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A class of cutoff functions for non-trivial Riemann zeros

Yu Li

Abstract

Abstract: In this paper we construct a pseudo ermitian operator \tilde{H} on the Hilbert space $L^2[\frac{1}{2},1]$ and show that the Hurwitz zeta functions over the interval $[\frac{1}{2},1]$ are eigenfunctions of \tilde{H} for all complex numbers.

Keywords. Riemann hypothesis, non-trivial Riemann zero, Pseudo hermitian operator, Berry-Keating operator

1 Introduction

One of the most attracting problems in Mathematics is Riemann hypothesis and it states that all the non-trivial zeros of Riemann zeta function

$$\zeta(s) = \sum_{n>1} \frac{1}{n^s} = \frac{1}{\Gamma(s)} \int_0^\infty \frac{t^{s-1}}{e^t - 1} dt$$

have the real part $\frac{1}{2}$.

One proposed way of proving Riemann hypothesis, known as Hilbert-Pólya conjecture, is to find an unbounded self-adjoint operator such that the imaginary part of Riemann zeros corresponds to the eigenvalues of this operator. It has been suggested [5] that the possible approach to solve Hilbert-Póplya conjecture lies in a quantization of the classical Hamiltonian H = XP, where P is the canonical momentum operator associated with position operator X. Inspired of the previous works, an operator \tilde{H} similar to quantum analogy was proposed [4] and asymptotic analysis was performed [3]. However, the momentum operator P in this case has no self-adjoint extension [2].

2 Hurwitz Zeta functions

The Hurwitz zeta function is originally defined for complex arguments s with $\mathcal{R}(s) > 1$ and x > 0 by

$$\zeta(s,x) = \sum_{n \ge 0} \frac{1}{(n+x)^s}$$

Hurwitz zeta function has an analytic continuation on \mathbb{C} with the only singularity s=1 and it can be represented [1, 10]

$$\zeta(s,x) = -\frac{\Gamma(1-s)}{2\pi i} \int_{\mathcal{C}} \frac{\xi^{s-1} e^{x\xi}}{e^{\xi} - 1} d\xi, \quad s \neq 1$$
 (1)

where C is the contour around the negative real axis and it starts at $-\infty$, encircles the origin once in the positive direction without enclosing any of the points $\xi = \pm 2ni, \pm 4ni, ...$ and returns to $-\infty$ [10]. The famous Riemann zeta function is the special case of Hurwitz zeta function with x = 1

$$\zeta(s,1) = \zeta(s)$$

and also with $x = \frac{1}{2}$

$$\zeta(s, \frac{1}{2}) = (2^s - 1)\zeta(s) \tag{2}$$

The partial derivative of Hurwitz zeta function with respect to x

$$\frac{\partial \zeta(s,x)}{\partial x} = -s\zeta(s+1,x), \quad s \neq 0, 1; x > 0$$
(3)

Let s denote a point in critical strip $0 < \mathcal{R}(s) < 1$ and we define a cutoff function ζ_s of Hurwitz function on interval $[\frac{1}{2}, 1]$

$$\zeta_s(x) := \zeta(s, x), \quad \frac{1}{2} \le x \le 1 \tag{4}$$

The functions ζ_s is continuous and even smooth on the closed interval $\left[\frac{1}{2},1\right]$ and have finite norm and belong to the Hilbert space, $\zeta_s \in L^2\left[\frac{1}{2},1\right]$. If a complex number z is a non-trivial Riemann zero,

$$\zeta_z(1) = 0$$

then the function ζ_z vanishes at $x = \frac{1}{2}$ from (2)

$$\zeta_z(\frac{1}{2}) = 0$$

3 Operators in Hilbert space $L^2[\frac{1}{2}, 1]$

We confine our discussion over the region $[\frac{1}{2},1]$. Let X be the position operator on $L^2[\frac{1}{2},1]$ and P_0 the momentum operator $P_0=-i\frac{d}{dx}$ densely defined on $L^2[\frac{1}{2},1]$ with domain

$$D(P_0) = \left\{ f \in C^{\infty}[\frac{1}{2}, 1] : f(\frac{1}{2}) = f(1) \right\}$$
 (5)

[6] and it has a unique self-adjoint extension $\overline{P_0} = P$. The closure of the graph of P_0 is the graph of P

$$\Gamma(P) = \overline{\Gamma(P_0)} = L^2[\frac{1}{2}, 1] \times L^2[\frac{1}{2}, 1]$$

From the Stone's theorem on one-parameter unitary groups, $\{e^{itP}\}_t$ is a strongly continuous one-parameter group of unitary operators on $L^2[\frac{1}{2},1]$. Thus, both operators $e^{\pm iP}-I$ are bijective and bounded, and then the bounded inverses $(e^{\pm iP}-I)^{-1}$ exist, from the bounded inverse theorem in functional analysis.

