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The Baker-Campbell-Hausdorff formula via mould calculus

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Abstract

The well-known Baker-Campbell-Hausdorff theorem in Lie theory says that the logarithm of a noncommutative product $e^X e^Y$ can be expressed in terms of iterated commutators of X and Y . This paper provides a gentle introduction to Écalle's mould calculus and shows how it allows for a short proof of the above result, together with the classical Dynkin explicit formula [Dy47] for the logarithm, as well as another formula recently obtained by T. Kimura [Ki17] for the product of exponentials itself. We also analyse the relation between the two formulas and indicate their mould calculus generalization to a product of more exponentials.

Contents

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1 Introduction

Let $\mathscr A$ be a noncommutative associative algebra with unit. In the associative algebra $\mathscr{A}[t]$ of all power series in an indeterminate t with coefficients in \mathscr{A} , one can take the exponential of any series without constant term in t and the logarithm of any series with constant term $1_{\mathscr{A}}$. In this context, the famous Baker-Campbell-Hausdorff theorem (BCH theorem, for short) can be phrased as

$$
\log(e^{tX}e^{tY}) \in \text{Lie}(X, Y)[[t]] \text{ for any } X, Y \in \mathcal{A}, \tag{1}
$$

where $Lie(X, Y)$ is the Lie subalgebra of $\mathscr A$ generated by X and Y, *i.e.* the smallest subspace which contains X and Y and is stable under commutator (see e.g. [BF12] and references therein).

In fact, using the notation [A, B] or ad_A B for a commutator $AB - BA$, one has

$$
\log(e^{tX}e^{tY}) = t(X+Y) + \frac{t^2}{2}[X,Y] + \frac{t^3}{12}([X,[X,Y]] + [Y,[Y,X]]) - \frac{t^4}{24}[Y,[X,[X,Y]]] + \cdots,
$$

where the coefficient of each power of t can be written in terms of nested commutators involving X and Y only, and there is a remarkable explicit formula due to Dynkin $[Dy47]$:

$$
\log(e^{X}e^{Y}) = \sum \frac{(-1)^{k-1}}{k} \frac{t^{\sigma}}{\sigma} \frac{[X^{p_1}Y^{q_1} \cdots X^{p_k}Y^{q_k}]}{p_1!q_1! \cdots p_k!q_k!}
$$
 (2)

with summation over all $k \in \mathbb{N}^*$ and $(p_1, q_1), \cdots, (p_k, q_k) \in \mathbb{N} \times \mathbb{N} \setminus \{(0, 0)\}\)$, where $\sigma \coloneqq$ $p_1 + q_1 + \cdots + p_k + q_k$ and $[X^{p_1}Y^{q_1} \cdots X^{p_k}Y^{q_k}] := \text{ad}_X^{p_1} \text{ad}_Y^{q_1} \cdots \text{ad}_X^{p_k} \text{ad}_Y^{q_k-1} Y$ if $q_k \ge 1$ and $\mathrm{ad}_X^{p_1} \mathrm{ad}_Y^{q_1} \cdots \mathrm{ad}_X^{p_k-1} X$ if $q_k = 0$ (in which case $p_k \geq 1$). Of course, the contribution of the terms with $q_k \geq 2$, or with $p_k \geq 2$ and $q_k = 0$, is zero.

Our aim is to revisit the BCH theorem and the Dynkin formula in the light of Ecalle's ´ so-called "mould calculus". We will show how mould calculus allows one to prove these results with little effort, as well as an interesting formula which was recently obtained by T. Kimura [Ki17] in relation to the BCH theorem and the Zassenhaus formula and reads

$$
e^{tX}e^{tY} = 1_{\mathscr{A}} + \sum_{r=1}^{\infty} \sum_{n_1, \dots, n_r=1}^{\infty} \frac{1}{n_r(n_r + n_{r-1}) \cdots (n_r + \cdots + n_1)} D_{n_1} \cdots D_{n_r}
$$

with $D_n := \frac{t^n}{(n-1)!} ad_X^{n-1}(X+Y)$ for each $n \ge 1$. (3)

We will also show how formula (3) and a little knowledge of mould calculus immediately imply the BCH theorem, and how the results can be generalized to a product of more than two exponentials. It seems hard to prove all these facts using the methods of [Ki17], which rely on a lot of explicit combinatorial computations, whereas almost no computation is needed when using a tiny part of mould machinery. In a nutshell, the point is that the rational coefficients in (3) make up a "symmetral mould"—in fact, a very classical one in

mould calculus—and that Dynkin's formula (2) is in essence a typical "Lie mould expansion" involving an "alternal" mould; we will explain in due time what "mould expansions", "symmetrality" and "alternality" are and how they relate to the Lie theory. We will also define a new operation in mould calculus, which gives the relation between the rational coefficients appearing in formulas (2) and (3).

Mould calculus was set up by J. Ecalle in the 1980s as part of his resurgence theory ($[Ec81]$, $[Ec92]$). Originally, Ecalle developed resurgence theory as a tool to study analytic classification problems within dynamical system theory, first for one-dimensional holomorphic germs, and then for much larger classes of discrete dynamical systems or vector fields, allowing him to tackle the Dulac conjecture about the finiteness of limit cycles of planar analytic vector fields. It soon turned out that resurgence theory has its own merits not only in mathematics but also in physics. For example, quantum resurgence was developed by Ecalle himself ($[Ec84]$) and Voros ($[Vo83]$) to study the spectrum of Schrödinger operators, and it was continued by Pham and his collaborators $(e.g.$ [DDP93]) as an essential aspect of exact WKB analysis. The mathematical side of resurgence theory has evolved steadily ([Sa16]). Recently, resurgence theory has been at the forefront in such diverse topics in mathematical physics as BPS spectrum ([GMN13]), supersymmetric field theories ([BD16] and references therein), resurgence and quantization as Riemann-Hilbert correspondence ([Ko17]), topological strings and Gromov-Witten theory ([CMS17], [CSV17]), to name a few.

Resurgence theory deals with analytic functions which enjoy a certain property of analytic continuation ("endlessly continuable functions"), which form an algebra, and which typically appear as Borel transforms of certain divergent series. In his systematic study of the singularities of these functions, their monodromies and Stokes data, Ecalle discovered ´ an infinite family of derivations acting on them, which generate a free Lie algebra. Mould calculus first appeared as a convenient combinatorial tool to manipulate these derivations. Later on, Ecalle also used mould calculus to study formal classification problems in dynamical system theory, without any relation to resurgence theory. Mould calculus has since been used in various branches of mathematics, for example in the theory of multiple zeta values ([Ec03], [Sc12], [BE17], [BS17]), in conjugacy problems for formal or analytic differential equations [Me09], [Sa09], in combinatorial Hopf algebras related to symmetric functions [Th11], in conjugacy problems in Lie algebras motivated by classical and quantum dynamics [PS17], in the study of Rayleigh-Schrödinger series [NP18].

In the present paper, we do not assume any familiarity with mould calculus on the part of the reader, and we introduce the most basic ideas about moulds. The BCH formula can be seen as an application, and we hope that readers can find other interesting applications in mathematics or physics.

The paper is organized as follows.

- Section 2 is a gentle introduction to mould calculus, containing the basic definitions and properties that we will require in our applications.
- Section 3 gives short proofs of the BCH theorem (Theorem A) and Dynkin's formula (Theorem B) based on mould calculus.
- Section 4 gives a short proof of Kimura's formula (Theorem C) via mould calculus, as well as another derivation of the BCH theorem (Corollary 4.5).
- Section 5 indicates how to generalize the previous results to the case of a product of more factors $e^{tX_1} \cdots e^{tX_N}$, with arbitrary $N \geq 2$ (Theorems B' and C').
- Section 6 defines a new operation in mould calculus, that we call σ -composition, which allows us to relate the mould used for Dynkin's formula and the one used for Kimura's formula.

2 Mould calculus for pedestrians

Throughout the article we use the notation

$$
\mathbb{N} = \{0, 1, 2, \ldots\}, \quad \mathbb{N}^* = \{1, 2, 3, \ldots\}.
$$

In this section, we denote by k a field of characteristic zero (it will be \mathbb{O} in our later applications) and by $\mathcal N$ a nonempty set (in our applications, it will be either a finite set or \mathbb{N}^*).

2.1 The mould algebra

Viewing $\mathcal N$ as an alphabet (the elements of which we call "letters"), we denote by $\mathcal N$ the corresponding set of "words" (or "strings"):

$$
\underline{\mathcal{N}} \coloneqq \{ \underline{n} = n_1 \cdots n_r \mid r \in \mathbb{N}, n_1, \ldots, n_r \in \mathcal{N} \}.
$$

The concatenation law $(a_1 \cdots a_r, b_1 \cdots b_s) \in \underline{\mathcal{N}} \times \underline{\mathcal{N}} \mapsto a_1 \cdots a_r b_1 \cdots b_s \in \underline{\mathcal{N}}$ yields a monoid structure, with the empty word \varnothing as unit.

Definition 2.1. A k-valued mould on \mathcal{N} is a function $\underline{\mathcal{N}} \to \mathbf{k}$. The set of all moulds is denoted by $k^{\underline{\mathcal{N}}}$.

Given a mould M, it is customary to denote by $M^{\underline{n}}$ the value it takes on a word \underline{n} . Mould multiplication is defined by the formula

$$
(M \times N)^{\underline{n}} := \sum_{(\underline{a}, \underline{b}) \text{ such that } \underline{n} = \underline{a}\underline{b}} M^{\underline{a}} N^{\underline{b}} \quad \text{for } \underline{n} \in \underline{\mathcal{N}},
$$
 (4)

for any two moulds $M, N \in \mathbf{k}^{\underline{N}}$. For instance,

$$
(M \times N)^{n_1 n_2} = M^{\emptyset} N^{n_1 n_2} + M^{n_1} N^{n_2} + M^{n_1 n_2} N^{\emptyset}.
$$

It is immediate to check that $k^{\mathcal{N}}$ is an associative k-algebra, noncommutative if N has more than one element, whose unit is the mould 1 defined by $\mathbb{1}^{\emptyset} = 1$ and $\mathbb{1}^{\mathbb{Z}} = 0$ for $n \neq \varnothing$.

