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# Logarithmic Decay of a Wave Equation with Kelvin-Voigt Damping \*

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#### Abstract

In this paper we analyze the long time behavior of a wave equation with local Kelvin-Voigt Damping. Through introducing proper class symbol and pseudo-differential calculus, we obtain a Carleman estimate, and then establish an estimate on the corresponding resolvent operator. As a result, we show the logarithmic decay rate for energy of the system without any geometric assumption on the subdomain on which the damping is effective.

keywords. Carleman estimate; wave equation; Kelvin-Voigt damping; logarithmic stability

AMS 2010 subject classification: Primary 93B05; Secondary 93B07, 35B37

#### 1 Introduction

In this paper, we consider a wave equation with local Kelvin-Voigt damping and analyze long time behaviour for the -solution of the system. Let  $\Omega \subset \mathbb{R}^d$  be a bounded domain with smooth boundary  $\Gamma = \partial \Omega$ . Denote by  $\partial_n$  the unit outward normal vector on boundary  $\Gamma$ . The PDE model is as follows.

$$\begin{cases} y_{tt}(t,x) - \operatorname{div}\left[\nabla y(t,x) + a(x)\nabla y_t(t,x)\right] = 0 & \text{in} \quad (0,\infty) \times \Omega, \\ y(t,x) = 0 & \text{on} \quad (0,\infty) \times \Gamma, \\ y(0,x) = y^0, \quad y_t(0,x) = y^1 & \text{in} \quad \Omega, \end{cases}$$

$$(1.1)$$

where the coefficient function  $a(\cdot) \in L^1(\Omega)$  is nonnegative and not identically null.

The natural energy of system (1.1) is

$$E(t) = \frac{1}{2} \left[ \int_{\Omega} (|\nabla y(t)|^2 + |y_t(t)|^2) dx \right].$$
 (1.2)

A direct computation gives that

$$\frac{d}{dt}E(t) = -\int_{\text{supp }a} a(x) |\nabla y_t(t)|^2 dx. \tag{1.3}$$

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Formula (1.3) shows that the only dissipative mechanism acting on the system is the viscoelastic damping div  $[a\nabla y_t]$ , which is only effective on supp a.

To rewrite the system as an evolution equation, we set the energy space as

$$\mathcal{H} = H_0^1(\Omega) \times L^2(\Omega), \tag{1.4}$$

with norm

$$||Y||_{\mathcal{H}} = \sqrt{||\nabla y_1||_{L^2(\Omega)}^2 + ||y_2||_{L^2(\Omega)}^2}, \quad \forall Y = (y_1, y_2) \in \mathcal{H}.$$
 (1.5)

Define an unbounded operator  $\mathcal{A}:D(\mathcal{A})\subset\mathcal{H}\to\mathcal{H}$  by

$$\mathcal{A}Y = (y_2, \operatorname{div}(\nabla y_1 + a\nabla y_2)), \quad \forall Y = (y_1, y_2) \in D(\mathcal{A}),$$

and

$$\mathcal{D}(\mathcal{A}) = \left\{ \left( y_1, y_2 \right) \in \mathcal{H} : y_2 \in H_0^1(\Omega), \operatorname{div} \left( \nabla y_1 + a \nabla y_2 \right) \right) \in L^2(\Omega) \right\}.$$

Let  $Y(t) = (y(t), y_t(t))$ . Then system (1.1) can be written as

$$\frac{d}{dt}Y(t) = AY(t), \quad \forall t > 0, \quad Y(0) = (y^0, y^1).$$
 (1.6)

It is known from [20] that if supp a is non-empty, the operator  $\mathcal{A}$  generates a contractive  $C_0$ semigroup  $e^{tA}$  on  $\mathcal{H}$  and  $i\mathbb{R} \subset \rho(\mathcal{A})$ , the resolvent of  $\mathcal{A}$ . Consequently, the semigroup  $e^{tA}$  is strongly stable. Moreover, if the entire medium is of the viscoelastic type (i.e. supp  $a = \overline{\Omega}$ ), the damping for the wave equation not only induces exponential energy decay, but also restricts the spectrum of the associated semigroup generator to a sector in the left half plane, and the associated semigroup is analytic ([13]). When the Kelvin-Voigt damping is localized on a subdomain of  $\Omega$ , the properties of system is quite complicated. First, it has been proved that properties of regularity and stability of 1-d system (1.1) depend on the continuousness of coefficient function  $a(\cdot)$ . More precisely, assume that  $\Omega = (-1, 1)$  and a(x) behaviours like  $x^{\alpha}$  with  $\alpha > 0$  in supp a = [0, 1]. Then the solution of system (1.1) is eventually differentiable for  $\alpha > 1$ , exponentially stable for  $\alpha \geq 1$ , polynomially stable of order  $\frac{1}{1-\alpha}$  for  $0 < \alpha < 1$ , and polynomially stable of optimal decay rate 2 for  $\alpha = 0$  (see [19, 22, 25, 27]). For the higher dimensional system, the corresponding semigroup is exponentially stable when  $a(\cdot) \in \mathscr{C}^2(\Omega)$  and supp  $a \supset \Gamma$  ([21]). However, when the Kelvin-Voigt damping is local and the material coefficient  $a(\cdot)$  is a positive constant on supp a, the energy of system (1.1) does not decay exponentially for any geometry of  $\Omega$  and supp a ([9, 28]). The reason is that the strong damping and non-continuousness of the coefficient function lead to reflection of waves at the interface  $\gamma \doteq \partial(\text{supp } a) \setminus \Gamma$ , which then fails to be effectively damped because they do not enter the region of damping. It turns out that the Kelvin-Voigt damping does not follow the principle that "geometric optics" condition implies exponential stability, which is true for the wave equation with local viscous damping ([2]).

Recently, [29] proves the polynomial stability of system (1.1) when  $a(\cdot) \equiv a_0 > 0$  on supp a and supp a satisfies certain geometry conditions. Then, a natural problem is: how about the decay rate if supp  $a \neq \emptyset$  is arbitrary? In [1], a is assumed discontinuous along a (d-1)-manifold, supp a is arbitrary and the rate of the decay of semi-group is estimated by  $(\log t)^{-k}$  for a data in  $D(\mathcal{A}^k)$ . In this paper, we analyze the logarithmic decay properties of the solution to (1.1) when a is smooth and supp a is arbitrary. The main result reads as follows.

**Theorem 1.1.** Suppose that the coefficient function  $a(\cdot) \in \mathscr{C}_0^{\infty}(\Omega)$  is nonnegative and supp  $a \subset \Omega$  is non-empty. Then the energy of the solution of (1.1) decays at logarithmic speed. More precisely,

one has that there exists a positive constant C such that

$$\|e^{t\mathcal{A}}Y_0\|_{\mathcal{H}} \le \frac{C}{[\log(t+2)]^{\frac{4k}{5}}} \|Y_0\|_{D(\mathcal{A}^k)}, \quad \forall \ t > 0, \ Y_0 = (y^0, y^1) \in D(\mathcal{A}^k).$$
 (1.7)

Our approach is based on the results duo to [7], which reduced the problem of determining the rate of energy decay to estimating the norm of the resolvent operator along the imaginary axis, see also [10, 18], etc. Our argument divides naturally three steps. First, in Section 2, we show some preliminaries including definitions and classical results about symbol, pseudo-differential calculus, and commutator estimate, etc. Then in Section 3, we prove corresponding Carleman estimates. Finally, in Section 4, we present a resolvent estimate and obtain Theorem 1.1. This theorem is a consequence of a resolvent estimate. The proof is given in Section 4. The method was developed in ([4, 5, 6, 16, 18, 26] and the references cited therein).

Throughout this paper, we use  $\|\cdot\|_V$  and  $(\cdot | \cdot)_V$  to denote the norm and inner product on  $L^2(V)$ , where  $V \subset \mathbb{R}^d$  if there is no further comments. When writing  $f \lesssim g$  (or  $f \gtrsim g$ ), we mean that there exists a positive constant C such that  $f \leq Cg$  (or  $f \geq Cg$ ). For  $j = 1, 2, \cdots$ , define operators  $D_j = -i\partial_{x_j}$ ,  $D = (D_1, \cdots, D_d)$ ,  $D^2 = \sum_{i=1}^d D_j^2$  and  $Da(x)D = \sum_{i=1}^d D_j a(x)D_j$ .

#### 2 Preliminaries

We shall prove Theorem 1.1 by Weyl-Hörmander calculus, which was introduced Hörmander ([12, 15]). In this section, some definitions and results on the class of symbol and pseudo-differential calculus are given.

#### 2.1 Symbol and Symbolic calculus

For any  $(x,\xi) \in \mathbb{R}^d \times \mathbb{R}^d$ ,  $\lambda \in \mathbb{R}$  and  $\tau > 0$ , we introduce the metric

$$g = g_{x,\xi} = \lambda dx^2 + \mu^{-2} d\xi^2$$
, where  $\mu^2 = \mu(\tau,\xi)^2 = \tau^2 + |\xi|^2$ , (2.1)

and the weight

$$\nu = \nu(x,\lambda) = \sqrt{1 + \lambda^2 a(x)^2}.$$
(2.2)

Note that and  $g_{x,\xi}(X,\Xi) = \lambda |X|^2 + \mu^{-2}(\tau,\xi) |\Xi|^2$  for all  $X,\Xi \in \mathbb{R}^d$ . Then we have the following results.

**Lemma 2.1.** Assume that there exist positive constants C and  $\lambda_0$  such that  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{C\lambda, 1\}$ . It holds

- (i) The metric  $g = g_{x,\xi}$  defined by (2.1) is admissible, i.e., it is slowly varying and temperate.
- (ii) The weight  $\nu = \nu(x, \lambda)$  defined by (2.2) is admissible, i.e., it is g-continuous and g-temperate.

*Proof.* (i) From Definition 18.4.1 in [12], the metric  $g_{x,\xi}$  defined by (2.1) is slowly varying if there exist  $\delta > 0$  and C > 0 such that

$$g_{x,\xi}(y-x,\eta-\xi) \le \delta \text{ implies } g_{y,\eta}(X,\Xi) \le Cg_{x,\xi}(X,\Xi), \ \ \forall \ x,y,\xi,\eta,X,\Xi \in \mathbb{R}^d,$$

where the constants  $\delta$  and C are independent on the parameters  $\lambda$  and  $\tau$ .

Suppose  $0 < \delta \le 1/4$  and

$$g_{x,\xi}(y-x,\eta-\xi) = \lambda |y-x|^2 + (\tau^2 + |\xi|^2)^{-1} |\eta-\xi|^2 \le \delta.$$

Then, we have

$$\tau^{2} + |\xi|^{2} \leq \tau^{2} + 2|\xi - \eta|^{2} + 2|\eta|^{2}$$
  
$$\leq \tau^{2} + 2\delta(\tau^{2} + |\xi|^{2}) + 2|\eta|^{2}.$$

This implies that  $\tau^2 + |\xi|^2 \le 4(\tau^2 + |\eta|^2)$ . Consequently,

$$g_{y,\eta}(X,\Xi) = \lambda |X|^2 + (\tau^2 + |\eta|^2)^{-1} |\Xi|^2 \le \lambda |X|^2 + \frac{1}{4} (\tau^2 + |\xi|^2)^{-1} |\Xi|^2 \le g_{x,\xi}(X,\Xi).$$

Therefore, g is slowly varying.

