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On a Feynman-Kac approach to growth-fragmentation semigroups and their asymptotic behaviors

Jean Bertoin*

This work develops further a probabilist approach to the asymptotic behavior of growth-fragmentation semigroups via the Feynman-Kac formula, which was introduced in a joint article with A.R. Watson [6]. Here, it is first shown that the sufficient condition for a Malthusian behavior which was established in [6], is also necessary. We then provide a simple criterion to ensure exponential speed of convergence, which enables us to treat cases than were not covered previously in the literature.

Keywords: Growth-fragmentation equation, Feynman-Kac formula, Malthus exponent, Krein-Rutman theorem, exponential ergodicity.

2010 Mathematics Subject Classification: 35Q92, 37A30, 47D06, 60G46, 60J25.

1 Introduction

Imagine a population, for instance of cells or of bacteria, where individuals grow and split as time passes, and such that the evolution of each individual only depends on its own mass, without interaction between different individuals. Assume also that when a fragmentation event happens, the sum of the masses of daughters resulting from the split equals the mass of the mother before splitting. In other words, the total mass is a preserved quantity when a fragmentation occurs, and grows between consecutive fragmentation events. The rate of growth of an individual may depend on its mass, and the rate at which a mother produces daughters may also depend both on the mass of the mother just before splitting and on the masses of its daughters right after the split.

Growth-fragmentation equations provide a mathematical model for such dynamics, by describing the evolution of concentrations $u(t, y)$ of individuals as a function of time

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$t \geq 0$ and mass $y > 0$. They can be expressed in the form

$$\partial_t u(t, y) + \partial_y(c(y)u(t, y)) + K(y)u(t, y) = \int_y^\infty k(z, y)u(t, z)dz, \quad (1)$$

where $c(y)$ describes the growth rate as a function of the mass, $k(z, y)$ the rate at which a daughter particle with mass y appears as the result of the fragmentation of a mother particle with mass $z > y$. Finally, $K(y)$ is the total fragmentation rate of individuals with mass y , and the assumption of conservation of mass at fragmentation events thus translates into

$$xK(x) = \int_0^x yk(x, y)dy \quad \text{for all } x > 0. \quad (2)$$

In this work, we rather deal with the weak form of (1) and focus on initial conditions of Dirac type. Namely, one introduces the operator

$$\mathcal{A}f(x) = c(x)f'(x) + \int_0^x f(y)k(x, y)dy - K(x)f(x), \quad x > 0, \quad (3)$$

which is defined on some domain $\mathcal{D}_\mathcal{A}$ of smooth functions $f : (0, \infty) \rightarrow \mathbb{R}$. Under fairly simple general assumptions on the rates c and k that will be introduced later on, \mathcal{A} is the infinitesimal generator of a unique strongly continuous positive semigroup $(T_t)_{t \geq 0}$, that is

$$\frac{dT_t f(x)}{dt} = \mathcal{A}(T_t f)(x), \quad f \in \mathcal{D}_\mathcal{A}. \quad (4)$$

In this setting, the measure $\mu_t(x, dy)$ on $(0, \infty)$ such that

$$T_t f(x) = \int_{(0, \infty)} f(y)\mu_t(x, dy) := \langle \mu_t(x, \cdot), f \rangle$$

should be interpreted informally as $\mu_t(x, dy) = u(t, y)dy$ and (4) as the weak version of (1) for the initial condition $\mu_0(x, dy) = u(0, y)dy = \delta_x(dy)$.

In general, there is of course no explicit expression for the growth-fragmentation semigroup $(T_t)_{t \geq 0}$, and many works in this area are concerned with its large time asymptotic behavior. See in particular [1, 2, 5, 8, 17, 19, 26, 27, 28, 30] and further references therein. Typically, one expects that under adequate assumptions on the rates of growth and of fragmentation, there exists a principal eigenvalue $\rho \in \mathbb{R}$ such that

$$\lim_{t \rightarrow \infty} e^{-\rho t} T_t f(x) = h(x)\langle \nu, f \rangle, \quad x > 0, \quad (5)$$

at least for every continuous and compactly supported function $f : (0, \infty) \rightarrow \mathbb{R}$. Here, $\nu(dy)$ is a Radon measure on $(0, \infty)$, which is often referred to as the asymptotic profile, and h some positive function. We stress that (5) may well fail; see for instance Doumic and Escobedo [15] and Gabriel [19].

When (5) holds, it is further important to be able to estimate the speed of convergence. Indeed, say for $\rho > 0$, an indefinite exponential growth is of course unrealistic in practice, and the growth-fragmentation equation can only be pertinent for describing rather early

stages of the evolution of a population when certain effects such as competition between individuals for space or resources can be neglected. As a consequence, the notions of principal eigenvalue and of asymptotic profile are only relevant for applications when the convergence in (5) occurs fast enough.

Spectral theory for semigroups and generators yields a well-established and classical framework for establishing the validity of (5), again provided that the growth and fragmentation rates are properly chosen. In short, if one can find positive eigenelements, namely a Radon measure ν and a positive function h on $(0, \infty)$, such that for some $\rho \in \mathbb{R}$:

$$\mathcal{A}h = \rho h, \quad \mathcal{A}^*\nu = \rho\nu, \quad \text{and} \quad \langle \nu, h \rangle = 1,$$

where \mathcal{A}^* denotes the dual of \mathcal{A} , then the so-called general relative entropy method (see in particular Chapter 6 in Perthame [29] and Michel et al. [27]) shows that (5) holds. In turn, explicit criteria in terms of the rates of growth c and of fragmentation k that ensure the existence of positive eigenelements, have been obtained by Michel [26] and by Doumic and Gabriel [17]. These works rely crucially on the Krein-Rutman theorem, a version of the Perron-Frobenius theorem for positive compact operators. On the other hand, exponential rate of convergence in (5) is essentially equivalent to the existence of a spectral gap. This has been obtained under specific assumptions on the growth and fragmentation rates notably by Perthame and Ryzhik [30], Laurençot and Perthame [23], Cáceres et al. [8] and Mischler and Scher [28].

Quite recently, together with A.R. Watson [6], we devised a probabilistic approach to (5), which circumvents spectral theory of semigroups and further provides probabilistic expressions for the various quantities of interest. This requires some assumptions on the growth rate c and the fragmentation k that we now introduce. First,

$$\text{the function } x \mapsto c(x)/x \text{ is continuous, positive and bounded on } (0, \infty). \quad (6)$$

In particular $c(x)/x$ is bounded away from 0 for x in any compact subset of $(0, \infty)$, but may tend to 0 as x goes to 0 and/or ∞ . Second, writing $\bar{k}(x, y) := x^{-1}yk(x, y)$ for every $0 < y < x$,

$$\text{the map } \begin{cases} (0, \infty) & \rightarrow L^1(dy) \\ x & \mapsto \bar{k}(x, \cdot) \end{cases} \quad \text{is continuous and } \sup_{x>0} \|\bar{k}(x, \cdot)\|_{L^1(dy)} = \|K\|_\infty < \infty. \quad (7)$$

Our probabilistic approach relies on an “instrumental” Markov process $X = (X_t)_{t \geq 0}$ with infinitesimal generator

$$\mathcal{G}f(x) := c(x)f'(x) + \int_0^x (f(y) - f(x))\bar{k}(x, y) dy. \quad (8)$$

In words, X describes a mass that grows continuously with rate $c(\cdot)$ and makes negative jumps at random times which should be thought of as fragmentation events. More precisely, the rate of negative jumps is $K(\cdot)$, and if a jump occurs when X is about to reach x , then its value immediately after the jump has the law $K(x)^{-1}\bar{k}(x, y)dy$. So the

dynamics of X can be thought of as that of the mass of a distinguished individual in the population, such that at each fragmentation event, the distinguished daughter is picked from her sibling by size-biased sampling.

Assumption (7) guaranties that the total jump rate (2) remains bounded, so the jump times of X never accumulate. One says X is piecewise deterministic (see [12] and references therein), in the sense that the trajectory $t \mapsto X_t$ is driven by the steady flow velocity c between two consecutive jumps, and jump times and locations are the sole source of randomness. We finally assume that

$$\text{the Markov process } X \text{ is irreducible,} \quad (9)$$

that is, for every $x, y > 0$, the probability that the Markov process started from x visits $y > 0$ is strictly positive. Roughly speaking, this means that there are no strict subintervals I of $(0, \infty)$ that form traps for X , in the sense that once the path enters I , it cannot exit from I . It can be seen that (9) is equivalent to a simple non-degeneracy assumption on the fragmentation rate kernel k , namely

$$\text{for every } x > 0, \text{ there exist } \alpha < x < \beta \text{ with } \int_{\alpha}^x \bar{k}(\beta, y) dy > 0; \quad (10)$$

see the forthcoming Lemma 3.1(i). We also recall that the absolute continuity of the fragmentation rates ensures aperiodicity of X (which might fail otherwise, as for instance in the case of equal mitosis; see e.g. [16]).

