using $y_2 = i\lambda y_1 + f_1$, we obtain,

$$
\lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} = \int_{\Omega} \nabla y_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx + i\lambda \int_{\Omega} a(x) \nabla y_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx \n+ i\lambda \int_{\Omega} f_{1} \chi^{2} \overline{y}_{1} dx + \int_{\Omega} f_{2} \chi^{2} \overline{y}_{1} dx + \int_{\Omega} a(x) \nabla f_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx.
$$
\n(4.4)

Since $\nabla(\chi^2 \overline{y}_1) = \chi^2 \nabla \overline{y}_1 + 2 \overline{y}_1 \chi \nabla \chi$, we have

$$
\left| \int_{\Omega} a(x) \nabla y_1 \cdot \nabla (\chi^2 \overline{y}_1) dx \right| \leq \|\nabla y_1\|_{\{x \in \Omega : a(x) > \eta\}}^2 + \|\chi y_1\|_{\Omega} \|\nabla y_1\|_{\{x \in \Omega : a(x) > \eta\}} \leq 2 \|\nabla y_1\|_{\{x \in \Omega : a(x) > \eta\}}^2 + \frac{1}{4} \|\chi y_1\|_{\Omega}^2.
$$
\n(4.5)

By the similar argument, we can deal with the rest terms on the right hand side of (4.4). Combining these with (4.4) , (4.5) yields

$$
\lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} \leq \frac{1}{2}\lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} + (|\lambda|+1) \|\nabla y_{1}\|_{\{x \in \Omega : a(x) > \eta\}}^{2} + \|f_{1}\|_{\Omega}^{2} + \lambda^{-2} \|f_{2}\|_{\Omega}^{2} + (\lambda^{-2}+1) \|\nabla f_{1}\|_{H^{1}(\Omega)}^{2}.
$$
\n(4.6)

Then, (i) is reached.

(ii) Multiplying the first equation in (4.3) by \overline{y}_1 and using $y_2 = i\lambda y_1 + f_1$, we obtain,

$$
\int_{\Omega} \left[-(\iota \lambda a(x) + 1) |\nabla y_1|^2 + \lambda^2 |y_1|^2 \right] dx = \int_{\Omega} \left[a(x) \nabla f_1 \cdot \nabla \overline{y}_1 + (\iota \lambda f_1 + f_2) \overline{y}_1 \right] dx. \tag{4.7}
$$

Taking the imaginary part of (4.7) yields

$$
\int_{\Omega} a(x)|\nabla y_1|^2 dx = -\lambda^{-1} \operatorname{Im} \int_{\Omega} [a(x)\nabla f_1 \cdot \nabla \overline{y_1} + (i\lambda f_1 + f_2)\overline{y_1}] dx.
$$
 (4.8)

Thus, by the Cauchy-Schwarz inequality, one can conclude from (4.8) that

$$
\int_{\Omega} a(x)|\nabla y_1|^2 dx \leq |\lambda|^{-1} \left(\int_{\Omega} a(x)|\nabla y_1|^2 dx \right)^{\frac{1}{2}} \|f_1\|_{H^1(\Omega)} + (\|f_1\|_{\Omega} + |\lambda|^{-1} \|f_2\|_{\Omega}) \|y_1\|_{\Omega}
$$

$$
\leq \frac{1}{4} \int_{\Omega} a(x)|\nabla y_1|^2 dx + |\lambda|^{-2} \|f_1\|_{H^1(\Omega)}^2 + (\|f_1\|_{\Omega} + |\lambda|^{-1} \|f_2\|_{\Omega}) \|y_1\|_{\Omega}.
$$

The proof of lemma is finished.

Proof of Theorem 4.1. Due to Lemma 4.1 (i) and (ii), we have

$$
||y_1||_{\{x \in \Omega : a(x) > 2\eta\}}^2 \lesssim ||\nabla y_1||_{\{x \in \Omega : a(x) > \eta\}}^2 + ||f_1||_{H^1(\Omega)}^2 + ||f_2||_{\Omega}^2,
$$
\n(4.9)

 \Box

and

$$
\|\nabla y_1\|_{\{\alpha \in \Omega : a(x) > \eta\}}^2 \lesssim \|f_1\|_{H^1(\Omega)}^2 + (\|f_1\|_{\Omega} + \|f_2\|_{\Omega}) \|y_1\|_{\Omega}.
$$
 (4.10)

