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Direct numerical simulation of nucleate boiling in micro-layer regime

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A B S T R A C T

The physical mechanisms associated with the evolution of a micro-layer beneath a bubble and the transition between contact line and micro-layer regimes are investigated with fully resolved numerical simulations, in the framework of nucleate pool boiling. Capturing the transition between these two regimes has been possible for the first time using very refined grids and parallel computations. Indeed, grids with a cell size under 1 μm must be used in order to capture thermal and dynamical effects in the micro-layer. Such multiscale computations require advanced code capabilities. The present simulations are used to analyse the physical processes involved in the formation and depletion of a micro-layer. A parametric study is carried out to investigate the impact of the main parameters affecting the presence of the micro-layer. From these results, the limit conditions between nucleate boiling in micro-layer and contact line regimes are deduced. Neglecting the micro-layer would lead to erroneous results because it has a strong influence on the overall bubble growth. Therefore the present results could be of major interest for designing models of nucleate pool boiling in larger scales computations, when the micro-layer cannot be resolved.

1. Introduction

Nucleate boiling is a physical phenomenon encountered in a large number of engineering applications. Its features are sometimes required, like in some heat transfer devices, and in other circumstances not desirable, like for instance in cryogenics tanks in microgravity conditions. Despite the large amount of studies concerning nucleate boiling, there is still a number of unanswered questions, in particular related to the region beneath a bubble where for some conditions a micro-layer is formed and persists during the first part of the bubble growth. It is a tiny film, which has a thickness smaller than ten microns and a length up to the millimetre. Actually, we can distinguish between two regimes of nucleate boiling, like schematically illustrated in Fig. 1 where the region around the bubble foot is represented: the contact line (CL) regime and the micro-layer (ML) regime [26]. In contact line regime the bubble maintains a quasi-static equilibrium shape during its growth [10,19]. The contact line moves while the bubble grows and the interface maintains a slope with respect to the wall given by the contact angle θ_1 . On the other hand, when the bubble growth is fast enough, it may happen that the contact line cannot move as fast as the bubble: consequently the interface is bended and a micro-layer is formed. We can thus distinguish between

two angles: the contact angle θ_1 formed at the contact line between the interface and the wall, and an apparent contact angle θ_2 which is given by the inclination of the interface outside the micro-layer.¹ By locally inducing high temperature gradients, the presence of a micro-layer strongly enhances the heat transfer between the solid heater and the bubble [21,9]. Therefore, one major difference between the two regimes is the wall heat flux profile q_w which has only a peak at the contact line in the CL regime whereas it exhibits a peak at the contact line followed by a plateau along the micro-layer in the ML regime [9], as schematically shown in Fig. 1. As a consequence, the micro-layer strongly affects the growth of the bubble, its departure diameter and frequency as well as its shape. Micro-layer effects have been observed experimentally since the pioneer work of Cooper and Lloyd [5]. However, micro-layers are thin and long and have a life time of only a few milliseconds which make the measurements tedious. More recently, the advance in techniques like high speed infrared thermography [25] and laser interferometry [12] together with high-speed cameras, have permitted to obtain interesting experimental analysis showing the formation and depletion of micro-layers with quantitative data on their thickness and length temporal evolutions [3,21,55]. Experimental

¹ Note that θ_1 is also an apparent contact angle at the outlet of a micro-region and it is related to a smaller microscopic angle θ_{mic} [26,27,38,11]. The micro region around the contact line, smaller than the micron, is not the focus of the present work and will not be considered in our discussion. See Appendix A for more details.

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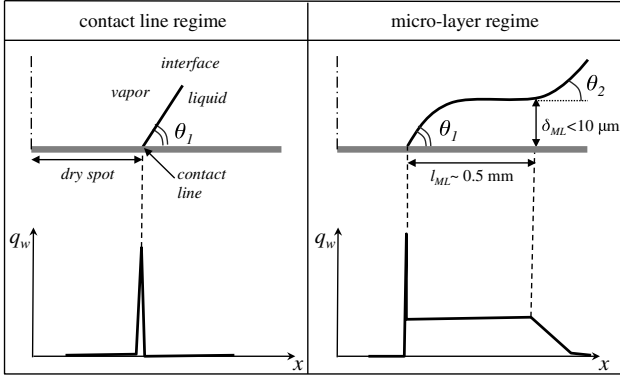


Fig. 1. Scheme of the different nucleate boiling regimes, contact line (left) and micro-layer (right), in terms of interface shape at the bubble foot (up) and wall heat flux profile (down).

evidences of the transition between contact line and micro-layer regime have been reported recently in [9] for a configuration involving a liquid meniscus moving on a heated wall. All these works have emphasised the importance of the micro-layer in the heat transfer process and the evidence that if a micro-layer exists it has to be taken into account in the nucleate boiling calculations [16]. Indeed numerical models have been developed to account for the micro-layer effects without directly solving it and are usually introduced in the numerical computations as boundary conditions [4,55] or as local source terms [43,44]. Recently, Hänsh and Walker [17] carried out a direct numerical simulation of the nucleation of a bubble with a resolved micro-layer beneath a bubble. However, they consider a purely hydrodynamic modelling without accounting for heat and mass transfer in their computation, but evaluating the evaporative mass flow rate from an offline computation. Differently, for what concerns CL regime, for which the grid refinement requirements are less stringent, fully resolved DNS simulations are available in the literature and have demonstrated their capability in correctly reproducing the phenomenon [27,45,48].

Several points remain unclear about the physical processes involved in the formation and depletion of micro-layers. In particular, it is not fully understood for which conditions nucleate boiling occurs in contact line or in micro-layer regime. The main motivation of the present work is to provide detailed informations in order to characterise the transition between the regimes, using fully resolved Direct Numerical Simulations (DNS). The in-house solver *DNVA* has been used to solve the incompressible Navier-Stokes equations for two phase flows and to account for phase change. Axisymmetric configurations have been employed, with very refined grids allowing the description of a possible micro-layer having a thickness of the order of the micron, beneath a bubble having a radius of the order of the millimetre. Simulations of this multi-scale phenomenon are challenging and require using massively parallel supercomputers and suitable numerical tools to verify that the results are not affected by grid dependence effects. Moreover, it is noteworthy that micro-layers involve locally very high heat fluxes in the range of several MW, that could induce stability issues for numerical solvers if the thermal gradients are not sufficiently resolved. To the authors knowledge, there are no other DNS results with phase change available in the literature showing the formation and depletion of a micro-layer and the present work is the first attempt to carry out a fully resolved DNS of nucleate boiling in micro-layer regime.

