

Pitman transforms and Brownian motion in the interval viewed as an affine alcove

Philippe Bougerol, Manon Defosseux

▶ To cite this version:

Philippe Bougerol, Manon Defosseux. Pitman transforms and Brownian motion in the interval viewed as an affine alcove. Annales Scientifiques de l'École Normale Supérieure, 2022, 55 (2), pp.429–472. 10.24033/asens.2499 . hal-01861862v4

HAL Id: hal-01861862 https://hal.science/hal-01861862v4

Submitted on 8 Jun2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

PITMAN TRANSFORMS AND BROWNIAN MOTION IN THE INTERVAL VIEWED AS AN AFFINE ALCOVE

PHILIPPE BOUGEROL AND MANON DEFOSSEUX

ABSTRACT. Pitman's theorem states that if $\{B_t, t \ge 0\}$ is a one dimensional Brownian motion, then $\{B_t - 2\inf_{0 \le s \le t} B_s, t \ge 0\}$ is a three dimensional Bessel process, i.e. a Brownian motion conditioned to remain forever positive. This paper gives a similar representation for the Brownian motion conditioned to remain in a given interval. Due to the double barrier condition, this representation is more involved and only asymptotic. One uses the fact that the interval is an alcove of the Kac-Moody affine Lie algebra $A_1^{(1)}$, the Littelmann path approach of representation theory and a dihedral approximation.

RÉSUMÉ. Le théorème de Pitman affirme que si $\{B_t, t \ge 0\}$ est un mouvement brownien unidimensionnel, alors $\{B_t - 2\inf_{0\le s\le t} B_s, t\ge 0\}$ est un processus de Bessel de dimension trois, c'est-à-dire un brownien conditionné à rester positif. Nous donnons dans cet article une représentation analogue pour le brownien conditionné à rester dans un intervalle donné. En raison de la présence de deux extrémités, cette représentation est plus compliquée que celle du théorème original. Nous utilisons le fait que l'intervalle est une alcôve pour l'algèbre de Kac-Moody affine $A_1^{(1)}$, l'approche par le modèle de chemins de Littelmann de la théorie des représentations et une approximation diédrale.

> couleurs de la corde dépôt de cette image cristaux du temps traces d'espace

Raymond Queneau

Contents

1.	Introduction	2
2.	Conditioned Brownian motions	6
2.1.	The conditioned Brownian motion in $[0, 1]$	6
2.2.	The conditioned space-time Brownian motion in C_{aff}	7
3.	Approximation by planar Brownian motions in dihedral cones	9
3.1.	Conditioned Brownian motion in a dihedral cone	9
3.2.	Pitman representation in a dihedral cone	11
3.3.	Approximation of the conditioned space time process	12
4.	String parameters in the dihedral case	13
5.	Affine string parameters of the space-time Brownian motion	17
5.1.	Pitman and Lévy transforms	17

5.2. Affine string parameters	17
5.3. Affine Verma weight $L^{(\mu)}(\infty)$	21
6. The highest weight process $\Lambda^{(\mu)}$	27
7. The representation theorem using Pitman and Lévy transforms	30
7.1. Representation of the conditioned space-time Brownian motion $A^{(\mu)}$	30
7.2. Representation of the conditioned Brownian motion in $[0, 1]$	31
7.3. A property of Pitman transform for piecewise C^1 paths	32
8. Some probability distributions	33
8.1. The conditional distribution of $B_t^{(\mu)}$ given $\sigma(\Lambda^{(\mu)}(s), s \leq t)$	33
8.2. Distribution of $L^{(\mu)}(\infty)$	35
8.3. Distribution of $\xi_1(\infty)$	35
9. Asymptotics for representations of the affine algebra $A_1^{(1)}$	37
9.1. The Kac-Moody algebra $A_1^{(1)}$	37
9.2. An analogue of the Duistermaat Heckman measure	38
9.3. Asymptotics for the crystal $B(\infty)$ of $A_1^{(1)}$	39
References	

1. INTRODUCTION

1.1. The probability transition of the Brownian motion conditioned to stay positive forever is the Doob transform of the difference of two heat kernels. This is a consequence of the reflection principle at 0. Pitman's theorem [33] of 1975 gives the path representation of this process as

$$\mathcal{P}B(t) = B_t - 2\inf_{0 \le s \le t} B_s,$$

where B is a standard Brownian motion with $B_0 = 0$. The transform $\mathcal{P}B$ is written with the reflection at 0. Consider now a Brownian motion conditioned to stay in the interval [0, 1] forever. Is it possible to write it as a path transform of a Brownian motion B by some kind of folding? The conditioned process can be seen as the Doob transform of the Brownian motion killed at 0 and at 1. Its probability transition is an alternating infinite sum which can be obtained by applying successive reflection principles at 0 and at 1 (method of images). It is therefore natural to ask if Pitman's theorem has an analogue for the conditioned process in the interval written with two similar transforms at 0 and 1, maybe repeated an infinite number of times. The main result of this article is to show that, to our surprise, this is not exactly the case. A small correction (a Lévy transform) has to be added: this is due to the non differentiability of the Brownian motion. Interestingly, the same correction also occurs in an asymptotic property of the highest weight representations of the affine Lie algebra $A_1^{(1)}$. Hence there are deep links between the trajectories of the Brownian motion and representation theory of Kac Moody algebras.

1.2. Let us state our main result. We suppose that $\mu \in [0, 1]$. We will give in Section 2.1 a precise definition of the following process.

Notation 1.1. Z^{μ} is a Brownian motion conditioned to stay in [0,1] forever such that $Z_0^{\mu} = \mu$.

 $\mathbf{2}$

We consider, for a continuous real path $\varphi : \mathbb{R}_+ \to \mathbb{R}$ such that $\varphi(0) = 0$, for $t \ge 0$,

$$\mathcal{L}_1 \varphi(t) = \varphi(t) - \inf_{\substack{0 \le s \le t}} \varphi(s),$$

$$\mathcal{P}_1 \varphi(t) = \varphi(t) - 2 \inf_{\substack{0 \le s \le t}} \varphi(s).$$

We call them the classical Lévy and Pitman transforms of φ . We introduce

$$\mathcal{L}_0\varphi(t) = \varphi(t) + \inf_{\substack{0 \le s \le t}} (s - \varphi(s)),$$

$$\mathcal{P}_0\varphi(t) = \varphi(t) + 2\inf_{\substack{0 \le s \le t}} (s - \varphi(s)).$$

For $n \in \mathbb{N}$, we let $\mathcal{P}_{2n} = \mathcal{P}_0$, $\mathcal{L}_{2n} = \mathcal{L}_0$ and $\mathcal{P}_{2n+1} = \mathcal{P}_1$, $\mathcal{L}_{2n+1} = \mathcal{L}_1$. The aim of this paper is the following representation theorem (see Theorem 7.5).

Theorem. Let $\mu \in [0, 1]$ and let $B_t^{\mu} = B_t + t\mu$ be the real Brownian motion with drift μ starting from 0. For any t > 0, almost surely,

$$\lim_{n \to \infty} t \mathcal{L}_{n+1} \mathcal{P}_n \cdots \mathcal{P}_1 \mathcal{P}_0 B^{\mu}(1/t) = \lim_{n \to \infty} t \mathcal{L}_{n+1} \mathcal{P}_n \cdots \mathcal{P}_2 \mathcal{P}_1 B^{\mu}(1/t) = Z_t^{\mu}.$$

1.3. Briefly, the strategy of the proof is as follows. The first step is to linearize the problem (the reflection at 1 is not linear): we introduce the process $A^{(\mu)}$ which is the space-time Brownian motion $B_t^{(\mu)} = (t, B_t^{\mu}), t \ge 0$, conditioned to stay in the affine cone

$$C_{\text{aff}} = \{ (t, x) \in \mathbb{R}^2; 0 < x < t \}.$$

We show that Z^{μ} is the space component of the time inverted process of $A^{(\mu)}$, i.e. $A_t^{(\mu)} = (t, tZ_{1/t}^{\mu})$ in distribution. So we will work essentially with $A^{(\mu)}$.

We define a sequence of non-negative random processes $\xi_n(t), t \ge 0, n \in \mathbb{N}$, by

(1.1)
$$\xi_n(t) = -\inf_{0 \le s \le t} \{ s \mathbb{1}_{2\mathbb{N}}(n) + (-1)^{n-1} \mathcal{P}_{n-1} \cdots \mathcal{P}_0 B^{\mu}(s) \}.$$

Then

(1.2)
$$\mathcal{P}_n \cdots \mathcal{P}_0 B^{\mu}(t) = B_t^{\mu} + 2 \sum_{k=0}^n (-1)^{k+1} \xi_k(t),$$

and

(1.3)
$$\mathcal{L}_{n+1}\mathcal{P}_n\cdots\mathcal{P}_0B^{\mu}(t)=\mathcal{P}_n\cdots\mathcal{P}_0B^{\mu}(t)+(-1)^n\xi_{n+1}(t).$$

We first suppose that $\mu \neq 0, 1$. Then the random variables

$$\xi_n(\infty) = \lim_{t \to +\infty} \xi_n(t)$$

are finite a.s. and their distributions have a simple explicit representation with independent exponential random variables. The properties of $\xi_n(t), n \in \mathbb{N}$, can be deduced by conditioning arguments. This allows us to show that for all $t \geq 0$ the limit of (1.3) exists a.s. and has the distribution of the space component of $A^{(\mu)}$. We also prove that for t > 0, $\xi_n(t)$ tends to 2 almost surely when n tends to $+\infty$. This shows that the limit of (1.2) itself does not exist. The boundary cases $\mu = 0, 1$, are dealt with using the Cameron–Martin–Girsanov (CMG) theorem.

To prove these results, we approximate the space-time Brownian motion $B^{(\mu)}$ by planar Brownian motions with proper drifts and we approximate $A^{(\mu)}$ by these planar Brownian motions conditioned to remain in a wedge in \mathbb{R}^2 of dihedral angle π/m . The dihedral case has been dealt with in Biane et al. [3] and we use their results. Due to the need of the correction term, the approximation is not immediate.

1.4. The article is organized as follows. We always suppose that $0 \leq \mu \leq 1$. In Section 2 we first define rigorously Z^{μ} . Then we define $A^{(\mu)}$ the conditioned space-time Brownian motion with drift μ in the affine Weyl cone C_{aff} . We prove in Theorem 2.4 that, as processes, $A_t^{(\mu)} = (t, tZ_{1/t}^{\mu})$ in distribution. In Section 3 we recall the Pitman representation theorem for planar Brownian motions in a dihedral cone and show how they approximate $A^{(\mu)}$. In Section 4 we introduce the string parameters in the dihedral case. The analogous parameters $\{\xi_n(t), n \in \mathbb{N}, t \geq 0\}$ for the space time Brownian motion $B^{(\mu)}$ are defined in Section 5 and called the affine string parameters. When $\mu \neq 0, 1$, then $\xi_n(\infty), n \in \mathbb{N}$, are finite and are called the Verma affine parameters. In Section 6 we study the highest weight process $\Lambda^{(\mu)}$ which is the limit of the image of $B^{(\mu)}$ under the transformation (1.3). It is shown in Section 7 that $\Lambda^{(\mu)}$ equals $A^{(\mu)}$ in distribution. The representation theorem for the Brownian motion Z^{μ} in [0, 1] follows. In Section 8 we first compute the conditional distribution of $B_t^{(\mu)}$ given the sigma-algebra $\sigma\{\Lambda^{(\mu)}(s), s \leq t\}$ and then the distributions of $L^{(\mu)}(\infty)$ and $\xi_1(\infty)$. Up to this point we only use probabilistic arguments with no reference to algebra.

In Section 9, we introduce the infinite dimensional affine Lie algebra $A_1^{(1)}$ and show how our results are related to its highest weight representations. We show that the conditional distribution of the Brownian motion is a Duistermaat Heckman measure for a circle action. It describes the semiclassical behaviour of the weights of a representation when its highest weight is large. The Lévy correction term also occurs in the behaviour of the elements of large weight of the Kashiwara crystal $B(\infty)$, which is of independent interest.

1.5. We have chosen to present the proof of our probabilistic results without explicit reference to Kac-Moody algebra, so that it can be read by probabilists. But let us now explain the ideas of representation theory behind the scenes because this has been a source of inspiration. This may be helpful for some readers. This also explain our choice of terminolgy (affine cone, string parameter, highest weight, ...).

At the heart of our approach is the fact that the interval [0, 1] is an alcove for the Kac-Moody affine Lie algebra $A_1^{(1)}$ and that $A^{(\mu)}$ can be seen as a process conditioned to remain in a Weyl chamber. This is linked to highest weight representations of $A_1^{(1)}$ through Littelmann path approach (see [29]).

To give some details, let us first recall the link between Littelmann path theory for the Lie algebra $\mathfrak{sl}_2(\mathbb{C})$ and the classical Pitman theorem as explained in Biane et al. [3]. Consider a real line $V = \mathbb{R}\alpha$ where α is the positive root. A path η in V is a continuous function $\eta : \mathbb{R}^+ \to V$ such that $\eta(0) = 0$. It can be written as $\eta(s) = \varphi(s)\alpha$ with $\varphi(s) \in \mathbb{R}$, for $s \ge 0$. A dominant path is a path with values in the Weyl chamber, which is here $\mathbb{R}_+\alpha$, so that η is dominant when $\varphi(s) \ge 0$ for all $s \ge 0$. We fix some t > 0. An integral path on [0, t] is a piecewise linear path such that $2\varphi(t)$ and $2\min_{0\le s\le t}\varphi(s)$ are in \mathbb{Z} . Consider an irreducible highest weight $\mathfrak{sl}_2(\mathbb{C})$ -module with highest weight λ . It has a combinatorial description given by a Kashiwara crystal (see Kashiwara [26]). Littelmann gives a path realization of this crystal (see Littelmann [29]). One chooses a dominant integral path π on [0, t], such that $\pi(t) = \lambda$. The Littelmann module is the set of integral paths η on [0, t] such that $\mathcal{P}_{\alpha}\eta(s) = \pi(s)$ for all $s \in [0, t]$, where \mathcal{P}_{α} is the path transform defined by

(1.4)
$$\mathcal{P}_{\alpha}\eta(s) = (\varphi(s) - 2\inf_{0 \le u \le s}\varphi(u))\alpha,$$

when $\eta(s) = \varphi(s)\alpha$. One recognizes the Pitman transform.

Let us consider now the Littelmann path theory for the affine Lie algebra $A_1^{(1)}$. This Kac Moody algebra corresponds to the group of loops of $Sl(2, \mathbb{C})$. Let $V = \mathbb{R}\Lambda_0 \oplus \mathbb{R}\Lambda_1 \oplus \mathbb{R}\delta$ be a 3 dimensional real vector space with basis $\Lambda_0, \Lambda_1, \delta$. We let $\alpha_0 = \Lambda_0 - \Lambda_1 + \delta, \alpha_1 = \Lambda_1 - \Lambda_0$. In the dual space \tilde{V} we define $\tilde{\alpha}_0, \tilde{\alpha}_1$ by requiring that $\tilde{\alpha}_0(\delta) = \tilde{\alpha}_1(\delta) = 0$ and that $(\tilde{\alpha}_i(\alpha_j))_{0 \leq i,j \leq 1}$ is the Cartan matrix

$$\begin{pmatrix} 2 & -2 \\ -2 & 2 \end{pmatrix}.$$

The Weyl group W is generated by the linear reflections s_0, s_1 on V defined by

$$s_i(v) = v - \tilde{\alpha}_i(v)\alpha_i,$$

for $v \in V$ and $i \in \{0, 1\}$. They are reflections along the walls of the Weyl chamber

$$C_W = \{ t\Lambda_0 + x\Lambda_1 + y\delta, 0 < x < t, y \in \mathbb{R} \} \simeq C_{\text{aff}} \times \mathbb{R}.$$

In Littlemann's theory, a path is now a continuous map $\eta : \mathbb{R}^+ \to V$ such that $\eta(0) = 0$. One defines path transforms $\mathcal{P}_{s_i}, i \in \{0, 1\}$, by

$$\mathcal{P}_{s_i}\eta(t) = \eta(t) - \inf_{0 \le s \le t} \tilde{\alpha}_i(\eta(s))\alpha_i.$$

A dominant path is a path with values in the closure of C_W and one can define integral paths on [0, t]. For a fixed t > 0 and an integral dominant path π on [0, t], the Littlemann module generated by π is the set of integral paths η on [0, t] for which there exists $n \in \mathbb{N}$ such that

$$\mathcal{P}_{s_n}\mathcal{P}_{s_{n-1}}\cdots\mathcal{P}_{s_1}\mathcal{P}_{s_0}\eta(s)=\pi(s)$$

when $0 \leq s \leq t$, where $s_{2k} = s_0$ and $s_{2k+1} = s_1$ for $k \in \mathbb{N}$. This gives a description of the Kashiwara crystal of highest weight $\pi(t)$ (see Kashiwara [27], Littelmann [29]). For an integral path η on [0, t] and more generally for a continuous piecewise C^1 path, there is a $k \in \mathbb{N}$ such that for all $n \geq k$,

$$\mathcal{P}_{s_n}\cdots\mathcal{P}_{s_0}\eta(s)=\mathcal{P}_{s_k}\cdots\mathcal{P}_{s_0}\eta(s),$$

for $0 \le s \le t$ (see Proposition 7.7). This new path is dominant. One can ask if, similarly to Pitman's theorem, at least the limit, when *n* tends to infinity, of $\mathcal{P}_{s_n} \cdots \mathcal{P}_{s_0} \eta$ exists when η is replaced by a space-time Brownian motion.

We show in this paper that this is not the case, but a slight modification does converge in $V/\mathbb{R}\delta$ to a space-time Brownian motion conditioned to remain in the affine cone C_{aff} . It will be enough for us to work in the space $V/\mathbb{R}\delta$, so we will use a space V without δ .

We call the sequence $\{\xi_k(t), k \in \mathbb{N}\}$ defined in (1.1) the affine string parameters by analogy with the string parameters in the Littelmann model. For integral paths, they are the string parameters of the corresponding element in the highest weight crystal (see Kashiwara [26], Littelmann [30]). Likewise, we call $\xi_k(\infty), k \in \mathbb{N}$, the Verma affine string parameters, by analogy with the string parameters of the crystal $B(\infty)$ of Kashiwara associated with the Verma module of highest weight 0. The fact that we are able to study the affine string parameters from the Verma ones is reminiscent of the fact that the highest weight crystals are obtained from $B(\infty)$.

1.6. In conclusion we see that the conditioned Brownian motion in the interval, a priori simple, can be studied here thanks to its links with the infinite dimensional affine Lie algebra $A_1^{(1)}$. It is an example of integrable probability in the sense of Borodin and Petrov [5]. It is interesting to remark that we use an approximation by processes which has no direct relation with group theoretic diffusions, since the dihedral group is not the Weyl group of a semi-simple Lie algebra (when $m \ge 7$). It would be interesting to study the higher rank case $A_n^{(1)}$, $n \ge 2$, which occurs in the analysis of n+1 non colliding Brownian motions on a circle (see Hobson and Werner [21]) or of the eigenvalues of the Brownian motion on SU(n+1). This requires new ideas.

1.7. We always suppose that $0 \le \mu \le 1$.

1.8. We thank Philippe Biane, Persi Diaconis and the referees for their advice.

2. Conditioned Brownian motions

2.1. The conditioned Brownian motion in [0, 1]. We recall some known facts about the Brownian motion conditioned to stay forever in the interval [0, 1]. It was first introduced by Knight [28] when the starting point is in (0, 1)and called the *taboo process*. It was defined there as the limit when t tends to infinity of a standard Brownian motion starting in the interval, conditioned to reach the boundary after time t. To define it rigorously in general, let us first consider the Brownian motion in (0, 1) killed at the boundary. Its generator is $\frac{1}{2}\frac{d^2}{dx^2}$ with Dirichlet boundary conditions. Its maximal eigenvalue is $-\pi^2/2$ with positive eigenvector

$$h(x) = \sin(\pi x),$$

called the ground state. We consider the associated h-Doob process $\{Z_t, t \ge 0\}$. It is the Markov process with transition probability density $q_t(x, y)$ given by

(2.1)
$$q_t(x,y) = \frac{\sin(\pi y)}{\sin(\pi x)} e^{\pi^2 t/2} u_t(x,y),$$

for $x, y \in (0, 1)$, where $u_t(x, y)$ is the transition probability density of the killed Brownian motion. It is the diffusion in [0, 1], with infinitesimal generator given by

(2.2)
$$\frac{1}{2}\frac{d^2}{dx^2} + \pi \cot(\pi x)\frac{d}{dx}.$$

By the reflection principle

(2.3)
$$u_t(x,y) = \sum_{k \in \mathbb{Z}} (p_t(x,y+2k) - p_t(x,-y+2k)),$$

for $x, y \in (0, 1)$, where p_t is the standard heat kernel (see Ito and McKean [23], p.30). Using for instance the Poisson formula (see Bellman [1]) one has,

(2.4)
$$u_t(x,y) = \sum_{n \in \mathbb{N}} \sin(n\pi x) \sin(n\pi y) e^{-\pi^2 n^2 t/2}.$$

One sees that 0 and 1 are entrance non-exit boundaries, by scale function techniques for instance. In other words, (Z_t) can start from the boundaries 0 and 1 and does not touch them at positive time. Let us remark that Z can also be defined by the latitude of the Brownian motion on the 3-dimensional sphere (see Ito and McKean [23], Section 7.15) or by the argument of an eigenvalue of the Brownian motion in SU(2). The behaviour at the boundaries is also clear from these descriptions. When $Z_0 = 0$, the entrance density measure starting from 0 is the limit of $q_t(x, y)$ when x tends to 0; hence, by (2.4)

(2.5)
$$q_t(0,y) = \sin(\pi y) \sum_{n \in \mathbb{Z}} n \sin(n\pi y) e^{-\frac{t}{2}\pi^2(n^2 - 1)},$$

for any $y \in (0, 1)$. The case $Z_0 = 1$ is symmetric.

