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▶ To cite this version:

Dominique Bakry, Ivan Gentil, Grégory Scheffer. Sharp Beckner-type inequalities for Cauchy and spherical distributions. 2018. hal-01761215v1

HAL Id: hal-01761215 https://hal.science/hal-01761215v1

Preprint submitted on 8 Apr 2018 (v1), last revised 15 Jan 2019 (v3)

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Sharp Beckner-type inequalities for Cauchy and spherical distributions

Dominique Bakry^{*}, Ivan Gentil[†]and Grégory Scheffer[‡]

April 8, 2018

Abstract

Using some harmonic extensions on the upper-half plane, and probabilistic representations, and curvature-dimension inequalities with some negative dimensions, we obtain some new opimal functional inequalities of the Beckner type for the Cauchy type distributions on the Euclidean space. These optimal inequalities appear to be equivalent to some non tight optimal Beckner inequalities on the sphere, and the family appear to be a new form of the Sobolev inequality.

Key words: Cauchy distribution, Beckner inequality, curvature-dimension condition, Poincaré inequality, Spherical analysis, Bessel processes, Stopping time.

1 Introduction

The so-called Γ_2 criterium is a way to prove functional inequalities such as Poincaré or logarithmic Sobolev inequalities. For instance, let $\Psi: \mathbb{R}^d \to \mathbb{R}$ be a smooth function such that $\nabla^2 \psi \geqslant \rho I$ with $\rho > 0$, then the probability measure

$$d\mu_{\psi} = \frac{e^{-\psi}}{Z} dx,$$

where Z is the normalization constant which turns μ_{ψ} into a probability, satisfies both the Poincaré inequality,

$$\int f^2 d\mu_{\psi} - \left(\int f d\mu_{\psi}\right)^2 \le \frac{1}{\rho} \int |\nabla f|^2 d\mu_{\psi},$$

and the logarithmic Sobolev inequality

$$\int f^2 \log f^2 d\mu_{\psi} - \int f^2 d\mu_{\psi} \log \int f^2 d\mu_{\psi} \leq \frac{2}{\rho} \int |\nabla f|^2 d\mu_{\psi},$$

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for any smooth function f. More generally, the same hypothesis leads to a family of interpolation inequalities between Poincaré and logarithmic Sobolev, namely, for any $p \in (1, 2]$,

$$\frac{p}{p-1} \left(\int f^2 d\mu_{\psi} - \left(\int f^{2/p} d\mu_{\psi} \right)^p \right) \le \frac{2}{\rho} \int |\nabla f|^2 d\mu_{\psi},$$

where the case p=2 corresponds to the Poincaré inequality and the limit $p\to 1$ to the logarithmic Sobolev inequality. The technics used to prove such inequalites are based on the Γ_2 calculus (the $CD(\rho,+\infty)$ condition) introduced by the first author with M. Émery in [BÉ85] (see also [BGL14]). This family is the so-called Beckner inequalities, proved for the usual Gaussian measure, that is when $\psi=\frac{\|x\|^2}{2}$, by W. Beckner in [Bec89] (in that case $\rho=1$). Some improvements of such inequalities can be found in [AD05, DNS08, BG10].

On the other hand, on the sphere $\mathbb{S}^d \subset \mathbb{R}^d$, with similar arguments, the following Beckner inequalities can be stated, $p \in (1,2]$,

$$\frac{p}{p-1} \left(\int f^2 d\mu_{\mathbb{S}} - \left(\int f^{2/p} d\mu_{\mathbb{S}} \right)^p \right) \le \frac{2}{d} \int |\nabla f|^2 d\mu_{\mathbb{S}},$$

see [BGL14], where in this situation, $\rho = d - 1$.

It is important to notice that all Beckner inequalities (for the spherical and the Gaussian models) are optimal in the sense that constants in front of the right hand side are optimal. However, only the constant functions saturate these inequalities when $p \neq 2$.

The aim of this article is two-fold. First, we explain and improve in a general context the method introduced by the last author [Sch03] to obtain some new optimal inequalities. This method is mainly probabilist, and makes strong use of various processes associated with the underlying structure behind the inequalities. He used this method to obtain an optimal Poincaré inequality under a $CD(\rho, n)$ condition, which is a stronger form of the condition $\nabla^2 \psi \geqslant \rho Id$, and without integration by parts, which apply even for non symmetric operators. This method is rather technical and is based on the $CD(\rho, n)$ condition with a negative dimension n, together with a construction of some sub-harmonic functional. Stochastic calculs is also a main tool of our approach.

Secondly, revisiting unpublished results of the last author, we apply this method to obtain Beckner inequalities for the generalized Cauchy distribution, that is the probability measure on \mathbb{R}^d with density

$$d\nu_b(x) = \frac{1}{Z(1+|x|^2)^b} dx$$

where b>d/2 and Z is the normalization constant. This measure does not satisfy any $CD(\rho, +\infty)$ condition with $\rho \geqslant 0$ but some CD(0, n) condition with n < 0, as will be explained with in detail in the paper. We prove the following new optimal family of inequalities,

$$\frac{p}{p-1} \left[\int f^2 d\nu_b - \left(\int f^{2/p} d\nu_b \right)^p \right] \le \frac{1}{(b-1)} \int \Gamma(f) (1+|y|^2) d\nu_b.$$

where $b \ge d + 1$ and $p \in [1 + 1/(b - d), 2]$.

When p=2 we recover a weighted Poincaré inequality proved by Nguyen [Ngu14] (see also others approaches in [BL09, BJM16, ABJ16]).

More surprisingly, this family is equivalent to the following one on the sphere \mathbb{S}^d ,

$$\int h^2 d\mu_{\mathbb{S}} \leq A \left(\int h^{2/p} d\mu_{\mathbb{S}} \right)^p + \frac{16}{(m+2-d)(3d-2+m)} \int \Gamma_{\mathbb{S}}(h) d\mu_{\mathbb{S}},$$

where

$$m \geqslant d+2, \quad p=1+\frac{2}{m-d} \in [1,2]$$

and $A \geqslant 1$ is some explicit constant to be described below, converging to 1 as p converges to 1, or equivalently when m converges to infinity. This inequality appears to be a non-tight inequality, that is constant functions are not optimal. On the other hand, this last inequality is optimal with explicit extremal functions. Moreover this new family of inequalities, or more precisely its behavior when p converges to 1, appears as a new form of a Sobolev inequality on the sphere, as many other ones described for example in [BCLS95].

The paper is organized as follows. In the next section, we describe the curvature-dimension condition $CD(\rho, n)$, even for negative n, and derive a general form of associated sub-harmonic functionals. In Section 3, we define an operator related to the generalized Cauchy distribution, and show that it satisfies the generalized curvature-dimension condition. In Section 4, we prove the new Beckner inequalities for these generalized Cauchy distribution. Finally, in Section 5, we prove the non-tight Becker inequalities on the d-dimensional sphere and show how they relate to the Sobolev inequality.

Notations: In all this paper, d will be the dimension of the main space, and satisfies $d \ge 1$. For every $X, Y \in \mathbb{R}^d$, $X \cdot Y = \sum_{i=1}^d X_i Y_i$ is the usual scalar product and $|X|^2$ is the Euclidean norm in \mathbb{R}^d .

2 General properties on curvature-dimension condition

Let (M^d,g) be a smooth connected d-dimensional Riemannian manifold and Δ_g its Laplace-Beltrami operator. The associated Ricci tensor is denoted Ric $_g$. For any diffusion operator L on M (that is a second order semi-elliptic differential operator with no zero-order term, see [BGL14]) we define the so-called carré du champ operator

$$\Gamma^L(f,f) = \Gamma^L(f) = \frac{1}{2}L(f^2) - fLf,$$

and its iterated operator

$$\Gamma_2^L(f, f) = \Gamma_2^L(f) = \frac{1}{2}L(\Gamma_L(f)) - \Gamma^L(f, Lf),$$

for any smooth function f. We note Γ (resp. Γ_2) instead Γ^L (resp. Γ_2^L) when there is no possible confusion.

2.1 Definitions

Definition 1 ($CD(\rho, n)$ condition) An operator $L = \Delta_g + X$ with X a smooth vector field satisfies a $CD(\rho, n)$ condition with $\rho \in \mathbb{R}$ and $n \in \mathbb{R} \setminus [0, d)$ if

$$\Gamma_2(f) \geqslant \rho \Gamma(f) + \frac{1}{n} (Lf)^2,$$
 (1)

for any smooth function f.

As we shall see of the proof in Lemma 2, this condition will never hold when $n \in [0, d)$ but we will extend it for n = 0.

Lemma 2 ([Bak94]) For any $\rho \in \mathbb{R}$ and $n \notin [0,d]$ the operator $L = \Delta + X$ satisfies a $CD(\rho,n)$ if and only if

 $\frac{n-d}{n}(\operatorname{Ric}(L) - \rho\Gamma) \geqslant \frac{X \otimes X}{n},\tag{2}$

and when n=d it reduces only to X=0 and the condition becomes $\mathrm{Ric}_g \geqslant \rho \Gamma$. In the definition, $\mathrm{Ric}(L)=\mathrm{Ric}_q-\nabla_S X$ where $\nabla_S X$ is the symmetrized tensor of ∇X .

