



HAL
open science

Determination of the magnetic structure of CePt 2 In 7 by means of neutron diffraction

Matthias Raba, Eric Ressouche, N. Qureshi, Claire Colin, V. Nassif, S. Ota,
Y. Hirose, R. Settai, Pierre Rodière, I. Sheikin

► **To cite this version:**

Matthias Raba, Eric Ressouche, N. Qureshi, Claire Colin, V. Nassif, et al.. Determination of the magnetic structure of CePt 2 In 7 by means of neutron diffraction. *Physical Review B*, 2017, 95 (16), pp.161102. 10.1103/PhysRevB.95.161102 . hal-01828364

HAL Id: hal-01828364

<https://hal.science/hal-01828364v1>

Submitted on 3 Jul 2018

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Determination of the magnetic structure of CePt₂In₇ by means of neutron diffraction

M. Raba,^{1,2,3,*} E. Ressouche,⁴ N. Qureshi,⁵ C. V. Colin,^{2,3} V. Nassif,^{2,3} S. Ota,⁶ Y. Hirose,⁷
R. Settai,⁷ P. Rodière,^{2,3,†} and I. Sheikin^{1,‡}

¹Laboratoire National des Champs Magnétiques Intenses (LNCMI-EMFL), CNRS, UGA, F-38042 Grenoble, France

²Institut Néel, Université Grenoble Alpes, F-38000 Grenoble, France

³CNRS, Institut Néel, F-38000 Grenoble, France

⁴INAC, CEA and Université Grenoble Alpes, CEA Grenoble, F-38054 Grenoble, France

⁵Institut Laue Langevin, 71 rue des Martyrs, BP156, F-38042 Grenoble Cedex 9, France

⁶Graduate School of Science and Technology, Niigata University, Niigata 950-2181, Japan

⁷Department of Physics, Niigata University, Niigata 950-2181, Japan

(Received 8 February 2017; revised manuscript received 20 March 2017; published 3 April 2017)

The magnetic structure of the heavy fermion antiferromagnet CePt₂In₇ is determined using neutron diffraction. We find a magnetic wave vector $\mathbf{q}_M = (1/2, 1/2, 1/2)$, which is temperature independent up to $T_N = 5.5$ K. A staggered moment of $0.45(1)\mu_B$ at 2 K resides on the Ce ion. The nearest-neighbor moments in the tetragonal basal plane are aligned antiferromagnetically. The moments rotate by 90° from one CeIn₃ plane to another along the c axis. A much weaker satellite peak with an incommensurate magnetic wave vector $\mathbf{q}_M = (1/2, 1/2, 0.47)$ seems to develop at low temperature. However, the experimental data available so far are not sufficient to draw a definitive conclusion about the possible coexistence of commensurate and incommensurate magnetic structures in this material.

DOI: [10.1103/PhysRevB.95.161102](https://doi.org/10.1103/PhysRevB.95.161102)

CePt₂In₇ is a recently discovered heavy fermion compound that belongs to the same family as the well-studied CeIn₃ and CeMIn₅ ($M = \text{Co, Rh, Ir}$). The spacing between Ce-In planes in CePt₂In₇ is drastically increased [1] as compared to its CeMIn₅ counterparts, implying a more two-dimensional crystal structure. Expectedly, the Fermi surface of CePt₂In₇ is also much more two dimensional [2]. This compound crystallizes in the body-centered-tetragonal structure (space group $I4/mmm$) with a unit cell considerably elongated along the c axis.

CePt₂In₇ undergoes an antiferromagnetic (AF) transition at $T_N = 5.5$ K [3–6]. Recent electrical resistivity and calorimetry measurements under pressure on single crystals of CePt₂In₇ revealed a quantum critical point at a critical pressure $P_c \approx 3.2$ GPa, where the AF order is completely suppressed [7]. A superconducting dome with the highest transition temperature $T_c \simeq 2.1$ K is observed around P_c [4,7,8], suggesting that critical AF fluctuations may mediate the Cooper pairing. Nuclear magnetic and quadrupole resonance (NMR and NQR, respectively) measurements on single crystals under pressure [9] suggest that P_c is not the only relevant pressure for this material. Indeed, a localized to itinerant crossover of the $4f$ electron of Ce occurs within the AF state at $P^* \approx 2.4$ GPa, approximately the pressure where superconductivity first emerges in single crystals.

While the magnetic structure of the cubic CeIn₃ is characterized by a simple commensurate ordering wave vector $(1/2, 1/2, 1/2)$ [10,11], that of the more two-dimensional CeRhIn₅ is more complicated. Its magnetically ordered ground state is an incommensurate helicoidal phase with the propagation vector $\mathbf{q}_M = (1/2, 1/2, 0.297)$ and the magnetic moment in the basal plane of the tetragonal structure [12–16].