For a given metric  $g_{x,\xi}$ , the associated metric  $g_{x,\xi}^{\sigma}$  is defined by  $g_{x,\xi}^{\sigma} = (\tau^2 + |\xi|^2)dx^2 + \lambda^{-1}d\xi^2$ . The metric  $g_{x,\xi}$  is **temperate** if there exist C > 0 and N > 0, such that

$$g_{x,\xi}(X,\Xi) \le Cg_{y,\eta}(X,\Xi) \left(1 + g_{x,\xi}^{\sigma}(x - y, \xi - \eta)\right)^{N}, \quad \forall \ x, y, \xi, \eta, X, \Xi \in \mathbb{R}^{d}, \tag{2.3}$$

where the constants C and N are independent on the parameters  $\lambda$  and  $\tau$  (Definition 18.5.1 in [12]). For the metric  $g = g_{x,\xi}$  defined by (2.1), (2.3) is equivalent to

$$\lambda |X|^{2} + (\tau^{2} + |\xi|^{2})^{-1} |\Xi|^{2}$$

$$\leq C(\lambda |X|^{2} + (\tau^{2} + |\eta|^{2})^{-1} |\Xi|^{2}) (1 + (\tau^{2} + |\xi|^{2}) |x - y|^{2} + \lambda^{-1} |\xi - \eta|^{2})^{N}.$$
(2.4)

First, assume that  $\tau^2 + |\eta|^2 \le 4(\tau^2 + |\xi|^2)$ . It follows that

$$(\tau^2 + |\eta|^2) \le C(\tau^2 + |\xi|^2) (1 + \lambda^{-1} |\xi - \eta|^2)^N, \quad C > 0, \ N > 0.$$
 (2.5)

Then it is easy to obtain (2.4) from (2.5).

Secondly, consider the case  $\tau^2 + |\eta|^2 > 4(\tau^2 + |\xi|^2)$ . Then

$$|\eta| > 2|\xi|, \quad |\eta| > \sqrt{3}\tau,$$
 (2.6)

and

$$|\xi - \eta| > \frac{1}{2} |\eta| > \frac{\sqrt{3}}{2} \tau > \frac{\sqrt{3}}{2} \mathcal{C}\lambda. \tag{2.7}$$

It follows from (2.6) and (2.7) that

$$\lambda^{-1} |\xi - \eta|^2 > \frac{\sqrt{3}}{2} \mathcal{C} |\xi - \eta| > \frac{\sqrt{3}}{4} \mathcal{C} |\eta|.$$

Consequently,

$$(1 + \lambda^{-1} | \xi - \eta |^2)^2 > \frac{3}{16} C^2 | \eta |^2 > \frac{3}{32} C^2 (| \eta |^2 + 3\tau^2).$$

This together with  $\tau^2 + |\xi|^2 \ge 1$  yields that there exists a positive constant C such that (2.5) holds with N = 2.

(ii) It is known from Definition 18.4.2 in [12] that a weight  $\nu(x)$  is g-continuous if there exist  $\delta > 0$  and C > 0 such that

$$g_{x,\xi}(y-x,\eta-\xi) \le \delta$$
 implies  $C^{-1}\nu(x) \le \nu(y) \le C\nu(x), \ \forall \ x,y,\xi,\eta \in \mathbb{R}^d$ .

where the constants  $\delta$  and C are independent on the parameters  $\lambda$  and  $\tau$ . Since the weight  $\nu(x)$  defined by (2.2) does not depend on  $\xi$ , the above condition is reduced to

$$\lambda |x - y|^2 \le \delta$$
 implies  $C^{-1}\nu(x) \le \nu(y) \le C\nu(x), \ \forall x, y \in \mathbb{R}^d$ .

The weight  $\nu(x)$  is g-temperate if there exist C>0 and N>0 such that

$$\nu(y) \le C\nu(x) \left(1 + g_{y,\eta}^{\sigma}(x - y, \xi - \eta)\right)^{N}, \quad \forall \ x, y, \xi, \eta \in \mathbb{R}^{d}, \tag{2.8}$$

where the constants C and N do not depend on the parameters  $\lambda$  and  $\tau$  (Definition 18.5.1 in [12]). The weight  $\nu(x)$  is admissible if it is g-continuous and g-temperate. When a weight is admissible, all the powers of this weight are g-continuous and g-temperate. Therefore, it suffice to prove that  $1 + \lambda a(x)$  is admissible.

Let  $s \in [0,t]$  and  $t \in [0,1]$ . Define  $f(s) = \lambda a(x+s(y-x))$  and  $F(t) = \sup_{s \in [0,t]} f(s)$  where  $x, y \in \Omega$  satisfying  $\lambda |x-y|^2 \le \delta$ . It is clear that  $f'(s) = \lambda a'(x+s(y-x))(y-x)$ . Combining this with the following inequality

$$|a'(x)|^2 \le 2a(x)||a''||_{L^{\infty}(\Omega)}, \quad \forall \ x \in \Omega,$$
 (2.9)

we obtain

$$|f'(s)| \le \lambda |a'(x+s(y-x))| |y-x| \le 2\lambda ||a''||_{\infty}^{\frac{1}{2}} [a(x+s(y-x))]^{\frac{1}{2}} |y-x|.$$

The proof of (2.9) will be given later. Consequently,

$$\sup_{s \in [0,t]} |f'(s)| \le 2\lambda^{\frac{1}{2}} ||a''||_{L^{\infty}(\Omega)}^{\frac{1}{2}} F(t)^{\frac{1}{2}} |y - x|.$$

Since  $f(t) \leq f(0) + t \sup_{s \in [0,t]} |f'(s)|$ , F is non-decreasing and  $\lambda |x-y|^2 \leq \delta$ , we obtain that for all  $t \in [0,1]$ ,

$$f(t) \leq f(0) + C\lambda^{\frac{1}{2}}F(t)^{\frac{1}{2}} \left| \, y - x \, \right| \\ \leq f(0) + C\sqrt{\delta}F(t)^{\frac{1}{2}} \leq f(0) + C\sqrt{\delta}F(\alpha)^{\frac{1}{2}},$$

where  $C = 2||a''||_{L^{\infty}(\Omega)}^{\frac{1}{2}}$  and  $\alpha \in [t,1]$ . Note that f(0) = F(0). It follows that

$$F(\alpha) = \sup_{t \in [0,\alpha]} f(t) \le F(0) + C\sqrt{\delta}F(\alpha)^{\frac{1}{2}}.$$

This yields

$$1 + F(\alpha) \le 1 + F(0) + C\sqrt{\delta}(1 + F(\alpha))^{\frac{1}{2}} \le 1 + F(0) + C\sqrt{\delta}(1 + F(\alpha)). \tag{2.10}$$

By choosing  $\delta$  sufficiently small such that  $C\sqrt{\delta} \leq 1/2$ , one can deduce from (2.10) that

$$1 + F(\alpha) < 2(1 + F(0)), \forall \alpha \in [t, 1].$$

In particular, we have

$$1 + \lambda a(y) < 2(1 + \lambda a(x)).$$

The above inequality remains true if we exchange x and y. Therefore, the weight  $1 + \lambda a(x)$  is g-continuous.

On the other hand, note that  $1 + \lambda a(x)$  is independent to  $\xi$ . Then, to obtain the weight  $1 + \lambda a(x)$  is  $\sigma$ -temperate, it is sufficient to prove that

$$1 + \lambda a(y) \le C(1 + \lambda a(x))(1 + \tau^2 | x - y |^2)^N.$$
(2.11)

In fact, it is clear that  $1 + \lambda a(y) \le 1 + \lambda (a(x) + C | x - y |)$  where  $C = ||a'||_{L^{\infty}(\Omega)}$ . Therefore, there exists positive constant  $C' = CC^{-1}$  such that

$$1 + \lambda a(y) \le (1 + \lambda a(x))(1 + C'\tau \mid x - y \mid) \le (1 + \lambda a(x))(2 + 2(C'\tau \mid x - y \mid)^{2})^{\frac{1}{2}}.$$

Thus, we obtain (2.11) with  $N = \frac{1}{2}$ ,  $C = 2 \max\{1, C'\}$ .

**Remark 2.1.** We claim that (2.9) holds for any compactly supported and nonnegative function  $a \in \mathcal{C}^2(\Omega)$ . In fact, from the following identity

$$a(x+h) = a(x) + a'(x)h + \int_0^1 (1-t)a''(x+th)h^2dt, \ \forall \ h \in \mathbb{R},$$

one can get

$$a(x) + a'(x)h + \frac{1}{2}||a''||_{L^{\infty}(\Omega)}|h|^2 \ge 0.$$

Let h = y a'(x), where  $y \in \mathbb{R}$  and  $x \in \Omega$  are arbitrary. It follows from the above inequality that

$$a(x) + |a'(x)|^2 y + \frac{1}{2} ||a''||_{L^{\infty}(\Omega)} |a'(x)|^2 y^2 \ge 0, \quad \forall x \in \Omega, y \in \mathbb{R}.$$

Then,

$$|a'(x)|^4 - 2a(x)||a''||_{L^{\infty}(\Omega)} |a'(x)|^2 \le 0,$$

and (2.9) is proved.

**Definition 2.1.** (Section 18.4.2 in [12]) Assume the weight  $m(x,\xi)$  is admissible and the metric g is defined by (2.1). Let  $q(x,\xi,\lambda,\tau)$  be a  $\mathscr{C}^{\infty}$  function with respect to  $(x,\xi)$  and  $\lambda$ ,  $\tau$  be parameters satisfying conditions in Lemma 2.1. The symbol  $q(x,\xi,\lambda,\tau)$  is in class S(m,g) if for all  $\alpha,\beta\in\mathbb{N}^d$  there exists  $C_{\alpha,\beta}$  independent of  $\tau$  and  $\lambda$  such that

$$|\partial_x^{\alpha} \partial_{\xi}^{\beta} q(x,\xi,\lambda,\tau)| \leq C_{\alpha,\beta} m(x,\xi) \lambda^{|\alpha|/2} \mu(\tau,\xi)^{-|\beta|}.$$

**Remark 2.2.** (i) It is clear that  $\mu = \sqrt{\tau^2 + |\xi|^2} \in S(\mu, g)$  since  $|\partial_{\xi}^{\beta} \mu(\tau, \xi)| \lesssim \mu^{1-|\beta|}$  for all  $\beta \in \mathbb{R}^d$ .

(ii) Let  $\nu$  be the weight defined by (2.2). It is easy to get that  $\lambda a \in S(\nu, g)$ . In fact, if  $|\alpha| \geq 2$ , it holds that  $|\partial_x^{\alpha}(\lambda a(x))| \leq C_{\alpha}\lambda \leq C_{\alpha}\lambda^{|\alpha|/2}\nu(x)$ , where  $C_{\alpha} > 0$ . For the case  $|\alpha| = 1$ , it follows from (2.9) that

$$|\partial_x^{\alpha}(\lambda a(x))| \leq \sqrt{2} \|a''\|_{L^{\infty}(\Omega)}^{\frac{1}{2}} \lambda^{\frac{1}{2}} (\lambda a(x))^{\frac{1}{2}}$$

Note that  $|\lambda a(x)| < C\nu^2(x)$  for some C > 0. This together with the above inequality, we have that

$$|\partial_x^{\alpha}(\lambda a(x))| < \sqrt{2C} \|a''\|_{L^{\infty}(\Omega)}^{\frac{1}{2}} \lambda^{\frac{1}{2}} \nu(x).$$

(iii) It is known from Lemma 18.4.3 of [12] that if the metric g and weights  $m_1$ ,  $m_2$  are admissible, symbols  $a \in S(m_1, g)$  and  $b \in S(m_2, g)$ , then  $ab \in S(m_1m_2, g)$ . In particular,  $(\lambda a)^j \mu^k \in S(\nu^j \mu^k, g)$  for all  $j, k \in \mathbb{N} \cup \{0\}$ .

**Definition 2.2.** Let  $b \in S(m,g)$  be a symbol and  $u \in \mathcal{S}(\mathbb{R}^d)$ , we set

$$b(x, D, \tau)u(x) = \text{Op}(b)u(x) := (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ix\cdot\xi} b(x, \xi, \tau) \,\widehat{u}(\xi) \,d\xi.$$

It is known that  $Op(b): \mathscr{S}(\mathbb{R}^d) \to \mathscr{S}(\mathbb{R}^d)$  is continuous and Op(b) can be uniquely extended to  $\mathscr{S}'(\mathbb{R}^d)$  continuously. The following two lemmas are consequences of Theorem 18.5.4 and 18.5.10 in [12].