In the next section the numerical solver is introduced, recalling the mathematical formalism and giving some details on the numerical methods that are implemented in the code. A simple axisymmetric domain is considered: details on the boundary con-

ditions and mesh requirements are given. A convergence study and a first comparison with experimental data in micro-layer regime are presented. The simulation of a bubble nucleation in micro-layer regime, showing the formation and depletion of the micro-layer, is deeply analysed and the basic physical mechanisms driving the phenomenon are interpreted. The influence of the main parameters is thus investigated with a parametric study varying the wall superheat, the contact angle, the surface tension and the liquid viscosity, and looking for the limit between CL and ML regimes. Eventually, a correlation for the frontier between the two regimes is proposed.

2. Numerical solver

2.1. Mathematical formalism

To describe the nucleate boiling at atmospheric pressure, in a isobaric environment, we consider an incompressible two-phase flow of a single component liquid-vapour system, with spatially uniform fluid densities ρ and thermophysical properties in each phase (viscosity μ , thermal conductivity k and specific heat at constant pressure c_p). Moreover, we assume that the heat produced by viscous dissipation is negligible. Under these hypothesis the flow can be described by the following form of the conservation equations of mass, momentum and energy in each phase:

$$\nabla \cdot \vec{V} = 0 \quad (1)$$

$$\rho \frac{D\vec{V}}{Dt} = -\nabla p + \nabla \cdot (2\mu\mathbf{D}) + \rho \vec{g} \quad (2)$$

$$\rho c_p \frac{DT}{Dt} = \nabla \cdot (k\nabla T) \quad (3)$$

where \vec{V} is the velocity field, p is the pressure, T is the temperature, \mathbf{D} is the deformation tensor and \vec{g} is the gravity acceleration. The corresponding mass, momentum and energy jump conditions that must be satisfied at the interface Γ , accounting for the phase change, are:

$$[\vec{V}]_{\Gamma} = \dot{m} \left[\frac{1}{\rho} \right]_{\Gamma} \vec{n} \quad (4)$$

$$[p]_{\Gamma} = \sigma \kappa + 2 \left[\mu \frac{\partial V_n}{\partial n} \right]_{\Gamma} - \dot{m}^2 \left[\frac{1}{\rho} \right]_{\Gamma} \quad (5)$$

$$[-k\nabla T \cdot \vec{n}]_{\Gamma} = \dot{m} L_{vap} \quad (6)$$

where σ is the surface tension, κ denotes the local interface curvature, \vec{n} is the normal vector at the interface, pointing towards the liquid phase, V_n is the velocity component in the \vec{n} direction, \dot{m} is the boiling mass flow rate and L_{vap} is the latent heat of vaporisation. The operator $[\cdot]_{\Gamma}$ indicates the jump across the interface Γ and is defined by:

$$[f]_{\Gamma} = f_v - f_l \quad (7)$$

with the subscripts l and v referring respectively to the liquid and vapour phases. Note that the present form of the energy jump condition Eq. (6) assumes that the interface temperature is equal to the saturation temperature ($T_{\Gamma} = T_{sat}$) in accordance with the second law of thermodynamics for a pure liquid-vapour system at local thermodynamic equilibrium. Moreover, for the sake of simplicity, we assume that T_{sat} is constant, which implies that the surface tension σ is constant. As a consequence, Marangoni convection effects (due to Kelvin effects for instance) are neglected. Interested readers can find more details on this specific point in [22].

2.2. Numerical schemes

The in house solver *DTVA* is used to carry out the numerical analysis in the present work. It is a structured incompressible solver, for two-phase flows simulations which is based on the Ghost Fluid Method [8,23,29,31,51], allowing a sharp discretisation of the jump conditions across the interface, and on the Level Set method to compute the motion of the interface [36]. A preserving distance algorithm ensures the signed distance property of the Level Set function in the whole domain [50]. Following the general guidelines of the Ghost Fluid Method, phase change simulations can be considered by imposing additional jump conditions to maintain mass conservation [35], energy conservation [14,15,54] and chemical species conservation [37,53,41]. Moreover, as demonstrated in [54], global accuracy of the numerical solver can be improved with second order extrapolations for the temperature field performed with the numerical methods described in [1]. Fifth order WENO-Z schemes are used for the convective terms [2], second order finite differences for the diffusion terms. A Black-Box MultiGrid solver [7] permits a fast and stable resolution of the pressure Poisson equation which comes out from the projection method, with a computational time only weakly dependent on the density ratio [32]. The temporal integration is performed with a second order Runge-Kutta scheme. Our in house solver also permits computations involving compressible two-phase flows [18] and the simulation of complex geometries on structured grid making use of an Immersed Interface Method for both single phase flows [34] and two-phase flows [30]. The solver has been extensively validated with experimental data both for isothermal two-phase flows [28,52] and for two-phase flows with phase change for impacting droplets in the Leidenfrost regime [42] and for nucleate pool boiling [19]. In particular, Huber et al. [19] have presented detailed comparisons between numerical results and experimental data on nucleate boiling in contact line regime, showing the ability of the solver in retrieving the correct departure diameter and departure frequency of a single bubble detaching from a nucleation site. This work showed, with a grid sensitivity analysis, that a grid cell size of about $10 \mu\text{m}$ is required in the vicinity of the contact line in order to achieve well resolved simulations of nucleate boiling in contact line regime. The simulations of boiling flows carried out in the present paper and in the former studies [14,19,54] are based on a Ghost Fluid Thermal Solver for Boiling (GFTSB) for computing the heat transfer and the phase change around the interface [41].

3. Configuration and numerical validations

3.1. Computational domain and initial conditions

The axisymmetric computational domain that is considered in the present study is shown in Fig. 2(a). It is a square of L dimension, limited by an isothermal wall in $y = 0$, an adiabatic wall in $x = L$ and opened in $y = L$. A symmetry condition is inferred in $x = 0$. Dirichlet conditions on the temperature are applied in $y = 0$, and Neumann conditions in $x = L$ and $y = L$. For the velocity field, no slip conditions are considered at the walls, whereas a free boundary condition is used at the upper boundary ($y = L$).

