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FIRST PASSAGE SETS OF THE 2D CONTINUUM GAUSSIAN FREE FIELD

JUHAN ARU, TITUS LUPU, AND AVELIO SEPÚLVEDA

ABSTRACT. We introduce the first passage set (FPS) of constant level $-a$ of the two-dimensional continuum Gaussian free field (GFF) on finitely connected domains. Informally, it is the set of points in the domain that can be connected to the boundary by a path on which the GFF does not go below $-a$. It is, thus, the two-dimensional analogue of the first hitting time of $-a$ by a one-dimensional Brownian motion. We provide an axiomatic characterization of the FPS, a continuum construction using level lines, and study its properties: it is a fractal set of zero Lebesgue measure and Minkowski dimension 2 that is coupled with the GFF Φ as a local set A so that $\Phi + a$ restricted to A is a positive measure. One of the highlights of this paper is identifying this measure as a Minkowski content measure in the non-integer gauge $r \mapsto |\log(r)|^{1/2}r^2$, by using Gaussian multiplicative chaos theory.

1. INTRODUCTION

The continuum Gaussian free field (GFF) is a canonical model of a Gaussian field satisfying a spatial Markov property. It first appeared in Euclidean quantum field theory, where it is known as bosonic massless free field [Sim74, Gaw96]. In the probability community the study of the 2D continuum GFF has reflowered in the 2000's due to its various connections to Schramm's SLE processes [SS13, She05, Dub09b, MS16a, MS16b, MS16c, MS17], Liouville quantum gravity measures [DS11, BSS14] and Brownian loop-soups [LJ10, LJ11].

The seminal papers that connected SLE processes and the free field showed that SLE_4 can be seen as a level line of the GFF - more precisely, in [SS09] Schramm and Sheffield showed that the level lines of the discrete GFF converge in law to SLE_4 and in [SS13] they gave a purely continuum definition of the limiting coupling, giving rise to the study of level lines of the continuum GFF.

In [ASW17], it was further shown that level lines give a way to define other geometric subsets of the GFF. More precisely, in [ASW17] the authors introduce two-valued local sets $\mathbb{A}_{-a,b}$. Heuristically, the set $\mathbb{A}_{-a,b}$ corresponds to the points in the domain that can be connected to the boundary by some path along which the height of the GFF remains in $[-a, b]$. The mathematical definition of these sets is based on thinking of the 2D GFF as a generalization of the Brownian motion, and it relies on the strong Markov property of the free field. In the case of the BM the same geometric heuristic defines the set $[0, T_{-a,b}]$, where $T_{-a,b}$ is the first exit time from the interval $[-a, b]$, and in [ASW17] it is proved that $\mathbb{A}_{-a,b}$ satisfies many properties expected from the analogy with these exit times.

In the current article, we introduce a further geometric subset of the GFF: the first passage set (FPS) \mathbb{A}_{-a} . Heuristically, it corresponds to the points in the domain that can be connected to the boundary by paths along which the height of the GFF is greater or equal to $-a$, i.e., it is the analogue of $[0, T_{-a}]$, where T_{-a} is the one-sided first passage time of a BM. We provide an axiomatic characterization of the continuum FPS, a construction using iterations of level lines as in [ASW17] and study several of its properties.

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There are two key aspects which make the FPS interesting to study. First of all, compared to most of the geometric subsets of the GFF studied so far, this set is large in the sense that the restriction of the GFF to this set is not zero. This not only requires new ways of working with this set, but also introduces interesting phenomena - as one of the key results we show that even though the GFF on this set is non-trivial, it is measurable with respect to the underlying set. Even more, we show that in fact this measure can be identified with the Minkowski content measure of the underlying set in the gauge $r \mapsto |\log(r)|^{1/2}r^2$. Secondly, in the case of the FPS the geometric definition given above can be made precise in the following sense: in a follow-up article [ALS18c], we show that the first passage sets of the metric graph GFF introduced in [LW16] converge to the continuum FPS. This will, among other things, allow us to identify the FPS with the trace of a clusters of Brownian loops and excursions, and to prove convergence results for the level lines of the GFF. Moreover, in a subsequent work we will use the results of these two papers to prove an excursion decomposition of the 2D continuum GFF [ALS18a].

Finally, let us mention that the FPS have already proved useful in studying Gaussian multiplicative chaos (GMC) measures of the GFF: in [APS17], the authors confirm that a construction of [Aid15] converges to the GMC measure. This was done using the fact that GFF can be approximated by its FPS of increasing levels. Moreover, in [APS18], and using the same FPS based construction, the authors prove that a “derivative” of subcritical GMC measures coincides with a multiple of the critical measure. This confirmed a conjecture of [DRSV14].

1.1. Overview of results. Let us give a more detailed overview of the results presented in the paper. To do this, first recall the local set coupling of a random set A with the Gaussian free field Φ in a domain D . It is a coupling (Φ, A) that induces a Markovian decomposition of Φ . That is to say Φ can be written as a sum between Φ_A and Φ^A , where Φ_A is a random distribution that is a.s. harmonic on $D \setminus A$ and conditional on (A, Φ_A) , Φ^A is a zero boundary GFF on $D \setminus A$. We denote by h_A the harmonic function corresponding to Φ_A outside of A . The local set condition implies that conditional on A and Φ_A , the GFF Φ restricted to $D \setminus A$ is given by the sum of h_A and Φ^A .

The two valued local sets (TVS) $\mathbb{A}_{-a,b}$ for a simply connected domain D , studied in [ASW17], can be then defined as the only thin local sets of the GFF such that $h_A \in \{-a, b\}$. Here, thin means that at Φ_A contains no extra mass on A , i.e. for any smooth function f , we have that $(h_A, f) = (\Phi_A, f)$.

Our first task is to generalize two-valued sets to multiply-connected domains and to more general boundary conditions u , and to show that the main properties of TVS proved in [ASW17] remain true also in this setup. The generalization to more general boundary conditions requires only slight modifications in the definitions, and slight extensions of the proofs. The case of multiply-connected domains, however, requires both new ideas and technical work. In particular, as technical result of independent interest we prove in Proposition 3.16 *that level lines in multiply-connected domains are continuous up to and at hitting a continuation threshold.*

We next introduce the first passage set (FPS). For a zero boundary GFF, the FPS of level $-a$, denoted \mathbb{A}_{-a} with $a \geq 0$ is then defined as a local set of the Gaussian free field on D satisfying the following properties:

- *Conditional on \mathbb{A}_{-a} , the law of the restriction of the GFF on D to $D \setminus \mathbb{A}_{-a}$ is that of a GFF on $D \setminus \mathbb{A}_{-a}$ with boundary condition $-a$, or in other words $h_{\mathbb{A}_{-a}} = -a$.*
- *The GFF on \mathbb{A}_{-a} is larger than $-a$, in the sense that for any positive test function f we have that $(\Phi_{\mathbb{A}_{-a}} + a, f) \geq 0$ - that is to say $\Phi_{\mathbb{A}_{-a}} + a$ is a positive measure.*

The full definition for general boundary conditions u is given in Definition 4.1. There, we also define the FPS in the other direction: \mathbb{V}_b will heuristically correspond to the local set A such that $h_A = b$ and the GFF on A is smaller than b . As proved in Theorem 4.3 in the setting of more general boundary conditions, the first passage set

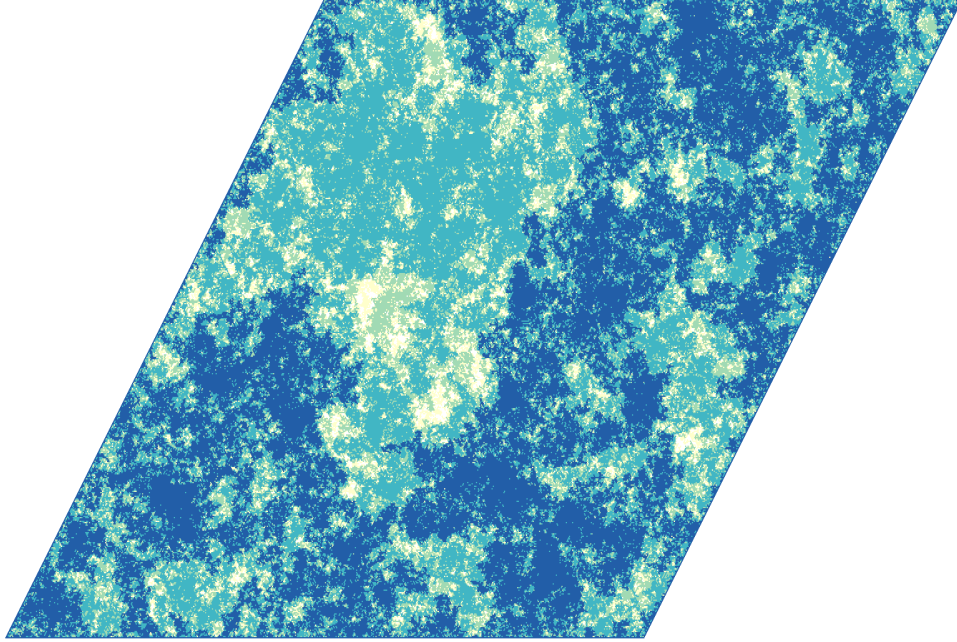


FIGURE 1. A simulation of four nested First passage sets. The first passage set $\mathbb{A}_{-\lambda}$ with $\lambda = \sqrt{\pi/8}$ is in dark blue. The difference between $\mathbb{A}_{-2\lambda}$ and $\mathbb{A}_{-\lambda}$ is in lighter blue, difference between $\mathbb{A}_{-2\lambda}$ and $\mathbb{A}_{-3\lambda}$ in green and yellow depicts the missing part of $\mathbb{A}_{-4\lambda}$. Image done by B. Werness.

- is unique in the sense that any other local set with the above conditions is a.s. equal to the FPS, and thus, it is a measurable with respect to the GFF it is coupled with;
- is monotone in the sense that for all $a \leq a'$ almost surely $\mathbb{A}_{-a} \subset \mathbb{A}_{-a'}$;
- as in the case of the Brownian motion can be constructed as a limit of two-valued local sets, $\mathbb{A}_{-a} = \lim_{b \rightarrow \infty} \mathbb{A}_{-a,b}$.

In fact the relation to two-valued sets is even stronger: we will show that *the intersection of two FPS \mathbb{A}_{-a} and \mathbb{V}_b is precisely $\mathbb{A}_{-a,b}$ in the simply connected case.* Proposition 4.14 generalizes this result from multiply connected domains and it is the key to show the uniqueness of TVS.

One can show that \mathbb{A}_{-a} has zero Lebesgue measure but, contrary to $\mathbb{A}_{-a,b}$, its Minkowski dimension is 2 and \mathbb{A}_{-a} is not a thin local set, i.e., Φ “charges” \mathbb{A}_{-a} . Also, quite surprisingly $\Phi_{\mathbb{A}_{-a}}$ is measurable function of just the set \mathbb{A}_{-a} itself. Even more, this measure is in fact equal to one halves times its Minkowski content measure in the gauge $r \mapsto |\log(r)|^{1/2}r^2$ (Proposition 5.1). Notice that in fact both of these statements are non-trivial! Our proof uses the recent construction of Liouville quantum gravity measures via the local sets [APS17], the fact that GFF is a measurable function of its Gaussian multiplicative chaos measures [BSS14], and a deterministic argument to link different measures on fractal sets, that could turn out to be useful in a more general setting. As a cute consequence we observe in Corollary 5.3 that *the GFF can be seen as a limit of recentered Minkowski content measures of a sequence of growing random sets.*

Finally, our techniques allow us also to compute explicitly the laws of several observables. In Propositions 4.8 and 4.11, we compute the extremal distance between FPS or TVS started from a given boundary to the rest of the boundary. This is the continuum analogue of some of the results obtained in the metric graph setting [LW16], where the extremal distance replaces the effective resistance. In an upcoming paper [ALS18b], we further use these techniques to calculate the law

of the extremal distance between the CLE_4 loop surrounding zero and the boundary, and the joint laws between different nested loops.

1.2. Outline of the paper. The rest of the article is structured as follows:

Section 2 contains the preliminaries: a summary of general potential theory objects, two-dimensional continuum GFF, local sets and basic results about Gaussian multiplicative chaos. The only novel parts are Propositions 2.7, 2.8 that heuristically allow us to parametrize local set processes using their distance to a part of the boundary.

In Section 3, we extend the theory of two-valued local set to the finitely-connected case. This will require a detailed study of the generalized level lines in multiply-connected domains. After that, in Section 4, we define and characterize the continuum FPS and prove several of its properties. Finally, in Section 5 we show that the measure $\Phi_{\mathbb{A}-a} + a$ corresponds to a constant times the Minkowski content measure (in a certain gauge) of the underlying set.

2. PRELIMINARIES

In this section, we describe the underlying objects and their key properties. First, we go over the conformally invariant notion of distance in complex analysis - the extremal length; then we discuss the continuum two-dimensional GFF and its local sets. The only new contribution of this section is Proposition 2.7.

We denote by $D \subseteq \mathbb{C}$ an open planar bounded domain with a non-empty and non-polar boundary. By conformal invariance, we can always assume that D is a subset of the unit disk \mathbb{D} . The most general case that we work with are domains D such that the complement of D has at most finitely many connected component and no complement being a singleton. Recall that by the Riemann mapping for multiply-connected domains [Koe22], such domains D are known to be conformally equivalent to a circle domain (i.e. to $\mathbb{D} \setminus K$, where K is a finite union of closed disjoint disks, disjoint also from $\partial\mathbb{D}$).

2.1. Extremal distance. In multiply-connected domains, the natural way to measure distances between the components is the *extremal length* (it is a particular case of *extremal distance*) and its reciprocal *conformal modulus*. Both of the quantities are conformally invariant and extremal distance is the analogue of the effective resistance on electrical networks [Duf62]. We introduce it shortly here and refer to [Ahl10], Section 4 for more details.

If $\rho(z)|dz|$ is a metric on D conformal equivalent to the Euclidean metric, we will denote by

$$\text{Length}_\rho(\gamma) := \int_\gamma \rho(z)|dz|$$

the ρ -length of a path γ , and by

$$\text{Area}_\rho(D) := \int_D \rho(z)^2 dz$$

the ρ -area of D . The extremal distance between \mathcal{B}_1 and \mathcal{B}_2 is defined as

$$\text{EL}(\mathcal{B}_1, \mathcal{B}_2) = \sup_\rho \inf_{\substack{\gamma \text{ connecting} \\ \mathcal{B}_1 \text{ and } \mathcal{B}_2}} \frac{\text{Length}_\rho(\gamma)^2}{\text{Area}_\rho(D)}.$$

The conformal modulus $M(\mathcal{B}_1, \mathcal{B}_2)$ is then defined as $\text{EL}(\mathcal{B}_1, \mathcal{B}_2)^{-1}$. We state here also a theorem giving an explicit formula for the extremal distance using the Dirichlet energy .

Theorem 2.1 (Theorem 4-5 of [Ahl10]). *Let D be finitely connected, \mathcal{B}_1 and \mathcal{B}_2 be unions of finitely many boundary arcs, such that $d(\mathcal{B}_1, \mathcal{B}_2) > 0$. Then $M(\mathcal{B}_1, \mathcal{B}_2) = \text{EL}(\mathcal{B}_1, \mathcal{B}_2)^{-1}$ is given by the Dirichlet energy $\int_D |\nabla \bar{u}|^2$ of the harmonic function \bar{u} equal to 0 on \mathcal{B}_1 , 1 on \mathcal{B}_2 , and having*

zero normal derivative on $\partial D \setminus (\mathcal{B}_1 \cup \mathcal{B}_2)$. If \mathcal{B}_1 , resp. \mathcal{B}_2 , has piecewise smooth boundary $\int_D |\nabla \bar{u}|^2$ is equal to $-\int_{\mathcal{B}_1} \partial_n \bar{u}$, resp. $\int_{\mathcal{B}_2} \partial_n \bar{u}$, where ∂_n is the outward derivative.

This theorem gives in particular a relation between the extremal distance and the boundary Poisson kernel. To explain this, we define the Green's function G_D of the Laplacian (with Dirichlet boundary conditions) in D . It is often useful to write

$$(2.1) \quad G_D(z, w) = (2\pi)^{-1} \log(1/|z - w|) + g_D(z, w),$$

where $g_D(z, \cdot)$ is the bounded harmonic function with boundary values given by $(2\pi)^{-1} \log(|z - x|)$ for $x \in \partial D$. It can be shown that the Green's function is conformally invariant. Additionally, note that in simply connected domains, $g_D(z, z)$ equals the log conformal radius:

$$g_D(z, z) = \frac{1}{2\pi} \log(\text{CR}(z, D)).$$

The Green's function can be used to define the Poisson kernel. In the case of domains with locally analytic boundary, the boundary Poisson kernel is defined as

$$H_D(x, y) = \partial_{n_x} \partial_{n_y} G_D(x, y), \quad x, y \in \partial D,$$

where ∂_{n_x} respectively ∂_{n_y} are the normal derivatives at x respectively y . If D and D' are domains with locally analytic boundaries and f is a conformal transformation from D to D' , then

$$(2.2) \quad H_{D'}(f(x), f(y)) = |f'(x)| |f'(y)| H_D(x, y).$$

One can see the boundary Poisson kernel as a measure on $\partial D \times \partial D$ rather than a function, by setting

$$H_D(dx, dy) = H_D(x, y) dx dy,$$

where on the right-hand side dx and dy denote the length measure on ∂D . This measure is conformal invariant by (2.2). Also note that it has infinite total mass due to diagonal divergence. For a domain D with general boundary, we can define the boundary Poisson kernel $H_D(dx, dy)$ as the push-forward measure of the boundary Poisson kernel on a domain D' with locally analytic boundary, under a conformal transformation taking $D' \rightarrow D$. This is true even in the case ∂D has locally infinite length, e.g. in the case where the boundary "looks like" an SLE₄ curve.

Finally, notice that from the definition using the Green's function and Theorem 2.1 we see that the extremal length introduced above can be expressed using the boundary Poisson kernel. Indeed, let \mathcal{B} be a union of finitely many boundary components. Then

$$M(\mathcal{B}, \partial D \setminus \mathcal{B}) = \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B})^{-1} = \iint_{\mathcal{B} \times \partial D \setminus \mathcal{B}} H_D(dx, dy).$$

In general, if $\mathcal{B}_1, \mathcal{B}_2 \subseteq \partial D$ are disjoint,

$$M(\mathcal{B}_1, \mathcal{B}_2) = \text{EL}(\mathcal{B}_1, \mathcal{B}_2)^{-1} \geq \iint_{\mathcal{B}_1 \times \mathcal{B}_2} H_D(dx, dy).$$

2.2. The continuum GFF. The (zero boundary) Gaussian Free Field (GFF) in a domain D can be viewed as a centered Gaussian process Φ indexed by the set of continuous functions with compact support in D , with covariance given by the Green's function:

$$\mathbb{E}[(\Phi, f_1)(\Phi, f_2)] = \iint_{D \times D} f_1(z) G_D(z, w) f_2(w) dz dw.$$

In this paper Φ always denotes the zero boundary GFF. We also consider GFF-s with non-zero Dirichlet boundary conditions - they are given by $\Phi + u$ where u is some bounded harmonic function that is piecewise constant¹ boundary data on ∂D .

¹Here and elsewhere this means piecewise constant that changes only finitely many times

Because the covariance kernel of the GFF blows up on the diagonal, it is impossible to view Φ as a random function. However, it can be shown that the GFF has a version that lives in the Sobolev space $H^{-1}(D)$ of generalized functions, justifying the notation (Φ, f) for Φ acting on functions f (see for example [Dub09b]). In fact, one can explicitly calculate that if Φ is a GFF in $D \subseteq \mathbb{D}$ then:

$$(2.3) \quad \mathbb{E} \left[\|\Phi\|_{H^{-1}(\mathbb{D})}^2 \right] = \iint_{\mathbb{D} \times \mathbb{D}} G_D(z, w) G_{\mathbb{D}}(z, w) dz dw \leq \iint_{\mathbb{D} \times \mathbb{D}} G_{\mathbb{D}}^2(z, w) dz dw < \infty.$$

Moreover, let us remark that it is in fact possible and useful to define the random variable (Φ, μ) for any fixed Borel measure μ , provided the energy $\iint \mu(dz) \mu(dw) G_D(z, w)$ is finite.