In CePt₂In₇, the magnetic structure of its AF ground state is still an open question. The existing reports on this matter are controversial. Indeed, NQR studies performed on polycrystalline samples [5] suggest that antiferromagnetism is commensurate in this material. The same conclusion was drawn from positive muon-spin rotation and relaxation (μ^+ SR) measurements also performed on polycrystalline samples [17]. On the contrary, the NQR spectra obtained on single crystals are consistent with the coexistence of an incommensurate and commensurate AF component of the magnetic structure [9,18]. The commensurate AF order first occurs just below T_N , then, at lower temperatures of about 3 K, incommensurate AF order gradually grows in. This probably accounts for the observation of only commensurate magnetism in polycrystalline NQR measurements, which were performed at temperatures down to 4 K [5]. At 1.6 K, the volume fraction of the incommensurate order is about 75%. However, the commensurate AF order is stabilized by hydrostatic pressure: Its volume fraction becomes nearly 100% at 2.4 GPa, the pressure where superconductivity first occurs and f electrons change from localized to itinerant. All the NQR experiments [5,9,18] lead to the same conclusion: The magnetic propagation vector is $(1/2, 1/2, \delta)$, although the value of δ is not predicted.

In this Rapid Communication, we report the magnetic structure of CePt₂In₇ determined from neutron diffraction, which is a bulk probe contrary to NQR and μ^+ SR measurements. The magnetic structure is found to be commensurate with a magnetic propagation vector $(1/2, 1/2, 1/2)$. The magnetic moments are aligned antiferromagnetically in the basal plane, and rotate by 90° from one CeIn₃ plane to another. As expected, the magnetic structure is more two dimensional than those of either CeIn₃ or CeRhIn₅.

Single crystals of CePt₂In₇ were grown by the In self-flux method, as explained in more detail elsewhere [19]. The high quality of the samples is confirmed by specific heat

*matthias.raba@lncmi.cnrs.fr

†pierre.rodriere@neel.cnrs.fr

‡ilya.sheikin@lncmi.cnrs.fr

measurements that show a clear and unique second-order phase transition at the Néel temperature $T_N = 5.5$ K [20] and de Haas–van Alphen effect measurements exhibiting quantum oscillations starting from about 2 T [21]. Two neutron diffraction experiments were performed at the Institut Laue-Langevin (ILL) of Grenoble (France) to determine the AF structure of CePt_2In_7 . At first, a powder neutron diffraction was carried out on the high-flux diffractometer D1B by using a pyrolytic graphite (002) monochromator providing a beam with a wavelength of 2.52 Å and a 128° multidetector. A great number of single crystals was ground into 1.6 g of fine powder. These measurements did not reveal any magnetic peaks below T_N in spite of 10-h-long acquisitions at both 1.5 and 10 K. This puts an upper limit for the magnetic moment at about $0.8\mu_B$. Both the crystal structure and the lattice parameters remain unchanged in the AF state. Similarly, no temperature-dependent magnetic peaks were detected in the previous powder neutron diffraction measurements on CeRhIn_5 [12]. Following this, a preliminary single-crystal neutron diffraction experiment was performed on the D23 beamline. This test measurement revealed a magnetic peak with $\mathbf{q}_M = (1/2, 1/2, 1/2)$, leading to more detailed measurements on the D10 beamline. For this experiment, we examined the biggest available sample from the same batch with the dimensions $1.9 \times 1.9 \times 1.0$ mm³ and the c axis perpendicular to the platelet surface. The instrument was used in a four-circle configuration with an additional triple-axis energy analysis. The latter was used in the elastic mode with a single detector to increase the signal-to-noise ratio. A vertically focusing pyrolytic graphite monochromator and analyzer was employed, fixing the incident and analyzed wavelength to 2.36 Å. A pyrolytic graphite filter was used to reduce the higher harmonic contamination to 10^{-4} of the primary beam intensity. In order to reach temperatures down to 2 K, we used a four-circle cryostat with helium circulation. The measured lattice parameters at $T = 2$ K are $a = 4.595(2)$ Å and $c = 21.558(5)$ Å, as obtained from the analysis of 30 nuclear Bragg peaks.