**Lemma 2.2.** Let  $b \in S(m,g)$  where m is an admissible weight and g is defined by (2.1). Then there exists  $c \in S(m,g)$  such that  $Op(b)^* = Op(c)$  and  $c(x,\xi) = \overline{b(x,\xi)} + r(x,\xi)$  where the remainder  $r \in S(\lambda^{\frac{1}{2}}\mu^{-1}m,g)$ .

**Lemma 2.3.** Let  $b \in S(m_1, g)$  and  $c \in S(m_2, g)$  where  $m_j$  are admissible weights for j = 1, 2 and g is defined by (2.1). Denote by  $[Op(b), Op(c)] = Op(b) \circ Op(c) - Op(c) \circ Op(b)$  and Poisson bracket  $\{b, c\}(x, \xi, \tau) = \sum_{1 \le j \le d} (\partial_{\xi_j} b \, \partial_{x_j} c - \partial_{x_j} b \, \partial_{\xi_j} c)(x, \xi, \tau)$ . Then,

- (i) there exists  $d \in S(m_1m_2, g)$  such that Op(b) Op(c) = Op(d) and  $d(x, \xi) = b(x, \xi)c(x, \xi) + r(x, \xi)$ where  $r \in S(\lambda^{\frac{1}{2}}\mu^{-1}m_1m_2, g)$ .
- (ii) for commutator i[Op(b), Op(c)] = Op(f), it holds that  $f \in S(\lambda^{\frac{1}{2}}\mu^{-1}m_1m_2, g)$  and  $f(x, \xi) = \{b, c\}(x, \xi) + r(x, \xi) \text{ where } r \in S(\lambda\mu^{-2}m_1m_2, g).$

The operators in  $S(\nu^j \mu^k, g)$  act on Sobolev spaces adapted to the class of symbol. Let  $b \in S(\nu^j \mu^k, g)$ , where  $\mu$  and g are defined by (2.1). Then there exists C > 0 such that

$$\|\operatorname{Op}(b)u\|_{\mathbb{R}^d} \le C\|\nu^j\operatorname{Op}(\mu^k)u\|_{\mathbb{R}^d}, \quad \forall \ u \in \mathscr{S}(\mathbb{R}^d).$$

By symbolic calculus, the above estimate is equivalent to  $\operatorname{Op}(\mu^{-k}\nu^{-j})\operatorname{Op}(b)$  acts on  $L^2(\mathbb{R}^d)$  since the operators associated with symbol in S(1,g) act on  $L^2(\mathbb{R}^d)$ . In particular, if  $b \in S(\nu^j \mu, g)$ , then for any  $\lambda \geq \lambda_0$ ,  $\tau \geq \max\{\mathcal{C}\lambda, 1\}$  and  $u \in \mathscr{S}(\mathbb{R}^d)$ , it holds

$$\|\operatorname{Op}(b)u\|_{\mathbb{R}^d} \le C\tau \|\nu^j u\|_{\mathbb{R}^d} + C\|\nu^j Du\|_{\mathbb{R}^d},$$

where C > 0 depends on positive constants  $\lambda_0$  and C.

#### 2.2 Commutator estimate

In this subsection, we suppose that  $\lambda = 1$  since the symbol does not depend on  $\lambda$ . The metric in (2.1) becomes

$$\tilde{g} = dx^2 + \mu^{-2}d\xi^2$$
, where  $\mu$  is defined by (2.1). (2.12)

To get the commutator estimate, we shall use the following Gårding inequality ([12, Theorem 18.6.7]).

**Lemma 2.4.** Let  $b \in S(\mu^{2k}, \tilde{g})$  be real valued.  $\mu$  and  $\tilde{g}$  are defined by (2.12). We assume there exists C > 0 such that  $b(x, \xi, \tau) \ge C\mu^{2k}$ . Then there exist  $\tilde{C} > 0$  and  $\tau_0 > 0$  such that

$$\operatorname{Re}(\operatorname{Op}(b)w \mid w)_{\mathbb{R}^d} \ge \widetilde{C} \|\operatorname{Op}(\mu^k)w\|_{\mathbb{R}^d}^2, \qquad \forall w \in \mathscr{S}(\mathbb{R}^d) \text{ and } \tau \ge \tau_0.$$
 (2.13)

**Definition 2.3.** Let V be a bounded open set in  $\mathbb{R}^d$ . We say that the weight function  $\varphi \in \mathscr{C}^{\infty}(\mathbb{R}^d; \mathbb{R})$  satisfies the sub-ellipticity condition in  $\overline{V}$  if  $|\nabla \varphi| > 0$  in  $\overline{V}$  and there exists constant C > 0,

$$\mathbb{P}_{\varphi}(x,\xi,\tau) = 0, \ \forall (x,\xi) \in \overline{V} \times \mathbb{R}^{d}, \ \tau > 0 \quad \Rightarrow \quad \{\mathbb{q}_{2},\mathbb{q}_{1}\}(x,\xi,\tau) \ge C(|\xi|^{2} + \tau^{2})^{3/2}, \tag{2.14}$$

where  $\mathbb{P}_{\varphi}(x,\xi,\tau) = |\xi + i\tau \nabla \varphi(x)|^2 = \mathbb{q}_2(x,\xi,\tau) + i\mathbb{q}_1(x,\xi,\tau)$  and  $\mathbb{q}_1$ ,  $\mathbb{q}_2$  are real valued.

**Lemma 2.5.** ([12]) Let V be a bounded open set in  $\mathbb{R}^d$  and  $\psi \in \mathscr{C}^{\infty}(\mathbb{R}^d; \mathbb{R})$  be such that  $|\nabla \psi| > 0$  in  $\overline{V}$ . Then, for  $\gamma > 0$  sufficiently large,  $\varphi = e^{\gamma \psi}$  fulfills the sub-ellipticity property in V.

**Lemma 2.6.** Assume that  $\varphi$  satisfies the sub-ellipticity in Definition 2.3. For all  $w \in \mathscr{C}_0^{\infty}(V)$ , there exist  $C_1, C_2 > 0$  and  $\tau_0 > 0$  such that the following inequality holds for all  $\tau \geq \tau_0$ ,

$$C_{1}\tau^{3}\|w\|_{V}^{2} + C_{1}\tau\|Dw\|_{V}^{2} \leq \operatorname{Re}\left(\operatorname{Op}(\{|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2}, 2\tau\xi \cdot \nabla\varphi(x)\})w \mid w\right)_{V} + C_{2}\tau^{-1}\|\operatorname{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})w\|_{V}^{2} + C_{2}\tau^{-1}\|\operatorname{Op}(2\tau\xi \cdot \nabla\varphi(x))w\|_{V}^{2}.$$

$$(2.15)$$

*Proof.* First, by homogeneity in  $(\xi, \tau)$ , compactness arguments and sub-ellipticity condition, we claim that there exist constants C,  $\delta > 0$  such that

$$C[|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2] + \{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\xi \cdot \nabla \varphi(x)\} \ge \delta \mu^2.$$
 (2.16)

The proof of (2.16) is classical. In fact, set

$$\mathcal{K} = \{ (x, \xi, \tau) \in \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R} : x \in \overline{V}, |\xi|^2 + \tau^2 = 1, \tau \ge 0 \},$$

and for  $(x, \xi, \tau) \in \mathcal{K}, \ \kappa > 0$ ,

$$G(x, \xi, \tau, \kappa) = \kappa \left[ |2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2 \right] + \{ |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, \ 2\xi \cdot \nabla \varphi(x) \}.$$

If  $|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2 = 0$  for  $(x, \xi, \tau) \in \mathcal{K}$ , it is clear that there exists a positive constant  $\delta$  such that (2.16) holds due to the fact that  $\phi$  is sub-elliptic. When  $|2\xi \cdot \nabla \varphi(x)|^2 + \mu^{-2} (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2 > 0$ , there exists a positive constant  $\kappa_{x,\xi,\tau}$  such that  $G(x,\xi,\tau,\kappa) > 0$  for every  $\kappa \geq \kappa_{x,\xi,\tau}$  since  $\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\xi \cdot \nabla \varphi(x)\}$  is bounded on  $\mathcal{K}$ . By continuity of  $G(x,\xi,\tau,\kappa)$ , there exists a neighborhood of  $(x,\xi,\tau)$ , denoted by  $V_{x,\xi,\tau}$ , such that  $G(x,\xi,\tau,\kappa) > 0$  for all  $(x,\xi,\tau) \in V_{x,\xi,\tau}$  and  $\kappa \geq \kappa_{x,\xi,\tau}$ . Since  $\mathcal{K}$  is compact, there exist finite sets  $V_j = V_{x_j,\xi_j,\tau_j}$  and corresponding constants  $\kappa_j = \kappa_{x_j,\xi_j,\tau_j}$   $(j=1,2,\cdots,n)$ , such that  $\mathcal{K} \subset \cup_{j=1}^n V_j$  and  $G(x,\xi,\tau,\kappa) > 0$  for all  $(x,\xi,\tau) \in V_j$  and  $\kappa > \kappa_j$ . Let  $\tilde{\kappa} = \max\{\kappa_j : j=1,2,\cdots,n\}$ . It follows that  $G(x,\xi,\tau,\tilde{\kappa}) > 0$  for all  $(x,\xi,\tau) \in \mathcal{K}$  and  $\kappa \geq \tilde{\kappa}$ . Finally, using the compactness of  $\mathcal{K}$  again, we conclude that there exists  $\delta > 0$  such that  $G(x,\xi,\tau,\tilde{\kappa}) \geq \delta$ . Thus, (2.16) is reached since g is a homogeneous function of degree 2 with respect to variables  $(\xi,\tau)$ .

By Gårding inequality (2.13), there exists a constant C>0 such that, for  $\tau \geq \tau_0$  with  $\tau_0$  sufficiently large,

$$C\|\operatorname{Op}(\mu)w\|_{V}^{2} \leq \operatorname{Re}\left(\operatorname{Op}\left(|2\xi \cdot \nabla \varphi(x)|^{2} + \mu^{-2}\left(|\xi|^{2} - \tau^{2}|\nabla \varphi(x)|^{2}\right)^{2} + \{|\xi|^{2} - \tau^{2}|\nabla \varphi(x)|^{2}, 2\xi \cdot \nabla \varphi(x)\}\right)w \mid w\right)_{V}.$$
(2.17)