Finally, there is an important numerical constant λ , that depends on the normalization of the GFF and is in the current setting (where $G_D(z, w) \sim -(2\pi)^{-1} \log(|z - w|)$ as $w \rightarrow z$) is given by

$$\lambda = \sqrt{\pi/8},$$

where 2λ is the *height gap* of the GFF [SS09, SS13]. Sometimes, other normalizations are used in the literature: if $G_D(z, w) \sim c \log(1/|z - w|)$ as $z \rightarrow w$, then λ should be taken to be $(\pi/2) \times \sqrt{c}$.

2.3. Local sets: definitions and basic properties. Let us now introduce more thoroughly the local sets of the GFF. We only discuss items that are directly used in the current paper. For a more general discussion of local sets and thin local sets (not necessarily of bounded type), we refer to [SS13, Wer16, Sep17].

Even though, it is not possible to make sense of (Φ, f) when $f = \mathbf{1}_A$ is the indicator function of an arbitrary random set A , local sets form a class of random sets where this is (in a sense) possible.

Definition 2.2 (Local sets). *Consider a random triple (Φ, A, Φ_A) , where Φ is a GFF in D , A is a random closed subset of \bar{D} and Φ_A a random distribution that can be viewed as a harmonic function when restricted to $D \setminus A$. We say that A is a local set for Φ if conditionally on (A, Φ_A) , $\Phi^A := \Phi - \Phi_A$ is a GFF in $D \setminus A$.*

Throughout this paper, we use the notation $h_A : D \rightarrow \mathbb{R}$ for the function that is equal to Φ_A on $D \setminus A$ and 0 on A .

Let us list a few properties of local sets (see for instance [SS13, Aru15, AS18a] for derivations and further properties):

Lemma 2.3.

- (1) *Any local set can be coupled in a unique way with a given GFF: Let $(\Phi, A, \Phi_A, \widehat{\Phi}_A)$ be a coupling where (Φ, A, Φ_A) and $(\Phi, A, \widehat{\Phi}_A)$ satisfy the conditions of this Definition 2.2. Then, a.s. $\Phi_A = \widehat{\Phi}_A$. Thus, being a local set is a property of the coupling (Φ, A) , as Φ_A is a measurable function of (Φ, A) .*
- (2) *If A and B are local sets coupled with the same GFF Φ , and (A, Φ_A) and (B, Φ_B) are conditionally independent given Φ , then $A \cup B$ is also a local set coupled with Φ . Additionally, $B \setminus A$ is a local set of Φ^A with $(\Phi^A)_{B \setminus A} = \Phi_{B \cup A} - \Phi_A$.*
- (3) *Let (Φ, A_n) be such that for all $n \in \mathbb{N}$, (Φ, A_n) is a local set coupling, the sets A_n are increasing in n and there exists $k \in \mathbb{N}$ such that the cardinal of connected components of $A_n \cup \partial D$ is bounded by k . Then, $\overline{\bigcup A_n}$ is also a local set and $\Phi_{A_n} \rightarrow \Phi_{\overline{\bigcup A_n}}$ in probability as $n \rightarrow \infty$.*
- (4) *Let (Φ, A_n) be such that for all $n \in \mathbb{N}$, (Φ, A_n) is a local set coupling and the sets A_n are decreasing in n . Then, $\bigcap A_n$ is also a local set and $\Phi_{A_n} \rightarrow \Phi_{\bigcap A_n}$ a.s. as $n \rightarrow \infty$.*

The property (3) follows from the fact that under the conditions on A_n Beurling estimate (Theorem 3.76 of [Law08]) ensures that $G_{D \setminus A_n} \rightarrow G_{D \setminus A}$ as $n \rightarrow \infty$. Property (4) just follows from inverse martingale theorem.

Often one is interested in a growing sequence of local sets, which we call local set processes.

Definition 2.4 (Local set process). *We say that a coupling $(\Phi, (\eta_t)_{t \geq 0})$ is a local set process if Φ is a GFF in D , $\eta_0 \subseteq \partial D$, and η_t is an increasing continuous family of local sets such that for all stopping time τ of the filtration $\mathcal{F}_t := \sigma(\eta_s : s \leq t)$, (Φ, η_t) is a local set.*

Let us note that in our definition η_t is actually a random set. In the rest of the paper, we are mostly interested in local set processes that are equal to the trace of a continuous curve. In those cases, we are going to denote by $\eta(t)$ the tip of the curve at time t . In other words, in our notation $\eta_t = \eta([0, t])$.

Local processes can be naturally parametrized from the viewpoint of any interior point z : the expected height $h_{\eta_t}(z)$ then becomes a Brownian motion. More precisely, we have that:

Proposition 2.5 (Proposition 6.5 of [MS16a]). *For any $z \in D$ if $(\eta_t)_{t \geq 0}$ is parametrized such that $(G_D - G_{D \setminus \eta_t})(z, z) = t$, then $(h_{\eta_t}(z))_{t \geq 0}$ has the law of a Brownian motion.*

Remark 2.6. *Notice that whereas G_D diverges on the diagonal, the difference of Green's functions can be given a canonical sense, using (2.1). In fact when D , and $D \setminus \eta_t$ are simply connected domains, it is a difference of logarithms of conformal radii:*

$$(G_D - G_{D \setminus \eta_t})(z, z) = \frac{1}{2\pi} \log(\text{CR}(z, D)) - \frac{1}{2\pi} \log(\text{CR}(z, D \setminus \eta_t)).$$

In fact, one can also parametrize local set processes η_t using their distance to the boundary. As the boundary values of the GFF away from η_t do not change, it is natural to look at normal derivatives. In order to obtain a conformally invariant quantity, notice that if $\mathcal{B} \subseteq \partial D$ and h is a harmonic function, then by Green's identities the quantity $\int_{\mathcal{B}} \partial_n h$ can be given a conformally invariant meaning: $\int_{\mathcal{B}} \partial_n h = \int_D \nabla h \nabla \bar{u}$ where \bar{u} is the harmonic extension of the function that takes the value 1 on \mathcal{B} and 0 on $\partial D \setminus \mathcal{B}$.

We will first consider the case, where the local set process is parametrized by its extremal distance to a whole boundary component:

Proposition 2.7. *Let D be finitely connected circle-domain and (Φ, η_t) be a local set process with Φ a GFF in D . Take $\mathcal{B} \subseteq \partial D$ be a union of finitely many boundary components. Then, if η_t is parametrized by its conformal modulus, i.e. such that*

$$t = M(\mathcal{B}, (\partial D \cup \eta_t) \setminus \mathcal{B}) - M(\mathcal{B}, \partial D \setminus \mathcal{B})$$

then $W_t := \int_{\mathcal{B}} \partial_n h_{\eta_t}$ has the law of a standard Brownian motion started from 0.

Equivalently, when parametrized by the extremal length

$$t = \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) - \text{EL}(\mathcal{B}, (\partial D \cup \eta_t) \setminus \mathcal{B}),$$

the process

$$\widehat{W}_t := \text{EL}(\mathcal{B}, (\partial D \cup \eta_t) \setminus \mathcal{B}) \int_{\mathcal{B}} \partial_n h_{\eta_t},$$

has the law of a Brownian bridge from 0 to 0 with length $\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B})$.

Moreover, the same holds (with the appropriate definitions) for any D finitely connected domain with all boundary components larger than a point.

Proof. As the conformal modulus is the reciprocal of the extremal length, a simple calculation shows that it is enough to prove the first claim.

Using the conformal invariance both of the quantity W_t , the Gaussian free field and the extremal length, it suffices to work in a circle domain and consider the case where \mathcal{B} is equal to the union of circles. In fact, for simplicity, we will only prove the case where \mathcal{B} is equal to a circle of radius 1. We can moreover assume that the circles of radius $1 + \varepsilon$, denoted by \mathcal{B}_ε are contained in D for all $\varepsilon > 0$ small enough. In this case $\varepsilon^{-1} \int_{\mathcal{B}_\varepsilon} h_{\eta_t}$ converges a.s. to $\partial_n h_{\eta_t}$.

Now write $h_{\eta_t} = \Phi - \Phi^{\eta_t}$, and note that for any $\varepsilon > 0$ we have that:

$$\begin{aligned}\sigma_\varepsilon^2 &:= \mathbb{E} \left[(\varepsilon^{-1} \int_{\mathcal{B}_\varepsilon} \Phi)^2 \right] = \varepsilon^{-2} \iint_{\mathcal{B}_\varepsilon \times \mathcal{B}_\varepsilon} G_D(z, w) dz dw, \\ \sigma_{\varepsilon, t}^2 &:= \mathbb{E} \left[(\varepsilon^{-1} \int_{\mathcal{B}_\varepsilon} \Phi^{\eta_t})^2 \mid \eta_t \right] = \varepsilon^{-2} \iint_{\mathcal{B}_\varepsilon \times \mathcal{B}_\varepsilon} G_{D \setminus \eta_t}(z, w) dz dw.\end{aligned}$$

Note that when $\varepsilon \rightarrow 0$, both terms individually diverge. However, their difference converges to

$$(2.4) \quad \iint_{\mathcal{B} \times \mathcal{B}} \partial_{n_y} \partial_{n_x} (G_D(x, y) - G_{D \setminus \eta_t}(x, y)) dx dy.$$

Moreover, this convergence is uniform in all sets η_t that are at extremal distance of at least δ to \mathcal{B} for some $\delta > 0$,

Now, by Theorem 2.1 and the fact that $\partial_{n_x} G_D(x, y) dx$ is the Poisson kernel we have that (2.4) is equal to $M(\mathcal{B}, (\partial D \cup \eta_t) \setminus \mathcal{B}) - M(\mathcal{B}, \partial D \setminus \mathcal{B})$. Further, note that a.s. $\varepsilon^{-1} \int_{\mathcal{B}} h_{\eta_t}(z) \rightarrow \int_{\mathcal{B}} \partial_n h_{\eta_t}(z) dz$, and that for all $\lambda \in \mathbb{R}$, by first conditioning on η_t ,

$$e^{-\frac{\lambda^2}{2}(\sigma_\varepsilon^2 - \sigma_{\varepsilon, t}^2)} \mathbb{E} \left[\exp \left(\lambda i \varepsilon^{-1} \int_{\mathcal{B}} h_{\eta_t}(z) dz \right) \mid \eta_t \right] = e^{-\sigma_\varepsilon^2} \mathbb{E} \left[(\varepsilon^{-1} \int_{\mathcal{B}_\varepsilon} \Phi)^2 \mid \eta_t \right].$$

We conclude by bounding $\sigma_\varepsilon^2 - \sigma_{\varepsilon, t}^2$ by $M(\mathcal{B}, \partial D \setminus \mathcal{B}) - M(\mathcal{B}, (\partial D \cup \eta_t) \setminus \mathcal{B}) \pm \delta$ and taking expected value. \square

Now, let us see how to extend this proposition to the case where the local set process is parametrized by its ‘‘distance’’ to a part of the boundary. One of the obstacles here is that when the growing set that we want to parametrize is of 0 Euclidean distance from this part of the boundary, then the naive conformal modulus between the set and this boundary part diverges. However, similarly to the reduced extremal distance (see Chapter 4.14 of [Ahl66]), the difference between the moduli is still non-trivial.

Let us explain this a bit more precisely: assume that ∂D can be partitioned as $\mathcal{B}_1 \cup \mathcal{B}_2$, where \mathcal{B}_2 is a connected subset of ∂D (that is not necessarily a whole boundary component). Let $B \subseteq \bar{D}$ a closed set that remains at positive distance from \mathcal{B}_2 and that intersects \mathcal{B}_1 . Define as u_0 to be the bounded harmonic function that takes value 0 on \mathcal{B}_2 and 1 on \mathcal{B}_1 , and u_B to be the bounded harmonic function taking again value 0 on \mathcal{B}_2 but is equal to 1 on $\mathcal{B}_1 \cup B$. Then $u_0 - u_B$ is a bounded harmonic function that takes values 0 in $\mathcal{B}_1 \cup \mathcal{B}_2$ and $u_0 - 1$ in B . In particular,

$$\int_{\mathcal{B}_2} \partial_n (u_0 - u_B) < \infty.$$

even when the conformal modulus between \mathcal{B}_1 and \mathcal{B}_2 is 0. Thus, we can define $M(\mathcal{B}_2, \mathcal{B}_1) - M(\mathcal{B}_2, \mathcal{B}_1 \cup B)$ as $\int_{\mathcal{B}_2} \partial_n (u_0 - u_B)$ even when both terms individually are infinite.

Using this observation, it is possible to use a proof similar to that of Proposition 2.7 to parametrize local sets using only a part of a boundary component:

Proposition 2.8. *Let D be a finitely connected circle domain, (Φ, η_t) a local set process with Φ a GFF in D . Now, let us partition ∂D in two sets: \mathcal{B}_1 and \mathcal{B}_2 a connected subset of $\partial D \setminus \mathcal{B}_1$ of positive length and suppose that $\eta_0 \subseteq \mathcal{B}_1$.*

Then, if η_t is parametrized by the difference of conformal moduli (as the individual conformal moduli may not exist)

$$t = M(\mathcal{B}_2, \mathcal{B}_1) - M(\mathcal{B}_2, \mathcal{B}_1 \cup \eta[0, t]) := \int_{\mathcal{B}_2} \partial_n (\bar{u}_0 - \bar{u}_t),$$

and \bar{u}_t is defined as the harmonic function taking the value 1 on \mathcal{B}_2 and the value 0 on $\mathcal{B}_1 \cup \eta[0, t]$. Then, the process $W_t = \int_{\mathcal{B}_2} \partial h_{\eta[0, t]}$ has the law of a Brownian motion.

Moreover, the same holds (with the appropriate definitions) for any finitely connected domain with all boundary components larger than a point.

2.4. Gaussian multiplicative chaos. Finally, let us summarize the definition and some properties of the Gaussian multiplicative chaos (GMC) associated to the GFF. The GMC measures were first introduced in the realm of self-interacting Euclidean field theories [HK71], and named by Kahane in his seminal article [Kah85]. We refer to e.g. [Aru17] for more detailed proofs and properties for the GMC measures, and to [RV14] for an overview of the GMC measures and their applications.

To define the GMC measure, one usually passes through an approximation procedure. Denote by $\Phi_\varepsilon(z)$ the circle-average process of the GFF: i.e. the GFF tested against the unit measure on the circle of radius ε around z . For $\varepsilon > 0$ and $\gamma \in \mathbb{R}$, we then set

$$d\mu_\gamma^\varepsilon := \varepsilon^{\frac{\gamma^2}{2} 2\pi} \exp\left(\gamma\sqrt{2\pi}\Phi_\varepsilon(z)\right) dz.$$

The $\sqrt{2\pi}$ factor comes from the fact that in the GMC literature the GFF is normalized differently (i.e. usually with covariance of that behaves like $-\log|z-w|$ near the diagonal).

In [DS11] it was shown that for $\gamma \in (-2, 2)$, as ε goes to 0 along the dyadics, μ_γ^ε converges towards a measure μ_γ a.s. and in \mathbb{L}^1 .

For us an important property is that in the subcritical regime the GMC measures depend analytically on γ . In fact, the following theorem (which can be, for example, found in Section 1 of [Aru17]) suffices our needs:

Proposition 2.9. *For any $\gamma \in (-\sqrt{2}, \sqrt{2})$ and any continuous compactly supported function f on D , there exists a modification of $((\mu_\gamma^\varepsilon, f))_{|\gamma| < \sqrt{2}}$, and a deterministic sequence $\varepsilon_k \rightarrow 0$ such that a.s. and in \mathbb{L}^2 , $((\mu_\gamma^{\varepsilon_k}, f))_{|\gamma| < \sqrt{2}}$ converges in the space $\mathcal{C}^1((-\sqrt{2}, \sqrt{2}))$ to $((\mu_\gamma, f))_{|\gamma| < \sqrt{2}}$.*

In particular, the map $\gamma \mapsto (\mu_\gamma, f)$ is \mathcal{C}^1 on $(-\sqrt{2}, \sqrt{2})$ and furthermore

$$(2.5) \quad (\Phi, f) = \lim_{k \rightarrow +\infty} (\Phi_{\varepsilon_k}, f) = \frac{1}{\sqrt{2\pi}} \lim_{k \rightarrow +\infty} \partial_\gamma(\mu_\gamma^{\varepsilon_k}, f) |_{\gamma=0} = \frac{1}{\sqrt{2\pi}} \partial_\gamma(\mu_\gamma, f) |_{\gamma=0},$$

where all the limits are a.s. and in \mathbb{L}^2 .

Remark 2.10. *In fact, in [Aru17] this is explained for a different normalization, instead of μ_γ , one considers $\tilde{\mu}_\gamma$ defined by*

$$d\tilde{\mu}_\gamma^\varepsilon := \exp\left(\gamma\sqrt{2\pi}\Phi_\varepsilon(z) - \frac{\gamma^2}{2} 2\pi \mathbb{E}[\Phi_\varepsilon(z)^2]\right) dz.$$

Since $\mu_\gamma(dz) = \text{CR}(z)^{\frac{\gamma^2}{2} 2\pi} \tilde{\mu}_\gamma(dz)$, Proposition 2.9 for μ_γ is equivalent to that for $\tilde{\mu}_\gamma$.

Remark 2.11. *Note that from (2.5) it follows that Φ is a deterministic measurable function of $(\mu_\gamma)_{|\gamma| < \sqrt{2}}$ (see e.g. Corollary 1.6 in [Aru17]). The measurability of the field w.r.t. its GMC measure was further studied in [BSS14], where the authors show that in fact Φ is a deterministic function of μ_γ for any fixed $\gamma \in (-2, 2)$.*

3. TWO-VALUED LOCAL SETS

Next, we discuss a specific type of local sets introduced in [ASW17]: two-valued local sets. In [ASW17], these sets were defined and studied in the case of the zero boundary GFF; we will extend this definition to general boundary conditions and to n -connected domains. We will also calculate the size of the set seen from interior points and from boundary components that do not intersect the set.

First, it is convenient to review a larger setting, that of bounded type local sets (BTLS) introduced in [ASW17]. These sets are thin local set A , for which its associated harmonic function h_A remains bounded. Here, by a thin local set (see [Wer16, Sep17]) we mean the following condition:

- For any smooth test function $f \in \mathcal{C}_0^\infty$, the random variable (Φ, f) is almost surely equal to $(\int_{D \setminus A} h_A(z) f(z) dz) + (\Phi^A, f)$.

This definition assumes that h_A belongs to $\mathbb{L}^1(D \setminus A)$ which is the case in our paper. For the general definition see [Sep17]. In order to say that the union of two thin sets is thin, it is more convenient to use a stronger condition. Indeed, it is not hard to show that (see Proposition 4.3 of [Sep17] for a proof):

- If h_A is $\mathbb{L}^1(D \setminus A)$ and for any compact set $K \subseteq D$, the Minkowski dimension of $A \cap K$ is strictly smaller than 2 then A is thin.

Now, we can define the bounded type local sets.

Definition 3.1 (BTLS). *Consider a closed subset A of \bar{D} and Φ a GFF in D defined on the same probability space. Let $K > 0$, we say that A is a K -BTLS for Φ if the following four conditions are satisfied:*

- (1) A is a thin local set of Φ .
- (2) Almost surely, $|h_A| \leq K$ in $D \setminus A$.
- (3) Almost surely, each connected component of A that does not intersect ∂D has a neighborhood that does intersect no other connected component of A .

If A is a K -BTLS for some K , we say that it is a BTLS.