Lemma 2.2 in [6] shows that the growth-fragmentation semigroup T_t can then be given by a Feynman-Kac formula (we refer to [14] for treatise on this topic in discrete time):

$$T_t f(x) = x \mathbb{E}_x \left(\mathcal{E}_t \frac{f(X_t)}{X_t} \right), \quad \text{with } \mathcal{E}_t := \exp \left(\int_0^t \frac{c(X_s)}{X_s} ds \right), \quad (11)$$

where \mathbb{E}_x stands for the expectation when X starts at $X_0 = x$. The first hitting time of $y > 0$ by X ,

$$H(y) := \inf \{ t > 0 : X_t = y \}, \quad (12)$$

and the Laplace transform¹

$$L_{x,y}(q) := \mathbb{E}_x \left(e^{-qH(y)} \mathcal{E}_{H(y)}, H(y) < \infty \right), \quad q \in \mathbb{R}, \quad (13)$$

then play a key role for the asymptotic behavior of T_t as we shall now explain. Note that $L_{x,y}$ is a non-increasing convex function with values in $(0, \infty]$, with $\lim_{q \rightarrow \infty} L_{x,y}(q) = 0$ and $\lim_{q \rightarrow -\infty} L_{x,y}(q) = \infty$. In particular, it possesses a right-derivative $L'_{x,y}(q)$ at every point q of its effective domain, i.e. with $L_{x,y}(q) < \infty$. Defining the Malthus exponent by

$$\lambda := \inf \{ q \in \mathbb{R} : L_{x,x}(q) < 1 \} \quad (14)$$

¹In (13) and in many further instances in this text, we use the notation $\mathbb{E}(Z, \Lambda) = \mathbb{E}(Z \mathbb{1}_{\Lambda})$ when Z is a nonnegative random variable and Λ an event.

(actually, this definition does not depend on $x > 0$; see Proposition 3.1 in [6]), the main results of [6] can be summarized as follows; see Theorem 5.3 in [6]. First, if

$$L_{x,x}(\lambda) = 1 \quad \text{and} \quad -L'_{x,x}(\lambda) < \infty \quad (15)$$

(again, this condition does not depend on x ; see Lemma 5.2 in [6]), then (5) holds with $\rho = \lambda$. Moreover, the asymptotic profile ν and the function h are given for some arbitrarily chosen $x_0 > 0$ by

$$h(y) = yL_{y,x_0}(\lambda) \quad \text{and} \quad \nu(dy) = \frac{dy}{h(y)c(y)|L'_{y,y}(\lambda)|}, \quad y > 0. \quad (16)$$

Second, if

$$\text{there exists some } q < \lambda \text{ and } x > 0 \text{ with } L_{x,x}(q) < \infty, \quad (17)$$

then the convergence (5) takes place exponentially fast. Specifically, there exists $\beta > 0$ such that

$$e^{-\lambda t}T_t f(x) = h(x)\langle \nu, f \rangle + o(e^{-\beta t}) \quad \text{as } t \rightarrow \infty \quad (18)$$

for every continuous function f with compact support and every $x > 0$. We stress that, by convexity of $L_{x,x}$, (17) is of course a stronger requirement than (15).

The idea of representing a non-Markovian semigroup in terms of some (possibly time-inhomogeneous) Markov process and then deriving the asymptotic behavior of the former from ergodic properties of the latter, has been used previously in the literature; see e.g. [4, 9, 10, 12] and references therein. In the setting of branching processes, the so-called many-to-one formulas and spine decompositions enable to express expectations of linear functionals of the system in terms of a single ‘‘typical’’ individual, where typical may be taken in the sense of an individual chosen uniformly at random in the system (like e.g. in [3, 4, 24]), or according to some well-chosen bias (like e.g. in [18, 11]). These approaches are mostly efficient when some information about the non-Markovian semigroup or its eigenelements is already available. The approach via the Feynman-Kac formula in [6] and further developed in the present work differs conceptually, in the sense that the ‘‘instrumental’’ Markov process X is not meant to describe the evolution of a ‘‘typical’’ individual. An important feature is that its evolution is given directly in terms of the parameters c and k of the model, without requiring *a priori* information about the eigenelements. The price to be paid is that the analysis of the Feynman-Kac formula is more involved; however it can be completed in a fairly great generality.

We have here two main purposes. First, we shall observe that the condition (15) is also necessary for the Malthusian behavior (5), and in particular, whenever the latter holds, the principal eigenvalue ρ is always given by the Malthus exponent defined by (14). We shall actually establish an even slightly stronger result.

Theorem 1.1. *Suppose that (6), (7) and (9) hold, and further that for some $\rho \in \mathbb{R}$:*

- (i) *there exist $x_1 > 0$ and a continuous function $f : (0, \infty) \rightarrow \mathbb{R}_+$ with compact support and $f \not\equiv 0$, such that*

$$\limsup_{t \rightarrow \infty} e^{-\rho t} T_t f(x_1) < \infty,$$

(ii) there exist $x_2 > 0$ and a continuous function $g : (0, \infty) \rightarrow \mathbb{R}_+$ with compact support, such that

$$\liminf_{t \rightarrow \infty} e^{-\rho t} T_t g(x_2) > 0.$$

Then $\rho = \lambda$, (15) holds and thus also the Malthusian behavior (5) with (16).

The assumptions in (15) and (17) are given in terms of the Laplace transform $L_{x,x}$ rather than directly in terms of the coefficients c and k as one might have wished, and the second purpose of the present work is to remedy (at least partly) this problem by providing the following simple criterion.

Theorem 1.2. *Suppose that (6), (7) and (9) hold. If the Malthus exponent λ defined in (14) and the growth rate c fulfill*

$$\limsup_{x \rightarrow 0+} \frac{c(x)}{x} < \lambda \quad \text{and} \quad \limsup_{x \rightarrow \infty} \frac{c(x)}{x} < \lambda, \quad (19)$$

then the exponential convergence (18) holds.

Theorem 1.2 might seem unsatisfactory, as its requirements are not given only in terms of the rates c and k , but also involve the Malthus exponent λ . However, there are simple explicit conditions in terms of c and k only that ensure (19). In particular, it is easily seen that $\lambda > \inf_{x>0} c(x)/x$ when X is recurrent and c is not linear, cf. Proposition 3.4(ii) in [6]. Thus (19) is then fulfilled whenever

$$\lim_{x \rightarrow 0+} \frac{c(x)}{x} = \lim_{x \rightarrow \infty} \frac{c(x)}{x} = \inf_{x>0} \frac{c(x)}{x}. \quad (20)$$

In turn, explicit conditions in terms of c and k guarantying recurrence for X are easy to obtain, as we shall further discuss in Section 3.5(i-ii). This yields explicit criteria for (18) that enables us to treat cases than were not covered previously in the literature.

It may be interesting to discuss a bit further Criterion (19). Requesting an upper-bound for the growth rate at infinity should not come as a surprise as similar assumptions are made in the literature to prevent the formation of too large particles. For instance, Doumic Jauffret and Gabriel [17] request (among other assumptions) that $\lim_{x \rightarrow \infty} xK(x)/c(x) = \infty$, which forces in our setting $\lim_{x \rightarrow \infty} c(x)/x = 0$ since we also assumed in (7) that the total rate of fragmentation K remains bounded; see Equation (13) in [17], and also Equation (11) in Balagué et al. [1]. On the other hand, imposing an upper-bound for the growth rate at 0+ may be more surprising, as on the contrary, it is often assumed in the literature that the growth for small particles should be strong enough in order to prevent shattering (see notably Equation (11) in [17] and Equation (10) in [1]). One might be further puzzled by the fact that the fragmentation rate k does not appear explicitly in (19); however, the value of the Malthus exponent λ depends of course both on c and k .

Let us also try to offer a rather informal interpretation of (19). The Feynman-Kac formula provides a representation of the growth-fragmentation semigroup $(T_t)_{t \geq 0}$ in

terms of a weighted particle (X_t, \mathcal{E}_t) , where X_t is the location of the particle at time t and \mathcal{E}_t its weight. The weight thus increases at rate $c(x)/x$ when the particle is located at x , and in this setting, the Malthus exponent λ can be interpreted as the long-time average rate of increase of the weight. Then (19) means that the weight of the particle increases more slowly than on average when the particle is either close to 0 or close to ∞ . Informally, the particle has thus a more important contribution to the Feynman-Kac formula when it stays away from 0 and from ∞ , that is essentially when it remains confined in a compact interval. And it is precisely for processes staying in compact spaces that exponential ergodicity is expected.

In this direction, we observe that Criterion (19) bears some striking similarity with Theorem 3.3 of Steinsaltz and Evans [31] and Theorem 2.6 of Kolb and Steinsaltz [22] about Yaglom convergence for linear diffusions with killing; see also Theorem 4.3 in [22] which states that convergence to a quasi-stationary distribution always holds when the killing rate at infinity is larger than the principal eigenvalue. Even though these works do not specifically consider spectral gap or exponential rate of convergence, they suggest that Criterion (19) could be rather sharp. We stress however that exponential convergence may hold even when (19) fails; for instance counter-examples can be built from Proposition 7.1 in [6].

The rest of this article is organized as follows. The two theorems are established in the next two sections, where the main ideas of the proofs are sketched first. We also gather in Section 3.5 miscellaneous comments about Theorem 1.2, notably discussing further the connection with earlier results in the literature.

We conclude this introduction by recalling that the Feynman-Kac functional \mathcal{E} defined in (11) is multiplicative, in the sense that for every $s, t \geq 0$, there is the identity

$$\mathcal{E}_{t+s} = \mathcal{E}_t \times (\mathcal{E}_s \circ \theta_t),$$

where $\mathcal{E}_s \circ \theta_t$ stands for the functional \mathcal{E}_s evaluated for the shifted path $X \circ \theta_t = X_{t+\cdot}$. In the sequel, this basic property will be often used without specific mentions, notably in combination with the Markov property.