Proof. — For completeness, we give a sketch of proof of this result which can be helpful for the rest of the paper. For any smooth function f, the Bochner-Lichnerowicz-Weitzenbock formula states that (cf. [BGL14, Sec C.5]),

$$\Gamma_2(f) = ||\nabla^2 f||_{H.S.}^2 + \operatorname{Ric}_q(\nabla f, \nabla f) - \nabla X_S(\nabla f, \nabla f) = ||\nabla^2 f||_{H.S.}^2 + \operatorname{Ric}(L)(\nabla f, \nabla f), \quad (3)$$

where $||\nabla^2 f||_{H.S.}$ is the Hilbert-Schmidt norm of f. On a fixed point $x \in M$, we choose a local chart such that the metric $g = (g_{i,j})$ is the identity matrix at the point x. In that case, $\nabla^2 f = (\nabla_{i,j} f)$ is a d-dimensional symmetric matrix denoted Y, $\Gamma(f) = \sum_{i=1}^d (\partial_i f)^2$ and let denote $Z = \nabla f$. Since $\sum_{i,j=1}^d Y_{i,j}^2 = \operatorname{tr}(Y^2)$, the $CD(\rho, n)$ condition (1) becomes

$$\operatorname{tr}(Y^2) + Z \cdot RZ \geqslant \rho Z \cdot Z + \frac{1}{n} (\operatorname{tr}(Y) + X \cdot Z)^2,$$

where R is a matrix representing the coordinates of the tensor Ric(L) at x. One can choose a new base in \mathbb{R}^d such that Y is diagonal. If we denote again Z in the new base in \mathbb{R}^d , we have to prove such inequality, for every vectors $(Y_{ii})_{1 \le i \le d}$ and $(Z_i)_{1 \le i \le d}$,

$$\sum_{i=1}^{d} Y_{ii}^{2} + Z \cdot RZ \geqslant \rho Z \cdot Z + \frac{1}{n} (\sum_{i=1}^{d} Y_{ii} + X \cdot Z)^{2}.$$

By taking Y_{ii} all equals, from $\sum_{i=1}^d Y_{ii}^2 \ge \frac{1}{d} (\sum_{i=1}^d Y_{ii})^2$, the previous inequality is equivalent to

$$\frac{1}{d} (\sum_{i=1}^{d} Y_{ii})^2 + Z \cdot RZ \geqslant \rho Z \cdot Z + \frac{1}{n} (\sum_{i=1}^{d} Y_{ii} + X \cdot Z)^2.$$

Let $y = \sum_{i=1}^{d} Y_{ii}^2$, then the inequality becomes

$$\frac{y^2}{d} + Z \cdot RZ \geqslant \rho Z \cdot Z + \frac{1}{n} (y + X \cdot Z)^2.$$

In other words, the $CD(\rho, n)$ condition is satisfied if and only if, for every $y \in \mathbb{R}$ and, $Z \in \mathbb{R}^d$

$$y^{2}(\frac{1}{d} - \frac{1}{n}) - y\frac{2}{n}X \cdot Z + Z \cdot (R - \rho \operatorname{Id})Z - \frac{1}{n}(X \cdot Z)^{2} \geqslant 0.$$

Then, either n=d, then we need to assume that X=0 and $Z\cdot(R-\rho\mathrm{Id})Z\geqslant 0$, either $\frac{1}{d}-\frac{1}{n}>0$ (with translates to $n\notin[0,d]$) and for any vector Z,

$$\frac{(X \cdot Z)^2}{n} \le \frac{n - d}{n} Z \cdot (R - \rho \mathrm{Id}) Z,$$

that is condition (2).

On may know define the tensor curvature-dimension condition, $TCD(\rho, n)$.

Definition 3 ($TCD(\rho, n)$ condition and quasi-models $QM(\rho, n)$) Let $\rho \in \mathbb{R}$, $n \in \mathbb{R} \setminus \{d\}$ and an operator $L = \Delta + X$ with X a smooth vector field.

• L satisfies a $TCD(\rho, n)$ condition if

$$\frac{n-d}{n}\left(\operatorname{Ric}(L) - \rho\Gamma\right) \geqslant \frac{X \otimes X}{n},\tag{4}$$

when $n \neq 0$, and

$$-d(\operatorname{Ric}(L) - \rho\Gamma) \le X \otimes X,$$

when n = 0.

• The operator L is a quasi model $QM(\rho, n)$ if

$$(n-d)(\operatorname{Ric}(L) - \rho\Gamma) = X \otimes X. \tag{5}$$

Indeed, this new definition allows us to extend (1) when $n \in [0, d)$, and we shall see later crucial examples of such operators, which are even quasi-models.

2.2 Sub-harmonic functionals

In this section, we construct sub-harmonic functionals from harmonic functions under $TCD(\rho, n)$. We systematically explore the most general ones, in view of further use.

The next result is rather technical, and we shall see in the next sections how it leads to Poincaré inequalities or Beckner inequalities.

Theorem 4 (Sub-harmonic functionals under $TCD(\rho, n)$ or $QM(\rho, n)$)

Let $\Phi: I \times [0, \infty) \to \mathbb{R}$ be a smooth function $(I \subset \mathbb{R})$, and denote Φ_i and Φ_{ij} its first and second derivatives, $i, j \in \{1, 2\}$. Let $L = \Delta_g + X$, where X is a smooth vector field.

We assume that L satisfies $TCD(\rho, n)$ with $\rho \in \mathbb{R}$ and $n \in \mathbb{R} \setminus (0, d)$ (or $QM(\rho, n)$ with $n \neq d$). Then, if $F: M \mapsto I$ is smooth and satisfies L(F) = 0, then $L\Phi(F, \Gamma(F)) \geqslant 0$ as soon as, at any point $(y, z) \in I \times [0, \infty)$, Φ satisfies, either

$$\begin{cases}
\Phi_{2} > 0; \\
\Phi_{2} \frac{n+1-d}{n-d} + 2\Phi_{22}z > 0; \\
(n-d)\left(n\Phi_{2} + 2(n-1)z\Phi_{22}\right) > 0; \\
2\rho\Phi_{2} + \Phi_{11} - 2(n-1)\frac{\Phi_{12}^{2}z}{n\Phi_{2} + 2(n-1)\Phi_{22}z} \geqslant 0,
\end{cases}$$
(6)

either

$$\begin{cases}
\Phi_2 = 0; \\
\Phi_{22}z = 0; \\
\Phi_{12}z = 0; \\
\Phi_{11}z \ge 0,
\end{cases}$$
(7)

either

$$\begin{cases}
\Phi_2 = 0; \\
\Phi_{22}z > 0; \\
\Phi_{11}\Phi_{22} - \Phi_{12}^2 \geqslant 0.
\end{cases}$$
(8)

Remarks 5 When $n \notin [0,d]$, condition $\Phi_2 \frac{n+1-d}{n-d} + 2\Phi_{22}z > 0$ is contained into $\Phi_2 + 2\frac{n-1}{n}z\Phi_{22} > 0$. Hence, provided that $n \notin [0,d]$, then (6) boils down to

$$\begin{cases}
\Phi_2 > 0, \\
\Phi_2 + 2\frac{n-1}{n}z\Phi_{22} > 0, \\
2\rho\Phi_2 + \Phi_{11} - 2(n-1)\frac{\Phi_{12}^2 z}{n\Phi_2 + 2(n-1)\Phi_{22}z} \geqslant 0.
\end{cases}$$
(9)

It is interesting to notice that when $n \notin [0,d]$ the dimension d does not appear in the previous conditions (9).

Proof. — The proof is an extension of the proof of Lemma 2. If LF = 0, then from the diffusion property of L (cf. [BGL14, Sec 1.14]),

$$L(\Phi(F,\Gamma(F))) = \Phi_2 L\Gamma(F) + \Phi_{11}\Gamma(F) + 2\Phi_{12}\Gamma(F,\Gamma(F)) + \Phi_{22}\Gamma(\Gamma(F),\Gamma(F))$$

= $2\Phi_2\Gamma_2(F) + \Phi_{11}\Gamma(F) + 2\Phi_{12}\Gamma(F,\Gamma(F)) + \Phi_{22}\Gamma(\Gamma(F),\Gamma(F)), \quad (10)$

since $L\Gamma(F) = 2\Gamma_2(F)$ when LF = 0, for simplicity we note $\Phi_{ij} = \Phi_{ij}(F, \Gamma(F))$. At some point $x \in M$, we write all the expressions appearing in $L(\Phi(F, \Gamma(F)))$. From (3), they make use only of $\nabla \nabla F$ and ∇F . We chose a system of coordinates in which at this point the metric g is identity, $\nabla \nabla F$ is diagonal, with eigenvalues $(\lambda_i)_{1 \leq i \leq d}$. The components of ∇F are denoted Z_i , and the components of X are denoted X_i .