Figure 1 shows the \mathbf{Q} scan performed along the [001] direction at $T = 2$ K. A clear peak is observed at $\mathbf{Q} = (1/2, 1/2, 1/2)$. As expected, the $(1/2, 1/2, 1/2)$ Bragg peak disappears above the Néel temperature, as shown in the inset of Fig. 1 for $T = 10$ K. Below T_N , the q_l value of the propagation vector does not seem to change with temperature. However, the full width at half maximum (FWHM) of the magnetic peak increases with temperature (see Fig. 2), suggesting a smaller size of magnetic domains along the c axis for temperatures close to T_N . We did not observe any other obvious peaks with an incommensurate magnetic wave vector. However, a small satellite peak at $(1/2, 1/2, 1/2 - \delta)$ with $\delta = 0.03$ seems to emerge above the background noise at $T = 2$ K, as shown in Fig. 2(a). The integrated intensity of this peak is 21% of the commensurate magnetic Bragg peak. The satellite peak seems to disappear at $T = 4$ K, as an attempt to fit the data by two Gaussian peaks with the fixed position of the incommensurate peak and integrated intensity ratio of the two peaks does not yield a satisfactory result [see Fig. 2(b)]. Remarkably, the NQR measurements performed on single crystals suggest the presence of only a commensurate order at $T = 4$ K, whereas an incommensurate order starts to develop below about 3 K [9,18]. However, these measurements also suggest that the

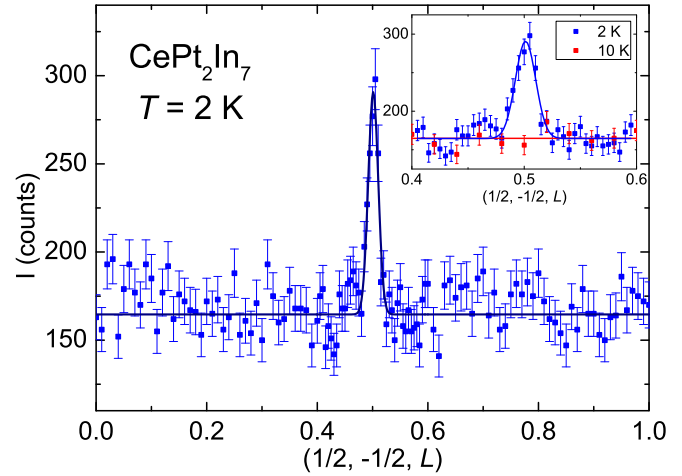


FIG. 1. \mathbf{Q} scan performed along the $[0,0,1]$ direction at 2 K. The inset shows a zoom of the same \mathbf{Q} scan near $\mathbf{Q} = (1/2, -1/2, 1/2)$ at 2 and 10 K. The intensity is in number of counts per 4×10^6 monitor counts, which corresponds roughly to 7 min. The solid lines are Gaussian fits of the peaks.

volume fraction of the incommensurate order is almost 3/4 at 2 K, and that the internal field due to the incommensurate order is larger than that due to its commensurate counterpart. This is difficult to reconcile with our results, if the small satellite

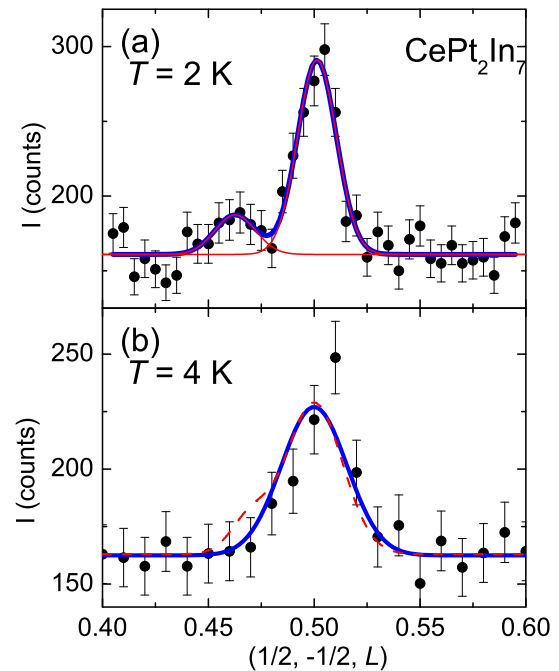


FIG. 2. Detailed \mathbf{Q} scans performed along the $[0,0,1]$ direction at (a) 2 K and (b) 4 K. The intensity is in number of counts per 4×10^6 monitor counts, which corresponds roughly to 7 min. The lines are Gaussian fits of the peaks. The data obtained at 2 K are well fit by two Gaussian peaks centered at $(1/2, -1/2, 1/2)$ and $(1/2, -1/2, 0.47)$. At 4 K, an attempt to fit the data by two Gaussian peaks while fixing the position of the second peak and the integrated intensity ratio of the two peaks (dashed line) does not yield a satisfactory result. The data are much better fitted by a single Gaussian peak (solid line).