Now we are going to estimate the terms  $\tau$  Op( $|2\xi \cdot \nabla \varphi(x)|^2$ ) and  $\tau$  Op( $\mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2$ ). Firstly, it follows from Lemma 2.3 that

$$\tau^{-1}\operatorname{Op}(|2\tau\xi\cdot\nabla\varphi(x)|^2) = \tau^{-1}\operatorname{Op}(2\tau\xi\cdot\nabla\varphi(x))^*\operatorname{Op}(2\tau\xi\cdot\nabla\varphi(x)) + \tau\operatorname{Op}(r_1), \tag{2.18}$$

where  $r_1 \in S(\mu, \tilde{g})$  and  $\tilde{g}$  is defined by (2.12). Therefore, for any  $\varepsilon > 0$ , there exists a positive constant  $C_{\varepsilon}$  such that

$$\left| \left( \tau^{-1} \operatorname{Op}(|2\tau\xi \cdot \nabla\varphi(x)|^{2}) w \mid w \right)_{V} \right|$$

$$\leq \tau^{-1} \| \operatorname{Op}(2\tau\xi \cdot \nabla\varphi(x)) w \|_{V}^{2} + \tau | \left( \operatorname{Op}(r_{1}) w \mid w \right)_{V} |$$

$$\leq \tau^{-1} \| \operatorname{Op}(2\tau\xi \cdot \nabla\varphi(x)) w \|_{V}^{2} + \varepsilon \tau \| \operatorname{Op}(\mu) w \|_{V}^{2} + C_{\varepsilon}\tau \| w \|_{V}^{2}.$$

$$(2.19)$$

Substituting (2.19) into (2.17) and choosing  $\varepsilon$  small enough, we have

$$C\tau \|\operatorname{Op}(\mu)w\|_V^2$$

$$\leq \operatorname{Re} \left( \operatorname{Op} \left( \tau \mu^{-2} \left( |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2 \right)^2 + \tau \left\{ |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, \ 2\xi \cdot \nabla \varphi(x) \right\} \right) w \mid w \right)_V$$

$$+ \tau^{-1} \| \operatorname{Op} (2\tau \xi \cdot \nabla \varphi(x)) w \|_V^2 + C_{\varepsilon} \tau \|w\|_V^2.$$

$$(2.20)$$

Secondly, by symbolic calculus, we have that

$$\tau \operatorname{Op}(\mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)^2) = \tau \operatorname{Op}(r_0) \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) + \tau \operatorname{Op}(r_2), \tag{2.21}$$

where  $r_0(x,\xi) = \mu^{-2}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) \in S(1,\tilde{g})$  and  $r_2 \in S(\mu,\tilde{g})$ . Therefore, for all  $\varepsilon > 0$ , there exists  $C_{\varepsilon} > 0$  such that

$$| (\tau \operatorname{Op}(r_0) \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) w | w)_V |$$

$$\leq C_{\varepsilon} \tau^{-1} || \operatorname{Op}(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2) w ||_V^2 + \varepsilon \tau^3 ||w||_V^2.$$
(2.22)

We choose  $\varepsilon$  small enough and combine (2.21)-(2.22) with (2.20) to get

$$C\tau \| \operatorname{Op}(\mu) w \|_{V}^{2} \leq \operatorname{Re} \left( \tau \operatorname{Op} \left\{ |\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}, \ 2\xi \cdot \nabla \varphi(x) \right\} w \mid w \right)_{V}$$

$$+ C_{\varepsilon} \tau^{-1} \| \operatorname{Op}(|\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}) w \|_{V}^{2}$$

$$+ \tau^{-1} \| \operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x)) w \|_{V}^{2} + C_{\varepsilon} (\tau + \varepsilon \tau^{3}) \| w \|_{V}^{2}.$$

$$(2.23)$$

Finally, it is clear that there exist positive constant C such that

$$\tau^{3} \|w\|_{V}^{2} + \tau \|Dw\|_{V}^{2} \le C\tau \|\operatorname{Op}(\mu)w\|_{V}^{2}. \tag{2.24}$$

Thus, we obtain (2.15) by using (2.23)-(2.24), choosing  $\varepsilon$  small enough and letting  $\tau > \tau_0$  large enough.

### 3 Carleman Estimate

In this section, we shall prove several Carleman inequalities. Define the operator

$$P(x, D, \lambda) = D^2 + i\lambda Da(x)D - \lambda^2$$
.

Let the weight function  $\varphi \in \mathscr{C}^{\infty}(\mathbb{R}^d; \mathbb{R})$ . The associated conjugate operator of  $P(x, D, \lambda)$  is  $P_{\varphi}(x, D, \lambda) = e^{\tau \varphi} P(x, D, \lambda) e^{-\tau \varphi}$ . Then,

$$P_{\varphi} = (D + i\tau \nabla \varphi(x))^{2} + i\lambda(D + i\tau \nabla \varphi(x))a(x)(D + i\tau \nabla \varphi(x)) - \lambda^{2}.$$

By setting  $Q_2 = \frac{1}{2}(P_{\varphi} + P_{\varphi}^*)$  and  $Q_1 = \frac{1}{2i}(P_{\varphi} - P_{\varphi}^*)$ , we have  $P_{\varphi} = Q_2 + iQ_1$ . We denote by  $p(x, \xi, \lambda)$ ,  $p_{\varphi}(x, \xi, \lambda)$  the associated symbol of  $P(x, D, \lambda)$ ,  $P_{\varphi}(x, D, \lambda)$ , respectively.

Let the metric g and weight  $\nu$  be defined by (2.1) and (2.2). Due to the results Remark 2.2, we know that  $D+i\tau\nabla\varphi(x)$  is an operator with symbol in  $S(\mu,g)$  class,  $\lambda a$  is in  $S(\nu,g)$ , and  $(1+i\lambda a(x))|\xi+i\tau\nabla\varphi(x)|^2$ , the principal symbol of  $P_{\varphi}$  belongs to  $S(\nu\mu^2,g)$ . Furthermore, from Lemma 2.2 and 2.3, one has that

$$P_{\varphi} = \operatorname{Op}\left((1+i\lambda a(x))|\xi+i\tau\nabla\varphi(x)|^{2}\right) - \lambda^{2} + R_{3},$$

$$Q_{2} = \operatorname{Op}(q_{2}) - \lambda^{2} + R_{2},$$

$$Q_{1} = \operatorname{Op}(q_{1}) + R_{1},$$
(3.1)

where  $q_2 = |\xi|^2 - \tau^2 |\nabla \varphi(x)|^2 - 2\lambda \tau a(x) \xi \cdot \nabla \varphi(x)$ ,  $q_1 = 2\tau \xi \cdot \nabla \varphi(x) + \lambda a(x) (|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2)$  belong to  $S(\nu \mu^2, g)$  and the symbols of  $R_j$  is in  $S(\lambda^{\frac{1}{2}} \nu \mu, g)$  for j = 1, 2, 3. It is clear that

$$||P_{\varphi}v||_{V}^{2} = ||Q_{2}v||_{V}^{2} + ||Q_{1}v||_{V}^{2} + 2\operatorname{Re}(Q_{2}v | iQ_{1}v)_{V}, \quad \forall v \in \mathscr{C}^{\infty}(V).$$
(3.2)

In what follows, several Carleman estimates are introduced. First, we give an estimation on the subdomain which is far away from the boundary  $\Gamma$ .

**Theorem 3.1.** Suppose  $\varphi$  satisfies sub-ellipticity condition in  $V \subset \Omega$ . Then, there exist positive constants C,  $\tilde{K}$  and  $\lambda_0$ , such that for every  $v \in \mathscr{C}_0^{\infty}(V)$ , it holds

$$\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} + \|Q_{1}v\|_{V}^{2} + \|Q_{2}v\|_{V}^{2} \lesssim \|P_{\varphi}v\|_{V}^{2}, \tag{3.3}$$

where  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{\tilde{K}|\lambda|^{\frac{5}{4}}, 1\}$ .

*Proof.* Let  $w = \nu(x)v$  in (2.15). We obtain

$$C_{1}\tau^{3}\|\nu(x)v\|_{V}^{2} + C_{1}\tau\|D(\nu(x)v)\|_{V}^{2}$$

$$\leq \operatorname{Re}\left(\operatorname{Op}(\left\{|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2}, 2\tau\xi \cdot \nabla\varphi(x)\right\})\nu(x)v|\nu(x)v\right)_{V}$$

$$+ C_{2}\tau^{-1}\|\operatorname{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})\nu(x)v\|_{V}^{2}$$

$$+ C_{2}\tau^{-1}\|\operatorname{Op}(2\tau\xi \cdot \nabla\varphi(x))\nu(x)v\|_{V}^{2}, C_{1}, C_{2} > 0.$$

$$(3.4)$$

Since the symbol of  $[D, \nu(x)]$  is in  $S(\lambda^{\frac{1}{2}}\nu, g)$ , we have

$$\|\nu(x)Dv\|_V^2 \le \|D(\nu(x)v)\|_V^2 + \|[D,\nu(x)]v\|_V^2 \le \|D(\nu(x)v)\|_V^2 + C\lambda\|\nu(x)v\|_V^2.$$

Consequently, for  $\tau \geq C\lambda$  with C > 0 sufficiently large, it holds

$$\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} \lesssim \tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|D(\nu(x)v)\|_{V}^{2}. \tag{3.5}$$

It follows from (3.4) and (3.5) that

$$C'_{1}\left(\tau^{3}\|\nu(x)v\|_{V}^{2} + \tau\|\nu(x)Dv\|_{V}^{2}\right)$$

$$\leq \operatorname{Re}\left(\operatorname{Op}(\left\{|\xi|^{2} - \tau^{2}\varphi(x)^{2}, 2\tau\xi \cdot \nabla\varphi(x)\right\})\nu(x)v \mid \nu(x)v\right)_{V}$$

$$+C_{2}\tau^{-1}\|\operatorname{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})\nu(x)v\|_{V}^{2}$$

$$+C_{2}\tau^{-1}\|\operatorname{Op}(2\tau\xi \cdot \nabla\varphi(x))\nu(x)v\|_{V}^{2}, \quad C'_{1}, C_{2} > 0.$$

$$(3.6)$$

Now, we estimate the first term on the right side hand of (3.6). Note that

$$2\operatorname{Re}(Q_{2}v \mid iQ_{1}v)_{V} = (Q_{2}v \mid iQ_{1}v)_{V} + (iQ_{1}v \mid Q_{2}v)_{V} = (i[Q_{2}, Q_{1}]v \mid v)_{V}, \tag{3.7}$$

where the principal symbol of  $i[Q_2, Q_1]$  is  $\{q_2, q_1\}$ . Due to Lemma 2.3, we obtain that

$$i[Q_2, Q_1] = \text{Op}(\{q_2, q_1\}) + R_4,$$
 (3.8)

where  $\{q_2, q_1\} \in S(\lambda^{\frac{1}{2}}\nu^2\mu^3, g)$  and the symbol of  $R_4$  is in  $S(\lambda\nu^2\mu^2, g)$ . A direct computation gives that

$$\begin{split} \left\{q_2,q_1\right\} &= (1+a^2(x)\lambda^2) \big\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, 2\tau\xi\cdot\nabla\varphi(x)\big\} \\ &+ \big(\big\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, \lambda a(x)\big\} - \lambda a(x) \big\{2\tau\xi.\nabla\varphi(x), \lambda a(x)\big\}\big) (|\xi|^2 - \tau^2|\nabla\varphi(x)|^2) \\ &- 2\tau\xi\cdot\nabla\varphi(x) \big(\lambda a(x) \big\{\lambda a(x), |\xi|^2 - \tau^2|\nabla\varphi(x)|^2\big\} + \big\{\lambda a(x), 2\tau\xi\cdot\nabla\varphi(x)\big\}\big). \end{split}$$

From the definition of  $q_1$  and  $q_2$ , we have

$$q_2(x,\xi) + \lambda a(x)q_1(x,\xi) = (1 + \lambda^2 a(x)^2)(|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2),$$
  

$$q_1(x,\xi) - \lambda a(x)q_2(x,\xi) = 2\tau \xi \cdot \nabla \varphi(x)(1 + \lambda^2 a^2(x)).$$
(3.9)

Consequently,

$$\begin{aligned}
& \left\{q_{2}, q_{1}\right\} \\
&= \nu^{2}(x) \left\{|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2}, 2\tau\xi \cdot \nabla\varphi(x)\right\} \\
&+ \nu^{-2}(x) \left(q_{2}(x, \xi) + \lambda a(x)q_{1}(x, \xi)\right) \left(\left\{|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2}, \lambda a(x)\right\} - \lambda a(x) \left\{2\tau\xi \cdot \nabla\varphi(x), \lambda a(x)\right\}\right) \\
&- \nu^{-2}(x) \left(q_{1}(x, \xi) - \lambda a(x)q_{2}(x, \xi)\right) \left(\lambda a(x) \left\{\lambda a(x), |\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2}\right\} + \left\{\lambda a(x), 2\tau\xi \cdot \nabla\varphi(x)\right\}\right).
\end{aligned}$$