Since L(F) = 0, the constraint can be written,

$$\sum_{i=1}^{d} \lambda_i = -\sum_{i=1}^{d} Z_i X_i.$$

and

$$\Gamma(F,\Gamma(F)) = 2\sum_{i=1}^d \lambda_i Z_i^2, \quad \Gamma(\Gamma(F),\Gamma(F)) = 4\sum_{i=1}^d \lambda_i^2 Z_i^2, \quad \Gamma(F) = \sum_{i=1}^d Z_i^2.$$

From the $TCD(\rho, n)$ conditions and (3), the Γ_2 operator has a lower bound,

$$\Gamma_2(F) = \sum_{i=1}^d \lambda_i^2 + \text{Ric}(L)(Z, Z) \geqslant \sum_{i=1}^d \lambda_i^2 + \rho \sum_{i=1}^d Z_i^2 + \frac{1}{n-d} \Big(\sum_{i=1}^d Z_i X_i\Big)^2.$$

Under the $TCD(\rho, n)$ this is an inequality and we need to assume that $n \notin (0, d)$ and when the operator is a quasi-model $QM(\rho, n)$ we only need to assume that $n \neq d$ and this is an equality. Since $\Phi_2 \geq 0$, from (10), we just have to minimize the following expression,

$$\frac{L(\Phi(F,\Gamma(F)))}{2} \geqslant 2\Phi_2 \Big(\sum_{i=1}^d \lambda_i^2 + \rho \sum_{i=1}^d Z_i^2 + \frac{1}{n-d} \Big(\sum_{i=1}^d Z_i X_i \Big)^2 \Big) + \Phi_{11} \sum_{i=1}^d Z_i^2 + 4\Phi_{12} \sum_{i=1}^d \lambda_i Z_i^2 + 4\Phi_{22} \sum_{i=1}^d \lambda_i^2 Z_i^2,$$

for all (λ_i) and Z_i , under the constraint $\sum_{i=1}^d \lambda_i = \sum_{i=1}^d X_i Z_i$.

Let $z = \sum_{i=1}^d Z_i^2 = \Gamma(F)$ and $Z_i^2 = z_i$, then the expression becomes by using the constraint (after dividing by 2)

$$\Phi_2\left(\sum_{i=1}^d \lambda_i^2 + \rho z + \frac{1}{n-d} \left(\sum_{i=1}^d \lambda_i\right)^2\right) + \frac{\Phi_{11}}{2} z + 2\Phi_{12} \sum_{i=1}^d \lambda_i z_i + 2\Phi_{22} \sum_{i=1}^d \lambda_i^2 z_i.$$

We first optimise on the simplex $\{z_i \geq 0, \sum_{i=1}^d z_i = z\}$ where z > 0 is fixed. We first observe that the expression is affine in z_i . The vector $(z_i)_{1 \leq i \leq d}$ belongs to the simplex $\{y_i \geq 0, \sum_{i=1}^d z_i = z\}$, so that the extremals of this expression $((\lambda_i)_{1 \leq i \leq d} \text{ and } z > 0 \text{ are fixed})$ is obtained at the one of the extremal point of the simplex. That is one of the z_i is z and all the other are 0. Assuming that $z_1 = z \neq 0$, one is now led to minimize

$$\begin{split} \lambda_1^2(\Phi_2 + 2\Phi_{22}z) + 2\Phi_{12}\lambda_1 z + \Phi_2 \sum_{i=2}^d \lambda_i^2 + \frac{\Phi_2}{n-d} \Big(\sum_{i=1}^d \lambda_i\Big)^2 + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) \\ &= \lambda_1^2(\Phi_2 \frac{n+1-d}{n-d} + 2\Phi_{22}z) + \lambda_1 \Big(2\Phi_{12}z + 2\frac{\Phi_2}{n-d}\sum_{i=2}^d \lambda_i\Big) \\ &+ \Phi_2 \Big(\sum_{i=2}^d \lambda_i^2 + \frac{1}{n-d} \Big(\sum_{i=2}^d \lambda_i\Big)^2\Big) + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) \end{split}$$

Then, we minimize in λ_1 . This imposes the second condition of (6), that is

$$\Phi_2 \frac{n+1-d}{n-d} + 2\Phi_{22}z > 0,$$

recall that $z = \Gamma(F)$. And the minimizer is given by

$$\frac{L(\Phi(F,\Gamma(F)))}{2} \geqslant \Phi_2\Big(\sum_{i=2}^d \lambda_i^2 + \frac{1}{n-d}\Big(\sum_{i=2}^d \lambda_i\Big)^2\Big) + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) - \frac{(\Phi_{12}z + \frac{\Phi_2}{n-d}\sum_{i=2}^d \lambda_i)^2}{\Phi_2\frac{n+1-d}{n-d} + 2\Phi_{22}z}$$

Let assume that $\Phi_2 > 0$. Decomposing the vector $\Lambda = (\lambda_i)_{2 \leq i \leq d}$ into a vector parallel to $(1, \dots, 1)$ and a vector $u = (u_i)_{2 \leq i \leq d}$ orthogonal to it, that is $\Lambda = \lambda(1, \dots, 1) + u$ with $\sum_{i=2}^{d} u_i = 0$. From this decomposition, we have $\sum_{i=2}^{d} \lambda_i = (d-1)\lambda$ and

$$\sum_{i=2}^{d} \lambda_i^2 = \sum_{i=2}^{d} (\lambda + u_i)^2 \geqslant (d-1)\lambda^2.$$

We see that we are bound to find the minimum in λ of

$$\begin{split} \Phi_2\Big((d-1)\lambda^2 + \frac{(d-1)^2}{n-d}\lambda^2\Big) + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) - \frac{(\Phi_{12}z + \frac{(d-1)\Phi_2}{n-d}\lambda)^2}{\Phi_2\frac{n+1-d}{n-d} + 2\Phi_{22}z} \\ &= \lambda^2\frac{\Phi_2(d-1)\Big(n\Phi_2 + 2(n-1)z\Phi_{22}\Big)}{(n-d)(\Phi_2\frac{n+1-d}{n-d} + 2\Phi_{22}z)} + \lambda\frac{2(d-1)\Phi_2\Phi_{12}z}{(n-d)(\Phi_2\frac{n+1-d}{n-d} + 2\Phi_{22}z)} \\ &\qquad \qquad - \frac{\Phi_{12}^2z^2}{\Phi_2\frac{n+1-d}{n-d} + 2\Phi_{22}z} + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) \end{split}$$

Then, once again, to get a minimizer in λ , we also need

$$(n-d)(n\Phi_2 + 2(n-1)\Phi_{22}z) > 0,$$

which is the third assumption in (6). Then, the minimizer is

$$\begin{split} \frac{L(\Phi(F,\Gamma(F)))}{2} \geqslant -\frac{(d-1)\Phi_2\Phi_{12}^2z^2}{(n-d)(n\Phi_2+2(n-1)z\Phi_{22})(\Phi_2\frac{n+1-d}{n-d}+2\Phi_{22}z)} \\ -\frac{\Phi_{12}^2z^2}{\Phi_2\frac{n+1-d}{n-d}+2\Phi_{22}z} + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) = \frac{-z^2(n-1)\Phi_{12}^2}{n\Phi_2+2z(n-1)\Phi_{22}} + z\Big(\Phi_2\rho + \frac{\Phi_{11}}{2}\Big) \geqslant 0, \end{split}$$

which is the fourth assumption in (6) since $z \ge 0$.

When $\Phi_2 = 0$ as some point (y, z), we follow the inequality to check the assumptions needed.

Corollary 6 (Sub-harmonic functionals under TCD(0,n)) Let assume that $L = \Delta + X$ satisfies a TCD(0,n) with n < 0. Let θ be a positive and smooth function on an interval $I \subset \mathbb{R}$ such that

$$2\frac{n-1}{n}(\theta')^2 \le \theta\theta''. \tag{11}$$

Then for any function $F: M \mapsto I$ such that LF = 0,

$$L(\theta(F)\Gamma(F)) \geqslant 0. \tag{12}$$

In particular, for any $\beta \in [n/(2-n), 0]$ and nonnegative harmonic function F,

$$L(F^{\beta}\Gamma(F)) \geqslant 0.$$

When n goes to $-\infty$, the condition (11) on θ degenerates into $2\theta' \leq \theta\theta''$. This means that the function Φ such that $\Phi'' = \theta$ is admissible : Φ is convex and $1/\Phi''$ is concave.

Moreover, when $n \in (-\infty, 0)$, the extremal case is given when $\theta(x) = x^{n/(2-n)}$, where (11) is an equality.

Although we shall mainly use Corollary 6 instead the general result Theorem 4, the proof in this last case is not really simpler than the general one.