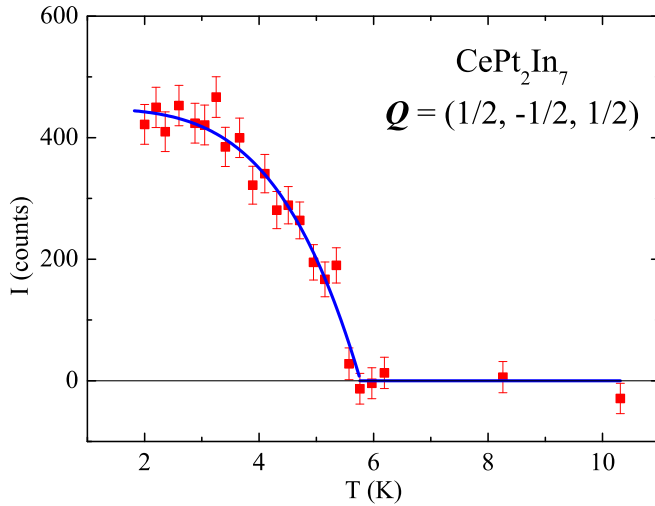


FIG. 3. Temperature dependence of the $(1/2, -1/2, 1/2)$ magnetic Bragg peak intensity after subtracting the background. The intensity is in number of counts per 1.6×10^7 monitor counts, which corresponds roughly to 27 min. The line is a phenomenological fit as explained in the text.

peak was a signature of an additional incommensurate order. Furthermore, if this was the case, a peak at $(1/2, 1/2, 1/2 + \delta)$ should have also been observed due to the tetragonal symmetry, although not necessarily with the same intensity. Such a peak, however, is absent in our data. Therefore, for the rest of this Rapid Communication, we will neglect this tiny contribution, whose possible magnetic origin remains to be clarified.

Figure 3 shows the temperature dependence of the $(1/2, -1/2, 1/2)$ magnetic Bragg peak intensity I proportional to the square of the magnetic phase transition order parameter. To determine the Néel temperature, the data were fitted by a phenomenological function $I/I_0 = 1 - (T/T_N)^\alpha$, with α a free parameter, which does not have a particular physical meaning. This function was successfully used to fit the temperature dependence of the magnetic Bragg peak intensity in other heavy fermion compounds, such as CePd_2Si_2 [22,23] and Sn-doped CeRhIn_5 [15]. The best fit is obtained with $\alpha = 4.1 \pm 0.7$ and $T_N = 5.7 \pm 0.1$ K. The latter value is consistent with T_N determined from previous measurements [3–6,20].

In order to determine the magnetic structure and the staggered moment, we measured 24 magnetic peaks at 2 K. For each peak, the measured neutron Bragg intensity was corrected for extinction, absorption, and Lorentz factor; the resulting intensities are shown in Table I. Only two arrangements of the magnetic moments are allowed by the group theory analysis: They can be either aligned in the basal plane or along the c axis. The best refinement, taking into account two magnetic domains and assuming they are equally populated, is obtained for magnetic moments aligned antiferromagnetically in the basal plane. The moments rotate by 90° from one CeIn_3 plane to its nearest-neighbor along the c axis. For this structure, the comparison between the observed intensities and the calculated ones is shown in Table I. Given the tetragonal symmetry of CePt_2In_7 , we cannot determine the orientation of the magnetic moments in the basal plane.

TABLE I. Magnetic refinement for the magnetic structure discussed in the text and shown in Fig. 4. (Note: $\chi^2 = 1.84$).

Q	I_{obs}	I_{calc}
$(1/2, -1/2, -1/2)$	8(1)	8.2
$(-1/2, -1/2, -1/2)$	7(1)	8.2
$(-1/2, -1/2, 1/2)$	7(1)	8.2
$(1/2, -1/2, 1/2)$	7.0(9)	8.2
$(1/2, -1/2, -5/2)$	10(1)	10.6
$(-1/2, -1/2, -5/2)$	9(1)	10.6
$(-1/2, -1/2, 5/2)$	11(1)	10.6
$(1/2, -1/2, 5/2)$	9(1)	10.6
$(1/2, -1/2, -3/2)$	10(1)	9.3
$(-1/2, -1/2, -3/2)$	9(1)	9.3
$(-1/2, -1/2, 3/2)$	8(1)	9.3
$(1/2, -1/2, 3/2)$	7(1)	9.3
$(1/2, -1/2, -9/2)$	12(2)	12.0
$(-1/2, -1/2, -9/2)$	19(3)	12.0
$(-1/2, -1/2, 9/2)$	13(2)	12.0
$(1/2, -1/2, 9/2)$	15(2)	12.0
$(1/2, -1/2, -7/2)$	14(2)	11.5
$(-1/2, -1/2, -7/2)$	13(2)	11.5
$(-1/2, -1/2, 7/2)$	12(1)	11.5
$(1/2, -1/2, 7/2)$	11(2)	11.5
$(1/2, -1/2, -11/2)$	15(2)	12.0
$(-1/2, -1/2, -11/2)$	18.8(25)	12.0
$(1/2, 1/2, 11/2)$	13(2)	12.0
$(-1/2, 1/2, 11/2)$	18(3)	12.0

