open science

# Boundary null-controllability of semi-discrete coupled parabolic systems in some multi-dimensional geometries 

Damien Allonsius, Franck Boyer

## To cite this version:

Damien Allonsius, Franck Boyer. Boundary null-controllability of semi-discrete coupled parabolic systems in some multi-dimensional geometries. 2018. hal-01827044v1

HAL Id: hal-01827044
https://hal.science/hal-01827044v1
Preprint submitted on 1 Jul 2018 (v1), last revised 13 Apr 2019 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# BOUNDARY NULL-CONTROLLABILITY OF SEMI-DISCRETE COUPLED PARABOLIC SYSTEMS IN SOME MULTI-DIMENSIONAL GEOMETRIES 

DAMIEN ALLONSIUS* AND FRANCK BOYER ${ }^{\dagger}$


#### Abstract

The main goal of this paper is to investigate the controllability properties of semidiscrete in space coupled parabolic systems with less controls than equations, in dimension greater than 1 . We are particularly interested in the boundary control case which is notably more intricate that the distributed control case, even though our analysis is more general.

The main assumption we make on the geometry and on the evolution equation itself is that it can be put into a tensorized form. In such a case, following [4] and using an adapted version of the Lebeau-Robbiano construction, we are able to prove controllability results for those semidiscrete systems (provided that the structure of the coupling terms satisfies some necessary Kalman condition) with uniform bounds on the controls.

To achieve this objective we actually propose an abstract result on ordinary differential equations with estimates on the control and the solution whose dependence upon the system parameters are carefully tracked. When applied to an ODE coming from the discretization in space of a parabolic system, we thus obtain uniform estimates with respect to the discretization parameters.


Key words. Control theory, Lebeau-Robbiano strategy, parabolic systems, semi-discretization in space.

AMS subject classifications. $35 \mathrm{~K} 10-65 \mathrm{M} 06-93 \mathrm{~B} 05$

## 1. Introduction.

1.1. Motivating example. In this paper, we are interested in null-controllability properties at any time $T>0$ of coupled linear parabolic equations at the continuous level as well as at the semi-discrete in space level. In this introduction, we will focus on the following prototype system (which is in the so-called cascade form)

$$
\left\{\begin{align*}
\partial_{t} \alpha-\Delta \alpha & =0, & & \text { in }(0, T) \times \Omega  \tag{1}\\
\partial_{t} \beta-\Delta \beta+\alpha & =0, & & \text { in }(0, T) \times \Omega \\
\alpha & =1_{\Gamma} v, & & \text { on }(0, T) \times \partial \Omega \\
\beta & =0, & & \text { on }(0, T) \times \partial \Omega
\end{align*}\right.
$$

where $\Omega$ is a bounded domain of $\mathbb{R}^{d}(d \geq 1), \Gamma$ is a non empty part of the boundary $\partial \Omega, \alpha$ and $\beta$ are the two components of the system, and $v$ is the boundary control we are looking for. The main difficulty in the analysis of the controllability of such system comes from the fact that we only have one boundary control $v$ to drive the two components $(\alpha, \beta)$ to 0 at the final time. The coupling terms (here the term $\alpha$ in the equation for $\beta$ ) plays a key role in the problem and it can be seen as a kind of indirect control for the second component of the system. Note that this indirect controllability issue arises even if $\Gamma=\partial \Omega$.

Actually, it appears that the results for such systems may be quite different from the case of scalar equations or from the case of coupled systems with a distributed control. We refer for instance to the survey [2] for a review on that topic. In particular, it is explained in that reference that usual techniques based on Carleman estimates

[^0]are useless on those problems. This is mainly because those techniques naturally gives the controllability of the system when there is as many controls as components of the system (for (1) it would consist in another boundary control for $\beta$ on $\Gamma$ ) and, in a second step, it is needed to prove that only on control is necessary. This is done, at the observability inequality level (see the discussion in Section 2.3) by removing one observation term thanks to the PDE itself. This last step cannot be done for boundary controls (or observations). This is why other approaches have to be developped.

Most of the results available up to now for such controllability problems are only proved in dimension $d=1$ by using the so-called moments method. This is a quite powerful method but, unfortunately, restricted to autonomous problem in space dimension 1. In the multi-dimensional case, one of the more advanced result available in the literature is proved in [4], in the case where the geometry and the diffusion operator can be tensorized and this is also the case we shall consider in the present work.


Figure 1: Typical geometric situation

For problem (1), this structure assumption amounts to assume that the geometry is as follows: $\Omega=\Omega_{1} \times \Omega_{2}$ with $\Omega_{1}=(0, a)$ and $\Omega_{2} \subset \mathbb{R}^{d-1}$, and $\Gamma=\{0\} \times \omega_{2}$ (see Figure 1). Then, we can rewrite the system as follows

$$
\left\{\begin{align*}
\partial_{t} \alpha-\partial_{x_{1}}^{2} \alpha-\Delta_{2} \alpha & =0, & & \text { in }(0, T) \times \Omega_{1} \times \Omega_{2},  \tag{2}\\
\partial_{t} \beta-\partial_{x_{1}}^{2} \beta-\Delta_{2} \beta+\alpha & =0, & & \text { in }(0, T) \times \Omega_{1} \times \Omega_{2}, \\
\alpha=\beta & =0, & & \text { on }(0, T) \times \Omega_{1} \times \partial \Omega_{2}, \\
\alpha & =1_{\{0\} \times \omega_{2}} v, & & \text { on }(0, T) \times\{0, a\} \times \Omega_{2}, \\
\beta & =0, & & \text { on }(0, T) \times\{0, a\} \times \Omega_{2},
\end{align*}\right.
$$

where $\Delta_{2}$ is the $(d-1)$-dimensional Laplace operator in $\Omega_{2}$. The fact that the diffusion operator is split into two parts, each of them acting on different sets of variables, is crucial in the analysis. That is the reason why we shall adopt a tensor product formalism that consists essentially in identifying $L^{2}(\Omega)$ to $L^{2}\left(\Omega_{1}\right) \otimes L^{2}\left(\Omega_{2}\right)$ and to write the two equations above in the following equivalent form

$$
\begin{cases}\partial_{t} \alpha+\left(-\partial_{x_{1}}^{2}\right) \otimes \mathcal{I} \alpha+\mathcal{I} \otimes\left(-\Delta_{2}\right) \alpha=0, & \text { in }(0, T),  \tag{3}\\ \partial_{t} \beta+\left(-\partial_{x_{1}}^{2}\right) \otimes \mathcal{I} \beta+\mathcal{I} \otimes\left(-\Delta_{2}\right) \beta+\alpha=0, & \text { in }(0, T),\end{cases}
$$

where the same symbol $\mathcal{I}$ is used for the identity operator in $L^{2}\left(\Omega_{1}\right)$ and $L^{2}\left(\Omega_{2}\right)$. All the necessary notations concerning tensor products will be recalled in Section 3.

By exploiting this tensor product structure, even though the tensor product formalism was not explicitly used, it was proved in [4], that the null-controllability of (3) holds at any time $T>0$.
1.2. Passing to the discrete world. We are now interested in semi-discrete versions of the controllability result for (1) mentioned just before. To simplify the presentation in this introduction, we assume that $d=2$, that $\Omega_{1}=\Omega_{2}=(0,1)$ and that the computation grid is made of $N \times N$ uniformly distributed points $(i h, j h)_{1 \leq i, j \leq N}$ with $h=\frac{1}{N+1}$. The semi-discrete system we consider is obtained by the finite difference method and reads
(4)
$\left\{\begin{aligned} \partial_{t} \alpha_{i, j}+\frac{4 \alpha_{i, j}-\alpha_{i-1, j}-\alpha_{i+1, j}-\alpha_{i, j-1}-\alpha_{i, j+1}}{h^{2}} & =0, & & \forall 1 \leq i, j \leq N, \\ \partial_{t} \beta_{i, j}+\frac{4 \beta_{i, j}-\beta_{i-1, j}-\beta_{i+1, j}-\beta_{i, j-1}-\beta_{i, j+1}}{h^{2}}+\alpha_{i, j} & =0, & & \forall 1 \leq i, j \leq N, \\ \alpha_{i, 0}=\beta_{i, 0}=\alpha_{i, N+1}=\beta_{i, N+1} & =0, & & \forall 1 \leq i \leq N, \\ \alpha_{0, j} & =v_{j} 1_{\omega_{2}}(j h), & & \forall 1 \leq j \leq N, \\ \alpha_{N+1, j}=\beta_{0, j}=\beta_{N+1, j} & =0, & & \forall 1 \leq j \leq N .\end{aligned}\right.$
The grid geometry is essentially the one described in Figure 2, where the control $v=\left(v_{j}\right)_{j}$ is only appearing in the first equation of the system (the one for $\alpha_{i, j}$ ) and only on the boundary points represented by the symbol $\square$ corresponding to the subdomain $\omega_{2}$.

At each time $t$, both components $\alpha_{h}=\left(\alpha_{i, j}\right)_{i, j} \in \mathbb{R}^{N \times N}$ and $\beta_{h}=\left(\beta_{i, j}\right)_{i, j} \in$ $\mathbb{R}^{N \times N}$ of the system are now considered as elements of the tensor product $\mathbb{R}^{N} \otimes \mathbb{R}^{N}$ and we observe that the five-point discrete Laplace operator can be written as the tensor product $\mathrm{A}_{h} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{A}_{h}$ with the usual definition of the three-point discrete Laplace matrix


We finally end up with the following equivalent form of our semi-discrete system

$$
\begin{cases}\partial_{t} \alpha_{h}+\left(\mathrm{A}_{h} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{h}\right) \alpha_{h} & =\mathrm{B}_{h} v_{h}  \tag{5}\\ \partial_{t} \beta_{h}+\left(\mathrm{A}_{h} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{h}\right) \beta_{h}+\alpha_{h} & =0\end{cases}
$$

where $\mathrm{B}_{h}$ is a matrix that accounts for the influence of the control $v_{h}$ in the system through the boundary conditions in (4), the precise definition of which will be given in Section 4.1. This is the semi discrete version of (3).

For this particular system, the main result of this paper is be the following (see Section 4.1 for the precise statement and definition of the norms involved).

THEOREM. There exists $C>0$ and $h_{0}>0$ such that for any $h<h_{0}$, any time $T>0$ and any initial data $\alpha_{h}^{0}, \beta_{h}^{0} \in \mathbb{R}^{N} \otimes \mathbb{R}^{N}$, (with $h=\frac{1}{N+1}$ ) there exists a control $v_{h} \in L^{2}\left(0, T, \mathbb{R}^{N}\right)$ such that

$$
\llbracket v_{h} \rrbracket_{L^{2}\left(0, T, \mathbb{R}^{N}\right)} \leq C e^{\frac{C}{T}}\left(\left\|\alpha_{h}^{0}\right\|+\left\|\beta_{h}^{0}\right\|\right)
$$

and the associated solution to (5) satisfies

$$
\left\|\alpha_{h}(T)\right\|+\left\|\beta_{h}(T)\right\| \leq C e^{-C / h^{2}} e^{\frac{C}{T}}\left(\left\|\alpha_{h}^{0}\right\|+\left\|\beta_{h}^{0}\right\|\right)
$$

It is well known, see $[6,16]$, that we cannot expect in general to achieve exactly $\alpha_{h}(T)=\beta_{h}(T)=0$ since the semi-discrete system may be not controllable. In this sense, achieving exponentially small targets with respect to $h$ is an optimal result.

Our aim will be to provide similar results for more general semi-discrete systems. That is the reason why, in order to formulate them more conveniently for any number of coupled equations and to ease the reading of the proofs, we shall actually gather the two components $\left(\alpha_{h}, \beta_{h}\right) \in\left(\mathbb{R}^{N} \otimes \mathbb{R}^{N}\right)^{2}$ into a single unknown $y_{h} \in \mathbb{R}^{N} \otimes \mathbb{R}^{N} \otimes \mathbb{R}^{n}$, with $n=2$ in the present case, in such a way that the considered system (5) will finally be written in the compact form

$$
\partial_{t} y_{h}+\left(\mathrm{A}_{h} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{h}\right) \otimes \mathbf{I} y_{h}+\mathrm{I} \otimes \mathrm{I} \otimes\left(\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right) y_{h}=\mathrm{B}_{h} \otimes\binom{1}{0} v_{h}
$$

In the sequel of this paper we shall not explicitly mention the subscript $h$ for the notation of quantities related to the discretization process but taking into account the fact that the considered spaces, norms and operators are grid-dependent is a central point in the analysis.
1.3. Main results and outline of the paper. Considering the previous discussion we shall analyse in this paper the controllability of parabolic systems of $n$ components and $m$ controls of the following tensorized form

$$
\begin{equation*}
\partial_{t} y+\mathcal{A}_{1} \otimes \mathcal{I} \otimes \mathbf{I} y+\mathcal{I} \otimes \mathcal{A}_{2} \otimes \mathbf{I} y+\mathcal{I} \otimes \mathcal{I} \otimes \mathrm{C} y=\mathcal{B}_{1} \otimes \mathcal{B}_{2} \otimes \mathrm{~B} v, \quad \text { in }(0, T) \tag{6}
\end{equation*}
$$

where $\mathcal{A}_{i}$ is a diffusion operator in $\Omega_{i}, \mathcal{B}_{i}$ is a (boundary or distributed) control operator in $\Omega_{i}$, I is the $n \times n$ identity matrix, C is a $n \times n$ coupling matrix and B a $n \times m$ control matrix.

Our main aim being to analyse semi-discrete versions of (6), we shall also consider linear ordinary differential equations of the similar form

$$
\begin{equation*}
\partial_{t} y+\mathrm{A}_{1} \otimes \mathrm{I} \otimes \mathrm{I} y+\mathrm{I} \otimes \mathrm{~A}_{2} \otimes \mathrm{I} y+\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C} y=\mathrm{B}_{1} \otimes \mathrm{~B}_{2} \otimes \mathrm{~B} v, \quad \text { in }(0, T) \tag{7}
\end{equation*}
$$

where the unknown $y$ and the control $v$ belong to a finite dimensional tensor space, $\mathrm{A}_{i}$ and $\mathrm{B}_{i}$ are linear operators on those space. All those objects depend, by nature, on discretization parameters, as in the example of Section 1.1. Therefore, it will be crucial to take care of all the constants in the estimates so as to obtain, at the end, controllability results for (7), that will not depend on those parameters.

Remark 1.1. We have used in (6) and (7) a convention that will be used all along the paper:

- Operators acting in infinite dimensional function spaces are written with calligraphic letters: $\mathcal{A}, \mathcal{B}, \mathcal{I}, \ldots$
- Operators (matrices) acting in finite dimensional spaces coming from discretisation issues (whose dimension may be large and depends on the discretization parameters) are written with upright letters: A, B, I, ...
- Matrices acting in finite dimensional spaces coming the number of components or controls in the system (their dimension is fixed and independent of discretization parameters) are written with sans serif letters: B, C, I, ...

The outline of the paper is the following. Section 2 is dedicated to recalling the main results in the controllability theory for ODEs while paying a particular attention to the discrete functional setting that will be adapted to the analysis of finite difference approximations of parabolic PDEs. In section 3, we first review the material we need concerning tensorized operators that are central in the present work, then we state the precise assumptions we need and our main abstract result (Theorem 3.1). In short, we will assume for our tensorized system (6) (in the continuous setting) or (7) (in the discrete setting):

- that the associated sub-problem concerning only the first coordinate of the tensor product which is

$$
\begin{equation*}
\partial_{t} y+\mathcal{A}_{1} \otimes \mathbf{I} y+\mathcal{I} \otimes \mathrm{C} y=\mathcal{B}_{1} \otimes \mathrm{~B} v, \quad \text { in }(0, T) \tag{8}
\end{equation*}
$$

or

$$
\begin{equation*}
\partial_{t} y+\mathrm{A}_{1} \otimes \mathrm{I} y+\mathrm{I} \otimes \mathrm{C} y=\mathrm{B}_{1} \otimes \mathrm{~B} v, \quad \text { in }(0, T) \tag{9}
\end{equation*}
$$

is null-controllable (or a relaxed version of this) at any time $T$ with a precise control of the cost of the control.
If we come back to our motivating example, this subsystem reads

$$
\left\{\begin{align*}
\partial_{t} \alpha-\partial_{x_{1}}^{2} \alpha & =0, & & \text { in }(0, T) \times \Omega_{1}  \tag{10}\\
\partial_{t} \beta-\partial_{x_{1}}^{2} \beta+\alpha & =0, & & \text { in }(0, T) \times \Omega_{1} \\
\alpha & =1_{\{0\}} v, & & \text { on }(0, T) \times \partial \Omega_{1} \\
\beta & =0, & & \text { on }(0, T) \times \partial \Omega_{1}
\end{align*}\right.
$$

Note that this system does not depend on the control set $\omega_{2}$.

- that the diffusion operator $\mathcal{A}_{2}$ (resp. $\mathrm{A}_{2}$ ) and the control operator $\mathcal{B}_{2}$ (resp. $B_{2}$ ) in the other direction satisfy a suitable spectral estimate similar to the Lebeau-Robbiano spectral inequality, except that we allow the inequality to hold only for a portion of the spectrum.
For system (2), this amounts to ask that the Lebeau-Robbiano spectral inequality holds, relative to the control set $\omega_{2}$, for the eigenfunctions of the operator $-\Delta_{2}$ in $\Omega_{2}$ with homogeneous boundary condition.
The complete proof of the theorem is given in Section 3.3. It consists, following the strategy developped in [4], to implement a construction of the control similar to the one originally proposed by Lebeau and Robbiano in [13]. We split the time interval into a suitable number of subintervals whose length is carefully chosen and, on each of those subintervals, we construct a partial control (obtained by combining the two assumptions above) that is able to damp out exponentially the part of the solution corresponding to the frequencies less than some threshold. Note that, contrary to the usual construction, we are not necessary able to drive this part of the solution exactly to zero at this stage. This threshold is then increased while the construction progresses towards the final time $T$; this eventually gives the expected control.

The main novelties in the present work are that: we allow relaxed controllability and spectral inequalities in our assumptions and moreover we precisely take care of the dependance of all the quantities of interest (norms, constants, ...) with respect to parameters on which the problem may depend. Those two refinements of the proof in [4] are mandatory since we want to apply this abstract result to systems obtained by semi-discretization processes.

To conclude the paper, in section 4, we precsiely explain how to use the abstract formalism developed here to achieve the uniform controllability results of semi-discrete coupled parabolic systems as announced in this introduction. As another example, we also show how to deduce the result of [4] (slightly generalized to variable coefficients operators) from the present abstract result and give some insights on other possible applications.
2. Controllability for linear ODEs. Before studying discrete versions of System (6), which is the main aim of the paper, we start by introducing the main notations and results that we shall use in the sequel concerning the controllability of linear ODEs. Most of this material is already well-known, however we propose a specific point of view adapted to our needs.
2.1. Framework. Let $\left(E,\langle\bullet, \bullet\rangle_{0}\right)$ and $\left(U,[\bullet, \bullet]_{0}\right)$ be two (finite dimensional) Euclidean spaces (each of them being identified with its own dual space). The corresponding norms are denoted by $\|\bullet\|_{0}$ and $\llbracket \bullet \rrbracket_{0}$. The presence of a subscript 0 in the notation is related to the fact that, in Section 2.2, those norms will be embedded in a scale of Sobolev-like norms.

We consider for the moment a general linear autonomous controlled system of the form

$$
\left\{\begin{align*}
y^{\prime}+\mathrm{L} y & =\mathrm{B} v, \text { on }(0, T)  \tag{11}\\
y(0) & =y_{0}
\end{align*}\right.
$$

where $\mathrm{L}: E \rightarrow E$ and $\mathrm{B}: U \rightarrow E$ are two linear operators, $y:[0, T] \rightarrow E$ is the state and $v:[0, T] \rightarrow U$ is the control we are looking for.

In the sequel of this paper, different such systems will be considered, coming in particular from the discretization of multi-D parabolic control problems such as (7) or their reduced version (9). In particular, the spaces $E, U$ and the operators L and B will depend on some discretization parameter $h$. We will be interested in properties of those systems that are uniform with respect to $h$, that is the reason why we will pay, in this section, a particular attention to the various constants appearing in the estimates. In section 4 we will propose a suitable framework ensuring that all those constants will be uniform with respect to $h$.
2.2. Well-posedness. It is clear that (11) is well-posed for any choice of $y_{0}$ and $v$ and that

$$
\begin{equation*}
\sup _{t \in[0, T]}\|y(t)\|_{0} \leq e^{T\|\mathrm{~L}\|}\left(\left\|y_{0}\right\|_{0}+\sqrt{T}\|\mathrm{~B}\| \llbracket v \rrbracket_{L^{2}(0, T ; U)}\right), \tag{12}
\end{equation*}
$$

where $\|\mathrm{L}\|$ (resp. $\|\mathrm{B}\|$ ) is the operator norm of $\mathrm{L}: E \rightarrow E$ (resp. B: $U \rightarrow E$ ). However, in the framework we are interested in which comes from the semi-discretization in space of an evolution PDEs, those operator norms will not be bounded in general with respect to the discretization parameter. This is the consequence at the discrete level of the fact that differential operators are naturally unbounded operators in Sobolev spaces. For example if L the discrete Laplace operator on a uniform mesh of size $h$ and B the boundary control operator as defined in section 4.1, then $\|\mathrm{B}\|$ and $\|\mathrm{L}\|$ both behave like $C / h^{2}$. Thus, inequality (12) will not give usable estimates.