The considered fluid is water at atmospheric conditions and the corresponding thermophysical properties are reported in Table 1. The sea level value for the gravity acceleration is used, $g = 9.81 \text{ m} \cdot \text{s}^{-2}$. The initial thermal field is constant and equal to $T_{\text{sat}} = 373.12 \text{ K}$ everywhere except in a layer near the isothermal wall where a linear profile is applied from T_{sat} to T_w . The thickness of the thermal layer δ_{KC} is computed from the correlation of Kays and Crawford [24] for free convection given by Eq. (8):

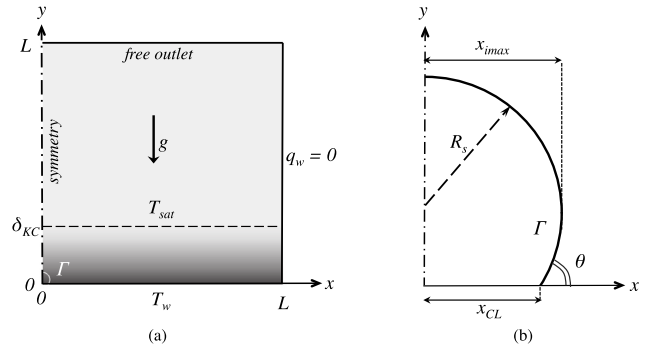


Fig. 2. Configuration used for the axisymmetric numerical simulations. (a) Computational domain, initial condition and boundary conditions. (b) Zoom view of the vapour seed which is used in the initial condition.

$$\delta_{\text{KC}} = 7.14 \left(\frac{\mu_l \alpha_l}{\rho_l g \beta_l (T_w - T_{\text{sat}})} \right)^{1/3} \quad (8)$$

In Eq. (8), α_l denotes the liquid thermal diffusivity and β_l the isobaric thermal expansion coefficient of the liquid, which is equal to $\beta_l = 7.52 \times 10^{-4} \text{ K}^{-1}$ for water at atmospheric conditions.

A vapour seed is initially deposited in $x = y = 0$ as shown in Fig. 2(a). The initial shape of the seed is reported in the zoom view of Fig. 2(b): it is a circle of radius R_s shifted upward in order to have the desired contact angle θ at the wall. The same contact angle will then be imposed during the computation with the Neumann boundary condition on the Level Set function² (see Huber et al. [19] for more details).

3.2. Mesh requirements

Preliminary simulations have been carried out to find the required level of grid refinement to capture hydrodynamic and thermal effects leading to the formation of a micro-layer. For this study we have considered a square of side $L = 2 \text{ mm}$ (see Fig. 2). The initial seed has a radius $R_s = 60 \mu\text{m}$ and the imposed contact angle is $\theta = 50^\circ$. The wall temperature is $T_w = 393.12 \text{ K}$ which corresponds to a Jakob number $\text{Ja} = 60$ defined by Eq. (9):

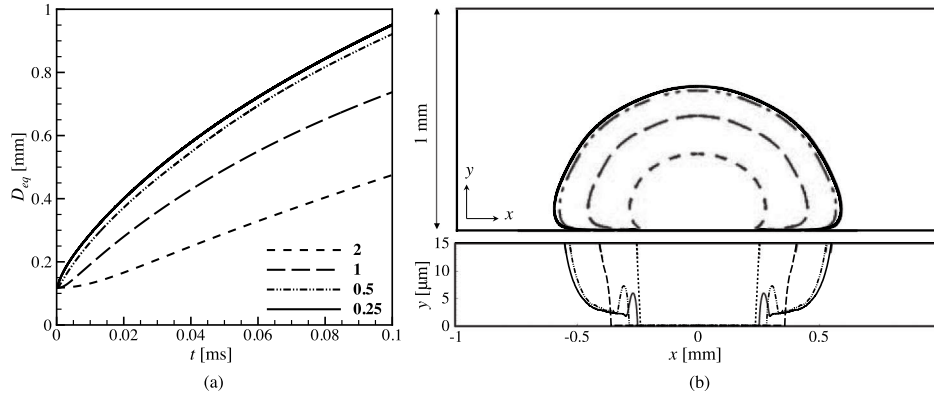
$$\text{Ja} = \frac{\rho_l c_{pl} (T_w - T_{\text{sat}})}{\rho_v L_{\text{vap}}} \quad (9)$$

Four uniform grids, with equal resolutions in both Cartesian directions, have been used. The cell size of the grids are equal to 2, 1, 0.5 and $0.25 \mu\text{m}$ which correspond respectively to 1024^2 , 2048^2 , 4096^2 and 8192^2 elements. Simulations on the latter two grids have been performed respectively on 64 and 256 processors by using the MPI library for parallel computations. A total time of 0.1 ms has been computed, which is sufficient to observe the formation of the micro-layer. Results are shown in Fig. 3 in terms of temporal evolution of the equivalent vapour diameter D_{eq} (a) and bubble shapes at $t = 0.1 \text{ ms}$ (b) for the different mesh resolutions. Also a detail of the interface in the region near the wall is reported in Fig. 3(c) to help the visualisation of the region where the micro-layer is formed (note that this image is not scaled to help the visualisation of the micro-layer). First, we observe that a micro-layer is present in the simulations with the two finer grids (cell size equal to 0.25 and $0.5 \mu\text{m}$), whereas no micro-layer is found with the two coarser grids (1 and $2 \mu\text{m}$). This can be explained considering that in order to simulate the thin micro-layer, the grid must be fine enough to correctly describe the hydrodynamic scales associated

² The imposed angle θ is an apparent contact angle corresponding to θ_1 in Fig. 1. See Appendix A for more details.

Table 1Thermophysical conditions for water in liquid and vapour phase at atmospheric pressure and $T_{sat} = 373.12$ K.

	ρ [kg/m ³]	μ [Pa·s]	k [W/K/m]	c_p [J/kg/K]	σ [N/m]	L_{vap} [J/kg]
Liquid	958	2.82×10^{-4}	0.677	4216	0.058	2.256×10^6
Vapour	0.5974	1.228×10^{-5}	0.024	2034	–	–

**Fig. 3.** Temporal evolution of the equivalent vapour diameter (a), bubble shape at $t = 0.1$ ms (b, up) and zoom of the bubble shape nearby the heated wall at $t = 0.1$ ms (b, down) for several levels of grid refinement (cell size equal to 0.25/0.5/1 and 2 μ m).

with its formation. Otherwise, if the simulation is under resolved, a contact line regime is retrieved and this strongly affects the results. In fact, as illustrated in Fig. 3(a), with the 1 and 2 μ m meshes, the growth rate is strongly underestimated with differences with respect to the 0.25 μ m mesh up to 54% and 24% respectively for D_{eq} . On the other hand, only slight differences, smaller than 4%, are observed between the two finest grids, which demonstrates the good spatial convergence of these simulations when sufficiently refined grids are employed. Both the finest grids permit to compute the formation of the micro-layer despite slight differences around the contact line position. The micro-layer has a minimum thickness around 2 μ m, which corresponds to 4 grid points for the 0.5 μ m mesh. The instantaneous field of velocity vectors at $t = 0.1$ ms, obtained with the 0.5 μ m mesh, is plotted in Fig. 4 around the whole bubble, (a), and around the micro-layer, (b) and (c). These figures allow the visualisation of the strong velocity jump across the interface, induced by the very high thermal flux in the micro-layer region which leads to a strong evaporation of the liquid in the layer. We can conclude from these results that the 0.5 μ m mesh is suitable to simulate the formation of the micro-layer during nucleate boiling of a bubble, and it will be used to carry out the simulations of the present study.