For $\mu \in [0, 1]$, we write (Z_t^{μ}) for the process Z when $Z_0 = \mu$. and we call it the conditioned Brownian motion process in [0, 1].

2.2. The conditioned space-time Brownian motion in C_{aff} . Let (B_t) be the standard real Brownian motion starting from 0 and

$$B_t^{(\mu)} = (t, B_t + t\mu)$$

be the space-time Brownian motion with drift μ . We now define rigorously the process $A^{(\mu)}(t), t \geq 0$, which is the process $B^{(\mu)}$ conditioned to stay forever in the affine cone

$$C_{\text{aff}} = \{(t, x) \in \mathbb{R}_+ \times \mathbb{R}_+ : 0 < x < t\},\$$

starting from (0,0). It has been introduced and studied in Defosseux [11, 12], or also [10]. Let $\{K_t^{(\mu)}, t \ge 0\}$ be the space-time process $B_t^{(\mu)}$ killed at the boundary of \bar{C}_{aff} . This is the process in the cone with generator $\frac{\partial}{\partial t} + \frac{1}{2} \frac{\partial^2}{\partial x^2} + \mu \frac{\partial}{\partial x}$ and Dirichlet boundary conditions.

We define a theta function φ_{μ} on $\mathbb{R}^*_+ \times \mathbb{R}$ first when $\mu \neq 0, 1$, by

(2.6)
$$\varphi_{\mu}(t,x) = \frac{e^{-\mu x}}{\sin(\mu \pi)} \sum_{k \in \mathbb{Z}} \sinh(\mu(2kt+x)) e^{-2(kx+k^2t)}$$

for $t > 0, x \in \mathbb{R}$. It follows from (2.3) that

(2.7)
$$\varphi_{\mu}(t,x) = \frac{\sqrt{\pi}}{\sqrt{2t}\sin(\mu\pi)} u_{1/t}(\mu,x/t) e^{\frac{t}{2}(\mu-x/t)^2}$$

which implies by (2.4) that $\varphi_{\mu}(t,x) = \varphi_{1-\mu}(t,t-x)$. By continuity,

$$\varphi_0(t,x) = \varphi_1(t,t-x) = \frac{1}{\pi} \sum_{k \in \mathbb{Z}} (2kt+x)e^{-2(kx+k^2t)}$$

Proposition 2.1. The function φ_{μ} is a space-time non negative harmonic function for the killed process $K^{(\mu)}$ on C_{aff} , vanishing on its boundary. *Proof.* The fact that φ_{μ} is harmonic is clear by computation. The boundary condition $\varphi_{\mu}(t,0) = 0$, resp. $\varphi_{\mu}(t,t) = 0$, follows from the change of variable from k to -k, resp. k to -1-k. It is non negative by (2.7).

Let us write $K_t^{(\mu)} = (t, K_t^{\mu})$ and let $w_t^{\mu}((r, x), (t + r, .))$ be the density of K_{t+r}^{μ} given $K_r^{\mu} = x$. Using the Cameron-Martin-Girsanov formula (see [25], Theorem 3.5.1) one has

(2.8)
$$w_t^{\mu}((r,x),(t+r,y)) = e^{\mu(y-x) - \frac{t}{2}\mu^2} w_t^0((r,x),(t+r,y))$$

and, by Defosseux ([11], Proposition 2.2),

$$w_t^{\mu}((r,x),(t+r,y)) = \sum_{k\in\mathbb{Z}} e^{-2(kx+k^2r)} (e^{\mu(2kr+x)} p_t^{\mu}(x+2kr,y) - e^{-\mu(2kr+x)} p_t^{\mu}(-x-2kr,y)),$$

where p_t^{μ} is the standard heat kernel with drift μ . We define the Markov process $\{A^{(\mu)}(t), t \ge 0\}$ in $C_{\text{aff}} \cup \{(0,0)\}$ as the Doob transform of $K^{(\mu)}$ in the following way: $A^{(\mu)}(t) = (t, A_t^{\mu})$ where the distribution of A_{r+t}^{μ} given $A_r^{\mu} = x$ has the density

(2.9)
$$s_t^{\mu}((r,x),(r+t,y)) = \frac{\varphi_{\mu}(r+t,y)}{\varphi_{\mu}(r,x)} w_t^{\mu}((r,x),(r+t,y)),$$

for $(r, x), (r + t, y) \in C_{\text{aff}}, A^{(\mu)}(0) = (0, 0)$, and the entrance density is given by

(2.10)
$$s_t^{\mu}((0,0),(t,y)) = \varphi_{\mu}(t,y)\sin(\frac{y}{t}\pi)e^{-\frac{1}{2t}(y-\mu t)^2}$$

for $(t, y) \in C_{aff}$.

Definition 2.2. The process $\{A^{(\mu)}(t), t \ge 0\}$ is the space-time Brownian motion conditioned to stay in C_{aff} .

Recall that u_t , resp. q_t , is the transition probability density of the killed Brownian motion in [0, 1], resp. of Z^{μ} (see (2.1) and (2.3)).

Lemma 2.3. For $0 < x \le r \le t$ and $0 \le y \le t$,

$$u_{1/r-1/t}(y/t, x/r)e^{-\frac{1}{2t}y^2} = w_{t-r}^0((r, x), (t, y))e^{-\frac{1}{2r}x^2},$$

and for $0 \le x, y \le 1$,

$$q_t(x,y) = \frac{1}{\sqrt{2\pi t}} e^{\pi^2 t/2} \sin(\pi y) e^{-\frac{1}{2t}(y-x)^2} \varphi_x(1/t,y/t).$$

Proof.

$$\begin{split} u_{1/r-1/t}(y/t, x/r) &= \sum_{k \in \mathbb{Z}} (p_{1/r-1/t}(y/t+2k, x/r) - p_{1/r-1/t}(-y/t-2k, x/r)) \\ &= \sum_{k \in \mathbb{Z}} \frac{p_r(0, x)}{p_t(0, y+2kt)} (p_{t-r}(x, y+2kt) - p_{t-r}(x, -y-2kt)) \\ &= \frac{p_r(0, x)}{p_t(0, y)} \sum_{k \in \mathbb{Z}} e^{-2kx-2k^2r} (p_{t-r}(x+2kt, y) - p_{t-r}(-x-2kr, y)), \end{split}$$

gives the first equality. The second one follows from (2.7) and (2.1).

Theorem 2.4. The processes $\{A^{(\mu)}(t), t \ge 0\}$ and $\{(t, tZ^{\mu}_{1/t}), t \ge 0\}$ have the same distribution.

Proof. For $0 < t_1 < \cdots < t_n$, $\mathbb{P}(t_1 Z_{1/t_1}^{\mu} \in dx_1, \cdots, t_n Z_{1/t_n}^{\mu} \in dx_n)$ equals

$$q_{\frac{1}{t_n}}(\mu, \frac{x_n}{t_n})q_{\frac{1}{t_{n-1}}-\frac{1}{t_n}}(\frac{x_n}{t_n}, \frac{x_{n-1}}{t_{n-1}})\cdots q_{\frac{1}{t_1}-\frac{1}{t_2}}(\frac{x_2}{t_2}, \frac{x_1}{t_1})\,dx_1\cdots dx_n.$$

Identity (2.1) and the first identity of Lemma 2.3 imply that

$$\prod_{k=2}^{n} q_{\frac{1}{t_{k-1}} - \frac{1}{t_k}} \left(\frac{x_k}{t_k}, \frac{x_{k-1}}{t_{k-1}}\right) = c e^{\frac{1}{2t_n} x_n^2 - \frac{1}{2t_1} x_1^2} \frac{\sin(\pi x_1/t_1)}{\sin(\pi x_n/t_n)} \prod_{k=2}^{n} w_{t_k - t_{k-1}}^0 \left((t_{k-1}, x_{k-1}), (t_k, x_k)\right),$$

where c is independent of x_1, \dots, x_n . We conclude by using the second identity of Lemma 2.3 for $q_{\frac{1}{t_n}}(\mu, \frac{x_n}{t_n})$ and (2.8),(2.9), and (2.10).

This implies that $A^{(\mu)}(t)$ is really in the interior of the cone for t > 0.

3. Approximation by planar Brownian motions in dihedral cones

In this section we approach the conditioned process $A^{(\mu)}$ by planar Brownian motions conditioned to stay in a dihedral cone.

Definition 3.1. For $m \in \mathbb{N}^*$, the convex dihedral cone C_m , with closure \bar{C}_m , is $C_m = \{(r \cos \theta, r \sin \theta) \in \mathbb{R}^2; r > 0, 0 < \theta < \pi/m\}.$

3.1. Conditioned Brownian motion in a dihedral cone. We will use the results of Biane et al. [3],[4], that we first recall. We consider $V = \mathbb{R}^2$ identified with its dual \tilde{V} , equipped with the usual scalar product $\langle ., . \rangle$. The dihedral group I(m) is the finite group generated by two involutions s_0^m, s_1^m with the only relation $(s_0^m s_1^m)^m = 1$. Notice that the exponent m in s_0^m and s_1^m is not a power but only refers to the angle of the cone. One chooses two pairs $(\alpha_0^m, \tilde{\alpha}_0^m), (\alpha_1^m, \tilde{\alpha}_1^m)$ in $V \times \tilde{V}$, associated with the matrix $(\tilde{\alpha}_i^m(\alpha_j^m))_{0 \le i,j \le 1}$ given by

$$\begin{pmatrix} 2 & -2\cos(\pi/m) \\ -2\cos(\pi/m) & 2 \end{pmatrix}.$$

Namely, one takes

$$\alpha_0^m = (2\sin(\pi/m), -2\cos(\pi/m)), \alpha_1^m = (0, 2),$$

and $\tilde{\alpha}_0^m = \alpha_0^m/2$, $\tilde{\alpha}_1^m = \alpha_1^m/2$. The following two linear reflections s_0^m, s_1^m of \mathbb{R}^2 ,

$$s_i^m(v) = v - \tilde{\alpha}_i^m(v)\alpha_i^m, \ v \in \mathbb{R}^2,$$

generate the group I(m). With our convention $\tilde{\alpha}_i^m(v) = \langle \tilde{\alpha}_i^m, v \rangle$ for $v \in V$. One has

$$C_m = \{ v \in \mathbb{R}^2; \tilde{\alpha}_i^m(v) > 0, i = 0, 1 \},\$$

and \bar{C}_m is a fundamental domain for the action of I(m) on \mathbb{R}^2 . Let $\gamma \in \bar{C}_m$. The conditioned planar Brownian motion in C_m with drift γ is intuitively given by the two-dimensional Brownian motion with drift starting from 0 conditioned to stay in the cone C_m forever. It is rigorously defined in 5.1 of Biane et al. [4] (see Proof of Theorem 5.5). The definition is as follows. Let $\kappa_t^{(\gamma)}$ be the semi group density of the standard planar Brownian motion in C_m with drift γ killed at the boundary. The function

(3.1)
$$\psi_{\gamma}^{m}(v) = \sum_{w \in I(m)} (-1)^{\det(w)} e^{\langle w(\gamma) - \gamma, v \rangle},$$

for $\gamma \in \overline{C}_m$, $v \in \mathbb{R}^2$, is harmonic positive in C_m , null on the boundary. We introduce the Doob transform

$$r_t^{(\gamma)}(v_1, v_2) = \frac{\psi_{\gamma}^m(v_2)}{\psi_{\gamma}^m(v_1)} \kappa_t^{(\gamma)}(v_1, v_2),$$

where $v_1, v_2 \in C_m$. Let

(3.2)
$$h_m(v) = \Im((x+iy)^m)$$

for $v = (x, y) \in \mathbb{R}^2$ (this is the alternating polynomial associated to the dihedral group I(m), equal to the product of the roots, see Dunkl and Xu [14], 6.2.3).

Definition 3.2. The conditioned planar Brownian motion $A_m^{(\gamma)}(t), t \ge 0$, in the dihedral cone C_m with drift $\gamma \in \overline{C}_m$ is defined as follows. It is the continuous Markov process with values in $C_m \cup \{0\}$ with transition probability densities $r_t^{(\gamma)}, t \ge 0$, such that $A_m^{(\gamma)}(0) = 0$, and with entrance probability density at time t proportional to

$$h_m(v)\frac{\psi_{\gamma}^m(v)}{h_m(\gamma)}e^{-\frac{1}{2t}\langle v,v\rangle}\mathbf{1}_{C_m}(v).$$

Here, $\psi_{\gamma}^{m}(v)/h_{m}(\gamma)$ is obtained by analytical continuation when $h_{m}(\gamma) = 0$ (see Lemma 5.11 in [4]). In particular for $\gamma = 0$ the entrance distribution is proportional to

$$h_m^2(v)e^{-\frac{1}{2t}\langle v,v\rangle}\mathbf{1}_{C_m}(v)$$

(see (5.1) of [4]).

Another quick manner to define rigorously $A_m^{(\gamma)}$ is to use Dunkl processes and the approach given by Gallardo and Yor in [17]. One considers the Dunkl Laplacian with multiplicity one associated with I(m). it is given by

$$\tilde{\Delta}f(v) = \Delta f(v) + 2\sum_{\alpha \in \mathbb{R}^+} (\frac{\langle \alpha, \nabla f(v) \rangle}{\langle \alpha, v \rangle} - \frac{f(v) - f(\sigma_\alpha v)}{\langle \alpha, v \rangle^2}),$$

for $f \in C^2(\mathbb{R}^2), v \in \mathbb{R}^2$, where σ_{α} is the orthogonal reflection with respect to $\{v \in \mathbb{R}^2, \langle \alpha, v \rangle = 0\}$, and where R^+ is the set of positive roots. Here Δ is the usual Laplacian and ∇ is the usual gradient. Roesler and Voit ([37], Section 3) show that for any $\gamma \in \mathbb{R}^2$, there exists a càd-làg process $D_{\gamma}(t), t \geq 0$, with generator $\tilde{\Delta}/2$ such that $D_{\gamma}(0) = \gamma$. Let $p : \mathbb{R}^2 \to \bar{C}_m$ be the projection where \bar{C}_m is identified with $\mathbb{R}^2/I(m)$. Then $Y_{\gamma}(t) = p(D_{\gamma}(t))$ is a continuous process called the radial Dunkl process. Its norm is a Bessel process of dimension 2(m+1) ([37], Theorems 4.10, 4.11). According to Gallardo and Yor ([17], Section 3.3), the conditioned Brownian motion in C_m with drift $\gamma \in \bar{C}_m$ starting from 0 is defined as $tY_{\gamma}(1/t)$. Its probability transitions are given by (25) of [17], which shows that it coincides in distribution with $A_m^{(\gamma)}$.

3.2. Pitman representation in a dihedral cone. Let $C_0(\mathbb{R}^2)$ be the set of continuous path $\eta : \mathbb{R}^+ \to \mathbb{R}^2$ such that $\eta(0) = 0$. The following path transforms are introduced in [3].

Definition 3.3. The Pitman transforms $\mathcal{P}^m_{s_i^m}$, i = 0, 1, are defined on $\mathcal{C}_0(\mathbb{R}^2)$ by

$$\mathcal{P}^m_{s^m_i}\eta(t)=\eta(t)-\inf_{0\leq s\leq t}\tilde{\alpha}^m_i(\eta(s))\alpha^m_i,$$

for $\eta \in \mathcal{C}_0(\mathbb{R}^2), t \ge 0$.

For each $i \in \mathbb{N}$, we take *i* modulo 2 when we write $\alpha_i^m, \tilde{\alpha}_i^m, s_i^m, \cdots$. Any $w \in I(m)$ can be written (maybe not uniquely) as

$$w = s_{i_r}^m \cdots s_{i_1}^m$$

for some $r \leq m$. It is called reduced when r is as small as possible.

Theorem 3.4 (Biane et al. [3]). Let $w = s_{i_r}^m \cdots s_{i_1}^m$ be a reduced decomposition of $w \in I(m)$ where $i_1, \cdots, i_r \in \{0, 1\}$. Then

$$\mathcal{P}^m_w := \mathcal{P}^m_{s_{i_r^m}} \cdots \mathcal{P}^m_{s_{i_1^m}}$$

depends only on w and not on its reduced decomposition.

There is a unique longest element w_m in I(m). It has two reduced decompositions, namely

(3.4)
$$w_m = s_{m-1}^m \cdots s_1^m s_0^m = s_m^m \cdots s_2^m s_1^m.$$

Proposition 3.5 ([3]). For any path $\eta \in C_0(\mathbb{R}^2)$, the path $\mathcal{P}_{w_m}^m \eta$ takes values in the closed dihedral cone \bar{C}_m .

The following important theorem is from Biane et al. [4] when $\gamma = 0$.

Theorem 3.6. Let $W^{(\gamma)}$ be the planar Brownian motion with drift $\gamma \in \overline{C}_m$, identity covariance matrix, starting from the origin. The process $\mathcal{P}_{w_m}^m W^{(\gamma)}$ has the same distribution as $A_m^{(\gamma)}$.

Proof. For $0 < t_1 < \cdots < t_n$, and a bounded measurable function $F : \mathbb{R}^n \to \mathbb{R}$, the Cameron–Martin–Girsanov formula (see [25], Theorem 3.5.1) gives

$$\mathbb{E}(F(\mathcal{P}_{w_m}^m W^{(\gamma)}(t_1),\ldots,\mathcal{P}_{w_m}^m W^{(\gamma)}(t_n)))$$

= $\mathbb{E}(F(\mathcal{P}_{w_m}^m W^{(0)}(t_1),\ldots,\mathcal{P}_{w_m}^m W^{(0)}(t_n))e^{-\frac{1}{2}\langle\gamma,\gamma\rangle t_n+\langle\gamma,W_{t_n}^{(0)}\rangle}).$

By Theorems 5.1 and 5.5 of [4] and (3.3) this equals

$$C \mathbb{E}(F(A_m^{(0)}(t_1), \dots, A_m^{(0)}(t_n))e^{-\frac{1}{2}\langle \gamma, \gamma \rangle t_n} \frac{\psi_{\gamma}^m(A_m^{(0)}(t_n))}{h_m(\gamma)h_m(A_m^{(0)}(t_n))}e^{\langle \gamma, A_m^{(0)}(t_n) \rangle}),$$

where C is a constant. This implies the theorem by using Definition 3.2. \Box

3.3. Approximation of the conditioned space time process. We consider the space-time process

$$B_t^{(\mu)} = (t, B_t + t\mu)$$

where B is a standard real Brownian motion starting from 0. Recall that we always suppose that $0 \le \mu \le 1$. For $m \in \mathbb{N}^*$, let $W^{(\frac{m}{\pi},\mu)}$ be the planar Brownian motion with drift $(\frac{m}{\pi}, \mu)$. Hence we can write

$$W^{(\frac{m}{\pi},\mu)}(t) = (\beta_t + t\frac{m}{\pi}, B_t + t\mu),$$

where β is a standard Brownian motions independent of B. Let $\tau_m : \mathbb{R}^2 \to \mathbb{R}^2$ be defined by

$$\tau_m(t,x) = (\frac{\pi t}{m}, x),$$

for $(t, x) \in \mathbb{R}^2$.

Our first approximation theorem is the following. Recall that $A^{(\mu)}$ is the conditioned space-time Brownian motion in C_{aff} . We use the standard definition of convergence in distribution for sequence of continuous processes as in Revuz and Yor ([36], XIII.1).

Theorem 3.7. In distribution, when m tends to $+\infty$, (i) $\tau_m W^{(\frac{m}{\pi},\mu)}$ converges to $B^{(\mu)}$, (ii) $\tau_m \mathcal{P}^m_{w_m} W^{(\frac{m}{\pi},\mu)}$ converges to $A^{(\mu)}$.

The first statement is clear. We will prove the second one after the following proposition. Let $\{Z_t^{\mu}, t \geq 0\}$ be the conditioned Brownian motion in [0, 1] starting from μ .

Proposition 3.8. For $m \in \mathbb{N}^*$, let $X_t^{(m)}, t \ge 0$, be the \mathbb{R}^2 -valued continuous process such that $X_{m^2t}^{(m)}$ is the conditioned planar Brownian motion in the cone C_m , with drift $\gamma_m = m^2 e^{i\pi\mu/m}$. One writes in polar coordinates

$$X_t^{(m)} = R_t^{(m)} \exp i\pi\theta_t^{(m)},$$

 $R_t^{(m)} \ge 0, \theta_t^{(m)} \in [0, 1/m]$. Then the process $m\theta_t^{(m)}, t > 0$, tends to $Z_{1/\pi^2 t}^{\mu}$ and the process $R_t^{(m)}$, $t \ge 0$, tends to t in distribution when m tends to $+\infty$.