2 Proof of Theorem 1.1

Throughout this section, we always assume that (6), (7) and (9) hold. The arguments for proving Theorem 1.1 belong to the same vein as in [6], with the difference that the role of remarkable martingales there is rather played here by supermartingales. Specifically, we shall first establish some properties of the first hitting time $H(y)$ and of its Laplace transform $L_{x,y}$, which are then applied to introduce supermartingales related to the Feynman-Kac formula. Then we shall use the latter and introduce another one-parameter family of (possibly defective) Markov process $Y^{(q)}$ by probability tilting. This yields a more direct probabilistic representation the growth-fragmentation semigroup, and analyzing the behavior of $Y^{(q)}$ via the regeneration property at return times then readily yields the conclusion.

We start by considering the motion $t \mapsto x(t)$ of a Lagrangian particle in the steady flow velocity c , viz.

$$dx(t) = c(x(t))dt,$$

which governs the dynamics of the piecewise deterministic process X between consecutive jump times, and introduce some notation in this setting that will be useful in several parts of this work. For $0 < x < y$, denote the travel time from x to y by

$$s(x, y) := \int_x^y \frac{dz}{c(z)},$$

which obviously decreases in the first variable and increases in the second one. Consider also the event $\Lambda_{x,y}$ that process X started at x reaches y before making any jump. Since $K(z)$ is the total jump rate when the process is located at z , we have

$$p(x, y) := \mathbb{P}_x(\Lambda_{x,y}) = \exp\left(-\int_0^{s(x,y)} K(x(t))dt\right) = \exp\left(-\int_x^y \frac{K(z)}{c(z)}dz\right). \quad (21)$$

This is a positive quantity which increases with x and decreases with y .

We proceed with the following uniform lower-bound for the cumulative distribution functions of first hitting times (12).

Lemma 2.1. *For every $0 < a < b$, there exists $t(a, b) \in \mathbb{R}_+$ such that*

$$\inf_{x,y \in [a,b]} \mathbb{P}_x(H(y) < t(a, b)) > 0.$$

Proof. Consider first the process X started from b . The irreducibility assumption (9) ensures that we can find two real numbers $q(a, b) \in (0, 1)$ and $h(a, b) > 0$ such that

$$\mathbb{P}_b(H(a) < h(a, b)) > q(a, b).$$

Next consider the process X started from an arbitrary point $x \in [a, b]$. By focusing on trajectories which first hit b before having any jump, then need an amount of time less than $h(a, b)$ for traveling from b to a , and finally hit $y \in [a, b]$ before having any further jump, we deduce from an application of the strong Markov property that there is the lowerbound

$$\mathbb{P}_x(H(y) < 2s(a, b) + h(a, b)) > p(a, b)^2 q(a, b) > 0.$$

This proves our claim with $t(a, b) = 2s(a, b) + h(a, b)$. □

Next, recall the notation (13) for the Laplace transform $L_{x,y}$, (14) for the Malthus exponent, and fix $x_1 > 0$ arbitrarily.

Lemma 2.2. *For every $q \geq \lambda$, the function*

$$\ell_q : x \mapsto L_{x,x_1}(q), \quad x > 0,$$

is bounded away from 0 and from ∞ on every compact interval of $(0, \infty)$.

Proof. We first consider the lower-bound and assume to start with that $q > 0$. We then have plainly

$$\ell_q(x) \geq v_q(x, x_1) := \mathbb{E}_x \left(e^{-qH(x_1)}, H(x_1) < \infty \right).$$

Fix $a > 0$ arbitrarily small and $b > 0$ arbitrarily large, with $a < x_1 < b$. Lemma 2.1 ensures the existence of $t(a, b) > 0$ and $p > 0$ such that, for all $x, y \in [a, b]$,

$$v_q(x, y) \geq \mathbb{E}_x \left(e^{-qH(y)}, H(y) < t(a, b) \right) \geq pe^{-qt(a,b)}.$$

A fortiori, $\inf_{a \leq x \leq b} \ell_q(x) > 0$.

The case when $q \leq 0$ is somewhat simpler: we can use

$$\ell_q(x) \geq \mathbb{P}_x (H(x_1) < \infty),$$

so Lemma 2.1 applies directly and yields $\inf_{a \leq x \leq b} \ell_q(x) > 0$.

We next turn our attention to the upper-bound; the argument does not depend on the sign of q . The assumption $q \geq \lambda$ and the very definition (14) entail that $L_{x_1, x_1}(q') < 1$ for every $q' > q$. Equation (16) in [6] states that then $L_{x, x_1}(q')L_{x_1, x}(q') < 1$ for every $x > 0$. By right-continuity of the functions $L_{x, y}$, we have

$$\ell_q(x) = L_{x, x_1}(q) \leq 1/L_{x_1, x}(q),$$

and since the first part of the proof shows that $x \mapsto L_{x_1, x}(q)$ is bounded away from 0 for $x \in [a, b]$, we conclude that $\sup_{a \leq x \leq b} \ell_q(x) < \infty$. \square

Theorem 4.4 in [6], which states that if $L_{x, x}(\lambda) = 1$, then the process $(e^{-\lambda t} \ell_\lambda(X_t) \mathcal{E}_t)_{t \geq 0}$ is a martingale, is a cornerstone of the probabilistic approach which is developed there. Here is a version of the latter in terms supermartingales.

Lemma 2.3. *For every $q \geq \lambda$, the process*

$$\mathcal{S}_t^{(q)} := e^{-qt} \ell_q(X_t) \mathcal{E}_t, \quad t \geq 0$$

is a \mathbb{P}_x -supermartingale for every $x > 0$.

Proof. Write

$$N_t := \#\{0 < s \leq t : X_s = x_1\}, \quad t \geq 0$$

for the process which counts the number of visits of X to x_1 as time passes, and

$$R_n := \inf\{t > 0 : N_t = n\}, \quad n \geq 1$$

for the instant when X returns to x_1 for the n -th time. Write also $(\mathcal{F}_t)_{t \geq 0}$ for the natural filtration of X and recall that the return times R_n are (\mathcal{F}_t) -stopping times. Further, writing $\mathcal{G}_n := \mathcal{F}_{R_n}$, we know that for every $t \geq 0$, $N_t + 1$ is a (\mathcal{G}_n) -stopping time and the first return to x_1 after time t can be expressed as

$$D_t := \inf\{s > t : X_s = x_1\} = R_{N_t+1}.$$

On the one hand, we see from the Markov property at time t and the definition of the function ℓ_q in Lemma 2.2 that for every $x > 0$,

$$\mathcal{S}_t^{(q)} = \mathbb{E}_x \left(e^{-qD_t} \mathcal{E}_{D_t} \mathbb{1}_{\{D_t < \infty\}} \mid \mathcal{F}_t \right) = \mathbb{E}_x \left(e^{-qR_{N_t+1}} \mathcal{E}_{R_{N_t+1}} \mathbb{1}_{\{R_{N_t+1} < \infty\}} \mid \mathcal{F}_t \right). \quad (22)$$

On the other hand, the strong Markov property and the fact that $L_{x_1, x_1}(q) \leq 1$ (from the definition (14) and our assumption $q \geq \lambda$) entail that for every $x > 0$,

$$e^{-qR_n} \mathcal{E}_{R_n} \mathbb{1}_{\{R_n < \infty\}}, \quad n \geq 1$$

is a \mathbb{P}_x -supermartingale in the filtration $(\mathcal{G}_n)_{n \geq 1}$. Since for $s \leq t$, $N_s + 1 \leq N_t + 1$ are two (\mathcal{G}_n) -stopping times, it follows from the optional sampling theorem for nonnegative supermartingales that

$$\mathbb{E}_x \left(e^{-qR_{N_t+1}} \mathcal{E}_{R_{N_t+1}} \mathbb{1}_{\{R_{N_t+1} < \infty\}} \mid \mathcal{G}_{N_s+1} \right) \leq e^{-qR_{N_s+1}} \mathcal{E}_{R_{N_s+1}} \mathbb{1}_{\{R_{N_s+1} < \infty\}}.$$

Then on both sides, take the conditional expectation given \mathcal{F}_s , which is a sub-algebra of $\mathcal{F}_{D_s} = \mathcal{G}_{N_s+1}$. We get from (22) (applied at time s rather than t)

$$\mathbb{E}_x \left(e^{-qR_{N_t+1}} \mathcal{E}_{R_{N_t+1}} \mathbb{1}_{\{R_{N_t+1} < \infty\}} \mid \mathcal{F}_s \right) \leq \mathcal{S}_s^{(q)}.$$

We conclude the proof by using once again (22) and the so-called tower property of conditional expectations on the left-hand. \square

The supermartingale $\mathcal{S}^{(q)}$ in Lemma 2.3 enables us to introduce a possibly defective (i.e. possibly with finite lifetime ζ) càdlàg Markov process $Y^{(q)} = (Y_t^{(q)})_{0 \leq t < \zeta}$ with distribution denoted by $\mathbb{P}^{(q)}$ as follows. For every $t \geq 0$ and every nonnegative functional F defined on Skorokhod's space $\mathcal{D}_{[0, t]}$ of càdlàg paths $\omega : [0, t] \rightarrow (0, \infty)$, one sets

$$\mathbb{E}_x^{(q)} [F((Y_s^{(q)})_{0 \leq s \leq t}), \zeta > t] = \frac{1}{\ell_q(x)} \mathbb{E}_x [\mathcal{S}_t^{(q)} F((X_s)_{0 \leq s \leq t})], \quad x > 0.$$

We stress that the distribution of $(Y_s^{(q)})_{0 \leq s \leq t}$ under the conditional law $\mathbb{P}_x^{(q)}(\cdot \mid \zeta > t)$ is absolutely continuous with respect to that of $(X_s)_{0 \leq s \leq t}$ under \mathbb{P}_x , and as a consequence, $Y^{(q)}$ inherits irreducibility from (9).