It is noteworthy that the simulations of this work are the first to fully resolve the hydrodynamic and thermal effects initiating the formation of a micro-layer at very low scale. This has been made possible by using the same numerical solver in the whole computational domain thanks to recent advances in the field of numerical methods for DNS of two-phase flows with phase change.

3.3. Micro-layer and contact line regimes

We are now concerned about the ability of the solver to differentiate between nucleate boiling in contact line or micro-layer regime, in the configuration presented in the previous section, if the 0.5 μ m grid is used. Another case has thus been simulated, keeping the same conditions of the above convergence study case, except for the wall temperature which has been lowered to $T_w = 380.12$ K which corresponds to $Ja = 21$. These are the conditions of the experimental data of Son et al. [48] for nucleate boiling in contact line regime, for which the *DDVA* solver has been previ-

ously validated in the recently published work by Huber et al. [19].³ As expected nucleate boiling occurs in contact line regime for $Ja = 21$. The instantaneous bubble interface, the temperature field and the wall heat flux profile at $t = 0.2$ ms are shown in Fig. 5(a) for $Ja = 21$ in contact line regime and in Fig. 5(b) for $Ja = 60$ in micro-layer regime (which is the case of the convergence study).

It appears that the micro-layer formation strongly modifies the bubble shape which becomes more flat than the bubble in contact line regime, which keeps an almost spherical shape. Another major consequence of the micro-layer is observable in the wall heat flux profile. In *CL* regime q_w only shows a peak at the contact line, whereas in *ML* regime q_w shows a peak at the contact line and then remains high all along the micro-layer extension, as was schematically introduced in Section 1 (see Fig. 1). These results show that the present setup is suitable to investigate on the transition from contact line to micro-layer regime.

3.4. Micro-layer thickness and length

To quantitatively verify our results in *ML* regime, we have carried out a simulation with the same configuration (see Fig. 2) but in conditions that are in the range of the experimental data of Jung and Kim [21], that is $T_w = 396.12$ K corresponding to $Ja = 69$, and a contact angle of $\theta = 65^\circ$. The aim of this analysis is to verify that the micro-layer characteristics that we compute are in agreement with experimental observations in similar conditions. The experimental observations report the formation of a micro-layer with a thickness $\delta_{ML,exp} < 3$ μ m and an extension $l_{ML,exp} \sim 0.15 - 0.6$ mm. These values have been deduced from Fig. 7 and 8 in [21]. The numerical results also show the formation of a micro-layer as observable in Fig. 6 where the instantaneous shape of the micro-layer and the temperature field are reported. As a measure of δ_{ML} we consider the minimum micro-layer thickness and l_{ML} is evaluated as the distance between the radial position where the slope of the interface ψ becomes greater than 10° and the contact line

³ Note that the grid resolution used in the validations carried out in [19] is 10 μ m, which is a sufficient resolution to capture the detachment diameter in quasi-static conditions which arise from a force balance.

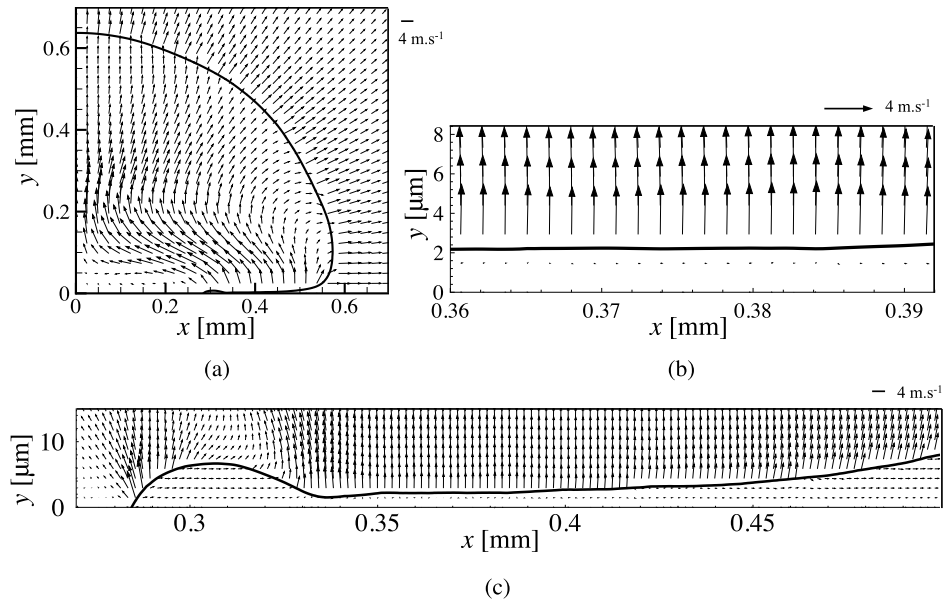
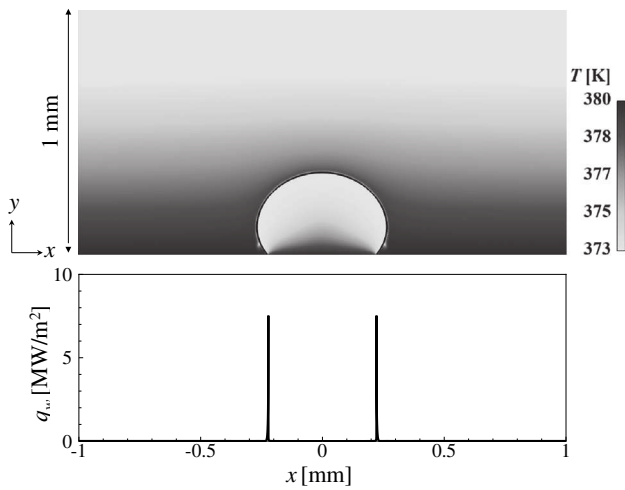
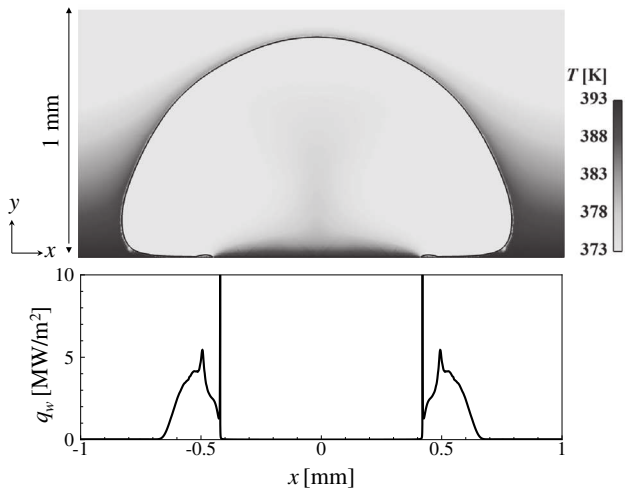


Fig. 4. Velocity vectors for the whole bubble (a) and zoom views around the micro-layer region (b and c) for the results of the $0.5 \mu\text{m}$ grid at $t = 0.1 \text{ ms}$.