3 The operators $Q_t^{(m)}$

Definition 7 (The operators $Q_t^{(m)}$) For any $t \ge 0$, m > 0, $x \in \mathbb{R}^d$, and any bounded function $f : \mathbb{R}^d \mapsto \mathbb{R}$, let

$$Q_t^{(m)}(f)(x) = \frac{1}{c(m,d)} \int \frac{f(yt+x)}{(1+|y|^2)^{\frac{m+d}{2}}} dy,$$
(13)

where

$$c(m,d) = \int \frac{1}{(1+|y|^2)^{\frac{m+d}{2}}} dy = \frac{\Gamma(\frac{m}{2})}{\Gamma(\frac{m+d}{2})} \pi^{d/2}$$
(14)

is the normalization constant, such that $Q_t^{(m)}(1) = 1$. In other words, for any t > 0,

$$Q_t^{(m)}(f)(x) = \int f(y)q_t(x, dy),$$

where the kernel q_t is given by

$$q_t(x,dy) = \frac{1}{c(m,d)} \frac{t^m}{(t^2 + |x - y|^2)^{\frac{m+d}{2}}} dy.$$
 (15)

The probability measure

$$\frac{1}{c(m,d)}(1+|y|^2)^{-\frac{m+d}{2}}dy$$

in \mathbb{R}^d , with m > 0, is known in statistics as the multivariate t-distribution with m degrees of freedom. There is a huge literature on the subject, see for instance [KN04].

Through an integration by parts (or from the computation of normalization constants) we may notice the useful formula,

$$\frac{c(m,d)}{c(m-2,d)} = \frac{m-2}{m-2+d},\tag{16}$$

for any m > 2.

Proposition 8 (Harmonicity of $Q_t^{(m)}f$ **)** Let m > 0. For any smooth and compactly supported function f, the map $(0, +\infty) \times \mathbb{R}^d \ni (t, x) \mapsto Q_t^{(m)}(f)(x)$ is solution of the elliptic equation $\Delta^{(m)}Q_t^{(m)}(f) = 0$ where

$$\Delta^{(m)} = \Delta_{\mathbb{R}^d} + \partial_{tt}^2 + \frac{1-m}{t} \partial_t. \tag{17}$$

Proof. — This can be proved directly on the expression of the kernel (15).

From Proposition 8 (or from a direct computation) the operators $(Q_t^{(m)})_{t\geqslant 0}$ admit an another formulation. Let $(Z_s)_{s\geqslant 0}=(X_s,Y_s)_{s\geqslant 0}\in\mathbb{R}^d\times\mathbb{R}$ be a diffusion process with generator $\Delta^{(m)}$ starting from $(x,t)\in\mathbb{R}^d\times(0,\infty)$. That is $(X_s)_{s\geqslant 0}$ is a Brownian motion up to a factor $\sqrt{2}$ and $(Y_s)_{s\geqslant 0}$ is independent of $(X_s)_{s\geqslant 0}$ with generator $\partial_t^2+\frac{1-m}{t}\partial_t$.

The process $(Y_s)_{s\geqslant 0}$ is a Bessel process with parameter m starting from t>0. When m>0, the stopping time $S=\inf\{u>0,\ Y_s=0\}$ is finite a.s., and since $Q_t^{(m)}(f)$ is a solution of the elliptic equation (13), from Itô's Lemma, or Dynkin's formula $[\emptyset k03, Sec 7.3]$

$$Q_t^{(m)}(f)(x) = E_{x,t}(f(X_S)), \tag{18}$$

for every smooth and bounded function f. In other words since, S and $(X_s)_{s\geqslant 0}$ are independent, if $\sigma_m(s,t)ds$ denotes the law of the hitting time S starting from t>0,

$$Q_t^{(m)}(f)(x) = \int_0^\infty P_s f(x) \sigma_m(s, t) ds, \tag{19}$$

where $(P_s)_{s\geqslant 0}$ is the heat semi-group in \mathbb{R}^d associated to the Euclidean Laplacian Δ . From the definition of the operators $Q_t^{(m)}$ and the definition of the the equation $(P_s)_{s\geqslant 0}$,

$$P_s f(x) = \int f(y) \exp\left(-\frac{|x-y|^2}{4s}\right) \frac{dy}{(4\pi s)^{d/2}},$$

and this gives a way to compute the law σ_m of the hitting time. One obtains, for any t > 0, the probability measure

$$\sigma_m(s,t)ds = \frac{1}{2^m \Gamma(m/2)} \frac{t^m \exp(-t^2/4s)}{s^{m/2+1}} ds.$$
 (20)

This formula can also be found for instance in [BS96].

The map $(t,s) \mapsto \sigma_m(t,s)$ satisfies

$$\partial_s = \partial_{tt}^2 + \frac{1 - m}{t} \partial_t,$$

with Dirichlet boundary conditions, $\sigma_m(t,0) = 0$ for every t > 0, and $\sigma_m(0,s)ds = \delta_0(s)$. We recognize in the RHS the generator of the Bessel process. This phenomena is universal and has nothing to do in particular with Bessel processes, see for instance [Pav14, Sec 7.2].

From this observation and an integration by parts, one recovers that the map $(t,x) \mapsto Q_t^{(m)} f(x)$ satisfies the elliptic equation $\Delta^{(m)} Q_t^{(m)}(f) = 0$.

Lemma 9 Let m > 0 and $p \in (0, m/2)$, then for any smooth and bounded function g,

$$E_{x,t}(S^p g(X_S)) = t^{2p} \frac{\Gamma(\frac{m}{2} - p)}{4^p \Gamma(\frac{m}{2})} Q_t^{(m-2p)}(g)(x).$$

In particular, when p = 1 and m > 2,

$$E_{x,t}(Sg(X_S)) = \frac{t^2}{2(m-2)} Q_t^{(m-2)}(g)(x).$$
(21)

Proof. — The formula comes from a direct computation since we have

$$E_{x,t}(S^p g(X_S)) = \frac{1}{2^m \Gamma(m/2)(4\pi)^{d/2}} \int g(y) \int_0^\infty s^{p - \frac{m+d}{2} - 1} \exp\left(-\frac{|x - y|^2}{4s} - \frac{t^2}{4s}\right) ds \, dy$$
$$= t^{2p} \frac{\Gamma(\frac{m}{2} - p)}{2^m \Gamma(m/2)(4\pi)^{d/2}} \int P_s g(y) \sigma_{m-2p}(s) ds.$$

Lemma 10 (TCD and QM conditions for $\Delta^{(m)}$) For any $m \neq 1$, the operator $\Delta^{(m)}$ on $\mathbb{R}^d \times (0,\infty)$ is a quasi-model QM(0,d-m+2). And for any $m \leq -1$ or $m \geq d+2$, $\Delta^{(m)}$ satisfies the TCD(0,d-m+2) condition.

Proof. — For that model,

$$\Delta^{(m)} = \Delta_{\mathbb{R}^{d+1}} + \frac{1-m}{t} \partial_t,$$

that is $X = \frac{1-m}{t}\partial_t$. Then $\operatorname{Ric}(L) = \frac{1-m}{t^2}(\partial_t)^2$ and $X \otimes X = \frac{(1-m)^2}{t^2}(\partial_t)^2$ and the result follows from the definition (5).

4 Beckner and Φ -entropy inequalities for Cauchy type distribution

4.1 Beckner inequalities for Cauchy type distribution

Let $f: \mathbb{R} \mapsto [0, \infty)$ be a smooth and compactly supported function, and let $F(t, x) = Q_t^{(m)}(f)(x)$ with m > 0.

Let $\beta \in \mathbb{R}$. From proposition (8), F is $\Delta^{(m)}$ -harmonic, and then

$$\Delta^{(m)}(F^{\beta+2}) = (\beta+2)(\beta+1)F^{\beta}\Gamma^{\Delta^{(m)}}(F).$$

and, as in Section 3, we can write,

$$E_{x,t}(F(Z_S)^{\beta+2}) = E_{x,t}(F(Z_0)^{\beta+2}) + (\beta+2)(\beta+1) \int_0^\infty E_{x,t}(F^{\beta}(Z_s)\Gamma^{\Delta^{(m)}}(F)(Z_s)\mathbb{1}_{s\leq S})ds,$$

where $(Z_s)_{s\geqslant 0}$ is the Markov process in \mathbb{R}^{d+1} with generator $\Delta^{(m)}$. Since $F(Z_S)=f(X_S)$, from (18) we have the general equation,

$$Q_t^{(m)}(f^{\beta+2}) = Q_t^{(m)}(f)^{\beta+2} + (\beta+2)(\beta+1) \int_0^\infty E_{x,t}(F^{\beta}(Z_s)\Gamma^{\Delta^{(m)}}(F)(Z_s)\mathbb{1}_{s \le S})ds.$$
 (22)

Again, from Dynkin's formula applied to $F^{\beta}\Gamma^{\Delta^{(m)}}(F)$, we have

$$E_{x,t}(F^{\beta}(Z_s)\Gamma^{\Delta^{(m)}}(F)(Z_s)\mathbb{1}_{s \le S}) = E_{x,t}(F^{\beta}(Z_S)\Gamma^{\Delta^{(m)}}(F)(Z_S)\mathbb{1}_{s \le S}) - E_{x,t}\left(\int_{s}^{S} L(F^{\beta}\Gamma^{\Delta^{(m)}}(F))(Z_u)du\mathbb{1}_{s \le S}\right).$$