Therefore, we need to introduce adapted estimates and some kind of discrete Sobolev norms to take into account the particular geometry of the (discrete) control operators under study. To this end, we introduce $\mathrm{D}: E \rightarrow E$ a self-adjoint definite
positive operator on $E$ (one can think of the discrete Laplace operator for instance) and we define a scale of inner products in $E$ defined, for any $s \in \mathbb{R}$, by

$$
\langle u, v\rangle_{s, \mathrm{D}}=\left\langle\mathrm{D}^{s} u, v\right\rangle_{0},
$$

and $\|\bullet\|_{s, \mathrm{D}}$ is the associated norms. Observe that $\|\bullet\|_{0}=\|\bullet\|_{0, \mathrm{D}}$ and that $\left(E,\langle\bullet, \bullet\rangle_{-s, \mathrm{D}}\right)$ is naturally isometric to the dual of $\left(E,\langle\bullet, \bullet\rangle_{s, \mathrm{D}}\right)$ since

$$
\begin{equation*}
\|u\|_{-s, \mathrm{D}}=\sup _{\psi \in E} \frac{\langle u, \psi\rangle_{0}}{\|\psi\|_{s, \mathrm{D}}} \tag{13}
\end{equation*}
$$

We shall now define, for given $s \in \mathbb{R}$, the two constants $M_{s, \text { adm }}>0$ and $M_{s, \text { cont }}>$ 0 that satisfy

$$
\begin{gather*}
\sup _{t \in[0, T]}\left\|e^{-t \mathrm{~L}^{*}} \psi\right\|_{s, \mathrm{D}} \leq M_{s, \text { cont }}\|\psi\|_{s, \mathrm{D}}, \quad \forall \psi \in E,  \tag{14}\\
\left.\left(\int_{0}^{T} \llbracket \mathrm{~B}^{*} e^{-\tau \mathrm{L}^{*}} \psi\right]_{0}^{2} d \tau\right)^{\frac{1}{2}} \leq M_{s, \mathrm{adm}}\|\psi\|_{s, \mathrm{D}}, \quad \forall \psi \in E \tag{15}
\end{gather*}
$$

where the adjoint operators $L^{*}$ and $\mathrm{B}^{*}$ are relative to the ambiant inner product on $E$ and $U$. Observe that (14) and (15) automatically holds since we consider finite dimension spaces, and the only interesting point is the uniformity (or not) of the constants with respect to the spaces and the operators involved. Depending on the targeted application (distributed control, Dirichlet boundary control or Neumann boundary control for instance) we will need to choose a convenient value of $s$ and of the operator D to ensure that those constants are actually uniform with respect to the discretization parameter.

Proposition 2.1. For any $v \in L^{2}(0, T ; U)$ and any $y_{0} \in E$, there exists a unique solution $y$ to (11) and it satisfies

$$
\sup _{t \in[0, T]}\|y(t)\|_{-s, \mathrm{D}} \leq M_{s, \text { cont }}\left\|y_{0}\right\|_{-s, \mathrm{D}}+M_{s, \mathrm{adm}} \llbracket v \rrbracket_{L^{2}(0, T ; U)}
$$

Proof. We write the Duhamel formula

$$
y(t)=e^{-t \mathrm{~L}} y_{0}+\int_{0}^{t} e^{-(t-s) \mathrm{L}} \mathrm{~B} v(s) d s
$$

then we take the inner product with any $\psi \in E$

$$
\langle y(t), \psi\rangle_{0}=\left\langle e^{-t \mathrm{~L}} y_{0}, \psi\right\rangle_{0}+\int_{0}^{t}\left\langle e^{-(t-\tau) \mathrm{L}} \mathrm{~B} v(\tau), \psi\right\rangle_{0} d \tau
$$

It follows that

$$
\langle y(t), \psi\rangle_{0}=\left\langle y_{0}, e^{-t \mathrm{~L}^{*}} \psi\right\rangle_{0}+\int_{0}^{t}\left[v(\tau), \mathrm{B}^{*} e^{-(t-\tau) \mathrm{L}^{*}} \psi\right]_{0} d \tau
$$

and then

$$
\left|\langle y(t), \psi\rangle_{0}\right| \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}\left\|e^{-t \mathrm{~L}^{*}} \psi\right\|_{s, \mathrm{D}}+\int_{0}^{t} \llbracket v(\tau) \rrbracket_{0} \llbracket \mathrm{~B}^{*} e^{-(t-\tau) \mathrm{L}^{*}} \psi \rrbracket_{0} d \tau
$$

By (14) and (15) we deduce that

$$
\left|\langle y(t), \psi\rangle_{0}\right| \leq\left(M_{s, \text { cont }}\left\|y_{0}\right\|_{-s, \mathrm{D}}+M_{s, \mathrm{adm}} \llbracket v \rrbracket_{L^{2}(0, T ; U)}\right)\|\psi\|_{s, \mathrm{D}}
$$

Since this is valid for any $\psi \in E$, we deduce the expected estimate by the duality property (13).

Remark 2.1. During this work we will often use the following very standard duality formula that was given in the proof above

$$
\begin{equation*}
\langle y(T), \psi\rangle_{0}-\left\langle y_{0}, e^{-T \mathrm{~L}^{*}} \psi\right\rangle_{0}=\int_{0}^{T}\left[v(t), \mathrm{B}^{*} e^{-(T-t) \mathrm{L}^{*}} \psi\right]_{0} d t, \quad \forall \psi \in E \tag{16}
\end{equation*}
$$

2.3. Relaxed observability inequalities. It is well-known (see [10, 15] for instance) that System (11) is null controllable at time $T$ for any initial condition $y_{0}$ if and only if there exists $C>0$, such that the following observability inequality for the adjoint problem is satisfied

$$
\begin{equation*}
\left\|e^{-T \mathrm{~L}^{*}} q_{T}\right\|_{0}^{2} \leq C^{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} e^{-(T-t) \mathrm{L}^{*}} q_{T} \rrbracket_{0}^{2} d t, \quad \forall q_{T} \in E . \tag{17}
\end{equation*}
$$

The value of the constant $C$ in this inequality is crucial since it appears in the measure of the control cost.

It happens that, when (11) comes from a discretization of a parabolic equation with the finite difference method then the null controllability of the semi-discrete system may not hold (in particular in a multi dimensional setting, see the example given by Kavian and reported in [16]). To tackle this problem, it was proposed (in $[11,7,8,6]$ for instance) to relax the controllability requirements by considering instead the $\varphi(h)$-null-controllability of (11). It consists in constructing uniformly bounded controls such that the solution $y(T)$ does not identically vanish but is small enough with respect to the discretization parameter $h$. This approach is based on the penalized HUM construction, where the penalization parameter is a given function $h \rightarrow \varphi(h)$ of the discretization parameter, given its name to this notion. Note that, the spaces $E, U$ and the operators L and B all depend on $h$, in particular the dimensions of $E$ and $U$ may increase when $h$ tends to zero. This is one of the main difficulty that we need to take care of in the analysis.

The $\varphi(h)$-null-controllability property is equivalent to a relaxed version of inequality (17) and the following Lemma 2.1, whose proof is given in appendix A, aims at establishing such an equivalence. This lemma is somehow related to [14, Lemma 3.4] or [3, Proposition 1] and is stated in a quite general framework : the constant $s$ and the operator D can be chosen arbitrarily. Moreover using appropriate spaces $F_{0}$ and $F_{T}$ (which are defined below), one can show that Lemma 2.1 encompasses some already known situations (see Remark 2.3). Roughly speaking, this Lemma is about controlling the components in the final state space $F_{T}$ of the solution of system (24) which starts from an initial condition $y_{0}$ in the initial state space $F_{0}$. Even though we only state it in a finite dimensional setting, it is clear that infinite dimensional versions also hold, as in the references quoted above.

Let $F_{0}$ and $F_{T}$ be two subspaces of $E$ and $\mathrm{P}_{F_{0}}$ (resp. $\mathrm{P}_{F_{T}}$ ) the orthogonal projection onto $F_{0}\left(\right.$ resp. $\left.F_{T}\right)$ with respect to the inner product $\langle\bullet, \bullet\rangle_{-s, \mathrm{D}}$.

We will denote the adjoint operators of the projectors $\mathrm{P}_{F_{0}}$ and $\mathrm{P}_{F_{T}}$ for the inner product $\langle\bullet \bullet \bullet\rangle_{0, \mathrm{D}}$ by $\mathrm{P}_{F_{0}}^{*}$ and $\mathrm{P}_{F_{T}}^{*}$. Observe that $\mathrm{P}_{F_{0}}^{*}$ and $\mathrm{P}_{F_{T}}^{*}$ are also the orthogonal
projectors in $\left(E,\langle\bullet, \bullet\rangle_{s, \mathrm{D}}\right)$ onto $\mathrm{D}^{-s} F_{0}$ and $\mathrm{D}^{-s} F_{T}$ respectively. In particular, we have

$$
\begin{aligned}
\left\|\mathrm{P}_{F_{0}}^{*} y\right\|_{s, \mathrm{D}} \leq\|y\|_{s, \mathrm{D}}, \quad \forall y \in E, \\
\left\|\mathrm{P}_{F_{T}}^{*} y\right\|_{s, \mathrm{D}} \leq\|y\|_{s, \mathrm{D}}, \quad \forall y \in E .
\end{aligned}
$$

REmARK 2.2. Observe that, if $F_{0}$ (resp. $F_{T}$ ) is stable by D then $\mathrm{P}_{F_{0}}$ (resp. $\mathrm{P}_{F_{T}}$ ) does not depend on $s$. In particular, those projections are orthogonal for the inner product $\langle\bullet, \bullet\rangle_{0}$.

Dealing with such subspaces will be crucial in the sequel when we will apply the Lebeau-Robbiano strategy since it requires to be able to control some precise components of the solution at each step (depending on eigenspaces of $\mathrm{A}_{2}$ ), see Section 3.3.

Lemma 2.1 (Relaxed observability inequalities and controllability). We use the above notations and assume that s, D , $F_{0}$ and $F_{T}$ are given. Let $M_{\mathrm{obs}}>0$ and $M_{\mathrm{rel}} \geq 0$ be two given numbers.

The following two propositions are equivalent.

1. For any $y_{0} \in F_{0}$ there exists a $v \in L^{2}(0, T ; U)$ satisfying

$$
\begin{equation*}
\frac{1}{M_{\mathrm{obs}}^{2}} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+\frac{1}{M_{\mathrm{rel}}^{2}}\left\|\mathrm{P}_{F_{T}}(y(T))\right\|_{-s, \mathrm{D}}^{2} \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}^{2} \tag{18}
\end{equation*}
$$

where $y$ is the corresponding solution of (11).
2. For any $q_{T} \in \mathrm{D}^{-s} F_{T}$, the following relaxed observability inequality holds:

$$
\begin{equation*}
\left\|\mathrm{P}_{F_{0}}^{*}\left(e^{-T \mathrm{~L}^{*}} q_{T}\right)\right\|_{s, \mathrm{D}}^{2} \leq M_{\mathrm{obs}}^{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} e^{-(T-t) \mathrm{L}^{*}} q_{T}\left\|_{0}^{2} d t+M_{\mathrm{rel}}^{2}\right\| q_{T} \|_{s, \mathrm{D}}^{2} \tag{19}
\end{equation*}
$$

Of course, for $M_{\text {rel }}=0$, the inequality (18) should be understood as

$$
\left\{\begin{array}{l}
\frac{1}{M_{\mathrm{obs}}^{2}} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2} \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}^{2}, \\
\left\|\mathrm{P}_{F_{T}}(y(T))\right\|_{-s, \mathrm{D}}=0
\end{array}\right.
$$

REMARK 2.3. Throughout this paper $F_{0}$ will always be equal to the whole space $E$. However it is worth noticing that by specifying spaces $F_{0}$ and $F_{T}$, one can recover usual inequalities related to different notions of controllability.

- When $F_{0}=F_{T}=E$, inequality (19) is the usual relaxed inequality. If (19) holds with $M_{\mathrm{rel}}=0$ then system (11) is null controllable.
As explained above, we cannot always expect $M_{\mathrm{rel}}$ to be equal to zero when system (11) is discretized by finite differences method with a space domain of dimension greater than one. However, if this inequality holds with $M_{\text {rel }}^{2}=$ $\varphi(h)$ and with $M_{\text {obs }}$ independent of $h$, we recover the $\varphi(h)$-null-controllability notion briefly described above.
- When $\operatorname{dim}\left(F_{0}\right)=1$ and $F_{T}=E$, then proving (19) amounts to drive only one given initial condition to zero. This question is tackled for instance in [6].
- The partial null-controllability consists in driving to zero only some components of the solution of a system of parabolic PDEs. It amounts to prove inequality (19) for $M_{\mathrm{rel}}=0$ and to choose an appropriate subspace $F_{T}$. This kind of controllability is studied for instance in [3] where related $\varphi(h)$-partial-null-controllability is also investigated.


## 3. Controllability of tensorized systems.

3.1. Notations. Let $\left(E_{i},\langle\bullet \bullet \bullet\rangle_{0, i}\right), i=1,2$ two finite dimensional Euclidean spaces of dimensions $N_{1}$ and $N_{2}$. The associated norms are denoted by $\|\bullet\|_{0, i}$. Let $\mathrm{D}_{i}$ be two positive definite self-adjoint operators in those spaces and for any $s \in \mathbb{R}$, we introduce the following scalar products

$$
\left\langle u_{i}, v_{i}\right\rangle_{s, \mathrm{D}_{i}}=\left\langle\mathrm{D}_{i}^{s} u_{i}, v_{i}\right\rangle_{0, i}, \quad \forall u_{i}, v_{i} \in E_{i}
$$

and the associated norms $\|\bullet\|_{s, \mathrm{D}_{i}}$.
In the case of a finite difference approximate system, the two spaces $E_{i}$ have to be understood as the space of discrete in space (scalar) functions defined on a grid of $\Omega_{i}$. We will then consider the tensor product space $E_{1} \otimes E_{2}$ as a natural discretization space for functions defined on the tensor product grid of $\Omega$. This space is equipped with the natural Euclidean structure defined by

$$
\begin{equation*}
\left\langle u_{1} \otimes u_{2}, v_{1} \otimes v_{2}\right\rangle_{0}=\left\langle u_{1}, v_{1}\right\rangle_{0,1}\left\langle u_{2}, v_{2}\right\rangle_{0,2} \tag{20}
\end{equation*}
$$

On $E_{1} \otimes E_{2}$ we consider the operator $\mathrm{D}=\mathrm{D}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{D}_{2}$, where I stands for the identity operator in $E_{1}$ and $E_{2}$. This is a positive definite self-adjoint operator, from which we can define the following natural inner products and associated norms

$$
\langle u, v\rangle_{s, \mathrm{D}}=\left\langle\mathrm{D}^{s} u, v\right\rangle_{0}, \quad \forall u, v \in E_{1} \otimes E_{2}
$$

Note, in particular that we have

$$
\begin{align*}
\left\langle u_{1} \otimes u_{2}, v_{1} \otimes v_{2}\right\rangle_{1, \mathrm{D}}=\left\langle u_{1}, v_{1}\right\rangle_{1, \mathrm{D}_{1}}\left\langle u_{2}, v_{2}\right\rangle_{0,2}+ & \left\langle u_{1}, v_{1}\right\rangle_{0,1}\left\langle u_{2}, v_{2}\right\rangle_{1, \mathrm{D}_{2}}  \tag{21}\\
& \forall u_{1}, v_{1} \in E_{1}, \forall u_{2}, v_{2} \in E_{2} .
\end{align*}
$$

Since we will be interested in vector-valued discrete functions that aim at being approximations of the solution (at any time $t$ ) of (6), we will naturally work with the space

$$
E=E_{1} \otimes E_{2} \otimes \mathbb{R}^{n}
$$

which corresponds to the fact that all the $n$ components of the system are approximated at each grid point.

We recall that if $L_{i}$ is a linear operator in $E_{i}$ and $L$ a linear operator in $\mathbb{R}^{n}$, the tensor product $\mathrm{L}_{1} \otimes \mathrm{~L}_{2} \otimes \mathrm{~L}$ is a linear operator on $E$ defined by

$$
\left(\mathrm{L}_{1} \otimes \mathrm{~L}_{2} \otimes \mathrm{~L}\right)\left(u_{1} \otimes u_{2} \otimes z\right)=\mathrm{L}_{1}\left(u_{1}\right) \otimes \mathrm{L}_{2}\left(u_{2}\right) \otimes \mathrm{L}(z), \quad \forall u_{i} \in E_{i}, \forall z \in \mathbb{R}^{n}
$$

Let $(\bullet, \bullet)$ be the Euclidean inner product in $\mathbb{R}^{n}$ and $|\bullet|$ the associated norm. The previous definitions are naturally extended to the space $E$ as follows

$$
\langle u \otimes y, v \otimes z\rangle_{s, \mathrm{D}}=\langle u, v\rangle_{s, \mathrm{D}}(y, z), \quad \forall u, v \in E_{1} \otimes E_{2}, \quad \forall y, z \in \mathbb{R}^{n}
$$

and we do similar extensions for spaces $E_{1} \otimes \mathbb{R}^{n}$ and $E_{2} \otimes \mathbb{R}^{n}$.
Recall that $E$ can be identified to $E_{1} \otimes \mathbb{R}^{n} \otimes E_{2}$ and $E_{2} \otimes \mathbb{R}^{n} \otimes E_{1}$. Therefore, if $u:=u_{i} \otimes y \in E_{i} \otimes \mathbb{R}^{n}$ and $v_{j} \in E_{j}$ where $i \neq j \in\{1,2\}$, one is allowed to consider $u \otimes v_{j}$ as an element of $E$, although one should instead deal with the element $u_{i} \otimes v_{j} \otimes y$.

We now introduce two other finite dimensional Euclidean spaces $\left(U_{i},[\bullet, \bullet]_{0, i}\right)$, $i=1,2$, with the associated norms $\llbracket \bullet \rrbracket_{0, i}$ that correspond to the control space for each subproblem.

We define $U:=U_{1} \otimes U_{2} \otimes \mathbb{R}^{m}$ and its inner product $[\bullet, \bullet]_{0}$, whose definition on $U_{1} \otimes U_{2}$ is analogous to (20) :

$$
\begin{equation*}
\left[u_{1} \otimes u_{2}, v_{1} \otimes v_{2}\right]_{0}=\left[u_{1}, v_{1}\right]_{0,1}\left[u_{2}, v_{2}\right]_{0,2}, \tag{22}
\end{equation*}
$$

and is extended to $U$ as before.
3.2. Main Theorem. Let $\mathrm{A}_{i}$ be a symmetric definite positive operator in $E_{i}$. One can think of $\mathrm{A}_{i}$ as an approximation of the continuous operator $\mathcal{A}_{i}$ but the statement of our result is generic and does not explicitly make use of such an assumption. The eigenvalues of $\mathrm{A}_{i}$ will be denoted by $\left(\lambda_{i, k}\right)_{k=1}^{N_{i}}$ and the corresponding eigenfunctions are $\left(\phi_{i, k}\right)_{k=1}^{N_{i}}$. Those form an orthonormal family for $\langle\bullet, \bullet\rangle_{0, i}$.

In our estimates, the following discrete Poincaré inequality for $\mathrm{A}_{1}$, will be needed

$$
\begin{equation*}
\|u\|_{0}=\|u\|_{0, \mathrm{~A}_{1}} \leq M_{\mathrm{P}, 1}\|u\|_{1, \mathrm{~A}_{1}}, \forall u \in E_{1}, \tag{23}
\end{equation*}
$$

for some $M_{\mathrm{P}, 1}$. Without loss of generality, we will assume that $M_{\mathrm{P}, 1} \geq 1$. Note that the best possible value for $M_{\mathrm{P}, 1}$ is $\lambda_{1,1}^{-1 / 2}$ but it is not sure that this value is greater than 1 . Note also that the following generalised Poincaré estimate holds

$$
\|u\|_{0} \leq M_{\mathrm{P}, 1}^{s}\|u\|_{s, \mathrm{~A}_{1}}, \forall u \in E_{1}, \forall s \geq 0
$$

For each $i=1,2$, we consider a control operator $\mathrm{B}_{i}: U_{i} \rightarrow E_{i}$, a coupling matrix $C$ and a control matrix $B$.

Our goal is to analyse the controllability properties (uniform with respect to any parameter on which the system may depend) of the following tensorized ODE

$$
\left\{\begin{array}{l}
\partial_{t} y+L y=\mathrm{B}_{1} \otimes \mathrm{~B}_{2} \otimes \mathrm{~B} v,  \tag{24}\\
y(0)=y_{0} \in E
\end{array}\right.
$$

where the control $v$ belongs to $L^{2}(0, T ; U)$ and

$$
L:=\mathrm{A}+\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C} \text { and } \mathrm{A}:=\mathrm{A}_{1} \otimes \mathrm{I} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2} \otimes \mathrm{I} .
$$

Here, the symbol I stands for the identity operator on $E_{1}$ or $E_{2}$ and the symbol $I$ is the identity operator of $\mathbb{R}^{n}$.