(a) $Ja = 21$, contact line regime



(b) $Ja = 60$, micro-layer regime

Fig. 5. Bubble interface, temperature field and wall heat flux profile at $t = 0.2 \text{ ms}$ in contact line regime for $Ja = 21$ (a) and in micro-layer regime for $Ja = 60$ (b).

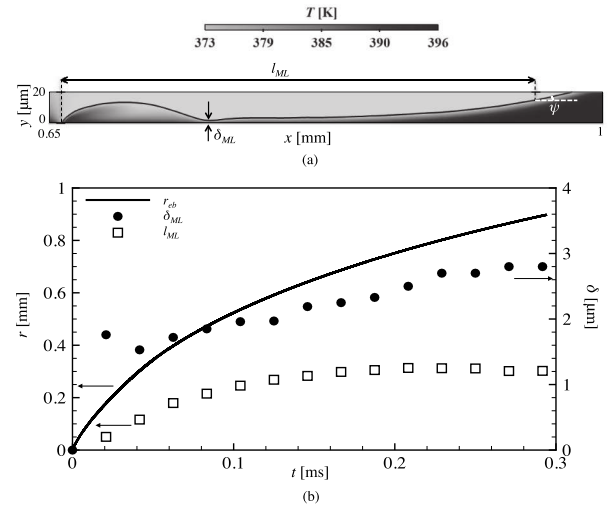


Fig. 6. Results for $Ja = 69$ and $\theta = 65^\circ$. Micro-layer visualisation and temperature field at $t = 0.29 \text{ ms}$ (a). Temporal evolution of the equivalent bubble radius r_{eb} , thickness δ_{ML} and axial extension l_{ML} of the micro-layer (b).

position (see Fig. 6). The micro-layer thickness δ_{ML} varies between 2 and 3 μm , its extension l_{ML} is around 0.3 mm and their temporal evolutions are reported in Fig. 6 together with the bubble equivalent radius. These values are in the range of the experimental data which demonstrate that the present numerical simulations are able to reproduce the micro-layer regime if it occurs.⁴

4. Numerical analysis of the growth and depletion of a micro-layer

This section is dedicated to the analysis of the formation and depletion of a micro-layer. For this purpose we have selected a square domain of dimension $L = 3 \text{ mm}$, which is greater than the previous domain, to allow the bubble to grow more, until the

⁴ Note that the temporal evolutions of the experimental data are slower than the numerical evolutions, which can be explained by the fact that, in a first approximation, we are imposing an isothermal boundary condition without accounting for the wall temperature variation due to heat conduction in the solid heater.

micro-layer depletion. The mesh has a uniform resolution of $0.5 \mu\text{m}$ in a 1 mm square from the nucleation site ($0 < x < 1 \text{ mm}$ and $0 < y < 1 \text{ mm}$), and it is coarsened towards $y = L$ and $x = L$ up to $1.5 \mu\text{m}$ (note that the most coarsened regions are not reached by the bubble). The same boundary conditions of the convergence case are inferred that is $\text{Ja} = 60$ and $\theta = 50^\circ$. A total physical time of 1.5 ms has been simulated. A micro-layer is formed and depleted in less than 1 ms , as can be observed in Fig. 7(a) where 5 instantaneous fields of temperature and bubble interfaces are reported, in a temporal interval that goes from 0.1 to 0.9 ms . Zoomed views in a fixed rectangular box of 1.1 mm length and $20 \mu\text{m}$ height are shown in Fig. 7(b) for the same 5 temporal instants together with the corresponding wall heat flux profiles. The wall heat flux initially has a profile similar to that expected in a *ML* regime (c.f. Fig. 1). The micro-layer length increases ($t = 0.3 \text{ ms}$) and then decreases up to its disappearance. At $t = 0.9 \text{ ms}$ the heat flux profile has the typical peak shape of the *CL* regime.

An amplified view of the micro-layer formation and depletion is shown in the velocity magnitude fields for the same five temporal instants in Fig. 8(a): streamlines coloured with the axial velocity (v_y) are also reported. At $t = 0.1 \text{ ms}$, the interface is bended just after the contact line and forms a rim followed by a very tiny film of liquid, with a thickness approximately equal to $3 \mu\text{m}$. Note that a rim has also been observed in the numerical results of Hänsh et al. [17], which are based on an hydrodynamic formulation (no heat transfer). Actually, this interface shape is typical of dewetting

phenomena: when a hole is nucleated on a film substrate the liquid retracts and a rim is formed which collects the liquid removed from the dry region, like shown for instance in the experiments of Redon et al. [39,40]. Similar observations have also been reported in the experimental work by Fourgeaud et al. [11] on the dynamic of a dewetting film with evaporation. For what concerns the micro-layer, the shape of the rim impacts the heat flux profile of Fig. 7(b), showing a secondary peak at the radial position corresponding to the minimum height of the interface between the rim and the film. In the rim the flow induced by the contact line movement is parallel to the wall, whereas, in the liquid film the flow is almost at rest. The micro-layer length increases from $t = 0.1 \text{ ms}$ to $t = 0.3 \text{ ms}$ and then decreases. Actually, when the rim grows, it pushes out the liquid film, which eventually disappears. At $t = 0.7 \text{ ms}$ the interface appears just slightly bended and at $t = 0.9 \text{ ms}$ it has a constant slope like in the *CL* regime. These observations suggest that the micro-layer depletion is an hydrodynamic phenomenon, which is different than the usually adopted hypothesis of the micro-layer being depleted by evaporation of the film [5,43]. The micro-layer shape is such that the slope of the interface is equal to θ at the contact line position, and then decreases along the radial direction, goes down around zero along the micro-layer extension and then grows up again after the micro-layer end. Therefore, a good indicator of the micro-layer length l_{ML} is the distance between the radial position where the interface slope is equal to the contact angle ($\psi = \theta$) and the contact-line