We say that a mould M has order $\geq p$ if $M^{\underline{n}} = 0$ for each word \underline{n} of length $\lt p$. Clearly, if ord $M \geq p$ and ord $N \geq q$, then ord $(M \times N) \geq p + q$. In particular, if $M^{\emptyset} = 0$, then ord $M^{\times k} \geq k$ for each $k \in \mathbb{N}^*$, hence the moulds

$$
e^M := \sum_{k \in \mathbb{N}} \frac{1}{k!} M^{\times k}
$$
 and $\log(1 + M) := \sum_{k \in \mathbb{N}^*} \frac{(-1)^{k-1}}{k} M^{\times k}$ (5)

are well-defined (because, for each $\underline{n} \in \underline{\mathcal{N}}$, only finitely many terms contribute to $(e^M)^{\underline{n}}$ or $(\log(1 + M))^n$). We thus get mutually inverse bijections

$$
\{ M \in \mathbf{k}^{\underline{\mathcal{N}}} \mid M^{\varnothing} = 0 \} \quad \underset{\log}{\overset{\text{exp}}{\rightleftharpoons}} \quad \{ M \in \mathbf{k}^{\underline{\mathcal{N}}} \mid M^{\varnothing} = 1 \}.
$$

2.2 Comoulds and mould expansions

Moulds are meant to provide the coefficients of certain multi-indexed expansions in an associative algebra A . To deal with infinite expansions, we require this A to be a complete filtered associative algebra, *i.e.* there is an order function ord: $\mathcal{A} \to \mathbb{N} \cup {\infty}$ compatible with sum and product,¹ such that every family $(A_i)_{i\in I}$ of A is formally summable provided, for each $p \in \mathbb{N}$, all the A_i 's have order $\geq p$ except finitely many of them. See [Sa09] or [PS17] for the details. For the present paper, the reader may think of

$$
\mathcal{A} = \mathscr{A}[[t]]
$$

with the order function relative to powers of t, where $\mathscr A$ is an associative algebra as in the introduction.

Assumption 2.2. We suppose that we are given a family $(B_n)_{n \in \mathcal{N}}$ in A such that all the B_n 's have order ≥ 1 and, for each $p \in \mathbb{N}$, only finitely many of them are not of order $\geq p$.

Definition 2.3. We call associative comould generated by $(B_n)_{n \in \mathcal{N}}$ the family $(B_n)_{n \in \mathcal{N}}$ defined by $B_{\emptyset} := 1_A$ and

 $B_{n_1\cdots n_r} \coloneqq B_{n_1} \cdots B_{n_r}$ for all $r \ge 1$ and $n_1, \ldots, n_r \in \mathcal{N}$.

¹We assume ord $(A + B) \ge \min\{\text{ord }A, \text{ord }B\}$ and ord $(AB) \ge \text{ord }A + \text{ord }B$ for any $A, B \in \mathcal{A}$, and ord $A = \infty$ iff $A = 0$.

Lemma 2.4. The formula

$$
M \in \mathbf{k}^{\underline{\mathcal{N}}} \mapsto MB \coloneqq \sum_{\underline{n} \in \underline{\mathcal{N}}} M^{\underline{n}} B_{\underline{n}} \in \mathcal{A} \tag{6}
$$

defines a morphism of associative algebras. Moreover,

$$
M^{\varnothing} = 0 \Rightarrow (e^M)B = e^{MB}, \qquad M^{\varnothing} = 1 \Rightarrow (\log M)B = \log(MB). \tag{7}
$$

Proof. Observe that the family $(M^{\underline{n}}B_{\underline{n}})_{\underline{n}\in\mathcal{N}}$ is formally summable in A thanks to our assumption on the B_n 's. The property $B_{\underline{a}\underline{b}} = B_{\underline{a}}B_{\underline{b}}$ for all $\underline{a}, \underline{b} \in \underline{\mathcal{N}}$ entails

$$
(M \times N)B = (MB)(NB),\tag{8}
$$

whence $M^{\times k}B = (MB)^k$ for all $k \in \mathbb{N}$, and (7) follows.

It is the right-hand side in (6) that is called a mould expansion.

Example 2.5. Suppose we are given $X, Y \in \mathcal{A}$, an associative algebra. Take $k = \mathbb{Q}$, $\mathcal{N} = \Omega := \{x, y\}$, a two-letter alphabet, and $\mathcal{A} = \mathscr{A}[[t]]$. We then consider the associative comould generated by

$$
B_x := tX, \qquad B_y := tY. \tag{9}
$$

Trivially, $tX = I_xB$ and $tY = I_yB$, where $I_x, I_y \in \mathbb{Q}^{\Omega}$ are defined by

$$
I_x^{\underline{\omega}} := \begin{cases} 1 & \text{if } \underline{\omega} \text{ is the one-letter word } x \\ 0 & \text{else,} \end{cases} \qquad I_y^{\underline{\omega}} := \begin{cases} 1 & \text{if } \underline{\omega} \text{ is the one-letter word } y \\ 0 & \text{else.} \end{cases}
$$

We thus get $e^{tX} = e^{I_x}B$, $e^{tY} = e^{I_y}B$, and

$$
e^{tX}e^{tY} = S_{\Omega}B
$$
 with $S_{\Omega} := e^{I_x} \times e^{I_y}$, $\log(e^{tX}e^{tY}) = T_{\Omega}B$ with $T_{\Omega} := \log S_{\Omega}$. (10)

By (4) and (5) , we get

$$
S_{\overline{\Omega}}^{\underline{\omega}} = \begin{cases} \frac{1}{p!q!} & \text{if } \underline{\omega} \text{ is of the form } x^p y^q \text{ with } p, q \in \mathbb{N} \\ 0 & \text{else,} \end{cases}
$$
(11)

thus the first part of (10) is just another way of writing $e^{tX}e^{tY} = \sum \frac{t^{p+q}}{p!q!}X^pY^q$.

In the general case, retaining from the associative algebra structure of A only the underlying Lie algebra structure, *i.e.* using only commutators (with the notation $ad_A B =$ $[A, B]$, one can define another kind of mould expansion:

Definition 2.6. We call Lie comould generated by $(B_n)_{n \in \mathcal{N}}$ the family $(B_{[\,\underline{n}\,]})_{\underline{n} \in \underline{\mathcal{N}}}$ of A defined by $B_{[\varnothing]} \coloneqq 0$ and

$$
B_{[n_1\cdots n_r]} := \mathrm{ad}_{B_{n_1}}\cdots \mathrm{ad}_{B_{n_{r-1}}}B_{n_r} = [B_{n_1},[\cdots [B_{n_{r-1}},B_{n_r}]\cdots]].
$$

 \Box

We define the Lie mould expansion associated with a mould $M \in \mathbf{k}^{\underline{N}}$ by the formula

$$
M[B] := \sum_{\underline{n} \in \underline{\mathcal{N}} \backslash \{\varnothing\}} \frac{1}{r(\underline{n})} M^{\underline{n}} B_{[\underline{n}]} \in \mathcal{A},\tag{12}
$$

where $r(n)$ denotes the length of a word n.

Division by $r(\underline{n})$ is just a normalization choice whose convenience will appear in Section 2.3. In Section 3, we will prove the BCH theorem by showing how to pass from the second part of (10) to a Lie mould expansion.

2.3 Symmetrality and alternality

One can get a morphism property for Lie mould expansions analogous to (8) by imposing restrictions to the moulds that we use: they must be "alternal". A tightly related notion is that of "symmetral" mould. The definition of both notions relies on word shuffling.

Recall that the shuffling of two words $\underline{a} = \omega_1 \cdots \omega_\ell$ and $\underline{b} = \omega_{\ell+1} \cdots \omega_r$ is the set of all the words \underline{n} which can be obtained by interdigitating the letters of \underline{a} and those of \underline{b} while preserving their internal order in \underline{a} or \underline{b} , *i.e.* the words which can be written $n = \omega_{\tau(1)} \cdots \omega_{\tau(r)}$ with a permutation τ such that $\tau^{-1}(1) < \cdots < \tau^{-1}(\ell)$ and $\tau^{-1}(\ell +$ 1) $\langle \cdots \langle \tau^{-1}(r) \rangle$. We define the *shuffling coefficient* sh $\left(\frac{a,b}{n}\right)$ to be the number of such permutations τ , and we set $\operatorname{sh}(\frac{a,b}{n}) \coloneqq 0$ whenever \underline{n} does not belong to the shuffling of \underline{a} and \underline{b} . For instance, if n, m, p, q are four distinct elements of N ,

$$
\operatorname{sh}\left(\frac{nmp, mq}{nmqpm}\right) = 0, \qquad \operatorname{sh}\left(\frac{nmp, mq}{m nqmp}\right) = 1, \qquad \operatorname{sh}\left(\frac{nmp, mq}{n mmqp}\right) = 2.
$$

We also define, for arbitrary words \underline{n} and \underline{a} , $\mathrm{sh}\left(\frac{a, \varnothing}{n}\right) = \mathrm{sh}\left(\frac{\varnothing, \underline{a}}{n}\right) = 1$ if $\underline{a} = \underline{n}$, 0 else.