3.1. Generalized level lines. One of the simplest family of BTLS are the generalized level lines, first described in [SS13], that correspond to $\text{SLE}_4(\rho)$ processes.

Let $D := \mathbb{H} \setminus \bigcup_{k=1}^n C_k$, where $(C_k)_{1 \leq k \leq n}$ is a finite family of disjoint closed disks, be a circle domain in the upper half plane. Further, let u be a harmonic function in D . We say that $\eta(\cdot)$, a curve parametrized by half plane capacity, is the generalized level line for the GFF $\Phi + u$ in D up to a stopping time τ if for all $t \geq 0$:

- (**): The set $\eta_t := \eta[0, t \wedge \tau]$ is a BTLS of the GFF Φ , with harmonic function satisfying the following properties: $h_{\eta_t} + u$ is a harmonic function in $D \setminus \eta_t$ with boundary values $-\lambda$ on the left-hand side of η_t , $+\lambda$ on the right side of η_t , and with the same boundary values as u on ∂D .

The first example of level lines comes from [SS13]: let u_0 be the unique bounded harmonic function in \mathbb{H} with boundary condition $-\lambda$ in \mathbb{R}^- and λ in \mathbb{R}^+ . Then it is shown in [SS13] that there exists a unique η satisfying (**) for $\tau = \infty$, and its law is that of an SLE_4 .

Several subsequent papers [SS13, MS16a, WW16, PW17b] have studied more general boundary data in simply-connected case and also level lines in a non-simply connected setting [ASW17]. The following lemma is a slight variant of the latter, stating existence of level lines until it either accumulate at another component, or hit the continuation threshold² on \mathbb{R} . It is a consequence of Theorem 1.1.3 of [WW16] and Lemma 15 of [ASW17].

Lemma 3.2 (Existence of generalized level line targeted at ∞). *Let u be a bounded harmonic function with piecewise constant boundary data such that $u(0^-) < \lambda$ and $u(0^+) > -\lambda$. Then, there exists a unique law on random simple curves $(\eta(t), t \geq 0)$ coupled with the GFF such that (**) holds for the function u and possibly infinite stopping time τ that is defined as the first time when η hits or accumulates at a point $x \in \partial D \setminus \mathbb{R}$ or hits a point $x \in \mathbb{R}$ such that $x \geq 0$ and $u(x^+) \leq -\lambda$ or*

²The continuation threshold is the first time τ in which the level line hits a boundary point, $x \in \partial D$, such that there is no level line of $\phi + u + h_{\eta_\tau}$ starting from x in O . Here O is the non bounded connected component of $D \setminus \eta_\tau$. This condition can be described explicitly using boundary values. See, for example, Definition 2.14 in [WW16].

$x \leq 0$ and $u(x^-) \geq \lambda$. We call η the generalized level line for the GFF $\Phi + u$ and it is measurable w.r.t Φ . Moreover, η_t if $\eta(\tau) \in \mathbb{R}$, then η is continuous on $[0, \tau]$.

Remark 3.3. In Section 3.3 we will be able to show that η is continuous up to τ even if $\eta(\tau) \notin \mathbb{R}$.

In simply-connected domains Theorem 1.1.3 of [WW16] and Lemma 15 of [Dub09a] also give us precise information on the subset of the boundary where the level line can hit:

Proposition 3.4 (Hitting of level lines in simply-connected domains). *Let $D = \mathbb{H}$ and let u be a bounded harmonic function with piecewise constant boundary data such that $u(0^-) < \lambda$ and $u(0^+) > -\lambda$. Let η be a generalized level line in D starting from 0. If either $u \geq \lambda$ or $u \leq -\lambda$ on some open interval $J \subset \partial\mathbb{D}$, then η_t stays at a positive distance of any $x \in J$. Moreover, if $u \leq -\lambda$ on a neighborhood in \mathbb{R} to the left of x , or $u \geq \lambda$ on a neighborhood in \mathbb{R} to the right of x , then almost surely η_t stays at a positive distance of x .*

Remark 3.5. Observe that the boundary points described by this lemma correspond exactly to points, from where one cannot start a generalized level line of $-\Phi - u$.

A simple, but important corollary of this result allows us to check whether a level line can enter a connected component of the complement of a bounded-type local set. This observation was key in [ASW17], where it was used for simply-connected domains and it followed just from the facts that 1) a generalized level line does not hit itself; 2) it has to exit such a component in finite time. We will prove the generalization of this lemma to finitely-connected setting below in Lemma 3.17; the same proof could be used in the simply-connected setting.

Lemma 3.6. *Let η be a generalized level line of a GFF $\Phi + u$ in D as above and A a BTLS of Φ conditionally independent of η . Take $z \in D$ and define $O(z)$ the connected component of $D \setminus A$ containing z . On the event where on any connected component of $\partial O(z)$ the boundary values of $(h_A + u)|_{O(z)}$ are either everywhere $\geq \lambda$ or everywhere $\leq -\lambda$, we have that a.s. $\eta([0, \infty]) \cap O(z) = \emptyset$.*

In Section 3.3, we extend all these results to finitely-connected domains, in particular, we extend the definition of generalized level lines by showing that they remain continuous until the stopping time τ . That is to say, that level lines remain continuous up to its accumulation point, even if it is on other boundary component. To do this, we will however first have to gain a better understanding of certain type of BTLS in simply-connected domains, called two-valued local sets.

3.2. Two-valued local sets in simply connected domains. Another family of useful BTLS is that of two-valued local sets. In [ASW17], two-valued local sets of the zero boundary GFF were introduced in the simply connected case, which we assume to be \mathbb{D} for convenience. Two-valued local sets are thin local sets A such that the harmonic function h_A takes precisely two values. More precisely, take $a, b > 0$, and consider thin local sets $\mathbb{A}_{-a,b}$ coupled with the GFF such that $h_{\mathbb{A}_{-a,b}}$ is constant in each connected component of $\mathbb{D} \setminus \mathbb{A}_{-a,b}$ and for all $z \in \mathbb{D} \setminus \mathbb{A}_{-a,b}$, $h_{\mathbb{A}_{-a,b}} \in \{-a, b\}$. It is somewhat more convenient to assume that the two-valued local sets and first passage sets introduced later also by convention contain the boundary.

[ASW17] dealt with the construction, measurability, uniqueness and monotonicity of two-valued local sets in the case of the zero boundary GFF in simply connected domains. Here we state a slight generalization of this main theorem for more general boundary values.

In this respect, let u be a bounded harmonic function with piecewise constant boundary values. Take $a, b > 0$ and define $u^{-a,b}$ to be the part of the boundary where the values of u are outside of $[-a, b]$. As long as $u^{-a,b}$ is empty, the harmonic function h_A still takes only two values $-a$ and b . Otherwise, we also allow for components where some of the boundary data for h_A (corresponding to $u^{-a,b}$) is not equal to $-a$ or b . More precisely, the complement of the two-valued set $\mathbb{A}_{-a,b}^u$ has two types of components O :

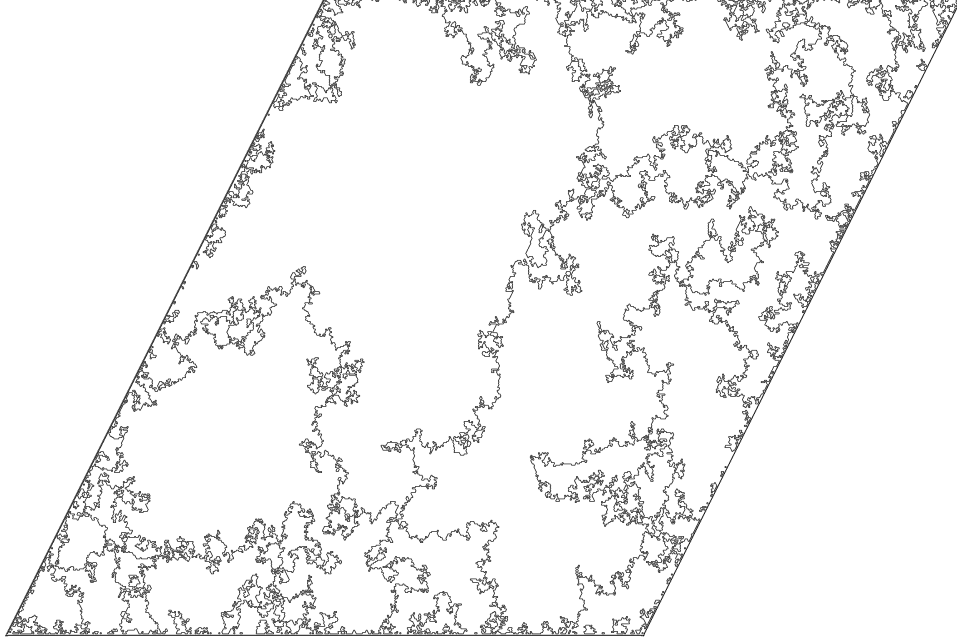


FIGURE 2. Simulation of $\mathbb{A}_{-\lambda, \lambda}$ done by B. Werness.

- (1) Those where $\partial O \cap \partial \mathbb{D}$ is a totally disconnected set. In these components $h_{\mathbb{A}_{-a, b}^u} + u$ takes the constant value: $-a$ or b .
- (2) Those where $\partial O \cap \partial \mathbb{D} \subset u^{-a, b}$. In these components $h_{\mathbb{A}_{-a, b}^u} + u$ takes boundary values u on the part $\partial O \cap \partial u^{-a, b}$ and has either constant boundary value $-a$ or b on the rest of ∂O , in such a way that $h_{\mathbb{A}_{-a, b}^u} + u$ is a bounded harmonic function that is either greater or equal to b or smaller or equal to $-a$ throughout the whole component.

The next proposition basically says that all the properties of the zero-boundary case generalize to the general boundary:

Proposition 3.7. *Consider a bounded harmonic function u as above. If $|a + b| \geq 2\lambda$ and $[\min(u), \max(u)] \cap (-a, b) \neq \emptyset$, then it is possible to construct $\mathbb{A}_{-a, b}^u \neq \emptyset$ coupled with a GFF Φ . Moreover, the sets $\mathbb{A}_{-a, b}^u$ are*

- *unique in the sense that if A' is another BTLS coupled with the same Φ , such that a.s. it satisfies the conditions above, then $A' = \mathbb{A}_{-a, b}^u$ almost surely;*
- *measurable functions of the GFF Φ that they are coupled with;*
- *monotone in the following sense: if $[-a, b] \subset [-a', b']$ with $b + a \geq 2\lambda$, then almost surely, $\mathbb{A}_{-a, b}^u \subset \mathbb{A}_{-a', b'}^u$.*

The proof is an extension of the proof of Proposition 2 and the arguments in Sections 6.1 and 6.2 of [ASW17]:

Proof. CONSTRUCTION: We know from [ASW17] that the condition $|a + b| \geq 2\lambda$ is necessary. Also, if $[\min(u), \max(u)] \cap (-a, b) \neq \emptyset$ does not hold, then the empty set satisfies our conditions. Thus, suppose that $|a + b| \geq 2\lambda$ and $[\min(u), \max(u)] \cap (-a, b) \neq \emptyset$. Notice that as soon as we have constructed the basic sets with $|b + a| = 2\lambda$, the rest of the proof follows exactly as in the construction in Section 6.2 of [ASW17] - one just iterates inside the components. Moreover, in this basic case, one can only concentrate on $\mathbb{A}_{-\lambda, \lambda}^u$ as for any other (a, b) with $b = -a + 2\lambda$ it is enough to construct $\mathbb{A}_{-\lambda, \lambda}^{u+a-\lambda}$.

Let us now build $\mathbb{A}_{-\lambda,\lambda}^u$. To do this, partition the boundary $\partial\mathbb{D} = \bigcup_{k=1}^n \mathcal{B}_k$ such that each \mathcal{B}_k is a finite segment, throughout each \mathcal{B}_k the function u is either larger or equal to λ , smaller or equal to $-\lambda$, or is contained in $(-\lambda, \lambda)$, and n is as small as possible. Call n the *boundary partition size*. Notice that n is finite by our assumption. We will now show the existence by induction on n .

In fact, the heart of the proof is the case $n = 2$, so we will start from this. If u is, say, larger than λ on \mathcal{B}_1 and smaller than $-\lambda$ on \mathcal{B}_2 , then by Lemma 3.2 we can draw a generalized level line from one point in $\partial\mathcal{B}_1$ to the other one, by Proposition 3.4 it almost surely finishes at the other point of $\partial\mathcal{B}_1$ and decomposes the domain into components satisfying (2).

So suppose u is larger than λ on \mathcal{B}_1 but in $(-\lambda, \lambda)$ on \mathcal{B}_2 . Then, we can similarly start a generalized level line from one point in $\partial\mathcal{B}_1$ targeted to the other one. Again, we know that it finishes there almost surely. It will decompose the domain into one piece that satisfies the condition (2) and possibly infinitely many simply connected pieces that have a boundary partition size equal to 2. We can iterate the level line in each of these components. Now for any $z \in \mathbb{D} \cap \mathbb{Q}^2$, denoting the local set process arising from the construction and continuing always in the connected component containing z by A_t^z , we have (say from Proposition 2.5) that $h_{A_t}(z)$ is a martingale. We claim that from this it follows that any z is in a component satisfying (1) or (2) above after drawing a finite number of generalized level lines. Indeed, fix some $z \in \mathbb{D} \cap \mathbb{Q}^2$; then any level line iterated in a component containing z that stays on the same side of z than the previous level line will have a larger harmonic measure than the previous one; as the sign of the level line facing z changes, we see that $h_{A_t}(z)$ changes by a bounded amount. This can happen only a finite number of times and thus the claim. Hence we have shown the construction in the case $n = 2$.

Now, if $n = 1$, then the only possible case is that u takes values in $(-\lambda, \lambda)$. In this case the generalized level lines can be started and ended at all points of the boundary. In choosing any two different points on the boundary and drawing a level line, we will decompose D into simply-connected components such that their boundary partition size equal to 2 (See Figure 3).

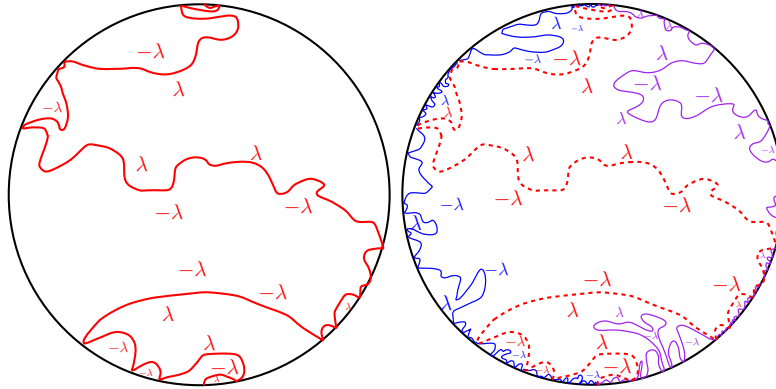


FIGURE 3. An illustration of the first two steps in the construction of $\mathbb{A}_{-\lambda,\lambda}^u$. Does it look similar to Figure 2?

For $n \geq 3$, we must have at least two \mathcal{B}_k , say \mathcal{B}_1 and \mathcal{B}_2 (not necessarily adjacent) such that $|u| \geq \lambda$ on them. We then start our generalized level line from a possible starting point in $\partial\mathcal{B}_1$ towards a possible target point in $\partial\mathcal{B}_2$. By Lemmas 3.2 and Proposition 3.4 it stops at either at its target point or at a point between two $\mathcal{B}_i, \mathcal{B}_j$ such that on both $|u| \geq \lambda$. One can verify that in each of these cases, in each component cut out the boundary partition size is strictly smaller than n . \square

Let us now make the following remarks:

- (i) In the construction we only need to use level lines of $\Phi + u$ whose boundary values are in $[-a, b]$.
- (ii) For a fixed point $z \in \mathbb{D}$ a.s. we only need a finite number of level lines to construct the connected component of $\mathbb{D} \setminus \mathbb{A}_{-a,b}^u$ containing z .
- (iii) As none of the level lines is started inside $u^{-a,b}$ nor can touch $u^{-a,b}$, any connected component of $u^{-a,b}$ belongs entirely to the boundary of a single connected component of $\mathbb{D} \setminus \mathbb{A}_{-a,b}^u$. In particular, each connected component O of $\mathbb{D} \setminus \mathbb{A}_{-a,b}^u$ with $\partial O \cap \partial \mathbb{D} \neq \emptyset$ has only finitely many intersection points and by Lemma 10 in [ASW17] we can assign them any values, in particular those that u already takes on $\partial \mathbb{D}$.

Proof. UNIQUENESS, MEASURABILITY AND MONOTONICITY: follow exactly as in the zero-boundary case, i.e. one follows the proof of the Proposition 2 in [ASW17]. The only small difference is in the uniqueness part: The first step is to show that $\mathbb{A}_{-a,b}^u \subset A'$. To do this one uses the construction and Lemma 3.3 as in the zero boundary case. Now, to show the opposite inclusion, notice that by conditions (1) and (2) we have that in any connected component O of $\mathbb{D} \setminus \mathbb{A}_{-a,b}^u$ the boundary values of $h_{\mathbb{A}_{-a,b}^u} + u$ are either larger or equal to b or smaller or equal to $-a$. In particular,

$$|h_{\mathbb{A}_{-a,b}^u} + u + a/2 - b/2| \geq |h_{A'} + u + a/2 - b/2|.$$

Thus, given that $A' = A' \cup \mathbb{A}_{-a,b}^u$ we can just use Lemma 9 of [ASW17], where instead of k we use $u + a/2 - b/2$ (the proof is exactly the same in this case). □

Remark 3.8. *In fact in the monotonicity statement, one could also include the changes in the harmonic function: if u, u' are two bounded harmonic function with piecewise constant boundary data, take $\delta = \sup_{\partial \mathbb{D}} |u - u'|$ if $[-a, b] \subset [-a' + \delta, b' - \delta]$ with $b + a \geq 2\lambda$, then almost surely, $\mathbb{A}_{-a,b}^u \subset \mathbb{A}_{-a',b'}^{u'}$. The proof just follows from the construction and Lemma 3.3.*

In [AS18b], the authors studied further properties of the TVS of the GFF in a simply connected domain with constant boundary condition. Let us mention two results that are important for us in analyzing generalized level lines in finitely-connected domains. First, it was proved in [AS18b] that $\mathbb{A}_{-\lambda,\lambda}^u$ is a.s. equal to the union of all level lines of $\Phi + u$ (see Lemma 3.6 of [AS18b]). The exact same proof works in this context and implies that:

Lemma 3.9. *Let Z be any countably dense set of points on ∂D . Then $\mathbb{A}_{-\lambda,\lambda}^u$ is equal to the closure of the union of all the level lines of $\Phi + u$ going between two different points of Z .*

Remark 3.10. *Note that the level line of $\Phi + u$ going from x to y is equal to the level line of $-\Phi - u$ going from y to x (Theorem 1.1.6 of [WW16]).*

Second, the TVS $\mathbb{A}_{-a,b}$ are locally finite for $a, b \leq 2\lambda$, again for zero boundary GFF (Proposition 3.15 of [AS18b]). A rather direct generalization is as follows:

Lemma 3.11. *In simply connected domains, $\mathbb{A}_{-\lambda,\lambda}^u$ is locally finite, i.e., a.s. for each $\varepsilon > 0$ there are only finitely many connected components of $D \setminus \mathbb{A}_{-\lambda,\lambda}^u$ with diameter larger than ε .*

This result will be proved in more generality in Proposition 4.15 (relying on the results of the paper [ALS18c]), but we will sketch a direct argument here too:

Proof sketch: First assume that the boundary condition changes twice and that one boundary value corresponds to either $-\lambda$ or λ , and the other to a constant $v \in (-\lambda, \lambda)$. Let us argue that in this case $\mathbb{A}_{-\lambda,\lambda}^u$ is locally finite, by reducing it to the case of $\mathbb{A}_{-\lambda,\lambda}^v$ (that has the same law as $\mathbb{A}_{-\lambda-v,\lambda-v}$). Indeed, consider Φ a GFF in \mathbb{D} and η a generalized level line of $\Phi + v$ joining two different boundary points. Then, $\mathbb{A}_{-\lambda,\lambda}^v$ restricted to any connected component O of $D \setminus \eta^v$ is given by $\mathbb{A}_{-\lambda,\lambda}^u$ of Φ^O .