Lemma 2.4. (i) Suppose that the assumption (i) of Theorem 1.1 holds. Then $\rho \geq \lambda$.

(ii) Suppose further that the assumption (ii) of Theorem 1.1 also holds, and set $Y := Y^{(\rho)}$. Then there exists $b > 0$ and $s > 0$ sufficiently large, such that

$$\liminf_{t \rightarrow \infty} \mathbb{P}_b^{(\rho)} (Y_r = b \text{ for some } r \in [t, t + s]) > 0.$$

Proof. (i) Suppose $\rho < \lambda$ and pick any $q \in (\rho, \lambda)$. On the one hand, since $q > \rho$, assumption (i) of Theorem 1.1 entails that

$$\int_0^\infty e^{-qt} T_t f(x_1) dt < \infty$$

for some $x_1 > 0$ and some continuous function $f : (0, \infty) \rightarrow \mathbb{R}_+$ with $f \not\equiv 0$. On the other hand, as $q < \lambda$, we have $L_{x_1, x_1}(q) \geq 1$, and the assertion above contradicts the fact that then for every $f : (0, \infty) \rightarrow [0, \infty)$ continuous with $f \not\equiv 0$, we must have

$$\int_0^\infty e^{-qt} T_t f(x) dt = \infty,$$

see Proposition 3.3 in [6].

(ii) Thanks to (i), we may now take $q = \rho$. Note from the very definition of Y that the Feynman-Kac formula (11) can be translated as follows: for every $x > 0$ and every continuous and compactly supported function $f : (0, \infty) \rightarrow \mathbb{R}$, there is the identity

$$\frac{e^{-\rho t}}{x\ell(x)} T_t f(x) = \mathbb{E}_x^{(\rho)} \left(\frac{f(Y_t)}{Y_t \ell(Y_t)}, \zeta > t \right), \quad (23)$$

with $\ell := \ell_\rho$. Combining this with assumption (ii) of Theorem 1.1 and the fact that, thanks to Lemma 2.2, ℓ remains bounded away from 0 on compact intervals of $(0, \infty)$, we deduce that

$$\liminf_{t \rightarrow \infty} \mathbb{P}_{x_2}^{(\rho)}(Y_t \in [a, b], \zeta > t) > 0, \quad (24)$$

for any $0 < a < b$ such that $\text{Supp}(g) \subseteq [a, b]$.

On the other hand, recall the notation from the first paragraph in the proof of Lemma 2.1 and notably the definition of $p(x, y)$ in (21), and for every $x \in [a, b]$, consider the probability $p^{(\rho)}(x)$ that process Y started at x reaches b at time $s(x, b)$, before dying or making any jump. The obvious bound $\mathcal{S}_t^{(\rho)} \geq \ell(X_t) \exp(-\rho^+ t)$ (where ρ^+ stands for the positive part of ρ) yields

$$\begin{aligned} p^{(\rho)}(x) &= \frac{1}{\ell(x)} \mathbb{E}_x \left(\mathcal{S}_{s(x,b)}^{(\rho)}, X \text{ has no jump before time } s(x, b) \right) \\ &\geq \frac{\ell(b)}{\ell(x)} \exp(-\rho^+ s(x, b)) p(x, b) \\ &\geq \frac{\ell(b)}{\sup_{[a,b]} \ell} \exp(-\rho^+ s(a, b)) p(a, b). \end{aligned}$$

Again by Lemma 2.2 and the first paragraph in the proof of Lemma 2.1, the right-hand above is positive, hence

$$\inf_{[a,b]} p^{(\rho)} > 0.$$

We now see from the Markov property of Y that

$$\mathbb{P}_{x_2}^{(\rho)}(Y_r = b \text{ for some } r \in [t, t + s(a, b)]) \geq \inf_{[a,b]} p^{(\rho)} \times \mathbb{P}_{x_2}^{(\rho)}(Y_t \in [a, b], \zeta > t),$$

and then, combining with (24), that

$$\liminf_{t \rightarrow \infty} \mathbb{P}_{x_2}^{(\rho)}(Y_r = b \text{ for some } r \in [t, t + s(a, b)]) > 0.$$

Recalling that Y is irreducible and applying the strong Markov property at time $H_Y(b)$ completes the proof. \square

We readily deduce from Lemma 2.4 the following

Corollary 2.5. *Under the assumptions (i) and (ii) of Theorem 1.1, the Markov process $Y = Y^{(\rho)}$ is point-recurrent and positive, that is*

$$\mathbb{E}_x^{(\rho)}(H_Y(y)) < \infty \quad \text{for every } x, y > 0,$$

where $H_Y(y) := \inf\{t \in (0, \zeta) : Y_t = y\}$ stands for the first hitting time of y by the process Y , with the usual convention that $\inf \emptyset = \infty$.

Proof. We start by observing from Lemma 2.4(ii) that b must be a recurrent state for Y , i.e.

$$\mathbb{P}_b^{(\rho)}(H_Y(b) < \infty) = 1.$$

Indeed, if we had $\mathbb{P}_b^{(\rho)}(H_Y(b) = \infty) > 0$, then by the strong Markov property at time $H_Y(b)$, the total number of visits to b would have a geometric distribution under $\mathbb{P}_b^{(\rho)}$, and in particular would be finite $\mathbb{P}_b^{(\rho)}$ -a.s., in contradiction with Lemma 2.4(ii).

By the strong Markov property at return times, the set of passage times at b , $\{t \geq 0 : Y_t = b\}$ is regenerative, i.e. it can be expressed as the set of partial sums of a sequence of independent copies of the variable $H_Y(b)$. Its so-called renewal function is given by

$$U(t) := \mathbb{E}_b^{(\rho)}(\text{Card}\{r \leq t : Y_r = b\}), \quad t \geq 0.$$

Choose $s > 0$ as in Lemma 2.4(ii), and observe that

$$U(t+s) - U(t) \geq \mathbb{P}_b^{(\rho)}(\exists r \in (t, t+s] : Y_r = b).$$

Recall from Blackwell's renewal theorem that

$$U(t+s) - U(t) \sim s/\mathbb{E}_b^{(\rho)}(H_Y(b)) \quad \text{as } t \rightarrow \infty,$$

and Lemma 2.4(ii) therefore implies that $\mathbb{E}_b^{(\rho)}(H_Y(b)) < \infty$. Since Y is irreducible, the conclusion of the statement follows. \square

We now have all the ingredients needed to prove Theorem 1.1. Indeed, Corollary 2.5 implies that under the assumptions (i) and (ii) of Theorem 1.1, Y cannot be defective, i.e. $\mathbb{P}_x^{(\rho)}(\zeta = \infty) = 1$. That is, equivalently, $\mathbb{E}_x(\mathcal{S}_t^{(\rho)}) = \ell(x)$ for all $t \geq 0$, and since a supermartingale with a constant expectation must be a martingale, $\mathcal{S}^{(\rho)}$ is a \mathbb{P}_x -martingale for every $x > 0$.

Then, point-recurrence for Y gives for every $x > 0$

$$1 = \lim_{t \rightarrow \infty} \mathbb{P}_x^{(\rho)}(H_Y(x) \leq t) = \lim_{t \rightarrow \infty} \frac{1}{\ell(x)} \mathbb{E}_x[\mathcal{S}_t^{(\rho)}, H(x) \leq t].$$

On the other hand, the martingale property of $\mathcal{S}^{(\rho)}$ under \mathbb{P}_x and the optional sampling theorem yield

$$\mathbb{E}_x[\mathcal{S}_t^{(\rho)}, H(x) \leq t] = \mathbb{E}_x[\mathcal{S}_{H(x)}^{(\rho)}, H(x) \leq t] = \ell(x) \mathbb{E}_x[e^{-\rho H(x)} \mathcal{E}_{H(x)}, H(x) \leq t].$$

We deduce by monotone convergence that

$$L_{x,x}(\rho) = \mathbb{E}_x[e^{-\rho H(x)} \mathcal{E}_{H(x)}, H(x) < \infty] = 1,$$

which implies both that $\lambda = \rho$ and the first condition in (15) holds.

We now see that the function $\ell = \ell_\rho = \ell_\lambda$ here is the same as that in Section 4 of [6], the martingale $\mathcal{S}^{(\rho)}$ coincides with the martingale \mathcal{M} there, and finally, the Markov process Y here is the same as that in Section 5 of [6]. Recall from Corollary 2.5 that Y is positive recurrent, and we conclude from Lemma 5.2(i) in [6] that the right-derivative of the Laplace transform $L_{x,x}(\cdot)$ at λ is necessarily finite, which is the second condition in (15). The proof of Theorem 1.1 is now completed.

3 Proof of Theorem 1.2

Our goal in this section is to check that, when the assumptions of Theorem 1.2 are fulfilled, then (17) holds, as the exponential convergence then follows from Theorem 1.1 in [6]. This will be achieved in three main steps.

To start with, we work on a compact interval $[a, b]$, where $0 < a < b$ are given, and consider the first exit-time

$$\sigma(a, b) := \inf\{t > 0 : X_t \notin [a, b]\}.$$

We first discuss irreducibility for the process killed when exiting from $[a, b]$, which is a necessary preamble for the rest of our analysis.