Then, it follows from (3.8) and the above equation that

$$i[Q_2, Q_1] = \nu(x) \operatorname{Op} (\{|\xi|^2 - \tau^2 |\nabla \varphi(x)|^2, 2\tau \xi \cdot \nabla \varphi(x)\}) \nu(x) + B_1 \nu^{-1}(x) \operatorname{Op} (q_2(x, \xi) + \lambda a(x) q_1(x, \xi)) - B_2 \nu^{-1}(x) \operatorname{Op} (q_1(x, \xi) - \lambda a(x) q_2(x, \xi)) + R_5,$$

where

$$B_1 = \operatorname{Op}\left(\nu^{-1}(x)\left(\left\{|\xi|^2 - \tau^2|\nabla\varphi(x)|^2, \lambda a(x)\right\} - \lambda a(x)\left\{2\tau\xi \cdot \nabla\varphi(x), \lambda a(x)\right\}\right)\right),$$
  

$$B_2 = \operatorname{Op}\left(\nu^{-1}(x)\left(\lambda a(x)\left\{\lambda a(x), |\xi|^2 - \tau^2|\nabla\varphi(x)|^2\right\} + \left\{\lambda a(x), 2\tau\xi \cdot \nabla\varphi(x)\right\}\right)\right),$$

symbols of  $B_1$  and  $B_2$  belong to  $S(\lambda^{\frac{1}{2}}\nu\mu, g)$  and the symbol of  $R_5$  is in  $S(\lambda\nu^2\mu^2, g)$ . Combining this with (3.1) yields

$$i[Q_{2}, Q_{1}] = \nu(x) \operatorname{Op} \left( \left\{ |\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}, 2\tau \xi \cdot \nabla \varphi(x) \right\} \right) \nu(x)$$

$$+ B_{1} \nu^{-1}(x) \left[ Q_{2}(x, \xi) + \lambda a(x) Q_{1}(x, \xi) + \lambda^{2} \right]$$

$$- B_{2} \nu^{-1}(x) \left[ Q_{1}(x, \xi) - \lambda a(x) Q_{2}(x, \xi) - \lambda^{3} a(x) \right] + R_{5},$$
(3.10)

We refer to Section 2 where the rules on symbolic calculus are given and precise. Therefore, by the continuity of pseudo-differential operator, we have that for j = 1, 2, k = 0, 1 and  $\ell = 1, 2,$ 

$$|(B_{j}\nu(x)^{-1}(\lambda a(x))^{k}Q_{\ell}v|v)_{V}|$$

$$= |(\nu(x)^{-1}(\lambda a(x))^{k}Q_{\ell}v|B_{j}^{*}v)_{V}| \lesssim ||Q_{\ell}v||_{V}||B_{j}^{*}v||_{V}$$

$$\leq \frac{1}{10}||Q_{\ell}v||_{V}^{2} + C\lambda\tau^{2}||\nu(x)v||_{V}^{2} + C\lambda||\nu(x)Dv||_{V}^{2}, \quad C > 0,$$
(3.11)

$$| (B_{j}\nu^{-1}(x)\lambda^{2}(\lambda a(x))^{k}v | v)_{V} | \lesssim \lambda^{\frac{5}{2}} (\tau \|\nu(x)v\|_{V} + \|\nu(x)Dv\|_{V}) \|\nu(x)v\|_{V}$$

$$\lesssim \lambda^{\frac{5}{2}}\tau \|\nu(x)v\|_{V}^{2} + \lambda^{\frac{5}{2}}\tau^{-1} \|\nu(x)Dv\|_{V}^{2},$$

$$(3.12)$$

and

$$|(R_5 v | v)_V| \lesssim \lambda \tau^2 ||\nu(x)v||_V^2 + \lambda ||\nu(x)Dv||_V^2.$$
 (3.13)

Due to (3.10)–(3.13), there exists a positive constant C such that

$$\left( \operatorname{Op}(\left\{ |\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}, 2\tau \xi \cdot \nabla \varphi(x) \right\}) \nu(x) v |\nu(x) v \rangle_{V} \right) 
\leq \left( i[Q_{2}, Q_{1}] v |v\rangle_{V} + \frac{1}{4} \sum_{\ell=1,2} \|Q_{\ell} v\|_{V}^{2} + C \left( \lambda \tau^{2} \|\nu(x) v\|_{V}^{2} \right) 
+ \lambda \|\nu(x) D v\|_{V}^{2} + \lambda^{\frac{5}{2}} \tau \|\nu(x) v\|_{V}^{2} + \lambda^{\frac{5}{2}} \tau^{-1} \|\nu(x) D v\|_{V}^{2} \right).$$
(3.14)

Next, we are going to estimate the last two terms on the right side hand of (3.6). It follows from (3.1) and (3.9) that

$$\|\operatorname{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})\nu(x)v\|_{V}^{2} + \|\operatorname{Op}(2\tau\xi\cdot\nabla\varphi(x))\nu(x)v\|_{V}^{2}$$

$$= \|\operatorname{Op}(\nu(x)^{-2}(q_{2} + \lambda a(x)q_{1})\nu(x)v\|_{V}^{2} + \|\operatorname{Op}(\nu(x)^{-2}(q_{1} - \lambda a(x)q_{2})\nu(x)v\|_{V}^{2}$$

$$= \|\nu(x)^{-1}[Q_{2} + \lambda^{2} - R_{2} + \lambda a(x)(Q_{1} - R_{1})]v\|_{V}^{2}$$

$$+ \|\nu(x)^{-1}[Q_{1} - R_{1} - \lambda a(x)(Q_{2} + \lambda^{2} - R_{2})]v\|_{V}^{2}.$$
(3.15)

Combining this with the fact that the symbols of  $R_j$  are in  $S(\lambda^{\frac{1}{2}}\nu\mu, g)$  for j=1,2, we have

$$\|\operatorname{Op}(|\xi|^{2} - \tau^{2}|\nabla\varphi(x)|^{2})\nu(x)v\|_{V}^{2} + \|\operatorname{Op}(2\tau\xi\cdot\nabla\varphi(x))\nu(x)v\|_{V}^{2}$$

$$\lesssim \sum_{\ell=1,2} (\|Q_{\ell}v\|_{V}^{2} + \|R_{\ell}v\|_{V}^{2}) + \lambda^{4}\|v\|_{V}^{2}$$

$$\lesssim \sum_{\ell=1,2} \|Q_{\ell}v\|_{V}^{2} + \lambda\tau^{2}\|\nu(x)v\|_{V}^{2} + \lambda\|\nu(x)Dv\|_{V}^{2} + \lambda^{4}\|v\|_{V}^{2}.$$
(3.16)

Consequently, for  $\tau \geq C\lambda$  with C large enough, it holds

$$\tau^{-1} \|\operatorname{Op}(|\xi|^{2} - \tau^{2} |\nabla \varphi(x)|^{2}) \nu(x) v\|_{V}^{2} + \tau^{-1} \|\operatorname{Op}(2\tau \xi \cdot \nabla \varphi(x)) \nu(x) v\|_{V}^{2}$$

$$\lesssim \tau^{-1} \sum_{\ell=1,2} \|Q_{\ell} v\|_{V}^{2} + \lambda \tau \|\nu(x) v\|_{V}^{2} + \|\nu(x) D v\|_{V}^{2} + \lambda^{3} \|v\|_{V}^{2}.$$
(3.17)

Finally, by (3.6), (3.14) and (3.17), one can choose  $\tau \geq \max\{\tilde{K}|\lambda|^{\frac{5}{4}}, 1\}$  with  $\tilde{K}$  sufficiently large such that for some C > 0,

$$C(\tau^{3} \| \nu(x)v \|_{V}^{2} + \tau \| \nu(x)Dv \|_{V}^{2}) \leq (i[Q_{2}, Q_{1}]v \| v)_{V} + \frac{1}{2} \sum_{\ell=1,2} \|Q_{\ell}v\|_{V}^{2} + \varepsilon \tau^{3} \|\nu(x)v\|_{V}^{2} + \varepsilon \tau \|\nu(x)Dv\|_{V}^{2},$$

$$(3.18)$$

where  $\varepsilon > 0$  is arbitrary. Choosing  $\varepsilon$  small with respect to C, using (3.2), (3.7) and (3.18), we obtain

$$C_1(\tau^3 \|\nu(x)v\|_V^2 + \tau \|\nu(x)Dv\|_V^2) + \frac{1}{2}(\|Q_1v\|_V^2 + \|Q_2v\|_V^2)$$

$$\leq (i[Q_2, Q_1]v, v)_V + \|Q_1v\|_V^2 + \|Q_2v\|_V^2 = \|P_{\varphi}v\|_V^2,$$

which implies (3.3).

Remark 3.1. The estimates in (3.12) impose the assumption  $\tau \geq \tilde{K}|\lambda|^{\frac{5}{4}}$ . The other remainder terms only impose the condition  $\tau \geq C\lambda$ . This condition is related with the principal normal condition. Indeed for a complex operator, with symbol  $p_1 + ip_2$  where  $p_1, p_2$  are both real valued, the Carleman estimate is only true if  $\{p_1, p_2\} = 0$  on  $p_1 = p_2 = 0$ . Here the symbol of operator before conjugation by weight is  $|\xi|^2 - \lambda^2 + i\lambda a(x)|\xi|^2$ , and the Poisson bracket is  $\{|\xi|^2 - \lambda^2, \lambda a(x)|\xi|^2\} = 2\lambda(\xi \cdot \nabla a(x))|\xi|^2$ . We can estimate this term, uniformly in a neighborhood of a(x) = 0, by  $C\lambda a^{\frac{1}{2}}(x)|\xi|^3$ . This explanation does not justify the power  $|\lambda|^{\frac{5}{4}}$  found at the end of computations but shows the difficulties.

**Theorem 3.2.** Suppose  $\varphi$  satisfies sub-ellipticity condition in  $V \subset \Omega$ . Then, there exist positive constants  $\tilde{K}$  and  $\lambda_0$ , such that for every  $u \in \mathscr{C}_0^{\infty}(V)$ , it holds

$$\tau^{3} \| (1 + \lambda^{2} a(x)^{2})^{\frac{1}{2}} e^{\tau \varphi} u \|_{V}^{2} + \tau \| (1 + \lambda^{2} a(x)^{2})^{\frac{1}{2}} e^{\tau \varphi} D u \|_{V}^{2} \lesssim \| e^{\tau \varphi} P u \|_{V}^{2}, \tag{3.19}$$

where  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{\tilde{K}|\lambda|^{\frac{5}{4}}, 1\}$ .

*Proof.* Set  $v = e^{\tau \varphi} u$ . From Theorem 3.1, it suffices to prove that (3.19) is equivalent to

$$\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} \lesssim \|P_{\varphi}v\|_{V}^{2}. \tag{3.20}$$

First, assume (3.20) holds. Then,  $Dv = e^{\tau \varphi}(Du - i\tau \nabla \varphi u)$  and  $e^{\tau \varphi}Du = Dv + i\tau \nabla \varphi v$ . Then there exist positive constants  $c_1$ ,  $c_2$  such that

$$c_{1} (\| \nu(x)Dv\|_{V} + \tau \|\nu(x)v\|_{V}) \leq \|\nu(x)e^{\tau\varphi}Du\|_{V} + \tau \|\nu(x)e^{\tau\varphi}u\|_{V}$$

$$\leq c_{2} (\|\nu(x)Dv\|_{V} + \tau \|\nu(x)v\|_{V}).$$
(3.21)

Combining this with (3.20), we conclude that

$$\tau^3 \|\nu(x) e^{\tau \varphi} u\|_V^2 + \tau \|\nu(x) e^{\tau \varphi} D u\|_V^2 \lesssim (\tau \|\nu(x) D v\|_V^2 + \tau^3 \|\nu(x) v\|_V^2) \lesssim \|e^{\tau \varphi} P u\|_V^2.$$

On the other hand, (3.19) implies that

$$\tau^{3} \| \nu(x) e^{\tau \varphi} u \|_{V}^{2} + \tau \| \nu(x) e^{\tau \varphi} D u \|_{V}^{2} \lesssim \| P_{\varphi} v \|_{V}^{2}.$$

Then, we proved (3.20) from the above estimate and (3.21).