To this end, we will take benefit from the tensor product structure of the system and make two main assumptions:

1. the first one concerns the vector-valued sub-control-system in $E_{1}$ defined by

$$
\left\{\begin{array}{l}
\partial_{t} y+\mathrm{A}_{1} \otimes \mathrm{I} y+\mathrm{I} \otimes \mathrm{C} y=\mathrm{B}_{1} \otimes \mathrm{~B} v,  \tag{25}\\
y(0)=y_{0} \in E_{1} \otimes \mathbb{R}^{n},
\end{array}\right.
$$

with $y \in \mathcal{C}^{0}\left([0, T], E_{1} \otimes \mathbb{R}^{n}\right)$ and $v \in L^{2}\left(0, T ; U_{1} \otimes \mathbb{R}^{m}\right)$.
Assumption 3.1. For some $M_{\text {obs }, 1}>0, \mu_{E_{1}} \in(0,+\infty]$, we have that for any $T>0$, for any $y_{0} \in E_{1} \otimes \mathbb{R}^{n}$, there exists a control $v \in L^{2}\left(0, T ; U_{1} \otimes \mathbb{R}^{m}\right)$ for (25) such that

$$
\left\{\begin{aligned}
\llbracket v \rrbracket_{L^{2}\left(0, T ; U_{1} \otimes \mathbb{R}^{m}\right)} & \leq e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left\|y_{0}\right\|_{-s, \mathbf{A}_{1}} \\
\|y(T)\|_{-s, \mathrm{~A}_{1}} & \leq e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} T}\left\|y_{0}\right\|_{-s, \mathrm{~A}_{1}}
\end{aligned}\right.
$$

2. the second one concerning a spectral property related with the pair of operators $\left(\mathrm{A}_{2}, \mathrm{~B}_{2}\right)$ similar to the well-known Lebeau-Robbiano inequality, excepted that it only holds for a certain portion of the spectrum of $\mathrm{A}_{2}$.
ASSUMPTION 3.2. For some $M_{\mathrm{LR}, 2}>0, \mu_{E_{2}} \in(0,+\infty]$, we have the following partial Lebeau-Robbiano spectral inequality

$$
\begin{equation*}
\|\psi\|_{0} \leq e^{M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \llbracket \mathrm{B}_{2}^{*} \psi \rrbracket_{0,2}, \quad \forall \psi \in \operatorname{span}\left(\phi_{2, j}, \quad \lambda_{2, j} \leq \mu\right) \tag{26}
\end{equation*}
$$

for any $0<\mu<\mu_{E_{2}}$.
The main theorem of this paper is the following.
THEOREM 3.1. Let $\mu^{*}=\min \left(\mu_{E_{1}}, \mu_{E_{2}}\right)$. There exist a $M_{\mathrm{obs}}>0$ depending only on $M_{\mathrm{obs}, 1}, M_{\mathrm{LR}, 2}, M_{s, \mathrm{cont}}, M_{s, \mathrm{adm}}$ and $\|\mathrm{C}\|$ such that for any $y_{0} \in E$ and any $T>0$ there exists a control $v \in L^{2}(0, T ; U)$ satisfying

$$
\left\{\begin{array}{l}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)} e^{-\mu^{*} T / 4}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{array}\right.
$$

Note that in this theorem the constant $M_{\text {obs, } 1}$ is built upon the Sobolev norms associated with $\mathrm{D}_{1}=\mathrm{A}_{1}$ and the constants $M_{s, \text { cont }}, M_{s, \text { adm }}$ with the Sobolev norms associated with $\mathrm{D}=\mathrm{A}$. Those norms are thus problem dependent.

In the case where $A_{i}$ are discrete version of diffusion operators, it can be tempting to use instead usual discrete Sobolev norms, that are defined by using for $\mathrm{D}_{1}$ and D the discrete Laplace operators. However, the equivalence between those norms, for large values of $s$, is uniform with respect to the discretization parameter if and only if the diffusion coefficients $\gamma_{i}$ are smooth enough. In the case of non smooth diffusion coefficients, using the norms given in the theorem is mandatory.
3.3. Proof of the main result. In this section, we will prove Theorem 3.1.
3.3.1. Preliminary estimates. We start with some preliminary results.

Lemma 3.1 (Norms comparison). For any $s \geq 0$, and any

$$
q=\sum_{j=1}^{N_{2}} q_{j} \otimes \phi_{2, j}, \quad q_{j} \in E_{1} \otimes \mathbb{R}^{n}
$$

we have

$$
\sum_{j=1}^{N_{2}}\left\|q_{j}\right\|_{s, \mathrm{~A}_{1}}^{2} \leq\|q\|_{s, \mathrm{~A}}^{2} \leq 2^{s} M_{\mathrm{P}, 1}^{2 s} \sum_{j=1}^{N_{2}}\left(1+\lambda_{2, j}^{s}\right)\left\|q_{j}\right\|_{s, \mathrm{~A}_{1}}^{2}
$$

Proof. We write each $q_{j}$ under the form

$$
q_{j}=\sum_{k=1}^{N_{1}} \phi_{1, k} \otimes q_{k, j}
$$

with $q_{k, j} \in \mathbb{R}^{n}$ in such a way that

$$
\mathrm{A}^{s} q=\sum_{\substack{1 \leq k \leq N_{1} \\ 1 \leq j \leq N_{2}}}\left(\lambda_{1, k}+\lambda_{2, j}\right)^{s} \phi_{1, k} \otimes \phi_{2, j} \otimes q_{k, j}
$$

It follows that

$$
\|q\|_{s, \mathrm{~A}}^{2}=\sum_{\substack{1 \leq k \leq N_{1} \\ 1 \leq j \leq N_{2}}}\left(\lambda_{1, k}+\lambda_{2, j}\right)^{s}\left|q_{k, j}\right|^{2}
$$

and we immediately obtain

$$
\|q\|_{s, \mathrm{~A}}^{2} \geq \sum_{j=1}^{N_{2}}\left(\sum_{k=1}^{N_{1}} \lambda_{1, k}^{s}\left|q_{k, j}\right|^{2}\right)=\sum_{j=1}^{N_{2}}\left\|q_{j}\right\|_{s, \mathrm{~A}_{1}}^{2} .
$$

Moreover, we have

$$
\begin{aligned}
\|q\|_{s, \mathrm{~A}}^{2} & \leq 2^{s} \sum_{\substack{1 \leq k \leq N_{1} \\
1 \leq j \leq N_{2}}}\left(\lambda_{1, k}^{s}+\lambda_{2, j}^{s}\right)\left|q_{k, j}\right|^{2} \\
& =2^{s} \sum_{j=1}^{N_{2}}\left\|q_{j}\right\|_{s, \mathrm{~A}_{1}}^{2}+2^{s} \sum_{j=1}^{N_{2}} \lambda_{2, j}^{s}\left\|q_{j}\right\|_{0, \mathrm{~A}_{1}}^{2} \\
& \leq 2^{s} M_{\mathrm{P}, 1}^{2 s} \sum_{j=1}^{N_{2}}\left(1+\lambda_{2, j}^{s}\right)\left\|q_{j}\right\|_{s, \mathrm{~A}_{1}}^{2}
\end{aligned}
$$

and the claim follows.
We define, for any $\mu \geq 0$, the following subspaces of $E_{2}$

$$
E_{\mu, 2}:=\operatorname{span}\left\{\phi_{2, j}: \lambda_{2, j} \leq \mu\right\}, \text { and } E_{\mu, 2}^{\perp}:=\operatorname{span}\left\{\phi_{2, j}: \lambda_{2, j}>\mu\right\}
$$

and the corresponding subspaces of $E$

$$
E_{\mu}:=E_{1} \otimes E_{\mu, 2} \otimes \mathbb{R}^{n}, \text { and } E_{\mu}^{\perp}:=E_{1} \otimes E_{\mu, 2}^{\perp} \otimes \mathbb{R}^{n}
$$

Let $\mathrm{P}_{\mu}$ (resp. $\mathrm{P}_{\mu}^{\perp}$ ) be the orthogonal projection in $E$ onto $E_{\mu}$ (resp. $E_{\mu}^{\perp}$ ) for the inner product $\langle\bullet, \bullet\rangle_{-s, \mathrm{~A}}$. By construction we have $\mathrm{P}_{\mu}+\mathrm{P}_{\mu}^{\perp}=\mathrm{I}$.

Let us now state some properties of the uncontrolled system

$$
\left\{\begin{array}{l}
\partial_{t} y+L y=0  \tag{27}\\
y(0)=y_{0} \in E
\end{array}\right.
$$

Proposition 3.1 (Dissipation estimates). 1. For any $\mu \geq 0$, if $y_{0} \in E_{\mu}^{\perp}$, then the unique solution to (27) satisfies $y(t) \in E_{\mu}^{\perp}$ for any $t$ and moreover, for any $s \in \mathbb{R}$, we have

$$
\|y(t)\|_{-s, \mathrm{~A}} \leq e^{\|\mathrm{C}\| t} e^{-t \mu}\left\|y_{0}\right\|_{-s, \mathrm{~A}}, \quad \forall t \geq 0
$$

2. For any $\mu \geq 0$, and any $y_{0} \in E$, the unique solution to (27) satisfies, for any $s \in \mathbb{R}$,

$$
\begin{gathered}
\left\|\mathrm{P}_{\mu} y(t)\right\|_{-s, \mathrm{~A}} \leq e^{\|\mathrm{C}\| t}\left\|\mathrm{P}_{\mu} y_{0}\right\|_{-s, \mathrm{~A}} \\
\left\|\mathrm{P}_{\mu}^{\perp} y(t)\right\|_{-s, \mathrm{~A}} \leq e^{\|\mathrm{C}\| t-\mu t}\left\|\mathrm{P}_{\mu}^{\perp} y_{0}\right\|_{-s, \mathrm{~A}} \leq e^{\|\mathrm{C}\| t-\mu t}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{gathered}
$$

Proof. 1. The first point is just a consequence of the fact that $E_{\mu}$ and $E_{\mu}^{\perp}$ are stable by $\mathrm{L}=\mathrm{A}+\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C}$. A straightforward computation, using that A commutes with $I \otimes I \otimes C$ shows that

$$
\frac{1}{2} \partial_{t}\left\langle\mathrm{~A}^{-s} y, y\right\rangle_{0}+\left\langle\mathrm{AA}^{-\frac{s}{2}} y, \mathrm{~A}^{-\frac{s}{2}} y\right\rangle_{0}=-\left\langle\mathrm{A}^{-s}(\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C}) y, y\right\rangle_{0}
$$

Then we observe that, $\mathrm{A}^{-\frac{s}{2}} y(t) \in E_{\mu}^{\perp}$ for any $t$, and that by definition of $E_{\mu}^{\perp}$ we have

$$
\langle\mathrm{A} z, z\rangle_{0} \geq \mu\langle z, z\rangle_{0}, \quad \forall z \in E_{\mu}^{\perp}
$$

It follows that

$$
\frac{1}{2} \partial_{t}\left\langle\mathrm{~A}^{-s} y, y\right\rangle_{0} \leq(\|\mathrm{C}\|-\mu)\left\langle\mathrm{A}^{-s} y, y\right\rangle_{0}
$$

and the claim follows by the differential form of Gronwall's Lemma.
2. We simply observe that $\mathrm{P}_{\mu} y$ and $\mathrm{P}_{\mu}^{\perp} y$ solve the same equation as $y$ with initial conditions $\mathrm{P}_{\mu} y_{0} \in E_{0}^{\perp}, \mathrm{P}_{\mu}^{\perp} y_{0} \in E_{\mu}^{\perp}$. It is then enough to use the inequality of the first point to conclude.
3.3.2. Partial controllability results. The following proposition is a partial version of Theorem 3.1. More precisely, we establish the existence of a control $v$ that reduces the norm of the projection on $E_{\mu}$ of the final state $y(T)$ as much as possible. However, the bound on the control is not uniform with respect to $\mu$ yet.

Proposition 3.2. There exists a positive number $M_{\mathrm{part}}$ which depends only on $s$, $M_{\mathrm{P}, 1}, M_{\mathrm{obs}, 1}, M_{\mathrm{LR}, 2}, M_{s, \mathrm{cont}}, M_{s, \mathrm{adm}}$ such that for any $\mu \in\left(0, \mu_{E_{2}}\right)$ and any $y_{0} \in E$, there exists a control $v \in L^{2}(0, T ; U)$ that satisfies

$$
\left\{\begin{array}{l}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq e^{M_{\mathrm{part}}\left(1+\frac{1}{T}+\sqrt{\mu}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}}  \tag{28}\\
\left\|\mathrm{P}_{\mu}(y(T))\right\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{part}}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} T}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{array}\right.
$$

and

$$
\begin{equation*}
\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{part}}\left(1+\frac{1}{T}+\sqrt{\mu}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \tag{29}
\end{equation*}
$$

Proof. Let $q_{T} \in E_{\mu}$ and $q$ be the solution of the backward equation

$$
\left\{\begin{aligned}
-\partial_{t} q+\mathrm{L}^{*} q & =0 \\
q(T) & =q_{T} .
\end{aligned}\right.
$$

We recall that $\mathrm{A}_{1}\left(\right.$ resp. $\left.\mathrm{A}_{2}\right)$ is self-adjoint in $E_{1}$ (resp. $E_{2}$ ) so that

$$
\mathrm{L}^{*}=\mathrm{A}_{1} \otimes \mathrm{I} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C}^{*}
$$

Since $E_{\mu}$ is stable by $\mathrm{L}^{*}$, we can decompose $q$ in the following way

$$
q(t)=\sum_{\lambda_{2, j} \leq \mu} q_{j}(t) \otimes \phi_{2, j}, \text { with } q_{j}(t) \in E_{1} \otimes \mathbb{R}^{n}
$$

where $q_{j}$ satisfies

$$
\left\{\begin{aligned}
-\partial_{t} q_{j}+\left(\mathrm{A}_{1} \otimes \mathbf{I}\right) q_{j}+\lambda_{2, j} q_{j}+\left(\mathrm{I} \otimes \mathrm{C}^{*}\right) q_{j} & =0 \\
q_{j}(T) & =q_{T, j} \in E_{1} \otimes \mathbb{R}^{n}
\end{aligned}\right.
$$

The new variable $z_{j}(t):=q_{j}(t) e^{\lambda_{2, j}(T-t)}$ satisfies the following system

$$
\left\{\begin{aligned}
-\partial_{t} z_{j}+\left(\mathrm{A}_{1} \otimes \mathrm{I}\right) z_{j}+\left(\mathrm{I} \otimes \mathrm{C}^{*}\right) z_{j}(t) & =0 \\
z_{j}(T) & =q_{T, j}
\end{aligned}\right.
$$

Thanks to Assumption 3.1, we can apply Lemma 2.1 in the space $E_{1} \otimes \mathbb{R}^{n}$, with $F_{0}=F_{T}=E_{1} \otimes \mathbb{R}^{n}, \mathrm{~L}=\mathrm{A}_{1} \otimes \mathbf{I}+\mathrm{I} \otimes \mathrm{C}, \mathrm{B}=\mathrm{B}_{1} \otimes \mathrm{~B}, \mathrm{D}=\mathrm{A}_{1} \otimes \boldsymbol{I}, M_{\mathrm{obs}}=e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}$ and $M_{\mathrm{rel}}=e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} T}$. Inequality (19) with $q_{T}=q_{T, j}$ gives

$$
\left\|z_{j}(0)\right\|_{s, \mathrm{~A}_{1}}^{2} \leq e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \int_{0}^{T} \llbracket \mathrm{~B}_{1}^{*} \otimes \mathrm{~B}^{*} z_{j}(t) \rrbracket_{0,1}^{2} d t+e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} 2 T}\left\|q_{T, j}\right\|_{s, \mathrm{~A}_{1}}^{2}
$$

and therefore, coming back to the variable $q_{j}$, we get

$$
\begin{aligned}
\left\|q_{j}(0)\right\|_{s, \mathrm{~A}_{1}}^{2} e^{2 \lambda_{2, j} T} \leq e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \int_{0}^{T} & \llbracket \mathrm{~B}_{1}^{*}
\end{aligned} \begin{aligned}
& \mathrm{B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} e^{2 \lambda_{2, j}(T-t)} d t \\
& +e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} 2 T}\left\|q_{T, j}\right\|_{s, \mathrm{~A}_{1}}^{2}
\end{aligned}
$$

and thus,

$$
\begin{aligned}
\left\|q_{j}(0)\right\|_{s, \mathrm{~A}_{1}}^{2} \leq e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \int_{0}^{T} \llbracket \mathrm{~B}_{1}^{*} & \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} d t \\
& +e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} 2 T} e^{-2 \lambda_{2, j} T}\left\|q_{T, j}\right\|_{s, \mathrm{~A}_{1}}^{2}
\end{aligned}
$$

Using the second inequality in Lemma 3.1, we get

$$
\begin{aligned}
\|q(0)\|_{s, \mathrm{~A}}^{2} & \leq 2^{s} M_{\mathrm{P}, 1}^{2 s} e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \sum_{\lambda_{2, j} \leq \mu}\left(\left(1+\lambda_{2, j}^{s}\right) \int_{0}^{T} \llbracket \mathrm{~B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} d t\right. \\
& \left.+\left(1+\lambda_{2, j}^{s}\right) e^{-2 \lambda_{2, j} T} e^{-\mu_{E_{1}} 2 T}\left\|q_{T, j}\right\|_{s, \mathrm{~A}_{1}}^{2}\right)
\end{aligned}
$$

Applying now the first inequality in Lemma 3.1 and the following inequality $\lambda^{s} e^{-2 \lambda T} \leq$ $\frac{s^{s}}{(2 T)^{s}}$, we get

$$
\begin{aligned}
\|q(0)\|_{s, \mathrm{~A}}^{2} & \leq 2^{s} M_{\mathrm{P}, 1}^{2 s}\left(1+\mu^{s}\right) e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \sum_{\lambda_{2, j} \leq \mu} \int_{0}^{T} \llbracket \mathrm{~B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} d t \\
& +2^{s} M_{\mathrm{P}, 1}^{2 s} e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left(1+\frac{s^{s}}{(2 T)^{s}}\right) e^{-\mu_{E_{1}} 2 T}\left\|q_{T}\right\|_{s, \mathrm{~A}}^{2}
\end{aligned}
$$

We shall now apply Assumption 3.2. To this end we choose any orthonormal basis $\left(\Psi_{k}\right)_{k=1, \ldots, K}$ of $U_{1} \otimes \mathbb{R}^{n}$ and we decompose each $\mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t)$ in this basis, the coefficients being denoted by $a_{k, j}(t)$. It follows

$$
\sum_{\lambda_{2, j} \leq \mu} \llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2}=\sum_{\lambda_{2, j} \leq \mu} \llbracket \sum_{k=1}^{K} a_{k, j}(t) \Psi_{k} \rrbracket_{0,1}^{2}=\sum_{k=1}^{K} \sum_{\lambda_{2, j} \leq \mu}\left(a_{k, j}(t)\right)^{2} .
$$

For any $k \in\{1, \ldots, K\}$, given that $0<\mu<\mu_{E_{2}}$, we can apply the discrete LebeauRobbiano inequality given by Assumption 3.2 to the vector $\psi=\sum_{\lambda_{2, j} \leq \mu} a_{k, j}(t) \phi_{2, j}$, to
obtain

$$
\begin{aligned}
& \left.\sum_{\lambda_{2, j} \leq \mu} \llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} \leq \sum_{k=1}^{K} e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \llbracket \sum_{\lambda_{2, j} \leq \mu} a_{k, j}(t) \mathrm{B}_{2}^{*} \phi_{2, j}\right]_{0,2}^{2} \\
& \leq e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \sum_{k=1}^{K} \sum_{\lambda_{2, j} \leq \mu} \sum_{\lambda_{2, j^{\prime}} \leq \mu}\left[a_{k, j}(t) \mathrm{B}_{2}^{*} \phi_{2, j}, a_{k, j^{\prime}}(t) \mathrm{B}_{2}^{*} \phi_{2, j^{\prime}}\right]_{0,2}
\end{aligned}
$$

Apply now (22) and the fact that for any $k, \llbracket \Psi_{k} \rrbracket_{0,1}=1$ :

$$
\begin{aligned}
& \sum_{\lambda_{2, j} \leq \mu} \llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j}(t) \rrbracket_{0,1}^{2} \\
& \leq e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \sum_{k=1}^{K} \sum_{\lambda_{2, j} \leq \mu} \sum_{\lambda_{2, j^{\prime}} \leq \mu}\left[a_{k, j}(t) \Psi_{k} \otimes \mathrm{~B}_{2}^{*} \phi_{2, j}, a_{k, j^{\prime}}(t) \Psi_{k} \otimes \mathrm{~B}_{2}^{*} \phi_{2, j^{\prime}}\right]_{0} \\
& \leq e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \sum_{\lambda_{2, j} \leq \mu} \sum_{\lambda_{2, j^{\prime}} \leq \mu}\left[\mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} \otimes \mathrm{~B}_{2}^{*}\left(q_{j}(t) \otimes \phi_{2, j}\right), \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} \otimes \mathrm{~B}_{2}^{*}\left(q_{j^{\prime}}(t) \otimes \phi_{2, j^{\prime}}\right)\right]_{0} \\
& \leq e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})} \llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} .
\end{aligned}
$$

Hence,

$$
\begin{align*}
\|q(0)\|_{s, \mathrm{~A}}^{2} & \leq 2^{s} M_{\mathrm{P}, 1}^{2 s}\left(1+\mu^{s}\right) e^{2 M_{\mathrm{LR}, 2}(1+\sqrt{\mu})+2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \int_{0}^{T} \llbracket \mathrm{~B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t  \tag{30}\\
& +2^{s} M_{\mathrm{P}, 1}^{2 s} e^{2 M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left(1+\frac{s^{s}}{T^{s}}\right) e^{-\mu_{E_{1}} 2 T}\left\|q_{T}\right\|_{s, \mathrm{~A}}^{2}
\end{align*}
$$

Note that, by construction, $E_{\mu}$ is stable by A so that (30) is valid for any $q_{T} \in F_{T}=$ $\mathrm{D}^{-s} E_{\mu}$. Thus, we can apply Lemma 2.1 in the space $E$ with $\mathrm{D}=\mathrm{A}, F_{0}=E$ and $F_{T}=E_{\mu}$. It follows that there exists a control $v \in L^{2}(0, T ; U)$ such that

$$
\left\{\begin{aligned}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} & \leq 2^{\frac{s}{2}} M_{\mathrm{P}, 1}^{s} \sqrt{1+\mu^{s}} e^{M_{\mathrm{LR}, 2}(1+\sqrt{\mu})+M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\left\|\mathrm{P}_{\mu}(y(T))\right\|_{-s, \mathrm{~A}} & \leq 2^{\frac{s}{2}} M_{\mathrm{P}, 1}^{s} e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)} \sqrt{1+\frac{s^{s}}{T^{s}}} e^{-\mu_{E_{1}} T}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{aligned}\right.
$$

Finally, inequality (29) comes from Proposition 2.1 and the estimate on the control cost just proved which give

$$
\begin{aligned}
& \|y(T)\|_{-s, \mathrm{~A}} \\
& \quad \leq\left(M_{s, \mathrm{cont}}+M_{s, \mathrm{adm}} 2^{\frac{s}{2}} M_{\mathrm{P}, 1}^{s} \sqrt{1+\mu^{s}} e^{M_{\mathrm{LR}, 2}(1+\sqrt{\mu})+M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{aligned}
$$

Those estimates can easily be put into the expected form for a suitable choice of $M_{\text {part }}$.

