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Characterization of partial derivatives with respect to material parameters in a fluid-solid interaction problem.

Izar Azpiroz^a, Hélène Barucq^a, Rabia Djellouli^b, Ha Pham^{a,*}

^a*Inria Bordeaux Sud-Ouest Research Center, Project-Team Magique-3D, LMA – UMR CNRS 5142 – E2S, Université de Pau et des pays de l'Adour, France.*

^b*IRIS, Department of Mathematics, California State University, Northridge.*

Abstract

For a fluid-solid interaction problem with Lipschitz interface, we investigate the partial Fréchet differentiability of the solutions and the approximate far-field-pattern with respect to solid material parameters. Differentiability is shown in standard Sobolev framework, and the derivatives are characterized as solutions to inhomogeneous fluid-solid transmission problems. To validate the accuracy of the characterization, we compare analytical values with numerical ones given by Interior Penalty Discontinuous Galerkin (IPDG) in a setting with circular obstacles. Our comparisons also show that IPDG gives results with high precision and incurs almost no effect of discretization error accumulation.

Keywords: sensitivity with respect to Lamé parameters, partial Fréchet derivative, fluid-solid interaction, approximate far-field-pattern.

1. Introduction

In this work, for a fluid-solid interaction problem in \mathbb{R}^m , $m = 2, 3$, with Lipschitz fluid-solid interface, we investigate the partial Fréchet differentiability of the solutions and the approximate far-field-pattern (a-FFP) with respect to the solid material parameters. These include the Lamé parameters λ , μ and the solid density ρ_s . The original fluid-solid interaction denoted by \mathbf{OP}_∞ is posed in infinite space, *cf.* Figure 1(a). For numerical evaluation using volumed-based discretization, the infinite domain is truncated using absorbing boundary condition (ABC), and the problem we focus on, denoted by \mathbf{OP} , is defined on a finite convex domain Ω_{finite} ,

*Corresponding author

Email addresses: izar.azpiroz@inria.fr (Izar Azpiroz), helene.barucq@inria.fr (Hélène Barucq), rabia.djellouli@csun.edu (Rabia Djellouli), ha.howard@inria.fr (Ha Pham)

cf. Figure 1(b). Because of domain truncation, we work with the approximated far-field-pattern (a-FFP), cf. (36), instead of the true far-field-pattern (FFP), cf. (11). Although defined on a finite domain, our problem is in the category of transmission

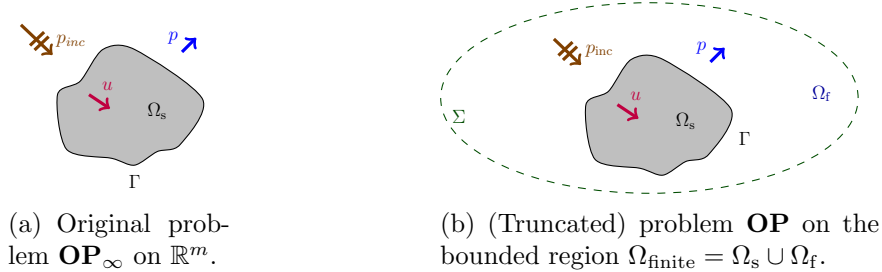


Figure 1: In fluid-solid interaction problems, an incident wave is diffracted by an immersed solid body, and generates scattered fluid pressure p and transmitted solid displacement u .

scattering with penetrable obstacles¹, to be distinguished with elastic problems on bounded domains such as [8] for elastostatics, [9] in time-harmonic elastodynamics, or [10, 11] in elastodynamics.

For problem \mathbf{OP} , we carry out the following tasks. We first extend the well-posedness results in [12] for a generic family of ABCs (*Task 1*). We then introduce the auxiliary problems and make sense of their PDE forms in the ‘standard’ Sobolev framework (see Notations in subsection 2.1), and study their well-posedness via variational formulation (*Task 2*). We next show the separate continuity and partial Fréchet differentiability of the solution operator and a-FFP (associated to \mathbf{OP}), and characterize their partial derivatives as solutions to the auxiliary problems (*Task 3*). This justifies their numerical evaluation by discretizing the corresponding variational formulations. The accuracy of this characterization is validated by an analytical-numerical comparison, using IPDG [13], for a setting with circular obstacle and circular artificial boundary (*Task 4*).

We limit ourselves to showing partial differentiability, which is sufficient for iterative inversion procedures which reconstruct separately each material parameter². What concerns us more is the Lipschitz interface. Note that most references assume at least \mathcal{C}^2 regularity, cf. [6, 2], with the exception of [5] for acoustic scattering with Lipschitz obstacles, and [14, 15, 16] for acoustic-elastic transmission (in fact \mathbf{OP}_∞) with polygonal-shaped obstacles, and the aforementioned [11, 8]

¹ In the category of penetrable obstacles, our problem studies acoustic-elastic transmission, to be distinguished with acoustic-acoustic cf. [1] and elastic-elastic cf. [2]. This is contrast with impenetrable obstacles, e.g. [3, 4, 5] in acoustics, and [6, 7] in elastics.

²Joint continuity and Fréchet differentiability can be obtained with slight modification of the current work.

for elasticity on bounded Lipschitz domain. Under Lipschitz assumption, the auxiliary problems can contain singular boundary or interface terms that do not fit in the canonical boundary Sobolev spaces, and to give sense to the problem is not trivial (*cf.* [14] and further discussion in subsection 2.3). Furthermore, for fluid-solid problems, existence of solutions is not always guaranteed due to the existence of solid resonance modes called Jones frequencies (defined in (26)).

The main feature of our work is that we make sense of the auxiliary problem in its PDE form within the standard Sobolev framework, then prove its equivalence to the variational formulation³. These justifications seem to be overlooked in literature, when one works directly with the variational forms, *e.g.* the aforementioned references [10, 11, 8], *cf.* Remark 5. Although [14, 15, 16] also study the auxiliary problems in PDE form, they use implicit function theorem to show differentiability, while we work directly with the solution operator and use standard analysis⁴. Other approaches to show Fréchet differentiability can be *e.g.* boundary integral operator, and factorization technique. Note that the literature investigating this topic is mostly concerned with domain derivatives⁵.

The rest of the paper contains two parts, mathematical analysis is in Section 2, and the numerical experiments in Section 3. The main results of *Task 1* are in Theorem 2 and Corollary 3, those of *Task 2* in Proposition 5 and Corollary 6, and those of *Task 3* in Proposition 8 and Corollary 9. The analysis for solid density ρ_s is straightforward and follows from the same analysis, with only results listed in subsection 2.5.

2. Mathematical analysis

2.1. Notations

Geometry. In Figure 1(b), we have introduced the solid obstacle Ω_s with Lipschitz boundary Γ , the finite convex domain Ω_{finite} with exterior boundary Σ , and the finite fluid region Ω_f . We have $\Omega_f = (\mathbb{R}^m \setminus \overline{\Omega_s}) \cap \Omega_{\text{finite}}$. Along any closed curve, the sign convention for the unit normal vector⁶ ν is outward.

³We work with the original PDE form which gives physical meaning to the solutions. Variational formulation is used for well-posedness and for numerical evaluation.

⁴*i.e.* by showing that the operators verify the definition of continuity and differentiability.

⁵We cite some representative references which study domain derivatives in obstacle scattering. Using *boundary integral operator* in acoustics *cf.* [17], elastics [2], electromagnetics [18, 19, 20]. Using *factorization technique* by [21], in elastics *cf.* [6], acoustics and electromagnetics [22]. Using *implicit function theorem* for acoustics *cf.* [23, 5], and acoustic-elastic transmission [14, 15, 16]. Using *variational methods* in acoustics *cf.* [3, 1], and elastics [7, 2]. The factorization technique is more adapted to inversion with Linear Sampling method, *cf.* [6, 24] for discussion.

⁶By definition, a Lipschitz boundary can be represented locally as a graph of a Lipschitz function. Such a boundary inherits a (surface) measure σ and a L^∞ unit normal vector ν which

Media. The fluid is assumed to be homogeneous compressible inviscid with constant density ρ_f and sound velocity c_f . The solid is supposed to be linearly isotropic homogeneous elastic of constant density ρ_s . This means its constitutive law relating the stress tensor $\boldsymbol{\sigma}$ to the strain tensor $\boldsymbol{\epsilon}(u)$,

$$\boldsymbol{\epsilon}(u) \triangleq \frac{\nabla u + (\nabla u)^t}{2},$$

is given by linear Hooke's law

$$\boldsymbol{\sigma}(u) = \mathbf{C} : \boldsymbol{\epsilon}(u),$$

with the elastic stiffness tensor \mathbf{C} invariant under rotations and reflections,

$$\mathbf{C}_{ijkl} = \lambda \delta_{ij} \delta_{kl} + \mu \delta_{il} \delta_{jk} \quad , \quad 1 \leq i, j, k, l \leq m.$$

Here λ and μ are the Lamé parameters. The material is in addition assumed to be pure elastic, *i.e.* $\rho_s, \lambda, \mu \in \mathbb{R}$, *cf.* [27]. We use the dot \cdot and double dot $:$ notation to denote single and double contraction, *cf.* [28]. For tensor $\mathbf{C} = (\mathbf{C}_{ijkl})$ of order 4, matrix $\mathbf{a} = (\mathbf{a}_{ij})$, and vector $u = (u_i)$, $(\mathbf{C} : \mathbf{a})_{ij} = \sum_{k,l} \mathbf{C}_{ijkl} \mathbf{a}_{kl}$; $(\mathbf{a} \cdot u)_i = \sum_j \mathbf{a}_{ij} u_j$.

Function spaces.

- For a function space \mathbb{F} , denote by \mathbb{F}' its dual, *i.e.* the set of bounded linear functionals from \mathbb{F} to \mathbb{C} .
- Following from the usual convention, $\mathcal{D}(\Omega)$ is the space of smooth functions with compact support in domain Ω .
- For $s \in \mathbb{R}$, $H^s(\mathbb{R}^m)$ is the usual Sobolev space on \mathbb{R}^m , *cf.* [25, p.76].
- We follow [25, p.77] to define the Sobolev space for a bounded domain $\Omega \subset \mathbb{R}^m$,

$$H^s(\Omega) \triangleq \{u \in \mathcal{D}'(\Omega) : u = U|_{\Omega} \text{ for some } U \in H^s(\mathbb{R}^m)\}, \quad s \in \mathbb{R}^+. \quad (1)$$

- For a Lipschitz boundary Λ , the Sobolev spaces $H^s(\Lambda)$ are well-defined for $|s| \leq 1$ *cf.* [25, p.98], having the property $H^{-s}(\Lambda) = (H^s(\Lambda))'$.
- We denote $\langle \cdot, \cdot \rangle_{s,\Lambda}$ as the duality pairing between $H^{-s}(\Lambda)$ and $H^s(\Lambda)$ with pivot $L^2(\Lambda)$, with convention

$$\langle f, \phi \rangle_{s,\Lambda} = (f, \phi)_{L^2(\Lambda)} = \int_{\Lambda} f \bar{\phi} \, ds.$$

exists σ -almost everywhere, [25, p96] or [26, Thm 2.7.1].

- Following [25, p.106], vector- or matrix-valued Sobolev spaces are written as,

$$\begin{aligned} H^s(\Omega)^m &= H^s(\Omega; \mathbb{C}^m) \quad , \quad H^s(\Omega)^{m \times m} = H^s(\Omega; \mathbb{C}^{m \times m}) \quad , \quad s \in \mathbb{R}^+ ; \\ H^s(\Lambda)^m &= H^s(\Lambda, \mathbb{C}^m) \quad , \quad |s| \leq 1 . \end{aligned}$$

- Our framework for well-posedness is the Gelfand triple $\mathbf{H} \subset \mathbf{U} \subset \mathbf{H}'$,

$$\mathbf{H} \triangleq H^1(\Omega_f) \times H^1(\Omega_s)^m \quad ; \quad \mathbf{U} \triangleq L^2(\Omega_f) \times L^2(\Omega_s)^m . \quad (2)$$

Trace operator. For $\frac{1}{2} < s \leq 1$, the trace operator $\gamma_0 : \mathcal{D}(\bar{\Omega}) \rightarrow \mathcal{D}(\partial\bar{\Omega})$, $u \mapsto \gamma_0 u = u|_{\partial\bar{\Omega}}$ has a unique extension to a bounded operator from $H^s(\Omega) \rightarrow H^{s-1/2}(\partial\bar{\Omega})$, with a continuous right inverse, see [25, Thm 3.37].