Proof. As seen above in Section 3.1, it follows from Gallardo and Yor [17] that $Y_{\gamma_m}(t) = t X_{m^2/t}^{(m)}$ is a radial multidimensional Dunkl process starting from γ_m . Therefore $R_{m^2t}^{(m)}$ is a Bessel process of dimension 2(m+1) with drift m^2 , starting from 0. In other words, one can write

$$(R_{m^2t}^{(m)})^2 = (m^2t + B_t(1))^2 + \sum_{k=2}^{2(m+1)} (B_t(k))^2$$

where $B_t(1), \dots, B_t(2(m+1))$ are independent standard real Brownian motions. Since, for any t > 0,

$$\mathbb{E}(\sup_{0 \le s \le t} \sum_{k=1}^{2(m+1)} B_{s/m^2}(k)^2) \le 8t(m+1)/m^2$$

tends to 0 as m tends to $+\infty$, the process $R_t^{(m)}, t \ge 0$ converges to t in distribution. Using the skew product decomposition of Y_{γ_m} , one can write

$$\pi \theta_{m^2 t}^{(m)} = \sigma^{(m)}(a_t^{(m)}), \text{ with } a_t^{(m)} = \int_t^{+\infty} \frac{1}{(R_{m^2 s}^{(m)})^2} \, ds,$$

where the process $\sigma^{(m)}(t), t \ge 0$, is a solution of the following stochastic differential equation

$$d\sigma^{(m)}(t) = dB_t + m\cot(m\sigma^{(m)}(t))dt$$

Here *B* is a Brownian motion independent of $R^{(m)}$, and $\sigma_0^{(m)} = \mu \pi/m$ (see Demni [13]). One remarks that $Z_t = \frac{m}{\pi} \sigma^{(m)} (\pi^2 t/m^2)$ satisfies

(3.5)
$$dZ_t = d\beta_t + \pi \cot(\pi Z_t) dt$$

where β is another Brownian motion and $Z_0 = \mu$. Therefore Z coincides with Z^{μ} (see (2.2)). When m tends to $+\infty$, $a_{t/m^2}^{(m)}$ is equivalent to $1/tm^2$; hence $m\theta_t^{(m)} = \frac{m}{\pi}\sigma^{(m)}(a_{t/m^2}^{(m)}), t > 0$, converges in distribution to $Z_{1/t\pi^2}^{\mu}$.

Proof of (ii) of Theorem 3.7. Let

$$X_t^{(m)} = \frac{1}{m\pi} \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(\pi^2 t)$$

Since $\frac{1}{m\pi}W_{m^2\pi^2t}^{(\frac{m}{\pi},\mu)}$ is a planar Brownian motion with drift $(m^2, m\pi\mu)$, one sees that $X_{m^2t}^{(m)}$ is the conditioned planar Brownian motion in the cone C_m with a drift equivalent to $m^2 e^{i\pi\mu/m}$ (see Theorem 3.6). One writes its polar decomposition as $X_t^{(m)} = R_t^{(m)} \exp i\pi\theta_t^{(m)}$. Using the continuity of the solution of (3.5) with respect to the initial condition, we see that Proposition 3.8 also holds when the drift γ_m is only equivalent to $m^2 e^{i\pi\mu/m}$. Therefore $m\theta_t^{(m)}$ tends to $Z_{1/t\pi^2}^{\mu}$ and $R_t^{(m)}$ tends to t in distribution. As a consequence the process

$$\tau_m \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(t) = (\pi^2 R_{t/\pi^2}^{(m)} \cos(\pi \theta_{t/\pi^2}^{(m)}), m\pi R_{t/\pi^2}^{(m)} \sin(\pi \theta_{t/\pi^2}^{(m)})), t \ge 0,$$

converges in distribution to $(t, tZ_{1/t}^{\mu})$ which is equal in distribution to $A^{(\mu)}$ by Theorem 2.4.

4. String parameters in the dihedral case

For simplicity of notations, and without loss of generality, one chooses one of the two decompositions of the longest element w_m in the dihedral group I(m), namely,

$$w_m = s_{m-1}^m \cdots s_1^m s_0^m.$$

For $\eta \in \mathcal{C}_0(\mathbb{R}^2)$, $0 \le k \le m-1$ and $0 \le t \le +\infty$, let

$$x_k^m(t) = -\inf_{0 \le s \le t} \tilde{\alpha}_k^m(\mathcal{P}_{s_{k-1}^m}^m \dots \mathcal{P}_{s_1^m}^m \mathcal{P}_{s_0^m}^m \eta(s)).$$

We call

$$x^{m}(t) = (x_{0}^{m}(t), x_{1}^{m}(t), \cdots, x_{m-1}^{m}(t))$$

the string parameters of the path η on [0, t] and we call $x^m(\infty)$ its Verma string parameters. One has, when $t < +\infty$,

(4.1)
$$\mathcal{P}_{w_m}^m \eta(t) = \eta(t) + \sum_{k=0}^{m-1} x_k^m(t) \alpha_k^m$$

Proposition 4.1. When

(4.2)
$$\lim_{s \to +\infty} \tilde{\alpha}_0^m(\eta(s)) = \lim_{s \to +\infty} \tilde{\alpha}_1^m(\eta(s)) = +\infty,$$

then for all k, $x_k^m(\infty) < +\infty$ and $x_k^m(t) = x_k^m(\infty)$ for t large enough.

Proof. Let $\mathfrak{t}_0 = \max\{t \ge 0, \eta(t) \notin C_m\}$, and, for $n \ge 0$,

$$\mathfrak{t}_{n+1} = \max\{t \ge \mathfrak{t}_n, \eta(t) + \sum_{k=0}^n x_k^m(t)\alpha_k \notin C_m\}.$$

We see by induction that $\mathfrak{t}_n < +\infty$ and that for $t > \mathfrak{t}_n$, $x_n^m(t) = x_n^m(\infty)$. \Box

Let, for $1 \leq k < m$,

$$a_k^m = \sin(k\pi/m).$$

Definition 4.2. The cone Γ_m in \mathbb{R}^m is defined as

$$\Gamma_m = \{ (x_0, \cdots, x_{m-1}) \in \mathbb{R}^m; \frac{x_1}{a_1^m} \ge \frac{x_2}{a_2^m} \ge \cdots \ge \frac{x_{m-1}}{a_{m-1}^m} \ge 0, \ x_0 \ge 0 \}.$$

For $\lambda \in \overline{C}_m$, the polytope $\Gamma_m(\lambda)$ is

$$\Gamma_m(\lambda) = \{(x_0, \cdots, x_{m-1}) \in \Gamma_m; 0 \le x_r \le \tilde{\alpha}_r^m (\lambda - \sum_{n=r+1}^{m-1} x_n \alpha_n^m), 0 \le r \le m-1\}.$$

Remark the particular role of x_0 . The following proposition is proved in [4], Propositions 4.4 and 4.7.

Proposition 4.3. For t > 0, the set of string parameters of paths on [0, t] is Γ_m .

Let $W^{(\gamma)}$ be a planar Brownian motion with drift γ in \overline{C}_m starting from 0.

Definition 4.4. We define $\xi^m(t)$ as the string parameters of $W^{(\gamma)}$ on [0, t] and $\xi^m(\infty)$ as its Verma string parameters.

Notice that, a.s., when $\gamma \in C_m$, $W^{(\gamma)}$ satisfies (4.2) by the law of large numbers, since $W_t^{(\gamma)}/t$ converges to an element of C_m when $t \to +\infty$. So its Verma string parameters are finite by Proposition 4.1. We will need the following proposition.

Proposition 4.5. For $\nu = 0$ and γ , let $\mathcal{F}_t^{\nu} = \sigma(\mathcal{P}_{w_m}^m W^{(\nu)}(s), s \leq t)$ and ψ : $\mathcal{C}_0([0,t], \mathbb{R}^2) \to \mathbb{R}$ be a bounded measurable function. There exists a function $\varphi : \mathbb{R}^2 \to \mathbb{R}$ such that, for both $\nu = 0$ and $\nu = \gamma$, a.s.,

$$\varphi(\mathcal{P}_{w_m}^m W^{(\nu)}(t)) = \frac{\mathbb{E}(\psi(W_s^{(\nu)}, 0 \le s \le t)e^{\langle \gamma - \nu, W_t^{(\nu)} \rangle} | \mathcal{F}_t^{\nu})}{\mathbb{E}(e^{\langle \gamma - \nu, W_t^{(\nu)} \rangle} | \mathcal{F}_t^{\nu})}.$$

Proof. By Biane et al. [4], Theorem 5.5, we know that when $\nu = 0$, the conditional expectation is $\sigma(\mathcal{P}_{w_m}^m W^{(\nu)}(t))$ measurable. So the formula defines φ in this case. We have to prove that for this φ the formula also holds when $\nu = \gamma$, namely that

(4.3)
$$\varphi(\mathcal{P}_{w_m}^m W^{(\gamma)}(t)) = \mathbb{E}(\psi(W_s^{(\gamma)}, 0 \le s \le t) | \mathcal{F}_t^{\gamma})$$

Let $F : \mathcal{C}_0([0,t],\mathbb{R}^2) \to \mathbb{R}$ be a bounded measurable function. We define, for $\nu = 0, \gamma,$

$$U_{\nu} = F(\mathcal{P}_{w_m}^m W_s^{(\nu)}, 0 \le s \le t).$$

By Cameron-Martin-Girsanov,

$$\mathbb{E}(U_{\gamma}\psi(W_s^{(\gamma)}, 0 \le s \le t)) = \mathbb{E}(U_0\psi(W_s^{(0)}, 0 \le s \le t)e^{\langle \gamma, W_t^{(0)} \rangle}) / \mathbb{E}(e^{\langle \gamma, W_t^{(0)} \rangle}).$$

Besides, since U_0 is \mathcal{F}_t^0 -measurable, using the definition of φ ,

$$\mathbb{E}(U_0\psi(W_s^{(0)}, 0 \le s \le t)e^{\langle \gamma, W_t^{(0)} \rangle}) = \mathbb{E}(U_0\mathbb{E}(\psi(W_s^{(0)}, 0 \le s \le t)e^{\langle \gamma, W_t^{(0)} \rangle} | \mathcal{F}_t^0))$$
$$= \mathbb{E}(U_0\varphi(\mathcal{P}_{w_m}^m W^{(0)}(t))\mathbb{E}(e^{\langle \gamma, W_t^{(0)} \rangle} | \mathcal{F}_t^0))$$
$$= \mathbb{E}(U_0\varphi(\mathcal{P}_{w_m}^m W^{(0)}(t))e^{\langle \gamma, W_t^{(0)} \rangle}).$$

Hence, by CMG again,

$$\mathbb{E}(U_{\gamma}\psi(W_{s}^{(\gamma)}, 0 \le s \le t)) = \mathbb{E}(U_{\gamma}\varphi(\mathcal{P}_{w_{m}}^{m}W^{(\gamma)}(t)))$$

which gives (4.3).

Theorem 4.6. Let $\mathcal{E}^m = (\mathcal{E}_0^m, \dots, \mathcal{E}_{m-1}^m)$ be a random vector in \mathbb{R}^m whose components \mathcal{E}_k^m are independent exponential random variables with respective parameters $\gamma_k^m = \langle \gamma, \alpha_k^m \rangle$.

(i) When $\gamma \in \overline{C}_m$, for any t > 0, conditionally on the σ -algebra

$$\sigma(\mathcal{P}^m_{w_m}W^{(\gamma)}(s), s \le t),$$

the random vector $(\xi_0^m(t), \ldots, \xi_{m-1}^m(t))$ is distributed as the vector \mathcal{E}^m conditioned

on the event $\{\mathcal{E}^m \in \Gamma_m(\mathcal{P}^m_{w_m}W^{(\gamma)}(t))\}$. (ii) When $\gamma \in C_m$, the random vector $(\xi_0^m(\infty), \ldots, \xi_{m-1}^m(\infty))$ is distributed as \mathcal{E}^m conditioned on the event $\{\mathcal{E}^m \in \Gamma_m\}$.

Proof. Let $F : \mathbb{R}^m \to \mathbb{R}$ be a bounded and continuous function. Let us denote by $\xi^{0,m}(t)$ the string parameters of $W^{(0)}$. It is shown in Biane et al. [4], Theorem 5.2, that the distribution of $\xi^{0,m}(t)$ conditionally on \mathcal{F}_t^0 is the normalized Lebesgue measure on $\Gamma_m(\mathcal{P}^m_{w_m}W^{(0)}(t))$. Using that, from (4.1),

$$W_t^{(0)} = \mathcal{P}_{w_m}^m W^{(0)}(t) - \sum_{k=0}^{m-1} \xi_k^{0,m}(t) \alpha_k^m,$$

one has,

$$\mathbb{E}(F(\xi^{0,m}(t))e^{\langle\gamma,W_t^{(0)}\rangle}|\mathcal{F}_t^0) = e^{\langle\gamma,\mathcal{P}_{w_m}^m W_t^{(0)}\rangle} \mathbb{E}(F(\xi^{0,m}(t))e^{-\langle\gamma,\sum_{k=0}^{m-1}\xi_k^{0,m}(t)\alpha_k^m\rangle}|\mathcal{F}_t^0)$$
$$= \frac{1}{c_t} \int \mathbf{1}_{\Gamma_m(\mathcal{P}_{w_m}^m W^{(0)}(t))}(x)F(x)e^{-\langle\gamma,x\rangle} dx$$

where dx is the Lebesgue measure on \mathbb{R}^m and c_t is the Lebesgue measure of $\Gamma_m(\mathcal{P}^m_{w_m}W^{(0)}(t)))$. If we apply Proposition 4.5 to the function $\psi(W_s^{(\gamma)}, 0 \le s \le t) = F(\xi^m(t))$ we see that

$$\mathbb{E}(F(\xi^m(t))|\mathcal{F}_t^{\gamma}) = \frac{\int \mathbf{1}_{\Gamma_m(\mathcal{P}_{w_m}^m W^{(\gamma)}(t))}(x)F(x)e^{-\langle \gamma, x \rangle} \, dx}{\int \mathbf{1}_{\Gamma_m(\mathcal{P}_{w_m}^m W^{(\gamma)}(t))}(x)e^{-\langle \gamma, x \rangle} \, dx},$$

which proves (i). We now suppose that $\gamma \in C_m$, so

$$\begin{split} \mathbb{E}(F(\xi^m(\infty))) &= \lim_{t \to +\infty} \mathbb{E}(F(\xi^m(t))) \\ &= \lim_{t \to +\infty} \mathbb{E}(\mathbb{E}(F(\xi^m(t))|\mathcal{F}_t^{\gamma})) \\ &= \lim_{t \to +\infty} \mathbb{E}(\frac{\int \mathbf{1}_{\Gamma_m}(\mathcal{P}_{w_m}^m W^{(\gamma)}(t))(x)F(x)e^{-\langle \gamma, x \rangle} \, dx}{\int \mathbf{1}_{\Gamma_m}(\mathcal{P}_{w_m}^m W^{(\gamma)}(t))(x)e^{-\langle \gamma, x \rangle} \, dx}) \\ &= \frac{\int \mathbf{1}_{\Gamma_m}(x)F(x)e^{-\langle \gamma, x \rangle} \, dx}{\int \mathbf{1}_{\Gamma_m}(x)e^{-\langle \gamma, x \rangle} \, dx}, \end{split}$$

since $\mathcal{P}_{w_m}^m W^{(\gamma)}(t)/t$ converges to γ as t tends to $+\infty$ and thus $\Gamma_m(\mathcal{P}_{w_m}^m W^{(\gamma)}(t))$ tends to Γ_m for $\gamma \in C_m$, a.s.. This proves (ii).

A direct consequence of the theorem is the following:

Corollary 4.7. The conditional distribution of $\xi^m(t)$ given $\sigma(\mathcal{P}^m_{w_m}W^{(\gamma)}(s), s \leq t)$ is the one of $\xi^m(\infty)$ given $\{\xi^m(\infty) \in \Gamma_m(\mathcal{P}^m_{w_m}W^{(\gamma)}(t))\}.$

Corollary 4.8. We suppose that $\gamma \in C_m$. Let $\xi^m(\infty)$ be the Verma string parameters of $W^{(\gamma)}$. If we let

$$(\xi_0^m(\infty), \frac{\xi_1^m(\infty)}{a_1^m} - \frac{\xi_2^m(\infty)}{a_2^m}, \frac{\xi_2^m(\infty)}{a_2^m} - \frac{\xi_3^m(\infty)}{a_3^m}, \cdots, \frac{\xi_{m-1}^m(\infty)}{a_{m-1}^m})$$

equals to

$$(\frac{\varepsilon_0^m}{\gamma_0^m}, \frac{\varepsilon_1^m}{\gamma_1^m a_1^m}, \frac{\varepsilon_2^m}{\gamma_1^m a_1^m + \gamma_2^m a_2^m}, \cdots, \frac{\varepsilon_{m-1}^m}{\gamma_1^m a_1^m + \cdots + \gamma_{m-1}^m a_{m-1}^m})$$

then the ε_n^m , $n = 0, 1, \dots, m-1$ are independent exponential random variables with parameter 1.

Proof. This follows from Theorem 4.6 since, on \mathbb{R}^m_+ ,

$$\mathbb{P}(\xi_0^m(\infty) \in dx_0, \cdots, \xi_{m-1}^m(\infty) \in dx_{m-1})$$

$$= Ce^{-\gamma_0^m x_0} dx_0 \prod_{k=2}^m e^{-\gamma_{k-1}^m x_{k-1}} \mathbb{1}_{\{\frac{x_{k-1}}{a_{k-1}^m} \ge \frac{x_k}{a_k^m}\}} dx_{k-1}$$

$$= Ce^{-\gamma_0^m x_0} dx_0 \prod_{k=2}^m e^{-(\gamma_1^m a_1^m + \dots + \gamma_{k-1}^m a_{k-1}^m)(\frac{x_{k-1}}{a_{k-1}^m} - \frac{x_k}{a_k^m})} \mathbb{1}_{\{\frac{x_{k-1}}{a_{k-1}^m} \ge \frac{x_k}{a_k^m}\}} dx_{k-1}$$

where by convention $x_m/a_m^m = 0$ and C is a normalizing constant.

Notice that this is similar to Renyi's representation of order statistics [35].

5. Affine string parameters of the space-time Brownian motion

5.1. **Pitman and Lévy transforms.** The infinite dihedral group $I(\infty)$ is the infinite group generated by two involutions s_0, s_1 with no relation. In $V = \mathbb{R}^2$ identified with its dual \tilde{V} , let $(\alpha_0, \tilde{\alpha}_0), (\alpha_1, \tilde{\alpha}_1)$ in $V \times \tilde{V}$ be given by

$$\begin{cases} \alpha_0 = (0, -2), \\ \alpha_1 = (0, 2), \\ \tilde{\alpha}_0 = (1, -1), \\ \tilde{\alpha}_1 = (0, 1). \end{cases}$$

The matrix $(\tilde{\alpha}_i(\alpha_j))_{0 \le i,j \le 1}$ is the Cartan matrix

$$\begin{pmatrix} 2 & -2 \\ -2 & 2 \end{pmatrix}.$$

The two linear reflections s_0, s_1 of \mathbb{R}^2 defined by

 $s_i(v) = v - \tilde{\alpha}_i(v)\alpha_i,$

for $v \in \mathbb{R}^2$, generate the group $I(\infty)$. Notice that s_0 is a non orthogonal reflection. For each $k \in \mathbb{N}$, when we write $\alpha_k, \tilde{\alpha}_k, s_k, \cdots$, we take k modulo 2, as above. Thus $\alpha_k = (-1)^k \alpha_0$ (but $\tilde{\alpha}_k \neq (-1)^k \tilde{\alpha}_0$). Notice that

$$C_{\text{aff}} = \{(t, x) \in \mathbb{R}^2; 0 < x < t\} = \{v \in V; \tilde{\alpha}_i(v) > 0, i = 0, 1\}.$$

We define the Pitman transforms \mathcal{P}_{s_0} and \mathcal{P}_{s_1} on the paths of $\mathcal{C}_0(\mathbb{R}^2)$ by

$$\mathcal{P}_{s_i}\eta(t) = \eta(t) - \inf_{0 \le s \le t} \tilde{\alpha}_i(\eta(s))\alpha_i,$$

where $\eta \in \mathcal{C}_0(\mathbb{R}^2), i = 0, 1$. Let $T : \mathbb{R}^2 \to \mathbb{R}^2$ be the involutive transformation T(t, x) = (t, t - x). Then

(5.1)
$$\mathcal{P}_{s_1}T = T\mathcal{P}_{s_0}$$

We will use mainly space-time paths, i.e. paths which can be written as $\eta(t) = (t, \varphi(t)), t \ge 0$, where $\varphi(t) \in \mathbb{R}$. In this case,

$$\mathcal{P}_{s_0}\eta(t) = (t,\varphi(t) + 2\inf_{s \le t}(s - \varphi(s)))$$
$$\mathcal{P}_{s_1}\eta(t) = (t,\varphi(t) - 2\inf_{s \le t}\varphi(s)).$$

One recognizes in the second component the transforms $\mathcal{P}_0, \mathcal{P}_1$ defined in Section 1.2.