The following corollary contains the main idea of Lebeau and Robbiano's strategy: during the first half of a given time interval, we control the lowest frequencies then, during the second half of the time interval, we turn the control to zero to take advantage of the natural dissipation of the problem.

Corollary 3.1. There exists a positive number $M_{\mathrm{LR}}$ which depends only on $M_{\mathrm{obs}, 1}, M_{\mathrm{LR}, 2}, M_{s, \mathrm{cont}}, M_{s, \mathrm{adm}}$ and $\|\mathrm{C}\|$ such that for any $\tau \in(0, T), \mu \in\left(0, \mu_{E_{2}}\right)$ and any $y_{0} \in E$, there exists a control $v \in L^{2}(0, \tau ; U)$ that satisfies

$$
\left\{\begin{array}{l}
\llbracket v \rrbracket_{L^{2}(0, \tau ; U)} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau}+\sqrt{\mu}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\|y(\tau)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau}+\tau\right)}\left(e^{-\mu_{E_{1}} \tau / 2}+e^{-\mu \tau / 2+M_{\mathrm{LR}} \sqrt{\mu}}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{array}\right.
$$

Proof. We apply Proposition 3.2 on the time interval $(0, \tau / 2)$. We get a control $v$ and a solution $y$ which satisfy (28) (with $T$ replaced by $\tau / 2$ ). Then, on the interval $(\tau / 2, \tau)$, we set the control to zero. Thus, we have constructed the following control

$$
\bar{v}(t):= \begin{cases}v(t) & \text { for } t \in(0, \tau / 2) \\ 0 & \text { for } t \in(\tau / 2, \tau)\end{cases}
$$

and the associated solution of the system is still denoted by $y$. Clearly, the $L^{2}$ norm of $\bar{v}$ on $(0, \tau)$ is equal to that of $v$ on $(0, \tau / 2)$.

1. Since the control is 0 on $(\tau / 2, \tau)$, the dissipation properties given in Proposition 3.1 give

$$
\left\|\mathrm{P}_{\mu}(y(\tau))\right\|_{-s, \mathrm{~A}} \leq e^{\|\mathrm{C}\| \frac{\tau}{2}}\left\|\mathrm{P}_{\mu}(y(\tau / 2))\right\|_{-s, \mathrm{~A}}
$$

Then, we use (28) to get

$$
\begin{equation*}
\left\|\mathrm{P}_{\mu}(y(\tau))\right\|_{-s, \mathrm{~A}} \leq e^{2 M_{\mathrm{part}}\left(1+\frac{1}{\tau}\right)} e^{\|\mathrm{C}\| \frac{\tau}{2}} e^{-\mu_{E_{1}} \frac{\tau}{2}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} . \tag{31}
\end{equation*}
$$

2. By the dissipation properties given in Proposition 3.1 we get

$$
\left\|\mathrm{P}_{\mu}^{\perp}(y(\tau))\right\|_{-s, \mathrm{~A}} \leq e^{(\|\mathrm{C}\|-\mu) \frac{\tau}{2}}\|y(\tau / 2)\|_{-s, A}
$$

Then inequality (29) of Proposition 3.2 (with still $T$ replaced by $\tau / 2$ ) leads to

$$
\begin{equation*}
\left\|\mathrm{P}_{\mu}^{\perp}(y(\tau))\right\|_{-s, \mathrm{~A}} \leq e^{2 M_{\mathrm{part}}\left(1+\frac{1}{\tau}+\sqrt{\mu}\right)} e^{(\|\mathrm{C}\|-\mu) \frac{\tau}{2}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \tag{32}
\end{equation*}
$$

We combine (31) and (32) to get the result for a suitable value of $M_{\mathrm{LR}}$ depending on $M_{\text {part }}$.
Observe that Corollary 3.1 is slightly different from the similar result in the classical Lebeau and Robbiano's strategy. Indeed, the modes corresponding to frequencies less than $\mu$ of the final state $y(\tau)$ are not cancelled; they still exist but are controlled by the small term $e^{-\mu_{E_{1}} \frac{\tau}{2}}$. When applying the complete strategy on interval $(0, T)$, one has to make sure that the term $e^{-\mu_{E_{1}} \tau / 2}$ is smaller than $e^{-\mu \tau / 2}$. This constraint is fulfilled in Theorem 3.1 by dealing with frequencies $\mu$ smaller than $\mu_{E_{1}}$.
3.3.3. Conclusion of the proof of the main theorem. We can now apply Lebeau-Robbiano's strategy, with a well chosen finite number of steps, and prove Theorem 3.1.

Proof. Without loss of generality we suppose that $M_{\mathrm{LR}} \geq \ln (2)$. Let $\alpha>0$ such that

$$
\begin{equation*}
\frac{\alpha}{8 M_{\mathrm{LR}}}-3 \frac{\sqrt{\alpha}}{2}-5>0 . \tag{33}
\end{equation*}
$$

Note that $\alpha$ only depends on $M_{\mathrm{LR}}$. Let $\mu_{j}=\frac{\alpha}{T^{2}}\left(2^{j}\right)^{2}, \tau_{j}=\frac{T}{2} \frac{1}{2^{j}}$ and $T_{j}=\sum_{k=1}^{j} \tau_{k}$.
Suppose that $\mu^{*} \leq \mu_{1}$. In this case we apply Corollary 3.1 with $\mu=\mu^{*}$ on the whole time interval $(0, T)$. There exists a control $v \in L^{2}(0, T, U)$ that satisfies

$$
\left\{\begin{array}{l}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{T}+\sqrt{\mu^{*}}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)}\left(e^{-\mu_{E_{1}} T / 2}+e^{-\mu^{*} T / 2+M_{\mathrm{LR}} \sqrt{\mu^{*}}}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{array}\right.
$$

Since $e^{-\mu_{E_{1}} T / 2} \leq e^{-\mu^{*} T / 2}$ we have

$$
\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)} e^{-\mu^{*} T / 2}\left(1+e^{M_{\mathrm{LR}} \sqrt{\mu^{*}}}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}}
$$

and moreover $e^{M_{\mathrm{LR}} \sqrt{\mu^{*}}} \leq e^{M_{\mathrm{LR}} \sqrt{\mu_{1}}} \leq e^{\frac{2}{T} M_{\mathrm{LR}} \sqrt{\alpha}}$. Therefore,

$$
\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}(1+2 \sqrt{\alpha})\left(1+\frac{1}{T}+T\right)} e^{-\mu^{*} T / 2}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
$$

Thus, if $\mu^{*} \leq \mu_{1}$, the proof of Theorem 3.1 is complete since $\alpha$ depends only on $M_{\mathrm{LR}}$, which itself depends only on $M_{\mathrm{obs}, 1}, M_{\mathrm{LR}, 2}, M_{s, \text { cont }}, M_{s, \text { adm }}$ and $\|\mathrm{C}\|$.

Suppose now that $\mu_{1}<\mu^{*}$ and define $j^{*} \in \mathbb{N}^{*}$ such that $\mu_{j^{*}}<\mu^{*} \leq \mu_{j^{*}+1}$.

- During the time interval $\left(0, T_{1}\right)=\left(0, \tau_{1}\right)$, we apply a control $v_{1}$ as given by Corollary 3.1 with $\mu=\mu_{1}$ (which applies here since $\mu_{1}<\mu^{*} \leq \mu_{E_{2}}$ ) and we get

$$
\left\{\begin{array}{l}
\llbracket v_{1} \rrbracket_{L^{2}\left(0, T_{1} ; U\right)} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau_{1}}+\sqrt{\mu_{1}}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\left.\left\|y\left(T_{1}\right)\right\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau_{1}}+\tau_{1}\right.}\right)\left(e^{-\mu_{E_{1}} \tau_{1} / 2}+e^{-\mu_{1} \tau_{1} / 2+M_{\mathrm{LR}} \sqrt{\mu_{1}}}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{array}\right.
$$

- For any index $j \leq j^{*}$, we continue this procedure by applying Corollary 3.1 on time interval $\left(T_{j-1}, T_{j}\right)=\left(T_{j-1}, T_{j-1}+\tau_{j}\right)$ with $\mu=\mu_{j}$. We get a control $v_{j}$ that satisfies

$$
\left\{\begin{align*}
\llbracket v_{j} \rrbracket_{L^{2}\left(T_{j-1}, T_{j} ; U\right)} & \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau_{j}}+\sqrt{\mu_{j}}\right)}\left\|y\left(T_{j-1}\right)\right\|_{-s, \mathrm{~A}}  \tag{34}\\
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} & \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{\tau_{j}}+\tau_{j}\right)} \times \\
& \left(e^{-\mu_{E_{1}} \tau_{j} / 2}+e^{-\mu_{j} \tau_{j} / 2+M_{\mathrm{LR} \sqrt{\mu_{j}}}}\right)\left\|y\left(T_{j-1}\right)\right\|_{-s, \mathrm{~A}}
\end{align*}\right.
$$

Let us focus on the estimation of the term $y\left(T_{j}\right)$. First, we use that $e^{-\mu_{E_{1}} \tau_{j} / 2} \leq$ $e^{-\mu_{j} \tau_{j} / 2}$ and the relations

$$
\left\{\begin{align*}
\mu_{j} \tau_{j} & =\frac{\alpha}{2 T} 2^{j}=\frac{\alpha}{4} \frac{1}{\tau_{j}}  \tag{35}\\
\sqrt{\mu_{j}} & =\frac{\sqrt{\alpha}}{2 \tau_{j}}
\end{align*}\right.
$$

It follows

$$
e^{-\mu_{j} \tau_{j} / 2+M_{\mathrm{LR}} \sqrt{\mu_{j}}} \leq e^{\frac{1}{\tau_{j}}\left(-\frac{\alpha}{8}+M_{\mathrm{LR}} \frac{\sqrt{\alpha}}{2}\right)}
$$

Hence

$$
\begin{equation*}
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} \leq 2 e^{M_{\mathrm{LR}}} e^{M_{\mathrm{LR}}\left(\frac{1}{\tau_{j}}+\tau_{j}\right)+\frac{1}{\tau_{j}}\left(-\frac{\alpha}{8}+M_{\mathrm{LR}} \frac{\sqrt{\alpha}}{2}\right)}\left\|y\left(T_{j-1}\right)\right\|_{-s, \mathrm{~A}} \tag{36}
\end{equation*}
$$

since $\ln (2) \leq M_{\mathrm{LR}}$, and $1 \leq \frac{1}{\tau_{j}}+\tau_{j}$ leads to

$$
2 e^{M_{\mathrm{LR}}} e^{M_{\mathrm{LR}}\left(\frac{1}{\tau_{j}}+\tau_{j}\right)} \leq e^{2 M_{\mathrm{LR}}} e^{M_{\mathrm{LR}}\left(\frac{1}{\tau_{j}}+\tau_{j}\right)} \leq e^{3 M_{\mathrm{LR}}\left(\frac{1}{\tau_{j}}+\tau_{j}\right)}
$$

which, associated with (36), finally yields

$$
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} \leq e^{3 M_{\mathrm{LR}}\left(\frac{1}{\tau_{j}}+\tau_{j}\right)+\frac{1}{\tau_{j}}\left(-\frac{\alpha}{8}+M_{\mathrm{LR}} \frac{\sqrt{\alpha}}{2}\right)}\left\|y\left(T_{j-1}\right)\right\|_{-s, \mathrm{~A}}
$$

Now we set $\beta:=\frac{\alpha}{8 M_{\mathrm{LR}}}-\frac{\sqrt{\alpha}}{2}-3$, in such a way that the previous estimate reads

$$
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} \leq e^{3 M_{\mathrm{LR}} \tau_{j}} e^{-M_{\mathrm{LR}} \frac{\beta}{\tau_{j}}}\left\|y\left(T_{j-1}\right)\right\|_{-s, \mathrm{~A}}
$$

According to (33), we have $\beta>0$ since $\frac{\alpha}{8 M_{\mathrm{LR}}}-\left(\frac{\sqrt{\alpha}}{2}+3\right)>\frac{\alpha}{8 M_{\mathrm{LR}}}-\left(\frac{3 \sqrt{\alpha}}{2}+5\right)>0$.
Thus, we can estimate $y\left(T_{j}\right)$ by using the previous upper-bound recursively,

$$
\begin{aligned}
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} & \leq\left(e^{3 M_{\mathrm{LR}} \sum_{k=1}^{j} \tau_{k}}\right)\left(e^{-\beta M_{\mathrm{LR}} \sum_{k=1}^{j} \frac{1}{\tau_{k}}}\right)\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
& \leq e^{3 M_{\mathrm{LR}} T_{j}} e^{-\frac{2}{T} \beta M_{\mathrm{LR}} \sum_{k=1}^{j} 2^{k}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
& \leq e^{3 M_{\mathrm{LR}} T_{j}} e^{-\frac{2 \beta M_{\mathrm{LR}} 2^{j}}{T}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} .
\end{aligned}
$$

Therefore,

$$
\begin{equation*}
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} \leq e^{3 M_{\mathrm{LR}} T_{j}} e^{-\frac{\beta M_{\mathrm{LR}}}{\tau_{j}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \tag{37}
\end{equation*}
$$

and with (35),

$$
\left\|y\left(T_{j}\right)\right\|_{-s, \mathrm{~A}} \leq e^{3 M_{\mathrm{LR}} T_{j}} e^{-2 \frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu_{j}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
$$

When $j=j^{*}$, we end up with (recall that $\mu^{*}<\mu_{j^{*}+1}=4 \mu_{j^{*}}$ )

$$
\begin{equation*}
\left\|y\left(T_{j^{*}}\right)\right\|_{-s, \mathrm{~A}} \leq e^{3 M_{\mathrm{LR}} T_{j^{*}}} e^{-\frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu^{*}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \tag{38}
\end{equation*}
$$

Let us now estimate the control $v$. Taking back (34) for $j \geq 2$, combined with (37) and (35),

$$
\llbracket v_{j} \rrbracket_{L^{2}\left(T_{j-1}, T_{j} ; U\right)} \leq e^{M_{\mathrm{LR}}\left[1+\frac{1}{\tau_{j}}\left(1+\frac{\sqrt{\alpha}}{2}-\frac{\beta}{2}\right)+3 T_{j-1}\right]}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
$$

Note that, (33) gives that $\bar{\beta}:=\beta / 2-1-\frac{\sqrt{\alpha}}{2}=\frac{1}{2}(\beta-2-\sqrt{\alpha})=\frac{1}{2}\left(\frac{\alpha}{8 M_{\mathrm{LR}}}-3 \frac{\sqrt{\alpha}}{2}-5\right)>0$.
This implies that

$$
\begin{aligned}
\llbracket v \rrbracket_{L^{2}\left(T_{1}, T_{j^{*}} ; U\right)}^{2} & =\sum_{j=2}^{j^{*}} \llbracket v_{j} \rrbracket_{L^{2}\left(T_{j-1}, T_{j} ; U\right)}^{2} \\
& \leq e^{2 M_{\mathrm{LR}}(1+3 T)} \sum_{j \geq 0} e^{-\frac{M_{\mathrm{LR}} \bar{\beta}}{T} 2^{j+2}}\left\|y_{0}\right\|_{-s, \mathrm{~A}}^{2} \\
& \leq\left(1+\frac{T}{M_{\mathrm{LR}} \bar{\beta}}\right) e^{2 M_{\mathrm{LR}}(1+3 T)}\left\|y_{0}\right\|_{-s, \mathrm{~A}}^{2}
\end{aligned}
$$

and the last inequality came from $\frac{1}{1-e^{-x}} \leq 1+\frac{1}{x}$, for $x>0$.
Using that

$$
\llbracket v_{1} \rrbracket_{L^{2}\left(0, T_{1} ; U\right)}^{2} \leq e^{2 M_{\mathrm{LR}}\left(1+\frac{1}{\tau_{1}}+\tau_{1}+\sqrt{\mu_{1}}\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}}^{2},
$$

we have built a control $v$ on $\left(0, T_{j^{*}}\right)$ that satisfies

$$
\begin{equation*}
\llbracket v \rrbracket_{L^{2}\left(0, T_{j^{*} *} ; U\right)} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}}, \tag{39}
\end{equation*}
$$

with $M_{\text {obs }}$ depending on $M_{\text {LR }}$ but not on $T$.
Notice that $T_{j^{*}}=\sum_{j=1}^{j^{*}} \tau_{j}=\frac{T}{2} \sum_{j=1}^{j^{*}} \frac{1}{2^{j}} \leq \frac{T}{2}$. The last step of the proof consists in applying Corollary 3.1 on the interval $\left(T_{j^{*}}, T\right)$ to recover the estimates of Theorem 3.1.