Definition 2.7. A mould $M \in \mathbf{k}^{\underline{N}}$ is said to be *alternal* if $M^{\emptyset} = 0$ and

$$
\sum_{\underline{n}\in\underline{\mathcal{N}}} \text{sh}\left(\frac{\underline{a}, \underline{b}}{\underline{n}}\right) M^{\underline{n}} = 0 \quad \text{for any two nonempty words } \underline{a}, \underline{b}.
$$
 (13)

A mould $M \in \mathbf{k}^{\mathcal{N}}$ is said to be *symmetral* if $M^{\emptyset} = 1$ and

$$
\sum_{\underline{n}\in\underline{\mathcal{N}}} \text{sh}\left(\frac{\underline{a}}{\underline{n}}, \frac{\underline{b}}{\underline{n}}\right) M^{\underline{n}} = M^{\underline{a}} M^{\underline{b}} \quad \text{for any two words } \underline{a}, \underline{b}.
$$
 (14)

Example 2.8. It is obvious that any mould M whose support is contained in the set of one-letter words (*i.e.* $r(\underline{n}) \neq 1 \Rightarrow M^{\underline{n}} = 0$) is alternal. For instance, the moulds I_x and I_y of Example 2.5 are alternal. An elementary example of symmetral mould is E defined by

²Indeed, $\tau^{-1}(i)$ is the position in <u>n</u> of ω_i , the *i*-th letter of <u>ab</u>.

 $E^{\underline{n}} := \frac{1}{r(\underline{n})!}$. Indeed, since the total number of words obtained by shuffling of any $\underline{a}, \underline{b} \in \underline{\mathcal{N}}$ (counted with multiplicity) is $\binom{r(\underline{a}\underline{b})}{r(\underline{a})}$,

$$
\sum_{n \in \mathcal{N}} \operatorname{sh}\left(\frac{a}{n}, \frac{b}{n}\right) E^{n} = \frac{r(\underline{a}\underline{b})!}{r(\underline{a})!r(\underline{b})!} \cdot \frac{1}{r(\underline{a}\underline{b})!} = E^{\underline{a}} E^{\underline{b}}.
$$

We shall see later that the moulds e^{I_x} , e^{I_y} and S_{Ω} involved in (10) are symmetral, and that T_{Ω} is alternal.

In this paper,³ we are interested in the shuffling coefficients because of the following classical relation between the Lie comould and the associative comould:

$$
B_{[\underline{n}]} = \sum_{(\underline{a}, \underline{b}) \in \underline{\mathcal{N}} \times \underline{\mathcal{N}}} (-1)^{r(\underline{b})} r(\underline{a}) \text{ sh}\left(\frac{a}{n}, \frac{b}{n}\right) B_{\underline{\tilde{b}}\underline{a}} \quad \text{for all } \underline{n} \in \underline{\mathcal{N}},\tag{15}
$$

where, for an arbitrary word $\underline{b} = b_1 \cdots b_s$, we denote by $\underline{\widetilde{b}}$ the reversed word: $\underline{\widetilde{b}} = b_s \cdots b_1$ (we omit the proof—see [vW66], [Re93], [PS17]). An immediate and useful consequence is

Lemma 2.9. If M is an alternal mould, then $M[B] = MB$, i.e.

$$
\sum_{n \in \mathcal{N}\backslash\{\varnothing\}} \frac{1}{r(n)} M^n B_{[n]} = \sum_{n \in \mathcal{N}} M^n B_n.
$$

Proof. Putting together (12) and (15), we get $M[B] = \sum$ n≠∅ \sum $\underline{a},\underline{b}$ $(-1)^{r(\underline{b})}\frac{r(\underline{a})}{r(\underline{n})}\,\operatorname{sh}\left(\frac{a,b}{\underline{n}}\right)M^{\underline{n}}B_{\underline{\widetilde{b}}\,\underline{a}}.$ Now, $\operatorname{sh}(\frac{a,b}{n}) \neq 0 \Rightarrow r(\underline{n}) = r(\underline{a}) + r(\underline{b})$, hence

$$
M[B] = \sum_{r(\underline{a})+r(\underline{b})\geq 1} (-1)^{r(\underline{b})} \frac{r(\underline{a})}{r(\underline{a})+r(\underline{b})} \bigg(\sum_{\underline{n}\in \underline{\mathcal{N}}} \text{sh}\left(\frac{a}{n}\underline{b}\right) M^{\underline{n}}\bigg) B_{\underline{\tilde{b}}\underline{a}} = \sum_{\underline{a}\neq \varnothing} M^{\underline{a}} B_{\underline{a}} = MB
$$

(the internal sum is $M^{\underline{a}}$ when $\underline{b} = \emptyset$ and it does not contribute when \underline{a} or $\underline{b} \neq \emptyset$ because of (13), nor when $\underline{a} = \emptyset$ because of the factor $r(\underline{a})$. \Box

Any mould expansion associated with an alternal mould thus belongs to the (closure of the) Lie subalgebra of A generated by the B_n 's, since it can be rewritten as a Lie mould expansion, involving only commutators of the B_n 's.

Lemma 2.9 is related to the classical Dynkin-Specht-Wever projection lemma in the context of free Lie algebras (see $e.g.$ [Re93]). One should also mention that the concepts

 3 In Écalle's work, the initial motivation for the definition of alternality and symmetrality is the situation when A is an algebra of operators (acting on an auxiliary algebra) and each B_n acts as a derivation: in that case, the B_{n} satisfy a modified Leibniz rule which involves the shuffling coefficients, whence it follows that MB is itself a derivation if M is an alternal mould, and an algebra automorphism if M is symmetral. Here we do not assume anything of that kind on A and the B_n 's but rather follow the spirit of "Lie mould calculus" as advocated in [PS17].

of symmetrality and alternality are related to certain combinatorial Hopf algebras, as emphasized by F. Menous in his work on the renormalization theory in perturbative quantum field theory—see e.g. [Me09] and footnote 4. Hopf-algebraic aspects of mould calculus are also touched upon in [Sa09], [PS17] and [NP18].

For our applications, we require a last general result from mould calculus (see e.g. [Sa09] for a proof):

Lemma 2.10.

- The product of two symmetral moulds is symmetral.
- The logarithm of a symmetral mould is alternal.
- The exponential of an alternal mould is symmetral.

Example 2.11. The mould I defined by

$$
I^{\underline{n}} = \begin{cases} 1 & \text{if } r(\underline{n}) = 1 \\ 0 & \text{else,} \end{cases}
$$
 (16)

is alternal (being supported in one-letter words). The symmetral mould E of Example 2.8 is e^I .

In fact, the set of all symmetral moulds is a group for mould multiplication, the set of all alternal moulds is a Lie algebra for mould commutator, and we get the analogue of (8) for Lie mould expansions:

$$
M, N \text{ alternal} \Rightarrow [M, N][B] = [M[B], N[B]].
$$

Let us also mention a manifestation of the antipode of the Hopf algebra related to moulds:⁴

M alternal $\Rightarrow S(M) = -M$, M symmetral $\Rightarrow S(M) =$ multiplicative inverse of M,

where $S(M)^{n_1\cdots n_r} \coloneqq (-1)^r M^{n_r\cdots n_1}$.

All these facts are mentioned in Ecalle's works and can be proved by Hopf-algebraic ´ techniques or by direct computation.

⁴ Denote by kN the linear span of the set of words, *i.e.* the k-vector space consisting of all formal sums $c = \sum c_n \underline{n}$ with finitely many nonzero coefficients $c_n \in \mathbf{k}$. Now, $\mathbf{k} \underline{\mathcal{N}}$ is a Hopf algebra if we define multiplication by extending $(\underline{a}, \underline{b}) \mapsto \underline{a} \Delta \underline{b} := \sum \text{sh}(\frac{a, b}{n}) \underline{n}$ by bilinearity, comultiplication by extending $\underline{n} \mapsto \sum_{n=a} \underline{b} \otimes \underline{b}$ by linearity, and antipode by extending $n_1 \cdots n_r \mapsto (-1)^r n_r \cdots n_1$ by linearity (the unit is \varnothing and the counit is $c \mapsto c_{\varnothing}$. The set of moulds can be identified with the set of linear forms on $k \underline{\mathcal{N}}$, if we identify $M \in \mathbf{k}^{\underline{N}}$ with $c \mapsto \sum M^n c_n$. The associative algebra structure (4) of $\mathbf{k}^{\underline{N}}$ is then dual to the coalgebra structure of kN , and alternal moulds appear as infinitesimal characters of kN (linear forms M such that $M(c\Delta c') = M(c)c'_{\mathcal{O}} + c_{\mathcal{O}}M(c')$ and symmetral moulds as characters (linear forms M such that $M(\emptyset) = 1$ and $M(c \Delta c') = M(c)M(c')$.

3 The BCH Theorem and Dynkin's formula

Let $\mathscr A$ be an associative algebra. We now use mould calculus to prove

Theorem A. Suppose $X, Y \in \mathscr{A}$. Let $\Psi = e^{tX} e^{tY} \in \mathscr{A}[[t]]$. Then

 $\log \Psi \in \mathrm{Lie}(X,Y)[[t]],$

where $Lie(X, Y)$ is the Lie subalgebra of $\mathscr A$ generated by X and Y.