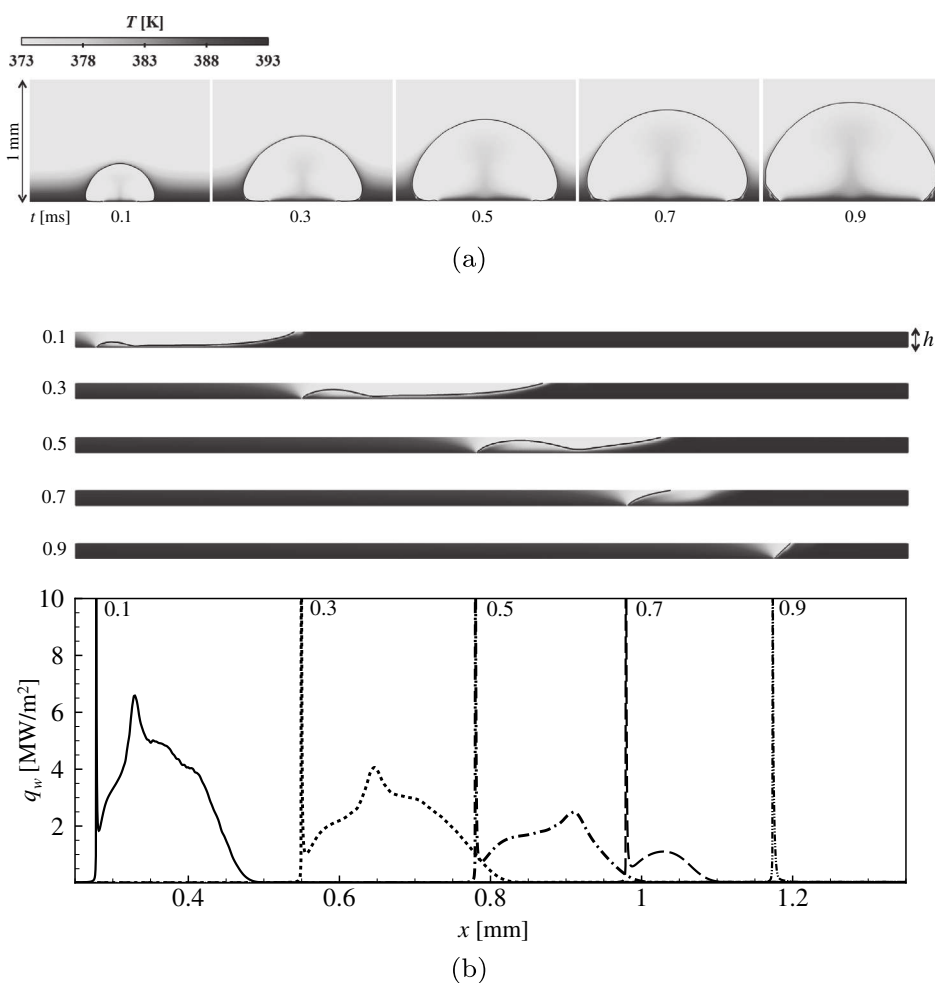


Fig. 7. (a) Instantaneous interface shapes and temperature fields of a bubble nucleated in micro-layer regime, with $\text{Ja} = 60$ and $\theta = 50^\circ$ (same grey scale of Fig. 3), for 5 temporal instants. (b) Zoom on a fixed box of dimensions $h = 20 \mu\text{m}$ and $\Delta x = 1.1 \text{ mm}$, and corresponding wall heat flux profiles, for the same 5 instants (saturated scale at 10 MW/m^2).

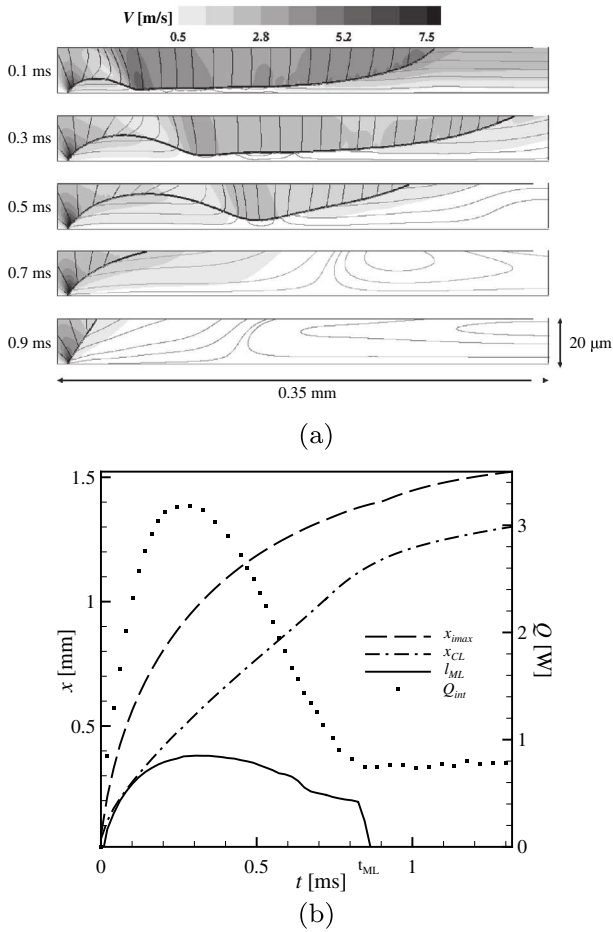


Fig. 8. Velocity fields and streamlines coloured by the axial velocity v_y for five temporal instants, in a constant dimension box moving with the contact line (a). Temporal evolution of the contact line position x_{CL} , the bubble maximal interface position x_{imax} , the micro-layer length l_{ML} and the integrated wall heat flux Q_{int} (b).

position (see Fig. 6(a)). To follow the formation and depletion of the micro-layer, the temporal evolution of l_{ML} is shown in Fig. 8 (b) where the radial position of the contact line x_{CL} and the maximum interface position x_{imax} (see Fig. 2) have also been reported. Initially the velocity of the bubble growth is higher than the contact line speed: as a consequence the interface is stretched in the radial direction and ends up forming a micro-layer as shown by the l_{ML} growth. After $t = 0.3$ ms, l_{ML} decreases, as a consequence of the micro-layer depletion and goes to zero around $t_{ML} = 0.85$ ms, which is the lifetime of the present micro-layer, in agreement with the instantaneous field shown in Fig. 8(a). The wall heat flux is strongly affected by the micro-layer and its integrated value from $x = 0$ to $x = L$ is also reported in the same figure (Q_{int}). It increases during the micro-layer growth and decreases during its depletion towards an almost constant value when a contact line regime is retrieved.⁵

5. Frontier between contact line and micro-layer regime

We may now investigate on the physical parameters that favour the formation of a micro-layer during the bubble growth. The development of the micro-layer seems to be originated by the difference between the velocity of the bubble expansion and the contact line velocity. Therefore, we can speculate that the micro-layer

creation results from a balance between the strength of the bubble growth rate, the capillary actions and the viscous dissipation acting on the contact line in the liquid region. The growth rate of the bubble can be related to the Ja number (defined by Eq. (9)), which is expected to be a driving parameter. The viscous dissipative force in the liquid rim is related to both the liquid viscosity μ_l and the liquid velocity, also related to the Ja number. Therefore, we expect a higher liquid viscosity to increase the dissipation and to facilitate the micro-layer formation. The capillary actions participate to the braking of the contact line on the one hand and against the bending of the interface, and thus the formation of the micro-layer, on the other hand. Actually, we expect the wettability of the fluid to have a major influence and this can be summarised by the spreading coefficient S which, for partial wetting conditions, is given by:

$$S = \sigma(\cos \theta - 1) \quad (10)$$

The spreading coefficient is negative, and tends towards zero for complete wetting [13]. Higher σ and θ characterise a lower wettability and thus are expected to prevent the micro-layer formation. On the contrary, lower σ and θ should help the micro-layer formation. In conclusion, we expect the most significant parameters affecting the micro-layer formation to be the Ja number, the contact angle θ , the liquid viscosity μ_l and the surface tension σ .