Conormal derivative. The conormal derivative operator $\frac{\partial}{\partial\nu} : \mathcal{D}(\bar{\Omega}) \rightarrow \mathcal{D}(\partial\bar{\Omega})$, $q \mapsto \frac{\partial q}{\partial\nu} = \gamma_0(\nabla q) \cdot \nu$ can be extended to a bounded map from $\{q \in H^1(\Omega) \mid \Delta q \in (H^1(\Omega))'\} \rightarrow H^{-1/2}(\partial\bar{\Omega})$, cf. [26, Thm 2.7.7 p69]. In particular, if $p \in H^1(\Omega_f)$ satisfies $\Delta p + \kappa^2 p = 0$ in $L^2(\Omega_f)$, its conormal derivative is well-defined with $\frac{\partial p}{\partial\nu} \in H^{-1/2}(\Gamma)$, and we have the integration-by-parts (IP) identity: for $q \in H^1(\Omega_f)$,

$$\int_{\Omega_f} -(\Delta p)\bar{q} \, dx = \int_{\Omega_f} (\nabla p \cdot \nabla \bar{q}) \, dx + \langle \frac{\partial p}{\partial\nu}, q \rangle_{1/2, \Gamma} - \langle \frac{\partial p}{\partial\nu}, q \rangle_{1/2, \Sigma} . \quad (3)$$

Normal trace. Following [29], the normal trace operator for matrices

$$\gamma_n : \mathcal{D}(\bar{\Omega})^{m \times m} \rightarrow \mathcal{D}(\partial\bar{\Omega})^m \quad , \quad A \mapsto A \cdot \nu$$

can be extended to a continuous map from

$$\mathcal{H}^0(\text{div}, \Omega) \triangleq \{\mathbf{v} \in L^2(\Omega)^{m \times m} \mid \text{div } \mathbf{v} \in L^2(\Omega)^m\} \quad \longrightarrow \quad H^{-1/2}(\partial\bar{\Omega}) \quad (4)$$

by using density arguments and the Green's identity,

$$\int_{\Omega} (\text{div } \mathbf{u}) \cdot v \, dx = - \int_{\Omega} \mathbf{u} : \nabla \mathbf{v} \, dx + \int_{\partial\bar{\Omega}} (\mathbf{u} \cdot \nu) \cdot v \, ds \quad , \quad \mathbf{u}, \mathbf{v} \in \mathcal{D}(\bar{\Omega})^{m \times m} . \quad (5)$$

2.2. Fluid-solid transmission problems

In this part, we realize *Task 1* and study the well-posedness for **OP**. We also provide a bound for the norm of the solution operator with an explicit bounding constant. This estimate will be used in a later subsection to show separate continuity and partial differentiability.

Original problem. When the configuration is excited by an exterior incident wave $p_{\text{inc}}(x)$ of frequency $\omega/2\pi$, cf. Figure 1(a), the scattered fluid pressure p and the

transmitted solid displacement u satisfy the following transmission problem in \mathbb{R}^m , with the fluid wavenumber $\kappa = \frac{\omega}{c_f}$,

$$(\mathbf{OP}_\infty) \quad \begin{cases} \operatorname{div} \boldsymbol{\sigma} + \omega^2 \rho_s u = 0 & \text{in } L^2(\Omega_s)^m & (6) \\ \Delta p + \kappa^2 p = 0 & \text{in } L^2(\Omega_f) & (7) \\ \boldsymbol{\sigma} \cdot \boldsymbol{\nu} + p \boldsymbol{\nu} = -p_{\text{inc}} \boldsymbol{\nu} & \text{in } H^{-1/2}(\Gamma)^m & (8) \\ \omega^2 \rho_f u \cdot \boldsymbol{\nu} - \frac{\partial p}{\partial \boldsymbol{\nu}} = \frac{\partial p_{\text{inc}}}{\partial \boldsymbol{\nu}} & \text{in } H^{-1/2}(\Gamma) & (9) \\ \lim_{r \rightarrow +\infty} r^{\frac{m-1}{2}} (\partial_r p - i \kappa p) = 0 & r = |x|. & (10) \end{cases}$$

Propagation in the fluid is governed by the Helmholtz equation (7), and in the solid by the time-harmonic elastodynamic equation (6). The interaction between the fluid and the solid is described by the kinematic interface condition (8), and the dynamic one (9), both defined along the interface Γ . To guarantee uniqueness of solutions, the outgoing Sommerfeld radiation condition (10) is imposed at infinity.

The scattered fluid pressure in \mathbf{OP}_∞ has the asymptotic behavior of an outgoing spherical wave, *cf.* [30, Thm 2.6],

$$p(x) = \frac{e^{i\kappa|x|}}{|x|^{(m-1)/2}} (p_\infty(\hat{x}) + \mathcal{O}(|x|^{-1})), \quad |x| \rightarrow \infty; \quad \hat{x} = \frac{x}{|x|}. \quad (11)$$

In the above expression, p_∞ is called the far-field pattern and satisfies the identity,

$$p_\infty(\hat{x}) = \mathbf{c}(m) \int_{\tilde{\Gamma}} \frac{\partial e^{-i\kappa \hat{x} \cdot y}}{\partial \boldsymbol{\nu}(y)} (\gamma_0 q) \, ds(y) - \mathbf{c}(m) \left\langle \frac{\partial q}{\partial \boldsymbol{\nu}}, e^{i\kappa \hat{x} \cdot y} \right\rangle_{1/2, \tilde{\Gamma}}, \quad (12)$$

for any simple (at least Lipschitz) closed curve $\tilde{\Gamma}$ enclosing Ω_s , *i.e.* $\tilde{\Gamma} \subset \overline{\mathbb{R}^m} \setminus \Omega_s$, and with the defined constant

$$\mathbf{c}(m) \triangleq \begin{cases} \frac{e^{i\pi/4}}{8\pi\kappa} & , m = 2 \\ \frac{1}{4\pi} & , m = 3 \end{cases}. \quad (13)$$

Truncated problem. For numerical computations, the region exterior to the solid is truncated to the finite domain Ω_{finite} , *cf.* Figure 1(b). The radiation condition (10) is replaced by absorbing boundary conditions (ABC) along Σ . We work with those of Robin type $\frac{\partial p}{\partial \boldsymbol{\nu}} = -\mathcal{B} \gamma_0 p$ in $H^{-1/2}(\Sigma)$, where operator \mathcal{B} satisfies assumptions,

$$\begin{aligned} \mathcal{B} : H^{1/2}(\Sigma) &\rightarrow H^{-1/2}(\Sigma) \quad \text{is bounded;} \\ \operatorname{Re} \langle \mathcal{B} \gamma, \gamma \rangle_{1/2, \Sigma} &\geq 0, \quad \operatorname{Im} \langle \mathcal{B} \gamma, \gamma \rangle_{1/2, \Sigma} < 0, \quad \text{for all } \gamma \neq 0 \text{ in } H^{1/2}(\Sigma). \end{aligned} \quad (14)$$

The resulting problem is denoted by $\mathbf{OP}_{\mathcal{B}}$ (or simply \mathbf{OP} for generic operator \mathcal{B}),

$$(\mathbf{OP}) \quad \begin{cases} \operatorname{div} \boldsymbol{\sigma} + \omega^2 \rho_s u = 0 & \text{in } L^2(\Omega_s)^m & (15) \\ \Delta p + \kappa^2 p = 0 & \text{in } L^2(\Omega_f) & (16) \\ \boldsymbol{\sigma} \cdot \boldsymbol{\nu} + p \boldsymbol{\nu} = -p_{\text{inc}} \boldsymbol{\nu} & \text{in } H^{-1/2}(\Gamma)^m & (17) \\ \omega^2 \rho_f u \cdot \boldsymbol{\nu} - \frac{\partial p}{\partial \boldsymbol{\nu}} = \frac{\partial p_{\text{inc}}}{\partial \boldsymbol{\nu}} & \text{in } H^{-1/2}(\Gamma) & (18) \\ \frac{\partial p}{\partial \boldsymbol{\nu}} = -\mathcal{B} p & \text{in } H^{-1/2}(\Sigma). & (19) \end{cases}$$

Assumptions (14) are needed for well-posedness, see discussion below. For the problem \mathbf{OP} to be a good approximation of the original one \mathbf{OP}_{∞} , operator \mathcal{B} is chosen to be the exterior Dirichlet-to-Neumann map (D-t-N) \mathbf{T} , see [31, Section 3.1.2 p.63] for definition, or a good approximation of this.

- When Σ is a circle, we can take \mathcal{B} as \mathbf{T} , as done in [12]. With this choice, the ABC is exact, *i.e.* the original problem \mathbf{OP}_{∞} is equivalent to \mathbf{OP} , *cf.* [12, Thm 1]. The resulting problem is referred to as $\mathbf{OP}_{\mathbf{T}}$.
- Under the general assumption that Σ is convex, we use the classical zero-th and first order absorbing boundary condition⁷, in which \mathcal{B} equals to operators

$$\mathbf{T}_0 \stackrel{\Delta}{=} -i\kappa \quad , \quad \mathbf{T}_1 \stackrel{\Delta}{=} -i\kappa + \frac{\varkappa_{\Sigma}}{2} \quad , \quad (20)$$

respectively. The resulting problems are denoted by $\mathbf{OP}_{\mathbf{T}_0}$ and $\mathbf{OP}_{\mathbf{T}_1}$. Here, \varkappa_{Σ} is the curvature of the curve Σ when $m = 2$, and $\frac{\varkappa_{\Sigma}}{2}$ is the mean curvature of the surface Σ when $m = 3$.

Lemma 1. *The following choices of \mathcal{B} satisfy properties (14): $\mathcal{B} = \mathbf{T}$ when Σ is spherical (or circular), $\mathcal{B} = \mathbf{T}_0$ for general Σ , or $\mathcal{B} = \mathbf{T}_1$ when Σ is convex with bounded positive curvature.*

Proof. We have $\mathbf{T} : H^{1/2}(\Sigma) \rightarrow H^{-1/2}(\Gamma)$ is bounded with a bounded inverse⁸, and satisfies (14), *cf.* [12, Thm 1 and Lem 3] or [31, Section 3.2] for $m = 3$, and [34, Thm 5.22] for $m=2$.

⁷For the derivation of these conditions, see *e.g.* [32, Eqn 21] for $m = 3$ and [33, Eqn 1.31] for $m = 2$. These conditions are also listed in [31, Section 3.3.3] for circular Σ ; in this case $\varkappa_{\Sigma} = R^{-1}$, with R the radius of Σ .

⁸The sign convention for \mathbf{T} in [34] is different from [12, 31]. In either convention, the properties of \mathbf{T} are shown by using the multipole expansions of solutions to the Helmholtz equation. For the boundedness of \mathbf{T} on general curves or surfaces, see [30, Thm 3.13] for \mathcal{C}^2 boundaries, and [35, Thm 1.4] for nontrapping polygons and star-shaped Lipschitz domains.

The boundedness for \mathbf{T}_0 follows immediately from the fact that it is a multiple of the identity. That for \mathbf{T}_1 follows additionally from the boundedness of the curvature of Σ . We next verify the second property in (14). For $\phi \in H^{1/2}(\Sigma)$, we have

$$\langle \mathcal{B}\phi, \phi \rangle_{1/2, \Sigma} = \int (\mathcal{B}\phi) \bar{\phi} \, dx = \begin{cases} -i \|\phi\|_{L^2(\Sigma)}^2 & , \quad \mathcal{B} = \mathbf{T}_0 \\ -i \|\phi\|_{L^2(\Sigma)}^2 + \int_{\Sigma} \frac{\varkappa_{\Sigma}}{2} |\phi|^2 \, ds & , \quad \mathcal{B} = \mathbf{T}_1 \end{cases}.$$

Under the assumption $\varkappa_{\Sigma} > 0$, we have $\operatorname{Re} \langle \mathcal{B}\phi, \phi \rangle_{1/2, \Gamma} \geq 0$ and $\operatorname{Im} \langle \mathcal{B}\phi, \phi \rangle_{1/2, \Gamma} \leq 0$. In addition, $\operatorname{Im} \langle \mathcal{B}\phi, \phi \rangle_{1/2, \Gamma} = 0$ implies $\phi = 0$. \square

Variational formulation. We follow [12] to introduce the following sesquilinear forms. The volume sesquilinear forms corresponding to the (reduced) elastodynamic equation in Ω_s and the Helmholtz equation in Ω_f are

$$\mathbf{a}_s(u, v) \triangleq \int_{\Omega_s} \boldsymbol{\sigma}(u) : \nabla \bar{v} \, dx - \int_{\Omega_s} \omega^2 \rho_s u \bar{v} \, dx ; \quad u, v \in H^1(\Omega_s)^m, \quad (21)$$

$$\mathbf{a}_f(p, q) \triangleq \int_{\Omega_f} \nabla p \cdot \nabla \bar{q} \, dx - \int_{\Omega_f} \kappa^2 p \bar{q} \, dx ; \quad p, q \in H^1(\Omega_f). \quad (22)$$

The boundary sesquilinear forms corresponding to the transmission conditions on the solid-fluid interface Σ are

$$\begin{aligned} \mathbf{b} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] &\triangleq \langle u \cdot \nu, \gamma_0 q \rangle_{1/2, \Gamma} ; \quad \begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H} ; \\ \mathbf{b}^* \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] &\triangleq \overline{\mathbf{b} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right]} ; \quad \begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}. \end{aligned} \quad (23)$$

Note that $\mathbf{b}^* \left[\begin{pmatrix} q \\ v \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix} \right] = \overline{\langle v \cdot \nu, p \rangle_{1/2, \Gamma}} = \langle p, v \cdot \nu \rangle_{1/2, \Gamma} = \langle p \nu, v \rangle_{1/2, \Gamma}$. The sesquilinear form associated to the absorbing boundary condition along Σ is,

$$\mathbf{t}(p, q) \triangleq \langle \mathcal{B} \gamma_0 p, \gamma_0 q \rangle_{1/2, \Gamma}. \quad (24)$$

The main sesquilinear form in all of our variational problems is

$$\begin{aligned} \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] &\triangleq \mathbf{a}_s(u, v) + \mathbf{b}^* \left[\begin{pmatrix} q \\ v \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix} \right] \\ &+ \frac{1}{\omega^2 \rho_f} \mathbf{a}_f(p, q) + \mathbf{b} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] + \frac{1}{\omega^2 \rho_f} \mathbf{t}(p, q) ; \quad \begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}. \end{aligned} \quad (25)$$

Denote by $\mathcal{A} = \mathcal{A}(\lambda, \mu, \rho_s)$ the operator associated to \mathbf{a} at (λ, μ, ρ_s) .

