





Optimization and control of magnetic high fields

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Introduction

Magnetic Fields (T:Tesla)

Static Field:

- Grenoble: 36T, soon 43T,
- Tallahassee, USA: 45T

Applications

- Semiconductors and nanophysics,
- Chemistry, Biochemistry,

Pulsed Field:

- Toulouse: 99T,
- Los Alamos, USA: 100T
- Magneto-science, Levitation,
- Applied Superconductivity



(a) Exterior





(c) Bitter



Applications

Optimisation

- Use experimental measurements to **identify unknown parameters**, such as heat exchange coefficient or conductivities
- Parametrize the form of the magnet to reach a **better homogeneity** of the magnetic field in the region of interest

Output and control quantity of interest

- Use different experimental and geometric parameters to control the magnetic field to create a magnetic structure to confine plasma
- Find parameters such as the magnetic field allow the **levitation** of an object inside the region of interest

Real time computation and visualization of the magnetic field

• Allow the **real time visualization** of the effect of different parameters on the magnetic field



Reduced Order Method

Weak formulation of the model : $a(u(\mu), v; \mu) = f(v; \mu)$

FEM Approximation

$$\underbrace{X^{\mathcal{N}} = \operatorname{span}\{\phi_{1}, \dots, \phi_{\mathcal{N}}\}}_{\text{FEM approximation space}} \longrightarrow \underbrace{u^{\mathcal{N}}(\mu) = \sum_{i=1}^{\mathcal{N}} u^{\mathcal{N}}_{i}\phi_{i}}_{\mathcal{N} \approx 10^{6}} \longrightarrow \underbrace{A^{\mathcal{N}}(\mu)u^{\mathcal{N}}(\mu) = F^{\mathcal{N}}(\mu)}_{\mathcal{N} \times \mathcal{N} \text{ system expensive to solve}}$$

RB Approximation : $u^N(\mu) \approx u^N(\mu)$: linear combination of FEM solution

$$\underbrace{W^{N} = span\{u^{\mathcal{N}}(\boldsymbol{\mu}_{1}), \dots, u^{\mathcal{N}}(\boldsymbol{\mu}_{N})\}}_{\text{RB approximation space}} \rightarrow \underbrace{u^{N}(\boldsymbol{\mu}) = \sum_{i=1}^{N} u_{i}^{N}(\boldsymbol{\mu}) \underline{u}^{\mathcal{N}}(\boldsymbol{\mu}_{i})}_{10 \leqslant N \leqslant 100} \rightarrow \underbrace{A^{N}(\boldsymbol{\mu})u^{N}(\boldsymbol{\mu}) = F^{N}(\boldsymbol{\mu})}_{N \times N \text{ system cheaper to solve}}$$

The reduced basis method

Ingredients

- Set of parameters : \mathcal{D}^{μ}
- FEM 'truth" approximation
 - $X^{\mathcal{N}}$: finite element approximation space of dimension $\mathcal{N}>>1$
 - $u^{\mathcal{N}}(\mu) \in X^{\mathcal{N}}$ is solution of $a(u^{\mathcal{N}}(\mu), v; \mu) = f(v; \mu) \quad \forall v \in X^{\mathcal{N}}$
- RB approximation
 - Sample : $S_N = \{\mu_1 \in \mathcal{D}^{\mu}, ..., \mu_N \in \mathcal{D}^{\mu}\}$
 - Approximation space : $W_N = span\{u^N(\mu_1), ..., u^N(\mu_N)\}$ with N << N
 - Galerkin projection on W_N to determine RB coefficients



$$u^{N}(\boldsymbol{\mu}) = \sum_{i=1}^{N} u_{i}^{N}(\boldsymbol{\mu}) \underline{u}^{\mathcal{N}}(\boldsymbol{\mu}_{i})$$

 $N \times N$ system to solve :

$$\sum_{i=1}^{N} a(\underline{u}^{\mathcal{N}}(\mu_i), v_k; \mu) u_i^{\mathcal{N}}(\mu) = f(v_k; \mu), \ 1 \leq k \leq N$$