However, previous NQR data [5,18] strongly suggest that they are parallel to the $[100]$ direction in the commensurate phase. The resulting magnetic structure (magnetic space group $I_c\bar{4}2m$) is schematically shown in Fig. 4. The staggered

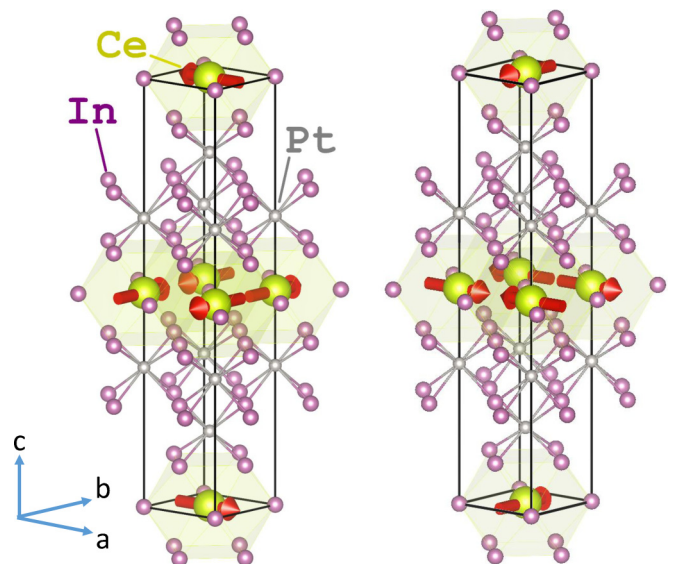


FIG. 4. Magnetic structure of CePt_2In_7 in a structural unit cell. The magnetic moment, schematically shown by arrows, is $0.45\mu_B$ per Ce and it is aligned in the basal plane. Both magnetic domains are shown.

TABLE II. Crystallographic and magnetic properties of CeT_mIn_{3+2m} ($T = Rh, Pt$) compounds. Here, d is the distance between $CeIn_3$ planes, α is the angle between magnetic moments on neighboring $CeIn_3$ planes, and M is the staggered moment.

	m	d (Å)	α	M (μ_B/Ce)
$CeIn_3$	0	4.6	180°	0.48 [11], 0.65 [10]
$CeRhIn_5$	1	7.5	107°	0.54 [16], 0.59 [14,15], 0.75 [12,13]
$CePt_2In_7$	2	10.8	90°	0.45

magnetic moment is determined at 2 K to be $M = 0.45(1)\mu_B$ per Ce.

Having determined the magnetic structure of $CePt_2In_7$, we now consider the relationship between the magnetic and crystal structures within the CeT_mIn_{3+2m} ($T = Rh, Pt$) family, where m TIn_2 layers separate a single $CeIn_3$ layer (see Table II). The building block of the family, $CeIn_3$ ($m = 0$), crystallizes into a simple cubic structure (space group $Pm\bar{3}m$). For $CeRhIn_5$ ($m = 1$) with alternating layers of $CeIn_3$ and $RhIn_2$, the crystal structure is primitive tetragonal (space group $P4/mmm$). Finally, $CePt_2In_7$ ($m = 2$), where two $PtIn_2$ layers separate each $CeIn_3$ layer, crystallizes into a body-centered-tetragonal structure with $I4/mmm$ space group. In all three materials, the magnetic moments of the Ce ions form a square lattice, surrounded by In ions in the a - b plane. They all are simple, nearest-neighbor antiferromagnets in the plane. In $CeIn_3$, the magnetic moments are also arranged antiferromagnetically along the c axis. In $CeRhIn_5$, magnetic correlations across the $RhIn_2$ layer are incommensurate, with neighboring magnetic moments being rotated by approximately 107° [12,13]. In $CePt_2In_7$, the magnetic moments are also aligned in the a - b plane. In this case, however, they rotate by 90° from one plane to another, as shown in Fig. 4. Clearly, the rotation angle is reduced with increasing the number of TIn_2 layers between the $CeIn_3$ layers (see Table II). The magnetic structure of $CePt_2In_7$ is expected to be the most two dimensional among the three compounds. Indeed, a greater separation between the $CeIn_3$ planes leads to a weaker coupling between them. In addition, geometrical frustration due to the body-centered-tetragonal crystal structure is likely to further reduce the effective dimensionality, as was experimentally observed in $BaCuSi_2O_6$ [24]. The 90° rotation angle of the moments from one plane to another observed in $CePt_2In_7$ is probably a signature of its most two-dimensional magnetic structure.