We then verify that the Krein-Rutman theorem can be applied in this compact setting, by analyzing the trajectories of X . This yields a principal eigenvalue $\rho_{a,b}$ for the system where particles are killed when exiting $[a, b]$, and a corresponding positive eigenfunction $h_{a,b}$. We then construct useful martingales from the latter, which in turn will enable us to compute certain expectations by application of optional sampling.

For the next step, we fix the upper-boundary point b large enough, and let the lower-boundary point a tend to $0+$. We shall establish the existence of $\gamma < \lambda$ and a non-degenerate function $g : (0, b] \rightarrow \mathbb{R}_+$ such that the process $g(X_t) \mathcal{E}_t e^{-\gamma t} \mathbb{1}_{\{t < H(b)\}}$ is a supermartingale.

Finally, (17) is established by putting the pieces together. In short, we pick a large enough interval $[a, b]$, $q < \lambda$ close enough to the Malthus exponent λ , decompose the excursion of the process away from its starting point at certain first-exit times, and estimate the various pieces using the preceding steps.

3.1 Irreducibility in compact intervals

We start by addressing the slightly technical question of irreducibility. Even when X is irreducible, it may happen that for some $0 < a < b$, there exist two states $x < y$ both

in (a, b) such that no path of X started from y can reach x without exiting first from $[a, b]$. When this occurs, the process killed at time $\sigma(a, b)$ is then no longer irreducible, and this creates an obstacle for our analysis. Recall however that the probability that X started from x follows the flow velocity without having jumps until it reaches y is always positive, so the problem can only arise when the starting point is larger than the target.

We call an interval (a, b) with $0 < a < b$ good, if the process killed at time $\sigma(a, b)$ remains irreducible, in the sense that

$$\mathbb{P}_x(H(y) < \sigma(a, b)) > 0 \quad \text{for all } x, y \in (a, b). \quad (25)$$

In the next statement, we first justify the criterion for irreducibility in terms of the kernel \bar{k} which has been asserted in the Introduction, and then argue that we can always find good intervals (a, b) with $a > 0$ as small as we wish and b as large as we wish.

Lemma 3.1. (i) *The conditions (9) and (10) are equivalent.*

(ii) *Assuming that either (9) or (10) hold, for every $\varepsilon \in (0, 1)$, there exists a good interval (a, b) with $a < \varepsilon$ and $b > 1/\varepsilon$.*

Proof. (i) We call a pair (β, α) with $0 < \alpha < \beta$ a possible jump if α belongs to the two-sided support of $\bar{k}(\beta, \cdot)$, in the sense that

$$\int_{\alpha}^{\alpha+\eta} \bar{k}(\beta, y) dy > 0 \quad \text{and} \quad \int_{\alpha-\eta}^{\alpha} \bar{k}(\beta, y) dy > 0 \quad \text{for all } 0 < \eta < \alpha.$$

Note that thanks to (7), the condition above remains fulfilled when we replace α , respectively β , by any α' close enough to α , respectively any β' close enough to β , so the set of possible jump pairs is open in $\{(y, x) : 0 < x < y\}$.

Because the sample paths of X increase between consecutive jump times and the jump rates of X are given by \bar{k} , we readily deduce from the observation above that if (β, α) is a possible jump pair, then the probability under \mathbb{P}_{β} that X makes a single jump straddling the interval (α, β) before hitting α , is strictly positive.

Now assume that (10) holds, so that the family of intervals (α, β) associated to possible jumps pairs (β, α) form a covering of $(0, \infty)$, and consider any $0 < x < y$. By compactness, there exists a finite covering of $[x, y]$ by intervals associated to possible jumps pairs, and it is now easy to check from the strong Markov property that $\mathbb{P}_y(H_x < \infty) > 0$. So (9) holds.

Conversely, if (10) fails for some $x > 0$, the process X can never make a jump across x . Because paths increase between two consecutive jump times, we see that the process started for any $y > x$ will never visit x , and hence (9) also fails.

(ii) We first observe that the argument in the proof of (i) actually yields a slightly better result. Namely, if (β, α) is a possible jump pair, then for any $\alpha < x < y < \beta$, the probability under \mathbb{P}_y that X makes a single jump straddling the interval $[x, y]$ before hitting x and also stays in the interval (α, β) up to time H_x , is strictly positive. Using

again the fact that on any finite time interval, the probability that X follows the flow velocity without having jump is positive, we now readily see that (25) holds for $a = \alpha$ and $b = \beta$. In other words, if (β, α) is a possible jump pair, then the interval (α, β) is good. As intervals associated to possible jump pair form a covering of $(0, \infty)$, there exists a finite covering of the compact interval $[\varepsilon, 1/\varepsilon]$, say $\{(\alpha_i, \beta_i) : i = 1, \dots, n\}$, and it is then immediate to check from the strong Markov property that $a = \min \alpha_i$ and $b = \max \beta_i$ fulfill the requirements of the statement. \square

Throughout the rest of this section, we henceforth assume that (6), (7) and (9) hold.

3.2 Applying the Krein-Rutman theorem in a compact interval

We next consider the Banach space $\mathcal{C}_0[a, b)$ of continuous functions $f : [a, b) \rightarrow \mathbb{R}$ with $f(b) = 0$, endowed with the usual norm $\|f\| = \sup_{x \in [a, b)} |f(x)|$. The reason for imposing $f(b) = 0$ is that the two boundary points of the interval $[a, b)$ have a different status for the Markov process killed at time $\sigma(a, b)$. Specifically, a is an entrance boundary, in the sense that the process started at a then stays in $[a, b)$ for a strictly positive amount of time \mathbb{P}_a -a.s., whereas b is an exit boundary, meaning that the process started at b leaves $[a, b)$ instantaneously \mathbb{P}_b -a.s. We refer to [21] on pages 108 and 130-131 for more about these notions introduced by Feller in the setting of diffusion processes. We do not assume right now that the interval (a, b) is good, but this assumption will of course be essential in a later part of our analysis.

Recall our assumption (6) and define $q_c := 1 + \sup_{x>0} c(x)/x$, so that

$$\mathcal{E}_t e^{-tq_c} \leq e^{-t} \quad \text{for all } t \geq 0.$$

We introduce for every bounded measurable function $f : [a, b) \rightarrow \mathbb{R}$

$$U_{a,b}f(x) := \mathbb{E}_x \left(\int_0^{\sigma(a,b)} f(X_t) \mathcal{E}_t e^{-tq_c} dt \right), \quad x \in [a, b).$$

Lemma 3.2. *The operator $U_{a,b}$ maps $\mathcal{C}_0[a, b)$ into itself. More precisely the family of functions $\{U_{a,b}f : \|f\| \leq 1\}$ is equicontinuous.*

Proof. We first note that $U_{a,b}f(b) = 0$ (since $\sigma(a, b) = 0$, \mathbb{P}_b -a.s.), and also that $U_{a,b}$ is a contraction, i.e. $\|U_{a,b}f\| \leq \|f\|$. Then recall that for $0 < x < y$, $s(x, y)$ denotes the travel time from x to y for a Lagrangian particle driven by the flow velocity c , and that $\Lambda_{x,y}$ stands for the event that X starts from x and reaches y before making any jump. Observe that on that event, we have

$$\ln \mathcal{E}_{s(x,y)} = \int_0^{s(x,y)} \frac{c(x(s))}{x(s)} ds = \ln y - \ln x.$$

Take $a \leq x < y \leq b$ and any f with $\|f\| \leq 1$. On the event $\Lambda_{x,y}$, we decompose the trajectory at time $s(x, y)$, set

$$X'_t := X_{s(x,y)+t} \quad , \quad \sigma'(a, b) := \inf\{t > 0 : X'_t \notin [a, b]\} \quad , \quad \mathcal{E}'_t := \exp\left(\int_0^t \frac{c(X'_s)}{X'_s} ds\right),$$

and get the bounds

$$\left| \int_0^{\sigma(a,b)} f(X_t) \mathcal{E}_t e^{-tq_c} dt - \frac{y}{x} e^{-q_c s(x,y)} \int_0^{\sigma'(a,b)} f(X'_t) \mathcal{E}'_t e^{-tq_c} dt \right| \leq s(x, y).$$

On the complementary event $\Lambda_{x,y}^c$, we use the trivial bounds

$$\left| \int_0^{\sigma(a,b)} f(X_t) \mathcal{E}_t e^{-tq_c} dt \right| \leq 1.$$

The Markov property shows that conditionally on the event $\Lambda_{x,y}$, X' has the law \mathbb{P}_y ; we deduce that $U_{a,b}f(x)$ is bounded from below by

$$-s(x, y) + \frac{y}{x} e^{-q_c s(x,y)} U_{a,b}f(y) \mathbb{P}_x(\Lambda_{x,y}) - (1 - \mathbb{P}_x(\Lambda_{x,y}))$$

and from above by

$$s(x, y) + \frac{y}{x} e^{-q_c s(x,y)} U_{a,b}f(y) \mathbb{P}_x(\Lambda_{x,y}) + 1 - \mathbb{P}_x(\Lambda_{x,y}).$$

Recalling that $|U_{a,b}f(y)| \leq 1$, this yields the inequality

$$|U_{a,b}f(x) - U_{a,b}f(y)| \leq s(x, y) + \left| \frac{y}{x} e^{-q_c s(x,y)} - 1 \right| + 2(1 - \mathbb{P}_x(\Lambda_{x,y})).$$