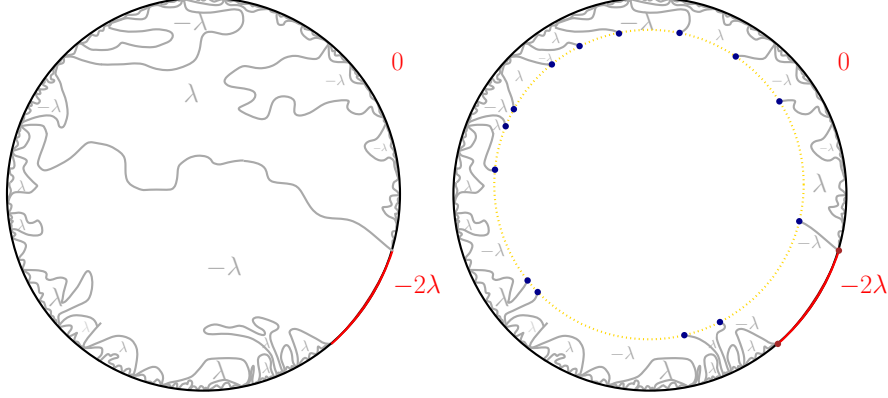


FIGURE 4. To the left the set $\mathbb{A}_{-\lambda, \lambda}^u$ where u is equal to -2λ on an arc, and equal to 0 on the rest of the boundary. To the right the set A described in Lemma 3.12. Note that the boundary values of $h_A + u$ on A are λ , $-\lambda$, equal to u on the parts of the boundary where $u \notin (-\lambda, \lambda)$ and can only change its sign at the points drawn in blue, i.e. at the points where A intersects $\partial(1 - \varepsilon)\mathbb{D}$.

But by Proposition 3.15 of [AS18b], we know that $\mathbb{A}_{-\lambda-v, \lambda-v}$ is locally finite, and moreover by following the proof of that proposition, we can deduce that if $\mathbb{A}_{-a, b}^u$ is locally finite inside some simply-connected set, it is locally finite in all simply-connected sets.

Now, the proof follows from two observations:

- By uniform continuity, a continuous curve in \mathbb{D} parametrized by $[0, 1]$ only separates finitely many components of diameter larger than ε for any $\varepsilon \geq 0$;
- From the construction of TVS with piecewise-constant boundary conditions given above, we see that after drawing a finite number of level lines (that are continuous up to their endpoint by Proposition 3.4) we can construct a local set $A \subset \mathbb{A}_{-\lambda, \lambda}^u$ such that the connected components of $D \setminus A$ either already correspond to those of $D \setminus \mathbb{A}_{-\lambda, \lambda}^u$, or where the boundary condition is as above: it changes twice between a part that is $\pm\lambda$, and a part where it takes a value in $(-\lambda, \lambda)$. \square

These two lemmas allow us to explicitly describe a part of $\mathbb{A}_{-\lambda, \lambda}$ in an ε -neighborhood of the boundary with a local set:

Lemma 3.12. *Fix $\varepsilon > 0$ and let Z be any countably dense set of points on $\partial\mathbb{D}$. Define A as closed union of the level lines of $\Phi + u$ and $-\Phi - u$ starting in \mathbb{D} and stopped the first time they reach distance ε from $\partial\mathbb{D}$. Let \widehat{A} be equal to the union of the connected components of $\mathbb{A}_{-\lambda, \lambda}^u \setminus ((1 - \varepsilon)\mathbb{D})$ that are connected to $\partial\mathbb{D}$. Then almost surely, A is equal to \widehat{A} . Furthermore, A is a local set such that $A \cap \partial((1 - \varepsilon)\mathbb{D})$ is a finite set of points.*

If we define O to be the connected component of $\mathbb{D} \setminus A$ containing 0 , then the boundary values of $h_A + u$ are, in absolute value bigger than or equal to λ . Furthermore, they are piece-wise constant, change their value only finitely many times, and change their sign only at points situated on $\partial((1 - \varepsilon)\mathbb{D})$.

Proof. A is a local set as it can be written as the closed union of countable many local sets. Now, by Lemma 3.9, the set $\mathbb{A}_{-\lambda, \lambda}^u$ is equal to the union of all level lines with starting and endpoints in Z . But taking the intersection of this union of level lines with $\mathbb{D} \setminus (1 - \varepsilon)\mathbb{D}$, and throwing away parts that are not connected to the boundary, gives exactly the union of level lines stopped at distance ε from ∂D .

Thus, to show that $A \cap \partial(1 - \varepsilon)\mathbb{D}$ is a finite set of points, it suffices to prove this claim for \widehat{A} , the connected component of $\mathbb{A}_{-\lambda, \lambda}^u \cap (\mathbb{D} \setminus ((1 - \varepsilon)\mathbb{D}))$ that is connected to $\partial\mathbb{D}$. Notice that this number of intersection points is bounded by twice the number of “excursions” of $\mathbb{A}_{-\lambda, \lambda}^u$ between two boundary points that intersect $(1 - \varepsilon)\mathbb{D}$, where by an excursion we mean a connected set of $\mathbb{A}_{-\lambda, \lambda}^u \cap \mathbb{D}$ that intersects $\partial\mathbb{D}$ only at its two endpoints. However, by Lemma 3.11 $\mathbb{A}_{-\lambda, \lambda}^u$ is a.s. locally finite. This implies that there are a.s. also only finitely many excursions of $\mathbb{A}_{-\lambda, \lambda}^u$ intersecting $(1 - \varepsilon)\mathbb{D}$, as one can associate to each such excursion a unique connected component of $\mathbb{D} \setminus \mathbb{A}_{-\lambda, \lambda}^u$ (for example the component that is separated from the point 1 by this excursion).

□

3.3. Generalized level lines in finitely connected domains. In this section, we prove that the generalized level lines are continuous up to their stopping time τ , i.e. that any accumulation point described in Lemma 3.2 is in fact a hitting point. WLoG, let us work in the circle domain $D = \mathbb{H} \setminus \bigcup_{k=1}^n C_k$, where $(C_k)_{1 \leq k \leq n}$ is a finite family of disjoint closed disks.

We first need to extend Proposition 3.4 to finitely-connected domains, showing that we cannot accumulate near the points that cannot be hit by a generalized level line:

Lemma 3.13 (Non-intersecting regime). *Let η be a generalized level line in D starting from 0 and $x \in D \setminus \mathbb{R}$. If either $u \geq \lambda$ or $u \leq -\lambda$ on some open interval $J \subset \partial D$, then a.s. η_t stays at a positive distance of any $x \in J$.*

Moreover, if $u \geq \lambda$ on a neighborhood in ∂D clockwise of x , or $u \leq -\lambda$ on a neighborhood in ∂D counter-clockwise of x , then almost surely η_t stays at a positive distance of x .

Proof. Let us start from the first claim. To do this, consider some compact interval $J' \subset J$ and parametrize η_t using the conformal modulus $t = M(J', \eta_t \cup \partial\mathbb{H}) - M(J', \partial\mathbb{H})$ as in Proposition 2.8. Let \bar{u}_t be the harmonic function that equals to 1 on J' and to 0 on $\eta_t \cup \partial\mathbb{H}$.

WLoG assume that the boundary values on J' satisfy $u \geq \lambda$. Suppose for contradiction that η comes arbitrarily close to J' . This means by Proposition 2.8 that $W_t = \int_{J'} \partial_n h_{\eta_t}$ is a Brownian motion on $[0, \infty)$. We argue that $\int_{J'} \partial_n h_{\eta_t}$ is bounded from below. Notice that this would give the desired contradiction, as then W_t would be a Brownian motion that remains bounded from below for all times.

Let u' be the harmonic function that is equal to λ in J and to u in $\partial D \setminus J$. Because $u' \leq u$, the harmonic function $h_{\eta_t} + u'$ has boundary values on η_t that are smaller than or equal to λ . Let \bar{h}_t be the bounded harmonic function with boundary values equal to λ in $\eta_t \cup J$ and to $\sup |u| + 1 > \lambda$ in the rest of the boundary. As the minimum of $\bar{h}_t - (h_{\eta_t} + u')$ is attained in J , we have that $\int_{J'} \partial_n (h_{\eta_t} + u') \geq \int_{J'} \partial_n \bar{h}_t$. To finish the first claim, we note that $\int_{J'} \partial_n u' < \infty$, and $\partial_n \bar{h}_t$ is increasing in t .

For the second claim, assume WLoG that $u \geq \lambda$ on some neighborhood in ∂D clockwise of x . Let us denote this neighborhood by J_L . If $u \geq \lambda$ also on a neighborhood counter-clockwise, then we are done by the first claim. Otherwise, let $\varepsilon > 0$ be very small, in particular smaller than the distance between any boundary components. Notice that we can start a generalized level line $\hat{\eta}$ of $\Phi + u$ from x stopped at $\hat{\tau}_\varepsilon$ the first time the level line reaches distance ε from the connected component it started from, or has reached its continuation threshold on the same connected component of ∂D that contains x . Note that $\hat{\eta}$ cannot hit J_L before time $\hat{\tau}_\varepsilon$ by absolute continuity w.r.t. to the simply-connected case (see Lemma 16 of [ASW17]). On the other hand, it a.s. hits the counter-clockwise boundary neighborhood if $u \in (-\lambda, \lambda)$ on this neighborhood, and a.s. does not hit it if $u \leq -\lambda$.

In both cases, by the first claim of this proposition, we have that η cannot intersect nor accumulate in $J_L \cup \hat{\eta}_t$ before either accumulating or hitting $\hat{\eta}(\hat{\tau})$, or finishing by accumulating somewhere on the boundary ∂D at positive distance from x . We now argue that η cannot accumulate at $\hat{\eta}(\hat{\tau})$

without hitting. Indeed, because Lemma 3.2 allows us to further grow $\hat{\eta}$ for a positive amount s after time $\hat{\tau}$, the first claim of this proposition implies again that η cannot accumulate near $\hat{\eta}(\hat{\tau})$, before hitting or accumulating $\hat{\eta}(\hat{\tau} + s)$. So it remains to argue that η cannot hit $\hat{\eta}(\hat{\tau})$. This now follows, as hitting $\hat{\eta}(\hat{\tau})$ would in particular mean that η is continuous up to this hitting time τ_1 and thus stay in some open neighborhood of $\hat{\eta}(\hat{\tau})$ after some time $s < \tau_1$. Thus, by the absolute continuity of the GFF in this neighborhood, i.e., the same proof as for Lemma 16 of [ASW17], one can show that η cannot hit $\hat{\eta}(\hat{\tau})$. \square

The following lemma tells us that when level lines approach the boundary they hit it in only one point and thus they are continuous until that time. This implies that, it is not possible for level lines to accumulate in the boundary.

Proposition 3.14. *Let η be a generalized level line of GFF $\Phi + u$ in D starting from 0. Let T_p be the set of boundary points on $\partial D \setminus \mathbb{R}$ that η could potentially hit, i.e. set of points x for which $u > -\lambda$ in a neighborhood in ∂D counter-clockwise of x and $u < \lambda$ in some neighbourhood in ∂D clockwise of x . Then a.s. either η hits T_p or stays at a positive distance from it.*

It is useful to observe that the set T_p is precisely the set of points for which one can start a level line of $-\Phi - u$. The following lemma says that when we start simultaneously a generalized level line of $\Phi + u$ and a level line of $-\Phi - u$ from different boundary components, then these level lines either agree on a continuous curve joining two boundary conditions or stay at a positive distance from each other.

Lemma 3.15. *Let η be a level line of $\Phi + u$ started at $x \in \mathbb{R}$ and let η' be a level line of $-\Phi - u$ started at $y \in D \setminus \mathbb{R}$, stopped at times τ, τ' respectively, that correspond to the first times they intersect a connected component of ∂D different to their starting point (or when they hit the continuation threshold on their own component). Then, either $\eta \cap \eta' = \emptyset$ or $\eta \cap \eta'$ is a connected set that intersects \mathbb{R} and the connected component of ∂D that contains y . In the latter case, the level lines η and η' are continuous until the first time they intersect a connected component of the boundary that does not contain their respective starting point.*

Proof. Let us note that for any fixed $s \in \mathbb{R}$, η stopped at the first time it hits η'_s (or at τ) is a level line of $\Phi + u + h_{\eta'_s}$. Thus, by Lemma 3.13 the only point it can hit or accumulate in η'_s is in its tip: $\eta'(s)$. To argue it cannot accumulate at $\eta'(s)$, we can use the same argument as in the proof of the second claim in Lemma 3.13: we can continue the level line η' for another positive time u to see that $\eta(t)$ cannot accumulate at $\eta'(s)$ before accumulating or hitting at $\eta(s+u)$. Thus, η can either hit $\eta'(s)$, or stay at a positive distance of η'_s .

Now, let us note that there is no rational s such that that $\eta'(s) \notin \eta$, but $\eta'_s \cap \eta \neq \emptyset$. Indeed, if there were such a time s , then η would hit a point in η'_s different from $\eta'(s)$. It is clear that the same holds when we switch the roles of η' and η . Thus, we see that if η, η' intersect at some time-points s, s' respectively, then $\eta'([s', \tau']) \subset \eta$ and $\eta([s, \tau]) \subset \eta'$. But this means in fact that in this case $\eta([s, \tau]) = \eta'([s', \tau'])$, where σ, σ' are respectively the last times before τ, τ' where the level lines η, η' touch the component of the boundary containing their starting point. Moreover, we see that if η and η' intersect they hit the same points but in inverse orders, and are thus both continuous up to and including the hitting time of the other boundary. \square

We can now prove Proposition 3.14.

Proof of Proposition 3.14. Fix $\varepsilon > 0$ very small (say much smaller than the minimal distance between two connected components of $\mathbb{C} \setminus D$). By Lemma 3.12 and the absolute continuity of the GFF (Proposition 13 and Corollary 14 of [ASW17]), for all C_k we can construct a local set A_k in an ε -neighbourhood of C_k such that its boundary values are either $\geq \lambda$ or $\leq -\lambda$ in some open interval around any boundary point of A_k that is of distance smaller than ε to the boundary. By Lemma

3.13 the generalized level line η of $\Phi + u$ started at 0 will stay at a positive distance from all these points. Thus, it can only get infinitely close to one of C_k by first hitting or accumulating at one of the points of A_k that is exactly at distance ε from a boundary component C_k . Notice that by Lemma 3.12 there are only finitely many of such points. Moreover, each such point belongs to either a level line of $\Phi + u$ or a level line of $-\Phi - u$ started from the boundary component C_k . By Lemma 3.13, we know that the level line η will stay at a positive distance from the first type of points. Finally, by Lemma 3.15 we know that if it gets infinitely close to one of the other type of points, it will actually agree with this level line until hitting the boundary component containing its starting point. □

For further reference, let us resume the existence and continuity of level lines in finitely-connected domains in a proposition - it follows directly by combining Lemmas 3.2 and 3.13 with Proposition 3.14.

Proposition 3.16 (Existence and continuity of generalized level lines). *Let u be a bounded harmonic function with piecewise constant boundary data such that $u(0^-) < \lambda$ and $u(0^+) > -\lambda$. Then, there exists a unique law on random simple curves $(\eta(t), t \geq 0)$ coupled with the GFF such that (**) holds for the function u and possibly infinite stopping time τ that is defined as the first time when η hits a point $x \in \partial D \setminus \mathbb{R}$ or hits a point $x \in \mathbb{R}$ such that $x \geq 0$ and $u(x^+) \leq -\lambda$ or $x \leq 0$ and $u(x^-) \geq \lambda$. We call η the generalized level line for the GFF $\Phi + u$, and it is measurable w.r.t Φ .*

Moreover, η_t is continuous on $[0, \tau]$, and it can only hit points from which we can start a generalized level line of $-\Phi - u$ (on $\partial D \setminus \mathbb{R}$ these are the points x such that $u < \lambda$ on some interval of ∂D clockwise of x and $u > -\lambda$ on some interval of ∂D counterclockwise of x).

As a consequence, we can prove the generalization of Lemma for finitely-connected domains, the only additional input being the continuity up to its stopping time, and the precise description of the hitting points stated in the proposition above.

Lemma 3.17. *Let η be a generalized level line of a GFF $\Phi + u$ in D as above and A a BTLS of Φ conditionally independent of η . Take $z \in D$ and define $O(z)$ the connected component of $D \setminus A$ containing z . On the event where on any connected component of $\partial O(z)$ the boundary values of $(h_A + u)|_{O(z)}$ are either everywhere $\geq \lambda$ or everywhere $\leq -\lambda$, we have that a.s. $\eta([0, \infty]) \cap O(z) = \emptyset$.*

Notice that we allow for the situation where on some boundary components the value is $\geq \lambda$ and on some it is $\leq -\lambda$. One of the key ingredients is Lemma 2.3 (2), that reads as follows: if A, B are conditionally independent local sets, then the boundary values on $A \cup B$ do not change at any point that is of positive distance of the boundary of $\partial A \cap \partial B$.

Proof. Define as E^z the event where on any connected component of $\partial O(z)$ the boundary values of $(h_A + u)|_{O(z)}$ are either everywhere $\geq \lambda$ or everywhere $\leq -\lambda$. Suppose for contradiction that on the event E^z , $\eta([0, \infty]) \cap O(z) \neq \emptyset$ with positive probability.

Take $\varepsilon > 0$ and $\tau := \tau(\varepsilon)$ the first time such that $\eta(\tau) \in O(z)$ and is at distance $\varepsilon > 0$ of $\partial O(z)$. Note that under our assumption for small enough ε , the event $\{\tau < \infty\} \cap E^z$ has non-zero probability. One can verify that $\bar{\eta}(t) := \eta(t + \tau)$ is a generalized level line of $\Phi^{A \cup \eta_\tau} + (\Phi_{A \cup \eta_\tau} + u)$. As the generalized level-line is a simple continuous curve, it stays at a positive distance of $\partial O(z) \cap \eta_\tau$. Additionally, from Lemma 2.3 (2) it follows that the boundary values of $(\Phi_{A \cup \eta_\tau} + u)$ are $\geq \lambda$ or $\leq -\lambda$ around any point on $\partial O(z)$ that is at positive distance from η_τ . Thus, by Lemma 16 of [ASW17], $\bar{\eta}$ cannot exit $O(z)$ through any point that is at positive distance of η_τ . But, by Proposition 3.16, we have that $\bar{\eta}(\infty)$ ends at a point on ∂D that is different from any of its previously visited points, staying continuous up to and at the moment it hits the boundary, giving a contradiction. □

A particularly useful corollary is the following.

Corollary 3.18. *Let A be a local set and η a generalized level line, both coupled with the same GFF. Then, η cannot enter a connected component of $D \setminus A$ inside which h_A is smaller than the boundary values of h_η .*

3.4. Two-valued local sets in n -connected domains. In order to study the FPS in n -connected domains, it is useful to first extend the definition of two-valued sets to this setting. In this respect, let u be a bounded harmonic function with piece-wise constant boundary values. Recall that by $u^{-a,b}$ we denoted the part of ∂D where the values of u are outside of $[-a, b]$. The two-valued set $\mathbb{A}_{-a,b}^u$ in n -connected domains is then a BTLS such that in each connected component O of $D \setminus \mathbb{A}_{-a,b}^u$ the bounded harmonic function $h_{\mathbb{A}_{-a,b}^u}$ satisfies the following conditions:

- (♣) On every boundary component of $\partial O \setminus u^{-a,b}$ the harmonic function $h_{\mathbb{A}_{-a,b}^u} + u$ takes constant value a or $-b$, and in $\partial O \cap u^{-a,b}$ it takes the value u . Additionally, we require that in every connected component of ∂O either $h_{\mathbb{A}_{-a,b}^u} + u \leq -a$ or $h_{\mathbb{A}_{-a,b}^u} + u \geq b$ holds.

