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HARPO, a gas TPC active target for high-performance γ -ray astronomy; demonstration of the polarimetry of MeV γ -rays converting to e^+e^- pair

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Abstract

No γ -ray polarimeter sensitive above 1 MeV has ever flown space. γ -ray polarimetry would be a new window on the radiative processes at work in cosmic sources, processes that produce linearly polarised emission, each of which with different polarisation fractions. The HARPO Collaboration has designed, built and characterised on beam a gas-TPC active target with which we have demonstrated for the first time the polarimetry of a linearly polarised MeV γ -ray beam, from the analysis of the conversions to e^+e^- pairs.

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Keywords: y-rays, polarimeter, TPC, pair production, telescope, gas detector

1. γ -ray polarimetry of cosmic sources: a science case

Today's γ -ray astronomy is polarisation-blind as no significant polarimetry of a cosmic source has been performed. Po-25 3 larimetry would provide the additional observables needed to 26 solve the pulsar (angle configuration, magnetic field, emission 27 location) parameter degeneracy. In particular in MeV pulsars, it 6 would enable to decipher the high-energy γ -ray emission mech- 28 7 anism (curvature radiation or synchrotron radiation) [1, 2]. Po-8 larimetry would enable to identify the nature of the emitting 29 9 particles in γ -ray emitting blazar jets (pure leptons or lepton-10 hadron mixture) [3]. Polarimetry could enable the indirect de-11 tection of dark matter [4]; the detection of Lorentz-invariance 12 violation due to "new" physics beyond the standard model, with 32 13 a potential larger than that of past searches in the X-ray band as 14 the sensitivity varies with the square of the photon energy [5]; ³³ 15 the discovery of the axion, i.e. of the pseudo-scalar believed to ³⁴ 16 be the pseudo-Goldstone boson associated to a breaking of the ³⁵ 17 U(1) symmetry devised to solve the QCD CP problem [6]. 18 36 37

¹⁹ 2. γ -ray polarimetry of cosmic sources: how achieve it ?

For the lowest part of the electromagnetic spectrum, polarimetry is performed by the measurement and the analysis

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of light intensities or of electric fields. At high energies, the (photo-electric, Compton scattering or conversion to pair) interaction of each single photon in a detector is recorded, and an estimate of the azimutal angle, ϕ , of the particles in the final-state is measured. The expression of the cross section, differential with respect to ϕ , is

$$d\sigma/d\phi \propto (1 + A \times P\cos(2(\phi - \phi_0))), \qquad (1)$$

where

- *A* is the polarisation asymmetry of the process at work. Its value depends on the photon energy.
- *P* is the linear polarisation fraction of the radiation.
- ϕ_0 is the angle between the polarisation direction of the radiation and the reference axis with respect to which the azimutal angle is measured.

The product $A \times P$ is sometimes named the modulation factor.

The polarimetry of a cosmic source in the MeV-GeV energy range has never been performed. For Compton polarimeters, both the total cross-section and the polarisation asymmetry decrease with energy (Fig. 2 of [7]). For nuclear pair-conversion, $\gamma Z \rightarrow e^+e^- Z$, the multiple scattering undergone by the electron and by the positron in the tracker blur the azimuthal-angle information carried by the pair within $\approx 10^{-3}$ radiation lengths [8, 9, 10].

Hopes have been put in the use of triplet conversions, that is, 45 of conversions to pair in the field of an electron of the detector, 46 $\gamma e^- \rightarrow e^+ e^- e^-$, as the target electron oftens recoils at a large 47 polar angle with respect to the direction of the incoming pho-48 ton, making it easier the measurement of its azimuthal angle. 49 Alas, the cross section is small and the fraction of it for high-50 enough recoil electron momentum is even much smaller (Fig. 6 51 of [11]), so the potential for a measurement on a cosmic source 52 ends up to be miserable (Sect. 5.3 of [11]). 53

A way to mitigate the multiple scattering issue is the use of 54 high-imaging-resolution active targets such as emulsion detec-55 tors [12]. Zooming in the first microns of the event just down-56 stream the conversion vertex, polarimetry with conversions to 57 pairs was demonstrated with a prototype detector in a GeV γ -58 ray beam [13] but the ability to take data in the MeV energy 59 range, where most of the statistics lies for cosmic sources (Fig. 60 2 of [11]), remains an issue. 61