Well-posedness discussion. The well-posedness of \mathbf{OP}_{∞} and $\mathbf{OP}_{\mathbf{T}}$ in the framework $\mathbf{H} \subset \mathbf{U} \subset \mathbf{H}'$ is addressed in [12]. In this case ($\mathcal{B} = \mathbf{T}$), the sesquilinear form

\mathbf{a} is shown to satisfy the Gårding's inequality on \mathbf{H} (*i.e.* \mathbf{H} -coercive with respect to \mathbf{U}), *cf.* [12, Proposition 6],

$$\operatorname{Re} \mathbf{a} \left[\begin{pmatrix} q \\ v \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] + c \left\| \begin{pmatrix} q \\ v \end{pmatrix} \right\|_{\mathbf{U}}^2 \geq \alpha \left\| \begin{pmatrix} q \\ v \end{pmatrix} \right\|_{\mathbf{H}}^2, \quad \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}.$$

This allows applying the Fredholm principle, *cf.* [26, Thm 2.1.60] to study the well-posedness of the variational problem equivalent to $\mathbf{OP}_{\mathbf{T}}$. Even when existence (of solutions) is given, uniqueness is not guaranteed due to Jones modes defined as

$$\omega \text{ is called a Jones frequency and } \begin{cases} \operatorname{div} \boldsymbol{\sigma} + \omega \rho_s u = 0 \text{ in } \Omega_s \\ \boldsymbol{\sigma} \cdot \boldsymbol{\nu} = 0; u \cdot \boldsymbol{\nu} = 0 \text{ on } \Gamma \end{cases} \text{ if } (u, \omega) \text{ solves } \quad (26)$$

We refer to [36] and the references therein for a detailed discussion of Jones modes.

We now extend the result of [12, Thm 7] to arbitrary operator \mathcal{B} satisfying (14), and obtain the well-posedness for $\mathbf{OP}_{\mathcal{B}}$. The proof follows [12] with minor modification. The completely new feature is an energy estimate with an explicit bounding constant⁹. For the rest of the subsection, notations are simplified and reflect only the dependence with respect to λ and μ . In particular, we write $\mathcal{A} = \mathcal{A}(\lambda, \mu)$ instead of $\mathcal{A}(\lambda, \mu, \rho_s)$.

Theorem 2. *Consider the sesquilinear form \mathbf{a} defined in (25) with operator \mathcal{B} satisfying (14), and the material parameters satisfying*

$$\lambda > 0, \quad \mu \geq 0, \quad \omega > 0, \quad \rho_f > 0, \quad \rho_s > 0. \quad (27)$$

For arbitrary functional $\mathbf{l} \in \mathbf{H}'$, the generic variational problem

$$\text{Find } \begin{pmatrix} p \\ u \end{pmatrix} \in \mathbf{H} \text{ so that } \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] = \mathbf{l} \left[\begin{pmatrix} q \\ v \end{pmatrix} \right], \quad \forall \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H} \quad (28)$$

has the following dichotomy.

1. *If ω is not a Jones frequency, there exists a unique solution for (28). In addition, the solution depends continuously on the right hand side, described by the energy bound: for a fixed ρ_s , and in a small enough neighborhood of a pair (λ, μ) with h satisfying (32),*

$$\begin{aligned} \left\| \begin{pmatrix} p_{\lambda+h, \mu} \\ u_{\lambda+h, \mu} \end{pmatrix} \right\|_{\mathbf{H}} &\leq 2 \left\| \mathcal{A}(\lambda, \mu)^{-1} \right\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})} \left\| \mathbf{l} \right\|_{\mathbf{H}'}, \\ \left\| \begin{pmatrix} p_{\lambda, \mu+h} \\ u_{\lambda, \mu+h} \end{pmatrix} \right\|_{\mathbf{H}} &\leq 2 \left\| \mathcal{A}(\lambda, \mu)^{-1} \right\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})} \left\| \mathbf{l} \right\|_{\mathbf{H}'}. \end{aligned} \quad (29)$$

⁹Estimate (60) in [12, Thm 7] needs some clarification. It should hold with a different bounding constant (not 1 as stated) and under the assumption that ω is not a Jones frequency.

2. If ω is a Jones frequency, then there exists a solution to (28) if and only if functional $\mathbf{1}$ satisfies the compatibility condition

$$\mathbf{I}\left[\begin{pmatrix} 0 \\ v \end{pmatrix}\right] = 0 \quad , \quad \forall \text{ Jones mode } v. \quad (30)$$

If (30) is satisfied, a solution $(p, u) \in \mathbf{H}$ to (28) exists, with the fluid part p being unique, and the solid displacement u being unique modulo Jones modes.

Remark 1. The assumption on the Lamé parameters is needed for the $H^1(\Omega_s)^m$ -coercivity of \mathbf{a}_s , see [12, Eqn (49)–(50)], see also [25, Thm 10.3].

Proof. Well-posedness statement: The proof is identical to that for [12, Thm 7] with the following modifications to adapt to a general \mathcal{B} .

- In the proof to show that \mathbf{a} satisfies the Gårding inequality, Eqn (57) in [12] is replaced by the property $\operatorname{Re} \langle \mathcal{B}\gamma, \gamma \rangle_{1/2, \Sigma} \geq 0$.
- In the investigation of the uniqueness of the pressure fluid, we consider $(p, u) \in \mathbf{H}$ such that $\mathbf{a}\left(\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix}\right) = 0$. This means that $(p, u) \in \mathbf{H}$ satisfies (15)–(19) with $p_{\text{inc}} = 0$. Eqn (63) in [12] is replaced with

$$\operatorname{Im} \mathbf{a}\left(\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix}\right) = \operatorname{Im} \langle \mathcal{B}\gamma_0 p, \gamma_0 p \rangle_{1/2, \Sigma} = 0.$$

From here, we use the third property in (14),

$$\operatorname{Im} \langle \mathcal{B}\gamma, \gamma \rangle_{1/2, \Sigma} < 0 \quad , \quad \text{for all } \gamma \in H^{1/2}(\Sigma), \gamma \neq 0,$$

to imply that $\gamma_0 p = 0$. Combining with the boundary condition (18), we obtain $\frac{\partial p}{\partial \nu} = -\mathcal{B}\gamma_0 q = 0$. The rest of the proof uses analytic continuation principle and follows as in [12, Section 4.2.1] to show that $p = 0$.

Energy estimate when ω is not a Jones frequency: At (λ, μ) , $(\lambda + h, \mu)$ and $(\lambda, \mu + h)$ with $h > 0$, the inverse of \mathcal{A} exists and is bounded, *i.e.*

$$\mathcal{A}(\lambda, \mu)^{-1} \quad , \quad \mathcal{A}(\lambda + h, \mu)^{-1} \quad , \quad \mathcal{A}(\lambda, \mu + h)^{-1} \in \mathcal{L}(\mathbf{H}', \mathbf{H}).$$

We will bound the norm of the last two inverses by that of $\mathcal{A}(\lambda, \mu)^{-1}$. Since the proof for $\mu + h$ is exactly the same, only that for $\lambda + h$ is written.

The map $(\lambda, \mu) \mapsto \mathcal{A}$ is affine separately in λ and μ , the only term that depends on λ or μ is \mathbf{a}_s . As a result, its partial derivative with respect to λ is independent of λ , and we denote this by $\mathcal{D}_1(\mu) \in \mathcal{L}(\mathbf{H}, \mathbf{H}')$. We can write

$$\begin{aligned} \mathcal{A}(\lambda + h, \mu) &= \mathcal{A}(\lambda, \mu) + h \mathcal{D}_1(\mu) \\ &= \mathcal{A}(\lambda, \mu)_{\mathbf{H} \rightarrow \mathbf{H}'} \left(\operatorname{Id}_{\mathbf{H} \rightarrow \mathbf{H}} - h \mathcal{A}(\lambda, \mu)_{\mathbf{H}' \rightarrow \mathbf{H}}^{-1} \mathcal{D}_1(\mu)_{\mathbf{H} \rightarrow \mathbf{H}'} \right) \\ \Rightarrow \mathcal{A}(\lambda + h, \mu)^{-1} &= \left(\operatorname{Id} - h \mathcal{K}(\lambda, \mu) \right)^{-1} \mathcal{A}(\lambda, \mu)^{-1}. \end{aligned}$$

Here, we have defined the bounded operator $\mathcal{K}(\lambda, \mu) = \mathcal{A}(\lambda, \mu)^{-1} \mathcal{D}_1(\mu)$. When $0 < h < \frac{1}{\mathfrak{c}(\lambda, \mu)}$, the inverse $(\text{Id} + h \mathcal{K}(\lambda, \mu))^{-1}$ can be written as an operator-valued Neumann series, and from this, we obtain the estimate

$$\begin{aligned} \|(\text{Id} + h \mathcal{K}(\lambda, \mu))^{-1}\|_{\mathcal{L}(\mathbf{H})} &\leq \frac{1}{1 - h \mathfrak{c}(\lambda, \mu)}, \\ \text{with } \mathfrak{c}(\lambda, \mu) &\triangleq \|\mathcal{A}(\lambda, \mu)^{-1}\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})} \|\mathcal{D}_1(\mu)\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})}. \end{aligned}$$

This is next used to bound,

$$\|\mathcal{A}(\lambda + h, \mu)^{-1}\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})} \leq \frac{\|\mathcal{A}(\lambda, \mu)^{-1}\|}{1 - h \mathfrak{c}(\lambda, \mu)}. \quad (31)$$

One has the same estimate for $\mathcal{A}(\lambda, \mu + h)^{-1}$, with \mathfrak{c} in (31) replaced by

$$\mathfrak{d}(\lambda, \mu) \triangleq \|\mathcal{A}(\lambda, \mu)^{-1}\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})} \|\mathcal{D}_2(\lambda)\|_{\mathcal{L}(\mathbf{H}', \mathbf{H})}.$$

To obtain the announced estimates (29), we impose further that

$$h \mathfrak{c}(\lambda, \mu) < \frac{1}{2}, \quad h \mathfrak{d}(\lambda, \mu) < \frac{1}{2} \quad \Rightarrow \quad \frac{1}{1 - h \mathfrak{c}(\lambda, \mu)} < 2, \quad \frac{1}{1 - h \mathfrak{d}(\lambda, \mu)} < 2. \quad (32)$$

□

Remark 2. *The compatibility condition (30) will be satisfied automatically, if $\mathbf{1}$ contains only boundary operators, which is the case for **OP**.*

Following the same reasoning as in [12, Thm 2], the transmission problem **OP** (15)–(19) is equivalent to the variational problem **VOP**,

$$\begin{aligned} \text{Find } \begin{pmatrix} p \\ u \end{pmatrix} \in \mathbf{H} : \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] &= \mathbf{l}_{\text{OP}} \left[\begin{pmatrix} \bar{q} \\ \bar{v} \end{pmatrix} \right], \quad \forall \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}, \\ \text{with } \mathbf{l}_{\text{OP}} \in \mathbf{H}' , \mathbf{l}_{\text{OP}} \left[\begin{pmatrix} q \\ v \end{pmatrix} \right] &\triangleq \frac{1}{\omega^2 \rho_f} \langle \frac{\partial p_{\text{inc}}}{\partial \nu}, \overline{\gamma_0 q} \rangle_{1/2, \Gamma} - \langle p_{\text{inc}}, \overline{\gamma_0 v} \rangle_{1/2, \Gamma}, \quad \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}. \end{aligned} \quad (33)$$

By Remark 2, functional \mathbf{l}_{OP} satisfies the compatibility condition (30), see also [12, Thm 7]. As a result, existence of solutions for problem **VOP** is guaranteed in all cases. We summarize this result in the corollary below.

Corollary 3. *With \mathcal{B} satisfying (14) and the material parameters satisfying (27), there exists a solution $(p, u) \in \mathbf{H}$ to problem **VOP**, with p unique. Additionally,*

- *if ω is a Jones frequency, the solid displacement u is unique modulo Jones modes;*
- *if not, such (p, u) is unique in \mathbf{H} , and in a neighborhood of (λ, μ) with h satisfying (32), we have*

$$\begin{aligned} \|(p_{\lambda+h, \mu}, u_{\lambda+h, \mu})\|_{\mathbf{H}} &\leq c(\lambda, \mu) \|p_{\text{inc}}\|_{H^2(\Omega_f)}, \\ \|(p_{\lambda, \mu+h}, u_{\lambda, \mu+h})\|_{\mathbf{H}} &\leq c(\lambda, \mu) \|p_{\text{inc}}\|_{H^2(\Omega_f)}, \end{aligned} \quad (34)$$

with a constant $c(\lambda, \mu) > 0$ independent of h .