Note that in particular when $-a \leq \inf u \leq \sup u \leq b$, the condition (♣) simplifies to $h_{\mathbb{A}_{-a,b}^u} + u$ takes constant values $-a$ or b in every connected component of ∂O . Now, we announce the fundamental proposition for two-valued sets in n -connected domains.

Theorem 3.19. *Let D be an n -connected domain and u be a bounded harmonic function that has piecewise constant boundary values. Consider $-a < 0 < b$, with $a + b \geq 2\lambda$. Then, $\mathbb{A}_{-a,b}^u$ exist and are non-empty if $[\min(u), \max(u)] \cap (-a, b) \neq \emptyset$. Moreover, $\mathbb{A}_{-a,b}^u$ satisfy the following properties:*

- They are unique in the sense that if A' is another BTLS coupled with the same Φ , such A' satisfies (♣) almost surely, $\mathbb{A}_{-a,b}^u = A'$.
- They are all measurable functions of the GFF Φ that they are coupled with.
- They are monotonic in the following sense: if $[a, b] \subset [a', b']$ and $-a < 0 < b$ with $b + a \geq 2\lambda$, then almost surely, $\mathbb{A}_{-a,b}^u \subset \mathbb{A}_{-a',b'}^u$.
- For each $\mathbb{A}_{-a,b}^u$, there are at most n connected components of $D \setminus \mathbb{A}_{-a,b}^u$ which are not simply connected.

We will now prove the existence and measurability. To prove the uniqueness, we will first in fact prove the uniqueness of the FPS in the next section; it will be a consequence of Lemma 4.14. Finally, monotonicity follows from uniqueness as in the proof of Proposition 2 of [ASW17]. Until having proved uniqueness, we mean by $\mathbb{A}_{-a,b}^u$ always the set constructed just below.

The proof of existence is in its spirit very similar to the proof of Proposition 3.7, that itself is modeled after Section 6.1 of [ASW17]. However, we do need extra arguments to treat the multiply-connected setting.

Proof. CONSTRUCTION Again, we can assume that we are in the non-trivial case, in other words that $[\min(u), \max(u)] \cap (-a, b) \neq \emptyset$. As in the proof of Proposition 3.7, it suffices to construct $\mathbb{A}_{-\lambda,\lambda}^u$.

This time we need a double induction. Let N be the number of boundary components and as in proof of Proposition 3.7. We take the minimal partition of any boundary component \mathcal{B}_i as $\mathcal{B}_i = \bigcup_{k=1}^{n_i} \mathcal{B}_i^k$, such that each \mathcal{B}_i^k is a finite segment, throughout each \mathcal{B}_i^k the function u is either larger or equal to λ , smaller or equal to $-\lambda$, or contained in $(-\lambda, \lambda)$. Recall that we called n_i boundary partition size of \mathcal{B}_i . We will now use induction on pairs $(N, \max_{i \leq N} n_i)$.

The case $(1, n)$ is given by Proposition 3.7. The key case is $(N, 2)$, so let us prove this by inducting on the number of boundary components N .

On any \mathcal{B}_i satisfying $|u| \geq \lambda$, draw a generalized level line starting from one point of $\partial \mathcal{B}_i^1$ to the other. If it hits some other boundary component, we have reduced the number of boundary components in each of the domains cut out and we can use induction hypothesis. Otherwise, by

Proposition 3.16, it ends at the other point of $\partial\mathcal{B}_i^1$ and reduces the boundary partition size of this boundary component to 1. Hence we can suppose that the only boundary components with boundary partition size equal to 2 have one part with $u \in (-\lambda, \lambda)$.

Now, pick any such component, say, \mathcal{B}_1 and suppose u is larger than λ on \mathcal{B}_1^1 . Then, we can start a generalized level line from one points on $\partial\mathcal{B}_1^1$ towards the other one. If the generalized level line hits some other component or cuts the domain into subdomains with strictly less than m boundary components, we can use induction hypothesis. Otherwise, we have finished all components O such that $\partial O \cap u^{-a,b}$ is non-zero. It now remains to see that all ‘inner components’ are also finished in finite time. This follows similarly to the proof of Proposition 3.7 by using the fact that $h_{A_t}(z)$ is a bounded martingale and converges almost surely.

Suppose now (N, n) satisfy $N \geq 2$, $n \geq 3$. Then we can similarly to the proof of Proposition 3.7 pick a generalized level line on some boundary component with boundary partition size bigger than 2 such that by drawing it we either reduce the boundary partition size to 2 for any subdomain with $N \geq 2$, or reduce the number of boundary components in each subdomain. Using a finite number of such lines we have reduced to either $(N, 2)$, or $(N', 3)$ with N' smaller than 3.

It remains to treat the case $(N, 1)$, if all components satisfy $|u| \geq \lambda$, we are done. Otherwise, in any component with $u \in (-\lambda, \lambda)$ we can start a level line from any point for some short amount of time. This will either reduce the setting to $(N, 3)$, $(N, 2)$ or reduce the number of boundary components. □

Examining closely the proof the following holds:

- (i) In the construction we only need to use level lines whose boundary values are in $[-a, b]$.
- (ii) For a fixed point $z \in D$, a.s. we only need a finite number of level lines to construct the connected component of $D \setminus \mathbb{A}_{-a,b}^u$ containing z .
- (iii) As none of the level lines is started inside $u^{-a,b}$ nor can touch $u^{-a,b}$, any connected component of $u^{-a,b}$ belongs entirely to the boundary of a single connected component of $D \setminus \mathbb{A}_{-a,b}^u$. In particular each connected component O of the complement of $\mathbb{A}_{-a,b}$ with $\partial O \cap \partial D \neq \emptyset$ has only finitely many intersection points and by Lemma 10 in [ASW17] we can assign them any values, in particular those that u already takes on ∂D .
- (iv) Due to the construction $D \setminus \mathbb{A}_{-a,b}^u$ has at most n non-simply connected components.

Proof. MEASURABILITY of the sets $\mathbb{A}_{-a,b}^u$ with respect to the GFF just follows from the measurability of the level lines used in the construction and the measurability result of Proposition 3.7. □

4. FIRST PASSAGE SETS OF THE 2D CONTINUUM GFF

The aim of this section is to define the first passage sets of the 2D continuum GFF, prove its characterization and properties. We first state an axiomatic definition of the continuum FPS inspired by its heuristic interpretation: i.e. the FPS stopped at value $-a$ is given by all points in \bar{D} that can be connected to the boundary via a path on which the values of the GFF do not decrease below $-a$. From this description, it is clear that it induces a Markovian decomposition of the GFF: the field outside of it is just a GFF with boundary condition equal to $-a$. In other words, the FPS is a local set, that we denote by \mathbb{A}_{-a} , and its harmonic function has to satisfy $h_{\mathbb{A}_{-a}} = -a$ as we stop at value $-a$. Finally, the question is how to translate the property for the values, as the GFF is not defined pointwise. The right way is to ask the distribution $\Phi_{\mathbb{A}_{-a}} + a$ to be a positive measure.

The set-up is again as follows: D is a finitely-connected domain where no component is a single point and u is a bounded harmonic function with piecewise constant boundary conditions. Here is the definition for general boundary values:

Definition 4.1. Let $a \in \mathbb{R}$ and Φ be a GFF in the multiple-connected domain D . We define the first passage set of Φ of level $-a$ and boundary condition u as the local set of Φ such that $\partial D \subseteq \mathbb{A}_{-a}^u$, with the following properties:

- (1) Inside each connected component O of $D \setminus \mathbb{A}_{-a}^u$, the harmonic function $h_{\mathbb{A}_{-a}^u} + u$ is equal to $-a$ on $\partial \mathbb{A}_{-a}^u \setminus \partial D$ and equal to u on $\partial D \setminus \mathbb{A}_{-a}^u$ in such a way that $h_{\mathbb{A}_{-a}^u} + u \leq -a$.
- (2) $\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u} \geq 0$, i.e., for any smooth positive test function f we have $(\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}, f) \geq 0$.
- (3) Additionally, for any connected component O of the complement of \mathbb{A}_{-a}^u , for any $\varepsilon > 0$ and $z \in \partial O$ and for all sufficiently small open ball U_z around z , we have that a.s.

$$h_{\mathbb{A}_{-a}^u}(z) + u(z) \geq \min\{-a, \inf_{w \in U_z \cap \overline{O}} u(w)\} - \varepsilon.$$

Notice that if $u \geq -a$, then the conditions (1) and (2) correspond more precisely to the heuristic and are equivalent to

- (1)' $h_{\mathbb{A}_{-a}^u} + u = -a$ in $D \setminus \mathbb{A}_{-a}^u$.
- (2)' $\Phi_{\mathbb{A}_{-a}^u} + u + a \geq 0$.

Moreover, in this case the technical condition (3) is not necessary. This condition roughly says that nothing odd can happen at boundary values that we have not determined: those on the intersection $\partial \mathbb{A}_{-a}^u$ and ∂D . This condition enters in the case $u < -a$: in [ALS18c] we want to take the limit of the FPS on metric graphs and it comes out that it is easier not to prescribe the value of the harmonic function at the intersection of ∂D and $\partial \mathbb{A}_{-a}^u$. Notice that in contrast we did prescribe the values at intersection points for two-valued sets.

Remark 4.2. One could similarly define excursions sets in the other direction, i.e. stopping the sets from above. We denote these sets by \mathbb{V}_b^u . In this case the definition goes the same way except that (2) should now be replaced by $\Phi_{\mathbb{V}_b^u} \leq h_{\mathbb{V}_b^u}$ and (3) by

$$h_{\mathbb{V}_b^u}(z) + u(z) \leq \max\{b, \inf_{w \in U_z \cap \overline{O}} u(w)\} - \varepsilon.$$

Let us also remark that in [APS17] the sets \mathbb{V}_b are unluckily denoted by \mathbb{A}_b , and that they can be obtained as \mathbb{A}_{-b}^{-u} of $-\Phi$.

Theorem 4.3. The FPS \mathbb{A}_{-a}^u of level $-a$ exists and is unique in the sense that if (Φ, A') is an FPS of level $-a$ for the GFF with boundary condition u , then $A' = \mathbb{A}_{-a}^u$ a.s.

We start from the existence of the FPS. Here, we provide a purely continuum construction using the two-valued sets $\mathbb{A}_{-a,b}$. Another approach would be to consider the scaling limit of the metric graph FPS when the mesh size goes to zero, as is done in [ALS18c].

Proposition 4.4. Denote by $\mathbb{A}_{-a,n}^u$ the two-valued local sets coupled with the GFF Φ in the domain D . Then for every $a \geq 0$ the local set $\mathbb{A}_{-a}^u := \overline{\cup_{n \in \mathbb{N}} \mathbb{A}_{-a,n}^u}$ is an FPS of level $-a$.

Proof. Let \mathbb{A}_{-a}^u be as in the statement. Then \mathbb{A}_{-a}^u is the closed union of nested measurable local sets so it is a measurable local set: it is a local set by Lemma 2.3 and measurable as a limit of measurable functions.

We first prove the condition (1) of the Definition 4.1. Take a countable dense set in D , $(z_i)_{i \in \mathbb{N}}$, and note that almost surely for all $i \in \mathbb{N}$, $z_i \notin \mathbb{A}_{-a}^u$. Consider $n > \sup u$. It suffices to show that for any z_i , there will be some finite n such that the component of the complement of $\mathbb{A}_{-a,n}^u$ containing z_i does not take the value n . Indeed, when this happens, then by the definition and uniqueness of two-valued sets above, it would take a value as described in (1) for all $\tilde{n} \geq n$.

Now, as a process in n , $h_{\mathbb{A}_{-a,n}^u}(z_i)$ is a lower bounded martingale, and thus it converges almost surely. It can, however, only converge when for some n it belongs to the component of the complement of $\mathbb{A}_{-a,n}^u$ not taking the value n . Hence we deduce the condition (1).

The condition (3) just comes from the fact that the value at the intersection points is prescribed by the definition of two-valued sets and it satisfies the appropriate condition.

It remains to prove (2), i.e. that $\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u} \geq 0$. Note that for all positive $f \in \mathcal{C}_0^\infty$ and all $n \geq 2$, we have that $(\Phi_{\mathbb{A}_{-a,n}^u} - h_{\mathbb{A}_{-a}^u}, f) \geq 0$. Thus, we conclude using Lemma 2.3 (iii). \square

Let us make the following observation about the construction above:

- (i) In the construction, we only need to use generalized level lines whose boundary values are in $[-a, \infty)$, moreover these generalized level lines never hit themselves.
- (ii) For a fixed point $z \in D$, it will belong to a component of the complement of $\mathbb{A}_{-a,n}^u$ with value n only for a finite number of n . Thus, we need only a finite number of level lines to construct the loop of \mathbb{A}_{-a}^u surrounding z .

We now want to use these remarks and the techniques of [ASW17] to prove the uniqueness of the FPS:

Proposition 4.5 (Uniqueness and Monotonicity of the FPS). *Let (Φ, A') be a FPS of level $-a$ for the GFF with boundary condition u . Then $A' = \mathbb{A}_{-a}^u$ a.s. Additionally, if $a \leq a'$ and $u \leq u'$, then $\mathbb{A}_{-a}^u \subseteq \mathbb{A}_{-a'}^{u'}$*

Proof. First let us prove that if A is a local set such that almost surely $\Phi_A \geq 0$, then a.s. A is a polar set. Given this condition, we have that $(\Phi_A, 1) \geq 0$, and due to the Markov property we know that $\mathbb{E}[(\Phi_A, 1)] = 0$, thus a.s. $(\Phi_A, 1) = 0$. Additionally, we know that $G_D \geq G_{D \setminus A}$, using again the strong Markov property we get that

$$\iint_{D \times D} G_D(z, w) dz dw = \mathbb{E}[(\Phi, 1)^2] = \mathbb{E}[(\Phi_G^{D \setminus A}, 1)^2] = \iint_{D \setminus A \times D \setminus A} G_{D \setminus A}(z, w) dz dw.$$

Hence almost everywhere $G_{D \setminus A} = G_D$, and thus A is a polar set.

We now prove the uniqueness of the FPS. Assume first that $\mathbb{A}_{-a}^u \subseteq A'$. We claim that then $A' \setminus \mathbb{A}_{-a}^u$ is a polar set. Indeed, consider $B := A' \setminus \mathbb{A}_{-a}^u$. From Lemma 2.3 (2), B is a local set of the zero boundary GFF $\Phi^{\mathbb{A}_{-a}^u}$. Moreover, one can check that from our conditions on the FPS, it follows that $h_{\mathbb{A}_{-a}^u} + u \leq h_{A'} + u$ and hence $(\Phi^{\mathbb{A}_{-a}^u})_B \geq 0$. Thus, by the previous argument B is polar.

Now, it suffices to prove that a.s. $\mathbb{A}_{-a}^u \subseteq A'$. We prove the monotonicity using arguments similar to those of Section 6 in [ASW17]. Take A' an FPS for $u' \geq u$ and $a' \geq a$. Suppose by contradiction, \mathbb{A}_{-a}^u is not contained in A' . Then choosing a countable dense set in D , $(z_i)_{i \in \mathbb{N}}$, there must be some z_i such that with positive probability during the construction of \mathbb{A}_{-a}^u a generalized level line enters the component $O(z_i)$ of the complement of A' containing z_i . Thus, there should be some finite $n \in \mathbb{N}$ such that with positive probability, η , the n^{th} -level line pointed towards z_i , is the first one to enter $O(z_i)$ and $\eta \cap O(z_i) \neq \emptyset$. This is, however, in contradiction with Lemma 3.17 and the remark just after the proof: indeed, the boundary values of $h_{A'}$ inside $O(z)$ are equal to $-a' - u' \leq -a - u$ and by the remarks above the boundary values of η are in $[-a, \infty)$. Thus, the uniqueness follows and thus monotonicity just follows from the construction given in Lemma 4.4. \square

Remark 4.6. *Let us remark that in fact $\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}$ is almost surely a non-trivial positive measure, unless $u \leq -a$ on D . Indeed, suppose for example that we work in \mathbb{D} and with a zero boundary GFF. In this case, it is known that the circle average (ρ_{1-r}^0, Φ) around 0 of radius $1 - r$ converges to 0 as $r \rightarrow 0$. But the circle average w.r.t. to $h_{\mathbb{A}_{-a}^u}$ is constantly equal to $-a$, and the variance of $(\rho_{1-r}^0, \Phi^{\mathbb{A}_{-a}^u})$ also converges to 0.*

A different way of seeing this is the following. Since

$$\mathbb{E}[\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}] = \mathbb{E}[-h_{\mathbb{A}_{-a}^u}] > 0,$$

$\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}$ is a non-trivial positive measure with positive probability. Then, in order to sample \mathbb{A}_{-a}^u , we can first explore $\mathbb{A}_{-a/2}^u$ and then further explore $\mathbb{A}_{-a/2}$ conditional independently in each connected component of the complement of $\mathbb{A}_{-a/2}^u$. Using that, one can argue that the FPS is non-trivial on every dyadic square it intersects.

4.1. Distance to interior points and boundary for \mathbb{A}_{-a}^u . We will now give an exact description of the law of the distance of \mathbb{A}_{-a}^u to interior points and to boundary components. This can be seen as a continuum analogue of Corollary 1 of [LW16]. The proofs follow from Propositions 2.5 and 2.7, that give a way to parametrize the FPS using the distance to an interior point or a boundary component respectively.

Proposition 4.7. *Let $a > 0$ and D a n -connected domain. Moreover, let u be a bounded harmonic function with piecewise constant boundary data and $z \in D$. Take W_t to be a Brownian motion started from $u(z)$ and with life-time $g_D(z, z)$. If $u \geq -a$, then*

$$G_D(z, z) - G_{D \setminus \mathbb{A}_{-a}^u}(z, z)$$

is distributed like the one-sided hitting time of the level $-a$ by W_t .

Proof. This follows exactly as the proof of Proposition 20 in [ASW17]. \square

Similarly, we can calculate the extremal distance between boundary components, analogously to Proposition 5 in [LW16].

Proposition 4.8. *Let a be a positive number, D a finitely-connected domain, and \mathcal{B} a union of connected components of ∂D . Moreover, let u be a bounded harmonic function with piecewise constant boundary data changing finitely many times such that u on $\partial D \setminus \mathcal{B}$ is a constant equal to $u_e \leq -a$. Let \widehat{W}_t be a Brownian bridge with starting point:*

$$u_s := \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) \iint_{\mathcal{B} \times \partial D \setminus \mathcal{B}} u(x) H_D(dx, dy)$$

endpoint u_e , and with length $\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B})$. If $u \geq -a$ on \mathcal{B} , then

$$\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) - \text{EL}(\mathcal{B} \cup (\mathbb{A}_{-a}^u \setminus \partial D), \partial D \setminus \mathcal{B})$$

has the law of the first hitting time of $-a$ by \widehat{W}_t .

Proof. From the construction of first passage sets (see the observations after the proof of Proposition 4.4), we know that the first passage set can be constructed by using only level lines with boundary values in $\geq -a$ that do not touch $\partial D \setminus \mathcal{B}$. We can parametrize the part of the construction, that always continues in the connected component containing $\partial D \setminus \mathcal{B}$ on its boundary, using its extremal length to $\partial D \setminus \mathcal{B}$. We denote the resulting local set process by $(A_t)_{0 \leq t \leq \tau}$. Here τ is the first time that this component stops growing. In other words $\tau := \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) - \text{EL}(\mathcal{B} \cup (\mathbb{A}_{-a}^u \setminus \partial D), \partial D \setminus \mathcal{B})$ and moreover τ is the first time that satisfies the following property:

- ⊕ Restricted to the connected component of O of $A_\tau \setminus D$ such that $\partial D \setminus \mathcal{B} \subseteq \partial O$, $h_{A_\tau} + u$ is the bounded harmonic function with boundary value $-a$ in $\partial O \setminus \mathcal{B}$ and $u_e \in \mathcal{B}$.