Computational opportunities

If the parameter dependance can be expressed as an affine decomposition :

$$a(u, v; \mu) = \sum_{q}^{Q_a} \theta_q^a(\mu) \ a_q(u, v) \quad \text{and} \quad f(v; \mu) = \sum_{q}^{Q_f} \theta_q^f(\mu) \ f_q(v)$$
$$\Rightarrow \sum_{i=1}^{N} \left[\sum_{q=1}^{Q_a} \theta_a^q(\mu) \underbrace{a^q(\underline{u}^{\mathcal{N}}(\mu_i), \underline{u}^{\mathcal{N}}(\mu_j))}_{\text{precomputed}} \right] u_i^N(\mu) = \sum_{q=1}^{Q_f} \theta_f^q(\mu) \underbrace{f^q(\underline{u}^{\mathcal{N}}(\mu_j))}_{\text{precomputed}}$$

Strategy : Offline-Online decomposition

- Computationnally intensive initial preprocessing
- Greatly reduced marginal cost

Multi-Physics Problems

- Monolithic Resolution
- Use Block Affine Decomposition

Non Affine Problems

- Empirical Interpolation Method (EIM) : approximate an expression with an affine decomposition
- Discrete EIM (DEIM) : approximate a discrete operator (matrix or vector) with an affine decomposition

Non Linear Problems

- Use of EIM/DEIM
- Iterative methods : Picard / Newton

a and f depend on non affinely parametrized $\{w_i(u, \mathbf{x}, \boldsymbol{\mu})\}_i$:

$$\mathsf{a}(u(\boldsymbol{\mu}), \boldsymbol{v}; \boldsymbol{\mu}; \{w_i(u, \mathbf{x}; \boldsymbol{\mu})\}_i) = f(\boldsymbol{v}; \boldsymbol{\mu}; \{w_i(u, \mathbf{x}; \boldsymbol{\mu})\}_i) \ \forall \boldsymbol{v} \in X$$

 \Rightarrow Affine decomposition doesn't exists !

Solution : build an affine approximation using Empirical Interpolation Method $w_i(u, \mathbf{x}, \boldsymbol{\mu}) \approx w_{i,M_i}(u, \mathbf{x}; \boldsymbol{\mu}) = \sum_{m=1}^{M_i} \beta_{i,m}^{M_i}(u; \boldsymbol{\mu}) q_{i,m}(\mathbf{x})$ a and f depend on non affinely parametrized $\{w_i(u, \mathbf{x}, \boldsymbol{\mu})\}_i$:

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$$\sum_{q=1}^{Q_{a}^{elm}} \sum_{m=1}^{M_{q}^{q}} \gamma_{a,m}^{q}(u;\mu) \hat{a}_{m}^{q}(u,v) + \sum_{l=1}^{Q_{a}^{elf}} \theta_{a}^{l}(\mu) \hat{a}^{l}(u,v) = \sum_{q=1}^{Q_{f}^{elm}} \sum_{m=1}^{M_{q}^{q}} \gamma_{f,m}^{q}(u;\mu) f_{m}^{q}(v) + \sum_{l=1}^{Q_{f}^{elf}} \theta_{f}^{l}(\mu) f^{l}(v)$$

• Build an affine approximation of a tensor $T(\mu)$

$$T_M(\mu) = \sum_{m=0}^M \beta^m(\mu) \Phi^m \tag{1}$$

where

- $\beta^m(\mu)$ are scalar functions, $\beta^m : \mathcal{D} \to \mathbf{R}$
- Φ^m are tensors independent of μ
- Offline/Online strategy very close to EIM algoritm.
- Can automatically provide the affine decomposition of a parametrized operator
 - Geometric Parameters
 - Complex Operator, ex : stabilization parameter

DEIM : Computation of $\beta^m(\mu)$

At any rank *M*, the vector of coefficient $\Theta_M(\mu) = (\beta^1(\mu), .., \beta^M(\mu))^T$ is the solution of the $M \times M$ system

$$B_{\mathcal{I}_M}\Theta_M(\mu) = T_{\mathcal{I}_M}(\mu) \tag{2}$$

where

- (B_{IM})_{k,l} = (Φ_l)_{ik}: the columns of B_{IM} are composed by the vectors Φ_l in which you only keep the entry for the index i_k ∈ I_M
- $(T_{\mathcal{I}_M}(\mu))_k = (T(\mu))_{i_k} : T_{\mathcal{I}_M}(\mu)$ is the evaluation of the tensor $T(\mu)$ in which you only keep the entry for the index $i_k \in \mathcal{I}_M$

 $B_{\mathcal{I}_M}$ is updated during the greedy algorithm when a new couple (Φ_M, i_m) is selected.