The HARPO project explored an other technique, the use 62 of a low-density, homogeneous active target such as a high-63 pressure gas TPC (time projection chamber) [14]. We first of 64 wrote [11] a Monte-Carlo event generator¹ that samples the $_{06}$ 65 five-dimensional (5D) differential cross section at the first or- $_{\circ 7}$ 66 der of the Born development, i.e., the "Bethe-Heitler" differen-67 tial cross section (non polarised [19], polarised [20, 21, 22]). 68 With that tool we determined the (68 % containment) contribu-69 tion to the single-photon angular resolution due to the fact that 70 the momentum of the recoil nucleus cannot be measured, to be 71 ≈ 1.5 rad[E/1 MeV]^{-5/4} [23, 24]. We obtained the contribu-72 tion to the single-photon angular resolution due to the single-73 track angular resolution in the case of an optimal tracking à la 74 Kalman to be $\propto p^{-3/4}$ [23, 25]. 75 106

At high thicknesses, we found the dilution² of the polarisa-76 tion asymmetry due to multiple scattering to worsen less (Fig. 77 17 of [11]), with this full Monte-Carlo event generation, than 78 what was predicted by [8, 9, 10] who approximated the pair¹⁰⁸ 79 opening-angle, θ_{+-} , by its most-probable value. This is due to 80 the long tail of the θ_{+-} distribution. We demonstrated that the¹⁰⁹ 81 single-track angular resolution of gas detectors can be so good110 82 that polarimetry can be performed before the azimuthal infor-111 83 mation is lost [11]. Due to the $p^{-3/4}$ variation of the single-track¹¹² 84 angular resolution, and given the 1/E scaling of the θ_{+-} distri-¹¹³ 85 bution [26], a better dilution is within reach at lower energies¹¹⁴ 86 (Fig. 20 of [11]). The simulation shows that with a 1 m^3 , 5 bar^{115} 87 argon polarimeter observing a cosmic source such as the Crab¹¹⁶ 88 nebula for one year, full time, with full efficiency, the expected117 89 precision in the measurement of P is of $\approx 1.4 \%$ [11]. 118 90

We implemented the moments' method to obtain an optimal¹¹⁹ measurement of the polarisation modulation factor $A \times P$. In¹²⁰ the case of the simple 1D differential cross section, the method¹²¹ is equivalent to a maximum likelihood fit of the 1D differential¹²²



Figure 1: The two "maps", that is, the two *x*, *t* and *y*, *t* projections of a conversion event of a 3.93 MeV γ -ray provided by the NewSUBARU BL01 beam line converting to an e^+e^- pair in the 2.1 bar argon-isobutane (95-5%) gas of the HARPO TPC prototype [34]. The vertical dashed lines denote the physical limits of the detector, i.e. the cathode (left) and the anote (right).

cross section $d\sigma/d\phi \propto (1 + A \times P \cos(2(\phi - \phi_0)))$ [11, 27]. In the case of the full 5D differential cross section, an additional gain in the precision of the measurement of $A \times P$ of a factor of two-to-three is obtained [11, 27].

In the case of interest of a final state consisting of three particles (even when the recoil cannot be measured), the definition of the azimuthal angle, ϕ , is an issue. We determined the bisector $\phi_{+-} = (\phi_+ + \phi_-)/2$ of the azimuthal angles of the direction of the electron, ϕ_- , and of the positron, ϕ_+ , to be the optimal choice [27]. We demonstrated that with that angle definition, our generator is the only one, to our knowledge, to match the known asymptotic expressions for *A* at low and at high energies [24].

3. HARPO: the experiment

The HARPO Collaboration built a high-pressure (0.5 - 5 bar) gas TPC prototype [14] that uses an hybrid GEM-micromegas amplification system that we characterised precisely [28]. The anode plane is segmented into 2 orthogonal series (x, y) of strips, rather than into pads, so as to limit the number of electronic channels and therefore the power consumption, a key factor for a detector intended to fly a space mission. The ambiguity in the track assignment in the two (x, t and y, t) event projections, where *t* is the ionisation electron drift time, is easily solved by matching the charge time-distribution for each track (Fig. 6 of [29]).