A parametric study has been performed to fully investigate the influence of the above parameters. The selected domain is a square of dimension $L = 2$ mm, and a uniform grid with a cell size equal to $0.5 \mu\text{m}$ is used (see 3.2). The first group of simulations have been carried out considering the thermophysical properties of water at atmospheric pressure (see Table 1). The same configuration of the previous cases is adopted (Fig. 2). We have varied the Ja number from 20 to 100 (by varying the wall temperature T_w) and for each Ja number several simulations with different contact angles θ , from 25° to 90° , have been carried out. The simulated physical time for each computation is 0.04 ms, which is sufficient to observe if a micro-layer is formed or not. Results are reported in Fig. 9(a) where a circle indicates a micro-layer regime, a square a contact line regime. Intermediate situations (for instance a strongly bended interface without a film creation) have been identified with triangles. A solid curve has also been reported to identify the limit between the regions of micro-layer and contact line regimes. The results show that for a given Ja, a limit contact angle exists under which a micro-layer is formed. For higher Ja the limit contact angle increases, showing that higher Jakob numbers facilitate the micro-layer formation. In our simulations, a micro-layer is always present for $Ja > 70$ in the range of θ investigated (that is for the maximum investigated $\theta = 90^\circ$ a micro-layer regime is obtained).

To investigate the influence of μ_l and σ , the same simulations have been carried out doubling or halving μ_l and σ , keeping all the other parameters constant. Results are shown in Fig. 9(a) for the viscosity variation and in Fig. 9(c) for the surface tension variation. The limit curve obtained for water (Fig. 9(a)) has also been reported in these figures for comparison. These curves show that increasing μ_l facilitates the micro-layer regime which, for a given θ , is found for lower Ja. On the contrary, a higher σ plays against the micro-layer formation, and for a given contact angle, a micro-layer is formed for higher Ja. These results are in qualitative agreement with the suggested impact of the viscous forces and the surface tension on the micro-layer formation.

All these results can be summarised in one single curve if the proper dimensionless variables are selected. For that purpose, let us introduce a capillary number Ca for the present problem:

$$Ca = \frac{\mu_l v_{int}}{\sigma} \quad (11)$$

where v_{int} represents an indicator of the bubble growth velocity, normal to the interface. We are concerned with the first instants

⁵ Note that these are the first instants after the bubble nucleation: for longer times, Q_{int} is expected to grow with the bubble growth.

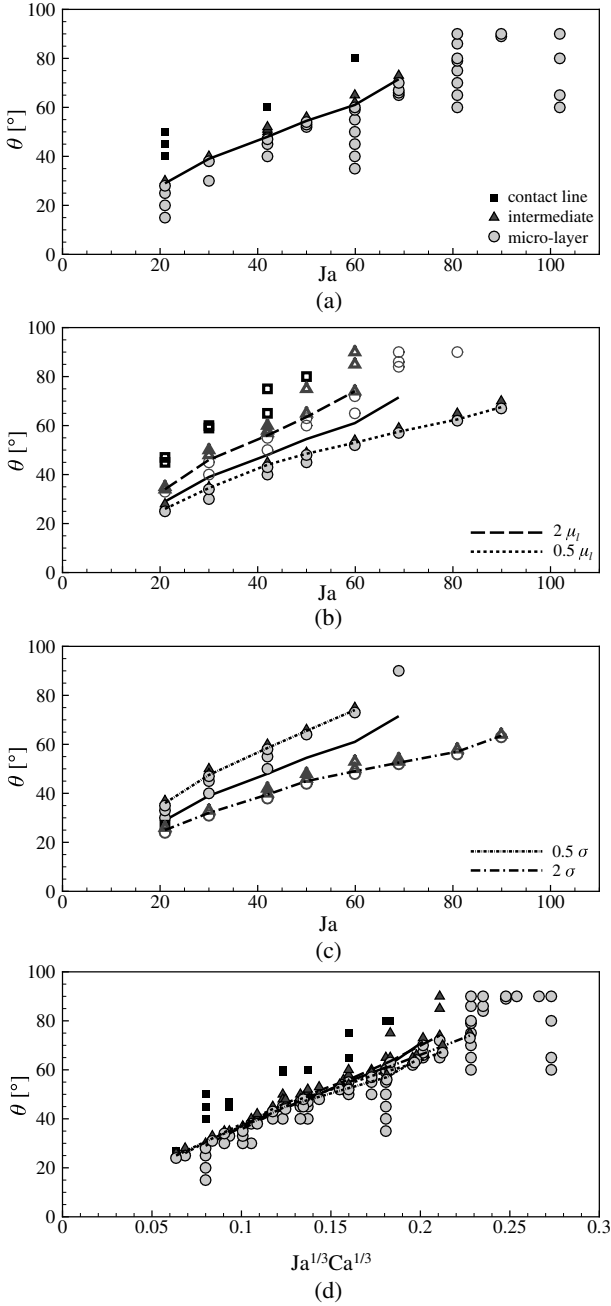


Fig. 9. Results of the parametric study on the transition from micro-layer to contact line regime. (a) Water, properties of Table 1. (b) Open symbols: $2\mu_l$, closed symbols: $0.5\mu_l$. (c) Open symbols: 2σ , closed symbols: 0.5σ . (d) All the results.