On the one hand, (21) shows that $1 - \mathbb{P}_x(\Lambda_{x,y})$ converges to 0 as $y - x \rightarrow 0+$, uniformly for $a \leq x < y \leq b$. On the other hand, it is easily checked that the same holds for $s(x, y)$ (because the flow velocity c is bounded away from 0 on $[a, b]$). This entails that

$$\sup_{x,y \in [a,b], |y-x| < \varepsilon} |U_{a,b}f(x) - U_{a,b}f(y)| \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0+,$$

uniformly for f with $\|f\| \leq 1$, and our claim is proven. \square

The subspace $\mathcal{C}_0^+[a, b]$ of nonnegative functions in $\mathcal{C}_0[a, b]$ is a reproducing cone, that is $\mathcal{C}_0^+[a, b]$ is a closed convex set which is stable by multiplication by nonnegative constants, and $f = f^+ - f^-$ is a decomposition of a generic function $f \in \mathcal{C}_0[a, b]$ as the difference of two functions in $\mathcal{C}_0^+[a, b]$. Obviously, the operator $U_{a,b}$ positive, i.e. maps the cone $\mathcal{C}_0^+[a, b]$ into itself. Recall also from the Riesz-Markov representation theorem that any linear functional on $\mathcal{C}_0[a, b]$ which is positive (in the sense that it maps $\mathcal{C}_0^+[a, b]$ into \mathbb{R}_+), can be represented by a finite Borel measure on $[a, b]$ that has no atom at b . Plainly,

the dual operator $U_{a,b}^*$ maps any such measure, say m , to another finite measure $U_{a,b}^*m$ on $[a, b]$ without atom at b , via the identity

$$\langle U_{a,b}^*m, f \rangle = \langle m, U_{a,b}f \rangle \quad \text{for all } f \in \mathcal{C}_0^+[a, b].$$

The Krein-Rutman theorem yields the first milestone for the proof of Theorem 1.2. Recall that the notion of good interval has been introduced in (25).

Proposition 3.3. *Let (a, b) be a good interval. Then:*

- (i) *the spectral radius $r(a, b)$ of $U_{a,b}$ is positive,*
- (ii) *there exist a function $h_{a,b} \in \mathcal{C}_0^+[a, b]$ with $h_{a,b}(x) > 0$ for every $x \in [a, b)$ and a finite Borel measure $\nu_{a,b}$ on $[a, b]$ with $\nu_{a,b}(\{b\}) = 0$, such that*

$$\langle \nu_{a,b}, h_{a,b} \rangle = 1, \quad U_{a,b}h_{a,b} = r(a, b)h_{a,b}, \quad \text{and } U_{a,b}^*\nu_{a,b} = r(a, b)\nu_{a,b},$$

- (iii) *the spectral gap holds, i.e. any $r \neq r(a, b)$ in the spectrum of $U_{a,b}$ has $|r| < r(a, b)$.*

Proof. We just need to verify the requirements of Theorem 19.5 in Deimling [13]. First, it follows from Lemma 3.2 by the Arzelà-Ascoli theorem that the operator $U_{a,b}$ is compact. Second, observe from the strong Markov property that for every $f \in \mathcal{C}_0^+[a, b]$ with $f \not\equiv 0$, (25) ensures that $U_{a,b}f(x) > 0$ for every $x \in [a, b)$. Then, $\langle m, U_{a,b}f \rangle > 0$ for every Borel measure on $[a, b]$ with no atom at b , which is Condition (8) on page 235 of [13]. \square

Proposition 3.3 enables us to introduce the following useful martingale.

Lemma 3.4. *Set $\rho_{a,b} := q_c - 1/r(a, b)$. The process*

$$\mathcal{M}_{a,b}(t) := \mathbb{1}_{\{t < \sigma(a,b)\}} h_{a,b}(X_t) \mathcal{E}_t e^{-t\rho_{a,b}}, \quad t \geq 0$$

is a \mathbb{P}_x -martingale for every $x \in [a, b]$.

Proof. Recall that $(\mathcal{F}_t)_{t \geq 0}$ denotes the natural filtration of X . By the Markov property, we can express the uniformly integrable martingale

$$M_t := \mathbb{E}_x \left(\int_0^{\sigma(a,b)} h_{a,b}(X_s) \mathcal{E}_s e^{-sq_c} ds \mid \mathcal{F}_t \right)$$

as

$$M_t = \int_0^{t \wedge \sigma(a,b)} h_{a,b}(X_s) \mathcal{E}_s e^{-sq_c} ds + \mathbb{1}_{\{t < \sigma(a,b)\}} \mathcal{E}_t e^{-tq_c} U_{a,b}h_{a,b}(X_t).$$

From the identity $U_{a,b}h_{a,b} = r(a, b)h_{a,b}$ and stochastic calculus, we deduce that

$$\mathcal{M}_{a,b}(t) = \mathcal{M}_{a,b}(0) + \frac{1}{r(a, b)} \int_0^t e^{s/r(a,b)} dM_s.$$

The stochastic integral in the right-hand side is a \mathbb{P}_x -martingale since both the integrand and the martingale M are bounded on $[0, t]$. \square

We now arrive at a second milestone of the proof of Theorem 1.2.

Proposition 3.5. *Take any good interval (a, b) . Then for all $x, y \in (a, b)$, there is the identity*

$$\mathbb{E}_x \left(\mathcal{E}_{H(y)} e^{-\rho_{a,b} H(y)}, H(y) < \sigma(a, b) \right) = h_{a,b}(x) / h_{a,b}(y).$$

Proof. Perhaps, it could be tempting to try to derive the statement from Lemma 3.4 by an application of optional sampling to the martingale $\mathcal{M}_{a,b}$ and the stopping time $H(y)$. Note however that this would not be legitimate as the latter is not bounded and the martingale $\mathcal{M}_{a,b}$ is not uniformly integrable (actually, $\mathcal{M}_{a,b}(\infty) = 0$ \mathbb{P}_x -a.s.).

We use the martingale $\mathcal{M}_{a,b}$ to introduce a new Markov process $(Z_t)_{t \geq 0}$ on $[a, b]$ with law $\mathbf{P}_x^{(a,b)}$, by setting

$$\mathbf{E}_x^{(a,b)} [F((Z_s)_{0 \leq s \leq t})] = \frac{1}{h_{a,b}(x)} \mathbb{E}_x [\mathcal{M}_{a,b}(t) F((X_s)_{0 \leq s \leq t})],$$

where F stands for a generic nonnegative functional defined on Skorokhod's space $\mathcal{D}_{[0,t]}$. We also write $H_Z(y) = \inf\{t > 0 : Z_t = y\}$ for the first hitting time of y by Z .

The purpose of introducing the process Z is that there are the identities

$$\begin{aligned} \mathbf{P}_x^{(a,b)}(H_Z(y) \leq t) &= \frac{1}{h_{a,b}(x)} \mathbb{E}_x [\mathcal{M}_{a,b}(t), H(y) \leq t] \\ &= \frac{1}{h_{a,b}(x)} \mathbb{E}_x [\mathcal{M}_{a,b}(H(y)), H(y) \leq t] \\ &= \frac{h_{a,b}(y)}{h_{a,b}(x)} \mathbb{E}_x [\mathcal{E}_{H(y)} e^{-\rho_{a,b} H(y)}, H(y) \leq t \wedge \sigma(a, b)], \end{aligned}$$

where the second line stems from Lemma 3.4 and Doob's optional sampling theorem applied at the bounded stopping time $t \wedge H(y)$, and the third from the definition of $\mathcal{M}_{a,b}$.

Our claim now amounts to checking that the left-hand side above converges to 1 as $t \rightarrow \infty$, that is that $H_Z(y) < \infty$, $\mathbf{P}_x^{(a,b)}$ -a.s. In this direction, we shall prove that Z is point recurrent using the full strength of Proposition 3.3 (whereas for Lemma 3.4, we only needed to know that $U_{a,b} h_{a,b} = r(a, b) h_{a,b}$).

Introduce the probability measure $m_{a,b}(dx) := h_{a,b}(x) \nu_{a,b}(dx)$ on $[a, b]$. We claim that $m_{a,b}$ is a stationary distribution for Z . In this direction, define first for every bounded measurable function $f : [a, b] \rightarrow \mathbb{R}$ and every $q > 0$

$$V^q f(x) := \mathbb{E}_x \left(\int_0^{\sigma(a,b)} f(X_t) \mathcal{E}_t e^{-(q+\rho_{a,b})t} dt \right), \quad x \in [a, b].$$

In particular, $V^0 = U_{a,b}$ and the resolvent equation reads $V^q = U_{a,b} - q U_{a,b} V^q$. With

this notation at hand, and recalling from Lemma 3.4 that $\rho_{a,b} = q_c - 1/r(a,b)$, we have

$$\begin{aligned} \mathbf{E}_{m_{a,b}}^{(a,b)} \left(\int_0^\infty e^{-qt} f(Z_t) dt \right) &= \mathbf{E}_{\nu_{a,b}} \left(\int_0^{\sigma(a,b)} e^{-(q+q_c-1/r(a,b))t} f(X_t) h_{a,b}(X_t) \mathcal{E}_t dt \right) \\ &= \langle \nu_{a,b}, V^{q-1/r(a,b)}(fh_{a,b}) \rangle \\ &= \langle \nu_{a,b}, U_{a,b}(fh_{a,b}) \rangle - (q-1/r(a,b)) \langle \nu_{a,b}, U_{a,b} V^{q-1/r(a,b)}(fh_{a,b}) \rangle \\ &= r(a,b) \langle \nu_{a,b}, fh_{a,b} \rangle - (r(a,b)q-1) \langle \nu_{a,b}, V^{q-1/r(a,b)}(fh_{a,b}) \rangle, \end{aligned}$$

where we used the resolvent equation at the third line, and that $U_{a,b}^* \nu_{a,b} = r(a,b) \nu_{a,b}$ at the fourth. This yields

$$\mathbf{E}_{m_{a,b}}^{(a,b)} \left(\int_0^\infty e^{-qt} f(Z_t) dt \right) = q^{-1} \langle \nu_{a,b}, fh_{a,b} \rangle = q^{-1} \langle m_{a,b}, f \rangle,$$

showing that indeed $m_{a,b}$ is a stationary law for Z .