Since there is higher order term  $\operatorname{div}(a(x)\nabla y_t)$  in system (1.1), it is necessary to deal with the term  $\operatorname{div}(a(x)\nabla f)$  for  $f\in H^1(\Omega)$  when proving the resolvent estimate. The following result is analogue to the work by Imanuvilov and Puel ([14]).

**Theorem 3.3.** Suppose  $\varphi$  satisfies sub-ellipticity condition on  $V \subset \Omega$ . Then, there exist  $C, \tilde{K}, \lambda_0 > 0$ , such that for all  $u \in \mathscr{C}_0^{\infty}(V)$  satisfying

$$P(x, D, \lambda)u = g_0 + \sum_{j=1}^{d} \partial_{x_j} g_j, \text{ where } g_j \in L^2(V), \ j = 0, 1, \dots, d,$$
 (3.22)

it holds

$$\tau \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} u \|_V^2 + \tau^{-1} \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} D u \|_V^2 \le C \sum_{j=0}^d \| e^{\tau \varphi} g_j \|_V^2, \tag{3.23}$$

where  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{\tilde{K}|\lambda|^{\frac{5}{4}}, 1\}$ .

*Proof.* First, from (3.1), we have  $D^2 = Q_2 + S_2$  and  $\lambda a(x)D^2 = Q_1 + S_1$  where  $S_1$  and  $S_2$  have symbols in  $S(\tau\nu\mu, g)$  if  $\tau \gtrsim \lambda$ . It follows that for any  $v \in \mathscr{C}_0^{\infty}(V)$ ,

$$||D^{2}v||_{V}^{2} \lesssim ||Q_{2}v||_{V}^{2} + \tau^{2}(\tau^{2}||\nu(x)v||_{V}^{2} + ||\nu(x)Dv||_{V}^{2}),$$

$$||\lambda a D^{2}v||_{V}^{2} \lesssim ||Q_{1}v||_{V}^{2} + \tau^{2}(\tau^{2}||\nu(x)v||_{V}^{2} + ||\nu(x)Dv||_{V}^{2}).$$
(3.24)

Using (3.3), (3.24) and the fact that  $\|\nu(x)D^2v\|_V^2 \le 2(\|D^2v\|_V^2 + \|\lambda a(x)D^2v\|_V^2)$ , we obtain

$$\tau^{3} \|\nu(x)v\|_{V}^{2} + \tau \|\nu(x)Dv\|_{V}^{2} + \tau^{-1} \|\nu(x)D^{2}v\|_{V}^{2} \lesssim \|P_{\varphi}v\|_{V}^{2}. \tag{3.25}$$

Let  $\widetilde{u}$  and  $\chi$  be in  $\mathscr{C}_0^{\infty}(\Omega)$  such that  $\chi=1$  on a neighborhood of supp  $\widetilde{u}$ . Similarly to (3.21), we obtain

$$\tau \| \nu(x) \widetilde{u} \|_V^2 + \tau^{-1} \| \nu(x) D \widetilde{u} \|_V^2 \lesssim \tau \| \nu(x) \widetilde{u} \|_V^2 + \tau^{-1} \| D(\nu(x) \widetilde{u}) \|_V^2.$$

Then, combining this with Fourier transform and the following inequality

$$\tau + \frac{|\xi|^2}{\tau} \lesssim \frac{\tau^3}{\tau^2 + |\xi|^2} + \frac{|\xi|^4}{\tau(\tau^2 + |\xi|^2)},$$

we conclude that

$$\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\operatorname{Op}(\mu^{-1})\nu(x)\chi\widetilde{u}\|_{V}^{2} + \tau^{-1} \|D^{2}\operatorname{Op}(\mu^{-1})\nu(x)\chi w\|_{V}^{2}.$$
(3.26)

From Lemma 2.3, we have  $\operatorname{Op}(\mu^{-1})\nu\chi = \nu\chi\operatorname{Op}(\mu^{-1}) + R_1$ , where  $R_1$  has a symbol in  $S(\mu^{-2}\nu\lambda^{\frac{1}{2}},g)$ , and  $D^2\operatorname{Op}(\mu^{-1})\nu\chi = \nu D^2\chi\operatorname{Op}(\mu^{-1}) + R_2$ , where  $R_2$  has a symbol in  $S(\nu\lambda^{\frac{1}{2}},g)$ . Then, it follows from (3.26) that

$$\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\nu(x)\chi\operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D^{2}\chi\operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2} + \lambda \|\nu(x)\widetilde{u}\|_{V}^{2}.$$

For  $\tau \geq \max\{\mathcal{C}\lambda, 1\}$  with  $\mathcal{C}$  large enough, one has the following result from the above inequality.

$$\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \tau^{3} \|\nu(x)\chi\operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu D^{2}\chi\operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2}. \tag{3.27}$$

Now, we apply (3.25) to  $v = \chi \operatorname{Op}(\mu^{-1})\widetilde{u}$  to have

$$\begin{split} \tau^{3} \| \nu(x) \chi \operatorname{Op}(\mu^{-1}) \widetilde{u} \|_{V}^{2} + \tau \| \nu(x) D\chi \operatorname{Op}(\mu^{-1}) \widetilde{u} \|_{V}^{2} + \tau^{-1} \| \nu(x) D^{2} \chi \operatorname{Op}(\mu^{-1}) \widetilde{u} \|_{V}^{2} \\ \lesssim & \quad \| P_{\varphi} \chi \operatorname{Op}(\mu^{-1}) \widetilde{u} \|_{V}^{2}. \end{split}$$

Thus, combining this with (3.27) yields

$$\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \|P_{\varphi}\chi \operatorname{Op}(\mu^{-1})\widetilde{u}\|_{V}^{2}. \tag{3.28}$$

Finally, note that  $P_{\varphi}$  has a symbol in  $S(\nu\mu^2, g)$ . Consequently,  $P_{\varphi}\chi \operatorname{Op}(\mu^{-1}) = \operatorname{Op}(\mu^{-1})P_{\varphi}\chi + R$ , where R has a symbol in  $S(\nu\lambda^{\frac{1}{2}}, g)$ . Then, we can deduce from (3.28) that

$$\tau \|\nu(x)\widetilde{u}\|_V^2 + \tau^{-1} \|\nu(x)D\widetilde{u}\|_V^2 \lesssim \|\operatorname{Op}(\mu^{-1})P_\varphi\widetilde{u}\|_V^2 + \lambda \|\nu\widetilde{u}\|_V^2.$$

When  $\tau \geq C\lambda$  with C large enough, the error term  $\lambda \|\nu \tilde{u}\|_{V}^{2}$  can be absorbed by the left hand side. Consequently,

$$\tau \|\nu(x)\widetilde{u}\|_{V}^{2} + \tau^{-1} \|\nu(x)D\widetilde{u}\|_{V}^{2} \lesssim \|\operatorname{Op}(\mu^{-1})P_{\varphi}\widetilde{u}\|_{V}^{2}. \tag{3.29}$$

For  $\widetilde{u} = e^{\tau \varphi} u$ , it follows from (3.29) and similar argument as (3.21) that

$$\tau \| \nu(x) e^{\tau \varphi} u \|_{V}^{2} + \tau^{-1} \| \nu(x) e^{\tau \varphi} D u \|_{V}^{2} \lesssim \tau \| \nu(x) e^{\tau \varphi} u \|_{V}^{2} + \tau^{-1} \| \nu(x) D e^{\tau \varphi} u \|_{V}^{2} \\
\lesssim \| \operatorname{Op}(\mu^{-1}) P_{\varphi} e^{\tau \varphi} u \|_{V}^{2}. \tag{3.30}$$

Obviously, one has that

$$P_{\varphi}\widetilde{u} = e^{\tau\varphi}Pu = e^{\tau\varphi}g_0 + \sum_{i=1}^{d} e^{\tau\varphi}\partial_{x_j}g_j = e^{\tau\varphi}g_0 + \sum_{i=1}^{d} \left(\partial_{x_j}(e^{\tau\varphi}g_j) - \tau e^{\tau\varphi}g_j\partial_{x_j}\varphi\right),$$

which yields

$$\|\operatorname{Op}(\mu^{-1})P_{\varphi}\widetilde{u}\|_{V}^{2} \lesssim \sum_{j=0}^{d} \|e^{\tau\varphi}g_{j}\|_{V}^{2}.$$
 (3.31)

Hence, we obtain Theorem 3.3 from (3.30) and (3.31).

**Remark 3.2.** Since  $a(\cdot)$  is nonnegative and not identically null, there exists  $\delta > 0$  such that  $\{x \in \Omega : a(x) > \delta\} \neq \emptyset$ . We introduce several sets as follows.

$$W_0 = \Omega \setminus \text{supp } a,$$

$$W_1 = \overline{\Omega} \setminus \mathcal{O}(\text{supp } a),$$

$$W_2 = \Omega \setminus (\{x \in \Omega : a(x) \ge \delta\} \cup \mathcal{O}(\Gamma)),$$

$$W_3 = \Omega \setminus \{x \in \Omega : a(x) \ge \frac{\delta}{2}\},$$

where  $\mathcal{O}(\Gamma)$  means the neighborhood of  $\Gamma$ .

It is known that there exists a function  $\psi \in \mathscr{C}^{\infty}(\Omega)$  such that ([11])

- 1)  $\psi(x) = 0$  for  $x \in \partial \Omega$ .
- **2)**  $\partial_n \psi(x) < 0$  for  $x \in \partial \Omega$ .
- **3)**  $\nabla \psi(x) \neq 0$  for  $x \in \overline{\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}}$ .

Let  $\varphi = e^{\gamma \psi}$ . It follows from Lemma 2.5 that  $\varphi$  satisfies the sub-ellipticity condition on  $x \in \Omega \setminus \{x \in \Omega : a(x) \geq \delta\}$  if  $\gamma > 0$  is sufficiently large. Then, in Theorem 3.2 and 3.3, one can choose V as  $\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}$ .

The following result is a classical Carleman estimate and corresponding to the Laplacian with Dirichlet boundary condition ([17], Proposition 2). We shall use it to deal with the terms on  $\Omega \setminus A$ 

**Lemma 3.1.** Suppose  $\varphi$  is chosen as in Remark 3.2. Then, there exist C > 0,  $\lambda_0 > 0$ , such that for all  $u \in \mathscr{C}^{\infty}(\overline{\Omega})$  satisfying supp  $u \subset W_1$  and u = 0 on  $\Gamma$ , it holds

$$\tau^{3} \| e^{\tau \varphi} u \|_{\Omega}^{2} + \tau \| e^{\tau \varphi} D u \|_{\Omega}^{2} \lesssim \| e^{\tau \varphi} (D^{2} - \lambda^{2}) u \|_{\Omega}^{2}, \tag{3.32}$$

where  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{\mathcal{C}\lambda, 1\}$ .

**Theorem 3.4.** Suppose  $\varphi$  is chosen as in Remark 3.2. Let  $u \in \mathscr{C}^{\infty}(\overline{\Omega})$  and satisfy

$$P(x, D, \lambda)u = f_0 + \sum_{j=1}^{d} \partial_{x_j} f_j$$
 in  $\Omega$ ,  
 $u = 0$  on  $\Gamma$ .

where  $f_j \in L^2(\Omega)$ , supp  $f_0 \subset \Omega$  and supp  $f_j \subset \mathcal{O}(\text{supp } a)$  for  $j = 1, \dots, d$ . Then, there exist  $\tilde{K} > 0$ ,  $\lambda_0 > 0$ , such that for all  $\lambda \geq \lambda_0$  and  $\tau \geq \max\{\tilde{K}|\lambda|^{\frac{5}{4}}, 1\}$ , it holds

$$\tau \|e^{\tau\varphi}u\|_{W_3}^2 + \tau^{-1} \|e^{\tau\varphi}Du\|_{W_3}^2 \lesssim \sum_{i=0}^d \|e^{\tau\varphi}f_j\|_{\Omega}^2 + \lambda \|e^{\tau\varphi}u\|_{\{x\in\Omega: a(x)\geq \delta/2\}}^2,$$

where the positive constant  $\delta$  is defined as in Remark 3.2.