The computation of $T_{\mathcal{I}_M}(\mu)$ has to be done online. It can not depend of the finite element dimension \mathcal{N} !

Applications

ThermoElectric CRB

Problem

From $V = \sum_{n=1}^{N} V_n(\mu) \xi_i^V$ and $T = \sum_{n=1}^{N} T_n(\mu) \xi_i^T$ such that: $\begin{cases}
-\nabla (\sigma_i(T) \nabla V) = 0 & \text{in } \Omega_i, \ i = 1, \dots, P_{mat} \\
-\nabla (k_i(T) \nabla T) = \mathbf{j} \cdot \mathbf{E} \text{ in } \Omega_i, \ i = 1, \dots, P_{mat} \\
V = g_{VD} & \text{on } \Gamma_{VD} \\
-\sigma_i(T) \nabla V \cdot \mathbf{n} = g_{VN} & \text{on } \Gamma_{VN} \\
-k_i(T) \nabla T \cdot \mathbf{n} = g_{TR} & \text{on } \Gamma_{TN} \\
-k_i(T) \nabla T \cdot \mathbf{n} + G_{TR}^1 T = g_{TR}^2 & \text{on } \Gamma_{TN}
\end{cases}$

All the configuration is made through a json file.

Empirical Interpolation Method

We need to use EIM for the conductivities that depends on T and for the right hand side:

$$\sigma(T) \approx \sum_{m_{\sigma}=1}^{M_{\sigma}} \beta_{m_{\sigma}}^{\sigma}(\mu) q_{m_{\sigma}}^{\sigma}, \quad k(T) \approx \sum_{m_{k}=1}^{M_{k}} \beta_{m_{k}}^{k}(\mu) q_{m_{k}}^{k}, \quad \mathbf{j} \cdot \mathbf{E} \approx \sum_{m_{f}=1}^{M_{f}} \beta_{m_{f}}^{f}(\mu) q_{m_{f}}^{f}$$

In the case of experimental parameters, non geometric (σ , k, h,...): $\mathbf{B}(\mathbf{x}; \boldsymbol{\mu}) = \frac{\mu_0}{4\pi} \int_{\Omega_{cond}} \frac{\mathbf{j}(\mathbf{r}; \boldsymbol{\mu}) \times (\mathbf{x} - \mathbf{r})}{|\mathbf{x} - \mathbf{r}|^3} d\mathbf{r} \quad \forall \mathbf{x} \in \Omega_{mgn}$

Using reduced basis for j:

$$\mathbf{j}(\mathbf{r};\boldsymbol{\mu}) = -\sigma(T)\nabla V = -\sum_{m_{\sigma}=1}^{M_{\sigma}} \beta_{m_{\sigma}}^{\sigma}(\boldsymbol{\mu}) q_{m_{\sigma}}^{\sigma}(\mathbf{r}) \sum_{i=1}^{N} V_{i}^{N}(\boldsymbol{\mu})\nabla \xi_{i}^{V}(\mathbf{r})$$

leads to:

$$\mathbf{B}(\mathbf{x};\boldsymbol{\mu}) = \sum_{i=1}^{N} \sum_{m_{\sigma}=1}^{M_{\sigma}} V_{i}^{N}(\boldsymbol{\mu}) \beta_{m_{\sigma}}^{\sigma}(\boldsymbol{\mu}) \underbrace{\frac{\mu_{0}}{4\pi} \int_{\Omega_{cond}} \frac{-q_{m_{\sigma}}(\mathbf{r}) \nabla \xi_{i}^{V}(\mathbf{r}) \times (\mathbf{x}-\mathbf{r})}{|\mathbf{x}-\mathbf{r}|^{3}} d\mathbf{r}}_{\xi_{i,m_{\sigma}}^{\mathsf{B}}(\mathbf{x}) \text{ precomputed}}$$

We can use the same parallel algorithm to compute $\xi_{i,m_{\sigma}}^{\mathbf{B}}(\mathbf{x})$.