In the above estimates, we have used,

$$\|\mathbf{l}_{\mathbf{OP}}\|_{\mathbf{H}'} = \hat{C}(\|\frac{\partial p_{\text{inc}}}{\partial \nu}\|_{H^{-1/2}(\Gamma)} + \|p_{\text{inc}}\|_{H^{-1/2}(\Gamma)}) \leq \tilde{C}\|p_{\text{inc}}\|_{H^2(\Omega_f)}.$$

As a result of Corollary 3, when ω is not a Jones frequency, we can define the solution operator

$$\begin{aligned} \mathcal{S} : \mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+ &\rightarrow \mathbf{H}, \quad (\lambda, \mu, \rho_s) \mapsto \mathcal{S}(\lambda, \mu, \rho_s) = (p_{\lambda, \mu, \rho_s}, u_{\lambda, \mu, \rho_s}), \\ &\text{with } (p_{\lambda, \mu, \rho_s}, u_{\lambda, \mu, \rho_s}) \text{ the unique solution in } \mathbf{H} \text{ to (15)–(19)}. \end{aligned} \quad (35)$$

Approximate far-field-pattern. Although \mathbf{OP} is defined on a finite domain, by using (12), we can generalize the notion of far-field-pattern to this problem. From (12), we define the integral operator \mathbf{I} : for $q \in H^1(\Omega_f)$ with $\frac{\partial q}{\partial \nu} \in H^{-1/2}(\Gamma)$,

$$(\mathbf{I}q)(\hat{x}) \triangleq \mathbf{c}(m) \int_{\Gamma} \frac{\partial e^{-i\kappa \hat{x} \cdot y}}{\partial \nu(y)} (\gamma_0 q) \, ds(y) - \mathbf{c}(m) \langle \frac{\partial q}{\partial \nu}, e^{i\kappa \hat{x} \cdot y} \rangle_{1/2, \Gamma}, \quad \hat{x} = \frac{x}{|x|}. \quad (36)$$

Recall that $\mathbf{c}(m)$ is defined in (13). The above definition exists in all cases, due to the uniqueness of the pressure field. Next, using the transmission conditions (17) and (18), we define the approximate FFP operator associated to \mathbf{OP} as follows,

$$\begin{aligned} \mathbf{F} : \mathbb{R}^+ \times \mathbb{R}^+ \times \mathbb{R}^+ &\longrightarrow \mathcal{C}^\infty(\mathbb{S}^1, \mathbb{C}), \quad (\lambda, \mu, \rho_s) \longrightarrow \mathbf{F}(\lambda, \mu, \rho_s) \\ (\mathbf{F}(\lambda, \mu, \rho_s))(\hat{x}) &= \mathbf{c}(m) \int_{\Gamma} \left(\frac{\partial e^{-i\kappa \hat{x} \cdot y}}{\partial \nu(y)} p - (\omega \rho_f^2 u \cdot \nu - \frac{\partial p_{\text{inc}}}{\partial \nu}) e^{-i\kappa \hat{x} \cdot y} \right) ds(y), \quad (37) \\ &\text{with } (p_{\lambda, \mu, \rho_s}, u_{\lambda, \mu, \rho_s}) = \mathcal{S}(\lambda, \mu, \rho_s). \end{aligned}$$

Remark 3. *This approach to compute approximate far-field-patterns from near fields is also employed in [37]. It is more common to define \mathbf{F} on a smooth curve $\tilde{\Gamma}$ outside of Ω_s and not on the boundary of the obstacle. However, using the latter provides some computation advantages, e.g. the possibility to use (37), which involves only the zero-th order traces, and avoids having to calculate the conormal derivative of p on a new curve.*

2.3. Auxiliary problem for partial derivatives with Lamé parameters

We now introduce the auxiliary problems for Lamé parameters λ and μ . In general, the auxiliary problems take the same form as the original problem, however they are inhomogeneous. They contain volume and/or interface (boundary) sources for penetrable (impenetrable) obstacles, all defined in terms of the solution to the original problem and its (possibly higher order) traces¹⁰. Difficulty arises

¹⁰ For examples of the auxiliary problems studying domain derivatives of FFP in obstacle scattering, see [6, Thm 3.5 and 4.4] for elastic scattering, see [2, Thm 3.1] for elastic-elastic transmission, and see [3, Thm 2.1] for acoustic scattering.

under low regularity of interfaces/boundaries, since interface (boundary) sources might not have well-defined traces in the canonical boundary spaces $H^s(\Gamma)$, $|s| \leq 1$.

In our case, the auxilliary problem for the solid density contains only a volume source, while that for the Lamé parameters has additionally a singular interface source in the form of the normal trace of $\operatorname{div} A$ with $A \in L^2(\Omega_s)^{m \times m}$, see also Remark 4. As such, we risk running into the same difficulty as [14] which studies domain derivatives for \mathbf{OP}_∞ . What differs in our problem is that the interface source is also the normal trace of the volume one, leading to cancellation in the more regular variational form, in which the ‘troublesome’ normal trace of $\operatorname{div} A$ no longer appears. This cancellation of singularities does not occur in the more difficult problem considered in [14] which only has interface sources¹¹.

Remark 4. *In our case, the above mentioned quantity A takes the form $\mathbf{C}' : \epsilon(u)$ where u is the solid displacement from \mathbf{OP} . When Γ is more regular, so is A , at least in $H^1(\Omega_s)$, and its normal trace exists in $H^{1/2}(\Gamma)$ (using the mapping property of the trace operator γ_0 [25, Thm 3.38]). When Γ is only Lipschitz, u is in at most $H^{3/2}(\Omega_s)^m$, cf. [12, Remark 2], and A is in at most $H^{1/2}(\Omega_s)^{m \times m}$.*

As mentioned in the Introduction, instead of passing directly to the (more regular) variational formulation of the auxiliary problem, we first make sense of its PDE form, see also Remark 5. The main idea is to group together terms which do not *a priori* have well-defined normal traces when considered separately, but do when grouped together. We will need the following lemma which generalizes the normal trace of a matrix-valued distribution.

Lemma 4. *If a matrix $A \in L^2(\Omega_s)^{m \times m}$ and $u \in H^1(\Omega_s)^m$ satisfy*

$$\operatorname{div} \boldsymbol{\sigma}(u) + \omega^2 \rho_s u = -\operatorname{div} A \text{ in } (H^1(\Omega_s)^m)', \quad (38)$$

then $(\boldsymbol{\sigma} + A) \in L^2(\Omega_s)^{m \times m}$, and its normal trace is well-defined with $(\boldsymbol{\sigma} + A) \cdot \nu \in H^{-1/2}(\Gamma)^m$. In addition, one has the generalized IP identity: for $v \in H^1(\Omega_s)^m$,

$$-\int_{\Omega_s} [\operatorname{div}(\boldsymbol{\sigma} + A)] \cdot \bar{v} \, dx = \int_{\Omega_s} (\boldsymbol{\sigma} + A) \cdot \nabla \bar{v} \, dx - \langle (\boldsymbol{\sigma} + A) \cdot \nu, v \rangle_{1/2, \Gamma}. \quad (39)$$

Proof. We rearrange (38) and obtain $\operatorname{div}(\boldsymbol{\sigma} + A) = -\omega^2 \rho_s u$ in $(H^1(\Omega_s)^m)'$. Since $u \in H^1(\Omega_s)^m$, $\operatorname{div}(\boldsymbol{\sigma} + A) \in H^1(\Omega_s)^m$ thus in $L^2(\Omega_s)^m$. On the other hand, $u \in H^1(\Omega_s)^m$, hence $\boldsymbol{\sigma}(u) \in L^2(\Omega_s)^{m \times m}$ and thus $\boldsymbol{\sigma}(u) + A \in L^2(\Omega_s)^{m \times m}$. This means $\boldsymbol{\sigma} + A \in \mathcal{H}^0(\operatorname{div}, \Omega_s)$ defined in (4). By the discussion in subsection 2.1, its normal trace is well-defined in $H^{-1/2}(\Gamma)^m$, and the integration-by-parts formula can be generalized from the Green’s identity (5). \square

¹¹In [14], the regularity is increased to polygonal-shaped obstacles, and the Sobolev framework had to be enlarged, in particular, the exponent $-\frac{3}{2}$ was considered. The well-posedness of the resulting auxilliary problem remains an open question.

From the above discussion, for $A \in L^2(\Omega_s)^{m \times m}$ and $(p, u) \in \mathbf{H}$, we can define

$$\begin{cases}
 \operatorname{div} \boldsymbol{\sigma} + \omega^2 \rho_s u = -\operatorname{div} A & \text{in } (H^1(\Omega_s)^m)' & (40) \\
 \Delta p + \kappa^2 p = 0 & \text{in } L^2(\Omega_f) & (41) \\
 (\boldsymbol{\sigma} + A) \cdot \boldsymbol{\nu} = -p \boldsymbol{\nu} & \text{in } H^{-1/2}(\Gamma)^m & (42) \\
 \omega^2 \rho_f u \cdot \boldsymbol{\nu} - \frac{\partial p}{\partial \boldsymbol{\nu}} = 0 & \text{in } H^{-1/2}(\Gamma) & (43) \\
 \frac{\partial p}{\partial \boldsymbol{\nu}} = -\mathcal{B} p & \text{in } H^{-1/2}(\Sigma) & (44)
 \end{cases}$$

Remark 5. *The quantity $\operatorname{div} A$ and its normal trace also appear in [8, Eqn (2.3)], and the right-hand-side of our variational formulation is the same as [8, Eqn (2.6)] and [11, Eqn (16),(17)]. Before passing to the variational formulation, [8] states the auxiliary problem in PDE form with the same assumed regularity of A , however without justifying its ‘troublesome’ normal trace. The PDE form of the auxiliary problem is not stated in [11]. Both works start their investigation directly with the variational formulation.*

We now study the well-posedness for **DP** by first showing that it is equivalent to variational problem called **VDP**,

$$\begin{aligned}
 & \text{Find } \begin{pmatrix} p \\ u \end{pmatrix} \in \mathbf{H} : \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] = \mathbf{l}_{\text{fr1}} \left[\begin{pmatrix} \bar{q} \\ \bar{v} \end{pmatrix} \right], \quad \forall \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}; \\
 & \text{with } \mathbf{l}_{\text{fr1}} \in \mathbf{H}', \quad \mathbf{l}_{\text{fr1}} \left[\begin{pmatrix} q \\ v \end{pmatrix} \right] \triangleq - \int_{\Omega_s} A : \nabla v \, dx.
 \end{aligned} \tag{45}$$

Note that,

$$\|\mathbf{l}_{\text{fr1}}\|_{\mathbf{H}'} = \|A\|_{L^2(\Omega_s)^{m \times m}}. \tag{46}$$

Proposition 5 (Variational equivalence for **DP**). *For matrix $A \in L^2(\Omega_s)^{m \times m}$ and $(p, u) \in \mathbf{H}$, the following statements are equivalent.*

(i) (p, u) satisfies (40)–(44).

(ii) (p, u) is a solution of (45).

Proof. **(i) \Rightarrow (ii):** Apply a test function $v \in H^1(\Omega_s)^m$ to both sides of (40), we obtain $-\langle \operatorname{div} \boldsymbol{\sigma} + \operatorname{div} A + \omega^2 \rho_s u, v \rangle_{1, \Omega_s} = 0$. By Lemma 4, $\operatorname{div}(\boldsymbol{\sigma} + A) \in L^2(\Omega_s)^m$, and the above equality can be written as

$$- \int_{\Omega_s} [\operatorname{div}(\boldsymbol{\sigma} + A)] \cdot \bar{v} \, dx - \omega^2 \rho_s \int_{\Omega_s} u \cdot \bar{v} \, dx = 0.$$

We next use the IP identity (39) from the same lemma, and transmission condition (42) to replace $(\boldsymbol{\sigma} + A) \cdot \boldsymbol{\nu}$ with $p \boldsymbol{\nu}$, and obtain

$$\int_{\Omega_s} \boldsymbol{\sigma} : \nabla \bar{v} \, dx + \langle p \boldsymbol{\nu}, v \rangle_{1/2, \Gamma} - \omega^2 \rho_s \int_{\Omega_s} u \cdot \bar{v} \, dx = - \int_{\Omega_s} A : \nabla \bar{v} \, dx, \quad \forall v \in H^1(\Omega_s)^m.$$

Here $\boldsymbol{\sigma}$ and A can be written separately, since each is in $L^2(\Omega_s)^{m \times m}$. In the terms of the sesquilinear forms defined in (21) and (23), the above equality is written as

$$\mathbf{a}_s(u, v) + \mathbf{b}^* \left[\begin{pmatrix} q \\ v \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix} \right] = \mathbf{l}_{\text{fr1}}(\bar{v}), \quad \forall (q, v) \in \mathbf{H}. \quad (47)$$

In dealing with the fluid part, the proof is more standard. By using IP identity (3) and transmission condition (43) and ABC (44), we obtain

$$\mathbf{a}_f(p, q) + \omega^2 \rho_f \mathbf{b} \left[\begin{pmatrix} q \\ v \end{pmatrix}, \begin{pmatrix} p \\ u \end{pmatrix} \right] + \mathbf{t}(p, q) = 0, \quad \forall (q, v) \in \mathbf{H}. \quad (48)$$

For more details, see also [12, Eqn (15)–(20)]. Statement (ii) is obtained by adding (47) to (48).

(ii) \Rightarrow (i): Suppose $(p, u) \in \mathbf{H}$ solves (45)

$$\mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] = \mathbf{l}_{\text{fr1}} \left(\frac{\bar{q}}{v} \right), \quad \forall (q, v) \in \mathbf{H}. \quad (49)$$

Step 1: Using test functions $\phi \in \mathcal{C}_0^\infty(\Omega_s)^m$, $\varphi \in \mathcal{C}_0^\infty(\Omega_f)$ in (49), all the boundary terms go to zero. Now set $\varphi = 0$ (also in (49)), we obtain

$$\int_{\Omega_s} \boldsymbol{\sigma} : \nabla \bar{\phi} \, dx - \omega^2 \rho_s \int_{\Omega_s} u \cdot \bar{\phi} \, dx = - \int_{\Omega_s} A : \nabla \bar{\phi} \, dx, \quad \forall \phi \in \mathcal{C}_0^\infty(\Omega_s)^m.$$

This means $-\text{div} \boldsymbol{\sigma} - \omega^2 \rho_s u = \text{div} A$ in $(\mathcal{C}_0^\infty(\Omega_s)^m)'$ *i.e.* in the sense of distributions. Since $u \in H^1(\Omega_s)^m$, $\text{div}(\boldsymbol{\sigma} + A) \in H^1(\Omega_s)^m$. As a result, the above equality holds in $L^2(\Omega_s)$,

$$\text{div}(\boldsymbol{\sigma} + A) + \omega^2 \rho_s u = 0, \quad \text{in } L^2(\Omega_s)^m. \quad (50)$$

Similarly, set $\phi = 0$ in (49), and we obtain that $\Delta p + \kappa^2 p = 0$ in the sense of distributions. Since $p \in H^1(\Omega_f)$, $\Delta p \in H^1(\Omega_f)$. As a result, we can write

$$\Delta p + \kappa^2 p = 0, \quad \text{in } L^2(\Omega_f). \quad (51)$$

We now use Lemma 4 to make sense of the normal trace $\text{div}(\boldsymbol{\sigma} + A) \cdot \nu$ and conormal trace $\frac{\partial p}{\partial \nu}$ and the integration by parts (39),(3).