The magnetic structure of both $CeIn_3$ and $CePt_2In_7$ is commensurate. The same conclusion was drawn for Ce_2RhIn_8 , in which two layers of $CeIn_3$ are separated by a single layer of $RhIn_2$ [25]. This suggests $CeRhIn_5$ with its incommensurate magnetic structure as a unique member of the whole $Ce_nT_mIn_{3n+2m}$ ($T =$ transition metal) family, where m TIn_2 layers separate n $CeIn_3$ layers. In all these compounds, superconductivity emerges in the vicinity of a

quantum critical point induced either by pressure [7,26–28] or chemical doping [29–31]. Interestingly, a commensurate magnetic order was observed to either coexist or compete with incommensurate ordering in $CeRhIn_5$ doped with either Ir [32] or Co [30,31]. Remarkably, in these compounds, commensurate antiferromagnetism emerges in the vicinity of a quantum critical point where superconductivity also appears. Furthermore, in Sn-doped $CeRhIn_5$, a drastic change in the magnetic order and a commensurate antiferromagnetism was observed in the proximity of the quantum critical point [15] where superconductivity is expected, but has not been observed so far. This suggests that a commensurate magnetic order might be favorable for the formation of superconductivity around a quantum critical point in this family of materials. On the other hand, several neutron diffraction experiments performed in $CeRhIn_5$ under pressure up to 1.7 GPa did not reveal the presence of a commensurate AF order [33–35]. This pressure, however, is considerably lower than the critical value, $P_c \approx 2.4$ GPa, although a pressure-induced bulk superconductivity is observed above about 1.5 GPa [28].

Regarding the staggered moment, its values are comparable for all three compounds of the CeT_mIn_{3+2m} family (see Table II). Given that the staggered moment of Ce_2RhIn_8 , $M = 0.55\mu_B/Ce$ [25], is also of the same order, there is no obvious correlation between the staggered moment and dimensionality in this family of heavy fermion materials.

Finally, the magnetic propagation vector $(1/2, 1/2, 1/2)$ observed both in $CeIn_3$ and $CePt_2In_7$ implies that the magnetic Brillouin zone in these compounds is eight times smaller than the crystallographic one. In $CeIn_3$, this naturally accounts for the observation of only small Fermi surfaces in quantum oscillation measurements performed at moderate magnetic fields up to 17 T [36]. Very high magnetic fields of about 55 T are required for the observation of large Fermi surfaces through magnetic breakdown tunneling [37]. Similarly, only small Fermi surface pockets, well within the interior of the AF Brillouin zone, are observed in $CePt_2In_7$ at moderately high magnetic fields up to about 25–30 T [21,38]. This is in contrast with $CeRhIn_5$, where large Fermi surfaces are observed at relatively low fields below 17–18 T [39,40].

In conclusion, we find the commensurate magnetic structure with $\mathbf{q}_M = (1/2, 1/2, 1/2)$ as depicted in Fig. 4 for $CePt_2In_7$. A magnetic moment of $0.45(1)\mu_B$ at 2 K resides on the Ce ion and the basal plane is its easy plane. Within the basal plane, magnetic moments form a simple nearest-neighbor antiferromagnet on a square lattice. The moments rotate by 90° from one $CeIn_3$ plane to another along the c axis.

We thank Y. Tokunaga, M. Horvatic, and J. Robert for fruitful discussions. We acknowledge scientific and technical support we received during the powder neutron diffraction experiment on the CRG-D1B (ILL) operated by the CNRS. This work was partially supported by the ANR-DFG grant “Fermi-NESt”.

[1] T. Klimczuk, O. Walter, L. MÜchler, J. W. Krizan, F. Kinnart, and R. J. Cava, *J. Phys.: Condens. Matter* **26**, 402201 (2014).