Results





Ν

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Results



- Parameters: σ, k, V_D, h, T_w
- Linear case: $\sigma(T) = \sigma(T_0)$ and $k(T) = k(T_0)$
- Dimensions:

| | Ω_{cond} | Ω_{mgn} |
|------|-----------------|----------------|
| dofs | 1350000 | 3000 |

• Times:

| | FE | RB |
|---|------|---------|
| V | 45s | 0.05s |
| В | 100s | 0.0005s |

In the case of geometric parameter, we need to compute:

$$\begin{split} \mathbf{B}(\mathbf{x};\boldsymbol{\mu}) &= \frac{\mu_0}{4\pi} \int_{\phi_{\boldsymbol{\mu}}(\Omega_{cond})} \frac{-\sigma \nabla V \times (\mathbf{x} - \mathbf{r})}{|\mathbf{x} - \mathbf{r}|^3} d\mathbf{r} \quad \forall \mathbf{x} \in \Omega_{mgn} \\ &= \frac{\mu_0}{4\pi} \int_{\Omega_{cond}^{ref}} \frac{-\sigma \nabla V \cdot J_{\phi_{\boldsymbol{\mu}}}^{-1} \times (\mathbf{x} - \phi_{\boldsymbol{\mu}}(\mathbf{r}))}{|\mathbf{x} - \phi_{\boldsymbol{\mu}}(\mathbf{r})|^3} d\mathbf{r} \quad \forall \mathbf{x} \in \Omega_{mgn} \end{split}$$

- How to handle the term $|\mathbf{x} \phi_{\mu}(\mathbf{r})|^3$? We could use EIM for the term $\phi_{\mu}(\mathbf{r})$, but the absolute value would prevent to extract the coefficients for the affine decomposition
 - And we need to compute it for *every* $\mathbf{x} \in \Omega_{mgn}$

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Discrete Empirical Interpolation Method

We approach the magnetic field with:

$$\mathbf{B}(\mathbf{x}; \boldsymbol{\mu}) pprox \sum_{m=1}^{M_B} eta_m(\boldsymbol{\mu}) \mathbf{Q}_m(\mathbf{x})$$

To compute $\beta(\mu)$, we need to solve the following problem:

$$\begin{pmatrix} \xi_{1}(t_{1}) & \dots & \xi_{1}(t_{M}) \\ \vdots & \ddots & \vdots \\ \xi_{M}(t_{1}) & \dots & \xi_{M}(t_{M}) \end{pmatrix} \begin{pmatrix} \beta_{1}(\mu) \\ \vdots \\ \beta_{M}(\mu) \end{pmatrix} = \begin{pmatrix} \mathsf{B}(t_{1};\mu) \\ \vdots \\ \mathsf{B}(t_{M};\mu) \end{pmatrix}$$

So we need to know $\mathbf{B}(t_i; \boldsymbol{\mu})$, where is the advantage ?

We approach the magnetic field with:

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We only need to know **B** on the interpolation points t_i , $i = 1, ..., M_B$

Optimization of the homogeneity

We apply ϕ a geometrical transformation to modify the shape of the helices to attain a better homogeneity

$$\phi = \begin{pmatrix} x \cos(\alpha(\mu)) + y \sin(\alpha(\mu)) \\ -x \sin(\alpha(\mu)) + y \cos(\alpha(\mu)) \\ z \end{cases}$$

where α is a Bezier curve which control points are the parameters μ of the optimization.

| | initial | optimal |
|---|---------|---------|
| $\max \mathbf{B}_z $ | 0.41337 | 0.4051 |
| $\min \mathbf{B}_z $ | 0.41076 | 0.4034 |
| $\frac{\max \mathbf{B}_z }{\min \mathbf{B}_z } - 1$ | 5381.9 | 3544.9 |





Conclusion and Perspectives

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- Non-Linear ThermoElectric model with or without geometric parameters
- Biot & Savart model with or without geometric parameters
- First implementation of optimization problems with or without geometric parameters

Perspectives

- Add Levitation and ions source applications
- Add Maxwell and Elasticity model
- Use HDG methods

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Thank you!

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