Theorem B (Dynkin, [Dy47]). In the above situation,

$$
\log \Psi = \sum \frac{(-1)^{k-1}}{k} \frac{t^{\sigma}}{\sigma} \frac{[X^{p_1} Y^{q_1} \cdots X^{p_k} Y^{q_k}]}{p_1! q_1! \cdots p_k! q_k!} \tag{17}
$$

with summation over all $k \in \mathbb{N}^*$ and $(p_1, q_1), \cdots, (p_k, q_k) \in \mathbb{N} \times \mathbb{N} \setminus \{(0, 0)\},\$ where $\sigma :=$ $p_1 + q_1 + \cdots + p_k + q_k$ and $[X^{p_1}Y^{q_1} \cdots X^{p_k}Y^{q_k}] := \text{ad}_X^{p_1} \text{ad}_Y^{q_1} \cdots \text{ad}_X^{p_k} \text{ad}_Y^{q_k-1} Y$ if $q_k \ge 1$ and $\text{ad}_X^{p_1} \text{ad}_Y^{q_1} \cdots \text{ad}_X^{p_k-1} X \text{ if } q_k = 0.$

Proof of Theorem A. Half of the work has already been done in Example 2.5! With the twoletter alphabet $\Omega = \{x, y\}$, $B_x = tX$ and $B_y = tY$, we have $\log \Psi = T_{\Omega}B$ with $T_{\Omega} = \log S_{\Omega}$, $S_{\Omega} = e^{I_x} \times e^{I_y}.$

The mould S_{Ω} is symmetral, because I_x and I_y are alternal (they are supported in the set of one-letter words) hence e^{I_x} and e^{I_y} are symmetral by Lemma 2.10 and so is their product. It follows, still by Lemma 2.10, that T_{Ω} is alternal. Lemma 2.9 then shows that

$$
\log \Psi = T_{\Omega}[B]. \tag{18}
$$

In particular, being expressed as a Lie mould expansion, $\log \Psi$ lies in Lie $(X, Y)[[t]]$. \Box

Proof of theorem B. With the same notation as previously, by definition,

$$
T_{\Omega}^{\underline{\omega}} = \sum_{k \ge 1} \frac{(-1)^{k-1}}{k} \sum_{\substack{\underline{\omega}^1, \dots, \underline{\omega}^k \in \underline{\Omega} \setminus \{\varnothing\} \\ \underline{\omega} = \underline{\omega}^1 \cdots \underline{\omega}^k}} S_{\Omega}^{\underline{\omega}^1} \cdots S_{\Omega}^{\underline{\omega}^k} \quad \text{for each word } \underline{\omega},
$$

hence (18) yields

$$
\log \Psi = \sum_{k \geq 1} \frac{(-1)^{k-1}}{k} \sum_{\underline{\omega}^1, \dots, \underline{\omega}^k \in \underline{\Omega} \backslash \{\varnothing\}} \frac{1}{r(\underline{\omega}^1) + \dots + r(\underline{\omega}^k)} S_{\Omega}^{\underline{\omega}^1} \dotsm S_{\Omega}^{\underline{\omega}^k} B_{[\underline{\omega}^1 \dotsm \underline{\omega}^k]}.
$$

Inserting (11) , we exactly get (17) .

Mould calculus also allows us to express the inner derivation associated with $\log \Psi$:

Corollary 3.1. The inner derivation of $\mathscr{A}[[t]]$ associated with $Z \coloneqq \log(e^{tX}e^{tY})$ is

$$
ad_Z = \sum \frac{(-1)^{k-1} t^{\sigma}}{k} \frac{ad_X^{p_1} ad_Y^{q_1} \cdots ad_X^{p_k} ad_Y^{q_k}}{p_1! q_1! \cdots p_k! q_k!} = \sum \frac{(-1)^{k-1}}{k} \frac{t^{\sigma}}{\sigma} \frac{[ad_X^{p_1} ad_Y^{q_1} \cdots ad_X^{p_k} ad_Y^{q_k}]}{p_1! q_1! \cdots p_k! q_k!}
$$
(19)

with summation over all $k \in \mathbb{N}^*$ and $(p_1, q_1), \cdots, (p_k, q_k) \in \mathbb{N} \times \mathbb{N} \setminus \{(0, 0)\},\$ where $\sigma :=$ $p_1 + q_1 + \cdots + p_k + q_k$ and with the same bracket notation as in Theorem A.

 \Box

Proof. Working in the associative algebra End $\mathscr{A}[[t]]$ with the comould and the Lie comould associated with $A_x := ad_{tX}$ and $A_y := ad_{tY}$, we get $ad_Z = T_{\Omega}[A]$ (*i.e.* the second part of (19)) from (17) because $A_{[\underline{\omega}]} = \text{ad}_{B_{[\underline{\omega}]}}$. Lemma 2.9 then entails $\text{ad}_Z = T_{\Omega}A$, *i.e.* the first part of (19) (which could have been obtained directly from $\text{ad}_Z = \log(\text{e}^{\text{ad}_t x} \text{e}^{\text{ad}_t y})$. \Box

4 Alternative formulas for $e^{tX}e^{tY}$ and its logarithm

In this section, we take $\mathcal{N} := \mathbb{N}^* = \{1, 2, 3, \ldots\}$ as our alphabet, and $\mathbf{k} := \mathbb{Q}$ as base field. We now show how to find Kimura's formula (3) from mould calculus.

4.1 An alternative mould expansion for $e^{tX}e^{tY}$

Theorem C ([Ki17]). Let $X, Y \in \mathcal{A}$ as in Theorem A. Then $\Psi = e^{tX}e^{tY}$ can be written

$$
\Psi = 1_{\mathscr{A}} + \sum_{r=1}^{\infty} \sum_{n_1, \dots, n_r=1}^{\infty} \frac{1}{n_r(n_r + n_{r-1}) \cdots (n_r + \dots + n_1)} D_{n_1} \cdots D_{n_r}
$$
(20)

with
$$
D_n := \frac{t^n}{(n-1)!} \operatorname{ad}_X^{n-1}(X+Y) \quad \text{for each } n \ge 1.
$$
 (21)

The rest of section 4.1 is devoted to a new proof of this formula.

Lemma 4.1. $\Psi = e^{tX}e^{tY}$ is the unique element of $\mathscr{A}[[t]]$ such that

$$
\Psi_{|t=0} = 1_{\mathscr{A}}, \qquad t\partial_t \Psi = D\Psi, \qquad \text{where } D := t e^{tX} (X + Y) e^{-tX}. \tag{22}
$$

Proof. The fact that Ψ satisfies (22) is straightforward. On the other hand, if $\tilde{\Psi} \in \mathscr{A}[[t]]$ is also solution to (22), then $\text{ord}(\tilde{\Psi}-\Psi) \geq 1$ and it is easy to see that in fact $\text{ord}(\tilde{\Psi}-\Psi) = \infty$ because $t\partial_t(\tilde{\Psi}-\Psi)=D(\tilde{\Psi}-\Psi)$ and ord $D\geq 1$; hence $\tilde{\Psi}-\Psi=0$. П

Let $\mathcal{N} := \mathbb{N}^*$ and consider the associative comould associated with the family $(D_n)_{n \in \mathcal{N}}$ defined by (21). We have

$$
D = \sum_{n \in \mathcal{N}} D_n = ID,\tag{23}
$$

where D in the left-hand side is the element of $\mathscr{A}[[t]]$ defined in (22), while the right-hand side is the mould expansion associated with the mould I defined by (16) . The proof of (23) is essentially the Hadamard lemma: ad_X can be written $L_X - R_X$, where L_X and R_X are the operators of left-multiplication and right-multiplication by X and they commute, hence $e^{t \cdot ad_X} = e^{t(L_X - R_X)} = e^{tL_X} e^{-tR_X}$, and e^{tL_X} and e^{-tR_X} are the operators of leftmultiplication and right-multiplication by e^{tX} and e^{-tX} , whence

$$
e^{t \operatorname{ad}_X} A = e^{tX} A e^{-tX} \quad \text{for any } A \in \mathscr{A}[[t]]. \tag{24}
$$

In particular, $e^{tX}(X+Y)e^{-tX} = \sum_{n \in \mathcal{N}} \frac{t^{n-1}}{(n-1)!} \operatorname{ad}_X^{n-1}(X+Y).$

Lemma 4.2. For any mould $S \in \mathbb{Q}^{\mathcal{N}}$,

$$
t\partial_t (SD) = (\nabla S)D,
$$

where ∇S is the mould defined by

$$
(\nabla S)^{n_1 \cdots n_r} := (n_1 + \cdots + n_r) S^{n_1 \cdots n_r} \quad \text{for each word } n_1 \cdots n_r \in \mathcal{N}.
$$

Proof. Obvious, since $D_n \in t^n \mathscr{A}$ for each $n \in \mathcal{N}$.

Lemma 2.4, formula (23) and Lemma 4.2 inspire us to look for a solution to (22) in the form of a mould expansion: $\Psi = SD$ will be solution to (22) if $S \in \mathbb{Q}^N$ is solution to the mould equation

$$
S^{\varnothing} = 1, \qquad \nabla S = I \times S \tag{25}
$$

(indeed: we have $(\nabla S)D = t\partial_t \Psi$ on the one hand, and $(I \times S)D = (ID)(SD) = D\Psi$ on the other hand, and $S^{\varnothing} = 1$ ensures $\text{ord}(\Psi - 1_{\mathscr{A}}) \geq 1$ because $\text{ord } D_n \geq 1$ for all nonempty word \underline{n}). Now the second part of (25) is equivalent to

$$
(n_1 + \dots + n_r)S^{n_1 \dots n_r} = S^{n_2 \dots n_r}
$$
 for each nonempty word $n_1 \dots n_r \in \mathcal{N}$, (26)

thus the mould equation (25) has a unique solution: the mould $S_{\mathcal{N}} \in \mathbb{Q}^{\mathcal{N}}$ defined by

$$
S_{\mathcal{N}}^{n_1\cdots n_r} \coloneqq \frac{1}{n_r(n_r + n_{r-1})\cdots(n_r + \cdots + n_1)} \quad \text{for each } n_1\cdots n_r \in \underline{\mathcal{N}}.\tag{27}
$$

In conclusion, S_N is a solution to (25), thus $S_N D$ is a solution to (22), thus

$$
S_{N}D = \Psi = e^{tX}e^{tY}
$$
\n(28)

and formula (20) is proved.

Remark 4.3. For any alphabet N and base field k, an arbitrary function $\phi: \mathcal{N} \to \mathbf{k}$ gives rise to a linear operator $\nabla_{\phi} \colon \mathbf{k}^{\underline{\mathcal{N}}} \to \mathbf{k}^{\underline{\mathcal{N}}}$ defined by the formula

$$
(\nabla_{\phi} M)^{n_1 \cdots n_r} = (\phi(n_1) + \cdots + \phi(n_r)) M^{n_1 \cdots n_r}
$$
\n(29)

(with the convention that an empty sum is 0). The reader can check that ∇_{ϕ} is a mould derivation, i.e. it satifies the Leibniz rule $\nabla_{\phi}(M \times N) = (\nabla_{\phi} M) \times N + M \times \nabla_{\phi} N$. Here, we have used the mould derivation associated with the inclusion map $\mathbb{N}^* \hookrightarrow \mathbb{Q}$.