Step 2: To derive the transmission and boundary conditions, we choose test functions $\phi \in \mathcal{C}^\infty(\overline{\Omega_s})^m$, $\varphi \in \mathcal{C}^\infty(\overline{\Omega_f})$. Since the arguments to obtain (44) and (43) are standard, *cf.* [12, Eqn 27–31], we will only write that for condition (42). Setting $\varphi = 0$ in (49), we get

$$\int_{\Omega_s} (\boldsymbol{\sigma} + A) : \nabla \bar{\phi} \, dx - \int_{\Omega_s} \omega^2 \rho_s u \cdot \bar{\phi} \, dx + \langle p \nu, \phi \rangle_{1/2, \Gamma} = 0.$$

This is further written, by keeping $\boldsymbol{\sigma}$ and A together and using IP identity (39) (given at the end of Step 1),

$$\int_{\Omega_s} -\operatorname{div}(\boldsymbol{\sigma} + A) \cdot \bar{\boldsymbol{\phi}} \, dx + \langle (\boldsymbol{\sigma} + A) \cdot \boldsymbol{\nu}, \boldsymbol{\phi} \rangle_{1/2, \Gamma} - \int_{\Omega_s} \omega^2 \rho_s u \bar{\boldsymbol{\phi}} \, dx + \langle p \boldsymbol{\nu}, \boldsymbol{\phi} \rangle_{1/2, \Gamma} = 0.$$

Due to (50), the above identity simplifies to

$$\langle (\boldsymbol{\sigma} + A) \cdot \boldsymbol{\nu} + p \boldsymbol{\nu}, \boldsymbol{\phi} \rangle_{1/2, \Gamma} = 0, \quad \forall \boldsymbol{\phi} \in \mathcal{C}^\infty(\overline{\Omega_s})^m.$$

By density arguments, we obtain the transmission condition (42) in $H^{-1/2}(\Gamma)$. \square

Given the equivalence in the above proposition, the well-posedness of problem **DP**, when ω is not a Jones frequency, is obtained via that of (45). The latter follows from Theorem 2 and (46). These results are gathered in the corollary below.

Corollary 6 (Existence and Uniqueness for **DP**). *With \mathcal{B} satisfying (14) and the material parameters satisfying (27), for a matrix $A \in L^2(\Omega_s)^{m \times m}$, when ω is not a Jones frequency, the problem **DP** (40)–(44) has a unique solution $(p, u) \in \mathbf{H}$. In addition, for each (λ, μ) and with h satisfying (32), we have*

$$\begin{aligned} \|\tilde{p}_{\lambda+h, \mu}, \tilde{u}_{\lambda+h, \mu}\|_{\mathbf{H}} &\leq d(\lambda, \mu) \|A\|_{L^2(\Omega_s)^{m \times m}}, \\ \|\tilde{p}_{\lambda, \mu+h}, \tilde{u}_{\lambda, \mu+h}\|_{\mathbf{H}} &\leq d(\lambda, \mu) \|A\|_{L^2(\Omega_s)^{m \times m}}, \end{aligned} \quad (52)$$

for a constant $d(\lambda, \mu) > 0$ independent of h .

2.4. Characterization of partial derivatives with respect to Lamé parameters

The goal of this section is to prove that the solution to **OP** (15)–(19) is separately continuous and partially Fréchet differentiable with respect to the Lamé parameters, and that their partial derivatives can be characterized as solutions to the corresponding auxiliary problem. We simplify the notation and only indicate the dependence with respect to λ and μ . In addition, we will assume that

$$\omega \text{ is not a Jones frequency.} \quad (53)$$

Recall that, under this assumption, we have defined in (35) the solution operator \mathcal{S} for problem **OP**, cf. Corollary 3. If we set $A = \mathbf{C}' : \boldsymbol{\epsilon}(u_{\lambda, \mu})$ with $\mathbf{C}' \in \{\partial_\lambda \mathbf{C}, \partial_\mu \mathbf{C}\}$ and $u_{\lambda, \mu} = \pi_2 \mathcal{S}(\lambda, \mu)$, with π_i the projection map onto the i -th component, then $A \in L^2(\Omega_s)^{m \times m}$. A version of Corollary 6 can be obtained for this particular A .

Corollary 7. Under (53) and (27), with $u_{\lambda,\mu} = \pi_2 \mathcal{S}(\lambda, \mu)$, the problem

$$(\text{LaDP}) \begin{cases} \operatorname{div} \boldsymbol{\sigma} + \omega^2 \rho_s u = -\operatorname{div} (\mathbf{C}' : \boldsymbol{\epsilon}(u_{\lambda,\mu})) & \text{in } (H^1(\Omega_s)^m)' & (54) \\ \Delta p + \kappa^2 p = 0 & \text{in } L^2(\Omega_f) & (55) \\ (\boldsymbol{\sigma} + \mathbf{C}' : \boldsymbol{\epsilon}(u_{\lambda,\mu})) \cdot \boldsymbol{\nu} = -p \boldsymbol{\nu} & \text{in } H^{-1/2}(\Gamma)^m & (56) \\ \omega^2 \rho_f u \cdot \boldsymbol{\nu} - \frac{\partial p}{\partial \boldsymbol{\nu}} = 0 & \text{in } H^{-1/2}(\Gamma) & (57) \\ \frac{\partial p}{\partial \boldsymbol{\nu}} = -\mathcal{B} p & \text{in } H^{-1/2}(\Sigma) & (58) \end{cases}$$

has a unique solution $(\tilde{u}_{\lambda,\mu}, \tilde{p}_{\lambda,\mu}) \in \mathbf{H}$. In addition, it is equivalent to the variational problem **VLaDP**

$$\begin{aligned} & \text{Find } \begin{pmatrix} p \\ u \end{pmatrix} \in \mathbf{H} : \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] = \mathbf{l}_{\text{fr2}} \left[\begin{pmatrix} \bar{q} \\ \bar{v} \end{pmatrix} \right], \quad \forall \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}; \\ & \text{with } \mathbf{l}_{\text{fr2}} \in \mathbf{H}', \quad \mathbf{l}_{\text{fr2}} \left[\begin{pmatrix} q \\ v \end{pmatrix} \right] \triangleq - \int_{\Omega_s} (\mathbf{C}' : \boldsymbol{\epsilon}(u_{\lambda,\mu})) : \nabla v \, dx. \end{aligned} \quad (59)$$

As a result, we can define the solution operators

$$\begin{aligned} & \tilde{\mathcal{S}}_\lambda, \tilde{\mathcal{S}}_\mu : \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbf{H}, \quad (\lambda, \mu) \mapsto (\tilde{p}_{\lambda,\mu}, \tilde{u}_{\lambda,\mu}), \\ & \text{with } (\tilde{p}_{\lambda,\mu}, \tilde{u}_{\lambda,\mu}) \text{ the unique solution in } \mathbf{H} \text{ to (54)–(58) for } \quad (60) \\ & \mathbf{C}' = \partial_\lambda \mathbf{C} \text{ and } \mathbf{C}' = \partial_\mu \mathbf{C} \text{ respectively.} \end{aligned}$$

Proposition 8. Under assumption (53), the mappings $\lambda \mapsto \mathcal{S}(\lambda, \mu)$ and $\mu \mapsto \mathcal{S}(\lambda, \mu)$ are continuous and Fréchet differentiable. Their derivatives are unique solutions to (54)–(58). In particular,

$$\partial_\lambda \mathcal{S}(\lambda, \mu) = \tilde{\mathcal{S}}_\lambda(\lambda, \mu), \quad \partial_\mu \mathcal{S}(\lambda, \mu) = \tilde{\mathcal{S}}_\mu(\lambda, \mu).$$

Proof. Since the proof for μ is identical, only that for λ is written. Since \mathbf{C}_λ is separately affine in λ and μ , we have

$$\mathbf{C}_{\lambda+h} - \mathbf{C}_\lambda = \mathbf{C}'_\lambda h \quad \text{with} \quad \mathbf{C}_\lambda \triangleq \mathbf{C}(\lambda, \mu), \quad \mathbf{C}'_\lambda \triangleq (\partial_\lambda \mathbf{C})(\lambda, \mu).$$

For a fixed μ , consider the unique solutions to **OP** at $\lambda + h$ and λ , and that to **LaDP** at λ ,

$$(p_{\lambda+h}, u_{\lambda+h}) = \mathcal{S}(\lambda + h, \mu); \quad (p_\lambda, u_\lambda) = \mathcal{S}(\lambda, \mu); \quad (\tilde{p}_\lambda, \tilde{u}_\lambda) = \tilde{\mathcal{S}}_\lambda(\lambda, \mu). \quad (61)$$

These quantities have the following regularity,

$$\begin{aligned} & (\tilde{p}_\lambda, \tilde{u}_\lambda), (p_{\lambda+h}, u_{\lambda+h}), (p_\lambda, u_\lambda) \in H^1(\Omega_f) \times H^1(\Omega_s)^m; \\ & \Delta p_{\lambda+h}, \Delta p_\lambda, \Delta \tilde{p}, \Delta \tilde{p}_\lambda \in L^2(\Omega_f), \quad (62) \\ & \operatorname{div}(\boldsymbol{\sigma}(\tilde{u}) + \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_\lambda)) \in L^2(\Omega_s)^m \text{ cf. Lemma 4.} \end{aligned}$$

We will show that their difference $(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)$, and $(\hat{p}_\lambda, \hat{u}_\lambda) \in \mathbf{H}$ with

$$\hat{u}_\lambda \stackrel{\Delta}{=} u_{\lambda+h} - u_\lambda - h \tilde{u}_\lambda \quad , \quad \hat{p}_\lambda \stackrel{\Delta}{=} p_{\lambda+h} - p_\lambda - h \tilde{p}_\lambda \quad (63)$$

satisfy,

$$\|(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)\|_{\mathbf{H}} = \mathcal{O}(|h|) \quad , \quad \|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}} = \mathfrak{o}(|h|) \quad , \quad h \rightarrow 0.$$

Step 1a: Consider the volume identity in Ω_s of each system,

$$\operatorname{div}(\mathbf{C}_{\lambda+h} : \boldsymbol{\epsilon}(u_{\lambda+h})) + \omega^2 \rho_s u_{\lambda+h} = 0 \quad \text{in } L^2(\Omega_s)^m \quad (64)$$

$$- \operatorname{div}(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_\lambda)) - \omega^2 \rho_s u_\lambda = 0 \quad \text{in } L^2(\Omega_s)^m \quad (65)$$

$$-h \left(\operatorname{div}(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\tilde{u}_\lambda)) + \omega^2 \rho_s \tilde{u}_\lambda + \operatorname{div} \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_\lambda) \right) = 0 \quad \text{in } (H^1(\Omega_s)^m)'. \quad (66)$$

In order to rewrite the above expressions, we use ‘as a buffer’

$$- \operatorname{div} \mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h}) + \operatorname{div} \mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h}) = 0 \quad \text{in } (H^1(\Omega_s)^m)'. \quad (67)$$

Adding the above four equations (64) + (65) + (66) + (67), replacing $\mathbf{C}_{\lambda+h} - \mathbf{C}_\lambda$ with $\mathbf{C}'_\lambda h$, and after some rearrangement, we obtain

$$\begin{aligned} \operatorname{div} \mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) + \omega^2 \rho_s \hat{u}_\lambda &= - \operatorname{div} \mathbf{E} \quad \text{in } (H^1(\Omega_s)^m)'; \\ \text{with } \mathbf{E} &\stackrel{\Delta}{=} h \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) \quad \text{in } L^2(\Omega_s)^{m \times m}. \end{aligned} \quad (68)$$

Since $\hat{u}_\lambda \in H^1(\Omega_s)$ and $\mathbf{E} \in L^2(\Omega_s)^{m \times m}$, by Lemma 4, the following normal trace is well-defined,

$$(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) + \mathbf{E}) \cdot \nu \in H^{-1/2}(\Gamma)^m. \quad (69)$$

Working in the same manner with (64) + (65) + (67), also replacing $\mathbf{C}_{\lambda+h} - \mathbf{C}_\lambda$, we obtain an equation for $(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)$,

$$\begin{aligned} \operatorname{div} \mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) + \omega^2 \rho_s (u_{\lambda+h} - u_\lambda) &= - \operatorname{div} \mathbf{e} \quad \text{in } (H^1(\Omega_s)^m)'; \\ \text{with } \mathbf{e} &\stackrel{\Delta}{=} h \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h}) \quad \text{in } L^2(\Omega_s)^{m \times m}. \end{aligned} \quad (70)$$

Similarly, since $u_{\lambda+h} - u_\lambda \in H^1(\Omega_s)$ and $\mathbf{e} \in L^2(\Omega_s)^{m \times m}$, Lemma 4 gives

$$(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) + \mathbf{e}) \cdot \nu \in H^{-1/2}(\Gamma)^m. \quad (71)$$