Corollary 3.1 with $\mu=\mu^{*}$ gives now a control $v \in L^{2}\left(T_{j^{*}}, T\right)$ such that

$$
\left\{\begin{align*}
& \llbracket v \|_{L^{2}\left(T_{j^{*}}, T ; U\right)} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{T-T_{j^{*}}}+T-T_{j^{*}}+\sqrt{\left.\mu^{*}\right)}\right.}\left\|y\left(T_{j^{*}}\right)\right\|_{-s, \mathrm{~A}}  \tag{40}\\
&\left.\|y(T)\|_{-s, \mathrm{~A}} \leq e^{M_{\mathrm{LR}}\left(1+\frac{1}{T-T_{j^{*}}}+T-T_{j^{*}}\right.}\right) \times \\
&\left(e^{-\mu_{E_{1}}\left(T-T_{j^{*}}\right) / 2}+e^{-\mu^{*}\left(T-T_{j^{*}}\right) / 2+M_{\mathrm{LR}} \sqrt{\mu^{*}}}\right)\left\|y\left(T_{j^{*}}\right)\right\|_{-s, \mathrm{~A}}
\end{align*}\right.
$$

First, with (38), we have

$$
\llbracket v \rrbracket_{L^{2}\left(T_{j^{*}}, T ; U\right)} \leq e^{M_{\mathrm{LR}}\left(1+\frac{2}{T}+T-T_{j^{*}}+\sqrt{\mu^{*}}\right)} e^{3 M_{\mathrm{LR}} T_{j^{*}}} e^{-\frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu^{*}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}},
$$

and by (33), we can check that

$$
\begin{equation*}
\beta>\sqrt{\alpha} . \tag{41}
\end{equation*}
$$

Hence

$$
\llbracket v \rrbracket_{L^{2}\left(T_{j^{*}}, T ; U\right)} \leq e^{2 M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
$$

Finally this estimation and (39) give the first inequality of Theorem 3.1.
According to (40) and (38),

$$
\begin{aligned}
\|y(T)\|_{-s, \mathrm{~A}} \leq & \left.e^{M_{\mathrm{LR}}\left(1+\frac{1}{T-T_{j^{*}}}+T-T_{j^{*}}\right.}\right)\left(e^{-\mu_{E_{1}}\left(T-T_{j^{*}}\right) / 2}+e^{-\mu^{*}\left(T-T_{j^{*}}\right) / 2+M_{\mathrm{LR}} \sqrt{\mu^{*}}}\right) \times \\
& e^{3 M_{\mathrm{LR}} T_{j^{*}}} e^{-\frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu^{*}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
& \leq e^{2 M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)}\left(e^{-\mu_{E_{1}} T / 4}+e^{-\mu^{*} T / 4+M_{\mathrm{LR}} \sqrt{\mu^{*}}}\right) e^{-\frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu^{*}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}} \\
\leq & 2 e^{2 M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)} e^{-\mu^{*} T / 4} e^{M_{\mathrm{LR}} \sqrt{\mu^{*}}} e^{-\frac{\beta M_{\mathrm{LR}}}{\sqrt{\alpha}} \sqrt{\mu^{*}}}\left\|y_{0}\right\|_{-s, \mathrm{~A}}
\end{aligned}
$$

Using (41), this reduces to

$$
\|y(T)\|_{-s, \mathrm{~A}} \leq 2 e^{2 M_{\mathrm{LR}}\left(1+\frac{1}{T}+T\right)} e^{-\mu^{*} T / 4}\left\|y_{0}\right\|_{-s, \mathrm{~A}},
$$

and this concludes the proof.

## 4. Applications.

4.1. Dirichlet boundary null control of a semi-discrete cascade system on a rectangle. In this section, we will apply the general framework introduced above to prove the $\phi(h)$ null controllability properties for semi-discrete versions of the boundary control problem of the coupled system (6) in the case of a cascade form (that is for particular B and C given below). The penalization term $\phi(h)$ will be exponentially small in $h$, just like in similar results obtained in the literature and quoted in the introduction section 1.1.

The results are valid for the finite difference discretization in space of the system in any dimension. However, for the simplicity of the presentation we will only state and prove our theorem in dimension $d=2$. Observe that other usual techniques (based on Carleman estimates, or on moments methods as in $[1,9,8,7]$ ) do not directly apply in this setting since we are considering boundary controls for multi-dimensional coupled systems with less controls than components in the system.
4.1.1. Presentation of the semi-discrete system. Let us rewrite the assumptions in the precise setting we consider here. Let $a_{1}, a_{2} \in \mathbb{R}_{+}^{*}, \Omega_{1}=\left(0, a_{1}\right)$, $\Omega_{2}=\left(0, a_{2}\right)$ and $\Omega:=\Omega_{1} \times \Omega_{2}$. Let $\omega_{2}$ be a non empty open subset of $\Omega_{2}$. Without loss of generality, we assume that

$$
\begin{equation*}
\overline{\omega_{2}} \subset \Omega_{2} \tag{42}
\end{equation*}
$$

Let $\gamma_{1} \in \mathcal{C}^{1}\left(\bar{\Omega}_{1}\right)$ and $\gamma_{2} \in \mathcal{C}^{1}\left(\bar{\Omega}_{2}\right)$ be two functions which satisfy:

$$
\gamma_{1, \min }:=\inf _{\bar{\Omega}_{1}} \gamma_{1}>0 \text { and } \gamma_{2, \min }:=\inf _{\bar{\Omega}_{2}} \gamma_{2}>0
$$

Let $\mathcal{A}_{1}=-\partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} \bullet\right)$ (resp. $\mathcal{A}_{2}=-\partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} \bullet\right)$ ) be the self-adjoint unbounded operator in $L^{2}\left(\Omega_{1}\right)$ (resp. $L^{2}\left(\Omega_{2}\right)$ ) whose domain is $D\left(\mathcal{A}_{1}\right)=H_{0}^{1}\left(\Omega_{1}\right) \cap$ $H^{2}\left(\Omega_{1}\right)$ (resp. $D\left(\mathcal{A}_{2}\right)=H_{0}^{1}\left(\Omega_{2}\right) \cap H^{2}\left(\Omega_{2}\right)$ ). Let $n \geq 1$, we define a self-adjoint unbounded operator in $L^{2}\left(\Omega_{1}\right) \otimes L^{2}\left(\Omega_{2}\right) \otimes \mathbb{R}^{n}$ on the domain $D(\mathcal{A})=D\left(\mathcal{A}_{1}\right) \otimes D\left(\mathcal{A}_{2}\right) \otimes \mathbb{R}^{n}$ as follows

$$
\mathcal{A}:=\mathcal{A}_{1} \otimes I \otimes \mathrm{I}+I \otimes \mathcal{A}_{2} \otimes \mathrm{I}
$$

Note that $\mathcal{A}$ can be seen as a self-adjoint operator in $\left(L^{2}(\Omega)\right)^{n}$ with domain $\left(H_{0}^{1}(\Omega) \cap\right.$ $\left.H^{2}(\Omega)\right)^{n}$. The control matrix $\mathrm{B} \in M_{n, 1}(\mathbb{R})$ and the coupling matrix $\mathrm{C} \in M_{n}(\mathbb{R})$ are chosen in the so-called cascade form, given as follows

$$
\mathrm{B}=\left(\begin{array}{l}
1 \\
0 \\
\mid \\
\mid \\
0
\end{array}\right), \quad \mathrm{C}=\left(\begin{array}{llll}
0 & & & 0 \\
1 & & & \\
0 & & & \mid \\
\mid & & & \\
0 & & \mid & 1 \\
0 & 0 & 1 & 0
\end{array}\right) .
$$

Let us now define the finite difference approximation that we consider for this system. For simplicity we present the results on uniform grids even though the same result will hold for regular enough families of meshes (see for instance the discussion in $[1,8]$ ). Let $N_{1}$ and $N_{2}$ be two integers greater or equal to 2 . For $i \in\{1,2\}$, we set $h_{i}:=\frac{a_{i}}{N_{i}+1}$ and consider $\left(x_{1, j}\right)_{j=0}^{N_{1}+1}$ and $\left(x_{2, j}\right)_{j=0}^{N_{2}+1}$ a discretization of $\Omega_{1}$ and $\Omega_{2}$ respectively : $x_{1, j}:=j h_{1}$ for $j \in\left\{0, \ldots, N_{1}+1\right\}, x_{2, j}:=j h_{2}$ for $j \in\left\{0, \ldots, N_{2}+1\right\}$. We define also, for $j \in\left\{0, \ldots, N_{1}\right\}, x_{1, j+1 / 2}=x_{1, j}+\frac{h_{1}}{2}$ and for $j \in\left\{0, \ldots, N_{2}\right\}$, $x_{2, j+1 / 2}=x_{2, j}+\frac{h_{2}}{2}$.

The geometry of the grid and of the control domain is summarized in Figure 2: the interior grid points are shown in black circles •, the homogeneous Dirichlet boundary points are shown in white circles $O$ and the boundary points in the control region are shown in gray squares $\square$.


Figure 2: Grid geometry

Associated with those grids, we can introduce the discrete functional spaces $E_{1}=$ $\mathbb{R}^{N_{1}}$ (resp. $E_{2}=\mathbb{R}^{N_{2}}$ ) equipped with the Euclidean inner product inherited from the $L^{2}$ inner product in $\Omega_{1}$ (resp. in $\Omega_{2}$ ) and defined by

$$
\langle y, z\rangle_{0, i}:=\sum_{j=1}^{N_{i}} h_{i} y_{j} z_{j}, \quad \forall y, z \in E_{i}, \forall i \in\{1,2\}
$$

For any function $f: \Omega_{i} \rightarrow \mathbb{R}$, we will use the same letter $f$ to denote the sampling of $f$ on the grid $\left(f\left(x_{i, j}\right)\right)_{j} \in E_{i}$ or the multiplication by $f$ operator in $E_{i}$ which means that, for any $y \in E_{i}$, the vector $f y \in E_{i}$ is defined by

$$
\begin{equation*}
(f y)_{j}=f\left(x_{i, j}\right) y_{j}, \quad \forall j \in\left\{1, \ldots, N_{i}\right\} \tag{43}
\end{equation*}
$$

Let $\mathrm{A}_{i}$ be the self-adjoint operator in $E_{i}$ (that can be seen as a $N_{i} \times N_{i}$ matrix) corresponding to the discretization of the scalar 1D operator $\mathcal{A}_{i}$ by the finite difference method, which is defined for any $y \in E_{i}$ by

$$
\begin{equation*}
\left(\mathrm{A}_{i} y\right)_{j}:=-\frac{1}{h_{i}}\left(\gamma_{i, j+1 / 2} \frac{y_{j+1}-y_{j}}{h_{i}}-\gamma_{i, j-1 / 2} \frac{y_{j}-y_{j-1}}{h_{i}}\right), \forall j \in \llbracket 1, N_{i} \rrbracket, \tag{44}
\end{equation*}
$$

with the usual convention that $y_{0}=y_{N_{i}+1}=0$. Here, we have used the notation $\gamma_{i, j+1 / 2}:=\gamma_{i}\left(x_{i, j+1 / 2}\right)$, for $i=1,2$ and $j \in\left\{0, \ldots, N_{i}-1\right\}$, for the sampling of the diffusion coefficient $\gamma_{i}$ on the dual grid of the mesh of $\Omega_{i}$. At some point we will also need to consider the discrete Laplace operators $\Delta_{i}$ defined by

$$
\left(-\Delta_{i} y\right)_{j}:=-\frac{y_{j+1}-2 y_{j}+y_{j-1}}{h_{i}^{2}}, \forall j \in \llbracket 1, N_{i} \rrbracket, \forall y \in E_{i}
$$

The discretization of the vector-valued 2 D operator $\mathcal{A}$ we will be interested in is thus given by

$$
\mathrm{A}=\mathrm{A}_{1} \otimes \mathrm{I} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2} \otimes \mathrm{I}
$$

and, if we take into account the coupling terms, the complete discretization of the operator appearing in (6) is given by

$$
\mathrm{L}:=\mathrm{A}+\mathrm{I} \otimes \mathrm{I} \otimes \mathrm{C}
$$

We introduce now the discrete control spaces and operators associated with our boundary control problem. We set $U_{2}=E_{2}$, and to avoid confusions, we will use the notation $[\bullet, \bullet]_{0,2}$ instead of $\langle\bullet, \bullet\rangle_{0,2}$ when dealing with objects in $U_{2}$.

Using the convention (43), we define the discrete control operator in $\Omega_{2}$ by

$$
\mathrm{B}_{2}=1_{\omega_{2}}
$$

This means that for any $y \in E_{2}$,

$$
\left(\mathrm{B}_{2} y\right)_{j}=\left\{\begin{array}{r}
y_{j}, \text { if } x_{2, j} \in \omega_{2} \\
0, \text { if } x_{2, j} \notin \omega_{2}
\end{array}\right.
$$

Observe that this operator is self-adjoint in $E_{2}$, that is $\mathrm{B}_{2}^{*}=\mathrm{B}_{2}$. Since the domain $\Omega_{1}$ is a 1 D interval, the corresponding boundary control $v$ is in fact a scalar control that lives in the space $U_{1}=\mathbb{R}$. However, the discretization $\mathrm{A}_{1}$ of the operator $\mathcal{A}_{1}$ is built upon the assumption that we are dealing with homogeneous Dirichlet boundary condition. Therefore, in order to include non-homogeneous Dirichlet boundary condition on the left end-point of $\Omega_{1}$ in the discretization, we need to add a source term in the discretization which is given by

$$
\mathrm{B}_{1} v:=\left(\begin{array}{c}
-\frac{\gamma_{1,1 / 2}}{h_{1}^{2}} v \\
0 \\
\mid \\
0
\end{array}\right) \in E_{1}
$$

For the analysis, it will be sightly more convenient to work with this operator in the following form

$$
\begin{equation*}
\mathrm{B}_{1} v=\gamma_{1,1 / 2}\left(\Delta_{1} r\right) v, \quad \forall v \in \mathbb{R} \tag{45}
\end{equation*}
$$

where $r$ is (the sampling of) the affine map defined by $x \in \Omega_{1} \mapsto 1-x / a_{1}$. This is the discrete counterpart of the boundary control operator as analysed in [15, Chapter 2].

A simple computation shows that the adjoint of $B_{1}$ is given by

$$
\mathrm{B}_{1}^{*} q=\gamma_{1,1 / 2} \partial_{l} q
$$

where $\partial_{l} q$ is the discrete normal derivative of $q \in E_{1}$ on the left boundary of the domain defined by

$$
\partial_{l} q=\frac{0-q_{1}}{h_{1}}=-\frac{q_{1}}{h_{1}}
$$

Note that this formula takes into account implicitely the homogeneous Dirichlet boundary condition for $q$.

We can now precisely write the semi-discrete control problem that we consider

$$
\left\{\begin{array}{l}
y^{\prime}(t)+\mathrm{L} y(t)=\mathrm{B}_{1} \otimes \mathrm{~B}_{2} \otimes \mathrm{~B} v(t)  \tag{46}\\
y(0)=y_{0} \in E
\end{array}\right.
$$

where $y(t) \in E=E_{1} \otimes E_{2} \otimes \mathbb{R}^{n}$ and $v(t) \in U=\mathbb{R} \otimes U_{2} \otimes \mathbb{R}$. Note that, in this particular case, the control space $U$ can be in fact identified with $E_{2}$.

The main theorem of this section is the following. The crucial point is that all the constants appearing in the estimates (47) do not depend on the discretization parameter $h$. In short, we prove that we can drive the semi-discrete system (46) to a target which is exponentially small with respect to $h$ with controls that are uniformly bounded. Up to a subsequence, this results imply the weak convergence of the semidiscrete controls towards a control of the continuous problem which leads the solution to zero, as soon as the discrete initial data converges towards the suitable initial data.

Theorem 4.1. There exist $C>0, \widetilde{C}>0$ and $h_{0}>0$, depending only on $\gamma_{1}, \gamma_{2}$ and $\omega_{2}$ such that, for any $T>0$, any mesh such that $h<h_{0}$, and any $y_{0} \in E$ there exists a control $v \in L^{2}(0, T ; U)$ satisfying

$$
\left\{\begin{array}{l}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq e^{\widetilde{C}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-1, \mathrm{~A}}  \tag{47}\\
\|y(T)\|_{-1, \mathrm{~A}} \leq e^{\widetilde{C}\left(1+\frac{1}{T}+T\right)} e^{-C / h^{2}}\left\|y_{0}\right\|_{-1, \mathrm{~A}}
\end{array}\right.
$$

where $y$ is the corresponding solution to the semi-discrete problem (46) with control $v$ and $h=\max \left(h_{1}, h_{2}\right)$ is the space discretization parameter.

Once such a theorem is proved, even with a non explicit/constructive proof, one can produce an optimization algorithm, based on the penalized HUM approach, that is able to compute a control $v$ satisfying (47). We can even relax the requirements by replacing the exponential factor $e^{-C / h^{2}}$ by any more convenient $\phi(h)$, such as $\phi(h)=h^{p}$ for some large enough integer $p$. Those questions are discussed in details for instance in [6] where some numerical illustrations are given.
4.1.2. Additional notations and properties. In order to simplify the presentation of the following proofs we need to introduce a few more notations. For any $i \in\{1,2\}$, we define $\gamma_{i}^{ \pm}$to be the translated sampling of the diffusion coefficient $\gamma_{i}$ defined by

$$
\left(\gamma_{i}^{ \pm}\right)_{j}=\gamma_{i, j \pm 1 / 2}, \quad \forall j \in\left\{1, \ldots, N_{i}\right\} .
$$

We also introduce the forward and backward difference operators $\nabla_{i}^{ \pm}$defined, for any $y \in E_{i}$, by

$$
\begin{aligned}
& \left(\nabla_{i}^{+} y\right)_{j}=\frac{y_{j+1}-y_{j}}{h_{i}}, \quad \forall j \in\left\{1, \ldots, N_{i}\right\} \\
& \left(\nabla_{i}^{-} y\right)_{j}=\frac{y_{j}-y_{j-1}}{h_{i}}, \quad \forall j \in\left\{1, \ldots, N_{i}\right\}
\end{aligned}
$$

For any $i \in\{1,2\}$, we define $\left(e_{i, j}\right)_{1 \leq j \leq N_{i}}$ to be the canonical basis of $E_{i}$ (each element corresponds to a point in the 1D grid of $\Omega_{i}$ ). At some point in the forthcoming analysis, we shall need to work with compactly supported discrete functions in order to justify some discrete integration by parts. To this end, we introduce the subspaces $E_{i}^{00} \subset E_{i}^{0} \subset E_{i}$ defined by

$$
\begin{aligned}
& E_{i}^{0}=\operatorname{span}\left(e_{i, j}, \quad j=2, \ldots, N_{i}-1\right) \\
& E_{i}^{00}=\operatorname{span}\left(e_{i, j}, \quad j=3, \ldots, N_{i}-2\right)
\end{aligned}
$$

With those notations, we observe that the gradient operators defined above satisfy the duality property

$$
\begin{equation*}
\left\langle\nabla_{i}^{+} y, z\right\rangle_{0, i}=-\left\langle y, \nabla_{i}^{-} z\right\rangle_{0, i}, \quad \forall y, z \in E_{i}^{0} \tag{48}
\end{equation*}
$$

that we obtain by a summation by parts.
Lemma 4.1. There exists a $C>0$ depending only on $\inf \gamma_{i}$, $\sup \gamma_{i},\left\|\gamma_{i}^{\prime}\right\|_{L^{\infty}}, a_{i}$ such that: for any $i \in\{1,2\}$, and any $y \in E_{i}$, we have

$$
\begin{gather*}
\|y\|_{0, i} \leq C\left\|\nabla_{i}^{ \pm} y\right\|_{0, i}  \tag{49}\\
\left\|\nabla_{i}^{+} y\right\|_{0, i}^{2}+\left\|\nabla_{i}^{-} y\right\|_{0, i}^{2} \leq C\left\langle\mathrm{~A}_{i} y, y\right\rangle_{0, i}  \tag{50}\\
\left\|\Delta_{i} y\right\|_{0, i} \leq C\left\|\mathrm{~A}_{i} y\right\|_{0, i} \tag{51}
\end{gather*}
$$

REMARK 4.1. By combining the above properties, we can obtain the following estimate

$$
\begin{equation*}
\|y\|_{0, i} \leq C\left\|\mathrm{~A}_{i} y\right\|_{0, i}, \quad \forall y \in E_{i} \tag{52}
\end{equation*}
$$

where $C$ depends only on $\gamma_{i}$ and $a_{i}$. This implies in particular that the Poincaré inequality (23) holds with a constant $M_{\mathrm{P}, 1}$ uniform with respect to the discretization parameters.

Proof. The first inequality is very classical. We recall the sketch of proof: for $i \in\{1,2\}$, we first write

$$
y_{j}=y_{j}-y_{0}=\sum_{k=1}^{j}\left(y_{k}-y_{k-1}\right)=\sum_{k=1}^{j} h_{i}\left(\nabla_{i}^{-} y\right)_{k}
$$

and by the Cauchy-Schwarz inequality, we obtain

$$
\left|y_{j}\right|^{2} \leq a_{i}\left\|\nabla_{i}^{-} y\right\|_{0, i}^{2}, \quad \forall j \in\left\{0, \ldots, N_{i}\right\}
$$

The claim follows by summing over $j$. The proof with the operator $\nabla_{i}^{+}$, is done in the same way but starting from the equality

$$
y_{j}=y_{j}-y_{N_{i}+1}=-\sum_{k=j}^{N_{i}}\left(y_{k+1}-y_{k}\right)=\sum_{k=j}^{N_{i}} h_{i}\left(\nabla_{i}^{+} y\right)_{k}
$$

For the third estimate, a simple calculation shows that

$$
\left\langle\mathrm{A}_{i} y, y\right\rangle_{0, i}=\sum_{j=0}^{N_{i}} h_{i} \gamma_{i, j+1 / 2}\left|\frac{y_{j+1}-y_{j}}{h_{i}}\right|^{2}, \forall y \in E_{i}
$$

and we deduce (50) with $C$ depending only on $\min _{\Omega_{i}}\left(\gamma_{i}\right)$.
Finally, a straightforward algebraic computation gives

$$
\mathrm{A}_{i}=\frac{\gamma_{i}^{+}+\gamma_{i}^{-}}{2}\left(-\Delta_{i}\right)-\frac{\gamma_{i}^{+}-\gamma_{i}^{-}}{h_{i}} \frac{\nabla_{i}^{+}+\nabla_{i}^{-}}{2}
$$

which implies

$$
\left\|-\Delta_{i} y\right\|_{0, i} \leq C\left\|\mathrm{~A}_{i} y\right\|_{0, i}+C\left\|\gamma_{i}^{\prime}\right\|_{L^{\infty}}\left(\left\|\nabla_{i}^{+} y\right\|_{0, i}+\left\|\nabla_{i}^{-} y\right\|_{0, i}\right)
$$

where $C$ only depends on $\min _{\Omega_{i}}\left(\gamma_{i}\right)$. By using (50), the Cauchy-Schwarz inequality and the Young inequality, we finally deduce that

$$
\left\|-\Delta_{i} y\right\|_{0, i} \leq C\left\|\mathrm{~A}_{i} y\right\|_{0, i}+C\|y\|_{0, i}, \quad \forall y \in E_{i}
$$

where $C$ only depends on the function $\gamma_{i}$ and we conclude by using (49) and (50).
4.1.3. Proof of the semi-discrete controllability result. The proof of Theorem 4.1 consists essentially in applying the analysis above to a suitable setting. To achieve the results, it is just needed to check that all the assumptions of Theorem 3.1 are satisfied with constants that do not depend on the discretization parameter.