In other words,

$$Q_{t}^{(m)}(f^{\beta+2}) - Q_{t}^{(m)}(f)^{\beta+2} = (\beta+2)(\beta+1)E_{x,t}(SF^{\beta}(Z_{S})\Gamma^{\Delta^{(m)}}(F)(Z_{S}))ds - (\beta+2)(\beta+1)\int_{0}^{\infty} E_{x,t}(\int_{s}^{S} L(F^{\beta}\Gamma^{\Delta^{(m)}}(F))(Z_{u})du\mathbb{1}_{s\leq S})ds.$$
 (23)

Let now assume that $m \geqslant d+2$, and let denote n=d-m+2 $(n \le 0)$. Let also assume that $\beta \in [\frac{n}{2-n},0]$ (that is $\beta \in [-1+\frac{2}{m-d},0]$, $\frac{2}{m-d}>0$). Since $n \le 0$, $\Delta^{(m)}$ satisfies TCD(0,n) (Lemma 10) and from Corollary 6,

$$L(F^{\beta}\Gamma^{\Delta^{(m)}}(F)) \geqslant 0.$$

The equation (23) becomes

$$Q_t^{(m)}(f^{\beta+2}) \leq Q_t^{(m)}(f)^{\beta+2} + (\beta+2)(\beta+1) \int_0^\infty E_{x,t}(F^{\beta}(Z_S)\Gamma^{\Delta^{(m)}}(F)(Z_S)\mathbb{1}_{s\leq S})ds.$$

The right-hand-term can be written

$$\int_0^\infty E_{s,t}(F^{\beta}(Z_S)\Gamma^{\Delta^{(m)}}(F)(Z_S)1\!\!1_{s\leq S})ds = E_{x,t}(SF^{\beta}(Z_S)\Gamma^{\Delta^{(m)}}(F)(Z_S)).$$

Starting form $Z_0 = (x, t)$,

$$F^{\beta}(Z_S)\Gamma^{\Delta^{(m)}}(F)(Z_S) = f^{\beta}(X_S)\Gamma^{\Delta^{(m)}}(F)(X_S, 0) = f^{\beta}(X_S)[\Gamma(f)(X_S) + (\partial_t F)^2(X_S, 0)],$$

since $\Gamma^{\Delta^{(m)}}(F) = \Gamma(F) + (\partial_t F)^2$ where Γ is the carré du champ operator associated to Euclidean Laplacian in \mathbb{R}^d . From Proposition 8, since m > 1, $\partial_t F(x, 0) = 0$ then

$$F^{\beta}(Z_S)\Gamma(F)(Z_S) = f^{\beta}(X_S)\Gamma(f)(X_S).$$

We can now apply Lemma 9, to get

$$Q_t^{(m)}(f^{\beta+2}) - Q_t^{(m)}(f)^{\beta+2} \le (\beta+2)(\beta+1)\frac{t^2}{2(m-2)}Q_t^{(m-2)}(f^{\beta}\Gamma(f)). \tag{24}$$

Changing $f^{\beta+2}$ by f^2 and $p=\beta+2$ we have obtained

Theorem 11 (Beckner inequalities for $Q_t^{(m)}$) For any $m \ge d+2$ and $p \in [1+\frac{2}{m-d},2]$, and for any smooth function $f \ge 0$,

$$\frac{p}{p-1} \left(Q_t^{(m)}(f^2) - Q_t^{(m)}(f^{2/p})^p \right) \le \frac{2t^2}{m-2} Q_t^{(m-2)}(\Gamma(f)). \tag{25}$$

When p=2 $(m \geqslant d+2)$, we obtain the Poincaré inequality for every smooth function f,

$$Q_t^{(m)}(f^2) - Q_t^{(m)}(f)^2 \le \frac{t^2}{(m-2)} Q_t^{(m-2)}(\Gamma(f)). \tag{26}$$

For t=1 and x=0 in the previous inequality, we have obtained the Beckner inequality for the Cauchy distribution $\frac{1}{c(m,d)} \frac{1}{(1+|y^2|)^{\frac{m+d}{2}}}$.

So, for any $m \geqslant d+2$ and $p \in [1+\frac{2}{m-d},2]$ one has,

$$\frac{p}{p-1} \left[\frac{1}{c(m,d)} \int \frac{f^2}{(1+|y|^2)^{\frac{m+d}{2}}} dy - \left(\frac{1}{c(m,d)} \int \frac{f^{2/p}}{(1+|y|^2)^{\frac{m+d}{2}}} dy \right)^p \right] \\
\leq \frac{2}{m-2} \frac{1}{c(m-2,d)} \int \frac{\Gamma(f)}{(1+|y|^2)^{\frac{m+d-2}{2}}} dy. \quad (27)$$

The inequality can be written as a weighted Poincaré-type inequality for a generalized Cauchy distribution. Let define for $b > \frac{d}{2}$, the probability measure

$$d\nu_b(y) = \frac{1}{c(2b-d,d)} \frac{1}{(1+|y|^2)^b} dy$$

and replacing b = (m + d)/2 in the previous inequality, we have obtained, using (16),

Theorem 12 (Beckner inequalities for Cauchy distribution) For any $b \ge d+1$ and $p \in [1+\frac{1}{b-d},2]$,

$$\frac{p}{p-1} \left[\int f^2 d\nu_b - \left(\int f^{2/p} d\nu_b \right)^p \right] \le \frac{1}{(b-1)} \int \Gamma(f) (1+|y|^2) d\nu_b, \tag{28}$$

for any smooth and nonnegative function f.

Moreover, the inequality is optimal, that is, fixing $p \in [1 + \frac{1}{b-d}, 2]$, the constant $\frac{1}{(b-1)}$ is the best possible in the inequality (28).

The case p=2 can be extended to any smooth function f (not only nonnegative function). Then we obtain the following result, proved in [Sch01], and also rediscover in [Ngu14].

Corollary 13 (Poincaré inequality for Cauchy distribution) For any $b \ge d+1$, the measure ν_b satisfies a Poincaré type inequality,

$$\int f^2 d\nu_b - \left(\int f d\nu_b \right)^2 \le \frac{1}{2(b-1)} \int \Gamma(f) (1+|y|^2) d\nu_b, \tag{29}$$

for any smooth function f. Moreover, functions $y \mapsto y_i$ with $1 \le i \le d$ are extremal functions.

Proof. — We only have to prove optimality. Apply the Poincaré inequality (29) to the function $y \mapsto y_i$ for some $1 \le i \le d$,

$$\nu_b(y_i^2) \le \frac{1}{2(b-1)}\nu_b(1+|y|^2) = \frac{1}{2(b-1)}(1+\nu_b(|y|^2)),$$

or equivalently

$$\frac{1}{d}\nu_b(|y|^2) \le \frac{1}{2(b-1)}(1+\nu_b(|y|^2)).$$

But this last one is an equality since, from (16), $\nu_b(|y|^2) = \frac{d}{2b-2-d}$.

Remarks 14

- Inequality (28) is equivalent to (25); indeed we only have to replace the map $y \mapsto f(y)$ by $y \mapsto f(ty + x)$.
- We cannot reach the logarithmic Sobolev inequality, since in our computation $p \ge 1 + 1/b$ and this inequality would require a limit p goes to 1.
- A Taylor expansion with $\varepsilon \to 0$ in inequality (28) with $f = 1 + \varepsilon g$ implies back the optimal Poincaré inequality (29), since

$$\frac{p}{p-1} \left[\int f^2 d\nu_b - \left(\int f^{2/p} d\nu_b \right)^p \right] = 2\varepsilon^2 \left[\int g^2 d\nu_b - \left(\int g d\nu_b \right)^2 \right] + o(\varepsilon^2).$$

We do not know if there are extremal functions apart of constant functions.

• Corollary 13 has also been proved by Nguyen [Ngu14, Cor. 14], using the Brunn-Minkowski theory (see also the approach in [BL09, ABJ16]).

In the one dimensional case, the exact value has been computed by Bonnefont, Joulin and Ma in [BJM16, Thm 3.1]. Their result is more general, since they compute the optimal constant for every b > 1/2. The constant of (29) with d = 1 is

$$\begin{cases} \frac{1}{2(b-1)} & \text{if} \quad b \geqslant 3/2; \\ \frac{4}{(2b-1)^2} & \text{if} \quad 1/2 < b \le 3/2. \end{cases}$$

It is interesting to notice that there are two regimes, depending on the range of the parameter b. Actually, we are not able with our method to reach this range, even in dimension 1.

• Applying inequality (28) to the function $x \mapsto f(\sqrt{2b}x)$, then, when $b \to +\infty$, the inequality becomes the optimal Beckner inequality for the Gaussian measure γ in \mathbb{R}^d ,

$$\frac{p}{p-1} \left[\int f^2 d\gamma - \left(\int f^{2/p} d\gamma \right)^p \right] \le 2 \int \Gamma(f) d\gamma,$$

for any smooth and positive function f. This inequality is proved by W. Beckner [Bec89]. In that case the inequality holds for any $p \in (1,2]$, and the limit case (when $p \to 1$) leads to the optimal Logarithmic Sobolev inequality for the Gaussian measure.

4.2 Tightness of the inequalities

We are interesting in the tightness of inequality (25), in the sense of the curvature-dimension condition.