Step 1b: We next work with the kinematic interface condition in each problem,

$$(\mathbf{C}_{\lambda+h} : \boldsymbol{\epsilon}(u_{\lambda+h})) \cdot \nu + p_{\lambda+h} \nu = -p_{\text{inc}} \nu \quad \text{in } H^{-1/2}(\Gamma)^m \quad (72)$$

$$- (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_\lambda)) \cdot \nu - p_\lambda \nu = p_{\text{inc}} \nu \quad \text{in } H^{-1/2}(\Gamma)^m \quad (73)$$

$$-h \left(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\tilde{u}) + \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_\lambda) \right) \cdot \nu - h \tilde{p} \nu = 0 \quad \text{in } H^{-1/2}(\Gamma)^m. \quad (74)$$

A ‘buffer’ in the form of normal trace is needed as before. However, to maintain the current regularity (in $H^{-1/2}(\Gamma)^m$), thanks to (69), we can use

$$- (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) + \mathbf{E}) \cdot \nu + (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) + \mathbf{E}) \cdot \nu = 0 \quad \text{in } H^{-1/2}(\Gamma)^m. \quad (75)$$

When adding the above four equations together, the sum of all first terms simplifies to zero¹², we obtain the kinematic interface condition,

$$(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) + \mathbf{E}) \cdot \nu + \hat{p}_\lambda \nu = 0 \quad \text{in } H^{-1/2}(\Gamma)^m. \quad (76)$$

To derive the interface condition for $u_{\lambda+h} - u_\lambda$, we use the same idea as before to maintain regularity in $H^{-1/2}(\Gamma)^m$. Using (71), we add (72) and (73) to

$$- (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) + \mathbf{e}) \cdot \nu + (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) + \mathbf{e}) \cdot \nu = 0 \in H^{-1/2}(\Gamma)^m. \quad (77)$$

After some rearrangement and simplification, we find

$$(\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda) + \mathbf{e}) \cdot \nu + (p_{\lambda+h} - p_\lambda) \nu = 0 \quad \text{in } H^{-1/2}(\Gamma)^m. \quad (78)$$

Step 1c: We next work with the Helmholtz equation, kinematic transmission condition (18), and ABC in each problem. Taking

$$((16), (18), (19))|_{\lambda+h} - ((16), (18), (19))|_\lambda - h((55), (57), (58))|_\lambda,$$

we obtain

$$\begin{aligned} \Delta \hat{p}_\lambda + \kappa^2 \hat{p}_\lambda &= 0 \in L^2(\Omega_f); \\ \omega^2 \rho_f \hat{u}_\lambda \cdot \nu &= \frac{\partial \hat{p}_\lambda}{\partial \nu} \in H^{-1/2}(\Gamma); \quad \frac{\partial \hat{p}_\lambda}{\partial \nu} = \mathcal{B} \hat{p}_\lambda \in H^{-1/2}(\Sigma). \end{aligned} \quad (79)$$

Using $((16), (18), (19))|_{\lambda+h} - ((16), (18), (19))|_\lambda$, we obtain the exact same system as above with \hat{u}_λ replaced by $u_{\lambda+h} - u_\lambda$ and \hat{p}_λ replaced by $p_{\lambda+h} - p_\lambda$.

At the end of Step 1a – 1c, we have obtained $\mathbf{E} \in L^2(\Omega_s)^{m \times m}$ and $(\hat{p}_\lambda, \hat{u}_\lambda) \in \mathbf{H}$ satisfying (40)–(44) with A replaced by \mathbf{E} . Similarly, we have $\mathbf{e} \in L^2(\Omega_s)^{m \times m}$ and $(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda) \in \mathbf{H}$ satisfying (40)–(44) with A replaced by \mathbf{e} .

¹² We first group the first terms of (74) and (75), and use the definition of \mathbf{E} ,

$$(-\mathbf{C}_\lambda : \boldsymbol{\epsilon}(h \tilde{u}) - h \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_\lambda) - \mathbf{C}_\lambda : \boldsymbol{\epsilon}(\hat{u}_\lambda) - h \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda)) \cdot \nu.$$

With the last three terms together grouped together, the above expression simplifies to

$$(-\mathbf{C}_\lambda : \boldsymbol{\epsilon}(h \tilde{u}) + \mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_\lambda + h \tilde{u}_\lambda) - \mathbf{C}_{\lambda+h} : \boldsymbol{\epsilon}(u_{\lambda+h})) \cdot \nu = (\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_\lambda) - \mathbf{C}_{\lambda+h} : \boldsymbol{\epsilon}(u_{\lambda+h})) \cdot \nu.$$

Now adding the first terms of (72) and (73), we get zero. Note that all the above expressions are in $H^{-1/2}(\Gamma)^m$.

Step 2: We next bound the differences. For the rest of the discussion, we assume that h satisfies (32). From Corollary 6, under assumption (53), we have existence and uniqueness of solutions, and the energy estimate (52) gives

$$\begin{aligned} \|(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)\|_{\mathbf{H}} &\leq d(\lambda, \mu) \|\mathbf{e}\|_{L^2(\Omega_s)^{m \times m}}, \\ \|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}} &\leq d(\lambda, \mu) \|\mathbf{E}\|_{L^2(\Omega_s)^{m \times m}}. \end{aligned} \quad (80)$$

Using its definition, we can bound \mathbf{e} by

$$\begin{aligned} \|\mathbf{e}\|_{L^2(\Omega_s)^{m \times m}} &\leq \|\mathbf{C}_{\lambda+h} - \mathbf{C}_\lambda\|_\infty \|\boldsymbol{\epsilon}(u_{\lambda+h})\|_{L^2(\Omega_s)^{m \times m}} \\ &\stackrel{(34)}{\leq} |h| \|\mathbf{C}'_\lambda\|_\infty c(\lambda, \mu) \|p_{\text{inc}}\|_{H^2(\Omega_f)}. \end{aligned} \quad (81)$$

With \mathbf{C}'_λ being constant, there exists $\tilde{c}(\lambda, \mu) > 0$ independent of h such that

$$\|(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)\|_{\mathbf{H}} \leq |h| \tilde{c}(\lambda, \mu) \|p_{\text{inc}}\|_{H^2(\Omega_f)}. \quad (82)$$

From this follows the continuity of the mapping $\lambda \mapsto \mathcal{S}(\lambda, \mu)$.

Similarly, the bound of \mathbf{E} follows from its definition,

$$\|\mathbf{E}\|_{L^2(\Omega_s)^{m \times m}} = \|h \mathbf{C}'_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h} - u_\lambda)\|_{L^2(\Omega_s)^{m \times m}} \leq |h| \|\mathbf{C}'_\lambda\|_\infty \|u_{\lambda+h} - u_\lambda\|_{H^1(\Omega_s)^m}.$$

Combining with (82), we obtain a constant $\tilde{d}(\lambda, \mu) > 0$ independent of h such that

$$\|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}} \leq |h|^2 \tilde{d}(\lambda, \mu) \|p_{\text{inc}}\|_{H^2(\Omega_f)}. \quad (83)$$

As a result, $\|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}} = \mathbf{o}(h)$. This finishes the proof that $\lambda \mapsto \mathcal{S}$ is Fréchet differentiable with derivative $\partial_\lambda \mathcal{S} = \tilde{\mathcal{S}}_\lambda(\lambda, \mu)$. \square

Remark 6. *A more natural buffer than that employed in (75) would have been the ‘formal’ normal trace of $-\mathbf{C}_\lambda : \boldsymbol{\epsilon}(u_{\lambda+h}) \in (H^1(\Omega_s)^m)'$, which is the buffer in (67). However, under Lipschitz assumption, this ‘formal’ normal trace does not make sense in $H^s(\Gamma)$, $|s| \leq 1$. That is why we work with the buffer terms given in (69) and (71), which have well-defined normal trace.*

Differentiability of approximate far-field-pattern (a-FFP) under assumption (53). We recall that the a-FFP operator $\mathbb{F}(\lambda, \mu)$ associated to \mathbf{OP} is given in (37). The a-FFP operator associated to \mathbf{LaDP} can be defined in a similar manner,

$$\begin{aligned} \tilde{\mathbb{F}}_\bullet : \mathbb{R}^+ \times \mathbb{R}^+ &\longrightarrow \mathcal{C}^\infty(\mathbb{S}^1, \mathbb{C}) \quad , \quad (\lambda, \mu) \longrightarrow \tilde{\mathbb{F}}_\bullet(\lambda, \mu) \\ (\tilde{\mathbb{F}}_\bullet(\lambda)(\hat{x})) &= \mathbf{c}(m) \int_\Gamma \left(\Phi(\hat{x}, y) \tilde{p}(y) - \omega^2 \rho_f \Psi(\hat{x}, y) \tilde{u}(y) \cdot \nu \right) ds(y). \quad (84) \\ \text{where } (\tilde{p}, \tilde{u}) &= \tilde{\mathcal{S}}_\bullet(\lambda, \mu) \quad , \quad \text{for } \bullet = \lambda, \mu. \end{aligned}$$

Here, we have written

$$\Phi(\hat{x}, y) \triangleq \frac{\partial \epsilon^{-i\kappa \hat{x} \cdot y}}{\partial \nu(y)} \quad ; \quad \Psi(\hat{x}, y) \triangleq e^{-i\kappa \hat{x} \cdot y}. \quad (85)$$

We proceed to show their separate continuity and differentiability.

Proposition 9. *The mappings $\lambda \mapsto F(\lambda, \mu)$ and $\mu \mapsto F(\lambda, \mu)$ are continuous and Fréchet differentiable with*

$$\partial_\lambda F(\lambda, \mu) = \tilde{F}_\lambda(\lambda, \mu) \quad , \quad \partial_\mu F(\lambda, \mu) = \tilde{F}_\mu(\lambda, \mu).$$

Proof. We will only write out the proof with respect to λ , that for μ is completely similar. We will show that

$$F(\lambda + h) - F(\lambda) = \mathcal{O}(|h|) \quad , \quad F(\lambda + h) - F(\lambda) - h\tilde{F}(\lambda) = \mathfrak{o}(|h|) \quad , \quad h \rightarrow 0.$$

In terms of the quantities defined in (61)–(63), the differences of the a-FFPs are

$$\begin{aligned} F(\lambda + h) - F(\lambda) - h\tilde{F}(\lambda) &= \mathbf{c}(m) \int_\Gamma \left(\Phi(\hat{x}, y) \hat{p}_\lambda(y) - \omega^2 \rho_f \Psi(\hat{x}, y) \hat{u}_\lambda(y) \cdot \nu \right) ds(y); \\ F(\lambda + h) - F(\lambda) &= \mathbf{c}(m) \int_\Gamma \left(\Phi(\hat{x}, y) (p_{\lambda+h} - p_\lambda)(y) - \right. \\ &\quad \left. \omega^2 \rho_f \Psi(\hat{x}, y) (u_{\lambda+h} - u_\lambda)(y) \cdot \nu \right) ds(y). \end{aligned}$$

Using $|\Phi(\hat{x}, y)| = 1$ and $|\Psi(\hat{x}, y)| = |i\kappa \hat{x} \cdot \hat{n}| = \kappa$, we obtain a constant $c = c(n, \Gamma, \omega, \rho_f, \kappa) > 0$ independent of h so that

$$|F(\lambda + h) - F(\lambda) - h\tilde{F}(\lambda)| \leq c \left(\|\hat{p}_\lambda\|_{L^2(\Gamma)} + \|\hat{u}_\lambda\|_{L^2(\Gamma)} \right) \leq c \|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}}.$$

Similarly, there exists a constant $\tilde{c} = \tilde{c}(n, \Gamma, \omega, \rho_f, \kappa) > 0$ independent of h so that

$$|F(\lambda + h) - F(\lambda)| \leq \tilde{c} \|(p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda)\|_{\mathbf{H}}.$$

Separate continuity is obtained by using the bound (82) for $\|p_{\lambda+h} - p_\lambda, u_{\lambda+h} - u_\lambda\|_{\mathbf{H}}$, and partial Fréchet differentiability by (83) for $\|(\hat{p}_\lambda, \hat{u}_\lambda)\|_{\mathbf{H}}$. \square

2.5. Characterization of the partial derivatives with respect to solid density

The analysis with respect to the solid density ρ_s is simpler. The corresponding auxiliary problem has better regularity; it has no interface source terms and, instead, only a very regular volume one (in $H^1(\Omega_s)^m$). As a result, its PDE form is readily well-defined in standard Sobolev spaces. For the current discussion, the

notation is simplified to indicate only the dependence with respect to ρ_s . With $u_{\rho_s} = \mathcal{S}(\rho_s) = \mathcal{S}(\lambda, \mu, \rho_s)$ cf. (35), the auxiliary problem with respect to ρ_s is

$$\begin{aligned}
 (\text{SoDP}) \quad & \begin{cases} \operatorname{div} \boldsymbol{\sigma} + \omega^2 \rho_s u = -\omega^2 u_{\rho_s} & \text{in } L^2(\Omega_s)^m & (86) \\ \Delta p + \kappa^2 p = 0 & \text{in } L^2(\Omega_f) & (87) \\ \boldsymbol{\sigma} \cdot \boldsymbol{\nu} = -p \boldsymbol{\nu} & \text{in } H^{-1/2}(\Gamma)^m & (88) \\ \omega^2 \rho_f u \cdot \boldsymbol{\nu} - \frac{\partial p}{\partial \boldsymbol{\nu}} = 0 & \text{in } H^{-1/2}(\Gamma) & (89) \\ \frac{\partial p}{\partial \boldsymbol{\nu}} = -\mathcal{B} p & \text{in } H^{-1/2}(\Sigma). & (90) \end{cases}
 \end{aligned}$$

The arguments to obtain the variational formulation equivalent to **SoDP** and the well-posedness for both problems are standard. The latter follows from Theorem 2.