[2] M. M. Altarawneh, N. Harrison, R. D. McDonald, F. F. Balakirev, C. H. Mielke, P. H. Tobash, J.-X. Zhu, J. D. Thompson, F. Ronning, and E. D. Bauer, *Phys. Rev. B* **83**, 081103 (2011).

- [3] P. H. Tobash, F. Ronning, J. D. Thompson, B. L. Scott, P. J. W. Moll, B. Batlogg, and E. D. Bauer, *J. Phys.: Condens. Matter* **24**, 015601 (2012).
- [4] E. D. Bauer, H. O. Lee, V. A. Sidorov, N. Kurita, K. Gofryk, J.-X. Zhu, F. Ronning, R. Movshovich, J. D. Thompson, and T. Park, *Phys. Rev. B* **81**, 180507 (2010).
- [5] N. apRoberts-Warren, A. P. Dioguardi, A. C. Shockley, C. H. Lin, J. Crocker, P. Klavins, and N. J. Curro, *Phys. Rev. B* **81**, 180403(R) (2010).
- [6] N. apRoberts-Warren, A. P. Dioguardi, A. C. Shockley, C. H. Lin, J. Crocker, P. Klavins, and N. J. Curro, *J. Phys.: Conf. Ser.* **344**, 012027 (2012).
- [7] V. A. Sidorov, X. Lu, T. Park, H. Lee, P. H. Tobash, R. E. Baumbach, F. Ronning, E. D. Bauer, and J. D. Thompson, *Phys. Rev. B* **88**, 020503 (2013).
- [8] E. D. Bauer, V. A. Sidorov, H. Lee, N. Kurita, F. Ronning, R. Movshovich, and J. D. Thompson, *J. Phys.: Conf. Ser.* **200**, 012011 (2010).
- [9] H. Sakai, Y. Tokunaga, S. Kambe, F. Ronning, E. D. Bauer, and J. D. Thompson, *Phys. Rev. Lett.* **112**, 206401 (2014).
- [10] J. M. Lawrence and S. M. Shapiro, *Phys. Rev. B* **22**, 4379 (1980).
- [11] A. Benoit, J. Boucherle, P. Convert, J. Flouquet, J. Palleau, and J. Schweizer, *Solid State Commun.* **34**, 293 (1980).
- [12] W. Bao, P. G. Pagliuso, J. L. Sarrao, J. D. Thompson, Z. Fisk, J. W. Lynn, and R. W. Erwin, *Phys. Rev. B* **62**, R14621 (2000).
- [13] W. Bao, P. G. Pagliuso, J. L. Sarrao, J. D. Thompson, Z. Fisk, J. W. Lynn, and R. W. Erwin, *Phys. Rev. B* **67**, 099903(E) (2003).
- [14] S. Raymond, E. Ressouche, G. Knebel, D. Aoki, and J. Flouquet, *J. Phys.: Condens. Matter* **19**, 242204 (2007).
- [15] S. Raymond, J. Buhot, E. Ressouche, F. Bourdarot, G. Knebel, and G. Lapertot, *Phys. Rev. B* **90**, 014423 (2014).
- [16] D. M. Fobes, E. D. Bauer, J. D. Thompson, A. Sazonov, V. Hutanu, S. Zhang, F. Ronning, and M. Janoschek, *J. Phys.: Condens. Matter* **29**, 17LT01 (2017).
- [17] M. Månsson, K. Prša, Y. Sassa, P. H. Tobash, E. D. Bauer, C. Rusu, D. Andreica, O. Tjernberg, K. Sedlak, M. Grioni *et al.*, *J. Phys.: Conf. Ser.* **551**, 012028 (2014).
- [18] H. Sakai, Y. Tokunaga, S. Kambe, H.-O. Lee, V. A. Sidorov, P. H. Tobash, F. Ronning, E. D. Bauer, and J. D. Thompson, *Phys. Rev. B* **83**, 140408 (2011).
- [19] S. Kurahashi, S. Ota, S. Tomaru, Y. Hirose, and R. Settai, *J. Phys.: Conf. Ser.* **592**, 012006 (2015).
- [20] Y. Krupko, A. Demuer, S. Ota, Y. Hirose, R. Settai, and I. Sheikin, *Phys. Rev. B* **93**, 085121 (2016).
- [21] K. Götze, Y. Krupko, J. Bruin, J. Klotz, S. Ota, Y. Hirose, H. Harima, R. Settai, A. McCollam, and I. Sheikin (unpublished).
- [22] N. H. van Dijk, B. Fåk, T. Charvolin, P. Lejay, and J. M. Mignot, *Phys. Rev. B* **61**, 8922 (2000).