*Proof.* Let  $\chi_1, \chi_2 \in \mathscr{C}^{\infty}(\Omega)$  be non-negative and satisfy the following assumption

- (i)  $0 \le \chi_1, \ \chi_2 \le 1$ ;  $\chi_1$  and  $\chi_2$  are supported on  $W_1$  and  $W_2$ , respectively.
- (ii)  $\chi_1 + \chi_2 \geq 1$  in  $W_3$ . In particular,  $\chi_1 \equiv 1$  on  $[\mathcal{O}(\partial\Omega) \cap \Omega] \setminus \mathcal{O}(\text{supp } a)$ , and  $\chi_2 \equiv 1$  on  $\mathcal{O}(\text{supp } a) \setminus \{x \in \Omega : a(x) \geq \frac{\delta}{2}\}$ .

First, it is clear that

$$P\chi_2 u = \chi_2 f_0 + \sum_{j=1}^d (\partial_{x_j} (\chi_2 f_j) - f_j \partial_{x_j} \chi_2) + [P, \chi_2] u.$$

Since  $[P, \chi_2]$  is a first order operator, we have there exist  $a_0, a_1, \dots, a_d$  and  $b_0, b_1, \dots, b_d$  such that

$$[P,\chi_2]u = \sum_{j=0}^d a_j \partial_{x_j} u + \lambda \sum_{j=0}^d b_j \partial_{x_j} u,$$

where supp  $a_j \subset \{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\} \cup [\Omega \setminus (\mathcal{O}(\text{supp } a) \cup \mathcal{O}(\text{supp } \partial\Omega))], \text{ supp } b_j \subset \{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\} \text{ for } j = 0, 1, \dots, d.$  Then, applying Theorem 3.3 with  $\chi_2 u$  instead of u,  $\Omega \setminus \{x \in \Omega : a(x) \geq \delta\}$  instead of V, we obtain

$$\tau \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} \chi_2 u \|_{\Omega}^2 + \tau^{-1} \| (1 + \lambda^2 a(x)^2)^{\frac{1}{2}} e^{\tau \varphi} D \chi_2 u \|_{\Omega}^2$$

$$\lesssim \sum_{j=0}^d \| e^{\tau \varphi} f_j \|_{\Omega}^2 + \| e^{\tau \varphi} u \|_{\{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\} \cup (\Omega \setminus \sup a)}^2 + \lambda \| e^{\tau \varphi} u \|_{\{x \in \Omega : \frac{\delta}{2} < a(x) < \delta\}}^2.$$

Consequently, due to  $(1 + \lambda^2 a(x)^2) \ge 1$ ,  $\tau \ge \{\mathcal{C}\lambda, 1\}$  and  $\lambda \ge \lambda_0$ , we obtain

$$\tau \| (1 + \lambda^{2} a(x)^{2})^{\frac{1}{2}} e^{\tau \varphi} \chi_{2} u \|_{\Omega}^{2} + \tau^{-1} \| (1 + \lambda^{2} a(x)^{2})^{\frac{1}{2}} e^{\tau \varphi} D \chi_{2} u \|_{\Omega}^{2}$$

$$\lesssim \sum_{j=0}^{d} \| e^{\tau \varphi} f_{j} \|_{\Omega}^{2} + \| e^{\tau \varphi} u \|_{\Omega \setminus \text{supp } a}^{2} + \lambda \| e^{\tau \varphi} u \|_{\{x \in \Omega : a(x) \geq \frac{\delta}{2}\}}^{2}.$$
(3.33)

On the other hand, by using  $\chi_1 u$  instead of u in Lemma 3.1, we obtain

$$\tau^{3} \| e^{\tau \varphi} \chi_{1} u \|_{\mathcal{O}}^{2} + \tau \| e^{\tau \varphi} D \chi_{1} u \|_{\mathcal{O}}^{2} \lesssim \| e^{\tau \varphi} (P - i\lambda D a(x) D) \chi_{1} u \|_{\mathcal{O}}^{2}. \tag{3.34}$$

Since  $\chi_1 \partial_{x_j} f_j = 0$  for  $j = 1, \dots, d$ , we have

$$P\chi_1 u = \chi_1 f_0 + [P, \chi_1] u.$$

Therefore, combining these with (3.34) and  $\mathcal{O}(\text{supp } a) \cap \text{supp } \chi_1 = \emptyset$  yields

$$\tau \|e^{\tau\varphi}\chi_1 u\|_{\Omega}^2 + \tau^{-1} \|e^{\tau\varphi}D\chi_1 u\|_{\Omega}^2 \lesssim \tau^{-2} (\|e^{\tau\varphi}f_0\|_{\Omega}^2 + \|e^{\tau\varphi}[P, \chi_1]u\|_{\Omega}^2). \tag{3.35}$$

Adding up (3.35) and (3.33), using  $(1 + \lambda^2 a(x)^2) \ge 1$  and  $\tau \ge 1$ , we obtain

$$\tau \| e^{\tau \varphi} (\chi_1 + \chi_2) u \|_{\Omega}^2 + \tau^{-1} \| e^{\tau \varphi} D(\chi_1 + \chi_2) u \|_{\Omega}^2$$

$$\lesssim \sum_{j=0}^{d} \|e^{\tau \varphi} f_j\|_{\Omega}^2 + \|e^{\tau \varphi} u\|_{\Omega \setminus \text{supp } a}^2 + \lambda \|e^{\tau \varphi} u\|_{\{x \in \Omega : a(x) \ge \frac{\delta}{2}\}}^2 + \tau^{-2} \|e^{\tau \varphi} [P, \chi_1] u\|_{\Omega}^2.$$

Note that  $[P, \chi_1]$  is a first order operator and supported on supp  $\chi_1 \setminus \{x \in \Omega : \chi_1 \equiv 1\}$ , which is an subset of  $\{x \in \Omega : \chi_1 + \chi_2 = 1\}$ . Then, for  $\tau$  sufficiently large, we have

$$\tau \|e^{\tau \varphi}(\chi_1 + \chi_2)u\|_{\Omega}^2 + \tau^{-1} \|e^{\tau \varphi}(\chi_1 + \chi_2)Du\|_{\Omega}^2 \lesssim \sum_{j=0}^d \|e^{\tau \varphi}f_j\|_{\Omega}^2 + \lambda \|e^{\tau \varphi}u\|_{\{x \in \Omega : a(x) \ge \frac{\delta}{2}\}}^2.$$

Then as  $1 \le \chi_1 + \chi_2 \le 2$  on  $W_3$ , we obtain the statement of Theorem 3.4.

#### 4 Resolvent estimate

In this section, we shall prove the main result. From the results in [3, 7, 10], the logarithmic decay of the energy in Theorem 1.1 can be obtained through the following theorem.

**Theorem 4.1.** Suppose conditions in Theorem 1.1 hold. Then, for every  $\lambda \in \mathbb{R}$  with  $|\lambda|$  large enough, there exists C > 0 such that

$$\left\| (\mathcal{A} - i\lambda)^{-1} \right\|_{\mathcal{L}(\mathcal{H})} \lesssim e^{C |\lambda|^{\frac{5}{4}}}. \tag{4.1}$$

Let  $\lambda$  be a real number such that  $|\lambda|$  is large enough. Consider the resolvent equation:

$$F = (A - i\lambda)Y$$
, where  $Y = (y_1, y_2) \in D(A)$ ,  $F = (f_1, f_2) \in \mathcal{H}$ , (4.2)

which yields

$$\begin{cases}
\operatorname{div}(\nabla y_1 + a\nabla y_2) + \lambda^2 y_1 = i\lambda f_1 + f_2, & \text{in } \Omega, \\
y_2 = i\lambda y_1 + f_1 & \text{in } \Omega, \\
y_1 \mid_{\Gamma} = 0
\end{cases} (4.3)$$

In what follows, we shall prove that there exists a positive constant C such that

$$\|(y_1, y_2)\|_{\mathcal{H}} \lesssim e^{C |\lambda|^{\frac{5}{4}}} \|(f_1, f_2)\|_{\mathcal{H}}.$$

First, let  $\eta > 0$  and  $\chi \in \mathscr{C}^{\infty}_{0}(\Omega)$  be real valued function such that supp  $\chi = \{x \in \Omega : a(x) > \eta\}$ ) and  $\chi \equiv 1$  in  $\{x \in \Omega : a(x) > 2\eta\}$ . The following lemma is helpful.

**Lemma 4.1.** For  $y_1$ ,  $f_1$ ,  $f_2$  satisfying (4.3) and  $|\lambda|$  large enough, it holds

(i)  $\|\chi y_1\|_{\Omega} \lesssim \|\nabla y_1\|_{\{x \in \Omega : a(x) > \eta\}} + \|f_1\|_{H^1(\Omega)} + \|f_2\|_{\Omega}$ .

(ii) 
$$\int_{\Omega} a(x) |\nabla y_1|^2 dx \lesssim ||f_1||_{H^1(\Omega)}^2 + (||f_1||_{\Omega} + ||f_2||_{\Omega}) ||y_1||_{\Omega}.$$

*Proof.* (i) Multiplying the first equation in (4.3) by  $\chi^2 \overline{y}_1$  and using  $y_2 = i\lambda y_1 + f_1$ , we obtain,

$$\lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} = \int_{\Omega} \nabla y_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx + i\lambda \int_{\Omega} a(x) \nabla y_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx + i\lambda \int_{\Omega} f_{1} \chi^{2} \overline{y}_{1} dx + \int_{\Omega} f_{2} \chi^{2} \overline{y}_{1} dx + \int_{\Omega} a(x) \nabla f_{1} \cdot \nabla(\chi^{2} \overline{y}_{1}) dx.$$

$$(4.4)$$

Since  $\nabla(\chi^2 \overline{y}_1) = \chi^2 \nabla \overline{y}_1 + 2 \overline{y}_1 \chi \nabla \chi$ , we have

$$\left| \int_{\Omega} a(x) \nabla y_{1} \cdot \nabla (\chi^{2} \overline{y}_{1}) dx \right| \lesssim \|\nabla y_{1}\|_{\{x \in \Omega : a(x) > \eta\}}^{2} + \|\chi y_{1}\|_{\Omega} \|\nabla y_{1}\|_{\{x \in \Omega : a(x) > \eta\}}$$

$$\leq 2 \|\nabla y_{1}\|_{\{x \in \Omega : a(x) > \eta\}}^{2} + \frac{1}{4} \|\chi y_{1}\|_{\Omega}^{2}.$$

$$(4.5)$$

By the similar argument, we can deal with the rest terms on the right hand side of (4.4). Combining these with (4.4), (4.5) yields

$$\lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} \lesssim \frac{1}{2} \lambda^{2} \|\chi y_{1}\|_{\Omega}^{2} + (|\lambda| + 1) \|\nabla y_{1}\|_{\{x \in \Omega : a(x) > \eta\}}^{2} + \|f_{1}\|_{\Omega}^{2} + \lambda^{-2} \|f_{2}\|_{\Omega}^{2} + (\lambda^{-2} + 1) \|\nabla f_{1}\|_{H^{1}(\Omega)}^{2}.$$

$$(4.6)$$

Then, (i) is reached.

