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Luminescence of Ho-doped lithium niobate crystals highlighted by Raman Spectroscopy

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ABSTRACT

Raman measurements were carried out on Ho³⁺ doped Lithium Niobate crystals. When the excitation wavelength of 532nm is used, in addition of expected Raman modes, forbidden bands are detected, while when exciting with 785nm, “classical” Raman spectrum was recorded with expected modes according to Raman selection rules. Additional lines are attributed to emission lines of Ho doped crystals. We detect, within a very good resolution, in the same Stokes spectrum, the transitions between the electronic states, and the vibrational states as well. We report on the analysis of these data as function of Ho-content, for different polarizations and wavelengths, of the incident laser beam.

Keywords: Lithium niobate, Raman spectroscopy, luminescence

1. INTRODUCTION

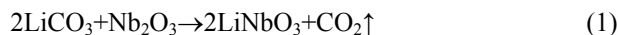
Lithium Niobate (LN) is one of the most used ferroelectrics having a wide number of applications such as phase-conjugation, holographic storage, frequency doubling, SAW sensors [1-3]. It is a versatile material, the properties of which can be changed according to the nature and concentration of doping [4-6]. Doping these crystals with rare earth impurity ions can be interesting for laser applications, and therefore characterization of emission spectra is required. Doping with holmium presents a special interest for up-conversion emission in green and red [7-9]. Moreover, recent works have shown that doping LN crystals by Ho³⁺ ions leads to the suppression of optical damage at 532nm [8] and it can be used for optical cooling based on the anti-Stokes luminescence in the wavelength range around 2μm [10].

Raman spectroscopy is a powerful spectroscopic technique providing a possibility to obtain a number of information about physico-chemical as well as optical properties of a given material [11,12]. This technique was proved to be efficient to detect the influence of introduction of dopant ion in the lattice. Here it is used in order to investigate the emission states beside vibrational modes.

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2. EXPERIMENTAL PART AND ANALYSIS

High purity materials in powder form from Johnson-Matthey (Nb_2O_5) and Merck (Li_2CO_3 and Ho_2O_3) were used to synthesize the initial material for the growth. LN powders doped with different concentrations (0.1, 1mol%) of Ho^{3+} ions were synthesized at 1000°C in two steps by solid reactions given below:



The obtained powders were investigated by the help of X-ray diffraction using Philips X'Pert diffractometer. The obtained results confirm existence of no any other phase than lithium niobate in above mentioned powders for all concentrations of impurity ions. In addition, the X-ray diffraction patterns allowed to determine the lattice parameters of the powders. In general, both a and c parameters of the Ho-doped powders under study increase, but if the behavior of the change of the a parameter with the change of the concentration of the impurity ion can be attributed to a monotonous one, in the case of the c parameter it is not monotonous. Most probably it is due to the fact that site occupation of Ho^{3+} ions in the lattice of the LN is different depending of its concentration.

Congruent composition LN crystals doped with different of Ho^{3+} impurity ions were grown by the modified Czochralski method. The growth was realized in air using a setup with a high-frequency heating furnace and a platinum crucible having a size of $50 \times 3 \times 50 \text{ mm}^3$. The crystals were pulled along the c -axis at a rate of 1 mm/h and a rotation speed of about 20 rpm.

Among additions conditions of the growth process it should be noted, that during the mentioned process, a direct electric field (electrical current of about 12 A/m^2) was applied to the crystal-melt system in order to obtain single domain crystals directly during the growth process and a rather homogeneous distribution of impurity ions within all the grown crystal.

Raman spectra were carried out utilizing a confocal Raman microscope LabRAM HR Evolution (*Horiba*) with exciting lines of 532 and 785nm. The laser light was focused through a $50\times$ objective (*Olympus*) at a distance above 1 cm from the sample. Raman system has a spectral resolution of 1 cm^{-1} using a grating with 1800 grooves per millimeter. In our series of measurements, we recorded backscattering Raman measurements in polarized configurations in order to obtain different Raman modes depending on their symmetry. The spectra were carried out in the configurations $\text{Y}(\text{ZZ})\text{Y}$ and $\text{Y}(\text{XZ})\text{Y}$ corresponding to $\text{A}_1(\text{TO})$ and $\text{E}(\text{TO})$ phonon modes respectively.

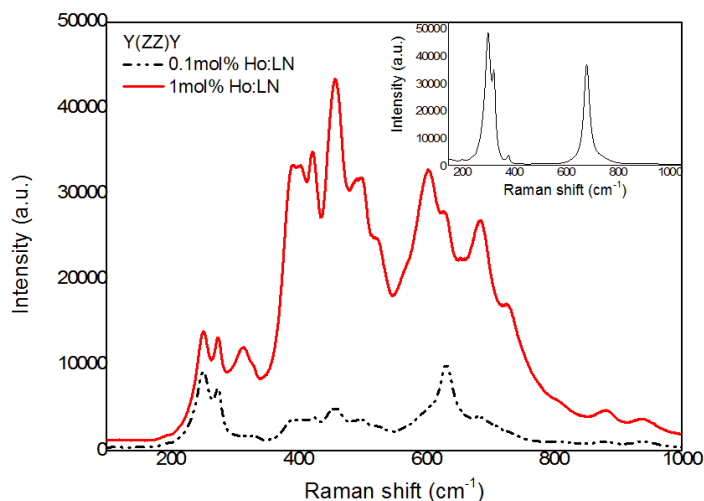


Fig. 1 – Spectra recorded with 532nm excitation wavelength on 0.1 and 1mol% Ho doped LN in $\text{Y}(\text{ZZ})\text{Y}$ configuration. The insert illustrates the spectrum of 1mol% Ho doped sample obtained with 785nm laser.