Combining (4.9) and (4.10) yields

$$
||y_1||_{H^1(\{x \in \Omega : a(x) > 2\eta\})}^2 \lesssim ||f_1||_{H^1(\Omega)}^2 + ||f_2||_{L^2(\Omega)}^2 + (||f_1||_{\Omega} + ||f_2||_{\Omega})||y_1||_{\Omega}.
$$
 (4.11)

On the other hand, by (4.3) , one has that y_1 satisfies

$$
\operatorname{div} (\nabla y_1 + i\lambda a(x)\nabla y_1) + \lambda^2 y_1 = i\lambda f_1 + f_2 - \operatorname{div}(a(x)\nabla f_1). \tag{4.12}
$$

Then, applying Theorem 3.4 to y_1 satisfying (4.12), we obtain

$$
\tau \|e^{\tau \varphi} y_1\|_{\Omega}^2 + \tau^{-1} \|e^{\tau \varphi} \nabla y_1\|_{\Omega}^2
$$

\n
$$
\lesssim \lambda^2 \|e^{\tau \varphi} f_1\|_{\Omega}^2 + \|e^{\tau \varphi} f_2\|_{\Omega}^2 + \sum_{j=1}^d \|e^{\tau \varphi} a \partial_{x_j} f_1\|_{\Omega}^2
$$

\n
$$
+ (\lambda + \tau) \|e^{\tau \varphi} y_1\|_{\{x \in \Omega : a(x) \ge \delta/2\}}^2 + \tau^{-1} \|e^{\tau \varphi} \nabla y_1\|_{\{x \in \Omega : a(x) \ge \delta/2\}}^2.
$$

Let $c_1 = \min_{x \in \Omega} \varphi(x)$ and $c_2 = \max_{x \in \Omega} \varphi(x)$. We conclude from the above inequality and $\tau \geq$ $\max{\{\widetilde{K}|\lambda|^{\frac{5}{4}}}, 1\}$ that

$$
\tau e^{2c_1 \tau} \|y_1\|_{\Omega}^2 + \tau^{-1} e^{2c_1 \tau} \|\nabla y_1\|_{\Omega}^2
$$
\n
$$
\lesssim \lambda^2 e^{2c_2 \tau} \|f_1\|_{H^1(\Omega)}^2 + e^{2c_2 \tau} \|f_2\|_{\Omega}^2 + \tau e^{2c_2 \tau} \|y_1\|_{H^1(\{x \in \Omega : a(x) \ge \delta/2\})}^2.
$$
\n(4.13)

Setting $\eta = \frac{\delta}{8}$ and substituting (4.11) into (4.13), we obtain

$$
\tau e^{2c_1\tau} \|y_1\|_{\Omega}^2 + \tau^{-1} e^{2c_1\tau} \|\nabla y_1\|_{\Omega}^2 \lesssim \tau^2 e^{2c_2\tau} \|f_1\|_{H^1(\Omega)}^2 + \tau e^{2c_2\tau} \|f_2\|_{\Omega}^2 + \tau e^{2c_2\tau} (\|f_1\|_{\Omega} + \|f_2\|_{\Omega}) \|y_1\|_{\Omega}.
$$

Let $c_3 = 2(c_2 - c_1) + 1$. For $\tau \ge \max{\{\widetilde{K}|\lambda|^{\frac{5}{4}}}, 1\}$, one has

$$
||y_1||_{\Omega}^2 + ||\nabla y_1||_{\Omega}^2 \lesssim e^{c_3 \tau} ||f_1||_{H^1(\Omega)}^2 + e^{c_3 \tau} ||f_2||_{\Omega}^2 + e^{c_3 \tau} (||f_1||_{\Omega} + ||f_2||_{\Omega}) ||y_1||_{\Omega}.
$$

For any $\varepsilon > 0$, using $e^{c_3 \tau} (\|f_1\|_{\Omega} + \|f_2\|_{\Omega}) \|y_1\|_{\Omega} \leq \varepsilon \|y_1\|_{\Omega}^2 + \frac{1}{4} \varepsilon^{-1} e^{2c_3 \tau} (\|f_1\|_{\Omega} + \|f_2\|_{\Omega})^2$ in the above estimate, we conclude that

$$
||y_1||_{\Omega}^2 + ||\nabla y_1||_{\Omega}^2 \lesssim e^{2c_3 \tau} (||f_1||_{H^1(\Omega)}^2 + ||f_2||_{\Omega}^2),
$$

which gives the desired result. \Box

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