The existence of a stationary law for Z now readily entails point recurrence. Take any $x \in (a,b)$; plainly Z inherits irreducibility from (25), and in particular $\mathbf{P}_{m_{a,b}}^{(a,b)}(H_Z(x) < \infty) := \alpha > 0$. By stationarity, we have for every $s > 0$ that

$$\mathbf{P}_{m_{a,b}}^{(a,b)}(\exists t \geq s : Z_t = x) \geq \alpha,$$

and *a fortiori*

$$\mathbf{P}_{m_{a,b}}^{(a,b)}(\sup\{t \geq 0 : Z_t = x\} = \infty) \geq \alpha.$$

We conclude from the strong Markov property that $H_Z(x) < \infty$, $\mathbf{P}_x^{(a,b)}$ -a.s., which is what we needed to verify. \square

Proposition 3.5 enables us to compare the eigenvalues $\rho_{a,b}$ for nested good intervals, and also with the Malthus exponent λ .

Lemma 3.6. *We have:*

(i) *Let (a,b) and (a',b') two good intervals with $a < a' < b' < b$. Then*

$$\rho_{a',b'} < \rho_{a,b} < \lambda.$$

(ii) $\lambda = \sup\{\rho_{a,b} : (a,b) \text{ is a good interval}\}$.

Proof. (i) Since $\sigma(a',b') \leq \sigma(a,b)$ and the inequality is strict with positive \mathbb{P}_x -probability for any $x \in [a',b']$, it follows from Proposition 3.5 that

$$\mathbb{E}_x \left(\mathcal{E}_{H(x)} e^{-\rho_{a',b'} H(x)}, H(x) < \sigma(a,b) \right) > 1 = \mathbb{E}_x \left(\mathcal{E}_{H(x)} e^{-\rho_{a,b} H(x)}, H(x) < \sigma(a,b) \right).$$

This forces $\rho_{a',b'} < \rho_{a,b}$. Similarly,

$$\mathbb{E}_x(\mathcal{E}_{H(x)} e^{-\lambda H(x)}, H(x) < \sigma(a,b)) < \mathbb{E}_x(\mathcal{E}_{H(x)} e^{-\lambda H(x)}, H(x) < \infty) \leq 1,$$

where the second inequality follows from the right-continuity of the Laplace transform $L_{x,x}$ and the definition (14) of the Malthus exponent. This yields $\rho_{a,b} < \lambda$.

(ii) Set $\bar{\rho} := \sup\{\rho_{a,b} : (a,b) \text{ is a good interval}\}$. Then, by monotone convergence and Lemma 3.1, we have

$$\begin{aligned} & \mathbb{E}_x \left(\mathcal{E}_{H(x)} e^{-\bar{\rho}H(x)}, H(x) < \infty \right) \\ &= \sup\{ \mathbb{E}_x \left(\mathcal{E}_{H(x)} e^{-\bar{\rho}H(x)}, H(x) < \sigma(a,b) \right) : (a,b) \text{ is a good interval} \} \\ &\leq \sup\{ \mathbb{E}_x \left(\mathcal{E}_{H(x)} e^{-\rho_{a,b}H(x)}, H(x) < \sigma(a,b) \right) : (a,b) \text{ is a good interval} \}, \end{aligned}$$

showing that $L_{x,x}(\bar{\rho}) \leq 1$. Hence $\bar{\rho} \geq \lambda$, and the converse inequality follows from (i). \square

3.3 Letting the lower-boundary point go to 0

We now fix the upper-boundary point b and let the lower-boundary point a tend to 0. Note that due to the absence of positive jumps, the identity $\lim_{a \rightarrow 0} \sigma(a,b) = H(b)$ holds \mathbb{P}_x -a.s. for all $x < b$.

Proposition 3.7. *Assume $\limsup_{x \rightarrow 0} c(x)/x < \lambda$. Then the following hold:*

- (i) *There exist a good interval (a,b) such that $\sup_{x \leq a} c(x)/x < \rho_{a,b}$.*
- (ii) *For every $a' \in (a,b)$ with $\sup_{x \leq a'} c(x)/x \leq \rho_{a,b}$ and every $b'' \in (a',b)$ sufficiently close to b , there exists then $\gamma < \lambda$ with*

$$\mathbb{E}_{a'} \left(\mathcal{E}_{H(a')} e^{-\gamma H(a')}, H(a') < H(b'') \right) \in (0, 1].$$

Proof. (i) We first use Lemma 3.6 and the assumption that $\limsup_{x \rightarrow 0} c(x)/x < \lambda$ to find a good interval (a',b') with $\limsup_{x \rightarrow 0} c(x)/x < \rho_{a',b'}$. Then consider

$$a'' := \inf\{x > 0 : c(x)/x \geq \rho_{a',b'}\} \in (0, \infty],$$

and choose a good interval (a,b) with $0 < a < a' \wedge a''$ and $b > b'$. By Lemma 3.6(i), we have then $\rho_{a,b} > \rho_{a',b'}$, and *a fortiori* $c(x)/x < \rho_{a,b}$ for all $x \leq a$.

(ii) The irreducibility (25) entails that $\mathbb{P}_{a'}(H(a') < H(b'')) > 0$ provided that $b'' \in (a',b)$ is chosen close enough to b . Then consider the convex and nonincreasing function $\Psi : \mathbb{R} \rightarrow (0, \infty]$ defined by

$$\Psi(q) := \mathbb{E}_{a'} \left(\mathcal{E}_{H(a')} e^{-qH(a')}, H(a') < H(b'') \right).$$

Since $\rho_{a,b} < \lambda$ by Lemma 3.6, we may pick $r \in (\rho_{a,b}, \lambda)$. We shall check that $\Psi(r) < \infty$ and our claim then follows. Indeed, if actually $\Psi(r) \leq 1$ then we simply take $\gamma = r$. Otherwise, since we have always $\Psi(\lambda) < L_{a',a'}(\lambda) \leq 1$, the equation $\Psi(q) = 1$ has a unique solution $\gamma \in (r, \lambda)$.

On the event $H(a') < H(b'')$, the process started from a' first stays in $[a', b'']$ until it makes a jump across a' at time $\sigma(a', b'')$, and then stays in $(0, a')$ until it eventually hits a' for the first time at the instant $H(a')$. Since

$$\sup_{x \leq a'} c(x)/x \leq \rho_{a,b} < r,$$

on the event that $H(a') < H(b'')$, there is the inequality

$$\mathcal{E}_{H(a')} e^{-rH(a')} \leq \mathcal{E}_{\sigma(a', b'')} e^{-r\sigma(a', b'')},$$

and we have therefore

$$\mathbb{E}_{a'} \left(\mathcal{E}_{H(a')} e^{-rH(a')}, H(a') < H(b'') \right) \leq \mathbb{E}_{a'} \left(\mathcal{E}_{\sigma(a', b'')} e^{-r\sigma(a', b'')}, H(a') < H(b'') \right).$$

Observe that there exists an instant $t \leq \sigma(a', b'')$ with $X_t < a'$ if and only if the process X stays in $[a, b]$ during the whole time-interval $[0, t]$ and exits from $[a', b'']$ at time t by jumping across the level a' , that is if and only if $t = \sigma(a', b'')$ and $H(a') < H(b'')$. Combining this observation with the fact that the predictable compensator of the jump process of X is $\bar{k}(X_{t-}, y) dy dt$, we deduce that

$$\begin{aligned} \mathbb{E}_{a'} \left(\mathcal{E}_{\sigma(a', b'')} e^{-r\sigma(a', b'')}, H(a') < H(b'') \right) &= \mathbb{E}_{a'} \left(\int_0^{\sigma(a', b'')} \mathcal{E}_t e^{-rt} \left(\int_0^{a'} \bar{k}(X_{t-}, y) dy \right) dt \right) \\ &\leq \|K\|_\infty \mathbb{E}_{a'} \left(\int_0^{\sigma(a', b'')} \mathcal{E}_t e^{-rt} dt \right), \end{aligned}$$

where $\|K\|_\infty = \sup_{x>0} K(x)$ is the maximal jump rate.

We then write $\delta := r - \rho_{a,b} > 0$ and use the inequality

$$\begin{aligned} \mathbb{E}_{a'} \left(\mathcal{E}_t e^{-rt}, t < \sigma(a', b'') \right) &\leq \frac{e^{-\delta t}}{\min_{[a', b'']} h_{a,b}} \mathbb{E}_{a'} \left(\mathcal{E}_t e^{-\rho_{a,b} t} h_{a,b}(X_t), t < \sigma(a, b) \right) \\ &= e^{-\delta t} \frac{h_{a,b}(a')}{\min_{[a', b'']} h_{a,b}}, \end{aligned}$$

where the equality is seen from Lemma 3.4. Since $\min_{[a', b'']} h_{a,b} > 0$ by Proposition 3.3, we now get that

$$\int_0^\infty \mathbb{E}_{a'} \left(\mathcal{E}_t e^{-rt}, t < \sigma(a', b'') \right) dt < \infty,$$

which entails our claim. \square

Proposition 3.7(ii) enables us to repeat the argument for the proof of Lemma 2.3, and this yields the following:

Corollary 3.8. *Notation and assumptions are as in Proposition 3.7. For $0 < x < b''$, we consider*

$$g(x) := \mathbb{E}_x \left(\mathcal{E}_{H(a')} e^{-\gamma H(a')}, H(a') < H(b'') \right).$$

The process

$$\mathcal{S}(t) := g(X_t) \mathcal{E}_t e^{-\gamma t} \mathbb{1}_{\{t < H(b'')\}}, \quad t \geq 0.$$

is then a \mathbb{P}_x -supermartingale for every $0 < x < b''$.