Figure 1 illustrates Raman spectra recorded on LN: 0.1 and 1 mol% Ho with 532nm exciting laser line. According to Raman selection rules only $A_1[TO]$ modes are expected for Y(ZZ)Y configuration while we notice multiple additional forbidden lines in both spectra. Only difference is the higher intensity for 1mol% Ho doped LN sample comparing to 0.1mol% doped one. In contrast, when 785nm is used solely 4 $A_1[TO]$ modes are detected. The same behavior is observed for Y(XZ)Y configuration where only $E[TO]$ modes are expected according to selection rules (figure 2).

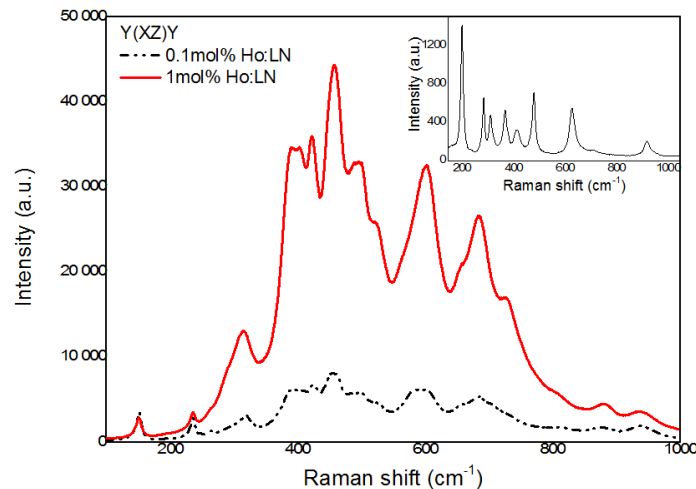


Fig. 2 – Spectra recorded with 532nm excitation wavelength on 0.1 and 1mol% Ho doped LN in Y(XZ)Y configuration. The insert illustrates the spectrum of 1mol% Ho doped sample obtained with 785nm laser.

It is clear that these additional bands are not attributed to Raman modes and are related to the emission by Ho^{3+} ions in the crystal lattice. It should be pointed out that here the emission lines are as narrow as vibrational ones. That means that we detect within a very good resolution, Raman lines and emission lines simultaneously in the same Stokes scattering spectrum. In figure 3 we present spectra as a function of wavelength for both samples in both configurations. Our results are in rather good agreement with luminescence results reported in literature [8,9]

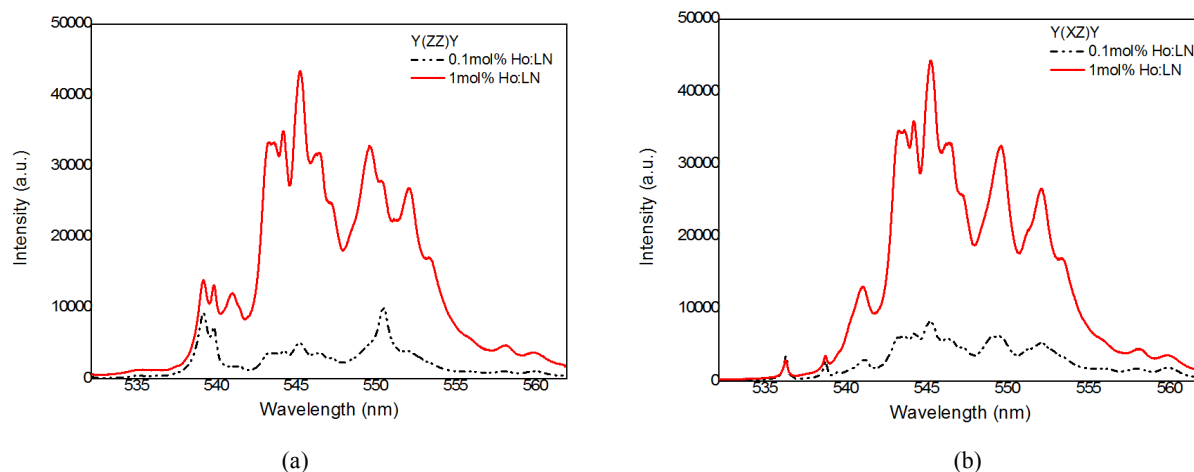


Fig. 3 – Spectra as a function of wavelength of 0.1 and 1% mol Ho doped LN (a) Y(ZZ)Y configuration; (b) Y(XZ)Y configuration

Emission lines of Ho doped LN in the range of 535nm-560nm are assigned to $^5S_2 + ^5F_4 \rightarrow ^5I_8$ Stark transitions [8,9]. One can notice that emission spectra in two configurations are not totally identical (without taking into account Raman

expected lines). This difference is related to π and σ emissions corresponding to Y(ZZ)Y and Y(XZ)Y configurations respectively. Our recent work [13] describes a methodology and data treatment procedure needed to separate Raman and emission lines in order to obtain purely the luminescence spectra.

3. CONCLUSION

Raman spectra of Ho doped LN were obtained with 532nm and 785nm excitation wavelengths and were found totally different from each other for both Ho concentrations (0.1 and 1mol%). Using 785nm incident laser beam only Raman modes were detected in accordance with selection rules. In contrast, when using 532nm laser in obtained spectra additional unexpected lines were detected. These spectra represent luminescence of Ho doped LN crystals superimposed on Raman modes. Obtained non-classical Raman spectra provide a possibility to acquire emission spectra with a rather good resolution. Our results are consonant with literature where emission spectra were obtained with optical absorption measurements.

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