5.2. Affine string parameters. We consider the space-time process $B_t^{(\mu)} = (t, B_t + t\mu)$ where B is a standard real Brownian motion starting from 0. We define, for $t \leq \infty$, $\xi(t) = \{\xi_k(t), k \geq 0\}$ by

$$\xi_k(t) = -\inf_{0 \le s \le t} \tilde{\alpha}_k(\mathcal{P}_{s_{k-1}} \dots \mathcal{P}_{s_1} \mathcal{P}_{s_0} B^{(\mu)}(s)).$$

Hence

$$\xi_0(t) = -\inf_{0 \le s \le t} \tilde{\alpha}_0(B^{(\mu)}(s)).$$

Definition 5.1. We call $\xi(t) = \{\xi_k(t), k \ge 0\}$ the affine string parameters of $B^{(\mu)}$ on [0, t], and $\xi(\infty) = \{\xi_k(\infty), k \ge 0\}$ its Verma affine string parameters.

We use a terminology inspired by the Kac-Moody affine algebra $A_1^{(1)}$ (see in particular Kac [24], Kashiwara [26], and Section 9 below). One has, for any $n \ge 0$,

(5.2)
$$\mathcal{P}_{s_n} \cdots \mathcal{P}_{s_0} B^{(\mu)}(t) = B_t^{(\mu)} + \sum_{k=0}^n \xi_k(t) \alpha_k.$$

When $\mu \neq 0, 1$, by the law of large numbers, a.s., $\lim_{t \to +\infty} B_t^{(\mu)}/t = (1, \mu) \in C_{\text{aff}}$, hence

$$\lim_{t \to +\infty} \tilde{\alpha}_0(B_t^{(\mu)}) = \lim_{t \to +\infty} \tilde{\alpha}_1(B_t^{(\mu)}) = +\infty,$$

and by the analogue of Proposition 4.1, $\xi_k(\infty) < +\infty$ for each $k \ge 0$. Recall that, for $v = (t, x) \in \mathbb{R}^2$,

$$\tau_m v = (\frac{\pi t}{m}, x).$$

We will use that for $v \in \mathbb{R}^2$,

(5.3)
$$\tau_m v = \left(\frac{\pi}{m\sin\frac{\pi}{m}} (\tilde{\alpha}_0^m(v) + \tilde{\alpha}_1^m(v)\cos\frac{\pi}{m}), \tilde{\alpha}_1^m(v)\right),$$

so the asymptotics of $\tau_m v$ and $(\tilde{\alpha}_0^m(v) + \tilde{\alpha}_1^m(v), \tilde{\alpha}_1^m(v))$ are the same as m tends to $+\infty$. It is easy to see that,

Lemma 5.2. For i = 0, 1,

$$\lim_{m \to +\infty} \tau_m \alpha_i^m = \alpha_i,$$
$$\lim_{m \to +\infty} \tau_m^{-1} \tilde{\alpha}_i^m = \tilde{\alpha}_i,$$
$$\lim_{m \to +\infty} \tau_m \circ s_i^m \circ \tau_m^{-1} = s_i$$

Definition 5.3. The dihedral highest weight process is

$$\Lambda_m^{(\mu)}(t) = \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(t), \ t \ge 0.$$

By Theorem 3.7, $\tau_m W^{(\frac{m}{\pi},\mu)}$ tends in distribution to $B^{(\mu)}$ and $\tau_m \Lambda_m^{(\mu)}$ tends to $A^{(\mu)}$. When $\mu \neq 0, 1$, we will always suppose that m is large enough to have $(\frac{m}{\pi}, \mu) \in C_m$. This ensures that each $\xi_k^m(\infty)$ is finite.

Proposition 5.4. Almost surely, for all $t \ge 0$ and $k \in \mathbb{N}$,

$$\lim_{m \to +\infty} \xi_k^m(t) = \xi_k(t),$$

and when $\mu \neq 0, 1$,

$$\lim_{m \to +\infty} \xi_k^m(\infty) = \xi_k(\infty).$$

Proof. Let $\eta_0^m = W^{(\frac{m}{\pi},\mu)}$, $\eta_0 = B^{(\mu)}$ and, for $k \ge 1$,

$$\eta_k^m = \mathcal{P}_{s_{k-1}^m}^m \dots \mathcal{P}_{s_0^m}^m \eta_0^m, \, \eta_k = \mathcal{P}_{s_{k-1}} \dots \mathcal{P}_{s_0} \eta_0.$$

We prove by induction on $k \in \mathbb{N}$, that for any T > 0, a.s., uniformly on [0, T],

$$\lim_{m \to +\infty} \tau_m \eta_k^m(t) = \eta_k(t)$$
$$\lim_{m \to +\infty} \xi_k^m(t) = \xi_k(t).$$

This is clear for k = 0. If it holds for k - 1, then, by Lemma 5.2,

$$\tau_m \eta_k^m(t) = \tau_m \mathcal{P}_{s_{k-1}^m}^m \eta_{k-1}^m(t) = \tau_m \eta_{k-1}^m(t) + \xi_{k-1}^m(t) \tau_m \alpha_{k-1}^m$$

converges uniformly to $\eta_k(t)$. Since one can write

$$\xi_k^m(t) = -\inf_{0 \le s \le t} \tilde{\alpha}_k^m(\eta_k^m(s)) = -\inf_{0 \le s \le t} \tau_m^{-1}(\tilde{\alpha}_k^m)(\tau_m \eta_k^m(s)),$$

 $\xi_k^m(t)$ converges to $\xi_k(t)$.

We now suppose that $\mu \neq 0, 1$. First we notice that, a.s., for any C > 0, there is a time $t_C > 0$ such that, for any $m \in \mathbb{N}$,

(5.4)
$$\tilde{\alpha}_i^m(\eta_0^m(t)) \ge C,$$

when $t \ge t_C$, i = 0, 1. Let us show that a.s., for t large enough (independent of m), $\xi_k^m(\infty) = \xi_k^m(t)$ for all $m \in \mathbb{N}$, large enough. This will imply that $\xi_k^m(\infty)$ tends to $\xi_k(\infty)$. Since $\tilde{\alpha}_k(\eta_k^m(0)) = 0$, it is enough to show that almost surely, there exists $T_k > 0$ such that for $t \ge T_k$, for all $m \in \mathbb{N}$, $\tilde{\alpha}_k^m(\eta_k^m(t)) \ge 0$. This follows from (5.4) and from the fact, easily proved by induction on k, that almost surely, there is a $C_k > 0$ such that, for all $t \ge 0, m \in \mathbb{N}$, $\|\eta_k^m(t) - \eta_0^m(t)\| \le C_k$. \Box

Theorem 5.5. For $\mu \neq 0, 1$, one can write $\xi_0(\infty) = \varepsilon_0/2(1-\mu)$, and for all $k \geq 1$,

$$\frac{\xi_k(\infty)}{k} = \sum_{n=k}^{+\infty} \frac{2\varepsilon_n}{n(n+1) + (1-2\mu)\nu_n}$$

where $\varepsilon_n, n \ge 0$, are independent exponential random variables with parameter 1, and where $\nu_n = n$ when n is even and $\nu_n = -(n+1)$ when n is odd.

Proof. We consider the independent exponentially distributed random variables ε_n^m of Corollary 4.8 for the drift $\gamma = (m/\pi, \mu)$. We have $\varepsilon_0^m = \gamma_0^m \xi_0^m(\infty)$ and for $1 \le k < m$,

$$\varepsilon_k^m = (\sum_{n=1}^k \gamma_n^m a_n^m) (\frac{\xi_k^m(\infty)}{a_k^m} - \frac{\xi_{k+1}^m(\infty)}{a_{k+1}^m}).$$

One has

$$\gamma_{2k}^m = \langle (\frac{m}{\pi}, \mu), \alpha_0^m \rangle = 2(\frac{m}{\pi}\sin(\frac{\pi}{m}) - \mu\cos(\frac{\pi}{m})),$$
$$\gamma_{2k+1}^m = \langle (\frac{m}{\pi}, \mu), \alpha_1^m \rangle = 2\mu.$$

When $m \to +\infty$, $\gamma^m_{2k} \to 2(1-\mu)$ and $a^m_k/m \to k\pi$ hence

$$\lim_{m \to +\infty} m \sum_{n=1}^{k} \gamma_n^m a_n^m = \frac{1}{2} (n(n+1) + (1-2\mu)\nu_n).$$

So it follows from Proposition 5.4 that the variables ε_n defined by, a.s., for $n \ge 1$,

(5.5)
$$\varepsilon_n := \lim_{m \to +\infty} \varepsilon_k^m = \frac{n(n+1) + (1-2\mu)\nu_n}{2} (\frac{\xi_n(\infty)}{n} - \frac{\xi_{n+1}(\infty)}{n+1})$$

and $\varepsilon_0 = 2(1-\mu)\xi_0(\infty)$ are independent exponentially distributed random variables with parameter 1. By definition of the cone Γ_m , the sequence $\xi_k^m(\infty)/a_k^m$, $k \ge 1$, is non negative and decreasing. Therefore $\xi_k(\infty)/k, k \ge 1$, is also non

negative and decreasing. We denote by S its limit, which is a non negative random variable. Thus one has, by (5.5),

$$\frac{\xi_k(\infty)}{k} = \sum_{n=k}^{+\infty} \frac{2\varepsilon_n}{n(n+1) + (1-2\mu)\nu_n} + S.$$

Moreover Fatou's lemma implies that

(5.6)
$$\mathbb{E}(\frac{\xi_k(\infty)}{k}) \le \liminf_m \mathbb{E}(\frac{\xi_k^m(\infty)}{k}) = \liminf_m \frac{\pi}{m} \sum_{n=k}^{m-1} \frac{1}{a_1^m \gamma_1^m + \dots + a_n^m \gamma_n^m}.$$

Let $c = \max\{1/\gamma_i^m, m \in \mathbb{N}, i = 0, 1\}$. Since $0 \le 2t \le \sin \pi t$ when $0 \le t \le 1/2$, one has, for $k \le [m/2]$,

$$\sum_{n=1}^{k} a_n^m \gamma_n^m \ge \frac{1}{c} \sum_{n=1}^{k} \sin(\frac{n\pi}{m}) \ge \sum_{n=1}^{k} \frac{2n}{cm} \ge \frac{k^2}{cm}$$

and, for k > [m/2],

$$\sum_{n=1}^k a_n^m \gamma_n^m \ge \sum_{n=1}^{\lfloor m/2 \rfloor + 1} a_n^m \gamma_n^m \ge \frac{m}{4c}.$$

Therefore, for $N \leq m - 1$,

$$\frac{\pi}{m} \sum_{n=N}^{m-1} \frac{1}{a_1^m \gamma_1^m + \dots + a_n^m \gamma_n^m} \le c\pi \left(\sum_{n=N}^{[m/2]} \frac{1}{n^2} + \sum_{n=[m/2]+1}^m \frac{4}{m^2}\right)$$

is as small as we want for N large enough. As one has for every fixed N > 0,

$$\lim_{m} \frac{\pi}{m} \sum_{n=k}^{N} \frac{1}{a_1^m \gamma_1^m + \dots + a_n^m \gamma_n^m} = \sum_{n=k}^{N} \frac{2}{n(n+1) + (1-2\mu)\nu_n},$$

the lim inf in (5.6) equals $\sum_{n=k}^{+\infty} \frac{2}{n(n+1)+(1-2\mu)\nu_n}$ and therefore

$$\mathbb{E}(\frac{\xi_k(\infty)}{k}) \le \sum_{n=k}^{+\infty} \frac{2}{n(n+1) + (1-2\mu)\nu_n}$$

This proves that $\mathbb{E}(S) = 0$, which implies that S = 0 a.s. and finishes the proof.

An important result for us is the following (it is maybe unexpected because the string coordinates $x_k(t)$ for C^1 paths vanish for k large enough, see Proposition 7.7).

Theorem 5.6. For $\mu \neq 0, 1$, almost surely, $\lim_{k \to +\infty} \xi_k(\infty) = 2$.

Proof. Since

$$\mathbb{E}(\xi_k(\infty)) = \sum_{n=k}^{+\infty} \frac{2k}{n(n+1) + (1-2\mu)\nu_n},$$

there is a C > 0 such that

$$|\mathbb{E}(\xi_k(\infty)) - \sum_{n=k}^{+\infty} \frac{2k}{n(n+1)}| \le Ck \sum_{n=k}^{+\infty} \frac{1}{n^3}.$$

Since $\sum_{n=k}^{+\infty} \frac{k}{n(n+1)} = 1$, $\mathbb{E}(\xi_k(\infty))$ tends to 2, as $k \to +\infty$. To prove that $\xi_k(\infty)$ tends to 2 almost surely it suffices to remark that if X_n is a sequence of i.i.d. centered random variables with a finite fourth moment then, almost surely, $Z_k = k \sum_{n=k}^{+\infty} X_n/n^2$ tends to 0. Indeed, it is easy to see that

$$\mathbb{E}(Z_k^4) = k^4 \mathbb{E}(X_1^4) \sum_{k \le n} \frac{1}{n^8} + 2k^4 \mathbb{E}(X_1^2)^2 \sum_{k \le n} \frac{1}{n^4} \sum_{n < m} \frac{1}{m^4}$$

For p > 1, $\sum_{n > k} n^{-p} \le C_p k^{1-p}$ for some $C_p > 0$, hence for some C > 0

$$\mathbb{E}(Z_k^4) \le Ck^4(\frac{1}{k^7} + \frac{1}{k^6}).$$

This shows that $\mathbb{E}(\sum_{k} Z_{k}^{4}) < +\infty$ which implies that $Z_{k} \to 0$, a.s.

5.3. Affine Verma weight $L^{(\mu)}(\infty)$. We consider, for $\mu \neq 0, 1$, and $k \geq 1$,

$$M_k(\infty) = \frac{1}{2}\xi_k(\infty)\alpha_k + \sum_{n=0}^{k-1}\xi_n(\infty)\alpha_n,$$

where $\xi(\infty) = \{\xi_k(\infty), k \ge 0\}$ are the affine Verma string parameters of $B^{(\mu)}$. Notice that, more simply, since $\alpha_n = (0, (-1)^{n+1}2)$,

$$M_k(\infty) = (0, (-1)^{k+1}\xi_k(\infty) + 2\sum_{n=0}^{k-1} (-1)^{n+1}\xi_n(\infty)).$$

The notation with the α_n 's may be strange for the reader; it is explained by its natural interpretation in $A_1^{(1)}$ (see 9.1).

Proposition 5.7. Let $\mu \neq 0, 1$. When k goes to infinity, $M_k(\infty)$ converges almost surely and in L^2 to

$$L^{(\mu)}(\infty) = \frac{1}{2} \sum_{n=0}^{\infty} \left(\frac{\varepsilon_{2n}}{n+1-\mu} \alpha_0 + \frac{\varepsilon_{2n+1}}{n+\mu} \alpha_1 \right).$$

Proof. One has the key relation

$$M_{2p+2}(\infty) - M_{2p}(\infty) = (0, -\xi_{2p}(\infty) + 2\xi_{2p+1}(\infty) - \xi_{2p+2}(\infty))$$
$$= (0, \frac{\varepsilon_{2p+1}}{p+\mu} - \frac{\varepsilon_{2p}}{p+1-\mu}).$$

Hence $\{M_{2p}(\infty) - \mathbb{E}(M_{2p}(\infty)), p \geq 0\}$ is a martingale bounded in L_2 . Since $\mathbb{E}(M_{2p}(\infty))$ converges, this shows that the sequence $\{M_{2p}(\infty), p \in \mathbb{N}\}$ converges a.s. and in L^2 and that its limit is $L^{(\mu)}(\infty)$. As

$$M_{2p+1}(\infty) = M_{2p}(\infty) + \frac{1}{2}(\xi_{2p}(\infty) - \xi_{2p+1}(\infty))\alpha_0$$

and $\xi_k(\infty)$ tends to 2 by Theorem 5.6, the sequence $\{M_{2p+1}(\infty), p \in \mathbb{N}\}$ has the same limit, namely $L^{(\mu)}(\infty)$.

Notice that $L^{(\mu)}(\infty)$ is not the limit of the series $\sum_{k=0}^{n} \xi_k(\infty) \alpha_k$ since $\xi_k(\infty)$ tends to 2 (Theorem 5.6).

Definition 5.8. We call $L^{(\mu)}(\infty)$ the affine Verma weight of $B^{(\mu)}$.

The terminology is due to the fact, as we will see in Section 9, that $-L^{(\mu)}(\infty)$ plays the role of an asymptotic weight in the Verma module of highest weight 0 in the case of the affine Lie algebra $A_1^{(1)}$, modulo δ . In coordinates, $L^{(\mu)}(\infty)$ can be written as $(0, D^{\mu}(\infty))$ where

(5.7)
$$D^{\mu}(\infty) = \sum_{n=0}^{\infty} \left(\frac{\varepsilon_{2n+1}}{n+\mu} - \frac{\varepsilon_{2n}}{n+1-\mu}\right).$$

For a sequence $x = (x_k) \in \mathbb{R}^{\mathbb{N}}_+$, we let

(5.8)
$$\omega(x) = \lim_{n \to +\infty} \sum_{k=0}^{n-1} x_k \alpha_k + \frac{1}{2} x_n \alpha_n,$$

and write $\omega(x) \in \mathbb{R}^2$ when this limit exists in \mathbb{R}^2 . When $\mu \neq 0, 1$,

$$\omega(\xi(\infty)) = L^{(\mu)}(\infty)$$

Definition 5.9. One defines, for $\lambda \in \bar{C}_{aff}$,

$$\Gamma = \{ x = (x_k) \in \mathbb{R}^{\mathbb{N}} : \frac{x_k}{k} \ge \frac{x_{k+1}}{k+1} \ge 0, \text{ for all } k \ge 1, x_0 \ge 0, \, \omega(x) \in \mathbb{R}^2 \},$$

$$\Gamma(\lambda) = \{ x \in \Gamma : x_k \le \tilde{\alpha}_k (\lambda - \omega(x) + \sum_{i=0}^k x_i \alpha_i), \text{ for every } k \ge 0 \}.$$

We remark that $x \in \Gamma$ is in $\Gamma(\lambda)$ if and only if for all $k \ge 0$,

(5.9)
$$\tilde{\alpha}_k(\omega(x) - \sum_{i=0}^{k-1} x_i \alpha_i - \frac{1}{2} x_k \alpha_k) \le \tilde{\alpha}_k(\lambda).$$

Notice the occurence of the coefficient 1/2 here which will explain the correction term in the representation theorem. The same coefficient already occurs in the description of the crystal $B(\lambda)$ of $A_1^{(1)}$ given by Nakashima [32] which is the discrete analogue of $\Gamma(\lambda)$. In a sense Γ will have the role of a continuous Verma crystal and $\omega(x)$ has the role of the weight of $x \in \Gamma$. In the same way $\Gamma(\lambda)$ will be a kind of continuous crystal of highest weight λ .

Proposition 5.10. When $\mu \neq 0, 1$, the Verma string parameters $\xi(\infty)$ of $B^{(\mu)}$ are a.s. in $\cup \{\Gamma(\lambda), \lambda \in C_{aff}\}$ and for each $\lambda \in C_{aff}$, $\mathbb{P}(\xi(\infty) \in \Gamma(\lambda)) > 0$.

Proof. Since $\omega(\xi(\infty)) = L^{(\mu)}(\infty)$ is the limit of $M_k(\infty)$, $\xi(\infty) \in \Gamma(\lambda)$ for some λ large enough. Now

$$L^{(\mu)}(\infty) - M_{2k}(\infty) = \sum_{p=k}^{\infty} (M_{2p+2}(\infty) - M_{2p}(\infty)) = (0, R_k)$$

where

$$R_k = \sum_{n=k}^{\infty} \left(\frac{\varepsilon_{2n+1}}{n+\mu} - \frac{\varepsilon_{2n}}{n+1-\mu}\right).$$

Let $T_k = \frac{\varepsilon_{2k+1}}{k+\mu}$ and

$$S_k = R_{k+1} - \sum_{n=2k+1}^{+\infty} \frac{2\varepsilon_n}{n(n+1) + (1-2\mu)\nu_n}$$

Then

$$L^{(\mu)}(\infty) - M_{2k+1}(\infty) = L^{(\mu)}(\infty) - M_{2k}(\infty) - (M_{2k+1}(\infty) - M_{2k}(\infty))$$

= $(0, R_k - (\xi_{2k+1}(\infty) - \xi_{2k}(\infty)))$
= $(0, S_k + T_k)$

since

$$\xi_{k+1}(\infty) - \xi_k(\infty) = \frac{\xi_{k+1}(\infty)}{k+1} - \frac{2k\varepsilon_k}{k(k+1) + (1-2\mu)\nu_k}.$$

Let a < 0 < b. As the sequences R_n, S_n, T_n converge almost surely to 0, we can choose $n_0 \ge 1$ such that

$$\mathbb{P}(R_n, S_n, T_n \in [a/2, b/2] \text{ for all } n \ge n_0) > 0.$$

One has

$$\mathbb{P}(R_n, S_n + T_n \in [a, b], \text{ for all } n \ge 0)$$

$$\ge \mathbb{P}(R_k - R_{n_0}, T_k + S_k - S_{n_0} \in [a/2, b/2], \text{ for } 0 \le k < n_0,$$

$$R_n, S_n, T_n \in [a/2, b/2], \text{ for } n \ge n_0)$$

$$\ge \mathbb{P}(R_k - R_{n_0}, T_k + S_k - S_{n_0} \in [a/2, b/2], \text{ for } 0 \le k < n_0)$$

$$\times \mathbb{P}(R_n, S_n, T_n \in [a/2, b/2], \text{ for } n \ge n_0) > 0.$$

This allows to finish the proof, since, by (5.9), $\xi(\infty) \in \Gamma(\lambda)$ if and only if, for all $k \ge 0$,

(5.10)
$$\tilde{\alpha}_k(L^{(\mu)}(\infty) - M_k(\infty)) \le \tilde{\alpha}_k(\lambda)$$

where for a given $\lambda \in C_{\text{aff}}$, $\tilde{\alpha}_k(\lambda)$ takes only two positive values.