Using Proposition 2.7 for the underlying GFF Φ , we deduce that

$$\widehat{W}_t := \text{EL}(\mathcal{B} \cup (A_t \setminus \partial D), \partial D \setminus \mathcal{B}) \left(- \int_{\partial D \setminus \mathcal{B}} \partial_n (h_{A_t} + u) \right) + u_e$$

is a Brownian bridge from u_s to u_e and of length $\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B})$. From an explicit calculation using Green's identities and Theorem 3.1 we see that:

$$\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) \iint_{\mathcal{B} \times \partial D \setminus \mathcal{B}} u(x) H_D(dx, dy) = - \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) \int_{\partial D \setminus \mathcal{B}} \partial_n u(x) dx + u_e.$$

Moreover, similarly we obtain that \ominus is equivalent to $\widehat{W}_\tau = -a$. Indeed, we can calculate using Theorem 2.1 that

$$\widehat{W}_\tau = (a + u_e) \text{EL}(\mathcal{B} \cup (A_\tau \setminus \partial D), \partial D \setminus \mathcal{B}) \int_{\partial D \setminus \mathcal{B}} \partial_n \bar{u}(x) + u_e = -a.$$

Here \bar{u} is the bounded harmonic function with values 0 in $\partial D \setminus \mathcal{B}$ and 1 in $\mathcal{B} \cup A_\tau$. The same calculation yields that for all times smaller than τ , $\widehat{W}_\tau > -a$, from where we conclude. \square

Let us point out the following corollary

Corollary 4.9. *Let \mathbb{A}_{-a}^u be an FPS with boundary condition u of Φ , where u is a bounded harmonic function with piecewise constant boundary data. Then $\mathbb{A}_{-a}^u \cap D$ is at positive distance of any connected component of ∂D where $u \leq -a$. Furthermore, for any $\varepsilon > 0$ and any allowed boundary condition u , there exists a strictly positive probability that the connected component of \mathbb{A}_{-a}^u started from a given boundary component remains at distance ε from that boundary component.*

Notice that if only a part of the boundary satisfies $u \leq -a$, then we also know that the FPS stays at a positive distance of any point on this interval. Indeed, this follows from the level line construction: we know from the proof of Theorem 3.19 and the remarks following the proof that for $\mathbb{A}_{-a,b}^u$ any connected set of $u^{-a,b}$ is entirely part of the boundary of a component of $D \setminus \mathbb{A}_{-a,b}^u$; on the other hand we also know that any component with $h_{\mathbb{A}_{-a,b}^u} + u \leq -a$ in the complement of $\mathbb{A}_{-a,b}^u$, will also be a component of the complement of \mathbb{A}_{-a}^u . Putting this together we conclude:

Corollary 4.10. *Let \mathbb{A}_{-a}^u be an FPS with boundary condition u of Φ , where u is a bounded harmonic function with piecewise constant boundary data. Then $\mathbb{A}_{-a}^u \cap D$ is at positive distance of any connected $J \subseteq \partial D$ where there is an open neighborhood J_ε such that $u(x) \leq -a$ for all $x \in J_\varepsilon \cap \partial D$.*

Finally, let us mention that one can, via the same proof, also prove an analogue of Proposition 4.7 for two-valued local sets:

Proposition 4.11. *Let a, b be positive with $a + b \geq 2\lambda$, D a n -connected domain, and \mathcal{B} a union of connected components of ∂D . Moreover, let u be a bounded harmonic function with piecewise constant boundary data changing finitely many times such that u on $\partial D \setminus \mathcal{B}$ is a constant equal to $u_e \notin (-a, b)$. Let \widehat{W}_t be a Brownian bridge with starting point:*

$$u_s := \text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) \iint_{\mathcal{B} \times \partial D \setminus \mathcal{B}} u(x) H_D(dx, dy)$$

endpoint u_e , and with length $\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B})$. If $u \in [-a, b]$ on \mathcal{B} , then

$$\text{EL}(\mathcal{B}, \partial D \setminus \mathcal{B}) - \text{EL}(\mathcal{B} \cup (\mathbb{A}_{-a}^u \setminus \partial D), \partial D \setminus \mathcal{B})$$

is stochastically bounded by the law of the first hitting time of $\{-a, b\}$ by \widehat{W}_t . Furthermore if D is an annular domain, then they have exactly the same law.

4.2. Level lines as boundaries of FPS. Now, let us see that level line can be identified with the boundary of certain FPS. Let D be finitely connected domain and $\partial_{\text{ext}} D$ be the outermost connected component of ∂D , that is to say the one that separates D from infinity. We consider two boundary points $x_0 \neq y_0 \in \partial_{\text{ext}} D$ that split $\partial_{\text{ext}} D$ in two boundary arcs, \mathcal{B}_1 and \mathcal{B}_2 such that y_0, \mathcal{B}_2, x_0 come in clockwise order. Assume that u is a bounded harmonic function with piecewise boundary values which are smaller than or equal to $-\lambda$ on \mathcal{B}_2 , $\inf_{\mathcal{B}_1} u > -\lambda$ and $\inf_{\partial D \setminus \partial_{\text{ext}} D} u \geq \lambda$. Note that thanks to Lemma 3.2 there is a generalized level line η of $\Phi + u$ starting at y_0 and targeted at x_0 .

Corollary 4.12. *Let u be as above and Φ be a GFF in D . Then, there exists a connected component O of $D \setminus \mathbb{A}_{-\lambda}^u$ such that $\partial_{\text{ext}} O$ is equal to the union of \mathcal{B}_2 and the generalized level line of $\Phi + u$ from y_0 to x_0 .*

Proof. This basically just follows from the uniqueness of the FPS, i.e. Proposition 4.5.

First, construct η , the generalized level line of $\Phi + u$ from y_0 to x_0 . By Lemma 3.13, it intersects \mathcal{B}_2 only at y_0 and x_0 and it does not intersect $\partial D \setminus \partial_{\text{ext}} D$. Denote by O the connected component of $D \setminus \eta$ containing \mathcal{B}_2 on its boundary: i.e. when we denote by $\partial_{\text{ext}} O$ the outermost connected component of its boundary, then it holds that $\partial_{\text{ext}} O = \mathcal{B}_2 \cup \eta$. To construct the FPS of level $-\lambda$ in D it remains to construct the FPS of level $-\lambda$ inside each connected component of $D \setminus \eta([0, \infty])$.

If all of the boundary components of $\partial D \setminus \partial_{\text{ext}} D$ are outside of O , then we can see that O remains fixed in the remaining construction and thus by the uniqueness of FPS is a component of the complement of $\mathbb{A}_{-\lambda}^u$.

If some of them are inside O , we have to make sure that the FPS of level $-\lambda$ constructed inside O does not intersect $\mathcal{B}_2 \cup \eta$. To do this, notice that the boundary values of $\partial_{\text{ext}} O = \mathcal{B}_2 \cup \eta$ are bounded from above by $-\lambda$. Thus, by Corollary 4.9, the FPS of level $-\lambda$ inside O stays at positive distance of $\partial_{\text{ext}} O$. □

Remark 4.13. *Let us remark that one can also obtain families of multiple level lines, as boundaries of the FPS $\mathbb{A}_{-\lambda}^u$. For instance, when u changes finitely many times between $-\lambda$ and λ in a simply connected domain, one can get in this way multiple commuting SLE₄ curves. In [PW17a] the authors provide an explicit expression for probabilities of different pairings.*

4.3. Uniqueness, monotonicity and local finiteness of $\mathbb{A}_{-a,b}^u$. In this subsection we will prove that $\mathbb{A}_{-a,b}^u$ is the only BTLS satisfying (\heartsuit) . The monotonicity then follows as in the proof of Proposition 2 of [ASW17].

To prove the uniqueness of $\mathbb{A}_{-a,b}^u$, we show that it can be represented as a simple function of \mathbb{A}_{-a}^u and \mathbb{V}_b^u , and then use the uniqueness of the FPS.

We say that a set $A \subseteq D$ is connected to the boundary if every connected component of $A \cup \partial D$ contains at least one boundary component of ∂D .

Proposition 4.14. *Let Φ be a GFF in an n -connected domain D . Then almost surely the union of the components of $\mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$ that are connected to boundary is equal to $\mathbb{A}_{-a,b}^u$. In particular, if D is simply connected we have that $\mathbb{A}_{-a,b}^u = \mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$*

Proof. From the monotonicity of two-valued sets (Theorem 3.19) and the construction of the FPS (Lemma 4.4) we see that for any $n \geq a \vee b$, we have that $\mathbb{A}_{-a,b}^u \subseteq \mathbb{A}_{-a,n}^u \subseteq \mathbb{A}_{-a}^u$ and, furthermore, $\mathbb{A}_{-a,b}^u \subseteq \mathbb{A}_{-n,b}^u \subseteq \mathbb{V}_b^u$. Hence $\mathbb{A}_{-a,b}^u \subseteq \mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$. Moreover, as by construction $\mathbb{A}_{-a,b}^u$ is connected to the boundary, we deduce that $\mathbb{A}_{-a,b}^u$ is contained on the union of connected components of $\mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$ that are connected to the boundary.

We will now prove the opposite inclusion. To do this, it suffices to show that for every connected component O of $D \setminus \mathbb{A}_{-a,b}^u$, either it is contained in $D \setminus \mathbb{A}_{-a}^u$ or $D \setminus \mathbb{V}_b^u$, or the closure $O \cap \mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$ does not intersect ∂O .

First, let us consider the case where $h_{\mathbb{A}_{-a,b}^u} + u$ is less than or equal to $-a$ in O . By construction of the FPS (Lemma 4.4), we see that O is also a connected component of $D \setminus \mathbb{A}_{-a}^u$. Similarly, if $h_{\mathbb{A}_{-a,b}^u} + u$ is greater than or equal to b in O , then O is a component of $D \setminus \mathbb{V}_b^u$.

Thus, it remains to deal with the case when O has more than one boundary component, and where u takes values that less than or equal to $-a$ on some of the boundary components, and values that are greater than or equal to b on all the other boundary components.

Now, by the uniqueness of the FPS (Proposition 4.5) we know that in order to construct the part of FPS \mathbb{A}_{-a}^u intersecting O , we can first take $\mathbb{A}_{-a,b}^u$, and then construct the FPS of level $-a$ in O . But now, by Corollary 4.9 we know that an FPS of level $-a$ inside O cannot touch any part of the boundary where u is smaller than or equal to $-a$. But similarly we see that the part of FPS \mathbb{V}_b^u intersecting O , cannot touch the boundary of O where u is larger than or equal to b . Hence

the closure of $\mathbb{A}_{-a}^u \cap \mathbb{V}_b^u \cap O$ cannot intersect the boundary of O and hence is not connected to the boundary of D . The proposition follows. \square

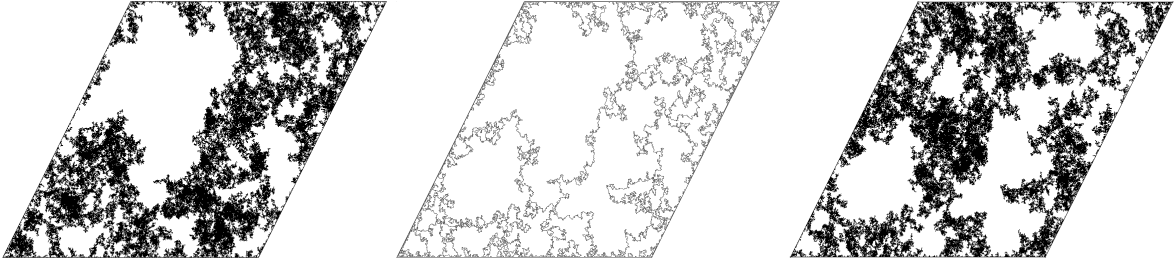


FIGURE 5. Graphical evidence of Lemma 4.14: On the left $\mathbb{A}_{-\lambda}$, on the right \mathbb{V}_λ and in the middle $\mathbb{A}_{-\lambda,\lambda}$. Simulation done by B. Werness.

We are now ready to prove the uniqueness of two-valued sets $\mathbb{A}_{-a,b}^u$ for general boundary data in n -connected domains. See Theorem 3.19 for the setting and precise statement.

Uniqueness of two-valued sets. In the proof of Theorem 3.19, we showed the existence of a two-valued set. Denote this set by $\mathbb{A}_{-a,b}^u$. Suppose A' is another two-valued set coupled with the same GFF, i.e. it satisfies the condition (\clubsuit) given just before Theorem 3.19.

First, notice that A' has to be connected to the boundary: indeed, suppose for contradiction that there is a component B of A' that is not connected to the boundary. Consider the component O of $D \setminus B$ that has B on part of its boundary, and let $\mathcal{B} = \partial O \cap B$. WLOG suppose that the boundary conditions on \mathcal{B} are equal to $-a$. Then, as in the last paragraph of proof of Lemma 4.14, we see that the FPS of height $-a$ also contains \mathcal{B} , and that moreover \mathcal{B} is also not connected to the boundary as a subset of the FPS. However, from the construction (Proposition 4.4) and uniqueness of the FPS (Proposition 4.5) we know that the FPS is connected to the boundary.

Now, inspecting the proof of Lemma 4.14, we conclude that exactly the same proof gives that almost surely the union of the components of $\mathbb{A}_{-a}^u \cap \mathbb{V}_b^u$ that are connected to boundary is equal to A' . However, we know from Lemma 4.14 that it is equal to $\mathbb{A}_{-a,b}^u$ and the claim follows. \square

Finally, let us state here the proposition about $\mathbb{A}_{-a,b}$ are locally finite. In fact this relies on a result from [ALS18c], that proves the local finiteness of first passage sets, using a proof that uses the relation of FPS to Brownian loop-soups. We should maybe stress that none of [ALS18c] depends on local finiteness of $\mathbb{A}_{-a,b}$.

Proposition 4.15. *Consider a finitely-connected domain with any piece-wise constant boundary condition u . For any $a, b \geq 0$, the two-valued set $\mathbb{A}_{-a,b}^u$ is almost surely locally finite.*

Proof. Proposition 5.7 of [ALS18c] gives us the local finiteness of \mathbb{A}_{-a}^u for any choice of $a \geq 0$, of a piece-wise constant boundary condition u and of a finitely-connected domain D . But now, note that if O is a simply connected component of $D \setminus \mathbb{A}_{-a}^u$, then it is also a connected component of either $D \setminus \mathbb{A}_{-a}^u$ or $D \setminus \mathbb{V}_b^u$ - indeed, this follows from the uniqueness of the FPS, as when constructing say, \mathbb{A}_{-a}^u from $\mathbb{A}_{-a,b}^u$ one keeps all the components that have boundary condition less than or equal to $-a$. As there are only finitely many non-simply connected components of $D \setminus \mathbb{A}_{-a}^u$, we conclude. \square

Remark 4.16. *Let us stress that if one restricts oneself to the simply-connected case, then one can prove local finiteness of \mathbb{A}_{-a}^u without relying on the construction or properties of FPS and TVS in multiply connected domains. In fact, the proof in [ALS18c] just uses properties of the Brownian loop-soup. In particular, this provides an alternative proof of Proposition 3.11.*

5. THE MINKOWSKI CONTENT MEASURE OF THE FPS

The aim of this section to identify the measure

$$\nu_{\mathbb{A}_{-a}^u} := \Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}$$

as a Minkowski content measure in a certain gauge:

Theorem 5.1. *The measure $\nu_{\mathbb{A}_{-a}^u}$ is a measurable function of \mathbb{A}_{-a}^u . Moreover, it is proportional to the Minkowski content measure in the gauge $r \mapsto |\log(r)|^{1/2}r^2$. More precisely, almost surely for any continuous f compactly supported in D ,*

$$\nu_{\mathbb{A}_{-a}^u} = \lim_{r \rightarrow 0} \frac{1}{2} |\log(r)|^{1/2} \int_D f(z) \mathbf{1}_{d(z, \mathbb{A}_{-a}^u) \leq r} dz.$$

Let us stress that there are two non-obvious statements in this theorem:

- (1) the fact that the measure $\nu_{\mathbb{A}_{-a}^u}$ can be obtained as a measurable function of \mathbb{A}_{-a}^u ;
- (2) and the identification of this function as the Minkowski content measure of \mathbb{A}_{-a}^u in a certain gauge.

As a simply corollary, we can deduce an almost sure statement on the dimension of \mathbb{A}_{-a}^u :

Corollary 5.2. *If the harmonic function u is not everywhere less or equal to $-a$ (i.e. if $\mathbb{A}_{-a}^u \setminus \partial D$ is non-empty), then \mathbb{A}_{-a}^u is a.s. of Minkowski dimension 2.*

Moreover, observe that the expected value of $\nu_{\mathbb{A}_{-a}^u}$ is equal to a in the sense that for any continuous function f , the expectation of $(f, \nu_{\mathbb{A}_{-a}^u})$ is equal to a times $\int_D f(z) dz$. Thus, the fact that the measures $\Phi_{\mathbb{A}_{-a}^u}$ converge to Φ as $a \rightarrow \infty$, can be interpreted as saying that the GFF is a limit of recentered Minkowski content measures of certain growing random sets.

Corollary 5.3. *For any $n \geq 1$, let A_n be random set with the law of the FPS \mathbb{A}_{-n} in \mathbb{D} , and denote by ν_{A_n} its Minkowski content measure in the gauge $|\log(r)|^{1/2}r^2$. Then, the recentered measures $\nu_{A_n} - \mathbb{E}[\nu_{A_n}]$ converge in law in $H^{-1}(\mathbb{D})$ to a zero boundary GFF on \mathbb{D} .*

Proof. Note that $\Phi - \nu_{A_n} - n$ has the law of a GFF in Φ^{A_n} . As $A_n \nearrow \mathbb{D}$ for the Hausdorff topology, Equation (2.3) and Dominated convergence theorem imply that $\mathbb{E}[\|\Phi^{A_n}\|_{H^{-1}(\mathbb{D})}^2]$ converges to 0, giving the result. \square

Let us also make two remarks.

Remark 5.4. *In Theorem 5.1, we restricted to functions supported compactly away from ∂D , because by construction, $\nu_{\mathbb{A}_{-a}^u}(\partial D) = 0$, yet ∂D might be irregular enough to have a positive, or even infinite Minkowski content measure. Thus, $\nu_{\mathbb{A}_{-a}^u}$ is rather the Minkowski content measure of $\mathbb{A}_{-a}^u \setminus \partial D$.*

Remark 5.5. *In fact, the theorem holds in a more general setting (with basically the same proof). Indeed, take A to be a local set of the GFF Φ such that there exists a deterministic K such that $\Phi_A = \nu_A + h_A \geq -K$, where h_A is harmonic of $D \setminus A$ and ν_A is a non-negative measure supported on A . Then, ν_A is given by the Minkowski content measure of $A \setminus \partial D$ as in Theorem 5.1.*

The rest of this section is devoted to the proof of Theorem 5.1. We start by looking at the decomposition of the Gaussian multiplicative chaos (GMC) measure induced by the Markovian decomposition of the field w.r.t the FPS. Then, we use this decomposition to see how the geometry of the FPS encodes its height and to give a short, but somewhat unexpected proof of why $\Phi_{\mathbb{A}_{-a}^u}$ is a measurable function of \mathbb{A}_{-a}^u . Next, we observe that a careful analysis of the proof gives in fact an explicit expression of $\Phi_{\mathbb{A}_{-a}^u} - h_{\mathbb{A}_{-a}^u}$ in terms of \mathbb{A}_{-a}^u , and finally we identify this expression with the Minkowski content measure of A in the gauge $r \mapsto |\log(r)|^{1/2}r^2$. \square In several of the arguments

it will be more convenient to work with \mathbb{V}_b^u , but notice that by the symmetry of the GFF this is equivalent.