Proposition 10. *Problem **SoDP** is equivalent to variational problem*

$$\begin{aligned}
 \text{Find } \begin{pmatrix} p \\ u \end{pmatrix} \in \mathbf{H} : \mathbf{a} \left[\begin{pmatrix} p \\ u \end{pmatrix}, \begin{pmatrix} q \\ v \end{pmatrix} \right] &= \mathbf{l}_{\text{fr3}} \left[\begin{pmatrix} \bar{q} \\ \bar{v} \end{pmatrix} \right], \quad \forall \begin{pmatrix} q \\ v \end{pmatrix} \in \mathbf{H}; \\
 \text{with } \mathbf{l}_{\text{fr3}} \in \mathbf{H}' , \mathbf{l}_{\text{fr3}} \left[\begin{pmatrix} q \\ v \end{pmatrix} \right] &\triangleq - \int_{\Omega_s} \omega^2 u_{\rho_s} : v \, dx. & (91)
 \end{aligned}$$

Given (53), there exist unique solutions to both problems in \mathbf{H} .

As a result of this, we can define the solution operator:

$$\begin{aligned}
 \tilde{\mathcal{S}} : \mathbb{R}^+ &\rightarrow \mathbf{H}, \quad \rho_s \mapsto \tilde{\mathcal{S}}(\rho_s), \\
 \text{with } \tilde{\mathcal{S}}(\rho_s) &\text{ the unique solution in } \mathbf{H} \text{ to (86)–(90)}. & (92)
 \end{aligned}$$

The a-FFP associated to **SoDP**, denoted by $\tilde{\mathbf{F}}(\rho_s)$, is defined using (84) with $(\tilde{p}, \tilde{u}) = \tilde{\mathcal{S}}(\rho_s)$.

The proofs for continuity and partial differentiability in this case are simpler than those for the Lamé parameters. In particular, we do not encounter the problem discussed in Remark 6. These results are gathered in the following proposition.

Proposition 11. *Given (53), at fixed (λ, μ) , the mapping $\rho_s \mapsto \mathcal{S}(\rho_s)$ is continuous and Fréchet differentiable with derivative given as the unique solution to (86)–(90). In particular, $\partial_{\rho_s} \mathcal{S}(\rho_s) = \tilde{\mathcal{S}}(\rho_s)$. Regarding a-FFP, the mapping $\rho_s \mapsto \mathbf{F}(\rho_s)$ (defined in (37)) is continuous and Fréchet differentiable with $\partial_{\rho_s} \mathbf{F}(\rho_s) = \tilde{\mathbf{F}}(\rho_s)$.*

3. Numerical validation

In this section, to evaluate the accuracy of this characterization, we consider a 2D setting with circular geometries, in which there are analytical expressions for

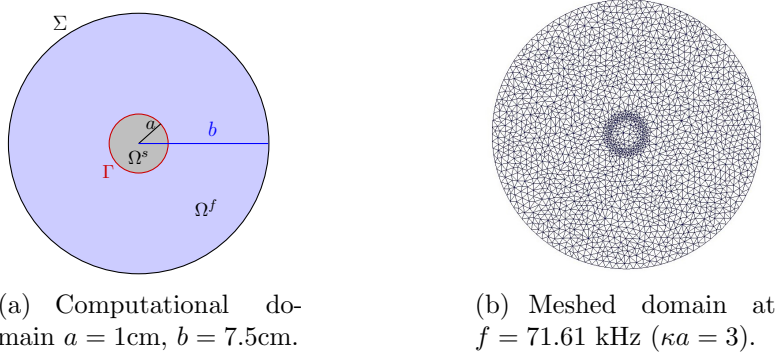


Figure 2: Circular computational domain Ω_{finite} with circular obstacle Ω_s . The mesh is used with IPDG method and satisfies 7 points per wavelength, *cf.* Remark 7. Boundary Σ and interface Λ are discretized using curved-edge elements *cf.* [13, Sec. 3.1].

the displacement fields, a-FFP and their derivatives. This allows for a comparison between analytical values and numerical ones. The latter are obtained by using IPDG method equipped with curved-boundary edge elements, *cf.* [13], to discretize the variational problems (**VOP** and **VLaDP**).

Parameters for experiments. We follow [38, 13] and consider a setup in which a time-harmonic planewave propagating along x -axis is diffracted by an infinitely long *steel* cylinder immersed in water. Here, the coordinates of \mathbb{R}^3 are written as (x, y, z) and of \mathbb{R}^2 as (x, y) . The cylinder has radius $a = 1\text{ cm}$ and principal axis parallel to z -direction. The properties of the involved materials are summarized in Table 1. This 3D phenomenon can be approximated by a 2D problem, in which the solid region Ω_s represents the circular cross-section of the cylinder in (x, y) -plane. The computation domain Ω_{finite} is chosen to be a disc of radius $b = 7.5\text{cm}$, see Figure 2(a), and the first-order ABC with $\mathcal{B} = \mathbf{T}_1$ defined in (20) is used.

<i>Medium</i>	<i>Density</i> (kg m^{-3})	<i>Propagation speed</i> (m s^{-1})		<i>Lamé parameters</i> (GPa)	
	ρ	c_P	c_S	λ	μ
Water	1000	1500	n/a	2.25	n/a
Steel	7900	5837	3100	115.40	76.9

Table 1: Properties of fluid and solid medium in the experiments.

In (x, y) - and polar coordinates, the spatial part of a time-harmonic planewave of frequency $f = \frac{\omega}{2\pi}$, is

$$p_{\text{pw}}(x, y) = e^{i \frac{\omega}{c_f} (x, y) \cdot (1, 0)} = p_{\text{pw}}(r, \theta) = e^{i \frac{\omega}{c_f} r \cos \theta}. \quad (93)$$

Experiments are carried out for frequency f and normalized wavenumber κa range

$$23 \text{ kHz} \leq f \leq 477 \text{ kHz} \quad \Leftrightarrow \quad 1 \leq \kappa a \leq 20.$$

We compare between the IPDG and analytical values of the quantities listed in Table 2. To quantify the difference between an analytical value g and the corresponding value g_h computed with IPDG, we use L^2 -relative error,

$$\mathbf{e}(g) = \frac{\|g - g_h\|_2}{\|g\|_2} \times 100. \quad (94)$$

	Fluid pressure	Solid displacement	a-FFP	Fréchet derivatives		
IPDG	p_h	$u_h = ((u_h)_x, (u_h)_y)$	$(p_h)_{a,\infty}$	∂p_h	∂u_h	$\partial(p_h)_{a,\infty}$
Exact	p	$u = (u_x, u_y)$	$p_{a,\infty}$	∂p	∂u	$\partial p_{a,\infty}$

Table 2: Compared quantities. For the Fréchet derivatives, we write ∂ to mean either ∂_λ or ∂_μ . Subscript x or y indicates corresponding components of the solid displacement.

Method of calculation. To calculate (p_h, u_h) , variational problem **VOP** (33) is discretized by IPDG of order \mathbb{P}_5 with curved-boundary edges also of order \mathbb{P}_5 , implemented in the software Hou10ni¹³. The displacement u_h is then used in **VLaDP** (59), also discretized by IPDG to give $(\partial p_h, \partial u_h)$. We use the same interior penalty terms as in [13, Eqn 4–5]. Note that our sesquilinear form differs from [13] in the ABC. Since our goal is not to study convergence order, parameters are chosen to obtain the best precision associated with \mathbb{P}_5 . In particular, the mesh is changed at each different frequency according to the criteria described in Remark 7.

Remark 7 (Mesh criteria). *In all of our experiments, at each frequency f , the mesh, cf. Figure 2(b), is created by the software `triangle`¹⁴ taking three variables: the number of discretization points on interface Γ and boundary Σ , denoted by N_Γ and N_Σ respectively, and elemental area \mathbf{a}_h . These are defined by*

$$N_\Gamma \triangleq \left\lceil \frac{2\pi a}{h} \right\rceil, \quad N_\Sigma \triangleq \left\lceil \frac{2\pi b}{h} \right\rceil, \quad \mathbf{a}_h \triangleq \frac{h^2}{2},$$

in terms of ‘mesh size’ h . The quantity ‘mesh size’ h defined as

$$h \triangleq \frac{\lambda_f}{N_{\text{ppwl}}} = \frac{2\pi c_f}{\omega N_{\text{ppwl}}},$$

¹³Degree \mathbb{P}_5 corresponds to 21 degree of freedoms per mesh element (triangle). Triangles with curved-boundary edges are used at the fluid-solid interface, cf. [13, Figure 1]. For Hou10ni, see team.inria.fr/magique3d/software/hou10ni/.

¹⁴<https://www.cs.cmu.edu/~quake/triangle.html>

is controlled by the frequency f , and the number of points per (shortest) wavelength¹⁵, denoted by N_{ppwl} . Unless indicated otherwise, $N_{\text{ppwl}} = 7$.

Due to circular geometries, the analytical values can be written in the form of multipole expansions, see [13, Eqn (A.1)] for the scattered field p and [13, Eqn (A.4)] for the solid displacement u . These expressions are defined in terms of the coefficients $X_n = (a_n, b_n, c_n, d_n)^t$ at each level $n \in \mathbb{N}$, which satisfy the linear system $\mathbb{E}_n X_n = E_n$, see [13, Eqn (17)]. Since our problem **OP** differs from [12] in the ABC (19), our entries \mathbb{E}_n^{41} and \mathbb{E}_n^{42} are different, with

$$\mathbb{E}_n^{4i} = \kappa H_n^{(i)'}(\kappa b) + \left(\frac{1}{2b} - i\kappa\right) H_n^{(i)}(\kappa b), \quad i = 1, 2. \quad (95)$$

Here, $H_n^{(\bullet)}$ are the cylindrical Hankel functions. The components of right-hand side $E_n \in \mathbb{C}^4$ are given in [12, Eqn (19)–(21)]. The derivatives $\partial p, \partial u$ are calculated in the same manner, however with coefficients $\partial X_n = (\partial a_n, \partial b_n, \partial c_n, \partial d_n)$, given by

$$\partial X_n = -\mathbb{E}_n^{-1} (\partial \mathbb{E}_n) X_n = -\mathbb{E}_n^{-1} (\partial \mathbb{E}_n) \mathbb{E}_n^{-1} E_n. \quad (96)$$

While the *true* analytical values are infinite sums¹⁶, those for numerical purposes are truncated versions with $N = 2\lceil \kappa b \rceil + 1$ modes (a standard choice *cf.* [13]). For illustrative purposes, we show the real part of $\partial_\lambda p$ at $\kappa a = 7$ in Figure 3(a).

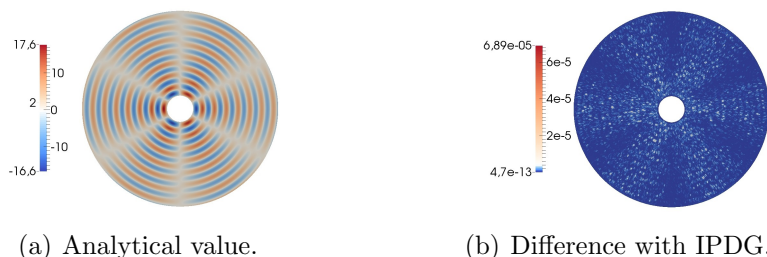


Figure 3: Analytical value of $\text{Re}(\partial_\mu p)$ at $\kappa a = 7$ is shown in (a). Absolute difference with IPDG $|\text{Re}(\partial_\mu p - \partial_\mu p_h)| = 1.95 \times 10^{-7}$ is shown in (b), corresponding to $2.38 \times 10^{-5}\%$ relative error.

With the analytical values $p, u, \partial p, \partial u$ and the IPDG values $p_h, u_h, \partial p_h, \partial u_h$ calculated, we now obtain the analytical $p_{a,\infty}$ and numerical $(p_h)_{a,\infty}$ a-FFP and their corresponding derivatives. We recall that a-FFPs associated to problem **OP** are given by (37),

$$\begin{bmatrix} (p_h)_{a,\infty}(\hat{\mathbf{x}}) \\ p_{a,\infty}(\hat{\mathbf{x}}) \end{bmatrix} = \mathbf{c}(2) \int_\Gamma \left\{ \Phi(\hat{\mathbf{x}}, \mathbf{y}) \begin{bmatrix} p_h(\mathbf{y}) \\ p(\mathbf{y}) \end{bmatrix} - (\omega \rho_f^2 \begin{bmatrix} u_h(\mathbf{y}) \\ u(\mathbf{y}) \end{bmatrix} \cdot \boldsymbol{\nu} - g(\mathbf{y})) \Psi(\hat{\mathbf{x}}, \mathbf{y}) \right\} ds(\mathbf{y}) \quad (97)$$

¹⁵ Wavelength is given by $\frac{\text{propagation speed}}{\text{source frequency } f}$. Water has the shortest wavelength, due to its speed c_f inferior to both the primary speed c_p and the shear speed c_s in steel as indicated in Table 1.

¹⁶ They are uniformly convergent on compact subsets.

with $g := \frac{\partial p_{pw}}{\partial \nu}$ and $\mathbf{c}(2)$ defined in (13), and their derivatives by (84),

$$\begin{bmatrix} \partial(p_h)_{a,\infty}(\hat{\mathbf{x}}) \\ \partial p_{a,\infty}(\hat{\mathbf{x}}) \end{bmatrix} = \mathbf{c}(2) \int_{\Gamma} \left\{ \Phi(\hat{\mathbf{x}}, \mathbf{y}) \begin{bmatrix} \partial p_h(\mathbf{y}) \\ \partial p(\mathbf{y}) \end{bmatrix} - \omega \rho_f^2 \begin{bmatrix} \partial u_h(\mathbf{y}) \\ \partial u(\mathbf{y}) \end{bmatrix} \cdot \nu \Psi(\hat{\mathbf{x}}, \mathbf{y}) \right\} ds(\mathbf{y}). \quad (98)$$

These quantities will be evaluated at 360 angles $\hat{\mathbf{x}} = \frac{\mathbf{x}}{|\mathbf{x}|}$, equally distributed on the unit circle. Here, square and curly braces group terms together, and do not have additional significance.