Lemma 15 (Taylor expansion of $(Q_t^{(m)})$ For any m > 4, we have

$$Q_t^{(m)} = \operatorname{Id} + \frac{\Delta}{2(m-2)}t^2 + \frac{\Delta^2}{8(m-2)(m-4)}t^4 + o(t^4).$$
 (30)

Proof. — One can check directly on the formula. On the other hand, one can deduce the formula with functional analysis, at least formally. Since $Q_t^{(m)}$ is $\Delta^{(m)}$ -harmonic from Proposition 8 then, from the functional analysis point of view, $Q_t^{(m)}$ is a function $G_m(t, \Delta)$. From the definition (17) of $\Delta^{(m)}$, the function G_m satisfies for $t \geq 0$, x < 0,

$$xG_m(t,x) + \frac{1-m}{t}\partial_t G_m(t,x) + \partial_{tt}^2 G_m(t,x) = 0,$$

with the boundary conditions $G_m(0,x) = 1$. Solution of such differential equation is given by, $G_m(t,x) = H_m(-t^2x)$ where H_m satisfies for any $\lambda \ge 0$,

$$4\lambda H_m''(\lambda) - 2(m-2)H_m'(\lambda) - H_m(\lambda) = 0.$$

An asymptotic development of $H_m(\lambda) = \sum_{p \geq 0} c_p \lambda^p$ gives $c_0 = 1$ (since $G_m(0, x) = 1$) and

$$c_{p+1} = \frac{1}{2(p+1)(2p-m+2)}c_p.$$

For the forth order Taylor expansion of $Q_t^{(m)}$, we get (30), that is $c_1 = \frac{1}{2(2-m)}$ and $c_2 = \frac{1}{8(m-4)(m-2)}$.

This Lemma is fundamental. The Taylor expansion in (25), or equivalently in (26), gives in return the CD(0,d) conditions. We say then that the inequalities are Γ_2 -tight and nothing has been lost in term of curvature-dimension condition (1), in other words the Taylor expansion of the inequality implies the CD(0,d) condition.

To see that, it is enough to compute the Taylor expansion of the inequality (25), when t goes to 0. It is important to fix $\beta = \frac{2-m+d}{m-d} = -1 + \frac{2}{m-d}$ the extremal bound of the admissible interval of β .

It is clear that the zero order terms vanish, it is also not difficult to observe that the second order terms also vanish. Therefor we only have to compute the fourth order terms and we obtain the following expression,

$$\begin{split} \frac{1}{8(m-4)(m-2)} \Delta^2(f^{\beta+2}) &- \frac{(\beta+2)}{8(m-4)(m-2)} f^{\beta-1} \Delta^2(f) - \frac{(\beta+2)(\beta+1)}{8(m-2)^2} f^{\beta-1} \Delta(f)^2 \leq \\ &\frac{\beta(\beta+1)(\beta+2)}{2(m-4)(m-2)} \Gamma(f,\Gamma(f)) f^{\beta-1} + \frac{\beta(\beta+1)(\beta+2)}{4(m-4)(m-2)} f^{\beta-1} \Delta(f) \Gamma(f) \\ &+ \frac{(\beta-1)\beta(\beta+1)(\beta+2)}{4(m-4)(m-2)} f^{\beta-2} \Gamma(f)^2 + \frac{(\beta+1)(\beta+2)}{4(m-4)(m-2)} f^{\beta} \Delta\Gamma(f). \end{split}$$

Expanding the term $\Delta^2(f^{\beta+2})$ (recall that $\beta = \frac{2-m+d}{m-d}$), after a bit of algebra, we finally get,

$$\Gamma_2(f) \geqslant \frac{\beta+1}{d(\beta+1)-2\beta} (\Delta f)^2 - \beta \frac{\Gamma(f,\Gamma(f))}{f} - \frac{\beta(\beta-1)}{2} \frac{\Gamma(f)^2}{f^2}.$$
 (31)

First, observe that if $\beta = 0$ that is m = d + 2 then the previous inequality becomes

$$\Gamma_2(f) \geqslant \frac{1}{d} (\Delta f)^2,$$

in others words we recover the CD(0,d) condition of the Laplacian Δ .

It is more subtle to notice that, for any $\beta \in (-1,0)$ fixed, inequality (31) implies the same condition CD(0,d). It is more tricky since, taking $f=1+\varepsilon g$ with $\varepsilon \to 0$, then (31) already implies $\Gamma_2(g) \geqslant \frac{\beta+1}{d(\beta+1)-2\beta}(\Delta g)^2$ which is the $CD(0,\frac{d(\beta+1)-2\beta}{\beta+1})$ condition, weaker than the CD(0,d) condition, since $\frac{d(\beta+1)-2\beta}{\beta+1} \geqslant d$. The way to recover the CD(0,d) condition is to replace f by $\Phi(f)$ in (31), where Φ is a smooth function.

The diffusion properties in (31) one has,

$$\begin{split} \Phi'^2(\Gamma_2(f)-a(\Delta f)^2) + \Phi''\Gamma(f)^2(1-a) - c\frac{\Phi'^4}{\Phi^2}\Gamma(f)^2 + \Phi'\Phi''(\Gamma(f,\Gamma(f)) - 2a\Delta f\Gamma(f)) \\ - 2b\frac{\Phi'^2\Phi''}{\Phi}\Gamma(f)^2 - b\frac{\Phi'^3}{\Phi}\Gamma(f,\Gamma(f)) \geqslant 0, \end{split}$$

where $a = \frac{\beta+1}{d(\beta+1)-2\beta}$, $b = -2\beta$ and $c = -\beta(\beta-1)$ and for simplicity we omit the variable f in the function Φ . It implies that the quadratic form with respect to the variables Φ' , Φ'^2/Φ and Φ'' is positive, then the determinant of its matrix is positive. We get

$$\begin{split} \Gamma_2(f) \geqslant \frac{1}{d} (\Delta f)^2 + \left(\frac{-ac - a^2}{c(a-1) - b^2} - \frac{1}{d} \right) (\Delta f)^2 + \frac{-b^2 - c - b^2 a}{4(c(a-1) - b^2)} \left(\frac{\Gamma(f, \Gamma(f))}{\Gamma(f)} \right)^2 \\ + \frac{b^2 + ac}{c(a-1) - b^2} \frac{\Gamma(f, \Gamma(f))}{\Gamma(f)} \Delta f. \end{split}$$

After an unpleasant computation, the inequality is independent of β , and may be written as follow

$$\Gamma_2(f) \geqslant \frac{1}{d}(\Delta f)^2 + \frac{d}{d-1} \left(\frac{\Gamma(f,\Gamma(f))}{2\Gamma(f)} - \frac{1}{d}\Delta f\right)^2,$$

which is a reinforced CD(0,d) condition.

Indeed, in all these computations, we could have replaced the Laplace operator Δ in \mathbb{R}^d by any operator satisfying the CD(0,d) condition. Then we could have constructed the associated $Q_t^{(m)}$ operator through (19), and obtain also a similar family of Beckner inequalities for this $Q_t^{(m)}$, still equivalent to the starting CD(0,d) condition. The main difference in the Euclidean case is that then the family of inequalities for $Q_t^{(m)}(x)$ reduces one for $Q_1^{(m)}(0)$, through dilations and translations.

4.3 Φ-entropy inequalities for Cauchy type distribution

We can extend Beckner's inequalities to a general Φ -entropy inequality. For any convex function Φ on an interval $I \subset \mathbb{R}$, let us define the Φ -entropy for measure μ (or a kernel) and a function f on I,

$$\operatorname{Ent}_{\mu}^{\Phi}(f) = \int \Phi(f) d\mu - \Phi\Big(\int f d\mu\Big).$$

Section 4.1 can be easily generalized to Φ -entropy.

Definition 16 (n-Admissible functions) Let $\Phi: I \mapsto \mathbb{R}$ be a smooth function on an interval $I \subset \mathbb{R}$ and let n < 0. We say that Φ is n-admissible if $\Phi'' > 0$ on I and $(\Phi'')^{\frac{2-n}{n}}$ is concave. In other words, Φ is n-admissible if and only if $\Phi'' > 0$ and moreover

$$2\frac{n-1}{n}(\Phi^{(3)})^2 \le \Phi''\Phi^{(4)},$$

that is $\theta = \Phi''$ satisfies condition (11).

When n goes to $-\infty$, an $(-\infty)$ -admissible is just a function Φ such that, $\Phi'' > 0$ on I and $(1/\Phi'')$ is concave. Many Φ -entropy inequalities have been proved for such function φ under the curvature-dimension condition $CD(\rho,\infty)$ (see [Cha04, BG10]). In our computations, the case $n=-\infty$ is then similar to the case $n=+\infty$.

We can state the following result.

Theorem 17 (Φ -entropy inequality for $Q_t^{(m)}$) For any function Φ (d-m+2)-admissible on an interval I, the following holds, for any $m \ge d+2$,

$$\operatorname{Ent}_{Q_t^{(m)}}^{\Phi}(f) \le \frac{t^2}{2(m-2)} Q_t^{(m-2)}(\Phi''(f)\Gamma(f)), \tag{32}$$

for any smooth function f on I.