- In a first step, we will check that the discrete diffusion operator $L$ and the discrete control operator $\mathrm{B}=\mathrm{B}_{1} \otimes \mathrm{~B}_{2} \otimes \mathrm{~B}$ are compatible in the sense that inequalities (14) and (15) are satisfied with $s=1$, and $\mathrm{D}=\mathrm{A}$, uniformly with respect to the mesh size.
- In a second step, using previous results in the literature, we shall prove that the discrete Lebeau-Robbiano spectral inequality (Assumption 3.2) and the semi-discrete controllability of the 1D system (Assumption 3.1) hold.
Let us start by proving the following result.
Lemma 4.2. There exist two constants $M_{\text {cont }}$ and $M_{a d m}$ independent of $h_{1}$ and $h_{2}$ that satisfy:

$$
\begin{gather*}
\sup _{t \in[0, T]}\left\|e^{-t \mathrm{~L}^{*}} \psi\right\|_{1, \mathrm{~A}} \leq M_{\text {cont }}\|\psi\|_{1, \mathrm{~A}}, \quad \forall \psi \in E,  \tag{53}\\
\left(\int_{0}^{T} \llbracket \mathrm{~B}^{*} e^{-\tau \mathrm{L}^{*}} \psi \rrbracket_{0}^{2} d \tau\right)^{\frac{1}{2}} \leq M_{a d m}\|\psi\|_{1, \mathrm{~A}}, \quad \forall \psi \in E,
\end{gather*}
$$

Proof. The proof of (53) is exactly the same as the one of the dissipation estimates of Proposition 3.1, for $\mu=0$ and $s=-1$, except that we deal with the adjoint matrix C*.

The more intricate part is now to prove (54). This will be a combination of discrete trace estimates and of discrete elliptic regularity properties for the operator A. In that proof the fact that we consider $s=1$ in the definition of the norms is crucial, this is the discrete counter-part of the fact that, in the continuous setting, we need to take $H_{0}^{1}$ initial data to ensure that the normal derivative of the solution of the backward heat equation belongs to $L^{2}$.

1. First, we prove a 1D trace inequality. More precisely, we show that there exists $C>0$ independent of $h_{1}$ such that:

$$
\begin{equation*}
\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0,1}^{2} \leq C\left\|\mathrm{~A}_{1} \otimes \mathrm{I} q\right\|_{0,1}^{2}, \forall q \in E_{1} \otimes \mathbb{R}^{n} \tag{55}
\end{equation*}
$$

To this end, we first use (45) and the fact that $\Delta_{1}$ is self-adjoint to get

$$
\mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*}=\gamma_{1,1 / 2}\left(r \otimes \mathrm{~B}^{*}\right)\left(\Delta_{1} \otimes \mathbf{I}\right)
$$

It follows that, for any $q \in E_{1} \otimes \mathbb{R}^{n}$,

$$
\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0,1} \leq \gamma_{1,1 / 2}\|r\|_{0,1}\left\|\mathrm{~B}^{*}\right\|\left\|\Delta_{1} \otimes \mathrm{I} q\right\|_{0,1} \leq C\left\|\Delta_{1} \otimes \mathrm{I} q\right\|_{0,1}
$$

where $C$ depends only on $\gamma_{1}$ and $\Omega_{1}$.
From (51) (that we apply on each component of the vector-valued unknown $q$ ), we deduce that

$$
\left\|\Delta_{1} \otimes \mathbf{I} q\right\|_{0,1} \leq C\left\|\mathrm{~A}_{1} \otimes \mathbf{I} q\right\|_{0,1}
$$

which proves (55).
2. In the second step, we prove that the inequality (55) can be extended to $\mathbf{2 D}$ discrete unknowns. More precisely, we show that there exists a constant $C>0$ independent of $h_{1}$ and $h_{2}$ such that

$$
\begin{equation*}
\llbracket \mathrm{B}^{*} q \rrbracket_{0}^{2}=\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0}^{2} \leq C\left\|\mathrm{~A}_{1} \otimes \mathrm{I} \otimes \mathrm{I} q\right\|_{0}^{2}, \forall q \in E . \tag{56}
\end{equation*}
$$

We define $f_{j}=e_{2, j} / \sqrt{h_{2}}$, in such a way that $\left(f_{2, j}\right)_{1 \leq j \leq N_{2}}$ is an orthonormal basis of $E_{2}$. Note that the particular structure of $\bar{B}_{2}$ implies that $\mathrm{B}_{2} f_{2, j}=$ $\alpha_{j} f_{2, j}$ with $\alpha_{j} \in\{0,1\}\left(\alpha_{j}=1\right.$ if and only if the mesh point $x_{2, j}$ lies in $\left.\omega_{2}\right)$. We can decompose any $q \in E$ into the unique form

$$
q=\sum_{j=1}^{N_{2}} q_{j} \otimes f_{2, j}
$$

with $q_{j} \in E_{1} \otimes \mathbb{R}^{n}$ for any $j$. We have, by orthogonality of $\left(f_{2, j}\right)_{j}$,

$$
\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0}^{2}=\llbracket \sum_{j=1}^{N_{2}} \alpha_{j}\left[\left(\mathrm{~B}_{1}^{*} \otimes \mathrm{~B}^{*}\right) q_{j}\right] \otimes f_{2, j} \rrbracket_{0}^{2}=\sum_{r=1}^{N_{2}} \llbracket \mathrm{~B}_{1}^{*} \otimes \mathrm{~B}^{*} q_{j} \rrbracket_{0,1}^{2} \alpha_{j}^{2} .
$$

Using (55) for each $j$, and the fact that $\alpha_{j}^{2} \leq 1$, we get

$$
\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0}^{2} \leq C \sum_{j=1}^{N_{2}}\left\|\mathrm{~A}_{1} \otimes \mathbf{I} q_{j}\right\|_{0,1}^{2}=C \sum_{j=1}^{N_{2}}\left\|\left[\left(\mathrm{~A}_{1} \otimes \mathrm{I}\right) \otimes \mathrm{I}\right] q_{j} \otimes f_{2, j}\right\|_{0}^{2}
$$

and, still by orthogonality of $\left(f_{2, j}\right)_{j}$, we obtain the claimed inequality (56), the constant $C$ being the same as in (55).
3. In the estimate (56) we only have the operator $\mathrm{A}_{1} \otimes \mathrm{I} \otimes \mathrm{I}$ that appears in the right-hand side and not the complete discrete 2 D elliptic operator A . The third step consists in proving a discrete elliptic regularity property that will allow us to get that there exists a constant $C>0$ independent of $h_{1}$ and $h_{2}$ such that

$$
\begin{equation*}
\llbracket \mathrm{B}^{*} q \rrbracket_{0}^{2} \leq C\|\mathrm{~A} q\|_{0}^{2}, \forall q \in E . \tag{57}
\end{equation*}
$$

The main idea is based on the following well-known computation: for any smooth and compactly supported function $f$ defined on $\mathbb{R}^{2}$, we can write

$$
\begin{aligned}
& \iint\left|\partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right)+\partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} f\right)\right|^{2} \\
= & \iint\left|\partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right)\right|^{2}+\left|\partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} f\right)\right|^{2} \\
& +2 \iint \partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right) \times \partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} f\right) .
\end{aligned}
$$

Moreover, by a double integration by parts, we find

$$
\iint \partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right) \times \partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} f\right)=\iint \gamma_{1}\left(x_{1}\right) \gamma_{2}\left(x_{2}\right)\left|\partial_{x_{1}} \partial_{x_{2}} f\right|^{2} \geq 0
$$

so that

$$
\iint\left|\partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right)\right|^{2} \leq \iint\left|\partial_{x_{1}}\left(\gamma_{1}\left(x_{1}\right) \partial_{x_{1}} f\right)+\partial_{x_{2}}\left(\gamma_{2}\left(x_{2}\right) \partial_{x_{2}} f\right)\right|^{2}
$$

We want to apply the same idea to prove roughly speaking that

$$
\left\|\mathrm{A}_{1} \otimes \mathrm{I} q\right\|_{0}^{2} \leq\left\|\mathrm{A}_{1} \otimes \mathrm{I} q+\mathrm{I} \otimes \mathrm{~A}_{2} q\right\|_{0}^{2}
$$

which, according to (56), would prove (57).
However, because of boundary terms, it is easier to prove this inequality for compactly supported discrete functions, that is for $q \in E_{1}^{00} \otimes E_{2}^{00}$. We will thus proceed by extension and truncation of the discrete functions under study.
Let $\xi_{1}:\left(0, a_{1}\right) \rightarrow \mathbb{R}$ be a smooth function such that $\xi_{1}=1$ on $\left(0, \frac{a_{1}}{3}\right)$, and $\xi_{1}=$ 0 on $\left(\frac{2 a_{1}}{3}, a_{1}\right)$, and $\xi_{2}:\left(0, a_{2}\right) \rightarrow \mathbb{R}$ be a smooth compactly supported function such that $\xi_{2}=1$ on $\omega_{2}$. Such a function exists thanks to the assumption (42). We introduce now the truncation operator on $E_{1} \otimes E_{2}$, defined by $\mathrm{T}=\xi_{1} \otimes \xi_{2}$. The choice of $\xi_{1}$ and $\xi_{2}$ implies that

$$
\left.\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0}^{2}=\llbracket\left(\mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*}\right)(\mathrm{T} \otimes \mathrm{I}) q\right) \rrbracket_{0}^{2}
$$

and thus by (56) we have

$$
\llbracket \mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q \rrbracket_{0}^{2} \leq C\left\|\left(\mathrm{~A}_{1} \otimes \mathrm{I} \otimes \mathbf{I}\right)(\mathrm{T} \otimes \mathbf{I}) q\right\|_{0}^{2}
$$

The discrete function $(\mathrm{T} \otimes \boldsymbol{I}) q$ vanish near all the boundaries of the domain, except the one corresponding to $\left\{x_{1}=0\right\}$ because of the choice of $\xi_{1}$. We will now introduce a symmetrization procedure that will let us work with a compactly supported discrete function.
We start by introducing the extended space in the first variable defined by

$$
\bar{E}_{1}:=\mathbb{R}^{\left\{-N_{1}, \ldots, N_{1}\right\}},
$$

which stands for discrete functions defined on a uniform discretisation of $\left(-a_{1}, a_{1}\right)$ with a mesh size $h_{1}$. If we denote by $\left(\bar{e}_{1, j}\right)_{-N_{1} \leq j \leq N_{1}}$ the canonical basis of $\bar{E}_{1}$, we can define the odd symmetrization operator $\mathrm{S}: E_{1} \rightarrow \bar{E}_{1}$ by

$$
\mathrm{Se}_{1, j}=\bar{e}_{1, j}-\bar{e}_{1,-j}, \quad \forall j \in\left\{1, \ldots, N_{1}\right\}
$$

With this notation, for $q \in E_{1} \otimes E_{2},(\mathrm{~S} \otimes \mathrm{I}) q$ is the odd symmetrization of $q$ with respect to the $x_{1}$ variable in the extended 2 D domain.
We consider $\bar{\gamma}_{1}$ to be the even extension of $\gamma_{1}$ to $\left(-a_{1}, a_{1}\right)$ and let $\overline{\mathrm{A}}_{1}$ be the discrete diffusion operator associated with $\bar{\gamma}_{1}$ and defined on $\bar{E}_{1}$ in the same way as in (44).
A simple computation shows that the symmetrization is compatible with the diffusion operator definitions, that is

$$
\begin{equation*}
\overline{\mathrm{A}}_{1} \mathrm{~S} q=\mathrm{SA}_{1} q, \quad \forall q \in E_{1} \tag{58}
\end{equation*}
$$

In particular, using the same notation for the norm in $\bar{E}_{1}$ as for the one in $E_{1}$, we have

$$
\left\|\overline{\mathrm{A}}_{1} \mathrm{~S} q\right\|_{0,1}^{2}=2\left\|\mathrm{~A}_{1} q\right\|_{0,1}^{2}, \quad \forall q \in E_{1}
$$

As a consequence, all the previous estimates lead to the following inequality

$$
\left\|\mathrm{B}_{1}^{*} \otimes \mathrm{~B}_{2}^{*} \otimes \mathrm{~B}^{*} q\right\|_{0,2}^{2} \leq C\left\|\overline{\mathrm{~A}}_{1} \otimes \mathrm{I} \otimes \mathrm{I} \bar{q}\right\|_{0}^{2}, \forall q \in E
$$

where, we have set $\bar{q}=(\mathbf{S} \otimes \mathbf{I} \otimes \mathbf{I})(\mathrm{T} \otimes \mathbf{I}) q$. Observe now that, by construction of the symmetrization and truncation operators, the discrete function $\bar{q}$ belongs to $\bar{E}_{1}^{00} \otimes E_{2}^{00} \otimes \mathbb{R}^{n}$.
Note that for such compactly supported discrete functions we have the algebraic identities

$$
\begin{aligned}
& \left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}\right) \bar{q}=\left(\left(-\bar{\nabla}_{1}^{-} \bar{\gamma}_{1}^{+} \bar{\nabla}_{1}^{+}\right) \otimes \mathrm{I}\right) \bar{q} \\
& \left(\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}=\left(\mathrm{I} \otimes\left(-\nabla_{2}^{-} \gamma_{2}^{+} \nabla_{2}^{+}\right)\right) \bar{q}
\end{aligned}
$$

We can now make the following computation

$$
\begin{aligned}
\left\|\left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}\right\|_{0}^{2}= & \left\|\overline{\mathrm{A}}_{1} \otimes \mathrm{I} \bar{q}\right\|_{0}^{2}+\left\|\mathrm{I} \otimes \mathrm{~A}_{2} \bar{q}\right\|_{0}^{2} \\
& +2\left\langle\overline{\mathrm{~A}}_{1} \otimes \mathrm{I} \bar{q}, \mathrm{I} \otimes \mathrm{~A}_{2} \bar{q}\right\rangle_{0}
\end{aligned}
$$

and the double product term can be evaluated as follows, the discrete integration by parts being justified by (48) and the fact that $\nabla_{i}^{ \pm} E_{i}^{00} \subset E_{i}^{0}$,

$$
\begin{aligned}
\left\langle\overline{\mathrm{A}}_{1} \otimes \mathrm{I} \bar{q}, \mathrm{I} \otimes \mathrm{~A}_{2} \bar{q}\right\rangle_{0} & =\left\langle\left(\bar{\nabla}_{1}^{-} \bar{\gamma}_{1}^{+} \bar{\nabla}_{1}^{+}\right) \otimes \mathrm{I} \bar{q}, \mathrm{I} \otimes\left(\nabla_{2}^{-} \gamma_{2}^{+} \nabla_{2}^{+}\right) \bar{q}\right\rangle_{0} \\
& =\left\langle\left(\bar{\gamma}_{1}^{+} \otimes \gamma_{2}^{+}\right)\left(\bar{\nabla}_{1}^{+} \otimes \nabla_{2}^{+}\right) \bar{q},\left(\bar{\nabla}_{1}^{+} \otimes \nabla_{2}^{+}\right) \bar{q}\right\rangle_{0} \geq 0
\end{aligned}
$$

At the end, we obtain that

$$
\left\|\overline{\mathrm{A}}_{1} \otimes \mathrm{I} \bar{q}\right\|_{0}^{2} \leq\left\|\left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}\right\|_{0}^{2}
$$

To conclude the proof of the final claim, we just need to show that

$$
\left\|\left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}\right\|_{0}^{2} \leq C\|\mathrm{~A} \bar{q}\|_{0}^{2}
$$

for a $C>0$ independent of the mesh.
Recall that $\bar{q}=(\mathrm{S} \otimes \mathrm{I}) \mathrm{T} q$ and by (58) we have

$$
\left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}=(\mathrm{S} \otimes \mathrm{I}) \mathrm{AT} q
$$

and in particular we have

$$
\left\|\left(\overline{\mathrm{A}}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right) \bar{q}\right\|_{0}^{2}=2\|\mathrm{AT} q\|_{0}^{2}
$$

It thus remains to evaluate the norm of $\mathrm{AT} q$ by the one of $\mathrm{A} q$. A simple algebraic computation leads to

$$
\begin{aligned}
\mathrm{AT} q= & \left(\mathrm{A}_{1} \otimes \mathrm{I}+\mathrm{I} \otimes \mathrm{~A}_{2}\right)\left(\xi_{1} \otimes \xi_{2}\right) q \\
= & \left(\xi_{1} \otimes \xi_{2}\right)(\mathrm{A} q)+\left(\left(\mathrm{A}_{1} \xi_{1}\right) \otimes \xi_{2}\right) q+\left(\xi_{1} \otimes\left(\mathrm{~A}_{2} \xi_{2}\right)\right) q \\
& +\left(\gamma_{1}^{+}\left(\nabla_{1}^{+} \xi_{1}\right) \otimes \mathrm{I}\right)\left(\nabla_{1}^{+} \otimes \mathrm{I}\right) q+\left(\gamma_{1}^{-}\left(\nabla_{1}^{-} \xi_{1}\right) \otimes \mathrm{I}\right)\left(\nabla_{1}^{-} \otimes \mathrm{I}\right) q \\
& +\left(\mathrm{I} \otimes \gamma_{2}^{+}\left(\nabla_{2}^{+} \xi_{2}\right)\right)\left(\mathrm{I} \otimes \nabla_{2}^{+}\right) q+\left(\mathrm{I} \otimes \gamma_{2}^{-}\left(\nabla_{2}^{-} \xi_{2}\right)\right)\left(\mathrm{I} \otimes \nabla_{2}^{-}\right) q
\end{aligned}
$$

Since $\xi_{1}$ and $\xi_{2}$ are smooth and $\gamma_{1}$ and $\gamma_{2}$ are bounded, we conclude, with the mean-value theorem that

$$
\begin{aligned}
&\|\mathrm{AT} q\|_{0} \leq C\|\mathrm{~A} q\|_{0}+C\|q\|_{0}+C\left\|\left(\nabla_{1}^{+} \otimes \mathrm{I}\right) q\right\|_{0} \\
&+C\left\|\left(\nabla_{1}^{-} \otimes \mathrm{I}\right) q\right\|_{0}+C\left\|\left(\mathrm{I} \otimes \nabla_{2}^{+}\right) q\right\|_{0}+C\left\|\left(\mathrm{I} \otimes \nabla_{2}^{-}\right) q\right\|_{0}
\end{aligned}
$$

The conclusion follows by (49), (50), (52) and the Cauchy-Schwarz and Young inequalities.
4. The last step of the proof consists in showing (54) by using (57) a classical energy estimate. Let $\psi \in E$ and let $q(t)=e^{-t \mathrm{~L}^{*}} \psi$. We set $z(t)=\mathrm{I} \otimes \mathrm{I} \otimes$ $e^{t \mathrm{C}^{*}} q(t)$. One can check that $z$ satisfies

$$
\left\{\begin{array}{r}
\partial_{t} z+\mathrm{A} z=0  \tag{59}\\
z(0)=\psi .
\end{array}\right.
$$

We multiply (59) by $\mathrm{A} z$ and integrate on $(0, T)$,

$$
\langle z(T), \mathrm{A} z(T)\rangle_{0}-\langle\psi, \mathrm{A} \psi\rangle_{0}+2 \int_{0}^{T}\|\mathrm{~A} z(t)\|_{0}^{2} d t=0,
$$

thus, since the operator A is positive,

$$
\int_{0}^{T}\|\mathrm{~A} z(t)\|_{0}^{2} d t \leq\|\psi\|_{1, \mathrm{~A}}^{2} .
$$

Since A commutes with $\mathrm{I} \otimes \mathrm{I} \otimes e^{t \mathrm{C}^{*}}$, we find that

$$
\int_{0}^{T}\left\|\mathrm{I} \otimes \mathrm{I} \otimes e^{t \mathrm{C}^{*}} \mathrm{~A} q(t)\right\|_{0}^{2} d t \leq\|\psi\|_{1, \mathrm{~A}}^{2},
$$

Moreover, $\left\|\mathrm{I} \otimes \mathrm{I} \otimes e^{t \mathrm{C}^{*}} \mathrm{~A} q(t)\right\|_{0}^{2} \geq e^{-2 T\|\mathrm{C}\|}\|\mathrm{A} q(t)\|_{0}^{2}$, for any $t \in(0, T)$ and so

$$
\int_{0}^{T}\|\mathrm{~A} q(t)\|_{0}^{2} d t \leq e^{2 T\|\mathrm{C}\|}\|\psi\|_{1, \mathrm{~A}}^{2}
$$

Applying (57) to $q(t)$ for any $t \in(0, T)$ and integrating in time this inequality, we finally get

$$
\int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t \leq C e^{C T}\|\psi\|_{1, \mathrm{~A}}^{2}
$$

and this concludes the proof of (54).
Remark 4.2. Note that we just proved that $M_{\text {adm }}$ depends on $T$ like $e^{C T}$ and this is consistant with estimates of $v$ and $y(T)$ given by Theorem 4.1.