The detector was exposed to a high-intensity γ -ray beam produced by the head-on inverse Compton scattering of a laser beam on the 0.6 - 1.5 GeV electron beam of the NewSUB-ARU storage ring [30]. The γ -ray beam was collimated on axis so as to select forward scattering and to obtain a quasimonochromatic beam with photon energies close to the Compton edge. After collimation, the linear polarisation of the laser beam is transferred almost perfectly to the γ -ray beam. By varying the laser wavelength and the electron beam energy we were able to take data between 1.7 and 74 MeV. We used a dedi-

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¹A C++ version of our event generator [15] has been deployed as the¹²⁵ G4BetheHeitler5DModel physics model of the 10.5beta release of Geant4_{126} [17, 18] (A fortran demonstrator is documented in [16]).

²The dilution factor *D* is the factor by which the polarisation asymmetry ¹²⁷ in an actual detector, A_{eff} , is decreased with respect to the theoretical value, ¹²⁸ A_{QED} , that is, $A_{eff} = D \times A_{QED}$. ¹²⁹



Figure 2: Distribution of the azimuthal angle of 11.8 MeV γ -rays (ratio of the₁₇₈ fully linearly polarised to the linearly non-polarised) converting to an e^+e^- pair₁₇₉ in the 2.1 bar argon-isobutane (95-5%) gas of the HARPO TPC prototype [33].₁₈₀

cated trigger system that enabled the data taking of γ -ray con-¹⁸³ versions in the TPC gas at rates of several 10 Hz, in the presence¹⁸⁴ of several 10 kHz of background noise, with a negligible dead¹⁸⁵ time [31]. We simulated the experiment with a Geant4-based¹⁸⁶ Monte-Carlo simulation, with a custom implementation of the¹⁸⁷ TPC-specific processes that we calibrated with respect to beam¹⁸⁹ data [32].

The strongly non-cylindrical-symmetric x, y, t structure of₁₉₁ 137 the detector induces a bias in the ϕ distribution: the distribu-192 138 tion angle of a non-polarised sources is not flat (Fig. 9 of [33]).193 139 This was analysed and mitigated by a two-step strategy: (a)194 140 chunks of data were taken with the detector rotated around the195 141 beam axis by multiples of 45°; (b) data were taken with ei-196 142 ther fully polarised (P = 1) or non polarised (P = 0) beam,¹⁹⁷ 143 so as to form the ratio of the ϕ distributions; On a space mis-144 sion, a non-polarised beam would not be available, and there-145 fore P = 0 simulated samples would be used to apply the bias₂₀₁ 146 correction. We demonstrated that this can be done without any₂₀₂ 147 additional systematic bias due to a possible difference between₂₀₃ 148 simulation and data, within uncertainty (Fig. 11 of [33]). The204 149 biases mentioned above would average out for an isotropic ex-205 150 posure anyway, as is the case for a long-duration data taking, so²⁰⁶ 151 the MC-based correction considered above would be a second-207 152 order correction (Fig. 12 of [33]). 153

The dilution factor is found to be of about 0.4 at 11.8 MeV^{209} and to decrease with photon energy, as expected (Fig. 13 of²¹⁰₂₁₁ [33]).

157 4. Results

We have demonstrated for the first time the polarimetry of a²¹⁶ 158 MeV γ -ray beam and with an excellent dilution factor [33] (Fig.²¹⁷ 159 2). In contrast to higher-density targets, gas detectors are able 160 to image the conversion to pairs of low-energy γ -rays (Fig. 1). 161 This is of utmost importance as the polarisation asymmetry of 162 conversions to pair does not decrease much on the energy range₂₂₂ 163 of a practical polarimeter [11, 27]: the precision of the measure-₂₂₃ 164 ment relies ultimately on the size of the collected statistics and224 165

therefore on the ability to collect data efficiently at low energies. An attempt to monitor the degradation of the TPC gas in a

sealed mode over half a year has shown that after the oxygen contamination (with an oxygen/nitrogen fraction found to be compatible with air) was removed by circulation through a purification cartridge, the nominal TPC fresh-gas parameters are restored [35], enabling a long-duration operation with the same gas on a space mission.

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