- [23] N. Kernavanois, S. Raymond, E. Ressouche, B. Grenier, J. Flouquet, and P. Lejay, *Phys. Rev. B* **71**, 064404 (2005).
- [24] S. E. Sebastian, N. Harrison, C. D. Batista, L. Balicas, M. Jaime, P. A. Sharma, N. Kawashima, and I. R. Fisher, *Nature (London)* **441**, 617 (2006).
- [25] W. Bao, P. G. Pagliuso, J. L. Sarrao, J. D. Thompson, Z. Fisk, and J. W. Lynn, *Phys. Rev. B* **64**, 020401 (2001).
- [26] N. D. Mathur, F. M. Grosche, S. R. Julian, I. R. Walker, D. M. Freye, R. K. W. Haselwimmer, and G. G. Lonzarich, *Nature (London)* **394**, 39 (1998).
- [27] M. Nicklas, V. A. Sidorov, H. A. Borges, P. G. Pagliuso, C. Petrovic, Z. Fisk, J. L. Sarrao, and J. D. Thompson, *Phys. Rev. B* **67**, 020506 (2003).
- [28] G. Knebel, D. Aoki, D. Braithwaite, B. Salce, and J. Flouquet, *Phys. Rev. B* **74**, 020501 (2006).
- [29] P. G. Pagliuso, C. Petrovic, R. Movshovich, D. Hall, M. F. Hundley, J. L. Sarrao, J. D. Thompson, and Z. Fisk, *Phys. Rev. B* **64**, 100503 (2001).
- [30] S. Ohira-Kawamura, H. Shishido, A. Yoshida, R. Okazaki, H. Kawano-Furukawa, T. Shibauchi, H. Harima, and Y. Matsuda, *Phys. Rev. B* **76**, 132507 (2007).
- [31] M. Yokoyama, N. Oyama, H. Amitsuka, S. Oinuma, I. Kawasaki, K. Tenya, M. Matsuura, K. Hirota, and T. J. Sato, *Phys. Rev. B* **77**, 224501 (2008).
- [32] A. Llobet, A. D. Christianson, W. Bao, J. S. Gardner, I. P. Swainson, J. W. Lynn, J.-M. Mignot, K. Prokes, P. G. Pagliuso, N. O. Moreno, J. L. Sarrao, J. D. Thompson, and A. H. Lacerda, *Phys. Rev. Lett.* **95**, 217002 (2005).
- [33] S. Majumdar, G. Balakrishnan, M. R. Lees, D. McK. Paul, and G. J. McIntyre, *Phys. Rev. B* **66**, 212502 (2002).
- [34] A. Llobet, J. S. Gardner, E. G. Moshopoulou, J.-M. Mignot, M. Nicklas, W. Bao, N. O. Moreno, P. G. Pagliuso, I. N. Goncharenko, J. L. Sarrao, and J. D. Thompson, *Phys. Rev. B* **69**, 024403 (2004).
- [35] S. Raymond, G. Knebel, D. Aoki, and J. Flouquet, *Phys. Rev. B* **77**, 172502 (2008).
- [36] M. Endo, N. Kimura, and H. Aoki, *J. Phys. Soc. Jpn.* **74**, 3295 (2005).
- [37] N. Harrison, S. E. Sebastian, C. H. Mielke, A. Paris, M. J. Gordon, C. A. Swenson, D. G. Rickel, M. D. Pacheco, P. F. Ruminer, J. B. Schillig, J. R. Sims, A. H. Lacerda, M.-T. Suzuki, H. Harima, and T. Ebihara, *Phys. Rev. Lett.* **99**, 056401 (2007).
- [38] A. Miyake, Y. Kohama, S. Ohta, Y. Hirose, R. Settai, K. Matsubayashi, Y. Uwatoko, A. Matsuo, K. Kindo, and M. Tokunaga, *J. Phys.: Conf. Ser.* **592**, 012149 (2015).
- [39] D. Hall, E. C. Palm, T. P. Murphy, S. W. Tozer, C. Petrovic, E. Miller-Ricci, L. Peabody, C. Q. H. Li, U. Alver, R. G. Goodrich, J. L. Sarrao, P. G. Pagliuso, J. M. Wills, and Z. Fisk, *Phys. Rev. B* **64**, 064506 (2001).
- [40] H. Shishido, R. Settai, D. Aoki, S. Ikeda, H. Nakawaki, N. Nakamura, T. Iizuka, Y. Inada, K. Sugiyama, T. Takeuchi *et al.*, *J. Phys. Soc. Jpn.* **71**, 162 (2002).