Now we can prove Assumption 3.2.
Proof (of Assumption 3.2). In [7] the authors proved a discrete Lebeau-Robbiano inequality on quite general meshes. We translate the statement of their Theorem 6.1 in our setting.

Theorem 4.2. There exists $C>0, \varepsilon>0$ and $h_{0}>0$, depending on $\gamma_{2}, \omega_{2}$ such that if $h_{2} \leq h_{0}$, we have for all $0<\mu \leq \varepsilon / h_{2}^{2}$,

$$
\sum_{k: \lambda_{2, k} \leq \mu}\left|\alpha_{k}\right|^{2}=\int_{\Omega_{2}}\left|\sum_{k: \lambda_{2, k} \leq \mu} \alpha_{k} \phi_{2, k}\right|^{2} \leq C e^{C \sqrt{\mu}} \int_{\omega_{2}}\left|\sum_{k: \lambda_{2, k} \leq \mu} \alpha_{k} \phi_{2, k}\right|^{2}, \quad \forall\left(\alpha_{k}\right)_{1 \leq k \leq N_{1}} \subset \mathbb{C}
$$

This result exactly yields that Assumption 3.2 is fulfilled in this setting with

$$
\mu_{E_{2}}=\frac{\varepsilon}{h_{2}^{2}} .
$$

Finally, we prove that Assumption 3.1 holds.
Proof (of Assumption 3.1). We base our proof on the strategy developped in [1], where the semi-discretized (on a uniform mesh) boundary null-control problem in space dimension 1 with operator $\mathrm{A}_{1}$ is tackled by applying the moments method. However, the explicit dependance in $T$ of the control cost by some bound in $e^{C / T}$ was not given in that work. This is crucial in the present analysis.

We can actually obtain this precised bound by using the expression of the control obtained with this method, in setting (S2) described in [1], and by using the refined bounds on biorthogonal families of exponential functions given by Theorem 1.5 of [4]. We will just describe here the new estimate that we need to adapt the results of [1] to our needs.

We first set $\varepsilon=1 / 2$ for instance (any value between 0 and 1 would be acceptable), and we follow [1], to define

$$
k_{\max , \varepsilon}^{h_{1}}:=\max \left\{k \in \llbracket 1, N_{1} \rrbracket ; \lambda_{1, k}<\frac{4}{h_{1}^{2}} \gamma_{1, \min }(1-\varepsilon)\right\}
$$

and :

$$
\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}:=\left\{\begin{array}{l}
\lambda_{1, k} \text { for } k \in \llbracket 1, k_{\max , \varepsilon}^{h_{1}} \rrbracket, \\
\lambda_{1, k_{\max , \varepsilon}^{h_{1}}}+4 \gamma_{1, \min } k^{2} \text { for } k \geq k_{\max , \varepsilon}^{h_{1}}+1
\end{array}\right.
$$

the sequence $\tilde{\Lambda}_{\varepsilon}^{h_{1}}$ satisfies the following items required to apply Theorem 1.5 of [4]

1. $\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k} \neq\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{n}$, for all $k, n \in \mathbb{N}$ with $k \neq n$;
$2,3\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k} \in \mathbb{R}_{+}^{*}$, for every $k \geq 1$;
2. $\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}$ is nondecreasing;
3. First, applying Theorem 3.2 of [1] there exists $C>0$ such that

$$
\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k+1}-\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k} \geq C \sqrt{\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}}, \forall k \geq 1 .
$$

Second, note that according to Lemma 3.3 of [1], there exist $c_{1}$ and $c_{2}$ two positive constants such that :

$$
\begin{equation*}
c_{1} k^{2} \leq \lambda_{1, k} \leq c_{2} k^{2}, \forall k \in \llbracket 1, N_{1} \rrbracket . \tag{60}
\end{equation*}
$$

Therefore,

$$
\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k+1}-\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k} \geq C k, \forall k \geq 1
$$

Take $k, n \in \mathbb{N}^{*}, k>n$, we have :

$$
\left|\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}-\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{n}\right| \geq C \sum_{i=n}^{k-1} i \geq C\left(k^{2}-n^{2}\right)
$$

6. Finally, let $\mathcal{N}$ the counting function associated with the sequence $\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k \geq 1}$, defined by

$$
\mathcal{N}(r)=\sharp\left\{k:\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k} \leq r\right\} .
$$

There exist $\alpha>0, p_{\min }>0$ and $p_{\max }>0$ such that for every $r>0$,

$$
-\alpha+p_{\min } \sqrt{r} \leq \mathcal{N}(r) \leq \alpha+p_{\max } \sqrt{r}
$$

Indeed, from item 4, it follows easily from the definition of $\mathcal{N}$ that $\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{\mathcal{N}(r)} \leq$ $r$ and $\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{\mathcal{N}(r)+1} \geq r$, and according to (60),

$$
c_{1} \mathcal{N}(r)^{2} \leq r \leq c_{2}(\mathcal{N}(r)+1)^{2}
$$

thus,

$$
\sqrt{\frac{r}{c_{2}}}-1 \leq \mathcal{N}(r) \leq \sqrt{\frac{r}{c_{1}}}
$$

We can now apply Theorem 1.5 of [4] to conclude that there exists $T_{0}>0$, such that for every $0<T<T_{0}$, there exists a family of functions in $L^{2}(0, T)$ denoted by $\left(q_{j, k}^{\tilde{\Lambda}_{\varepsilon}^{h_{1}}}\right)_{k \geq 1, j \in \llbracket 1, d-1 \rrbracket}$, satisfying for all $k, l \geq 1$, and $i, j \in \llbracket 0, d-1 \rrbracket$,

$$
\int_{0}^{T} q_{j, k}^{\tilde{\Lambda}_{\varepsilon}^{h_{1}}}(t)(t-T)^{i} \exp \left(-\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}(T-t)\right) d t=\delta_{k, l} \delta_{i, j}
$$

and for all $k \geq 1$ and $j \in \llbracket 0, d-1 \rrbracket$,

$$
\begin{equation*}
\left\|q_{j, k}^{\tilde{h}_{\varepsilon}^{h_{1}}}\right\|_{L^{2}(0, T)} \leq C \exp \left(C \sqrt{\left(\tilde{\Lambda}_{\varepsilon}^{h_{1}}\right)_{k}}+\frac{C}{T}\right) \tag{61}
\end{equation*}
$$

where $C$ does not depend on $h_{2}$.
Remark 4.3. In [4], Theorem 1.5, hypothesis of 6 on the counting function is slightly different from item 6 given above. Indeed, in this reference the authors require the following condition : for some $p, \alpha>0$,

$$
\begin{equation*}
|p \sqrt{r}-\mathcal{N}(r)| \leq \alpha, \forall r>0 \tag{62}
\end{equation*}
$$

Actually, looking carefully at the proof of Theorem 1.5, we realize that we can use item 6 above instead of hypothesis (62). Indeed, we can just replace $p$ by $p_{\max }$ in the definition of the number d at page 2988 and in the proofs of Theorem 4.3 at page 2987, Lemma A. 3 at page 2994 (note that the property $\lim _{t \rightarrow+\infty} \frac{\mathcal{N}(t)}{t}=0$ is preserved), and Lemma $A .4$ at page 2996 and change $p$ into $p_{\min }$ at page 2993 in the proof of Lemma A.1.

Now, in the proof of [1, Theorem 5.4], we can use the improved estimate (61) into the bound of the control (with $t_{0}=\frac{T}{2}$ ), and find that

$$
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq C_{\varepsilon}\left\|y_{0}\right\|_{-1, \mathrm{~A}_{1}} \sum_{k=1}^{k_{m a x, \varepsilon}^{h_{1}}} e^{-\lambda_{1, k} \frac{T}{2}} e^{C \sqrt{\lambda_{1, k}}+\frac{C}{T}}
$$

and using (60),

$$
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq C_{\varepsilon} e^{\frac{C}{T}}\left\|y_{0}\right\|_{-1, \mathrm{~A}_{1}} \sum_{k=1}^{+\infty} e^{-c_{2} k^{2} \frac{T}{2}+C \sqrt{c_{1}} k}
$$

thus there exists $C_{1}>0$ depending only on $\varepsilon$ (whose value has been arbitrarily set to $1 / 2$ ) and $\gamma_{1}$ such that

$$
\begin{equation*}
\llbracket v \rrbracket_{L^{2}(0, T ; U)} \leq C_{1} e^{\frac{C_{1}}{T}}\left\|y_{0}\right\|_{-1, \mathrm{~A}_{1}} \tag{63}
\end{equation*}
$$

Indeed if $k \geq k_{0}:=\left\lceil\frac{C \sqrt{c_{1}}+1}{c_{2} T / 2}\right\rceil$, we have $-c_{2} k^{2} \frac{T}{2}+C \sqrt{c_{1}} k \leq-k$, so

$$
\begin{aligned}
\sum_{k=1}^{+\infty} e^{-c_{2} k^{2} \frac{T}{2}+C \sqrt{c_{1}} k} & \leq \sum_{k=1}^{k_{0}} e^{-c_{2} k^{2} \frac{T}{2}+C \sqrt{c_{1}} k}+\sum_{k=k_{0}+1}^{+\infty} e^{-k} \\
& \leq k_{0} e^{C \sqrt{c_{1}} k_{0}}+\frac{1}{1-e}
\end{aligned}
$$

which gives (63) since $k_{0} \leq \frac{C_{\gamma_{1}, \varepsilon}}{T}$.
Still using the estimations in the proof of Theorem 5.4 of [1], we get

$$
\|y(T)\|_{-1, \mathrm{~A}_{1}} \leq e^{C_{2}\left(1+\frac{1}{T}\right)} e^{-\mu_{E_{1}} T}\left\|y_{0}\right\|_{-1, \mathrm{~A}_{1}}
$$

with

$$
\mu_{E_{1}}=\frac{C_{1}}{h_{1}^{2}}
$$

and $C_{1}>0$ and $C_{2}>0$ depending only on $\gamma_{1}$.
Now we can apply Theorem 3.1 with $\mu_{E_{1}}$ and $\mu_{E_{2}}$ both of the same order $\frac{C}{h^{2}}$. The proof of Theorem 4.1 is complete.
4.2. Dirichlet boundary null control of a continuous n-dimensional system on a cylindrical domain of dimension $d$. Our aim is to show how the finite dimensional framework developped in Section 3 actually applies to the study of nullcontrollability problems of tensorized parabolic systems.

As an illustration of this statement, we shall give a short proof of the main result of [4] by using Theorem 3.1 of the present article. This is of course not surprising since our approach is directly inspired from the one developped in [4]. However, it seems to us interesting to show how our general finite dimensional framework actually encompasses already known results through a spectral projection technique. This example should convince the reader that, by using the same strategy, one can easily adapt the proof to other kinds of tensorized controlled systems like, for instance, Neumann or Robin boundary controls, or even mixed (distributed and boundary) controls, as soon as we have in hand a suitable controllability result on the associated 1D system.

Finally, since we have taken care of all the constants in the proofs, this strategy can be used to derive controllability properties that are uniform with respect to some parameters present in the problem. As an illustration, in the case of a 1D Robin boundary control problem, we can prove estimates that are uniform in the Robin parameter (see [5]). By the present technique, those result will automatically be translated to the corresponding multi-D result, generalizing the one proved in [4].

Let us recall the statement of [4, Theorem 1.3] with the notation of the present paper.

Theorem 4.3. Let $\Omega_{1}=(0,1)$, and $\Omega_{2}$ be a smooth bounded connected domain of $\mathbb{R}^{d-1}$ and set $\Omega=\Omega_{1} \times \Omega_{2}$.

We suppose given a coupling matrix C and a control matrix B . Assume that the
following system of $n$ equations

$$
\left\{\begin{aligned}
& \partial_{t} y^{1}+\mathcal{A}_{1} \otimes \mathbf{I} y^{1}+\mathcal{I} \otimes \mathrm{C} y^{1}=0, \\
& y^{1}=1_{\{0\}} \otimes\left(\mathrm{B} v^{1}\right) \text { in }(0, T) \times \Omega_{1} \\
& y^{1}(0)=y_{0}^{1} \\
& \text { in } \Omega_{1},
\end{aligned}\right.
$$

is null controllable for any $y_{0}^{1} \in H^{-1}\left(\Omega_{1}\right) \otimes \mathbb{R}^{n}$ and any time $T>0$ with, in addition, the following bound:

$$
\llbracket v^{1} \rrbracket_{L^{2}\left(0, T ; \mathbb{R}^{m}\right)} \leq e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left\|y_{0}^{1}\right\|_{H^{-1}\left(\Omega_{1}\right) \otimes \mathbb{R}^{n}}
$$

Then, for any nonempty open set $\omega_{2} \subset \Omega_{2}$, the following system of $n$ equations:

$$
\left\{\begin{align*}
& \partial_{t} y+\mathcal{L} y=0, \text { in }(0, T) \times \Omega  \tag{64}\\
& y=\left(1_{\{0\} \times \omega_{2}} \otimes \mathrm{~B}\right) v \\
& \text { in }(0, T) \times \partial \Omega \\
& y(0)=y_{0} \\
& \text { in } \Omega,
\end{align*}\right.
$$

where $\mathcal{L}:=\mathcal{A}_{1} \otimes \mathcal{I} \otimes \mathbf{I}+\mathcal{I} \otimes \mathcal{A}_{2} \otimes \mathbf{I}+\mathcal{I} \otimes \mathcal{I} \otimes \mathbf{C}$, is null controllable for any $y_{0} \in$ $H^{-1}(\Omega) \otimes \mathbb{R}^{n}$ and any time $T>0$ with, in addition, the following bound:

$$
\begin{equation*}
\llbracket v \rrbracket_{L^{2}((0, T) \times \partial \Omega)^{m}} \leq e^{M_{\mathrm{obs}, 1}\left(1+\frac{1}{T}\right)}\left\|y_{0}\right\|_{H^{-1}(\Omega) \otimes \mathbb{R}^{n}} \tag{65}
\end{equation*}
$$

Proof. Let $y_{0} \in H^{-1}(\Omega) \otimes \mathbb{R}^{n}$, recall that the solution $y$ of (64) satisfies for any $\psi \in H_{0}^{1}(\Omega) \otimes \mathbb{R}^{n}$, the equality

$$
\begin{aligned}
\langle y(T), \psi\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} & -\left\langle y_{0}, e^{-T \mathcal{L}^{*}} \psi\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
& =-\left[\left(\left.\frac{\partial}{\partial x_{1}}\right|_{x_{1}=0} \otimes 1_{\omega_{2}} \otimes \mathrm{~B}^{*}\right) e^{-(T-\cdot) \mathcal{L}^{*}} \psi, v\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

Therefore, $v$ is a null-control for this system, if and only if, it satisfies

$$
\begin{aligned}
&-\left\langle y_{0}, e^{-T \mathcal{L}^{*}} \psi\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
&=-\left[\left(\left.\frac{\partial}{\partial x_{1}}\right|_{x_{1}=0} \otimes 1_{\omega_{2}} \otimes \mathrm{~B}^{*}\right) e^{-(T-\cdot) \mathcal{L}^{*}} \psi, v\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

for any $\psi \in H_{0}^{1}(\Omega) \otimes \mathbb{R}^{n}$. Actually, since the problem is linear, it is enough to check the previous equality for any $\psi$ belonging to a total family of $H_{0}^{1}(\Omega) \otimes \mathbb{R}^{n}$.

Let us consider the total family of $H_{0}^{1}(\Omega)$ made of the eigenfunctions of $\mathcal{A}_{1}$ and $\mathcal{A}_{2}$. For $i=1,2$, we denote by $\left(\phi_{i, j}, \lambda_{i, j}\right)_{j \geq 1}$ the eigenfunctions and eigenvalues of the operator $\mathcal{A}_{i}$ with homogeneous Dirichlet boundary conditions. We choose them to form an orthonormal basis of $L^{2}\left(\Omega_{i}\right)$. We will thus consider the total family of $H_{0}^{1}(\Omega)$ defined by $\left(\phi_{1, i} \otimes \phi_{2, j}\right)_{i, j}$ that will be tensorized with the canonical basis of $\mathbb{R}^{n},\left(e_{k}\right)_{k \in\{1, \ldots, n\}}$ to finally produce a total family of $H_{0}^{1}(\Omega) \otimes \mathbb{R}^{n}$.

We are thus led to find a control satisfying, $\forall i, j \geq 1$ and $\forall k \in\{1, \ldots, n\}$,

$$
\begin{aligned}
& e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)}\left\langle y_{0}, \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
& \quad=\left[e^{-(T-.)\left(\lambda_{1, i}+\lambda_{2, j}\right)}\left(\phi_{1, i}\right)^{\prime}(0) \otimes\left(1_{\omega_{2}} \phi_{2, j}\right) \otimes\left(\mathrm{B}^{*} e^{-(T-\cdot) \mathrm{C}^{*}} e_{k}\right), v\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

In order to apply Theorem 3.1, we need to consider a projection of (64) on a finite dimensional space. For $i=1,2$ and $J \geq 1$, we define the spaces

$$
F_{i, J}=\operatorname{span}\left(\phi_{i, j}, j \leq J\right),
$$

and the operator $\mathrm{A}_{i}: F_{i, J} \rightarrow F_{i, J}$ as the restriction to $F_{i, J}$ of $\mathcal{A}_{i}$, that is the unique isomorphism of $F_{i, J}$ satisfying $\mathrm{A}_{i} \phi_{i, j}=\lambda_{i, j} \phi_{i, j}$, for any $j \leq J$.

Let $\mathcal{P}_{i, J}$ be the orthogonal projection from the space $H^{-1}\left(\Omega_{i}\right)$ equipped with the inner product $\langle\bullet, \bullet\rangle_{-1, \mathcal{A}_{i}}$ onto the space $F_{i, J}$.

We define also $\mathcal{P}_{J}:=\mathcal{P}_{1, J} \otimes \mathcal{P}_{2, J} \otimes \mathrm{I}$ which is an orthogonal projection from $H^{-1}(\Omega) \otimes \mathbb{R}^{n}$ onto $F_{1, J} \otimes F_{2, J} \otimes \mathbb{R}^{n}$, with respect to $\langle\bullet, \bullet\rangle_{-1, \mathcal{A}}$.

REmARK 4.4. Since the eigenfunctions $\left(\phi_{i, j}\right)_{j \geq 1}$ are orthogonal with respect to the $L^{2}\left(\Omega_{i}\right), H^{-1}\left(\Omega_{i}\right)$ and $H_{0}^{1}\left(\Omega_{i}\right)$ inner products, we deduce that the projection $\mathcal{P}_{J}$ (resp. $\left.\mathcal{P}_{i, J}\right)$ is orthogonal with respect to $\langle\bullet, \bullet\rangle_{s, \mathcal{A}}\left(\right.$ resp. $\left.\langle\bullet, \bullet\rangle_{s, \mathcal{A}_{i}}\right)$, for any $s \in\{-1,0,1\}$.