We have now completed all the preliminary steps needed to establish Theorem 1.2.

3.4 Proof of (17)

We first pick two good intervals (a, b) and (a', b') with $0 < a < a' < b' < b$ sufficiently large such that

$$\sup_{(0, a'] \cup [b', \infty)} c(x)/x < \rho_{a, b}$$

That this is indeed possible is seen from the same argument as in the proof of Proposition 3.7(i), thanks to Lemmas 3.1(ii) and 3.6 and the assumptions of Theorem 1.2. The irreducibility of the process killed when exiting from $[a, b]$ shows that provided we choose $b'' \in (b', b)$ close enough to b , then $\mathbb{P}_{b'}(H(a') < H(b'')) > 0$. Next, in the notation introduced in Proposition 3.7, take any $\max\{\rho_{a, b}, \gamma\} < q < \lambda$. In particular,

$$c(x)/x < q \text{ for all } x \in (0, a'] \cup [b', \infty).$$

We shall prove that (17) holds with $x = b'$. We thus let X start from b' and split the excursion interval $(0, H(b'))$ at times $\sigma(b', \infty)$ and $\sigma(a', \infty)$. Recall that $\sigma(b', \infty)$ is the first instant when X jumps across b' , so plainly $\sigma(b', \infty) \leq \sigma(a', \infty)$ $\mathbb{P}_{b'}$ -a.s., and these two first-exit times may coincide with positive probability. Because X stays in $[b', \infty)$ until time $\sigma(b', \infty)$, we have

$$\mathcal{E}_{\sigma(b', \infty)} e^{-q\sigma(b', \infty)} \leq 1 \quad \mathbb{P}_{b'}\text{-a.s.},$$

and we see from the strong Markov property that (17) will follow from

$$\sup_{x < b'} \mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \infty \right) < \infty. \quad (26)$$

First consider the case $x \leq a'$. The process started from x stays in $(0, a']$ until time $H(a')$, thus

$$\mathcal{E}_{H(a')} e^{-qH(a')} \leq 1 \quad \mathbb{P}_x\text{-a.s.},$$

and therefore we have, again from the strong Markov property,

$$\mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \infty \right) \leq \mathbb{E}_{a'} \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \infty \right).$$

Next recall Proposition 3.7, and chose $b'' \in (b', b)$ close enough to b such that $\mathbb{P}_{b'}(H(a') < H(b'')) > 0$. In the notation of Corollary 3.8, Proposition 3.7(ii) ensures that $g(a') \in (0, 1]$, and we have also $g(b') > 0$. We now deduce from the optional sampling theorem applied under $\mathbb{P}_{a'}$ to the supermartingale \mathcal{S} in Corollary 3.8 and the stopping time $H(b') < H(b'')$, that

$$\mathbb{E}_{a'} \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \infty \right) \leq \frac{1}{g(b')} \mathbb{E}_{a'} \left(\mathcal{S}_{H(b')} \right) \leq \frac{g(a')}{g(b')},$$

and hence

$$\sup_{x \leq a'} \mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \infty \right) \leq \frac{g(a')}{g(b')} < \infty. \quad (27)$$

Next, we consider the case $x \in (a', b')$, and distinguish whether the process exits from $[a', b']$ through the upper or the lower boundary. On the one hand, Lemma 3.4 and the optional sampling theorem readily yield for every $x \in [a', b']$:

$$\mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') < \sigma(a, b) \right) \leq \frac{h_{a,b}(x)}{h_{a,b}(b')},$$

and *a fortiori*

$$\sup_{a' \leq x < b'} \mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, H(b') \leq \sigma(a', b') \right) \leq \frac{\max_{[a', b']} h_{a,b}}{h_{a,b}(b')} < \infty. \quad (28)$$

On the other hand, the same argument as in the proof of Proposition 3.7(ii) shows that for every $x \in [a', b']$:

$$\begin{aligned} \mathbb{E}_x \left(\mathcal{E}_{\sigma(a', b')} e^{-q\sigma(a', b')}, \sigma(a', b') < H(b') \right) &\leq \|K\|_\infty \mathbb{E}_x \left(\int_0^{\sigma(a', b')} \mathcal{E}_t e^{-qt} dt \right) \\ &\leq \frac{\|K\|_\infty h_{a,b}(x)}{\delta \min_{[a', b']} h_{a,b}}, \end{aligned}$$

where $\delta = q - \rho_{a', b'} > 0$ and $\|K\|_\infty$ is the maximal jump rate. Combining this with (27) and the strong Markov property, we conclude that

$$\sup_{a' \leq x < b'} \mathbb{E}_x \left(\mathcal{E}_{H(b')} e^{-qH(b')}, \sigma(a', b') < H(b') < \infty \right) < \infty. \quad (29)$$

Together, (27), (28) and (29) establish (26), and this completes the proof of Theorem 1.2.

3.5 Miscellaneous comments about Theorem 1.2

- (i) We have argued in the Introduction that (19) yields the more explicit criterion (20) when X is recurrent, and we now discuss simple conditions in terms of the growth and fragmentation rates that ensure recurrence. Since X is irreducible and its trajectories have no positive jumps, point-recurrence can only fail if sample paths converge either to 0 or to ∞ a.s., which can easily be impeded by Foster-type conditions (see for instance [25] and [20] for general references). Typically, consider the power function $f(x) = x^r$ for some $r > 0$, and assume that $\mathcal{G}f(x) \leq 0$ for all large enough x . That is, for some $x_\infty > 0$:

$$rc(x)x^r + \int_0^x (y^r - x^r)yk(x, y)dy \leq 0 \quad \text{for all } x \geq x_\infty. \quad (30)$$

Then the process $f(X_{t \wedge \sigma(x_\infty, \infty)})$ is a \mathbb{P}_x -supermartingale for every $x \geq x_\infty$, and it follows that $\mathbb{P}_x(\lim_{t \rightarrow \infty} X_t = \infty) = 0$. Similarly, working now with negative powers and considering $g(x) = x^{-q}$ for some $q > 0$, if for some $x_0 > 0$

$$-qc(x)x^{-q} + \int_0^x (y^{-q} - x^{-q})yk(x, y)dy \leq 0 \quad \text{for all } x \leq x_0, \quad (31)$$

then the process $g(X_{t \wedge H(x_0)})$ is a \mathbb{P}_x -supermartingale for every $x \leq x_0$, and it follows that $\mathbb{P}_x(\lim_{t \rightarrow \infty} X_t = 0) = 0$. We refer further to Bouguet [7], who considered specifically positive recurrence for the family of piecewise deterministic Markov processes that arise in our setting; see notably Theorem 4 and also Remark 3 there. Putting the pieces together and excluding implicitly the case of linear growth², we get that exponentially fast convergence (18) holds provided that (20), (30) and (31) are satisfied.

Recall that Doumic Jauffret and Gabriel [17] obtained conditions that ensure the existence of eigenelements and thus also the Malthusian behavior (5) by the general relative entropy method (however their approach does not yield exponential speed of convergence (18)). The comparison of those conditions displays certain resemblance and but also differences. For instance, (12) and (13) in [17] can be loosely related to (30) and (31) here; we do not need here to make assumptions such as (5,7,10) in [17]; on the other hand [17] also covers the situation where the growth rate c does not fulfill (6).

- (ii) If we assume self-similarity of the fragmentation kernel, that is

$$k(x, y) = x^{-1}K(x)p(y/x), \quad 0 < y < x,$$

where $p \in L^1_+([0, 1])$ with $\int_0^1 p(u)udu = 1$, then (30) and (31) translate respectively into

$$\frac{c(x)}{xK(x)} \leq r^{-1} \int_0^1 (1 - u^r)up(u)du \quad \text{for all } x \geq x_\infty \quad (32)$$

and

$$\frac{c(x)}{xK(x)} \geq q^{-1} \int_0^1 (u^{-q} - 1)up(u)du \quad \text{for all } x \leq x_0. \quad (33)$$

Putting pieces together, we thus see that in the self-similar case, exponentially fast convergence (18) holds whenever (20), (32) and (33) are fulfilled.

This should be compared with Theorem 1.11 of Balagué et al. [1] in which the existence of a spectral gap is asserted under more stringent conditions. We stress however that [1] also covers cases where the growth rate c does not fulfill (6).

- (iii) We point out that exponential speed of convergence (18) may hold without (19). For instance, in the case of linear growth rate $c(x) = ax$, Section 6 in [6] discusses situations where $\lambda = a$ and nonetheless (18) takes place.

²Linear growth rate $c(x) = ax$ was already discussed in details in Section 6 of [6]

- (iv) We also recall from Corollary 4.5 and Lemma 4.6 in [6], that when (19) holds, then h is an eigenfunction for the eigenvalue λ of the growth-fragmentation operator \mathcal{A} (in particular h belongs to the domain of \mathcal{A}), and further the function $x \mapsto h(x)/x$ is continuous and bounded.

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