5.1. Decomposition of the GMC using the FPS. In [APS17], the decomposition of the GMC measures w.r.t. to the FPS and w.r.t. more general local sets was used to view the GMC measure of the GFF as a multiplicative cascade. In particular, it was observed that in the case of the FPS one recovers a construction of [Aid15].

As in [APS17] the decomposition of the GMC measures was stated w.r.t. the sets \mathbb{V}_b^u instead of \mathbb{A}_{-a}^u (as this allows to consider positive values of γ), we will do the same here.

Define $\tilde{\gamma} = \sqrt{2\pi}\gamma$ and $\mathcal{F}_{\mathbb{V}_b^u} := \sigma(\mathbb{V}_b^u, \Phi_{\mathbb{V}_b^u})$. Then from the Proposition 4.1 of [APS17] (see also Theorem 2.5 in [APS18]) it follows that:

Proposition 5.6. *Take $0 \leq \gamma < 2$, i.e., $0 < \tilde{\gamma} < 2\sqrt{2\pi}$. Then, for all continuous function f compactly supported in D ,*

$$\mathbb{E}[(\mu_\gamma, f) \mid \mathcal{F}_{\mathbb{V}_b^u}] = \int_{D \setminus \mathbb{V}_b^u} f(z) \exp\left(\tilde{\gamma} h_{\mathbb{V}_b^u}(z) + \frac{\tilde{\gamma}^2}{2} g_{D \setminus \mathbb{V}_b^u}(z, z)\right) dz$$

where g_D was defined in (2.1). Moreover, conditioned on \mathbb{V}_b^u , we have that

$$(\mu_\gamma, f) = \int_{D \setminus \mathbb{V}_b^u} f(z) \exp(\tilde{\gamma} h_{\mathbb{V}_b^u}(z)) d\mu_\gamma^{\Phi_{\mathbb{V}_b^u}}(z),$$

where $\mu_\gamma^{\Phi_{\mathbb{V}_b^u}}$ is the GMC measure for the field $\Phi_{\mathbb{V}_b^u}$.

As from condition (1) of the definition of the FPS it follows that in the coupling $(\Phi, \mathbb{A}_{-\lambda}^u)$, the harmonic extension $h_{\mathbb{A}_{-a}^u}$ is a measurable function only of the set \mathbb{A}_{-a}^u , we infer a direct but useful corollary:

Lemma 5.7. *$(\mu_\gamma)_{0 \leq \gamma < 2}$ is a measurable function of \mathbb{V}_b^u and $\Phi_{\mathbb{V}_b^u}$.*

Observe that this means that in order to construct the Liouville measure μ_γ we do not need the positive measure $\nu_{\mathbb{A}_{-a}^u}$.

5.2. How to read the height of the FPS from its geometry. We have seen that the FPS is measurable w.r.t. the underlying free field. Moreover, we just saw that $h_{\mathbb{A}_{-a}^u}$ is a measurable function only of \mathbb{A}_{-a}^u itself. This tells us in particular that knowing the boundary conditions u and the set \mathbb{A}_{-a}^u is sufficient to implicitly find the value of $-a$: for example if $x \in D$, it is given by, for example, $u(x) + h_{\mathbb{A}_{-a}^u}(x)$. A natural question is whether one can also determine $-a$ in an explicit way by just looking at the geometry set \mathbb{A}_{-a} . The following proposition says that this is indeed the case, and it presents some of the basic tools later used to show how to construct $\Phi_{\mathbb{A}_{-a}^u}$ as an explicit function of \mathbb{A}_{-a}^u . As this is more of a little side-story, we state and prove it in the simplest setting of the unit disk and the zero boundary GFF Φ . A similar result can be stated and proved more generally for finitely-connected domains and general boundary conditions.

Proposition 5.8. *Consider the FPS \mathbb{A}_{-a} of the GFF Φ defined on \mathbb{D} and let $\gamma \in (0, 1)$. Then, we have that $\int_0^{2\pi} \left(\frac{\text{CR}(re^{i\theta}, D \setminus \mathbb{A}_{-a})}{\text{CR}(re^{i\theta}, \mathbb{D})}\right)^{\pi\gamma^2} d\theta$ converges to $2\pi e^{-\gamma\sqrt{2\pi}a}$ as $r \rightarrow 1$.*

Proof. It is slightly more convenient to work with \mathbb{V}_b instead of \mathbb{A}_{-a} . Denote by

$$A_r := \int_0^{2\pi} \left(\frac{\text{CR}(re^{i\theta}, \mathbb{D} \setminus \mathbb{V}_b)}{\text{CR}(re^{i\theta}, \mathbb{D})}\right)^{\pi\gamma^2} d\theta.$$

From Proposition 5.6 it follows that for any $z \in \mathbb{D}$,

$$\mathbb{E} \left[\left(\frac{\text{CR}(z, \mathbb{D} \setminus \mathbb{V}_b)}{\text{CR}(z, \mathbb{D})} \right)^{\pi\gamma^2} e^{\gamma\sqrt{2\pi}b} \right] = 1.$$

Thus, the mean of A_r is equal to $2\pi e^{-\gamma\sqrt{2\pi}b}$. It remains to show that its variance tends to zero as $r \rightarrow 1$. By the Markov decomposition w.r.t. \mathbb{V}_b and a direct Gaussian computation, we have that for any $z, w \in \mathbb{D}$

$$\mathbb{E} \left[e^{\gamma\sqrt{2\pi}(\Phi_\varepsilon(z) + \Phi_\varepsilon(w)) - \frac{\gamma^2}{2} 2\pi(\mathbb{E}(\Phi_\varepsilon(z)^2) + \mathbb{E}(\Phi_\varepsilon(w)^2))} \middle| \mathbb{V}_b \right]$$

converges almost surely as $\varepsilon \rightarrow 0$ to

$$e^{\gamma^2 4\pi G_{D \setminus \mathbb{V}_b}(z, w)} \left(\frac{\text{CR}(z, \mathbb{D} \setminus \mathbb{V}_b)}{\text{CR}(z, \mathbb{D})} \right)^{\pi\gamma^2} \left(\frac{\text{CR}(w, \mathbb{D} \setminus \mathbb{V}_b)}{\text{CR}(w, \mathbb{D})} \right)^{\pi\gamma^2} e^{2\gamma\sqrt{2\pi}b}.$$

Thus, as the Green's function is always positive, using bounded convergence theorem $\mathbb{E}[A_r^2]$ is upper bounded by

$$e^{-2\gamma\sqrt{2\pi}b} \lim_{\varepsilon \rightarrow 0} \int_0^{2\pi} \int_0^{2\pi} \mathbb{E} \left[e^{\gamma\sqrt{2\pi}(\Phi_\varepsilon(re^{i\theta_1}) + \Phi_\varepsilon(re^{i\theta_2})) - \frac{\gamma^2}{2} 2\pi(\mathbb{E}(\Phi_\varepsilon(re^{i\theta_1})^2) + \mathbb{E}(\Phi_\varepsilon(re^{i\theta_2})^2))} \right] d\theta_1 d\theta_2.$$

However, this expected value can be explicitly calculated from a Gaussian estimate and we can get that the right-hand term converges to

$$e^{2\gamma\sqrt{2\pi}b} \int_0^{2\pi} \int_0^{2\pi} e^{\gamma^2 G_D(re^{i\theta_1}, re^{i\theta_2})} d\theta_1 d\theta_2 < \infty$$

as $\gamma < 1$. Finally, using dominated convergence and the fact that a.e. $G_D(re^{i\theta_1}, re^{i\theta_2})$ converges to 0 as $r \rightarrow 1$, we have that $\mathbb{E}[A_r^2]$ converges as $r \rightarrow 1$ to $(2\pi e^{-\gamma\sqrt{2\pi}b})^2$. \square

5.3. Measurability of $\Phi_{\mathbb{A}_{-a}^u}$ w.r.t. \mathbb{A}_{-a}^u . We will now give a short argument to prove the measurability of $\Phi_{\mathbb{A}_{-a}^u}$ w.r.t. \mathbb{A}_{-a}^u :

Proposition 5.9. *Let Φ be a GFF in D . We have that $\Phi_{\mathbb{A}_{-a}^u}$ is a measurable function of the set \mathbb{A}_{-a}^u .*

Proof. In fact it is again clearer to prove the claim for the sets \mathbb{V}_b^u . To prove the measurability of $\Phi_{\mathbb{V}_b^u}$, take $\gamma \in (0, 2)$ and consider μ_γ the γ -GMC measure corresponding to $\sqrt{2\pi}\Phi$. The proof is based on three measurability statements:

- (1) $\Phi_{\mathbb{V}_b^u}$ is a measurable function of Φ : this follows from Lemma 4.4.
- (2) μ_γ is a measurable function of \mathbb{V}_b^u and $\Phi_{\mathbb{V}_b^u}$: this follows from Lemma 5.7.
- (3) Φ is a measurable function of μ_γ for a fixed $\gamma < 2$: this follows from Remark 2.11.

Thus, if F is a bounded measurable function, we have that

$$F(\Phi_{\mathbb{V}_b^u}) = \mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \Phi] = \mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \mu_\gamma(\Phi)] = \mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \mathbb{V}_b^u, \Phi_{\mathbb{V}_b^u}].$$

Note that when A is a local set, it follows from the definition that Φ^A and Φ_A are conditionally independent given A . Hence

$$\mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \mathbb{V}_b^u, \Phi_{\mathbb{V}_b^u}, h_{\mathbb{V}_{-a}^u}] = \mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \mathbb{V}_b^u, \Phi_{\mathbb{V}_b^u}] = \mathbb{E} [F(\Phi_{\mathbb{V}_b^u}) \mid \mathbb{V}_b^u].$$

This proves the proposition. \square

5.4. **Explicit expression for the measure $\nu_{\mathbb{A}_{-a}^u}$.** Let us now derive an explicit expression for the measure $\nu_{\mathbb{A}_{-a}^u}$. In fact, it is simpler to continue working with \mathbb{V}_b^u , and thus write $\nu_{\mathbb{V}_b^u} := h_{\mathbb{V}_b^u} - \Phi_{\mathbb{V}_b^u}$. Notice the difference in sign in the definition due to the fact of taking the FPS in the other direction. Our aim is to prove that:

Proposition 5.10. *Almost surely for all f continuous function compactly supported in D ,*

$$(5.1) \quad (\nu_{\mathbb{V}_b^u}, f) = \lim_{\gamma \rightarrow 0} \frac{\gamma}{\sqrt{2\pi}} \int_{D \setminus \mathbb{V}_b^u} f(z) |\log d(z, \mathbb{V}_b^u)| d(z, \mathbb{V}_b^u)^{\gamma^2/2} dz.$$

By symmetry, the same holds for $\nu_{\mathbb{A}_{-a}^u}$.

To prove this proposition we will first look at the step 3 in the proof of Proposition 5.9 in more detail. Indeed, this step stems from (2.5), which identifies the GFF as the derivative of the GMC measures at $\gamma = 0$, i.e. for any compactly supported continuous f :

$$(\Phi, f) = \frac{1}{\sqrt{2\pi}} \partial_\gamma (\mu_\gamma, f) |_{\gamma=0}.$$

Let us take conditional expectation with respect to $\mathcal{F}_{\mathbb{V}_b^u}$ and write

$$(\Phi_{\mathbb{V}_b^u}, f) = \mathbb{E} [(\Phi, f) | \mathcal{F}_{\mathbb{V}_b^u}] = \frac{1}{\sqrt{2\pi}} \mathbb{E} \left[\lim_{\gamma \rightarrow 0^+} \partial_\gamma (\mu_\gamma, f) | \mathcal{F}_{\mathbb{V}_b^u} \right].$$

Now, by Proposition 2.9 $\partial_\gamma (\mu_\gamma, f)$ converges in \mathbb{L}^1 as $\gamma \rightarrow 0^+$ towards $\sqrt{2\pi}(\Phi, f)$ and thus one can interchange the limit as $\gamma \rightarrow 0^+$ and the conditional expectation. For the same reason, one can also change the order of the γ -derivative and the conditional expectation. We obtain:

$$(5.2) \quad (\Phi_{\mathbb{V}_b^u}, f) = \frac{1}{\sqrt{2\pi}} \lim_{\gamma \rightarrow 0^+} \partial_\gamma \mathbb{E} [(\mu_\gamma, f) | \mathcal{F}_{\mathbb{V}_b^u}]$$

We claim that:

Lemma 5.11. *Almost surely, for all compactly supported functions f on D ,*

$$(5.3) \quad (\nu_{\mathbb{V}_b^u}, f) = \lim_{\gamma \rightarrow 0^+} -\sqrt{2\pi}\gamma \int_{D \setminus \mathbb{V}_b^u} f(z) (g_{D \setminus \mathbb{V}_b^u}(z, z)) e^{\frac{\tilde{\gamma}^2}{2} g_{D \setminus \mathbb{V}_b^u}(z, z)} dz.$$

By symmetry, the same holds for $\nu_{\mathbb{A}_{-a}^u}$.

Proof. By Proposition 5.6, we have that $\partial_\gamma \mathbb{E} [(\mu_\gamma, f) | \mathbb{V}_b^u]$ is equal to

$$\sqrt{2\pi} \int_{D \setminus \mathbb{V}_b^u} f(z) \left[h_{\mathbb{V}_b^u}(z) + \sqrt{2\pi}\gamma g_{D \setminus \mathbb{V}_b^u}(z, z) \right] e^{\tilde{\gamma} h_{\mathbb{V}_b^u}(z) + \frac{\tilde{\gamma}^2}{2} g_{D \setminus \mathbb{V}_b^u}(z, z)} dz.$$

Using the fact that $\tilde{\gamma} h_{\mathbb{V}_b^u}(z) + \tilde{\gamma}^2 g_{D \setminus \mathbb{V}_b^u}(z, z)/2$ converges to 0 as $\gamma \rightarrow 0$ and (5.2), we obtain the result for a fixed function f . Thus, this also holds for a countable dense family, for the uniform convergence on compact subsets, in the space of continuous functions on D . Now, notice that as $g_{D \setminus \mathbb{V}_b^u}(z, z) \leq g_D(z, z) < C_r$ for any compact subset $D_r \subset D$, the measure defined by the RHS of (5.3) on any such compact subset D_r is of bounded total mass. As the space of such measures is compact, they have a weak limit and (5.3) holds simultaneously for all f continuous on D_r . As we can take $D_r \rightarrow D$, the lemma follows. \square

In order to deduce Proposition 5.10, we will now make use of a simple estimate relating $g_D(z, z)$ as defined in (2.1) with the distance of z to ∂D . This result follows directly from Koebe's quarter theorem in simply connected domains, but in the case of n -connected domains requires an argument.

Lemma 5.12. *For any D open bounded connected finitely connected domain of \mathbb{C} with non-polar boundary components, there is a constant $c = c(D) > 1$ such that*

$$\forall z \in D, d(z, \partial D) \leq e^{2\pi g_D(z, z)} \leq cd(z, \partial D).$$

Proof. If D is simply connected, then $e^{2\pi g_D(z, z)}$ equals the conformal radius $\text{CR}(z, D)$, and the results follows from Koebe's quarter theorem, with $c = 4$.

For the other cases, note that $2\pi g_D(z, z)$ is the value at z of the harmonic extension of the boundary values $x \mapsto \log(|z - x|)$, $x \in \partial D$. This already implies that

$$2\pi g_D(z, z) \geq \log(d(z, \partial D)).$$

Take B^z a Brownian motion started at $z \in D$ and define $T_{\partial D}$ as the first time B^z hits ∂D . Let δ be the smallest diameter of an inner hole of D . With the Beurling's estimate (Theorem 3.76 of [Law08]) we get that

$$\mathbb{P}(|z - B_{T_{\partial D}}^z| \geq rd(z, \partial D)) \leq Cr^{-1/2},$$

where C is a constant not depending on z , as long as the open disk $B(z, rd(z, \partial D))$ does not entirely contain an inner hole of D . A sufficient condition for that is $rd(z, \partial D) \leq \delta$. Thus,

$$\begin{aligned} 2\pi g_D(z, z) &= \log(d(z, \partial D)) + \int_1^{\frac{\text{diam}(D)}{d(z, \partial D)}} \mathbb{P}(|z - B_{T_{\partial D}}^z| \geq e^s d(z, \partial D)) ds \\ &\leq \log(d(z, \partial D)) + \int_1^{\frac{\delta}{d(z, \partial D)}} \mathbb{P}(|z - B_{T_{\partial D}}^z| \geq e^s d(z, \partial D)) ds \\ &\quad + \mathbb{P}(|z - B_{T_{\partial D}}^z| \geq \delta) \log\left(\frac{\text{diam}(D)}{d(z, \partial D)}\right) \\ &\leq \log(d(z, \partial D)) + C \int_1^{+\infty} e^{-\frac{s}{2}} ds + C \frac{d(z, \partial D)^{1/2}}{\delta^{1/2}} \log\left(\frac{\text{diam}(D)}{d(z, \partial D)}\right). \quad \square \end{aligned}$$

We can now prove Proposition 5.10.

Proof of Proposition 5.10: First, note that $D \setminus \mathbb{A}_{-a}^u$ has a.s. finitely many non simply connected components. By 5.12, there is some $c > 1$, such that for all $z \in D \setminus \mathbb{A}_{-a}^u$,

$$d(z, \mathbb{A}_{-a}^u) \leq e^{2\pi g_{D \setminus \mathbb{A}_{-a}^u}(z, z)} \leq cd(z, \mathbb{A}_{-a}^u).$$

Plugging this into (5.3) and taking f positive, we obtain that on the one hand

$$\limsup_{\gamma \rightarrow 0^+} -\frac{\gamma}{\sqrt{2\pi}} \int_{D \setminus \mathbb{A}_{-a}^u} f(z) \log(d(z, \partial D) + c) d(z, \partial D)^{\frac{\gamma^2}{2}} dz \leq (\nu_{\mathbb{A}_{-a}^u}, f),$$

and on the other hand

$$(\nu_{\mathbb{A}_{-a}^u}, f) \leq \liminf_{\gamma \rightarrow 0^+} -\frac{\gamma}{\sqrt{2\pi}} c^{\frac{\gamma^2}{2}} \int_{D \setminus \mathbb{A}_{-a}^u} f(z) \log(d(z, \partial D)) d(z, \partial D)^{\frac{\gamma^2}{2}} dz.$$

As both limits are equal to (5.1), the lemma follows. \square

5.5. Identification with the Minkowski content measure. In this section, we finish the proof of Theorem 5.1, by showing that the measure $\nu_{\mathbb{A}_{-a}^u}$ agrees with the Minkowski content in the gauge $r \mapsto |\log(r)|^{1/2} r^2$. The proof is ‘‘deterministic’’ and gives a rather general strategy for identifying two measures defined on a fractal set and satisfying coherent scaling.

Let us first introduce some notation: given F a function on $(0, +\infty)$ and a compact subset A of \overline{D} , we define the measure

$$\mathfrak{M}(F) = \mathfrak{M}_A(F) = \mathbf{1}_{D \setminus A} F(d(z, A)) dz.$$

For $r \in (0, 1)$, denote

$$\mathcal{J}_r = |\log(r)|^{1/2} \mathbf{1}_{(0,r)}.$$

The Minkowski content measure A in the gauge $|\log(r)|^{1/2} r^2$ would be the weak limit of $\mathfrak{M}_A(\mathcal{J}_r)$ as $r \rightarrow 0$, if it exists. Finally, for $\gamma > 0$, denote

$$F_\gamma(s) = \frac{\gamma}{\sqrt{2\pi}} \mathbf{1}_{s \in (0,1)} |\log(s)| s^{\frac{\gamma^2}{2}}.$$

Proposition 5.10 can be then interpreted as saying that $\nu_{\mathbb{A}_a^u} = \lim_{\gamma \rightarrow 0} \mathfrak{M}(F_\gamma)$. Notice that in this limit, as we are on a bounded domain, the extra indicator function plays no role. The missing part of Theorem 5.1 then follows from:

Proposition 5.13. *Suppose that $(\mathfrak{M}(F_\gamma), f)$ converges as $\gamma \rightarrow 0$ for all f continuous functions compactly supported in D . Then, also $(\mathfrak{M}(\mathcal{J}_r), f)$ converges as $r \rightarrow 0$ and moreover*

$$\frac{1}{2} \lim_{r \rightarrow 0} (\mathfrak{M}(\mathcal{J}_r), f) = \lim_{\gamma \rightarrow 0} (\mathfrak{M}(F_\gamma), f),$$

for all continuous functions compactly supported in D .