4.2 An alternative Lie mould expansion for $log(e^{tX}e^{tY})$

The mould S_N that we have just constructed happens to be a very common and useful object of mould calculus (see e.g. [Ec81] or [Sa09, §13]). It is well-known that it is symmetral; we give the proof for the sake of completeness.

 \Box

Lemma 4.4. The mould S_N defined by the formula (27) is symmetral.

Proof. We prove the property (14) for $M = S_N$ by induction on $r(\underline{a}) + r(\underline{b})$. The property holds when $\underline{a} = \emptyset$ or $\underline{b} = \emptyset$ because $S_N^{\emptyset} = 1$. In particular it holds when $r(\underline{a}) + r(\underline{b}) = 0$. Suppose now that \underline{a} and \underline{b} are arbitrary nonempty words. Using the notation

 $|\underline{n}| := n_1 + \cdots + n_r$, $\underline{n} := n_2 \cdots n_r$ for any nonempty word $n_1 \cdots n_r$,

we multiply the right-hand side of (14) by $|\underline{a}| + |\underline{b}|$: we get

$$
(|\underline{a}|+|\underline{b}|)S_{\mathcal{N}}^{\underline{a}}S_{\mathcal{N}}^{\underline{b}} = |\underline{a}|S_{\mathcal{N}}^{\underline{a}}S_{\mathcal{N}}^{\underline{b}} + |\underline{b}|S_{\mathcal{N}}^{\underline{a}}S_{\mathcal{N}}^{\underline{b}} = S_{\mathcal{N}}^{'\underline{a}}S_{\mathcal{N}}^{\underline{b}} + S_{\mathcal{N}}^{\underline{a}}S_{\mathcal{N}}^{'\underline{b}} = \sum_{\underline{c}} \operatorname{sh}\left(\frac{a_{\underline{c}}b}{\underline{c}}\right)S_{\mathcal{N}}^{\underline{c}} + \sum_{\underline{c}} \operatorname{sh}\left(\frac{a_{\underline{c}}b}{\underline{c}}\right)S_{\mathcal{N}}^{\underline{c}},\tag{30}
$$

where we have used (26) and the induction hypothesis. On the other hand, multiplying the left-hand side of (14) by $|\underline{a}| + |\underline{b}|$, we get

$$
(|\underline{a}| + |\underline{b}|) \sum_{\underline{n}} \operatorname{sh}\left(\frac{a, \underline{b}}{\underline{n}}\right) S_{\mathcal{N}}^{\underline{n}} = \sum_{\underline{n}} |\underline{n}| \operatorname{sh}\left(\frac{a, \underline{b}}{\underline{n}}\right) S_{\mathcal{N}}^{\underline{n}} = \sum_{\underline{n}} \operatorname{sh}\left(\frac{a, \underline{b}}{\underline{n}}\right) S_{\mathcal{N}}^{\underline{n}} \tag{31}
$$

(using (26) again). The last sum can be split into two according to the first letter of n, which must come either from the first letter of <u>a</u> or from the first letter of <u>b</u> for sh $\left(\frac{a,b}{n}\right)$ to be nonzero: either $\underline{n} = a_1 \underline{c}$ and $\mathrm{sh}\left(\frac{a,b}{n}\right) = \mathrm{sh}\left(\frac{a,b}{\underline{c}}\right)$, or $\underline{n} = b_1 \underline{c}$ and $\mathrm{sh}\left(\frac{a,b}{n}\right) = \mathrm{sh}\left(\frac{a,b}{\underline{c}}\right)$, therefore (30) and (31) coincide, which proves (14) with $M = S_{\mathcal{N}}$. \Box

We are now in a position to obtain a new formula for $\log \Psi$, on which its Lie character is manifest—the new formula thus contains the BCH theorem:

Corollary 4.5. Let $T_N := \log S_N \in \mathbb{Q}^N$. Then, with the notation of Theorem C, we have $log \Psi = T_{\mathcal{N}}[D], i.e.$

$$
\log(e^{tX}e^{tY}) = \sum_{r\geq 1} \sum_{n_1,\dots,n_r=1}^{\infty} \frac{1}{r} T_{\mathcal{N}}^{n_1\cdots n_r} [D_{n_1},[\cdots [D_{n_{r-1}}, D_{n_r}]\cdots]] \in \text{Lie}(X,Y)[[t]].
$$

Proof. From Theorem C and Lemma 2.4 we deduce

$$
\log \Psi = \log(S_{\mathcal{N}}D) = T_{\mathcal{N}}D. \tag{32}
$$

 \Box

By Lemmas 2.10 and 4.4, T_N is alternal. We conclude by Lemma 2.9.

From the definition $T_{\mathcal{N}} = \sum_{n=1}^{\infty}$ $k=1$ $(-1)^{k-1}$ $\frac{\partial^{k-1}}{\partial k}(S_{\mathcal{N}}-1)^{\times k}$, we can write down the coefficients for words of small length:

$$
T^{n_1} = S^{n_1} = \frac{1}{n_1}
$$

\n
$$
T^{n_1 n_2} = S^{n_1 n_2} - \frac{1}{2} S^{n_1} S^{n_2} = \frac{n_1 - n_2}{2n_1 n_2 (n_1 + n_2)}
$$

\n
$$
T^{n_1 n_2 n_3} = S^{n_1 n_2 n_3} - \frac{1}{2} S^{n_1 n_2} S^{n_3} - \frac{1}{2} S^{n_1} S^{n_2 n_3} + \frac{1}{3} S^{n_1} S^{n_2} S^{n_3}
$$

\n
$$
T^{n_1 n_2 n_3 n_4} = S^{n_1 n_2 n_3 n_4} - \frac{1}{2} S^{n_1} S^{n_2 n_3 n_4} - \frac{1}{2} S^{n_1 n_2} S^{n_3 n_4} - \frac{1}{2} S^{n_1 n_2 n_3} S^{n_4}
$$

\n
$$
+ \frac{1}{3} S^{n_1} S^{n_2} S^{n_3 n_4} + \frac{1}{3} S^{n_1} S^{n_2 n_3} S^{n_4} + \frac{1}{3} S^{n_1 n_2} S^{n_3} S^{n_4} - \frac{1}{4} S^{n_1} S^{n_2} S^{n_3} S^{n_4}
$$

(omitting the subscript N to lighten notation). The low powers of t in log $\Psi = T_N[D]$ can then be extracted from the Lie mould expansion and we recover the classical BCH series:

· · · · · ·

$$
\log \Psi = \sum_{n_1=1}^{\infty} T^{n_1} D_{n_1} + \sum_{n_1,n_2=1}^{\infty} \frac{1}{2} T^{n_1 n_2} [D_{n_1}, D_{n_2}] + \sum_{n_1,n_2,n_3=1}^{\infty} \frac{1}{3} T^{n_1 n_2 n_3} [D_{n_1}, [D_{n_2}, D_{n_3}]]
$$

+
$$
\sum_{n_1,n_2,n_3,n_4=1}^{\infty} \frac{1}{4} T^{n_1 n_2 n_3 n_4} [D_{n_1}, [D_{n_2}, [D_{n_3}, D_{n_4}]]] + \cdots
$$

= $t(X+Y) + \frac{t^2}{2}[X,Y] + \frac{t^3}{3!}[X,[X,Y]] + \frac{t^4}{4!}[X,[X,[X,Y]]] + \frac{t^5}{5!}[X,[X,[X,Y]]]] + \cdots$
-
$$
\frac{t^3}{12}([(X+Y), [X,Y]]) - \frac{t^4}{24}([(X+Y), [X,[X,Y]]]) - \frac{t^5}{120}[[X,Y], [X,[X,Y]]]] + \cdots
$$

+
$$
\frac{t^5}{720}[(X+Y), [(X+Y), [X,[X,Y]]]] - \frac{t^5}{240}[[X,Y], [(X+Y), [X,Y]]] + \cdots
$$

+
$$
\frac{t^5}{720}[(X+Y), [(X+Y), [(X+Y), [X,Y]]]] + \cdots
$$

= $t(X+Y) + \frac{t^2}{2}[X,Y] + \frac{t^3}{12}([X,[X,Y]] + [Y,[Y,X]]) - \frac{t^4}{24}[Y,[X,[X,Y]]]$
-
$$
\frac{t^5}{720}[X,[X,[X,[X,Y]]]] - \frac{t^5}{720}[Y,[Y,[Y,[Y,[Y,X]]]] + \frac{t^5}{360}[X,[Y,[Y,[X,X]]]] + \cdots
$$

5 Generalization to an arbitrary number of factors

One of the merits of the mould calculus approach is that the formulas are easily generalized to the case of

$$
\Psi = e^{tX_1} \cdots e^{tX_N} \in \mathscr{A}[[t]],
$$

where $\mathscr A$ us an associative algebra and $X_1, \ldots, X_N \in \mathscr A$ for some $N \geq 2$.

5.1 Mould expansion of the first kind

Theorem B'. Let $\mathbb{N}_*^N \coloneqq \{ p \in \mathbb{N}^N \mid p_1 + \cdots + p_N \geq 1 \}$. We have

$$
\log \Psi = \sum \frac{(-1)^{k-1}}{k} \frac{t^{\sigma}}{\sigma} \frac{[X_1^{p_1^1} \cdots X_N^{p_N^1} \cdots X_1^{p_1^k} \cdots X_N^{p_N^k}]}{p_1^{1!} \cdots p_N^{1!} \cdots p_1^{k!} \cdots p_N^{k!}}
$$

with summation over all $k \in \mathbb{N}^*$ and $p^1, \dots, p^k \in \mathbb{N}^N_*$, where $\sigma \coloneqq \sum^k$ $i=1$ $\frac{N}{\sum}$ $j=1$ p^i_j and the bracket denote nested commutators as before.