(ii) Multiplying the first equation in (4.3) by  $\overline{y}_1$  and using  $y_2 = i\lambda y_1 + f_1$ , we obtain,

$$\int_{\Omega} \left[ -(i\lambda a(x) + 1)|\nabla y_1|^2 + \lambda^2 |y_1|^2 \right] dx = \int_{\Omega} \left[ a(x)\nabla f_1 \cdot \nabla \overline{y}_1 + (i\lambda f_1 + f_2)\overline{y}_1 \right] dx. \tag{4.7}$$

Taking the imaginary part of (4.7) yields

$$\int_{\Omega} a(x) |\nabla y_1|^2 dx = -\lambda^{-1} \operatorname{Im} \int_{\Omega} [a(x) \nabla f_1 \cdot \nabla \overline{y_1} + (i\lambda f_1 + f_2) \overline{y_1}] dx. \tag{4.8}$$

Thus, by the Cauchy-Schwarz inequality, one can conclude from (4.8) that

$$\int_{\Omega} a(x) |\nabla y_{1}|^{2} dx \leq |\lambda|^{-1} \left( \int_{\Omega} a(x) |\nabla y_{1}|^{2} dx \right)^{\frac{1}{2}} ||f_{1}||_{H^{1}(\Omega)} + (||f_{1}||_{\Omega} + |\lambda|^{-1} ||f_{2}||_{\Omega}) ||y_{1}||_{\Omega} 
\leq \frac{1}{4} \int_{\Omega} a(x) |\nabla y_{1}|^{2} dx + |\lambda|^{-2} ||f_{1}||_{H^{1}(\Omega)}^{2} + (||f_{1}||_{\Omega} + |\lambda|^{-1} ||f_{2}||_{\Omega}) ||y_{1}||_{\Omega}.$$

The proof of lemma is finished.

**Proof of Theorem 4.1.** Due to Lemma 4.1 (i) and (ii), we have

$$||y_1||_{\{x\in\Omega: a(x)>2\eta\}}^2 \lesssim ||\nabla y_1||_{\{x\in\Omega: a(x)>\eta\}}^2 + ||f_1||_{H^1(\Omega)}^2 + ||f_2||_{\Omega}^2, \tag{4.9}$$

and

$$\|\nabla y_1\|_{\{x\in\Omega: a(x)>\eta\}}^2 \lesssim \|f_1\|_{H^1(\Omega)}^2 + (\|f_1\|_{\Omega} + \|f_2\|_{\Omega})\|y_1\|_{\Omega}. \tag{4.10}$$

Combining (4.9) and (4.10) yields

$$||y_1||_{H^1(\{x\in\Omega: a(x)>2\eta\})}^2 \lesssim ||f_1||_{H^1(\Omega)}^2 + ||f_2||_{L^2(\Omega)}^2 + (||f_1||_{\Omega} + ||f_2||_{\Omega})||y_1||_{\Omega}. \tag{4.11}$$

On the other hand, by (4.3), one has that  $y_1$  satisfies

$$\operatorname{div}(\nabla y_1 + i\lambda a(x)\nabla y_1) + \lambda^2 y_1 = i\lambda f_1 + f_2 - \operatorname{div}(a(x)\nabla f_1). \tag{4.12}$$

Then, applying Theorem 3.4 to  $y_1$  satisfying (4.12), we obtain

$$\tau \|e^{\tau\varphi}y_1\|_{\Omega}^2 + \tau^{-1} \|e^{\tau\varphi}\nabla y_1\|_{\Omega}^2$$

$$\lesssim \lambda^2 \|e^{\tau\varphi}f_1\|_{\Omega}^2 + \|e^{\tau\varphi}f_2\|_{\Omega}^2 + \sum_{j=1}^d \|e^{\tau\varphi}a\,\partial_{x_j}f_1\|_{\Omega}^2$$

$$+(\lambda+\tau) \|e^{\tau\varphi}y_1\|_{\{x\in\Omega: a(x)\geq \delta/2\}}^2 + \tau^{-1} \|e^{\tau\varphi}\nabla y_1\|_{\{x\in\Omega: a(x)\geq \delta/2\}}^2.$$

Let  $c_1 = \min_{x \in \Omega} \varphi(x)$  and  $c_2 = \max_{x \in \Omega} \varphi(x)$ . We conclude from the above inequality and  $\tau \ge \max\{\widetilde{K}|\lambda|^{\frac{5}{4}}, 1\}$  that

$$\tau e^{2c_1\tau} \|y_1\|_{\Omega}^2 + \tau^{-1} e^{2c_1\tau} \|\nabla y_1\|_{\Omega}^2$$

$$\lesssim \lambda^2 e^{2c_2\tau} \|f_1\|_{H^1(\Omega)}^2 + e^{2c_2\tau} \|f_2\|_{\Omega}^2 + \tau e^{2c_2\tau} \|y_1\|_{H^1(\{x \in \Omega : a(x) > \delta/2\})}^2.$$
(4.13)

Setting  $\eta = \frac{\delta}{8}$  and substituting (4.11) into (4.13), we obtain

$$\tau e^{2c_1\tau}\|y_1\|_{\Omega}^2 + \tau^{-1}e^{2c_1\tau}\|\nabla y_1\|_{\Omega}^2 \lesssim \tau^2 e^{2c_2\tau}\|f_1\|_{H^1(\Omega)}^2 + \tau e^{2c_2\tau}\|f_2\|_{\Omega}^2 + \tau e^{2c_2\tau}(\|f_1\|_{\Omega} + \|f_2\|_{\Omega})\|y_1\|_{\Omega}.$$

Let  $c_3 = 2(c_2 - c_1) + 1$ . For  $\tau \ge \max\{\widetilde{K}|\lambda|^{\frac{5}{4}}, 1\}$ , one has

$$||y_1||_{\Omega}^2 + ||\nabla y_1||_{\Omega}^2 \lesssim e^{c_3\tau} ||f_1||_{H^1(\Omega)}^2 + e^{c_3\tau} ||f_2||_{\Omega}^2 + e^{c_3\tau} (||f_1||_{\Omega} + ||f_2||_{\Omega}) ||y_1||_{\Omega}.$$

For any  $\varepsilon > 0$ , using  $e^{c_3\tau}(\|f_1\|_{\Omega} + \|f_2\|_{\Omega})\|y_1\|_{\Omega} \le \varepsilon \|y_1\|_{\Omega}^2 + \frac{1}{4}\varepsilon^{-1}e^{2c_3\tau}(\|f_1\|_{\Omega} + \|f_2\|_{\Omega})^2$  in the above estimate, we conclude that

$$||y_1||_{\Omega}^2 + ||\nabla y_1||_{\Omega}^2 \lesssim e^{2c_3\tau} (||f_1||_{H^1(\Omega)}^2 + ||f_2||_{\Omega}^2),$$

which gives the desired result.

#### References

- [1] K. Ammari, F. Hassine and L. Robbiano, Stabilization for the wave equation with singular Kelvin-Voigt damping, arXiv:1805.10430
- [2] C. Bardos, G. Lebeau and J. Rauch, Sharp sufficient conditions for the observation, control and stabilization of waves from the boundary, SIAM J. Control Optim., 30(5): 1024–1065, 1992.
- [3] C. Batty, T. Duyckaerts, Non-uniform stability for bounded semi-groups on Banach spaces, J. Evol. Equ. 8(4): 765–780, 2008
- [4] M. Bellassoued, Distribution of resonances and decay rate of the local energy for the elastic wave equation, Comm. Math. Phys., 215(2):375–408, 2000.
- [5] M. Bellassoued, Carleman estimates and distribution of resonances for the transparent obstacle and application to the stabilization. Asymptot. Anal., 35(3-4):257–279, 2003.
- [6] M. Bellassoued, Decay of solutions of the elastic wave equation with a localized dissipation, Ann. Fac. Sci. Toulouse Math. (6), 12(3):267–301, 2003.
- [7] N. Burq, Décroissance de l'énergie locale de l'équation des ondes pour le problème extérieur et absence de résonance au voisinage du réel, Acta Math., 180(1):1–29, 1998.
- [8] G. Chen, S. A. Fulling, F. J. Narcowich, and S. Sun, Exponential decay of energy of evolution equation with locally distributed damping, SIAM J. Appl. Math., 51(1): 266–301, 1991.
- [9] S. Chen, K. Liu and Z. Liu, Spectrum and stability for elastic systems with global or local Kelvin-Voigt damping, SIAM J. Appl. Math., 59(2): 651–668, 1998.
- [10] T. Duyckaerts, Optimal decay rates of the energy of a hyperbolic-parabolic system coupled by an interface, Asymptot. Anal., 51(1):17–45, 2007.
- [11] A. Fursikov and O. Yu. Imanuvilov, Controllability of evolution equations, V. 34, Seoul National University, Korea, 1996, Lecture notes.
- [12] L. Hörmander, The Analysis of Linear Partial Differential Operators, Springer-Verlag, Berlin, 1983–85.
- [13] F. Huang, On the mathematical model for linear elastic systems with analytic damping, SIAM J. Control Optim., 26(3): 714–724, 1988.

- [14] O. Y. Imanuvilov, J.-P. Puel, Global Carleman estimates for weak solutions of elliptic nonhomogeneous Dirichlet problems, Int. Math. Res. Not., 16: 883–913, 2003.
- [15] J. Le Rousseau and L. Robbiano, Carleman estimate for elliptic operators with coefficients with jumps at an interface in arbitrary dimension and application to the null controllability of linear parabolic equations, HAL: hal-00193885.
- [16] G. Lebeau, Équation des ondes amorties, In algebraic and geometric methods in mathematical physics (Kaciveli, 1993), volume 19 of Math. Phys. Stud., 73-109, Kluwer Acad. Publ., Dordrecht, 1996.
- [17] G. Lebeau and L. Robbiano, Contrôle exact de l'équation de la chaleur, Comm. Partial Differential Equations, 20(1-2):335–356, 1995.
- [18] G. Lebeau and L. Robbiano, Stabilisation de l'équation des ondes par le bord, Duke Math. J., 86(3):465–491, 1997.
- [19] K. Liu, Z. Liu and Q. Zhang, Eventual differentiability of a string with local Kelvin-Voigt damping, to appear in ESAIM: Control, Optimization and the Caculus of Variations, 23(2): 443-454, 2017.
- [20] K. Liu and B. Rao, Exponential stability for the wave equation with local Kelvin-Voigt damping,Z. Angew. Math. Phys., 57(3): 419-432, 2006.
- [21] Z. Liu and B. Rao, Frequency domain characterization of rational decay rate for solution of linear evolution euqations, Z. Angew. Math. Phys., 56(4): 630–644, 2005.
- [22] Z. Liu and Q. Zhang, Stability of a string with local Kelvin-Voigt damping and non-smooth coefficient at interface, SIAM J. Control. Optim. 54(4):1859–1871, 2016.
- [23] A. Pazy, Semigroups of Linear Operators and Applications to Partial Differential Equations, Springer-Verlag, New York, 1983.
- [24] J. Rauch, X. Zhang and E. Zuazua, Polynomial decay for a hyperbolic-parabolic coupled system, J. Math. Pures Appl., 84(4): 407–470, 2005.
- [25] M. Renardy, On localized Kelvin-Voigt damping, Z. Angew. Math. Mech, 84(4): 280–283, 2004.
- [26] L. Robbiano, Fonction de coût et contrôle des solutions des équations hyperboliques, Asymptotic Anal., 10(2):95–115, 1995.
- [27] Q. Zhang, Exponential stability of an elastic string with local Kelvin-Voigt damping, Z. Angew. Math. Phys., 6(6):1009–1015, 2010.
- [28] Q. Zhang, On the lack of exponential stability for an elastic-viscoelastic waves interaction system, Nonlinear Analysis: Real World Applications 37: 387–411, 2017.
- [29] Q. Zhang, Polynomial decay of an elastic/viscoelastic waves interaction system, submitted.
- [30] X. Zhang and E. Zuazua, Long time behavior of a coupled heat-wave system arising in fluid-structure interaction, Arch. Rat. Mech. Anal. 184: 49–120, 2007.