Hypothesis (4.3) allows to apply Assumption 3.1 with $E_{1}=F_{1, J}, U_{1}=\mathbb{R}, \mathcal{B}_{1}=$ $\Delta_{1} r$, where we have used the affine function $r(x)=1-x$ (note that $\mathcal{B}_{1}^{*}=-1_{\{0\}} \partial_{x_{1}}$ ) and $\mu_{E_{1}}=+\infty$. Moreover, thanks to the Lebeau-Robbiano's spectral inequality in the domain $\Omega_{2}$, (see $[13,12]$ ), Assumption 3.2 is fulfilled with $E_{2}=F_{2, J}, U_{2}=F_{2, J}$, $\mathcal{B}_{2}: U_{2} \rightarrow E_{2}, \mathcal{B}_{2}=\mathcal{P}_{2, J} 1_{\omega_{2}}$, and $\mu_{E_{2}}=+\infty$. Indeed, the usual spectral inequality gives that there exists $C>0$ such that for any $\mu>0, J \in \mathbb{N}^{*}$

$$
\|\psi\|_{L^{2}\left(\Omega_{2}\right)}^{2} \leq C e^{C \sqrt{\mu}} \int_{\Omega_{2}}\left(1_{\omega_{2}} \psi\right)^{2}, \quad \forall \psi \in \operatorname{Span}\left(\phi_{2, j}, \quad \lambda_{2, j} \leq \mu\right) \cap E_{2}
$$

Moreover,

$$
\int_{\Omega_{2}}\left(1_{\omega_{2}} \psi\right)^{2}=\int_{\Omega_{2}}\left(1_{\omega_{2}} \psi\right) \psi
$$

and according to Remark 4.4 and given that $\psi \in E_{2}$,

$$
\int_{\Omega_{2}}\left(1_{\omega_{2}} \psi\right) \psi=\int_{\Omega_{2}}\left(\mathcal{P}_{2, J} 1_{\omega_{2}} \psi\right) \psi=\int_{\Omega_{2}}\left(\mathcal{B}_{2} \psi\right) \psi=\int_{\Omega_{2}}\left(\mathcal{B}_{2}^{*} \psi\right) \psi
$$

since $\mathcal{B}_{2}$ is self-adjoint. It follows from the Cauchy-Schwarz inequality that

$$
\|\psi\|_{L^{2}\left(\Omega_{2}\right)} \leq C e^{C \sqrt{\mu}}\left\|\mathcal{B}_{2}^{*} \psi\right\|_{L^{2}\left(\Omega_{2}\right)}, \quad \forall \psi \in \operatorname{Span}\left(\phi_{2, j}, \quad \lambda_{2, j} \leq \mu\right) \cap E_{2}
$$

Thus, for any fixed $J$, Theorem 3.1 applies with $\mu^{*}$ as large as we want. Thus, letting $\mu^{*}$ going to infinity, we finally build a null control for the projected finite dimensional problem denoted by $v_{J} \in L^{2}\left(0, T ; U_{2}\right) \otimes \mathbb{R}^{m}$ that drives solution $y$ of (24) from $P_{J} y_{0}$ to zero and which satisfies the following bound, uniformly with respect to $J$ :

$$
\begin{equation*}
\llbracket v_{J} \rrbracket_{L^{2}((0, T) \times \partial \Omega)^{m}} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)}\left\|P_{J} y_{0}\right\|_{-1, \mathrm{~A}} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-1, \mathcal{A}} \tag{66}
\end{equation*}
$$

where $\|\bullet\|_{-1, \mathrm{~A}}^{2}=\langle\bullet, \bullet\rangle_{-1, \mathrm{~A}}=\left\langle\bullet, \mathrm{A}^{-1} \bullet\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)}$.
From (66) we can infer that there exists a subsequence of $\left(v_{J}\right)_{J}$ denoted by $\left(v_{J^{\prime}}\right)_{J^{\prime}}$ which weakly converges in $L^{2}((0, T) \times \partial \Omega)^{m}$ to a limit denoted by $v$, satisfying also

$$
\llbracket v \rrbracket_{L^{2}((0, T) \times \partial \Omega)^{m}} \leq e^{M_{\mathrm{obs}}\left(1+\frac{1}{T}+T\right)}\left\|y_{0}\right\|_{-1, \mathcal{A}}
$$

We claim that this limit drives the solution $y$ of (64) from $y_{0}$ to zero. Indeed, let $k \in\{1, \ldots, N\}, i, j \geq 1$, and $J \geq \max (i, j)$.

First, we have by definition of a solution of (24)

$$
\begin{aligned}
& -\left\langle\mathcal{P}_{J} y_{0}, e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
& \quad=\left[e^{-(T-.)\left(\lambda_{1, i}+\lambda_{2, j}\right)}\left(\mathrm{B}_{1}^{*} \phi_{1, i}\right) \otimes\left(\mathrm{B}_{2}^{*} \phi_{2, j}\right) \otimes\left(\mathrm{B}^{*} e^{-(T-.) \mathrm{C}^{*}} e_{k}\right), v_{J}\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

Second, since $J \geq i$ and $J \geq j$, and by definition of $\mathcal{P}_{J}$, we have

$$
\begin{aligned}
- & \left\langle\mathcal{P}_{J} y_{0}, e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T C^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
& =-\left\langle\mathcal{P}_{J} y_{0},\left(\lambda_{1, i}+\lambda_{2, j}\right) e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{-1, \mathrm{~A}} \\
& =-\left\langle y_{0},\left(\lambda_{1, i}+\lambda_{2, j}\right) e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{-1, \mathcal{A}}
\end{aligned}
$$

Hence, we get

$$
\begin{aligned}
&-\left\langle\mathcal{P}_{J} y_{0}, e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
&=-\left\langle y_{0}, e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)} \phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)}
\end{aligned}
$$

Third, since $F_{2, J}$ is stable by $\mathrm{A}_{2}, \mathcal{P}_{2, J}$ is an orthogonal projection with respect to the $L^{2}\left(\Omega_{2}\right)$ norm, see Remark 2.2. So, since $v_{J} \in L^{2}\left(0, T ; U_{2}\right) \otimes \mathbb{R}^{m}$ and $U_{2}=F_{2, J}$, according to remark 4.4, we have for any $t$

$$
\left[\mathcal{B}_{2}^{*} \phi_{2, j}, v_{j}(t)\right]_{L^{2}\left(\Omega_{2}\right)^{m}}=\left[\mathcal{P}_{2, J} 1_{\omega_{2}} \phi_{2, j}, v_{j}(t)\right]_{L^{2}\left(\Omega_{2}\right)^{m}}=\left[1_{\omega_{2}} \phi_{2, j}, v_{j}(t)\right]_{L^{2}\left(\Omega_{2}\right)^{m}}
$$

Hence,

$$
\begin{aligned}
{\left[\left(\mathrm{B}_{1}^{*} \phi_{1, i}\right) \otimes\left(\mathrm{B}_{2}^{*} \phi_{2, j}\right) \otimes\left(\mathrm{B}^{*} e^{-(T-.) \mathrm{C}^{*}} e_{k}\right), v_{J}\right]_{L^{2}((0, T) \times \partial \Omega)^{m}} } \\
\quad=\left[-\left(\phi_{1, i}\right)^{\prime}(0) \otimes 1_{\omega_{2}} \phi_{2, j} \otimes\left(\mathrm{~B}^{*} e^{-(T-.) \mathrm{C}^{*}} e_{k}\right), v_{J}\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

Gathering (11), (11) and (11), we get

$$
\begin{aligned}
-\left\langle y_{0}, e^{-T\left(\lambda_{1, i}+\lambda_{2, j}\right)}\right. & \left.\phi_{1, i} \otimes \phi_{2, j} \otimes\left(e^{-T \mathrm{C}^{*}} e_{k}\right)\right\rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} \\
& =\left[-\phi_{1, i}^{\prime}(0) \otimes 1_{\omega_{2}} \phi_{2, j} \otimes\left(\mathrm{~B}^{*} e^{-(T-.) \mathrm{C}^{*}} e_{k}\right), v_{J}\right]_{L^{2}((0, T) \times \partial \Omega)^{m}}
\end{aligned}
$$

Considering the subsequence $v_{J^{\prime}}$ in (11) and letting $J^{\prime}$ go to infinity, we get (11) and this concludes the proof.

## Appendix A. Proof of Lemma 2.1.

$1 . \Rightarrow 2$. For any $q_{T} \in \mathrm{D}^{-s} F_{T}$, we set $q(t)=e^{-(T-t) \mathrm{L}^{*}} q_{T}$ and we consider the initial data $y_{0}:=\mathrm{P}_{F_{0}} \mathrm{D}^{s} \mathrm{P}_{F_{0}}^{*} q(0)$; by assumption there exists a control $v$ and an associated solution $y$ satisfying (18).
By (16) we find

$$
\begin{aligned}
& \left\langle\mathrm{P}_{F_{0}} \mathrm{D}^{s} \mathrm{P}_{F_{0}}^{*} q(0), q(0)\right\rangle_{0}=\left\langle y(T), q_{T}\right\rangle_{0}-\int_{0}^{T}\left[v(t), \mathrm{B}^{*} q(t)\right]_{0} d t \\
& \left\|\mathrm{P}_{F_{0}}^{*}(q(0))\right\|_{s, \mathrm{D}}^{2}=\left\langle y(T), \mathrm{D}^{s} q_{T}\right\rangle_{-s, \mathrm{D}}-\int_{0}^{T}\left[v(t), \mathrm{B}^{*} q(t)\right]_{0} d t
\end{aligned}
$$

Since $\mathrm{D}^{s} q_{T} \in F_{T}$ and $\mathrm{P}_{F_{T}}$ is orthogonal with respect to $\langle\bullet \bullet \bullet\rangle_{-s, \mathrm{D}}$, we deduce

$$
\begin{aligned}
\left\|\mathrm{P}_{F_{0}}^{*}(q(0))\right\|_{s, \mathrm{D}}^{2} & \leq\left\langle\mathrm{P}_{F_{T}} y(T), \mathrm{D}^{s} q_{T}\right\rangle_{-s, \mathrm{D}} \\
& +\llbracket v \rrbracket_{L^{2}(0, T ; U)}\left(\int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t\right)^{1 / 2} .
\end{aligned}
$$

The Cauchy-Schwarz inequality gives

$$
\begin{aligned}
\left\|\mathrm{P}_{F_{0}}^{*}(q(0))\right\|_{s, \mathrm{D}}^{2} \leq & \left(\frac{1}{M_{\mathrm{rel}}^{2}}\left\|\mathrm{P}_{F_{T}} y(T)\right\|_{-s, \mathrm{D}}^{2}+\frac{1}{M_{\mathrm{obs}}^{2}} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}\right)^{\frac{1}{2}} \\
& \times\left(M_{\mathrm{rel}}^{2}\left\|\mathrm{D}^{s} q_{T}\right\|_{-s, \mathrm{D}}^{2}+M_{\mathrm{obs}}^{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t\right)^{\frac{1}{2}}
\end{aligned}
$$

Using the estimate (18) given in the assumptions we get

$$
\left\|\mathrm{P}_{F_{0}}^{*}(q(0))\right\|_{s, \mathrm{D}}^{2} \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}\left(M_{\mathrm{rel}}^{2}\left\|\mathrm{D}^{s} q_{T}\right\|_{-s, \mathrm{D}}^{2}+M_{\mathrm{obs}}^{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t\right)^{\frac{1}{2}}
$$

Recall that we have chosen $y_{0}=\mathrm{P}_{F_{0}} \mathrm{D}^{s} \mathrm{P}_{F_{0}}^{*} q(0)$, and that $\mathrm{P}_{F_{0}}$ is the orthogonal projection onto $F_{0}$ with respect to the inner product $\langle\cdot, \cdot\rangle_{-s, \mathrm{D}}$. Therefore,

$$
\left\|y_{0}\right\|_{-s, \mathrm{D}} \leq\left\|\mathrm{D}^{s} \mathrm{P}_{F_{0}}^{*} q(0)\right\|_{-s, \mathrm{D}}=\left\|\mathrm{P}_{F_{0}}^{*} q(0)\right\|_{s, \mathrm{D}}
$$

and thus

$$
\left\|\mathrm{P}_{F_{0}}^{*}(q(0))\right\|_{s, \mathrm{D}} \leq\left(M_{\mathrm{rel}}^{2}\left\|q_{T}\right\|_{s, \mathrm{D}}^{2}+M_{\mathrm{obs}}^{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t\right)^{\frac{1}{2}}
$$

which gives the claim.
$2 . \Rightarrow 1$. Let $y_{0}$ be an element of $F_{0}$. We apply the quite usual penalised HUM approach (as described for instance in [6]) to find such a control $v$. To this end, we introduce the functional defined on $\mathrm{D}^{-s} F_{T}$ by

$$
J\left(q_{T}\right):=\frac{1}{2} \int_{0}^{T} \llbracket \mathrm{~B}^{*} q(t) \rrbracket_{0}^{2} d t+\frac{M_{\mathrm{rel}}^{2}}{2 M_{\mathrm{obs}}^{2}}\left\|q_{T}\right\|_{s, \mathrm{D}}^{2}+\left\langle y_{0}, q(0)\right\rangle_{0}
$$

It is clear that $J$ is continuous, strictly convex and coercive; we denote by $q_{T}^{\text {opt }} \in \mathrm{D}^{-s} F_{T}$ its minimizer and by $q^{\text {opt }}(t)=e^{-(T-t) \mathrm{L}^{*}} q_{T}^{\text {opt }}$ the solution to the backward equation with final condition $q_{T}^{\mathrm{opt}}$. We set $v(t):=\mathrm{B}^{*} q^{\mathrm{opt}}(t)$ and introduce the associated solution $y$ to (11). Let us show the relation

$$
\begin{equation*}
\mathrm{P}_{F_{T}}(y(T))=-\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}} \mathrm{D}^{s} q_{T}^{\mathrm{opt}} \tag{67}
\end{equation*}
$$

Take any $q_{T} \in \mathrm{D}^{-s} F_{T}$ and let $q(t)=e^{-(T-t) \mathrm{L}^{*}} q_{T}$ be the solution of the adjoint equation with final condition $q_{T}$. The associated Euler-Lagrange equation for $J$ at the optimal point $q_{T}^{\text {opt }}$ writes

$$
\begin{equation*}
\int_{0}^{T}\left[\mathrm{~B}^{*} q^{\mathrm{opt}}(t), \mathrm{B}^{*} q(t)\right]_{0} d t+\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}}\left\langle q_{T}^{\mathrm{opt}}, q_{T}\right\rangle_{s, \mathrm{D}}+\left\langle y_{0}, q(0)\right\rangle_{0}=0 \tag{68}
\end{equation*}
$$

Combining (68) with (16) gives that for any $q_{T} \in \mathrm{D}^{-s} F_{T}$,

$$
\left\langle y(T), q_{T}\right\rangle_{0}=-\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}}\left\langle q_{T}^{\mathrm{opt}}, q_{T}\right\rangle_{s, \mathrm{D}}
$$

and thus,

$$
\left\langle y(T), \mathrm{D}^{s} q_{T}\right\rangle_{-s, \mathrm{D}}=-\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}}\left\langle\mathrm{D}^{s} q_{T}^{\mathrm{opt}}, \mathrm{D}^{s} q_{T}\right\rangle_{-s, \mathrm{D}}
$$

and using that $\mathrm{D}^{s} q_{T}^{\text {opt }} \in F_{T}$, and that $\mathrm{D}^{s} q_{T}$ is an arbitrary element in $F_{T}$, we finally find the equality (67).
Now we apply (68) with $q_{T}=q_{T}^{\text {opt }}$ to get

$$
\begin{aligned}
\llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}}\left\|q_{T}^{\mathrm{opt}}\right\|_{s, \mathrm{D}}^{2} & =\left\langle y_{0}, q^{\mathrm{opt}}(0)\right\rangle_{0} \\
& =\left\langle y_{0}, \mathrm{P}_{F_{0}}^{*} q^{\mathrm{opt}}(0)\right\rangle_{0} \\
& \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}\left\|\mathrm{P}_{F_{0}}^{*} q^{\mathrm{opt}}(0)\right\|_{s, \mathrm{D}}
\end{aligned}
$$

where the second equality comes from the fact that $y_{0} \in F_{0}$. Using now the hypothesis (19), we obtain

$$
\llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+\frac{M_{\mathrm{rel}}^{2}}{M_{\mathrm{obs}}^{2}}\left\|q_{T}^{\mathrm{opt}}\right\|_{s, \mathrm{D}}^{2} \leq\left\|y_{0}\right\|_{-s, \mathrm{D}}\left(M_{\mathrm{obs}}^{2} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+M_{\mathrm{rel}}^{2}\left\|q_{T}^{\mathrm{opt}}\right\|_{s, \mathrm{D}}^{2}\right)^{1 / 2}
$$

and thus

$$
\begin{equation*}
M_{\mathrm{obs}}^{2} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+M_{\mathrm{rel}}^{2}\left\|q_{T}^{\mathrm{opt}}\right\|_{s, \mathrm{D}}^{2} \leq M_{\mathrm{obs}}^{4}\left\|y_{0}\right\|_{-s, \mathrm{D}}^{2} \tag{69}
\end{equation*}
$$

Note that taking the norm $\|\cdot\|_{s, \mathrm{D}}$ in (67) gives

$$
\left\|q_{T}^{\mathrm{opt}}\right\|_{s, \mathrm{D}}^{2}=\frac{M_{\mathrm{obs}}^{4}}{M_{\mathrm{rel}}^{4}}\left\|\mathrm{P}_{F_{T}} y(T)\right\|_{-s, \mathrm{D}}^{2}
$$

so that with (69) we end up with

$$
M_{\mathrm{obs}}^{2} \llbracket v \rrbracket_{L^{2}(0, T ; U)}^{2}+\frac{M_{\mathrm{obs}}^{4}}{M_{\mathrm{rel}}^{2}}\left\|\mathrm{P}_{F_{T}} y(T)\right\|_{-s, \mathrm{D}}^{2} \leq M_{\mathrm{obs}}^{4}\left\|y_{0}\right\|_{-s, \mathrm{D}}^{2}
$$

and the claim is proved.

## REFERENCES

[1] D. Allonsius, F. Boyer, and M. Morancey, Spectral analysis of discrete elliptic operators and applications in control theory, Numerische Mathematik, to appear (2018), https://hal. archives-ouvertes.fr/hal-01422168.
[2] F. Ammar-Khodja, A. Benabdallah, M. González-Burgos, and L. Teresa, Recent results on the controllability of linear coupled parabolic problems: A survey, Mathematical Control and Related Fields, 1 (2011), pp. 267-306, doi:10.3934/mcrf.2011.1.267.
[3] F. Ammar Khodja, F. Chouly, and M. Duprez, Partial null controllability of parabolic linear systems, Math. Control Relat. Fields, 6 (2016), pp. 185-216, https://doi.org/10.3934/mcrf. 2016001.
[4] A. Benabdallah, F. Boyer, M. González-Burgos, and G. Olive, Sharp estimates of the one-dimensional boundary control cost for parabolic systems and application to the $\$ n \$-$ dimensional boundary null controllability in cylindrical domains, SIAM Journal on Control and Optimization, 52 (2014), pp. 2970-3001, doi:10.1137/130929680, http://dx.doi.org/10. 1137/130929680.
[5] K. Bhandari and F. Boyer, Boundary null-controllability of one-dimensional coupled parabolic system with robin conditions, in preparation, (2018).
[6] F. Boyer, On the penalised HUM approach and its applications to the numerical approximation of null-controls for parabolic problems, in CANUM 2012, 41e Congrès National d'Analyse Numérique, vol. 41 of ESAIM Proc., EDP Sci., Les Ulis, 2013, pp. 15-58.
[7] F. Boyer, F. Hubert, and J. L. Rousseau, Discrete carleman estimates for elliptic operators and uniform controllability of semi-discretized parabolic equations, Journal de Mathématiques Pures et Appliquées, 93 (2010), pp. 240-276.
[8] F. Boyer, F. Hubert, and J. L. Rousseau, Discrete carleman estimates for elliptic operators in arbitrary dimension and applications, SIAM Journal on Control and Optimization, 48 (2010), pp. 5357-5397.
[9] F. Boyer and J. L. Rousseau, Carleman estimates for semi-discrete parabolic operators and application to the controllability of semi-linear semi-discrete parabolic equations, Annales de l'Institut Henri Poincaré (C) Non Linear Analysis, 31 (2014), pp. 1035-1078.
[10] J.-M. Coron, Control and nonlinearity, vol. 136 of Mathematical Surveys and Monographs, American Mathematical Society, Providence, RI, 2007.
[11] S. Labbé and E. Trélat, Uniform controllability of semidiscrete approximations of parabolic control systems, Systems Control Lett., 55 (2006), pp. 597-609.
[12] J. Le Rousseau and G. Lebeau, On carleman estimates for elliptic and parabolic operators. applications to unique continuation and control of parabolic equations, ESAIM: COCV, 18 (2011), pp. 712-747, doi:10.1051/cocv/2011168.
[13] G. Lebeau and L. Robbiano, Contrôle exact de l'équation de la chaleur, Séminaire Équations aux dérivées partielles (Polytechnique), (1994-1995), pp. 1-11, http://eudml.org/doc/ 112118.
[14] L. Miller, A direct Lebeau-Robbiano strategy for the observability of heat-like semigroups, Discrete Contin. Dyn. Syst. Ser. B, 14 (2010), pp. 1465-1485, https://doi.org/10.3934/ dcdsb.2010.14.1465.
[15] M. Tucsnak and G. Weiss, Observation and control for operator semigroups, Birkhäuser Advanced Texts: Basler Lehrbücher. [Birkhäuser Advanced Texts: Basel Textbooks], Birkhäuser Verlag, Basel, 2009, doi:10.1007/978-3-7643-8994-9, http://dx.doi.org/10. 1007/978-3-7643-8994-9.
[16] E. Zuazua, Control and numerical approximation of the wave and heat equations, International Congress of Mathematicians, Madrid, Spain, III (2006), pp. 1389-1417.


[^0]:    *Aix-Marseille Université, CNRS, Centrale Marseille, I2M, UMR 7373, 13453 Marseille, France (damien.allonsius@univ-amu.fr, morgan.morancey@univ-amu.fr).
    ${ }^{\dagger}$ Institut de Mathématiques de Toulouse, UMR 5219, Université de Toulouse, CNRS, UPS IMT, F-31062 Toulouse Cedex 9, France (franck.boyer@math.univ-toulouse.fr).