Proof. Let $\Omega := \{x_1, \ldots, x_N\}$ be an N-element set. We consider the associative comould generated by the family

$$
B_{x_1} \coloneqq tX_1, \ \dots, \ B_{x_N} \coloneqq tX_N \ \in \mathscr{A}[[t]]. \tag{33}
$$

We can write $tX_1 = I_1B, \ldots, tX_N = I_NB$, with moulds $I_1, \ldots, I_N \in \mathbb{Q}^{\Omega}$ defined by

$$
I_j^{\underline{\omega}}:=\begin{cases} 1 \ \ \text{if}\ \underline{\omega}\ \text{is the one-letter word}\ x_j \\ 0 \ \ \text{else} \end{cases}
$$

for $j = 1, \ldots, N$. Hence

$$
\Psi = S_{\Omega} B \text{ with } S_{\Omega} := e^{I_1} \times \cdots \times e^{I_N}, \quad \log \Psi = T_{\Omega} B \text{ with } S_{\Omega} := \log S_{\Omega}. \tag{34}
$$

The moulds I_1, \ldots, I_N are alternal (being supported in one-letter words), hence Lemma 2.10 entails that their exponentials are symmetral, and also S_{Ω} , while T_{Ω} is alternal. We deduce that

$$
\log \Psi = T_{\Omega}[B] = \sum_{k \ge 1} \frac{(-1)^{k-1}}{k} \sum_{\underline{\omega}^1, \dots, \underline{\omega}^k \in \underline{\Omega} \backslash \{\varnothing\}} \frac{1}{r(\underline{\omega}^1) + \dots + r(\underline{\omega}^k)} S_{\Omega}^{\underline{\omega}^1} \dotsm S_{\Omega}^{\underline{\omega}^k} B_{[\underline{\omega}^1 \dotsm \underline{\omega}^k]}.
$$

The conclusion stems from the fact that

$$
S_{\overline{\Omega}}^{\underline{\omega}} = \begin{cases} \frac{1}{p_1! \cdots p_N!} & \text{if } \underline{\omega} \text{ is of the form } x_1^{p_1} \cdots x_N^{p_N} \text{ with } (p_1, \dots, p_N) \in \mathbb{N}^N \\ 0 & \text{else.} \end{cases}
$$

 \Box

5.2 Mould expansion of the second kind

Theorem C'. In the above situation, $\Psi = e^{tX_1} \cdots e^{tX_N}$ can also be written

$$
\Psi = 1_{\mathscr{A}} + \sum_{r=1}^{\infty} \sum_{n_1,\dots,n_r=1}^{\infty} \frac{1}{n_r(n_r + n_{r-1})\cdots(n_r + \cdots + n_1)} \mathfrak{D}_{n_1}\cdots\mathfrak{D}_{n_r}
$$
(35)

with
$$
\mathfrak{D}_n := t^n \sum_{j=1}^N \sum_{\substack{m_1, \dots, m_{j-1} \in \mathbb{N} \\ m_1 + \dots + m_{j-1} = n-1}} \frac{\mathrm{ad}_{X_1}^{m_1} \cdots \mathrm{ad}_{X_{j-1}}^{m_{j-1}}}{m_1! \cdots m_{j-1}!} X_j
$$
 for each $n \ge 1$. (36)

Note that formula (35) involves exactly the same rational coefficients as in the case $N = 2$. The only difference in the formula is that the D_n 's of (21) have been generalized to the \mathfrak{D}_n 's which are defined in (36) and read

$$
\mathfrak{D}_n := \begin{cases} t(X_1 + \dots + X_N) & \text{for } n = 1, \\ t^n \frac{\text{ad}_{X_1}^{n-1}}{(n-1)!} X_2 + \dots + t^n & \sum_{m_1 + \dots + m_{N-1} = n-1} \frac{\text{ad}_{X_1}^{m_1} \dots \text{ad}_{X_{N-1}}^{m_{N-1}}}{m_1! \dots m_{N-1}!} X_N & \text{for } n > 1. \end{cases}
$$
(37)

Proof. We have $\Psi_{|t=0} = 1_{\mathscr{A}}$ and

$$
t\partial_t \Psi = tX_1 e^{tX_1} \cdots e^{tX_N} + t e^{tX_1} X_2 e^{tX_2} \cdots e^{tX_N} + \cdots + t e^{tX_1} \cdots e^{tX_{N-1}} X_N e^{tX_N}
$$

$$
= \mathfrak{D}\Psi, \quad \text{where } \mathfrak{D} := t \sum_{j=1}^N \text{Ad}_{e^{tX_1}} \cdots \text{Ad}_{e^{tX_{j-1}}} X_j
$$

with the notation $\text{Ad}_E A = E A E^{-1}$ for any $A \in \mathscr{A}[[t]]$ whenever E is an invertible element of $\mathscr{A}[[t]]$. Moreover, we observe that there is no other solution in $\mathscr{A}[[t]]$ to the system

$$
\Psi_{|t=0} = 1_{\mathscr{A}}, \qquad t\partial_t \Psi = \mathfrak{D}\Psi, \tag{38}
$$

because ord $\mathfrak{D} \geq 1$.

Thanks to (24), we compute $\mathfrak{D} = t \sum_{j=1}^{N} e^{ad_t X_1} \cdots e^{ad_t X_{j-1}} X_j = \sum_{n \geq 1} \mathfrak{D}_n$. Let us thus take $\mathcal{N} = \mathbb{N}^*$ as alphabet and consider the associative comould generated by $(\mathfrak{D}_n)_{n \in \mathcal{N}}$, so that $\mathfrak D$ can be rewritten as the mould expansion $I\mathfrak D$, with the same mould as in (16).

Lemmas 2.4 and 4.2 show that a mould expansion $\Psi = SD$ is solution to (38) if $S \in \mathbb{Q}^{\mathcal{N}}$ is solution to the mould equation (25) (indeed: $(\nabla S)\mathfrak{D} = t\partial_t \Psi$ on the one hand, and $(I \times S)$ $\mathfrak{D} = (I \mathfrak{D})(S \mathfrak{D}) = \mathfrak{D} \Psi$ on the other hand, and $S^{\varnothing} = 1$ ensures ord $(\Psi - 1_{\mathscr{A}}) \geq 1$ because ord $\mathfrak{D}_n \geq 1$ for all nonempty word \underline{n}). But we already know that $S = S_{\mathcal{N}}$ defined by (27) is the unique solution to (25), hence

$$
\Psi = S_{\mathcal{N}} \mathfrak{D},\tag{39}
$$

 \Box

which is equivalent to (35).

Notice that, in view of Section 4.2, the mould S_N is symmetral, the mould $T_N = \log S_N$ is alternal, whence

$$
\log \Psi = T_{\mathcal{N}} \mathfrak{D} = T_{\mathcal{N}}[\mathfrak{D}], \tag{40}
$$

i.e.

$$
\log(e^{tX_1}\cdots e^{tX_N}) = \sum_{r\geq 1} \sum_{n_1,\ldots,n_r=1}^{\infty} \frac{1}{r} T_N^{n_1\cdots n_r} \left[\mathfrak{D}_{n_1},\left[\cdots\left[\mathfrak{D}_{n_{r-1}},\mathfrak{D}_{n_r}\right]\cdots\right]\right]
$$

which thus belongs to $Lie(X_1, \ldots, X_N)[[t]]$, in accordance with the BCH theorem.

6 Relation between the two kinds of mould expansion

In our application to products of two or more exponentials, we have seen two different kinds of mould expansion. The first kind involves an N-element alphabet $\Omega := \{x_1, \ldots, x_N\}$ and the comould generated by the family $(B_{\omega})_{\omega \in \Omega}$ defined by (33). For the second one, the alphabet is $\mathcal{N} := \mathbb{N}^*$ and the comould is generated by the family $(\mathfrak{D}_n)_{n \in \mathcal{N}}$ which is defined by (37) and boils down to the D_n 's of (21) when $N = 2$. A natural question is: What is the relation between both kinds of mould expansion? *i.e.* can one pass from the representation of the product Ψ as $S_{\Omega}B$ in (34) to its representation as $S_{\mathcal{N}}\mathfrak{D}$ in (39), or from $\log \Psi = T_{\Omega}B$ in (34) to $\log \Psi = T_{\mathcal{N}} \mathfrak{D}$ in (40)?

In this section, we will answer this question by defining a new operation on moulds, which allows one to pass directly from S_N to S_Ω , or from T_N to T_Ω . We take $N = 2$ for simplicity but the generalization to arbitrary N is easy.

We start by giving a mould expansion of the first kind for the D_n 's themselves.

Lemma 6.1. Let $\Omega := \{x, y\}$. The formula

$$
\underline{\omega} \in \underline{\Omega} \mapsto U^{\underline{\omega}} := \begin{cases} 1 & \text{if } \underline{\omega} = x \\ \frac{(-1)^q}{p!q!} & \text{if } \underline{\omega} \text{ is of the form } x^p y x^q \text{ for some } p, q \in \mathbb{N} \\ 0 & \text{else} \end{cases} \tag{41}
$$

defines an alternal mould $U \in \mathbb{Q}^{\Omega}$ such that

$$
D_n = U_n B \quad \text{for each } n \in \mathbb{N}^*, \tag{42}
$$

where the left-hand side is defined by (21) and the right-hand side is the mould expansion (for the comould generated by (9)) associated with

 $U_n \coloneqq \text{restriction of } U$ to the words of length n.