In order to prove Proposition 5.14 below, and to study the string parameters $\xi_k(t)$ in the next section we need to show some uniform approximations by the dihedral case. We take here $\mu \neq 0, 1$, and m large enough to ensure that $(\frac{m}{\pi}, \mu) \in C_m$. Let

$$M_k^m(\infty) = \frac{1}{2} \xi_k^m(\infty) \alpha_k^m + \sum_{n=0}^{k-1} \xi_n^m(\infty) \alpha_n^m,$$

for $0 \le k \le m$, with the convention that $\xi_m^m(\infty) = 0$.

Proposition 5.11. When $\mu \neq 0, 1$, for any $k \geq 0$, $\tau_m M_k^m(\infty)$ converges to $M_k(\infty)$ a.s. when m goes to infinity.

Proof. This follows from Lemma 5.2 and Proposition 5.4.

Proposition 5.12. When $\mu \neq 0, 1$, one has, in probability,

(5.11)
$$\lim_{m \to +\infty} \sup_{1 \le k \le m} \|\tau_m M_k^m(\infty) - M_k(\infty)\| = 0.$$

Proof. We use Corollary 4.8 and its notations. For $1 \le 2p + 1 \le m$,

$$M_{2p+2}^{m}(\infty) - M_{2p}^{m}(\infty) = \frac{1}{2}\xi_{2p+2}^{m}(\infty)\alpha_{2p+2}^{m} + \xi_{2p+1}^{m}(\infty)\alpha_{2p+1}^{m} + \frac{1}{2}\xi_{2p}^{m}(\infty)\alpha_{2p}^{m}.$$

Therefore,

$$\begin{split} \tilde{\alpha}_{0}^{m}(M_{2p+2}^{m}(\infty) - M_{2p}^{m}(\infty)) &= \xi_{2p+2}^{m}(\infty) - 2\xi_{2p+1}^{m}(\infty) \cos \frac{\pi}{m} + \xi_{2p}^{m}(\infty) \\ &= -\frac{2\varepsilon_{2p+1}^{m}a_{2p+1}^{m}\cos \frac{\pi}{m}}{\gamma_{1}^{m}a_{1}^{m} + \gamma_{0}^{m}a_{2}^{m} + \dots + \gamma_{1}^{m}a_{2p+1}^{m}} \\ &+ a_{2p}^{m}(\frac{\varepsilon_{2p+1}^{m}}{\gamma_{1}^{m}a_{1}^{m} + \dots + \gamma_{1}^{m}a_{2p+1}^{m}} + \frac{\varepsilon_{2p}^{m}}{\gamma_{1}^{m}a_{1}^{m} + \dots + \gamma_{0}^{m}a_{2p}^{m}}), \end{split}$$

since $\$

$$a_{2p+2}^m - 2a_{2p+1}^m \cos\frac{\pi}{m} + a_{2p}^m = 0.$$

One deduces that

$$\tilde{\alpha}_{0}^{m}(M_{2p}^{m}(\infty)) = 2\sin\frac{\pi}{m} \sum_{l=0}^{p-1} \left(\frac{\varepsilon_{2l}^{m}\cos\frac{l\pi}{m}}{\gamma_{1}^{m}\sin\frac{l}{m}\pi + \gamma_{0}^{m}\sin\frac{l+1}{m}\pi} - \frac{\varepsilon_{2l+1}^{m}\cos\frac{(l+1)\pi}{m}}{\gamma_{1}^{m}\sin\frac{l+1}{m}\pi + \gamma_{0}^{m}\sin\frac{l}{m}\pi}\right)$$

by using the relations

$$\sum_{k=1}^{n} a_{2k}^{m} = \frac{\sin\frac{n\pi}{m}\sin\frac{(n+1)\pi}{m}}{\sin\frac{\pi}{m}}, \sum_{k=1}^{n} a_{2k-1}^{m} = \frac{\sin^{2}\frac{n\pi}{m}}{\sin\frac{\pi}{m}}, \sum_{k=1}^{n} a_{k}^{m} = \frac{\sin\frac{n\pi}{2m}\sin\frac{(n+1)\pi}{2m}}{\sin\frac{\pi}{2m}}.$$

Similarly, one finds

$$\tilde{\alpha}_{1}^{m}(M_{2p}^{m}(\infty)) = -2\xi_{0}^{m}(\infty)\cos\frac{\pi}{m} + \xi_{2p}^{m}(\infty)\sin\frac{\pi}{m}\tan\frac{p\pi}{m} + 2\sin\frac{\pi}{m}\sum_{l=1}^{p-1} \left(\frac{\varepsilon_{2l+1}^{m}\cos\frac{l\pi}{m}}{\gamma_{1}^{m}\sin\frac{l+1}{m}\pi + \gamma_{0}^{m}\sin\frac{l}{m}\pi} - \frac{\varepsilon_{2l}^{m}\cos\frac{(l+1)\pi}{m}}{\gamma_{1}^{m}\sin\frac{l}{m}\pi + \gamma_{0}^{m}\sin\frac{l+1}{m}\pi}\right).$$

On the other hand, one has

$$M_{2p+1}^{m}(\infty) - M_{2p}^{m}(\infty) = \frac{1}{2} (\xi_{2p}^{m}(\infty)\alpha_{0}^{m} + \xi_{2p+1}^{m}(\infty)\alpha_{1}^{m}).$$

Thus the proposition follows from the next lemma, from (5.5) and from Lemma 5.2. $\hfill \Box$

Lemma 5.13. Let $\gamma_0^m, \gamma_1^m, m \in \mathbb{N}$, be real numbers such that γ_0^m tends to $2(1-\mu)$ and γ_1^m tends to 2μ when m tends to $+\infty$. Let, for $p \in \{0, \cdots, [m/2]\}$,

$$S_p^m = \sin \frac{\pi}{m} \sum_{l=0}^{p-1} \left(\frac{\varepsilon_{2l}^m \cos \frac{l\pi}{m}}{\gamma_1^m \sin \frac{l}{m} \pi + \gamma_0^m \sin \frac{l+1}{m} \pi} - \frac{\varepsilon_{2l+1}^m \cos \frac{(l+1)\pi}{m}}{\gamma_1^m \sin \frac{l+1}{m} \pi + \gamma_0^m \sin \frac{l}{m} \pi} \right),$$
$$S_p = \sum_{l=0}^{p-1} \left(\frac{\varepsilon_{2l}}{2l+2(1-\mu)} - \frac{\varepsilon_{2l+1}}{2l+2\mu} \right),$$

where $(\varepsilon_n^m)_{n\geq 0}$ is a sequence of independent exponentially distributed random variables with parameter 1 that converges almost surely to $(\varepsilon_n)_{n\geq 0}$ when m goes to infinity. Then in probability

$$\lim_{m \to \infty} \sup_{1 \le p \le m/2} |S_p^m - S_p| = 0.$$

Proof. We have

$$\sup_{1 \le p \le m/2} |S_p^m - S_p| \le \sup_{1 \le p \le m/2} |S_p^m - S_p - \mathbb{E}(S_p^m - S_p)| + \sup_{1 \le p \le m/2} |\mathbb{E}(S_p^m - S_p)|.$$

Let us show that the first term of the right-hand side converges to 0 in probability. One has for $N \leq m/2$,

(5.13)
$$\sup_{1 \le p \le m/2} |S_p^m - S_p - \mathbb{E}(S_p^m - S_p)| \le 2 \sup_{1 \le p \le N} |S_p^m - S_p - \mathbb{E}(S_p^m - S_p)| + \sup_{N \le p \le m/2} |\tilde{S}_p^m - \tilde{S}_p - \mathbb{E}(\tilde{S}_p^m - \tilde{S}_p)|,$$

where $\tilde{S}_p^m = S_p^m - S_N^m$ and $\tilde{S}_p = S_p - S_N$, for $p \ge N$. The first term of the right hand-side of (5.13) converges to 0, for any N, when m goes to infinity. As $\{\tilde{S}_{N+k}^m - \tilde{S}_{N+k} - \mathbb{E}(\tilde{S}_{N+k}^m - \tilde{S}_{N+k}), k = 0, \cdots, [m/2] - N\}$ is a martingale, Doob's martingale inequality gives, for a > 0,

(5.14)
$$\mathbb{P}(\sup_{N \le p \le m/2} |\tilde{S}_p^m - \tilde{S}_p - \mathbb{E}(\tilde{S}_p^m - \tilde{S}_p)| \ge a) \le \frac{1}{a^2} \operatorname{Var}(\tilde{S}_{[m/2]}^m - \tilde{S}_{[m/2]})$$

Besides, one has

$$\operatorname{Var}(\tilde{S}^{m}_{[m/2]} - \tilde{S}_{[m/2]}) \le 2\operatorname{Var}(\tilde{S}^{m}_{[m/2]}) + 2\operatorname{Var}(\tilde{S}_{[m/2]})$$

One has

$$\operatorname{Var}(\tilde{S}_{[m/2]}^{m}) \leq \sum_{l=N+1}^{[m/2]-1} \frac{\sin^2 \pi/m}{(\gamma_1^m \sin \frac{l}{m}\pi + \gamma_0^m \sin \frac{l+1}{m}\pi)^2} + \frac{\sin^2 \pi/m}{(\gamma_1^m \sin \frac{l+1}{m}\pi + \gamma_0^m \sin \frac{l}{m}\pi)^2},$$

and

$$\operatorname{Var}(\tilde{S}_{[m/2]}) \le \sum_{l=N+1}^{[m/2]-1} \frac{1}{(2l+2(1-\mu))^2} + \frac{1}{(2l+2\mu)^2}.$$

As $\frac{2x}{\pi} \le \sin x \le x$ for $0 \le x \le \frac{\pi}{2}$,

$$\frac{\sin\frac{\pi}{m}}{\gamma_1^m \sin\frac{l}{m}\pi + \gamma_0^m \sin\frac{l+1}{m}\pi} \le \frac{\pi/2}{\gamma_1^m l + \gamma_0^m (l+1)}$$

and

$$\frac{\sin\frac{\pi}{m}}{\gamma_1^m \sin\frac{l+1}{m}\pi + \gamma_0^m \sin\frac{l}{m}\pi} \le \frac{\pi/2}{\gamma_1^m(l+1) + \gamma_0^m l},$$

for $N \leq l \leq m/2 - 1$. Thus for any $\varepsilon > 0$, there is a $N_0 > 0$ such that for all $m \geq 2(N_0 + 1)$, the right hand-side term of (5.14) is smaller than ε . Hence $\sup_{1\leq p\leq m/2} |S_p^m - S_p - \mathbb{E}(S_p^m - S_p)|$ converges in probability to 0. Let us deal with the second term of the right hand-side of (5.12). One has for $0 \leq q \leq p \leq m/2$,

$$|\mathbb{E}(S_p^m - S_q^m)| \le \sum_{l=q}^{p-1} \frac{\pi^2(\gamma_0^m + \gamma_1^m)}{4(\gamma_1^m l + \gamma_0^m (l+1))(\gamma_1^m (l+1) + \gamma_0^m l)}$$

and

$$\mathbb{E}(S_p - S_q) = \sum_{l=q}^{p-1} \frac{2(2\mu - 1)}{(2l + 2(1-\mu))(2l + 2\mu)}$$

As for a fixed q, $\mathbb{E}(S_q^m)$ converges to $\mathbb{E}(S_q)$ when m goes to infinity, one obtains that $\sup_{0 \le p \le m/2} |\mathbb{E}(S_p^m - S_p)|$ converges to 0, which finishes the proof of the lemma.

Recall that $\xi^m(\infty)$, resp. $\xi(\infty)$, are the Verma string parameters of $W^{(\frac{m}{\pi},\mu)}$, resp. $B^{(\mu)}$.

Proposition 5.14. Suppose that $\mu \neq 0, 1$. Let (λ_m) be a sequence of \mathbb{R}^2 such that $\lambda_m \in C_m$ and such that $\tau_m \lambda_m$ tends to λ when m tends to ∞ where $\lambda \in C_{aff}$. The random sets $\{\xi^m(\infty) \in \Gamma_m(\lambda_m)\}$ converge in probability to $\{\xi(\infty) \in \Gamma(\lambda)\}$.

Proof. To prove convergence in probability of a sequence, it is enough to show that each subsequence has a subsequence which converges almost surely. Therefore, by Propositions 5.12 and 5.7, working with a subsequence, we can suppose that the set of $\omega \in \Omega$ for which, as $m \to +\infty$,

$$\sup_{1 \le k \le m} \|\tau_m M_k^m(\infty)(\omega) - M_k(\infty)(\omega)\| \to 0,$$

and for which $M_m(\infty)(\omega)$ tends to $L^{(\mu)}(\infty)(\omega)$ has probability one. If

$$X_k^m = M_m^m(\infty) - M_k^m(\infty), \quad X_k = L^{(\mu)}(\infty) - M_k(\infty),$$

then, when m tends to $+\infty$ (see (5.3))

(5.15)
$$\sup_{0 \le k \le m} |\tilde{\alpha}_k^m(X_k^m(\omega)) - \tilde{\alpha}_k(X_k(\omega))| \to 0,$$

and

(5.16)
$$\tilde{\alpha}_m(X_m(\omega)) \to 0.$$

For these ω we will show that

$$\limsup_{m \to \infty} \{\xi^m(\infty) \in \Gamma_m(\lambda_m)\} \subset \{\xi(\infty) \in \Gamma(\lambda)\} \subset \liminf_{m \to \infty} \{\xi^m(\infty) \in \Gamma_m(\lambda_m)\}.$$

Notice that $\tilde{\alpha}_k^m(\lambda_m)$ tends to $\tilde{\alpha}_k(\lambda)$ for k = 0, 1 by Lemma 5.2. One has $\xi^m(\infty) \in \Gamma_m(\lambda_m)$ if and only if $\tilde{\alpha}_k^m(\lambda_m - X_k^m) \ge 0$ for $0 \le k < m$, and $\xi(\infty) \in \Gamma(\lambda)$ if and only if $\tilde{\alpha}_k(\lambda - X_k) \ge 0$ for every $k \in \mathbb{N}$. The left inclusion above follows from (5.15). Now, suppose that $\xi(\infty)(\omega) \in \Gamma(\lambda)$. Since λ is fixed and since the distribution of X_k is continuous,

$$\mathbb{P}(\tilde{\alpha}_k(\lambda - X_k) = 0) = 0,$$

and one can suppose that for all $k \ge 0$ (see 5.10),

$$\tilde{\alpha}_k(\lambda - X_k(\omega)) > 0.$$

We choose $\varepsilon > 0$ such that $\tilde{\alpha}_i(\lambda) > \varepsilon, i = 0, 1$. Using (5.15), one can choose m_0 such that $m \ge m_0$ implies that

$$\sup_{0 \le k \le m} \tilde{\alpha}_k^m(X_k^m(\omega)) - \tilde{\alpha}_k(X_k(\omega)) < \varepsilon.$$

Then, by (5.16), we choose k_0 such that $k \ge k_0$ implies that

$$\tilde{\alpha}_k(X_k(\omega)) < \tilde{\alpha}_k(\lambda) - \varepsilon$$

As for each k, $\tilde{\alpha}_k^m(X_k^m)(\omega)$ converges to $\tilde{\alpha}_k(X_k)(\omega)$ when m goes to infinity, one takes m_1 such that when $m \geq m_1$,

$$\tilde{\alpha}_k^m(X_k^m(\omega)) < \tilde{\alpha}_k(\lambda),$$

for $k \leq k_0$. For $m \geq \max(m_0, m_1)$ one has for $k \geq k_0$

$$\tilde{\alpha}_k^m(X_k^m(\omega)) = \tilde{\alpha}_k^m(X_k^m(\omega)) - \tilde{\alpha}_k(X_k(\omega)) + \tilde{\alpha}_k(X_k(\omega)) < \tilde{\alpha}_k(\lambda).$$

This proves the right inclusion above, and shows the proposition.

This proposition will be crucial to move from Verma affine string parameters to affine string parameters in the next section.

6. The highest weight process $\Lambda^{(\mu)}$

In this section we introduce the highest weight process $\Lambda^{(\mu)}$ associated with the space-time Brownian motion $B^{(\mu)}$ as the limit of its Pitman's transforms (with a correction) and show that it coincides in distribution with $A^{(\mu)}$. We define for $k \geq 0$ and $t \in \mathbb{R}^+$,

$$M_{k}(t) = \frac{1}{2}\xi_{k}(t)\alpha_{k} + \sum_{n=0}^{k-1}\xi_{n}(t)\alpha_{n},$$

where $\xi(t) = \{\xi_n(t), n \ge 0\}$ are the affine string parameters of $B^{(\mu)}$ on [0, t], and

$$M_{k}^{m}(t) = \frac{1}{2}\xi_{k}^{m}(t)\alpha_{k}^{m} + \sum_{n=0}^{k-1}\xi_{n}^{m}(t)\alpha_{n}^{m},$$

where $\xi^m(t) = \{\xi_n^m(t), 0 \le n < m\}$ are the string parameters of $W^{(\frac{m}{\pi},\mu)}$ on [0,t] (by convention $\xi_n^m(t) = 0$ when $n \ge m$). In particular,

(6.1)
$$M_m^m(t) = \sum_{n=0}^{m-1} \xi_n^m(t) \alpha_n^m = \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(t) - W^{(\frac{m}{\pi},\mu)}(t).$$

Below we will often use the Cameron Martin Girsanov theorem (see [25] Theorem 3.5.1), abreviated CMG, to reduce the case $0 \le \mu \le 1$ to the case where $0 < \mu < 1$. Each time the reason is the following: on the time interval [0, t] the processes considered for $\mu = 0$ and $\mu = 1$ are absolutely continuous with the ones for $0 < \mu < 1$.

Proposition 6.1. For each $t \ge 0$, $M_k(t)$ converges in probability when k goes to infinity. We denote by $L^{(\mu)}(t)$ the limit.

Proof. By CMG, one can suppose that $\mu \neq 0, 1$. For every $k \geq 1$ and $t \geq 0$, $\tau_m M_k^m(t)$ converges to $M_k(t)$ almost surely when m goes to infinity by Proposition 5.4 and Lemma 5.2. Let $\varepsilon > 0$. For $p, q \in \mathbb{N}$,

$$\mathbb{P}(\|M_{p+q}(t) - M_p(t)\| \ge \varepsilon) = \lim_{m \to \infty} \mathbb{P}(\|\tau_m M_{p+q}^m(t) - \tau_m M_p^m(t)\| \ge \varepsilon)$$
$$= \lim_{m \to \infty} \mathbb{E}(1_{\|\tau_m M_{p+q}^m(t) - \tau_m M_p^m(t)\| \ge \epsilon} |\sigma(\Lambda_m^{(\mu)}(s), s \le t))$$
$$= \lim_{m \to \infty} f_m(\Lambda_m^{(\mu)}(t)),$$

where $\Lambda_m^{(\mu)}(t) = \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(t)$ and where

$$f_m(\lambda) = \mathbb{P}(\|\tau_m M_{p+q}^m(\infty) - \tau_m M_p^m(\infty)\| \ge \varepsilon |\xi^m(\infty) \in \Gamma_m(\lambda)),$$

by Corollary 4.7. If $\lambda \in C_{\text{aff}}$ and $\tau_m \lambda_m$ tends to λ , then by Propositions 5.14 and 5.11, $f_m(\lambda_m)$ tends to $f(\lambda)$ where

$$f(\lambda) = \mathbb{P}(\|M_{p+q}(\infty) - M_p(\infty)\| \ge \varepsilon |\xi(\infty) \in \Gamma(\lambda)).$$

Notice that this is well defined by Proposition 5.10. On the other hand we have seen in Theorem 3.7 that when m tends to infinity, $\tau_m \Lambda_m^{(\mu)}(t)$ converges in distribution to $A^{(\mu)}(t)$, where $A^{(\mu)}$ is the Brownian motion conditioned to remain in C_{aff} defined in Section 2.2. Therefore $\mathbb{E}(f_m(\Lambda_m^{(\mu)}(t)))$ tends to $\mathbb{E}(f(A^{(\mu)}(t)))$. This shows that

$$\mathbb{P}(\|M_{p+q}(t) - M_p(t)\| \ge \varepsilon)$$

= $\int_{\Omega} \mathbb{P}(\|M_{p+q}(\infty) - M_p(\infty)\| \ge \varepsilon |\xi(\infty) \in \Gamma(A^{(\mu)}(t)(\omega))) d\mathbb{P}(\omega)$

Since $M_p(\infty)$ converges a.s. to $L^{(\mu)}(\infty)$ (Proposition 5.7) we see that $M_p(t), p \in \mathbb{N}$, is a Cauchy sequence for the convergence in probability, and thus converges in probability.