The inequality is optimal in the sense that if inequality (32) holds for a some function Φ , the constant $\frac{t^2}{2(m-2)}$ is the optimal one.

Corollary 18 (Φ -entropy inequality for Cauchy distribution) For any function Φ (d-m+2)-admissible on an interval I, the following holds, for $b \ge d+1$

$$\operatorname{Ent}_{\nu_b}^{\Phi}(f) \le \frac{1}{4(b-1)} \int \Phi''(f) \Gamma(f) (1+|x|^2) d\nu_b,$$

for any smooth function f on I. The inequality is optimal.

The proof is the same as the one of Theorem 11, using the sub-harmonicity inequality (12), so we skip it, the optimality being proved in a same way.

5 A Beckner-type inequality on the sphere

In this section we are looking for the link between the family of Beckner inequalities (12) for ν_b (or $Q_t^{(m)}$) and a family of Beckner inequalities on the sphere.

The unit sphere $\mathbb{S}^d \subset \mathbb{R}^{d+1}$ can be seen in \mathbb{R}^d through with the stereographic projection, with the carré du champ operator on the sphere

$$\Gamma_{\mathbb{S}}(f) = \frac{\rho^4}{4} \Gamma(f),$$

where $\rho(x) = \sqrt{1 + |x|^2}$ and the spherical measure

$$\mu_{\mathbb{S}}(dx) = \frac{1}{c(d,d)} \frac{1}{\rho^{2d}} dx,$$

where the normalization constant c(d, d) has been defined in (14) (see [BGL14, Sec 2.2]). The Laplace-Beltrami operator takes the form in the stereographic projection representation

$$\Delta_{\mathbb{S}} = \frac{(1+|x|^2)^2}{4} \Delta - \frac{d-2}{2} (1+|x|^2) \sum_{i=1}^d x_i \partial_i,$$

where Δ is the classical Laplacian in \mathbb{R}^d . The operator is symmetric in $L^2(\mu_{\mathbb{S}})$. It is interesting to recall that the map $\mathbb{S}^d \ni x \mapsto x_{d+1}$ is a eigenvector associated to the eigenvalue -d. This function, from the stereographic projection, takes the form

$$u(x) = \frac{1 - |x|^2}{1 + |x|^2}, \ x \in \mathbb{R}^d.$$
(33)

Then, u satisfies $\Delta_{\mathbb{S}}u = -du$ and $\Gamma_{\mathbb{S}^d}(u) = 1 - u^2$ since

$$\Gamma_{\mathbb{S}^d}(u) = \frac{(1+|x|^2)^2}{4} \Gamma\left(\frac{1-|x|^2}{1+|x|^2}\right) = \frac{1}{(1+|x|^2)^2} \Gamma(|x|^2) = \frac{4|x|^2}{(1+|x|^2)^2} = 1 - u^2.$$

Indeed this computation is much easer (and indeed completely elementary) using the standard representation in the sphere, see [BGL14, Sec 2.2].

We can now transform the Beckner inequality for $Q_t^{(m)}$ into a new Beckner inequality on the sphere. With these new notations, inequality (24) for t=1 and x=0 becomes

$$\frac{c(d,d)}{c(m,d)} \int f^{\beta+2} \rho^{d-m} d\mu_{\mathbb{S}} - \left(\frac{c(d,d)}{c(m,d)} \int f \rho^{d-m} d\mu_{\mathbb{S}}\right)^{\beta+2} \\
\leq \frac{2(\beta+1)(\beta+2)}{m-2} \frac{c(d,d)}{c(m-2,d)} \int f^{\beta} \Gamma_{\mathbb{S}}(f) \rho^{d-m-2} d\mu_{\mathbb{S}}.$$
(34)

We concentrate to the case $\beta = \frac{d+2-m}{m-d}$, the limit of the admissible interval. Let $f = \rho^{m-d}g$ and we compute the various terms. First, we have form diffusion properties of $\Gamma_{\mathbb{S}}$,

$$f^{\beta}\Gamma_{\mathbb{S}}(f)\rho^{d-m-2} = g^{\beta}\rho^{2(d-m)}\Gamma_{\mathbb{S}}(\rho^{m-d}g) = g^{\beta}\Gamma_{\mathbb{S}}(g) + 2(m-d)g^{\beta+1}\Gamma_{\mathbb{S}}(\log\rho, g) + (m-d)^2g^{\beta+2}\Gamma_{\mathbb{S}}(\log\rho),$$

so that

$$\begin{split} \int f^{\beta} \Gamma_{\mathbb{S}}(f) \rho^{d-m-2} d\mu_{\mathbb{S}} &= \\ &\int g^{\beta} \Gamma_{\mathbb{S}}(g) d\mu_{\mathbb{S}} + (m-d) \int \frac{2}{\beta+2} \Gamma_{\mathbb{S}}(\log \rho, g^{\beta+2}) d\mu_{\mathbb{S}} + (m-d)^2 \int g^{\beta+2} \Gamma_{\mathbb{S}}(\log \rho) d\mu_{\mathbb{S}}. \end{split}$$

With an integration by parts, we get

$$\int \Gamma_{\mathbb{S}}(\log \rho, g^{\beta+2}) d\mu_{\mathbb{S}} = -\int g^{\beta+2} \Delta_{\mathbb{S}} \log \rho d\mu_{\mathbb{S}},$$

and then

$$\int f^{\beta} \Gamma_{\mathbb{S}}(f) \rho^{d-m-2} d\mu_{\mathbb{S}} = \int g^{\beta} \Gamma_{\mathbb{S}}(g) d\mu_{\mathbb{S}} + \int g^{\beta+2} K d\mu_{\mathbb{S}},$$

where

$$K = (m - d)^{2} \left(2 \frac{\Delta_{\mathbb{S}} \log \rho}{d - m - 2} + \Gamma_{\mathbb{S}} (\log \rho) \right),$$

since $\beta + 2 = (d - m - 2)/(d - m)$. Secondly we have

$$\int f^{\beta+2} \rho^{d-m} d\mu_{\mathbb{S}} = \int g^{\beta+2} \rho^2 d\mu_{\mathbb{S}}$$

and

$$\int f \rho^{d-m} d\mu_{\mathbb{S}} = \int g d\mu_{\mathbb{S}}.$$

Then, inequality (34) can be written as

$$\frac{c(d,d)}{c(m,d)} \int g^{\beta+2} \rho^2 d\mu_{\mathbb{S}} - \left(\frac{c(d,d)}{c(m,d)} \int g d\mu_{\mathbb{S}}\right)^{\beta+2} \\
\leq \frac{2(\beta+1)(\beta+2)}{m-2} \frac{c(d,d)}{c(m-2,d)} \left(\int g^{\beta} \Gamma_{\mathbb{S}}(g) d\mu_{\mathbb{S}} + \int g^{\beta+2} K d\mu_{\mathbb{S}}\right),$$

or

$$\int g^{\beta+2} R d\mu_{\mathbb{S}} - \left(\frac{c(d,d)}{c(m,d)} \int g d\mu_{\mathbb{S}}\right)^{\beta+2} \leq \frac{2(\beta+1)(\beta+2)}{m-2} \frac{c(d,d)}{c(m-2,d)} \int g^{\beta} \Gamma_{\mathbb{S}}(g) d\mu_{\mathbb{S}},$$

with

$$R = \frac{c(d,d)}{c(m,d)}\rho^2 - \frac{2(\beta+1)(\beta+2)}{m-2} \frac{c(d,d)}{c(m-2,d)} K.$$

From (16), the definition of β and the fact that $\rho^2 = 2/(1+u)$ where u has been defined in (33), we can write

$$R = \frac{c(d,d)}{c(m,d)} \left(\frac{2}{1+u} - 2\frac{m-d+2}{(m-d)^2} \frac{K}{m-2+d} \right),$$

and the inequality

$$\int g^{\beta+2}Rd\mu_{\mathbb{S}} - \left(\frac{c(d,d)}{c(m,d)}\int gd\mu_{\mathbb{S}}\right)^{\beta+2} \leq \frac{4(m-d+2)}{(m-d)^2(m-2+d)}\frac{c(d,d)}{c(m,d)}\int g^{\beta}\Gamma_{\mathbb{S}}(g)d\mu_{\mathbb{S}},$$

We need now to compute $\Delta_{\mathbb{S}} \log \rho$ and $\Gamma_{\mathbb{S}}(\log \rho)$. We have

$$\begin{split} \Delta_{\mathbb{S}}(\log \rho) &= \frac{1}{2} \Delta_{\mathbb{S}} \log \rho^2 = -\frac{1}{2} \Delta_{\mathbb{S}} \log (1+u) = \frac{d}{2(1+u)} u + \frac{1}{2(1+u)^2} \Gamma_{\mathbb{S}^d}(u) \\ &= \frac{d}{2(1+u)} u + \frac{1}{2(1+u)^2} (1-u^2). \end{split}$$