Proof. In view of (8), we have $\text{ad}_{MB}(NB) = [MB, NB] = [M, N]B = (\text{ad}_{M} N)B$ for any $M, N \in \mathbb{Q}^{\Omega}$, hence (21) can be rewritten as $D_n = \frac{1}{(n-1)!} \operatorname{ad}_{I_x}^{n-1}((I_x + I_y)B) = U_nB$ with $U_n := \frac{1}{(n-1)!} \operatorname{ad}_{I_x}^{n-1}(I_x + I_y)$. Since I_x and I_y are alternal and the set of all alternal moulds is stable under mould commutator (as mentioned at the end of Section 2.3), we see that this mould U_n is alternal. Since the support of U_n is contained in the set of words of length n, the formula $U := \sum_{n \geq 1} U_n$ makes sense and defines an alternal mould (and U_n now appears as the restriction of this U to the set of words of length n). There only remains to check (41) .

Now, $ad_{I_x} = L - R$, where L and R are the operators of left-multiplication and rightmultiplication by I_x , which commute, hence the binomial theorem yields

$$
U_n = \sum_{p+q=n-1} \frac{(-1)^q}{p!q!} L^p R^q (I_x + I_y) = \sum_{p+q=n-1} \frac{(-1)^q}{p!q!} I_x^{\times p} \times (I_x + I_y) \times I_x^{\times q},
$$

i.e. $U_n^{\underline{\omega}} = 1$ if $\underline{\omega} = x$ and $n = 1$, $\frac{(-1)^q}{p!q!}$ if $\underline{\omega}$ is of the form $x^p y x^q$ for some $p, q \in \mathbb{N}$ such that $p + q = n - 1$ (in which case p and q are uniquely determined), and 0 else. Our U thus coincides with the mould defined by (41). \Box In fact the proof just given shows that

$$
U = e^{ad_{I_x}} (I_x + I_y) = e^{I_x} \times (I_x + I_y) \times e^{-I_x}.
$$
 (43)

This mould will allow us to relate D-mould expansions and B-mould expansions:

Theorem D. Let $\mathcal{N} := \mathbb{N}^*$. Define a linear map $M \in \mathbb{Q}^{\underline{N}} \mapsto M \odot U \in \mathbb{Q}^{\underline{\Omega}}$ by the formulas

$$
(M \odot U)^{\varnothing} \coloneqq M^{\varnothing},\tag{44}
$$

$$
(M \odot U)^{\underline{\omega}} := \sum_{s \ge 1} \sum_{\underline{\omega} = \underline{\omega}^1 \cdots \underline{\omega}^s \atop \underline{\omega}^1, \dots, \underline{\omega}^s \in \underline{\Omega} \setminus \{\varnothing\}} M^{r(\underline{\omega}^1) \cdots r(\underline{\omega}^s)} U^{\underline{\omega}^1} \cdots U^{\underline{\omega}^s} \quad \text{for } \underline{\omega} \in \underline{\Omega} \setminus \{\varnothing\}. \tag{45}
$$

Then

 $MD = (M \odot U)B$ for any $M \in \mathbb{Q}^{\mathcal{N}}$.

Recall that $r: \Omega \to \mathbb{N}^* = \mathcal{N}$ is our notation for the length function. In (45), $r(\omega^1) \cdots r(\omega^s)$ is to be understood as a word of length s of $\mathcal N$ (and the sum is finite because the words ω^j are nonempty, hence $s \leq r(\underline{\omega})$.

Proof. By direct computation, using (42) to express $D_{n_1\cdots n_s} = D_{n_1} \cdots D_{n_s}$,

$$
MD = \sum_{\underline{n} \in \underline{N}} M^{\underline{n}} D_{\underline{n}} = M^{\varnothing} 1_{\mathcal{A}} + \sum_{s \ge 1} \sum_{n_1, \dots, n_s \in \mathcal{N}} M^{n_1 \cdots n_s} \sum_{\underline{\omega}^1 \in \underline{\Omega} \atop r(\underline{\omega}^1) = n_1} U^{\underline{\omega}^1} B_{\underline{\omega}^1} \cdots \sum_{\underline{\omega}^s \in \underline{\Omega} \atop r(\underline{\omega}^s) = n_1} U^{\underline{\omega}^s} B_{\underline{\omega}^s}
$$

\n
$$
= M^{\varnothing} 1_{\mathcal{A}} + \sum_{s \ge 1} \sum_{n_1, \dots, n_s \in \mathcal{N}} M^{n_1 \cdots n_s} \sum_{\underline{\omega}^1, \dots, \underline{\omega}^s \in \underline{\Omega} \atop r(\underline{\omega}^1) = n_1, \dots, r(\underline{\omega}^s) = n_s} U^{\underline{\omega}^1} \cdots U^{\underline{\omega}^s} B_{\underline{\omega}^1 \cdots \underline{\omega}^s}
$$

\n
$$
= M^{\varnothing} 1_{\mathcal{A}} + \sum_{s \ge 1} \sum_{\underline{\omega}^1, \dots, \underline{\omega}^s \in \underline{\Omega} \setminus \{\varnothing\}} M^{r(\underline{\omega}^1) \cdots r(\underline{\omega}^s)} U^{\underline{\omega}^1} \cdots U^{\underline{\omega}^s} B_{\underline{\omega}^1 \cdots \underline{\omega}^s}
$$

\n
$$
= M^{\varnothing} 1_{\mathcal{A}} + \sum_{\underline{\omega} \in \underline{\Omega} \setminus \{\varnothing\}} \left(\sum_{s \ge 1, \underline{\omega}^1, \dots \underline{\omega}^s \in \underline{\Omega} \setminus \{\varnothing\}} M^{r(\underline{\omega}^1) \cdots r(\underline{\omega}^s)} U^{\underline{\omega}^1} \cdots U^{\underline{\omega}^s} \right) B_{\underline{\omega}} = (M \odot U) B.
$$

The relations $S_{\mathcal{N}}D = S_{\Omega}B$ (which coincides with Ψ according to (10) and (28)) and $T_{\mathcal{N}}D = T_{\Omega}B$ (which coincides with log Ψ according to (28) and (32)) now appear as a manifestation of Theorem D and the following

Theorem E.

$$
S_{\mathcal{N}} \odot U = S_{\Omega}, \qquad T_{\mathcal{N}} \odot U = T_{\Omega}.
$$

The proof of Theorem E is given at the end of this section.

Our definition (44)–(45) of the mould operation ' \odot ' is a variant of Écalle's mould composition '◦' which is defined for any alphabet that is a commutative semigroup ($[Ec84]$, [Sa09], [FFM17]). Here is a definition which encompasses both operations:

Definition 6.2. Given two alphabets Ω and \mathcal{N} , and a map $\sigma: \Omega \setminus \{\emptyset\} \to \mathcal{N}$, we define the σ -composition

$$
(M, U) \in \mathbf{k}^{\underline{N}} \times \mathbf{k}^{\underline{\Omega}} \mapsto M \circ_{\sigma} U \in \mathbf{k}^{\underline{\Omega}}
$$

by the formulas

$$
(M \circ_{\sigma} U)^{\varnothing} := M^{\varnothing},\tag{46}
$$

$$
(M \circ_{\sigma} U)^{\underline{\omega}} := \sum_{s \ge 1} \sum_{\substack{\underline{\omega} = \underline{\omega}^{1} \cdots \underline{\omega}^{s} \\ \underline{\omega}^{1} \cdots \underline{\omega}^{s} \in \Omega \setminus \{\varnothing\}}} M^{\sigma(\underline{\omega}^{1}) \cdots \sigma(\underline{\omega}^{s})} U^{\underline{\omega}^{1}} \cdots U^{\underline{\omega}^{s}} \quad \text{for } \underline{\omega} \in \underline{\Omega} \setminus \{\varnothing\}. \tag{47}
$$

Thus, we recover the ' \odot ' composition in the special case when $\mathcal{N} = \mathbb{N}^*$ and $\sigma(\underline{\omega}) = r(\underline{\omega})$ (with arbitrary Ω), and Écalle's composition ' \circ ' when $\mathcal{N} = \Omega$ is a commutative semigroup and $\sigma(n_1 \cdots n_r) = n_1 + \cdots + n_r$ for any nonempty word of $\underline{\mathcal{N}}$. Some classical properties of the latter operation can be generalized as follows:

(i) $(M \circ_{\sigma} U) \times (N \circ_{\sigma} U) = (M \times N) \circ_{\sigma} U$.

(ii)
$$
e^{M\circ_{\sigma}U} = (e^M)\circ_{\sigma} U
$$
 if $M^{\emptyset} = 0$, $\log(M \circ_{\sigma} U) = (\log M) \circ_{\sigma} U$ if $M^{\emptyset} = 1$.

- (iii) $I \circ_{\sigma} U = U U^{\emptyset} \mathbb{1}_{\Omega}$, where I is defined by (16) and $\mathbb{1}_{\Omega}$ is the unit of $\mathbf{k}_{\Omega}^{\Omega}$.
- (iv) Denote by $\iota_{\Omega} \colon \Omega \hookrightarrow \Omega \setminus \{\varnothing\}$ the inclusion map. If $\phi \colon \mathcal{N} \to \mathbf{k}$ is a function such that $\phi \circ \sigma$ maps the concatenation in Ω to the addition in **k**, then

$$
(\nabla_{\phi} M) \circ_{\sigma} U = \nabla_{\psi} (M \circ_{\sigma} U) \quad \text{for all } M \in \mathbf{k}^{\underline{N}}, \qquad \text{with } \psi := \phi \circ \sigma \circ \iota_{\Omega}, \qquad (48)
$$

where ∇_{ϕ} and ∇_{ψ} are the mould derivations defined by (29).

(v) If U is alternal and $\sigma(\omega_1 \cdots \omega_r) = \sigma(\omega_{\tau(1)} \cdots \omega_{\tau(r)})$ for every permutation τ and for any $\omega_1, \ldots, \omega_r \in \Omega$, then

M alternal $\Rightarrow M \circ_{\sigma} U$ alternal, M symmetral $\Rightarrow M \circ_{\sigma} U$ symmetral.