We will see in Theorem 6.9 that the convergence in the proposition actually also holds almost surely.

Proposition 6.2. When $\mu \neq 0, 1$, there is an almost surely finite random time $\mathfrak{t} \geq 0$ such that $L^{(\mu)}(t) = L^{(\mu)}(\infty)$ for $t \geq \mathfrak{t}$.

Proof. As in the proof of Proposition 4.1 we define $\mathfrak{t}_0 = \max\{t \ge 0, B_t^{(\mu)} \notin C_{\text{aff}}\}$, and, for $n \ge 0$,

$$\mathfrak{t}_{n+1} = \max\{t \ge \mathfrak{t}_n, B_t^{(\mu)} + \sum_{k=0}^n \xi_k(t)\alpha_k \notin C_{\text{aff}}\}.$$

We see by induction that $\mathfrak{t}_n < +\infty$ and that for $t > \mathfrak{t}_n$, $\xi_n(t) = \xi_n(\infty)$. Therefore

$$\mathfrak{t}_{n+1} = \max\{t \ge \mathfrak{t}_n, B_t^{(\mu)} + \sum_{k=0}^n \xi_k(\infty) \alpha_k \notin C_{aff}\}.$$

We know that a.s., $\sum_{k=0}^{n} \xi_k(\infty) \alpha_k$ is bounded. Hence $\mathfrak{t} = \sup \mathfrak{t}_n < +\infty$ and for $t \geq \mathfrak{t}, \, \xi_k(t) = \xi_k(\infty)$ for all $k \geq 0$. So, for $t \geq \mathfrak{t}, \, L^{(\mu)}(t) = L^{(\mu)}(\infty)$.

Proposition 6.3. For any $t \ge 0$, as m tends to ∞ , $\tau_m M_m^m(t)$ converges in probability to $L^{(\mu)}(t)$.

Proof. One can suppose that $\mu \neq 0, 1$. In that case, one has for $\varepsilon > 0$,

$$\mathbb{P}(\|\tau_m M_m^m(t) - L^{(\mu)}(t)\| > \varepsilon) \le \mathbb{P}(\|\tau_m M_m^m(t) - M_m(t))\| > \frac{\varepsilon}{2}) + \mathbb{P}(\|M_m(t) - L^{(\mu)}(t)\| > \frac{\varepsilon}{2})$$

For the first term, we condition by $\sigma(\Lambda_m^{(\mu)}(s), s \leq t)$ as in the proof of Proposition 6.1 and we use Proposition 5.12. Convergence in probability of Proposition 6.1 gives the convergence of the second term.

We introduce the following definition,

Definition 6.4. We define the highest weight process $\Lambda^{(\mu)}$ of $B^{(\mu)}$ by, for $t \ge 0$,

$$\Lambda^{(\mu)}(t) = B_t^{(\mu)} + L^{(\mu)}(t)$$

In the analogy with the Littelmann model, for t > 0 fixed, $\{B_s^{(\mu)}, 0 \le s \le t\}$ can be considered as a path with weight $B_t^{(\mu)}$, with highest weight $\Lambda^{(\mu)}(t)$, and $L^{(\mu)}(t)$ as the weight seen from the highest weight. Recall that

(6.2)
$$\Lambda_m^{(\mu)}(t) = \mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)}(t) = W_t^{(\frac{m}{\pi},\mu)} + M_m^m(t).$$

Proposition 6.5. In probability,

$$\lim_{m \to +\infty} \tau_m W_t^{(\frac{m}{\pi},\mu)} = B_t^{(\mu)},$$
$$\lim_{m \to +\infty} \tau_m \Lambda_m^{(\mu)}(t) = \Lambda^{(\mu)}(t).$$

Proof. The first statement is obvious. The second one is then a consequence of Proposition 6.3 and (6.2). \Box

Let $\{A^{(\mu)}(t), t \geq 0\}$ be the conditioned space-time Brownian motion in the affine Weyl chamber C_{aff} with drift μ starting from the origin, defined in 2.2.

Theorem 6.6. In distribution,

$$\{\Lambda^{(\mu)}(t), t \ge 0\} = \{A^{(\mu)}(t), t \ge 0\}$$

Proof. This follows from Proposition 6.5, and Theorem 3.7.

This shows in particular that $\Lambda^{(\mu)}$ has a continuous version. Recall that for any $\lambda \in C_{\text{aff}}$, $\mathbb{P}(\xi(\infty) \in \Gamma(\lambda)) > 0$ by Proposition 5.10.

Proposition 6.7. For $\mu \neq 0, 1$, and $f : \mathbb{R}^2 \times \Gamma \times \mathbb{R}^2 \to \mathbb{R}$, bounded and measurable,

$$\mathbb{E}(f(B_t^{(\mu)},\xi(t),\Lambda^{(\mu)}(t))|\sigma(\Lambda^{(\mu)}(s),s\leq t))=g(\Lambda^{(\mu)}(t))$$

for each $t \ge 0$, where

$$g(\lambda) = \mathbb{E}(f(\lambda - L^{(\mu)}(\infty), \xi(\infty), \lambda) | \xi(\infty) \in \Gamma(\lambda))$$

Proof. It is enough to prove this when f is a bounded continuous function depending on a finite number of variables. By Corollary 4.7,

$$\mathbb{E}(f(\tau_m \Lambda_m^{(\mu)}(t) - \tau_m M_m^m(t), \xi^m(t), \Lambda_m^{(\mu)}(t)) | \sigma(\Lambda_m^{(\mu)}(s), s \le t))) = g_m(\Lambda_m^{(\mu)}(t))$$

where

$$g_m(\lambda) = \mathbb{E}(f(\tau_m \lambda - \tau_m M_m^m(\infty), \xi^m(\infty), \lambda) | \xi^m(\infty) \in \Gamma_m(\lambda)).$$

Let h be a bounded continuous function depending only on a finite number of variables. Using the convergences in probability of the propositions 6.3 and 6.5,

$$\mathbb{E}(f(\Lambda^{(\mu)}(t) - L^{(\mu)}(t), \xi(t), \Lambda^{(\mu)}(t))h(\Lambda^{(\mu)}(s), s \le t))) = \lim_{m \to +\infty} \mathbb{E}(f(\tau_m \Lambda_m^{(\mu)}(t) - \tau_m M_m^m(t), \xi^m(t), \tau_m \Lambda_m^{(\mu)}(t))h(\tau_m \Lambda_m^{(\mu)}(s), s \le t))) = \lim_{m \to +\infty} \mathbb{E}(g_m(\Lambda_m^{(\mu)}(t))h(\tau_m \Lambda_m^{(\mu)}(s), s \le t))).$$

Since $\tau_m \Lambda_m^{(\mu)}(t)$ tends to $\Lambda^{(\mu)}(t)$, we also know from Proposition 5.14 that $\{\xi^m(\infty) \in \Gamma_m(\Lambda_m^{(\mu)}(t))\}$ tends to $\{\xi(\infty) \in \Gamma(\Lambda^{(\mu)}(t))\}$ in probability. Therefore the limit above equals

$$\mathbb{E}(g(\Lambda^{(\mu)}(t))h(\Lambda^{(\mu)}(s), s \le t)))$$

which proves the proposition.

Applying the proposition with $f(b,\xi,\lambda) = 1_{\Gamma(\lambda)}(\xi)$ we obtain

Corollary 6.8. Almost surely, $\xi(t)$ belongs to $\Gamma(\Lambda^{(\mu)}(t))$.

Theorem 6.9. $M_k(t)$ converges almost surely to $L^{(\mu)}(t)$ when $k \to +\infty$.

Proof. One can suppose that $\mu \neq 0, 1$. In that case, by Proposition 6.7, one has $\mathbb{E}(1_{\{\lim M_k(t)=L^{(\mu)}(t)\}} | \sigma(\Lambda^{(\mu)}(s), s \leq t)) = \mathbb{E}(1_{\{\lim M_k(\infty)=L^{(\mu)}(\infty)\}} | \xi(\infty) \in \Gamma(\Lambda^{(\mu)}(t))).$ As the right hand-side term equals 1 by Proposition 5.7, one obtains that

$$\mathbb{P}(\lim M_k(t) = L^{(\mu)}(t)) = \mathbb{E}(\mathbb{E}(1_{\{\lim M_k(t) = L^{(\mu)}(t)\}} | \sigma(\Lambda^{(\mu)}(s), s \le t))) = 1.$$

7. The representation theorem using Pitman and Lévy transforms

7.1. Representation of the conditioned space-time Brownian motion $A^{(\mu)}$. Let us remind where we stand. For $B_t^{(\mu)} = (t, B_t + t\mu)$ we have written

$$\mathcal{P}_{s_n}\cdots\mathcal{P}_{s_1}\mathcal{P}_{s_0}B^{(\mu)}(t) = B_t^{(\mu)} + \sum_{i=0}^n \xi_i(t)\alpha_i,$$

see (5.2), and

$$M_{n+1}(t) = \frac{1}{2}\xi_{n+1}(t)\alpha_{n+1} + \sum_{i=0}^{n}\xi_i(t)\alpha_i.$$

We have seen that when for t > 0, $\lim_{k \to +\infty} M_k(t) = L^{(\mu)}(t)$ a.s. and that the process $\Lambda^{(\mu)}(t) = B_t^{(\mu)} + L^{(\mu)}(t)$, $t \ge 0$, has the same distribution as the process $A^{(\mu)}(t), t \ge 0$. Hence,

Theorem 7.1. For $t \ge 0$, almost surely,

$$\lim_{n \to +\infty} \mathcal{P}_{s_n} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} B^{(\mu)}(t) + \frac{1}{2} \xi_{n+1}(t) \alpha_{n+1}$$

exists, and the limiting process has the same distribution as $\{A^{(\mu)}(t), t \ge 0\}$.

To interpret the correction term, it is worthwhile to introduce the Lévy transform (sometimes called the Skorokhod transform). A theorem of Lévy (see Revuz and Yor [36], VI.2) states that if β is the standard Brownian motion, then

$$\mathcal{L}\beta(t) = \beta_t - \inf_{0 \le s \le t} \beta_s, t \ge 0,$$

has the same distribution as $|\beta_t|, t \ge 0$, and that $-\inf_{0\le s\le t} \beta_s$ is the local time of $\mathcal{L}\beta$ at 0. We define here the following Lévy transform (sometimes the Lévy transform of β is defined as $\int_0^t \operatorname{sign}(\beta_s) d\beta_s$, this is related to the transform here, but different).

Definition 7.2. For $\eta \in C_0(\mathbb{R}^2)$ and i = 0, 1, the Lévy transform $\mathcal{L}_{s_i}\eta$ of η is

$$\mathcal{L}_{s_i}\eta(t) = \eta(t) - \frac{1}{2} \inf_{0 \le s \le t} \tilde{\alpha}_i(\eta(s))\alpha_i.$$

Another way to state the former theorem is

Theorem 7.1 bis. For $t \ge 0$,

$$\lim_{n \to +\infty} \mathcal{L}_{s_{n+1}} \mathcal{P}_{s_n} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} B^{(\mu)}(t)$$

exists a.s. and the limiting process has the same distribution as $A^{(\mu)}$.

The following proposition indicates that the presence of the Lévy transform is due to the bad behavior of $A^{(\mu)}(t)$ for t near 0.

Proposition 7.3. For all t > 0, a.s. $\lim_{k \to +\infty} \xi_k(t) = 2$.

Proof. When $\mu \neq 0, 1, \xi_k(\infty)$ tends to 2 almost surely (Theorem 5.6), and the proposition follows from Proposition 6.7 by conditioning by $\sigma(\Lambda^{(\mu)}(s), s \leq t)$. The case when $\mu \in \{0, 1\}$ follows from the CMG theorem.

This implies that, for t > 0,

$$\lim_{n \to +\infty} \mathcal{P}_{s_n} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} B^{(\mu)}(t) + (-1)^n \alpha_1 = A^{(\mu)}(t).$$

So without the Levy correction, the iterates of Pitman's transform do not converge.

In the whole paper we have chosen to begin by first applying \mathcal{P}_{s_0} to $B^{(\mu)}$ and then \mathcal{P}_{s_1} , in the iterations of Pitman's transforms. Let us show the non trivial fact that we obtain the same limit if we begin with \mathcal{P}_{s_1} . More precisely, if we denote with a tilde the quantities previously defined when we begin by \mathcal{P}_{s_1} instead of \mathcal{P}_{s_0} , one has,

Theorem 7.4. (i) Almost surely,

$$\lim_{n \to +\infty} \mathcal{L}_{s_{n+1}} \mathcal{P}_{s_n} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} B^{(\mu)}(t) = \lim_{n \to +\infty} \mathcal{L}_{s_{n+1}} \mathcal{P}_{s_n} \cdots \mathcal{P}_{s_2} \mathcal{P}_{s_1} B^{(\mu)}(t).$$

(ii) $\tilde{\xi}(t)$ defined for $B^{(\mu)}$ has the same distribution as $\xi(t)$ defined for $B^{(1-\mu)}$.

Proof. For (i) we have to prove that $L^{(\mu)}(t) = \tilde{L}^{(\mu)}(t)$. We have seen in Proposition 6.3 that $\tau_m M_m^m(t)$ converges in probability to $L^{(\mu)}(t)$. By the same proof $\tau_m \tilde{M}_m^m(t)$ converges in probability to $\tilde{L}^{(\mu)}(t)$. Now by (6.2),

$$M_{m}^{m}(t) = \mathcal{P}_{w_{m}}^{m} W^{(\frac{m}{\pi},\mu)}(t) - W_{t}^{(\frac{m}{\pi},\mu)}$$

and by Theorem 3.4

$$\mathcal{P}_{w_m}^m W^{(\frac{m}{\pi},\mu)} = \tilde{\mathcal{P}}_{w_m}^m W^{(\frac{m}{\pi},\mu)}.$$

since

$$w_m = s_{m-1}^m \cdots s_1^m s_0^m = s_m^m \cdots s_2^m s_1^m.$$

Hence $M_m^m(t) = \tilde{M}_m^m(t)$ and $L^{(\mu)}(t) = \tilde{L}^{(\mu)}(t)$ a.s.. (ii) follows from the relation (5.1) and the fact that $TB^{(\mu)}$ has the same distribution as $B^{(1-\mu)}$.

7.2. Representation of the conditioned Brownian motion in [0, 1]. For a continuous real path $\varphi : \mathbb{R}_+ \to \mathbb{R}$ such that $\varphi(0) = 0$, we have defined

$$\mathcal{L}_0\varphi(t) = \varphi(t) + \inf_{0 \le s \le t} (s - \varphi(s)), \ \mathcal{P}_0\varphi(t) = \varphi(t) + 2\inf_{0 \le s \le t} (s - \varphi(s)),$$
$$\mathcal{L}_1\varphi(t) = \varphi(t) - \inf_{0 \le s \le t} \varphi(s), \ \mathcal{P}_1\varphi(t) = \varphi(t) - 2\inf_{0 \le s \le t} \varphi(s).$$

For $B_t^{\mu} = B_t + t\mu, t \ge 0$, the theorem stated in the introduction is

Theorem 7.5. For any t > 0, almost surely,

$$\lim_{n \to \infty} t \mathcal{L}_{n+1} \mathcal{P}_n \cdots \mathcal{P}_1 \mathcal{P}_0 B^{\mu}(1/t) = \lim_{n \to \infty} t \mathcal{L}_{n+1} \mathcal{P}_n \cdots \mathcal{P}_2 \mathcal{P}_1 B^{\mu}(1/t) = Z_t^{\mu},$$

where Z_t^{μ} , $t \geq 0$, is a Brownian motion conditioned to stay in the interval [0,1] forever, starting from μ .

Proof. The proof is just the juxtaposition of the theorems 7.1, 7.4 and 2.4.

Remark that, for any t > 0 and a.s., by this theorem and Proposition 7.3,

$$\lim_{n \to \infty} t(\mathcal{P}_n \cdots \mathcal{P}_1 \mathcal{P}_0 B^{\mu} (1/t) + (-1)^n 2) = Z_t^{\mu}.$$

As an illustration let us show that: if we write $L^{(\mu)}(\infty) = (0, D^{\mu}(\infty))$ for some $D^{\mu}(\infty) \in \mathbb{R},$

Proposition 7.6. When $\mu \neq 0, 1$, there is a standard Brownian motion β and a stopping time τ with respect to the filtration $\{\sigma(\beta_s, s \leq t), t \geq 0\}$ such that $\tau > 0$ a.s. and for $0 \leq t \leq \tau$,

$$Z_t^{\mu} = \beta_t + \mu + tD^{\mu}(\infty).$$

Proof. Let $\mathfrak{t} \geq 0$ be the random time given by Proposition 6.2. For $t > \mathfrak{t}$, $L^{(\mu)}(t) = L^{(\mu)}(\infty)$ hence $A^{(\mu)}(t) = B_t^{(\mu)} + L^{(\mu)}(\infty)$. This implies the relation for $\tau = 1/\mathfrak{t}$ and $\beta_t = tB_{1/t}$, since $(t, tZ_{1/t}^{\mu}) = A^{(\mu)}(t)$.

Notice that, by (8.7), $\mathbb{E}(D^{\mu}(\infty)) = \pi \cot \pi \mu$ as expected from the expression of the drift of the generator $\frac{1}{2}d^2/dx^2 + (\pi \cot \pi x)d/dx$ of Z_t^{μ} .

7.3. A property of Pitman transform for piecewise C^1 paths. As mentioned in the introduction, the need of an infinite number of Pitman transforms to represent the space-time Brownian motion in C_{aff} is due to its wild behaviour and in particular to its non differentiability. The situation is much simpler for regular space-time curves, as shown by the following proposition (recall that paths in C_{aff} are fixed under the Pitman transforms).

Proposition 7.7. Let $\varphi : [0,T] \to \mathbb{R}$ be a continuous piecewise C^1 function such that $\varphi(0) = 0$ and let $\eta(t) = (t, \varphi(t))$. There is an n such that for all $t \in [0, T]$,

(7.1)
$$\mathcal{P}_{s_n} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} \eta(t) \in C_{aff}$$

and for all $m \geq n$,

(7.2)
$$\mathcal{P}_{s_m} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} \eta(t) = \mathcal{P}_{s_m} \cdots \mathcal{P}_{s_1} \mathcal{P}_{s_0} \eta(t)$$

We use the notations of 1.2 and Section 7.2. It is equivalent to proving that there is an n > 0 such that

$$0 \leq \mathcal{P}_n \cdots \mathcal{P}_1 \mathcal{P}_0 \varphi(t) \leq t$$

for all $t \in [0,T]$. Let $\tau_1(\varphi) = \inf\{t > 0; \varphi(t) < 0\}, \tau_0(\varphi) = \inf\{t > 0; \varphi(t) > t\},$ and let $|\varphi'|$ be the supremum of the left and right derivatives of φ on [0, T].

Lemma 7.8. (1) $|(\mathcal{P}_1 \varphi)'| \leq |\varphi'|$, and $|(\mathcal{P}_0 \varphi)'| \leq 2 + |\varphi'|$. (2) When $0 \le t \le \tau_1(\varphi) \land \tau_0(\varphi), \ \mathcal{P}_1\varphi(t) = \mathcal{P}_0\varphi(t) = \varphi(t).$

- (3) There is an n > 0 such that $0 \leq (\mathcal{P}_n \cdots \mathcal{P}_1 \mathcal{P}_0 \varphi)'(0) \leq 1$.

Proof. (1) and (2) are straightforward. For (3): if σ_0 and σ_1 are the real reflections given by $\sigma_0(x) = -x$ and $\sigma_1(x) = 2 - x$, then $(\mathcal{P}_i \varphi)'(0) = \sigma_i(\varphi'(0))$, for i = 0, 1 and it is well known that one can bring any real number into [0, 1] by repeated actions of σ_0 and σ_1 .

Lemma 7.9. Proposition 7.7 holds when $\tau_1(\varphi) \wedge \tau_0(\varphi) > 0$.