Then, finally

$$\Delta_{\mathbb{S}}(\log \rho) = \frac{1 + u(d-1)}{2(1+u)},$$

and

$$\Gamma_{\mathbb{S}}(\log \rho) = \frac{1 - u}{4(1 + u)}.$$

If we plug those values on the definition of K, then a miracle occurs. The function R is constant,

$$R = \frac{c(d,d)}{c(m,d)} \frac{3d + m - 2}{d + m - 2}.$$

Setting now $f^2 = g^{\beta+2} = g^{\frac{d-m-2}{d-m}}$, we obtain

Theorem 19 (Beckner-type inequalities on the sphere) For any smooth nonnegative function f on the sphere \mathbb{S}^d , and $m \ge d+2$, we have

$$\int f^2 d\mu_{\mathbb{S}} \le A \left(\int f^{2/p} d\mu_{\mathbb{S}} \right)^p + \frac{16}{(m+2-d)(3d-2+m)} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}, \tag{35}$$

where

$$p = 1 + \frac{2}{m - d} \in (1, 2]$$

and

$$A = \left(\frac{c(d,d)}{c(m,d)}\right)^{\frac{2}{m-d}} \frac{m+d-2}{m+3d-2}.$$

Remarks 20

- The function R is constant only for the parameter $p = 1 + \frac{2}{m-d}$.
- When m > d+2, A > 1. In that case inequality (35) is not tight, in the sense that when g = 1, we do not have an equality. But the inequality is still optimal since it is saturated for $f = \rho^{\frac{d-m-2}{2}}$.
- When m = d + 2, then A = 1, p = 2 and inequality (35) is nothing else than the optimal Poincaré inequality on the sphere,

$$\int f^2 d\mu_{\mathbb{S}} - \left(\int f^2 d\mu_{\mathbb{S}}\right)^2 \le \frac{1}{d} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}.$$

• The classical Beckner inequalities hold for the spherical model (see e.g. [BGL14, Rmk 6.8.4]): for any $p \in (1,2]$

$$\int f^2 d\mu_{\mathbb{S}} \le \left(\int f^{2/p} d\mu_{\mathbb{S}} \right)^p + \frac{p-1}{pd} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}. \tag{36}$$

Inequality (36) is optimal in the sense that $\frac{p-1}{pd}$ is the best constant. Only constant functions saturate inequality (36) for $p \in (1,2)$. And when p goes to 1, inequality (36) provides the Logarithmic Sobolev inequality on the sphere

$$\int f^2 \log \frac{f^2}{\int f^2 d\mu_{\mathbb{S}}} d\mu_{\mathbb{S}} \le \frac{2}{d} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}},$$

proved in [MW82] (see also [BGL14, Thm. 5.7.4]). On the other hand, this last logarithmic Sobolev inequality implies back the full family (36), $p \in (1,2]$.

Let us mention that inequality (36) has been generalized by in [DEKL14] in the following way,

$$\Psi\left(\frac{1-\int f^{2/p}d\mu_{\mathbb{S}}}{p-1}\right) \le \int \Gamma_{\mathbb{S}}(f)d\mu_{\mathbb{S}},$$

for every function f such that $\int f^2 d\mu_{\mathbb{S}} = 1$ and for some explicit function Ψ .

• So (35) appears as a new and optimal inequality on the sphere. It is probably worth to compare it to the del Pino-Dolbeault family of optimal Gagliardo-Nirenberg inequalities,

$$||f||_{2\frac{a-1}{a-2}} \le C ||\nabla f||_2^{\theta} ||f||_{\frac{2a}{a-2}}^{1-\theta}, \quad a \geqslant d,$$

where θ is a fixed by scaling properties, see [DD02].

• The Taylor expansion (when $m \to \infty$) of the constant A is given by

$$A = 1 + d \frac{\log(m)}{m} + \frac{C}{m} + o(1/m),$$

where C is a constant depending on the dimension d. This will be used in the rest of the section.

Even if we are not able from Theorem 19 to reach directly the Logarithmic Sobolev inequality, inequality (35) contains enough information to obtain a Sobolev inequality for the spherical model, however with a non optimal constant.

For this, we are going to obtain from (35) a Nash inequality on the sphere.

First, it is not difficult to see that there exits two constants $\alpha, \beta > 0$ depending only on d, such that, for any $m \ge d + 2$,

$$A \le \alpha^{\frac{2}{m-d}} m^{\frac{d}{m-d}}.$$

which is just a precised form of the previous Taylor expansion, and

$$\frac{16}{(m+2-d)(3d-2+m)} \le \frac{\beta}{m^2}.$$

From Hölder's inequality, if f is nonnegative and $p \in [1, 2]$,

$$\left(\int f^{2/p} d\mu_{\mathbb{S}}\right)^{p} \leq \left(\int f d\mu_{\mathbb{S}}\right)^{p-1} \left(\int f^{2} d\mu_{\mathbb{S}}\right)^{2-p}$$

with $p = 1 + 2/(m - d) \in [1, 2]$. Inequality (35) becomes, for any nonnegative function f such that $\int f d\mu_{\mathbb{S}} = 1$, and any $m \ge m + 2$,

$$\int f^2 d\mu_{\mathbb{S}} \leq \alpha^{\frac{2}{m-d}} m^{\frac{d}{m-d}} \left(\int f^2 d\mu_{\mathbb{S}} \right)^{1-\frac{2}{m-d}} + \frac{\beta}{m^2} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}.$$

Let define for $x \ge 0$, the map

$$\varphi_m(x) = \alpha^{\frac{2}{m-d}} m^{\frac{d}{m-d}} x^{1-\frac{2}{m-d}} + \frac{\beta}{m^2} E,$$

where $E = \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}$. The function φ_m as a unique fixed point $x_m > 0$, and the previous inequality implies that $\int f^2 d\mu_{\mathbb{S}} \leq x_m$.

Let us proved that for any $m \ge d + 2$,

$$x_m \le \alpha m^{d/2} \left(1 + \frac{\beta E}{2\alpha m^{1+d/2}} \right) = a_m.$$

To prove such an inequality, it is enough to prove that, for any $m \ge d + 2$, $\varphi_m(a_m) \le a_m$. We have

$$\varphi_m(a_m) = \alpha \, m^{d/2} \left(1 + \frac{\beta E}{2\alpha m^{1+d/2}} \right)^{1 - \frac{2}{m-d}} + \frac{\beta}{m^2} E$$

$$\leq \alpha \, m^{d/2} \left[\left(1 + \frac{\beta E}{2\alpha m^{1+d/2}} \right)^{1 - \frac{2}{m}} + \frac{\beta}{m^{2 + \frac{d}{2}}} E \right] \leq \alpha \, m^{d/2} \left(1 + \frac{\beta E}{2\alpha m^{1+d/2}} \right) = a_m,$$

where we used the inequality $(1+x)^u \le 1+ux$, for x>0 and $u\in(0,1)$. Modifying the numerical constants (depending only on d), we see that for any $m\geqslant d+2$,

$$\int f^2 d\mu_{\mathbb{S}} \le C m^{d/2} + \frac{C}{m^{d/2}} \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}.$$

Optimizing with respect to the parameter $m \ge d+2$, we have obtained, for some other constant C > 0,

$$\begin{split} \int f^2 d\mu_{\mathbb{S}} &\leq C \left(\left(\int f d\mu_{\mathbb{S}} \right)^2 + \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}} \right)^{\frac{d}{d+2}} \left(\int f d\mu_{\mathbb{S}} \right)^{\frac{2}{d+2}} \\ &\leq C \left(\int f^2 d\mu_{\mathbb{S}} + \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}} \right)^{\frac{d}{d+2}} \left(\int f d\mu_{\mathbb{S}} \right)^{\frac{2}{d+2}}, \end{split}$$

from Jensen inequality.

This last inequality is a so-called Nash inequality, and is known to be equivalent to the Sobolev inequality (up to the optimal constant) for the spherical model, see [BGL14, Prop. 6.2.3],

$$\left(\int f^{\frac{2d}{d-2}}d\mu_{\mathbb{S}}\right)^{\frac{d-2}{d}} \le C\int f^2d\mu_{\mathbb{S}} + C\int \Gamma_{\mathbb{S}}(f)d\mu_{\mathbb{S}}.$$

From the Poincaré inequality (m = d + 2) and [BGL14, Prop. 6.2.2], this inequality implies a tight Sobolev inequality. We have obtained

Theorem 21 (From Beckner inequalities to Sobolev inequality on the sphere) From the family of Beckner type inequalities (35), one can prove that there exists a constant C > 0 depending only on d such that, for any smooth and nonnegative function f on the sphere \mathbb{S}^d ,

$$\left(\int f^{\frac{2d}{d-2}} d\mu_{\mathbb{S}}\right)^{\frac{d-2}{d}} \le \int f^2 d\mu_{\mathbb{S}} + C \int \Gamma_{\mathbb{S}}(f) d\mu_{\mathbb{S}}.$$
 (37)

Conclusion: The family of optimal inequalities (35) appears as a new form of optimal inequalities on the sphere, implying a Sobolev inequality.

Acknowledgements. This work was partly written while the two first authors were visiting Institut Mittag-Leffler in Stockholm; it is a pleasure for them to thank this institution for its kind hospitality.

This research was supported by the French ANR-17-CE40-0030 EFI project.

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