(vi) Suppose $(B_{\omega})_{\omega \in \Omega}$ satisfies Assumption 2.2. Then the formula $D_n := \sum$ $\underline{\omega} \in \sigma^{-1}(n)$ $U^{\underline{\omega}}B_{\underline{\omega}}$ defines a family $(D_n)_{n \in \mathcal{N}}$ which also satisfies Assumption 2.2, and

$$
MD = (M \circ_{\sigma} U)B \quad \text{for any } M \in \mathbf{k}^{\underline{N}}.
$$

(vii) Suppose that $\tau: \mathcal{N} \setminus \{\varnothing\} \to \mathcal{M}$ is a map such that $\psi \coloneqq \tau \circ \iota_{\mathcal{N}} \circ \sigma$ satisfies

$$
\psi(\underline{\omega}^1 \cdots \underline{\omega}^s) = \tau(\sigma(\underline{\omega}^1) \cdots \sigma(\underline{\omega}^s)) \text{ for any } s \ge 1 \text{ and } \underline{\omega}^1, \dots \underline{\omega}^s \in \underline{\Omega} \setminus \{\varnothing\},
$$

then

$$
M \circ_{\psi} (N \circ_{\sigma} U) = (M \circ_{\tau} N) \circ_{\sigma} U
$$
 for any $M \in \mathbf{k}^{\underline{M}}, N \in \mathbf{k}^{\underline{N}}, U \in \mathbf{k}^{\underline{\Omega}}$.

(The proof of these properties is left to the reader.)

Proof of Theorem E. Here $\Omega = \{x, y\}, \mathcal{N} = \mathbb{N}^*$ and $\sigma = r: \Omega \to \mathcal{N}$ is word length. Since $T_{\mathcal{N}} = \log S_{\mathcal{N}}$ and $T_{\Omega} = \log S_{\Omega}$, in view of (ii) it is sufficient to prove $S_{\mathcal{N}} \circ_{\sigma} U = S_{\Omega}$.

As noticed in Section 4.1, S_N is a solution in \mathbb{Q}^N to equation (25), which involves $\nabla = \nabla_{\phi}$, with the notation $\phi \colon \mathcal{N} \hookrightarrow \mathbb{Q}$ for the inclusion map. Taking ' $\odot U$ ' of both sides of (25) , we get

$$
(\nabla_{\phi} S_{\mathcal{N}}) \circ_{\sigma} U = (I \times S_{\mathcal{N}}) \circ_{\sigma} U. \tag{49}
$$

We compute the left-hand side by means of (iv): $\phi \circ \sigma(\underline{\omega}) = r(\underline{\omega})$ is word length, in particular it maps concatenation in $\underline{\Omega}$ to addition in \mathbb{Q} , and $\phi \circ \sigma \circ \iota_{\Omega} \equiv 1$, hence the left-hand side is $\nabla_1(S_N \circ_{\sigma} U)$. Note that the mould derivation ∇_1 is given by $(\nabla_1 M)^{\underline{\omega}} = r(\underline{\omega})M^{\underline{\omega}}$.

By (i) and (iii), the right-hand side of (49) is $(I \circ_{\sigma} U) \times (S_{\mathcal{N}} \circ_{\sigma} U) = U \times (S_{\mathcal{N}} \circ_{\sigma} U)$. Therefore, $S_{\mathcal{N}} \circ_{\sigma} U$ is a solution to

$$
M^{\emptyset} = 1, \qquad \nabla_1 M = U \times M. \tag{50}
$$

 \Box

It is easy to see that (50) has no other solution in \mathbb{Q}^{Ω} .

On the other hand, by (10), $S_{\Omega} = e^{I_x} \times e^{I_y}$, and ∇_1 is a mould derivation which satisfies $\nabla_1 I_x = I_x$ and $\nabla_1 I_y = I_y$, thus

$$
\nabla_1 S_{\Omega} = \nabla_1 (e^{I_x}) \times e^{I_y} + e^{I_x} \times \nabla_1 (e^{I_y}) = I_x \times e^{I_x} \times e^{I_y} + e^{I_x} \times I_y \times e^{I_y}
$$

$$
= (I_x + e^{I_x} \times I_y \times e^{-I_x}) \times e^{I_x} \times e^{I_y} = U \times S_{\Omega}
$$

by (43). Therefore S_{Ω} is a solution to (50), hence it must coincide with $S_{\mathcal{N}} \circ_{\sigma} U$.

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References

- [BS17] S. Baumard and L. Schneps, On the derivation representation of the fundamental Lie algebra of mixed elliptic motives, Annales Mathématiques du Québec 41, 1 (2017), 43–62.
- [BD16] A. Behtash, G. V. Dunne, T. Schäfer, T. Sulejmanpasic and M. Unsal, *Com*plexified path integrals, exact saddles, and supersymmetry, Phys. Rev. Lett. 116 (2016),011601.
- [BF12] A. Bonfiglioli and R. Fulci, Topics in Noncommutative Algebra—The Theorem of Campbell, Baker, Hausdorff and Dynkin, Lecture Notes in Mathematics, 2034. Springer, Heidelberg, (2012). xxii+539 pp.
- $[BE17]$ O. Bouillot and J. Écalle, *Invariants of identity-tangent diffeomorphisms ex*panded as series of multitangents and multizetas, in Resurgence, physics and numbers, 109–232 (2017), CRM Series, 20, Ed. Norm., Pisa.
- [CSV17] R. Couso-Santamaría, R. Schiappa and R.Vaz, On asymptotices and resurgent structures of enumerative Gromov-Witten invariants, Commun. Num. Theor. Phys. 11 (2017) 707–790, [1605.07473].
- [CMS17] R. Couso-Santamaría1, M. Mariño and R. Schiappa, Resurgence matches quantization, J. Phys. A: Math. Theor. 50 (2017) 145402.
- [Dy47] E. B. Dynkin, Calculation of the coefficients in the Campbell-Hausdorff formula (Russian), Dokl. Akad. Nauk SSSR (N.S.), 57, (1947), 323–326.
- [DDP93] H. Dillinger, E. Delabaere and F. Pham, Résurgence de Voros et périodes des courbes hyperelliptiques, Annales de l'institut Fourier 43, 1 (1993), 163–199.
- [Ec81] J. Écalle, Les fonctions résurgentes, Publ. Math. d'Orsay [Vol.1: 81-05, Vol.2: 81-06,Vol.3: 85-05], (1981,1985).
- [Ec84] J. Écalle, *Cinq applications des fonctions résurgentes*, Publ. Math. d'Orsay, 84– 62, (1984).
- [Ec92] J. Écalle, *Introduction aux fonctions analysables et preuve constructive de la* conjecture de Dulac, Actualités Math., Hermann, Paris (1992).
- [Ec03] J. Écalle, $ARI/GARI$, la dimorphie et l'arithmétique des multizêtas: un premier $bilan$, Journal de Théorie des Nombres de Bordeaux 15, 2 (2003), 411–478.
- [FFM17] F. Fauvet, L. Foissy, and D. Manchon, The Hopf algebra of finite topologies and mould composition, Annales de l'Institut Fourier 67, 3 (2017), 911–945.
- [GMN13] D. Gaiotto, G. W. Moore and A. Neitzke, Wall-crossing, Hitchin Systems, and the WKB Approximation, Adv. in Math. 234, (2013), 239–403.
- [Ki17] T. Kimura, Explicit description of the Zassenhaus formula, Theor.Exp.Phys., (2017), 041A03.
- [Ko17] M. Kontsevich, Resurgence and Quantization, Course given at IHES, Paris in April, 2017.
- [Ma15] M. Matone, An algorithm for the Baker-Campbell-Hausdorff formula, J. High Energy Phys. 05 (2015) 113, [arXiv:1502.06589].
- [Me09] F. Menous, Formal differential equations and renormalization, in Renormalization and Galois theories, A. Connes, F. Fauvet, J.-P. Ramis (eds.), IRMA Lect.Math.Theor.Phys., 15 (2009), 229–246.
- [NP18] J.-C. Novelli, T. Paul, D. Sauzin, J.-Y. Thibon, "Rayleigh-Schrödinger series and Birkhoff decomposition", Letters in Mathematical Physics 108 (2018), 18 p. <https://doi.org/10.1007/s11005-017-1040-1>
- [PS17] T. Paul and D. Sauzin, Normalization in Lie algebras via mould calculus and applications, Regular and Chaotic Dynamics 22, 6 (2017), 616–649.
- [Re93] C. Reutenauer, Free Lie algebras, London Mathematical Society Monographs 7, Clarendon Press, Oxford University Press, New York 1993, xviii+269 pp.
- [Sa08] D. Sauzin, Initiation to mould calculus through the example of saddle-node singularities, Rev. Semin. Iberoam. Mat. 3 (2008), no. 5-6, 147–160.
- [Sa09] D. Sauzin, Mould expansions for the saddle-node and resurgence monomials, in Renormalization and Galois theories, A. Connes, F. Fauvet, J.-P. Ramis (eds.), IRMA Lectures in Mathematics and Theoretical Physics 15, Zürich: European Mathematical Society, 83–163 (2009).
- [Sa16] D. Sauzin, Introduction to 1-summability and resurgence, in Divergent series, summability and resurgence. I. Monodromy and resurgence, C. Mitschi, D. Sauzin. Lecture Notes in Mathematics, 2153. Springer, 2016. xxi+298 pp.
- [Sc12] L. Schneps, Double shuffle and Kashiwara-Vergne Lie algebras, Journal of Algebra 367 (2012), 54–74.
- [Th11] J.-Y. Thibon, Noncommutative symmetric functions and combinatorial Hopf algebras, in: Asymptotics in dynamics, geometry and PDEs; generalized Borel summation. Vol. I, (2011) 219–258, CRM Series, 12, Ed. Norm., Pisa.
- [Vo83] A. Voros, The return of the quartic oscillator. The complex WKB method, Annales de l'I. H. P., section A, tome 39, n° 3 (1983) 211−338.
- [vW66] W. von Waldenfels, "Zur Charakterisierung Liescher Elemente in freien Algebren," Arch. Math. (Basel) 17, 44–48 (1966).