Proof. We first suppose that $\tau_0(\varphi) < \tau_1(\varphi)$. Let $\varphi_1 = \mathcal{P}_0\varphi, \varphi_2 = \mathcal{P}_1\varphi_1, \varphi_3 = \mathcal{P}_0\varphi_2, \cdots$ and $a_1 = \tau_0(\varphi), a_2 = \tau_1(\varphi_1), a_3 = \tau_0(\varphi_2), \cdots$. By (1) of Lemma 7.8, $|\varphi'_n| \leq |\varphi'| + 2n$. Since $\varphi_{2n}(a_{2n}) = 0$, and $\varphi_{2n}(a_{2n+1}) = a_{2n+1}$, we obtain by the mean value theorem that

$$a_{2n+1} - a_{2n} \ge \frac{a_{2n+1}}{|\varphi'| + 4n}$$

So the increasing sequence a_n is bigger that T for n large enough. When $\tau_0(\varphi) > \tau_1(\varphi)$, then $\tau_0(\varphi_1) > \tau_1(\varphi_1)$ and we use the same proof for φ_1 instead of φ by beginning to apply \mathcal{P}_1 .

Proof of Proposition 7.7. Using (3) of Lemma 7.8 one can suppose that $0 \leq \varphi'(0) \leq 1$. When $0 < \varphi'(0) < 1$, then $\tau_1(\varphi) \wedge \tau_0(\varphi) > 0$ and we can apply Lemma 7.9. If $\varphi'(0) = 1$, let $\gamma = \mathcal{P}_0 \varphi$. Then $\gamma'(0) = 1$ and $\gamma(t) \leq t$ for $t \geq 0$. In that case $\tau_0(\gamma) \geq T$ and $\tau_1(\gamma) > 0$. Indeed otherwise there is a sequence $t_n > 0$ decreasing to 0 such that $\gamma(t_n) \leq 0$ and $\gamma'(0)$ cannot be 1. So we can also apply Lemma 7.9. The case $\varphi'(0) = 0$ is similar. This proves (7.1), which gives 7.2 since the action of the path transforms \mathcal{P}_{s_0} and \mathcal{P}_{s_1} are trivial for path in $\overline{C}_{\text{aff}}$.

8. Some probability distributions

8.1. The conditional distribution of $B_t^{(\mu)}$ given $\sigma(\Lambda^{(\mu)}(s), s \leq t)$. We will compute this conditional distribution by approaching it by the dihedral case.

Let $m \geq 1$. Recall the definitions of $\psi_{\gamma}^m(v)$ and h_m given in (3.1) and (3.2). Let $W^{(\gamma)}$ be the standard planar Brownian motion in \mathbb{R}^2 with drift $\gamma \in \overline{C}_m$.

Lemma 8.1. For $\zeta \in \mathbb{R}^2$ and $v = \mathcal{P}_{w_m}^m W^{(\gamma)}(t)$,

$$\mathbb{E}(e^{\langle \zeta, W_t^{(\gamma)} \rangle} | \sigma(\mathcal{P}_{w_m}^m W^{(\gamma)}(s), 0 \le s \le t)) = \frac{\psi_{\zeta+\gamma}^m(v)}{\psi_{\gamma}^m(v)} \frac{h_m(\gamma)}{h_m(\zeta+\gamma)} e^{\langle \zeta, v \rangle}.$$

Proof. Theorem 5.5 in [4] gives

$$\mathbb{E}(e^{\langle \zeta, W_t^{(0)} \rangle} | \sigma(\mathcal{P}_{w_m}^m W^{(0)}(s), 0 \le s \le t)) = C \frac{\psi_{\zeta}^m(v)}{h_m(v)h_m(\zeta)} e^{\langle \zeta, v \rangle}$$

for $v = \mathcal{P}_{w_m}^m W^{(0)}(t)$ and a constant C > 0 independent of v and ζ . We conclude with Proposition 4.5.

For any $l \in \mathbb{R}$ and $(t, x) \in C_{\text{aff}}$ we define φ_l as in (2.6) by

(8.1)
$$\varphi_l(t,x) = \frac{e^{-lx}}{\sin(l\pi)} \sum_{k \in \mathbb{Z}} \sinh(2ktl + xl)e^{-2(kx+k^2t)}$$

Lemma 8.2. Let $\mu_m \in \mathbb{R}^2, v_m \in \overline{C}_m$ be such that

$$\lim_{m \to +\infty} \tau_m v_m = (t, x), \lim_{m \to +\infty} \tau_m \mu_m = (1, \mu),$$

then

(8.2)
$$\lim_{m \to \infty} \psi^m_{\mu_m}(v_m) = \frac{\sin(\mu\pi)}{2} \varphi_\mu(t, x),$$

(8.3)
$$\lim_{m \to \infty} \left(\frac{\pi}{m}\right)^m h_m(\mu_m) = \sin(\mu\pi).$$

Proof. Let r be the rotation of \mathbb{R}^2 of angle $2\pi/m$ and s be the symmetry s(a, b) = (a, -b), for $(a, b) \in \mathbb{R}^2$. Let $J(m) = \{-[m/2], \cdots, [m/2]\}$ when m is odd and $J(m) = \{-[m/2] + 1, \cdots, [m/2] - 1\} \cup \{m/2\}$ when m is even. The dihedral group I(m) is $I(m) = \{r^k, k \in J(m)\} \rtimes \{Id, s\}, l(r) = 2$, and l(s) = 1. Let $\mu_m = (\mu_m(1), \mu_m(2))$ and $v_m = (v_m(1), v_m(2))$. One has

$$\psi_{\mu_m}^m(v_m) = e^{-\langle \gamma_m, \alpha_m \rangle} \sum_{k \in J(m)} (e^{\langle r^k(\mu_m), v_m \rangle} - e^{\langle r^k s(\mu_m), v_m \rangle})$$
$$= 2e^{-\mu_m(2)v_m(2)} \sum_{k \in J(m)} I(m, k)$$

where

$$I(m,k) = \exp[\mu_m(1)v_m(1)(\cos\frac{2k\pi}{m} - 1) + \mu_m(1)v_m(2)\sin\frac{2k\pi}{m}] \times \sinh[\mu_m(2)v_m(2)\cos\frac{2k\pi}{m} - \mu_m(2)v_m(1)\sin\frac{2k\pi}{m})].$$

For $\varepsilon > 0$, we choose $m_0 \ge 0$ such that, for $m \ge m_0$,

$$\sum_{1/4 \le |k| \le m/2} |I(m,k)| \le \varepsilon$$

by using the inequality $\cos(\frac{2k}{m}\pi) \leq 0$ when $\frac{m}{4} \leq |k| \leq \frac{m}{2}$. Besides, when $t \in [-\frac{\pi}{2}, \frac{\pi}{2}]$, then $\cos t \leq 1 - \frac{t^2}{\pi}$ and $|\sin(t)| \leq |t|$, hence, for $|k| \leq \frac{m}{4}$,

$$|I(m,k)| \le e^{-c_1 k^2 + c_2 k}$$

where $c_1, c_2 > 0$ do not depend on m. So, one can choose N such that for $m \ge 1$

$$\sum_{N \le |k| \le m/4} |I(m,k)| \le \epsilon$$

We obtain (8.2) since, for N fixed,

$$\lim_{m \to +\infty} \sum_{k=-N}^{N} I(m,k) = e^{-\mu x} \sum_{k=-N}^{N} \sinh(\mu(2kt+x)) e^{-2(kx+k^2t)}.$$

The relation (8.3) is immediate.

For $\lambda = (t, x) \in \bar{C}_{aff}$, we consider the probability measure $\nu_{\lambda}^{(\mu)}$ on \mathbb{R}^2 carried by the line $\{(t, y), y \in \mathbb{R}\}$ which has the Laplace transform

(8.4)
$$\int e^{\langle \zeta, v \rangle} d\nu_{\lambda}^{(\mu)}(v) = \frac{e^{\langle \zeta, \lambda \rangle} \varphi_{\zeta_2 + \mu}(\lambda)}{\varphi_{\mu}(\lambda)}$$

for $\zeta = (\zeta_1, \zeta_2) \in \mathbb{R}^2$. This measure was introduced by Frenkel [16] in his study of orbital measures (see below Section 9.2).

Theorem 8.3. The conditional distribution of $B_t^{(\mu)}$ given $\sigma(\Lambda^{(\mu)}(s), s \leq t))$ is the probability measure $\nu_{\Lambda^{(\mu)}(t)}^{(\mu)}$.

Proof. By Lemma 8.1 and for $\mu_m = (\frac{m}{\pi}, \mu)$ one has

$$\mathbb{E}(e^{\langle \zeta,\tau_m W_t^{(\mu_m)} \rangle} | \sigma(\Lambda_m^{(\mu)}(s), s \le t)) = \frac{\psi_{\tau_m \zeta + \mu_m}^m(\Lambda_m^{(\mu)}(t))}{\psi_{\mu_m}^m(\Lambda_m^{(\mu)}(t))} \frac{h_m(\mu_m)}{h_m(\tau_m \zeta + \mu_m)} e^{\langle \tau_m \zeta, \Lambda_m^{(\mu)}(t) \rangle}.$$

Hence the formula

$$\mathbb{E}(e^{\langle \zeta, B_t^{(\mu)} \rangle} | \sigma(\Lambda^{(\mu)}(s), s \le t)) = \frac{\varphi_{\zeta_2 + \mu}(\Lambda^{(\mu)}(t))}{\varphi_{\mu}(\Lambda^{(\mu)}(t))} e^{\langle \zeta, \Lambda^{(\mu)}(t) \rangle}$$

is obtained by letting m go to infinity using Proposition 6.5 and Lemma 8.2.

8.2. Distribution of $L^{(\mu)}(\infty)$. We suppose that $\mu \neq 0, 1$, and write $L^{(\mu)}(\infty) = (0, D^{\mu}(\infty))$. The Laplace transform of $D^{\mu}(\infty)$ is, by (5.7), for $\tau > 0$,

$$\mathbb{E}(e^{-\tau D^{\mu}(\infty)}) = \prod_{n=0}^{\infty} \left(\left(1 + \frac{\tau}{(n+\mu)}\right) \left(1 - \frac{\tau}{(n+1-\mu)}\right) \right)^{-1}.$$

Using the relation

(8.5)
$$\frac{\Gamma(\alpha)\Gamma(\beta)}{\Gamma(\alpha+\gamma)\Gamma(\beta-\gamma)} = \prod_{n=0}^{+\infty} (1+\frac{\gamma}{n+\alpha})(1-\frac{\gamma}{n+\beta})$$

(Formula 8.325.1 of [19]) and $\Gamma(z)\Gamma(1-z) = \pi/\sin{(\pi z)}$, we obtain that

(8.6)
$$\mathbb{E}(e^{-\tau D^{\mu}(\infty)}) = \frac{\sin(\pi\mu)}{\sin(\pi(\mu+\tau))}.$$

In particular,

(8.7)
$$\mathbb{E}(D^{\mu}(\infty)) = \pi \cot(\pi \mu).$$

Corollary 8.4. The density of $D^{1/2}(\infty)$ is $1/(2\pi \cosh(x/2))$.

Proof. One uses that the Fourier transform of $1/(\pi \cosh x)$ is $1/\cosh(\lambda \pi/2)$. \Box

Notice that the distribution of $\sum_{n=1}^{\infty} \left(\frac{\varepsilon_{2n+1}}{n} - \frac{\varepsilon_{2n}}{n+1}\right)$ appears in Diaconis et al. [9].

8.3. **Distribution of** $\xi_1(\infty)$. We denote the Bessel process of dimension 3 with drift $\nu \geq 0$ by $\rho^{(\nu)}$. It is the norm of a Brownian motion in \mathbb{R}^3 with a drift of length ν .

Corollary 8.5. $\xi_1(\infty)$ has the same distribution as $\sup_{t\geq 0}(\varrho_t^{(1-\mu)}-t)$.

Proof. By Pitman and Rogers [34], $\varrho_t^{(1-\mu)}$ has the same distribution as $\mathcal{P}_{s_1}B_t^{(1-\mu)}$, so the claim follows from (2) of Theorem 7.4 when $\mu \neq 0, 1$, and by continuity when $\mu = 1$. For $\mu = 0, \xi_1(\infty) = +\infty$.

When $\mu = 1$, $\xi_0(\infty) = +\infty$ and

$$\xi_1(\infty) = \frac{1}{2} \sum_{n=1}^{+\infty} \frac{\varepsilon_{2n-1} + \varepsilon_{2n}}{n^2}.$$

Its distribution is studied in Biane et al. [2] where it is symbolized $\pi^2 S_2/4$. Its Laplace transform is given by, for $\tau \ge 0$,

$$\mathbb{E}(e^{-2\tau\xi_1(\infty)}) = \frac{\pi^2\tau}{\sinh^2(\pi\sqrt{\tau})},$$

and its distribution function is (cf. Table 1 in [2])

$$F(x) = 1 + 2\sum_{k=1}^{+\infty} (1 - 4k^2 x)e^{-2k^2 x}.$$

In this case the corollary is also given in Example 20 of Salminen and Yor [38]. When $\mu = 1/2$,

$$\xi_1(\infty) = \sum_{n=1}^{+\infty} \frac{2\varepsilon_n}{n(n+1)}.$$

Proposition 8.6. When $\mu = 1/2$, the Laplace transform of $\xi_1(\infty)$ is, for $\tau \ge 0$,

$$\mathbb{E}(e^{-\tau\xi_1(\infty)}) = \frac{2\pi\tau}{\cosh(\pi\sqrt{2\tau - 1/4})},$$

and its density is

$$\sum_{n=0}^{+\infty} (-1)^{n+1} n(n+1)(2n+1)e^{-n(n+1)x/2}.$$

In distribution,

$$\xi_1(\infty) = \sup_{n>0, i=1,2,3} \frac{\varepsilon_n^{(i)}}{n},$$

where the $\varepsilon_n^{(i)}$ are exponential independent random variables with parameter 1. Proof. One has

$$\mathbb{E}(e^{-\tau\xi_1(\infty)}) = \prod_{n=1}^{+\infty} (1 + \frac{2\tau}{n(n+1)})^{-1}.$$

Using the formula

$$\cosh \pi x = (1+4x^2) \prod_{n=1}^{+\infty} (1+\frac{x^2}{(n+1/2)^2}),$$

one obtains

$$\mathbb{E}(e^{-\tau\xi_1(\infty)}) = \frac{2\pi\tau}{\cosh(\pi\sqrt{2\tau - 1/4})}.$$

Let

$$g(x) = 2\sum_{n=0}^{+\infty} (-1)^n (2n+1)e^{-\frac{x}{2}(n+\frac{1}{2})^2}.$$

Since $1/\cosh\sqrt{2\tau}$ is the Laplace transform of $\pi g(\pi^2 x)/2$, (e.g. [2]), one has

$$\frac{2\tau\pi}{\cosh\pi\sqrt{2\tau-1/4}} = \int_0^\infty \tau e^{-\tau x} e^{x/8} g(x) \, dx = \int_0^\infty e^{-\tau x} (e^{x/8} g(x))' \, dx,$$

by an integration by parts. Computing the derivative $(e^{x/8}g(x))'$ gives the density. By integration, the distribution function of $\xi_1(\infty)$ is

$$\sum_{n=0}^{+\infty} (-1)^n (2n+1) e^{-n(n+1)x/2}$$

which equals $\prod_{n=1}^{\infty} (1 - e^{-nx})^3$ by a formula of Jacobi ([18], Theorem 357). \Box

9. Asymptotics for representations of the affine algebra $A_1^{(1)}$

We will show that some of the results we have obtained have a direct interpretation in term of semi-classical limits of highest weight representations of the affine Kac Moody Lie algebra $A_1^{(1)}$. The probability measure $\nu_{\lambda}^{(\mu)}$ defined in (8.4) is interpreted as a kind of Duistermaat Heckman measure. The affine string coordinates $\xi(\infty)$ describe the asymptotic behaviour of the large weights of Kashiwara's infinity crystal $B(\infty)$.

9.1. The Kac-Moody algebra $A_1^{(1)}$. We consider the affine Lie algebra $A_1^{(1)}$. For our purpose, we only need to define and consider a realization of a real Cartan subalgebra. We introduce two copies of \mathbb{R}^3 in duality,

$$\mathfrak{h}_{\mathbb{R}} = \operatorname{Span}_{\mathbb{R}} \{ c, \tilde{\alpha}_1, d \}, \ \mathfrak{h}_{\mathbb{R}}^* = \operatorname{Span}_{\mathbb{R}} \{ \Lambda_0, \alpha_1, \delta \},$$

where c = (1, 0, 0), $\tilde{\alpha}_1 = (0, 1, 0)$, d = (0, 0, 1), and $\Lambda_0 = (1, 0, 0)$, $\alpha_1 = (0, 2, 0)$, $\delta = (0, 0, 1)$ in \mathbb{R}^3 . We let $\tilde{\alpha}_0 = (1, -1, 0)$ and $\alpha_0 = (0, -2, 1)$, so that $c = \tilde{\alpha}_0 + \tilde{\alpha}_1$ and $\delta = \alpha_0 + \alpha_1$. Notice that these α_0 and α_1 project on the ones given in Section 5 by the projection on $\mathfrak{h}_{\mathbb{R}}^*/\mathbb{R}\delta$, identified with $\mathbb{R}\Lambda_0 \oplus \mathbb{R}\alpha_1$. With the notations of the introduction $\mathbb{R}\Lambda_0 \oplus \mathbb{R}\alpha_1 = \mathbb{R}\Lambda_0 \oplus \mathbb{R}\Lambda_1$. These notations are frequently used in the litterature in this context (for instance in Kac [24]). Usually α_0, α_1 are called the two positive simple roots of $A_1^{(1)}$ and $\tilde{\alpha}_0, \tilde{\alpha}_1$ their coroots. One considers the set of integral weights

$$P = \{ \lambda \in \mathfrak{h}_{\mathbb{R}}^* : \lambda(\tilde{\alpha}_i) \in \mathbb{Z}, i = 0, 1 \},\$$

and the set of dominant integral weights

$$P_{+} = \{ \lambda \in \mathfrak{h}_{\mathbb{R}}^{*} : \lambda(\tilde{\alpha}_{i}) \in \mathbb{N}, i = 0, 1 \}.$$

For a dominant integral weight λ , the character of the irreducible representation $V(\lambda)$ of $A_1^{(1)}$ with highest weight λ is defined as the formal series

(9.1)
$$\operatorname{char}_{\lambda} = \sum_{\beta \in P} \dim(V_{\beta}(\lambda)) e^{\beta},$$

where $V_{\beta}(\lambda)$ is the weight space corresponding to the weight β . If we let $e^{\beta}(h) = e^{\beta(h)}$ for $h \in \mathfrak{h}_{\mathbb{R}}$, and evaluate this formal series at h, the series converges absolutely or diverges, and it converges when $\delta(h) > 0$. When $\lambda = n\Lambda_0 + m\frac{\alpha_1}{2}$,

with $(m,n) \in \mathbb{N}^2$ such that $0 \leq m \leq n$, $a \in \mathbb{R}$, and b > 0, the Weyl-Kac character formula for $A_1^{(1)}$ is

(9.2)

$$\operatorname{char}_{\lambda}(a\tilde{\alpha}_{1} + bd) = \frac{\sum_{k \in \mathbb{Z}} \sinh(a(m+1) + 2ak(n+2))e^{-b(k(m+1) + k^{2}(n+2))}}{\sum_{k \in \mathbb{Z}} \sinh(a + 4ak)e^{-b(k+2k^{2})}}$$

For more details about affine Lie algebras and their representations, see Kac [24].

9.2. An analogue of the Duistermaat Heckman measure. Let

$$\lambda = t\Lambda_0 + x\alpha_1/2$$

where $x \in (0, t)$ and $\{\lambda_r, r \in \mathbb{N}^*\}$ be a sequence of dominant integral weights such that

$$\lim_{r \to +\infty} \lambda_r / r = \lambda.$$

We consider the irreducible module $V(\lambda_r)$ of highest weight λ_r . Since $V(\lambda_r)$ is infinite dimensional, we put a Boltzman factor on each weight. More precisely let $h_r^{(\mu)} = \frac{1}{r}(\mu \tilde{\alpha}_1 + 2d)$. As before $0 \le \mu \le 1$. We denote by $\gamma_r^{(\mu)}$ the probability measure on $\mathfrak{h}_{\mathbb{R}}^*$ given by

$$\gamma_r^{(\mu)} = \frac{1}{\operatorname{char}_{\lambda_r}(h_r^{(\mu)})} \sum_{\beta \in P} \dim(V_\beta(\lambda_r)) e^{\beta(h_r^{(\mu)})} \varepsilon_{r^{-1}\beta},$$

where $\varepsilon_{r^{-1}\beta}$ is the Dirac measure at $r^{-1}\beta$,

Proposition 9.1. The sequence of the push forward probabilities of $\{\gamma_r^{(\mu)}, r \ge 1\}$ by the quotient map from $\mathfrak{h}_{\mathbb{R}}^*$ to $\mathfrak{h}_{\mathbb{R}}^*/\mathbb{R}\delta$ converges to $\nu_{\lambda}^{(\mu)}$, when $r \to +\infty$.