To get some insight into why this proposition might be true, notice that both families \mathcal{J}_r and F_γ satisfy the same scaling property:

$$\begin{aligned} \mathcal{J}_{r'}(s) &= \beta^{-1/2} \mathcal{J}_r(s^\beta), \quad \text{for } \beta = \frac{|\log(r)|}{|\log(r')|}; \\ F_{\gamma'}(s) &= \beta^{-1/2} F_\gamma(s^\beta), \quad \text{for } \beta = \frac{\gamma'^2}{\gamma^2}. \end{aligned}$$

Moreover, for $\gamma = |\log r|^{-1/2}$, one can notice that the maximum of F_γ is also of order $|\log r|^{-1/2}$ and this maximal value is taken at distance $O(r)$.

We prove the proposition in two steps: first we show that if both the limits of $\mathfrak{M}(F_\gamma)$ as $\gamma \rightarrow 0$ and of $\mathfrak{M}(\mathcal{J}_r)/2$ as $r \rightarrow 0$ exist, then they are equal. Then, we show that if the limit of $\mathfrak{M}(F_\gamma)$ exists then so does the limit of $\mathfrak{M}(\mathcal{J}_r)$. From now on, let us take A a closed set such that for all compactly supported f , $(\mathfrak{M}(F_\gamma), f) = (\mathfrak{M}_A(F_\gamma), f)$ converges as γ tends to 0.

5.5.1. Relationship between $\mathfrak{M}(F_\gamma)$ and $\mathfrak{M}(\mathcal{J}_r)$.

Lemma 5.14. *Assume that both $\lim_{r \rightarrow 0} \mathfrak{M}(\mathcal{J}_r)$ and $\lim_{\gamma \rightarrow 0} \mathfrak{M}(F_\gamma)$ exist, then*

$$\lim_{\gamma \rightarrow 0} (\mathfrak{M}(F_\gamma), f) = \frac{1}{2} \lim_{r \rightarrow 0} (\mathfrak{M}(\mathcal{J}_r), f).$$

Proof. Write $F_\gamma(s)$ as

$$\begin{aligned} F_\gamma(s) &= - \int_0^1 F'_\gamma(r) \mathbf{1}_{s \leq r} dr = - \int_0^1 F'_\gamma(r) |\log(r)|^{-1/2} \mathcal{J}_r(s) dr \\ &= \frac{\gamma}{\sqrt{2\pi}} \int_0^1 r^{\frac{\gamma^2}{2}-1} \left(1 - \frac{\gamma^2}{2} |\log(r)|\right) |\log(r)|^{-1/2} \mathcal{J}_r(s) dr. \end{aligned}$$

By changing the order of integration, we get

$$(\mathfrak{M}(F_\gamma), f) = \frac{\gamma}{\sqrt{2\pi}} \int_0^1 r^{\frac{\gamma^2}{2}-1} \left(1 - \frac{\gamma^2}{2} |\log(r)|\right) |\log(r)|^{-1/2} (\mathfrak{M}(\mathcal{J}_r), f) dr.$$

Further, by the change of variable $r = \exp(-t^2/\gamma^2)$ and symmetrization, this equals

$$\int_{-\infty}^{+\infty} (1 - t^2/2) \frac{e^{-t^2/2}}{\sqrt{2\pi}} (\mathfrak{M}(\mathcal{J}_{\exp(-t^2/\gamma^2)}), f) dt.$$

Now, if the limit $\lim_{r \rightarrow 0} \mathfrak{M}(\mathcal{J}_r)$ exists, then $\sup_{r \in (0,1)} |(\mathfrak{M}(\mathcal{J}_r), f)| < \infty$. Thus, from dominated convergence it follows that

$$(\nu_{\mathbb{A}_{-a}^u}, f) = \lim_{\gamma \rightarrow 0} (\mathfrak{M}_{\mathbb{A}_{-a}^u}(F_\gamma), f) = c \lim_{r \rightarrow 0} (\mathfrak{M}_{\mathbb{A}_{-a}^u}(\mathcal{J}_r), f),$$

where

$$c = \int_{-\infty}^{+\infty} (1 - t^2/2) \frac{e^{-t^2/2}}{\sqrt{2\pi}} dt = \frac{1}{2}. \quad \square$$

5.5.2. Convergence of $\mathfrak{M}(\mathcal{J}_r)$. In the last paragraph, we saw that a certain integrated version of the Minkowski content measure $\mathfrak{M}(\mathcal{J}_r)$ converges. We now strengthen it to full convergence. Throughout the section we suppose that $(\mathfrak{M}(F_\gamma), f)$ converges as $\gamma \rightarrow 0$ for all f continuous compactly supported in D .

The idea of the proof is to approximate \mathcal{J}_r by linear combinations $\sum_{k=1}^n c_k F_{\gamma_k}$. As \mathcal{J}_r is not continuous on $[0, 1]$ and not 0 in 0, one cannot expect a uniform approximation over the whole interval. Thus we have to use an uniform approximation over a subinterval of $[0, 1]$ for certain smoothed versions of \mathcal{J}_r , and to argue that the cut out parts do not matter.

As a first step let us show that the Minkowski content measure is bounded. In fact we will show something a tiny bit stronger, that is useful for us in the later approximations:

Lemma 5.15. *Let $q_r(s) = \mathbf{1}_{\{s \in (0,1)\}} \frac{|\log(r \vee s)|}{|\log r|^{1/2}}$. Then, for all f continuous compactly supported in D , $(\mathfrak{M}(q_r), f)$ stays bounded as $r \rightarrow 0$.*

Proof. It is enough to prove the result for f non-negative. Let $\gamma_0 > 0$. There is $C > 0$ such that

$$\mathbf{1}_{s \in (1/4, 1/2)} \leq CF_{\gamma_0}(s) \text{ and } \mathbf{1}_{s \in (1/2, 1)} |\log s| \leq CF_{\gamma_0}(s).$$

Thus, by the scaling of F_γ ,

$$\mathbf{1}_{(2^{-2k+1}, 2^{-2k})}(s) = \mathbf{1}_{(1/4, 1/2)}(s^{2^{-k}}) \leq CF_{\gamma_0}(s^{2^{-k}}) = C2^{-\frac{k}{2}} F_{2^{-\frac{k}{2}} \gamma_0}(s).$$

From where one can see that

$$q_{1/2} \leq C \sum_{k=1}^{+\infty} 2^{-\frac{k}{2}} F_{2^{-\frac{k}{2}} \gamma_0},$$

and as q_r satisfies the scaling $q_r(s) = \beta^{-1/2} q_{1/2}(s^\beta)$ with $\beta = \frac{\log 2}{|\log r|}$, it follows that

$$q_r \leq C \sum_{k=1}^{+\infty} 2^{-\frac{k}{2}} F_{2^{-\frac{k}{2}} \beta \gamma_0}.$$

Thus, the lemma follows from the convergence of $(\mathfrak{M}(F_\gamma), f)$. □

For $\alpha > 2$, let us now define $\mathcal{J}_{1/2, \alpha}$ to be a positive smooth function that is bounded by $|\log 2|^{1/2}$, equals $|\log 2|^{1/2}$ over the interval $[1/2^\alpha, 1/2 - 1/2^\alpha]$ and takes the value 0 outside of the interval $(1/2^{2\alpha}, 1/2)$. Finally, let $\mathcal{J}_{r, \alpha}(s) = \beta^{-1/2} \mathcal{J}_{1/2, \alpha}(s^\beta)$ with $\beta = \frac{\log 2}{|\log r|}$ as above.

We claim that the convergence of $\mathfrak{M}(\mathcal{J}_{r, \alpha})$ as $r \rightarrow 0$ implies that of $\mathfrak{M}(\mathcal{J}_r)$:

Claim 5.16. *Assume that for all $\alpha > 2$ and all continuous f compactly supported in D , $(\mathfrak{M}(\mathcal{J}_{r, \alpha}), f)$ converges as $r \rightarrow 0$. Then, also $(\mathfrak{M}(\mathcal{J}_r), f)$ converges as $r \rightarrow 0$.*

Proof. By definition $\mathcal{J}_{r, \alpha}(s) \leq \mathcal{J}_r(s)$ for all $s \in (0, 1)$. Moreover, similarly

$$\mathcal{J}_r(s) \leq \log(1 + r^\alpha) \mathcal{J}_{r+r^\alpha/2, \alpha}(s) + \alpha^{-1/2} q_{r^\alpha}(s).$$

Hence for all positive f ,

$$(\mathfrak{M}(\mathcal{J}_{r, \alpha}), f) \leq (\mathfrak{M}(\mathcal{J}_r), f) \leq (1 + r^\alpha) (\mathfrak{M}(\mathcal{J}_{r+r^\alpha/2, \alpha}), f) + \alpha^{-1/2} (\mathfrak{M}(q_{r^\alpha}), f).$$

Thus, by letting first $r \rightarrow 0$ and then $\alpha \rightarrow \infty$, we conclude by Lemma 5.15 \square

Finally, it remains to show that

Claim 5.17. *For all $\alpha > 2$ and all continuous f compactly supported in D , we have $(\mathfrak{M}(\mathcal{J}_{r,\alpha}), f)$ converges as $r \rightarrow 0$.*

Proof. Notice that for any n -th degree polynomial Q_n with 0 constant coefficient, we can write $|\log(s)|Q_n(s)$ over $(0, 1)$ as $\sum_{k=1}^n c_{k,n}F_{\gamma_k}(s)$, where $\gamma_k := \sqrt{2k}$ and $c_{k,n}\sqrt{2\pi}/\gamma_k$ is the coefficient in Q_n of the power s^k . Thus, our aim will be to approximate $\mathcal{J}_{1/2,\alpha}$ via polynomials. To do this, let

$$\widehat{\mathcal{J}}_{1/2,\alpha}(s) = |\log s|^{-1} \mathcal{J}_{1/2,\alpha}(s).$$

We claim there exists a sequence of polynomials Q_n such that

$$\lim_{n \rightarrow +\infty} \sup_{s \in (0,1)} |\log(s)| \frac{|\widehat{\mathcal{J}}_{1/2,\alpha}(s) - Q_n(s)|}{q_{1/2}(s)} = 0.$$

Indeed, this is just a consequence of the fact that the polynomials having zeros at 0 and 1 are dense in

$$\{f \in \mathcal{C}^1([0, 1]) : f(0) = f(1) = 0\}$$

for the norm $\|f\|_{\mathcal{C}^1} = \max_{[0,1]} |f| + \max_{[0,1]} |f'|$. One way to see this is to approximate f' by a polynomial P_n and to define $Q_n(x) = \int_0^x (P_n(y))dy - x \int_0^1 P_n(y)dy$.

Thus we obtain that for any $\varepsilon > 0$, there is some $n \in \mathbb{N}$ large enough such that for all $s \in (0, 1)$ we have

$$|\mathcal{J}_{1/2,\alpha}(s) - \sum_{k=1}^n c_{k,n}F_{\gamma_k}(s)| \leq \varepsilon q_{1/2}(s).$$

Now take $\beta = \frac{\log 2}{|\log r|}$ as before. Using the scaling of each of the terms, we see that over $s \in (0, 1)$,

$$|\mathcal{J}_{r,\alpha}(s^\beta) - \sum_{k=1}^n c_{k,n}F_{\beta^{1/2}\gamma_k}(s^\beta)| \leq \varepsilon q_r(s^\beta).$$

Hence for all $\varepsilon > 0$ there exists $n \in \mathbb{N}$ large enough such that

$$\left| (\mathfrak{M}(\mathcal{J}_{r,\alpha}), f) - \left(\mathfrak{M}\left(\sum_{k=1}^n c_{k,n}F_{\beta^{1/2}\gamma_k}\right), f \right) \right| \leq \varepsilon (\mathfrak{M}(q_r), |f|).$$

Combining now Lemma 5.15 and the fact that as $\beta \rightarrow 0$ the sequence $\left(\mathfrak{M}\left(\sum_{k=0}^n c_{k,n}F_{\beta^{1/2}\gamma_k}\right), f\right)$ converges, we conclude that $(\mathfrak{M}(\mathcal{J}_{r,\alpha}), f)$ has a finite limit as $r \rightarrow 0$. \square

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REFERENCES

- [Ahl66] Lars Valerian Ahlfors. *Complex Analysis*. McGraw-Hill Book Company, 2nd edition, 1966.
- [Ahl10] Lars Valerian Ahlfors. *Conformal invariants: topics in geometric function theory*, volume 371. American Mathematical Society, 2010.
- [Aid15] Elie Aïdékon. The extremal process in nested conformal loops. *Preprint available on the webpage of the author*, 2015.
- [ALS18a] Juhan Aru, Titus Lupu, and Avelio Sepúlveda. Excursion decomposition of the 2D Gaussian free field. In preparation, 2018.
- [ALS18b] Juhan Aru, Titus Lupu, and Avelio Sepúlveda. Extremal distance of CLE_4 . In preparation, 2018.
- [ALS18c] Juhan Aru, Titus Lupu, and Avelio Sepúlveda. The first passage sets of the 2D Gaussian free field: convergence and isomorphisms. *arXiv preprint arXiv:1805.09204*, 2018.
- [APS17] Juhan Aru, Ellen Powell, and Avelio Sepúlveda. Approximating Liouville measure using local sets of the Gaussian free field. *arXiv preprint arXiv:1701.05872*, 2017.
- [APS18] Juhan Aru, Ellen Powell, and Avelio Sepúlveda. Critical Liouville measure as a limit of subcritical measures. *arXiv preprint arXiv:1802.08433*, 2018.
- [Aru15] Juhan Aru. *The geometry of the Gaussian free field combined with SLE processes and the KPZ relation*. PhD thesis, Ecole normale supérieure de lyon-ENS LYON, 2015.
- [Aru17] Juhan Aru. Gaussian multiplicative chaos through the lens of the 2D Gaussian free field. *arXiv preprint arXiv:1709.04355*, 2017.
- [AS18a] Juhan Aru and Avelio Sepúlveda. Survey in preparation. 2018.
- [AS18b] Juhan Aru and Avelio Sepúlveda. Two-valued local sets of the 2d continuum gaussian free field: connectivity, labels, and induced metrics. *arXiv preprint arXiv:1801.03828*, 2018.
- [ASW17] Juhan Aru, Avelio Sepúlveda, and Wendelin Werner. On bounded-type thin local sets of the two-dimensional Gaussian free field. *Journal of the Institute of Mathematics of Jussieu*, pages 1–28, 2017.
- [BSS14] Nathanaël Berestycki, Scott Sheffield, and Xin Sun. Equivalence of Liouville measure and Gaussian free field. *arXiv preprint arXiv:1410.5407*, 2014.
- [DRSV14] Bertrand Duplantier, Rémi Rhodes, Scott Sheffield, and Vincent Vargas. Critical Gaussian multiplicative chaos: convergence of the derivative martingale. *The Annals of Probability*, 42(5):1769–1808, 2014.
- [DS11] Bertrand Duplantier and Scott Sheffield. Liouville quantum gravity and KPZ. *Inventiones mathematicae*, 185(2):333–393, 2011.
- [Dub09a] Julien Dubédat. Duality of Schramm Lxoewner evolutions. *Ann. Sci. Éc. Norm. Supér.(4)*, 42(5):697–724, 2009.
- [Dub09b] Julien Dubédat. SLE and the free field: partition functions and couplings. *Journal of the American Mathematical Society*, 22(4):995–1054, 2009.
- [Duf62] Richard Duffin. The extremal length of a network. *Journal of Mathematical Analysis and Applications*, 5(2):200–215, 1962.
- [Gaw96] Krzysztof Gawędzki. Lectures on conformal field theory. *Nuclear Physics B*, 328:733–752, 1996.
- [HK71] Raphael Hoegh-Krohn. A general class of quantum fields without cut-offs in two space-time dimensions. *Communications in Mathematical Physics*, 21(3):244–255, 1971.
- [Kah85] Jean-Pierre Kahane. Sur le chaos multiplicatif. *Annales des sciences mathématiques du Québec*, 9(2):105–150, 1985.
- [Koe22] Paul Koebe. Über die Konforme Abbildung Endlich-und Unendlich-Vielfach Zusammenhängender Symmetrischer Bereiche. *Acta Mathematica*, 43(1):263–287, 1922.
- [Law08] Gregory Lawler. *Conformally invariant processes in the plane*. Number 114. American Mathematical Society, 2008.
- [LJ10] Yves Le Jan. Markov loops and renormalization. *The Annals of Probability*, 38(3):1280–1319, 2010.
- [LJ11] Yves Le Jan. Markov paths, loops and fields. In *2008 St-Flour summer school, Lecture Notes in Mathematics*, volume 2026. Springer, 2011.
- [LW16] Titus Lupu and Wendelin Werner. The random pseudo-metric on a graph defined via the zero-set of the Gaussian free field on its metric graph. *Probability Theory and Related Fields*, pages 1–44, 2016.
- [MS16a] Jason Miller and Scott Sheffield. Imaginary geometry I: interacting sles. *Probability Theory and Related Fields*, 164(3-4):553–705, 2016.
- [MS16b] Jason Miller and Scott Sheffield. Imaginary geometry II: reversibility of $SLE_\kappa(\rho_1; \rho_2)$ for $\kappa \in (0, 4)$. *The Annals of Probability*, 44(3):1647–1722, 2016.
- [MS16c] Jason Miller and Scott Sheffield. Imaginary geometry III: reversibility of $SLE_\kappa(\rho_1; \rho_2)$ for $\kappa \in (4, 8)$. *Annals of Mathematics*, 184(2):455–486, 2016.

- [MS17] Jason Miller and Scott Sheffield. Imaginary geometry IV: interior rays, whole-plane reversibility, and space-filling trees. *Probability Theory and Related Fields*, 169:729–869, 2017.
- [PW17a] Eveliina Peltola and Hao Wu. Global and local multiple SLEs for $\kappa \leq 4$ and connection probabilities for level lines of GFF. *arXiv preprint arXiv:1703.00898*, 2017.
- [PW17b] Ellen Powell and Hao Wu. Level lines of the Gaussian free field with general boundary data. *Annales de l’Institut Henri Poincaré Probabilités et Statistiques*, 53(4):2229–2259, 2017.
- [RV14] Rémi Rhodes and Vincent Vargas. Gaussian multiplicative chaos and applications: a review. *Probability Surveys*, 11:315–392, 2014.
- [Sep17] Avelio Sepúlveda. On thin local sets of the Gaussian free field. *arXiv preprint arXiv:1702.03164*, 2017.
- [She05] Scott Sheffield. Local sets of the Gaussian free field: Slides and audio, 2005.
- [Sim74] Barry Simon. *The $P(\Phi)_2$ Euclidean (Quantum) Field Theory*. Princeton Series in Physics. Princeton University Press, 1974.
- [SS09] Oded Schramm and Scott Sheffield. Contour lines of the two-dimensional discrete Gaussian free field. *Acta Mathematica*, 202(1):21, 2009.
- [SS13] Oded Schramm and Scott Sheffield. A contour line of the continuum Gaussian free field. *Probability Theory and Related Fields*, 157(1-2):47–80, 2013.
- [Wer16] Wendelin Werner. Topics on the GFF and CLE(4), 2016.
- [WW16] Menglu Wang and Hao Wu. Level lines of Gaussian free field I: zero-boundary GFF. *Stochastic Processes and their Applications*, 2016.

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