Remark 8 (a-FFP vs FFP). *Recall that FFP can only be defined for the fluid pressure associated to problem \mathbf{OP}_{∞} . If we consider a disc-shaped obstacle in this problem, the true FFP denoted by p_{∞} has the following expansion,*

$$p_{\infty}(\hat{\mathbf{x}}) = \sum_n \sqrt{\frac{2}{\kappa\pi}} \frac{e^{i\pi/4}}{i^{n+1}} a_n \cos(n\theta), \quad \hat{\mathbf{x}} = \frac{\mathbf{x}}{|\mathbf{x}|}, \quad (99)$$

where the coefficients a_n are obtained by solving $\hat{\mathbb{E}}_n \hat{X}_n = \hat{E}_n$ with $\hat{X}_n = (a_n, c_n, d_n)^t$. These are obtained in the same way as \mathbb{E}_n but without the ABC¹⁷.

Experiments and observations. There are two main sets of comparison, both carried out for the normalized wavenumber range $4 \leq \kappa a \leq 12$. In the set 1, Figure 4 and 5 show the comparison for p , u , and a-FFP, while in the second set, Figure 6 and 7 show that for their partial derivatives.

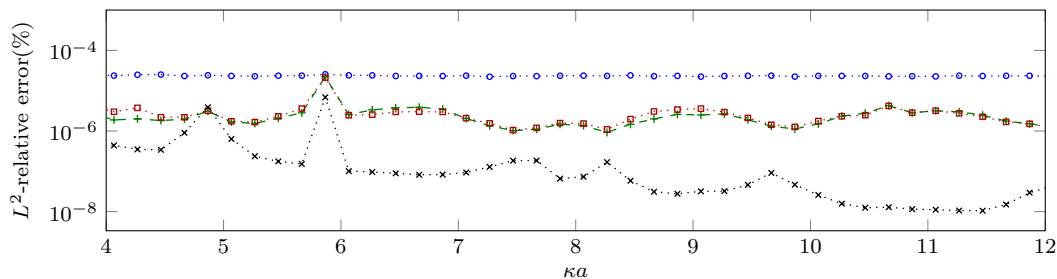


Figure 4: Comparison between the analytical and IPDG values for fluid pressure $\cdots\circ\cdots$ $\mathbf{c}(p)$, solid displacement $-+-$ $\mathbf{c}(u_x)$, $\cdots\square\cdots$ $\mathbf{c}(u_y)$, and a-FFP $\cdots\times\cdots$ $\mathbf{c}(p_{a,\infty})$.

Below are the observations for Figure 4.

¹⁷In particular, $\hat{\mathbb{E}}_n$ is obtained from \mathbb{E}_n removing the last row and second column, *i.e.*

$$\hat{\mathbb{E}}_n = \begin{pmatrix} \mathbb{E}_n^{11} & \mathbb{E}_n^{13} & \mathbb{E}_n^{14} \\ \mathbb{E}_n^{21} & \mathbb{E}_n^{23} & \mathbb{E}_n^{24} \\ \mathbb{E}_n^{31} & \mathbb{E}_n^{33} & \mathbb{E}_n^{34} \end{pmatrix}, \quad \hat{E}_n = (E_n^1, E_n^2, E_n^3)^t.$$

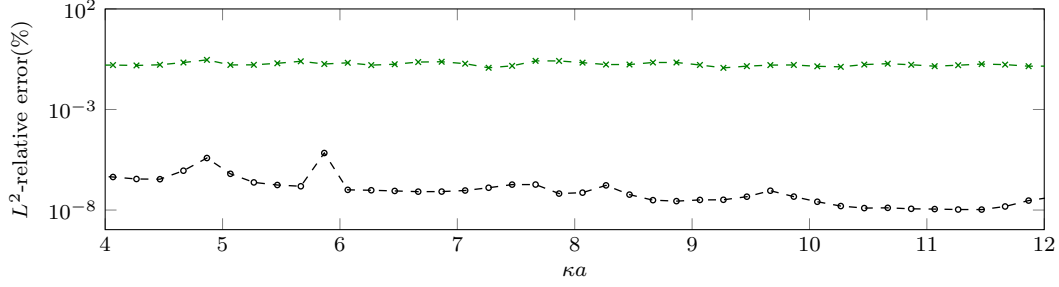


Figure 5: Error a-FFP - \circ - $\epsilon(p_{a,\infty}) = \frac{|(p_h)_{a,\infty} - p_{a,\infty}|}{|p_{a,\infty}|} \times 100$ vs. error FFP - $*$ - $\epsilon = \frac{|p_\infty - p_{a,\infty}|}{|p_\infty|} \times 100$. The error $\epsilon(p_{a,\infty})$ is an indication of the quality of IPDG and the chosen integration quadrature, while ϵ reflects only the quality of the absorbing boundary condition in approximating the original problem \mathbf{OP}_∞ on \mathbb{R}^2 .

- Figure 4 shows that IPDG gives a very high accuracy level, with errors $\epsilon(p)$, $\epsilon(u)$, $\epsilon(p_{a,\infty})$ around $10^{-5}\%$ for the current normalized wavenumber range. Although numerical a-FFP is calculated from one procedure of discretization (to calculate u_h and p_h) and integration (97), we do not see the effect of accumulation of numerical error. Instead, the error $\epsilon(p_{a,\infty})$ is lower than both $\epsilon(p)$ and $\epsilon(u)$.
- We see here the effectiveness of the mesh criteria defined in Remark 7, in maintaining the same error level, with the error curve either flat for $\epsilon(p)$ or varying in a fixed range for $\epsilon(u)$, instead of increasing with the frequency.
- With curved-boundary edge elements, the error curve $\epsilon(p)$ is flat and smooth, compared to those for solid displacement $\epsilon(u_x)$ and $\epsilon(u_y)$, which have spikes, especially around a Jones frequency¹⁸. This effect is also reported in [13].
- Since a-FFP is calculated in terms of both the traces of p and u , its error curve $\epsilon(p_{a,\infty})$ loses the flatness of $\epsilon(p)$, and takes on the fluctuation of $\epsilon(u)$.

The quality of the numerical solution also depends on the choice of the ABC. Figure 5 shows the satisfactory quality¹⁹ of \mathbf{T}_1 , by comparing the true FFP p_∞ , *cf.* (99), with the analytical a-FFP $p_{a,\infty}$, *cf.* (97), with 0.3% relative error.

We next comment on the second set of experiments given in Figures 6 – 7, which show the errors for the partial derivatives, $\epsilon(\partial p)$, $\epsilon(\partial u)$, and $\epsilon(\partial p_{a,\infty})$.

¹⁸For an illustration of Jones frequencies, see [13, Fig.8], [38, Fig. 7].

¹⁹ Although (satisfactorily) small, ϵ is much larger than the relative error $\epsilon(p_{a,\infty})$. To decrease further the effect of the ABC (currently first order), higher orders can be used, and the radius b of Ω_{finite} can be chosen larger.

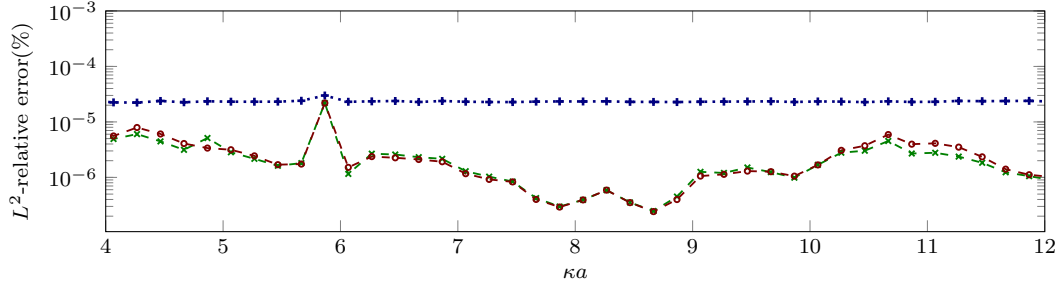
- Relative errors for both $\mathbf{e}(\partial p)$ and $\mathbf{e}(\partial u)$ are below than 10^{-4} , *cf.* Figure 6(a)–6(b). Although two discretization procedures are used (first to obtain u_h and then the partial derivatives), $\mathbf{e}(\partial p)$ and $\mathbf{e}(\partial u)$ are around the same magnitude of $\mathbf{e}(p)$ and $\mathbf{e}(u)$ (shown *cf.* Figure 4). This means no effect of discretization error accumulation.
- The error $\mathbf{e}(\partial p_{a,\infty})$ in Figure 7 is slightly higher, fluctuating around $10^{-4}\%$. This slight increase is expected, since IPDG values of derivatives of a-FFP are computed from two discretization procedures (to obtain ∂u_h and ∂p_h) and an integration procedure in (98). However, accumulation of discretization error is still very negligible.
- Compared with the flat and smooth curve of $\mathbf{e}(p)$, a small peak just before $\kappa a = 6$ is now observed in that of $\mathbf{e}(\partial p)$. Outside of this neighborhood, the error curve stays smooth and flat. For the solid displacement, this peak, already observed for the curve of $\mathbf{e}(u)$, is now more pronounced for $\mathbf{e}(\partial u)$.
- The error with respect to λ behaves in the same way as with respect to μ , apart from the following minute distinctions. Figure 7 shows that $\mathbf{e}(\partial_\mu p_{a,\infty})$ is always below $\mathbf{e}(\partial_\lambda p_{a,\infty})$, and in Figure 6(a)–6(b), the peak just before $\kappa a = 6$ is more pronounced in $\mathbf{e}(\partial_\mu p)$ than in $\mathbf{e}(\partial_\lambda p)$.

4. Conclusion

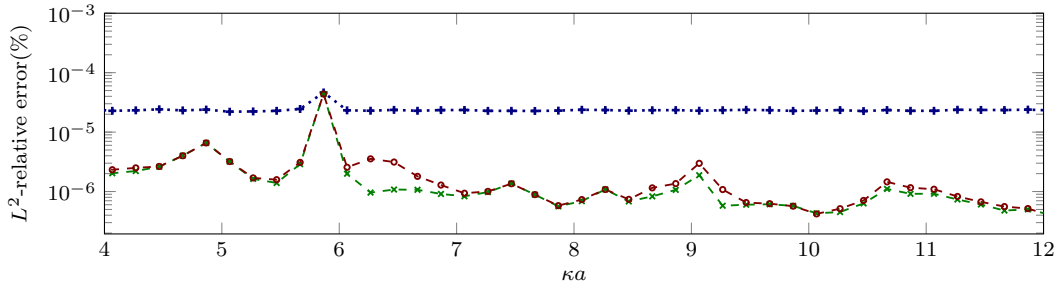
For the stated fluid-solid interaction problem with a Lipschitz interface, we have shown the separate continuity and partial Fréchet differentiability of the solution operator and a-FFP with respect to the solid materials. The main difficulty in treating Lipschitz interfaces is the presence, in the PDE form of the auxiliary problem (which is used to characterize the partial derivatives), of inhomogeneous interface terms whose normal traces cannot be defined in canonical boundary Sobolev spaces. This is overcome in our work, with well-posedness shown in the standard framework for both the PDE and variational form (of the auxiliary problems). Joint continuity and full Fréchet differentiability results can be obtained with slight modification of the current work. Our analytical-numerical comparisons show that using IPDG equipped with curved boundary edge elements and high enough discretization order, in this characterization, gives high precision and incurs almost no effect of discretization error accumulation.

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(a) Comparison for partial derivative with respect to λ .



(b) Comparison for partial derivative with respect to μ .

Figure 6: Comparison between the analytical and IPDG values of the Fréchet derivative of fluid pressure $\cdots+$ $\epsilon(\partial p)$ and solid displacement $-x-$ $\epsilon(\partial u_x)$, $-o-$ $\epsilon(\partial u_y)$.

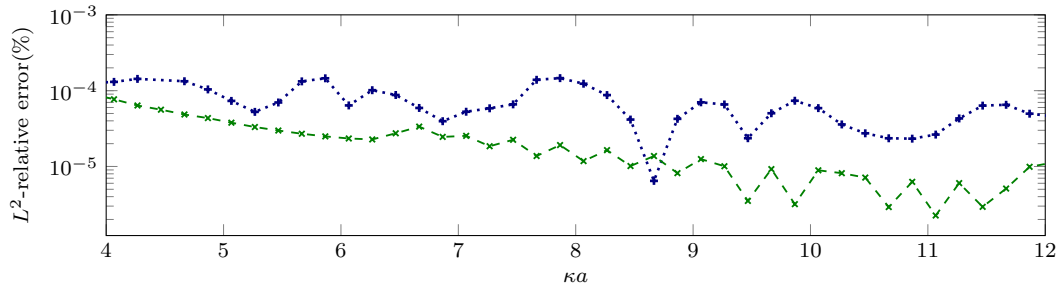


Figure 7: Comparison between the analytical and IPDG values of the Fréchet derivative of a-FFP. $\cdots+$ $\epsilon(\partial_\lambda p_{a,\infty})$, $-x-$ $\epsilon(\partial_\mu p_{a,\infty})$.

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