Proof. For any $v \in \mathfrak{h}_{\mathbb{R}}$,

$$\int_{\mathfrak{h}_{\mathbb{R}}^{*}} e^{\beta(v)} \gamma_{r}^{(\mu)}(d\beta) = \frac{\operatorname{char}_{\lambda_{r}}(r^{-1}v + h_{r}^{(\mu)})}{\operatorname{char}_{\lambda_{r}}(h_{r})}.$$

For $\tau \in \mathbb{R}$ and $v = \tau \tilde{\alpha}_1$, the numerator of the character $\operatorname{char}_{\lambda_r}(r^{-1}v + h_r^{(\mu)})$ given by (9.2) converges to

$$e^{(\tau+\mu)x}\sin(\pi(\tau+\mu))\varphi_{(\tau+\mu)}(t,x)$$

when r goes to infinity. Besides Lemma 2.7 implies that the denominator of this character is equivalent to

$$\frac{\sqrt{\pi r}}{2}e^{-\frac{1}{4}r\pi^2}\sin(\pi(\tau+\mu))$$

Therefore

$$\lim_{r \to +\infty} \int_{\mathfrak{h}_{\mathbb{R}}^*} e^{\tau\beta(\tilde{\alpha}_1)} \gamma_r^{(\mu)}(d\beta) = e^{\tau x} \frac{\varphi_{\tau+\mu}(t,x)}{\varphi_{\mu}(t,x)}.$$

Then the proposition follows from Theorem 8.3.

The Duistermaat Heckman measure for a compact connected Lie group is an approximation of the distribution of the weights of an irreducible finite dimensional representation when its highest weight is large (see Heckman [20]). So the theorem above shows that $\nu_{\lambda}^{(\mu)}$ is an analogue of a normalized DH-measure. For a compact connected Lie group, the Duistermaat Heckman measure also appears

as the image of the Liouville measure on a coadjoint orbit by the projection on the dual of a Cartan subalgebra. Frenkel has shown in [16] that the distribution of a Brownian motion on $\mathfrak{su}(2)$ indexed by the time in [0, 1], given the conjugacy class of the endpoint of its wrapping on SU(2), plays the role of a normalized Liouville measure on a coadjoint orbit of the loop group L(SU(2)). The image of this measure by the moment map associated to the action of a maximal torus of SU(2) on the coadjoint orbit is the measure $\nu_{\lambda}^{(\mu)}$ when $\mu = 0$ (see Frenkel [16] and also Defosseux [12]).

9.3. Asymptotics for the crystal $B(\infty)$ of $A_1^{(1)}$. The infinity crystal $B(\infty)$ of Kashiwara [26] is the crystal of the Verma module with highest weight 0 of $A_1^{(1)}$. It gives a combinatorial graph describing precisely this module. This crystal is important since any irreducible highest weight crystal may be obtained from $B(\infty)$. It is shown in Kashiwara [26], and more explicitly in Nakashima and Zelevinski [31], that using string parametrizations, a realization of $B(\infty)$ is given by

$$B(\infty) = \{x \in \mathbb{N}^{\mathbb{N}}; \text{ for some } n \in \mathbb{N}, \frac{x_1}{1} \ge \frac{x_2}{2} \ge \dots \ge \frac{x_n}{n} > 0, x_k = 0 \text{ for } k > n\}.$$

Notice that the only condition on x_0 is $x_0 \in \mathbb{N}$. For $x \in B(\infty)$, we let

$$\omega(x) = \sum_{k=0}^{+\infty} x_k \alpha_k \in \mathfrak{h}_{\mathbb{R}}^*$$

By [31], $-\omega(x)$ is the weight of x in the crystal $B(\infty)$. For $\tilde{\rho} = \tilde{\alpha}_1/2 + 2d$, we let

$$s(x) = \sum_{k=0}^{+\infty} x_k = \omega(x)(\tilde{\rho}).$$

The character $char_{\infty}$ of the Verma module of highest weight 0 is defined as in (9.1) and is given by (see Kac [24], (9.7.2))

(9.3)
$$\operatorname{char}_{\infty} = \prod_{\beta \in R_+} (1 - e^{-\beta})^{-1},$$

where

(9.4)
$$R_{+} = \{\alpha_0 + n\delta, \alpha_1 + n\delta, (n+1)\delta, n \in \mathbb{N}\}$$

is the set of positive roots of $A_1^{(1)}$. As previously, if we let $e^{\beta}(h) = e^{\beta(h)}$ for $h \in \mathfrak{h}_{\mathbb{R}}$, and evaluate the formal character at h, it converges when $\delta(h) > 0$. Let $r \in \mathbb{R}_+^*$. On each element x of the crystal $B(\infty)$ we put the Boltzman weight $e^{-s(x)/r}$. We introduce the probability distribution β_r on $B(\infty)$ by

$$\beta_r = \frac{1}{Z_r} \sum_{x \in B(\infty)} e^{-s(x)/r} \varepsilon_x$$

when ε_x is the Dirac mass at x and (by Kashiwara [26])

$$Z_r = \sum_{x \in B(\infty)} e^{-s(x)/r} = \operatorname{char}_{\infty}(\tilde{\rho}/r).$$

The following theorem indicates that the affine Brownian model describes a kind of continuous version of the infinity crystal $B(\infty)$ for the affine Lie algebra $A_1^{(1)}$. Let $j_r: B(\infty) \to \mathbb{R}^{\mathbb{N}}$ be given by $j_r((x_k, k \ge 0)) = (x_k/r, k \ge 0).$

Theorem 9.2. The image of β_r by the map j_r converges in distribution to the distribution of the Verma parameter $\xi(\infty)$ of $B^{(1/2)}$, when $r \to +\infty$.

Proof. We first notice that

$$A_{\infty} = \{ (\lambda_1, \lambda_2, \cdots); (0, \lambda_1, \lambda_2, \cdots) \in B(\infty) \}$$

is the set of anti-lecture hall compositions defined in Corteel and Savage [7]. For $\lambda = (\lambda_1, \lambda_2 \cdots, \lambda_n, 0, 0, \cdots) \in A_{\infty}$, let $|\lambda| = \sum_{k=1}^n \lambda_k$. Then, see Corteel et al. [7, 8] and (1.2) and (1.3) in Chen et al. [6], for $0 \le q < 1, k \in \mathbb{N}$,

(9.5)
$$\sum_{\lambda \in A_{\infty}} q^{|\lambda|} = \frac{(-q;q)_{\infty}}{(q^2;q)_{\infty}}$$

(9.6)
$$\sum_{\lambda \in A_{\infty}, \lambda_1 \le k} q^{|\lambda|} = \frac{(-q;q)_{\infty}(q;q^{k+2})_{\infty}(q^{k+1};q^{k+2})_{\infty}(q^{k+2};q^{k+2})_{\infty}}{(q;q)_{\infty}},$$

where $(a;q)_{\infty} = \prod_{n=0}^{\infty} (1 - aq^n)$. Let $X^{(r)} = (X_k^{(r)}; k \ge 0)$ be a random element with distribution β_r . First it is clear that $X_0^{(r)}/r$ converges in distribution to $\xi_0(\infty)$. Let us now prove that $X_1^{(r)}/r$ converges in distribution to $\xi_1(\infty)$. It follows from (9.5) and (9.6) that, for $a \ge 0$ and $q = e^{-1/r}$,

$$\mathbb{P}(X_1^{(r)} \le ar) = \frac{\sum_{\{x \in B(\infty); x_1 \le [ar]\}} q^{s(x)}}{\sum_{\{x \in B(\infty)\}} q^{s(x)}} = \frac{\sum_{\{\lambda \in A_\infty; \lambda_1 \le [ar]\}} q^{|\lambda|}}{\sum_{\{\lambda \in A_\infty\}} q^{|\lambda|}} = \frac{1}{1-q} (q; e^{-[ar]/r} q^2)_\infty (e^{-[ar]/r} q; e^{-[ar]/r} q^2)_\infty (e^{-[ar]/r} q^2)_\infty (e^{-$$

Since $(q; e^{-[ar]/r}q^2)_{\infty}$ is equivalent to $(1-q)(e^{-a}; e^{-a})_{\infty}$ when q tends to 1, one has, by Proposition 8.6,

$$\lim_{r \to \infty} \mathbb{P}(X_1^{(r)} \le ra) = (e^{-a}; e^{-a})_{\infty}^3 = \prod_{n=1}^{\infty} (1 - e^{-na})^3 = \mathbb{P}(\xi_1(\infty) \le a).$$

So $X_1^{(r)}/r$ converges to $\xi_1(\infty)$. We now consider the full sequence $X_k^{(r)}/r, k \in \mathbb{N}$. For any r and $n \ge 1$, one has

$$\frac{X_1^{(r)}}{1} \ge \frac{X_2^{(r)}}{2} \ge \dots \ge \frac{X_n^{(r)}}{n} \ge 0$$

which implies that for any $n \in \mathbb{N}$, the collection of distributions of

$$(\frac{1}{r}X_1^{(r)},\cdots,\frac{1}{r}X_n^{(r)})_{r>0}$$

is tight since $X_1^{(r)}/r$ converges in distribution. By Cantor's diagonal argument, we construct an increasing sequence $\varphi(r) \in \mathbb{N}, r \in \mathbb{N}$, such that the random variables $\frac{1}{\varphi(r)}X_k^{(\varphi(r))}, k \ge 0$, converge in finite dimensional distribution when r goes to infinity. Let us denote by $(R_k, k \ge 0)$ the limit, and let us prove that $\frac{R_k}{k}$ has the same distribution as $\sum_{n=k}^{+\infty} \frac{2\varepsilon_n}{n(n+1)}$ where the ε_n 's are independent exponential random variables with parameter 1. For $x \in B(\infty)$, one has, since $s(x) = x_0 + \frac{1}{2} \sum_{k=2}^{+\infty} (kx_{k-1} - (k-1)x_k)$,

$$\mathbb{P}(X^{(r)} = x) = \frac{e^{-\frac{1}{r}x_0}}{\operatorname{char}_{\infty}(r^{-1}\tilde{\rho})} \prod_{k=2}^{\infty} e^{-\frac{1}{2r}(kx_{k-1} - (k-1)x_k)} \mathbb{1}_{\{\frac{x_{k-1}}{k-1} \ge \frac{x_k}{k}\}}$$
$$= \frac{e^{-\frac{1}{r}x_0}}{\operatorname{char}_{\infty}(r^{-1}\tilde{\rho})} \prod_{k=1}^{\infty} e^{-\frac{k+1}{2r}(x_k - \lceil\frac{k}{k+1}x_{k+1}\rceil)} e^{-\frac{1}{2r}(k\lceil\frac{k-1}{k}x_k\rceil - (k-1)x_k)} \mathbb{1}_{\{x_k \ge \lceil k\frac{x_{k+1}}{k+1}\rceil\}},$$

where $\lceil \cdot \rceil$ is the ceiling function. Let $Y_k^{(r)}$, $k = 0, \dots, n$, be independent geometric random variables with values in \mathbb{N} , where $Y_0^{(r)}$ has the parameter 1 and $Y_k^{(r)}$ has the parameter $e^{-\frac{k+1}{2r}}$ when $k \ge 1$. As,

$$e^{-\frac{k}{2r}} \le e^{-\frac{1}{2r}(k\lceil \frac{k-1}{k}x_k\rceil - (k-1)x_k)} \le 1,$$

one obtains that for $t_0, \cdots, t_n \ge 0$,

$$C(r,n)(1+O(\frac{1}{r})) \le \frac{\mathbb{P}(X_0^{(r)} \le t_0, \cdots, X_n^{(r)} - \lceil \frac{n}{n+1} X_{n+1}^{(r)} \rceil \le t_n)}{\mathbb{P}(Y_0^{(r)} \le t_0, \cdots, Y_n^{(r)} \le t_n)} \le C(r,n),$$

where C(r, n) is independent of t_0, \dots, t_n , and tends to 1 when r tends to $+\infty$. This proves that for any n,

$$\frac{1}{r}X_0^{(r)}, \ \frac{1}{r}(X_k^{(r)} - \lceil \frac{k}{k+1}X_{k+1}^{(r)}\rceil), \quad 1 \le k \le n$$

converge jointly to $\varepsilon_0, \frac{2}{k+1}\varepsilon_k, 1 \le k \le n$, when r goes to infinity. Besides

$$\lim_{r \to \infty} \frac{1}{\varphi(r)} \left(X_k^{(\varphi(r))} - \left\lceil \frac{k}{k+1} X_{k+1}^{(\varphi(r))} \right\rceil \right) = R_k - \frac{k}{k+1} R_{k+1}.$$

Thus, for $k \ge 1$, $R_k - \frac{k}{k+1}R_{k+1}$ are independent random variables with the same distribution as $\frac{2}{k+1}\varepsilon_k$. The positive sequence R_k/k is decreasing. Let S be its limit. We have the identity in distribution, for all $k \ge 1$,

$$\frac{R_k}{k} = \sum_{n=k}^{+\infty} \frac{2\varepsilon_n}{n(n+1)} + S$$

We have proved that in distribution $R_1 = \xi_1(\infty)$, so S = 0 which finishes the proof.

Recall (Corollary 8.4) that $L^{(1/2)}(\infty) = D^{1/2}(\infty)\alpha_1/2$ where the density of $D^{1/2}(\infty)$ is $1/(\pi \cosh x)$, and that $-\omega(x)$ is the weight of $x \in B(\infty)$.

Proposition 9.3. When r goes to infinity, in distribution, if $X^{(r)}$ has the distribution β_r ,

- (1) in the quotient space $\mathfrak{h}_{\mathbb{R}}^*/\mathbb{R}\delta$, the normalized weights $\omega(X^{(r)})/r$ converge to $L^{(1/2)}(\infty)$,
- (2) the coordinate of $\omega(X^{(r)})/r$ along δ goes to $+\infty$.

Proof. One has for any $u \in \mathfrak{h}_{\mathbb{R}}$,

$$\mathbb{E}(e^{-\omega(X^{(r)})(u)}) = \frac{\operatorname{char}_{\infty}(u+r^{-1}\tilde{\rho})}{\operatorname{char}_{\infty}(r^{-1}\tilde{\rho})}.$$

In view of (9.4), the expression (9.3) of the character gives the Laplace transform of $\omega(X^{(r)})$ and shows that, in distribution,

$$\omega(X^{(r)}) = \sum_{n \ge 0} (G_0(n)(\alpha_0 + n\delta) + G_1(n)(\alpha_1 + n\delta) + G_2(n)(n+1)\delta).$$

Here $G_i(n)$, $i = 0, 1, 2, n \in \mathbb{N}$, are independent random variables such that $G_0(n), G_1(n)$ and $G_2(n)$, are geometrically distributed with respective parameter $e^{-(\alpha_0+n\delta)(r^{-1}\tilde{\rho})}$, $e^{-(\alpha_1+n\delta)(r^{-1}\tilde{\rho})}$, and $e^{-(n+1)\delta(r^{-1}\tilde{\rho})}$, i.e. $e^{-2(n+1/2)/r}$, $e^{-2(n+1/2)/r}$ and $e^{-2(n+1)/r}$. The proposition follows easily.

Due to the Lévy correction, the image by ω of the limit in law of $\frac{X^{(r)}}{r}$ is not equal to limit in law of $\omega(\frac{X^{(r)}}{r})$, modulo δ .

References

- [1] R. Bellman. A brief introduction to Theta functions, Holt, Rinehart and Winston, 1961.
- [2] Ph. Biane, J.W. Pitman and M. Yor. Probability laws related to the Jacobi Theta and Riemann Zeta functions, and Brownian excursions. Bulletin of the American Mathematical Society 38 (2001) 435–465.
- [3] Ph. Biane, Ph. Bougerol and N. O'Connell. Littlemann paths and Brownian paths. Duke Math. J. 130 (2005) 127–167.
- [4] Ph. Biane, Ph. Bougerol and N. O'Connell. Continuous crystal and Duistermaat-Heckmann measure for Coxeter groups. Adv. Maths. 221 (2009) 1522–1583.
- [5] A. Borodin and L. Petrov. Integrable probability: From representation theory to Macdonald processes. Probability Surveys 11 (2014) 1–58.
- [6] W.Y.C. Chen, D.D.M. Sang and D.Y.H. Shi. Anti-lecture hall compositions and overpartitions. J. Combin. Theory Ser. A 118 (2011) 1451–1464.
- [7] S. Corteel and C. Savage. Anti-lecture hall compositions. Discrete Math. 263 (2003) 275– 280.
- [8] S. Corteel, J. Lovejoy and C. Savage. Anti-lecture hall compositions and Andrews' generalization of the Watson-Whipple transformation. J. Combin. Theory Ser. A 134 (2015) 188–195.
- [9] P. Diaconis, S. Janson and R. Rhoades. Note on a partition limit theorem for rank and crank. Bull. London Math. Soc. 45 (2013) 551–553.
- [10] M. Defosseux. The affine Lie algebra $\hat{s}l_2$ and a conditioned space-time Brownian motion, arXiv:1401.3115 (2014).
- [11] M. Defosseux. Affine Lie algebras and conditioned space-time Brownian motions in affine Weyl chambers. Probab. Theory Relat. Fields 165 (2015) 1–17.
- [12] M. Defosseux. Kirillov-Frenkel character formula for loop groups, radial part and Brownian sheet. Ann. of Probab. 47 (2019) 1036–1055.
- [13] N. Demni. Radial Dunkl processes associated with dihedral systems. Séminaire de Probabilités, 43, Lecture Notes in Math. Springer (2009) 153–169.
- [14] C.F. Dunkl and Y. Xu. Orthogonal Polynomials of several variables. Encyclopedia of Mathematics and its Applications 155, 2nd Edition, Cambridge University Press, 2014.
- [15] J. J. Duistermaat and G. J. Heckman. On the variation in the cohomology of the symplectic form of the reduced phase space. Invent. Math. 69 (1982) 259–268.
- [16] I.B. Frenkel. Orbital theory for affine Lie algebras. Invent. Math. 77 (1984) 301–352.
- [17] L. Gallardo and M. Yor. Some new examples of Markov processes which enjoy the timeinversion property. Probab. Theory Relat. Fields 132 (2005) 150–162.
- [18] G.H. Hardy and E.M. Wright. An Introduction to the theory of numbers. Sixth Edition, Oxford University Press, 2008.
- [19] I.S. Gradshteyn and I.M. Ryzhik. Tables of integrals, series and products. Seventh Edition, Academic Press, 2015.
- [20] G. J. Heckman. Projections of orbits and asymptotic behavior of multiplicities for compact connected Lie groups. Invent. Math. 67 (1982) 333–356.

- [21] D.G. Hobson and W. Werner. Non-colliding Brownian motions on the circle. Bulletin of the London Mathematical Society 28 (1996) 643–650.
- [22] J. E. Humphreys. *Reflection groups and Coxeter groups*. Cambridge Studies in Advanced Mathematics 29. Cambridge University Press, 1990.
- [23] K. Ito and H.P. McKean. Diffusion processes and their sample paths, Springer Verlag, 1965.
- [24] V.G. Kac. Infinite dimensional Lie algebras. Third edition, Cambridge University Press, 1990.
- [25] I. Karatzas and S.E. Shreve Brownian motion and stochastic calculus, Springer Verlag, 1991.
- [26] M. Kashiwara. The crystal base and Littelmann's refined Demazure character formula. Duke Math. J. 71 (1993) 839–858.
- [27] M. Kashiwara. On Crystal Bases. Representations of groups (Banff, AB, 1994), CMS Conf. Proc., Amer. Math. Soc. 16 (1995) 155–191.
- [28] F.B. Knight. Brownian local times and taboo processes. Trans. Amer. Math. Soc. 143 (1969) 173–185.
- [29] P. Littelmann. Paths and root operators in representation theory. Annals of Mathematics 142 (1995) 499–525.
- [30] P. Littelmann. Cones, crystals, and patterns. Transform. Groups 3 (1998) 145–179.
- [31] T. Nakashima and A. Zelevinsky. Polyhedral realizations of crystal bases for quantized Kac-Moody algebras. Adv. Math. 131 (1997) 253–278.
- [32] T. Nakashima. Polyhedral realizations of crystal bases for integrable highest weight modules. Journal of Algebra 219 (1999) 571–597.
- [33] J.W. Pitman. One-dimensional Brownian motion and the three-dimensional Bessel process. Adv. Appl. Probab. 7 (1975) 511–526.
- [34] J.W. Pitman and L.C.G. Rogers. Markov functions. Ann. Prob. 9 (1981) 573–582.
- [35] A. Renyi. On the theory of order statistics. Acta Mathematica Academiae Scientiarum Hungaricae 4 (1953) 191–231.
- [36] D. Revuz and M. Yor. Continuous martingales and Brownian motion, Second edition Springer-Verlag. (1994).
- [37] M. Roesler and M. Voit. Markov processes related with Dunkl operators. Adv. in Appl. Math. 21 (1998) 575–643.
- [38] P. Salminen and M. Yor. On hitting times of affine boundaries by reflecting Brownian motion and Bessel processes. Periodica Mathematica Hungarica 62 (2011) 75–101.

PH.B. : SORBONNE UNIVERSITÉ, LPSM-UMR 8001, PARIS, FRANCE *Email address*: philippe.bougerol@upmc.fr

M. D. : UNIVERSITÉ PARIS DESCARTES, MAP5-UMR 8145, PARIS, FRANCE *Email address*: manon.defosseux@parisdescartes.fr