after nucleation, when a global estimation of v_{int} is given by the expansion velocity of the bubble induced by the phase change:

$$v_{int} \approx \frac{q_w}{\rho_v L_{vap}} \quad (12)$$

In Eq. (12) the evaporation mass flow rate \dot{m} has been expressed as a function of the wall heat flux q_w which can be related to the thickness δ_{KC} of the thermal boundary layer that has been initially imposed (see Eq. (8)):

$$q_w = k_l \frac{(T_w - T_{sat})}{\delta_{KC}} \quad (13)$$

Making use of Eqs. (12), (13) and (9), after some rearrangements, Eq. (11) gives:

$$Ca = Ja \frac{\alpha_l \mu_l}{\delta_{KC} \sigma} \quad (14)$$

where α_l is the liquid thermal diffusivity. Being Ja the other dimensionless number considered, we are looking for a proper combination $Ja^a Ca^b$ permitting to summarise the results. The best fit is found for $a = b = \frac{1}{3}$ as shown in Fig. 9(d), where all the numerical results of the present parametric study have been reported, and which corresponds to the following condition to have a micro-layer regime:

$$\frac{Ja Ca}{(\theta - \theta_0)^3} > \frac{1}{A^3} \quad (15)$$

where $\theta_0 = 5^\circ$ and $A = 313$, valid in the range of the interpolated results that is for $2.2 \times 10^{-4} < Ca Ja < 1.2 \times 10^{-2}$ and $15^\circ < \theta < 90^\circ$. Note that the correlation Eq. (15) has been deduced in wall isothermal conditions which correspond to the limit case of Biot number $Bo \rightarrow 0$. In terms of dimensional quantities, Eq. (15) corresponds to the following limit for the interface velocity v_{int} :

$$v_{int} > \frac{1}{A^{3/2}} \sqrt{\frac{\alpha_l \sigma}{\delta_{KC} \mu_l}} (\theta - \theta_0)^3 \quad (16)$$

The square root term in Eq. (16) can be interpreted as the product of two characteristic velocities:

- a capillary velocity v_d^* given by the following expression:

$$v_d^* = \frac{\sigma}{\mu_l} (\theta - \theta_0)^3 \quad (17)$$

which is the characteristic velocity related to the dewetting of a film when a hole is nucleated, and depends on a balance between the capillary forces and the dissipative viscous forces in the dewetting rim⁶[11,39,47];

- a velocity associated with the phase change v_m^* :

$$v_m^* = \frac{\alpha_l}{\delta_{KC}} \quad (19)$$

which is the velocity of heat diffusion through the thermal boundary layer.

Eq. (16) suggests that a micro-layer is formed beneath a nucleated bubble if the bubble growth is faster than the maximal velocity that the contact line can have, which depends on the thermophysical properties of the fluid, and corresponds to a combination of the dewetting velocity and the velocity locally induced by the phase change at the contact line.

6. Conclusion

The first DNS of a nucleated bubble showing the formation and depletion of a micro-layer has been carried out and analysed in the present work. The shape of the micro-layer, showing a rim followed by a film, is similar to a dewetting film. This has suggested a mechanism for the micro-layer depletion induced by capillary actions. It has been shown that fine grid resolutions, under the half μm , are necessary to capture the micro-layer, that only measures a few microns height. However, the presence of a micro-layer strongly influences both the shape of the bubble and the wall heat flux at larger scales. Therefore, knowing what is the nucleate boil-

⁶ Usually the velocity of the dewetting film is given by:

$$v_d = k v_d^* \quad (18)$$

where k is a constant depending on the molecular weight and the nature of the fluid [39].

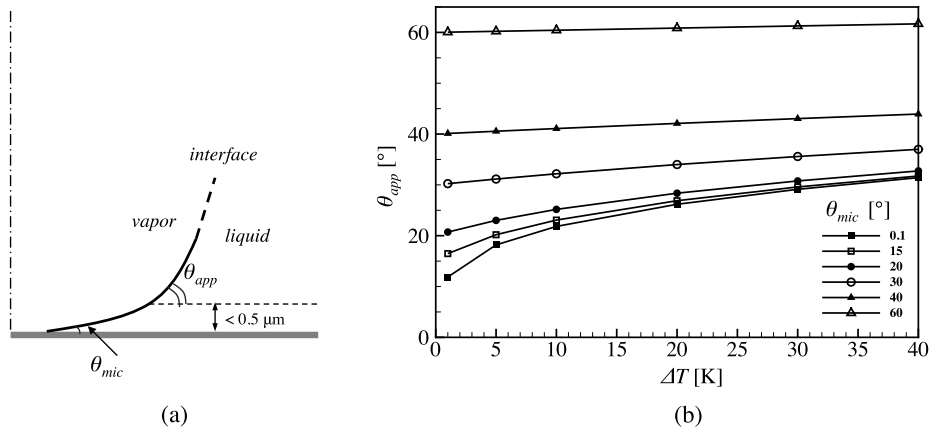


Fig. A.10. (a) Schematic of the micro-region near the contact line of a bubble. (b) Evolution of the apparent contact angle θ_{app} with the wall superheating ΔT for different microscopic contact angles θ_{mic} , for water at atmospheric pressure.

ing regime is of major interest to correctly predict the wall heat flux profile, either by fully solving the micro-layer or by using an integrated model for the micro-layer. A parametric study has thus been carried out to investigate on the transition between the two regimes. This has permit to find a correlation for the limit in terms of Ja number, Ca number and apparent contact angle θ . The correlation could help in knowing a priori the nucleate boiling regime and thus in improving the heat flux profile prediction. Moreover, the results have shown that the formation of the micro-layer occurs if the bubble grows faster than a limit velocity for the contact line which depends on the dewetting velocity, related to viscous and capillary forces, and on the liquid thermal diffusion through the thermal boundary layer.

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Appendix A. Contact angles

The static microscopic contact angle θ_{mic} formed between a liquid-vapour interface and a solid wall results from an energetic balance at the contact line, and depends on the nature of both the fluid and the solid [46]. Wettable fluids are characterised by static contact angles θ_{mic} close to zero. In the framework of bubble nucleated on a superheated wall a micro-region exists, having a characteristic dimension of the order of 100 nm, where the interface is bended. In the micro-region the interface slope varies between θ_{mic} at the wall and the apparent contact angle θ_{app} , as schematically shown in Fig. A.10(a) [27,38]. Micro-region models allow to compute θ_{app} as a function of the wall superheat $\Delta T = T_w - T_{sat}$ for a given fluid with a given θ_{mic} . As an example, in Fig. A.10(b) are reported the evolutions of θ_{app} obtained with

Mathieu model [33] which is a generalisation of Stephan model [49], and that has been analysed in the framework of nucleate boiling in contact line regime by Huber et al. [19]. In particular, these curves show that high values of θ_{mic} involve small variations of θ_{app} .

In the present simulations, we impose the angle between the interface and the wall with a boundary condition on the level set function (see [19] for the numerical details). The grid cell size is equal to $0.5 \mu m$ and therefore it does not allow to capture the micro-region that should be smaller than $0.1 \mu m$. As a consequence, the imposed contact angle θ must be considered an apparent contact angle. Moreover, note that the conditions analysed in the present work refer to Ca numbers, based on the velocity of the contact line, around 0.02. As a consequence of these small values, we have neglected the variation of the contact angle induced by dynamic effects [6,20,46].

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