We construct an operator \tilde{H} on $L^2[\frac{1}{2},1]$

$$\tilde{H} = (e^{iP} - I)^{-1} (XP + PX) (e^{iP} - I)$$

It holds obviously for $\frac{1}{2} \le x \le 1$ and $0 < \mathcal{R}(s) < 1$

$$(XP + PX)x^{-s} = i(2s - 1)x^{-s}$$
(6)

The operator \tilde{H} is pseudo hermitian [7, 8, 9], i.e. there exists an invertible and self-adjoint operator η such that the adjoint operator \tilde{H}^{\dagger} can be represented

$$\tilde{H}^{\dagger} = \eta \tilde{H} \eta^{-1}$$

Proposition 3.1. \tilde{H} is pseudo hermitian and all eigenvalues are real.

Proof. Let $\eta = (e^{iP} - I)^{\dagger}(e^{iP} - I)$ and it is obviously self-adjoint, then

$$\tilde{H}^{\dagger} = \eta \tilde{H} \eta^{-1}$$

The operator η is positive definite and self-adjoint by construction, the operator \tilde{H} is hermitian in the new inner product with respect to η

$$\langle \tilde{H}\phi,\varphi\rangle_{\eta}=\langle \tilde{H}\phi,\eta\varphi\rangle=\langle \phi,\eta\tilde{H}\eta^{-1}\eta\varphi\rangle=\langle \phi,\tilde{H}\varphi\rangle_{\eta},\quad \phi,\varphi\in L^{2}[\frac{1}{2},1]$$

Then all eigenvalues of \tilde{H} are real.

Proposition 3.2. For any complex number z, the function $\zeta_z(x) - \zeta_z(\frac{1}{2})$ is an eigenfunction of \tilde{H} with eigenvalue i(2z-1)

$$\tilde{H}\left(\zeta_z(x) - \zeta_z(\frac{1}{2})\right) = i(2z - 1)\left(\zeta_z(x) - \zeta_z(\frac{1}{2})\right)$$

Proof. We consider for non-trivial Riemann zeros z the term

$$\frac{1}{iP} \left(\frac{iP}{e^{iP} - I} \right) x^{-z}, \quad \frac{1}{2} \le x \le 1$$

where $\frac{1}{iP}$ is defined as an integral operator on $L^2[\frac{1}{2},1]$ with boundary $x=\frac{1}{2}$

$$\frac{1}{iP}f(x) := \int_{\frac{1}{2}}^{x} f(\xi)d\xi, \quad f \in L^{2}[\frac{1}{2}, 1]$$

and $\frac{iP}{e^{iP}-I}f$ is defined as a Bernoulli generating function

$$\frac{iP}{e^{iP} - I}f(x) := \sum_{n>0} \frac{B_n}{n!} (iP)^n f(x), \quad f \in L^2[\frac{1}{2}, 1]$$

Notice

$$(iP)^n x^{-z} = \frac{\Gamma(-z+1)}{\Gamma(-z+1-n)} x^{-z-n}$$

We use Hankel loop contour integral formula [10, 11]

$$\frac{1}{\Gamma(-z+1-n)} = \frac{1}{2\pi i} \int_{\mathcal{C}} \xi^{-(1-z-n)} e^{\xi} d\xi$$

where the integration is done along the Hankel-Bromwich contour and it begins at $-\infty$, circles the origin once in the positive direction, and returns to $-\infty$

$$(iP)^{n}x^{-z} = \Gamma(1-z)\left(\frac{1}{2\pi i}\int_{\mathcal{C}}\xi^{-(1-z-n)}e^{\xi}d\xi\right)x^{-z-n}$$
$$= \frac{\Gamma(1-z)}{2\pi i}\int_{\mathcal{C}}\left(\frac{\xi}{x}\right)^{z+n-1}e^{\xi}d\frac{\xi}{x}$$

and then

$$\frac{iP}{e^{iP} - I} x^{-z} = \sum_{n \ge 0} \frac{B_n}{n!} (iP)^n x^{-z}$$
$$= \frac{\Gamma(1-z)}{2\pi i} \sum_{n \ge 0} \int_{\mathcal{C}} \frac{B_n}{n!} \left(\frac{\xi}{x}\right)^{z+n-1} e^{\xi} d\frac{\xi}{x}$$

Let $u = \frac{\xi}{x}$ and $C' = \frac{C}{x}$ is still a loop around the negative real axis. Uniform convergence on closed interval justifies the interchange of integral and sum. Then Hurwitz zeta function can be reproduced (1)

$$\frac{iP}{e^{iP} - I}x^{-z} = \frac{\Gamma(1-z)}{2\pi i} \int_{\mathcal{C}'} \frac{u^z}{e^u - 1} e^{ux} du = -z\zeta(z+1, x)$$

and the inverse

$$\frac{iP}{e^{iP} - I}x^{-z} = -z\zeta(z+1, x)$$

Then from (3)

$$\frac{1}{iP} \left(\frac{iP}{e^{iP} - I} \right) x^{-z} = \int_{\frac{1}{2}}^{x} -z\zeta(z+1,\xi)d\xi$$

$$= \zeta(z,x) - \zeta(z,\frac{1}{2}) \tag{7}$$

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